

# Charm and beauty production in hadronic collisions via muon measurements at forward rapidity with ALICE

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## Outline

- ❖ Motivations
- ❖ The ALICE detector
- ❖ Latest measurements
- ❖ Run 3 perspectives



Heavy (charm and beauty) quarks: **sensitive probes of the quark-gluon plasma (QGP)**

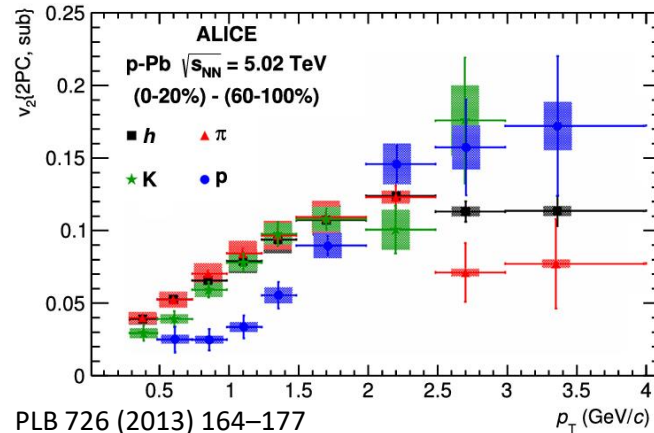
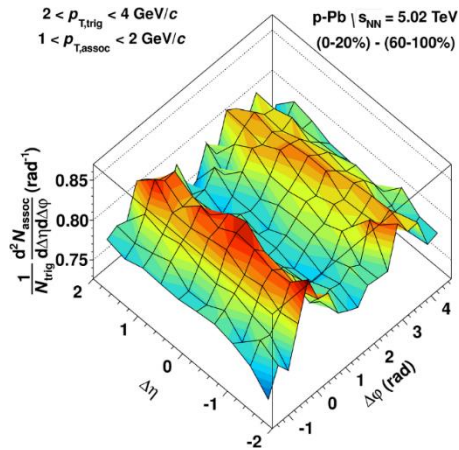
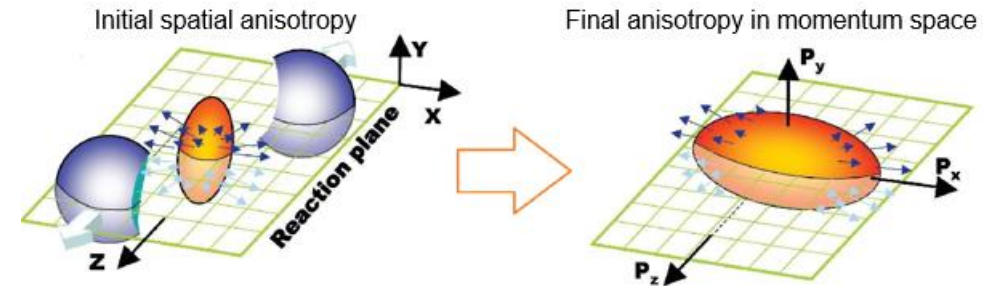
- $\tau_{c/b} \sim 0.01-0.1$  (fm/c)  $< \tau_{QGP} \sim 0.3$  (fm/c) [PRC 89 (2014) 034906]
- Experience the **whole collision evolution**

❖ **Key observable: azimuthal anisotropy** quantified by means of a Fourier expansion of azimuthal distributions of produced particles

$$\frac{dN}{d\varphi} \propto 1 + \sum_{n=1}^{\infty} 2v_n \cos[n(\varphi - \Psi_n)]$$

$$v_n = \langle \cos n(\varphi - \Psi_n) \rangle$$

$n = 2$ , elliptic flow



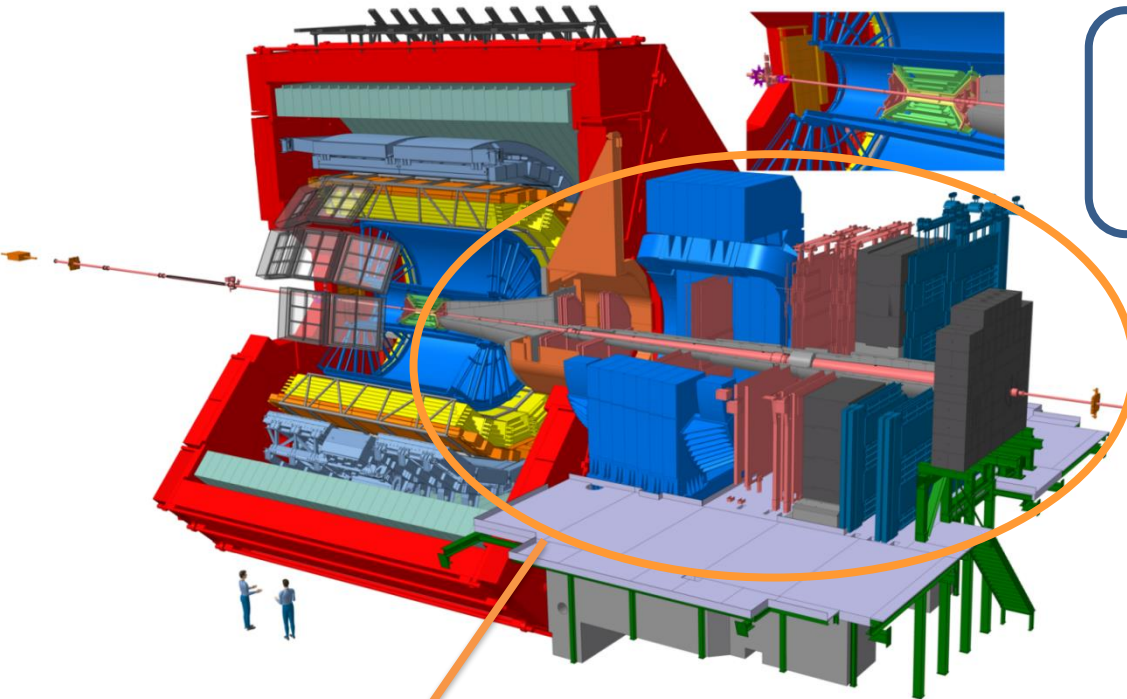
**Small collision systems (pp & p-Pb collisions):**

- ❖ Baseline for heavy-ion collisions
- ❖ pp collisions: test pQCD-based calculations and production mechanisms
- ❖ p-Pb collisions: cold nuclear matter effects and study of nuclear parton distribution functions

**No (or very tiny) QGP effect is expected**

❖ Long-range angular correlations and clear mass ordering observed in p-Pb collisions **at high multiplicity**, as in Pb-Pb collisions

# The ALICE detector (Run 2 layout)

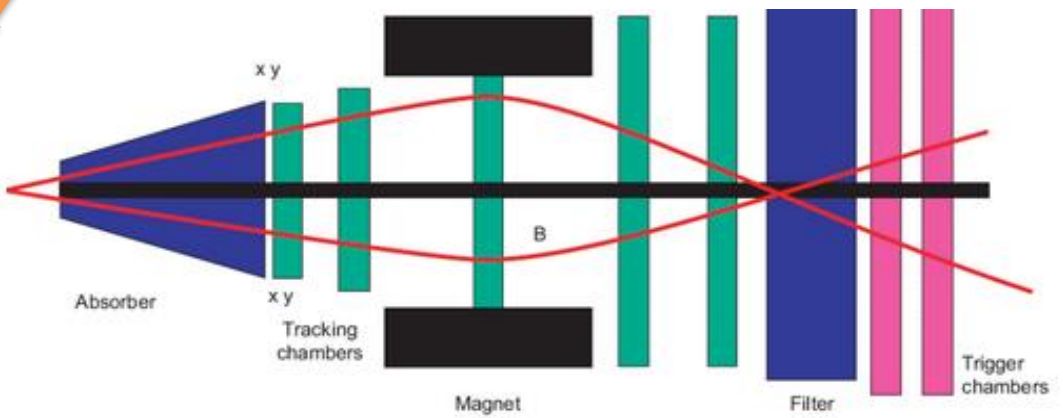


**Central Barrel,  $|\eta| < 0.9$**   
 Vertexing (ITS),  
 Tracking (ITS, TPC),  
 PID (ITS, TPC, TOF, TRD, HMPID,  
 Calorimeters)

**Forward/  
 Backward  
 detectors**  
 (V0, T0, ZDC):  
 Trigger,  
 Timing,  
 Multiplicity/  
 Centrality,  
 Event plane

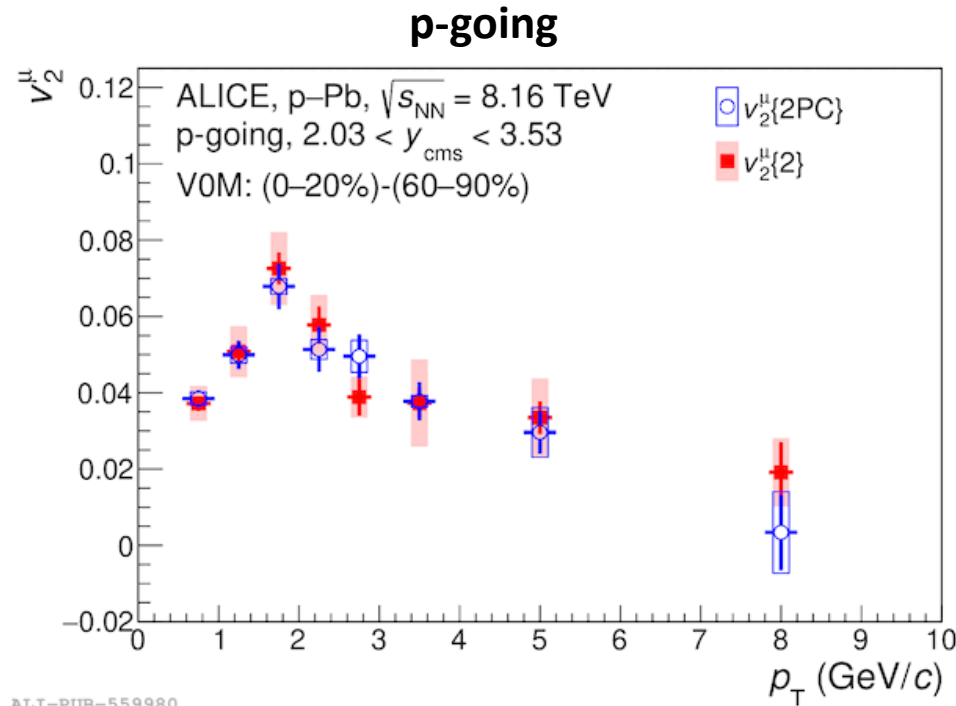
- Hadronic decays ( $|y| < 0.8$ )**
- $D^0 \rightarrow K^- \pi^+$
  - $D^+ \rightarrow K^- \pi^+ \pi^+$
  - $D^{*+} \rightarrow D^0 (\rightarrow K^- \pi^+) \pi^+$
  - $D_s^+ \rightarrow \phi (\rightarrow K^- K^+) \pi^+$
  - $D_{s1}^+ \rightarrow D^{*+} K_s^0$
  - $D_{s2}^{*+} \rightarrow D^+ K_s^0$
  - $\Lambda_c^+ \rightarrow p K_s^0, \Lambda_c^+ \rightarrow p K^- \pi^+$
  - $\Lambda_c^+ \rightarrow e^+ \Lambda \nu_e$
  - $\Xi_c^0 \rightarrow e^+ \Xi^- \nu_e, \Xi_c^0 \rightarrow \pi^+ \Xi^-$
  - $\Xi_c^+ \rightarrow \pi^+ \pi^+ \Xi^-$
  - $\Omega_c^0 \rightarrow \Omega^- \pi^+$
  - $\Sigma_c^{0,++}(2455) \rightarrow \Lambda_c^+ \pi^-$
  - $\Sigma_c^{0,++}(2520) \rightarrow \Lambda_c^+ \pi^-$
- Semi-leptonic decays**
- **$c, b \rightarrow \mu^\pm (2.5 < y < 4.0)$**
  - $c, b \rightarrow e^\pm (|y| < 0.8 \text{ or } 0.6)$

**Muon spectrometer**  
 $-4 < \eta < -2.5$   
 Tracking,  
 Triggering,  
 $\mu$  identification

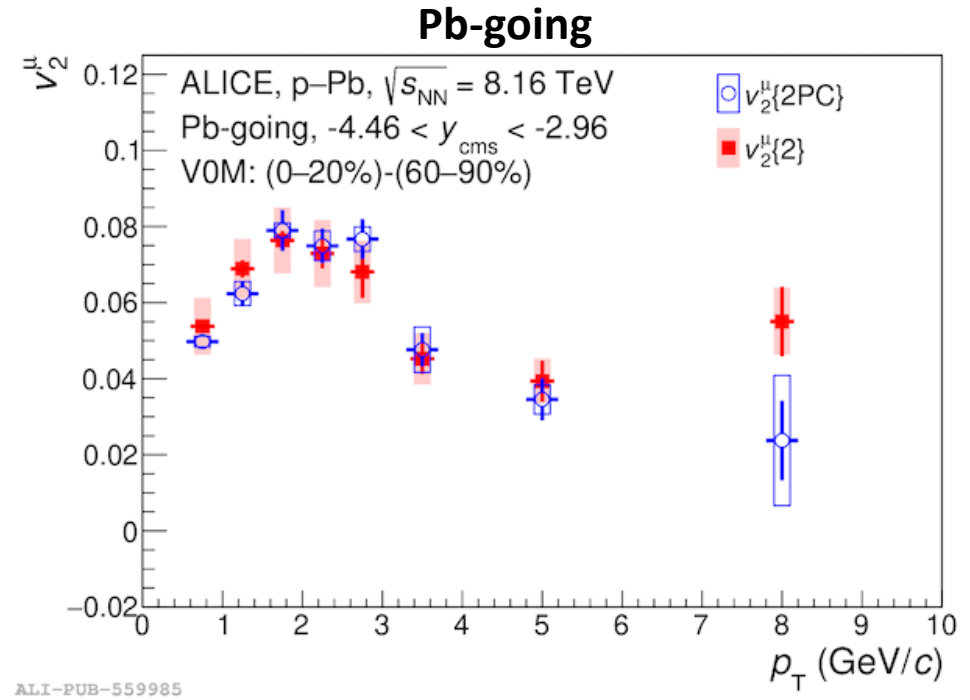


*Int. J. Mod. Phys. A29 (2014) 1430044*

# Heavy-flavour hadron decay muon $v_2$ in p–Pb collisions



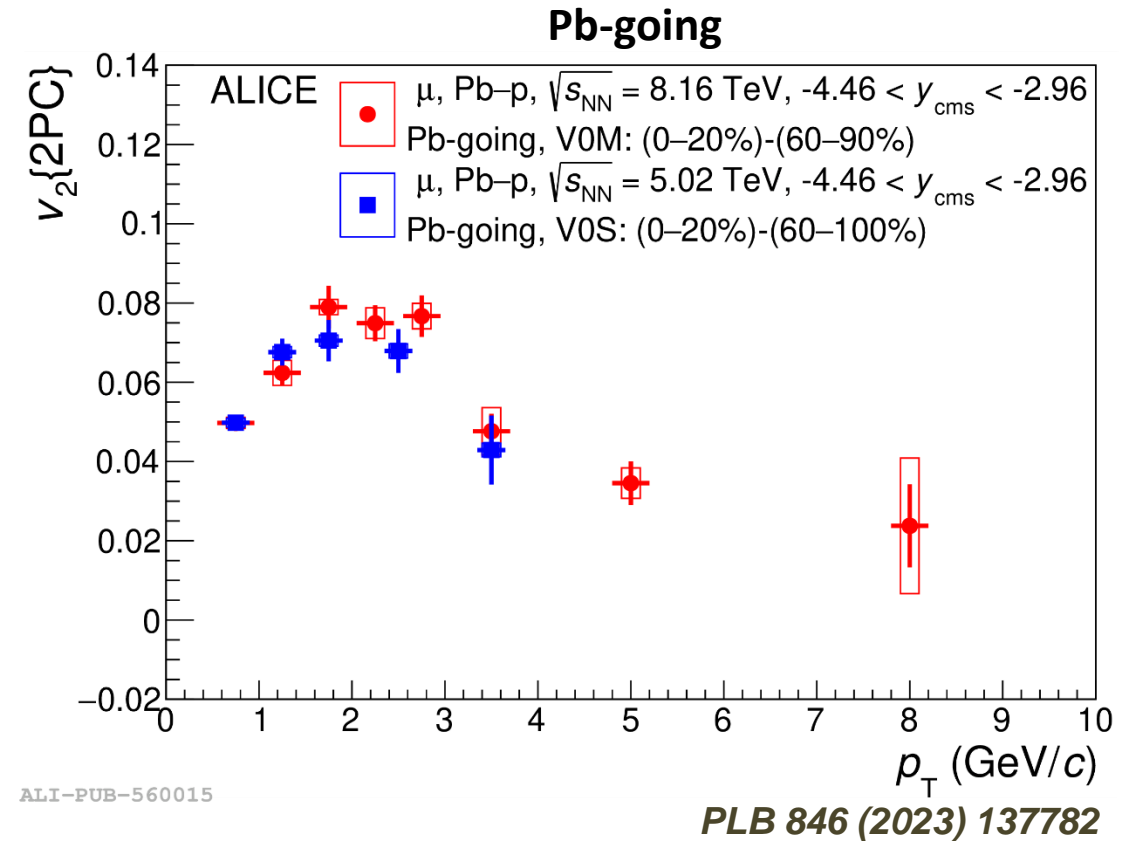
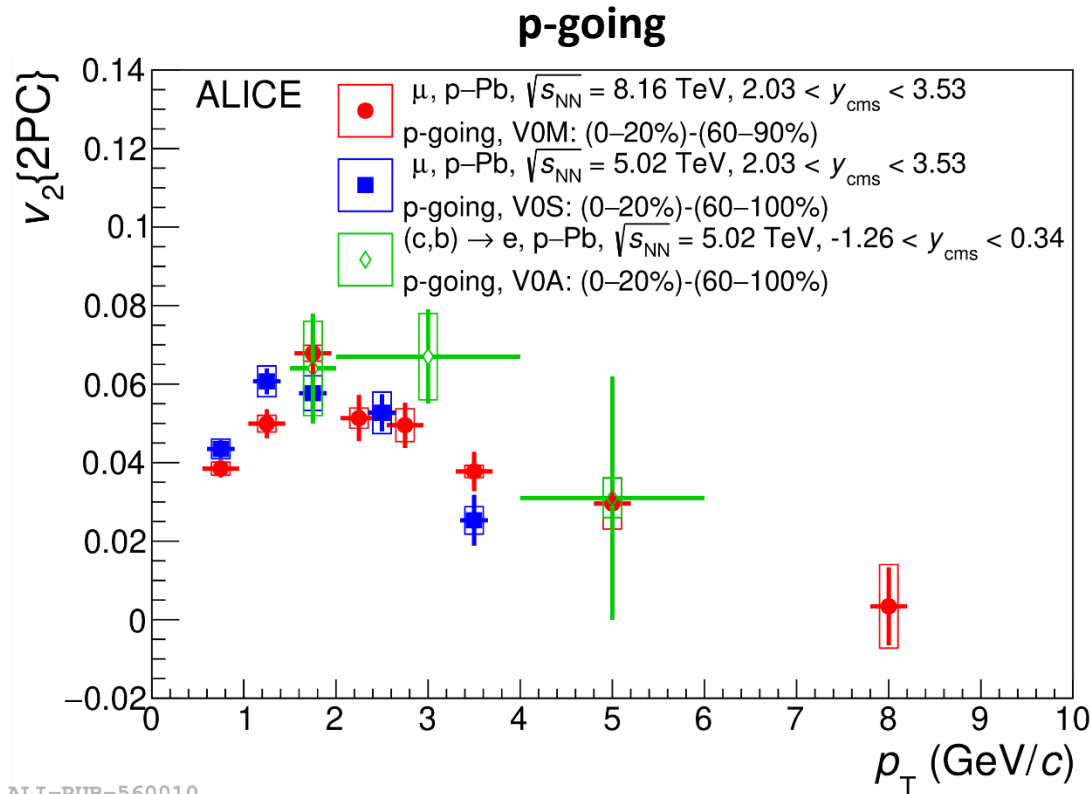
**HF- $\mu$  dominate at  $p_T > 2$  GeV/c**



**PLB 846 (2023) 137782**

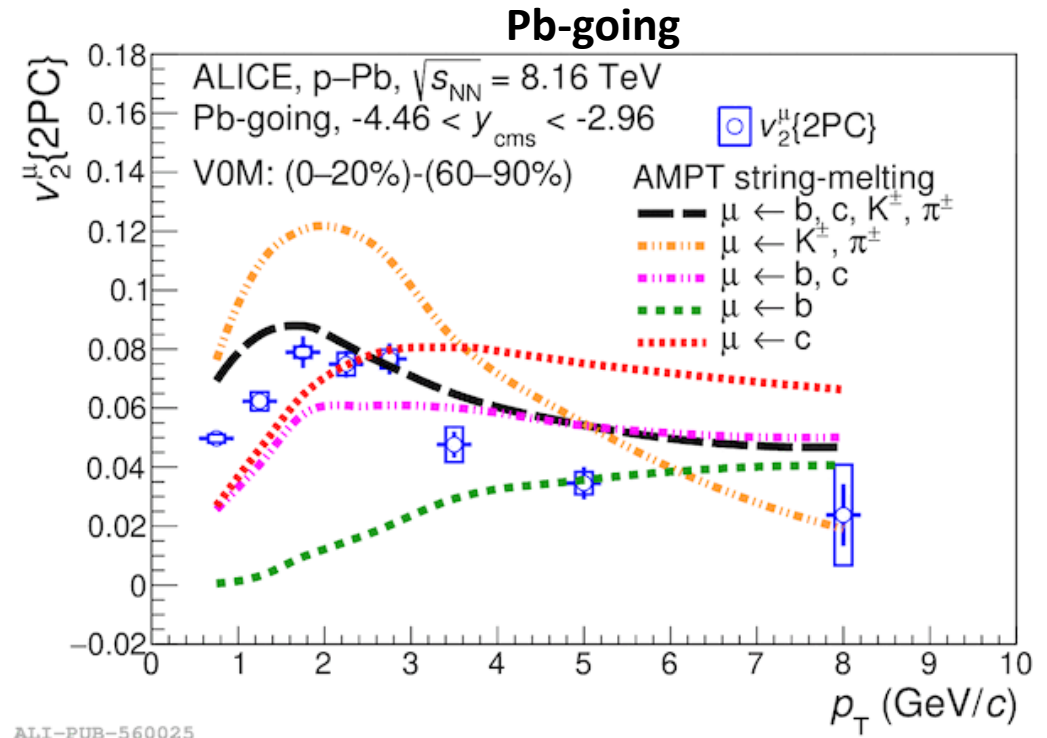
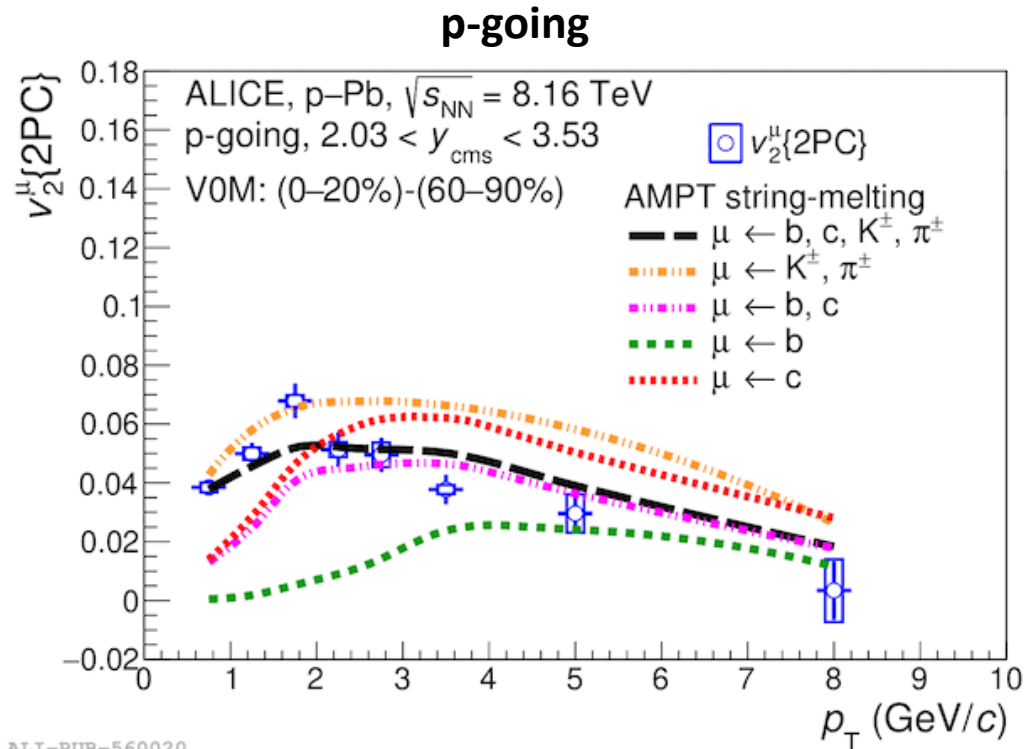
- ❖ Two different techniques: **forward-central two-particle correlations** and **two-particle cumulants**
- ❖ **Positive muon  $v_2$**  measured for the first time over a **wide  $p_T$  interval** at both forward and backward rapidities with a **significance of  $4.7\sigma$  and  $7.6\sigma$**  for  $2 < p_T < 6$  GeV/c, respectively
- ❖ Consistent  $v_2$  values with two-particle correlations and two-particle cumulants
- ❖ Hint for a **smaller  $v_2$  at forward** than backward rapidity: consequence of **decorrelation effects of flow vectors in different rapidity regions**

# Comparison with other published measurement



- ❖ **No significant energy dependence** on the muon  $v_2$  values obtained at  $\sqrt{s_{NN}} = 5.02$  TeV and 8.16 TeV
- ❖ Results in **good agreement** within uncertainties at midrapidity with the **heavy-flavour hadron decay electron  $v_2$**  at  $\sqrt{s_{NN}} = 5.02$  TeV

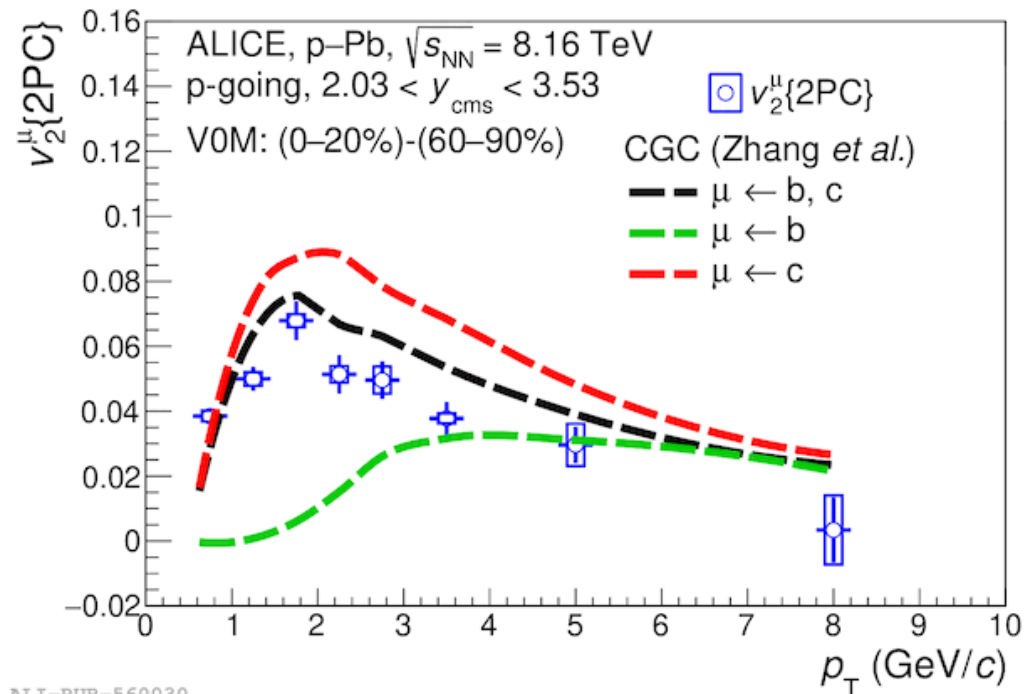
# Comparisons with AMPT calculations



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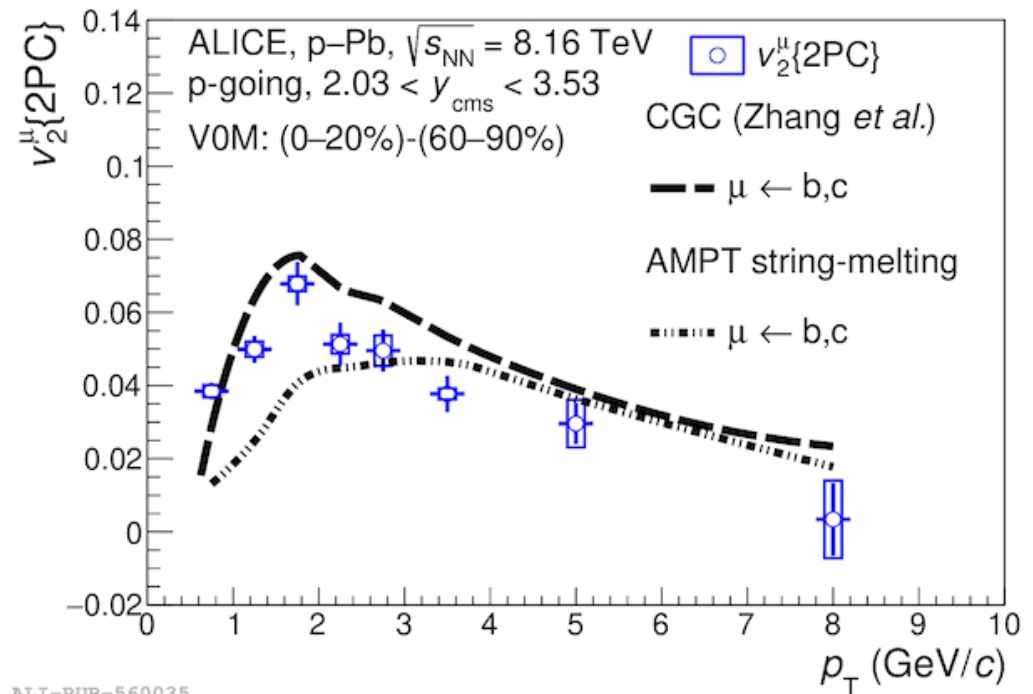
- ❖ AMPT calculations [Z. W. Lin, PRC 72 (2005) 064901]: non-equilibrium dynamics, **microscopic evolution** of parton interactions
  - Larger  $v_2$  for muons from light-flavor hadron decays than for heavy-flavor hadron decay muons
  - Larger  $v_2$  at backward rapidities: rapidity-dependent flow-vector fluctuations
- ❖ AMPT predictions in **fair agreement** with the measured  $v_2$ , although the model slightly overestimates the data at backward rapidities
  - Suggests that the  $v_2$  could be due to the **anisotropic parton escape mechanism**

# Comparisons with AMPT and CGC calculations



ALI-PUB-560030

**HF- $\mu$  dominate at  $p_T = 2$  GeV/c**



ALI-PUB-560035

**PLB 846 (2023) 137782**

❖ CGC calculations [C. Zhang *et al.*, PRL 122 (2019) 172303]: anisotropy generated from **parton interactions in the early stage of collisions**

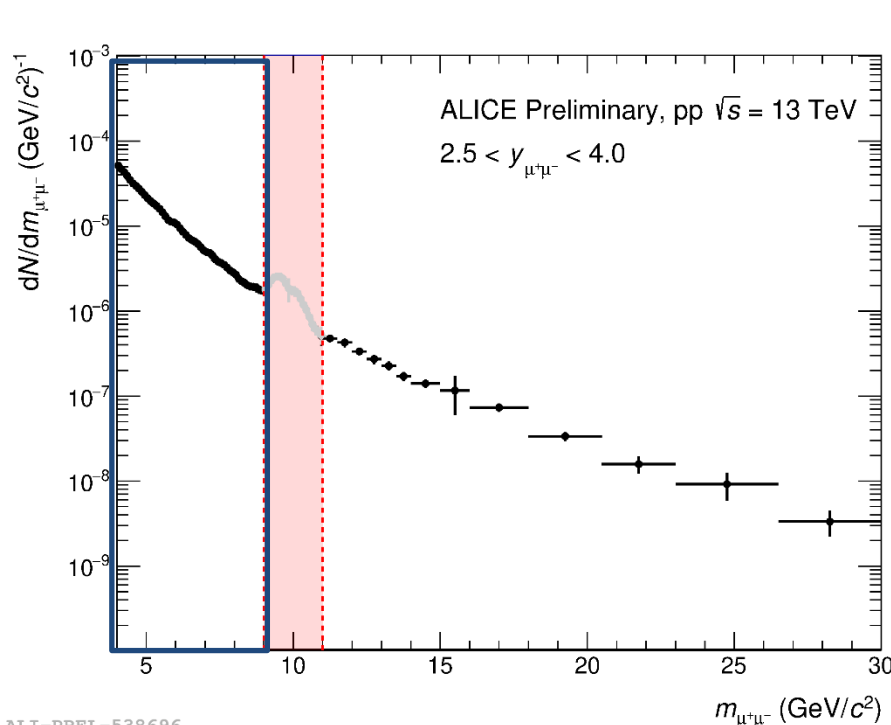
- Qualitative agreement with the measured  $v_2$
- Larger  $v_2$  for heavy-flavour hadron decay muons at low  $p_T$  compared to AMPT
- Compatible  $v_2$  at high  $p_T$  with AMPT: heavy-flavour dominate
- Possible **contributions from initial-state effects** not fully excluded

# Heavy-flavour production in pp collisions at forward $y$ via $\mu^+\mu^-$

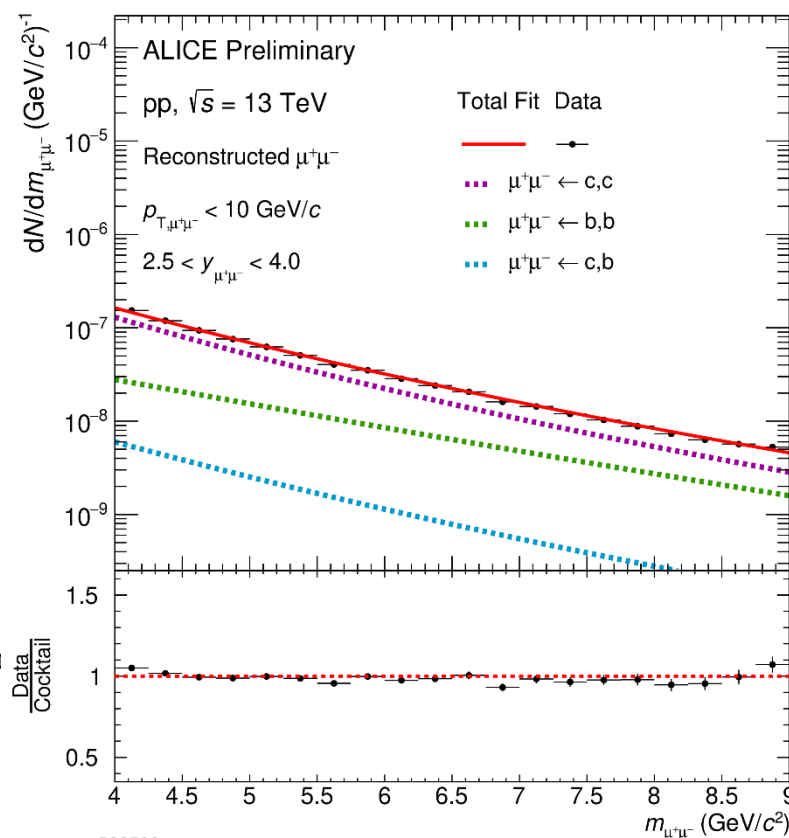


- ❖ Charm and beauty production measured at  $2.5 < y < 4.0$  in pp collisions at  $\sqrt{s} = 13$  TeV, exploiting the **dimuon high-mass continuum dominated by the semimuonic decays of heavy-flavour hadrons**
- ❖ Simultaneous fit to the mass and  $p_T$  distributions with a combined template of the main sources in the continuum
  - ❖ Templates extracted from the heavy-flavour enriched PYTHIA 8 simulations

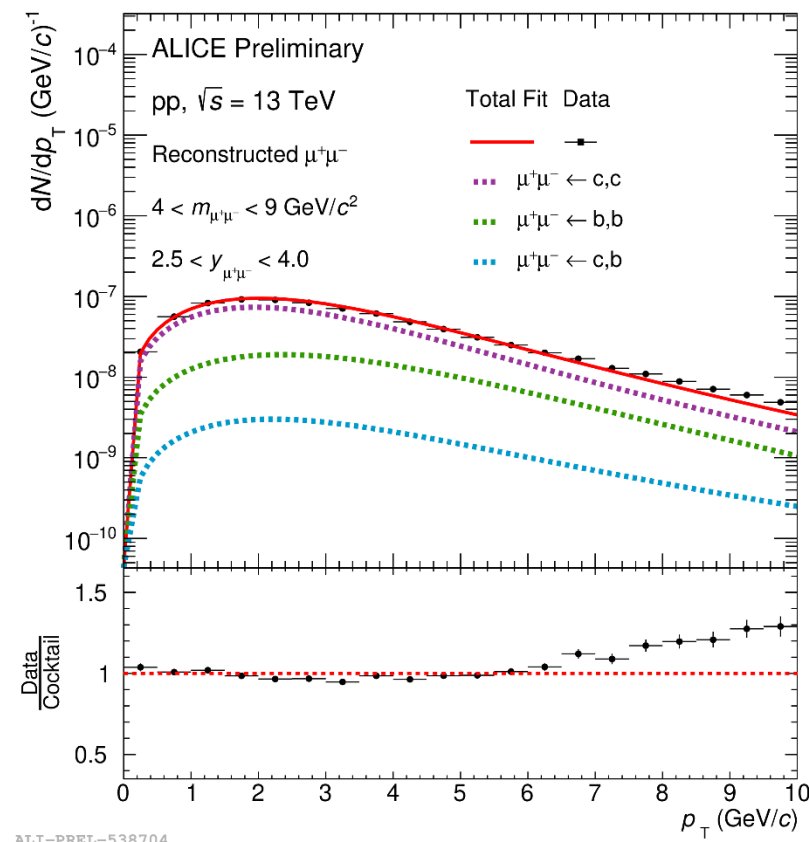
PYTHIA 8: Comput. Phys. Commun. 191 (2015) 159



ALI-PREL-538696



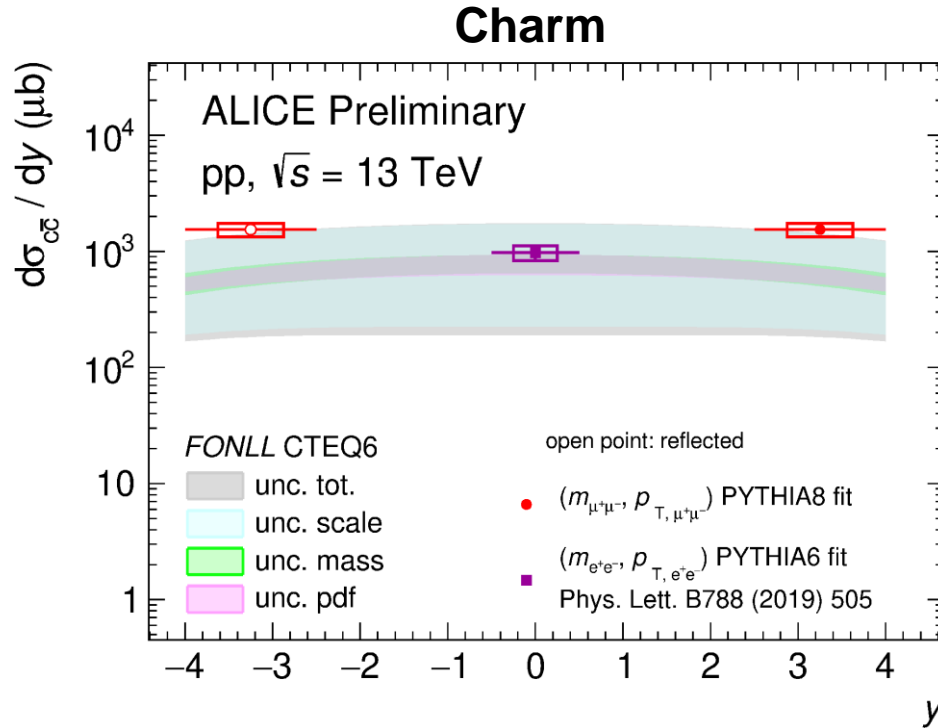
ALI-PREL-538700



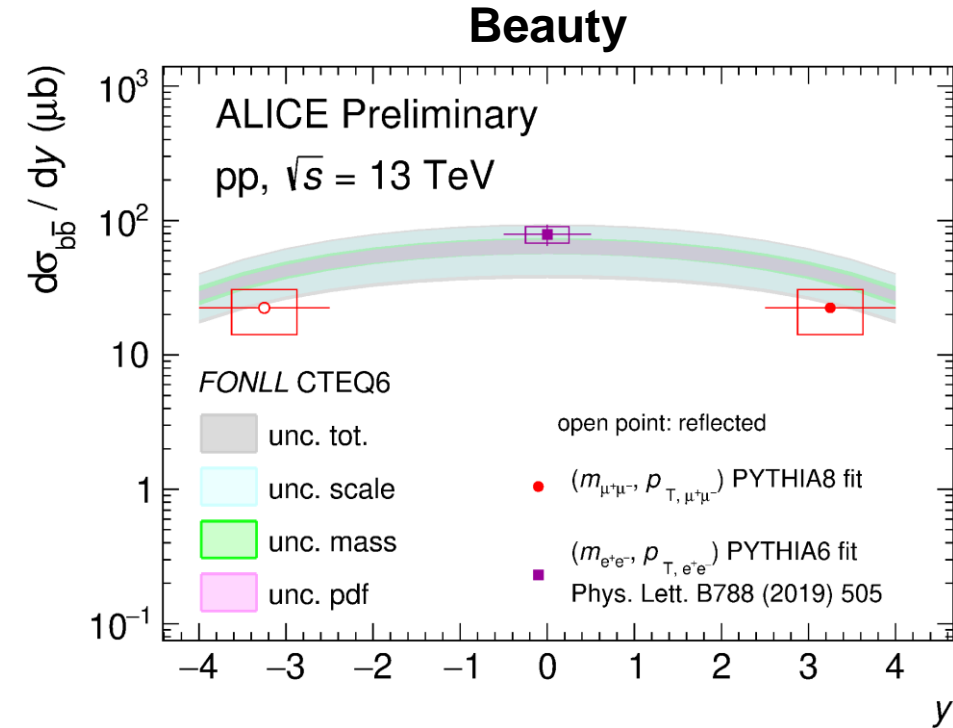
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# Charm and beauty production cross sections at forward $y$ in pp collisions at $\sqrt{s} = 13$ TeV



ALI-PREL-538716

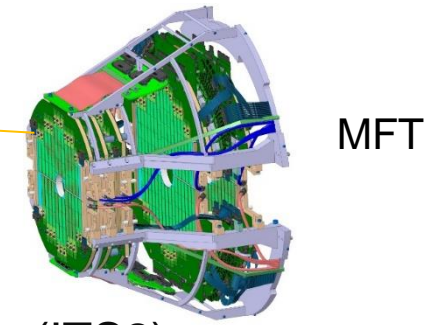
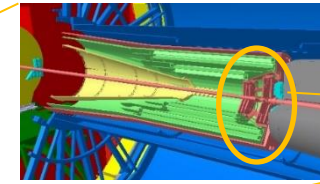
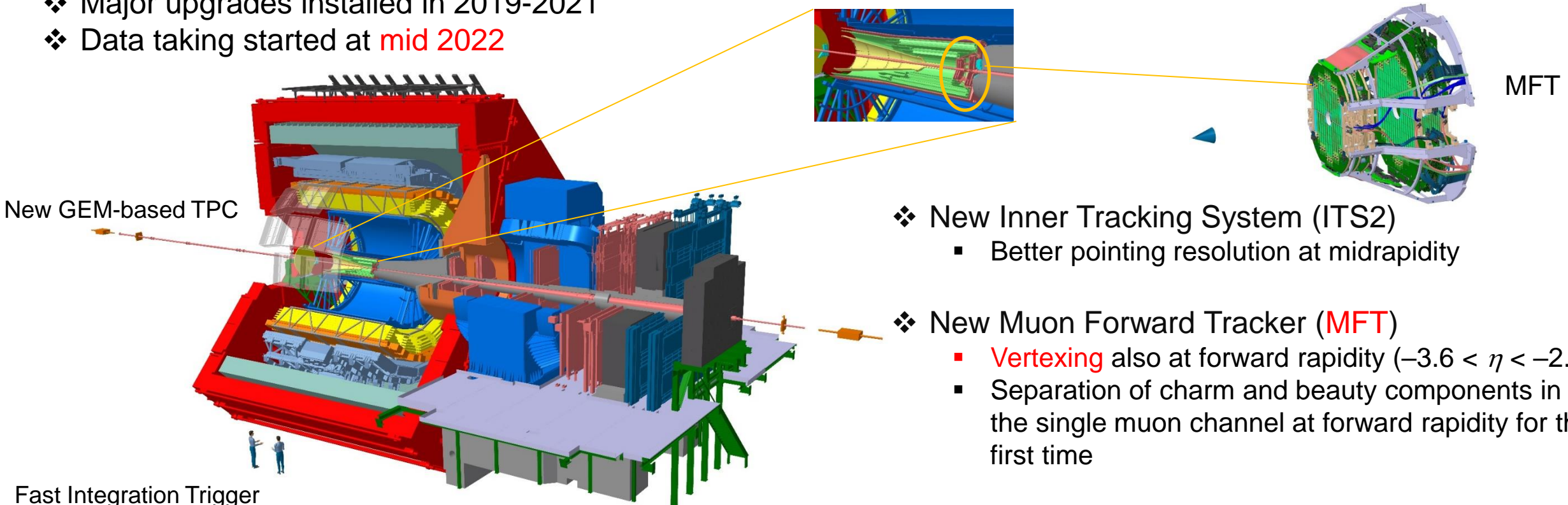


ALI-PREL-538708

- ❖ Charm and beauty production cross sections measured separately at forward rapidity via the dimuon continuum
- ❖ Results in agreement with FONLL predictions within uncertainties, although they lie at the upper and lower limit of the calculations for charm and beauty production cross section, respectively
- ❖ Complement the previously published results at midrapidity in the dielectron channel

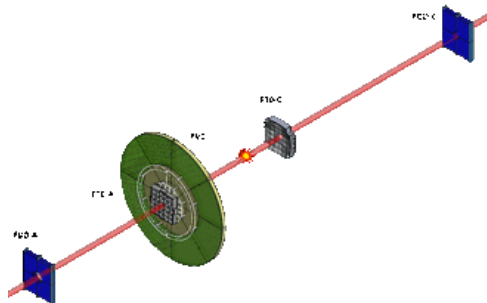
# The ALICE detector: Run 3 setup

- ❖ Major upgrades installed in 2019-2021
- ❖ Data taking started at **mid 2022**

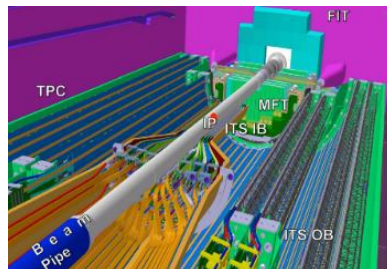


- ❖ New Inner Tracking System (ITS2)
  - Better pointing resolution at midrapidity
- ❖ New Muon Forward Tracker (**MFT**)
  - **Vertexing** also at forward rapidity ( $-3.6 < \eta < -2.5$ )
  - Separation of charm and beauty components in the single muon channel at forward rapidity for the first time
- ❖ New readout system of most subdetectors
  - Continuous readout and increased luminosity
  - Multidifferential measurements of various muon sources down to lower  $p_T$  and up to higher  $p_T$
- ❖ Extend **heavy flavour measurements** towards the **forward rapidity region** allowing for **c/b separation**

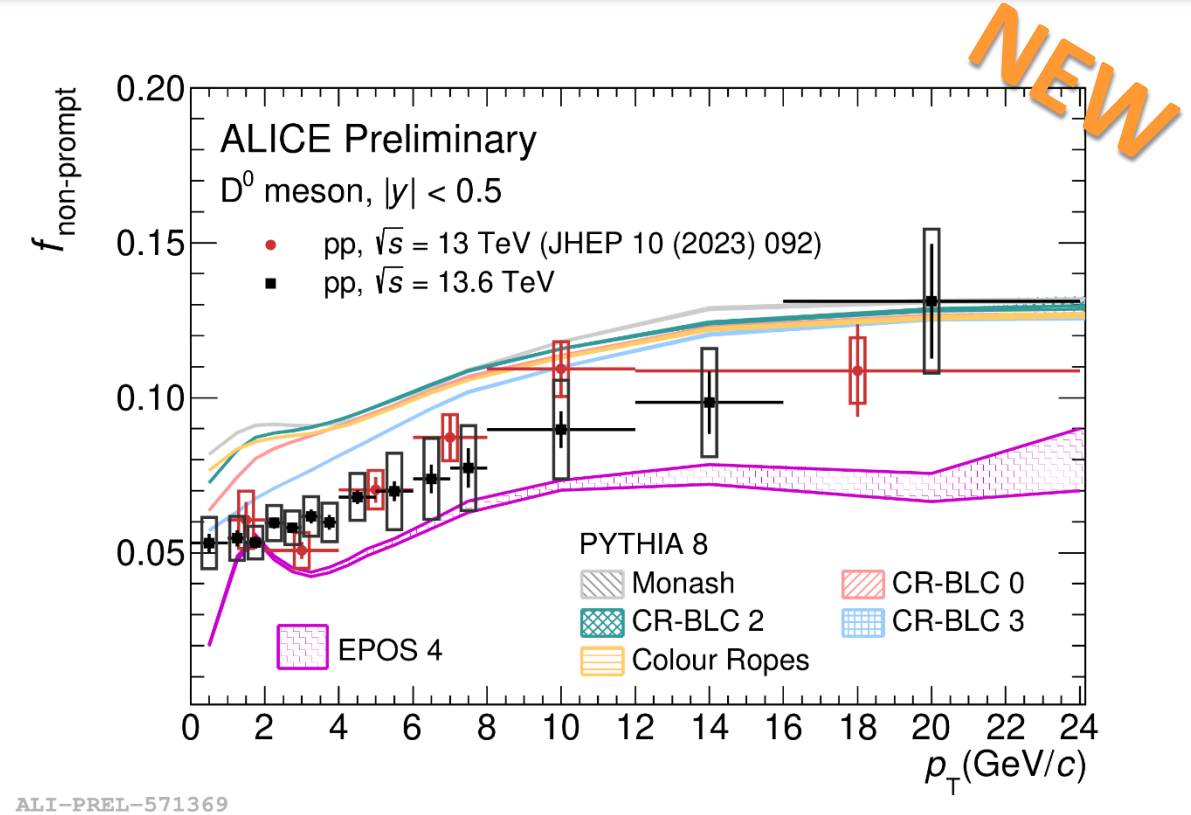
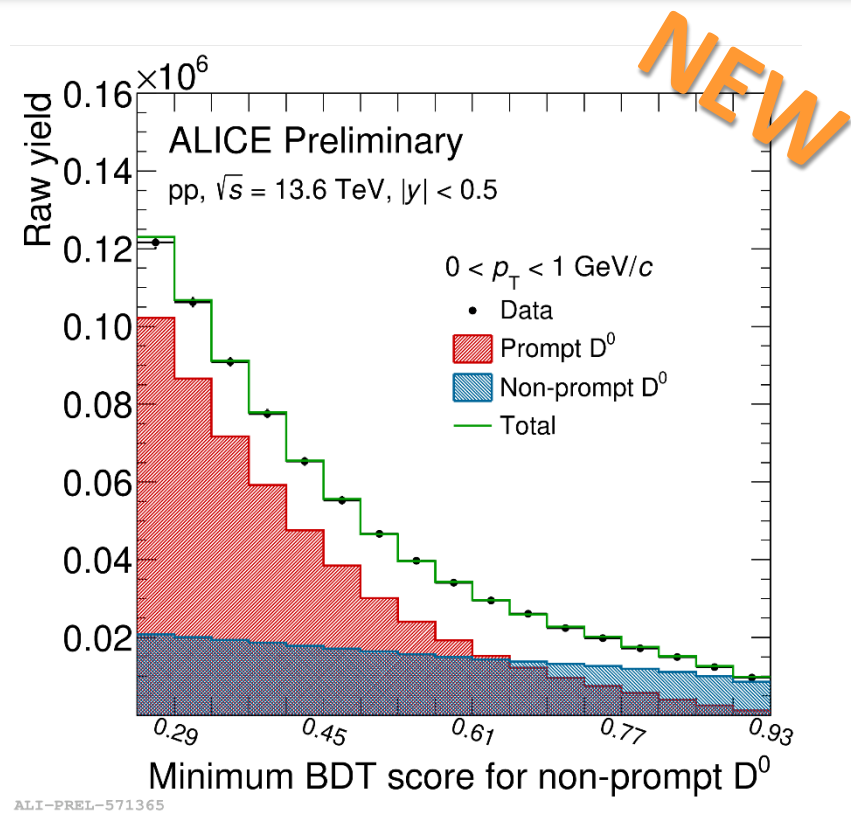
Fast Integration Trigger



New beam pipe

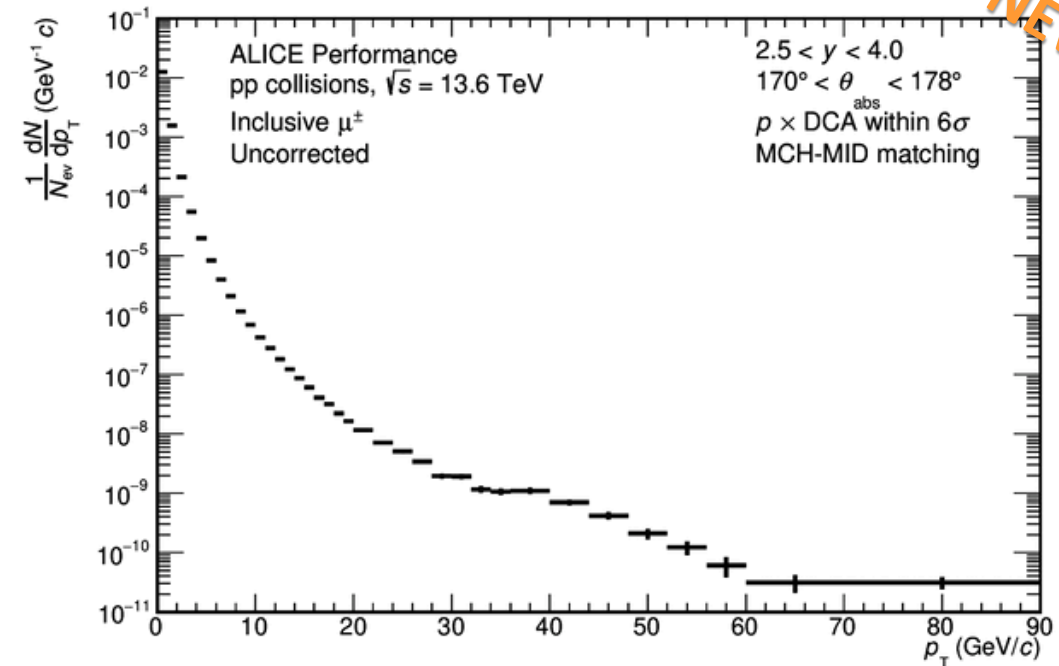
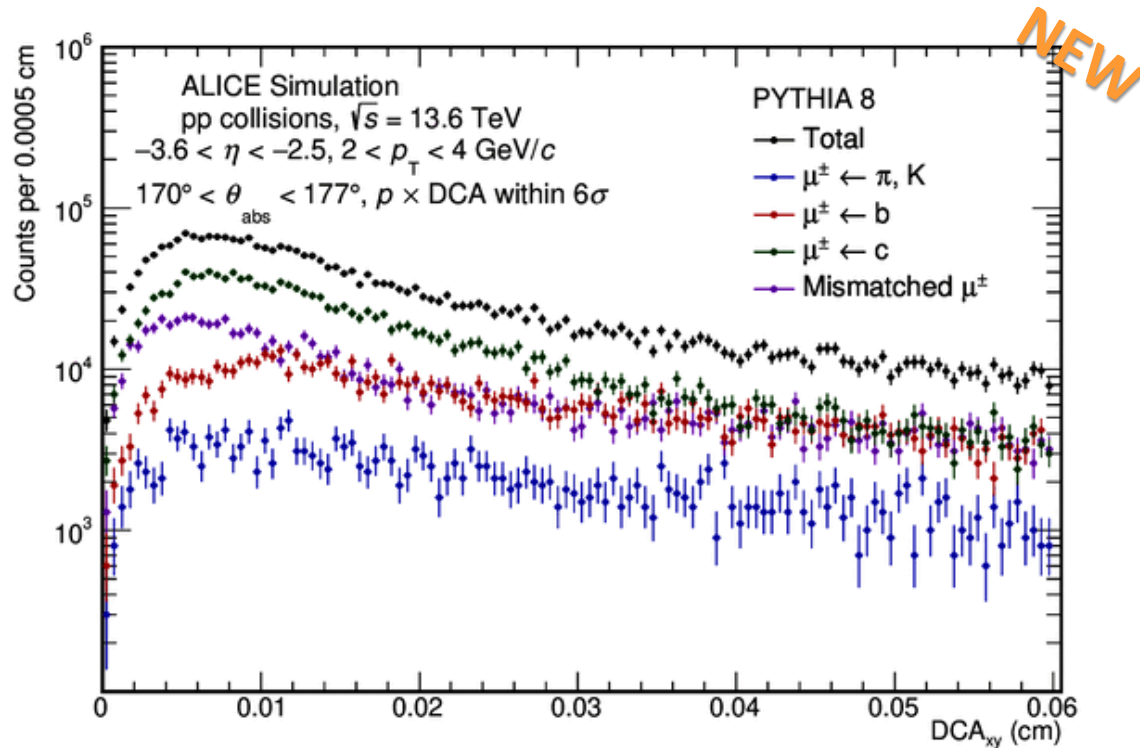


# Midrapidity: non-prompt $D^0$ -meson fraction in pp collisions at $\sqrt{s} = 13.6$ TeV



- ❖ First measurement of the non-prompt  $D^0$ -meson fraction down to  $p_T = 0$  in Run 3
  - In agreement with the measurements in pp collisions  $\sqrt{s} = 13$  TeV
  - Increased granularity and extended  $p_T$  reach
  - Strong constraints on model calculations

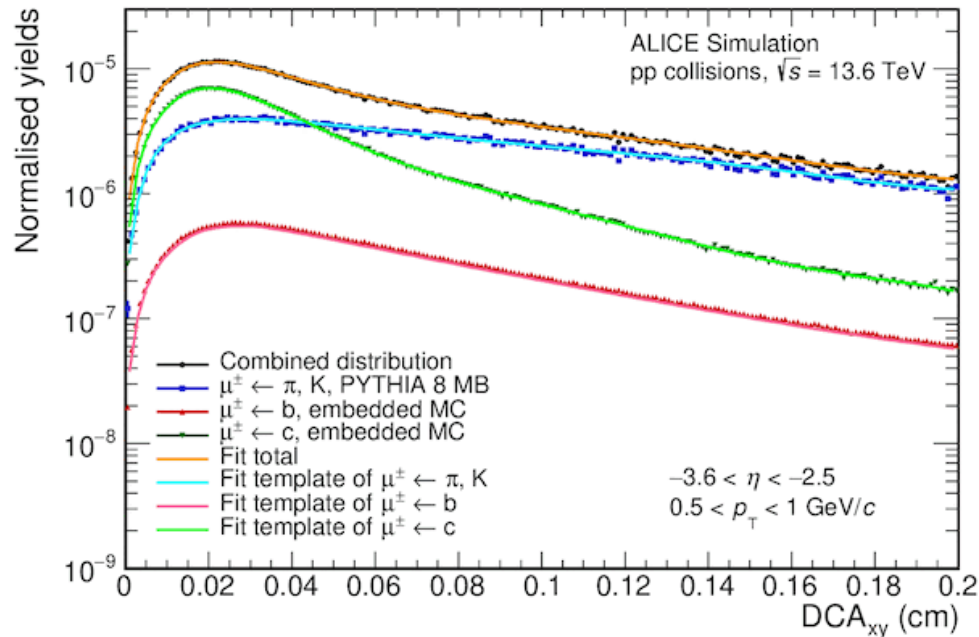
- ❖ High  $p_T$  region measurements feasible even with a fraction of the 2022-2023 pp sample at  $\sqrt{s} = 13.6$  TeV
- ❖ Muons from W-boson decays clearly observed at  $p_T \sim 40$  GeV/c
- ❖ Very promising to perform **multidifferential measurements of open heavy flavours in the semimuonic channel** with high precision, analysing the full pp sample



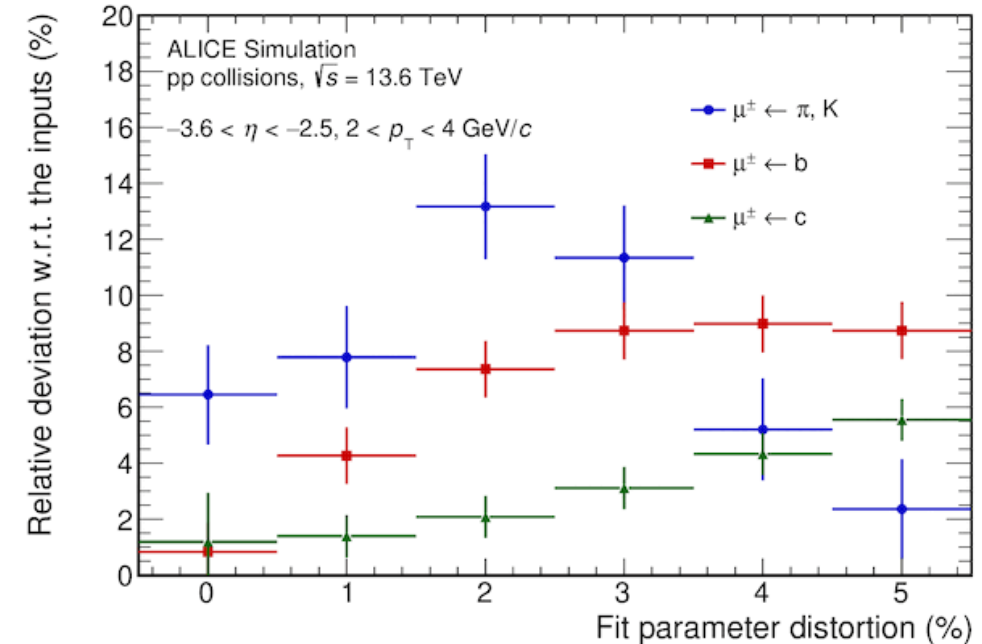
- ❖ Different **decay length** of charm and beauty hadrons
  - Observable: **DCA<sub>xy</sub>** (Distance of Closest Approach to the primary vertex in the transverse plane) of heavy-flavour decay muons
- ❖ Distinct features of various sources: **charm and beauty separation** at forward rapidity **can be achieved with the MFT**

# Strategy to separate charm- and beauty-decay muons

- ❖ Procedure based on Monte Carlo templates of  $DCA_{xy}$  for each source and in each  $p_T$  interval
  - Parametrized with a **variable-width Gaussian function**
  - Total  $DCA_{xy}$  distribution for each  $p_T$  interval fitted as a combination of the various templates
  - Closure test performed varying fit parameters to mimic detector effects
  
- ❖ Procedure validated with full realistic simulations



ALI-SIMUL-547324



ALI-SIMUL-547377

- ❖  $\mu \leftarrow c$  and  $\mu \leftarrow b$  can be measured separately in the semimuonic channel at forward rapidity in pp collisions at  $\sqrt{s} = 13.6$  TeV with the MFT coupled with the muon spectrometer, **down to the low  $p_T$  region**

- ❖ **Azimuthal anisotropies in small collision systems**
  - Collective-like behaviour of heavy quarks in high-multiplicity p–Pb collisions at both forward and backward rapidities
  - New constraints in the interpretation of the collective-like behaviour in small collision systems and to the model calculations
  
- ❖  $p_T$ -integrated production cross section of **charm and beauty measured via dimuons** in pp collisions at  $\sqrt{s} = 13$  TeV
  
- ❖ The Muon Forward Tracker allows us to inquire further physics channels
  - Multidifferential measurements of production of **charm- and beauty-decay muons separately** in the semimuonic channel down to low  $p_T$
  - Template fit method tested and validated with full Monte Carlo simulations

**Stay tuned: more to come soon in both pp and Pb–Pb collisions for the measurements in the semimuonic channel**



Thank you for your attention

Merci pour votre attention

感谢聆听



ALICE

Run 3 Pb-Pb

$\sqrt{s_{NN}} = 5.36 \text{ TeV}$

27 September 2023, 04:50

**BACKUP SLIDES**



# Backup slides: AMPT(A multiphase transport) model

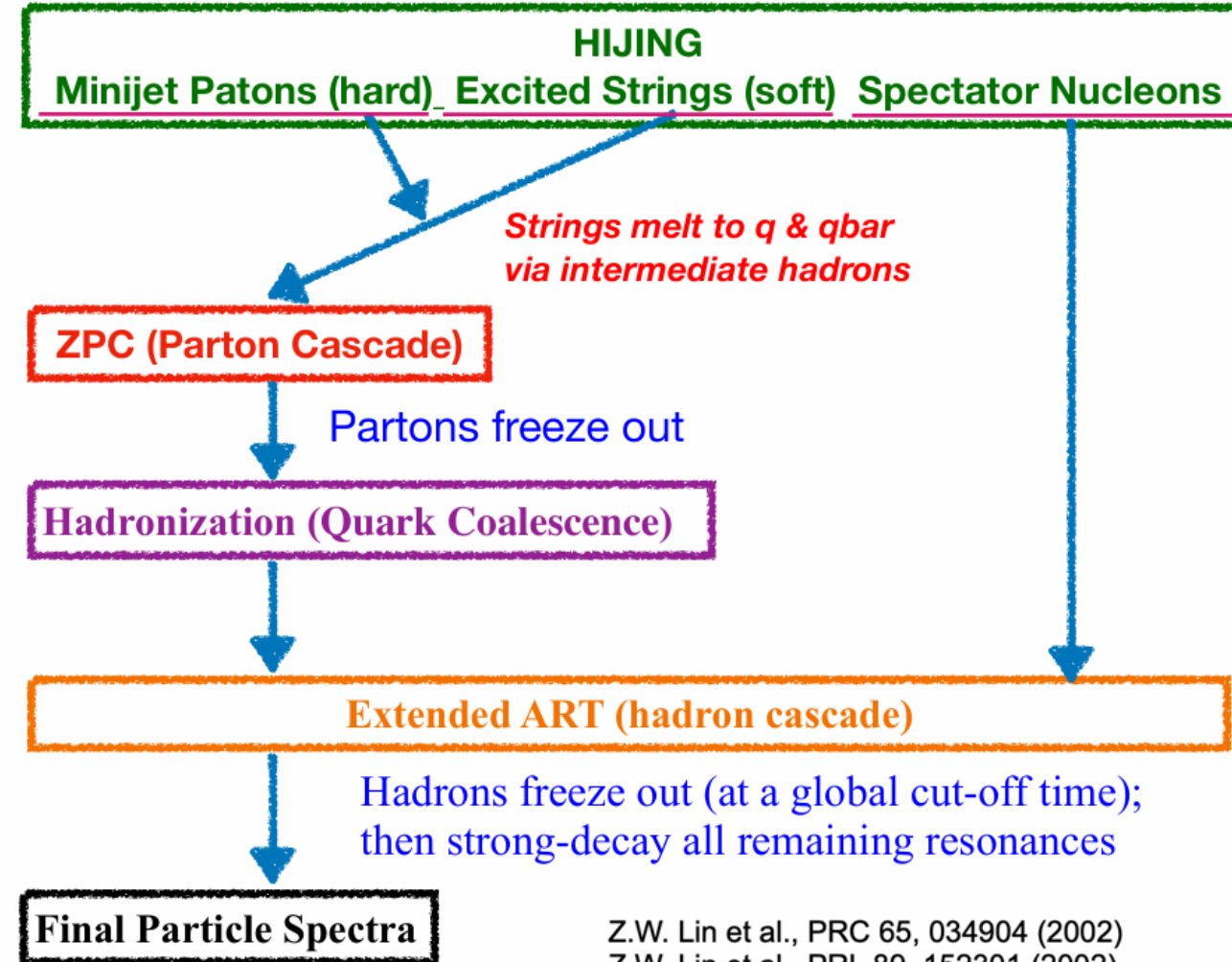
- **Initial conditions:** HIJING two-component model
- **String melting:** hadrons from string fragmentation are melted into primordial quarks and antiquarks
- Quark formation time:  $t_f = E_H/m_{T,H}^2$
- **Parton cascade:** two-body elastic scattering described by ZPC model

Debye screening mass

$$\frac{d\sigma}{dt} = \frac{9\pi\alpha_s^2}{2} \left(1 + \frac{\mu^2}{s}\right) \frac{1}{(t - \mu^2)^2}$$

- **Coalescence:** combine nearest quarks to meson/baryon
- **A Relativistic Transport (ART)** to describe hadron scatterings

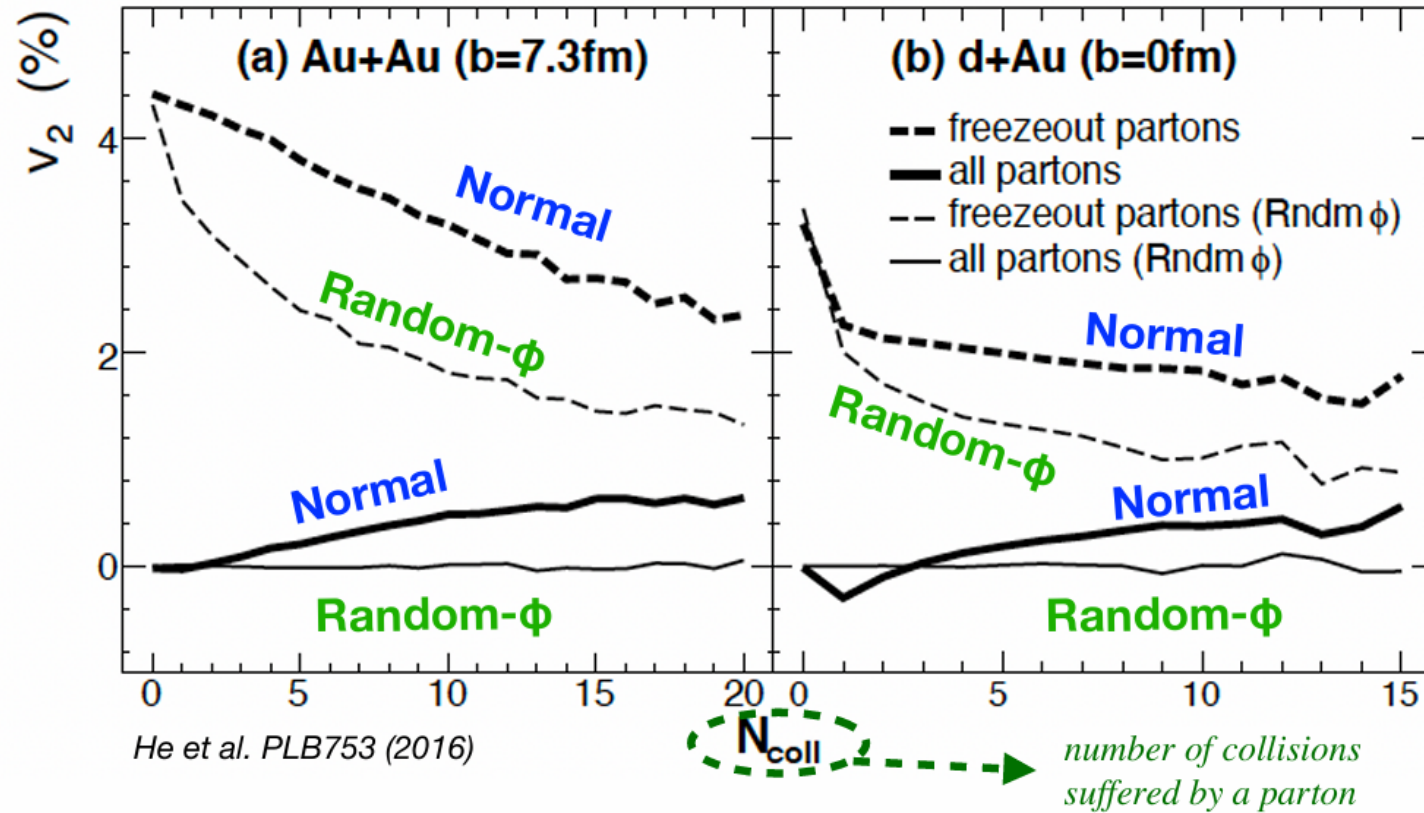
## Structure of AMPT v2.xx (String Melting Version)



Z.W. Lin et al., PRC 65, 034904 (2002)  
Z.W. Lin et al., PRL 89, 152301 (2002)

# Backup slides: escape mechanisms

Zi-Wei Lin, IHEP2016



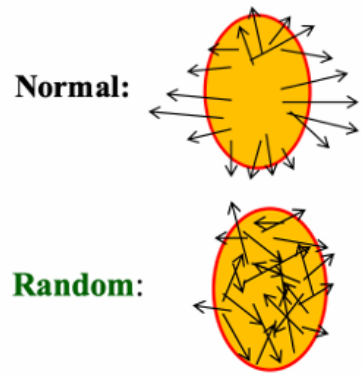
**Freeze out partons:** freeze out after exactly  $N_{\text{coll}}$  collisions;

**Active partons:** will collide further, freeze out after  $> N_{\text{coll}}$ ;

**All partons:** sum of the above two partons

$v_2$	Normal $v_2$	Random- $\phi$ $v_2$
Au+Au	3.9%	2.7%
d+Au	2.7%	2.5%

	Fraction from pure escape	$\langle N_{\text{coll}} \rangle$ all partons
Au+Au	69%	4.6(modest)
d+Au	93%	1.2(low)



- Few interactions are sufficient for anisotropy in AMPT
- Partons more likely to escape in the short direction
- Particle escape is the dominant mechanism for  $v_2$  in small collision systems
- The  $v_2$  from pure escape decrease with increasing  $N_{\text{coll}}$  → transport approach hydrodynamics

- The nonflow contribution is estimated in low-multiplicity (60–90%) events, and subtracted for both differential and reference cumulants

$$v_2^\mu\{2\}(p_T) = \frac{[d_2^\mu\{2\}(p_T)]_{(0-20\%)}}{\sqrt{[c_2\{2\}]_{(0-20\%)}}}$$



subtraction of long-range jet correlation

$$v_2^\mu\{2\}(p_T) = \frac{[d_2^\mu\{2\}(p_T)]_{(0-20\%)} - f \cdot [d_2^\mu\{2\}(p_T)]_{(60-90\%)}}{\sqrt{[c_2\{2\}]_{(0-20\%)} - f \cdot [c_2\{2\}]_{(60-90\%)}}}$$



subtraction of remains long-range and short-range jet correlation

$$v_2^\mu\{2\}(p_T) = \frac{[d_2^\mu\{2\}(p_T)]_{(0-20\%)} - f \cdot f' \cdot [d_2^\mu\{2\}(p_T)]_{(60-90\%)}}{f_{\text{RP}} \cdot \sqrt{[c_2\{2\}]_{(0-20\%)} - f \cdot [c_2\{2\}]_{(60-90\%)}}} \cdot f_{\Delta\eta}$$

$f$ : the ratio of mean SPD tracklet in low-multiplicity collisions to that in high-multiplicity collisions

$f_{\text{RP}}, f'$ : the factors to remove the **remaining long-range jet correlations**, obtained from two-particle correlations

$f_{\Delta\eta}$ : the factor to suppress the **short-range jet correlations**

Define  $n$  sets of selections with different prompt and non-prompt  $D^0$  contributions

For each selection set, the raw yield and the efficiencies are related to the corrected yields of prompt  $N_{prompt}$  and non-prompt  $N_{non-prompt}$   $D^0$

$$\begin{cases} (\text{Acc} \times \epsilon)_1^{prompt} \cdot N_{prompt} + (\text{Acc} \times \epsilon)_1^{non-prompt} \cdot N_{non-prompt} = Y_1 \\ \dots \\ (\text{Acc} \times \epsilon)_n^{prompt} \cdot N_{prompt} + (\text{Acc} \times \epsilon)_n^{non-prompt} \cdot N_{non-prompt} = Y_n \end{cases}$$

The algebraic equations can be represented by:

$$\begin{pmatrix} \epsilon_1^p & \epsilon_1^{np} \\ \vdots & \vdots \\ \epsilon_n^p & \epsilon_n^{np} \end{pmatrix} \times \begin{pmatrix} N_p \\ N_{np} \end{pmatrix} - \begin{pmatrix} Y_1 \\ \vdots \\ Y_n \end{pmatrix} = \begin{pmatrix} \delta_1 \\ \vdots \\ \delta_n \end{pmatrix}$$

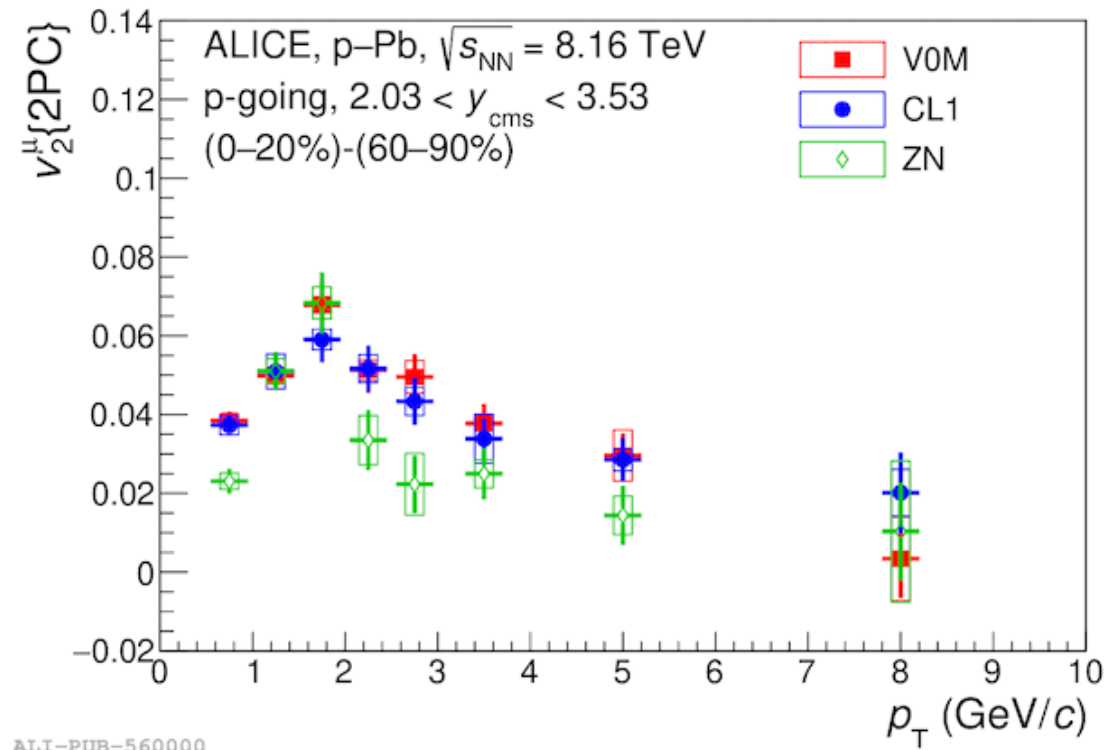
The  $\chi^2$  of the system is defined as:

$$\chi^2 = \delta^T C^{-1} \delta, \quad \text{where } C \text{ is the covariance matrix from the uncertainties}$$

Corrected yields of prompt and non-prompt  $D^0$  obtained from  $\chi^2$  minimization of the system

Non-prompt fraction  $f_{non-prompt}$  evaluated for a given set of selections as

$$f_{non-prompt}^i = \frac{(\text{Acc} \times \epsilon)_i^{non-prompt} \cdot N_{non-prompt}}{(\text{Acc} \times \epsilon)_i^{non-prompt} \cdot N_{non-prompt} + (\text{Acc} \times \epsilon)_i^{prompt} \cdot N_{prompt}}$$



- ❖ Compatible  $v_2$  values with different multiplicity estimators within uncertainties with a hint for a smaller  $v_2$  using the energy deposited in the neutron ZDC

- ❖ Procedure based on Monte Carlo templates of  $DCA_{xy}$  for each source and in each  $p_T$  interval
  - Parametrized with a **variable-width Gaussian function**, the width being a polynomial function of the  $DCA_{xy}$

$$f(x) = A e^{-(x-\mu)^2/2\sigma(x)^2}$$
$$\sigma(x) = \sigma_0^L + \sigma_1^L (\mu - x) + \dots + \sigma_3^L (\mu - x)^3 \text{ for } x \leq \mu$$
$$\sigma(x) = \sigma_0^R + \sigma_1^R (x - \mu) + \dots + \sigma_6^R (x - \mu)^6 \text{ for } x > \mu$$

- Total  $DCA_{xy}$  distribution for each  $p_T$  interval fitted as a combination of the various templates

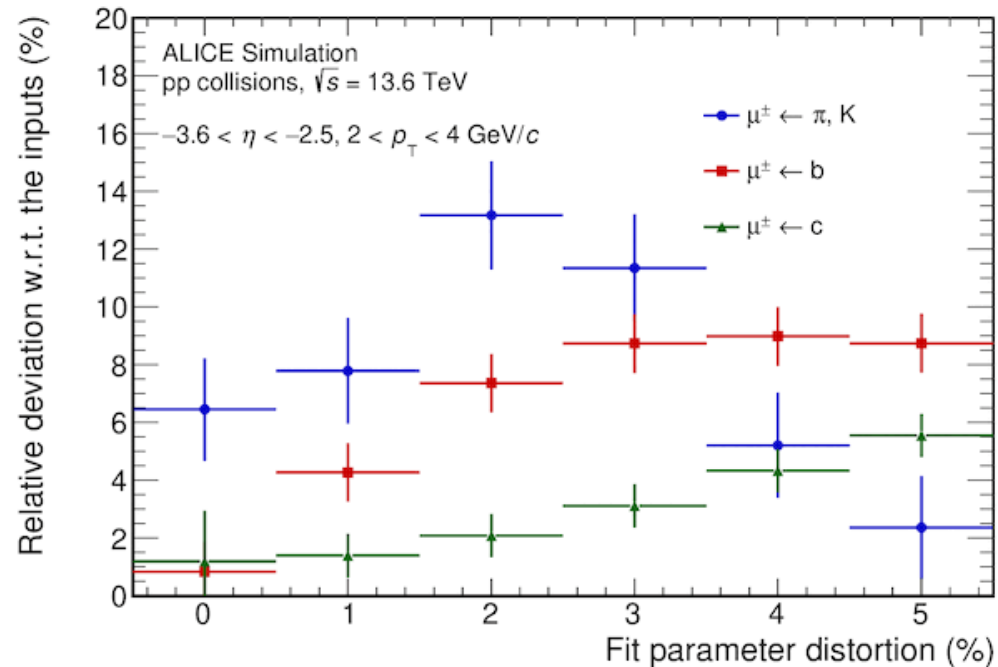
$$f(x) = B \cdot f_b(x) + C \cdot f_c(x) + D \cdot f_{bkg}(x)$$

- Extraction of yields of  $\mu \leftarrow c$  and  $\mu \leftarrow b$

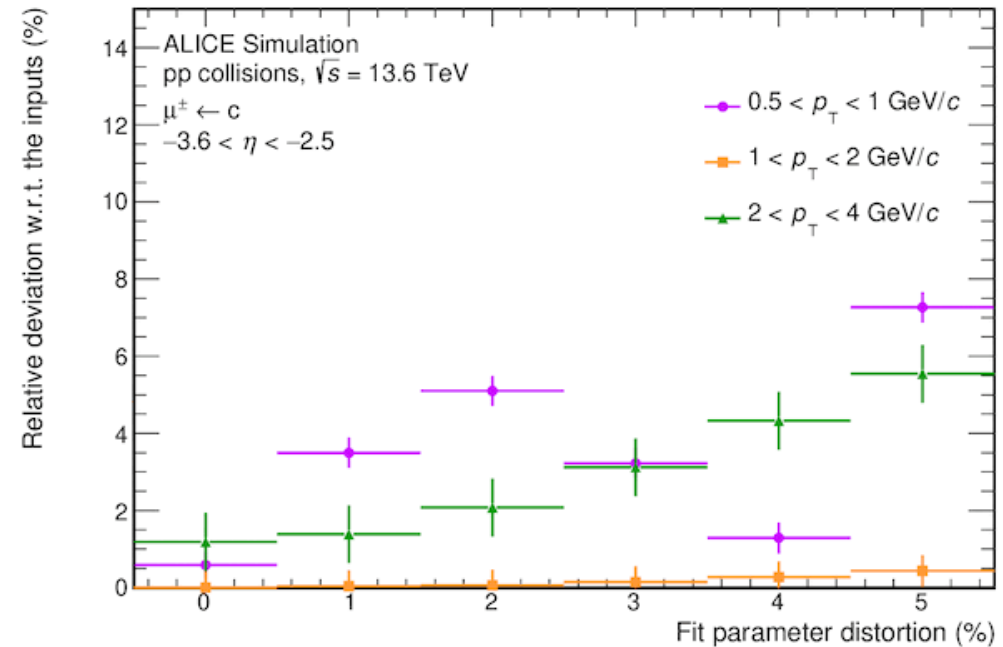
# Strategy to separate charm- and beauty-decay muons

## ❖ Closure test: varying fit parameters to mimic detector effects

- Randomly vary width parameters by 1% to 5% to generate distorted templates
- Refit  $DCA_{xy}$  using these distorted templates and examine the sensitivity through their relative deviations



ALI-SIMUL-547377



ALI-SIMUL-547380

- ❖  $\mu \leftarrow c$  and  $\mu \leftarrow b$  can be measured separately in the semimuonic channel at forward rapidity in pp collisions at  $\sqrt{s} = 13.6$  TeV with the MFT coupled with the muon spectrometer, **down to about 0.5 GeV/c and 2 GeV/c for charm and beauty, respectively**