

The gallium anomaly and quantum decoherence Thomas Schwetz — IRN Neutrino meeting, Nantes, France — 19 June 2023



KIT – Die Forschungsuniversität in der Helmholtz-Gemeinschaft









Outline

The gallium anomaly and why sterile neutrinos are problematic

- A QFT calculation of standard decoherence effects in reactor and radioactive source experiments w. Raphael Krüger, arXiv:2303.15524 (EPJC to appear)
- An explanation of the gallium anomaly in terms of new physics quantum decoherence without sterile neutrinos w. Yasaman Farzan, arXiv:2306.09422 (today)



The gallium anomaly

- Measurements of gallium solar neutrino experiments GALLEX and SAGE with radioactive ⁵¹Cr or ³⁷Ar sources lead to rates lower than expected ($\sim 2\sigma$) e.g. Giunti, Laveder, 2011
- possible explanation due to eV sterile neutrino oscillations?





The BEST experiment

$R_{in} = 0.79 \pm 0.05$

$R_{out} = 0.77 \pm 0.05$



V. V. Barinov et al., Phys. Rev. Lett. 128 (2022), no. 23 232501; Phys. Rev. C 105 (2022), no. 6 065502





The gallium anomaly

	$\chi^2_{\rm null}/{ m dof}$	<i>p</i> -value		
CS1, BEST	32.1/2	$1.1 \times 10^{-7} (5.3\sigma)$		
CS1, all	36.3/6	$2.4 \times 10^{-6} (4.7\sigma)$		
CS2, BEST	34.7/2	$2.9 \times 10^{-8} (5.5\sigma)$		
CS2, all	38.4/6	$9.4 \times 10^{-7} (4.9\sigma)$		

Farzan, TS, 2306.09422 cross sections CS1, CS2 from Haxton et al., 2303.13623







Can it be explained by eV sterile neutrino oscillations?





10^{-1} 10^{-3}





Can quantum decoherence have an effect on the gallium/reactor tension?

- Arguelles, Bertolez-Martinez, Salvado, Impact of Wave Packet Separation in Low-Energy Sterile Neutrino Searches [2201.05108]
- Arguelles, Conrad et al., New Clues About Light Sterile Neutrinos: Preference for Models with Damping Effects in Global Fits [2211.02610]
- Akhmedov, Smirnov, Damping of neutrino oscillations, decoherence and the lengths of neutrino wave packets, [2208.03736]; Jones, Comment on... [2209.00561], Akhmedov, Smirnov, Reply to comment... [2210.01547]

 Jones, Marzec, Spitz, The Width of a Beta-decay-induced Antineutrino Wavepacket [2211.00026]



Can quantum decoherence have an effect on the gallium/reactor tension?

• R. Krüger, T. Schwetz, 2303.15524 (EPJC to appear):

perform a QFT calculation of the oscillation amplitude, taking into account localization of external particles for reactor and gallium experiments

"first principle" calculation of decoherence effects







QFT approach to neutrino oscillations



- Feynman diagram for joint process of production, propagation, detection macroscopic separation of the two vertices
- early papers:
- Rich, 1993; Giunti, Kim, Lee, Lee, 1993; Grimus, Stockinger, 1996; Kiers, Weiss, 1998 review paper: M. Beuthe, Phys. Rept. 375 (2003) 105 [hep-ph/0109119]



QFT approach to neutrino oscillations



Figure 1: Feynman diagram for the total process in an oscillation experiment.

reactor experiments: 3 Neutrino oscillation amplitude and event rate b production: $Cr \rightarrow V + \nu_e$

We now move to the discussion of the amplitude relevant for neutrino oscillation experiments considertection: Gasterine poduction, propagation and detection: Gasterine, we consider neutrino production by the decay of a particle A into two final state particles and an anti-neutrino, $A \rightarrow 1+2+\bar{\nu}$, and anti-neutrino detection via the process $B + \bar{\nu} \rightarrow 3+4$. We have in mind reactor neutrinos, where the production process corresponds to the beta decay of a_{\bullet} nucleus (A), and the detection process is the inverse beta decay reaction on a proton (B), but



QFT approach to neutrino oscillations

- of external particles and their velocities
- assume Gaussian wave packets in momentum space:

$$|\phi\rangle = \int d\tilde{\mathbf{k}}\phi(\mathbf{k})|\mathbf{k}\rangle$$

• coherence properties of oscillation amplitude is determined by localization

$$\phi(\mathbf{k}) = \left(\frac{2\pi}{\sigma^2}\right)^{3/4} e^{-\frac{(\mathbf{k}-\mathbf{p})^2}{4\sigma^2}}$$



The transition amplitude

$$i\mathcal{A}_{\alpha\beta} \propto \sum_{j} U_{\alpha j} U_{\beta j}^{*} \int \frac{d^{4}p}{(2\pi)^{4}} i\tilde{\mathcal{M}}_{P} \frac{\not{p} - m_{j}}{p^{2} - m_{j}^{2} + i\epsilon} i\tilde{\mathcal{M}}_{D} e^{-ip(x_{D} - x_{P})}$$
$$\times \prod_{I=P,D} \frac{\pi^{2}}{\sigma_{pl}^{3}\sigma_{El}} \exp\left[-\frac{(\mathbf{p} - \mathbf{p}_{l})^{2}}{4\sigma_{pl}^{2}} - \frac{(p^{0} - E_{l} - \mathbf{v}_{l}(\mathbf{p} - \mathbf{p}_{l}))^{2}}{4\sigma_{El}^{2}}\right]$$

and effective momentum and energy spreads:

$$\sigma_p^2 \equiv \sum_{i,f} \sigma_{i,f}^2, \qquad \sigma_e^2 \equiv \sigma_p^2 (\Sigma - \mathbf{v}^2)$$

and a weighted velocity and velocity-squared:

$$\mathbf{v} \equiv \frac{1}{\sigma_p^2} \sum_{i,f} \sigma_{i,f}^2 \mathbf{v}_{i,f} , \qquad \Sigma \equiv \frac{1}{\sigma_p^2} \sum_{i,f} \sigma_{i,f}^2 \mathbf{v}_{i,f}^2 ,$$

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$$\mathbf{v}_{\mathbf{i}} \equiv \left. \frac{\partial E_i}{\partial \mathbf{k}_i} \right|_{\mathbf{k}_i = \mathbf{p}_i}$$

generalizations:

- finite life-time of decaying particles
- Interactions with environment (effect in analogy to collisional line broadening)





...after some algebra:

$$\overline{|\mathcal{A}_{\alpha\beta}|^{2}} \propto \exp\left[i\frac{\Delta m^{2}L}{2E_{0}}\right]$$

$$\times \exp\left[-\frac{1}{2}\left(\frac{\Delta m^{2}}{4E_{0}\sigma_{m}}\right)^{2}\right]$$

$$\times \exp\left[-\frac{1}{2}\left(\frac{\Delta m^{2}L\sigma_{\mathrm{en}}}{2E_{0}^{2}}\right)^{2}\right]$$

σ_m, σ_{en} are calculable from localization properties of initial and final state particles

standard oscillation phase

localization decoherence ξ_{loc}

energy decoherence ξ_{en}



• Localization decoherence:

$$\xi_{\rm loc} = \exp\left[-\frac{1}{2}\left(\frac{\Delta m^2}{4E_{\nu}\sigma_m}\right)^2\right] = \exp\left[-2\pi^2\left(\frac{\delta_{\rm loc}}{L_{\rm osc}}\right)^2\right] \quad \text{with} \quad \sigma_m\delta_{\rm loc} = \frac{1}{2}$$

- be identified: $\sigma_m \gg \Delta m^2 / E_{\nu}$

energy/momentum uncertainty be large enough that individual mass states cannot

• production/detection regions be localized better than oscillation length: $\delta_{loc} \ll L_{osc}$



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- be identified: $\sigma_m \gg \Delta m^2 / E_{\nu}$
- Energy decoherence:



- the neutrino energy needs to be well
- can be interpreted as neutrino wave p

• energy/momentum uncertainty be large enough that individual mass states cannot

• production/detection regions be localized better than oscillation length: $\delta_{loc} \ll L_{osc}$

$$\frac{1}{2} \left(\frac{\Delta m^2 L \sigma_{\rm en}}{2E_{\nu}^2} \right)^2 = \exp \left[-2\pi^2 \left(\frac{L}{L_{\rm osc}} \frac{\sigma_{\rm en}}{E_{\nu}} \right)^2 \right]$$

defined
$$\sigma_{\rm en} \ll E_{\nu}$$
 packet separation ($v_j \approx 1 - m_j^2/2E_{\nu}^2$)



Classical averaging

$$\overline{|\mathcal{A}_{\alpha\beta}|^2} \propto \exp\left[i\frac{\Delta m^2 L}{2E_0}\right] \exp\left[-2\pi^2\left(\frac{\delta_{\text{loc}}}{L_{\text{osc}}}\right)^2\right] \times \exp\left[-\frac{1}{2}\left(\frac{\Delta m^2 L\sigma_{\text{en}}}{2E_0^2}\right)^2\right]$$

• Classical averaging of $|\mathscr{A}_{\alpha\beta}|^2$ over space and energy have the same effect as intrinsic QM decoherence:

$$\delta_{\rm loc}^2 \to \delta_{\rm loc}^2 + \delta_{\rm clas}^2$$
, $\sigma_{\rm en}^2 \to \sigma_{\rm en}^2 + \sigma_{\rm clas}^2$

- Indistinguishable phenomenologically Kiers, Nussinov, Weiss, 1996; Stodolsky, 1998; Ohlsson, 2001
- the quantum level

• quantum mechanical uncertainties provide a fundamental lower bound on the uncertainty • to observe QM decoherence, classical averaging effects have to be suppressed down to



R. Krüger, TS, 2303.15524 Numerical estimates for gallium and reactors

- nuclei localization: inter-atomic distances in cristal lattice or fluid
- velocities: either thermal velocities or velocities for typical kin. energies of outgoing particles
- electron/positron localization: distance the particle travels until it deposits one mean excitation energy (integrating $\langle dE/dx \rangle$ Bethe equation)

	Particle	$\delta_x \text{ [nm]}$	$\sigma [{\rm eV}]$	v
Reactor (P)	N	0.24	410	1×10
$N \to N' + e^- + \overline{\nu}_e$	N'	0.24	410	4×10
	e^{-}	260	0.38	0.99
Reactor (D)	p	0.1	990	5×10
$p + \overline{\nu}_e \to n + e^+$	n	5×10^{6}	2×10^{-5}	5×10
	e^+	320	0.3	0.99
Gallium (P)	Cr	0.20	480	7×10
$\mathrm{Cr} \to \mathrm{V} + \nu_e$	V	0.20	480	2×10
Gallium (D)	Ga	0.27	370	6×10
$Ga + \nu_e \rightarrow Ge + e^-$	Ge	0.27	370	1×10
	e^{-}	310	0.32	0.83







Numerical estimates for gallium and reactors

Incalization decoherence:

for reactor and gallium experiments dominated by hadronic particles

$$\sigma_m \simeq (400 - 500) \,\mathrm{eV}\,, \qquad \delta_{\mathrm{loc}} = \frac{1}{2\sigma_m} \simeq 0.2 \,\mathrm{nm}$$
$$-\ln\xi_{\mathrm{loc}} = \frac{1}{2} \left(\frac{\Delta m^2}{4E_{\nu}\sigma_m}\right)^2 \approx 1.3 \times 10^{-19} \left(\frac{\Delta m^2}{1 \,\mathrm{eV}^2}\right)^2 \left(\frac{1 \,\mathrm{MeV}}{E_{\nu}}\right)^2 \left(\frac{500 \,\mathrm{eV}}{\sigma_m}\right)^2$$

 \Rightarrow QM localization decoherence irrelevant for all practical purposes: $\xi_{loc} = 1$ Note: classical spatial averaging is relevant and needs to be taken into account





Numerical estimates for gallium and reactors

• energy decoherence:

including phase space integration and Doppler broadening:

$$-\ln \xi_{\rm en} = 2\pi^2 \left(\frac{L}{L_{\rm osc}} \frac{\sigma_{\rm en}}{E_{\nu}}\right)^2 \approx 4.9 \times 10^{-12} \left(\frac{L}{L_{\rm osc}}\right)^2 \left(\frac{1\,{\rm MeV}}{E_{\nu}}\right)^2 \left(\frac{\sigma_{\rm en}}{0.5\,{\rm eV}}\right)^2$$

purposes: $\xi_{\rm en} = 1$

Note: classical energy averaging is relevant and needs to be taken into account

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 $\sigma_{\rm en} \approx 0.5 \, {\rm eV} \,, \qquad \delta_{\rm en} \approx 200 \, {\rm nm}$

 \Rightarrow QM energy decoherence ("wave packet separation") irrelevant for all practical





Summary energy decoherence

• our result: $\sigma_{en} \approx 0.5 \,\mathrm{eV}$, $\delta_{en} \approx 200 \,\mathrm{nm}$

energy resolution of typical reactor neutrino detectors:

 $(0.03 - 0.06) \,\mathrm{MeV}\sqrt{E/\mathrm{MeV}}$

about 6 orders of magnitude larger than σ_{en} !

• phenomenological constraint is dominated by classical energy resolution: $\sigma < 0.47 \,\mathrm{MeV}, \delta > 2.1 \times 10^{-4} \,\mathrm{nm}$ [Gouvea, Romeri, Ternes, 2005.03022, 2104.05806]

R. Krüger, TS, 2303.15524 s. also Akhmedov, Smirnov, 2208.03736

$$\Rightarrow \sigma_{clas} \simeq 0.1 \text{ MeV}$$

QM decoherence negligible — dominated by classical averaging (holds also for JUNO)







Any decoherence effect on top of classical averaging will point towards new physics!

Barenboim, Mavromatos, Sarkar, Waldron-Lauda, 2006 Fogli, Lisi, Marrone, Montanino, Palazzo, 2007 Farzan, TS, Smirnov, 2008; Bakhti, Farzan, TS, 2015 Guzzo, de Holanda, Oliveira, 2014 Hellmann, Pas, Rani, 2022 Banks, Kelly, McCullough, 2023



 three SM neutrinos (no steriles) modified QM evolution:

 $\frac{d\rho}{dt} = i[\rho, H] - \{\rho, D^2\} + 2D\rho D$

$$H = \frac{1}{2E_{\nu}} \operatorname{diag}(m_1^2, m_2^2, m_3^2) ,$$

 $D = \operatorname{diag}(d_1, d_2, d_3),$





 three SM neutrinos (no steriles) modified QM evolution:

 $\frac{d\rho}{dt} = i[\rho, H] - \{\rho, D^2\} + 2D\rho D$ • damping of interference terms:

$$P_{ee} = \sum_{i=1}^{3} |U_{ei}|^4 + \sum_{i \neq j} |U_{ei}|^2 |U_{ej}|^2$$

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$$H = \frac{1}{2E_{\nu}} \operatorname{diag}(m_1^2, m_2^2, m_3^2) ,$$

 $D = \operatorname{diag}(d_1, d_2, d_3),$

$$\gamma_{ij} = (d_i - d_j)^2$$
$$= \frac{1}{\lambda_{ij}} \left(\frac{E_{\text{ref}}}{E_{\nu}}\right)$$











$$r = 12$$

$$\frac{\chi^2_{\min}/\text{dof} \quad p\text{-val.} \quad \Delta\chi^2 \quad \#\sigma \quad \lambda_{12}}{2.2/1 \quad 0.14 \quad 32.5 \quad 5.4 \quad 1.}$$
all
$$9.2/5 \quad 0.10 \quad 29.2 \quad 5.0 \quad 1.$$

$$P_{ee}^{\text{gal}} \approx 1 - \frac{1}{2} \sin^2 2\theta_{13} - \frac{1}{2} \cos^4 \theta_{13} \sin^2 2\theta_{12} \left(1\right)$$









Consistency with other oscillation data





Consistency with other oscillation data







solar and reactor neutrinos require steep energy dependence of decoherence effects: $\gamma \propto E_{\nu}^{-r}$, $r \gtrsim 10 - 12 \Rightarrow$ other oscillation evidences not affected



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- predict distance dependence in gallium exps. (still compatible with BEST) slightly modified MSW transition region in $P_{ee}(E_{\nu})$ for solar neutrinos



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- can be made compatible with decoherence explanation of LSND Bakhti, Farzan, TS [1503.05374] either
 - different decoherence parameters for neutrinos and antineutrinos, or
 - play with energy dependence: $d_1 \approx 0, d_2$ peak @ 0.75 MeV, d_3 peak @ 30 MeV



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• no effect predicted in MiniBooNE or SBL reactors



Summary

• BEST results make gallium anomaly highly significant $\approx 5\sigma$

• standard QM localization/wave packet effects are negligible

 presented an explanation in terms of new-physics decoherence, explanation

- sterile neutrino explanation in tension with reactors and solar neutrinos

 - consistent with oscillation data, potentially compatible with an LSND



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