

# Light by light scattering in Ultra-peripheral collisions @ LHC

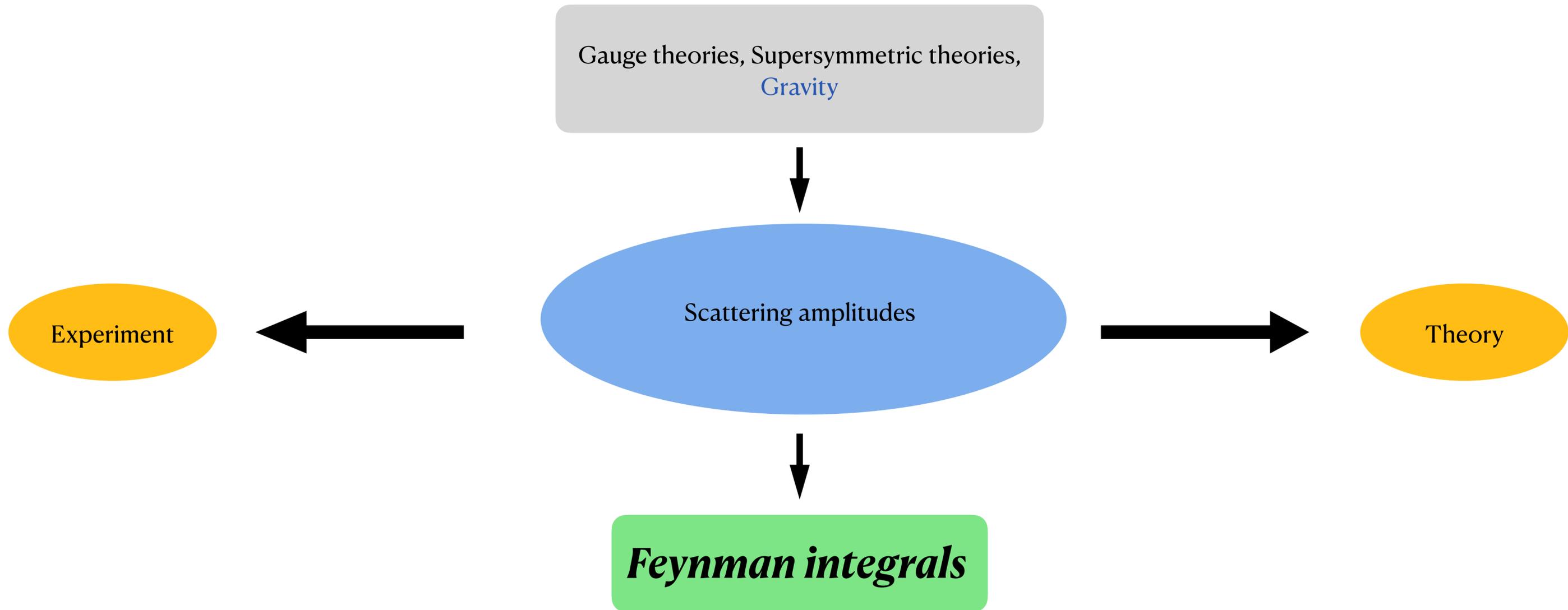
Ekta Chaubey

*Assemblée Générale 2023 du GDR LCD*  
*Strasbourg, France*  
29th September 2023



Bethe Center for  
Theoretical Physics





For precision physics, **precise** theoretical predictions are needed  
→ computation of **higher-order loop corrections**

# Analytic structure

Class of functions?

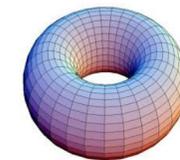
Rich Mathematical structure



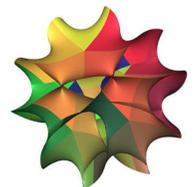
Multiple Polylogarithms



Elliptic Polylogarithms



K3 surfaces

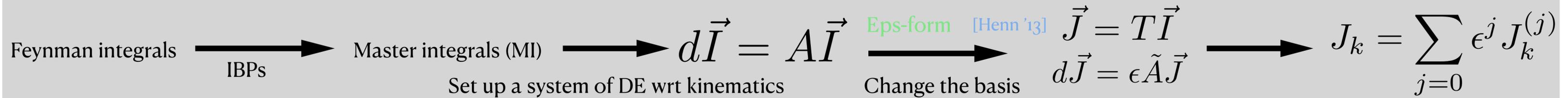


Higher-dimensional Calabi-Yau manifold

*Can we identify them?*

*Can we “use” them?*

# Modern techniques for solving Feynman integrals



The choice of basis not unique; no general methods fits all

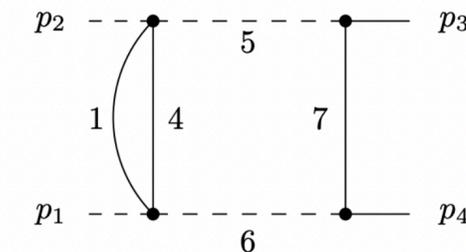
- [Lee '14] [Meyer '18] [Wasser '20]
- [Adams, E. Chaubey, Weinzierl, '17]
- [Dlapa, Henn, Wagner, '22]

## Maximal cuts

- [Baikov, '96, '97]
- [Frellesvig, Papadopoulos, '17]
- [Primo, Tancredi, '17]
- [Adams, E. Chaubey, Weinzierl, '17, '18]

## Square roots

Geometry of the Feynman integral often manifested by the square root  
[Festi, van Straten, '18] [Adams, E. Chaubey, Weinzierl '18]



4

$$\text{Maxcut}_C I_{2001111}(4-2\epsilon) \propto \int_C \frac{dP}{\sqrt{(P-t)(P-t+4m^2)(P^2+2m^2\frac{s+4t}{s-4m^2}P+m^2(m^2-4t)\frac{s}{s-4m^2}-\frac{4m^2t^2}{s-4m^2})}}$$

# Functional representation

Chen's definition of iterated integrals:  $\gamma : [0, 1] \rightarrow M$        $x_i = \gamma(0)$        $x_f = \gamma(1)$

[Chen '77]

$$\begin{aligned} I_\gamma(\omega_1, \dots, \omega_k; \lambda) &= \int_0^\lambda d\lambda_1 f_1(\lambda_1) \int_0^{\lambda_1} d\lambda_2 f_2(\lambda_2) \dots \int_0^{\lambda_{k-1}} d\lambda_k f_k(\lambda_k) \\ &= \int_0^\lambda d\lambda_1 f_1(\lambda_1) I_\gamma(\omega_2, \dots, \omega_k; \lambda_1), \end{aligned}$$

Goncharov's Polylogarithms (MPLs)

Elliptic Polylogarithms (eMPLs)

# Numerical evaluation

Numerical evaluation of the functions at various phase-space points

Multiple polylogarithms

Elliptic integrals

GINAC PolyLogTools

[Vollinga, Weinzierl, '05][Duhr, Dulat, '19]

Local series expansion methods

[Moriello, '19][Hidding, '21]

Iterated integrals  
Canonical DE

Ofcourse PySecDec, Fiesta, Feyntrop  
Auxiliary mass flow methods

[Liu, Ma, '22]

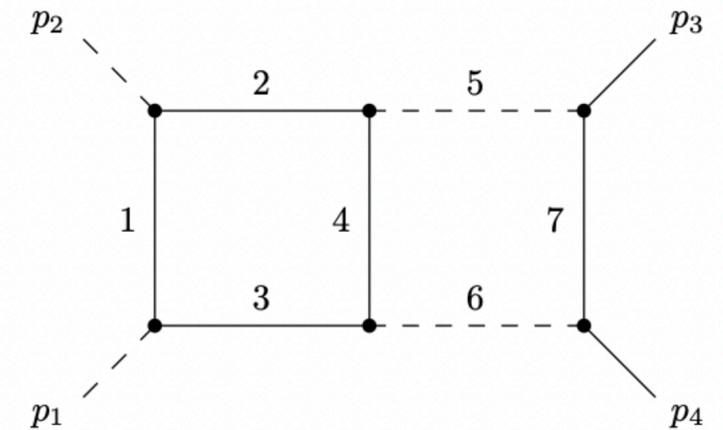
# Some selected examples

# Scattering amplitudes ( $t\bar{t}$ )

[Badger, E. Chaubey, Hartanto, Marzucca, '21]

Needed for precise understanding of top quarks

- Computation of helicity amplitudes for top-quark pair production involves elliptic integral contributions
- Topbox evaluates to 3 different elliptic curves
- Faster numerical evaluation of elliptic functions needs optimisation of the basis [E. Chaubey '21]

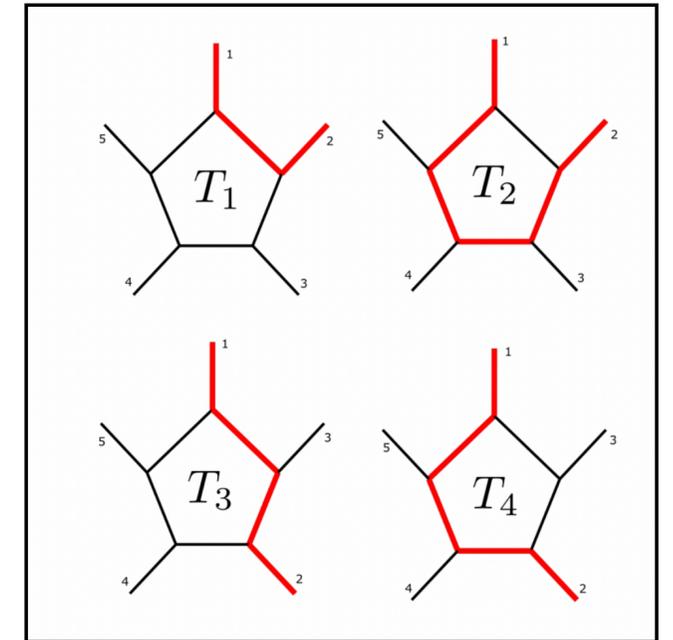


[Adams, E. Chaubey, Weinzierl, '18]

# 1-loop amplitude for tt-jet

[Badger, Becchetti, E. Chaubey, Marzucca, Sarandrea, '22]

- Analytic helicity amplitudes for 1-loop QCD corrections, Previously missing ingredient for NNLO, expansion of the 1-loop helicity amplitudes up to  $O(\epsilon^2)$



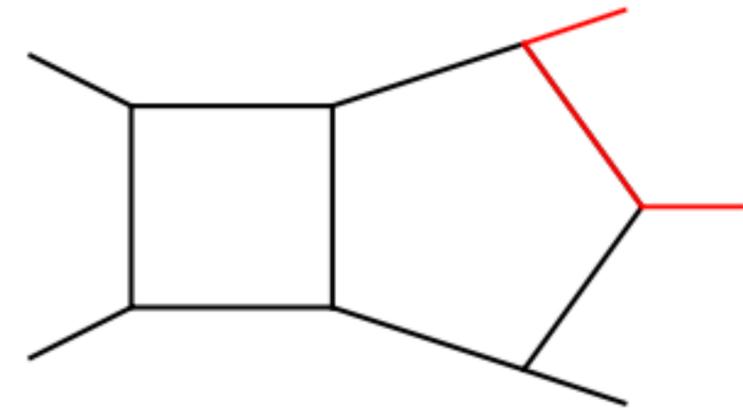
- Canonical form DE for all 130 MIs across 4 pentagon topologies
- Numerical solution using generalised power series expansion in DiffExp
- Analytic result of boundary constants

# 2-loop 5-point integrals

[Badger, Becchetti, E. Chaubey, Marzucca, '22]

- One of the topologies that appear in the NNLO corrections for tt-jet production
- 5 Mandelstam variables and mass dependence through top-quarks
- 88 master integrals
- Well-thought basis important to reconstruct analytically

$$\mu_{ij} = -k_i^{(-2\epsilon)} \cdot k_j^{(-2\epsilon)}$$



$$pp \rightarrow t\bar{t}j$$

- Evaluation at 1 boundary point with a precision of  $O(100)$  using AMFlow
- Integration of analytic differential equations using generalised power series expansion as implemented in DiffExp

# Two-loop QCD & QED corrections for light-by-light scattering

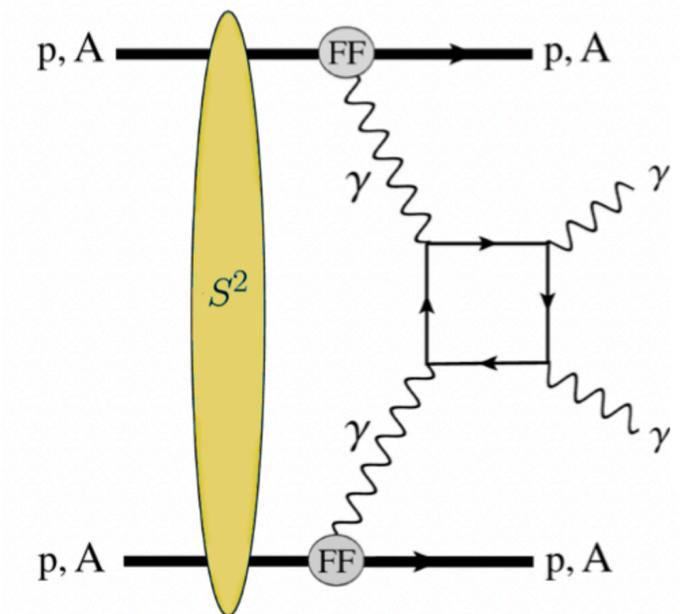
- The LHC can accelerate not just protons but heavy ions with charges up to  $Z=82$  for lead (Pb) ions. This enables many  $\gamma\gamma$  collision measurements in ultra-peripheral pp, pn and nn collisions (UPCs).
- Light-by-light scattering one of the few photon-fusion processes observed for the first time at the LHC.

[ATLAS collaboration '17, '19, '20]

[Klusek-Gawenda, Schaefer, Szczurek, '16]

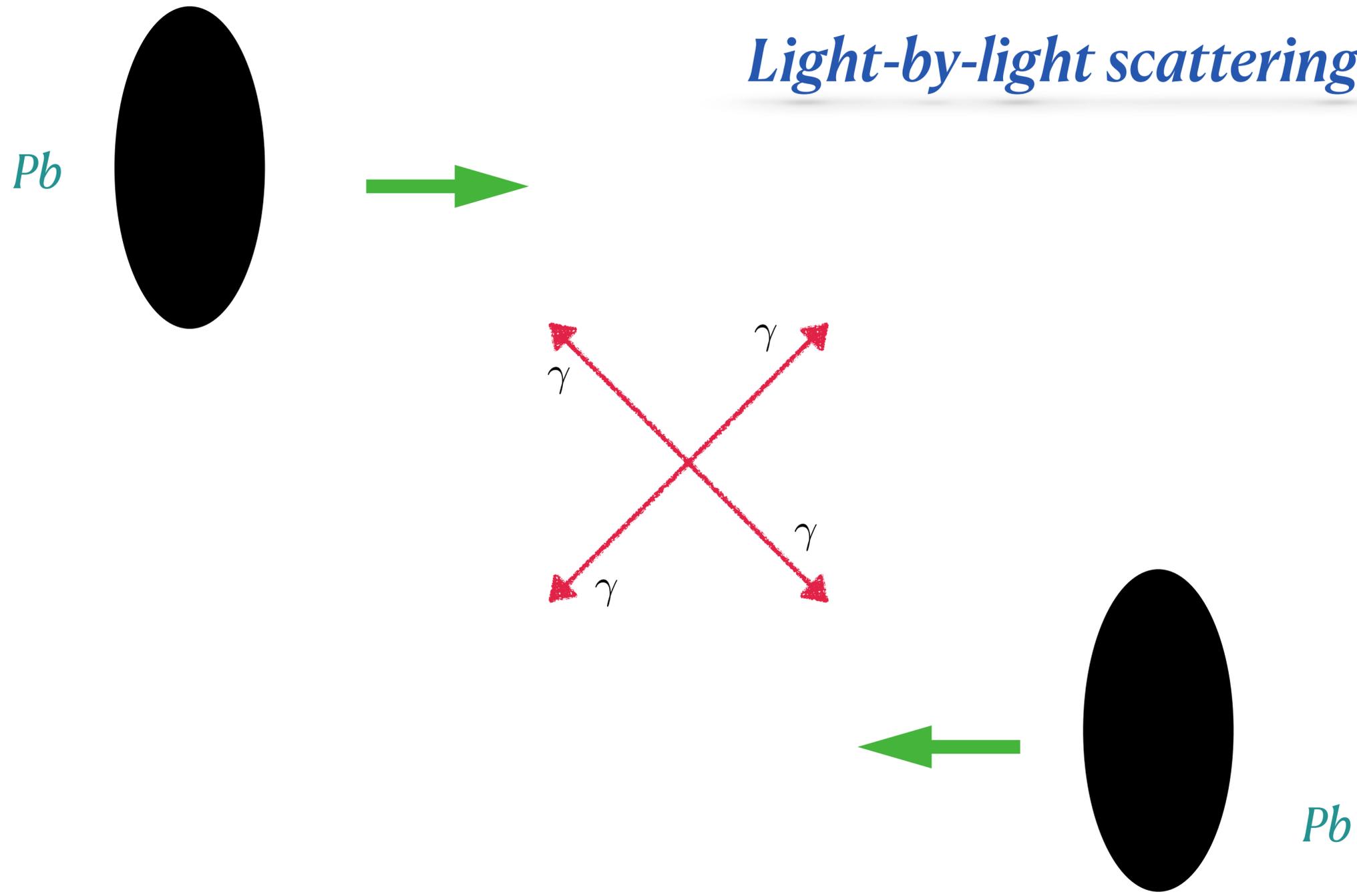
[Beloborodov, Kharlamova, Telnov, '23]

Important for studies of anomalous quartic gauge couplings, axion-like particles, Born-Infeld extensions of QED or anomalous tau electromagnetic moments as well as for important SM and BSM studies.



Light-by-light scattering at UPCs  
[Shao, d'Enterria, '22]

# *Light-by-light scattering*



Process	Physics motivation
$\gamma\gamma \rightarrow e^+e^-, \mu^+\mu^-$	“Standard candles” for proton/nucleus $\gamma$ fluxes, EPA calculations, and higher-order QED corrections
$\gamma\gamma \rightarrow \tau^+\tau^-$	Anomalous $\tau$ lepton e.m. moments [29–32]
$\gamma\gamma \rightarrow \gamma\gamma$	aQGC [25], ALPs [27], BI QED [28], noncommut. interactions [36], extra dims. [37],...
$\gamma\gamma \rightarrow \mathcal{T}_0$	Ditauonium properties (heaviest QED bound state) [38, 39]
$\gamma\gamma \rightarrow (c\bar{c})_{0,2}, (b\bar{b})_{0,2}$	Properties of scalar and tensor charmonia and bottomonia [40, 41]
$\gamma\gamma \rightarrow XYZ$	Properties of spin-even XYZ heavy-quark exotic states [42]
$\gamma\gamma \rightarrow VM VM$	(with VM = $\rho, \omega, \phi, J/\psi, \Upsilon$ ): BFKL-Pomeron dynamics [43–46]
$\gamma\gamma \rightarrow W^+W^-, ZZ, Z\gamma, \dots$	anomalous quartic gauge couplings [11, 26, 47, 48]
$\gamma\gamma \rightarrow H$	Higgs- $\gamma$ coupling, total H width [49, 50]
$\gamma\gamma \rightarrow HH$	Higgs potential [51], quartic $\gamma\gamma HH$ coupling
$\gamma\gamma \rightarrow t\bar{t}$	anomalous top-quark e.m. couplings [11, 49]
$\gamma\gamma \rightarrow \tilde{\ell}\tilde{\ell}, \tilde{\chi}^+\tilde{\chi}^-, H^{++}H^{--}$	SUSY pairs: slepton [11, 52, 53], chargino [11, 54], doubly-charged Higgs bosons [11, 55].
$\gamma\gamma \rightarrow a, \phi, MM, G$	ALPs [27, 56], radions [57], monopoles [58–61], gravitons [62–64],...

# Gamma UPC

## C. Light-by-light scattering in Pb-Pb UPCs at $\sqrt{s_{NN}} = 5.02$ TeV

The loop-induced LbL signal is generated with gamma-UPC plus MADGRAPH5\_AMC@NLO v2.6.6 [75, 138] with the virtual box contributions computed at leading order. Table XIV compares the integrated fiducial cross sections measured by ATLAS [15] with the gamma-UPC using EDFF and ChFF  $\gamma$  fluxes and the SUPERCHIC predictions. The measured cross section is about 2 standard deviations above the gamma-UPC and SUPERCHIC predictions.

TABLE XIV: Fiducial light-by-light cross sections measured in Pb-Pb UPCs at  $\sqrt{s_{NN}} = 5.02$  TeV (with  $E_T^\gamma > 2.5$  GeV,  $|\eta^\gamma| < 2.4$ ,  $m_{\gamma\gamma} > 5$  GeV,  $p_{T,\gamma\gamma} < 1$  GeV), compared to the theoretical gamma-UPC results obtained with EDFF and ChFF  $\gamma$  fluxes (and their average), as well as with the SUPERCHIC MC prediction.

Process, system	ATLAS data [15]	gamma-UPC $\sigma$			SUPERCHIC $\sigma$
		EDFF	ChFF	average	
$\gamma\gamma \rightarrow \gamma\gamma$ , Pb-Pb at 5.02 TeV	$120 \pm 22$ nb	63 nb	76 nb	$70 \pm 7$ nb	$78 \pm 8$ nb

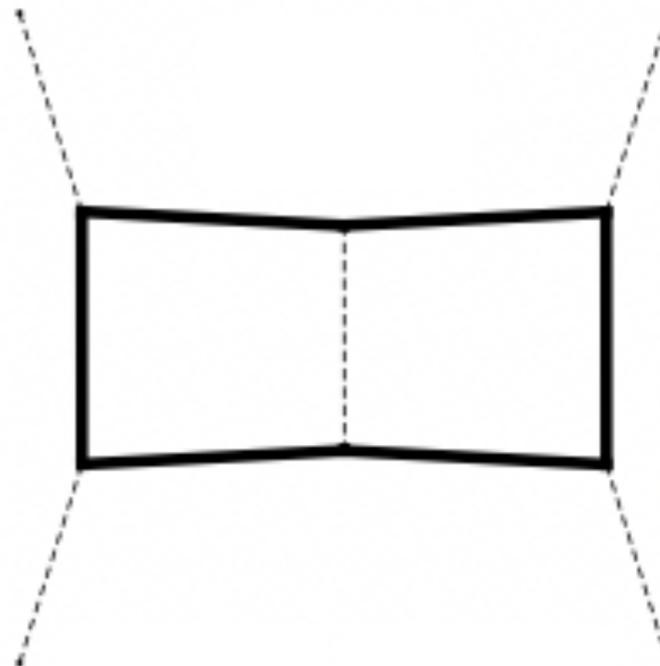
- Scale we are probing is 5 GeV (tau, lepton, bottom should be massive)
- Observed recently [ATLAS collaboration '17, '19, '20]
- Gamma\_UPC has been integrated into automated event generators Madgraph5\_aMC@NLO for NLO

# Light-by-light scattering

[AH, E. Chaubey, Shao, To appear]

$$\gamma(p_1, \lambda_1) + \gamma(p_2, \lambda_2) + \gamma(p_3, \lambda_3) + \gamma(p_4, \lambda_4) \rightarrow 0$$

- 2-loop QCD & QED corrections in the ultra-relativistic limit ( $s, t, u \gg m^2$ , massless internal lines)  
[Bern, De Freitas, Dixon, Ghinculov & Wong, 2001]
- Two-loop corrections to light-by-light scattering in supersymmetric QED  
[Binoth, Glover, Marquard, & van Der Bij, 2002]



# Amplitude computation

$$\mathcal{M} = \varepsilon_{1,\mu_1} \varepsilon_{2,\mu_2} \varepsilon_{3,\mu_3} \varepsilon_{4,\mu_4} \mathcal{M}^{\mu_1 \mu_2 \mu_3 \mu_4} (p_1, p_2, p_3, p_4)$$

$$\begin{aligned} \mathcal{M}^{\mu_1 \mu_2 \mu_3 \mu_4} = & A_1 g^{\mu_1 \mu_2} g^{\mu_3 \mu_4} + A_2 g^{\mu_1 \mu_3} g^{\mu_2 \mu_4} + A_3 g^{\mu_1 \mu_4} g^{\mu_2 \mu_3} + \sum_{j_1, j_2=1}^3 (B_{j_1 j_2}^1 g^{\mu_1 \mu_2} p_{j_1}^{\mu_3} p_{j_2}^{\mu_4} \\ & + B^2 j_1 j_2 g^{\mu_1 \mu_3} p_{j_1}^{\mu_2} p_{j_2}^{\mu_4} + B^3 j_1 j_2 g^{\mu_1 \mu_4} p_{j_1}^{\mu_2} p_{j_2}^{\mu_3} + B^4 j_1 j_2 g^{\mu_2 \mu_3} p_{j_1}^{\mu_2} p_{j_2}^{\mu_4} \\ & + B^5 j_1 j_2 g^{\mu_2 \mu_4} p_{j_1}^{\mu_1} p_{j_2}^{\mu_3} + B^6 j_1 j_2 g^{\mu_3 \mu_4} p_{j_1}^{\mu_1} p_{j_2}^{\mu_2}) \\ & + \sum_{j_1, j_2, j_3, j_4=1}^3 C_{j_1 j_2 j_3 j_4} p_{j_1}^{\mu_1} p_{j_2}^{\mu_2} p_{j_3}^{\mu_3} p_{j_4}^{\mu_4}. \end{aligned}$$

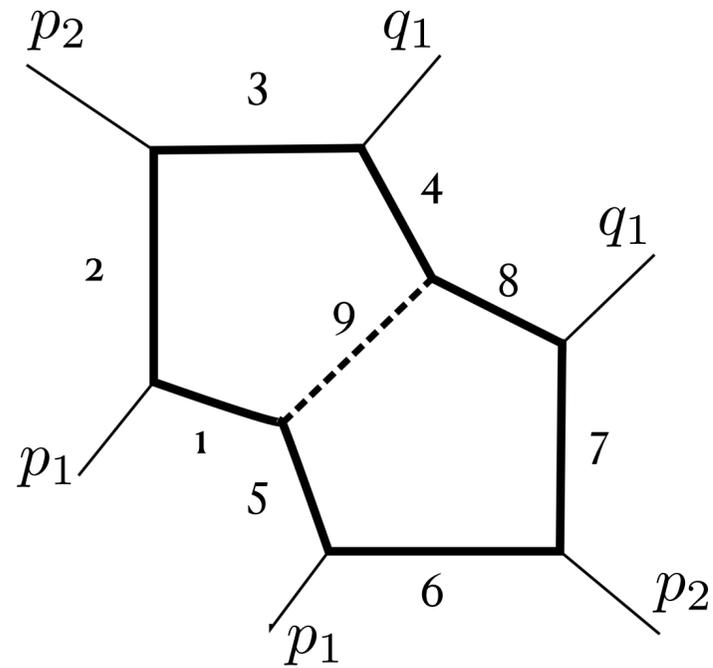
Number of independent functions

$$A_1(s, t, u) \quad B_{11}^1(s, t, u) \quad C_{2111}(s, t, u)$$

$\varepsilon_j \cdot p_j = 0$   
Bose symmetry  
Gauge symmetry

[Binoth, Glover, Marquard, & van Der Bij, 2002]

# Reduction to MIs & Simplification

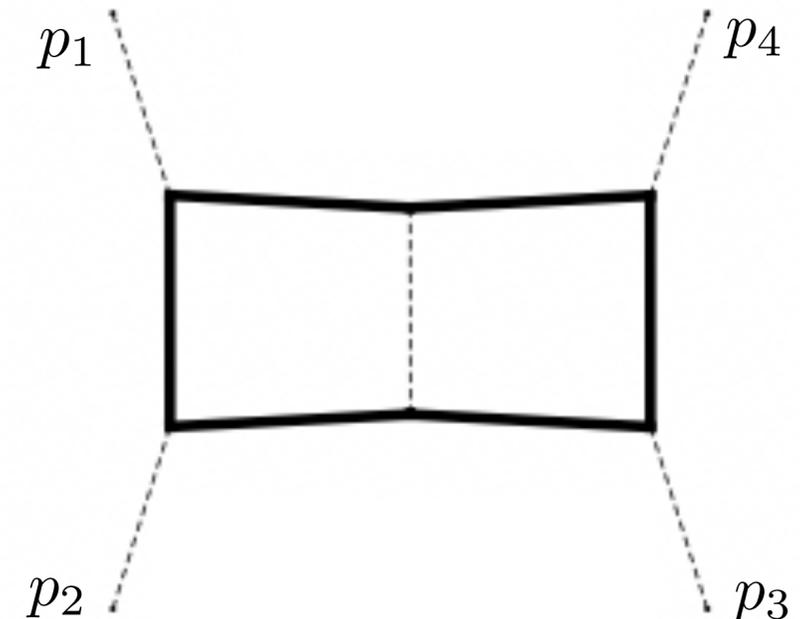


$$I_{a_1, \dots, a_9} = \left( \frac{e^{\epsilon \gamma_E}}{i\pi^{\frac{d}{2}}} \right)^2 \int \prod_{i=1}^2 d^d k_i \frac{D_4^{a_4} D_6^{a_6}}{D_1^{a_1} D_2^{a_2} D_3^{a_3} D_5^{a_5} D_7^{a_7} D_8^{a_8} D_9^{a_9}},$$

$$k_1^2 - m_t^2, (k_1 + p_1)^2 - m_t^2, (k_1 + p_1 + p_2)^2 - m_t^2, (k_1 + p_1 + p_2 + p_3)^2 - m_t^2, \\ (k_2)^2 - m_t^2, (k_2 + p_1)^2 - m_t^2, (k_2 + p_1 + p_2)^2 - m_t^2, (k_2 + p_1 + p_2 + p_3)^2 - m_t^2, \\ (k_2 - k_1)^2$$

LiteRed (FiniteFlow), KIRA  
 [Lee, '13] [Peraro, '19] [Klappert,  
 Lange, Maierhöfer, Usovitsch,

60 diagrams in total  
 7798 integrals before IBP  
 18 top-level sectors  
 Can be mapped into the 2-loop  
 diagram shown on the right



# Analytic computation of the MIs

- 29 MIs; use of differential equations;
- Choice of a canonical basis [\[Caron-huot, Henn, 14\]](#)

$$\begin{aligned}\partial_s \vec{f} &= \epsilon A_s \vec{f} \\ \partial_t \vec{f} &= \epsilon A_t \vec{f}\end{aligned}$$

- Square roots:

$$\sqrt{s(s - 4m^2)} \quad \sqrt{t(t - 4m^2)} \quad \sqrt{st(st - 4m^2(s + t))} \quad \sqrt{s(m^4s - 2m^2t(s + 2t) + st^2)}$$

- Choice of variables

$$s = -\frac{4(w - z)^2}{(1 - w^2)(1 - z^2)} \quad t = -\frac{(w - z)^2}{wz} \quad m^2 = 1$$

$$\sqrt{-2wz + z^2 + w^4z^2 - 2w^3z^3 + w^2(1 + z^2 + z^4)} \quad \text{Non-rationalizable??..}$$

# The alphabet

$$sq = \sqrt{-2wz + z^2 + w^4z^2 - 2w^3z^3 + w^2(1 + z^2 + z^4)} \quad [\text{Caron-huot, Henn, 14}]$$

$$1 - w, 1 + w, 1 - wz, w - z, w, w + z, 1 + w - z + wz, 1 - w + z + wz, 1 + wz, 1 - z, 1 + z,$$

$$z, \frac{-sq + w - z - wz + w^2z - wz^2}{sq + w - z - wz + w^2z - wz^2}, \frac{-sq + w^2 - 3wz + z^2}{sq + w^2 - 3wz + z^2}, \frac{-1 - sq + w^2z - wz^2}{-1 + sq + w^2z - wz^2},$$

$$\frac{1 - sq + w^2z - wz^2}{1 + sq + w^2z - wz^2}$$

*6 master integrals containing all 4 square roots in the integrand at weight-4.*

# Solving the canonical master integrals

We keep the analytic result in terms of iterated integrals with dog one-forms constructed using

[Heller, von Manteuffel, Schabinger, '20]

We convert all the integral in terms of 1-dimensional integrations

[Caron-huot, Henn, 14] [Chicherin, Sotnikov '21] [Chicherin, Sotnikov, Zoia '22]

We also express the first two orders in terms of logs and classical polylogs by matching symbols

[Duhr, Gangl, Rhodes, '11]

$$\int_{\gamma} I(\omega_1 \dots \omega_4; \lambda) = \int_0^{\lambda} d\lambda_1 f_1(\lambda_1) \int_0^{\lambda_1} d\lambda_2 f_2(\lambda_2) \underbrace{\int_0^{\lambda_2} d\lambda_3 f_3(\lambda_3) \int_0^{\lambda_3} d\lambda_4 f_4(\lambda_4)}_{\{\text{Log}^2(z), \text{Li}_2(z)\}}$$

# Numerical evaluation

*Physical phase-space regions of interest*

$$0 < w < 1 \ \& \ (0 < z < w \mid w < z < 1)$$

$$0 < w < 1 \ \& \ (1 < z < \frac{1}{w} \mid z > \frac{1}{w})$$

$$0 < w < 1 \ \& \ (-1 < z < -w \mid -w < z < 0)$$

*We obtain different analytic representations of the results valid in different regions*

**Outputs**



- *We obtained a completely analytic representation for the squared matrix element at 2-loop*
- *NLO cross section with massive contributions for light-by-light scattering at UPC within reach!*

# Take away

- Computation of higher order perturbative corrections important
- Analytic solution of multi-loop integrals requires understanding the **mathematical structure**
- With the inclusion of masses in the loop, the analytic structure starts becoming **more complicated**
- With the inclusion of more legs, often one needs go beyond current mathematical understand; often also desirable to **combine analytic as well as numerical** techniques