



# Direct and Indirect Measurements in Nuclear Astrophysics



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## Outline

- Nuclear Physics in Astrophysics
- Feature of thermonuclear reactions and experimental requirements in direct measurements
- Indirect measurements
- Direct vs. Indirect measurements:  $^{12}\text{C}+^{12}\text{C}$  fusion as case study

## Nuclear Astrophysics

Rich & Diverse Interdisciplinary Field bringing together

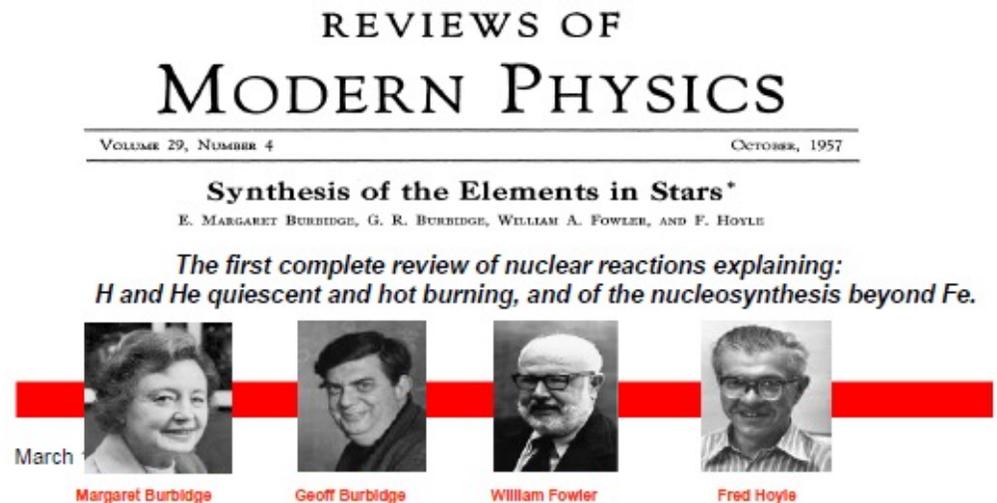
- Modelers
- Observers
- **Nuclear physicists**: Experimentalists as well as Theorists

Nucleosynthesis in 1900



... to the seminal **B<sup>2</sup>FH** review paper of 1957,  
the basis of the modern nuclear astrophysics

this work has been considered as the greatest gift of astrophysics to modern civilization



Nuclear reactions responsible for both ENERGY PRODUCTION and CREATION OF ELEMENTS in 4 ways/environments:

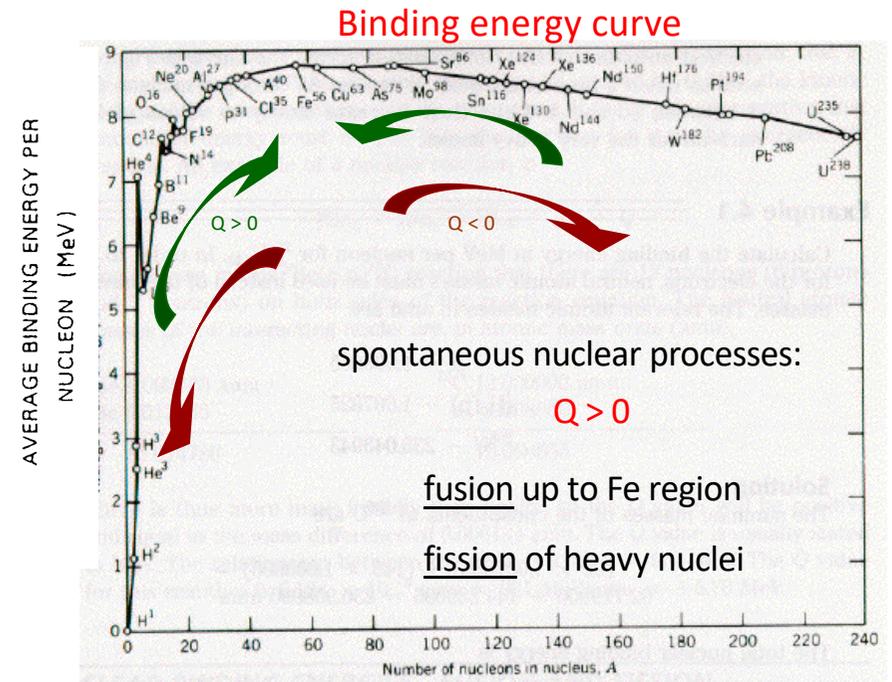
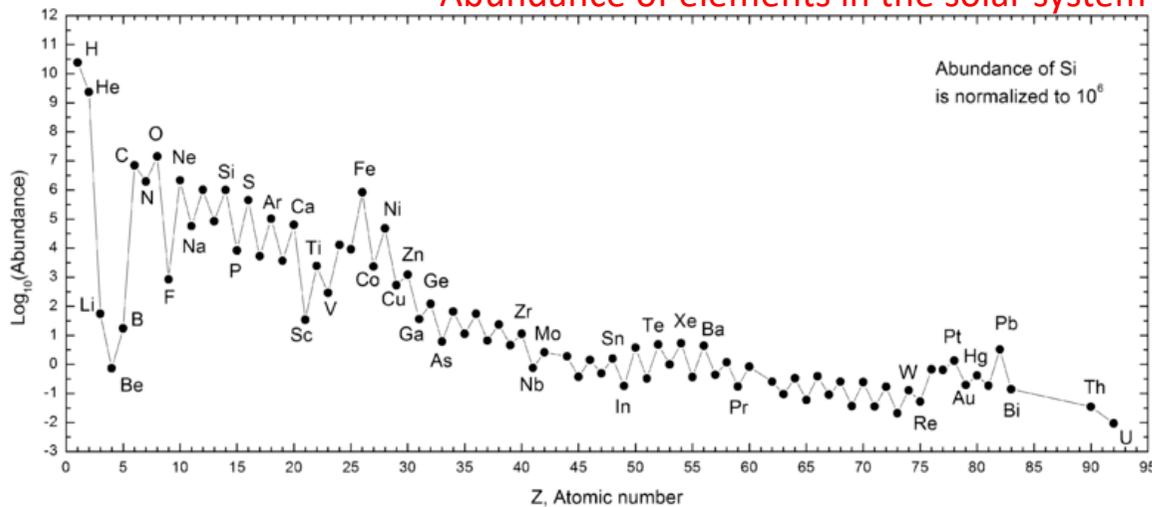


# Where's the Nuclear Physics?

- H burning → conversion of H to He
- He burning → conversion of He to C, O ...
- C, O and Ne burning → production of A: 16 to 28
- Si burning → production of A: 28 to 60
- s-, r- and p-processes → production of A>60
- Li, Be, and B from cosmic rays

- Big Bang Nucleosynthesis does not go beyond Li due to missing stable nuclei of mass number 5 or 8
- Odd-even staggering of abundances (Oddo-Harkins rule)
- Larger alpha-nuclei abundance, particularly those connected to particular values of Z and N (so called magic numbers, 2, 8, 20, 28, 50 ...) which are significant with regard to the structure of nuclei ... at least up to Fe
- Broad peak around Fe

Abundance of elements in the solar system

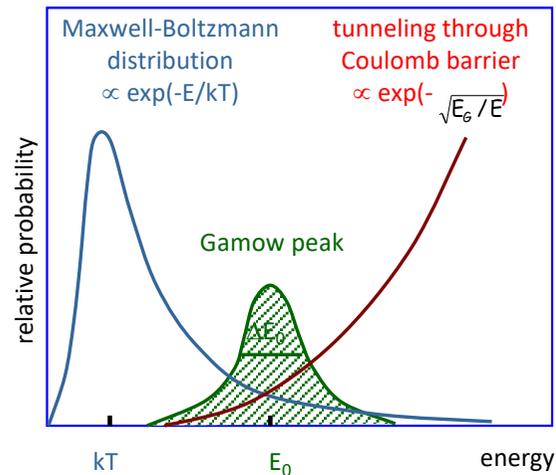


# Nuclear Astrophysics

Astrophysical energies for reactions between charged particles are determined by the Gamow peak: the most effective energy region for thermonuclear reactions

The Gamow energy  $E_0 = f(Z_1, Z_2, T)$

varies depending on the reaction and/or the temperature, usually from tens to hundreds of keV, but also MeV.



Examples:  $T \sim 15 \times 10^6 \text{ K}$  ( $T_6 = 15$ )

reaction	Coulomb barrier (MeV)	$E_0$ (keV)	area under Gamow peak $\sim \langle \sigma v \rangle$
p + p	0.182	5.9	$7.0 \times 10^{-6}$
$\alpha + {}^{12}\text{C}$	2.242	56	$5.9 \times 10^{-56}$
${}^{16}\text{O} + {}^{16}\text{O}$	10.349	237	$2.5 \times 10^{-237}$

$$kT \ll E_0 \ll E_{\text{coul}}$$

$10^{-18} \text{ barn} < \sigma < 10^{-9} \text{ barn}$  major experimental challenges

## General features: Quiescent vs. Explosive Burning

Feature	Quiescent Burning (e.g., Main Sequence, Red Giant)	Explosive Burning (e.g., Supernovae, X-ray Bursts)
Temperature (T)	$10^6 - 10^8$ K	$> 10^9$ K
Energy ( $E_0$ )	10 keV – 1 MeV ( $E_0 \ll E_{Coul}$ )	$\sim$ MeVs ( $E_0 \leq E_{Coul}$ )
Cross-Section ( $\sigma$ )	$10^{-18}$ barn $< \sigma < 10^{-9}$ barn	$10^{-6}$ barn $< \sigma < 10^3$ barn
Interaction Time ( $\tau$ )	$\sim \langle \sigma v \rangle^{-1} \sim 10^9$ years	$\sim \langle \sigma v \rangle^{-1} \sim$ seconds
Unstable Species	DO NOT play significant role	DO play significant role

## Shared Experimental Challenges

- **Low Signal-to-Noise Ratio** (Common across both due to low cross-sections or low beam intensities)
- **Unknown Nuclear Properties** (Especially for exotic, unstable nuclei)
- **Low Beam Intensities** (Particularly for Radioactive Ion Beams - RIBs, which are several orders of magnitude lower than stable beams)

## Experimental Requirements

- **Extrapolation Procedures** (Crucial for quiescent, less so for explosive) using  $S(E) = E\sigma(E) \exp(2\pi\eta)$
- **High Beam Intensities / RIBs production & acceleration** (High for quiescent, RIBs essential for explosive)
- **Ultra Pure Targets**
- **Long Measurements High Detection Efficiency** (Essential for both)
- **Storage Rings**

# Experimental approach: alternative solutions

- Underground experiments to reduce (cosmic) background: LUNA (LNGS Italy), Felsenkeller (Germany), CASPAR (USA), JUNA (China), particularly suited to perform gamma spectroscopy

- Surface experiments:

inverse kinematics: accelerating heavier nuclei into lighter targets.

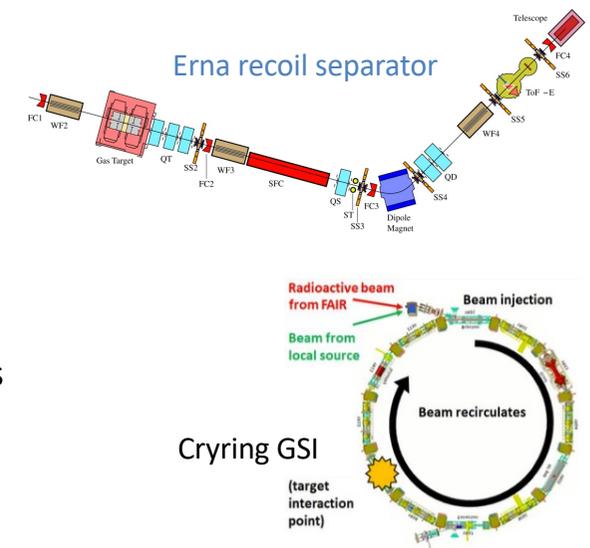
coincidence experiments (g-g, g-particle, ...);

recoil separators: to separate reaction products from unreacted beam and

disperse them according to their mass-to-charge-state ratio;

storage rings: to overcome beam intensity limitations. Recirculate the beam many times for repeated target interactions.

...



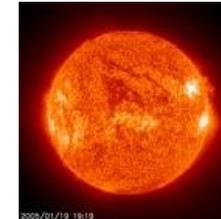
- Use indirect methods: Coulomb Dissociation (CD), Asymptotic Normalization Coefficients (ANC), Trojan Horse Method (THM) ...

Dedicated Talks for some of these items already this afternoon

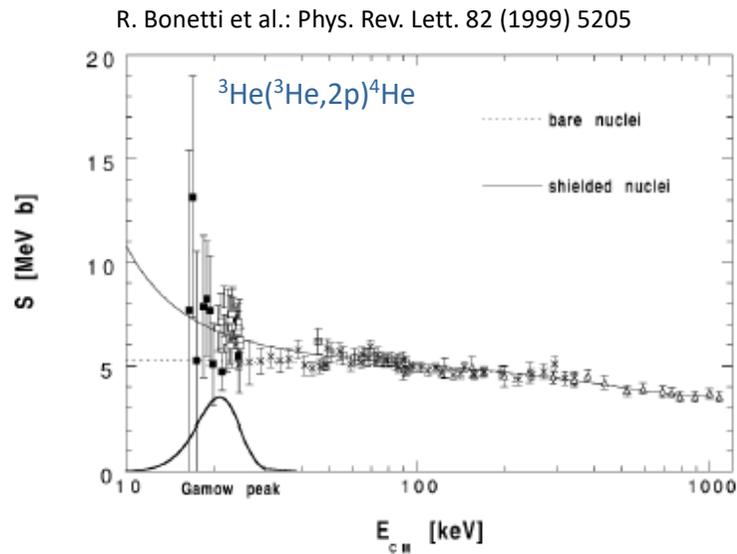
To give you the feeling what low signal-to-noise ratio means



LUNA – Phase I: 50 kV accelerator (1992-2001)

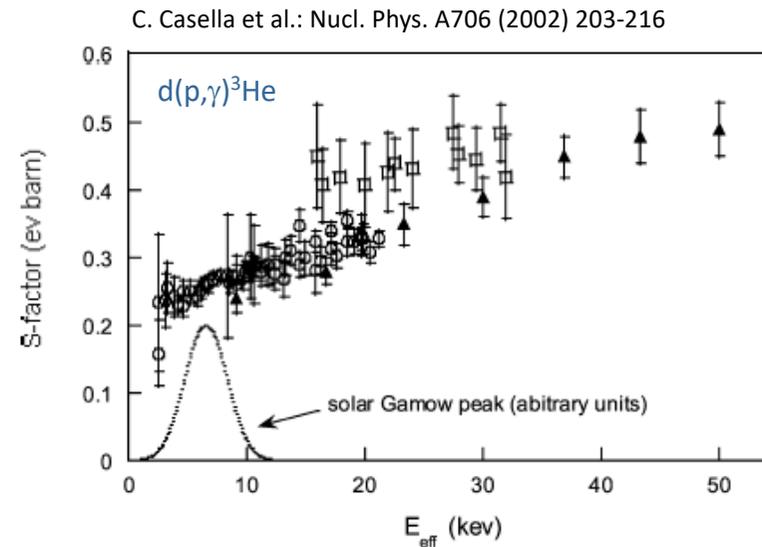


investigate reactions in solar pp chain



@ lowest energy:

$\sigma \sim 20 \text{ fb} \rightarrow 1 \text{ count/month}$



@ lowest energy:

$\sigma \sim 9 \text{ pb} \rightarrow 50 \text{ counts/day}$

only two reactions studied directly at the Gamow peak

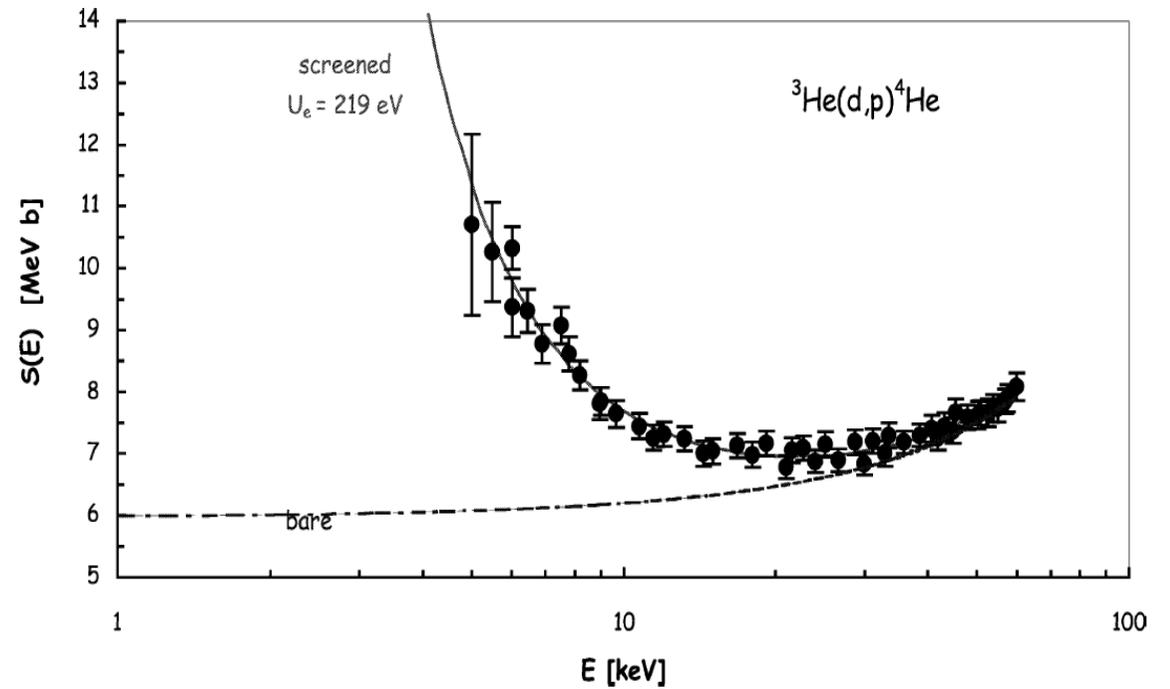
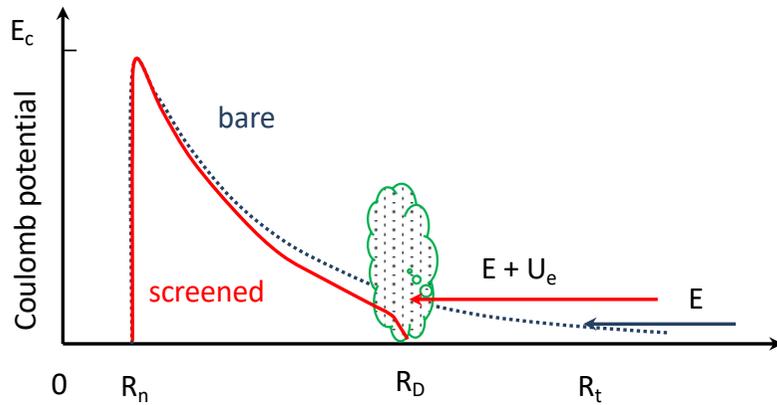
...but... intrinsic limitation at astrophysical energies

→ → → →

Electron Screening

$S(E)$  experimental enhancement due to the electron screening

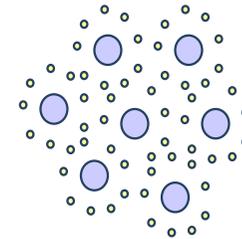
$$S(E)_s = S(E)_b \exp(\pi\eta U_e/E)$$



## Electron Screening

In astrophysical plasma:

- the screening, due to free electrons in plasma, can be different  
→ we need  $S(E)_b$  to evaluate reaction rates



Debye-Hückel radius

$$R_D \sim (kT/\rho)^{1/2}$$

A theoretical approach to extract the electron screening potential  $U_e$  in the laboratory is needed



... however, experimental studies of reactions involving light nuclides have shown that the observed exponential enhancement of the cross section at low energies were in all cases significantly larger (about a factor of 2) than it could be accounted for from available atomic-physics model, i.e. the adiabatic limit  $(U_e)_{ad}$  ... screening yet to be fully understood

→ No way to measure  $S(E)_b$  from direct experiments at energies where screening is important

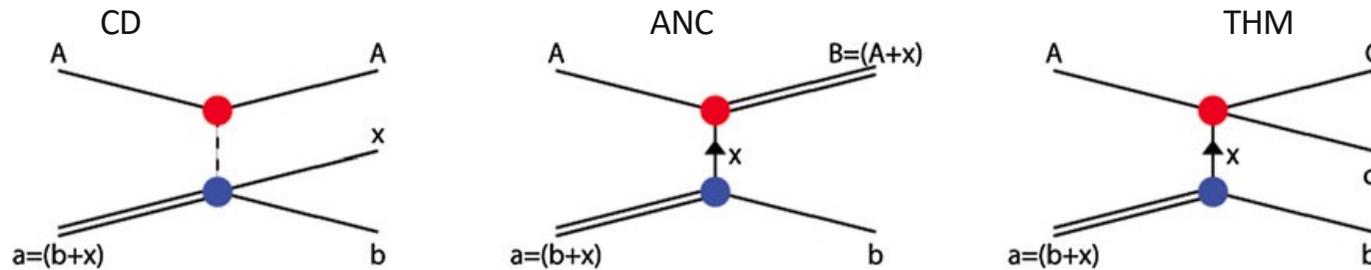
$S_b(E)$ -factor extracted from extrapolation of higher energy data

# Indirect Methods for Nuclear Astrophysics

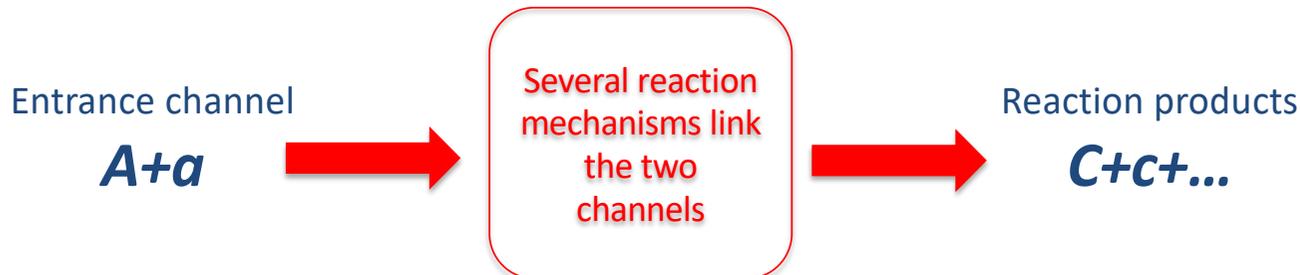
- to measure cross sections at never reached energies (no Coulomb suppression), where the **signal is below current detection sensitivity**
- to get independent information on  $U_e$
- to overcome difficulties in producing the beam or the target (radioactive ions, neutrons..)

## Coulomb Dissociation (CD), Asymptotic Normalization Coefficient (ANC) and Trojan Horse Methods (THM)

They share some common features: the astrophysically relevant two-body reaction at low energies is replaced by a high-energy reaction usually with a three-body final state. In all cases a virtual particle (photon  $\gamma$  or nucleus  $x$ ) is transferred between the two subsystems in the reaction.



Quite straightforward experiment, no Coulomb suppression, no electron screening but ...



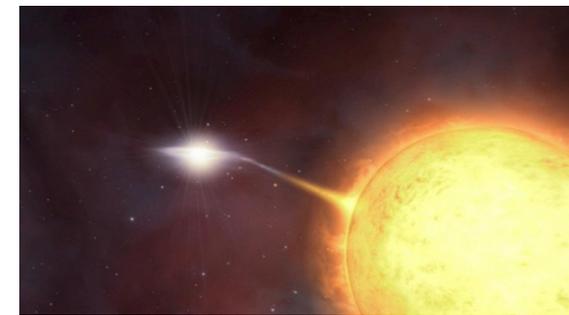
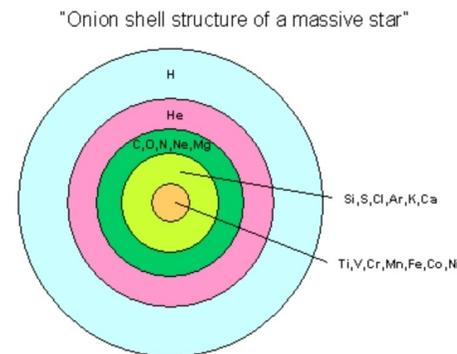
The reaction theory is needed to select only one reaction mechanism. However, nowadays powerful techniques and observables for careful data analysis and theoretical investigation.

# C-burning

Relevant temperature  $T \sim 0.1-1.2$  GK

It powers scenarios that influence the fate of stars, such as:

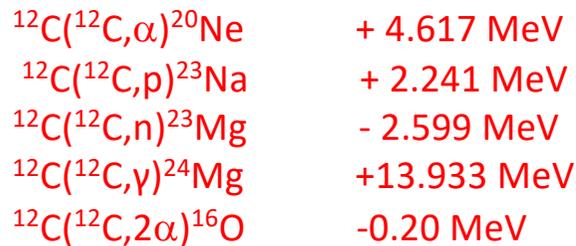
- **Late evolutionary stages of massive stars** ( $> 8 M_{\odot}$ ). It influences  $M_{\text{up}}$ , threshold in mass for hydrostatic carbon burning to occur (relevant  $T \sim 1$  GK)
  - **Superbursts** from accreting neutron stars ( $T \sim 0.15-0.7$  GK): C-fusion triggers (?) the ignition in the burning ashes of hydrogen and helium on the surface of accreting neutron stars
  - **Type Ia supernovae** ( $0.15-0.7$  GK and  $\rho \sim (2-5) 10^9$  g/cm<sup>3</sup>): interpreted as the consequence of explosive carbon burning ignited near the core of the white dwarf star in a binary system
- the C/O ratio influences the nucleosynthesis and, in turn, the resulting light curve



# 12C+12C fusion

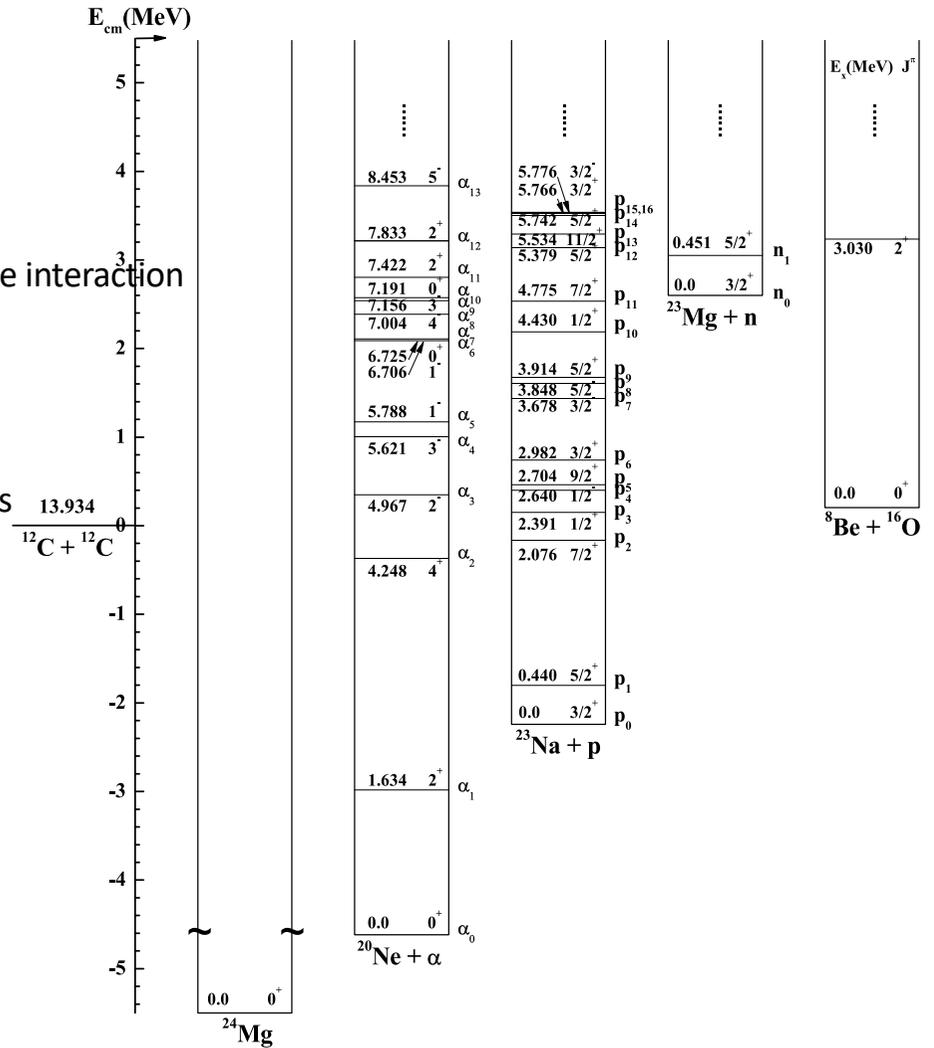
Carbon burning temperature from 0.1 to 1.2 GK →  $E_{cm}$  from 1 to 3 MeV

Principal reactions:



Considerable efforts to measure the  $^{12}\text{C}+^{12}\text{C}$  cross section at astrophysical energies

M.G. Mazarakis & W.E. Stephensen, Phys. Rev. C 7 1280 (1973)  
 K. U. Kettner *et al.*, Phys. Rev. Lett. 38, 337 (1977)  
 H.W. Becker *et al.*, Z. Phys. A 303, 305 (1981)  
 L. Barron-Palos *et al.*, Nucl. Phys. A 779, 318 (2006)  
 E.F. Aguleira *et al.*, Phys. Rev. C 73, 064601 (2006)  
 T. Spillane *et al.*, Phys. Rev. C 73, 064601 (2006)  
 T. Spillane *et al.*, Phys. Rev. Lett. 98, 122501 (2007)  
 B. Bucher *et al.*, Phys. Rev. Lett. 114, 251102 (2015)  
 C.L. Jiang *et al.*, Phys. Rev. C 97 012801 (2018)  
 J. Zickefoose *et al.*, Phys. Rev. C 97, 065806 (2018)  
 A. Tumino *et al.*, Nature (2018)  
 G. Fruet *et al.* PRL, 124 192701 (2020)  
 W.P. Tan *et al.* PRL, 124 192702 (2020), W.P. Tan *et al.* PRC (2024)  
 L. Morales-Gallegos *et al.* EPJ A (2024) ....



# $^{12}\text{C}+^{12}\text{C}$ : recent experiments

Jiang et al. 2018: down to  $E_{\text{c.m.}} = 2.84 \text{ MeV}$  and  $2.96 \text{ MeV}$  for the p and  $\alpha$  channels, respectively, using a sphere array of 100 Compton-suppressed Ge detectors in coincidence with silicon detectors.

Fruet et al. 2020: down to  $E_{\text{c.m.}}=2.16 \text{ MeV}$  using the particle- $\gamma$  coincidence technique and thin C-target. Charged particles were detected using annular silicon strip detectors, while  $\gamma$ -ray detection was accomplished with an array of LaBr3(Ce) scintillators. Only the  $p_1$  and  $\alpha_1$  channels. Total  $S^*$  factor reconstructed taking not observed branchings from the literature

Tan et al. 2020 and 2024: down to  $E_{\text{c.m.}}=2.2 \text{ MeV}$  using the particle- $\gamma$  coincidence technique and thick target. In particular, p and  $\alpha$ s were detected using a silicon detector array, and  $\gamma$ -rays with HPGe detectors. Only the  $p_1$  and  $\alpha_1$  channels. Total  $S^*$  factor reconstructed taking not observed branchings from the literature.

Tumino et al. 2018: THM measurement down to  $0.8 \text{ MeV}$  for the  $p_{0,1}$  and  $\alpha_{0,1}$  channels. Coincidence experiment using the  $^{14}\text{N}+^{12}\text{C}$  reaction at 30 MeV of beam energy.

Morales Gallegos et al. 2023: down to  $E_{\text{c.m.}} = 2.4 \text{ MeV}$  and  $2.9 \text{ MeV}$  for the p and  $\alpha$  channels, respectively, using telescopes (GASTLY modules) made up of Cl+Si detectors and thick target. Problems with background at the lowest energies.

Most of direct measurements: **particle- $\gamma$  coincidence technique** to overcome limitations due to target purity and beam-induced/cosmic rays background **but no access to the ground state transitions**

**ground state transitions are crucial, as these channels contribute significantly at stellar energies**

# Our Experiment with the THM

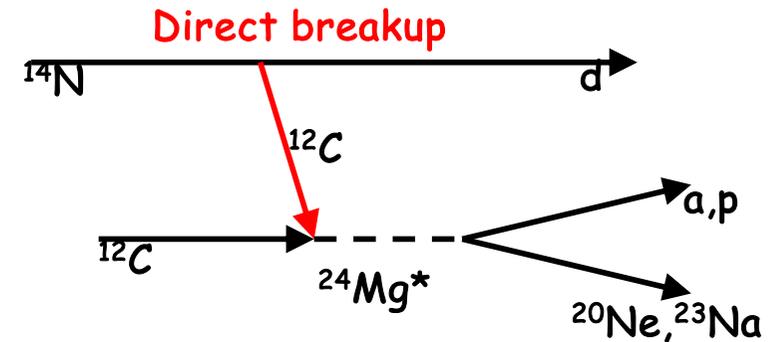
$^{12}\text{C}(^{12}\text{C},\alpha)^{20}\text{Ne}$  and  $^{12}\text{C}(^{12}\text{C},\text{p})^{23}\text{Na}$  reactions via the Trojan Horse Method applied to the  $^{12}\text{C}(^{14}\text{N},\alpha)^{20}\text{Ne}^2\text{H}$  and  $^{12}\text{C}(^{14}\text{N},\text{p})^{23}\text{Na}^2\text{H}$  three-body processes

Observation of  $^{12}\text{C}$  cluster transfer in the  $^{12}\text{C}(^{14}\text{N},\text{d})^{24}\text{Mg}^*$  reaction  $^2\text{H}$  from the  $^{14}\text{N}$  as spectator s  
(R.H. Zurmühle et al. PRC 49(1994) 5)

## QUASI-FREE MECHANISM

- ✓ only  $^{12}\text{C}$  -  $^{12}\text{C}$  interaction
- ✓ d = spectator

$$E_{\text{QF}} = E_{^{14}\text{N}} \frac{m_{^{12}\text{C}}}{m_{^{14}\text{N}}} \cdot \frac{m_{^{12}\text{C}}}{m_{^{12}\text{C}} + m_{^{12}\text{C}}} - 10.27 \text{ MeV}$$



$E_{^{14}\text{N}} = 30 \text{ MeV} > E_{\text{Coul}} \Rightarrow$  NO Coulomb barrier in the entrance channel  
NO electron screening

## PWIA

$$\frac{d^3\sigma}{d\Omega_B d\Omega_b dE_B} \propto \text{KF} \left| \Phi(p_{xs}) \right|^2 \left[ \frac{d^2\sigma_{xA \rightarrow bB}}{dE_{xA} d\Omega} \right]_{\text{HOES}}$$

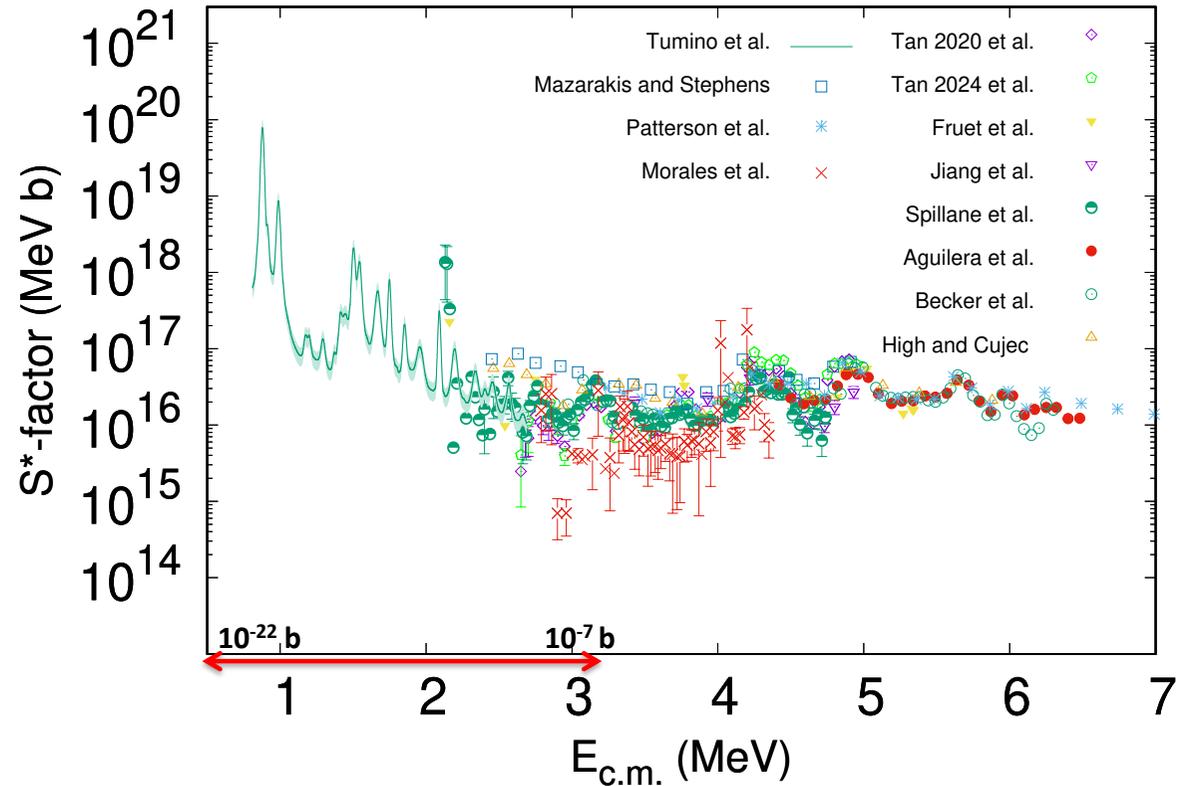
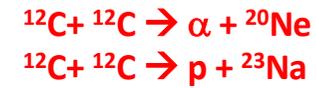
kinematical factor
momentum distribution of s inside a
Nuclear cross section for the  $A+x \rightarrow C+c$  reaction

$$\frac{d^2\sigma_{xA \rightarrow c'}}{dE_{xA} d\Omega_s} = \text{NF} \sum_i (2J_i + 1) \left| \frac{\sqrt{k_{c'}}}{\mu_{c'}} \frac{\sqrt{2P_{c'}} M_i(p_{xA} R_{xA}) \gamma_{xA}^i \gamma_{c'}^i}{D_i(E_{xA})} \right|^2$$

From the modified R-matrix approach assuming non-interfering resonances

R-matrix fits on all channels at the same time in the full energy range of interest

## 12C+12C comprehensive figure



Next step:

- direct data below 2 MeV (LUNA MV, Stella collaboration ...)
- improve the normalization of THM data to direct ones with larger overlap

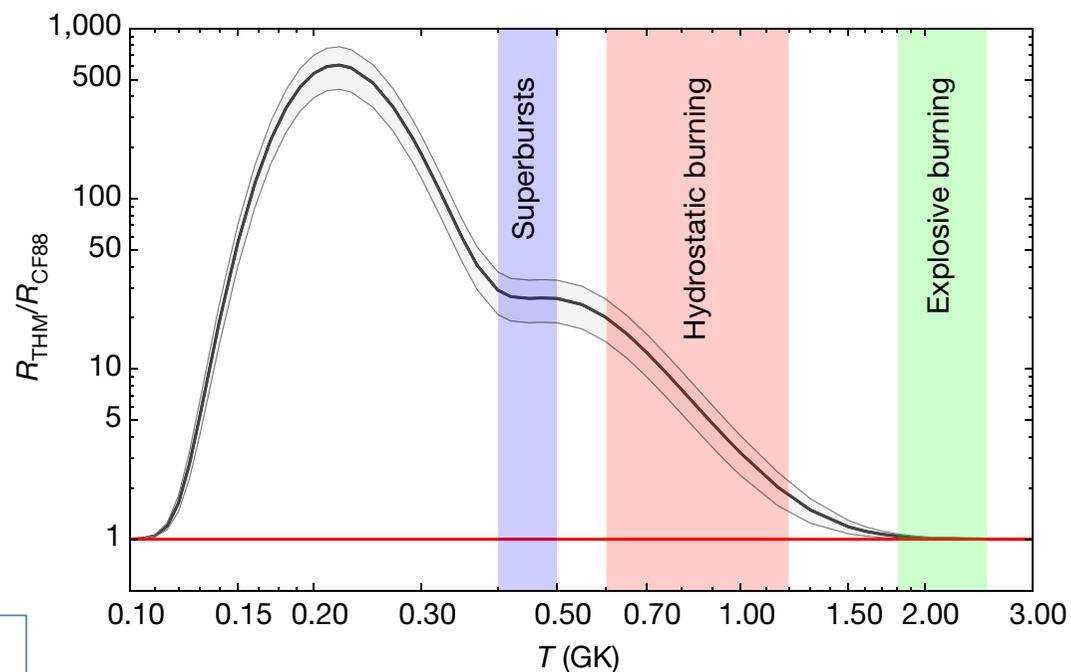
## An increase in the $^{12}\text{C} + ^{12}\text{C}$ fusion rate from resonances at astrophysical energies

A. Tumino , C. Spitaleri, M. La Cognata, S. Cherubini, G. L. Guardo, M. Gulino, S. Hayakawa, I. Indelicato, L. Lamia, H. Petrascu, R. G. Pizzone, S. M. R. Puglia, G. G. Rapisarda, S. Romano, M. L. Sergi, R. Spartá & L. Trache

Nature **557**, 687–690 (2018) | [Download Citation](#)

### 12C+12C Reaction Rate

Color shadings mark typical regions for C-burning



Compared to CF88, the present rate increases from a factor of 1.18 at 1.2 GK to a factor of more than 25 at 0.5 GK

# Conclusions

Nuclear Astrophysics a fascinating and active field

Nuclear physics is a fundamental pillar of Nuclear Astrophysics

Direct measurements are becoming more and more accurate

Indirect methods are unique tools to investigate reactions on energy ranges difficult to study otherwise.

However, when possible, a joint work with direct and indirect measurements is the best choice to ensure accurate normalization and reaction rates for astrophysical applications

**Thank you for your attention!**