



Search for a heavy scalar X decaying to a scalar S and a Higgs boson in the $X \rightarrow SH \rightarrow bb\gamma\gamma$ channel with ATLAS Run-2 data

Maxime Fernoux

Supervisors : Elisabeth Petit and Arnaud Duperrin

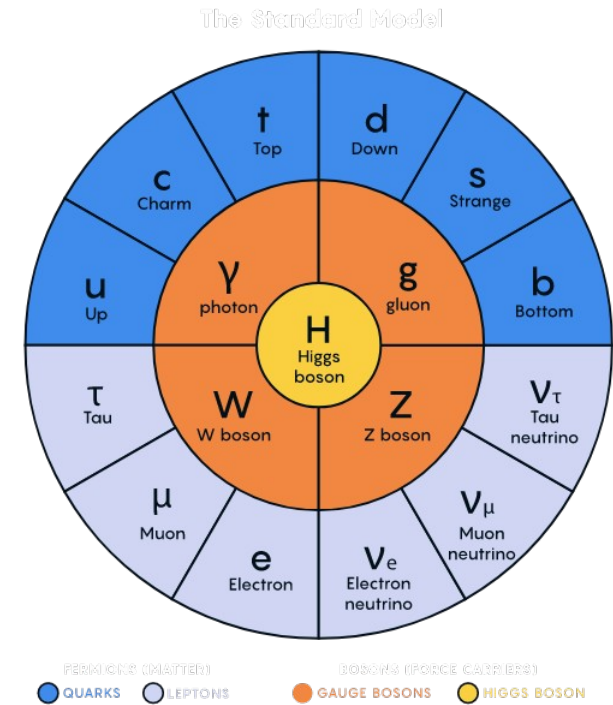
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Search for Beyond Standard Model physics

- SM is great but have some caveats :
 - Doesn't explain gravity
 - Doesn't explain neutrino oscillations and mass
 - No dark matter candidate

→ We need physics beyond standard model (BSM)

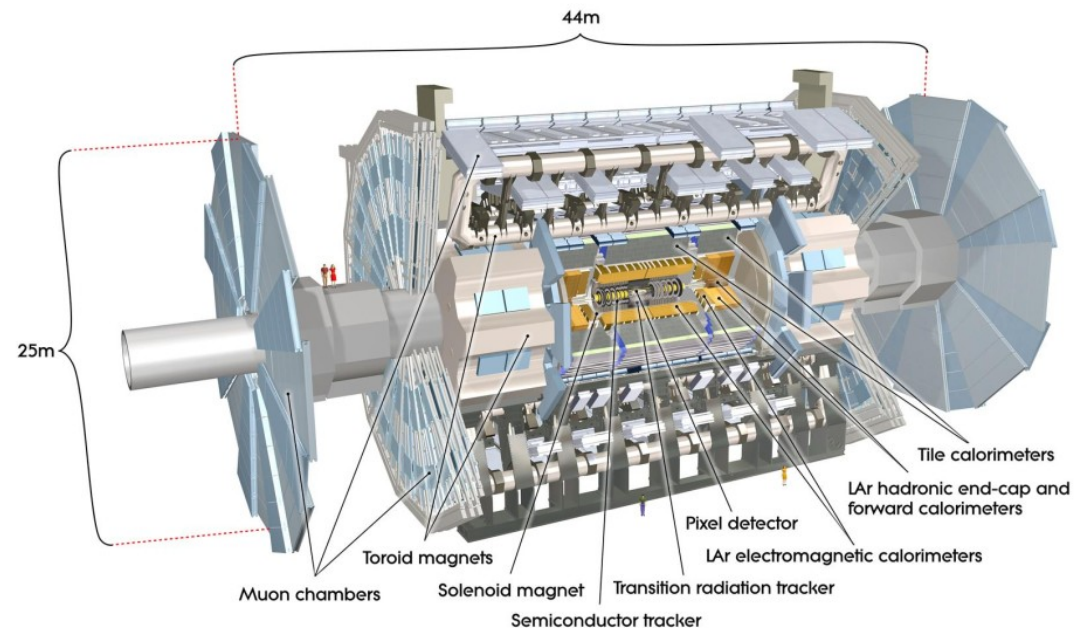
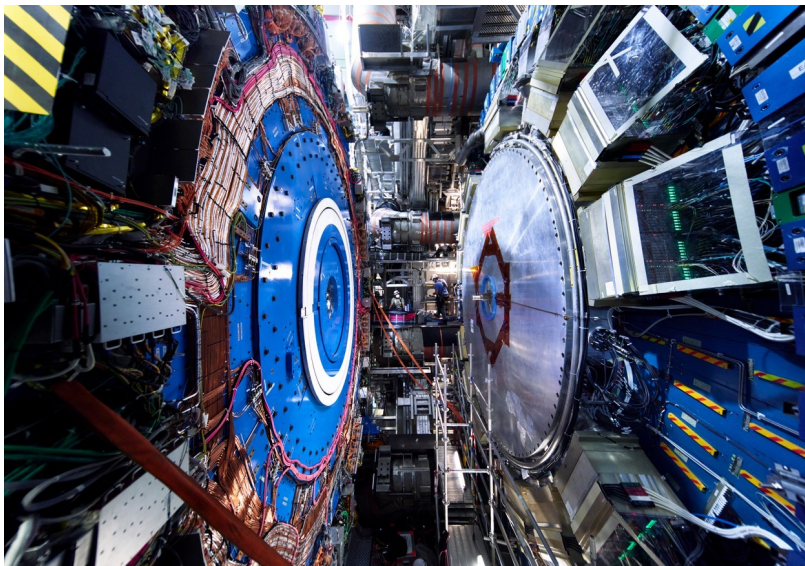


- Some BSM models predict additional scalar particles with different range of masses, notably in the Higgs sector
 - Double Higgs doublet model : 2HDM
 - Next-to-Minimal Supersymmetric Standard Model : NMSSM
 - ... etc
- Could be at reach of LHC and Atlas !



ATLAS detector

- ATLAS is a multipurposed detector using proton-proton collisions at the Large Hadron Collider (LHC) located at CERN
 - Higgs boson discovered in 2012
 - Make precision measurements to detect deviations from the SM
 - Search for BSM particles, especially in the Higgs sector

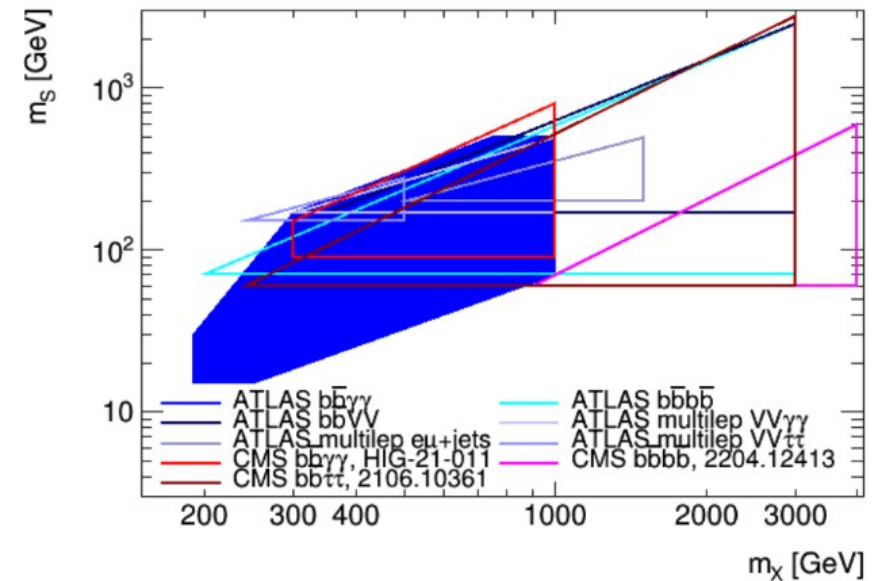
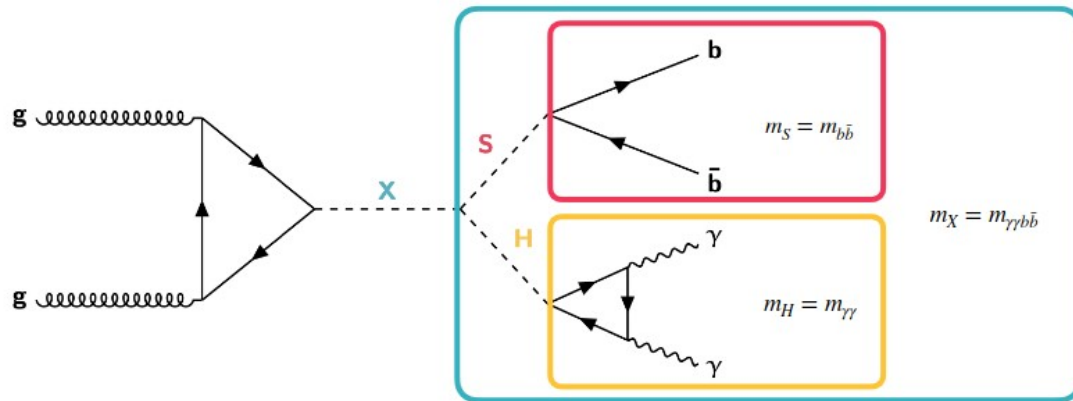


- This analysis uses data collected during Run-2 (from 2015 to 2018) with pp collisions at $\sqrt{s} = 13$ TeV

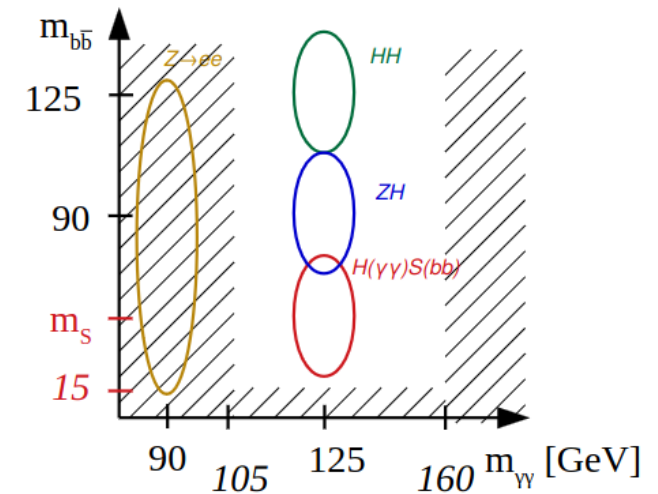


Targeted signal

- We search for a heavy scalar X decaying into a light scalar S and the SM Higgs where $H \rightarrow \gamma\gamma$ and $S \rightarrow b\bar{b}$
- The search is **model-independent** and a wide range of m_X , m_S signal is targeted :
 $15 < m_S < 500 \text{ GeV}$ and $170 \text{ GeV} < m_X < 1 \text{ TeV}$



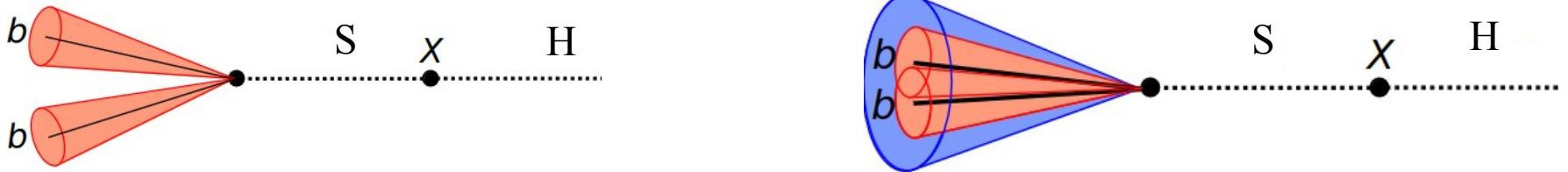
- Analysis is also heavily linked to $HH \rightarrow b\bar{b}\gamma\gamma$ analysis
 Di-Higgs is a major goal of LHC physics program
 - Can test new strategies with this final state
 - Could also help to remove potential background to HH signal



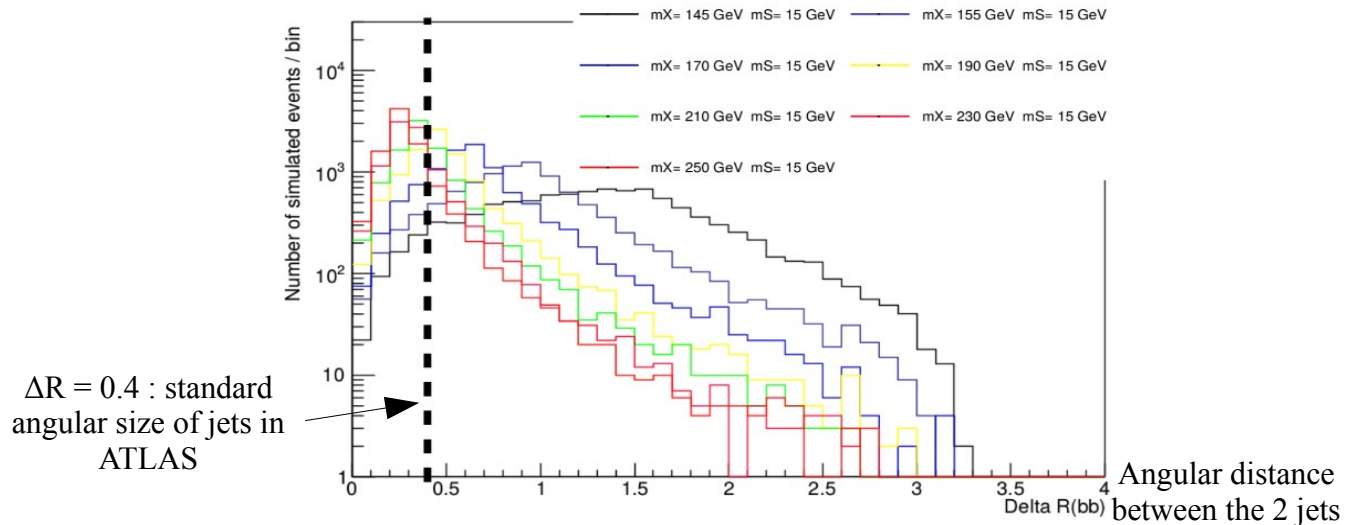


Different search regions

- A challenging situation arises when m_S is much smaller than m_X ($m_S/m_X < \text{around } 0.1$):
b-jets from the S decay are boosted and reconstructed as one b-tagged jet



→ We separate the search space in a **resolved region** with 2 b-tagged jets and a **merged region** with only one b-tagged jet





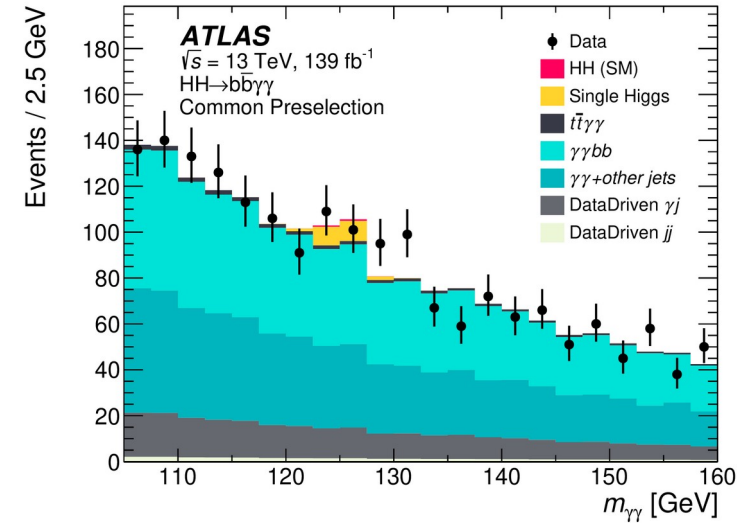
Event selection

- Selection

Table 4: The definitions of selections used in the analysis.

	2 <i>b</i> -tagged	1 <i>b</i> -tagged
Number of 'tight' and isolated photons	≥ 2	
$m_{\gamma\gamma}$ [GeV]	$\in [105, 160]$	
Number of leptons	$= 0$	
Number of central jets	$\in [2, 5]$	
Number of <i>b</i> -tagged jets @ 77% WP	$= 2$	$= 1$

Invariant photon mass distribution



- Predicted background yields (from theoretical cross sections) : main **non-resonant background** is $\gamma\gamma +$ jets, main **resonant** ones are $t\bar{t}H$, ggH , ZH and also $VBFH$ for 1 *b*-jet selection

2 *b*-tagged jets

	Selection
<i>HH</i> ggF+VBF	1.691 ± 0.004
<i>ZH</i>	3.691 ± 0.013
<i>WH</i>	0.207 ± 0.004
<i>VBFH</i>	0.685 ± 0.012
<i>bbH</i>	0.62 ± 0.023
<i>ggH</i>	5.453 ± 0.065
<i>tHjb</i>	0.969 ± 0.029
<i>tWH</i>	0.131 ± 0.005
<i>ttH</i>	8.313 ± 0.014
$Z(\rightarrow qq)\gamma\gamma$	20.345 ± 0.303
$t\bar{t}\gamma\gamma$	28.69 ± 0.109
$\gamma\gamma + \text{jets}$	1418.32 ± 4.596
Total SM	1489.116 ± 4.608

1 *b*-tagged jet

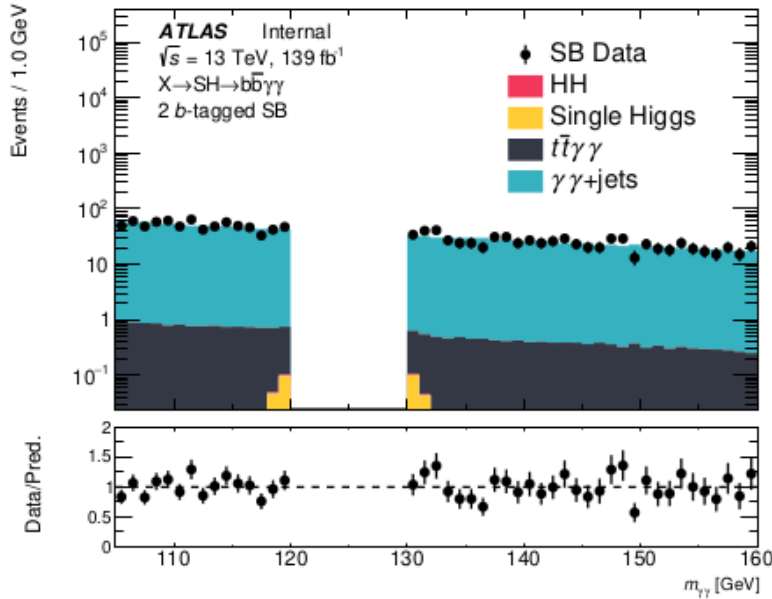
	Selection
<i>HH</i> ggF+VBF	1.827 ± 0.004
<i>WH</i>	5.97 ± 0.021
<i>VBFH</i>	8.333 ± 0.042
<i>ZH</i>	5.941 ± 0.015
<i>bbH</i>	2.973 ± 0.047
<i>ggH</i>	48.532 ± 0.202
<i>ggZH</i>	1.581 ± 0.01
<i>tHjb</i>	2.681 ± 0.05
<i>tWH</i>	0.572 ± 0.01
<i>ttH</i>	11.681 ± 0.017
$Z(\rightarrow qq)\gamma\gamma$	53.728 ± 0.8
$t\bar{t}\gamma\gamma$	49.78 ± 0.142
$\gamma\gamma + \text{jets}$	16298.8 ± 16.031
Total SM	16492.398 ± 16.053



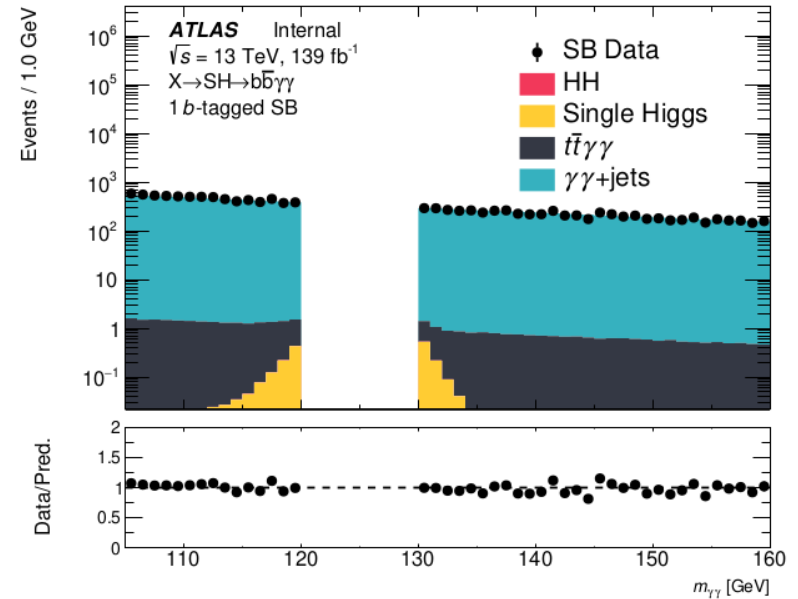
Analysis recap – Strategy

- The $m_{\gamma\gamma}$ distribution is used to split the events into a SR with $120 < m_{\gamma\gamma} < 130$ GeV and a sideband control region (CR)

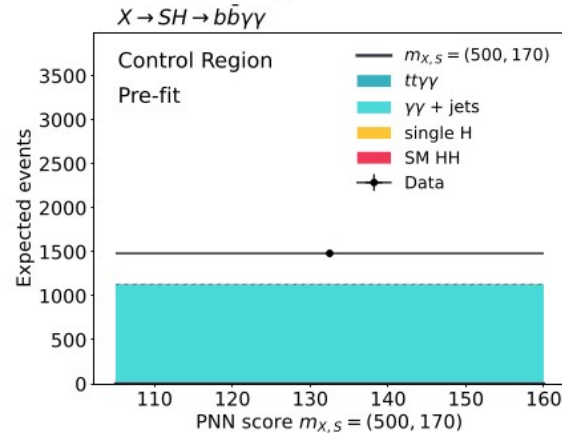
2 b-tagged jets



1 b-tagged jet



- The CR allows to correct the normalisation of the $\gamma\gamma$ + jets events using true data

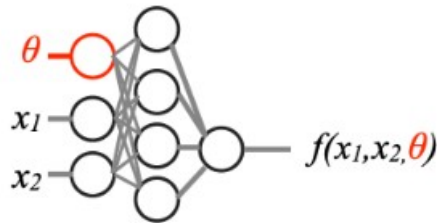
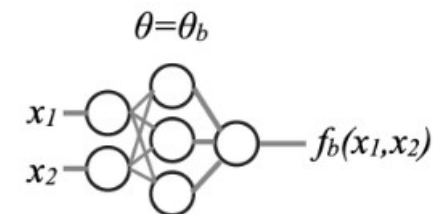
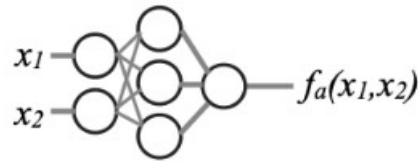




PNN discriminant

- A parametrised Neural Network (PNN) is used as discriminant in the SR. It is trained with simulated events that pass each selection (SR + SB)

$\theta = \theta_a$ Training variables = x_1, x_2 , Parameter = θ



1 PNN(θ) can act as M NNs optimised to target a specific θ

- Two separate PNNs :

- **2 b-tagged** :

Parameter $\theta = \mathbf{mX, mS}$

Training samples : signal, ttH, ggH, ZH and $\gamma\gamma$ + jets backgrounds (no HH – too signal-like and confuses the network)

Training variables : m_{jj} and $m_{\gamma\gamma j}^* = m_{\gamma\gamma j} - (m_{\gamma\gamma} - 125 \text{ GeV})$

- **1 b-tagged** : target low mS/mX values

Parameter $\theta = \mathbf{mX}$

Training samples : signal, VBFH, HH, ttH, ggH, ZH and $\gamma\gamma$ + jets backgrounds

Training variables : b-jet p_T and $m_{\gamma j}^* = m_{\gamma j} - (m_{\gamma} - 125 \text{ GeV})$

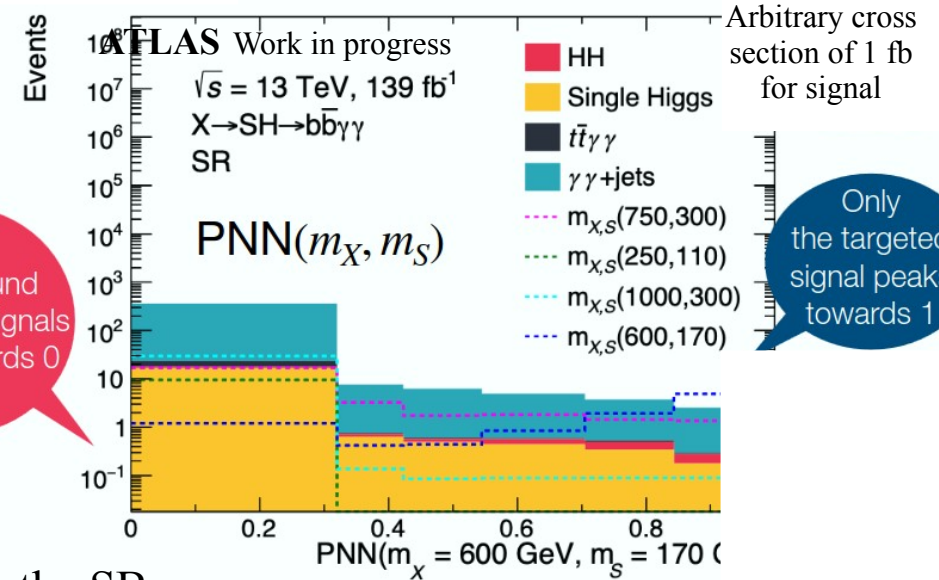


PNN distribution

- PNN shapes of backgrounds comes from MC samples
- Final results are computed with a binned log-likelihood fit on the PNN distribution

- Example of PNN distribution

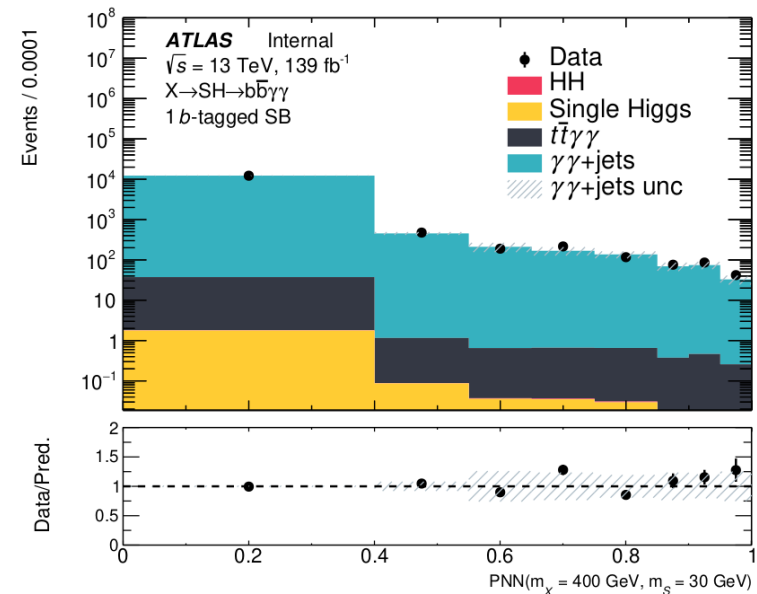
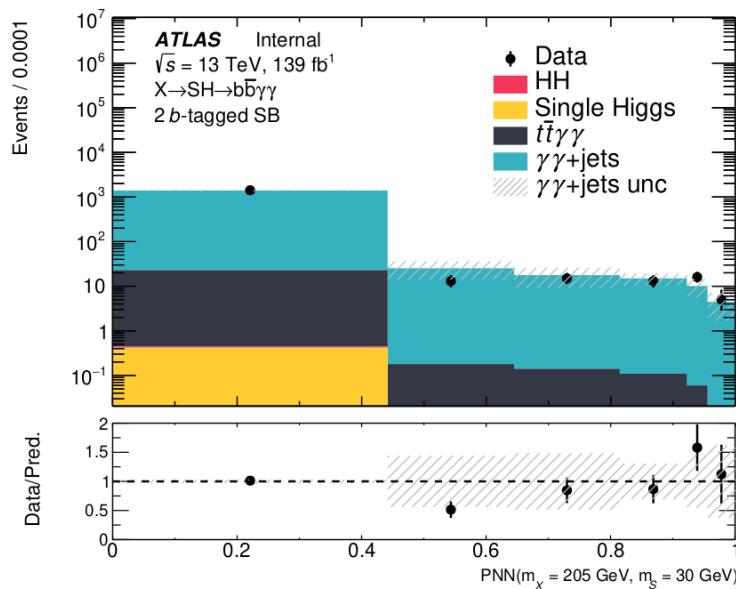
A constrain on the binning is to have at least 1 background event in every bin



Background and other signals peak towards 0

Only the targeted signal peaks towards 1

- Consistency between data and MC is checked in the SB

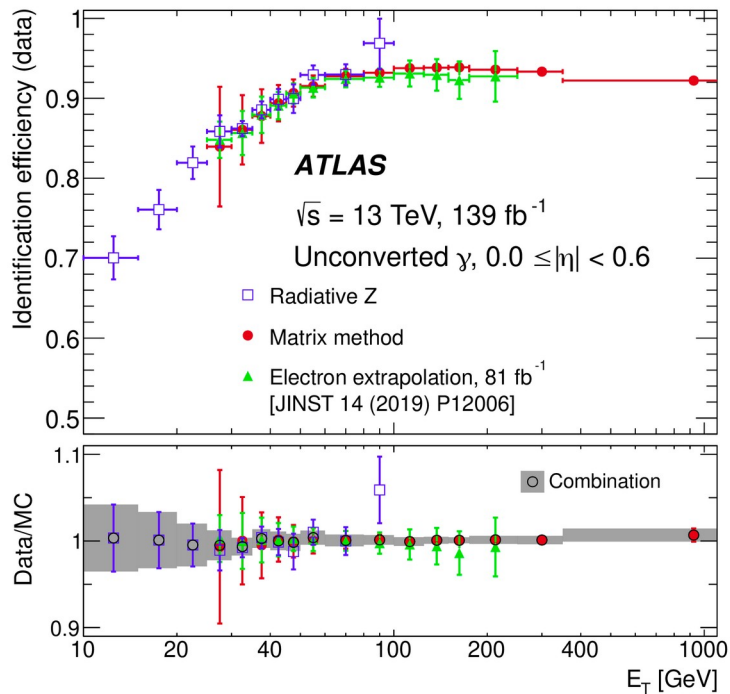




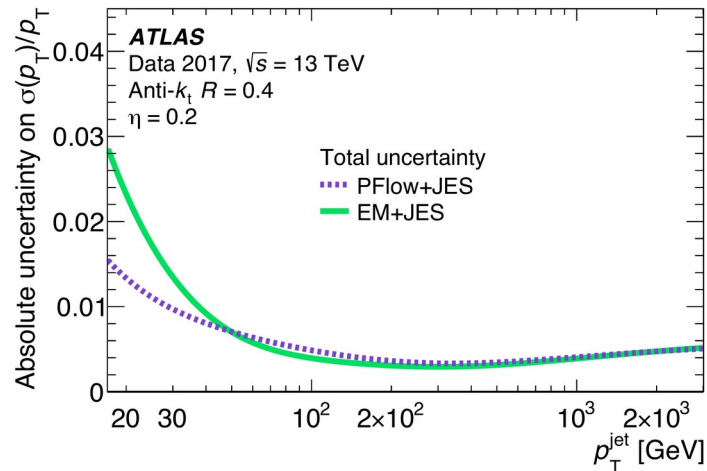
Experimental Systematics

- Physics is an experimental science → we have uncertainties affecting the measures
- Eventually at our analysis level it can have various impacts :
 - Particle identification can change the number of events in the CR and SR regions
 - Flavour tagging change the number of b-tagged jets
 - p_T and energy resolution can change the position and width of the peak in the $m_{\gamma\gamma}$, m_{bb} and $m_{bb\gamma\gamma}$ distributions and afterwards the shape of the PNN distribution

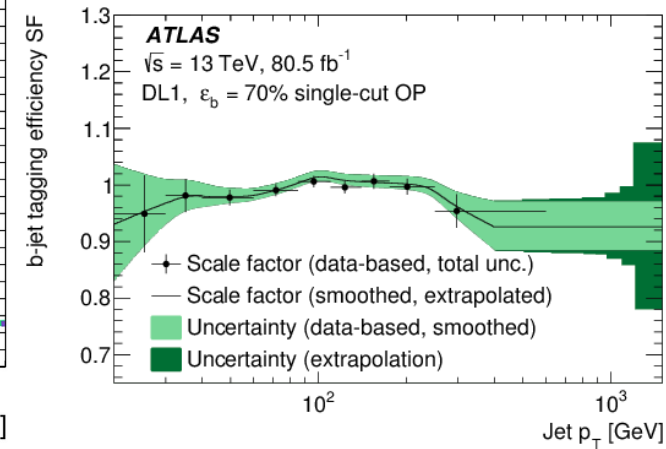
Photon identification efficiency as a function of E_T



Photon p_T resolution as a function of p_T



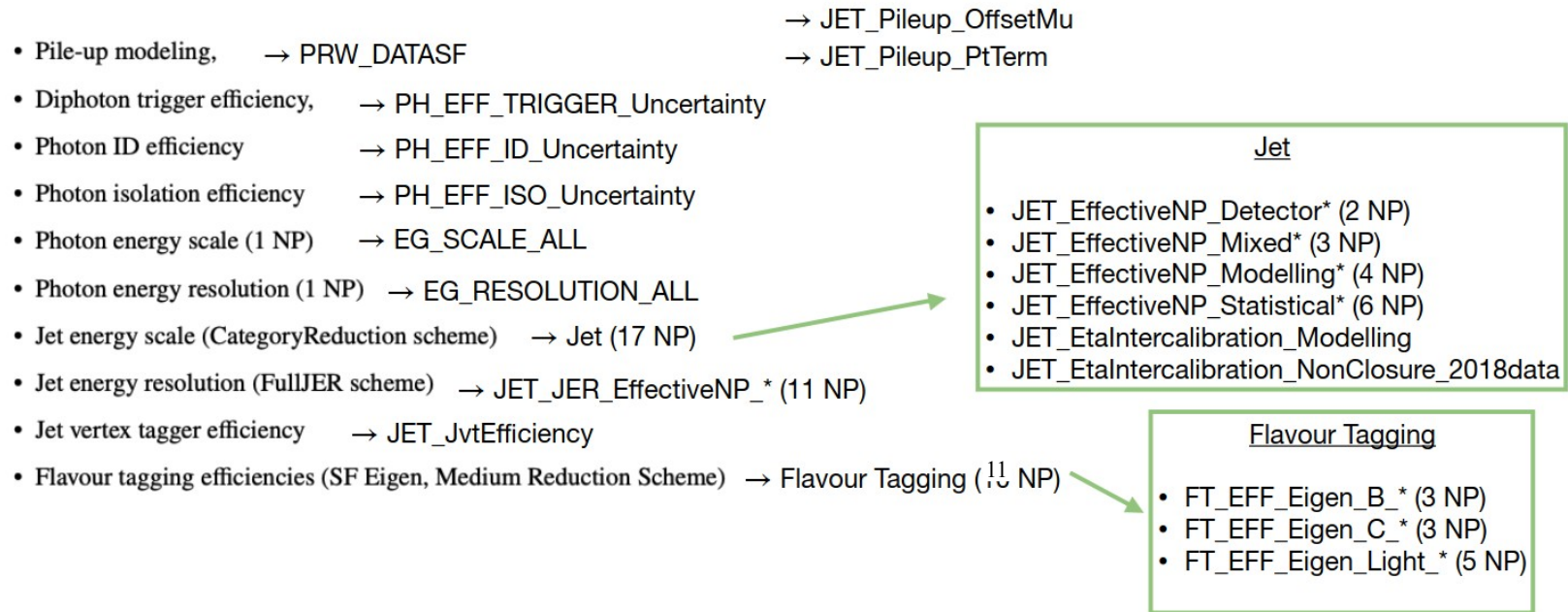
b-jet tagging efficiency as a function of p_T





Experimental Systematics

- In the analysis framework experimental systematics are studied through MC samples at $\pm 1\sigma$ away from nominal values from each effect
- They have been produced for major backgrounds only : ttH, ZH, ggH, ggHH and VBFHH, ggZH, $\gamma\gamma$ + jets and $Z(bb/qq)\gamma\gamma$
- Systematics are treated as nuisance parameters (NP) in the fit \rightarrow 47 in total !
Like the number of slices of a prefou from Vendee





Experimental Systematics

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- Systematics are treated as nuisance parameters (NP) in the fit \rightarrow 47 in total !
Like the number of slices of a prefou from Vendée

- Pile-up
- Diphoton
- Photon
- Photon i
- Photon e
- Photon e
- Jet ener
- Jet ener
- Jet verte
- Flavour



ata



Experimental Systematics – Yield change

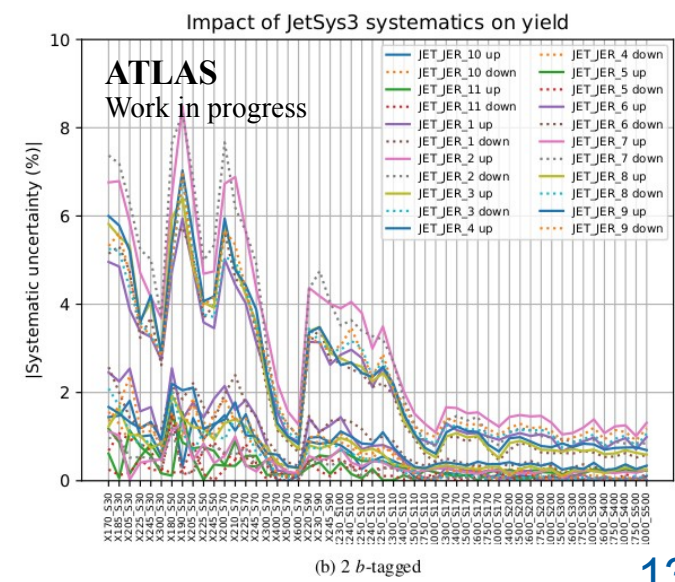
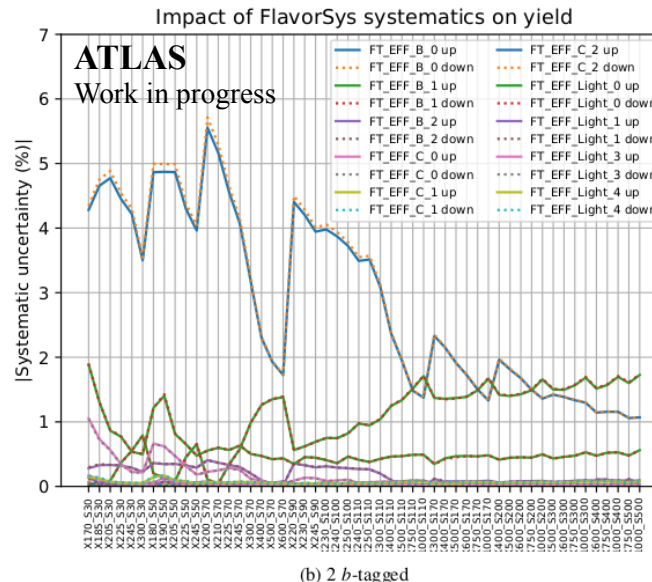
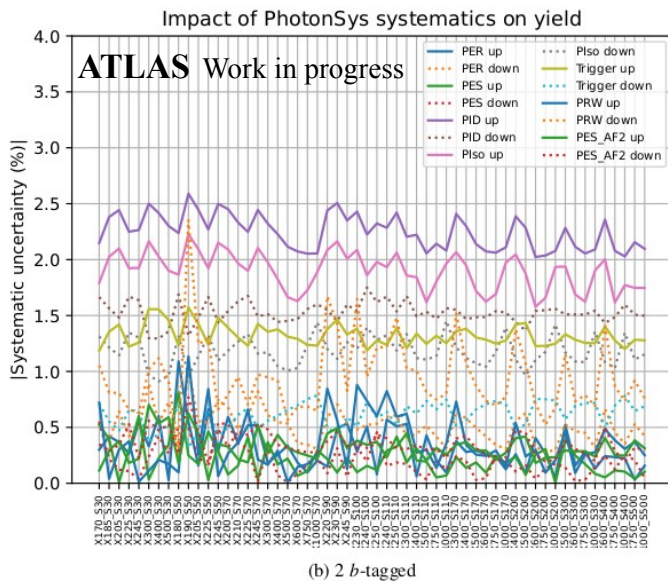
- Systematics change both the yields and the shape of the PNN distribution of the samples

Table 18: Uncertainty [in %] on the yield for backgrounds in the 2 *b*-tagged category.

Main uncertainties

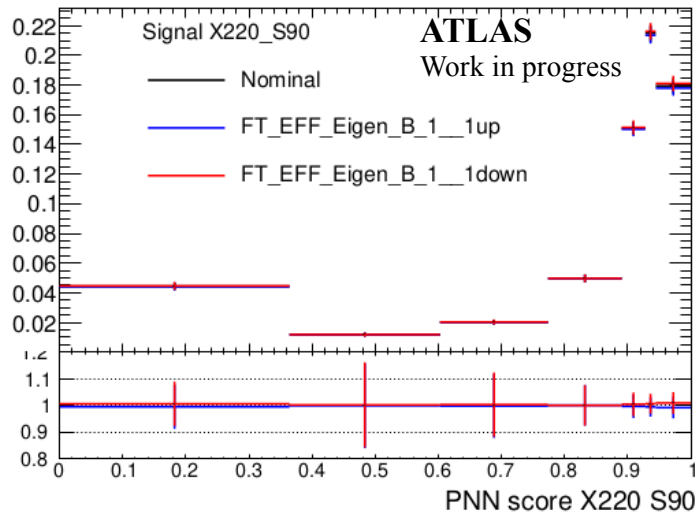
Source		Yield uncertainty [%]								
		ttH	ZH	ggHH	ggH	ggZH	tHjb	VBFHH	Zqqyy	Zbbyy
Event-based	Photon Trigger	1.01	1.02	1.00	1.04	0.98	1.03	1.02	1.00	1.25
	Pile-up reweighting	0.88	0.76	0.56	0.43	0.60	0.99	0.60	3.69	1.15
Photon	Photon Energy Resolution	0.42	0.42	0.34	0.42	0.43	0.63	0.43	2.79	0.78
	Photon Energy Scale	0.17	0.18	0.12	0.07	0.11	0.18	0.24	14.53	0.90
	Photon ID	1.59	1.61	1.44	1.64	1.49	1.60	1.59	1.72	1.96
	Photon Isolation	1.55	1.57	1.45	1.60	1.46	1.59	1.59	1.27	1.88
Jet	Jet Energy Scale	1.36	0.94	0.55	1.81	0.74	0.76	0.72	5.30	1.09
	Jet Energy Resolution	7.33	4.60	2.91	7.50	3.36	4.88	3.08	0.68	5.37
	b-jet efficiency	2.07	2.99	2.51	3.05	2.55	2.30	2.83	0.10	3.36
Flavour-tagging	c-jet efficiency	0.40	0.71	0.06	1.68	0.60	0.92	0.07	13.12	0.22
	light-jet efficiency	0.79	0.38	0.40	2.72	0.51	0.90	0.42	1.85	0.48

- For signal, yields changes are dependant on mS and mX
Jet energy resolution systematics are the most important (but below 10%)

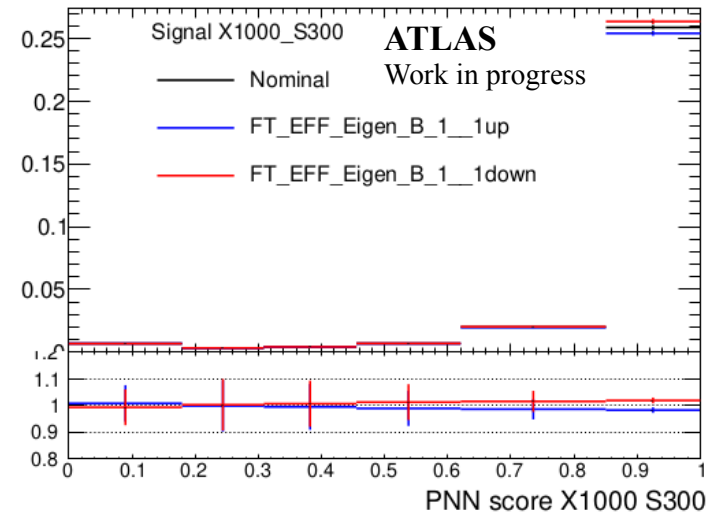




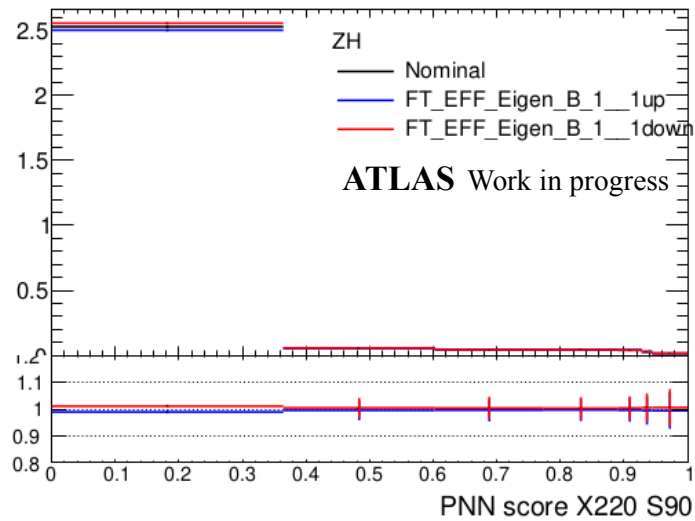
Experimental Systematics – Shape change



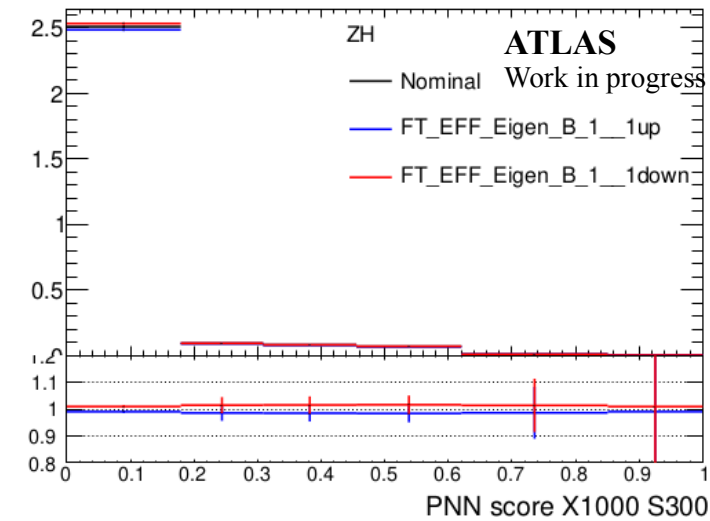
(a) SH signal, $m_X = 220$ GeV, $m_S = 90$ GeV



(b) SH signal, $m_X = 1000$ GeV, $m_S = 300$ GeV



(e) ZH, $m_X = 220$ GeV, $m_S = 90$ GeV



(f) ZH, $m_X = 1000$ GeV, $m_S = 300$ GeV

- NB : for $\gamma\gamma + \text{jets}$, only shape changes are used as normalisation is imposed by the sideband

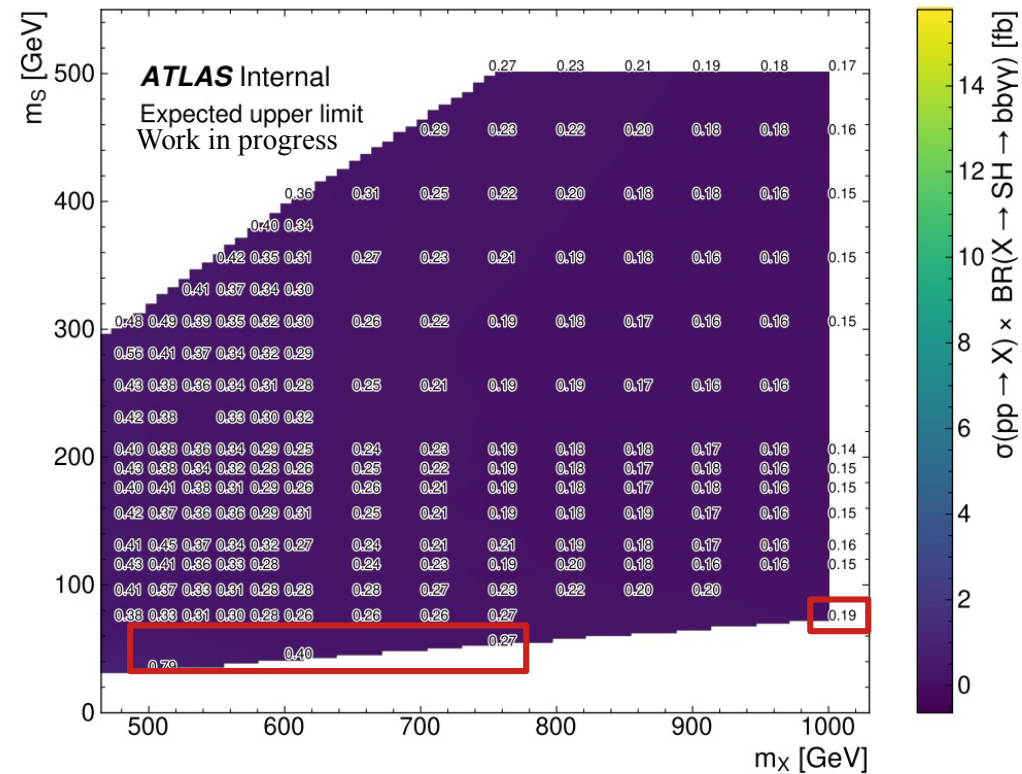
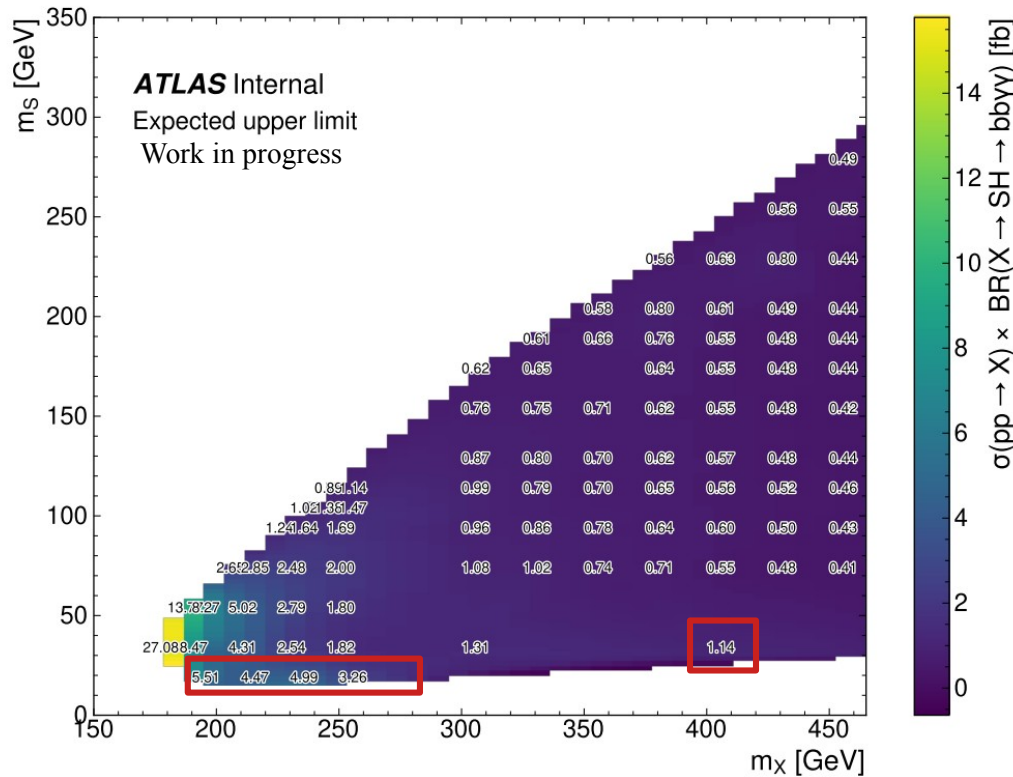


Blinded expected limits

- If no signal is observed, the goal of the analysis is to set upper limits on the cross section of the $X \rightarrow SH$ signal in the $b\bar{b}\gamma\gamma$ final state
- I can only show blinded results (i.e using Monte-Carlo and not true ATLAS data)
Not final results but gives an idea of the analysis sensibility



1 b-jet points

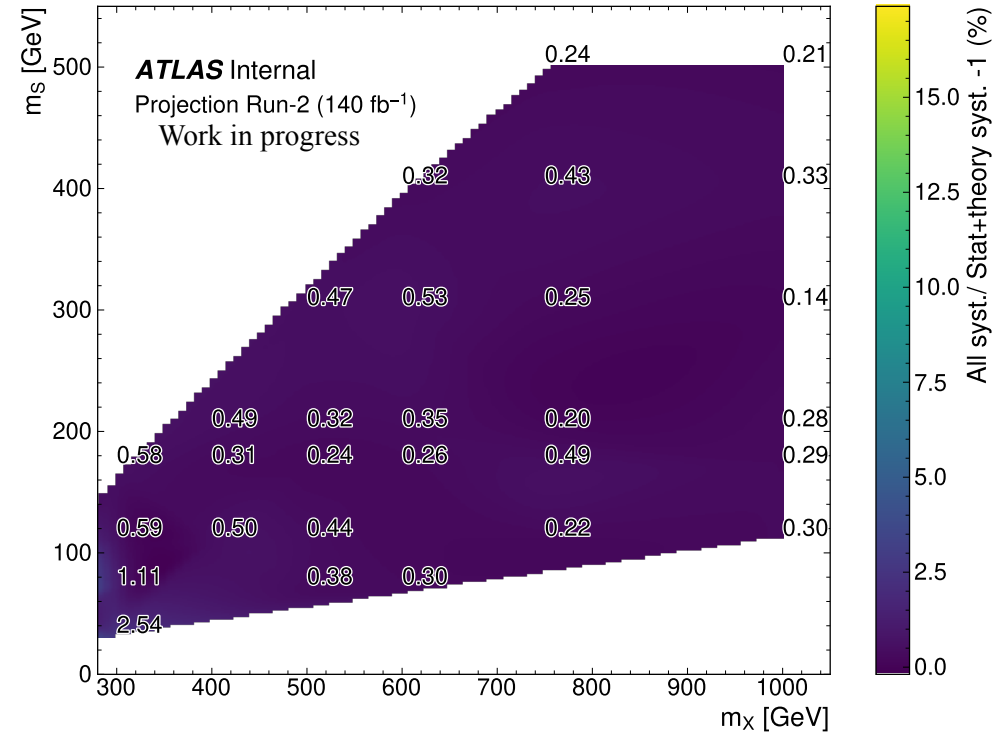
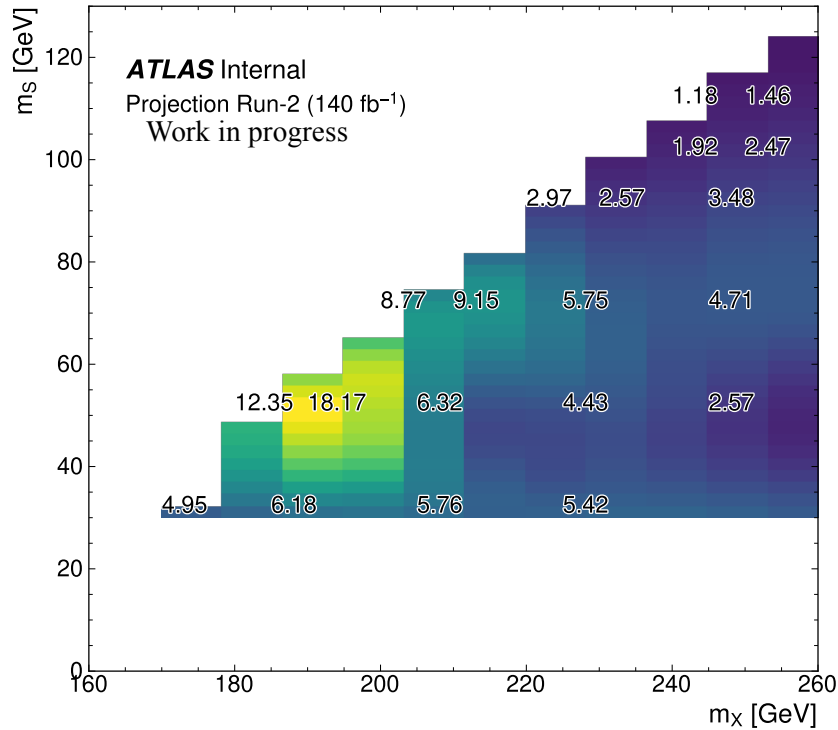


- Limits range from 0.15 to 30 fb
Sensibility is better in high mass region



Impact of experimental systematics uncertainty

- Here we plot the ratio between blinded expected limits and the limits obtained without taking into account experimental systematics to check their impact :

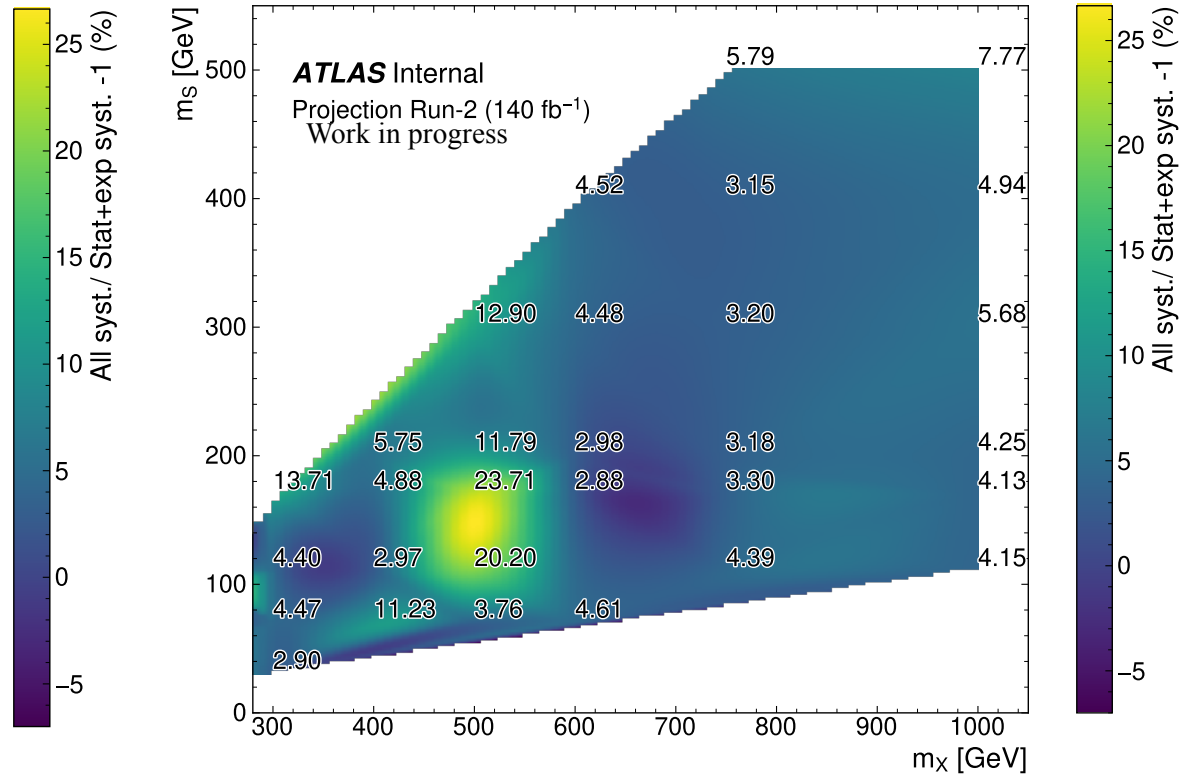
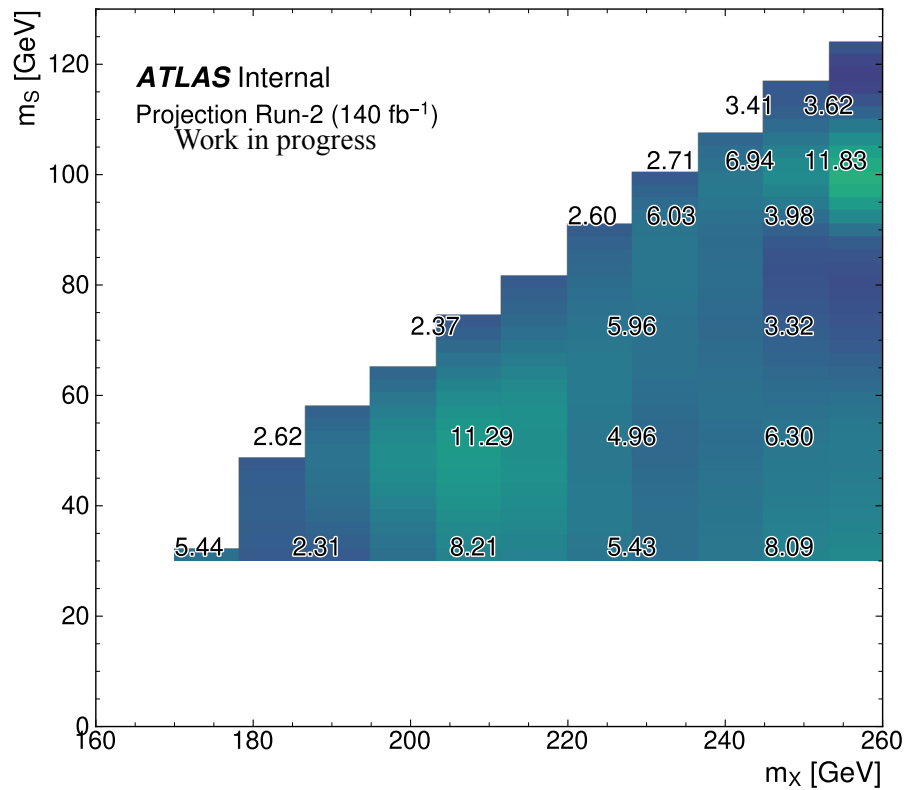


- Their impact on limit can reach 18% but are mostly between 5-10% at low mass and below 1% at high mass



Impact of theoretical systematics uncertainties

- Same plot as before but with theoretical uncertainties

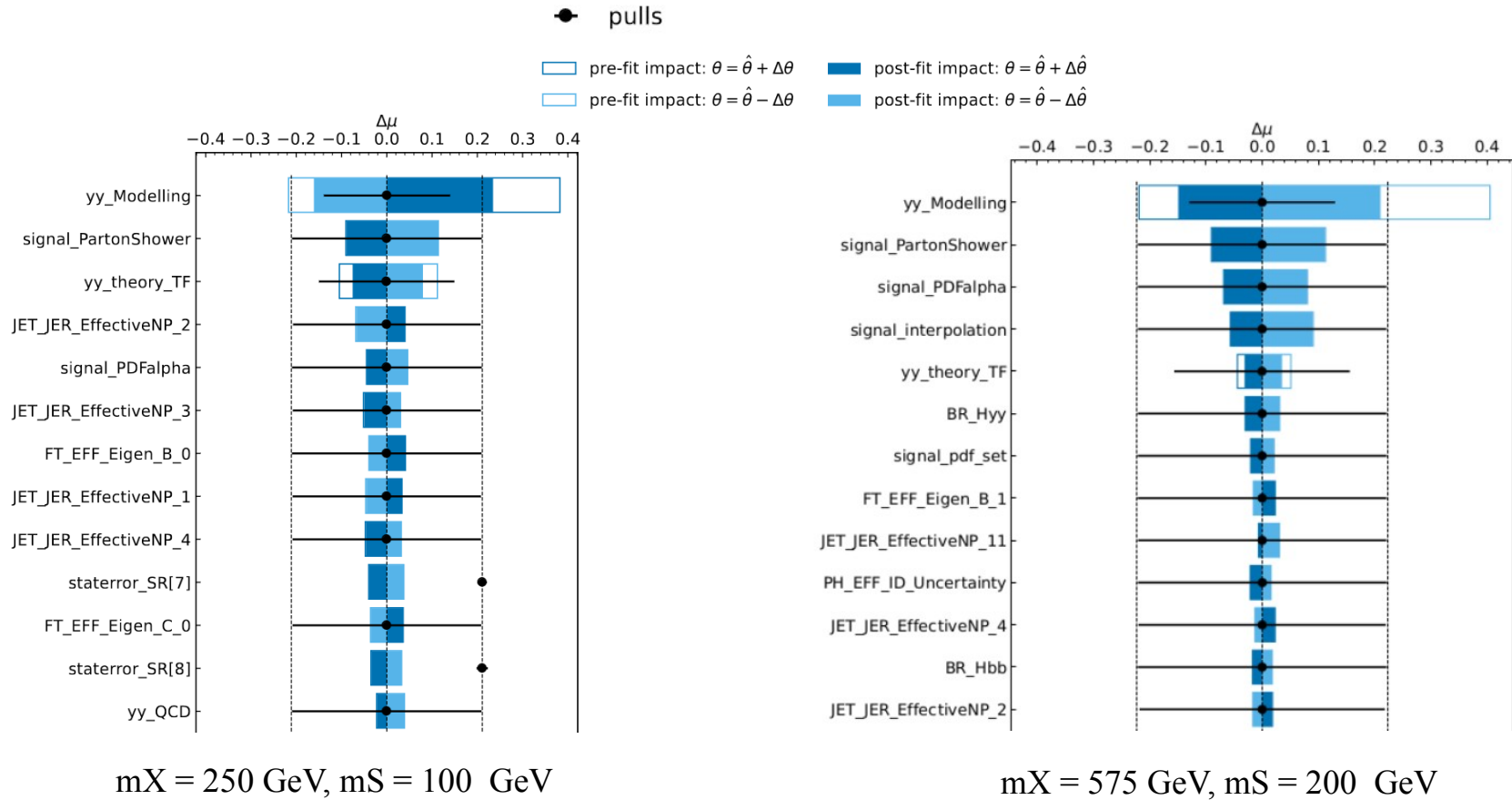


- Theory systematics impact vary a lot and can reach 20%
- It seem to be dominated by $\gamma\gamma$ modelling (in back-up slide)
Other theoretical systematics account for 3-4%



Impact of systematics uncertainty – ranking plots

- Ranking plots : class the uncertainties with the impact they have on the fit POI



- Largest systematics is the modelling of the $\gamma\gamma + \text{jets}$ background
- Largest experimental systematics are about flavour tagging and jet energy resolution



Summary

- A search for a resonant scalar particle X decaying into a scalar S and SM Higgs is performed on the $X \rightarrow SH \rightarrow b\bar{b}\gamma\gamma$ channel with the ATLAS Run-2 data
- Most interesting point in this analysis is the PNN that has been developed to target the signal for any values of m_X and m_S
- Analysis at the group internal review stage
More analyses to come with Run-3 data !



Back-up



All systematics uncertainties

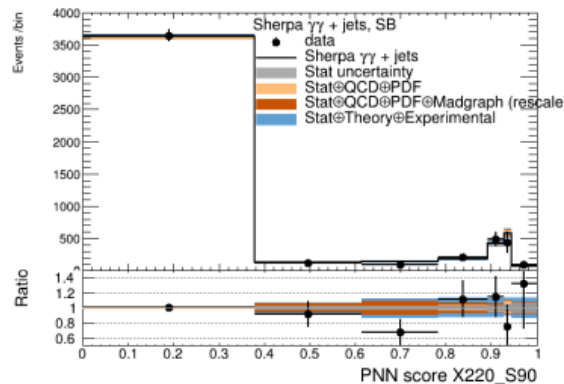
Main uncertainties

		Signal	HH ggF	HH VBF	ttH & ZH	Other Single Higgs	Continuum $\gamma\gamma$ +jets
Theory	Normalisation	$BR(H \rightarrow \gamma\gamma)$	$BR(H \rightarrow \gamma\gamma)$ $BR(H \rightarrow b\bar{b})$ PDF+ α_S Scales + m_t	$BR(H \rightarrow \gamma\gamma)$ $BR(H \rightarrow b\bar{b})$ PDF+ α_S Scales	$BR(H \rightarrow \gamma\gamma)$	$BR(H \rightarrow \gamma\gamma)$ PDF+ α_S Scales	$\gamma\gamma$ transfer factor
	Shape+Norm.	Scales, PDF+ α_S Parton shower Interpolation	Parton Shower		Scales, PDF+ α_S Parton Shower		Scales, PDF+ α_S Modelling
Exp.	Shape+Norm.	Pile-up modelling Diphoton trigger efficiency Photon identification and isolation efficiency Photon energy scale and resolution Jet energy scale and resolution Jet vertex tagger efficiency Flavour tagging efficiency (all exp. systematics are neglected for bbH , tH and $VBF H$)					

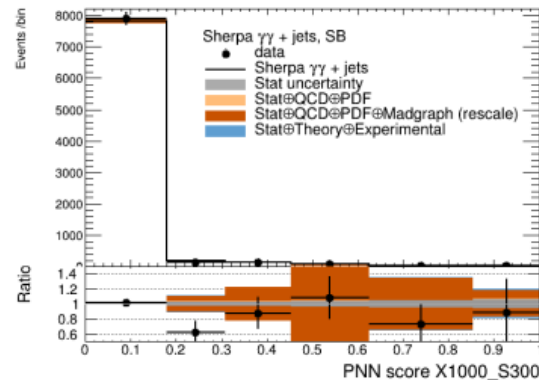


Theoretical systematics uncertainties

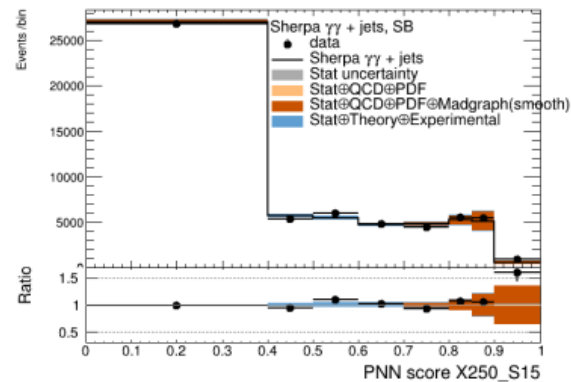
- Largest theoretical uncertainty is the modelling of $\gamma\gamma + \text{jets}$ events which is difficult to handle
- Normalization of the $\gamma\gamma + \text{jets}$ events is determined by a normalization factor from the sideband distribution
- The uncertainty regarding the modelling is evaluated by comparing simulated events from two different MC generators : Sherpa and MadGraph



(a) $m_X = 220 \text{ GeV}, m_S = 90 \text{ GeV}$



(b) $m_X = 1000 \text{ GeV}, m_S = 300 \text{ GeV}$

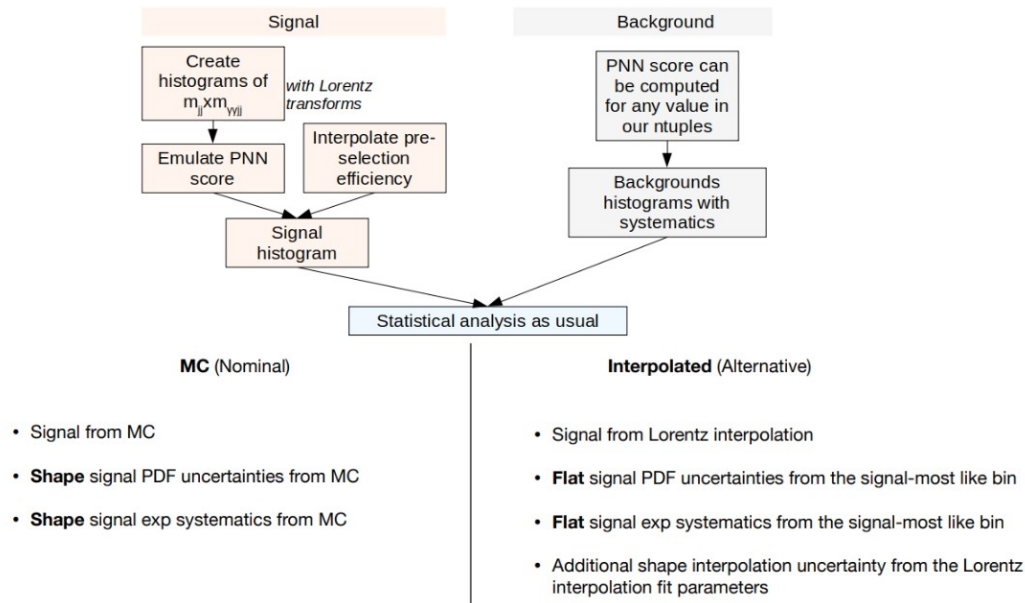


(c) $m_X = 250 \text{ GeV}, m_S = 15 \text{ GeV}$



Interpolation strategy

- Why interpolate ?
 - We need to be able to look for any signal in the region and granularity is not precise enough with MC samples
- Interpolation works separately for signal and background
 - PNN score can be computed for any m_X , m_S values in background samples
 - For signal we need to interpolate both the yields and the PNN shape
 - The shape is obtained with Lorentz transforms
 - The yields are obtained using Delaunay triangulation from the available MC samples

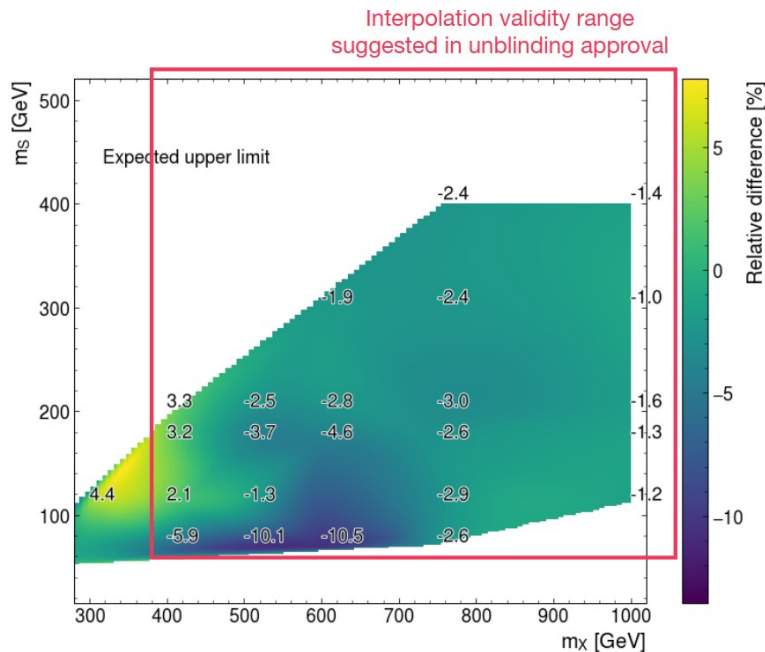
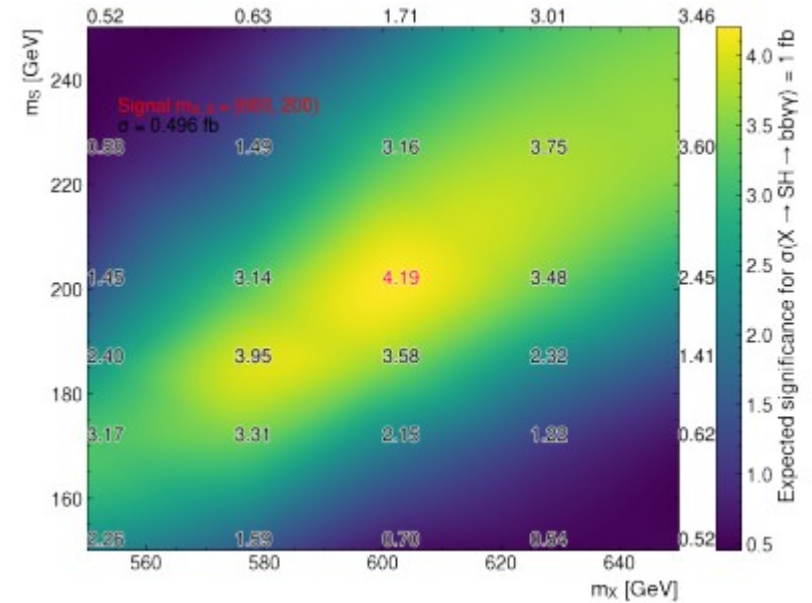




Interpolation range and validity

- Where should we interpolate ?
- Injection tests are made to ensure granularity is enough to allow us to be sensible to any signal in the probed region

→ we inject signal at $\sigma = 2 \times$ expected limit and want at least one neighbouring point to have an expected significance ≥ 3



- Validity of interpolation is evaluated by comparing interpolated and MC signal limits
 - Difference is below 5% for most points with a maximum of $\sim 10\%$
 - Interpolation is more difficult in low m_S regions where there are some jets overlap
- Interpolation is made for points with $m_X \geq 300 \text{ GeV}$ and $m_S \geq 70 \text{ GeV}$