

Quarkonium production in pp and heavy-ion collisions

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outline

- Introduction
- Quarkonium production in pp collision
- Quarkonium production in AA collision
- Summary

quarkonium production in p+p collisions



uncertainty principle

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Wigner projection

• Wigner density of S- & P-wave states: $\langle p, r | p_1, r_1; p_2, r_2 \rangle$

$$\Phi_{\rm S}^{W}(\mathbf{r},\mathbf{p}) = 8 \frac{D}{d_1 d_2} \exp\left[-\frac{r^2}{\sigma^2} - \sigma^2 p^2\right],$$

$$\Phi_{\rm P}^{W}(\mathbf{r},\mathbf{p}) = \frac{16}{3} \frac{D}{d_1 d_2} \left(\frac{r^2}{\sigma^2} - \frac{3}{2} + \sigma^2 p^2\right),$$

$$\times \exp\left[-\frac{r^2}{\sigma^2} - \sigma^2 p^2\right],$$

- D, d1, d2: color-spin degeneracy of quarkonium, heavy quark and antiquark
- r, p: relative distance and momentum in center-of-mass frame
- σ : the only parameter ~ quarkonium radius

J/ψ in pp @ 200 GeV

T. Song, J. Aichelin, E. Bratkovskaya, PRC96, 014907

direct: purely initial J/ ψ prompt: including the feeddown from χ_c , ψ' Inclusive: including χ_c , ψ' and B decay













Quarkonium production in heavy-ion collisions

- Different from in p+p collisions,
- Quarkonium cannot be formed above the dissociation temperature
- Quarkonium radius changes with time (temperature)
- Quarkonium dissociation and regeneration take place
- ...

• We use the Remler formalism

E. A. Remler and A. P. Sathe, Ann. Phys. 91, 295 (1975).
E. A. Remler, Ann. Phys. 95, 455 (1975).
E. A. Remler, Ann. Phys. 136, 293 (1981).
D. Y. A. Villar, J. Zhao, J. Aichelin, and P. B. Gossiaux, arXiv:2206.01308 (2022).

T. Song, J. Aichelin, and E. Bratkovskaya, Phys. Rev. C ${\bf 96},\,014907$ (2017), 1705.00046.

Remler's formalism



1. This is carried out for all $Q\bar{Q}$ pairs for all physical states: J/ ψ , χ_c , ψ' for $c\bar{c}$ and Υ (1S), Υ (2S), Υ (3S), χ_b (1P), χ_b (2P) for $b\bar{b}$ 2. Heavy quark potential and scattering are closely related to quarkonium production/dissociation 2-dimensional table is prepared for all possible combinations of heavy quark and heavy antiquark whether charm or bottom $(Q_i, \overline{Q}_j) = (c, \overline{c}), (c, \overline{b}), (b, \overline{c}), (b, \overline{b})$

	Q_1	Q_2	Q ₃	Q_4	Q_5	•••
\overline{Q}_1	d					
\bar{Q}_2		i		off	diag onal	
\overline{Q}_3			а			
\overline{Q}_4				g		
\overline{Q}_5					0	
•••						nal

Initially all cells are empty because of high temperature. When the local temperature at the middle of (Q_i, \overline{Q}_j) becomes lower than T_{diss} , the cell begins to be filled in \rightarrow initial production

	Q_1	Q_2	Q ₃	Q_4	Q_5	•••
\overline{Q}_1	d					
\bar{Q}_2		i		off	diag onal	
\bar{Q}_3			а			
${ar Q}_4$				g		
\overline{Q}_5					0	
•••						nal

When Q_2 scatters, Q_2 column is updated

	Q_1	Q_2	Q ₃	Q_4	Q_5	•••
$ar{Q}_1$	d					
\bar{Q}_2		i		off	diag onal	
\overline{Q}_3			а			
$ar{Q}_4$				g		
\bar{Q}_5					0	
•••						nal

When \bar{Q}_3 scatters, \bar{Q}_3 row is updated

	Q_1	Q_2	Q_3	Q_4	Q_5	•••
$ar{Q}_1$	d					
\bar{Q}_2		i		off	diag onal	
\bar{Q}_3			а			
$ar{Q}_4$				g		
\overline{Q}_5					0	
•••						nal



Parton-Hadron-String-Dynamics (PHSD)

Parton-Hadron-String Dynamics (PHSD) is a non-equilibrium microscopic transport approach for the description of dynamics of strongly-interacting hadronic and partonic matter produced in heavy-ion collisions

Dynamics: based on the solution of generalized off-shell transport equations derived from Kadanoff-Baym many-body theory (beyond semi-classical BUU)



time

PHSD provides a good description of 'bulk' hadronic and electromagnetic observables from SIS to LHC energies

PHSD: W. Cassing, E. Bratkovskaya, PRC 78 (2008) 034919; NPA831 (2009) 215; P. Moreau et al., PRC100 (2019) 014911



Heavy quark scattering in PHSD



$R_{AA}(y)$ of Y(1S)

T. Song, J. Aichelin, J. Zhao, P. Gossiaux, E. Bratkovskaya, PRC108, 054908



$$R_{AA} = \frac{\frac{dN^{AA}}{dydp_T}}{N_{coll}\left(\frac{dN^{pp}}{dydp_T}\right)}$$

Upper limit: Initial production at T_{diss} (Wigner projection at T_{diss})

Solid line: only 10 % of scatterings update the Wigner projection

Lower limit: final number at T_c (Wigner projection at T_c) Because of the spatial diffusion of heavy (anti)quark, Wigner density normally decreases with time

$R_{AA}(p_T)$ of Y(1S)

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There must be a suppression of heavy (anti)quark scattering, if it is bound



If $Q\bar{Q}$ form a color singlet, the size of $Q\bar{Q}$ pair will be small and the scattering cross section must be reduced due to interference

Y(1S) interaction rate \approx interaction rate of *b* and \overline{b} quarks/10



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 $R_{AA}(y)$ of Y(2S)

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The band is too thin, because $T_{diss} \approx T_c$ for Y(2S)

$R_{AA}(p_T)$ of Y(2S)

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The band is too thin, because $T_{diss} \approx T_c$ for Y(2S)

R_{AA} vs N_{part}

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- Diagonal: Upsilon only from initial pairs
- All: include upsilon from initial other pairs
- The contribution from offdiagonal is little, because b quark pairs are rare even at LHC

Off-diagonal quarkonium

Summary

- Relmer formalism is applied to quarkonium production in pp and AA collisions
- In pp collisions, only initial Wigner projection is carried out
- In AA collisions, the initial Wigner projection starts at T_{diss} and then Wigner density is updated whenever Q or \overline{Q} scatters until T_c
- Since the scattering cross section of color singlet (quarkonium) is smaller than that of separate Q and \overline{Q} , there must be a suppression of scattering
- R_{AA} (y), R_{AA} (p_T), R_{AA} (N_{part}) of Y(1S), Y(2S) in Pb+Pb collisions at LHC are well reproduced by the PHSD using free energy for heavy quark potential and applying that the quarkonium interaction rate \approx 10 % interaction rate of separate Q and \overline{Q}
- Off-diagonal bottomonium is little even at LHC

Thank you for your attention