# Measurement of strange baryon production in chargedparticle jets in pp and p-Pb collisions with ALICE

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**Abstract.** The enhancement of the ratio of strange to non-strange hadrons with increasing charged-particle multiplicity density is often interpreted as a signature of the medium formed in heavy-ion collisions. However, this effect has also been observed in smaller systems, such as pp and p–Pb collisions. Understanding this effect requires a detailed description of the production mechanisms of strange hadrons in and out of jets. In this manuscript published results on the production of  $K_S^0$ ,  $\Lambda(\bar{\Lambda})$ ,  $\Xi^{\pm}$  and  $\Omega^{\pm}$  hadrons in charged-particle jets and the underlying event in pp and p-Pb collisions measured by ALICE are presented. In addition, a new analysis of in-jet fragmentation into  $\Lambda$  and  $K_S^0$  hadrons in pp collisions at  $\sqrt{s} = 13.6$  TeV from LHC Run 3 data is shown.

## 1 Introduction

In heavy-ion collisions, a medium of deconfined quarks and gluons is formed. One signature of this medium is the enhancement of the ratio of strange to non-strange hadrons with increasing charged-particle multiplicity density. However, this effect has also been observed in smaller systems, such as pp and p–Pb collisions, where such a medium is not expected.

The yields of  $K_S^0$ ,  $\Lambda(\bar{\Lambda})$ ,  $\Xi^{\pm}$ , and  $\Omega^{\pm}$  hadrons were measured as a function of chargedparticle multiplicity in pp, p–Pb and Pb–Pb collisions by the ALICE Collaboration [1]. The yields of these strange hadrons were found to increase with multiplicity at a rate faster than that of the pion yields. Questions remain on the mechanisms driving the production of these hadrons and their increase with increasing charged-particle multiplicity. In these proceedings, possible ways to disentangle the strangeness production mechanisms in jets and in the underlying event (UE) are investigated. A published analysis [2] on the production of  $K_S^0$ ,  $\Lambda(\bar{\Lambda})$ ,  $\Xi^{\pm}$ , and  $\Omega^{\pm}$  in charged-particle jets and in the UE in pp and p–Pb collisions is presented. In particular, this is the first measurement of  $\Omega^{\pm}$  production in jets. Additionally, a new analysis is discussed on in-jet fragmentation into  $K_S^0$  and  $\Lambda(\bar{\Lambda})$  using Run 3 ALICE data.

# 2 Strange hadron yields in charged-particle jets

Jets are reconstructed using the anti- $k_{\rm T}$  algorithm [3, 4], with a jet resolution parameter R = 0.4, using tracks with transverse momentum  $p_{\rm T,track} > 0.15$  GeV/c and pseudorapidity  $|\eta_{\rm track}| < 0.8$ , across the full acceptance in azimuth  $\varphi$ . The pseudorapidity of reconstructed

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jets is required to be in the range  $|\eta_{jet}| < 0.35$ , ensuring the jets are fully contained within the acceptance region of reconstructed weak decays. In addition, a transverse momentum cut of  $p_{T,jet} > 10 \text{ GeV}/c$  is applied to ensure the jet originates from a hard scattering process.

This analysis considers the V0 particles ( $K_S^0$ ,  $\Lambda$ , and  $\bar{\Lambda}$ ), as well as the cascades (here meaning  $\Xi^{\pm}$  and  $\Omega^{\pm}$ ). The following decay channels were searched for, with their corresponding branching ratios (*BR*):

$$\begin{split} & K^0_{\rm S} \to \pi^+ + \pi^- & BR = (69.20 \pm 0.05)\%, \\ & \Lambda(\bar{\Lambda}) \to p(\bar{p}) + \pi^-(\pi^+) & BR = (63.9 \pm 0.5)\%, \\ & \Xi^-(\bar{\Xi}^+) \to \Lambda(\bar{\Lambda}) + \pi^-(\pi^+) & BR = (99.887 \pm 0.035)\%, \\ & \Omega^-(\bar{\Omega}^+) \to \Lambda(\bar{\Lambda}) + {\rm K}^-({\rm K}^+) & BR = (67.8 \pm 0.7)\%. \end{split}$$

The V0 and cascade candidates are reconstructed based on the topology of their weak decays, within a pseudorapidity range of  $|\eta| < 0.75$ . The purity is improved by identifying charged tracks based on their energy deposition in the Time Projection Chamber subdetector. The invariant mass is used to extract the signal and combinatorial background. In addition, candidates that satisfy the reconstruction criteria for multiple particles, are rejected.

Particles are defined as within the jet cone (JC) if their distance in the  $\eta - \varphi$  plane with respect to the jet axis is less than the jet resolution parameter. This includes both particles associated with the hard scattering and particles from the UE. The UE contribution is estimated with the perpendicular cone method (PC) [2], where the yields of strange hadrons are measured within a cone of size *R*, placed at the same  $\eta$  as the jet and at an angle  $\Delta \varphi = \pi/2$ .

The  $p_{\rm T}$ -differential yields  $d\rho/dp_{\rm T}$  of the particles are calculated by

$$\frac{\mathrm{d}\rho}{\mathrm{d}p_{\mathrm{T}}} = \frac{1}{N_{\mathrm{ev}}} \times \frac{1}{A_{\mathrm{acc}}} \times \frac{\mathrm{d}N}{\mathrm{d}p_{\mathrm{T}}},\tag{1}$$

where  $dN/dp_T$  is the  $p_T$ -differential particle production yields,  $N_{ev}$  is the number of events,  $A_{acc}$  is the area of the  $\eta - \varphi$  acceptance for a given selection. The  $p_T$ -differential yields of in-jet particles can then be extracted by subtracting the PC selected yields from the JC selected yields.

Figure 1 shows the baryon-to-meson and baryon-to-baryon ratios for the measured strange particles in pp collisions. For each of these ratios, the UE yields closely follow the inclusive yields. For the baryon-to-meson ratios and  $\Xi/\Lambda$ , an enhancement is seen in the inclusive and UE yields with respect to the in-jet yields. For the  $\Omega/\Lambda$  and  $\Omega/\Xi$  ratios, this enhancement is not seen, but the uncertainties are too large to make any definitive statements.

Figure 2 shows the  $\Xi/K_S^0$  yield ratio in p–Pb collisions for different multiplicity bins: high (0-10%), intermediate (10-40%) and low (40-100%). The UE yields show an enhancement with increasing multiplicity, whereas the in-jet particle yields are consistent within uncertainties across multiplicity. A similar enhancement is found for the ratios of  $\Lambda/K_S^0$  and  $\Xi/\Lambda$ . Due to the low number of  $\Omega$  candidates in the studied data sample, the ratios of  $\Omega/K_S^0$ ,  $\Omega/\Lambda$ , and  $\Omega/\Xi$  were only measured for Minimum Bias (MB) p–Pb collisions. The results are found to be consistent within uncertainties for MB pp and MB p–Pb collisions.

#### 3 Jet fragmentation

A different approach to studying the production of strange hadrons in jets is to measure the jet fragmentation functions. The jet fragmentation function for a hadron is defined as the number density of hadrons in a jet, that carry a fraction z of the longitudinal momentum of the jet. The



Figure 1:  $p_{\rm T}$ -dependent baryon-to-meson (top) and baryon-to-baryon (bottom) yield ratios for inclusive, in-jet, and PC selected particles in pp collisions.



Figure 2:  $\Xi/K_S^0$  yield ratios for inclusive, in-jet, and PC selected particles in p–Pb collisions in different multiplicity classes.

longitudinal momentum fraction is calculated by projecting the particle 3-momentum vector  $\mathbf{p}_i$  onto that of the jet ( $\mathbf{p}_{jet}$ ) and normalising to the jet momentum

$$z_i = \frac{\mathbf{p}_i \cdot \mathbf{p}_{jet}}{\mathbf{p}_{iet}^2}.$$
 (2)

Due to the relatively long lifetimes of  $K_S^0$  (2.6844 cm/*c*) and  $\Lambda$  (7.89 cm/*c*) [5], their decay daughters are removed from the jet clustering input by removing tracks with a large distance of closest approach to the primary vertex. Therefore, the V0 candidates are reconstructed separately and added to the jet clustering input, yielding a collection of charged-particle + V0 jets (ch+V0 jets).

To illustrate the performance of the ch+V0 jet approach in Run 3 ALICE data, figure 3 shows the invariant mass spectra of V0 particles in ch+V0 jets around the  $K_S^0$  invariant mass. V0 particles with 0.5 < z < 0.7 were selected in jets with  $10 < p_{T,jet} < 20$  GeV/c (left) and  $20 < p_{T,jet} < 40$  GeV/c (right). The clearly visible invariant mass peaks with large numbers of candidates in the Run 3 data enables precise measurements of V0 fragmentation distributions in jets and study of the dependence on the jet momentum.



Figure 3: Invariant mass spectrum of V0 candidates in ch+V0 jets around the K<sub>S</sub><sup>0</sup> mass, with  $0.5 < z_{V0} < 0.7$ . Left:  $10 < p_{T,jet} < 20$  GeV/c, right:  $20 < p_{T,jet} < 40$  GeV/c.

## 4 Conclusions

Studying the production of strange hadrons in jets can provide insight into the mechanisms that drive the increase of the strange to non-strange hadron ratio with charged-particle multiplicity density. In these proceedings, the first measurement of  $\Omega^{\pm}$  baryons in jets was presented, as well as results on  $K_S^0$ ,  $\Lambda$  ( $\bar{\Lambda}$ ), and  $\Xi^{\pm}$  yields in jets, in pp and p–Pb collisions. It is observed that the baryon-to-meson and the 2-strange-to-1-strange yield ratios show an enhancement with multiplicity in the UE, but not in jets. Also, an approach to measuring in-jet fragmentation into  $K_S^0$  and  $\Lambda$  ( $\bar{\Lambda}$ ) in pp collisions using jets constructed from both charged particles and reconstructed V0 candidates was shown. This approach showcases the possibility to identify V0 particles at high *z* in Run 3 ALICE data.

# References

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