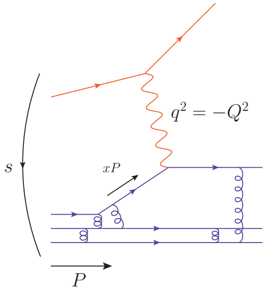
Probing the gluon orbital angular momentum with double spin asymmetries

Renaud Boussarie

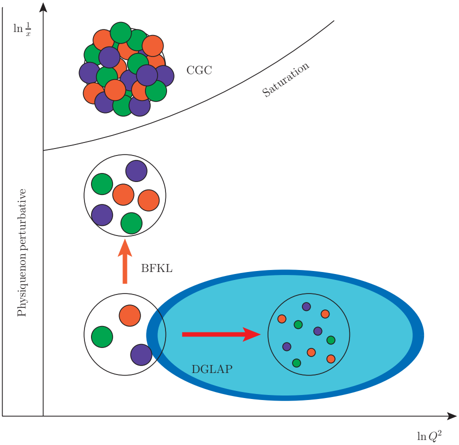
AG du GDR QCD 2023

[Bhattacharya, RB, Hatta, PRL 128 (2022) 18, 182002]

Accessing the partonic content of hadrons with an electromagnetic probe

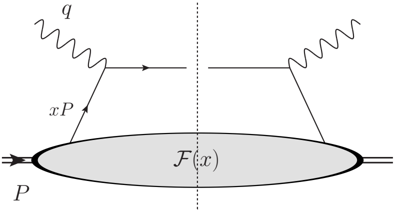


Electron-proton collision (parton model)



QCD factorization

processes with a hard scale $Q \gg \Lambda_{QCD}$



$$\sigma = \mathcal{F}(x, \mu) \otimes \mathcal{H}(x, \mu)$$

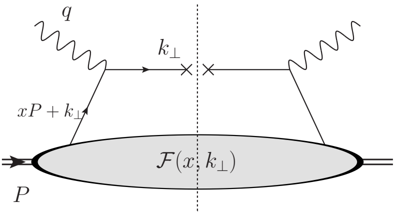
At a scale μ, the process is factorized into:

- A hard scattering subamplitude $\mathcal{H}(x, \mu)$
- A Parton Distribution Function (PDF) $\mathcal{F}(x, \mu)$

μ independence: DGLAP renormalization equation for $\mathcal{F}$

Transverse Momentum Dependent (TMD) factorization: semi-inclusive processes with one hard and one semihard scale

$$Q \sim \sqrt{s} \gg k_\perp$$



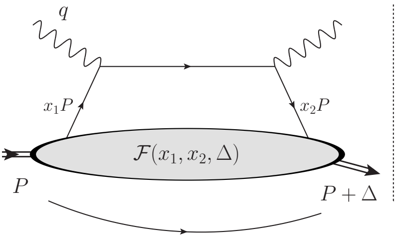
$$\sigma = \mathcal{F}(x, k_\perp, \zeta, \mu) \otimes \mathcal{H}(\mu) \otimes \hat{\mathcal{F}}(\hat{x}, \hat{k}_\perp, \hat{\zeta}, \mu)$$

At a scale μ, the process is factorized into:

- A hard scattering subamplitude $\mathcal{H}(\mu)$
- A TMD PDF $\mathcal{F}(x, k_\perp, \zeta, \mu)$
- A TMD FF $\hat{\mathcal{F}}(\hat{x}, \hat{k}_\perp, \hat{\zeta}, \mu)$

$\mu, \zeta, \hat{\zeta}$ independence: TMD evolution for $\mathcal{F}, \hat{\mathcal{F}}$

Factorization with Generalized Parton Distributions (GPD): exclusive processes with one hard scale $Q \sim \sqrt{s}$



$$\sigma = \mathcal{F}(x_1, x_2, |\Delta_\perp|, \mu) \otimes \mathcal{H}\,(x_1, x_2, \mu)$$

At a scale μ, the process is factorized into:

- A hard scattering subamplitude $\mathcal{H}(x_1, x_2, \mu)$
- A Generalized Parton Distribution (GPD) $\mathcal{F}(x_1, x_2, |\Delta_\perp|, \mu)$

μ independence: DGLAP/ERBL renormalization equation for $\mathcal{F}$

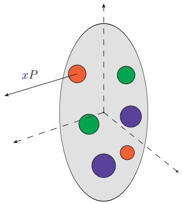
Operator definition for parton distributions

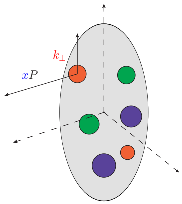
Parton distribution function

$$\mathcal{F}(x) \propto \int \mathrm{d}z^- \mathrm{e}^{ixP^+z^-} \left\langle P \left| F^{+i}(z^-) \left[z^-, 0^-\right] F^{+i}(0) \left[0^-, z^-\right] \right| P \right\rangle$$

Transverse Momentum Dependent distribution

$$\mathcal{F}(x, k_\perp) \propto \int \mathrm{d}^4 z \delta(z^+) \mathrm{e}^{ixP^+z^- + i(k_\perp \cdot z_\perp)} \left\langle P \left| F^{+i}(z) \mathcal{U}_{z,0} F^{+i}(0) \mathcal{U}_{0,z} \right| P \right\rangle$$





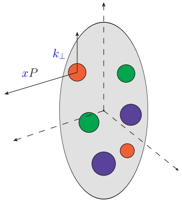
Operator definition for parton distributions

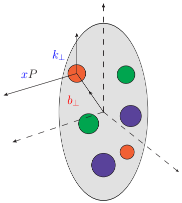
TMD distribution

$$\mathcal{F}(x, k_\perp) \propto \int \mathrm{d}^4 z \delta(z^+) \mathrm{e}^{ixP^+z^- + i(k_\perp \cdot z_\perp)} \left\langle P \left| F^{+i}(z) \mathcal{U}_{z,0} F^{+i}(0) \mathcal{U}_{0,z} \right| P \right\rangle$$

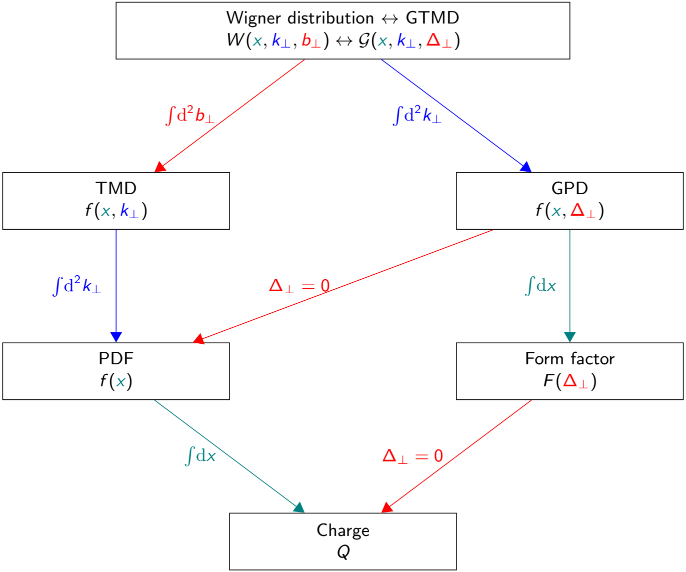
Generalized TMD distribution

$$\mathcal{F}(x, k_\perp, \Delta) \propto \int \mathrm{d}^4 z \delta(z^+) \mathrm{e}^{ixP^+z^- + i(k_\perp \cdot z_\perp)} \left\langle P+\Delta \left| F^{+i}(z) \mathcal{U}_{z,0} F^{+i}(0) \mathcal{U}_{0,z} \right| P \right\rangle$$





The family tree of parton distributions



Wigner distributions in NRQM

Wigner distributions in Quantum Mechanics [Wigner, 1932]

Defined via wavefunctions

$$W(x,k) = \int \frac{\mathrm{d}x'}{2\pi} \mathrm{e}^{-i(k\cdot x')} \psi\left(x+\frac{x'}{2}\right)\psi^*\left(x-\frac{x'}{2}\right)$$

Connection with probability densities:

$$|\psi(x)|^2 = \int \mathrm{d}k \, W(x,k)$$

and

$$|\psi(k)|^2 = \int \mathrm{d}x \, W(x,k)$$

Connection with observables:

$$\langle O \rangle = \int \mathrm{d}x \int \mathrm{d}k \, O(x,k) \, W(x,k)$$

Wigner distributions in QCD

QCD Wigner distributions

Defined as the Fourier transform of a GTMD

$$\mathcal{W}^g(x, k_\perp, b_\perp) = \int \mathrm{d}^2\Delta_\perp \mathrm{e}^{i(b_\perp \cdot \Delta_\perp)} \times \int \mathrm{d}^4 z \delta(z^+) \mathrm{e}^{ixP^+ z^- + i(k_\perp \cdot z_\perp)} \left\langle P+\Delta_\perp \left| F^{+i}(z)\mathcal{U}_{z,0} F^{+i}(0)\mathcal{U}_{0,z} \right| P \right\rangle$$

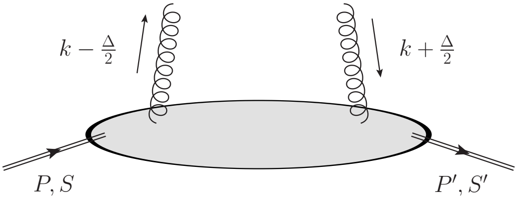
Connection with observables: e.g. orbital angular momentum of gluons inside a proton

$$\langle L_z^g \rangle = \int \mathrm{d}x \int \mathrm{d}^2 k_\perp \int \mathrm{d}^2 b_\perp (b_\perp \times k_\perp)_z W^g(x, k_\perp, b_\perp)$$

GTMD

Parametrization and coupling to the target hadron

[Meissner, Metz, Schlegel, 2009], [Lorcé, Pasquini, 2013]



$$\int \mathrm{d}^4 v \delta(v^+) e^{ix\bar{P}^+ v^- - i(\boldsymbol{k}\cdot\boldsymbol{v})} \langle P' S' | F^{i+}(-\frac{v}{2}) \mathcal{U}^{[+]}_{\frac{v}{2},-\frac{v}{2}} F^{i+}(\frac{v}{2}) | PS \rangle$$

$$= (2\pi)^3 \frac{\bar{P}^+}{2M} \bar{u}_{P'S'} \left[F^g_{1,1} + i \frac{\sigma^{i+}}{\bar{P}^+} (\boldsymbol{k}^i F^g_{1,2} + \boldsymbol{\Delta}^i F^g_{1,3}) + i \frac{\sigma^{ij} \boldsymbol{k}^i \boldsymbol{\Delta}^j}{M^2} F^g_{1,4} \right] u_{PS}$$

Leading twist GTMDs

Leading twist GTMDs in a spin 1/2 hadron

[Meissner, Metz, Schlegel, 2009], [Lorcé, Pasquini, 2013]

- Unpolarized parton pairs: $F_{1,1}$, $F_{1,2}$, $F_{1,3}$ and $F_{1,4}$
- Polarized parton pairs: $G_{1,1}$, $G_{1,2}$, $G_{1,3}$ and $G_{1,4}$
- Transversity distributions: $H_{1,1}$, $H_{1,2}$, $H_{1,3}$, $H_{1,4}$, $H_{1,5}$, $H_{1,6}$, $H_{1,7}$ and $H_{1,8}$

16 distributions with real and imaginary values at leading twist

Leading twist GTMDs vs GPDs

GTMDs span GPDs

- Unpolarized parton pairs:
  $\mathrm{Re}(F_{1,1})$, $\mathrm{Re}(F_{1,2})$, $\mathrm{Re}(F_{1,3}) \to H, E$
- Polarized parton pairs:
  $\mathrm{Re}(G_{1,2})$, $\mathrm{Re}(G_{1,3})$, $\mathrm{Re}(G_{1,4}) \to \tilde{H}, \tilde{E}$
- Transversity distributions:
  $\mathrm{Re}(H_{1,3})$, $\mathrm{Re}(H_{1,4})$, $\mathrm{Re}(H_{1,5})$, $\mathrm{Re}(H_{1,6})$, $\mathrm{Re}(H_{1,7})$, $\mathrm{Re}(H_{1,8}) \to H_T, E_T$

Most distributions have a null GPD limit

GTMDs and Orbital Angular Momentum

Spin sum rule

Jaffe Manohar spin sum rule

$$\frac{1}{2} = \frac{1}{2}\Delta\Sigma + \Delta G + L^q + L^g$$

- Quark helicity $\Delta\Sigma$
- Gluon helicity ΔG
- Quark OAM L^q
- Gluon OAM L^g

Gluon OAM

Gluon orbital angular momentum

'Intuitive' (canonical) definition for the OAM

$$L_z^g = \int \mathrm{d}^2 k_\perp (\Delta_\perp \times k_\perp)_z \, x f^g(x, \xi, k_\perp, \Delta_\perp)$$

with $xf^g(x, \xi, k_\perp, \Delta_\perp)$ the gluon GTMD correlator

Gluon OAM

Gluon orbital angular momentum

$$xf_g(x,\xi,\boldsymbol{k},\boldsymbol{\Delta}) = \frac{1}{2M}\bar{u}(p',h_{p'})\left[F_{1,1} + i\frac{\sigma^{j+}}{P^+}(\boldsymbol{k}^j F_{1,2} + \boldsymbol{\Delta}^j F_{1,3}) + i\frac{\sigma^{ij}\boldsymbol{k}^i\boldsymbol{\Delta}^j}{M^2}F_{1,4}\right]u(p,h_p)$$

Symmetry properties:

$$F^*_{1,(1,3,4)}(\xi,\boldsymbol{k}\cdot\boldsymbol{\Delta},\boldsymbol{k}^2,\boldsymbol{\Delta}^2) = F_{1,(1,3,4)}(-\xi,-\boldsymbol{k}\cdot\boldsymbol{\Delta},\boldsymbol{k}^2,\boldsymbol{\Delta}^2)$$

$$F^*_{1,2}(\xi,\boldsymbol{k}\cdot\boldsymbol{\Delta},\boldsymbol{k}^2,\boldsymbol{\Delta}^2) = -F_{1,2}(-\xi,-\boldsymbol{k}\cdot\boldsymbol{\Delta},\boldsymbol{k}^2,\boldsymbol{\Delta}^2)$$

Gluon OAM

Gluon orbital angular momentum

Symmetry properties:

$$F^*_{1,(1,3,4)}(\xi, \boldsymbol{k}\cdot\boldsymbol{\Delta}, \boldsymbol{k}^2, \boldsymbol{\Delta}^2) = F_{1,(1,3,4)}(-\xi, -\boldsymbol{k}\cdot\boldsymbol{\Delta}, \boldsymbol{k}^2, \boldsymbol{\Delta}^2)$$

$$F^*_{1,2}(\xi, \boldsymbol{k}\cdot\boldsymbol{\Delta}, \boldsymbol{k}^2, \boldsymbol{\Delta}^2) = -F_{1,2}(-\xi, -\boldsymbol{k}\cdot\boldsymbol{\Delta}, \boldsymbol{k}^2, \boldsymbol{\Delta}^2)$$

so we can write

$$F_{1,(1,3,4)} = \left[R^{(1)}_{1,(1,3,4)} + \xi(\boldsymbol{k}\cdot\boldsymbol{\Delta})R^{(2)}_{1,(1,3,4)}\right] + i\left[\xi I^{(1)}_{1,(1,3,4)} + (\boldsymbol{k}\cdot\boldsymbol{\Delta})I^{(2)}_{1,(1,3,4)}\right]$$

$$F_{1,2} = \left[\xi R^{(1)}_{1,2} + (\boldsymbol{k}\cdot\boldsymbol{\Delta})R^{(2)}_{1,2}\right] + i\left[I^{(1)}_{1,2} + \xi(\boldsymbol{k}\cdot\boldsymbol{\Delta})I^{(2)}_{1,2}\right]$$

Gluon OAM

Gluon orbital angular momentum

$$xf_g(x,\xi,\boldsymbol{k},\boldsymbol{\Delta})$$
$$= \frac{1}{2M}\bar{u}(p',h_{p'})\left[F_{1,1} + i\frac{\sigma^{j+}}{P^+}(\boldsymbol{k}^j F_{1,2} + \boldsymbol{\Delta}^j F_{1,3}) + i\frac{\sigma^{ij}\boldsymbol{k}^i\boldsymbol{\Delta}^j}{M^2}F_{1,4}\right]u(p,h_p)$$

with

$$F_{1,(1,3,4)} = \left[R^{(1)}_{1,(1,3,4)} + \xi(\boldsymbol{k}\cdot\boldsymbol{\Delta})R^{(2)}_{1,(1,3,4)}\right] + i\left[\xi I^{(1)}_{1,(1,3,4)} + (\boldsymbol{k}\cdot\boldsymbol{\Delta})I^{(2)}_{1,(1,3,4)}\right]$$

$$F_{1,2} = \left[\xi R^{(1)}_{1,2} + (\boldsymbol{k}\cdot\boldsymbol{\Delta})R^{(2)}_{1,2}\right] + i\left[I^{(1)}_{1,2} + \xi(\boldsymbol{k}\cdot\boldsymbol{\Delta})I^{(2)}_{1,2}\right]$$

Identify the terms which survive the convolution with $(\boldsymbol{k}\times\boldsymbol{\Delta})$

Gluon OAM

Gluon orbital angular momentum

$$\int \mathrm{d}^2\boldsymbol{k}(\boldsymbol{k}\times\boldsymbol{\Delta})xf_g(x,\xi,\boldsymbol{k},\boldsymbol{\Delta})$$

$$= \int \mathrm{d}^2\boldsymbol{k}\,\frac{i(\boldsymbol{k}\times\boldsymbol{\Delta})}{2M}\times \bar{u}(p',h_{p'})\left[\xi\frac{\boldsymbol{k}^i\sigma^{i+}}{P^+}R^{(1)}_{1,2}+\frac{\sigma^{ij}\boldsymbol{k}^i\boldsymbol{\Delta}^j}{M^2}F_{1,4}\right]u(p,h_p)$$

At small ξ, only $F_{1,4}$ contributes

Gluon OAM

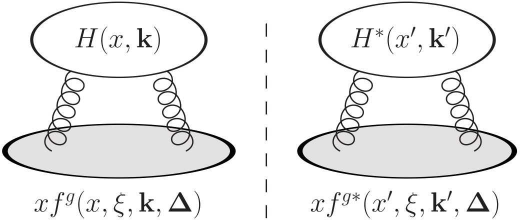
Gluon orbital angular momentum as a GTMD

$$\int \mathrm{d}^2\boldsymbol{k}\,\frac{\boldsymbol{k}^2}{M^2}F^g_{1,4}(x,0,\boldsymbol{k},\mathbf{0}) = -xL_g(x,0)$$

Where to constrain the OAM GTMD?

How to constrain the OAM GTMD

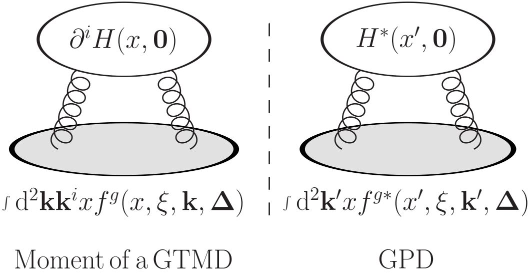
Exclusive process with general(-ish) gluon kinematics



Squared amplitude of a process involving a gluon GTMD

How to constrain the OAM GTMD

Power expand the cross section



We need to study the GTMD*GPD proton structure

How to constrain the OAM GTMD

GTMD*GPD correlators, summed over outgoing helicities

$$\sum_{h'_p} \frac{1}{2P^+} \left[x f_g(x,\xi,\boldsymbol{k},\boldsymbol{\Delta})\right] \left[x' F_g(x',\xi,\boldsymbol{\Delta})\right]^*$$

$$= H_g F_{1,1} + \frac{\xi^2}{1-\xi^2} E_g F_{1,1}$$

$$+ \frac{1}{2}\frac{\boldsymbol{k}\cdot\boldsymbol{\Delta}}{M^2} E_g F_{1,2} + \frac{1}{2}\frac{\boldsymbol{\Delta}^2}{M^2} E_g F_{1,3}$$

$$+ \frac{i S^+ \epsilon^{ij} \boldsymbol{k}^i \boldsymbol{\Delta}^j}{(1+\xi)P^+ M^2} \left(H_g F^g_{1,4} - \frac{\xi^2}{1-\xi^2} E_g F_{1,4} - \frac{1}{2} E_g F_{1,2} \right)$$

Longitudinal spin asymmetries select $F^g_{1,4}$

How to constrain the OAM GTMD

Target correlator for longitudinal target spin asymmetry

$$\sum_{h'_p} \frac{1}{2P^+} \left[x f_g(x,\xi,\boldsymbol{k},\boldsymbol{\Delta})\right] \left[x' F_g(x',\xi,\boldsymbol{\Delta})\right]^* \Bigg|_{S^+}$$

$$= \frac{iS^+ \epsilon^{ij} \boldsymbol{k}^i \boldsymbol{\Delta}^j}{(1+\xi)P^+ M^2} \left(H_g F^g_{1,4} - \frac{1}{2} E_g F_{1,2} - \frac{\xi^2}{1-\xi^2} E_g F_{1,4} \right)$$

- Phenomenologically, $E_g \ll H_g$
- The moment of $F^g_{1,2}$ is a genuine higher twist while that of $F^g_{1,4}$ has a kinematic part

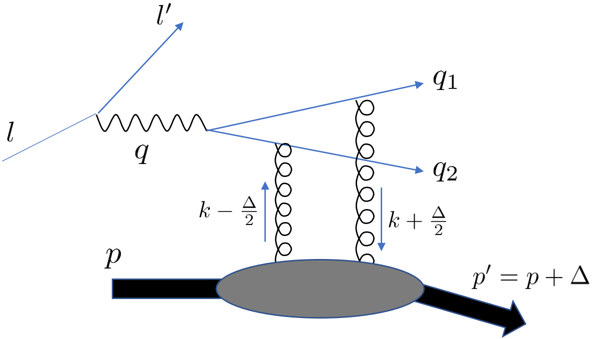
Longitudinal spin asymmetry selects the OAM GTMD $F^g_{1,4}$

How to constrain the gluon OAM

A process will contrain the gluon OAM if:

- It involves gluons
- It is exclusive enough
- It involves a longitudinal target spin asymmetry
- Its hard part involves $\boldsymbol{k} \times \boldsymbol{\Delta}$

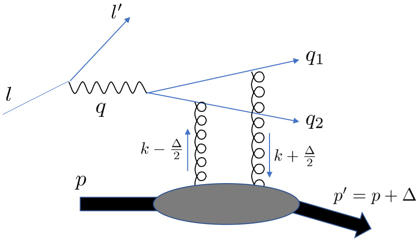
Exclusive dijet electroproduction and gluon OAM



[Ji, Yuan, Zhao], [Bhattacharya, RB, Hatta]

Kinematics

Dijet kinematics



Longitudinal fractions

$$z = p \cdot q_1 / p \cdot q$$
$$\bar{z} = p \cdot q_2 / p \cdot q$$

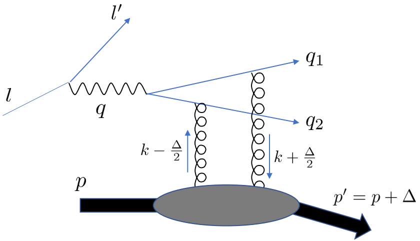
Transverse momenta

$$q_{1\perp} = q_\perp - z\Delta_\perp$$
$$q_{2\perp} = -q_\perp - \bar{z}\Delta_\perp$$

Consider longitudinal target spin asymmetry, try to select $(\mathbf{k} \times \mathbf{\Delta})$

SSA kinematics

Dijet kinematics



Longitudinal fractions

$$z = p \cdot q_1 / p \cdot q$$
$$\bar{z} = p \cdot q_2 / p \cdot q$$

Transverse momenta

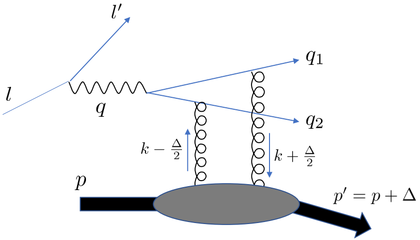
$$q_{1\perp} = q_\perp - z\Delta_\perp$$
$$q_{2\perp} = -q_\perp - \bar{z}\Delta_\perp$$

Single (target) spin asymmetry [Ji, Yuan, Zhao]

$$\mathrm{d}\sigma^{h_p} \sim h_p(z - \bar{z}) \sin(\phi_{q_\perp} - \phi_{\Delta_\perp}) \mathrm{Re}(A_2 A_3^*)$$

SSA kinematics

Dijet kinematics



Longitudinal fractions

$$z = p \cdot q_1 / p \cdot q$$
$$\bar{z} = p \cdot q_2 / p \cdot q$$

Transverse momenta

$$q_{1\perp} = q_\perp - z\Delta_\perp$$
$$q_{2\perp} = -q_\perp - \bar{z}\Delta_\perp$$

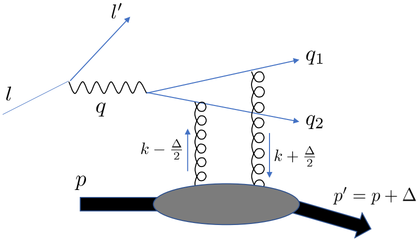
Single (target) spin asymmetry [Ji, Yuan, Zhao]

$$\mathrm{d}\sigma^{h_p} \sim h_p(z - \bar{z}) \sin(\phi_{q_\perp} - \phi_{\Delta_\perp}) \mathrm{Re}(A_2 A_3^*)$$

BUT A_3 has end point singularities (higher poles in $x \pm \xi$)

DSA kinematics

Dijet kinematics



Longitudinal fractions

$$z = p \cdot q_1 / p \cdot q$$
$$\bar{z} = p \cdot q_2 / p \cdot q$$

Transverse momenta

$$q_{1\perp} = q_\perp - z\Delta_\perp$$
$$q_{2\perp} = -q_\perp - \bar{z}\Delta_\perp$$

Double spin asymmetry [Bhattachary, RB, Hatta]

$$\mathrm{d}\sigma^{h_p h_l} \sim h_p h_l \cos(\phi_{l_\perp} - \phi_{\Delta_\perp})\mathrm{Re}(A_2' A_3'^*)$$

A_3' has end point singularities BUT they cancel at $z = 1/2$

DSA

Best gluon OAM observable so far: DSA

- Consider exclusive DIS dijet production
- Make sure gluons are involved (quark-induced jets, $c\bar{c}$ pair...)
- Measure double spin asymmetries
- Focus on $\cos(\phi_{l_\perp} - \phi_{\Delta_\perp})$ at $z = 1/2$

DSA

Best gluon OAM observable so far: DSA

- Consider exclusive DIS dijet production
- Make sure gluons are involved (quark-induced jets, $c\bar{c}$ pair...)
- Measure double spin asymmetries
- Focus on $\cos(\phi_{l_\perp} - \phi_{\Delta_\perp})$ at $z = 1/2$

HOWEVER

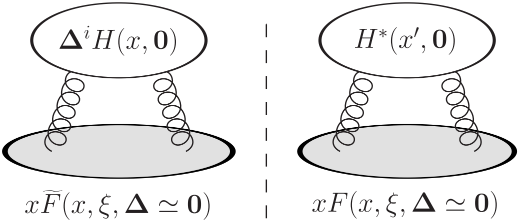
Helicity contribution

Another contribution to that observable

Polarized gluon GPDs contribute to

$$\mathrm{d}\sigma^{h_p h_l} \sim h_p h_l \cos(\phi_{l_\perp} - \phi_{\Delta_\perp}) \mathrm{Re}(A_2' A_3'^*)$$

at the same order as the OAM term



Cross section

Cross section for DSA in exclusive DIS dijet production

Neglecting GPD E wrt GPD H

$$\frac{d\sigma}{\mathrm{d}y\mathrm{d}Q^2\mathrm{d}\phi_{l_\perp}\mathrm{d}z\mathrm{d}\boldsymbol{q}^2\mathrm{d}^2\boldsymbol{\Delta}} = \frac{\alpha_s^2\alpha_{\mathrm{em}}^2 e_q^2 y}{2\pi^3 N_c Q^2}\xi\frac{h_l h_p|\boldsymbol{l}||\boldsymbol{\Delta}|\cos(\phi_{l_\perp}-\phi_{\Delta_\perp})}{(W^2+Q^2)(-\boldsymbol{q}^2+W^2/4)(\boldsymbol{q}^2+Q^2/4)^2}$$
$$\times\,\mathrm{Re}\left(\mathcal{H}_g^{(1)*}\widetilde{\mathcal{H}}_g^{(2)}-\mathcal{H}_g^{(1)*}\mathcal{L}_g-\frac{4\boldsymbol{q}^2}{\boldsymbol{q}^2+z\bar{z}Q^2}\mathcal{H}_g^{(2)*}\mathcal{L}_g\right)$$

Compton form factors for helicity ($\widetilde{\mathcal{H}}$) and for OAM ($\mathcal{L}$).

Positive or negative interference depending on $\boldsymbol{q}^2$ vs Q^2
$\Rightarrow$ insight on their interplay

Cross section estimates

Numerical estimation

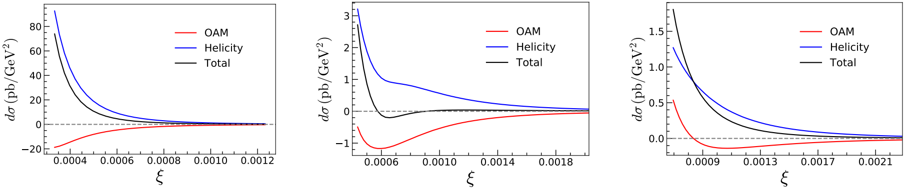
- $H_g(x,\xi)$, $\tilde{H}_g(x,\xi)$: double distribution ansatz $\Leftrightarrow xg(x), x\Delta g(x)$
- $xg(x), x\Delta g(x)$: use known PDF sets, such as JAM
- $L_g(x,\xi)$: double distribution ansatz $L_g(x)$
- $L_g(x)$: known in terms of $xg(x)$ and $x\Delta g(x)$ in the WW approximation

$$L_g(x) \simeq x \int_x^1 \frac{\mathrm{d}x'}{x'^2} \left[x'g(x') - 2\Delta g(x')\right]$$

[Hatta, Yoshida]

Cross section estimates

Numerical estimation



$Q^2 = 2.7\mathrm{GeV}^2$ $\quad Q^2 = 4.8\mathrm{GeV}^2$ $\quad Q^2 = 10\mathrm{GeV}^2$

$$\left.\frac{d\sigma}{\mathrm{d}y\mathrm{d}Q^2\mathrm{d}z\mathrm{d}\xi\mathrm{d}\delta\phi}\right|_{\delta\phi\simeq 0, y=0.7}$$

Summary

Conclusions

- GTMDs provide an intuitive definition of the OAM
- One GTMD in particular defines it
- Longitudinal SSA and DSA can constrain this GTMD
- DSA in exclusive DIS dijet probes both OAM and helicity