

# Extraction of TMD distributions

Valerio Bertone

IRFU, CEA, Université Paris-Saclay



September 27, 2023, Assemblée Générale 2023 du GDR QCD, Strasbourg

# Introduction

- 🍎 The **transverse momentum** ( $q_T$ ) distribution of a **high-mass** ( $Q$ ) system has two main regimes:

- 🍎 for  $q_T \gtrsim Q$  **collinear factorisation** at *fixed perturbative order* is appropriate:

$$\left( \frac{d\sigma}{dq_T} \right)_{\text{f.o.}} = \int_0^1 dx_1 \int_0^1 dx_2 f_1(x_1, Q) f_2(x_2, Q) \frac{d\hat{\sigma}}{dq_T} + \mathcal{O}\left[\left(\frac{\Lambda_{\text{QCD}}}{Q}\right)^n\right]$$

- 🍎 for  $q_T \ll Q$  **transverse-momentum-dependent (TMD) factorisation** at *fixed logarithmic order* is appropriate:

$$\left( \frac{d\sigma}{dq_T} \right)_{\text{res.}} \stackrel{\text{TMD}}{=} \sigma_0 H(Q) \int d^2 \mathbf{b}_T e^{i \mathbf{b}_T \cdot \mathbf{q}_T} F_1(x_1, \mathbf{b}_T, Q, Q^2) F_2(x_2, \mathbf{b}_T, Q, Q^2) + \mathcal{O}\left[\left(\frac{q_T}{Q}\right)^m\right]$$

- 🍎 Collinear and TMD factorisations may eventually be **matched** to produce accurate results over the full  $q_T$  spectrum.

# Introduction

- 🍎 The **transverse momentum** ( $q_T$ ) distribution of a **high-mass** ( $Q$ ) system has two main regimes:

- 🍎 for  $q_T \gtrsim Q$  **collinear factorisation** at *fixed perturbative order* is appropriate:

$$\left( \frac{d\sigma}{dq_T} \right)_{\text{f.o.}} = \int_0^1 dx_1 \int_0^1 dx_2 f_1(x_1, Q) f_2(x_2, Q) \frac{d\hat{\sigma}}{dq_T} + \mathcal{O}\left[\left(\frac{\Lambda_{\text{QCD}}}{Q}\right)^n\right]$$

- 🍎 for  $q_T \ll Q$  **transverse-momentum-dependent (TMD) factorisation** at *fixed logarithmic order* is appropriate:

Main subject of this talk

$$\left( \frac{d\sigma}{dq_T} \right)_{\text{res.}} \stackrel{\text{TMD}}{=} \sigma_0 H(Q) \int d^2 \mathbf{b}_T e^{i \mathbf{b}_T \cdot \mathbf{q}_T} F_1(x_1, \mathbf{b}_T, Q, Q^2) F_2(x_2, \mathbf{b}_T, Q, Q^2) + \mathcal{O}\left[\left(\frac{q_T}{Q}\right)^m\right]$$

- 🍎 Collinear and TMD factorisations may eventually be **matched** to produce accurate results over the full  $q_T$  spectrum.

# TMD factorisation

- 🍎 TMD factorisation introduces two independent scales:
  - 🍎 the **renormalisation scale  $\mu$** , originating from the UV renormalisation,
  - 🍎 the **rapidity scale  $\zeta$** , originating from the cancellation of rapidity divergences.

- 🍎 The respective **evolution equations** are:

$$\frac{\partial \ln F}{\partial \ln \sqrt{\zeta}} = K(\mu_0) - \int_{\mu_0}^{\mu} \frac{d\mu'}{\mu'} \gamma_K(\alpha_s(\mu'))$$

$$\frac{\partial \ln F}{\partial \ln \mu} = \gamma_F(\alpha_s(\mu)) - \gamma_K(\alpha_s(\mu)) \ln \frac{\sqrt{\zeta}}{\mu}$$

- 🍎 At small  $b_T$ , TMDs can be matched onto collinear distributions:

$$F(\mu, \zeta) = C(\mu, \zeta) \otimes f(\mu)$$

- 🍎 The solution final is:

$$F(\mu, \zeta) = \exp \left\{ K(\mu_0) \ln \frac{\sqrt{\zeta}}{\sqrt{\zeta_0}} + \int_{\mu_0}^{\mu} \frac{d\mu'}{\mu'} \left[ \gamma_F(\alpha_s(\mu')) - \gamma_K(\alpha_s(\mu')) \ln \frac{\sqrt{\zeta}}{\mu'} \right] \right\} C(\mu_0, \zeta_0) \otimes f(\mu_0)$$

- 🍎 Anomalous dims. and matching funcs. **perturbatively** computable. 4

# TMD factorisation

- 🍎 TMD factorisation introduces two independent scales:
  - 🍎 the **renormalisation scale  $\mu$** , originating from the UV renormalisation,
  - 🍎 the **rapidity scale  $\zeta$** , originating from the cancellation of rapidity divergences.

- 🍎 The respective **evolution equations** are:

$$\frac{\partial \ln F}{\partial \ln \sqrt{\zeta}} = K(\mu_0) - \int_{\mu_0}^{\mu} \frac{d\mu'}{\mu'} \gamma_K(\alpha_s(\mu'))$$

$$\frac{\partial \ln F}{\partial \ln \mu} = \gamma_F(\alpha_s(\mu)) - \gamma_K(\alpha_s(\mu)) \ln \frac{\sqrt{\zeta}}{\mu}$$

- 🍎 At small  $b_T$ , TMDs can be matched onto collinear distributions:

$$F(\mu, \zeta) = C(\mu, \zeta) \otimes f(\mu)$$

- 🍎 The solution final is:

$$\mu_b = b_0 / b_T$$

$$F(\mu, \zeta) = \exp \left\{ K(\mu_0) \ln \frac{\sqrt{\zeta}}{\sqrt{\zeta_0}} + \int_{\mu_0}^{\mu} \frac{d\mu'}{\mu'} \left[ \gamma_F(\alpha_s(\mu')) - \gamma_K(\alpha_s(\mu')) \ln \frac{\sqrt{\zeta}}{\mu'} \right] \right\} C(\mu_0, \zeta_0) \otimes f(\mu_0)$$

- 🍎 Anomalous dims. and matching funcs. **perturbatively** computable. 5

# TMD factorisation

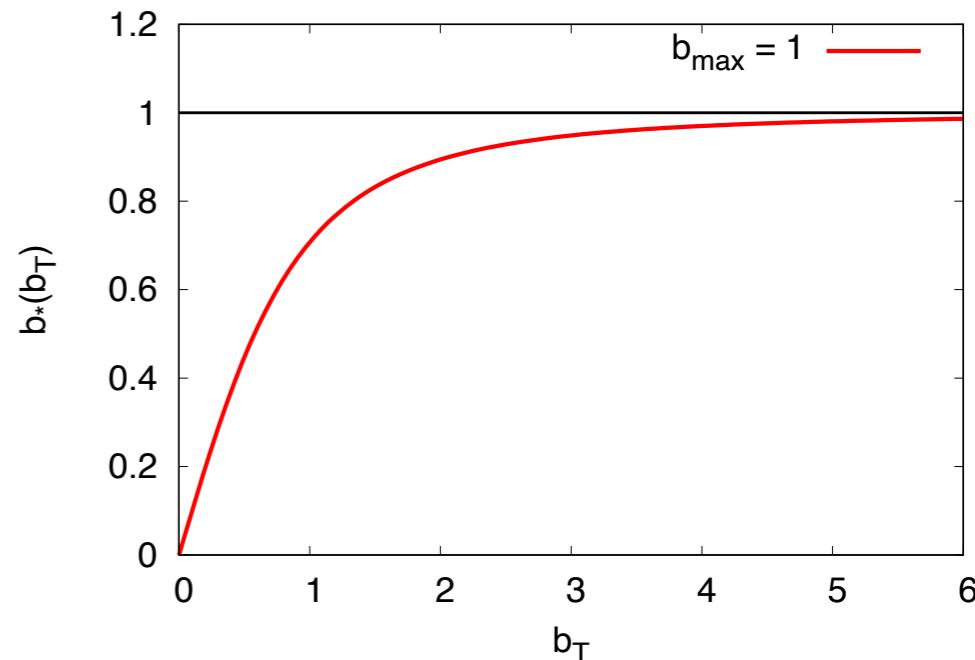
When integrating over  $b_T$ , **large values of  $b_T$**  give raise to low scales in the **non-perturbative** region.

Introduce the so-called  **$b_*$ -prescription**:

$$b_*(b_T) = \frac{b_T}{\sqrt{1 + b_T^2/b_{\max}^2}}$$

and rewrite:

$$F(x, b_T, \mu, \zeta) = \left[ \frac{F(x, b_T, \mu, \zeta)}{F(x, b_*(b_T), \mu, \zeta)} \right] F(x, b_*(b_T), \mu, \zeta) \equiv f_{\text{NP}}(x, b_T, \zeta) F(x, b_*(b_T), \mu, \zeta)$$



# TMD factorisation

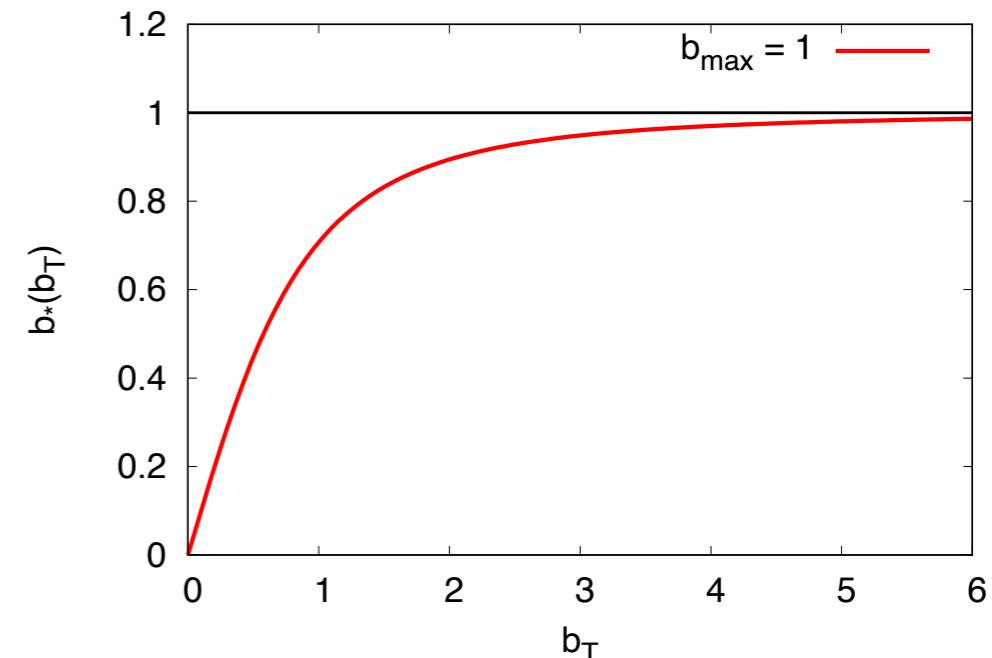
When integrating over  $b_T$ , **large values of  $b_T$**  give raise to low scales in the **non-perturbative** region.

Introduce the so-called  **$b_*$ -prescription**:

$$b_*(b_T) = \frac{b_T}{\sqrt{1 + b_T^2/b_{\max}^2}}$$

and rewrite:

$$F(x, b_T, \mu, \zeta) = \left[ \frac{F(x, b_T, \mu, \zeta)}{F(x, b_*(b_T), \mu, \zeta)} \right] F(x, b_*(b_T), \mu, \zeta) \equiv f_{\text{NP}}(x, b_T, \zeta) F(x, b_*(b_T), \mu, \zeta)$$



Purely perturbative

Non-perturbative,  
determine from data

Properties of  $f_{\text{NP}}$ :

has to go to **one** as  $b_T$  goes to zero: reproduce the fully perturbative regime,

has to go to **zero** as  $b_T$  becomes large: mimic the Sudakov suppression.

Bottom line: avoidance of the non-perturbative region upon integration in  $b_T$  implies the presence of **both**  $b_*$ -prescription and  $f_{\text{NP}}$ .

# TMD factorisation

🍎 Final expression:

$$\begin{aligned} F_{f/P}(x, \mathbf{b}_T; \mu, \zeta) &= \sum_j C_{f/j}(x, b_*; \mu_b, \mu_b^2) \otimes f_{j/P}(x, \mu_b) && : A \\ &\times \exp \left\{ K(b_*; \mu_b) \ln \frac{\sqrt{\zeta}}{\mu_b} + \int_{\mu_b}^{\mu} \frac{d\mu'}{\mu'} \left[ \gamma_F - \gamma_K \ln \frac{\sqrt{\zeta}}{\mu'} \right] \right\} && : B \\ &\times \exp \left\{ g_{j/P}(x, b_T) + g_K(b_T) \ln \frac{\sqrt{\zeta_F}}{\sqrt{\zeta_{F,0}}} \right\} && : C \end{aligned}$$

# TMD factorisation

Final expression:

$$F_{f/P}(x, \mathbf{b}_T; \mu, \zeta) = \sum_j C_{f/j}(x, b_*; \mu_b, \mu_b^2) \otimes f_{j/P}(x, \mu_b) : A$$
$$\times \exp \left\{ K(b_*; \mu_b) \ln \frac{\sqrt{\zeta}}{\mu_b} + \int_{\mu_b}^{\mu} \frac{d\mu'}{\mu'} \left[ \gamma_F - \gamma_K \ln \frac{\sqrt{\zeta}}{\mu'} \right] \right\} : B$$
$$\times \exp \left\{ \underline{g_{j/P}(x, b_T)} + g_K(b_T) \ln \frac{\sqrt{\zeta_F}}{\sqrt{\zeta_{F,0}}} \right\} : C$$

- matching onto the collinear region at  $b_T \ll 1/\Lambda_{\text{QCD}}$ ,
- factorises as *hard* (perturbative) and *longitudinal* (*i.e.* collinear, non-perturbative).

# TMD factorisation

Final expression:

$$\begin{aligned}
 F_{f/P}(x, \mathbf{b}_T; \mu, \zeta) &= \sum_j C_{f/j}(x, b_*; \mu_b, \mu_b^2) \otimes f_{j/P}(x, \mu_b) && : A \\
 &\times \exp \left\{ K(b_*; \mu_b) \ln \frac{\sqrt{\zeta}}{\mu_b} + \int_{\mu_b}^{\mu} \frac{d\mu'}{\mu'} \left[ \gamma_F - \gamma_K \ln \frac{\sqrt{\zeta}}{\mu'} \right] \right\} && : B \\
 &\times \exp \left\{ \underline{g_{j/P}(x, b_T)} + \underline{g_K(b_T) \ln \frac{\sqrt{\zeta_F}}{\sqrt{\zeta_{F,0}}}} \right\} && : C
 \end{aligned}$$

- matching onto the collinear region at  $b_T \ll 1/\Lambda_{\text{QCD}}$ ,
- factorises as *hard* (perturbative) and *longitudinal* (*i.e.* collinear, non-perturbative).
- CS and RGE evolution,
- evolution in  $\mu$  and  $\zeta$ ,
- perturbative.

# TMD factorisation

Final expression:

$$F_{f/P}(x, \mathbf{b}_T; \mu, \zeta) = \sum_j C_{f/j}(x, b_*; \mu_b, \mu_b^2) \otimes f_{j/P}(x, \mu_b)$$

: A

$$\times \exp \left\{ K(b_*; \mu_b) \ln \frac{\sqrt{\zeta}}{\mu_b} + \int_{\mu_b}^{\mu} \frac{d\mu'}{\mu'} \left[ \gamma_F - \gamma_K \ln \frac{\sqrt{\zeta}}{\mu'} \right] \right\}$$

: B

$$\times \exp \left\{ g_{j/P}(x, b_T) + g_K(b_T) \ln \frac{\sqrt{\zeta_F}}{\sqrt{\zeta_{F,0}}} \right\}$$

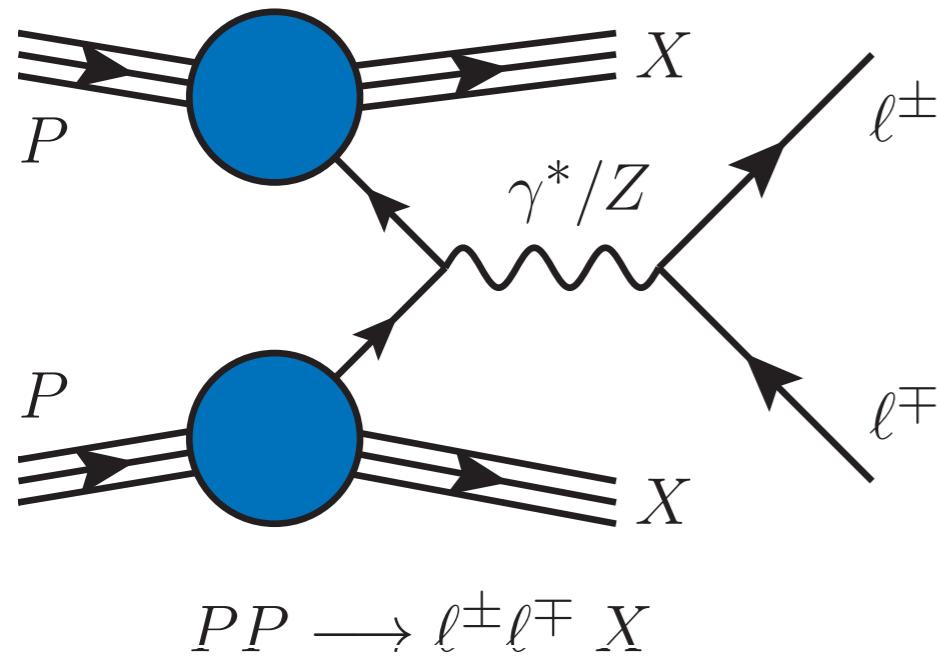
: C

- matching onto the collinear region at  $b_T \ll 1/\Lambda_{\text{QCD}}$ ,
- factorises as *hard* (perturbative) and *longitudinal* (*i.e.* collinear, non-perturbative).
  - avoid the Landau pole,
  - $f_{\text{NP}}$  accounts for the introduction of  $b_*$ ,
  - $f_{\text{NP}}$  is non-perturbative thus **fit** to data.
- CS and RGE evolution,  
evolution in  $\mu$  and  $\zeta$ ,  
perturbative.

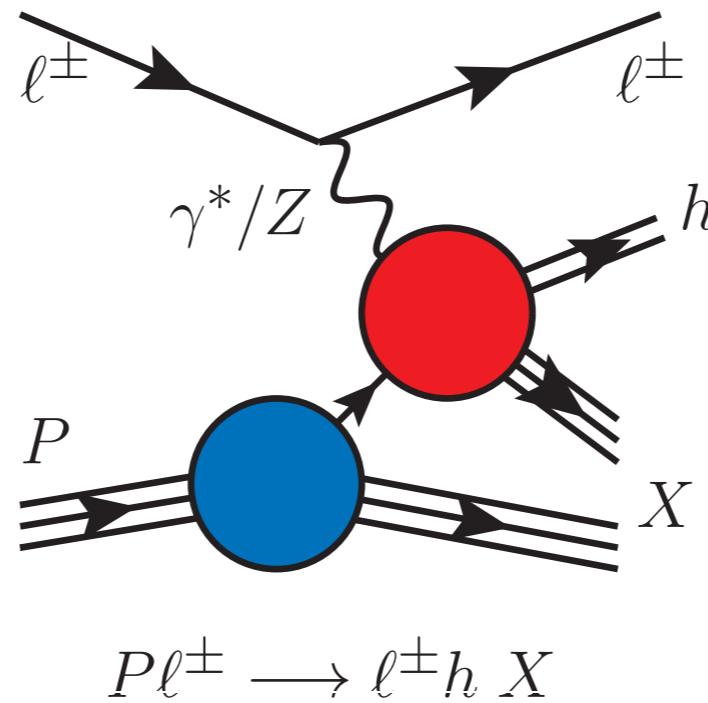
# Factorising processes

- Processes for which leading-power TMD factorisation has been **proven**:

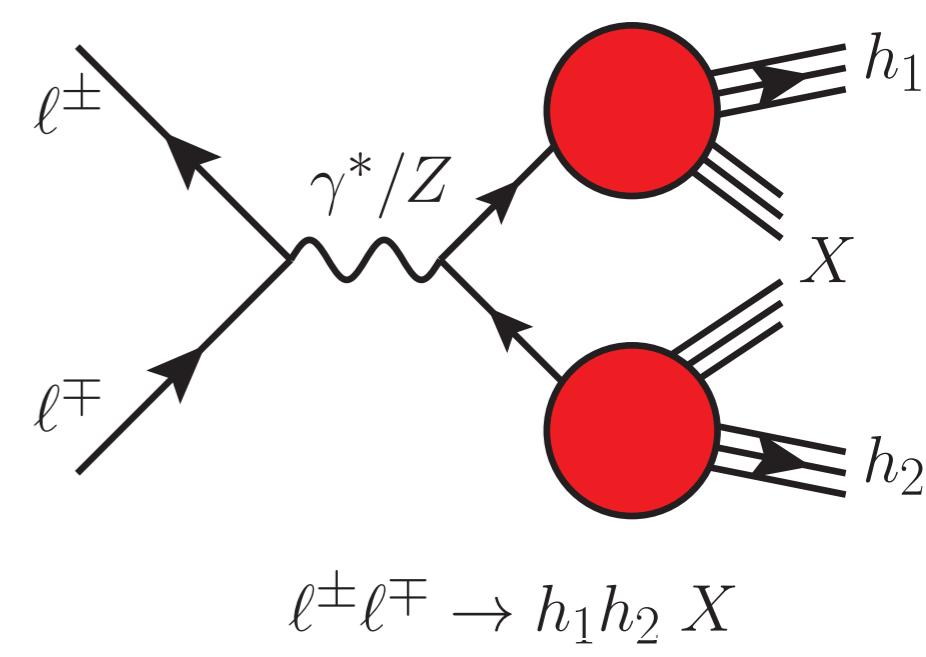
Drell-Yan



Semi-inclusive DIS



$e^+e^-$  annihilation



- Two TMD PDFs:

- Lots of data:

- low-energy: FNAL,

- mid-energy: RHIC,

- high-energy: Tevatron, LHC.

- Examples of other processes:

- thrust and  $p_{hT}$  distributions in single-hadron production in  $e^+e^-$ ,

- hadron-in-jet production,

- ...

- One TMD PDF one FF:

- many precise data points:

- HERMES at DESY,

- COMPASS at CERN.

- Two TMD FFs:

- di-hadron prod. from:

- BELLE at KEK,

- BABAR at SLAC.

# Logarithmic counting

$$\left( \frac{d\sigma}{dq_T} \right)_{\text{res.}} \stackrel{\text{TMD}}{=} \sigma_0 \textcolor{violet}{H}(Q) \int d^2 \mathbf{b}_T e^{i \mathbf{b}_T \cdot \mathbf{q}_T} F_1(x_1, \mathbf{b}_T, Q, Q^2) F_2(x_2, \mathbf{b}_T, Q, Q^2)$$

$$F_i = \sum_j (\textcolor{brown}{C}_{i/j} \otimes f_j) \exp \left\{ \textcolor{green}{K} \ln \frac{\sqrt{\zeta}}{\mu_b} + \int_{\mu_b}^{\mu} \frac{d\mu'}{\mu'} \left[ \textcolor{blue}{\gamma_F} - \textcolor{red}{\gamma_K} \ln \frac{\sqrt{\zeta}}{\mu'} \right] \right\}$$

Accuracy	$\gamma_K$	$\gamma_F$	$K$	$C_{f/j}$	$H$	FFs/PDFs/ $\alpha_s$
LL	$\alpha_s$	-	-	1	1	-
NLL	$\alpha_s^2$	$\alpha_s$	$\alpha_s$	1	1	LO
NLL'	$\alpha_s^2$	$\alpha_s$	$\alpha_s$	$\alpha_s$	$\alpha_s$	LO
$N^2LL$	$\alpha_s^3$	$\alpha_s^2$	$\alpha_s^2$	$\alpha_s$	$\alpha_s$	NLO
$N^2LL'$	$\alpha_s^3$	$\alpha_s^2$	$\alpha_s^2$	$\alpha_s^2$	$\alpha_s^2$	NLO
$N^3LL$	$\alpha_s^4$	$\alpha_s^3$	$\alpha_s^3$	$\alpha_s^2$	$\alpha_s^2$	NNLO
$N^3LL'$	$\alpha_s^4$	$\alpha_s^3$	$\alpha_s^3$	$\alpha_s^3$	$\alpha_s^3$	NNLO

# Unpolarised TMD extractions

## *A selection of fits*

	Accuracy	SIDIS	Drell-Yan	N. of points
DWS 1984, <a href="#">CERN-TH.3987/84</a>	NLL	✗	✓	a few
BLNY 2003, <a href="#">hep-ph/0212159</a>	NLL'-NNLL	✗	✓	116
Pavia 2013, <a href="#">1309.3507</a>	No evolution	✓	✗	1538 (HERMES)
Torino 2014, <a href="#">1312.6261</a>	No evolution	✓	✗	576 (H) 6284 (C)
DEMS 2014, <a href="#">1407.3311</a>	NNLL	✗	✓	223
Pavia 2017, <a href="#">1703.10157</a>	NLL	✓	✓	8059
SV 2017, <a href="#">1706.01473</a>	N <sup>3</sup> LL	✗	✓ (LHC)	309
BSV 2019, <a href="#">1902.08474</a>	N <sup>3</sup> LL	✗	✓ (LHC)	457
SV 2019, <a href="#">1912.06532</a>	N <sup>3</sup> LL <sup>(-)</sup>	✓	✓ (LHC)	1039
Pavia 2019, <a href="#">1912.07550</a>	N <sup>3</sup> LL	✗	✓ (LHC)	353
SV+ 2022, <a href="#">2201.07114</a>	N <sup>3</sup> LL	✗	✓ (LHC)	507/309
MAPTMD 2022, <a href="#">2206.07598</a>	N <sup>3</sup> LL <sup>(-)</sup>	✓	✓ (LHC)	2031

# Unpolarised TMD extractions

## *A selection of fits*

	Accuracy	SIDIS	Drell-Yan	N. of points
DWS 1984, <a href="#">CERN-TH.3987/84</a>	NLL	✗	✓	a few
BLNY 2003, <a href="#">hep-ph/0212159</a>	NLL'-NNLL	✗	✓	116
Pavia 2013, <a href="#">1309.3507</a>	No evolution	✓	✗	1538 (HERMES)
Torino 2014, <a href="#">1312.6261</a>	No evolution	✓	✗	576 (H) 6284 (C)
DEMS 2014, <a href="#">1407.3311</a>	NNLL	✗	✓	223
Pavia 2017, <a href="#">1703.10157</a>	NLL	✓	✓	8059
SV 2017, <a href="#">1706.01473</a>	N <sup>3</sup> LL	✗	✓ (LHC)	309
BSV 2019, <a href="#">1902.08474</a>	N <sup>3</sup> LL	✗	✓ (LHC)	457
SV 2019, <a href="#">1912.06532</a>	N <sup>3</sup> LL <sup>(-)</sup>	✓	✓ (LHC)	1039
Pavia 2019, <a href="#">1912.07550</a>	N <sup>3</sup> LL	✗	✓ (LHC)	353
SV+ 2022, <a href="#">2201.07114</a>	N <sup>3</sup> LL	✗	✓ (LHC)	507/309
MAPTMD 2022, <a href="#">2206.07598</a>	N <sup>3</sup> LL <sup>(-)</sup>	✓	✓ (LHC)	2031

# Unpolarised TMD extractions

*Many more studies and extractions...*

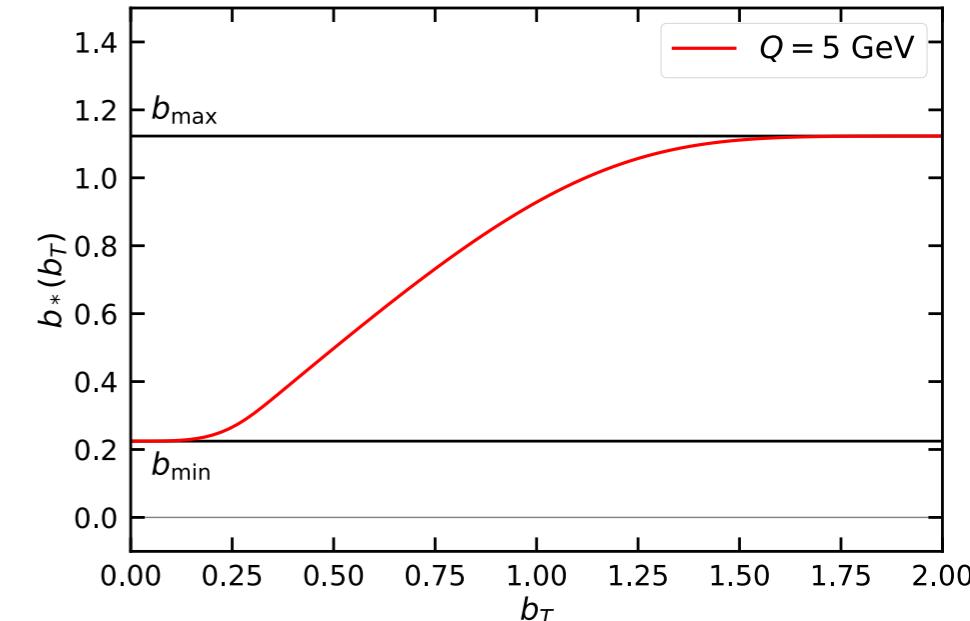
- 🍎 TMD fragmentation functions from  $e^+e^-$  data [[2108.04182](#), [1704.08882](#)]
- 🍎  $W$  production in  $pp$  collisions [[2011.05351](#)]
- 🍎 Di-jet and heavy-meson pair production in DIS [[2008.07531](#), [2111.03703](#)]
- 🍎 Dijet production in  $pp$  collisions [*e.g.* [1807.07573](#)]
- 🍎 hadron-in-jet production [[1612.04817](#)]
- 🍎 Model-independent prescription to extract TMDs [[2201.07237](#)]
- 🍎 Parton-branching methods [*e.g.* [1804.11152](#)]
- 🍎  $q_T$ -resummation based extractions [[2203.05394](#)]
- 🍎 Study of the Sivers TMDs [[1308.5003](#), [2004.14278](#), [2009.10710](#), [2103.03270](#),...]
- 🍎 Pion TMDs [[1907.10356](#), [2210.01733](#)]
- 🍎 TMD flavour dependence [[1807.02101](#), [2201.07114](#)]
- 🍎 ...

# MAPTMD 2022

## Main settings

- 🍎  $b_*$  prescription:

$$b_*(b_T) = b_{\max} \left( \frac{1 - e^{-b_T^4/b_{\max}^4}}{1 - e^{-b_T^4/b_{\min}^4}} \right)^{1/4} \quad \text{with} \quad \begin{cases} b_{\max} = 2e^{-\gamma_E} \\ b_{\min} = b_{\max}/Q \end{cases}$$



- 🍎 Non-perturbative function  $f_{\text{NP}}$ :

🍎 evolution (CS kernel):  $g_K(\mathbf{b}_T^2) = -g_2^2 \frac{\mathbf{b}_T^2}{2}$

- 🍎 PDFs:

$$f_{1NP}(x, \mathbf{b}_T^2; \zeta, Q_0) = \frac{g_1(x) e^{-g_1(x) \frac{\mathbf{b}_T^2}{4}} + \lambda^2 g_{1B}^2(x) \left[ 1 - g_{1B}(x) \frac{\mathbf{b}_T^2}{4} \right] e^{-g_{1B}(x) \frac{\mathbf{b}_T^2}{4}} + \lambda_2^2 g_{1C}(x) e^{-g_{1C}(x) \frac{\mathbf{b}_T^2}{4}}}{g_1(x) + \lambda^2 g_{1B}^2(x) + \lambda_2^2 g_{1C}(x)} \left[ \frac{\zeta}{Q_0^2} \right]^{g_K(\mathbf{b}_T^2)/2}$$

- 🍎 FFs:

$$D_{1NP}(z, \mathbf{b}_T^2; \zeta, Q_0) = \frac{g_3(z) e^{-g_3(z) \frac{\mathbf{b}_T^2}{4z^2}} + \frac{\lambda_F}{z^2} g_{3B}^2(z) \left[ 1 - g_{3B}(z) \frac{\mathbf{b}_T^2}{4z^2} \right] e^{-g_{3B}(z) \frac{\mathbf{b}_T^2}{4z^2}}}{g_3(z) + \frac{\lambda_F}{z^2} g_{3B}^2(z)} \left[ \frac{\zeta}{Q_0^2} \right]^{g_K(\mathbf{b}_T^2)/2}$$

$$g_{\{1,1B,1C\}}(x) = N_{\{1,1B,1C\}} \frac{x^{\sigma_{\{1,2,3\}}} (1-x)^{\alpha_{\{1,2,3\}}^2}}{\hat{x}^{\sigma_{\{1,2,3\}}} (1-\hat{x})^{\alpha_{\{1,2,3\}}^2}} \quad g_{\{3,3B\}}(z) = N_{\{3,3B\}} \frac{(z^{\beta_{\{1,2\}}} + \delta_{\{1,2\}}^2)(1-z)^{\gamma_{\{1,2\}}^2}}{(\hat{z}^{\beta_{\{1,2\}}} + \delta_{\{1,2\}}^2)(1-\hat{z})^{\gamma_{\{1,2\}}^2}}$$

- 🍎 11 (PDFs) + 9 (FFs) + 1 (evol): **21 free parameters** to fit to data.
- 🍎 Perturbative accuracies: **N<sup>3</sup>LL(-)**.
- 🍎 **Monte Carlo** method for the experimental error propagation.

# MAPTMD 2022

## Dataset



DY data:

- 🍎 fixed-target low-energy DY,
- 🍎 RHIC data,
- 🍎 LHC and Tevatron data,
- 🍎 selection cut  $q_T / Q < 0.2$ ,
- 🍎 484 data points.



SIDIS data:

- 🍎 HERMES and COMPASS,
- 🍎  $P_{hT}|_{\max} = \min[\min[0.2Q, 0.5zQ] + 0.3 \text{ GeV}, zQ]$

🍎  $Q > 1.4 \text{ GeV}$ ,  $0.2 < z < 0.7$ ,

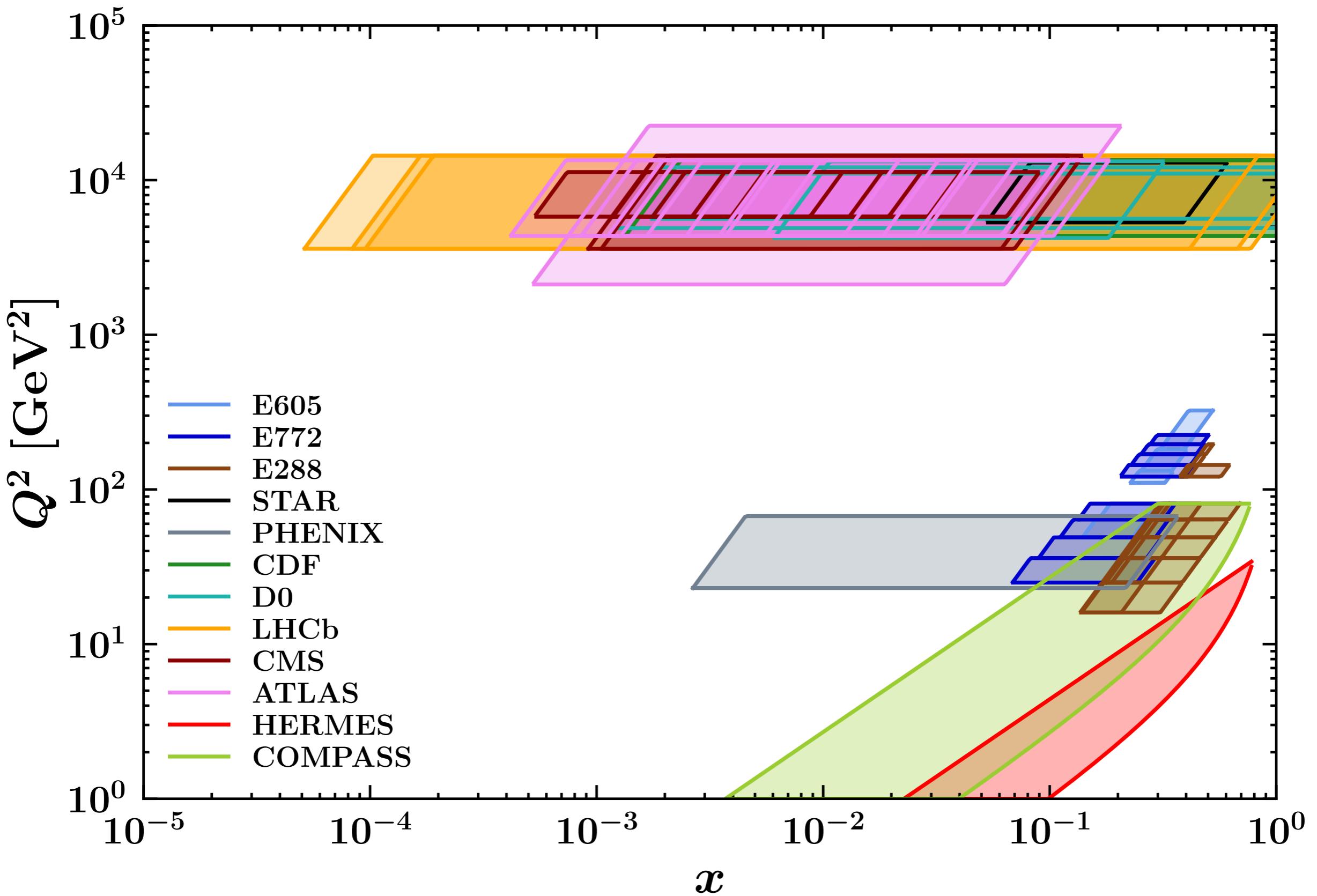
🍎 1547 points.

Experiment	$N_{\text{dat}}$	Observable	$\sqrt{s} [\text{GeV}]$	$Q [\text{GeV}]$	$y$ or $x_F$	Lepton cuts	Ref.
E605	50	$Ed^3\sigma/d^3q$	38.8	7 - 18	$x_F = 0.1$	-	[55]
E772	53	$Ed^3\sigma/d^3q$	38.8	5 - 15	$0.1 < x_F < 0.3$	-	[51]
E288 200 GeV	30	$Ed^3\sigma/d^3q$	19.4	4 - 9	$y = 0.40$	-	[56]
E288 300 GeV	39	$Ed^3\sigma/d^3q$	23.8	4 - 12	$y = 0.21$	-	[56]
E288 400 GeV	61	$Ed^3\sigma/d^3q$	27.4	5 - 14	$y = 0.03$	-	[56]
STAR 510	7	$d\sigma/d q_T $	510	73 - 114	$ y  < 1$	$p_{T\ell} > 25 \text{ GeV}$ $ \eta_\ell  < 1$	-
PHENIX200	2	$d\sigma/d q_T $	200	4.8 - 8.2	$1.2 < y < 2.2$	-	[52]
CDF Run I	25	$d\sigma/d q_T $	1800	66 - 116	Inclusive	-	[57]
CDF Run II	26	$d\sigma/d q_T $	1960	66 - 116	Inclusive	-	[58]
D0 Run I	12	$d\sigma/d q_T $	1800	75 - 105	Inclusive	-	[59]
D0 Run II	5	$(1/\sigma)d\sigma/d q_T $	1960	70 - 110	Inclusive	-	[60]
D0 Run II ( $\mu$ )	3	$(1/\sigma)d\sigma/d q_T $	1960	65 - 115	$ y  < 1.7$	$p_{T\ell} > 15 \text{ GeV}$ $ \eta_\ell  < 1.7$	[61]
LHCb 7 TeV	7	$d\sigma/d q_T $	7000	60 - 120	$2 < y < 4.5$	$p_{T\ell} > 20 \text{ GeV}$ $2 < \eta_\ell < 4.5$	[62]
LHCb 8 TeV	7	$d\sigma/d q_T $	8000	60 - 120	$2 < y < 4.5$	$p_{T\ell} > 20 \text{ GeV}$ $2 < \eta_\ell < 4.5$	[63]
LHCb 13 TeV	7	$d\sigma/d q_T $	13000	60 - 120	$2 < y < 4.5$	$p_{T\ell} > 20 \text{ GeV}$ $2 < \eta_\ell < 4.5$	[64]
CMS 7 TeV	4	$(1/\sigma)d\sigma/d q_T $	7000	60 - 120	$ y  < 2.1$	$p_{T\ell} > 20 \text{ GeV}$ $ \eta_\ell  < 2.1$	[65]
CMS 8 TeV	4	$(1/\sigma)d\sigma/d q_T $	8000	60 - 120	$ y  < 2.1$	$p_{T\ell} > 15 \text{ GeV}$ $ \eta_\ell  < 2.1$	[66]
CMS 13 TeV	70	$d\sigma/d q_T $	13000	76 - 106	$ y  < 0.4$ $0.4 <  y  < 0.8$ $0.8 <  y  < 1.2$ $1.2 <  y  < 1.6$ $1.6 <  y  < 2.4$	$p_{T\ell} > 25 \text{ GeV}$ $ \eta_\ell  < 2.4$	[53]
ATLAS 7 TeV	6 6 6	$(1/\sigma)d\sigma/d q_T $	7000	66 - 116	$ y  < 1$ $1 <  y  < 2$ $2 <  y  < 2.4$	$p_{T\ell} > 20 \text{ GeV}$ $ \eta_\ell  < 2.4$	[67]
ATLAS 8 TeV on-peak	6 6 6 6 6	$(1/\sigma)d\sigma/d q_T $	8000	66 - 116	$ y  < 0.4$ $0.4 <  y  < 0.8$ $0.8 <  y  < 1.2$ $1.2 <  y  < 1.6$ $1.6 <  y  < 2$ $2 <  y  < 2.4$	$p_{T\ell} > 20 \text{ GeV}$ $ \eta_\ell  < 2.4$	[68]
ATLAS 8 TeV off-peak	4 8	$(1/\sigma)d\sigma/d q_T $	8000	46 - 66 116 - 150	$ y  < 2.4$	$p_{T\ell} > 20 \text{ GeV}$ $ \eta_\ell  < 2.4$	[68]
ATLAS 13 TeV	6	$(1/\sigma)d\sigma/d q_T $	13000	66 - 113	$ y  < 2.5$	$p_{T\ell} > 27 \text{ GeV}$ $ \eta_\ell  < 2.5$	[54]
Total	484						

Experiment	$N_{\text{dat}}$	Observable	Channels	$Q [\text{GeV}]$	$x$	$z$	Phase space cuts	Ref.
HERMES	344	$M(x, z,  \mathbf{P}_{hT} , Q)$	$p \rightarrow \pi^+$ $p \rightarrow \pi^-$ $p \rightarrow K^+$ $p \rightarrow K^-$ $d \rightarrow \pi^+$ $d \rightarrow \pi^-$ $d \rightarrow K^+$ $d \rightarrow K^-$	$1 - \sqrt{15}$	$0.023 < x < 0.6$ (6 bins)	$0.1 < z < 1.1$ (8 bins)	$W^2 > 10 \text{ GeV}^2$ $0.1 < y < 0.85$	[46]
COMPASS	1203	$M(x, z, \mathbf{P}_{hT}^2, Q)$	$d \rightarrow h^+$ $d \rightarrow h^-$	$1 - 9$ (5 bins)	$0.003 < x < 0.4$ (8 bins)	$0.2 < z < 0.8$ (4 bins)	$W^2 > 25 \text{ GeV}^2$ $0.1 < y < 0.9$	[72]
Total	1547							

# MAPTMD 2022

## *Kinematic coverage*



# MAPTMD 2022

## *Fit quality*

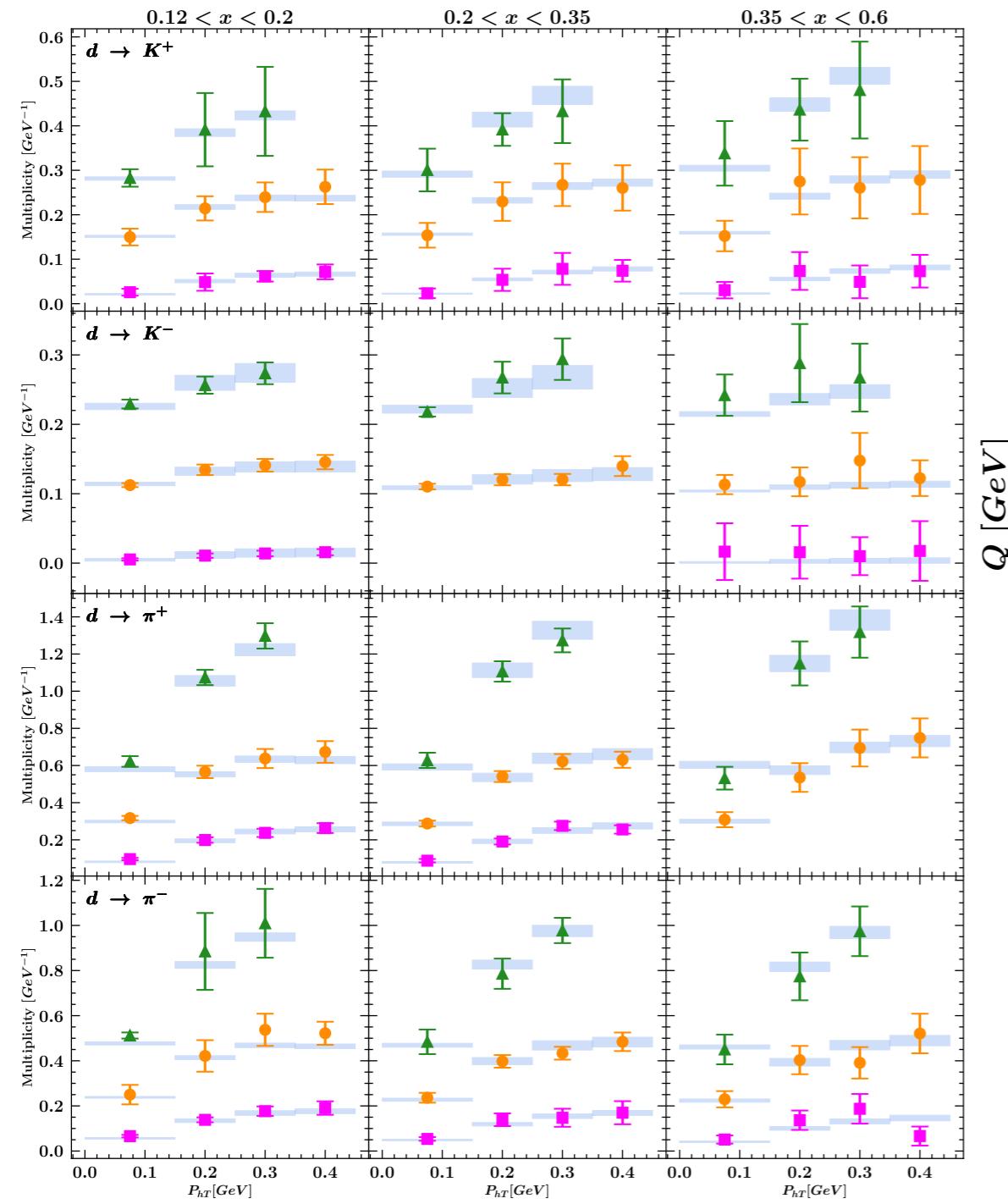
Data set	N <sup>3</sup> LL−			
	$N_{\text{dat}}$	$\chi^2_D$	$\chi^2_\lambda$	$\chi^2_0$
CDF Run I	25	0.45	0.09	0.54
CDF Run II	26	0.995	0.004	1.0
D0 Run I	12	0.67	0.01	0.68
D0 Run II	5	0.89	0.21	1.10
D0 Run II ( $\mu$ )	3	3.96	0.28	4.2
<i>Tevatron total</i>	71	0.87	0.06	0.93
LHCb 7 TeV	7	1.24	0.49	1.73
LHCb 8 TeV	7	0.78	0.36	1.14
LHCb 13 TeV	7	1.42	0.06	1.48
<i>LHCb total</i>	21	1.15	0.3	1.45
ATLAS 7 TeV	18	6.43	0.92	7.35
ATLAS 8 TeV	48	3.7	0.32	4.02
ATLAS 13 TeV	6	5.9	0.5	6.4
<i>ATLAS total</i>	72	4.56	0.48	5.05
CMS 7 TeV	4	2.21	0.10	2.31
CMS 8 TeV	4	1.938	0.001	1.94
CMS 13 TeV	70	0.36	0.02	0.37
<i>CMS total</i>	78	0.53	0.02	0.55
PHENIX 200	2	2.21	0.88	3.08
STAR 510	7	1.05	0.10	1.15
DY collider total	251	1.86	0.2	2.06

E288 200 GeV	30	0.35	0.19	0.54
E288 300 GeV	39	0.33	0.09	0.42
E288 400 GeV	61	0.5	0.11	0.61
E772	53	1.52	1.03	2.56
E605	50	1.26	0.44	1.7
DY fixed-target total	233	0.85	0.4	1.24
HERMES ( $p \rightarrow \pi^+$ )	45	0.86	0.42	1.28
HERMES ( $p \rightarrow \pi^-$ )	45	0.61	0.31	0.92
HERMES ( $p \rightarrow K^+$ )	45	0.49	0.04	0.53
HERMES ( $p \rightarrow K^-$ )	37	0.18	0.13	0.31
HERMES ( $d \rightarrow \pi^+$ )	41	0.68	0.45	1.13
HERMES ( $d \rightarrow \pi^-$ )	45	0.63	0.35	0.97
HERMES ( $d \rightarrow K^+$ )	45	0.2	0.02	0.22
HERMES ( $d \rightarrow K^-$ )	41	0.14	0.08	0.22
<i>HERMES total</i>	344	0.48	0.23	0.71
COMPASS ( $d \rightarrow h^+$ )	602	0.55	0.31	0.86
COMPASS ( $d \rightarrow h^-$ )	601	0.68	0.3	0.98
<i>COMPASS total</i>	1203	0.62	0.3	0.92
SIDIS total	1547	0.59	0.28	0.87
<b>Total</b>	<b>2031</b>	<b>0.77</b>	<b>0.29</b>	<b>1.06</b>

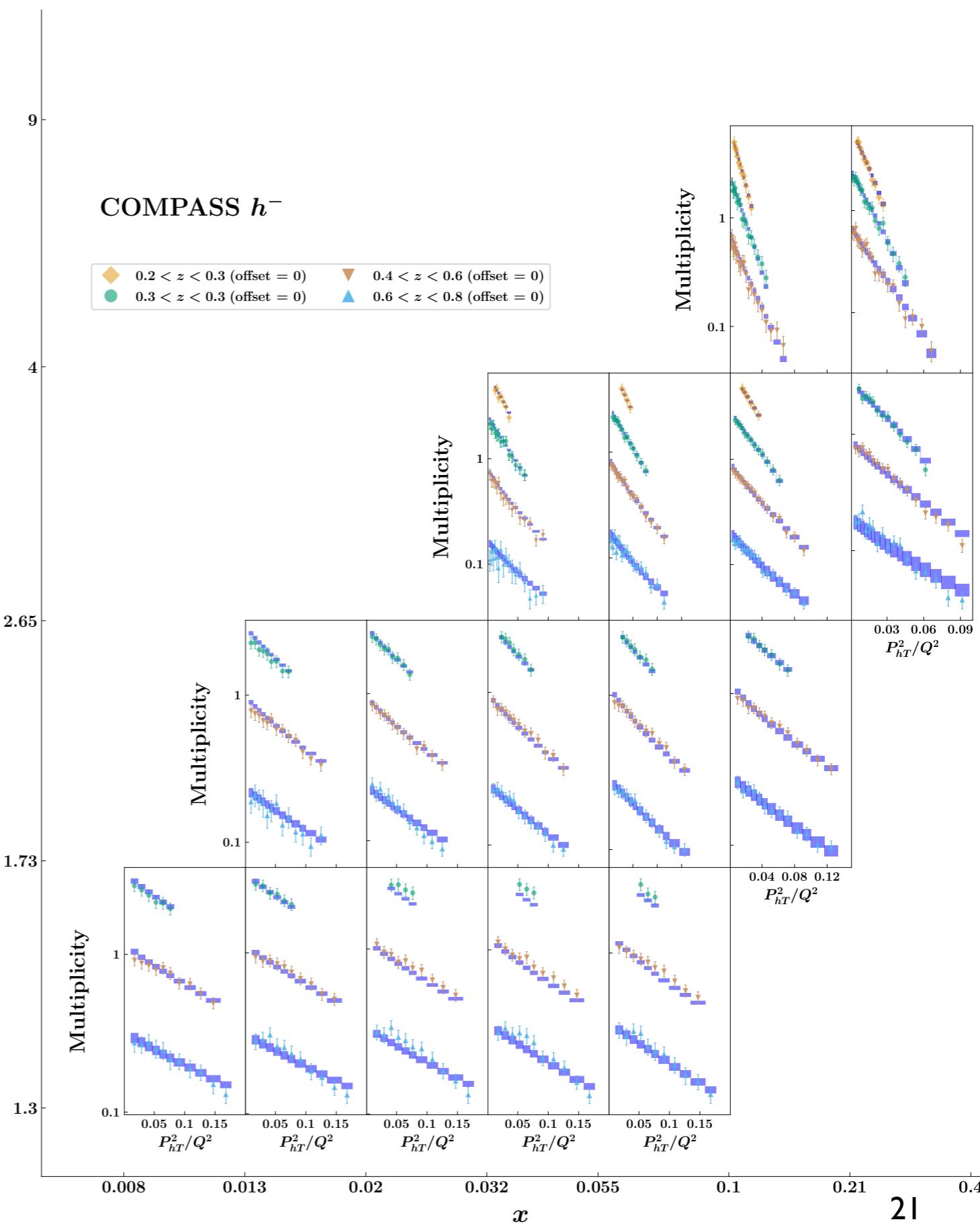
# MAPTMD 2022

## Fit quality: SIDIS

HERMES

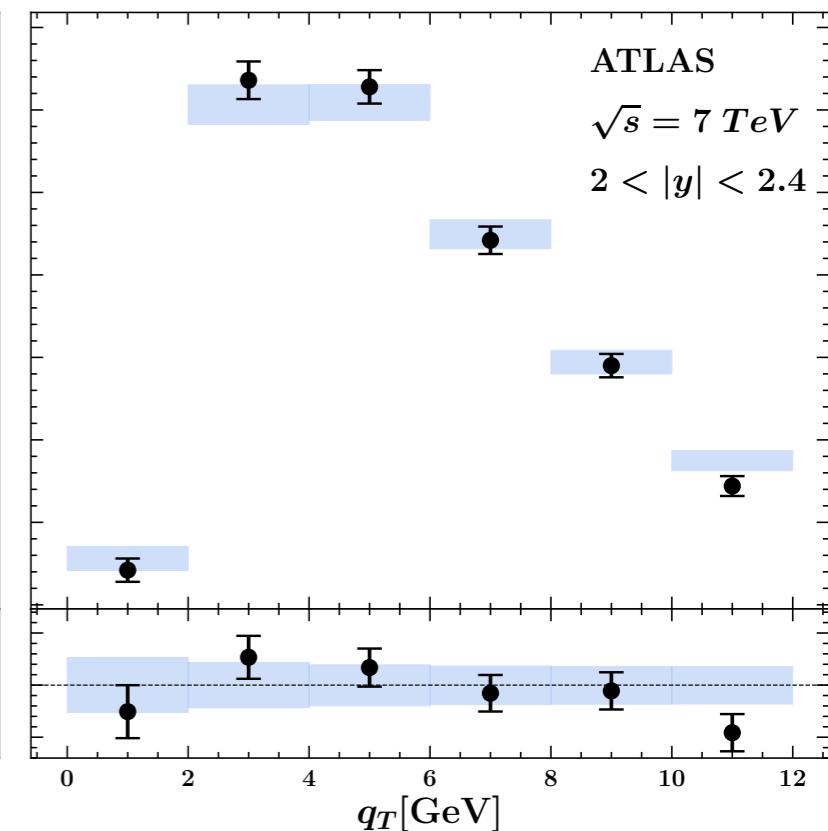
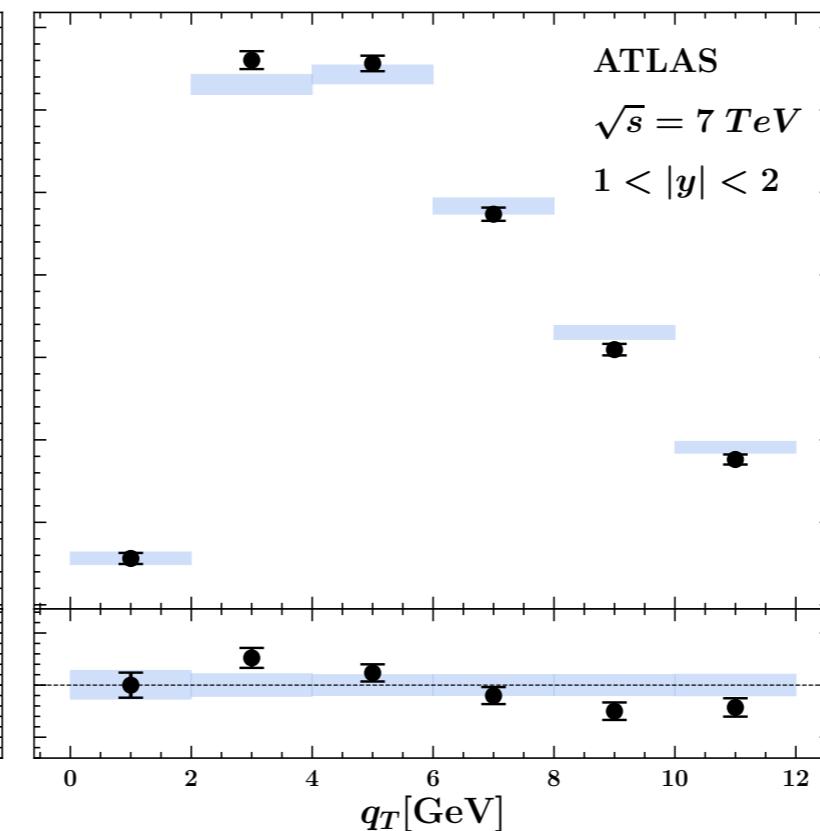
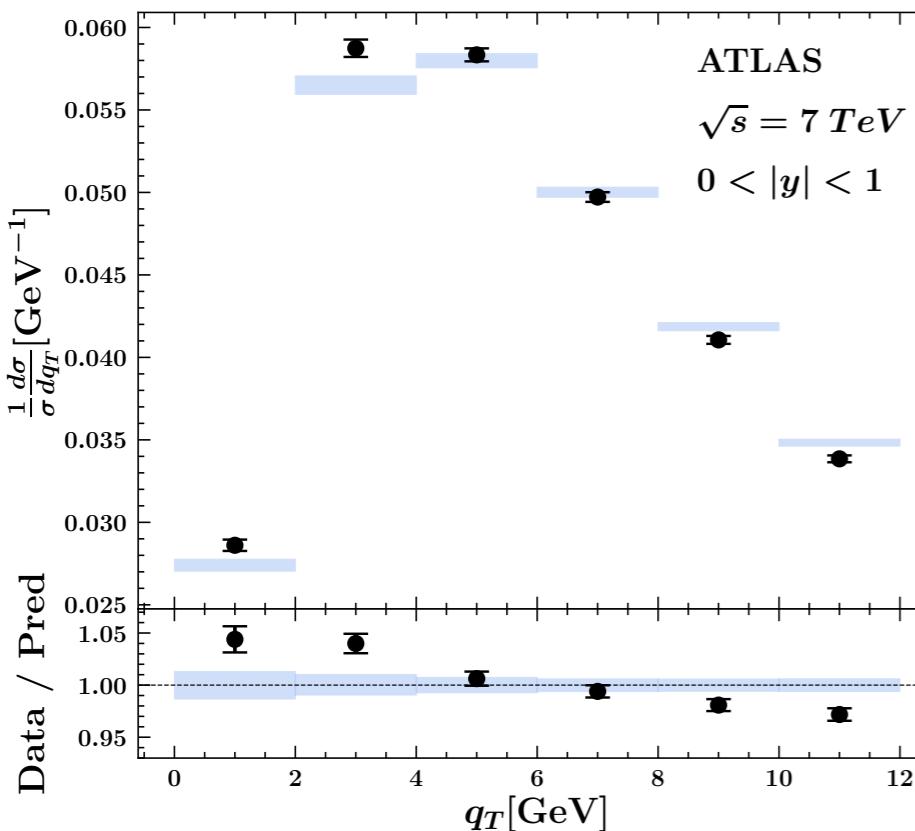
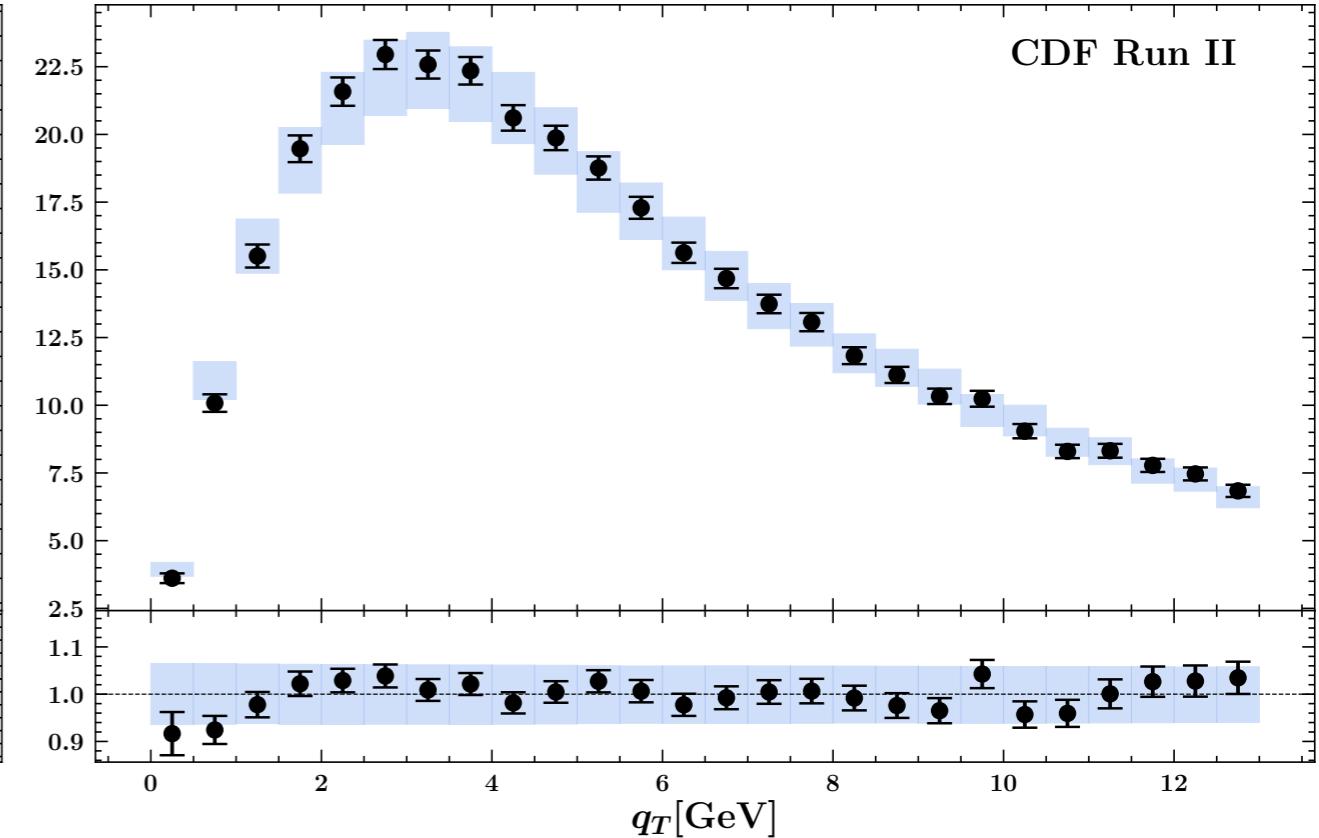
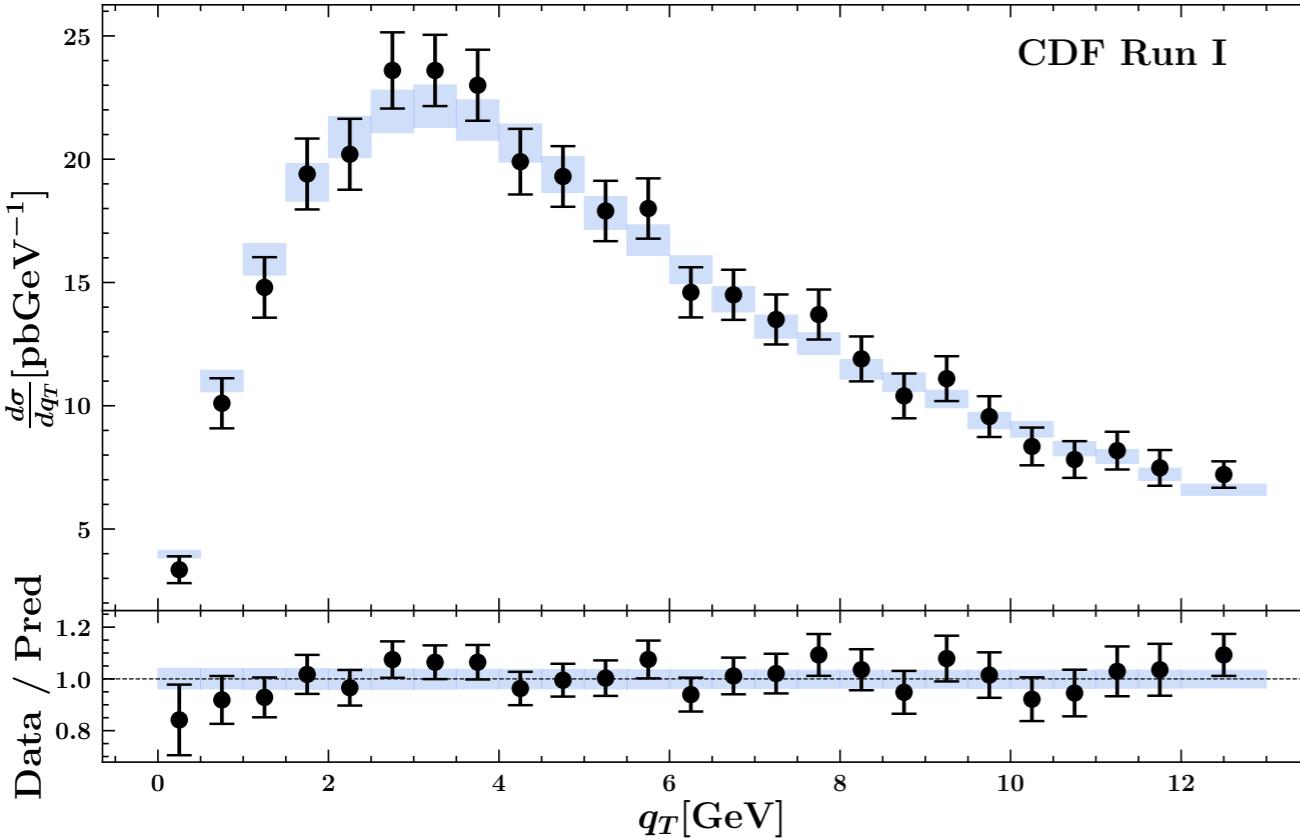


Legend:   
▲  $0.375 < z < 0.475$  (offset = 0.2)   ●  $0.475 < z < 0.6$  (offset = 0.1)   ■  $0.6 < z < 0.8$  (offset = 0)   ▲  $0.3 < z < 0.3$  (offset = 0)



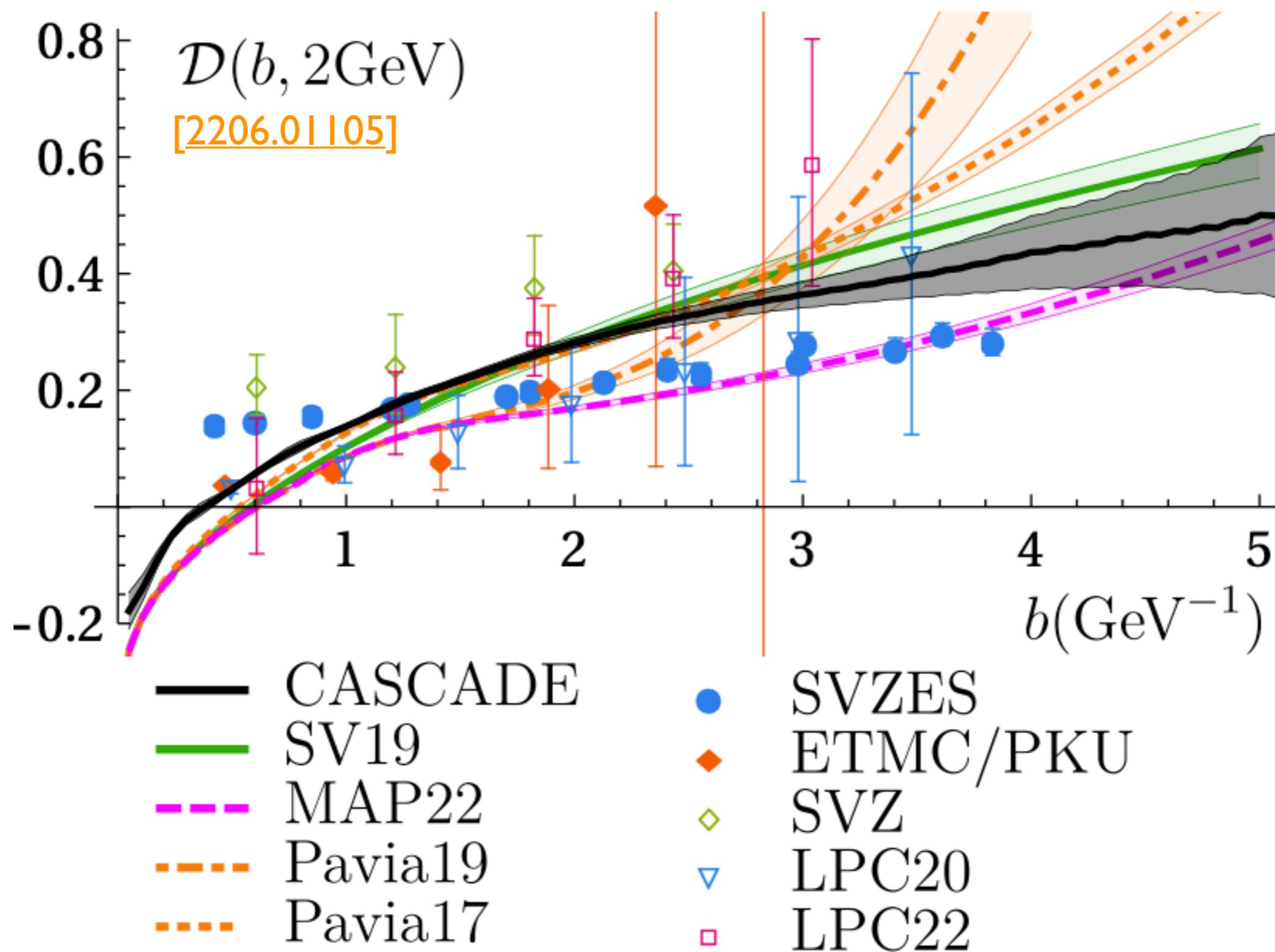
# MAPTMD 2022

*Fit quality:  $DY$*



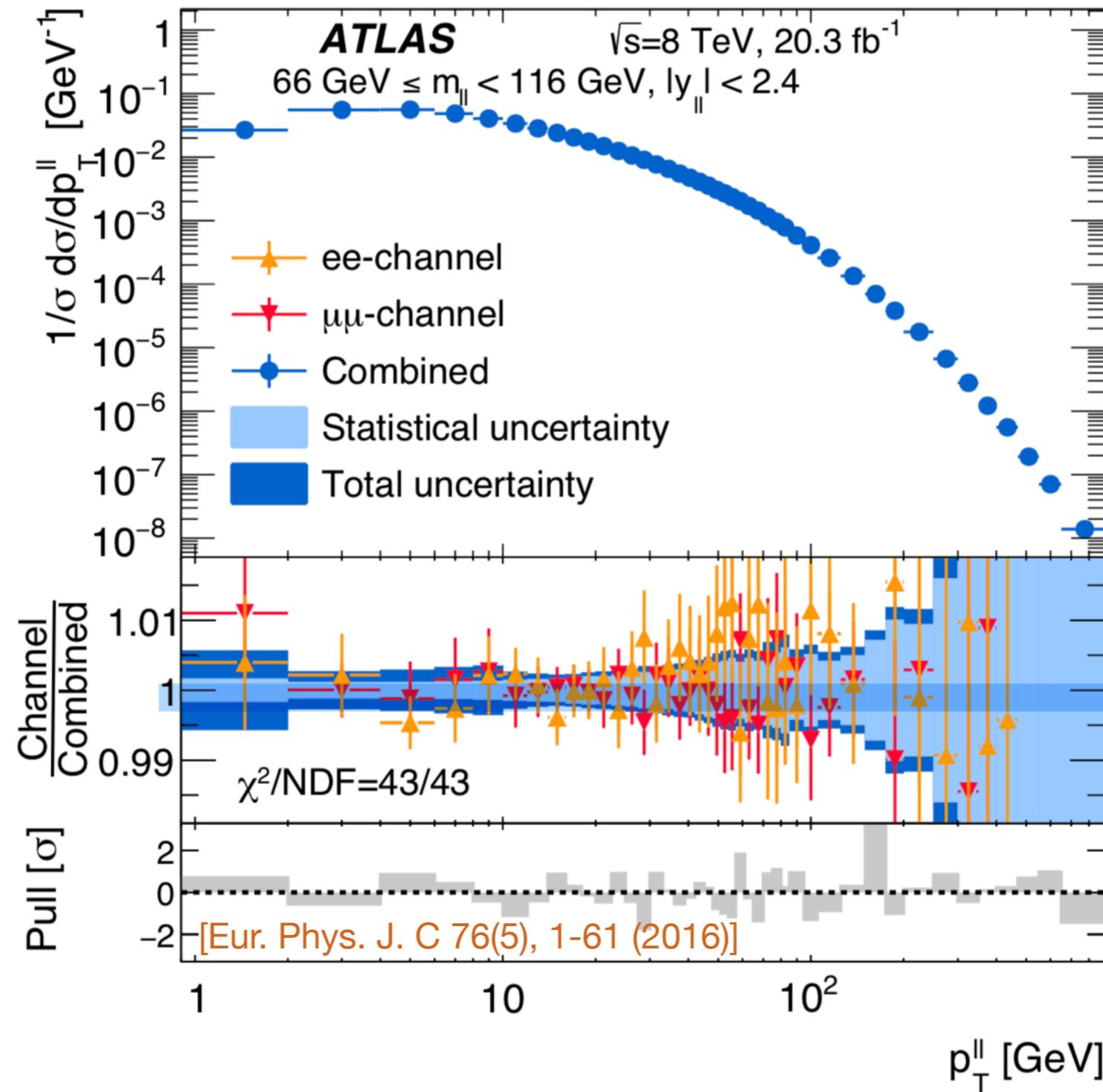
# MAPTMD 2022

*Collins-Soper kernel*



# TMDs at the LHC

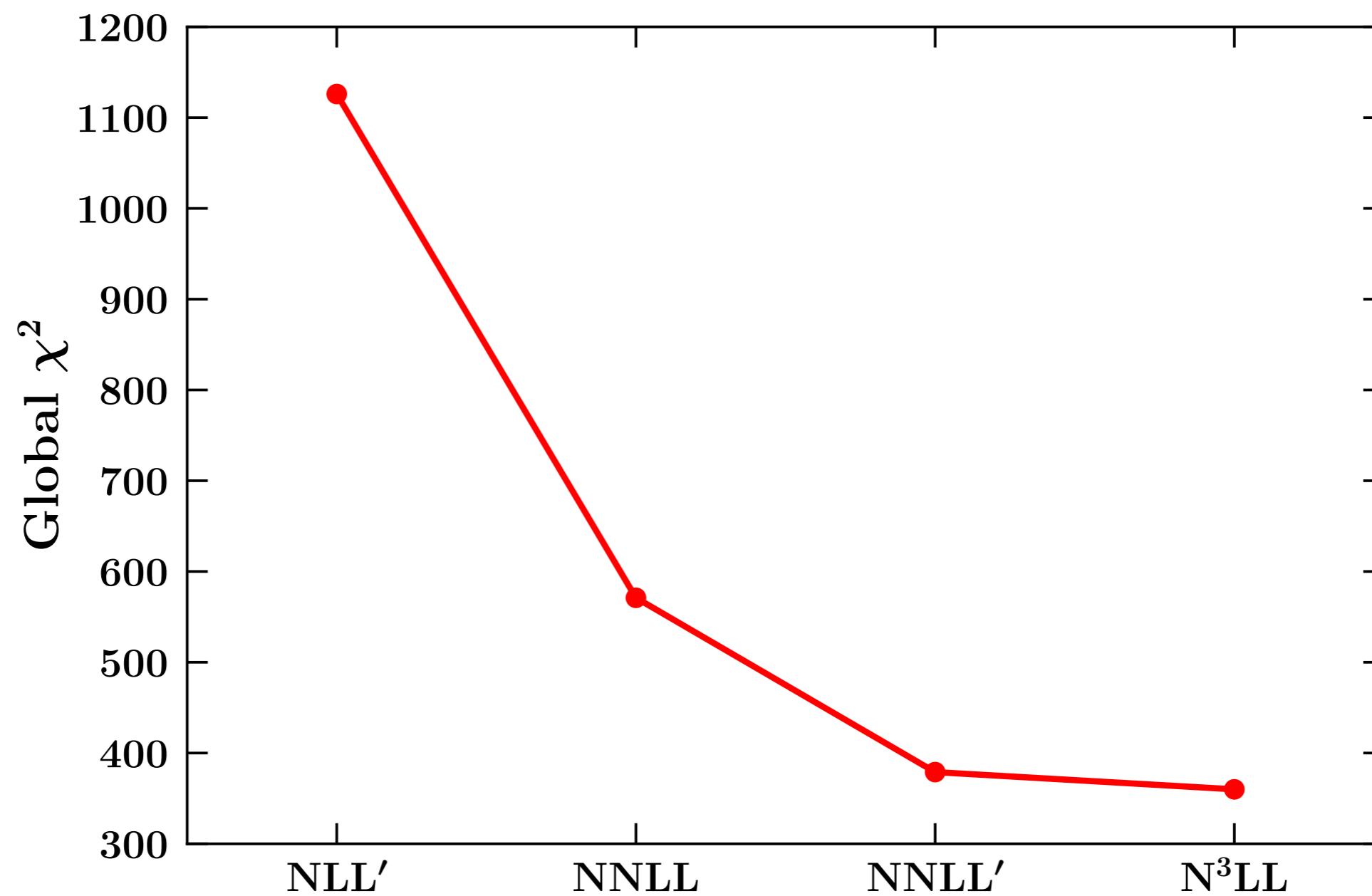
- Measurements of  $\mathcal{Z} q_T$  distributions have reached **sub-percent level** uncs.:



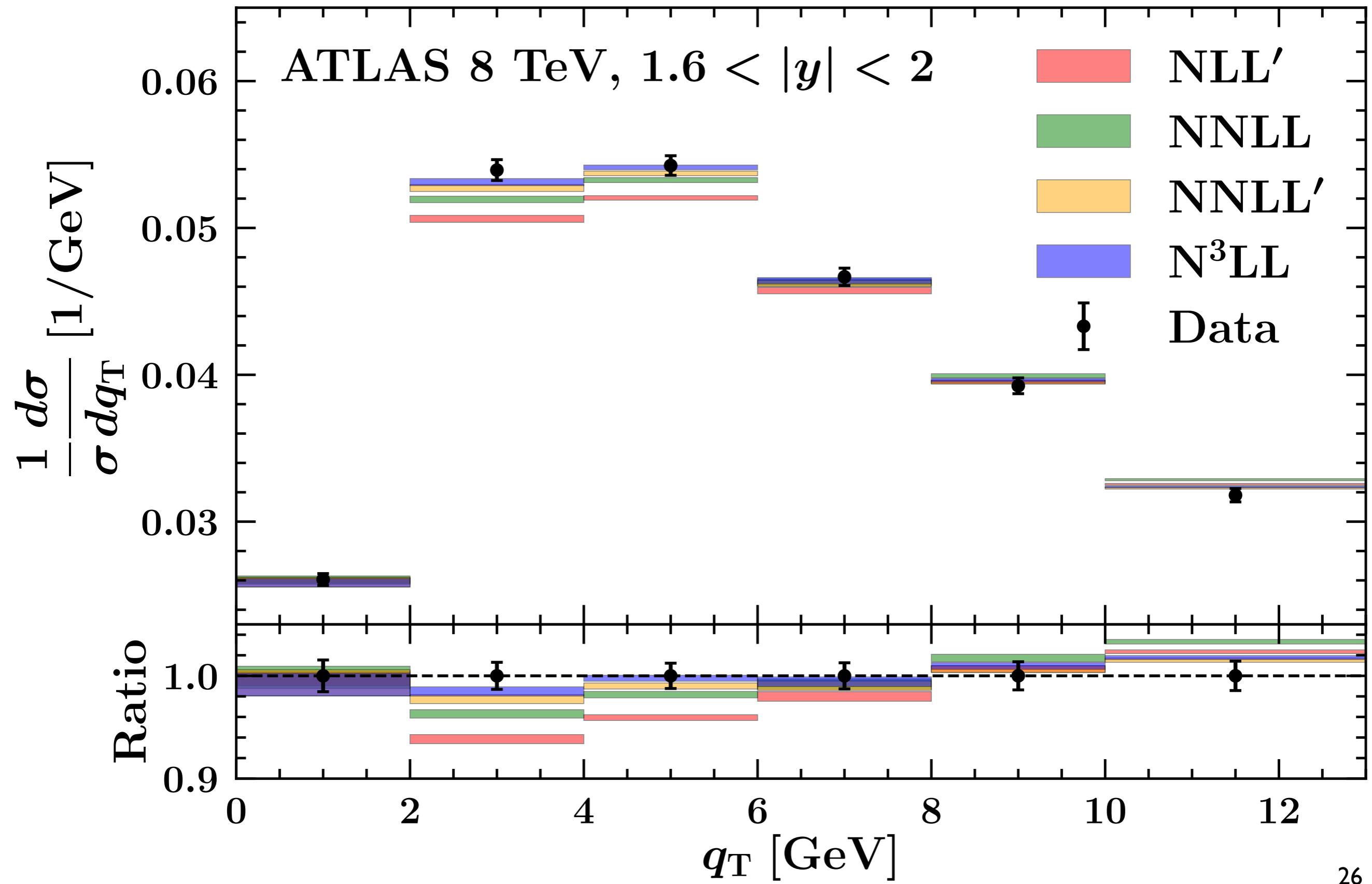
- State-of-the-art** calculations are thus necessary to describe this data.
- Non-perturbative** corrections are *very* relevant at very low  $q_T$ .

# Perturbative convergence

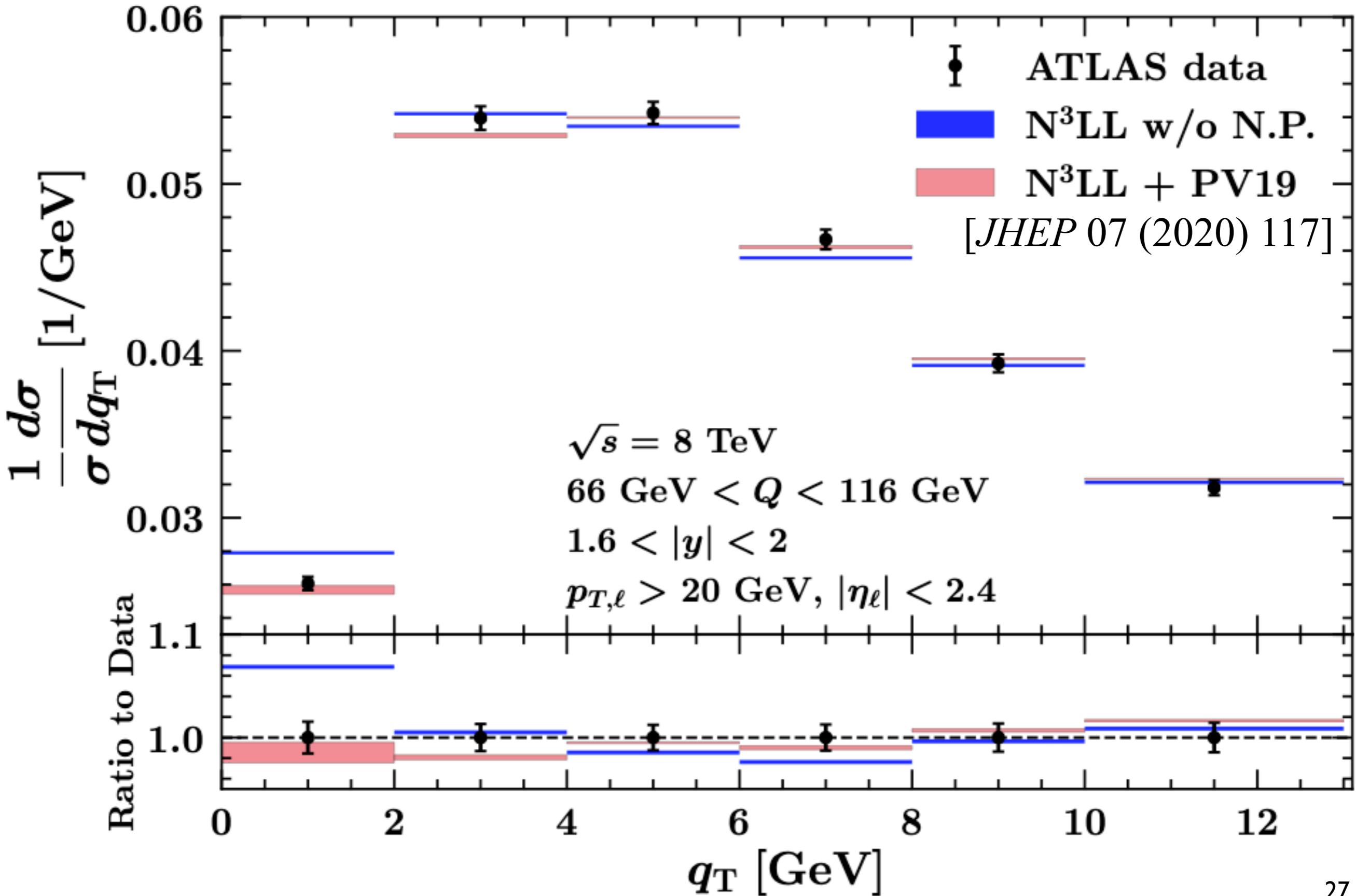
	NLL'	NNLL	NNLL'	$N^3LL$
Global $\chi^2$	1126	571	379	360



# Perturbative convergence



# TMDs at the LHC



# TMDs at the LHC

## *The W mass*

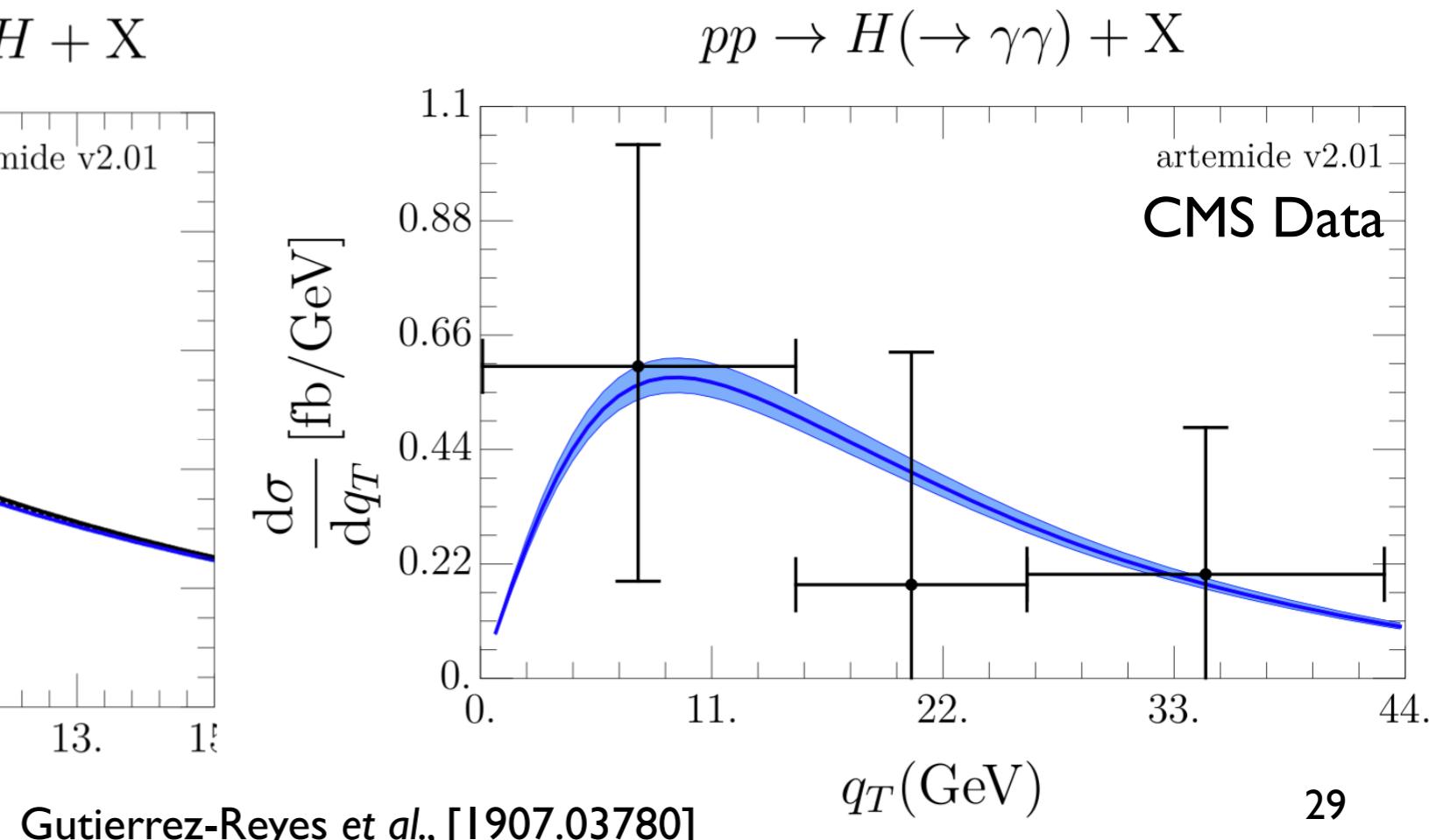
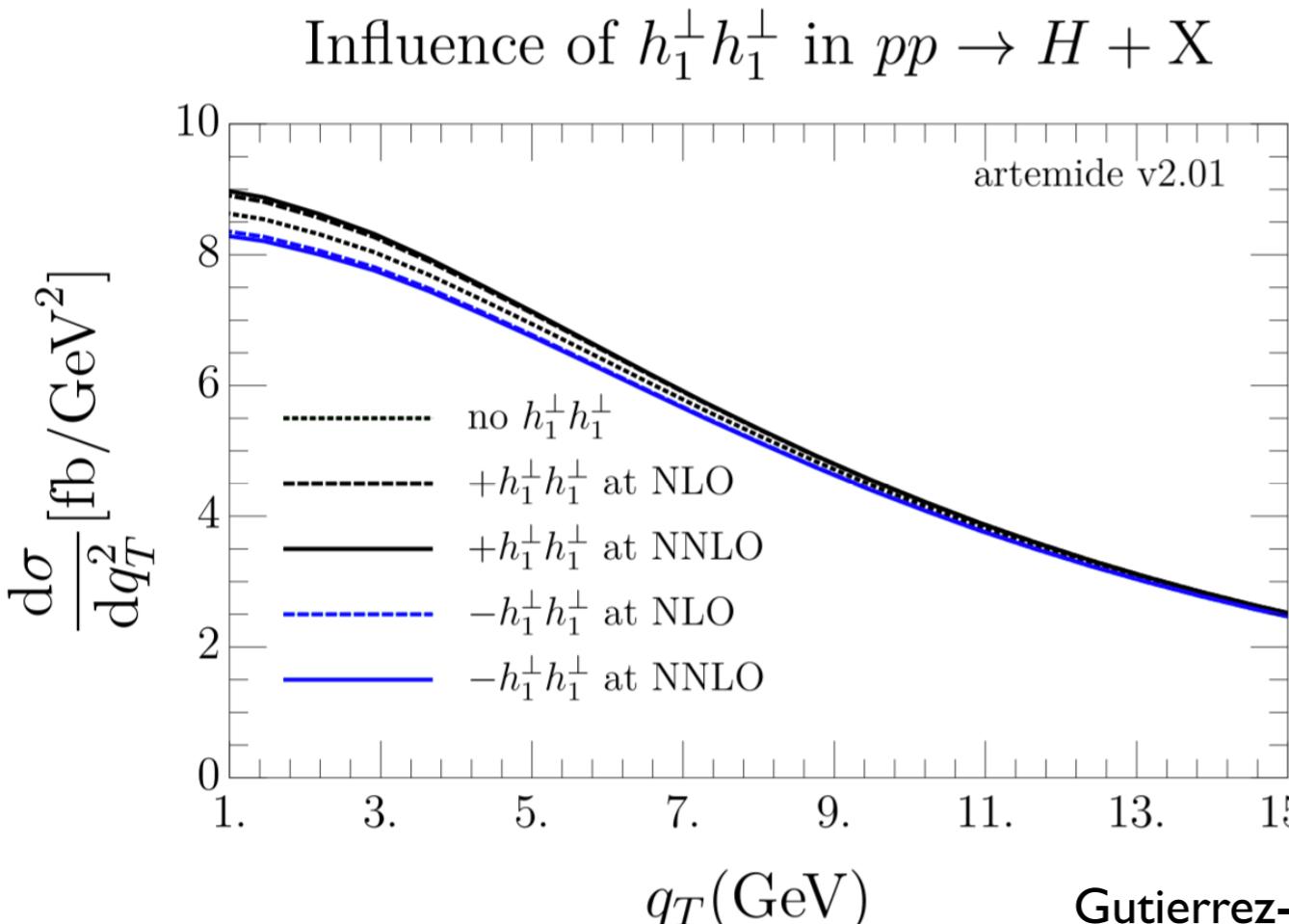
- A precise determination of the **W mass** plays an important role in testing the Standard Model and thus for **BSM** physics.
- This is a central task of the **LHC physics programme**.
- In order to minimise experimental systematic effects, the most promising procedure relies on the measurement of the  $W/Z$  ratio cross section:
  - the  $W$  mass is basically determined through template fits of:
$$\frac{d\sigma^W}{dq_T} = \left( \frac{d\sigma^W/dq_T}{d\sigma^Z/dq_T} \right)_{\text{exp.}} \left( \frac{d\sigma^Z}{dq_T} \right)_{\text{th.}}$$
- Therefore, an accurate and reliable prediction of the **Z spectrum** is essential.
- **TMD-based predictions** are currently playing an important role within the LHC electroweak working group along **other formalisms**.

# Gluon TMDs

- 🍎 The *linearly polarised* ( $h_T$ ) and *unpolarised* ( $f$ ) gluon TMD PDFs contribute to the low- $q_T$  spectrum of Higgs production in gluon fusion:

$$\frac{d\sigma}{dy d^2 \mathbf{q}_T} = \frac{\sigma_{gg \rightarrow H}}{(2\pi)^2} \int d^2 \mathbf{b} e^{-i(\mathbf{b}\mathbf{q}_T)} \left( f_{1,g}(x_A, \mathbf{b}) f_{1,g}(x_B, \mathbf{b}) + h_{1,g}^\perp(x_A, \mathbf{b}) h_{1,g}^\perp(x_B, \mathbf{b}) \right)$$

- 🍎 Despite linearly polarised gluons enter at NNLL, their effect on the cross section can be assessed but **below current LHC data accuracy**:

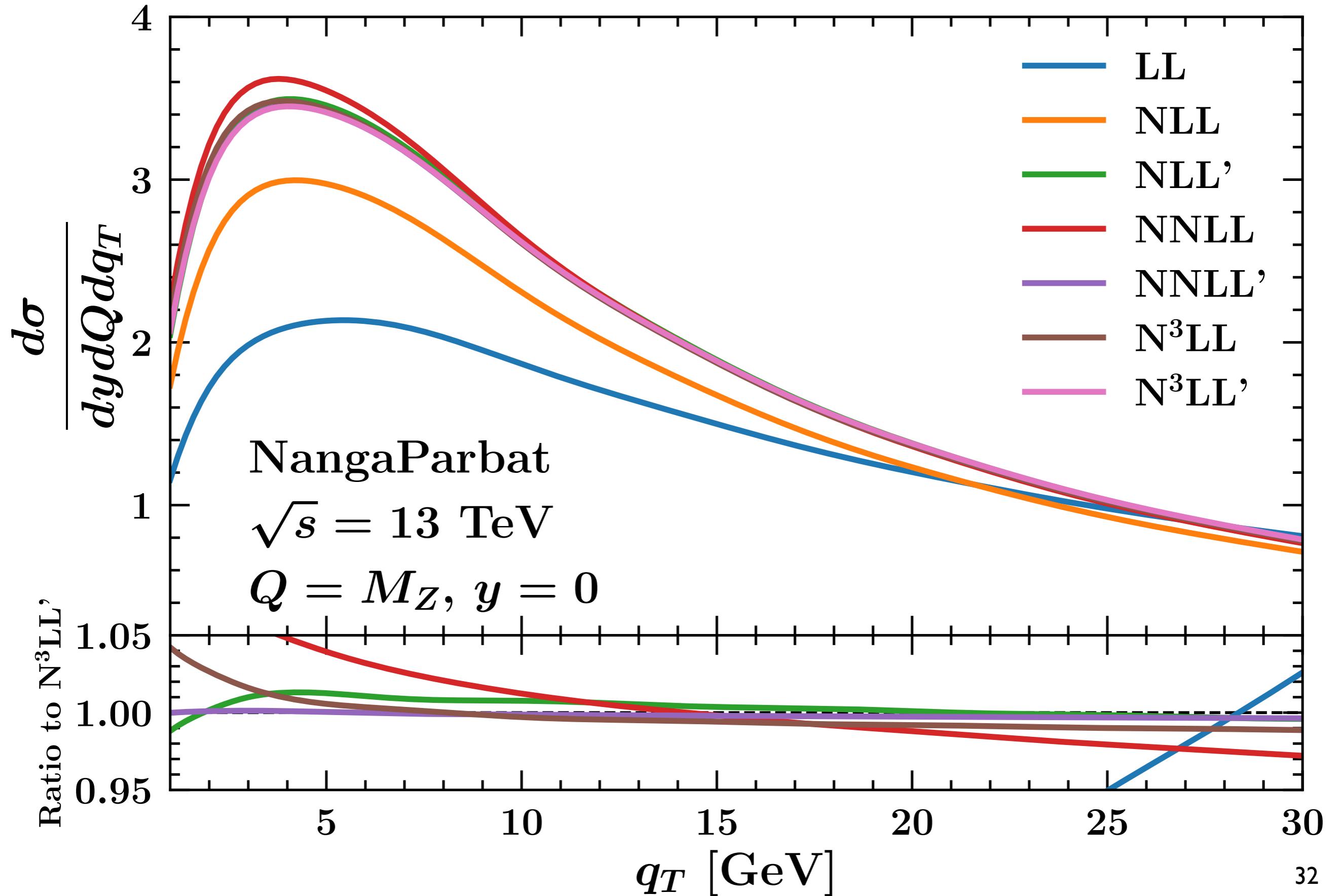


# Conclusions

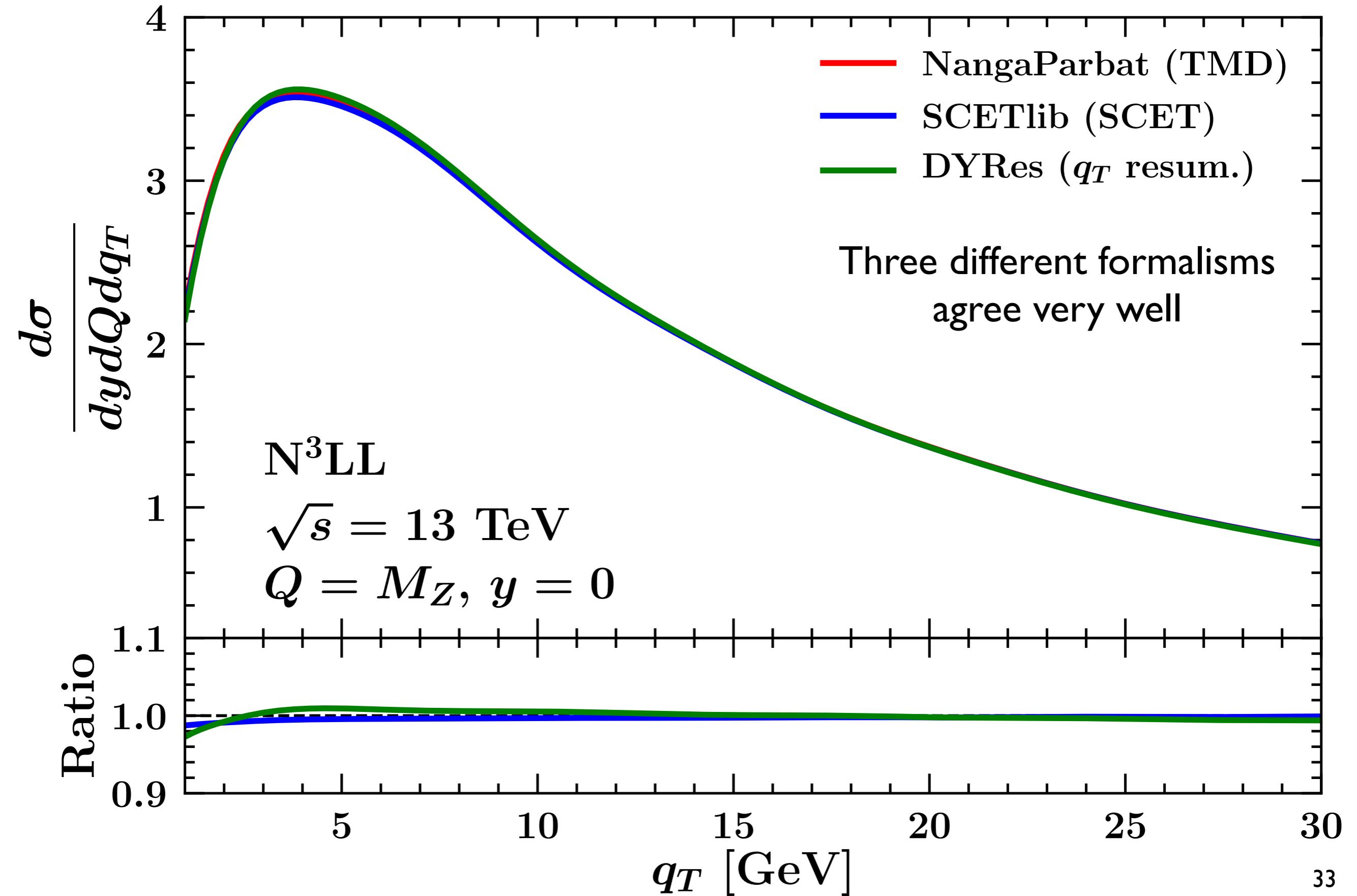
- 🍎 **TMD factorisation** provides a valuable tool to describe  $q_T$  distributions at small values of  $q_T$  (resummation of large logs),
  - 🍎 written in terms of **TMD distributions**,
- 🍎 Non-perturbative component of TMDs is to be determined from **data**.
- 🍎 A lot of effort is being invested on the extraction of TMD PDFs and FFs:
  - 🍎 tremendous progress made over the past few years,
  - 🍎 wide and precise **datasets** (COMPASS, HERMES, LHC and Tevatron exps.),
  - 🍎 more data to come from the LHC,
  - 🍎 state-of-the-art **theoretical computation** moving to even higher accuracy,
  - 🍎 looking forward to the **EIC** to pin down TMDs to unprecedented accuracy.

# Backup

# Perturbative convergence



# TMD, $q_T$ resummation, SCET



# Logarithmic counting

- apple TMD factorisation provides **resummation** of large logs  $L = \log(q_T/Q)$ :
  - apple implemented through the **Sudakov** form fact  $R$ .

- apple A **perturbative expansion** in powers of  $\alpha_s$  of  $R$  would give:

One Sudakov for each TMD  $\longrightarrow R^2 = \sum_{n=0}^{\infty} a_s^n \sum_{k=1}^{2n} \tilde{S}^{(n,k)} L^k$  Double-log expansion

- apple that can be rearranged as:

$$R^2 = \sum_{m=0}^{\infty} R_{N^m LL}^2 \quad \text{with} \quad R_{N^m LL}^2 = \sum_{n=[m/2]}^{\infty} \tilde{S}^{(n,2n-m)} a_s^n L^{2n-m}$$

$n=[m/2]$  Integer part of  $m/2$

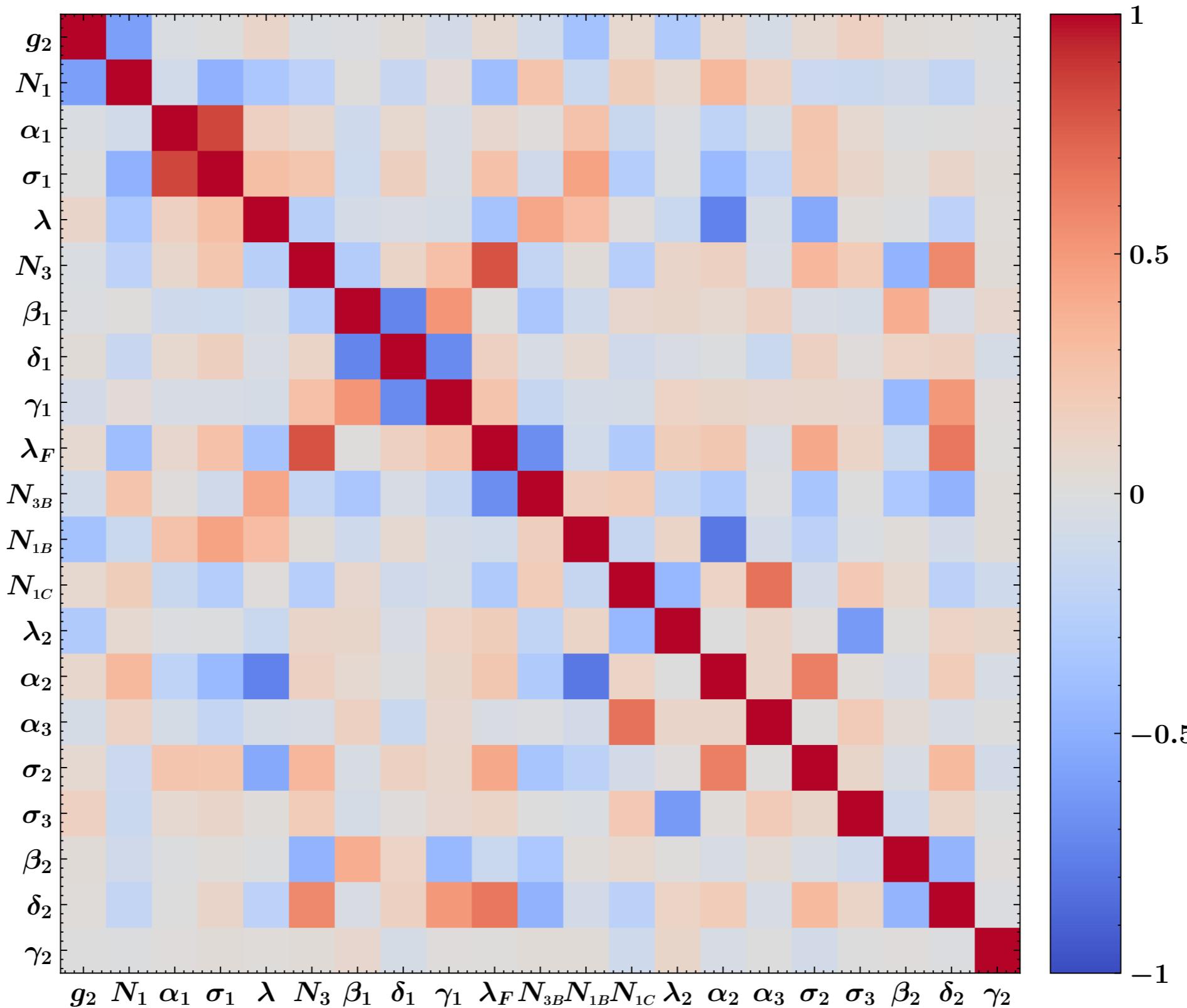
- apple Therefore, multiplying  $R$  by a power  $p$  of  $\alpha_s$  gives:

$$a_s^p R_{N^m LL}^2 = \sum_{j=[(m+2p)/2]}^{\infty} \tilde{S}^{(j-p,2j-(m+2p))} a_s^j L^{2j-(m+2p)} \sim R_{N^{m+2p} LL}^2$$

- apple Bottom line: any additional power of  $\alpha_s$  causes a shift of **two units** in the logarithmic ordering.

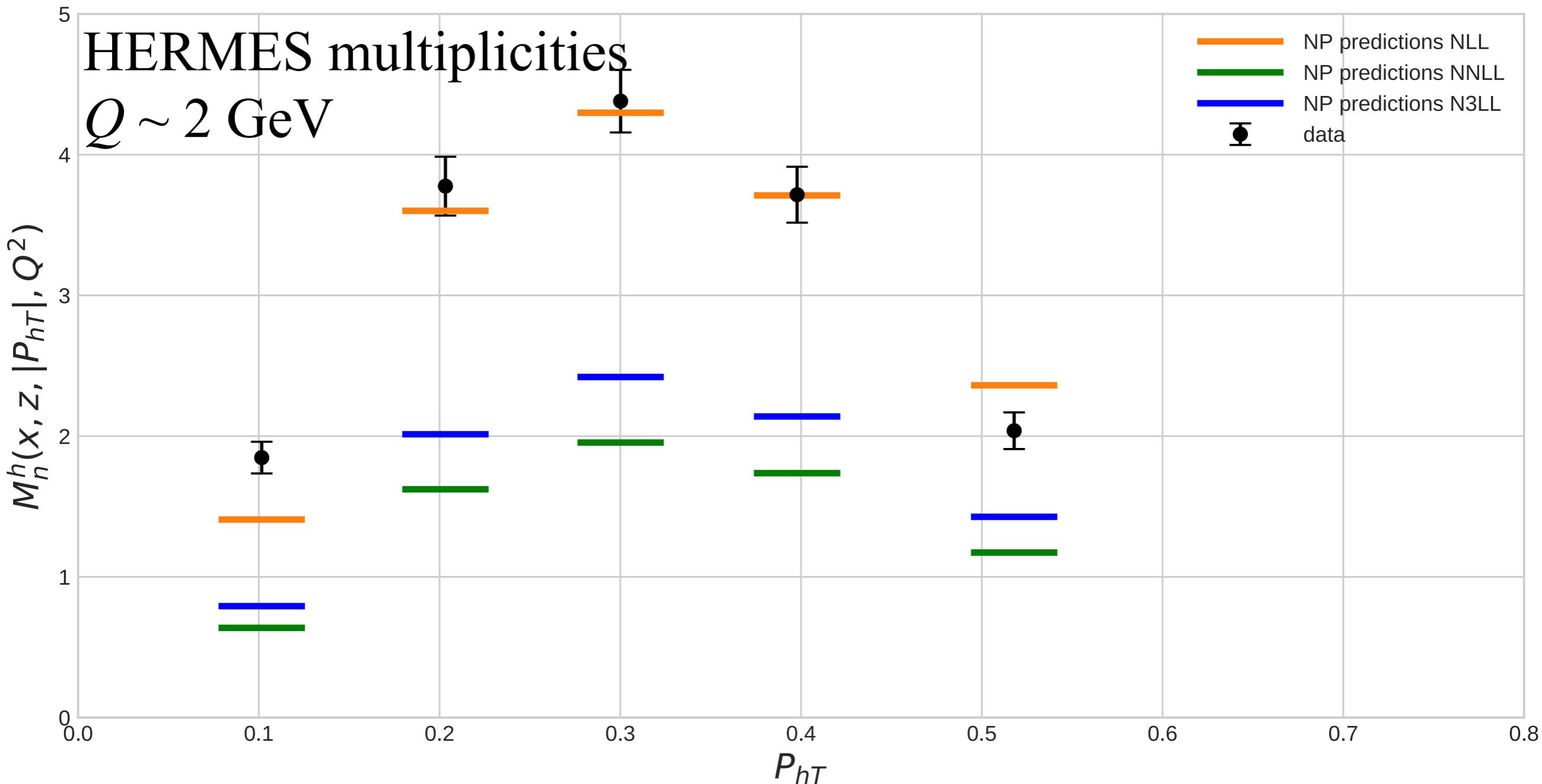
# MAPTMD 2022

## *Correlation between fit parameters*



# MAPTMD 2022

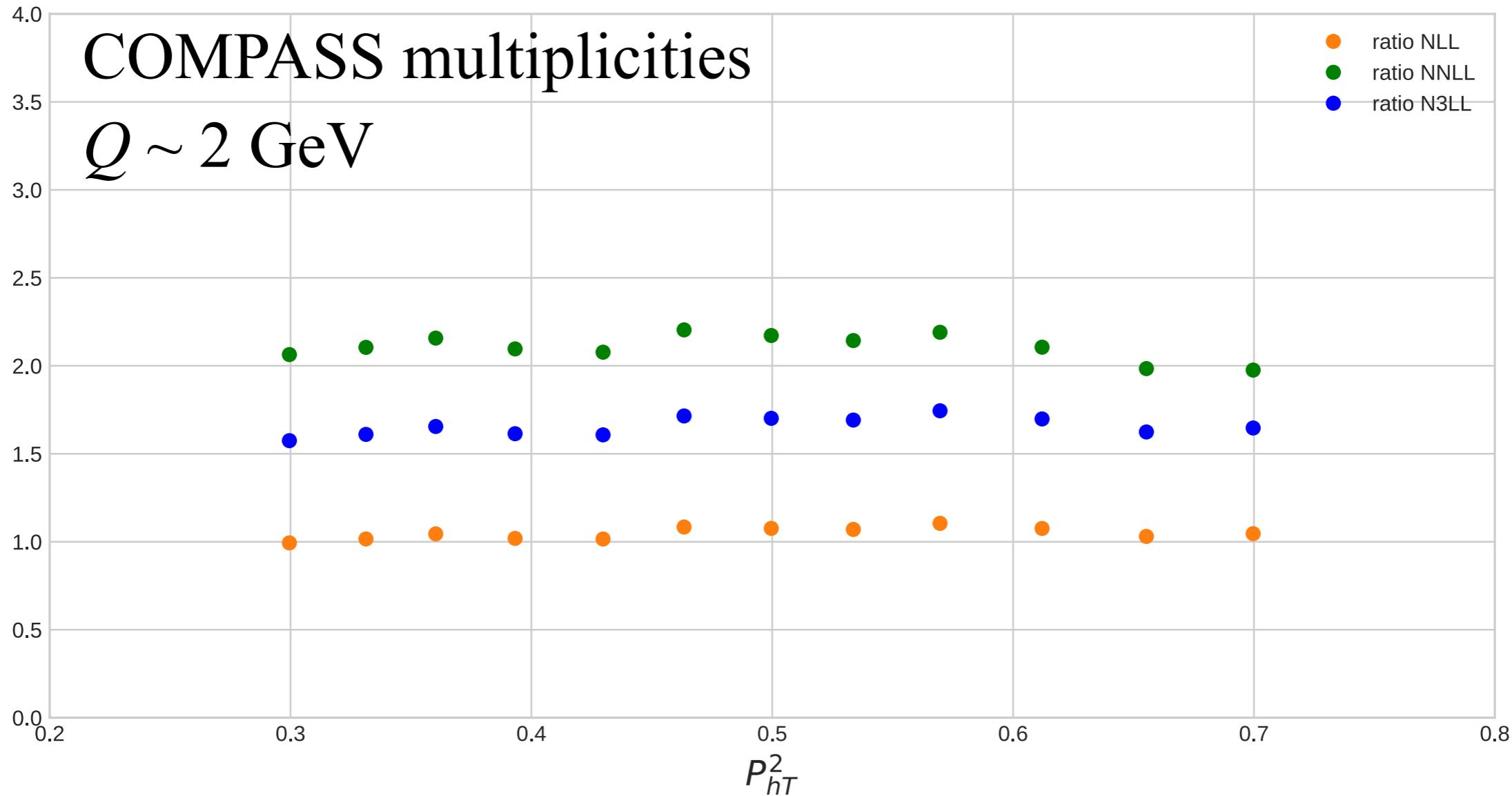
## *Normalisation of SIDIS*



- 🍎 Description of SIDIS multiplicities considerably worsens moving from NLL to higher perturbative orders.

# MAPTMD 2022

## *Normalisation of SIDIS*



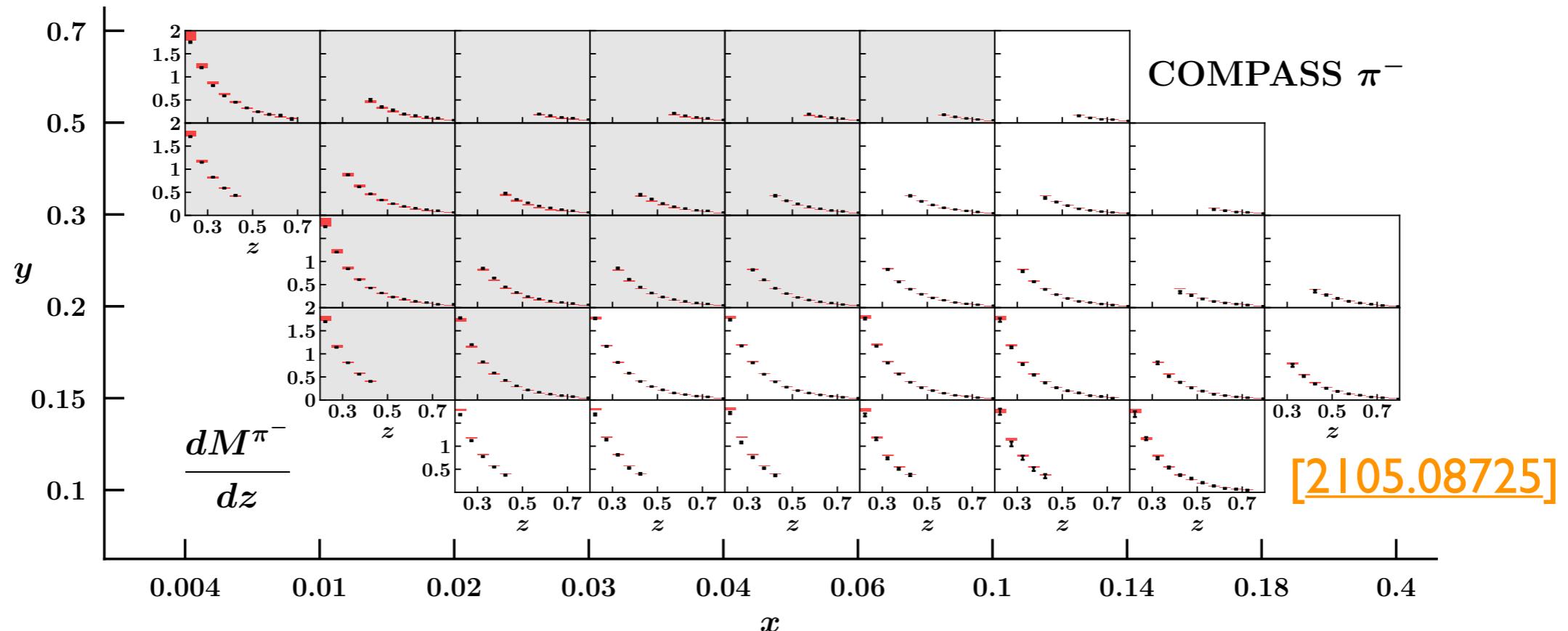
- 🍏 Normalisation problem already observed in the literature.
- 🍏 Large perturbative corrections particularly to the hard function:

$$H_{\text{SIDIS}}(Q) = 1 + \frac{\alpha_s(Q)}{4\pi} \underbrace{C_F \left( -16 + \frac{\pi^2}{3} \right)}_{\sim -17} + \mathcal{O}(\alpha_s^2)$$

# MAPTMD 2022

## *Normalisation of SIDIS*

- 🍎 SIDIS multiplicity:  $M(x, z, P_{hT}, Q) = \frac{d\sigma}{dx dQ dz dP_{hT}} / \frac{d\sigma}{dx dQ}$
- 🍎 The SIDIS cross section integrated over  $P_{hT}$  ( $d\sigma/dxdQdz$ ) is ok.



- 🍎 Normalise predictions such that integral over  $P_{hT}$  gives  $d\sigma/dxdQdz$ :

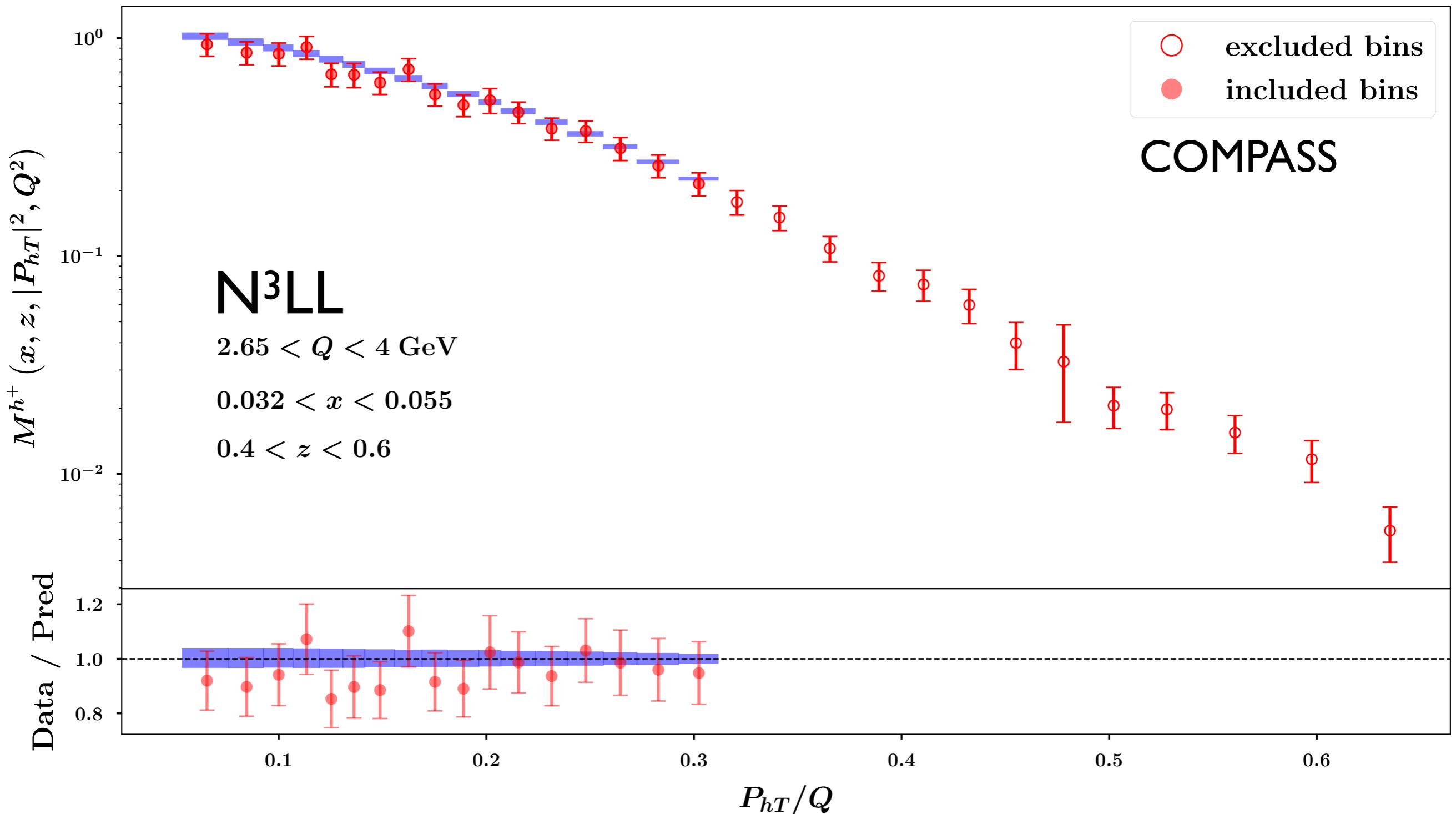
$$M(x, z, P_{hT}, Q) = \mathcal{N} \frac{\frac{d\sigma}{dx dQ dz dP_{hT}}}{\frac{d\sigma}{dx dQ}}$$

$$\mathcal{N} = \frac{\frac{d\sigma}{dx dQ dz}}{\int dP_{hT} \frac{d\sigma}{dx dQ dz dP_{hT}}}$$

- 🍎 Theoretically determined normalisation, **not fitted**.

# MAPTMD 2022

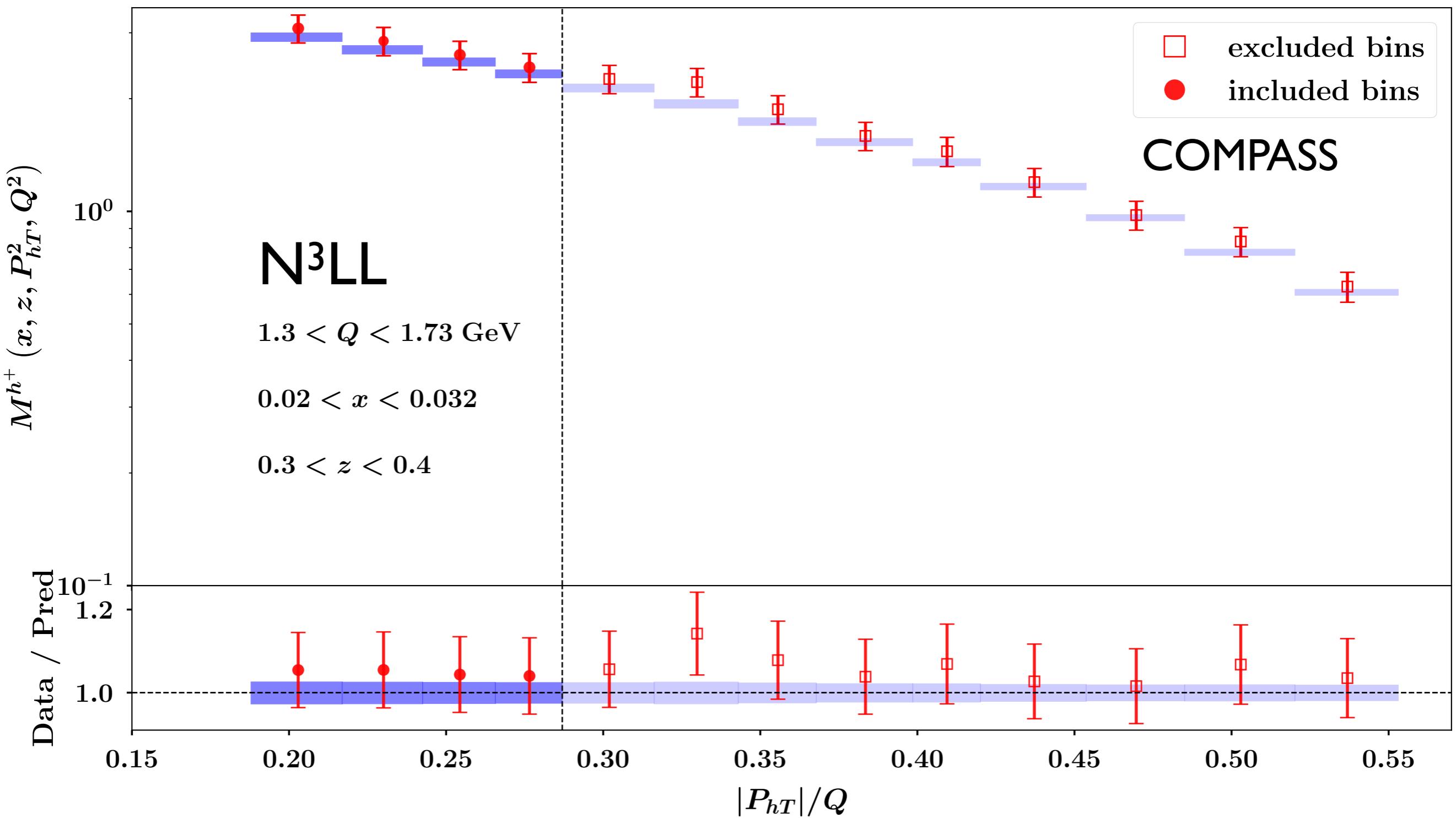
## *Normalisation of SIDIS*



Excellent agreement upon normalisation.

# MAPTMD 2022

## *Normalisation of SIDIS*



Agreement extends well above the expected validity region. To be clarified.

# Pavia2019

## Dataset

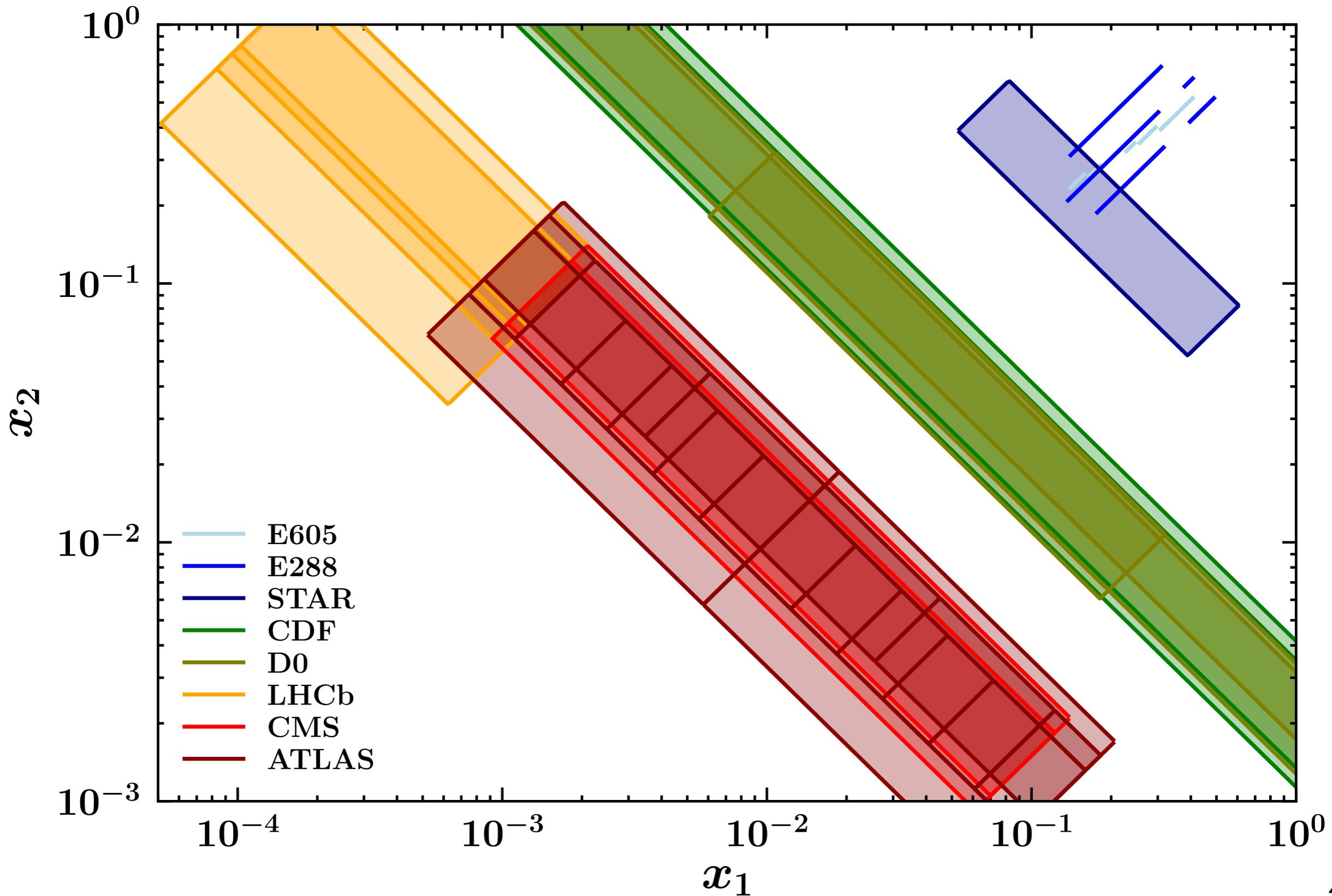
- 🍎 DY data only:
- 🍎 fixed-target low-energy DY,
- 🍎 STAR data
- 🍎 LHC and Tevatron data,
- 🍎 353 data points,
- 🍎 selection cut  $q_T / Q < 0.2$ .

Experiment	$N_{\text{dat}}$	Observable	$\sqrt{s}$ [GeV]	$Q$ [GeV]	$y$ or $x_F$	Lepton cuts	Ref.
E605	50	$Ed^3\sigma/d^3q$	38.8	7 - 18	$x_F = 0.1$	-	[79]
E288 200 GeV	30	$Ed^3\sigma/d^3q$	19.4	4 - 9	$y = 0.40$	-	[80]
E288 300 GeV	39	$Ed^3\sigma/d^3q$	23.8	4 - 12	$y = 0.21$	-	[80]
E288 400 GeV	61	$Ed^3\sigma/d^3q$	27.4	5 - 14	$y = 0.03$	-	[80]
STAR 510	7	$d\sigma/dq_T$	510	73 - 114	$ y  < 1$	$p_{T\ell} > 25 \text{ GeV}$ $ \eta_\ell  < 1$	-
CDF Run I	25	$d\sigma/dq_T$	1800	66 - 116	Inclusive	-	[81]
CDF Run II	26	$d\sigma/dq_T$	1960	66 - 116	Inclusive	-	[82]
D0 Run I	12	$d\sigma/dq_T$	1800	75 - 105	Inclusive	-	[83]
D0 Run II	5	$(1/\sigma)d\sigma/dq_T$	1960	70 - 110	Inclusive	-	[84]
D0 Run II ( $\mu$ )	3	$(1/\sigma)d\sigma/dq_T$	1960	65 - 115	$ y  < 1.7$	$p_{T\ell} > 15 \text{ GeV}$ $ \eta_\ell  < 1.7$	[85]
LHCb 7 TeV	7	$d\sigma/dq_T$	7000	60 - 120	$2 < y < 4.5$	$p_{T\ell} > 20 \text{ GeV}$ $2 < \eta_\ell < 4.5$	[86]
LHCb 8 TeV	7	$d\sigma/dq_T$	8000	60 - 120	$2 < y < 4.5$	$p_{T\ell} > 20 \text{ GeV}$ $2 < \eta_\ell < 4.5$	[87]
LHCb 13 TeV	7	$d\sigma/dq_T$	13000	60 - 120	$2 < y < 4.5$	$p_{T\ell} > 20 \text{ GeV}$ $2 < \eta_\ell < 4.5$	[92]
CMS 7 TeV	4	$(1/\sigma)d\sigma/dq_T$	7000	60 - 120	$ y  < 2.1$	$p_{T\ell} > 20 \text{ GeV}$ $ \eta_\ell  < 2.1$	[88]
CMS 8 TeV	4	$(1/\sigma)d\sigma/dq_T$	8000	60 - 120	$ y  < 2.1$	$p_{T\ell} > 15 \text{ GeV}$ $ \eta_\ell  < 2.1$	[89]
ATLAS 7 TeV	6 6 6	$(1/\sigma)d\sigma/dq_T$	7000	66 - 116	$ y  < 1$ $1 <  y  < 2$ $2 <  y  < 2.4$	$p_{T\ell} > 20 \text{ GeV}$ $ \eta_\ell  < 2.4$	[93]
ATLAS 8 TeV on-peak	6 6 6 6	$(1/\sigma)d\sigma/dq_T$	8000	66 - 116	$ y  < 0.4$ $0.4 <  y  < 0.8$ $0.8 <  y  < 1.2$ $1.2 <  y  < 1.6$ $1.6 <  y  < 2$ $2 <  y  < 2.4$	$p_{T\ell} > 20 \text{ GeV}$ $ \eta_\ell  < 2.4$	[90]
ATLAS 8 TeV off-peak	4 8	$(1/\sigma)d\sigma/dq_T$	8000	46 - 66 116 - 150	$ y  < 2.4$	$p_{T\ell} > 20 \text{ GeV}$ $ \eta_\ell  < 2.4$	[90]
Total	353	-	-	-	-	-	-

# Pavia2019

*Kinematic coverage*

$$x_{1,2} = \frac{Q}{\sqrt{s}} e^{\pm y}$$

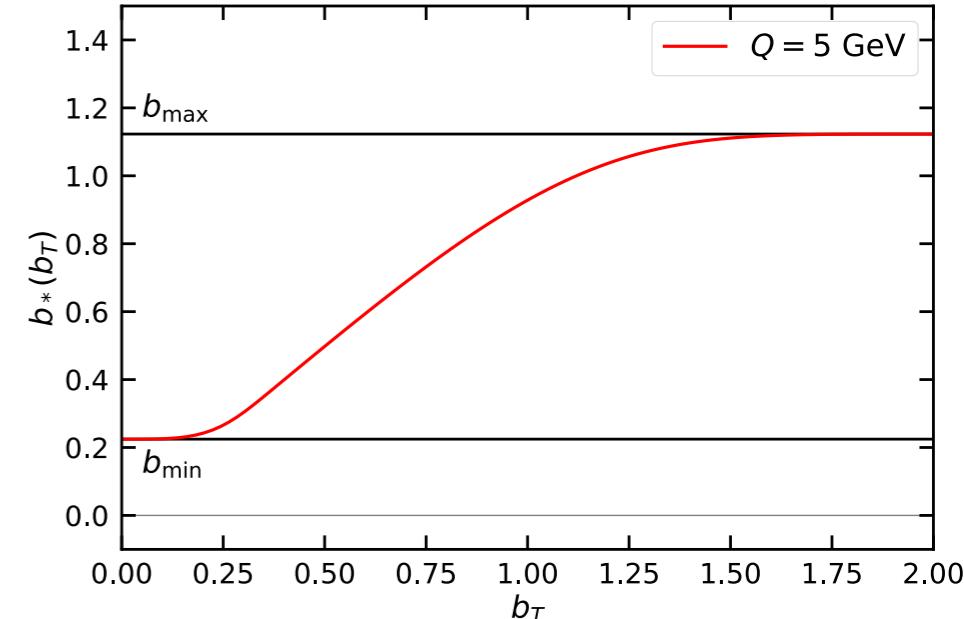


# Pavia2019

## Main settings

- 🍎  $b_*$  prescription:

$$b_*(b_T) = b_{\max} \left( \frac{1 - e^{-b_T^4/b_{\max}^4}}{1 - e^{-b_T^4/b_{\min}^4}} \right)^{1/4} \quad \text{with} \quad \begin{cases} b_{\max} = 2e^{-\gamma_E} \\ b_{\min} = b_{\max}/Q \end{cases}$$



- 🍎 Non-perturbative function  $f_{\text{NP}}$ :

- 🍎 evolution:

$$g_K(b_T) = - (g_2 + g_{2B} b_T^2) \frac{b_T^2}{2}$$

- 🍎 PDFs:

$$\tilde{f}_{\text{NP}}(x, b_T) = \left[ \frac{1 - \lambda}{1 + g_1(x) \frac{b_T^2}{4}} + \lambda \exp \left( -g_{1B}(x) \frac{b_T^2}{4} \right) \right]$$

$$g_1(x) = \frac{N_1}{x\sigma} \exp \left[ -\frac{1}{2\sigma^2} \ln^2 \left( \frac{x}{\alpha} \right) \right] \quad g_{1B}(x) = \frac{N_{1B}}{x\sigma_B} \exp \left[ -\frac{1}{2\sigma_B^2} \ln^2 \left( \frac{x}{\alpha_B} \right) \right]$$

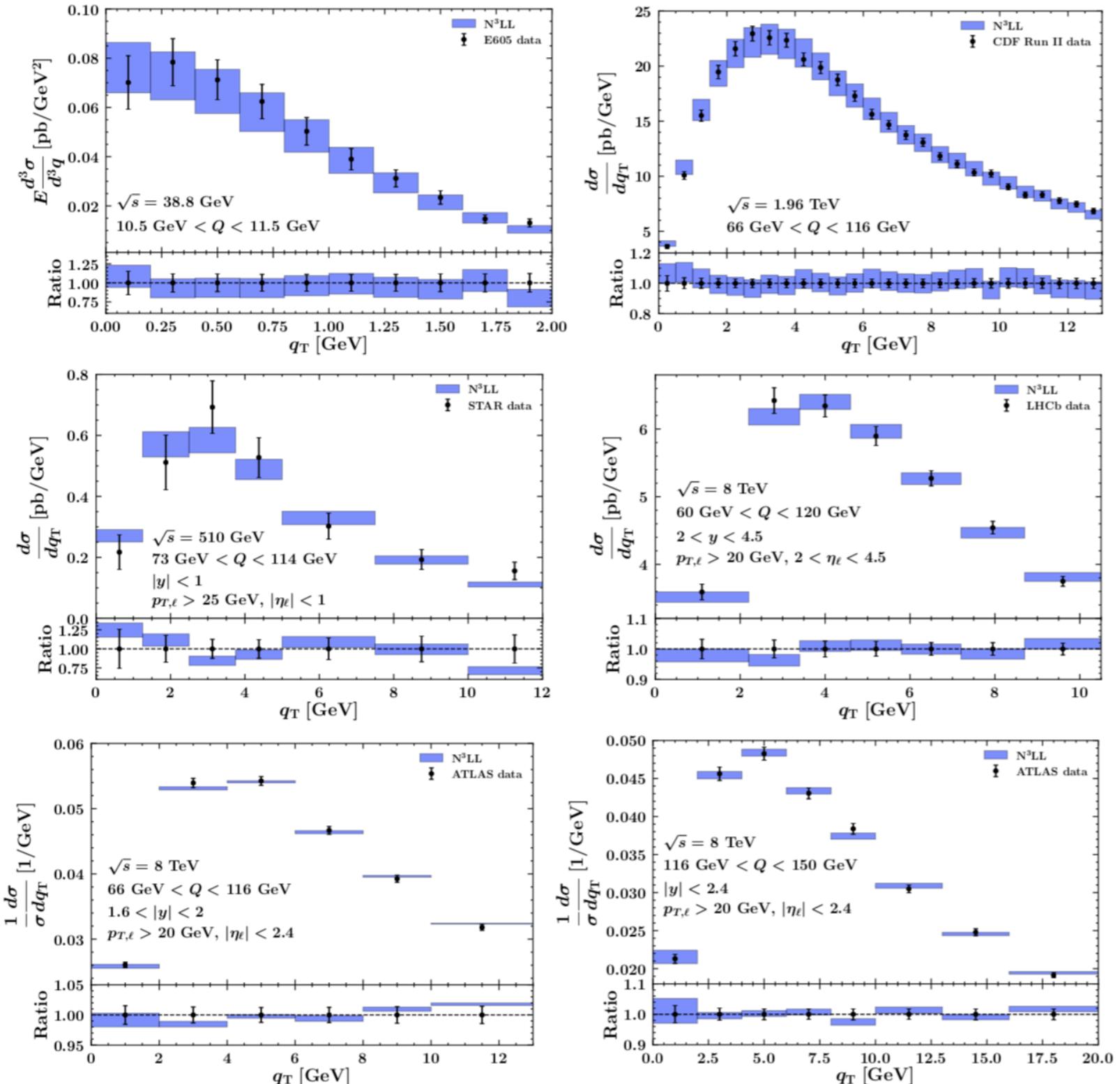
- 🍎 **9 free parameters** to fit to data.

- 🍎 Perturbative accuracies: **NLL'**, **NNLL**, **NNLL'**, **N<sup>3</sup>LL**

- 🍎 **Monte Carlo** method for the experimental error propagation.

# Pavia2019

## Fit quality



Experiment		$\chi_D^2/N_{\text{dat}}$	$\chi_\lambda^2/N_{\text{dat}}$	$\chi^2/N_{\text{dat}}$
E605	$7 \text{ GeV} < Q < 8 \text{ GeV}$	0.419	0.068	0.487
	$8 \text{ GeV} < Q < 9 \text{ GeV}$	0.995	0.034	1.029
	$10.5 \text{ GeV} < Q < 11.5 \text{ GeV}$	0.191	0.137	0.328
	$11.5 \text{ GeV} < Q < 13.5 \text{ GeV}$	0.491	0.284	0.775
	$13.5 \text{ GeV} < Q < 18 \text{ GeV}$	0.491	0.385	0.877
E288 200 GeV	$4 \text{ GeV} < Q < 5 \text{ GeV}$	0.213	0.649	0.862
	$5 \text{ GeV} < Q < 6 \text{ GeV}$	0.673	0.292	0.965
	$6 \text{ GeV} < Q < 7 \text{ GeV}$	0.133	0.141	0.275
	$7 \text{ GeV} < Q < 8 \text{ GeV}$	0.254	0.014	0.268
E288 300 GeV	$8 \text{ GeV} < Q < 9 \text{ GeV}$	0.652	0.024	0.676
	$4 \text{ GeV} < Q < 5 \text{ GeV}$	0.231	0.555	0.785
	$5 \text{ GeV} < Q < 6 \text{ GeV}$	0.502	0.204	0.706
	$6 \text{ GeV} < Q < 7 \text{ GeV}$	0.315	0.063	0.378
	$7 \text{ GeV} < Q < 8 \text{ GeV}$	0.056	0.030	0.086
E288 400 GeV	$8 \text{ GeV} < Q < 9 \text{ GeV}$	0.530	0.017	0.547
	$11 \text{ GeV} < Q < 12 \text{ GeV}$	1.047	0.167	1.215
	$5 \text{ GeV} < Q < 6 \text{ GeV}$	0.312	0.065	0.377
	$6 \text{ GeV} < Q < 7 \text{ GeV}$	0.100	0.005	0.105
	$7 \text{ GeV} < Q < 8 \text{ GeV}$	0.018	0.011	0.029
STAR	$8 \text{ GeV} < Q < 9 \text{ GeV}$	0.437	0.039	0.477
	$11 \text{ GeV} < Q < 12 \text{ GeV}$	0.637	0.036	0.673
	$12 \text{ GeV} < Q < 13 \text{ GeV}$	0.788	0.028	0.816
	$13 \text{ GeV} < Q < 14 \text{ GeV}$	1.064	0.044	1.107
		0.782	0.054	0.836
CDF Run I		0.480	0.058	0.538
CDF Run II		0.959	0.001	0.959
D0 Run I		0.711	0.043	0.753
D0 Run II		1.325	0.612	1.937
D0 Run II ( $\mu$ )		3.196	0.023	3.218
LHCb 7 TeV		1.069	0.194	1.263
LHCb 8 TeV		0.460	0.075	0.535
LHCb 13 TeV		0.735	0.020	0.755
CMS 7 TeV		2.131	0.000	2.131
CMS 8 TeV		1.405	0.007	1.412
ATLAS 7 TeV	$0 <  y  < 1$	2.581	0.028	2.609
	$1 <  y  < 2$	4.333	1.032	5.365
	$2 <  y  < 2.4$	3.561	0.378	3.939
ATLAS 8 TeV	$0 <  y  < 0.4$	1.924	0.337	2.262
	$0.4 <  y  < 0.8$	2.342	0.247	2.590
	$0.8 <  y  < 1.2$	0.917	0.061	0.978
	$1.2 <  y  < 1.6$	0.912	0.095	1.006
	$1.6 <  y  < 2$	0.721	0.092	0.814
	$2 <  y  < 2.4$	0.932	0.348	1.280
ATLAS 8 TeV	$46 \text{ GeV} < Q < 66 \text{ GeV}$	2.138	0.745	2.883
off-peak	$116 \text{ GeV} < Q < 150 \text{ GeV}$	0.501	0.003	0.504
Global		<b>0.88</b>	<b>0.14</b>	<b>1.02</b>

# SV2019

## Dataset

- 🍎 Both DY and SIDIS data:
- 🍎 fixed-target low-energy DY,
- 🍎 PHENIX data,
- 🍎 LHC and Tevatron data,
- 🍎 HERMES and COMPASS,
- 🍎  $457 + 582 = 1039$  data points.

**SIDIS**     $\langle Q \rangle \geq 2\text{GeV}$      $\delta \equiv \frac{\langle q_T \rangle}{\langle Q \rangle} < 0.25$

Experiment	Reaction	ref.	Kinematics	$N_{\text{pt}}$ after cuts
HERMES	$p \rightarrow \pi^+$	[58]	$0.023 < x < 0.6$ (6 bins) $0.2 < z < 0.8$ (6 bins) $1.0 < Q < \sqrt{20}\text{GeV}$ $W^2 > 10\text{GeV}^2$ $0.1 < y < 0.85$	24
	$p \rightarrow \pi^-$			24
	$p \rightarrow K^+$			24
	$p \rightarrow K^-$			24
	$D \rightarrow \pi^+$			24
	$D \rightarrow \pi^-$			24
	$D \rightarrow K^+$			24
	$D \rightarrow K^-$			24
COMPASS	$d \rightarrow h^+$	[59]	$0.003 < x < 0.4$ (8 bins) $0.2 < z < 0.8$ (4 bins) $1.0 < Q \simeq 9\text{GeV}$ (5 bins)	195
	$d \rightarrow h^-$			195
	Total			582

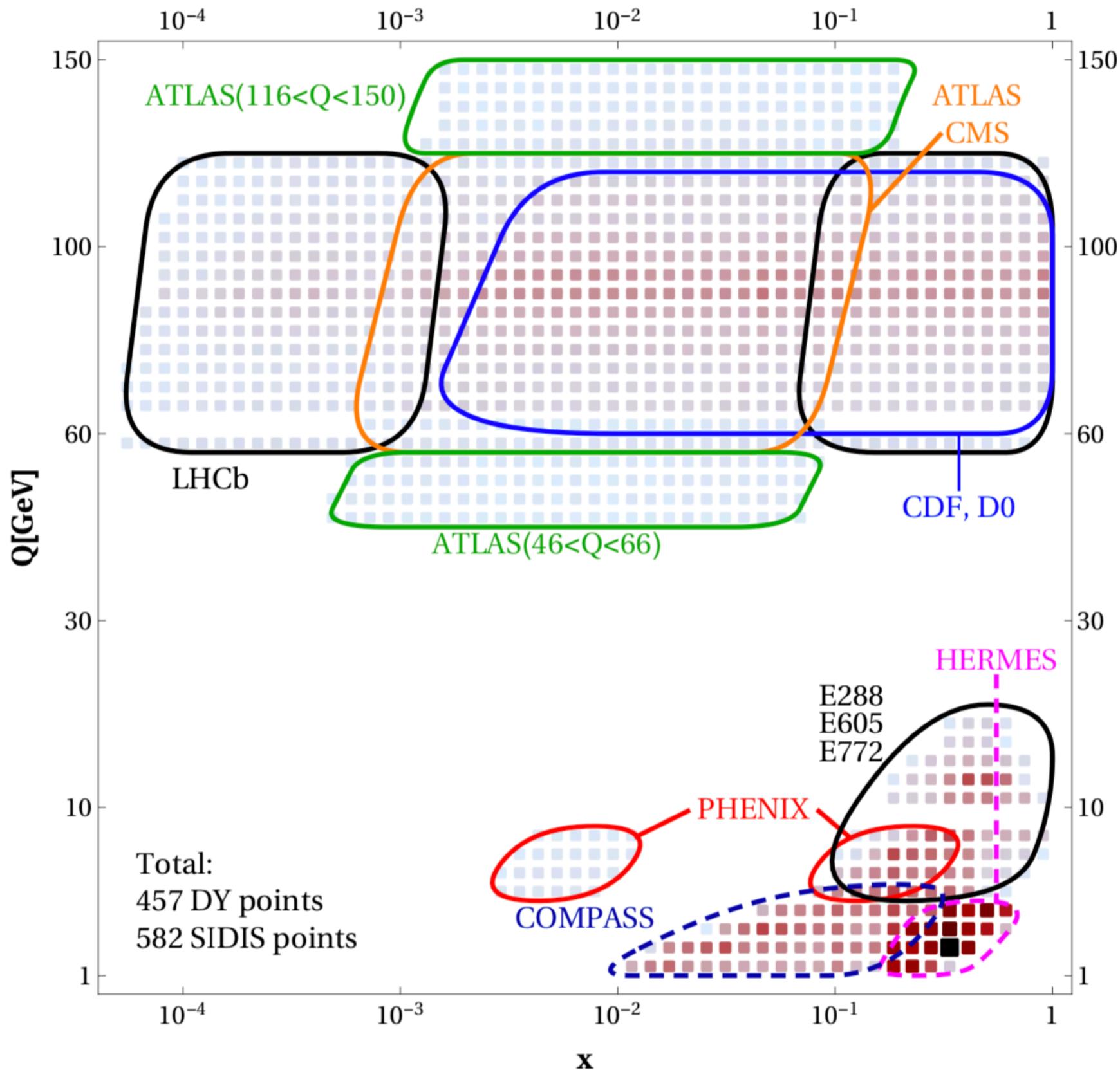
$$\text{DY} \quad \delta \equiv \frac{\langle q_T \rangle}{\langle Q \rangle} < 0.1 \quad \delta < 0.25 \quad \text{if} \quad \delta^2 < \sigma$$

Experiment	ref.	$\sqrt{s}$ [GeV]	$Q$ [GeV]	$y/x_F$	fiducial region	$N_{\text{pt}}$ after cuts
E288 (200)	[64]	19.4	4 - 9 in 1 GeV bins*	$0.1 < x_F < 0.7$	-	43
E288 (300)	[64]	23.8	4 - 12 in 1 GeV bins*	$-0.09 < x_F < 0.51$	-	53
E288 (400)	[64]	27.4	5 - 14 in 1 GeV bins*	$-0.27 < x_F < 0.33$	-	76
E605	[65]	38.8	7 - 18 in 5 bins*	$-0.1 < x_F < 0.2$	-	53
E772	[66]	38.8	5 - 15 in 8 bins*	$0.1 < x_F < 0.3$	-	35
PHENIX	[67]	200	4.8 - 8.2	$1.2 < y < 2.2$	-	3
CDF (run1)	[68]	1800	66 - 116	-	-	33
CDF (run2)	[69]	1960	66 - 116	-	-	39
D0 (run1)	[70]	1800	75 - 105	-	-	16
D0 (run2)	[71]	1960	70 - 110	-	-	8
D0 (run2) $_\mu$	[72]	1960	65 - 115	$ y  < 1.7$	$p_T > 15\text{ GeV}$ $ \eta  < 1.7$	3
ATLAS (7TeV)	[45]	7000	66 - 116	$ y  < 1$ $1 <  y  < 2$ $2 <  y  < 2.4$	$p_T > 20\text{ GeV}$ $ \eta  < 2.4$	15
ATLAS (8TeV)	[46]	8000	66 - 116	$ y  < 2.4$ in 6 bins	$p_T > 20\text{ GeV}$ $ \eta  < 2.4$	30
ATLAS (8TeV)	[46]	8000	46 - 66	$ y  < 2.4$	$p_T > 20\text{ GeV}$ $ \eta  < 2.4$	3
ATLAS (8TeV)	[46]	8000	116 - 150	$ y  < 2.4$	$p_T > 20\text{ GeV}$ $ \eta  < 2.4$	7
CMS (7TeV)	[47]	7000	60 - 120	$ y  < 2.1$	$p_T > 20\text{ GeV}$ $ \eta  < 2.1$	8
CMS (8TeV)	[48]	8000	60 - 120	$ y  < 2.1$	$p_T > 20\text{ GeV}$ $ \eta  < 2.1$	8
LHCb (7TeV)	[73]	7000	60 - 120	$2 < y < 4.5$	$p_T > 20\text{ GeV}$ $2 < \eta < 4.5$	8
LHCb (8TeV)	[74]	8000	60 - 120	$2 < y < 4.5$	$p_T > 20\text{ GeV}$ $2 < \eta < 4.5$	7
LHCb (13TeV)	[75]	13000	60 - 120	$2 < y < 4.5$	$p_T > 20\text{ GeV}$ $2 < \eta < 4.5$	9
Total						457

\*Bins with  $9 \lesssim Q \lesssim 11$  are omitted due to the  $\Upsilon$  resonance.

# SV2019

## *Kinematic coverage*



# SV2019

## Main settings

🍎  $b^*$  prescription:

$$b_*(b_T) = \sqrt{\frac{b_T^2 B_{\text{NP}}^2}{b_T^2 + B_{\text{NP}}^2}}$$

🍎 Non-perturbative function  $f_{\text{NP}}$ :

🍎 evolution:

$$g_K(b_T) = -c_0 b_T b_*(b_T) \rightarrow \begin{cases} -c_0 b_T^2 & \text{for } b_T \rightarrow 0 \\ -c_0 B_{\text{NP}} b_T & \text{for } b_T \rightarrow \infty \end{cases}$$

🍎 PDFs and FFs:

$$f_{NP}(x, b) = \exp\left(-\frac{\lambda_1(1-x) + \lambda_2 x + x(1-x)\lambda_5}{\sqrt{1 + \lambda_3 x^{\lambda_4}} \mathbf{b}^2} \mathbf{b}^2\right)$$

$$D_{NP}(x, b) = \exp\left(-\frac{\eta_1 z + \eta_2(1-z)}{\sqrt{1 + \eta_3(\mathbf{b}/z)^2}} \frac{\mathbf{b}^2}{z^2}\right) \left(1 + \eta_4 \frac{\mathbf{b}^2}{z^2}\right)$$

🍎 **11 free parameters** to fit to data.

🍎 Perturbative accuracies: **NNLL'(NNLO)**, **N<sup>3</sup>LL(-) (N<sup>3</sup>LO)**

🍎 **Monte Carlo** method for the experimental error propagation.

# SV2019

## *Fit quality*

🍎 Remarkably good total  $\chi^2$ ,

🍎 DY and SIDIS data are separately well described,

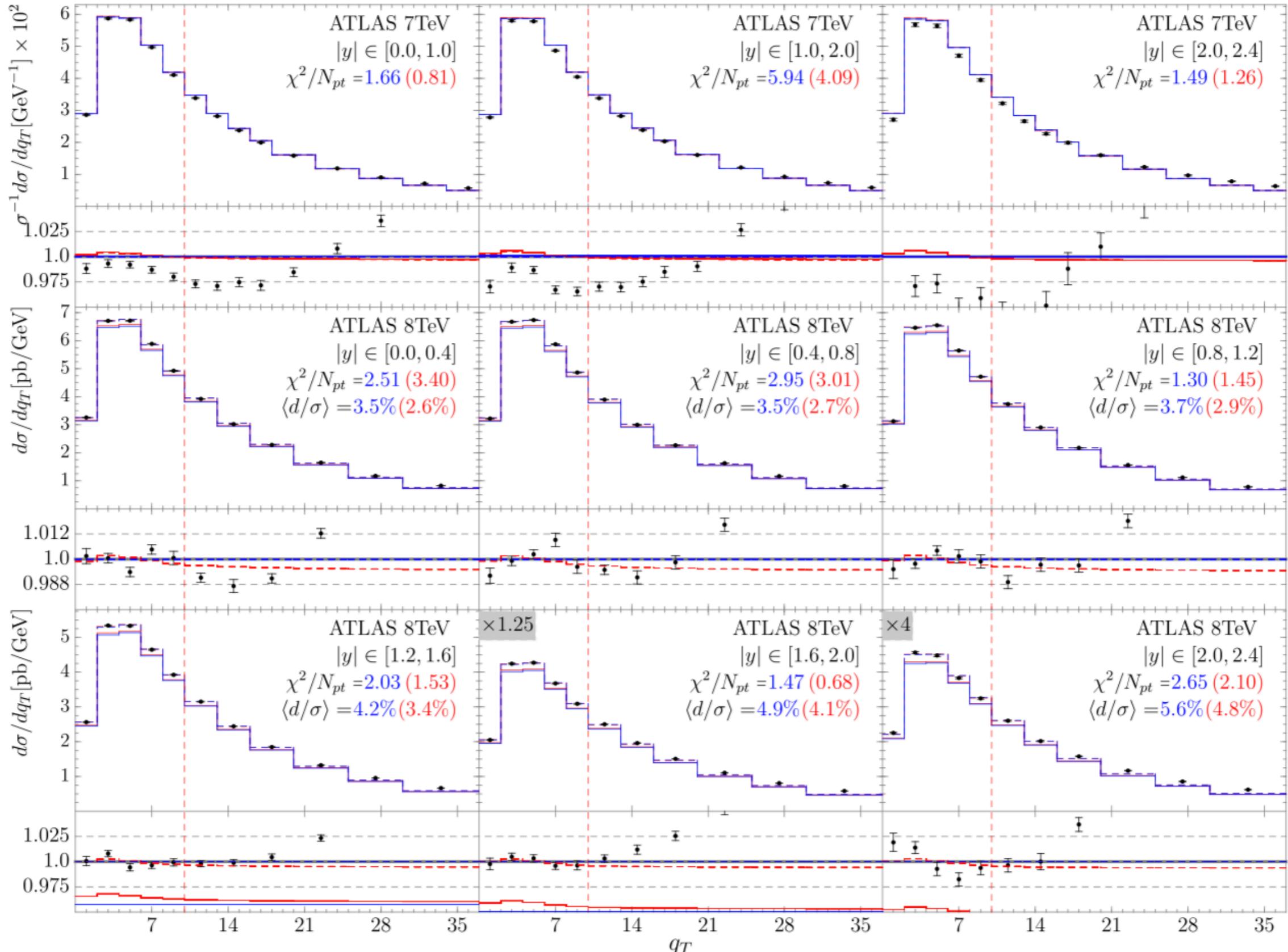
### 🍎 Important achievement:

🍎 simultaneous description of SIDIS and DY data within the same fit at high perturbative order.

Data set	$N_{pt}$	NNLO		$N^3\text{LO}$	
		$\chi^2/N_{pt}$	$\langle d/\sigma \rangle$	$\chi^2/N_{pt}$	$\langle d/\sigma \rangle$
CDF run1	33	0.66	8.4%	0.67	7.8%
CDF run2	39	1.28	2.8%	1.41	2.1%
D0 run1	16	0.72	0.1%	0.78	-0.5%
D0 run2	8	1.38	-	1.64	-
D0 run2 ( $\mu$ )	3	0.62	-	0.69	-
Tevatron total	99	0.97		1.06	
ATLAS 7TeV $0.0 <  y  < 1.0$	5	1.66	-	0.81	-
ATLAS 7TeV $1.0 <  y  < 2.0$	5	5.94	-	4.09	-
ATLAS 7TeV $2.0 <  y  < 2.4$	5	1.49	-	1.26	-
ATLAS 8TeV $0.0 <  y  < 0.4$	5	2.51	3.5%	3.40	2.8%
ATLAS 8TeV $0.4 <  y  < 0.8$	5	2.95	3.5%	3.03	2.7%
ATLAS 8TeV $0.8 <  y  < 1.2$	5	1.30	3.7%	1.45	2.9%
ATLAS 8TeV $1.2 <  y  < 1.6$	5	2.03	4.2%	1.53	3.4%
ATLAS 8TeV $1.6 <  y  < 2.0$	5	1.47	4.9%	0.70	4.1%
ATLAS 8TeV $2.0 <  y  < 2.4$	5	2.64	5.6%	2.10	4.8%
ATLAS 8TeV $46 < Q < 66\text{GeV}$	3	0.31	1.1%	0.31	0.2%
ATLAS 8TeV $116 < Q < 150\text{GeV}$	7	0.84	1.9%	0.97	1.2%
ATLAS total	55	2.12		1.82	
CMS 7TeV	8	1.25	-	1.24	-
CMS 8TeV	8	0.77	-	0.76	-
CMS total	16	1.01		1.00	
LHCb 7TeV	8	2.68	5.8%	2.37	5.2%
LHCb 8TeV	7	4.81	5.8%	4.16	5.1%
LHCb 13TeV	9	0.91	6.4%	0.81	5.7%
LHCb total	24	2.63		2.31	
<b>High energy DY total</b>	<b>194</b>	<b>1.51</b>		<b>1.42</b>	
PHE200	3	0.28	0.2%	0.29	-0.3%
E228-200	43	1.00	35.7%	1.12	35.0%
E228-300	53	0.90	29.2%	1.01	28.3%
E228-400	76	0.86	20.6%	0.96	19.5%
E772	35	1.84	9.5%	1.91	8.5%
E605	53	0.57	21.3%	0.60	20.1%
<b>Low energy DY total</b>	<b>263</b>	<b>0.96</b>		<b>1.04</b>	
HERMES ( $p \rightarrow \pi^+$ )	24	2.20	1.7%	3.06	2.2%
HERMES ( $p \rightarrow \pi^-$ )	24	1.12	0.6%	1.45	0.9%
HERMES ( $p \rightarrow K^+$ )	24	0.71	-0.1%	0.66	0.0%
HERMES ( $p \rightarrow K^-$ )	24	0.69	0.0%	0.66	0.0%
HERMES ( $d \rightarrow \pi^+$ )	24	0.57	0.3%	0.78	0.8%
HERMES ( $d \rightarrow \pi^-$ )	24	0.74	0.5%	0.96	0.7%
HERMES ( $d \rightarrow K^+$ )	24	0.52	-0.1%	0.53	0.0%
HERMES ( $d \rightarrow K^-$ )	24	1.27	0.0%	1.17	0.1%
HERMES total	192	0.98		1.16	
COMPASS ( $d \rightarrow h^+$ )	195	0.61	3.3%	0.76	5.1%
COMPASS ( $d \rightarrow h^-$ )	195	0.68	-2.3%	0.92	-0.5%
COMPASS total	390	0.65		0.84	
<b>SIDIS total</b>	<b>582</b>	<b>0.76</b>		<b>0.95</b>	
<b>Total</b>	<b>1039</b>	<b>0.95</b>		<b>1.06</b>	

# SV2019

## *Fit quality*



# SV2019

*Fit quality*

$$z^2 \times M(z, p_T)$$

$$d \rightarrow h^+$$

# COMPASS

