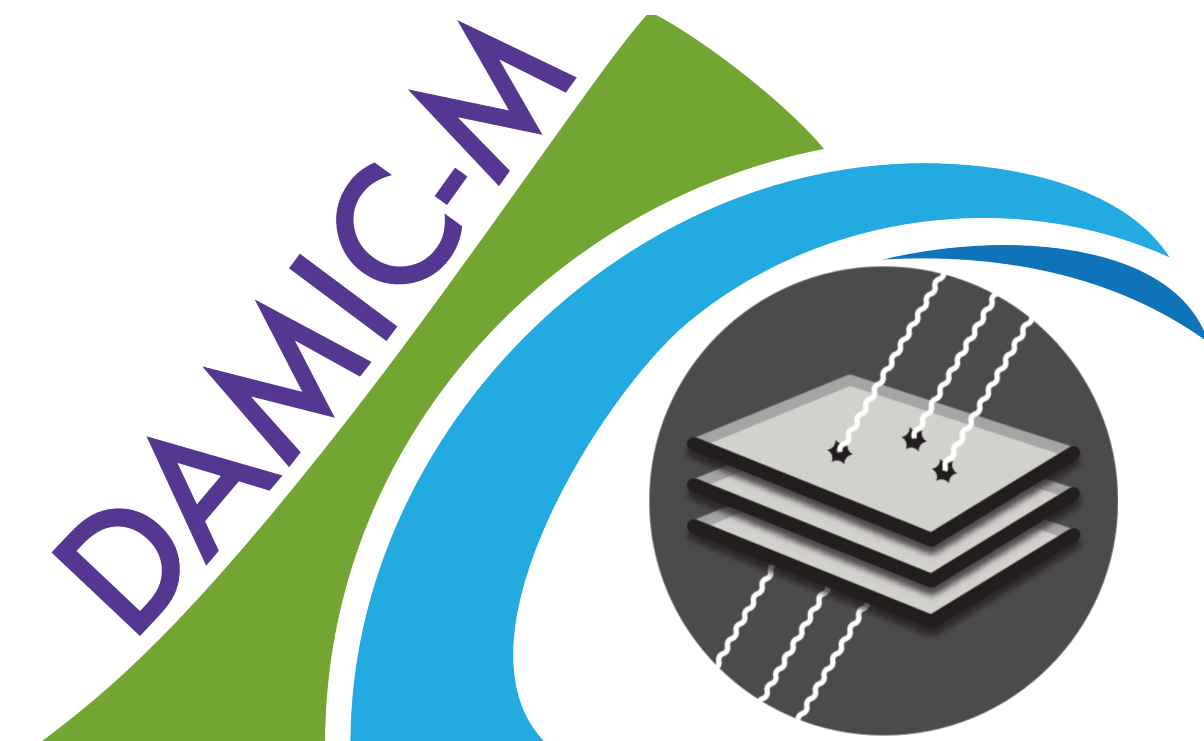# The DAMIC-M Experiment: First Results of the Low Background Chamber prototype

Jean-Philippe Zopounidis on behalf of the DAMIC-M collaboration
Sorbonne University, LPNHE

GDR Deep Underground Physics 2023
**06/21/2023**

DAMIC @ Modane

Location

Modane Underground Laboratory (LSM), France

CCD technology

- Use of skipper-readout silicon CCDs with thickness of 675 um and 9 Mpixels

- 200 CCDs in order to achieve a 700 gr target mass

- Custom electronics for the transfer of charge and a low noise read-out
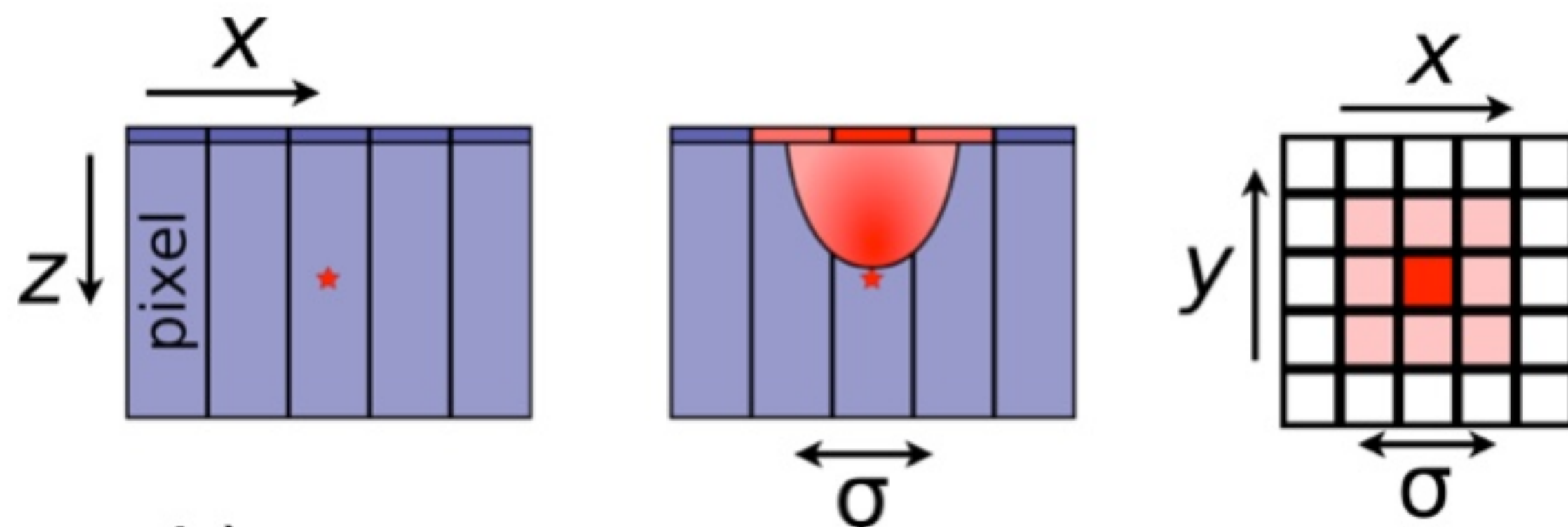
Main scientific objectives

- Sub-electron charge resolution together with low-noise readout electronics

- Operate ionization detector with energy threshold 5-10 eV

- Achieve a background rate of < 1 d.r.u.

- Search for light-mass DM candidates of the hidden sector or light WIMPs

cryostat

cold finger

CCD array

CCD module array
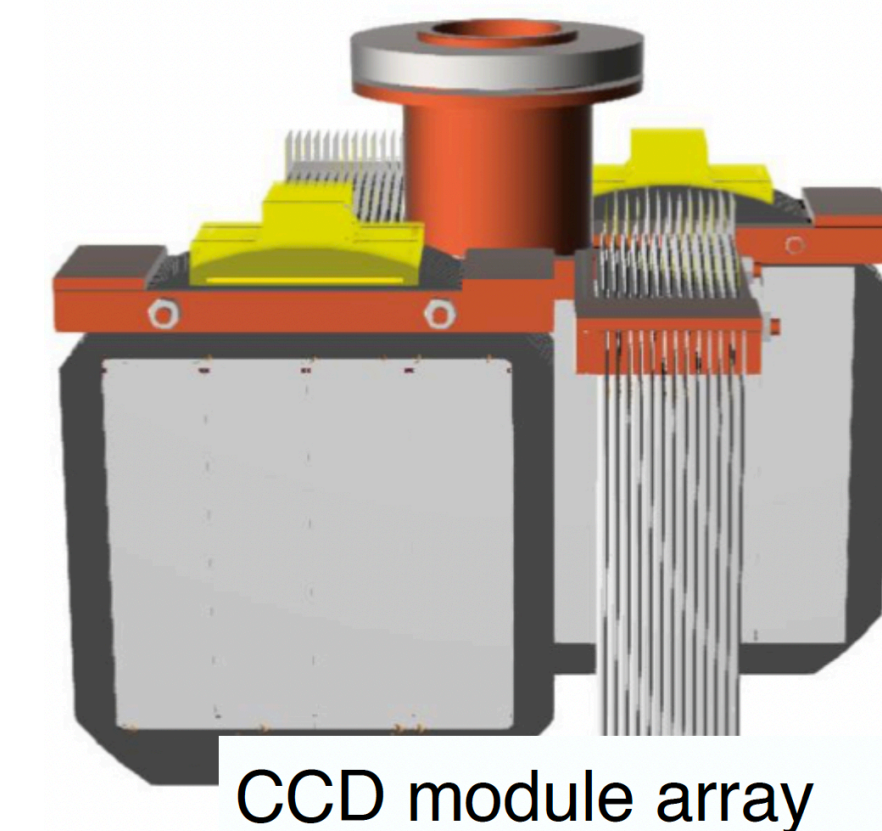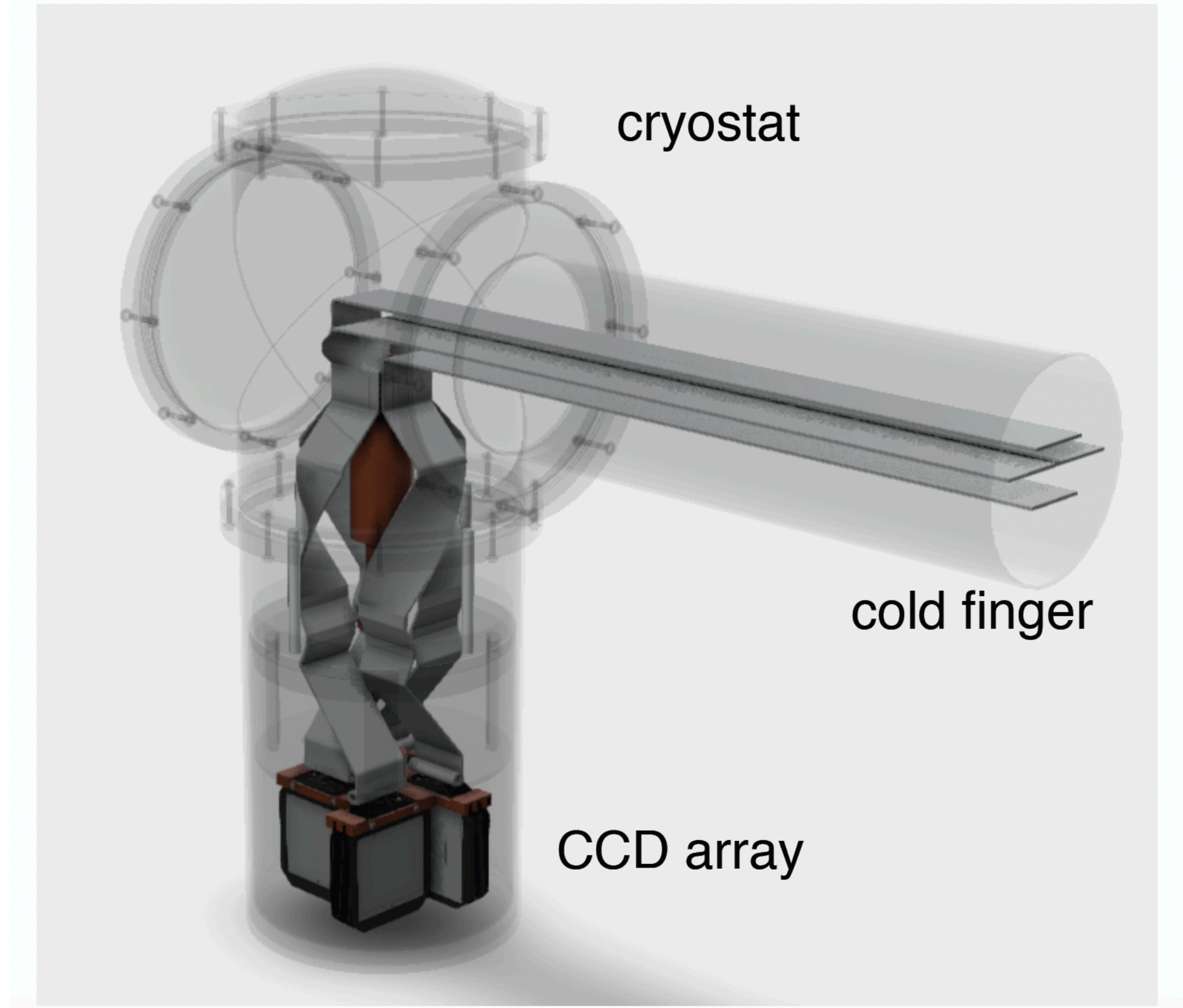
Detection principle
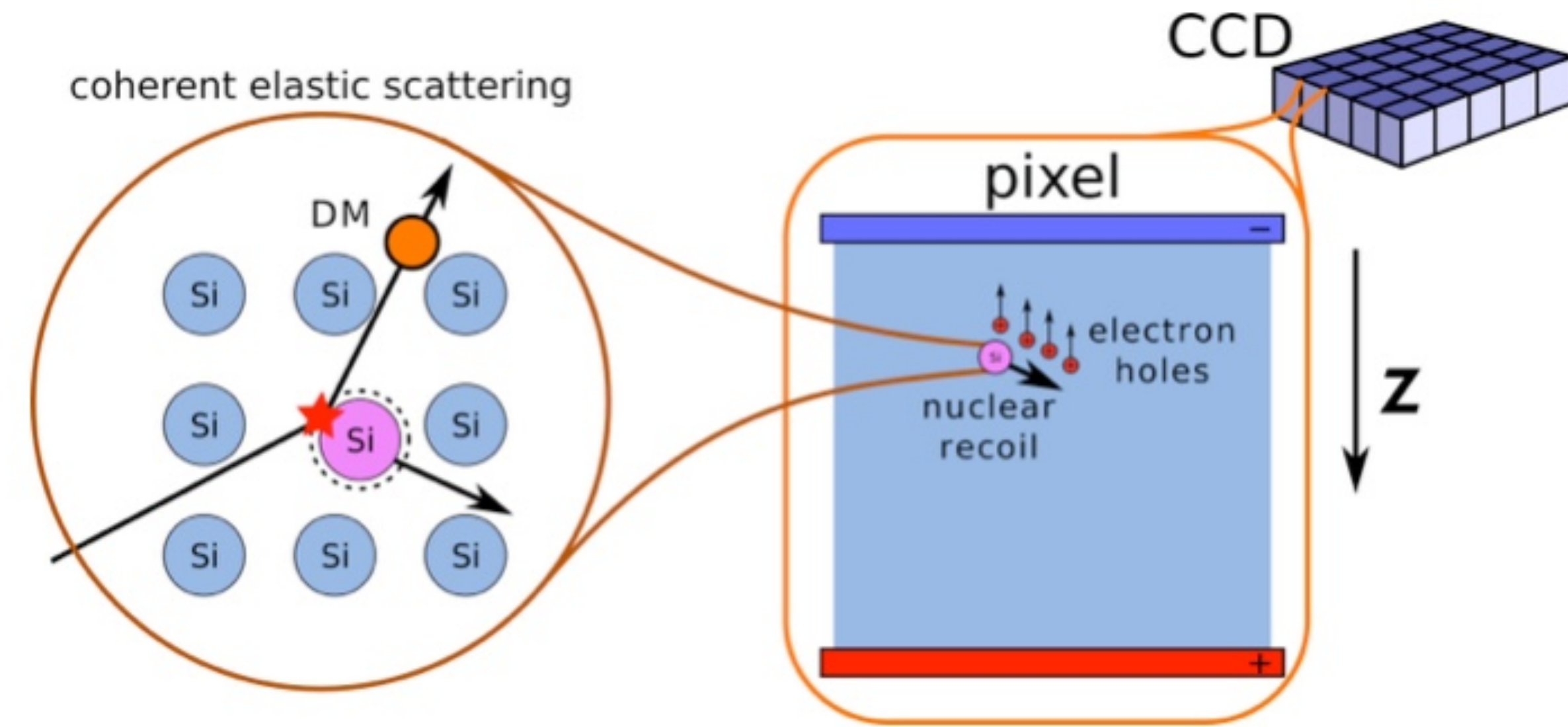
- A DM particle can scatter off a nucleus or a valence electron and create a point-like ionization event

- Charge will be drifted to the pixel array under a voltage bias

- Lateral spread of the charge cloud due to thermal diffusion

- The lateral spread is proportional to the drift time (depth of the interaction)

| 3D reconstruction of the interaction location |
| Identification of particle type via cluster pattern |

3

Charge transfer and skipper readout

- After exposure of the active target and charge generations the readout take place

- A series of 3 voltage clocks create potential walls that are used to move the charges through the pixels

- In a vertical transfer one row of pixels is moved one row closer to the horizontal register

- The charge in the horizontal register is moved pixel-by-pixel to a readout amplifier

- The charge fall into the SW gate and then to the SN where it is measured

4
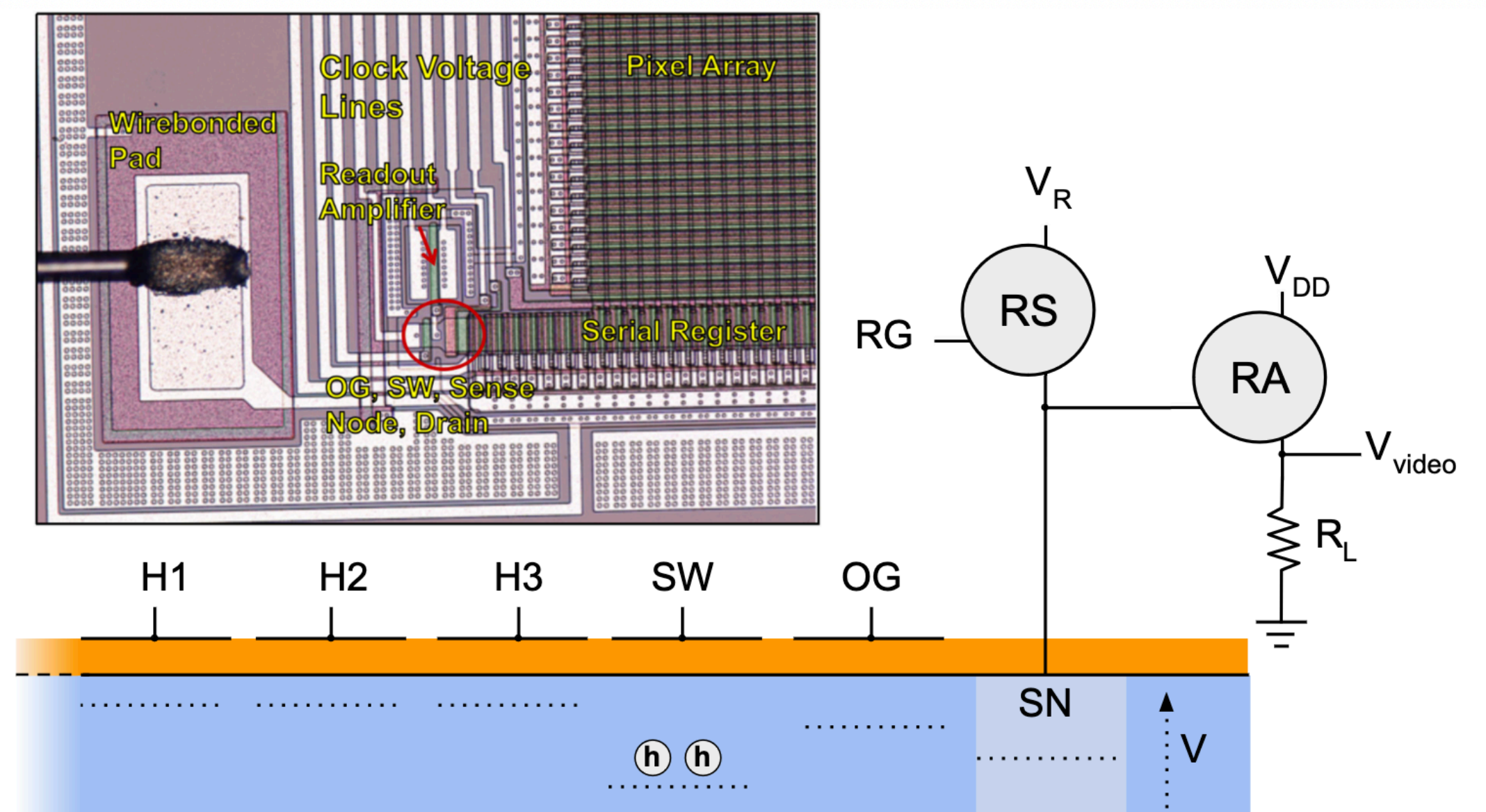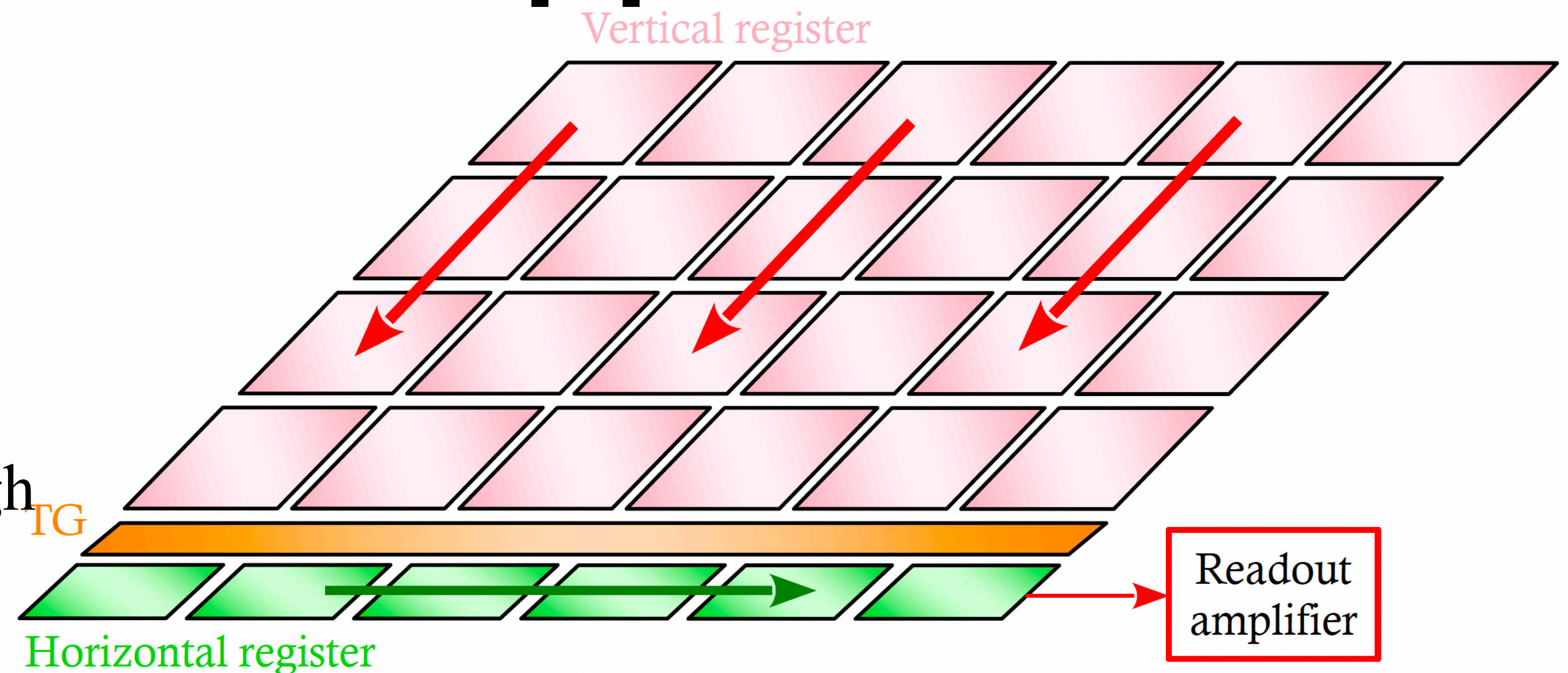
Charge transfer and skipper readout

• Using a floating gate as SN and replacing the bias VOG with a clock, permits a multiple non-destructive measurement of the charge packet
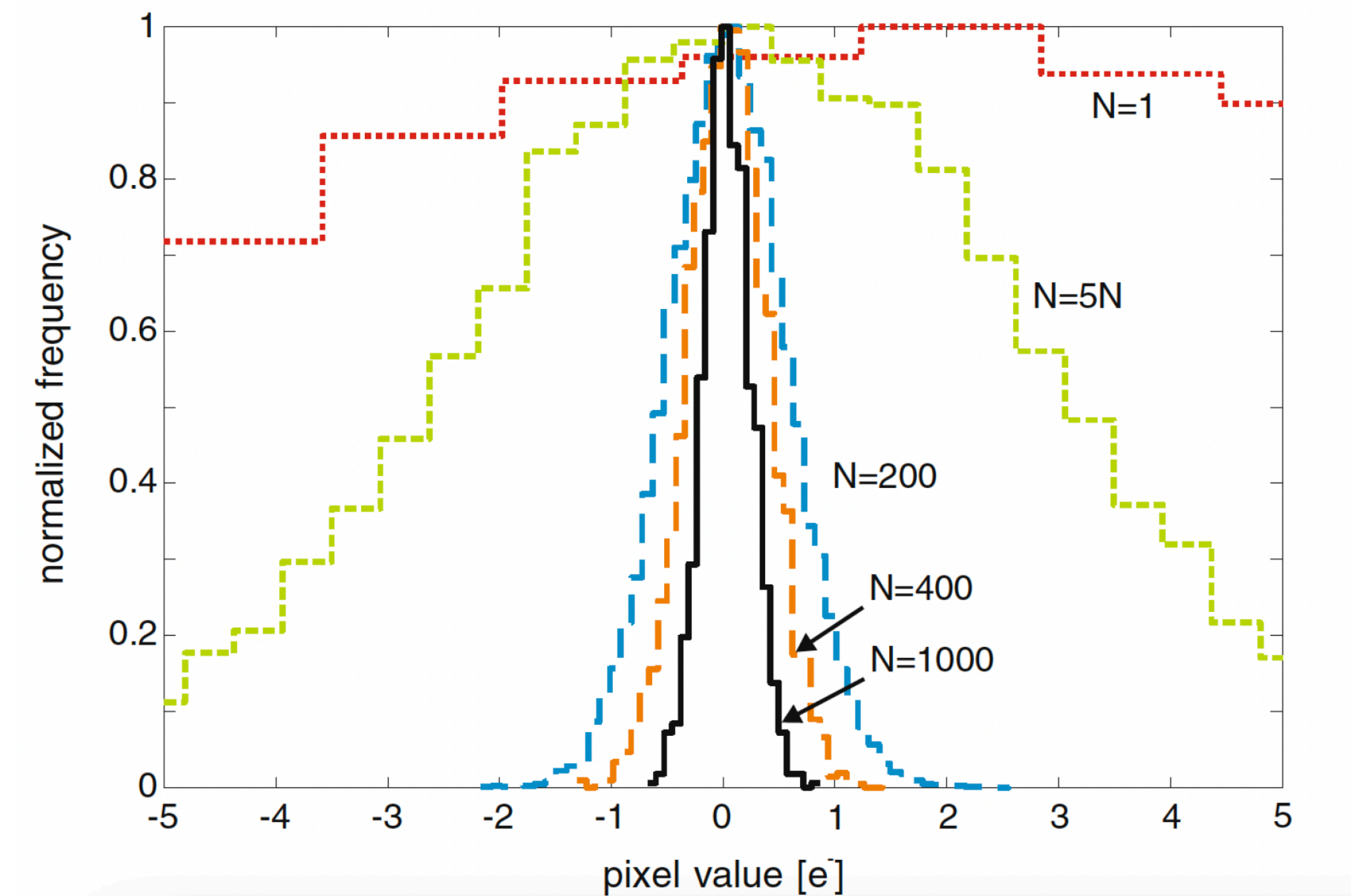
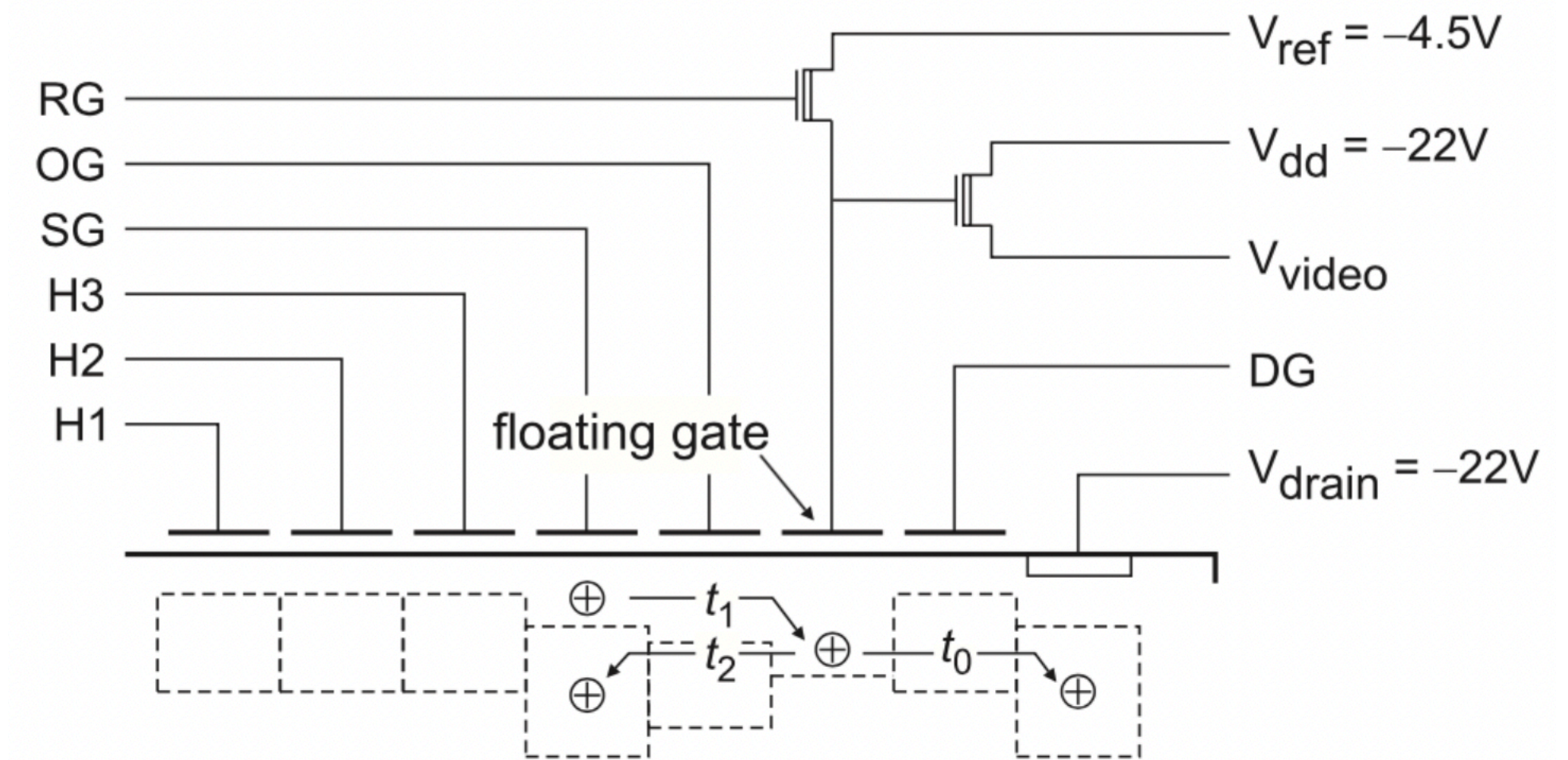• The measurement error σ will decrease as

$$\sim 1/\sqrt{N_{skip}}$$

• Thus the 1/f amplifier low frequency noise is now subdominant

• For a large number of N_{skip} the resolution reaches sub-electron values

• But $t_{readout} \sim N_{skip}$

Effort for Background mitigation

Radiogenic and cosmogenic background limit sensitivity for WIMP search (nuclear recoils of $\sim$ keV$_{ee}$ energy deposits)

- Use of Si wafers with low cosmogenic activation and limiting time above ground (fabrication, transport and storage)

- Careful material selection

- Oxygen-free cooper box for CCDs and further reduction by multiple layers of Pb Shield

- Analysis techniques for efficient identification of particle type from cluster shape
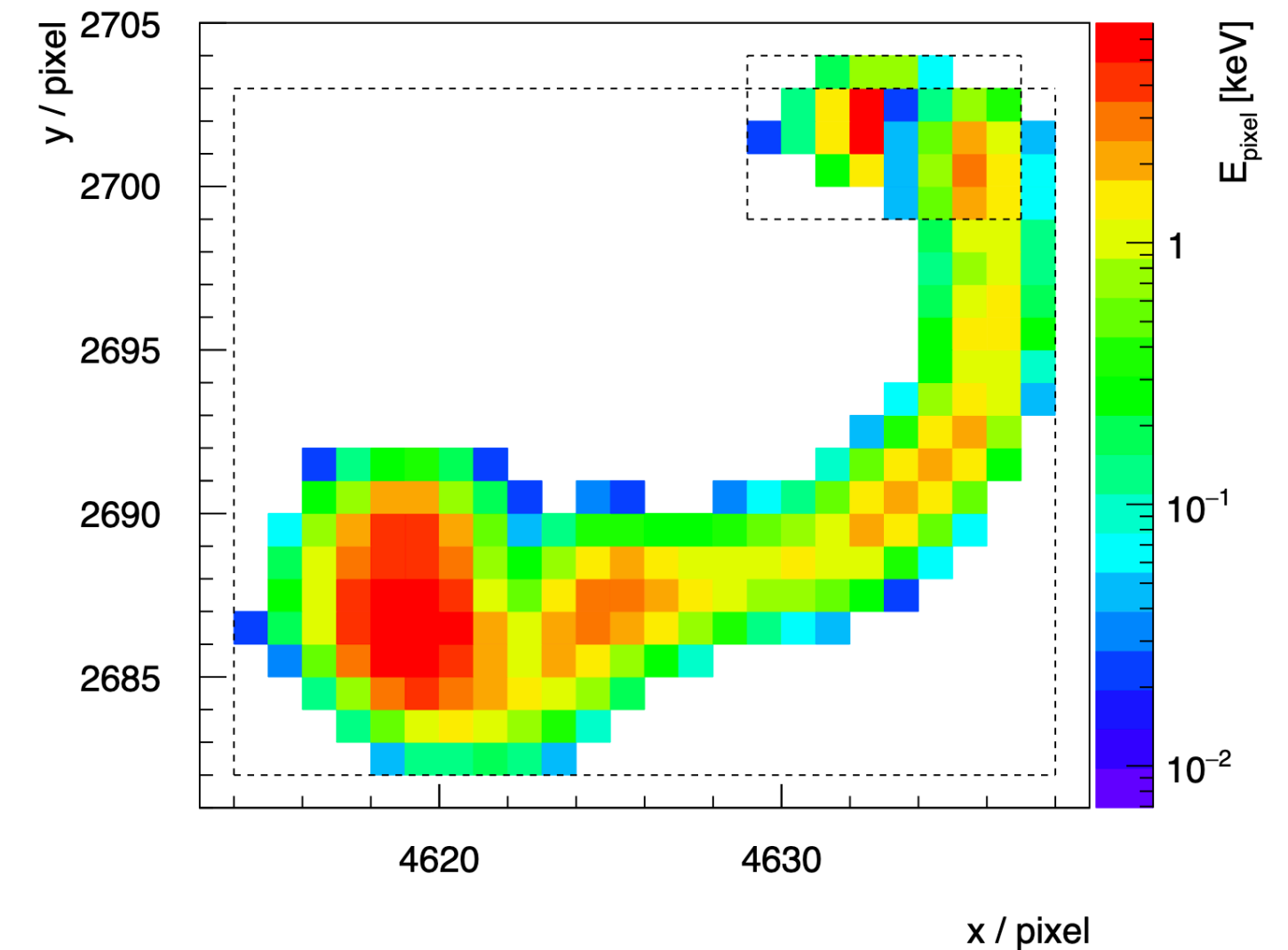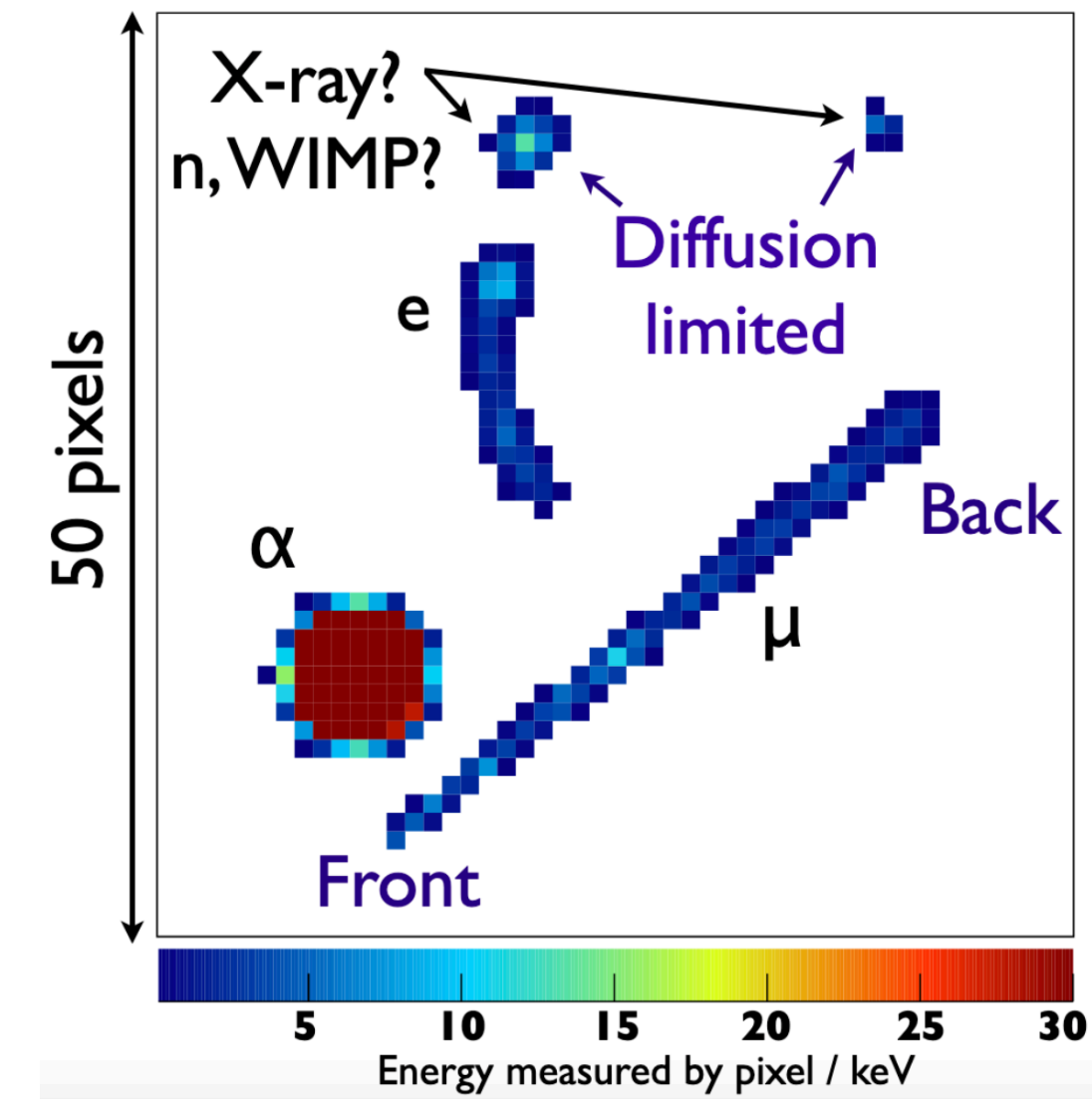
Low Background Chamber at LSM

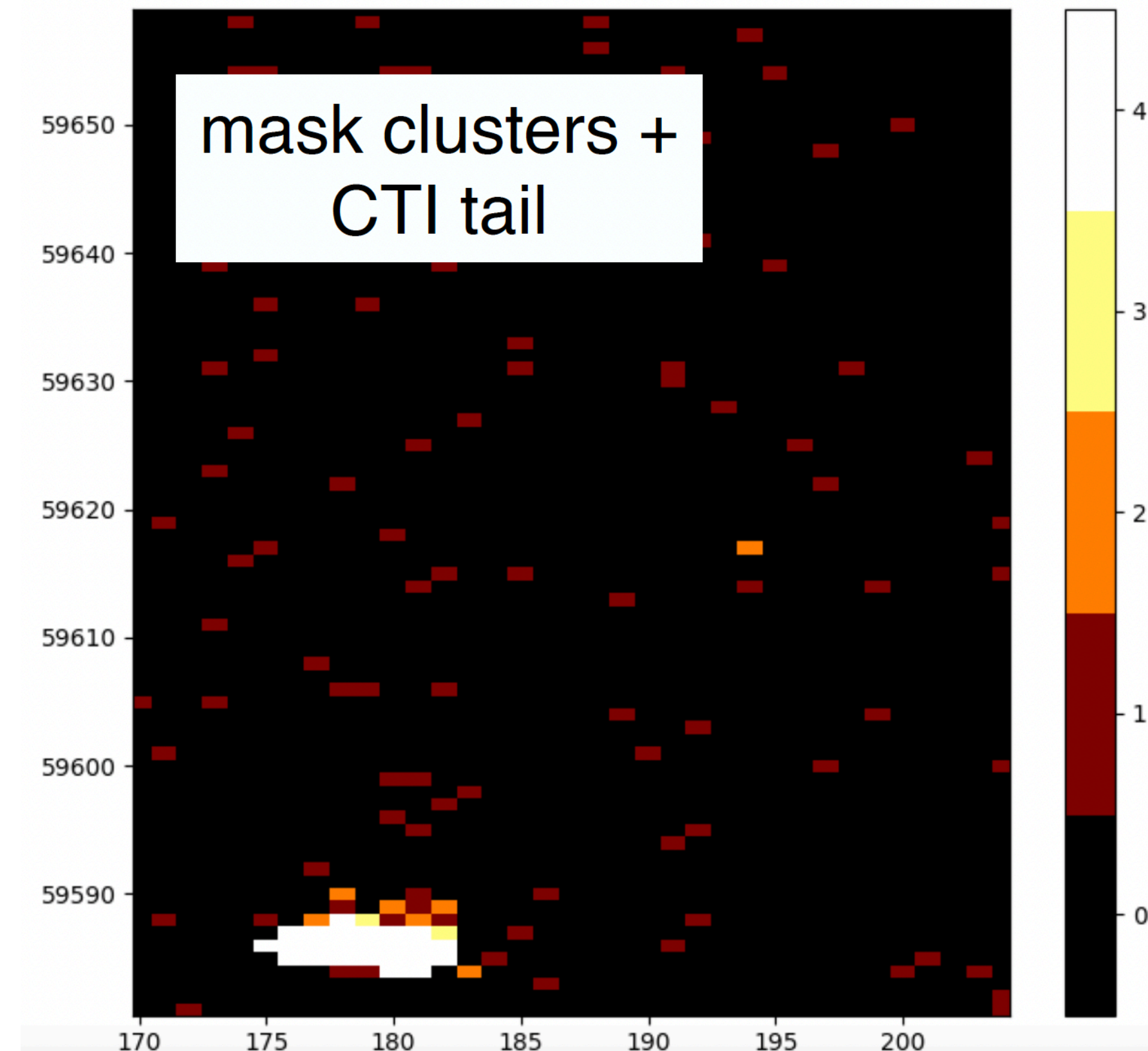# Low Background Chamber at LSM

LBC experimental configuration

- $6k \times 4k$ pixel skipper CCDs ($\times 2$)

- Total mass of active target ~ 18gr

- Background reduction with layered polyethylene+lead shielding, innermost layer of ancient lead

- Readout is done with the commercially available Astronomical Research Cameras electronics

Advancements

- Validation of detector components and subsystems (DAQ, and instrumentation, slow control monitoring)

- Reduction of the high-energy background levels to few d.r.u

- Operation with 650 NDCM (resolution of 0.2e)

- The level of DC is 3×10^{-3}e/pix/day (x10 higher than the initial goal)

- First results for hidden sector candidates with an exposure of 85.2 gr-day

Cryostat

External shielding

Vacuum system

Support structure

Box with two CCDs
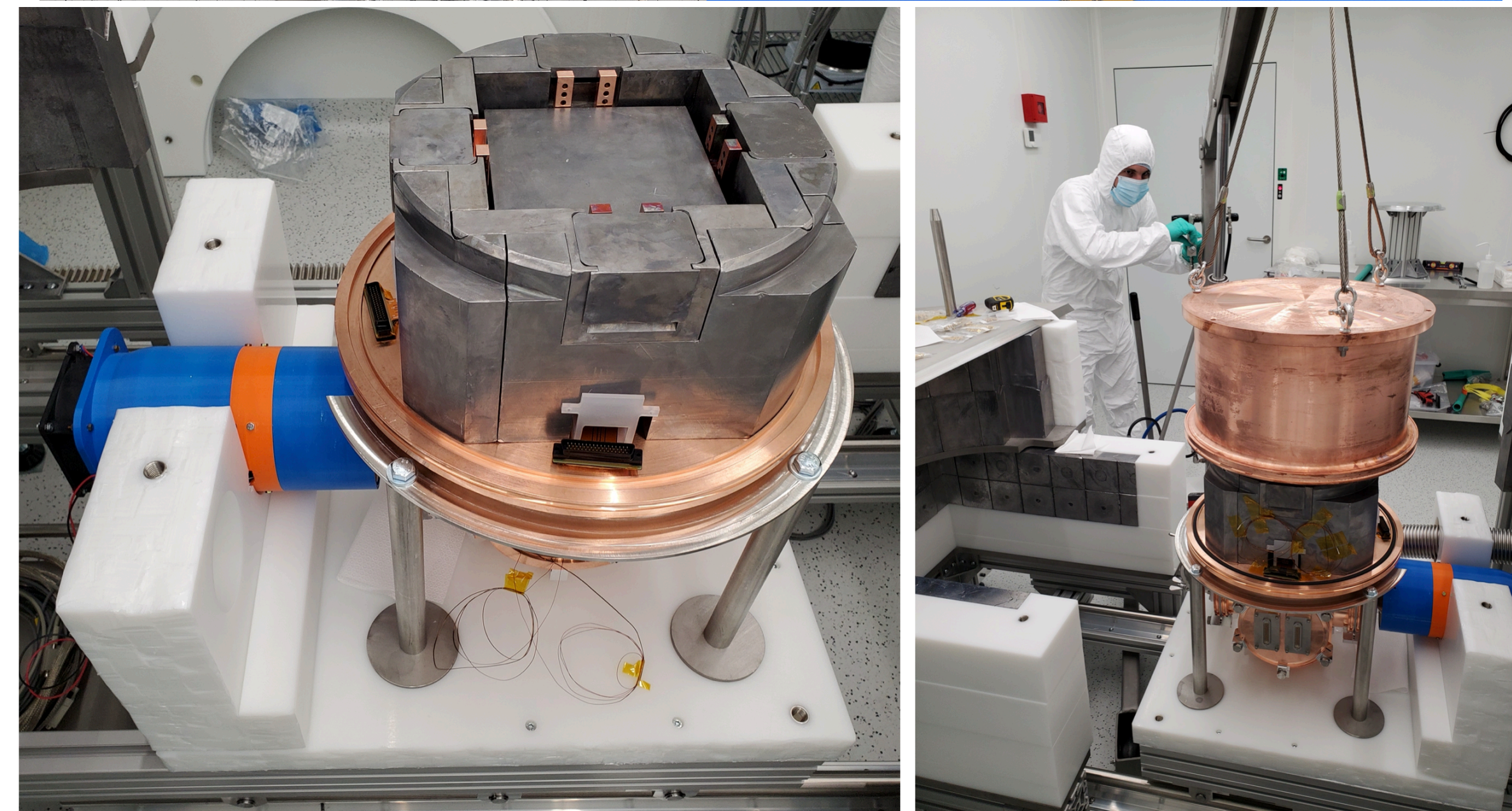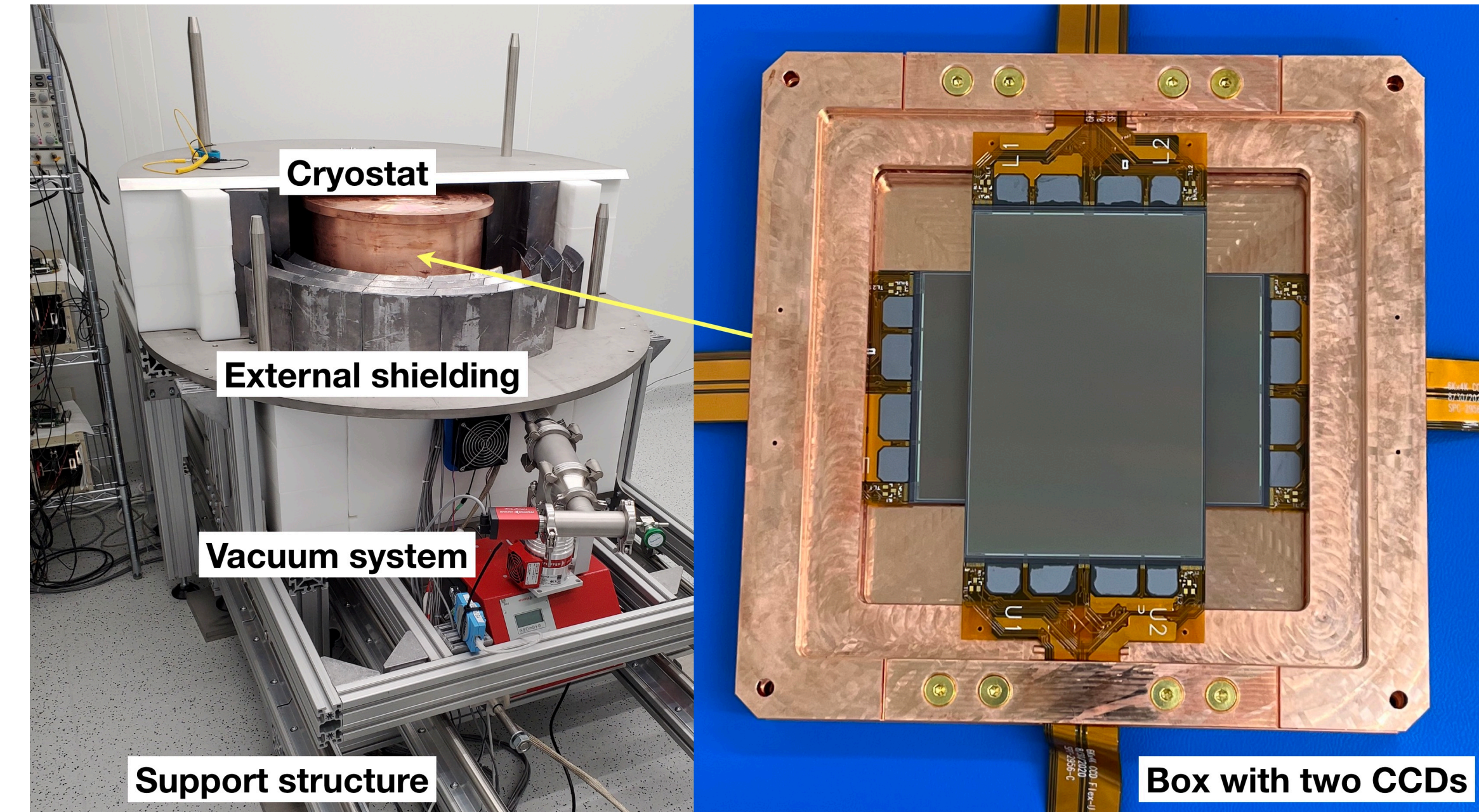
Data selection procedure

CCD **2**, amplifer L
- Baseline subtracted
— Masked > 2σ

mask clusters + CTI tail
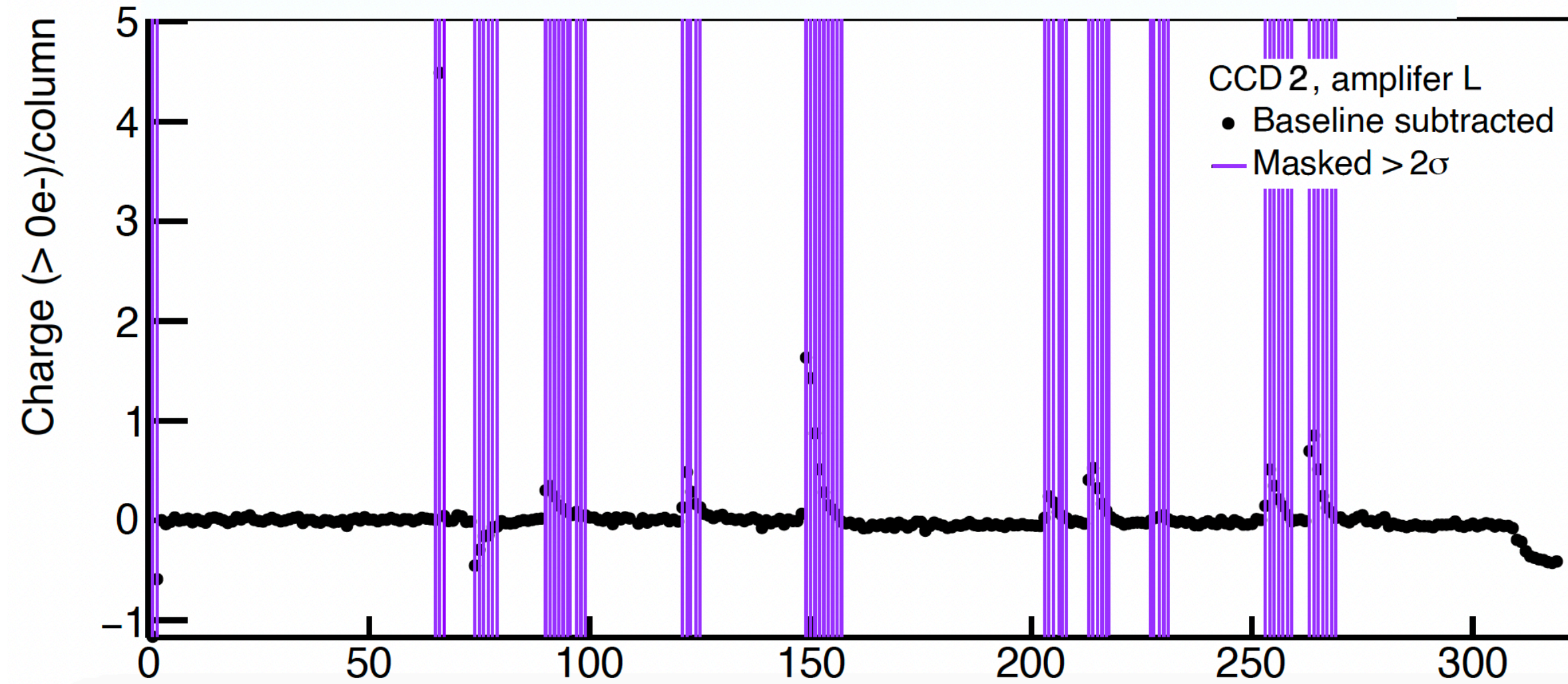
Identification and Rejection of high-energy events

clustering reconstruction algorithm, with threshold of $2\sigma_{avg}$ and seed of $2e^-$

Masking

Removal of pixel clusters and tails from CTI

Defect masking

Searching for "hot" columns with abnormally high charge

Artefacts in the serial register

They can distort the pixel charge distribution

Searching for Hidden Sector candidates

A hypothetical massive vector boson A' of a broken (dark) gauge group $U(1)_D$ that kinetically mix with the SM weak hypercharge. At low energies the mixing is between A' and a photon
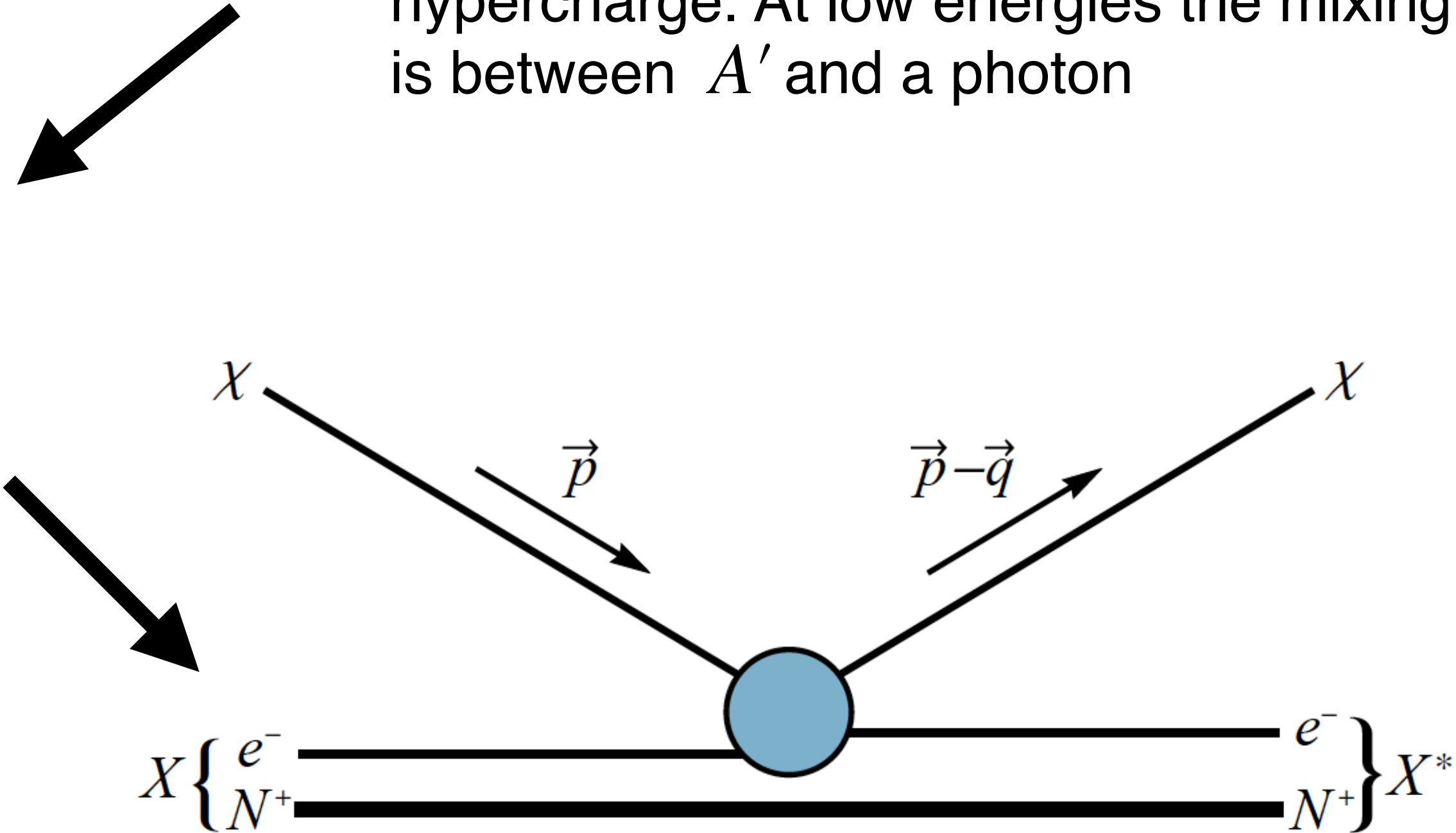
Particle candidates for a light dark matter

$$\mathscr{L} \supset -\frac{1}{4}F'^{\mu\nu}F'_{\mu\nu} - \frac{\epsilon}{2}F^{\mu\nu}F'_{\mu\nu} + \frac{1}{2}m_{A'}^2 A'^\mu A'_\mu$$

The Dark Sector interacts with the SM via the gauge boson A'.
DM particles can scatter off bound electrons of the Xe atom via A' exchange.

Two cases are of interest

- $F_{DM}(q) = 1$, heavy mediator, $(m_{A'} \gg \alpha m_e)$

- $F_{DM}(q) = (\alpha m_e/q)^2$, ultra-light vector mediator. $(m_{A'} \ll \alpha m_e)$

$$\frac{dR_{crystal}}{d \ln E_{er}} = \frac{\rho_\chi}{m_\chi} N_{cell} \bar{\sigma}_e \alpha \frac{m_e^2}{\mu_{\chi,e}^2} \int d \ln q \left(\frac{E_e}{q} \eta(u_{\min}) \right) |f_{crystal}(q, E_e)|^2 |F_{DM}(q)|^2 dq$$

χ $\vec{p}$ $\vec{p} - \vec{q}$ χ

$X \left\{ \begin{array}{l} e^- \\ N^+ \end{array} \right.$ $\left. \begin{array}{l} e^- \\ N^+ \end{array} \right\} X^*$

Cosmic Time

The rate depends on the initial and final state of the electron, the particular interaction and the Halo model
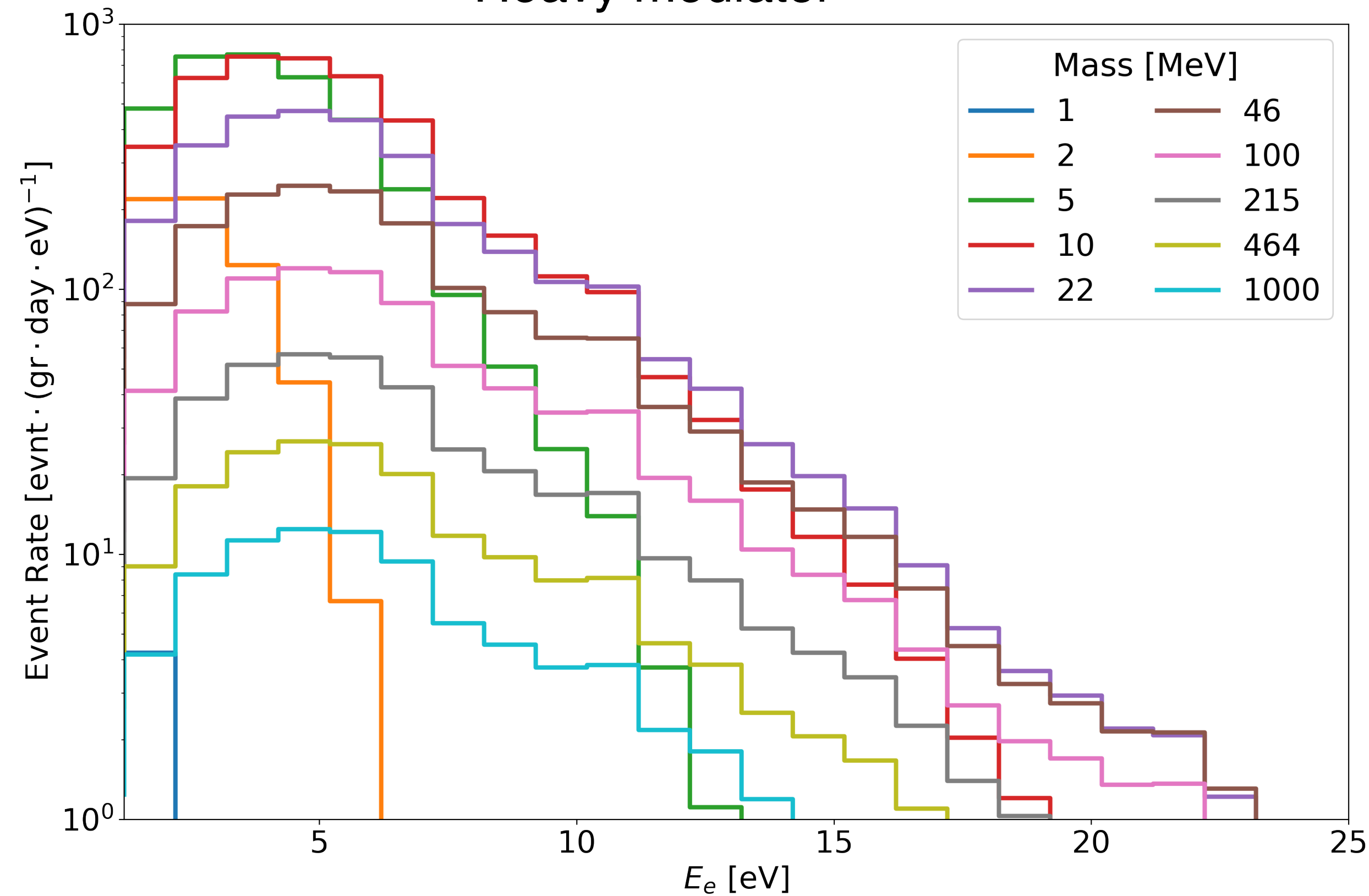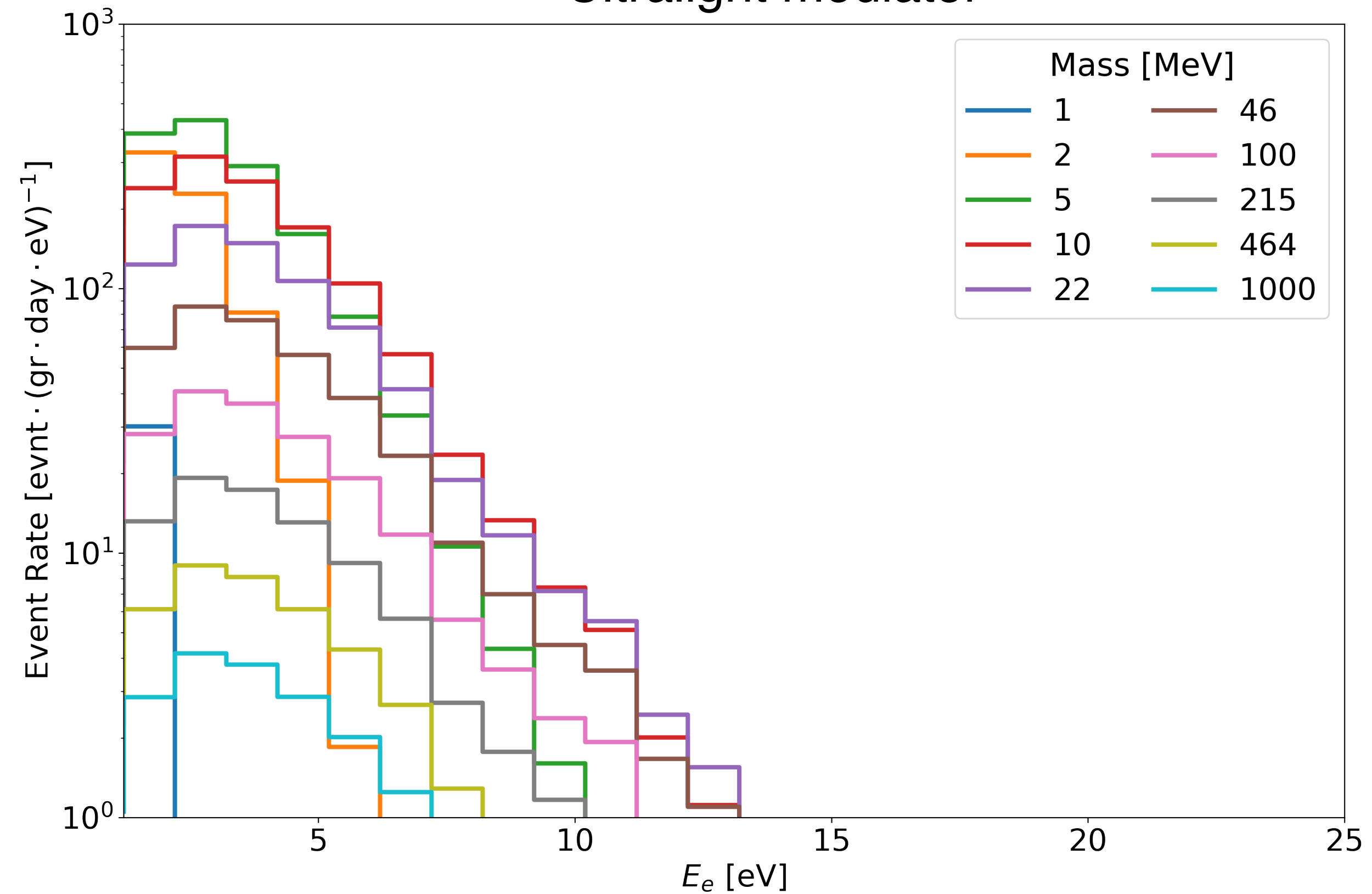
10

Computing the expected event rates

$$\frac{dR_{crystal}}{d \ln E_{er}} = \frac{\rho_\chi}{m_\chi} N_{cell} \bar{\sigma}_e \alpha \frac{m_e^2}{\mu_{\chi,e}^2} \int d \ln q \left(\frac{E_e}{q} \eta(u_{\min}) \right) |f_{crystal}(q, E_e)|^2 |F_{DM}(q)|^2 dq$$

Heavy mediator

Ultralight mediator

- We are using the QEDark (JHEP05(2016)046) framework to obtain the computed $f_{cryst}(q, E_e)$
- We use the PhystatDM conventions about the halo model and the local DM density parameters

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Detector response: Ionization

$$\frac{dR_{crystal}}{d \ln E_{er}} = \frac{\rho_\chi}{m_\chi} N_{cell} \bar{\sigma}_e \alpha \frac{m_e^2}{\mu_{\chi,e}^2} \int d \ln q \left(\frac{E_e}{q} \eta(u_{\min}) \right) |f_{crystal}(q, E_e)|^2 |F_{DM}(q)|^2 dq$$

Necessity to convert the deposited energy in resulted ionization Q

E_{er} and Q are related by a long chain of secondary scattering processes redistributing the deposited energy

Simple model: Extrapolation of the high energy understanding of ionization. If ε is the mean energy to create an electron hole pair then:

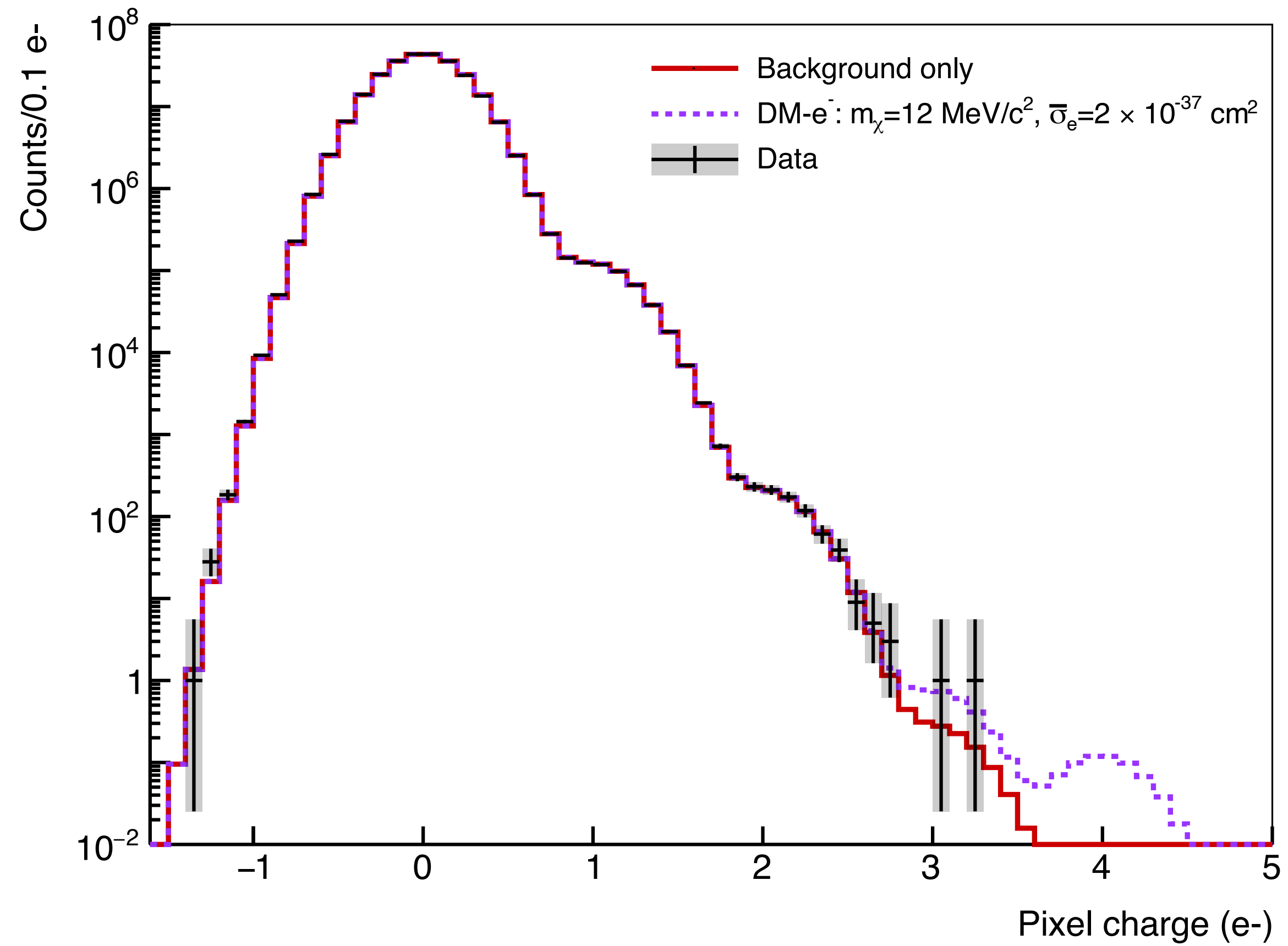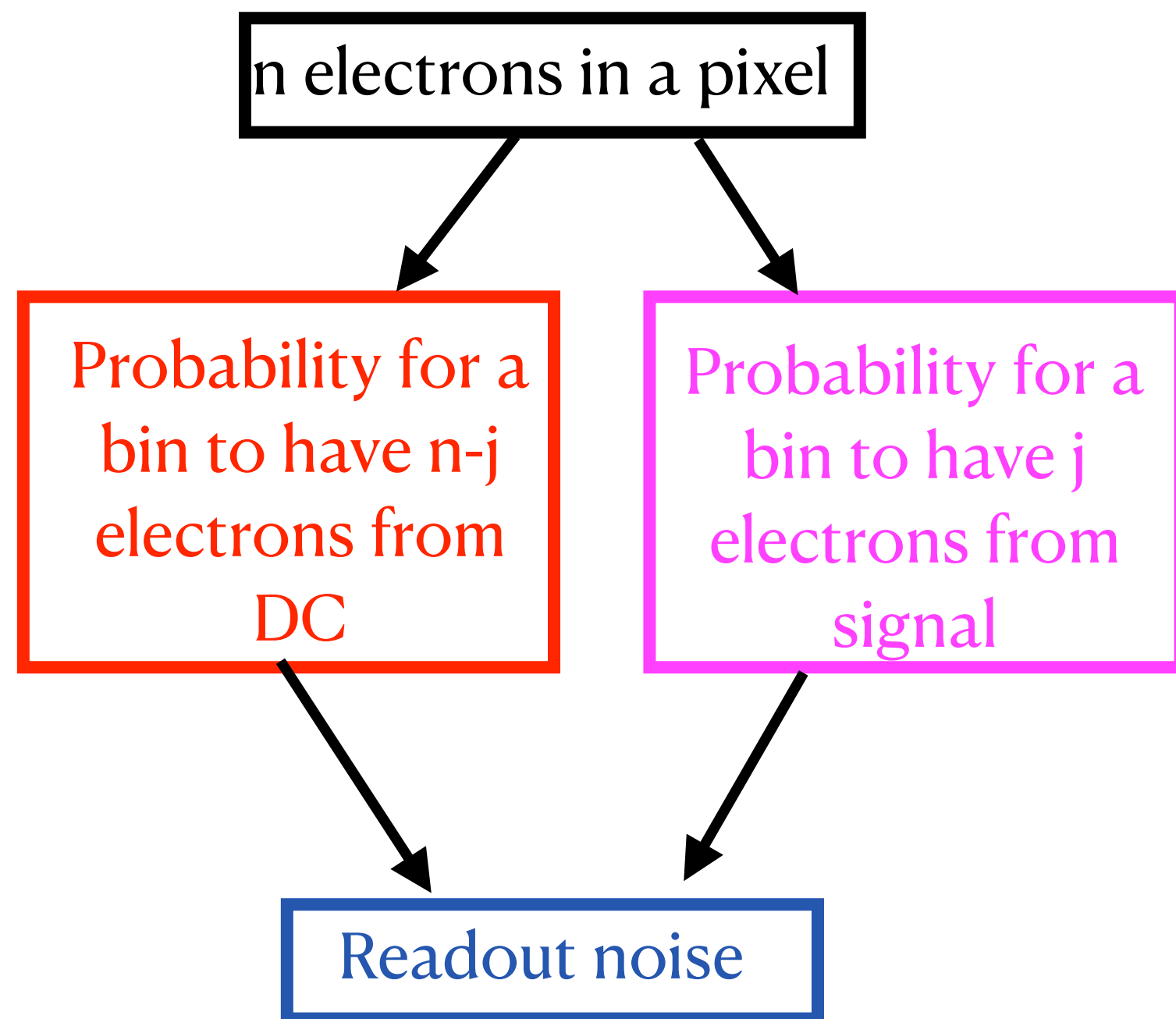$$Q(E_{er}) = 1 + [(E_{er} - E_{gap})/\epsilon]$$

Previous relation brake-down when $E_{er} = \mathcal{O}(E_{gap})$. We use a phenomenological model of impact ionization to explore the likely charge yield in this energy regime.
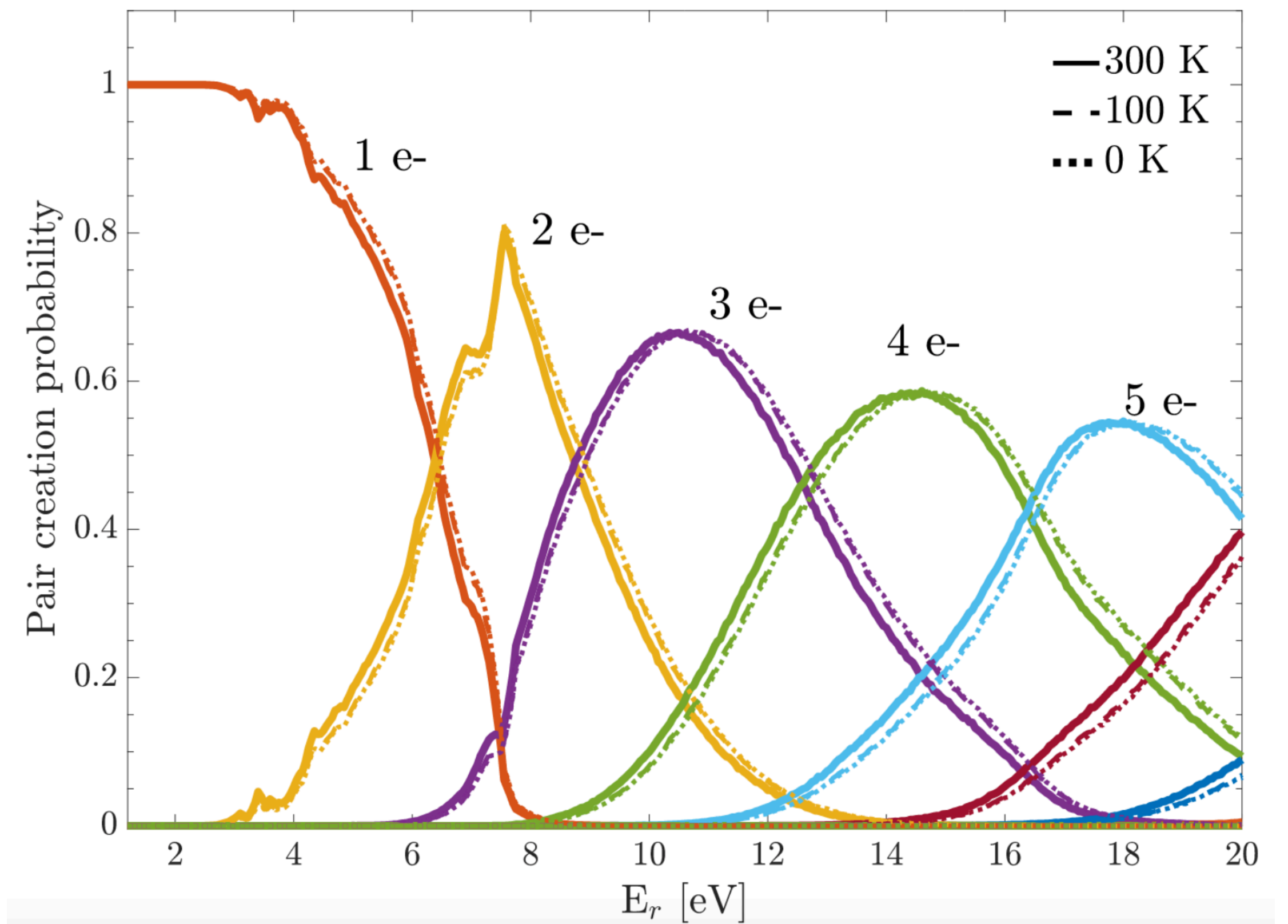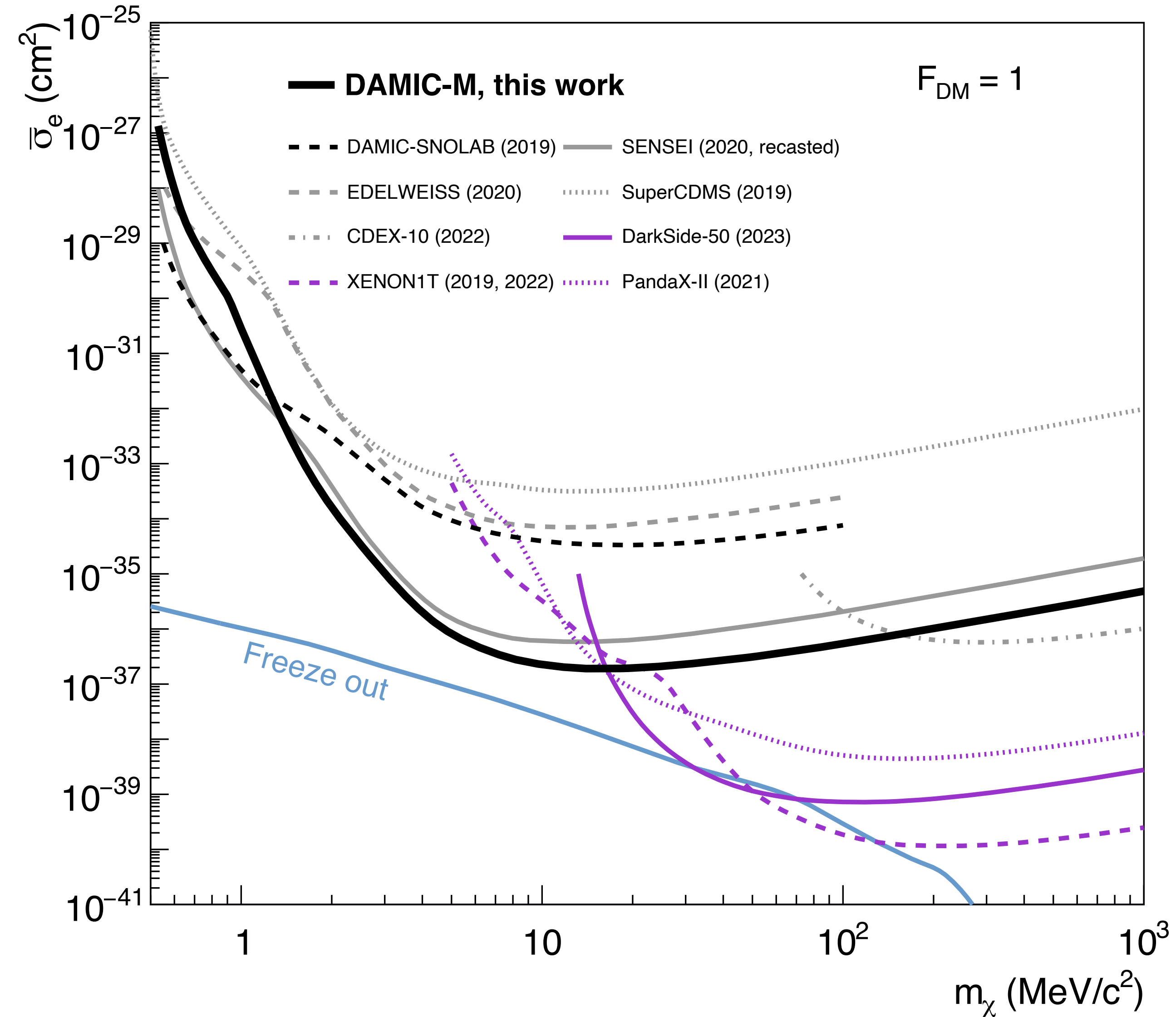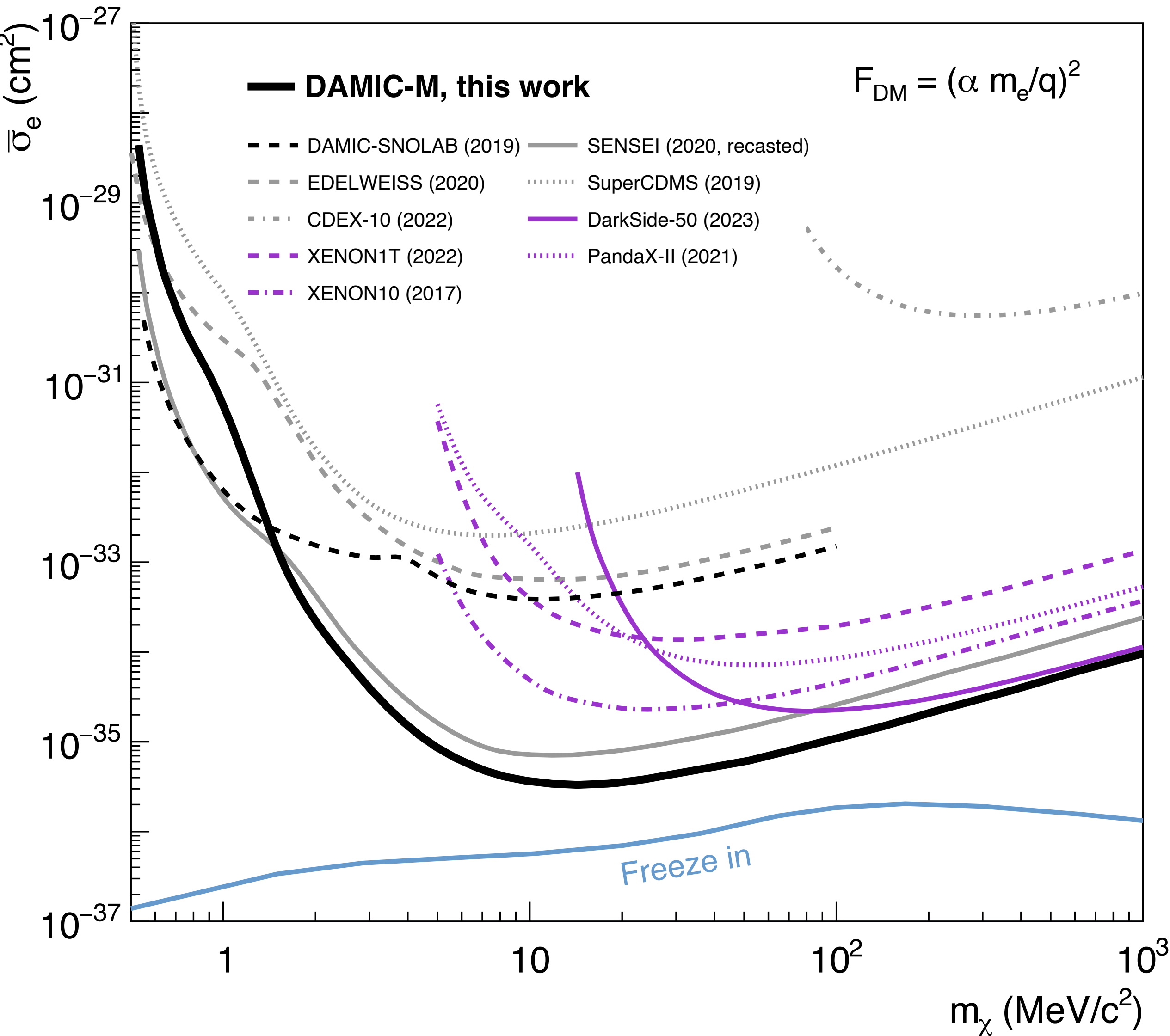
Final Signal and Background model

Then one can construct the pixel charge distribution
for any particular value for the dark current λ_{DC} and the cross section σ_e

$$\Pi(p) = N_{pixels} \sum_{n=0}^{\infty} \left(\sum_{j=0}^{n} \boxed{S(j\,|\,m_\chi, \sigma_e)} \boxed{\text{Poisson}(n-j\,|\,\lambda_{DC})} \right) \boxed{\text{Normal}(p\,|\,n, \sigma_{avg})}$$

n electrons in a pixel

Probability for a bin to have n-j electrons from DC

Probability for a bin to have j electrons from signal

Readout noise

90% CL upper limits results

Phys. Rev. Lett. 130, 171003 (2023).

90% CL upper limits results

$\bar{\sigma}_e$ (cm^2) 10^{-27}

DAMIC-M, this work $F_{DM} = (\alpha \, m_e/q)^2$

- - - DAMIC-SNOLAB (2019) —— SENSEI (2020, recasted)

$\bar{\sigma}_e$ (cm^2) 10^{-25}

10^{-27}

DAMIC-M, this work $F_{DM} = 1$

- - - DAMIC-SNOLAB (2019) —— SENSEI (2020, recasted)
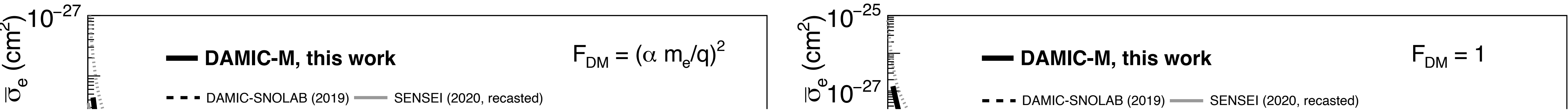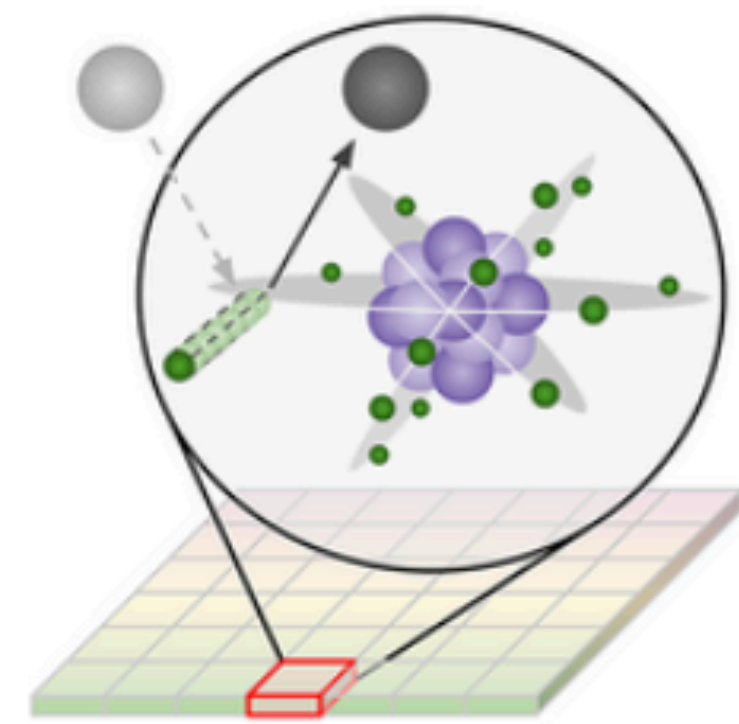
Editors' Suggestion

PDF   HTML

First Constraints from DAMIC-M on Sub-GeV Dark-Matter Particles Interacting with Electrons

I. Arnquist *et al.* (DAMIC-M Collaboration)

Phys. Rev. Lett. **130**, 171003 (2023) – Published 28 April 2023

World-leading constraints are placed on electron interactions with dark matter in the MeV to GeV range by the first underground operation of a new CCD detector.

Show Abstract +

Phys. Rev. Lett. 130, 171003 (2023).

Conclusion

- Prototype LBC is already installed at the LSM

- LBC is taking data under stable conditions of low background and optimised readout noise

- First science results with 85.2 gr-day exposure. World leading limits for DM-e scattering candidates

- DAMIC-M aim to explore a vast region of the parametric space in light supersymmetric and hidden sector DM candidates

- The final DAMIC-M detector will reach a kg-year scale exposure using skipper silicon-CCD technology with a fraction of a dru as goal of the background rate

- The actual experiment construction will start at 2024

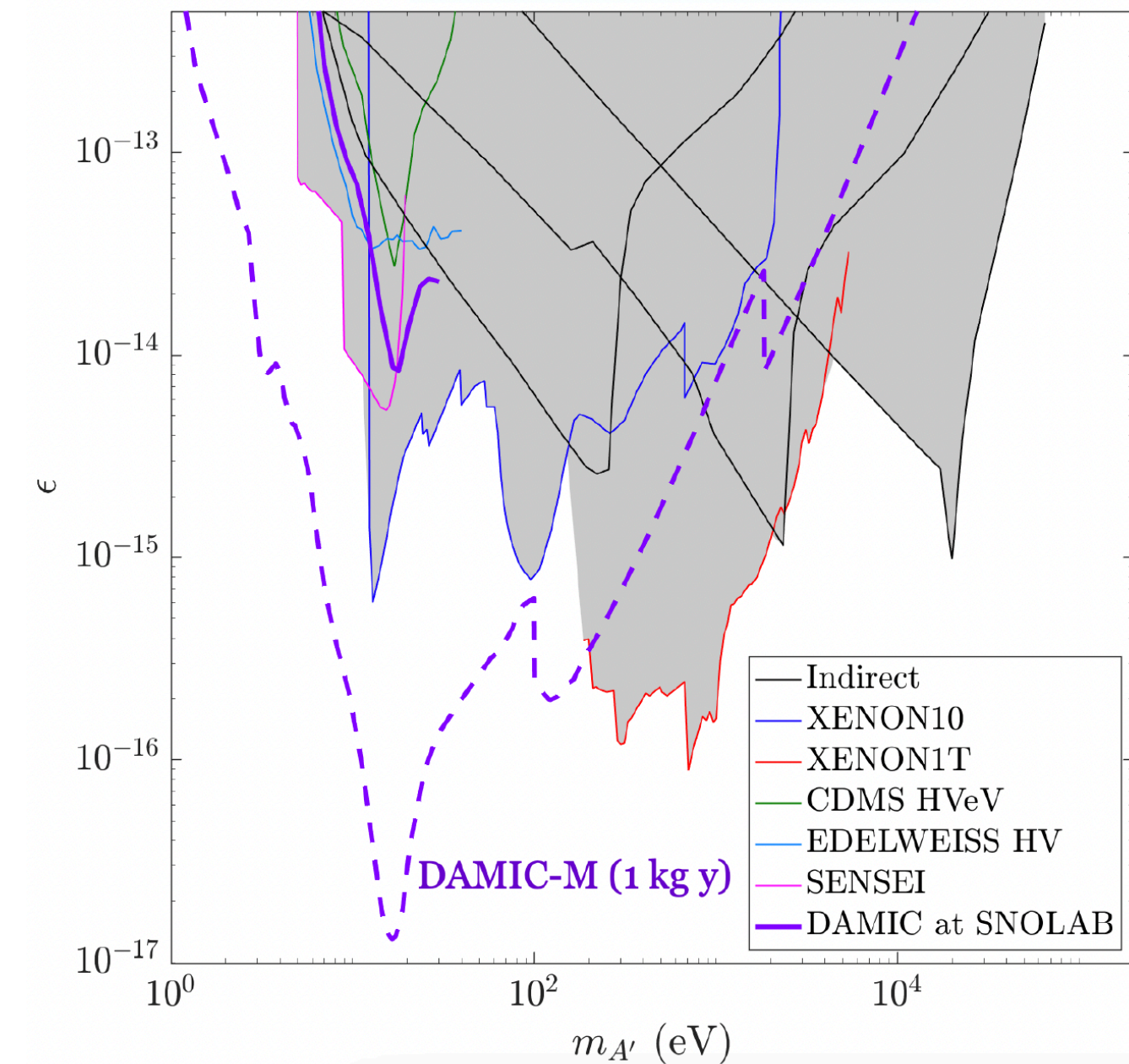# Supplementary slides
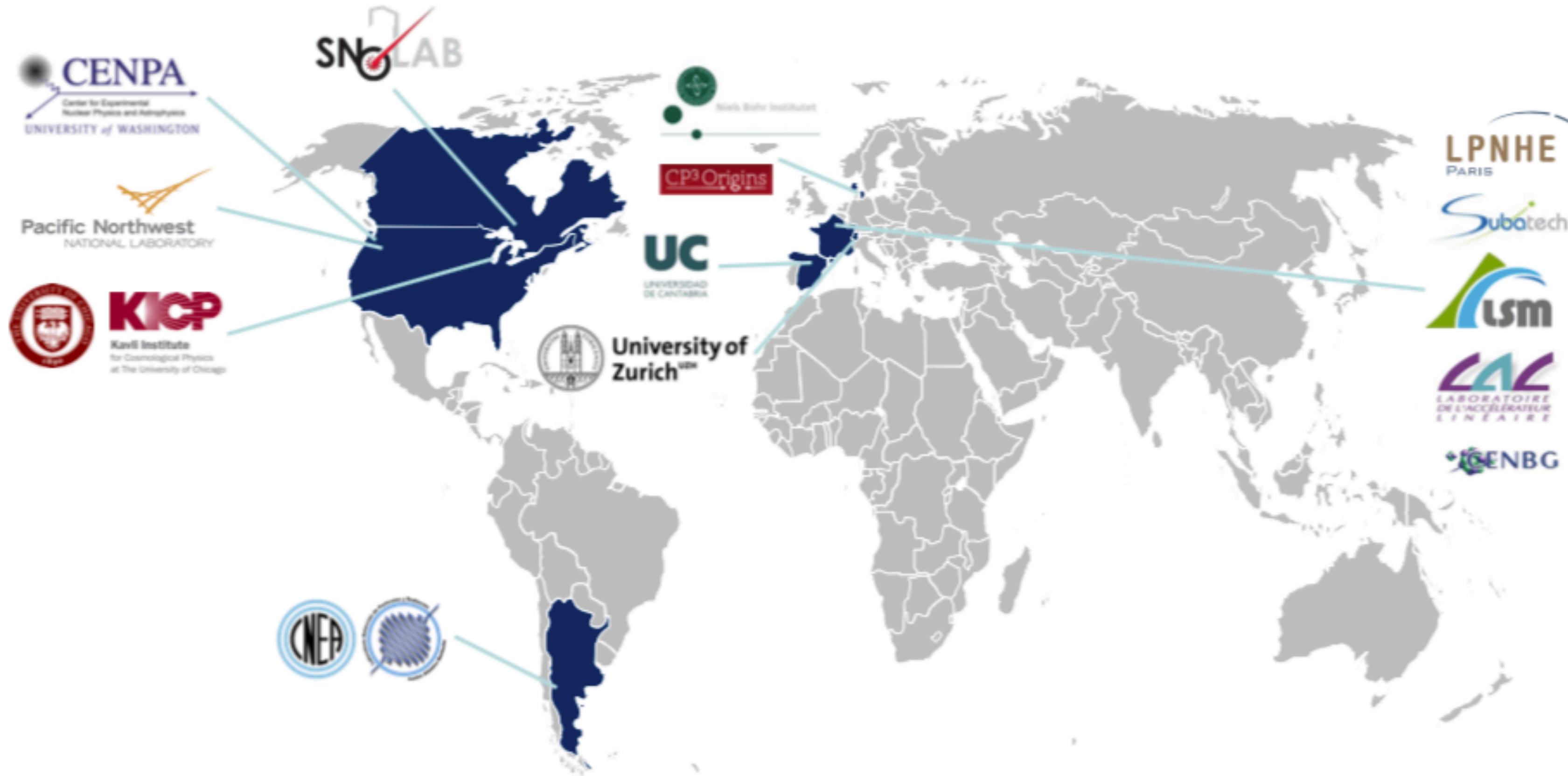
DAMIC-M Collaboration

Science goals

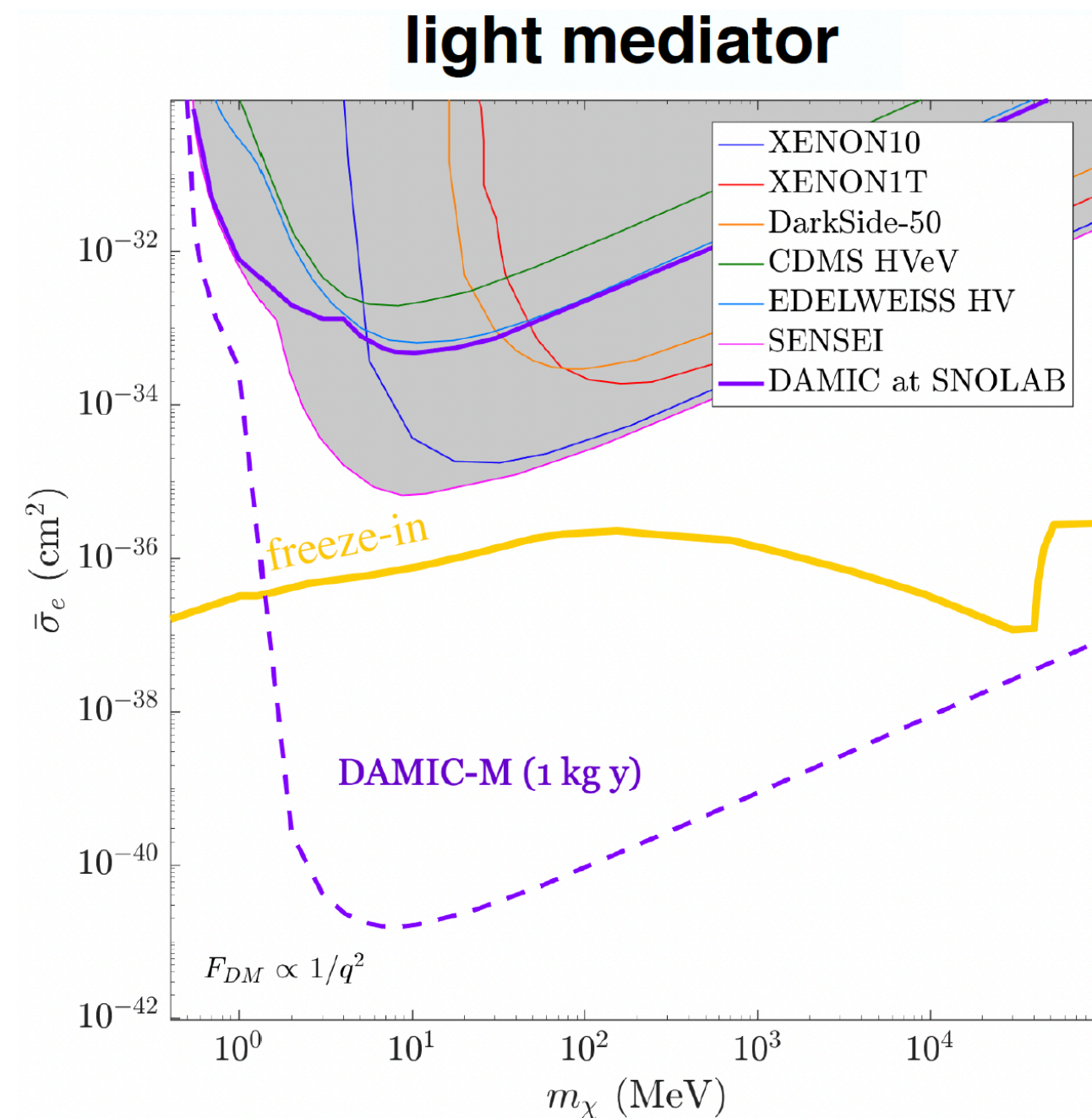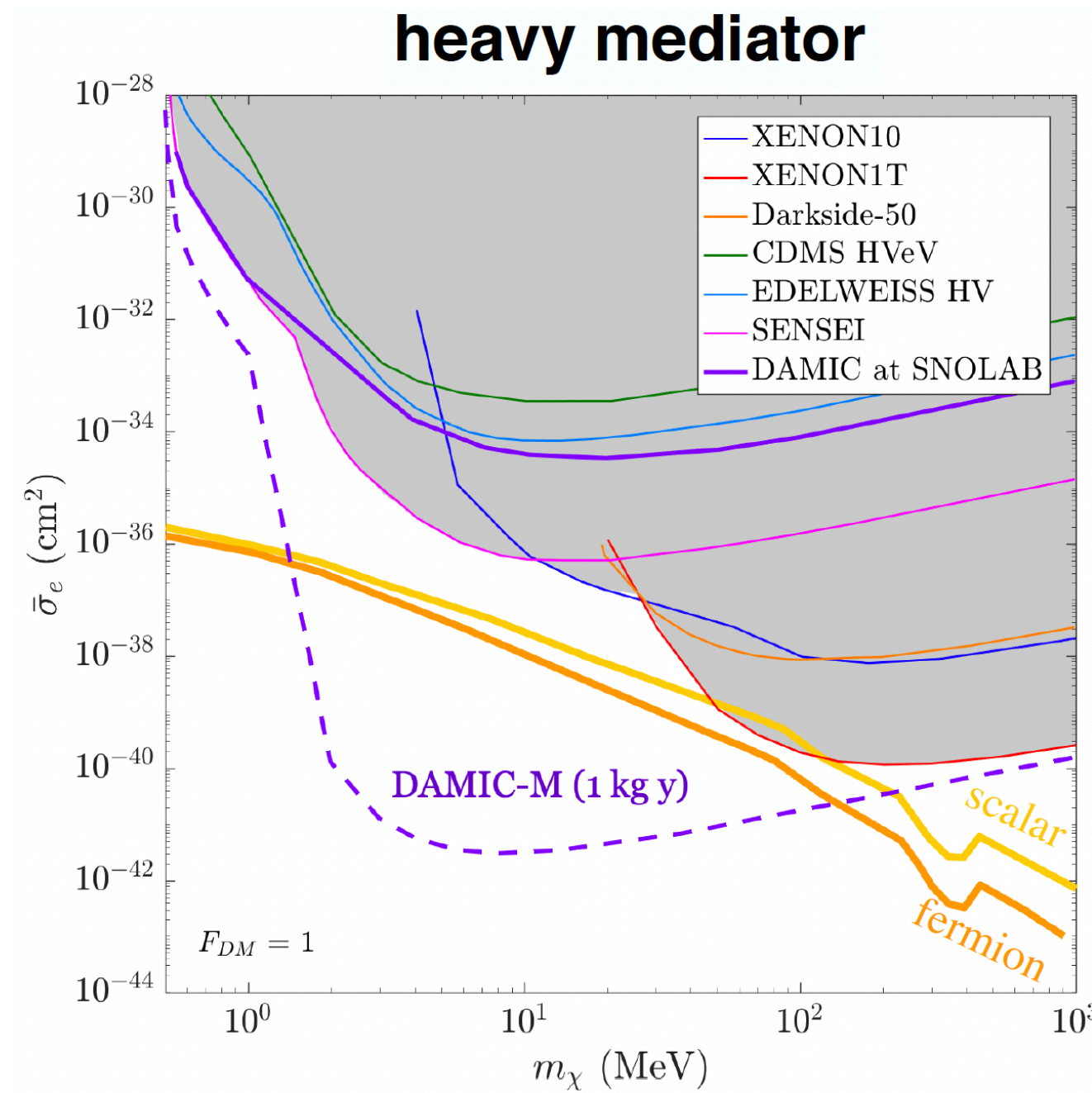- **Spin independent WIMP-nucleon elastic scattering**

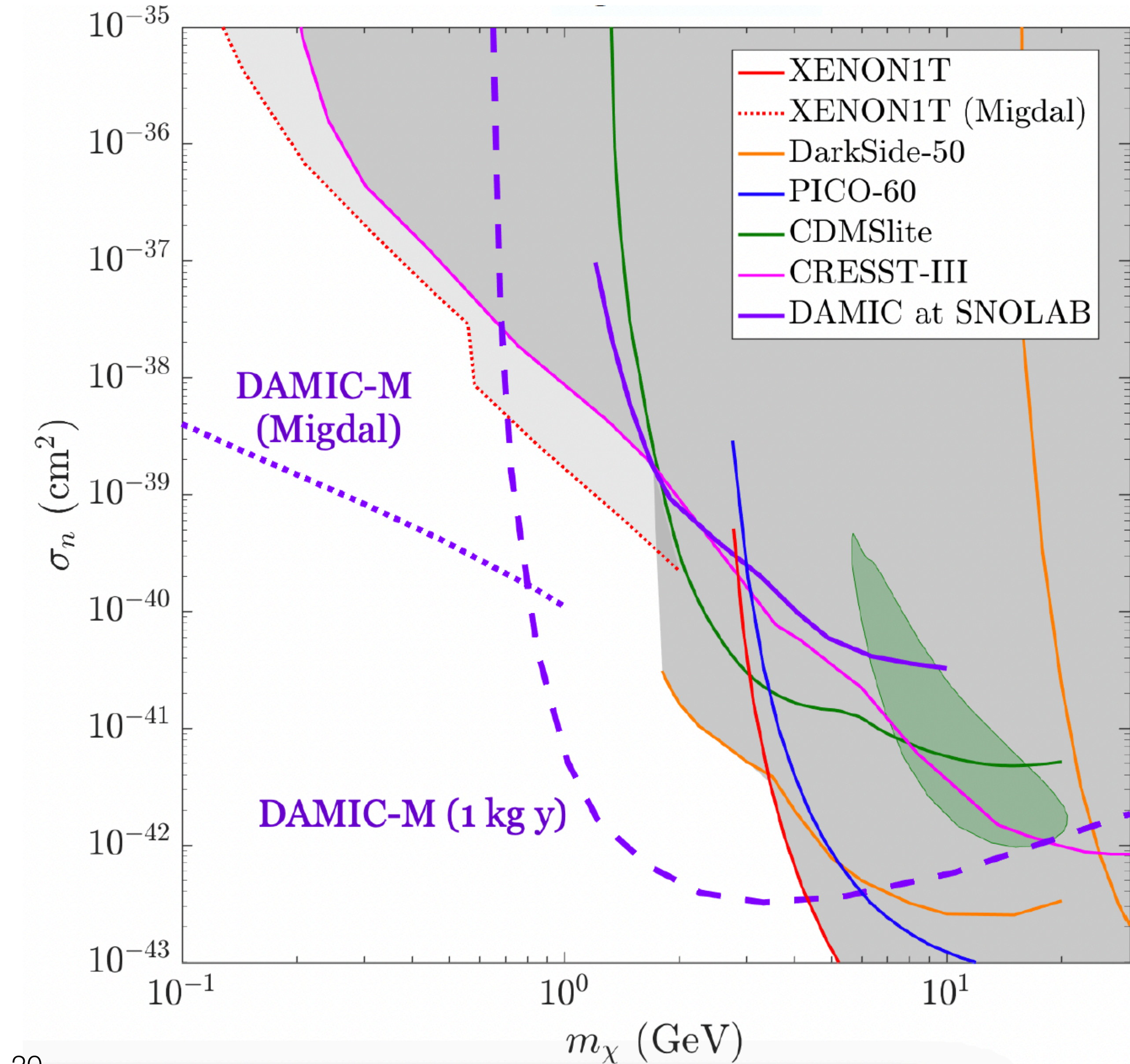- **Hidden sector candidates (DM-e scattering, Bosonic DM)**

With an exposure of 1 kg-year and low background and Dark Current conditions is possible to explore a vast region of the available parametric space
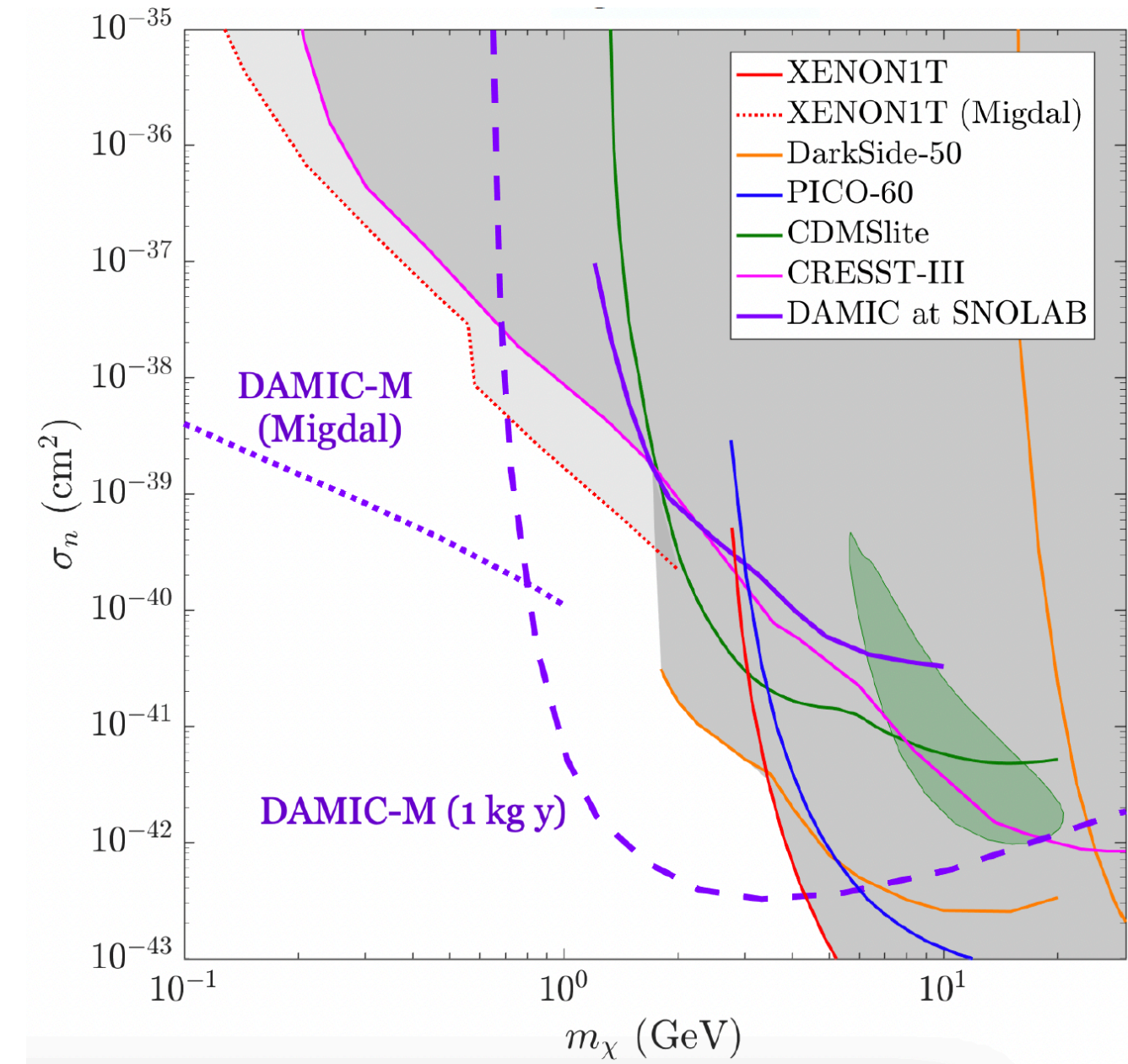
Science goals: Light WIMPs

Spin independent WIMP-nucleon elastic scattering

- Signal expects to be diffusion limited NR interaction

- Due to the small mass of Si atom the sensitivity is strong down to WIMPs with mass $\mathcal{O}(1\text{GeV})$

- Possibility to explore a vast region of the available parametric space with kg-year scale exposure

- Possibility to extend to lower DM masses via the Migdal channel
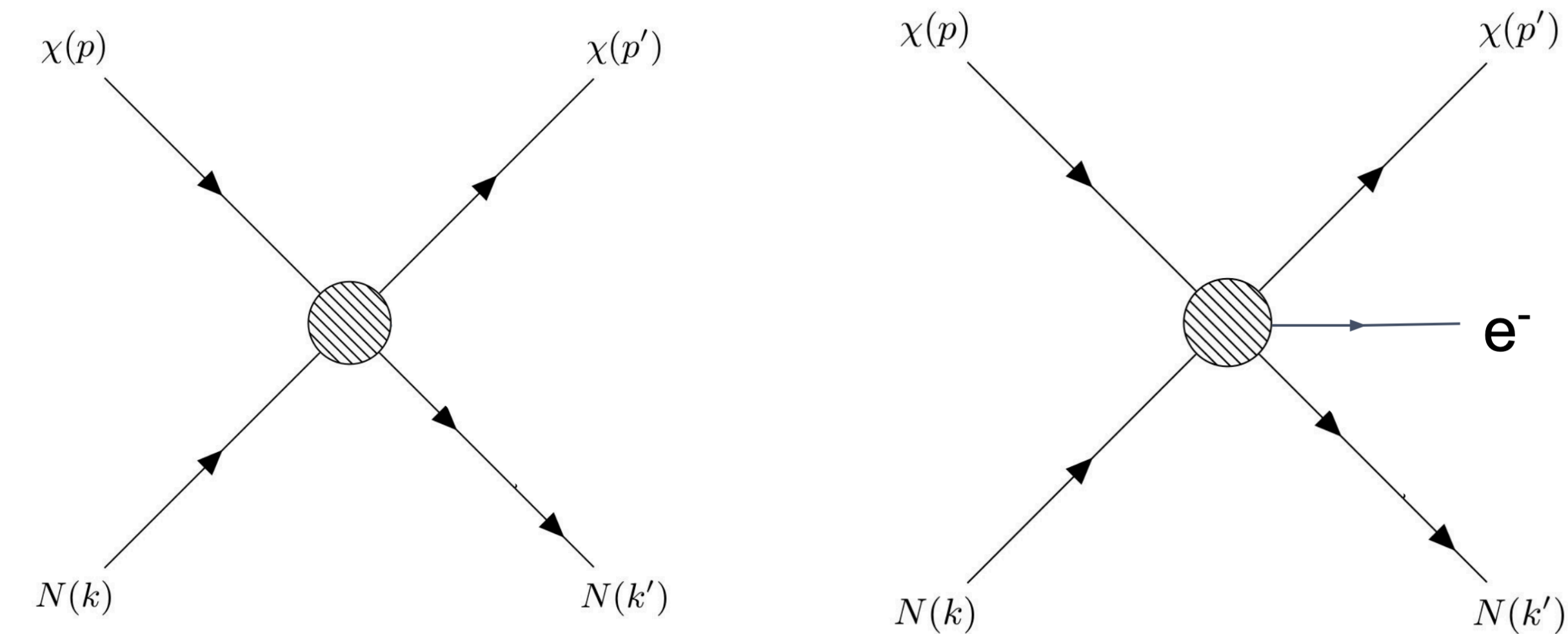
# Science goals: Hidden Sector candidates

DM scattering off electrons

- Signal expects to be diffusion limited ER interaction

- Due to the very small dark current and the single electron resolution the sensitivity is strong down to DM masses of $\mathcal{O}(10^{-1})\,MeV$

- Possibility to explore a vast region of the available parametric space with kg-day scale exposure

- Physics channels to be explored: DM-electron scattering via heavy or ultralight mediator, dark photon absorption
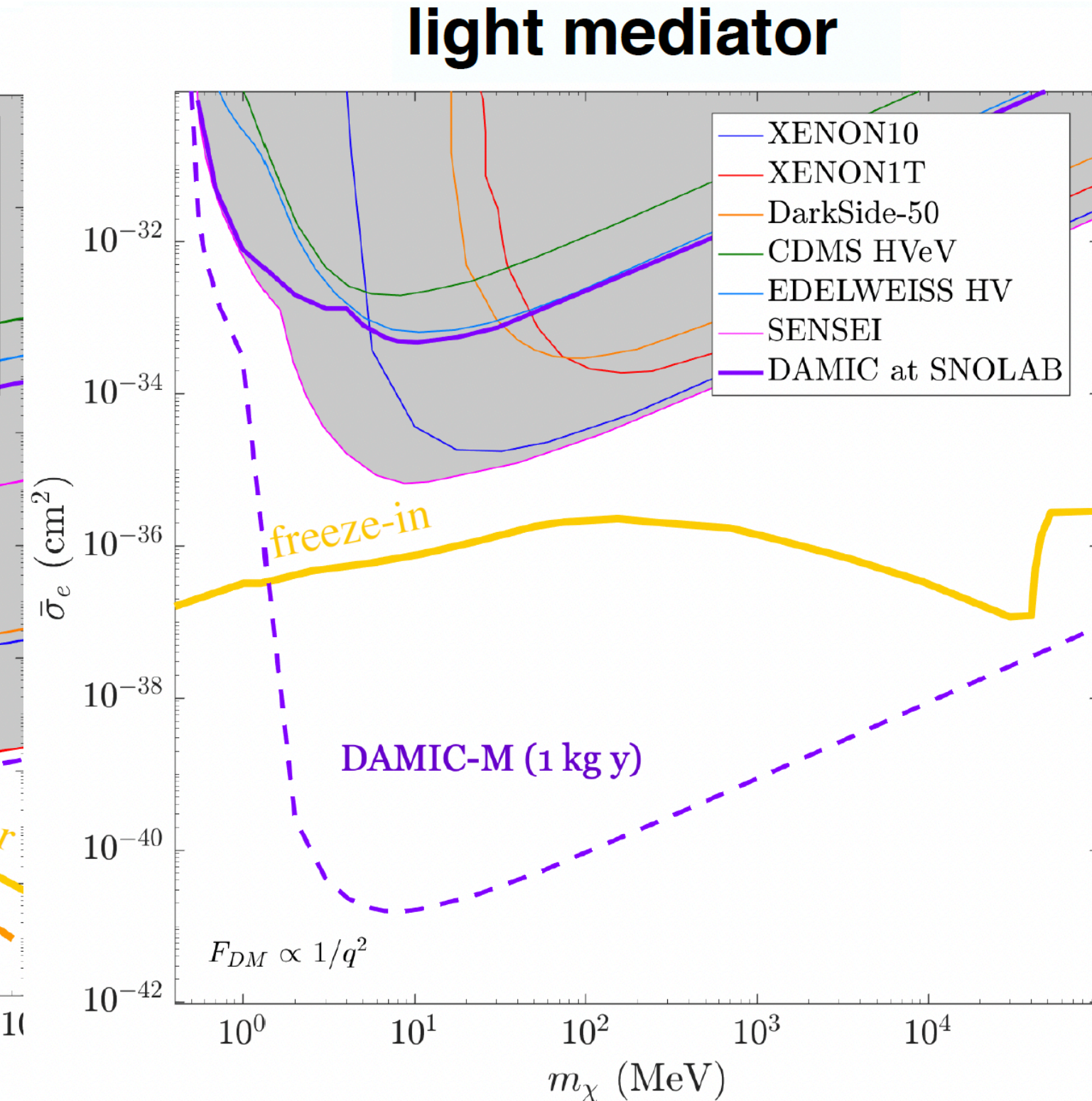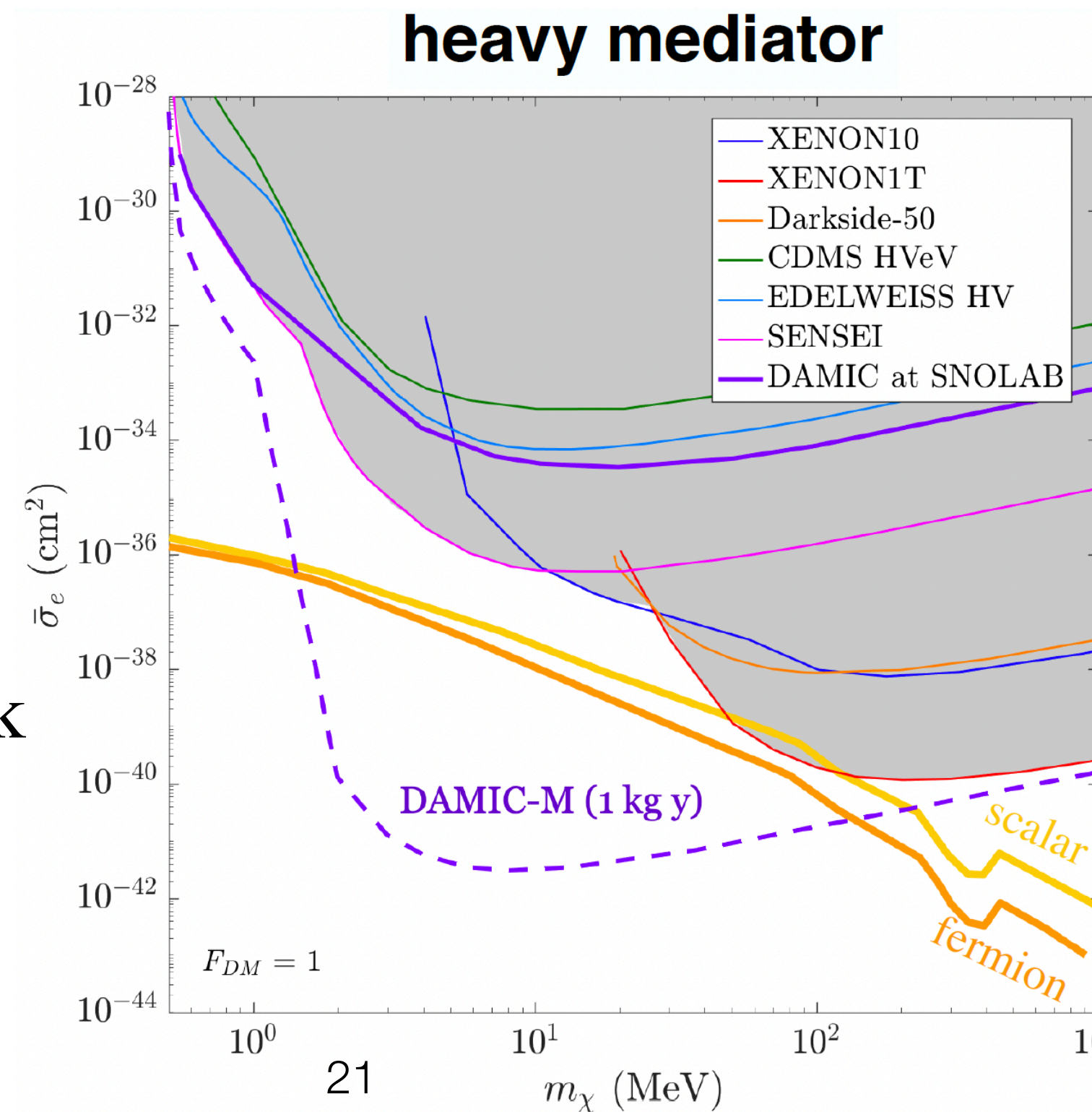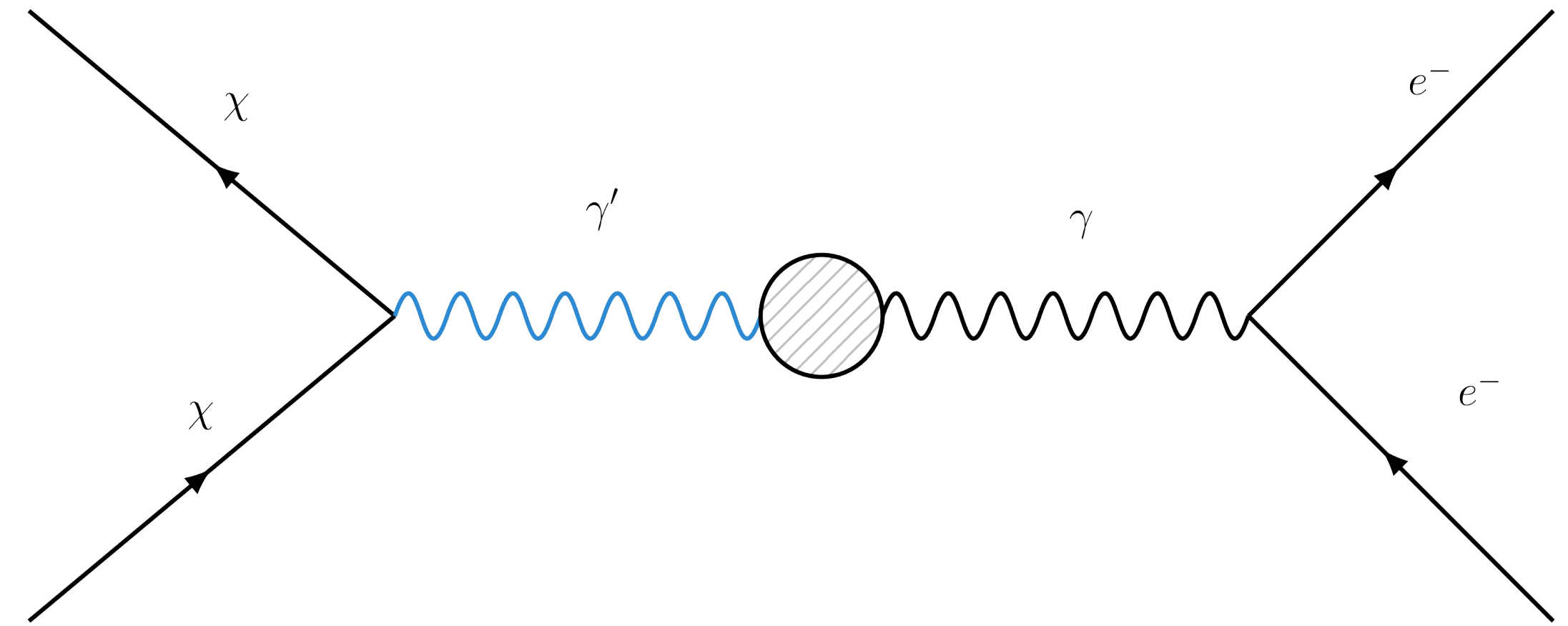
Detector response

Computation of the signal probability per pixel

- We simulate CCD images with DM events for a given cross section

- We take into account the pair creation probability in the determination of the secondary ionization

- We account for the diffusion of the ionization charges during drift to the pixel array

- We account for the 10x10 binning readout mode

For a detailed description of the signal model calculation see:
[Michelangelo Traina Ph.D. thesis, Sorbonne University, 2022](#)

Charge diffusion and binning

Legend:
- Pre-diffusion (1499760 pixels)
- Post-diffusion (2171260 pixels)
- Post-binning 10×10 (1602600 binned pixels)

x-axis: Pixel charge [e^-]
y-axis: Pixels per unit charge per kg-year
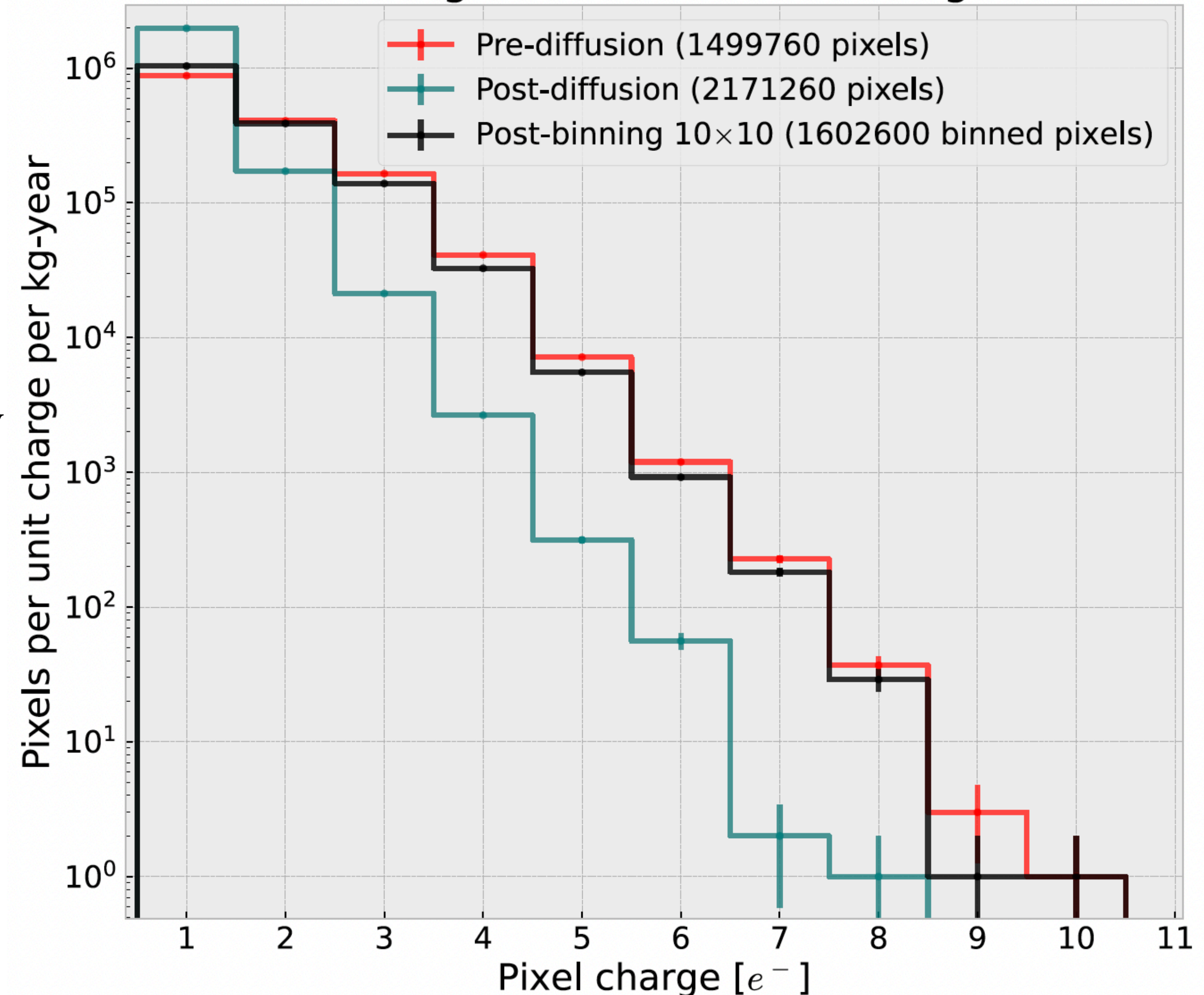
How DM interacts with the Si crystal?

It's necessary to account for the crystal lattice nature of the Si target.

Si crystal is a **multi-body system** with **delocalised valence electrons** occupying an energy **band-structure** with energy gap separating from the unoccupied conduction bands

QM Problem in a periodic potential can be reduced to the 1st Brillouin zone

$$\psi_{i,\mathbf{k}}(\mathbf{x}) = \frac{1}{V_{cryst}} \sum_{\mathbf{G}} u_i(\mathbf{k}+\mathbf{G}) \cdot \exp\left(i(\mathbf{k}+\mathbf{G})\mathbf{x}\right)$$

The wave function coefficients are obtained with **DFT approximations**. All the properties of the system are obtained from the ground state particle density. This is obtained by use of pseudo-potential approximation of independent electrons with the same ground state density
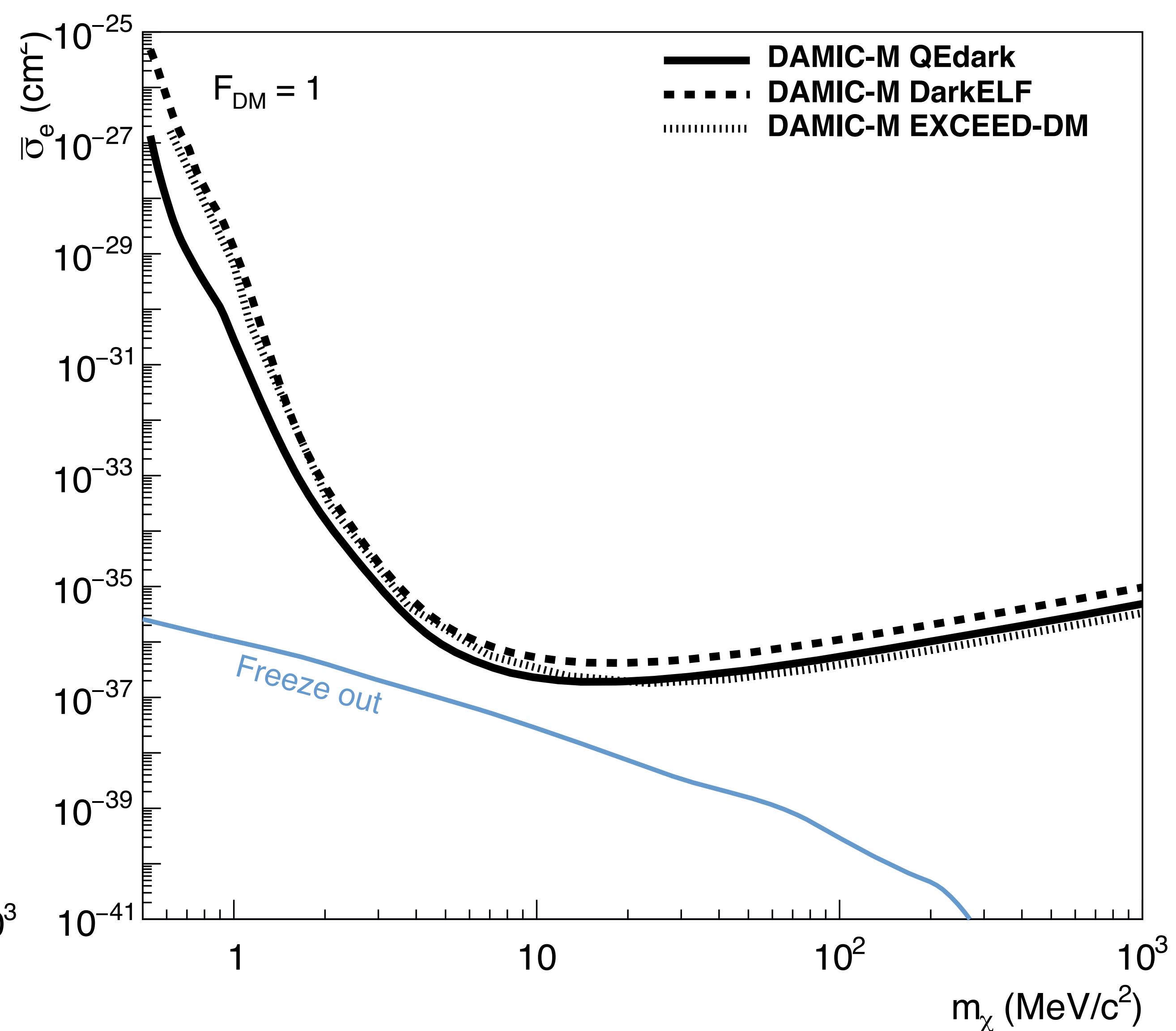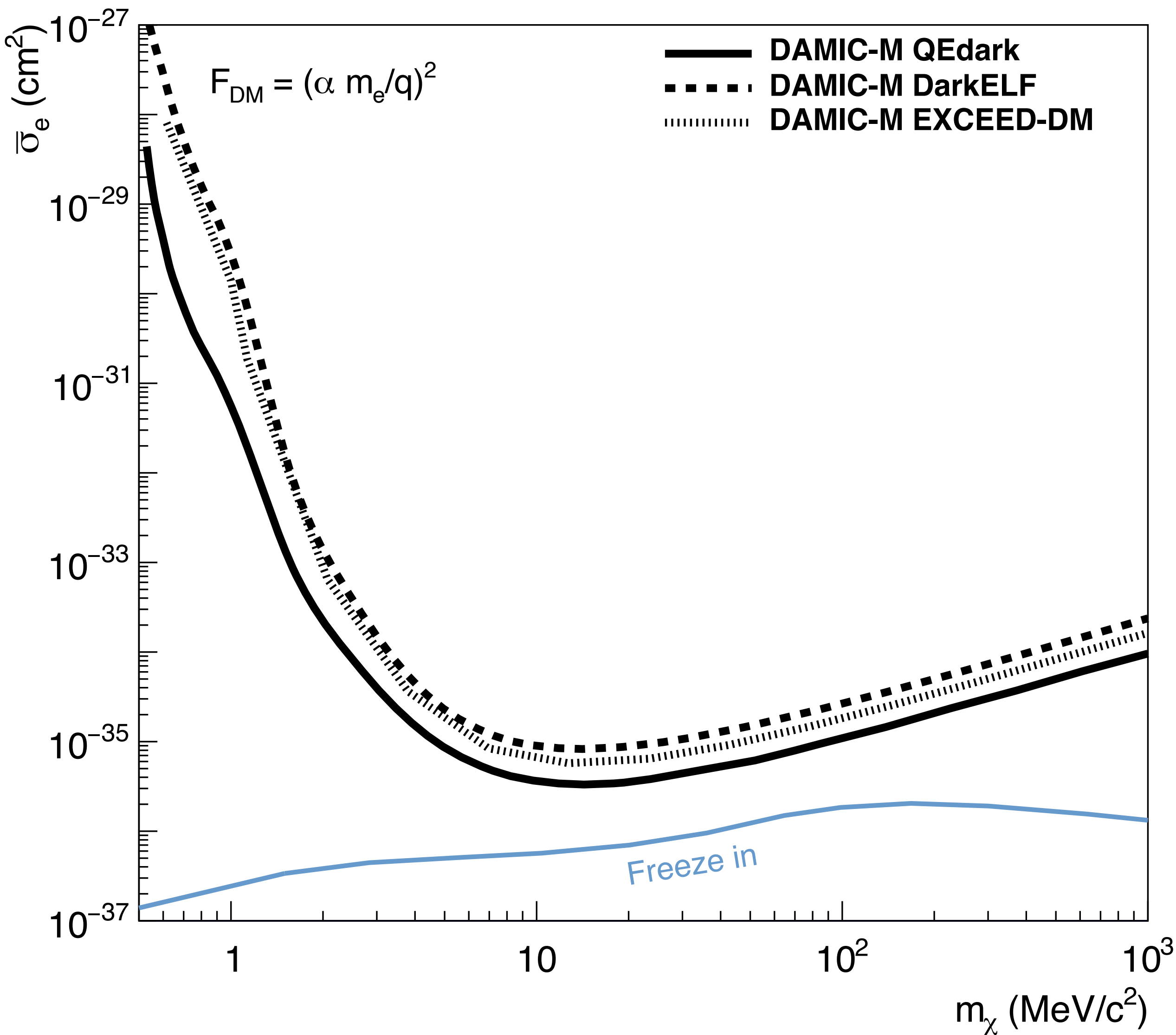
The form factor for the transition from occupied valence state to unoccupied conduction state is

$$f_{i\mathbf{k},i'\mathbf{k}',\mathbf{G}'} = \sum_{\mathbf{G}} u_{i'}^*(\mathbf{k}'+\mathbf{G}'+\mathbf{G}) \cdot u_i(\mathbf{k}+\mathbf{G})$$

The form factor for the transition from occupied valence state to unoccupied conduction state is

$$|f_{crystal}(q,E_e)|^2 = \frac{2\pi^2(\alpha m_e^2 V_{cell})^{-1}}{E_e} \sum_{i,i'} \int_{BZ} \frac{V_{cell}d\mathbf{k}}{(2\pi)^3} \frac{V_{cell}d\mathbf{k}'}{(2\pi)^3} E_e \delta(E_e - E_{i'\mathbf{k}'} - E_{i\mathbf{k}}) \sum_{\mathbf{G}'} q\delta(q - |\mathbf{k}'-\mathbf{k}+\mathbf{G}'|)|f_{i\mathbf{k},i'\mathbf{k}',\mathbf{G}'}|^2$$

Evaluation of theoretical uncertainties

Inference process

The signal framework can give us the histogram of the signal or the DC-only $\Pi_i, i \in bins$. This can be compared to the data histogram $D_i, i \in bins$ by a binned likelihood

$$\log \mathscr{L} = - (\theta - \bar{\theta})^{\top} \Sigma^{-1} (\theta - \bar{\theta}) + \sum_{i \in bins} \left(D_i \log(\Pi_i) - \Pi_i - \log(D_i!) \right)$$

If one want to set an upper limit to a value s of the signal strength then one can define the **log likelihood ratio** **(**Cowan et al.**)**

$$t_s = - 2 \log \frac{\mathscr{L}(s, \hat{\hat{\theta}})}{\mathscr{L}(\hat{s}, \hat{\theta})}$$

The asymptotic distribution of this test statistic can either be evaluate with toy MC either make use of Wilks' theorem
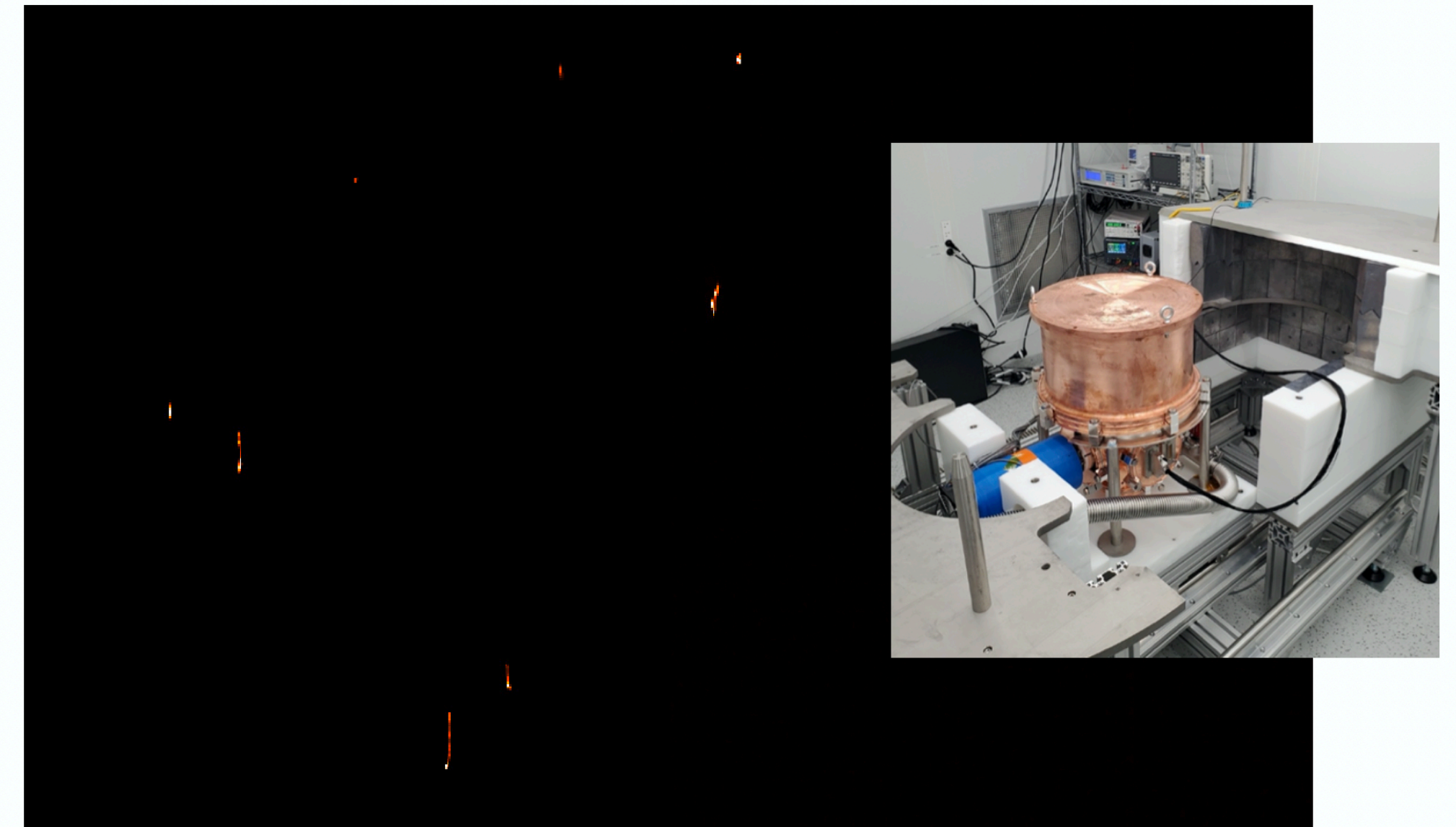
LBC commissioning and first science runs

Feb 22 - May 22: Commissioning runs

- Systematic effort for DC reduction

- High-energy BG at 3000 dru with the internal lead shield

- Various calibrations and analysis tools development

- Optimisation of CCD readout scheme

May 22 - now: Science runs

- Further suppression of high-energy BG at 10 dru with external polyethylene+lead shielding

- Data taking with 650 skips and resolution of 0.2 e

- Accumulation of 115 gr-days of exposure with a binning readout scheme of 10x10

- First science results on DM-e scattering

Internal shield

Internal + external shield