

Combined angular analysis of $B \rightarrow D^* e \nu_e$ and $B \rightarrow D^* \mu \nu_\mu$ at the LHCb detector

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Journées de Rencontre Jeunes Chercheurs 2023 - October 23, 2023



Content

1 Introduction

- Standard Model
- Lepton flavour universality
- Kinematics of $B \rightarrow D^* \ell \nu_\ell$
- Motivation for the analysis

2 Methodology

- Template fit
- Fit examples

3 New Physics contributions

- Template fit model independence
- New Physics sensitivity tests

4 Results

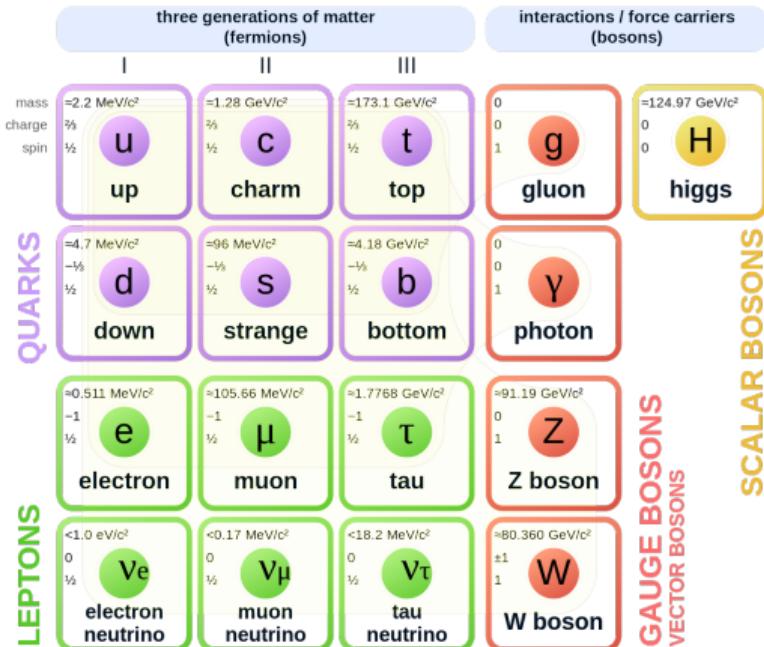
- Status and outlook

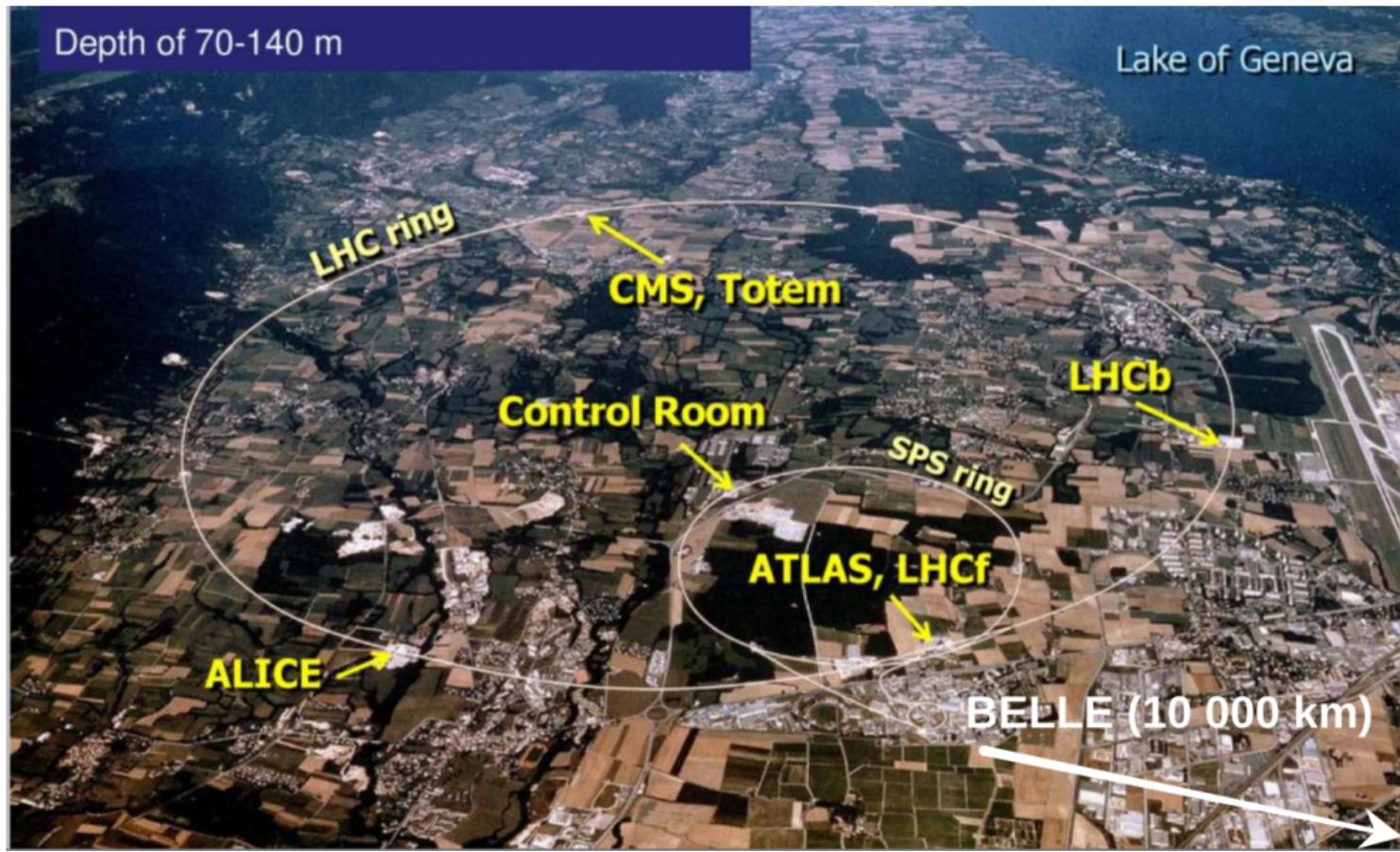
5 Backup Slides

Standard Model(SM)

- SM is a **Quantum Field Theory (QFT)** where particles are described as excitations of quantum fields
- The SM counts six "flavors" of **quarks** and six "flavours" of **leptons**
- Spin 1/2 particles called fermions:
 - Quarks
 - Three colours
 - Cannot be found isolated in nature. Are combined into colorless objects called hadrons, which are combination of:
two (Mesons)
three (Baryons)
 - ...
 - Leptons:
 - Electrically neutral - neutrino
 - Charged leptons

Standard Model of Elementary Particles

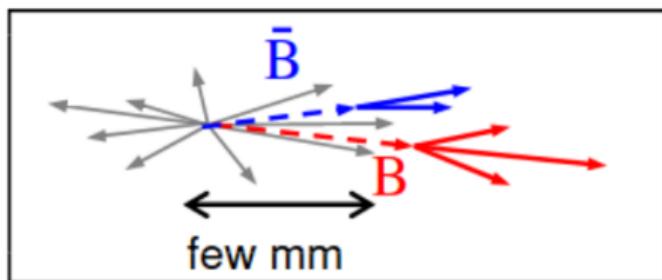




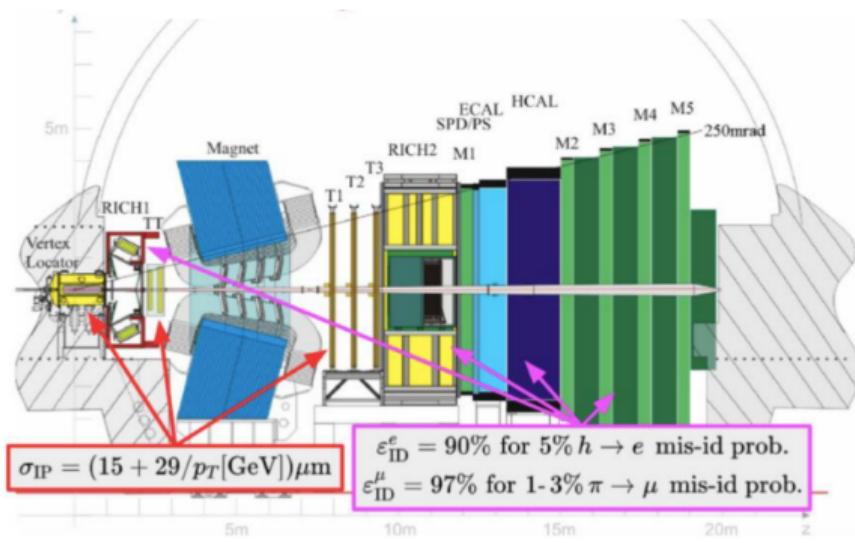
LHCb detector

Single-arm **forward** spectrometer:

- Forward angular coverage → b-hadrons produced in pairs ($b\bar{b}$) in the same direction
- Excellent performance of tracking, powerful particle identification allows to perform **high precision measurement of semileptonic B decays**



B mesons have displaced vertices. Fly a few mm before decaying in inner tracking detector - Vertex Locator

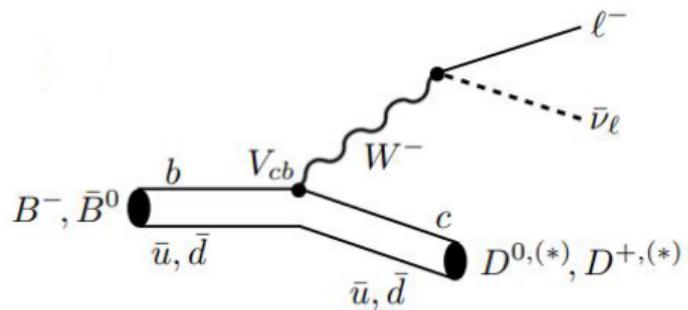


Introduction - LFU

- Semileptonic b-hadron decays provide powerful probes for testing the

Lepton Flavor Universality,

which states that the interactions of the electroweak bosons with the leptons are independent of the lepton flavor



- LFU can be tested with ratios of branching fractions to final states with different lepton flavors:

$$R(D^*)_{e/\mu} = \frac{Br(B^0 \rightarrow D^* e \nu_e)}{Br(B^0 \rightarrow D^* \mu \nu_\mu)}$$

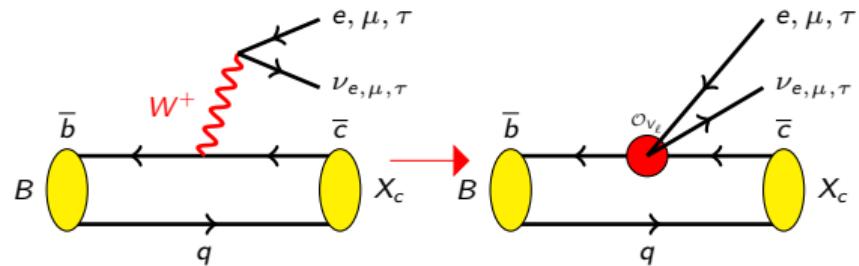
- New Physics (NP) can be detected in angular coefficients even if $R(D^*)_{e/\mu}$ is compatible with SM. Furthermore NP can be characterized in the angular analysis

Effective field theory

- The low-energy effective theory based on the assumption of three light lefthanded neutrino flavours below the electroweak scale for $B^0 \rightarrow D^* \ell \nu_\ell$ at dimension six can be written as

$$\mathcal{H}_{\text{eff}}(b \rightarrow c \ell) = \frac{4G_F}{\sqrt{2}} V_{cb} \sum_i (\mathcal{O}_{SM} + C_i \mathcal{O}_i)$$

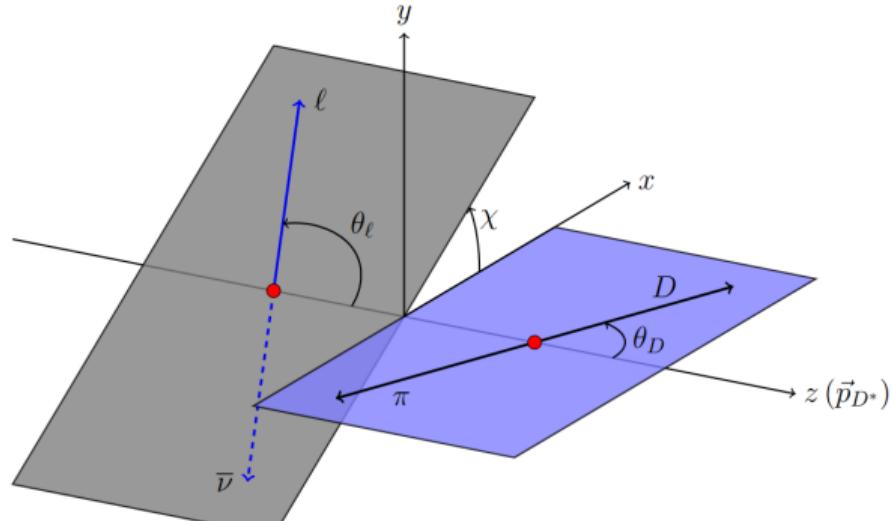
- Dominant SM. NP corrections parametrized by adding new operators and coefficients
- Wilson coefficients**
- Wilson operators**



$$\mathcal{H}_{\text{eff}} = \frac{4G_F}{\sqrt{2}} V_{cb} \{ [(1 + C_{V_L}) P_L + C_{V_R} P_R] \gamma_\mu P_L + [C_S + C_P \gamma^5] P_L + C_T \sigma^{\mu\nu} P_L \sigma_{\mu\nu} P_L + h.c. \}$$

- C_L, C_R, C_S, C_P, C_T are complex NP couplings ($\equiv 0$ in SM)
- Different NP models (W' , H^+ , LQ , ...) \Rightarrow different combinations of couplings

Introduction - Kinematics of $B \rightarrow D^*(\rightarrow D\pi)\ell\nu_\ell$



$$B \rightarrow D^* \ell \nu_\ell$$

$$D^* \rightarrow D\pi$$

$$D \rightarrow K\pi$$

Taken from [arXiv:1907.02257](https://arxiv.org/abs/1907.02257)

Kinematic variables : $\cos \theta_\ell, \cos \theta_D, \chi, q^2 = (p_\nu + p_\ell)^2$

$$\ell = \mu, e$$

Introduction - Analysis Idea

- The experimental fit of the full differential rate of $B \rightarrow D^* \ell \nu_\ell$ is rather complicated:
 - The analysis suffers from biases in the angular distributions caused by various reconstruction and acceptance effects at LHCb
 - Binned multidimensional template histograms must be used to resolve these biases
- For the simplification purposes, template histograms can be one dimensional if the analyses concentrate on single-differential distribution. In such case, 5(4 in CP averaged measurement) angular coefficients can be obtained

$$\frac{1}{\hat{\Gamma}^{(\ell)}} \frac{d\hat{\Gamma}^{(\ell)}}{d \cos \theta_\ell} = \frac{1}{2} + \langle A_{FB}^{(\ell)} \rangle \cos \theta_\ell + \frac{1}{4}(1 - 3\langle F_L^{(\ell)} \rangle) \frac{3 \cos^2 \theta_\ell - 1}{2}$$

$$\frac{1}{\hat{\Gamma}^{(\ell)}} \frac{d\hat{\Gamma}^{(\ell)}}{d \cos \theta_D} = \frac{3}{4}(1 - \langle F_L^{(\ell)} \rangle) \sin^2 \theta_D + \frac{3}{2}\langle F_L^{(\ell)} \rangle \cos^2 \theta_D$$

$$\frac{1}{\hat{\Gamma}^{(\ell)}} \frac{d\hat{\Gamma}^{(\ell)}}{d \chi} = \frac{1}{2\pi} + \frac{2}{3\pi}\langle S_3^{(\ell)} \rangle \cos 2\chi + \frac{2}{3\pi}\langle S_9^{(\ell)} \rangle \sin 2\chi$$

Introduction - Analysis Idea

- There are hints of LFU violation in Belle measurements - forward-backward assymetry of muon $\langle A_{FB}^{\mu} \rangle$ and electron $\langle A_{FB}^e \rangle$ are different within "4 σ "
- Idea : Make the one dimensional model independent template fit of single-differential distributions and measure 4 angular coefficients per lepton for $B \rightarrow D^* e \nu_e$ and $B \rightarrow D^* \mu \nu_\mu$ based on LHCb 2016-2018 Data

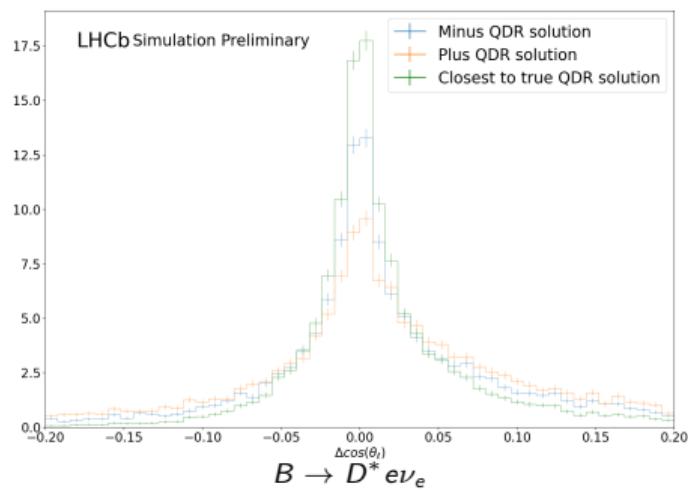
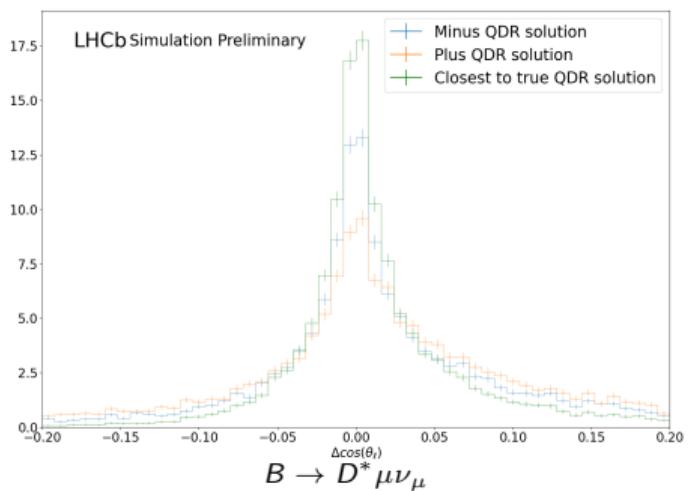
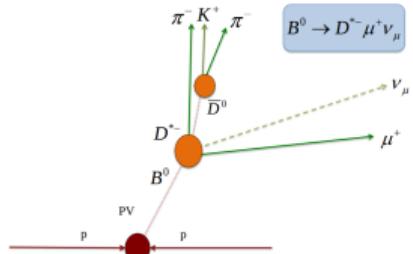
Problems:

- Charged particles, especially electrons, emit Bremsstrahlung photons when passing through the detector material degrading their momentum resolution and reconstruction efficiency
- The total and missing 4-momentum is unknown, so neutrinos cannot be reconstructed... Or they still can?

Reconstruction procedure

As the B^0 mass is well known, its momentum can be estimated using conservation of energy and momentum up to a two-fold ambiguity from its line of flight between the reconstructed primary and B^0 vertices by solving the quadratic equation (QDR):

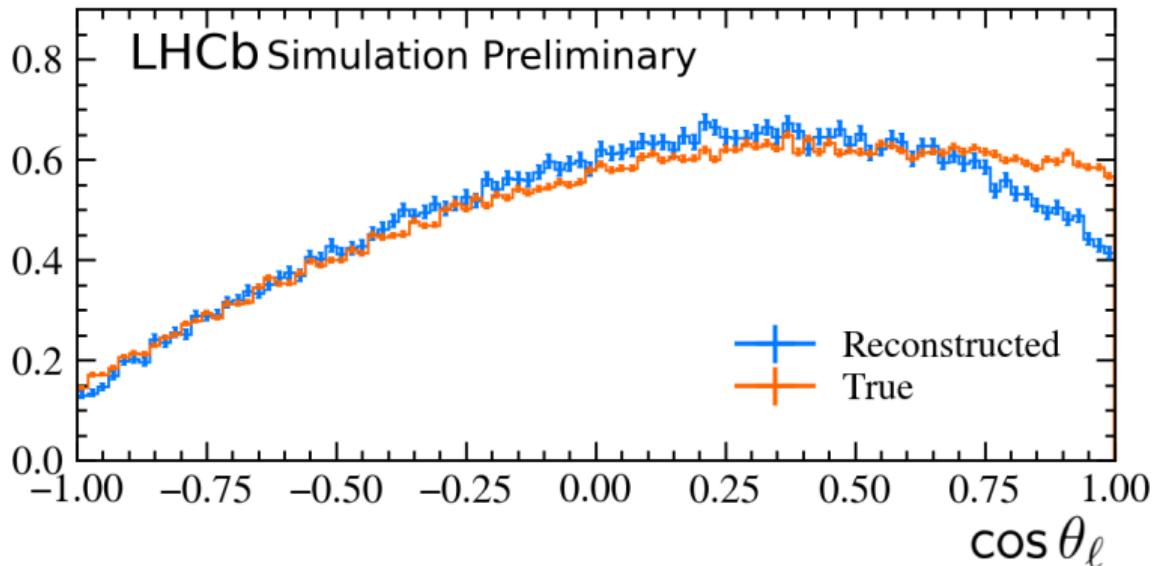
$$|\vec{p}_{B^0}| = \frac{(m_Y^2 + m_{B^0}^2)|\vec{p}_Y| \cos \theta_{B^0,Y} \pm E_Y \sqrt{(m_{B^0}^2 - m_Y^2)^2 - 4m_{B^0}^2 |\vec{p}_Y|^2 \sin^2 \theta_{B^0,Y}}}{2(E_Y^2 - |\vec{p}_Y|^2 \cos^2 \theta_{B^0,Y})}$$



Methodology

The following effects significantly change kinematic distributions:

- Non-linear detector efficiencies, such as the reconstruction resolution dependence on the transverse momentum of the lepton
- Neutrino reconstruction procedure



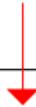
Template fit procedure

- To resolve these discrepancies between reconstructed and true angles the angular fit procedure described in [JHEP 11, \(2019\) 133](#) can be used
-

$$\frac{1}{\hat{\Gamma}^{(\ell)}} \frac{d\hat{\Gamma}^{(\ell)}}{d \cos \theta_\ell} = \frac{1}{2} \cdot [1 + \langle A_{FB}^{(\ell)} \rangle \cos \theta_\ell] + \frac{1}{4}(1 - 3\langle \tilde{F}_L^{(\ell)} \rangle) \frac{3 \cos^2 \theta_\ell - 1}{2}$$

$$\frac{1}{\hat{\Gamma}^{(\ell)}} \frac{d\hat{\Gamma}^{(\ell)}}{d \cos \theta_D} = \frac{3}{4}(1 - \langle F_L^{(\ell)} \rangle) \sin^2 \theta_D + \frac{3}{2}\langle F_L^{(\ell)} \rangle \cos^2 \theta_D$$

$$\frac{1}{\hat{\Gamma}^{(\ell)}} \frac{d\hat{\Gamma}^{(\ell)}}{d \chi} = \frac{1}{2\pi} \cdot [1 + \frac{2}{3\pi} \langle S_3^{(\ell)} \rangle \cos 2\chi + \frac{2}{3\pi} \langle S_9^{(\ell)} \rangle \sin 2\chi]$$



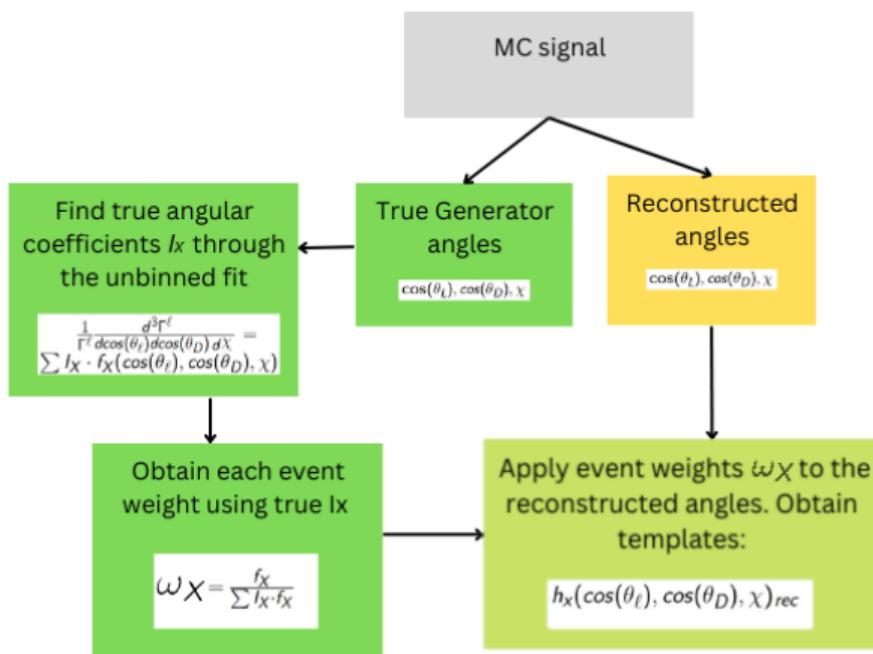
$$\frac{1}{\hat{\Gamma}^{(\ell)}} \frac{d\hat{\Gamma}^{(\ell)}}{d \cos \theta_\ell} = (\frac{1}{2} - \frac{1}{8}(1 - 3)\langle \tilde{F}_L^{(\ell)} \rangle) h_{const, \theta_\ell} + \langle A_{FB}^{(\ell)} \rangle h_{\cos \theta_\ell} + \frac{3}{8}(1 - 3\langle \tilde{F}_L^{(\ell)} \rangle) h_{\cos^2 \theta_\ell}$$

$$\frac{1}{\hat{\Gamma}^{(\ell)}} \frac{d\hat{\Gamma}^{(\ell)}}{d \cos \theta_D} = \frac{3}{4}((1 - \langle F_L^{(\ell)} \rangle) h_{const, \theta_D} + (3\langle F_L^{(\ell)} \rangle - 1) h_{\cos^2 \theta_D}$$

$$\frac{1}{\hat{\Gamma}^{(\ell)}} \frac{d\hat{\Gamma}^{(\ell)}}{d \chi} = (\frac{1}{2\pi} - \frac{2}{3\pi} \langle S_3^{(\ell)} \rangle) h_{const, \chi} + \frac{2}{3\pi} \langle S_3^{(\ell)} \rangle h_{(1+\cos 2\chi)}$$

Template fit procedure

- Truth-level angular coefficients are obtained from True MC distributions with a specific description of the SM for these types of decays

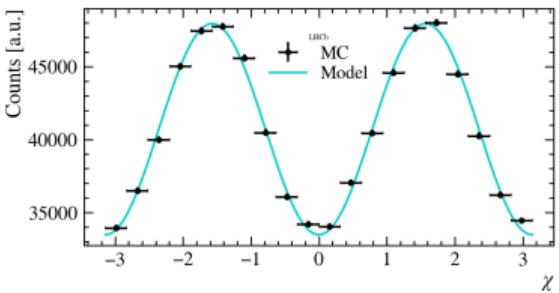
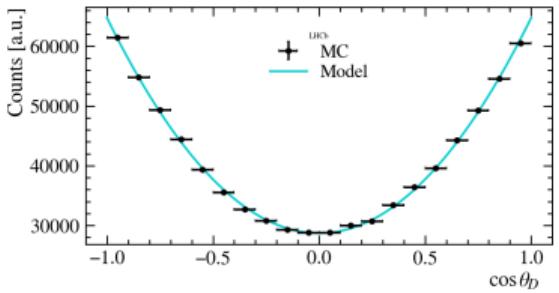
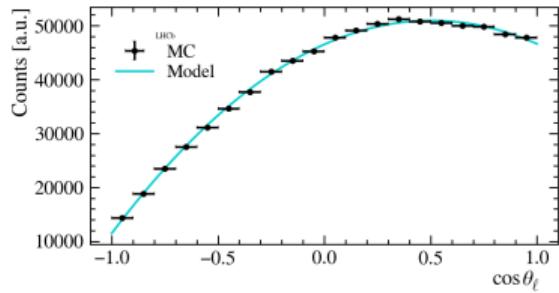
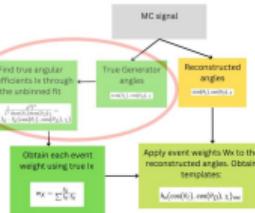


The préfou receipe:

- 1 Mix all the ingredients and knead the dough in a bread machine or by hand.
- 2 Let the dough rise for about an hour.
- 3 Arrange the dough in the shape of a baguette on baking paper. Bake for 20 minutes at 180C.
- 4 Take the bread out and let it cool, then cut it in half lengthwise.
- 5 Top the bread with the butter and cut it into small slices.
- 6 Return to the oven for 10 to 15 minutes until the bread is golden brown.
- 7 Serve while it's hot

True MC distributions $B \rightarrow D^* \mu \nu_\mu, \cos \theta_D, \cos \theta_\ell, \chi$

Generator SM:



(a) $\cos \theta_\ell$

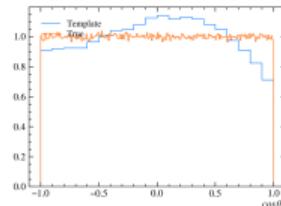
(b) $\cos \theta_D$

(c) χ

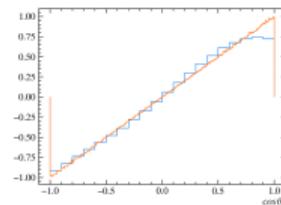
Fit of True MC distributions

Template fit $\cos\theta_\ell$, $B \rightarrow D^*\mu\nu_\mu$

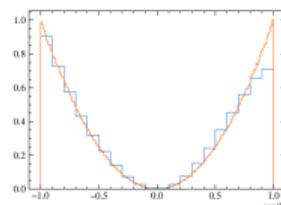
Template for a constant



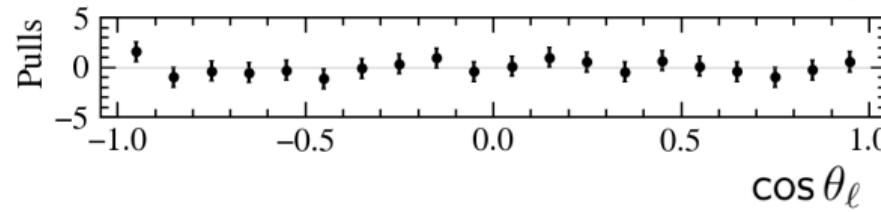
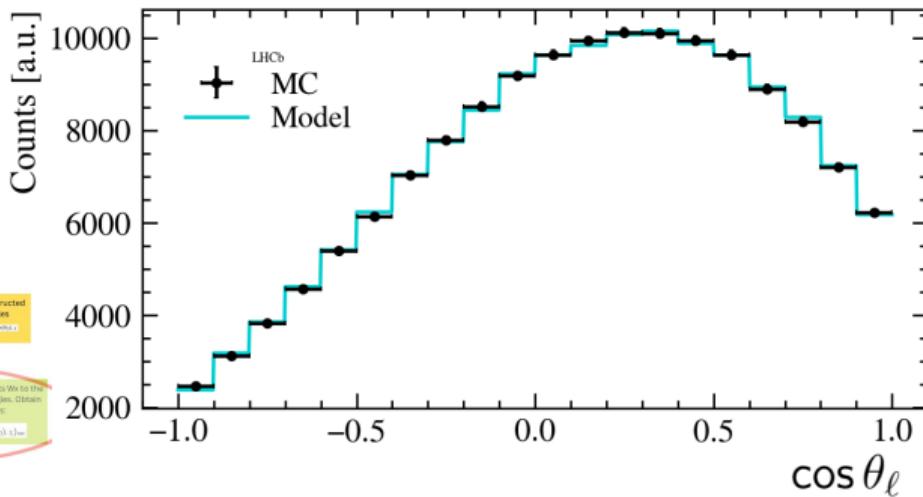
Template for a $\cos\theta_\ell$



Template for a $\cos^2\theta_\ell$



Template fit of MC with SM formfactor parametrization



New Physics contributions

The dependence of angular observables from **Wilson Coefficients**

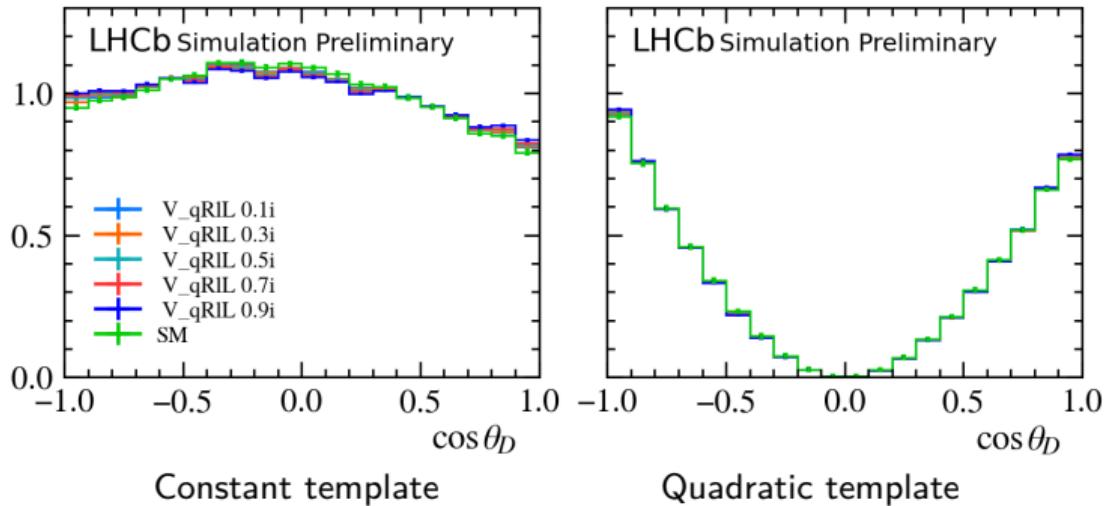
$$C_V^{\ell\ell'} = C_{V_R}^{\ell\ell'} + C_{V_L}^{\ell\ell'}, \quad C_A^{\ell\ell'} = C_{V_R}^{\ell\ell'} - C_{V_L}^{\ell\ell'}, \quad C_P^{\ell\ell'} = C_{S_R}^{\ell\ell'} - C_{S_L}^{\ell\ell'},$$

Observable	$ C_A ^2$	$ C_V ^2$	$ C_P ^2$	$ C_T ^2$	$\text{Re}(C_A C_V^*)$	$\text{Re}(C_A C_P^*)$	$\text{Re}(C_A C_T^*)$	$\text{Re}(C_V C_P^*)$	$\text{Re}(C_V C_T^*)$	$\text{Re}(C_P C_T^*)$
$d\Gamma/dq^2$	✓	✓	✓	✓	—	(m)	(m)	—	(m)	—
$\text{num}(A_{\text{FB}})$	(m^2)	—	—	(m^2)	✓	(m)	(m)	—	(m)	✓
$\text{num}(F_L)$	✓	—	✓	✓	—	(m)	(m)	—	—	—
$\text{num}(F_L-1/3)$	✓	✓	✓	✓	—	(m)	(m)	—	(m)	—
$\text{num}(\tilde{F}_L)$	✓	(m^2)	✓	✓	—	(m)	(m)	—	(m)	—
$\text{num}(\tilde{F}_L-1/3)$	✓	✓	—	✓	—	—	—	—	—	—
$\text{num}(S_3)$	✓	✓	—	✓	—	—	—	—	—	—

The table is taken from [C. Bobeth, et al](#)

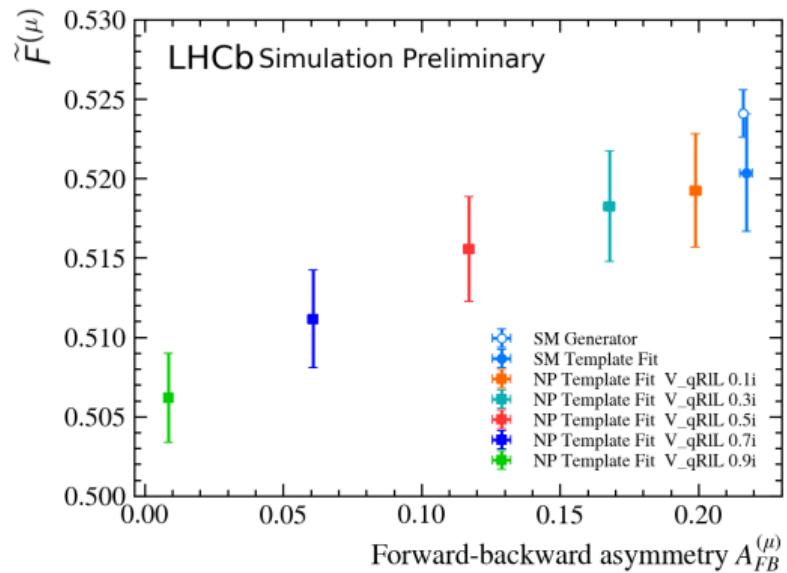
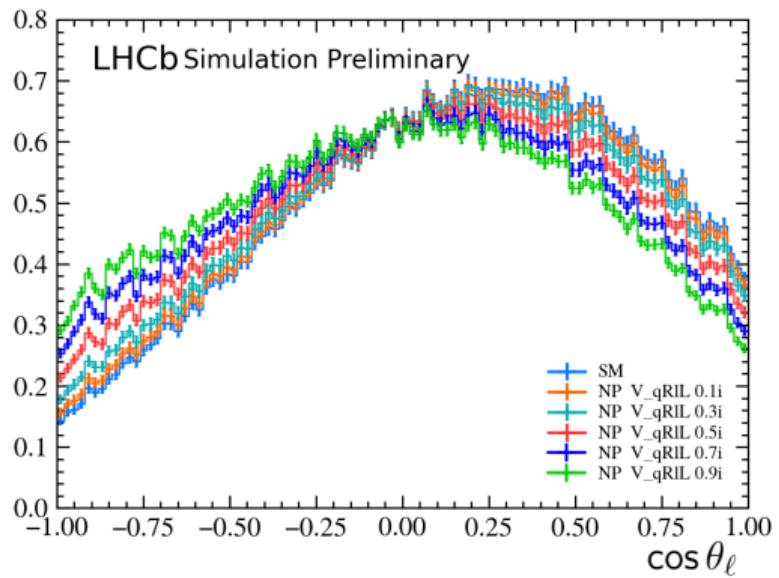
Template fit model independence

- The simplest way to prove the model Independence of 1D template fit is to remake templates with different NP contributions, i.e. to reweigh reconstructed and generator angles. If the approach is indeed model-independent, all templates are supposed to be the same within statistical uncertainty

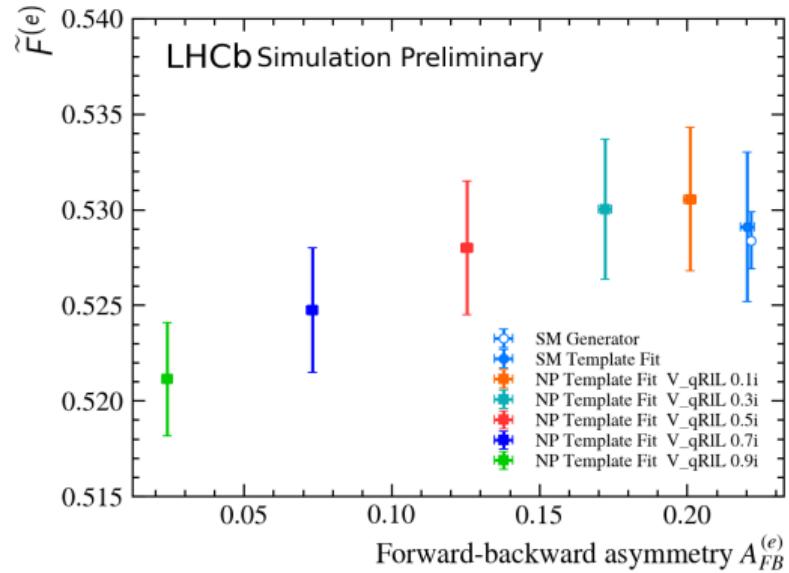
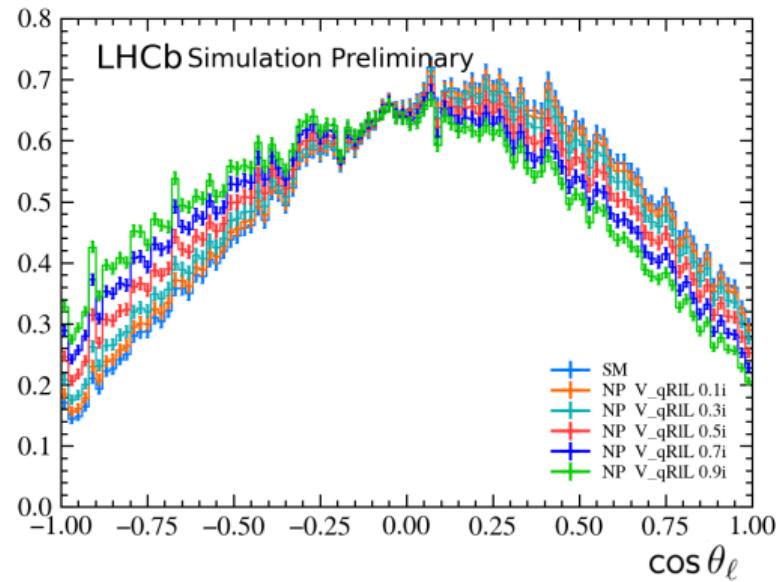


NP template fit $\cos \theta_\ell$, $B \rightarrow D^* \mu \nu_\mu$

- Example of sensitivity to NP of angular observable - $\langle A_{FB}^\mu \rangle$ for vector contribution with right helicity of b quark V_{qRIL} with different amplitudes:

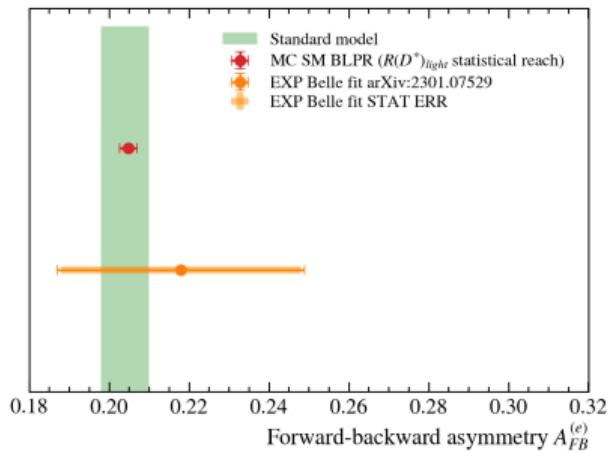
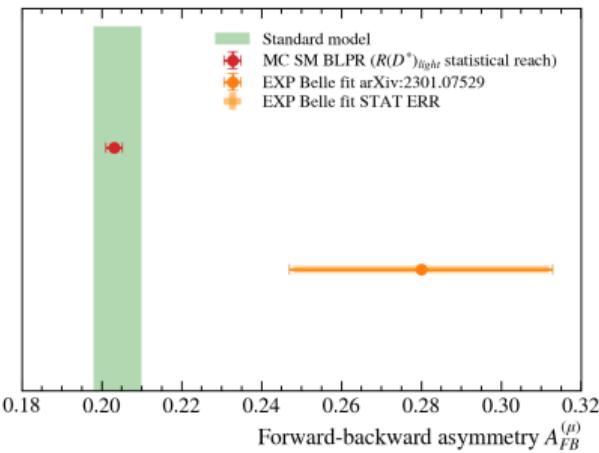


NP template fit $\cos \theta_\ell$, $B \rightarrow D^* e \nu_e$



Results on MC

- Expected statistical uncertainty is estimated using the same selection criteria as in the ongoing analysis for $R(D^*)_{e/\mu}$, which is currently in the advanced state
- All of the results obtained on the [simulation](#)
- The same plots were obtained for all four angular coefficients



- Competitive measurement to Belle II will be possible
- But we expect to be dominated by systematic uncertainties ($B \rightarrow D^{**} \ell \nu_\ell$ decay model, MC correction)

Status and outlook

- Precision measurements of flavor physics observables provide a great test of the SM
- LHCb allows to perform these high precision measurements
- Analysis of 1D angular distributions for $B \rightarrow D^* \mu \nu_\mu$ and $B \rightarrow D^* e \nu_e$ was conducted on MC samples
- This analysis is sensitive to New Physics (NP) couplings
- Expected statistical uncertainty for angular observables is estimated
- Measurements can be competitive with Belle
- It will be the first-ever angular analysis at LHCb for $B \rightarrow D^* e \nu_\ell$

Plans:

- Finalize data fit
- Obtain blinded results for data
- Evaluate systematic uncertainties contributions

Thank you!

Introduction - Analysis Idea

- One dimensional fit was performed based on Belle data, the discrepancy with SM was observed in lepton Forward-Backward asymmetry :

- ① Belle's 2018 untagged data ([4 angular coefficients per lepton were measured](#)) were analyzed by C. Bobeth et. al.

$$\langle A_{FB}^{\mu} \rangle = 0.2300 \pm 0.0059 \text{ (} 2.4\sigma \text{ higher than SM)}$$

$$\langle F_L^{\mu} \rangle = 0.5271 \pm 0.0046$$

$$\langle A_{FB}^e \rangle = 0.1951 \pm 0.0069 \text{ (} 1.6\sigma \text{ lower than SM)}$$

$$\langle F_L^e \rangle = 0.5336 \pm 0.0045$$

- ② Belle's 2022 hadronic tag-side reconstruction ([2 angular coefficients per lepton was measured](#)) analysis [arXiv:2301.07529](#) :

$$\langle A_{FB}^{\mu} \rangle = 0.280 \pm 0.032(\text{stat.}) \pm 0.009(\text{syst.}) \text{ (} 2.3\sigma \text{ higher than SM)}$$

$$\langle F_L^{\mu} \rangle = 0.503 \pm 0.023(\text{stat.}) \pm 0.007(\text{syst.})$$

$$\langle A_{FB}^e \rangle = 0.218 \pm 0.030(\text{stat.}) \pm 0.008(\text{syst.})$$

$$\langle F_L^e \rangle = 0.471 \pm 0.024(\text{stat.}) \pm 0.007(\text{syst.}) \text{ (} 2.6\sigma \text{ lower than SM)}$$

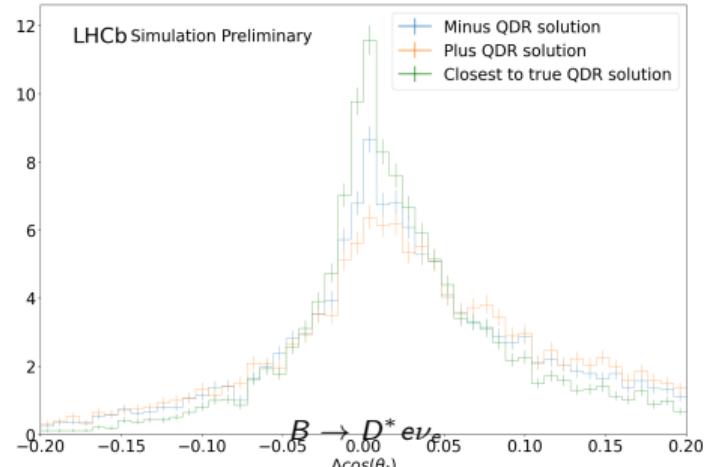
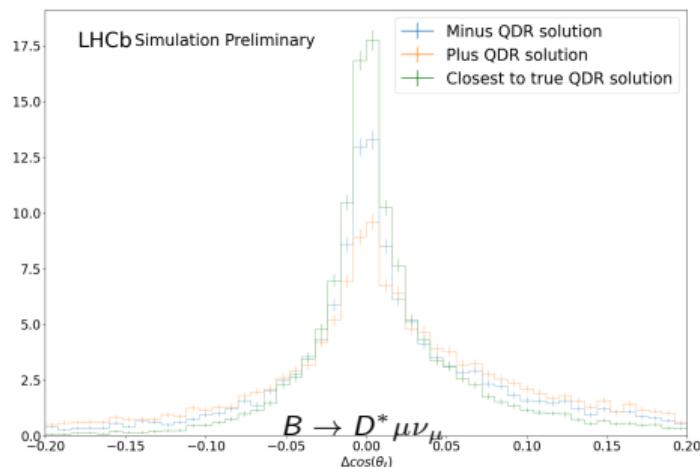
- Idea : Make the one dimensional model independent template fit of single-differential distributions and measure 4 angular coefficients per lepton for $B \rightarrow D^* e \nu_e$ and $B \rightarrow D^* \mu \nu_\mu$ based on LHCb Run 2 Data

Backup Slides

Solution of quadratic equation (QDR)

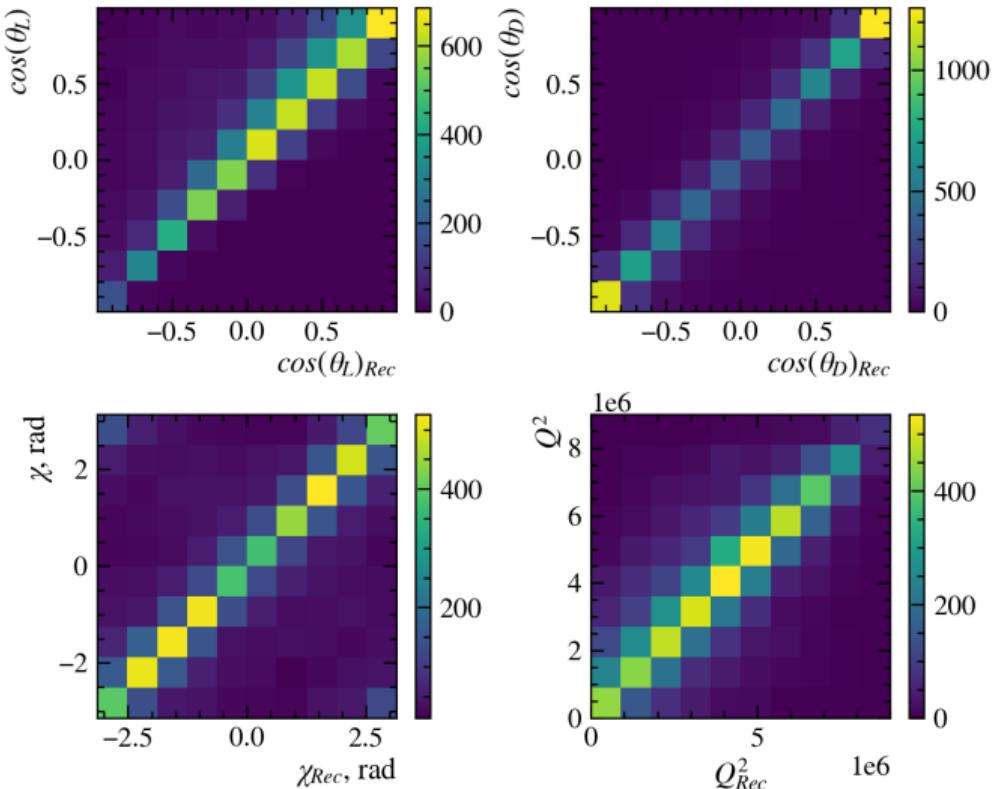
As the B^0 mass is well known, its momentum can be estimated up to a two-fold ambiguity from its line of flight between the reconstructed primary and B^0 vertices:

$$|\vec{p}_{B^0}| = \frac{(m_Y^2 + m_{B^0}^2)|\vec{p}_Y| \cos \theta_{B^0,Y} \pm E_Y \sqrt{(m_{B^0}^2 - m_Y^2)^2 - 4m_{B^0}^2 |\vec{p}_Y|^2 \sin^2 \theta_{B^0,Y}}}{2(E_Y^2 - |\vec{p}_Y|^2 \cos^2 \theta_{B^0,Y})}$$

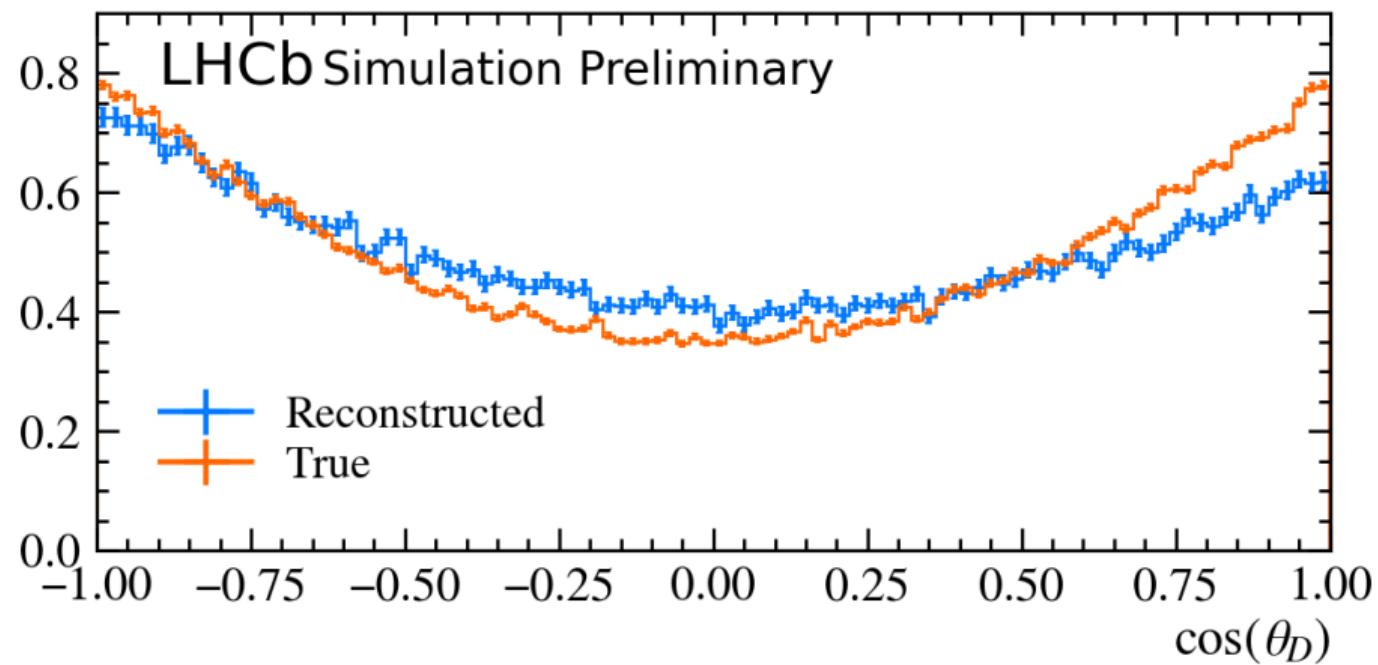


Since the "minus" solution has better resolution, it was chosen in all cases to resolve ambiguity

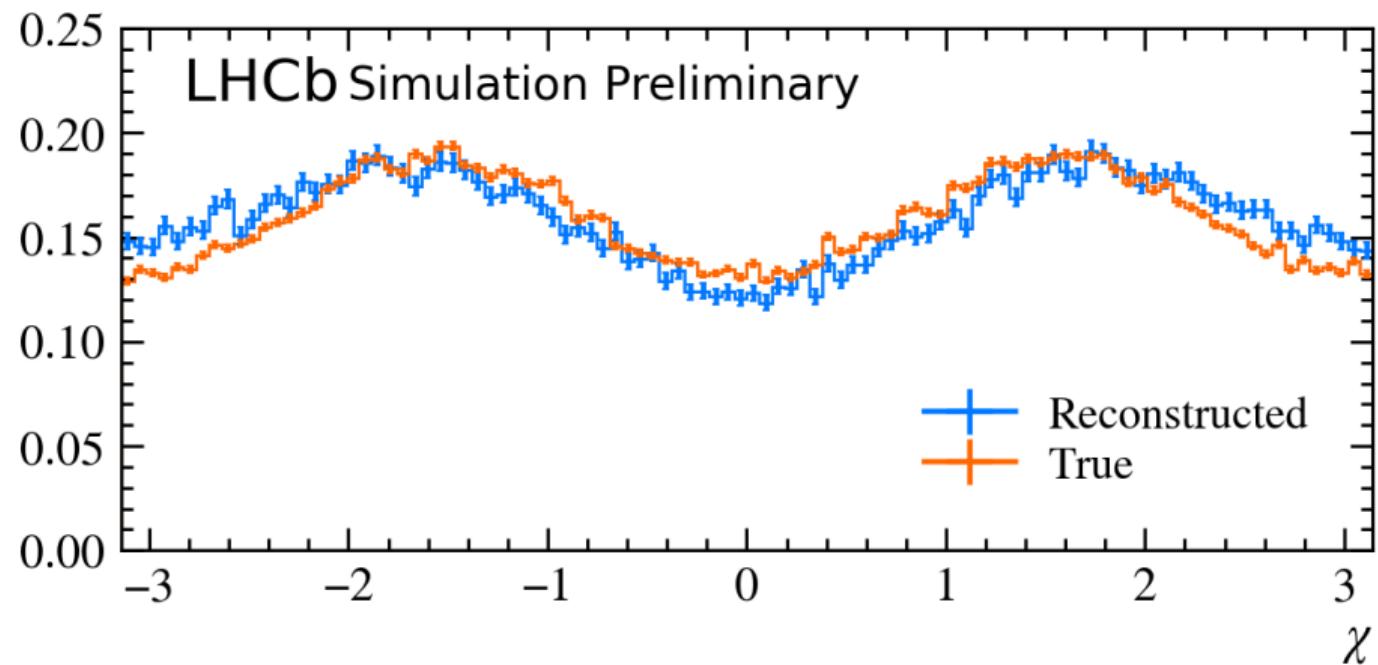
$B \rightarrow D^* \mu \nu_\mu$ correlations, QDR



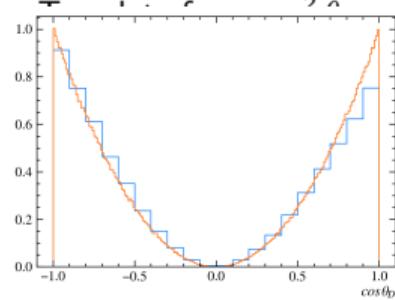
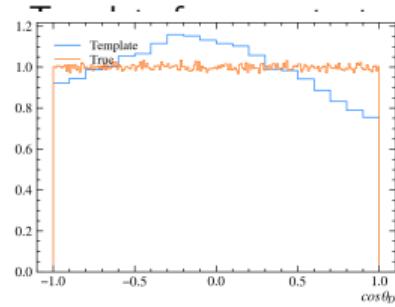
Reconstructed and True angles



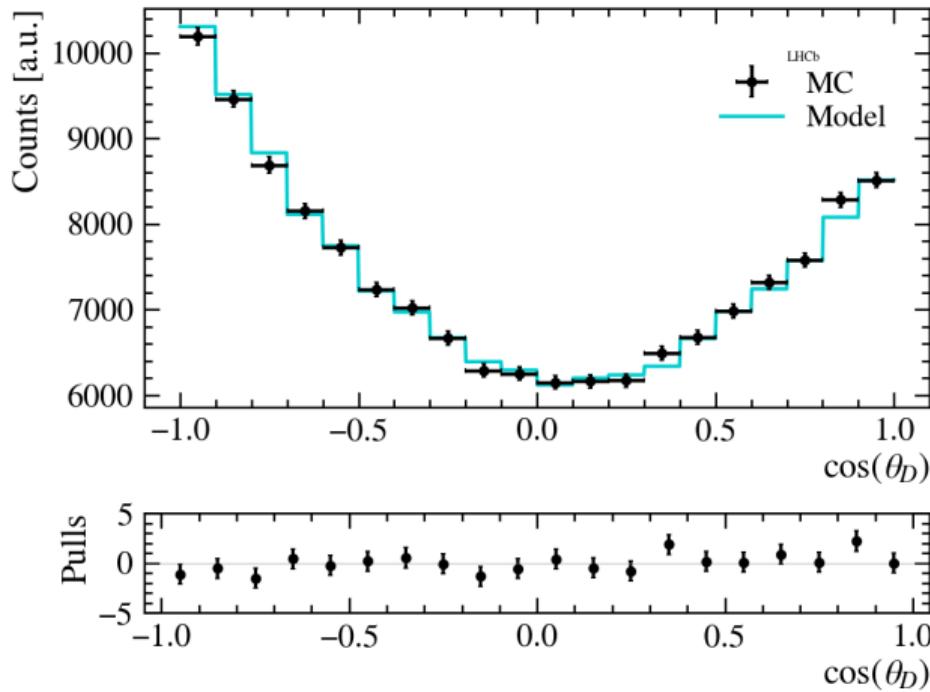
Reconstructed and True angles



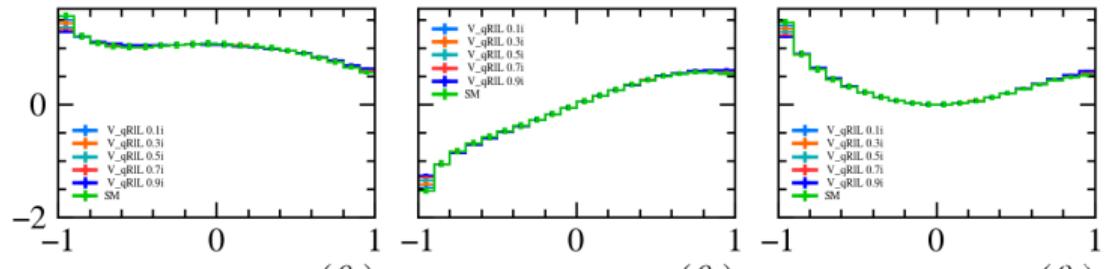
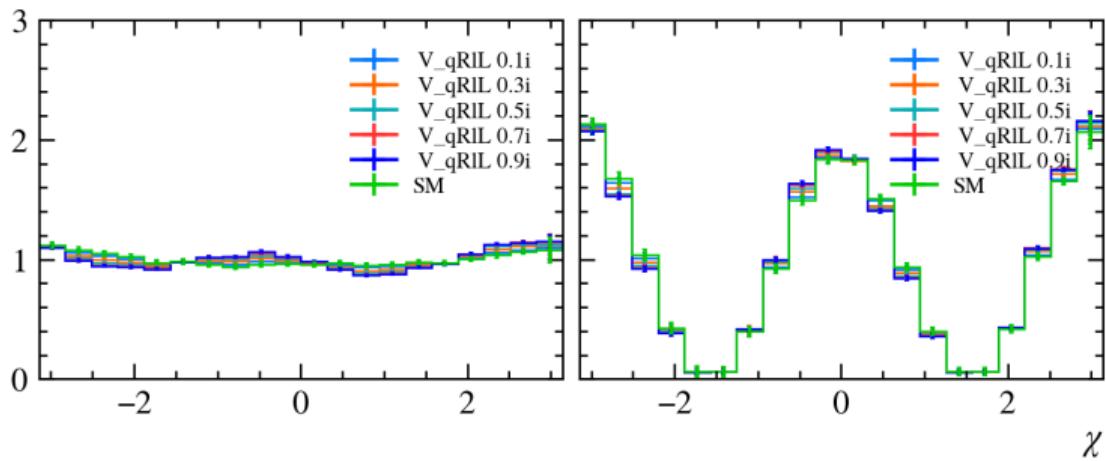
Template fit $\cos\theta_D$, $B \rightarrow D^*\mu\nu_\mu$



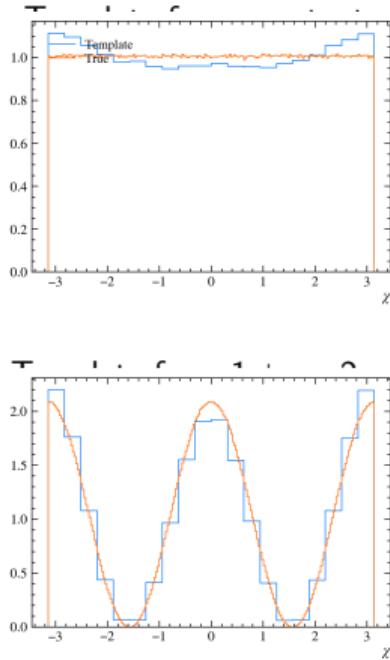
Template fit of MC with SM formfactor parametrization



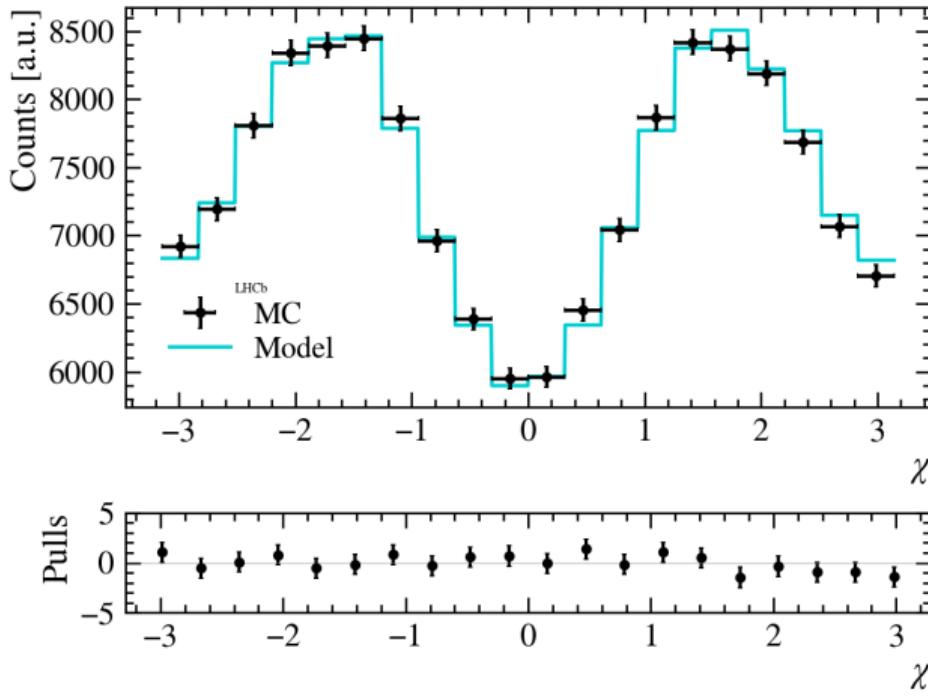
Template fit model independence



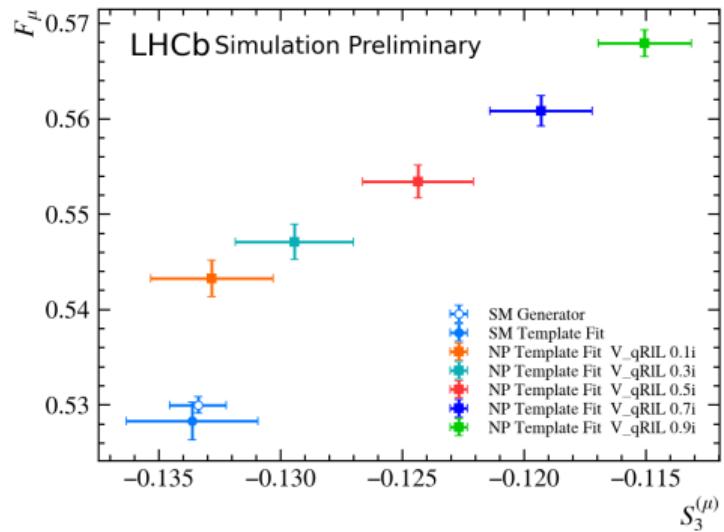
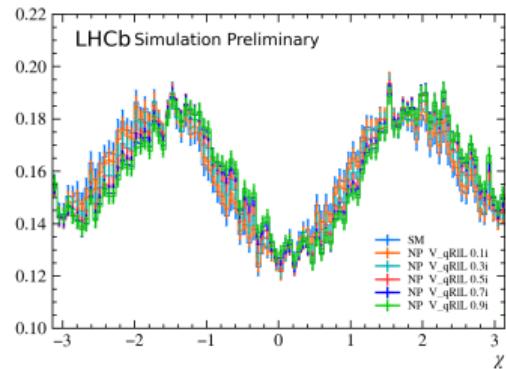
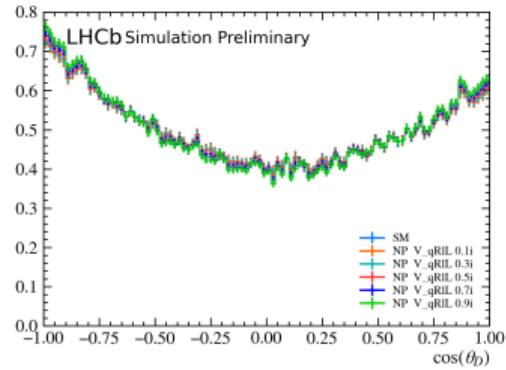
Template fit $\cos\chi$, $B \rightarrow D^*\mu\nu_\mu$



Template fit of MC with SM formfactor parametrization



New Physics template fit $\cos \theta_D, \chi, B \rightarrow D^* \mu \nu_\mu$



NP template fit $\cos \theta_D$, χ , $B \rightarrow D^* e \nu_e$

