

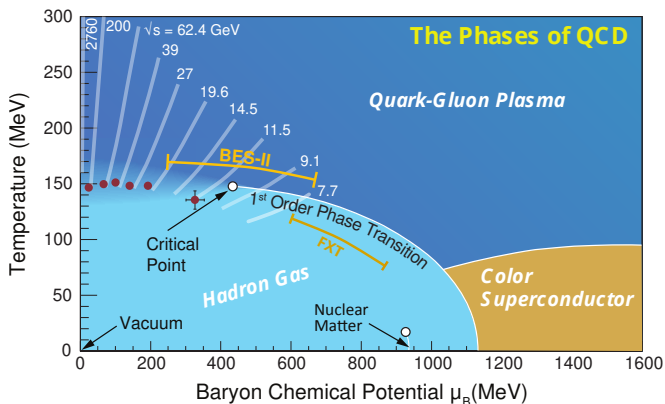
QCD EOS, critical point, and hydrodynamic fluctuations

M. Stephanov



QCD critical point

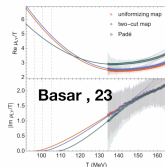
Where on the QCD phase boundary is the CP?



Motivation for BES at RHIC and BEST topical collaboration.

Latest theory developments on locating CP

From Maneesha Pradeep's talk at CPOD 2024:



G. Basar, Fri,
11:40 am

Extrapolations of Lee-Yang edge singularities to real axis

$$(\mu_{BC}, T_c) \approx (580, 100) \text{ MeV}$$

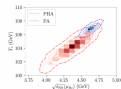
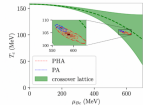
Bayesian Holography + Lattice input at $\mu = 0$

[Higbert et al. e-Print: 2309.00579 \[nucl-th\]](#)

Predict CFP (95% confidence level):

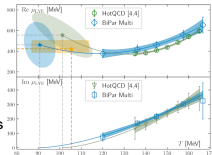
$$T_c = 101 - 108 \text{ MeV} \quad \mu_c = 560 - 625 \text{ MeV}$$

$$\sqrt{s} = 4.0 - 4.8 \text{ GeV}$$



J. Noronha,
Tue, 11:40
am

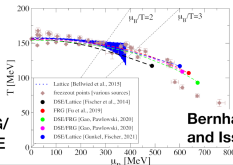
(Ratti's talk)



C. Schmidt,
Tue, 11:00
am

Clarke et al.,
24

$$(\mu_{BC}, T_c) = (422^{+80}_{-35}, 105^{+8}_{-18}) \text{ MeV}$$



C. Fischer,
Tue, 11:20
am

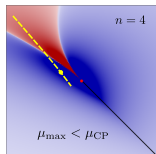
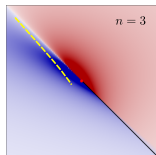
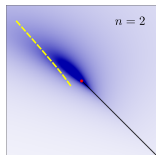
FRG/
DSE

Bernhardt, Fischer
and Isserstedt, 23

$$(\mu_{BC}, T_c) = (495 - 654, 108 - 119) \text{ MeV}$$

Theory vs BES-II data

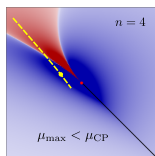
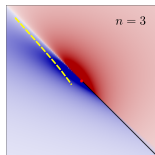
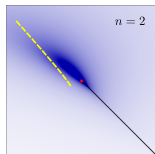
(universal EOS) critical χ_n :



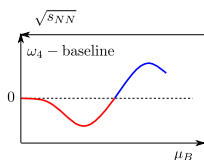
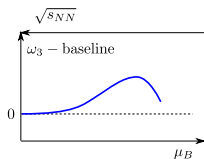
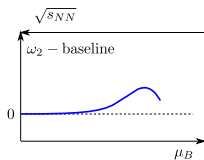
Bzdak et al review 1906.00936

Theory vs BES-II data

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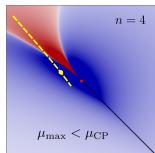
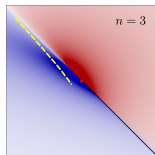
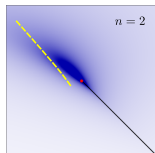
(irreducible correlations) $FC_n[N_p] \sim \chi_n$ (Pradeep, MS 2211.09142), $\omega_n \equiv FC_n/FC_1$



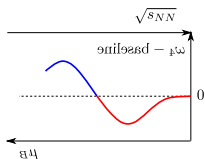
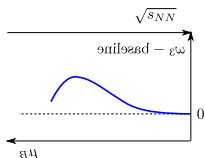
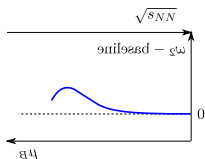
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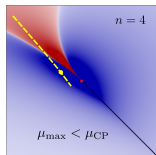
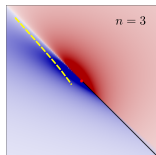
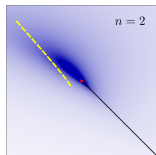


Bzdak et al review 1906.00936

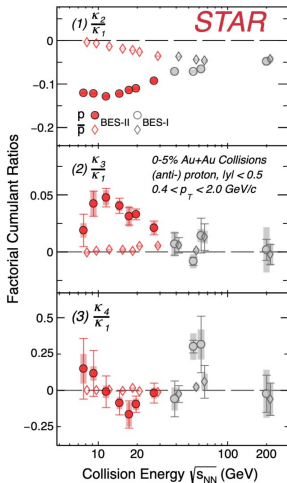
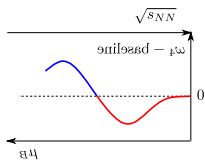
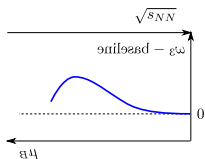
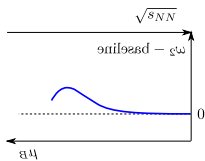
Expected signatures: **bump** in ω_2 and ω_3 , **dip** then **bump** in ω_4

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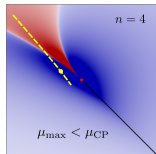
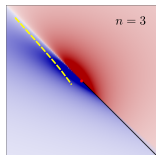
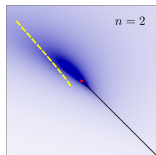


Bzdak et al review 1906.00936

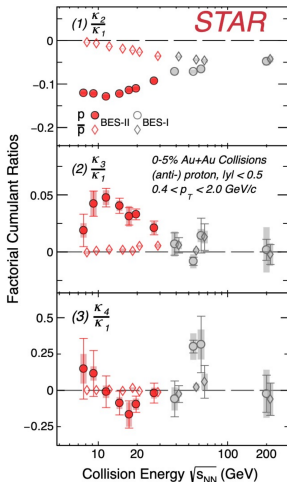
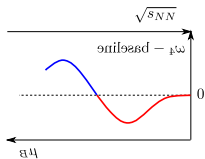
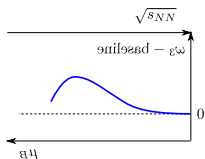
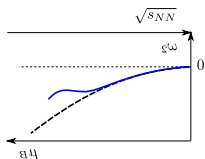
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Bzdak et al review 1906.00936

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BEST Framework

The goal of BES theory: connect observables to QCD phase diagram.

BEST framework: An et al (40+ authors, 100+ pp, 369 refs) [2108.13867](#)

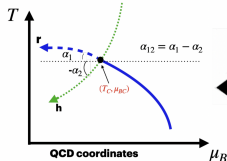
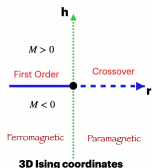
BES theory review: Du, Sorensen, MS [2402.10183](#)

- 🟢 Lattice EOS + CP \rightarrow parametric EOS
- 🔴 EOS \rightarrow Hydrodynamics with (non-gaussian) fluctuations.
- 🔴 Freezeout, including fluctuations. *reviewed in [2403.03255](#)*
- 🟢 Comparison with experiment. Bayesian analysis (ML).
Determine/constrain EOS, critical point parameters.

Parametric EOS (now with T' -expansion)

From Maneesha Pradeep's review talk at CPOD 2024:

$$P_{\text{QCD}}(\mu, T) = P_{\text{BG}}(\mu, T) + A G(r(\mu, T), h(\mu, T))$$



Independent & non-universal parameters

$$\mu_c, \alpha_{12}, \rho, w$$

Weakly constrained in the chiral limit

MP, Stephanov, 19

Kahangirwe et al., 24

7

Range of Validity improved

$$0 \leq \mu_B \leq 700 \text{ MeV}, 25 \text{ MeV} \leq T \leq 800 \text{ MeV}$$

The new construction is causal and stable for a larger range of ρ and w

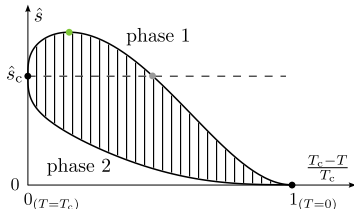
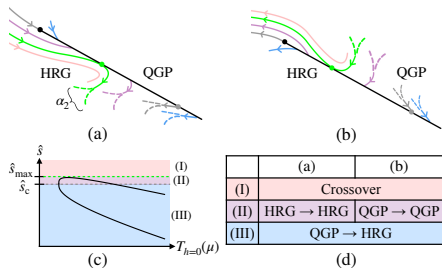
**M. Kahangirwe,
Wed, 12:10 pm**

*Kahangirwe et al 2402.08636, PRD
(also Jahan's and Karthein's talks)*

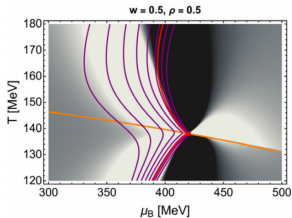
Critical point and non-trivial hydro trajectories

Pradeep, Sogabe, MS, Yee 2402.09519, PRC

- $\hat{s} \equiv s/n$ is non-monotonic along coexistence (1st order) line
- non-trivial deformation of trajectories



explains “lensing”, “cusp”



Critical lensing—Dore et al,22,
Nonaka&Asakawa, 05

depending on $(\partial P/\partial T)_n$ at CP

Deterministic approach to non-Gaussian fluctuations

non-Gaussian fluctuations are non-trivial and sensitive signatures of the critical point

● *Infinite* hierarchy of coupled equations *An et al 2009.10742, PRL*

for connected hydro correlators $H_n \equiv \underbrace{\langle \delta\psi \dots \delta\psi \rangle}_{n}^{\text{connected}}$.

$$\partial_t \psi = -\nabla \cdot \text{Flux}[\psi, H, H_3, H_4, \dots];$$

$$\partial_t H = \mathbf{F}[\psi, H, H_3, H_4, \dots];$$

$$\partial_t H_3 = \mathbf{F}_3[\psi, H, H_3, H_4, \dots];$$

⋮

Controlled perturbation theory

An et al [2009.10742](#), PRL

- Small fluctuations are *almost* Gaussian
- Introduce expansion parameter ε , so that $\delta\psi \sim \sqrt{\varepsilon}$.

Then $H_n \equiv \varepsilon^{n-1}$ and to leading order in ε :

$$\partial_t \psi = -\nabla \cdot (\text{Flux}[\psi] + \mathcal{O}(\varepsilon));$$

$$\partial_t H = -2\Gamma(H - \bar{H}[\psi]) + \mathcal{O}(\varepsilon^2);$$

\vdots

$$\partial_t H_n = -n\Gamma(H_n - \bar{H}_n[\psi, H, \dots, H^{n-1}]) + \mathcal{O}(\varepsilon^n);$$

To leading order, the equations are iterative and “linear”.

- In hydrodynamics the small parameter is $(q/\Lambda)^3$, i.e., fluctuation wavelength $1/q \gg$ size of hydro cell $1/\Lambda$ (UV cutoff).

Diagrammatic representation

An et al [2009.10742](#), [2212.14029](#), An's talk at CPOD 2024

Leading order in $\varepsilon \Leftrightarrow$ tree diagrams.

$$\left(\text{---}\bullet \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} \right)^* = \begin{array}{c} \text{---}\text{D} \begin{array}{c} \bullet \\ \vdots \\ \bullet \end{array} \\ \text{drift} \end{array} + \begin{array}{c} \text{---}\Delta \text{---} \\ \text{---}\bullet \text{---} \\ \text{---}\bullet \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \quad \text{all combinatorial configurations of trees}$$

Loops describe feedback of fluctuations (renormalization and long-time tails).

$$\left(\text{---}\bullet \right)^* = \text{---}\text{D} + \text{---}\text{D} \begin{array}{c} \text{---} \\ \text{---} \end{array} \quad \text{1-pt equation including leading loop}$$

conventional hydro equations *one loop (renormalization & long-time tails)*

Generalizing Wigner transform

An et al [2009.10742](#), PRL

Definition:

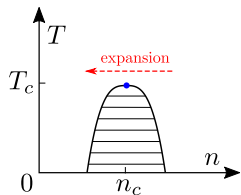
$$W_n(\mathbf{x}; \mathbf{q}_1, \dots, \mathbf{q}_n) \equiv \int d\mathbf{y}_1^3 \dots \int d\mathbf{y}_n^3 H_n(\mathbf{x} + \mathbf{y}_1, \dots, \mathbf{x} + \mathbf{y}_n) \delta^{(3)}\left(\frac{\mathbf{y}_1 + \dots + \mathbf{y}_n}{n}\right) e^{-i(\mathbf{q}_1 \cdot \mathbf{y}_1 + \dots + \mathbf{q}_n \cdot \mathbf{y}_n)},$$



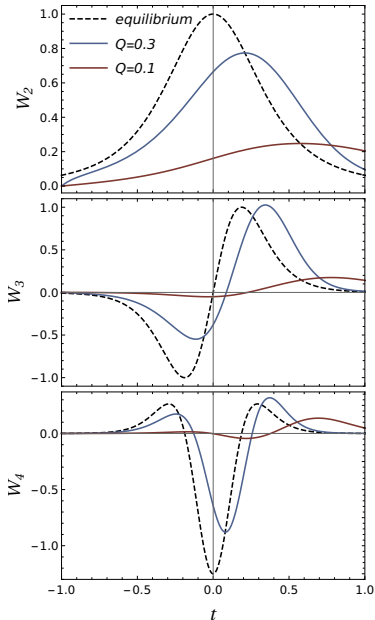
W_n 's quantify magnitude and non-gaussianity of fluctuation harmonics with wave-vectors \mathbf{q}_i .

Example: expansion through a critical region

An et al [2009.10742](#), PRL



- Two main features:
 - Lag, "memory".
 - Smaller Q – slower evolution.
- Conservation laws.



Freezeout of fluctuations

- Freezeout: translation of correlators of hydrodynamic fluctuations ($\psi = \epsilon, n_B, u$)

$$\langle \delta\psi \dots \delta\psi \rangle = H_n(\mathbf{x}_1, \dots, \mathbf{x}_n)$$

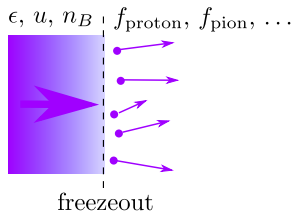
to particle correlators

$$\langle \delta f \dots \delta f \rangle = G_n(\mathbf{x}_1, \mathbf{p}_1, \dots, \mathbf{x}_n, \mathbf{p}_n).$$

- Conservation laws relate \mathbf{p} integrals of G_n to H_n .

- But the \mathbf{p} dependence in G_n is not constrained.

There are ∞ many possibilities/solutions (G_n) satisfying conservation laws.



Maximum entropy freezeout

Pradeep, MS, [2211.09142](#), PRL

- There is a unique solution which maximizes the entropy!
 - for $n = 1$ equivalent to Cooper-Frye
 - for critical fluctuations similar to the σ field coupling
 - but applies much more generally
 - model independent, i.e., determined by QCD EOS

$$\underbrace{\hat{\Delta}G_{ABC}}_{\text{irreducible particle correlations (FC)}} = \underbrace{\hat{\Delta}H_{abc}}_{\text{hydrodynamic correlations}} \underbrace{(\bar{H}^{-1}P\bar{G})_A^a (\bar{H}^{-1}P\bar{G})_B^b (\bar{H}^{-1}P\bar{G})_C^c}_{\text{kinematic factors}}$$

- Work in progress – implement in a hydro model and estimate *nonequilibrium* expectations for multiplicity cumulants in BES

Karthein, Pradeep, MS, Rajagopal, Yin; Karthein's talk

Summary

- BES-II data is in.
Qualitatively agrees with non-monotonic expectations from CP, not only in $n = 4$ factorial cumulant, but $n = 3$ and $n = 2$.
- To produce such signatures the CP has to be at $\mu_B > 420$ MeV. Agreement with recent theory estimates by different approaches.
- To convert these qualitative statements into quantitative ones, i.e., provide constraints on the QCD EOS from BES-II data more work is needed and is underway.

More

Factorial Cumulants are better experimental measures

Three reasons:

- Normal cumulants (NC) measure non-gaussianity;
Factorial cumulants (FC) measure non-poissonianity,
(irreducible particle correlations).

NCs are for densities (continuous);

FCs are for multiplicities (discrete).

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FCs are powers of Δy for small Δy ; NCs are polynomials.

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- Acceptance dependence:

FCs are powers of Δy for small Δy ; NCs are polynomials.

- Maximum Entropy freezeout (*Pradeep, MS 2211.09142*):

FCs of multiplicities are directly related to hydrodynamic correlators (or susceptibilities in thermodynamics).