

SQM 2024, Strasbourg (France)

Heavy flavor production and collectivity in high energy proton-proton collisions

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With: J. Aichelin, K. Werner, J. Zhao (main author)

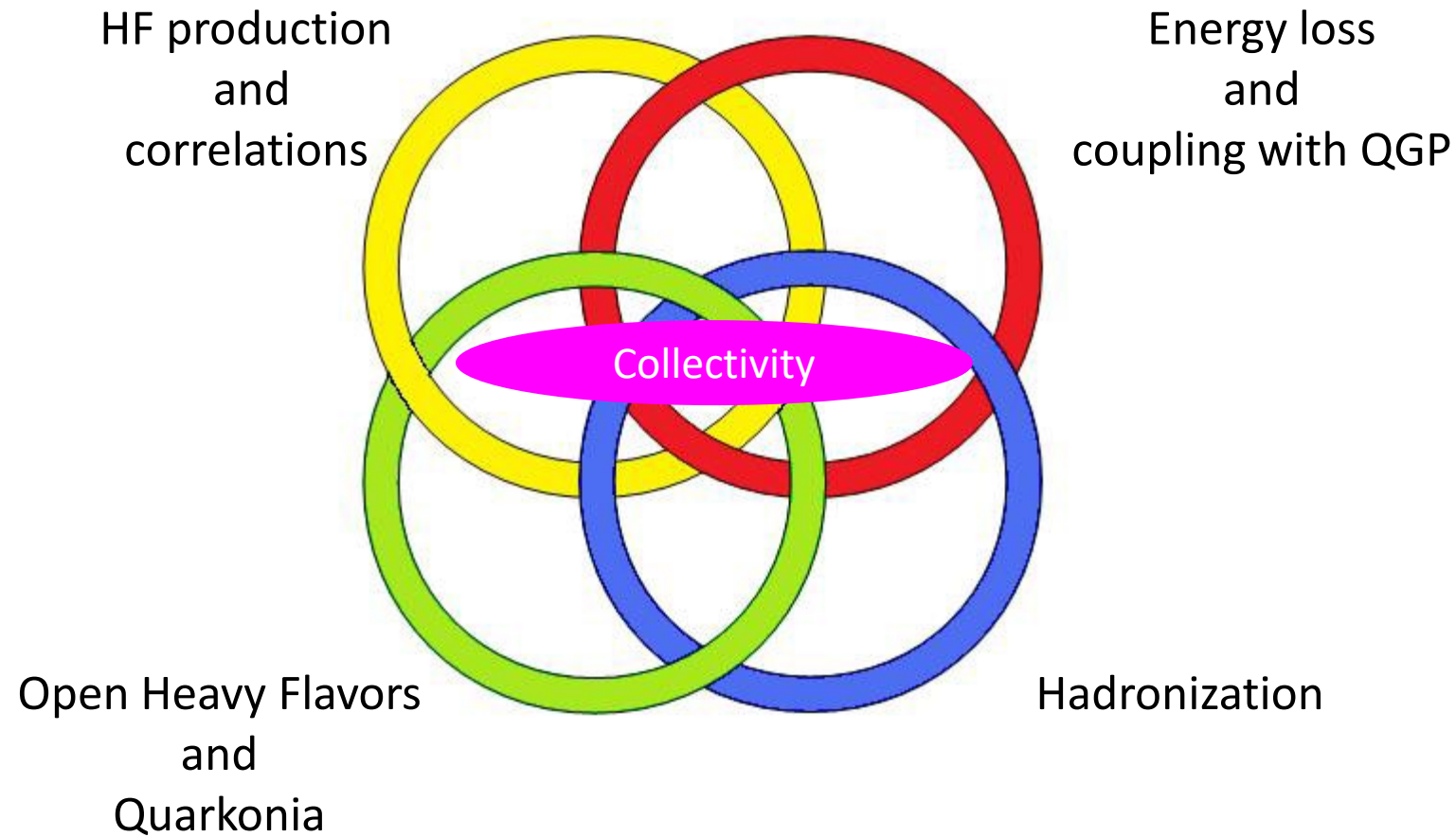


and Pays de la Loire



HF Borromeo's rings of collectivity

Motivation: Studying the collectivity effects in large and small systems as modeled by EPOS4-HQ



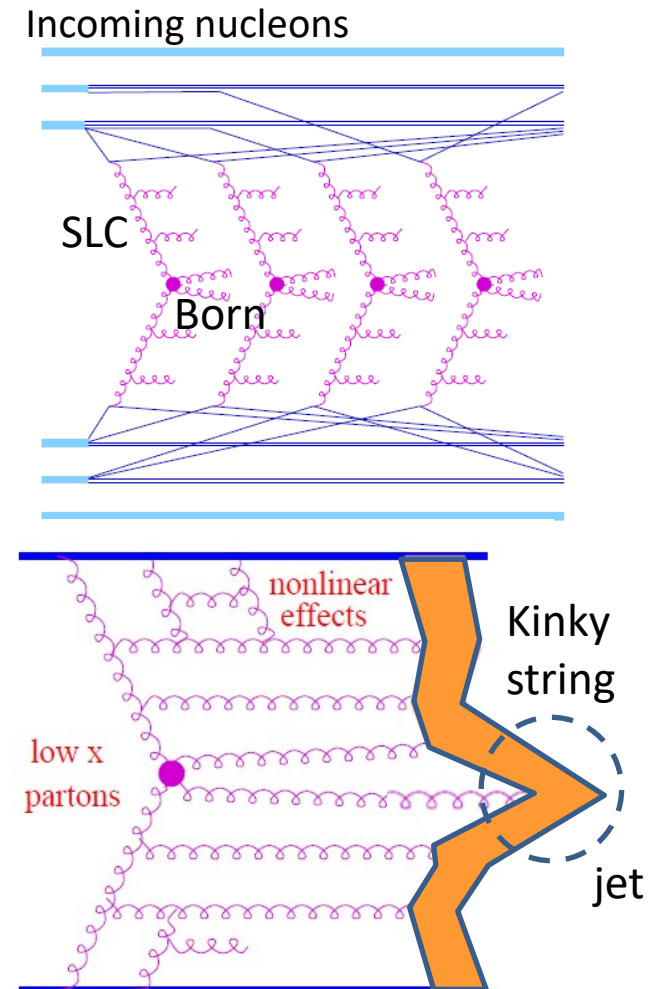
EPOS initial conditions

EPOS + Hydro : state of the art framework that encompass pp, pA and AA collisions

EPOS (initial conditions):

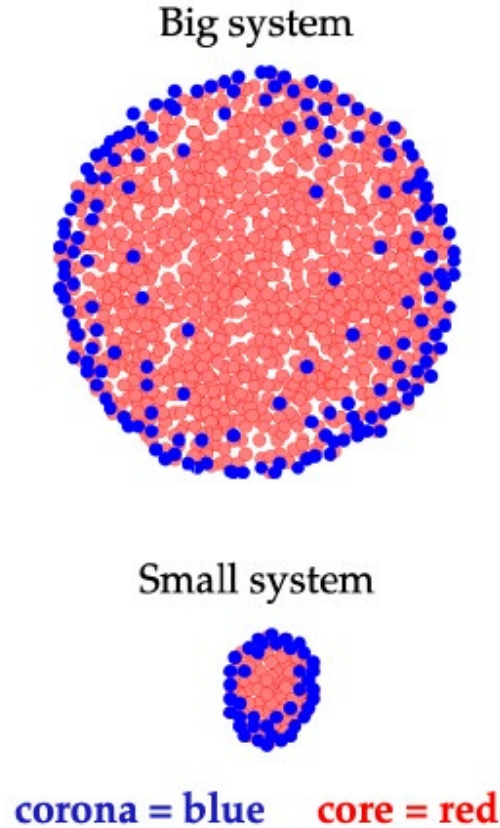
- Model based on Gribov-Regge multiple pomeron interactions
- Particle production (including HQ, from EPOS 3 on) in cut (semi-hard) pomerons, seen as partons ladder
- Space like DGLAP evolution with hard Born process
- Soft particles form a flux tube (string, with its own dynamics, incl. string breaking)... lots of them in A-A
- Slow string segments (pre-hadrons), far from the surface, are mapped to fluid dynamic fields
- Hard particles (kinky string) -> jets

Ref: K. Werner, Iu. Karpenko, M. Bleicher, T. Pierog, and S. Porteboeuf-Houssais Phys. Rev. C 85 (2012), 064907 + many refs in 2023 (<https://klaus.pages.in2p3.fr/epos4/>)

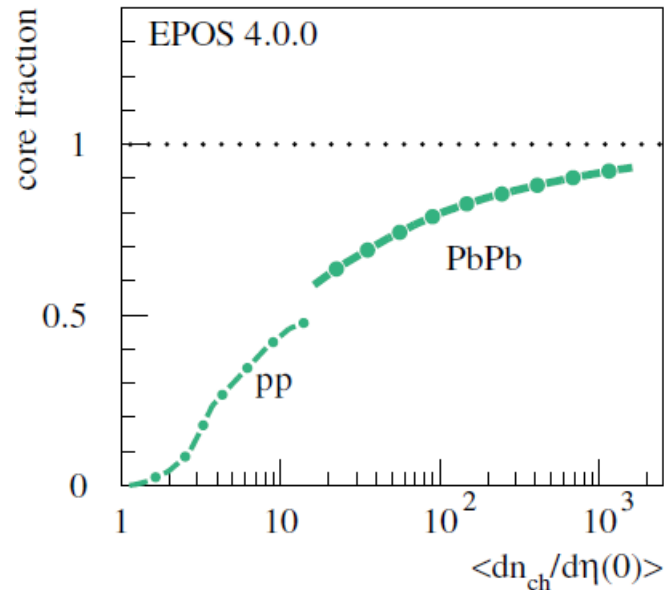


EPOS initial conditions

Simple but efficient initial-stage in EPOS : Core-Corona picture



- If the energy loss is bigger than the energy of the prehadron, it is considered to be part of the “core”
- If the energy loss is smaller than the energy, the prehadron escapes, it is called “corona”
- *Core: hydrodynamics; Corona: hadronic phase*



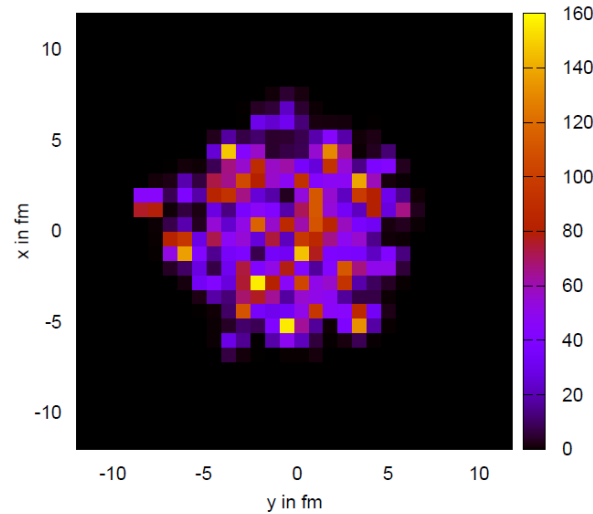
The energy density is larger than the critical energy density $\epsilon_0 \rightarrow$ deconfined QCD matter and hydro evolution

EPOS + VHHLE hydro as a background for HQ

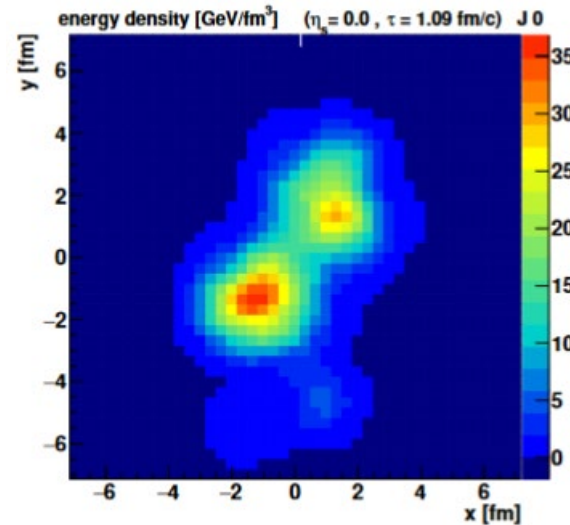
EPOS: state of the art framework that encompass pp, pA and AA collisions

Initial energy density

Beware: \neq color scales



EPOS Pb-Pb central



EPOS Pb-Pb peripheral

More realistic hydro and initial conditions => original HQ studies such as:

- 1) fluctuations in HQ observables
- 2) correlations between HF and light hadrons

≈1 year ago: EPOS4

<https://klaus.pages.in2p3.fr/epos4/>

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Code

EPOS4: A Monte Carlo tool for simulating high-energy scatterings

from electron-positron annihilation up to heavy ion collisions

Whereas elementary processes such as electron-positron annihilation and electron-nucleon scattering in EPOS4 are following standard procedures found in textbooks, the theoretical framework of both proton-proton and nucleus-nucleus scattering is quite unique, based on the (obvious!) fact that at very high energies, multiple (nucleonic or partonic) primary scatterings must happen in parallel, and not sequentially.

The EPOS4 approach brings together ancient knowledge about S-matrix theory (to deal with parallel scatterings) and modern concepts of perturbative QCD and saturation, going much beyond the usual factorization approach. The parallel scattering principle requires sophisticated Monte Carlo techniques, inspired by those used in statistical physics to investigate the Ising model.

| | | |
|--|---|--|
| Challenges and highlights in heavy ion physics | Why we need a parallel scattering framework | The essence of EPOS4: S-matrix, factorization, saturation (NEW: paper) |
| Heavy flavor in EPOS4 | Microcanonical hadronization and strangeness enhancement | Heavy ion collisions from 5TeV per nucleon down to 4GeV |
| A fast EPOS4 option, using parameterized fluid expansion | Flow harmonics in heavy ion collisions from 5TeV down to 7.7GeV | System size dependence of flow harmonics in pp and AA |

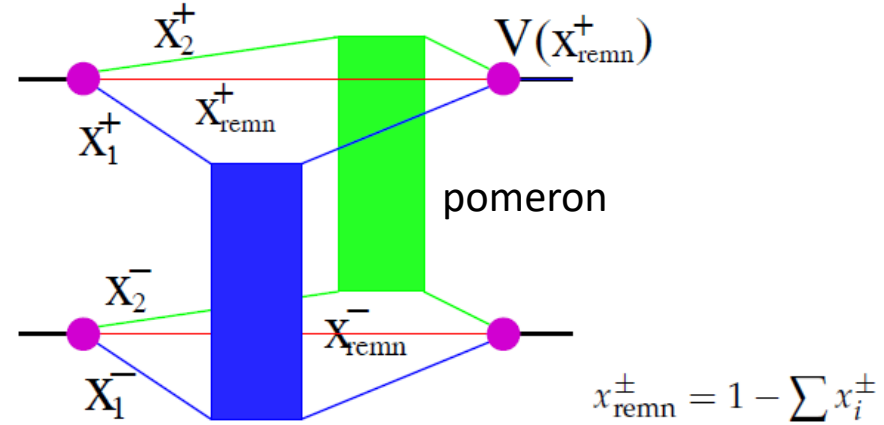
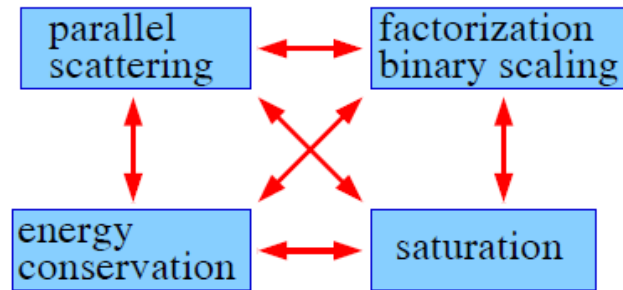
What is new in EPOS4?

The basic principle of treating (nucleonic or partonic) scatterings in parallel, based on S-matrix theory, has been used for two decades. But there is a major problem. For inclusive cross sections (but only for those!), important cancellations occur, leading to factorization (in pp) or binary scaling (in AA). In a parallel scattering scenario, these cancellations must come out (they cannot be imposed), which requires very high precision and good strategies -- and this is provided with EPOS4.

The treatment of parton ladders is completely redone, with unprecedented precision concerning the parton kinematics, in particular in case of heavy flavor partons being involved. Also, backward parton evolution in each of the

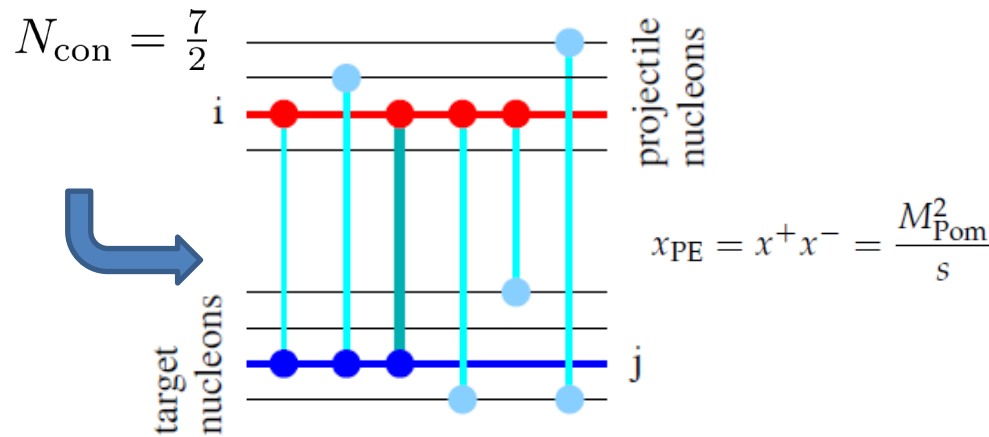
One Novelty in EPOS4: Curing the factorization issue

K. Werner. arXiv: 2301.12517

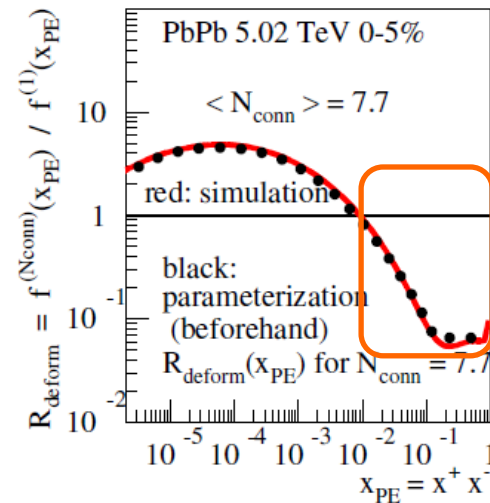


Double scattering diagram

// energy conservation

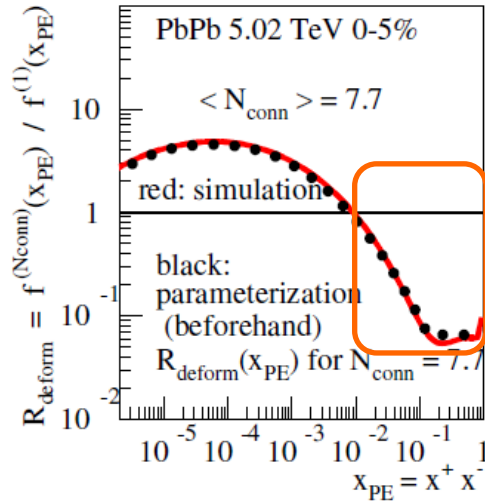


Invariant mass of the “hard pomeron” is reduced due to the connection of both projectile and target pomerons to other pomerons

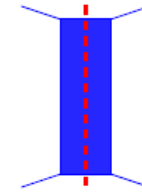


R_{deform} : Depletion of the high pomeron mass frequency due to many energy sharings

One Novelty in EPOS4: Curing the factorization issue



R_{deform} : Depletion of the high pomeron mass frequency due to many energy sharings



$$= G(x^+, x^-, s, b)$$

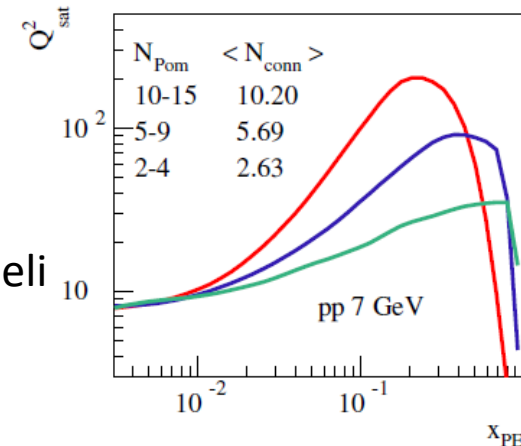
The cut pomeron

The “amazingly simple” solution in EPOS4: define $G(x^+, x^-, s, b) = \frac{n}{R_{\text{deform}}(x_{\text{PE}})} G_{\text{QCD}}(Q_{\text{sat}}^2, x^+, x^-, s, b)$

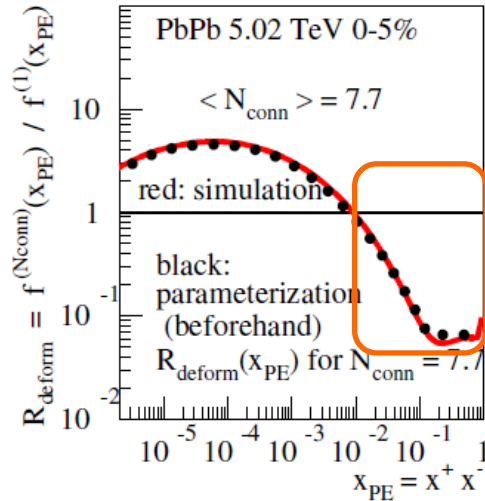
With $Q_{\text{sat}}(N_{\text{conn}}, x^+, x^-)$ chosen implicitly such that G does not depend on N_{conn} .

$$\frac{d^2\sigma_{\text{incl}}^{AB}(N_{\text{conn}})}{dx^+ dx^-} \propto \frac{d\sigma_{\text{incl}}^{\text{single Pom}}}{dx^+ dx^-} \left[Q_{\text{sat}}^2(N_{\text{conn}}, x^+, x^-) \right]$$

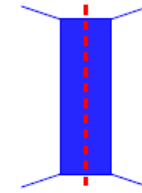
- which perfectly warrant the factorization at large p_T ; one recovers binary scaling (generalized Abramovskii Gribov Kancheli theorem).
- For large N_{conn} , low p_T is suppressed



One Novelty in EPOS4: Curing the factorization issue



R_{deform} : Depletion of the high pomeron mass frequency due to many energy sharings



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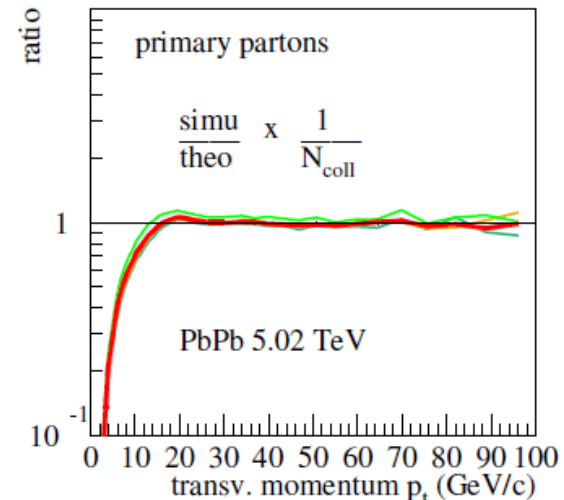
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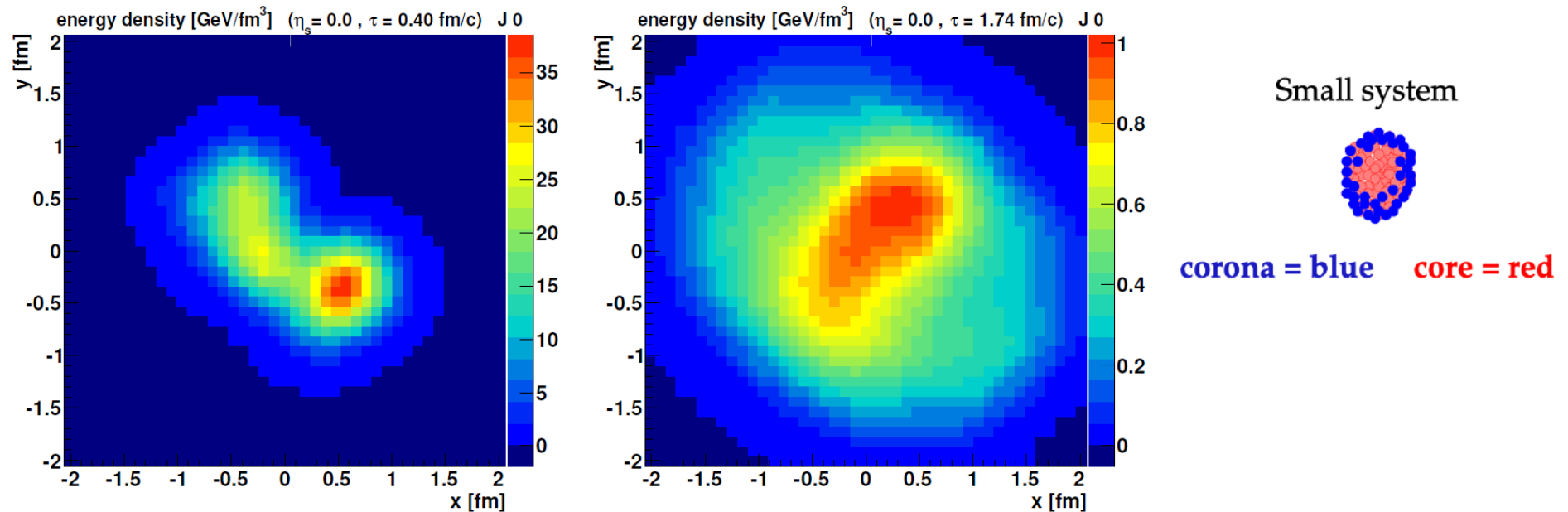
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EPOS4 for small systems

EPOS + Hydro : state of the art framework that encompass pp, pA and AA collisions

=> Go and look in pp



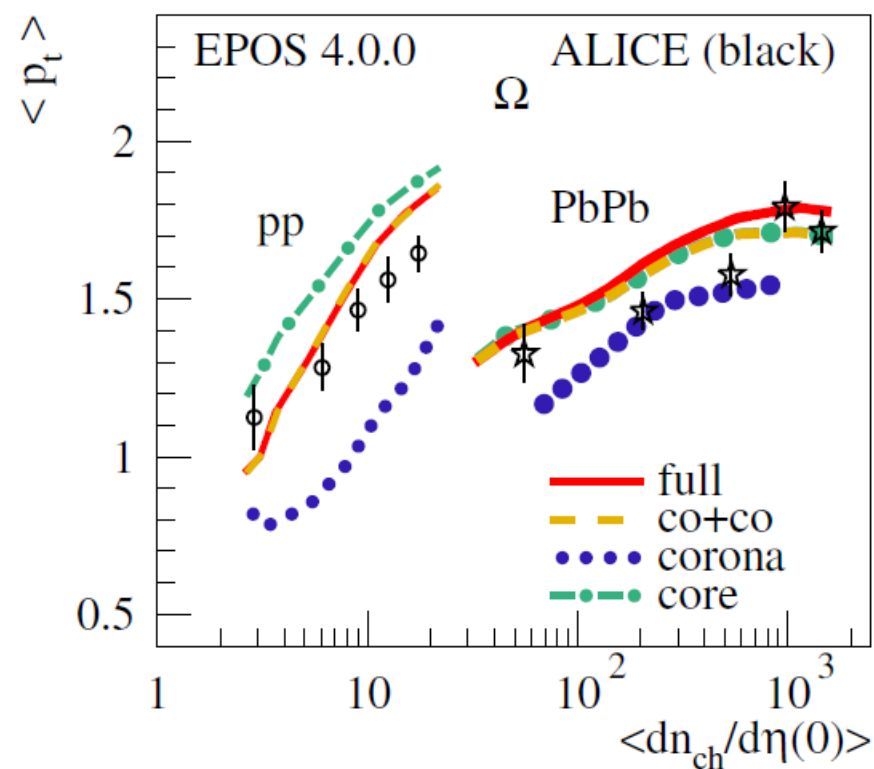
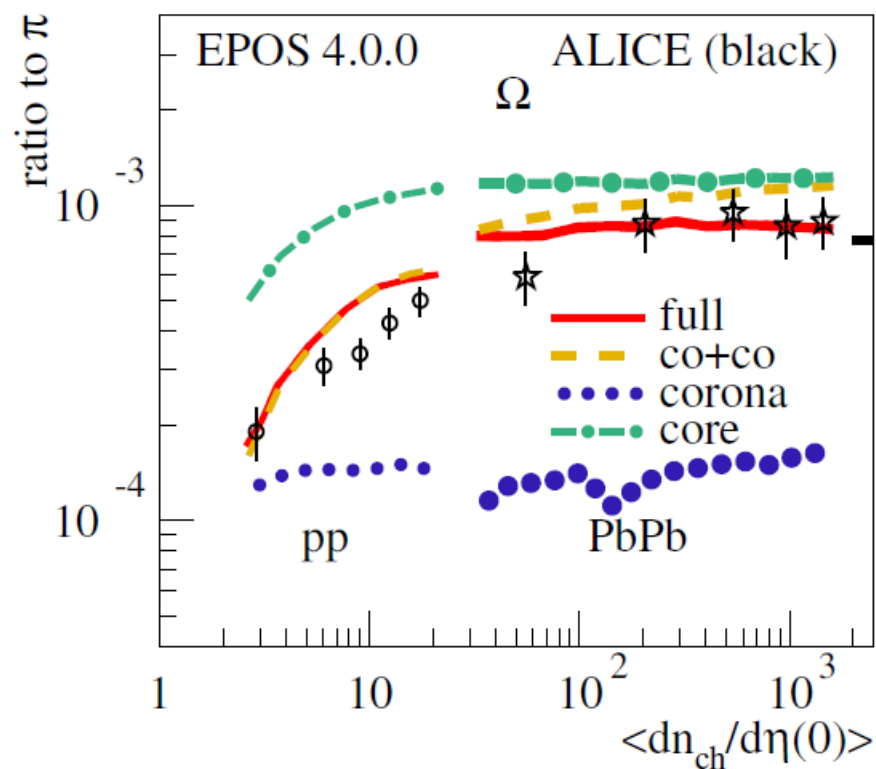
The energy density is larger than the critical energy density ϵ_0 —> deconfined QCD matter in pp as well !

=> In EPOS4, QGP droplet is one of the ingredients of collectivity

Full EPOS4: checking multiplicity dependencies and $\langle p_t \rangle$

continuous curve

jump



Affected by:

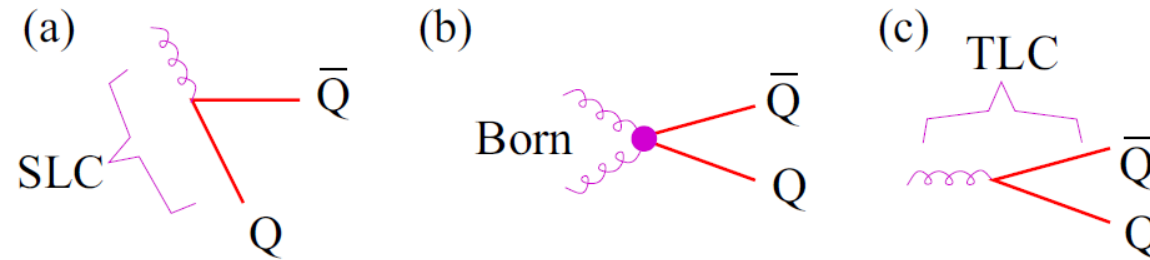
- core-corona
- microcanonical
- hadronic cascade (UrQMD)

- Saturation
- Flow
- core-corona

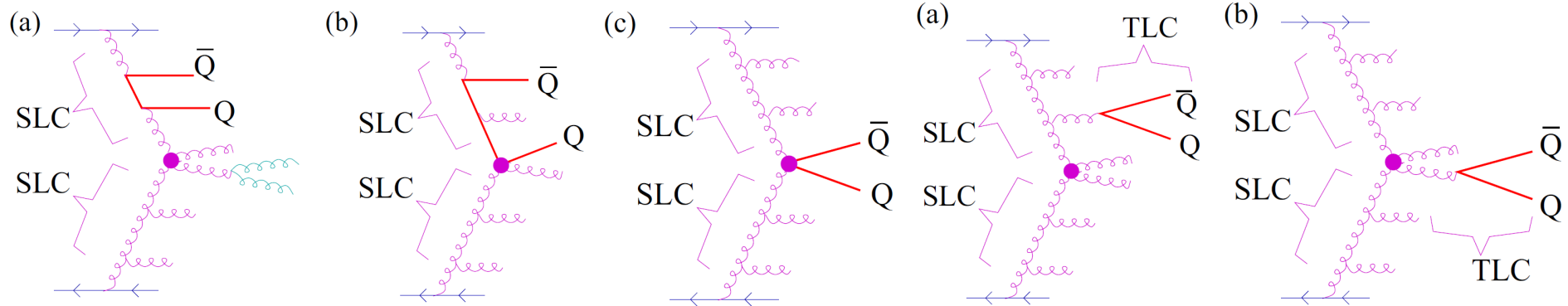
HQ sector: Improved HF production in EPOS4

Initial production of heavy quarks through

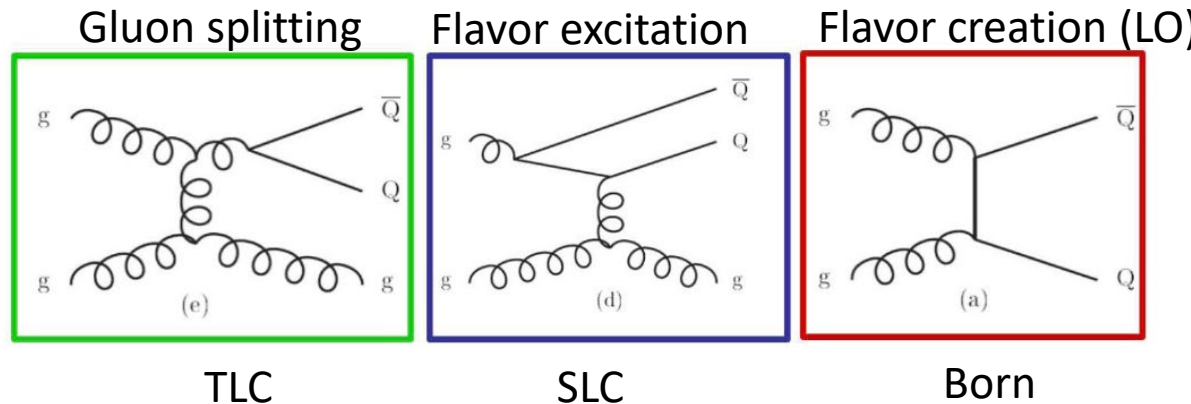
K. Werner, B. Guiot, Phys.Rev.C 108 (2023) 3, 034904



Found in the following “evolutions” :



Includes the 3 basic mechanisms present in other MC generator like Pythia



Later hadronized through string breaking in EPOS4

From EPOS4 to (MC@s)HQ

QGP Bulk (t) + initial
distribution of HQ

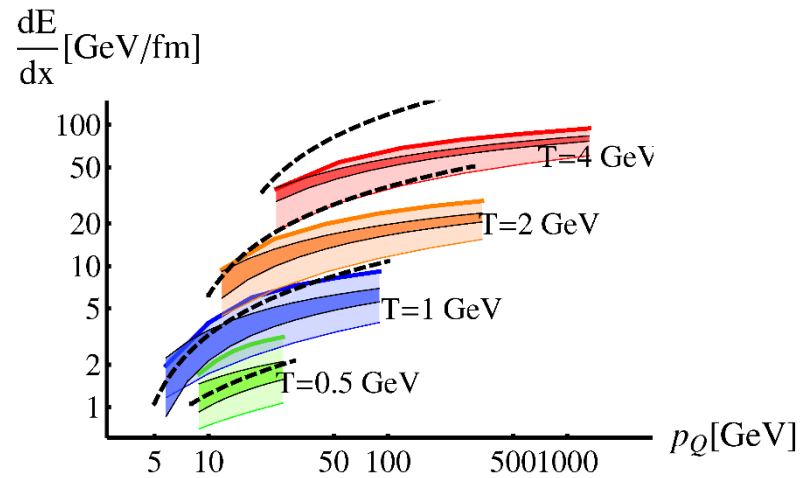
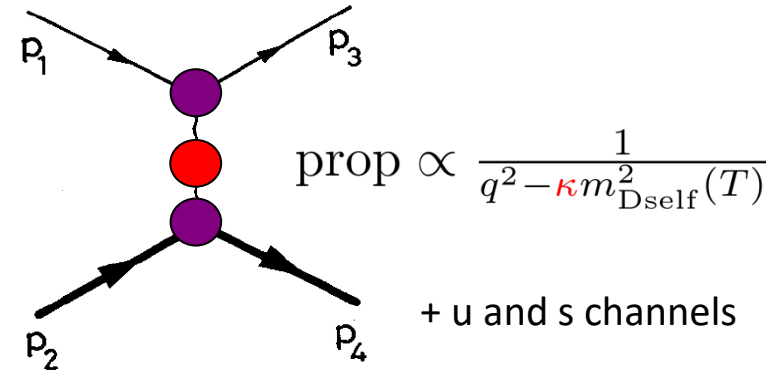


The core energy loss from the HQ part

Collisional component

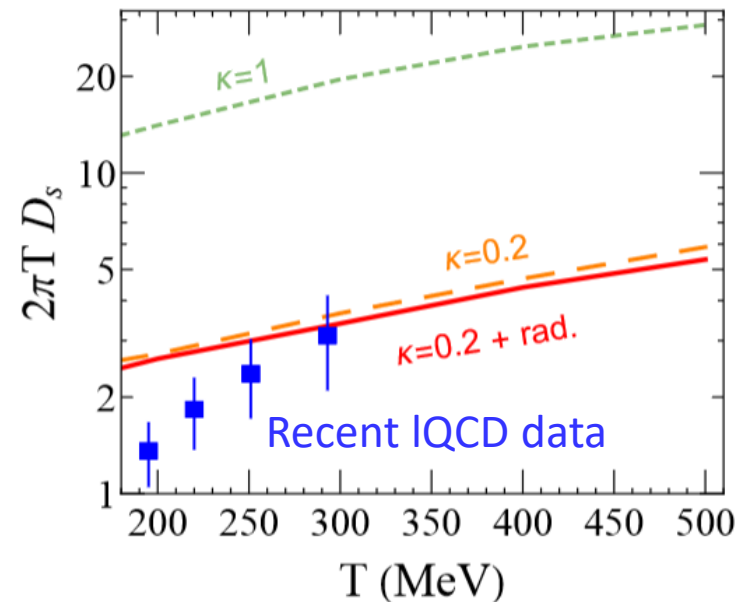
- One-gluon exchange model: reduced IR regulator $\kappa m_{\text{Dself}}^2$ in the hard propagator, fixed on HTL Energy loss (maximal insensitivity of dE/dx on q^*)
- Running coupling $\alpha_{\text{eff}}(t)$
- self consistent Debye mass

$$m_{\text{Dself}}^2(T) = (1+n_f/6) 4\pi\alpha_{\text{eff}}(m_{\text{Dself}}^2)T^2$$



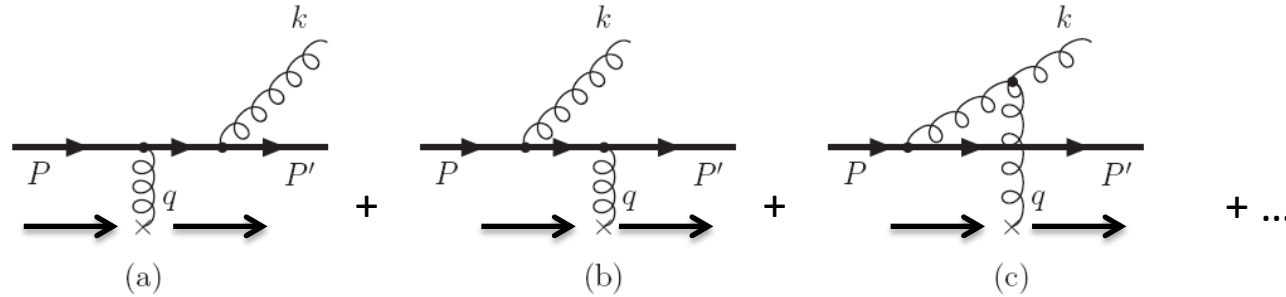
Comparison with Peigné-Peshier at finite momentum

$$D_s = \lim_{p_Q \rightarrow 0} T / (M_Q \eta_D) \quad \eta_D : \text{drag coefficient}$$



The core energy loss from the HQ part

Radiative component



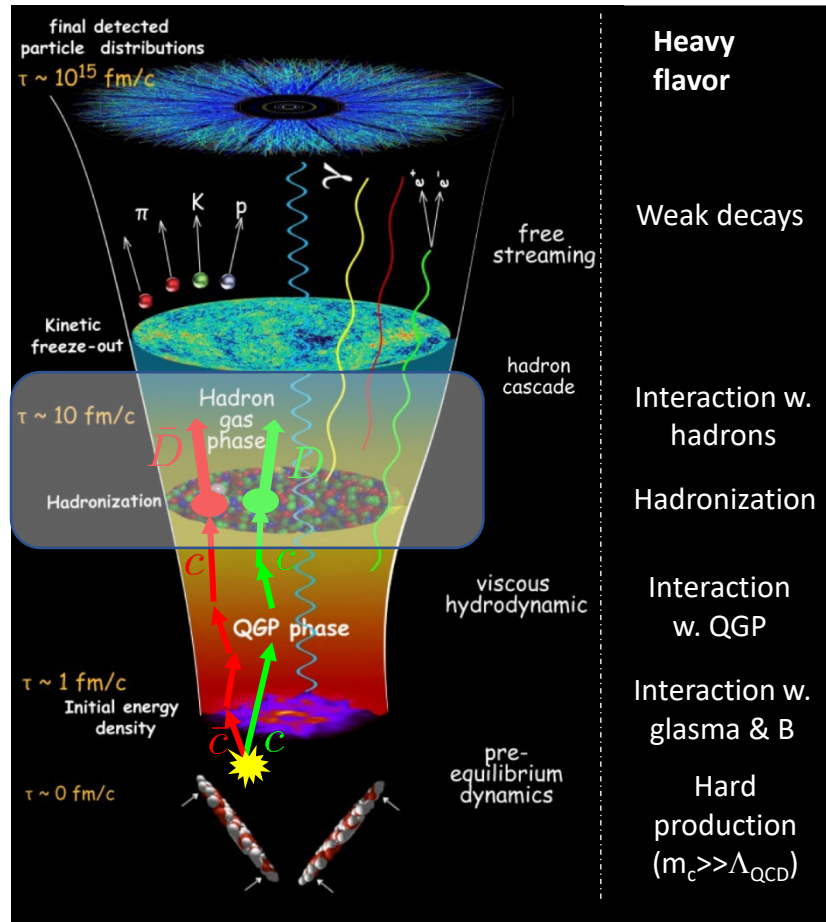
- Extension of Gunion-Bertsch approximation beyond mid-rapidity and to finite mass m_Q) distribution of induced gluon radiation per collision ($\Delta E_{\text{rad}} \propto E L$):

$$P_g(x, \mathbf{k}_\perp, \mathbf{q}_\perp, m_Q) = \frac{3\alpha_s}{\pi^2} \frac{1-x}{x} \left(\frac{\mathbf{k}_\perp}{\mathbf{k}_\perp^2 + x m_Q^2} - \frac{\mathbf{k}_\perp - \mathbf{q}_\perp}{(\mathbf{k}_\perp - \mathbf{q}_\perp)^2 + x m_Q^2} \right)^2$$

- LPM effect for moderate gluon energy

Implemented in EPOS4-HQ through Boltzmann transport

Heavy quarks (Q) as ideal hard probes:

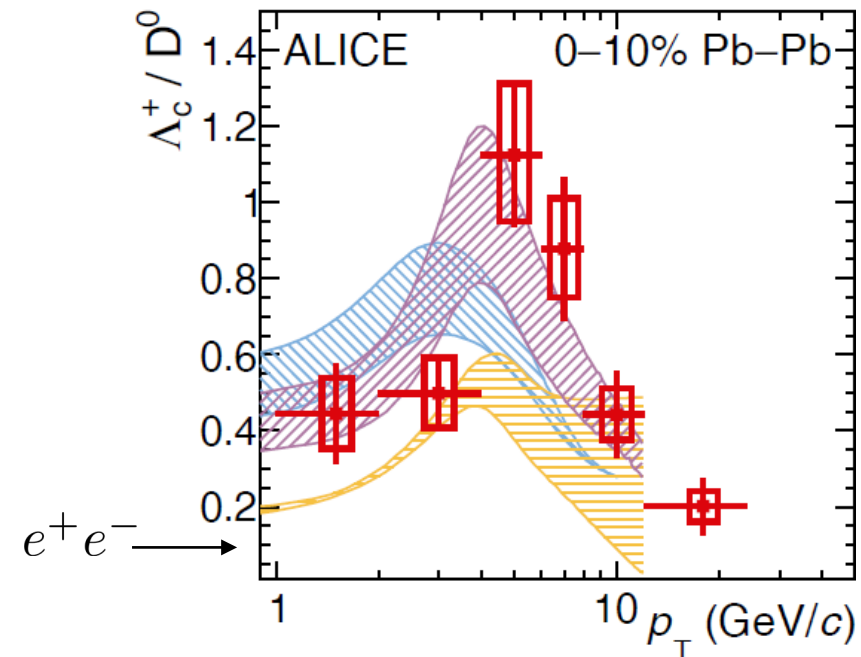


The recombination of heavy quark with some *existing* light quark(s) from the QGP is an essential mechanism at “low” $p_T < 5-10 \text{ GeV}/c$...

Mandatory to understand the Λ_c/D^0 ratio

- Evolution in the QGP DOES NOT modify the yield of initial Q and Qbar: Negligible annihilation rate !
- It only impacts their distribution in momentum
 - Hence, for usual observables like RAA and v_2 ... only the initial “1 body” distribution matters

Mostly like in elementary pp collisions...
- ... However hadronization is affected by the QGP
 - Other mechanism wrt usual fragmentation of HQ in elementary collisions : **coalescence / recombination**



The coalescence + fragmentation hadronization

When the local energy density is lower than the critical value ($T \sim 165 \text{ MeV}$)

Heavy quarks hadronize via **coalescence + fragmentation** in **EPOS4HQ!**



$$\frac{dN}{d^3\mathbf{P}} = g_H \sum_{N_Q} \int \prod_{i=1}^k \frac{d^3 p_i}{(2\pi)^3} f(\mathbf{p}_i) W_H(\mathbf{p}_1, \dots, \mathbf{p}_i) \delta^{(3)}\left(\mathbf{P} - \sum_{i=1}^N \mathbf{p}_i\right),$$

EPOS4 with only string fragmentation

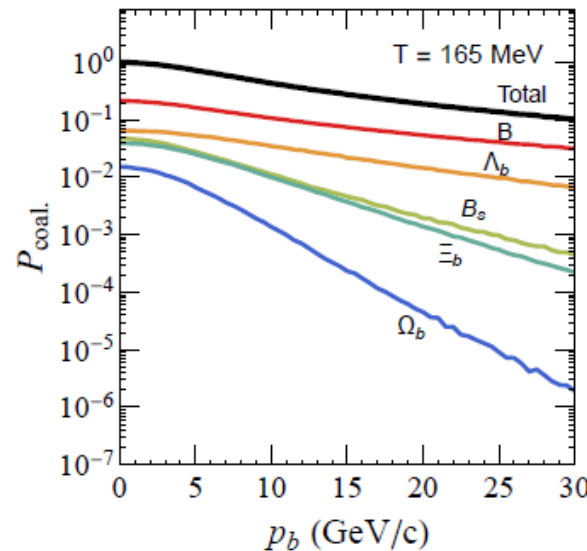
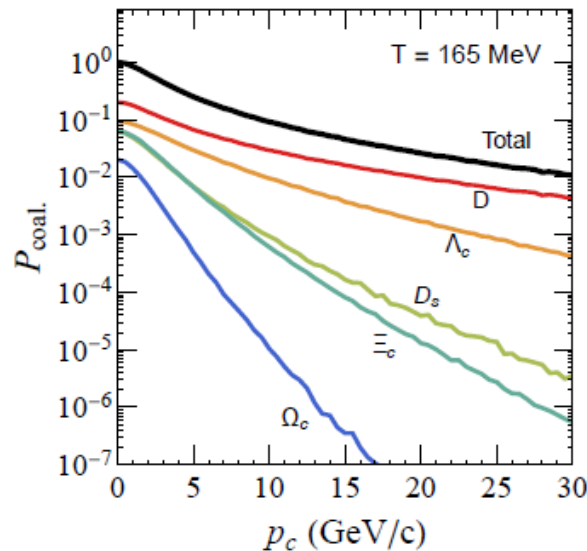
$1 - P_{\text{coal.}}$ for fragmentation (HQET based fragmentation function)

J. Zhao's work

We include almost all hadrons (missing baryons predicted by the potential model; $17D, 10D_s, 38\Lambda_c, 54\Sigma_c, 92\Xi_c, 54\Omega_c$; except the rare HF hadrons)

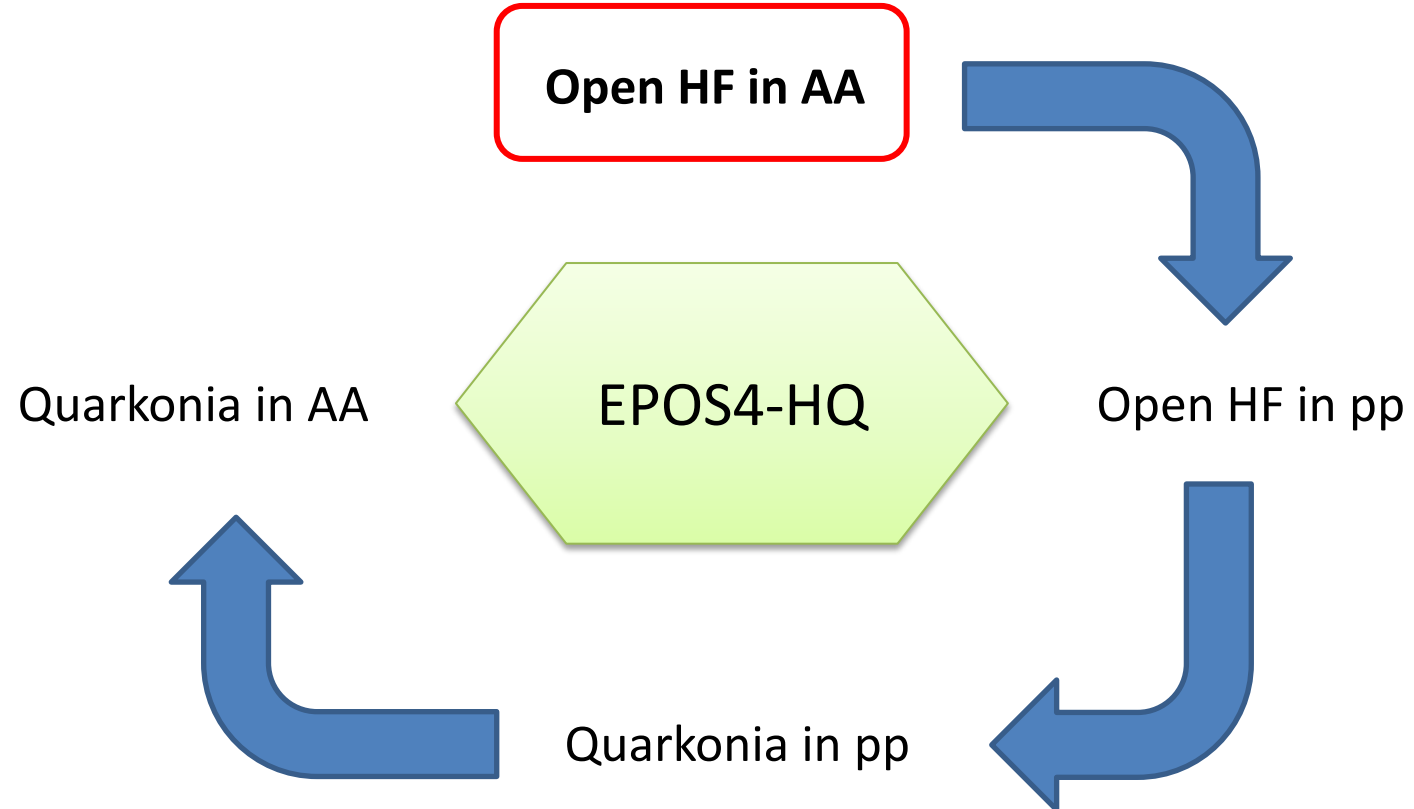
Ground states Wigner density: $W(p_r) = (2\sqrt{\pi}\sigma)^3 e^{-\sigma^2 p_r^2}$ Width is given by the potential model

Excited states are involved via the thermal ratio: $n_i = \frac{g_i}{2\pi^2} T_{\text{FO}} m_i^2 K_2\left(\frac{m_i}{T_{\text{FO}}}\right)$ $R^m = n_{\text{excited}}^m / n_{\text{ground}}$.



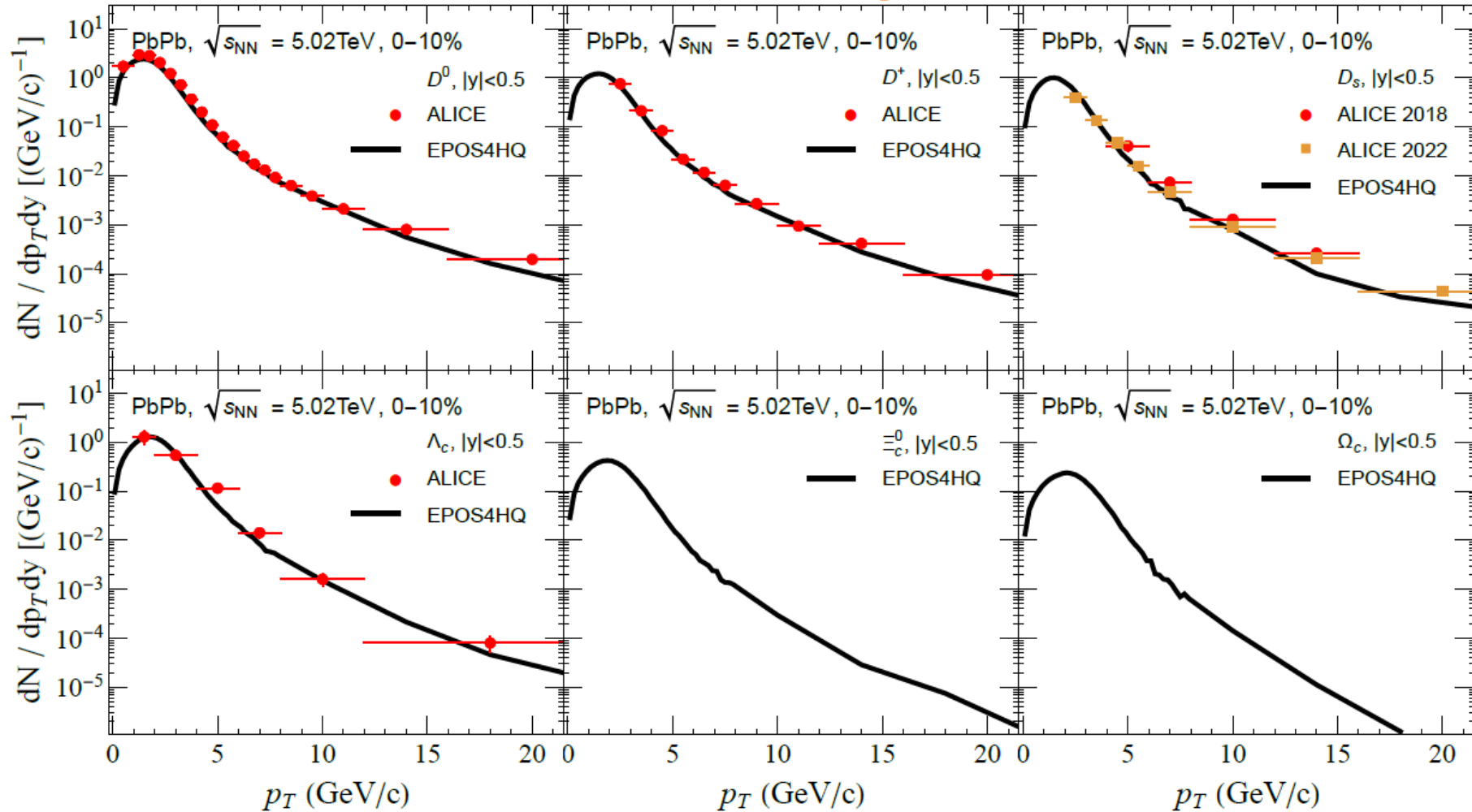
After hadronization, evolution in hadronic phase \rightarrow UrQMD

HF in EPOS4-HQ



OHF hadrons p_T distributions in EPOS4-HQ

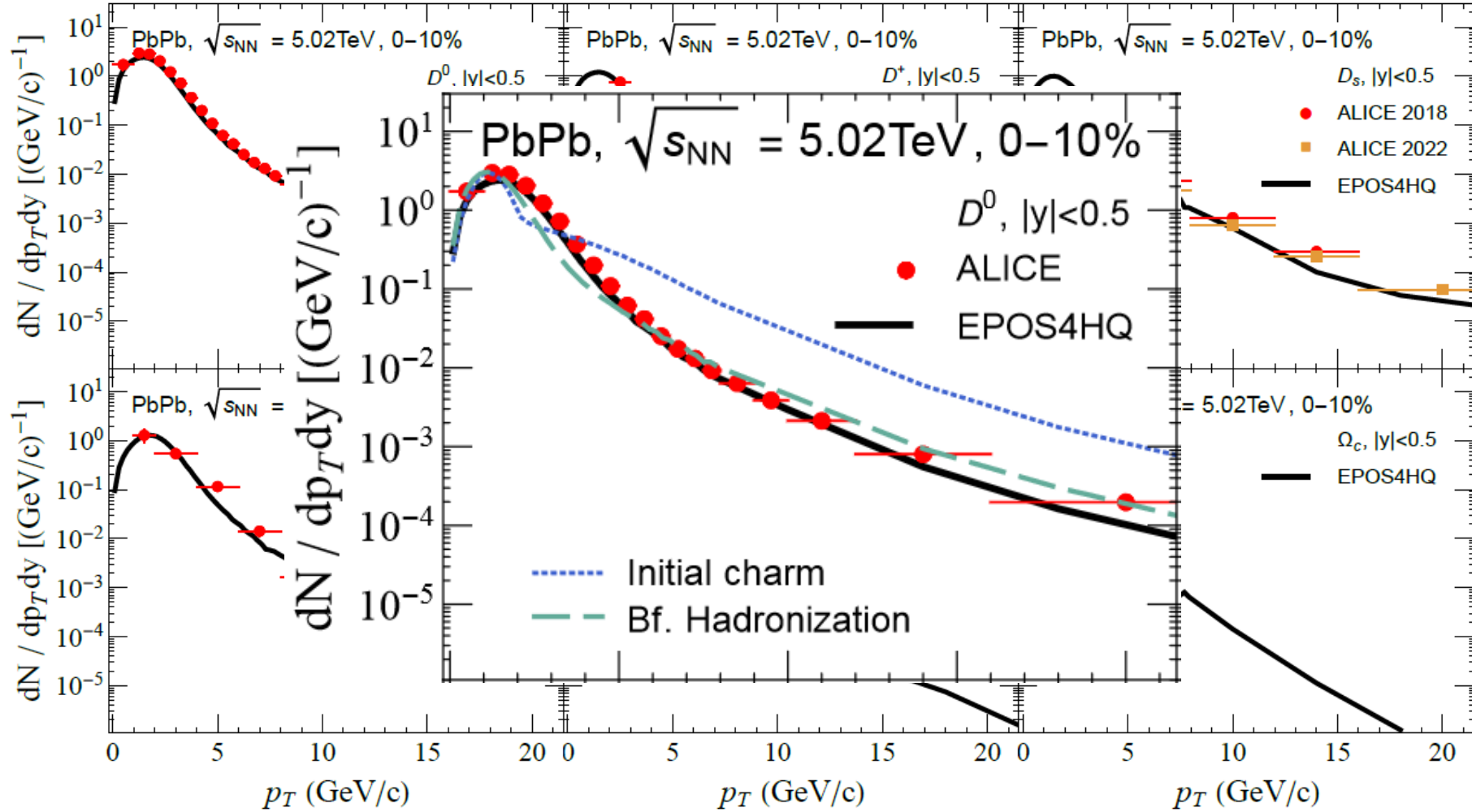
PbPb central @ LHC



Very good agreement for all resonances

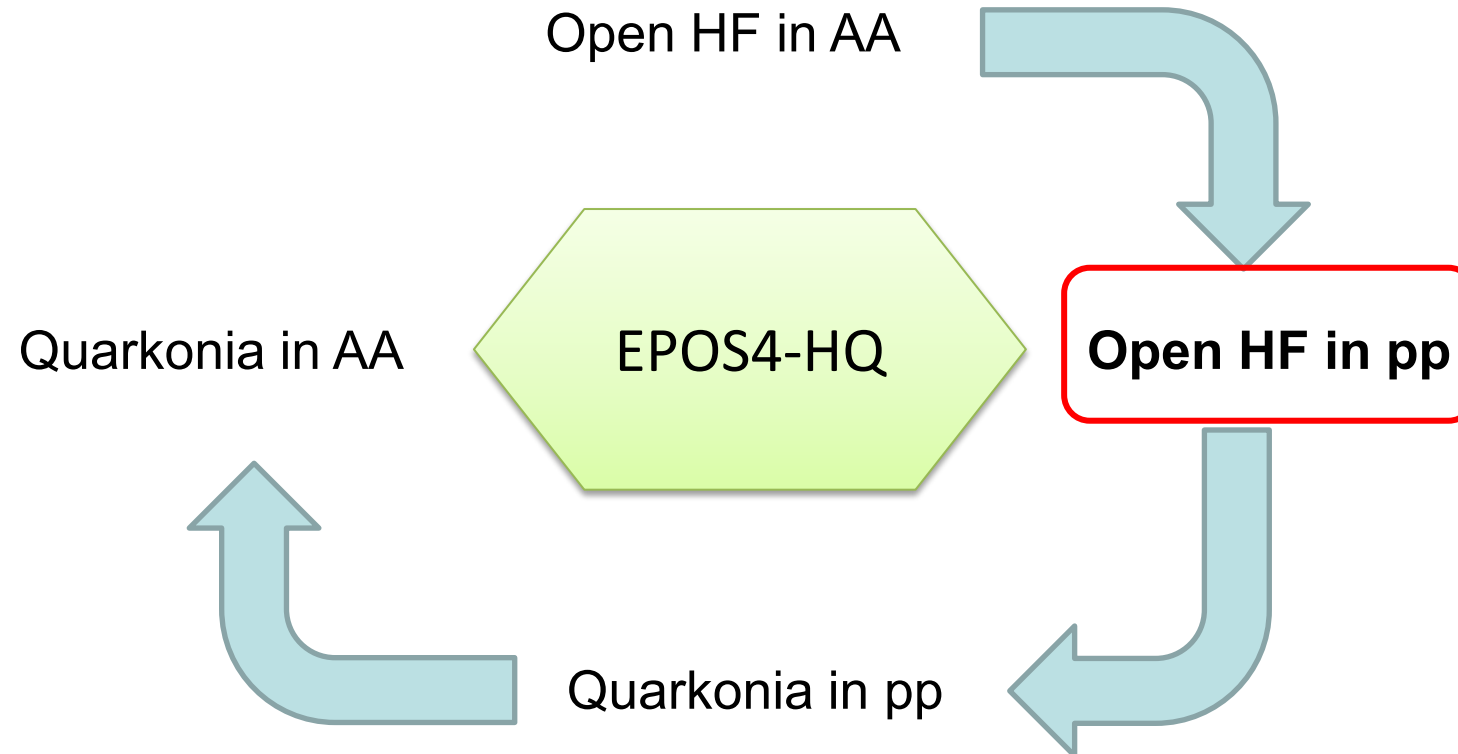
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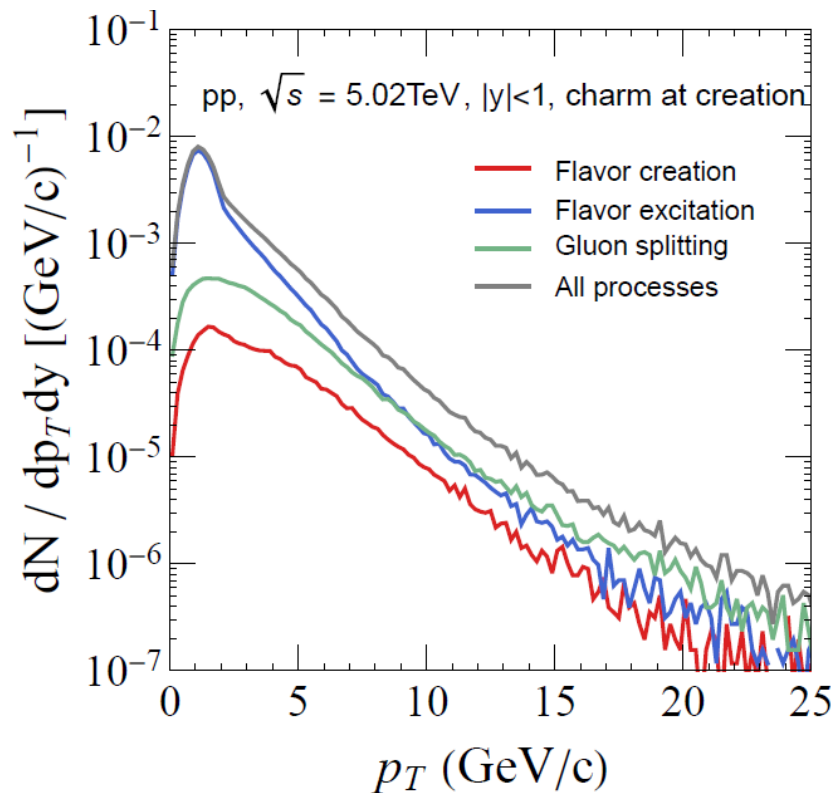
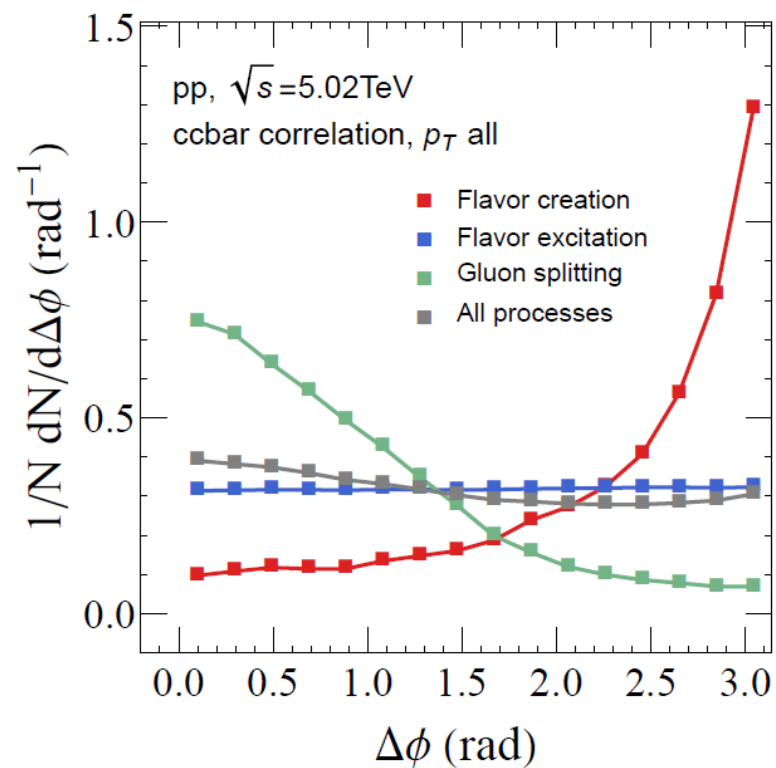
Largest effect in AA : Energy loss of c-quarks

HF in EPOS4-HQ

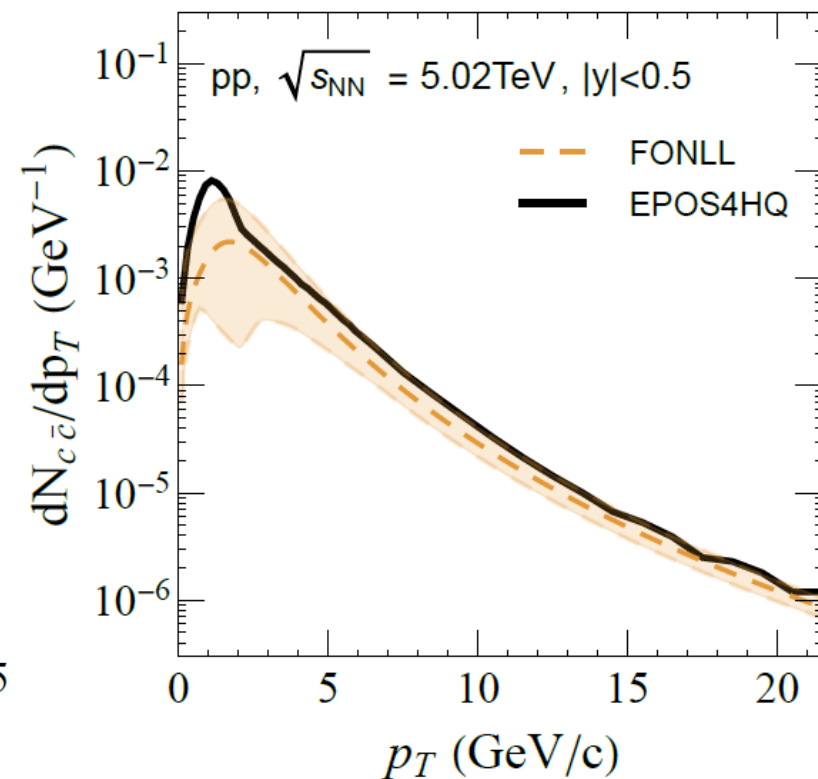


HF production in pp

See *K. Werner. arXiv: 2306.02396*



Dominance of the flavor excitation mechanism at small p_t

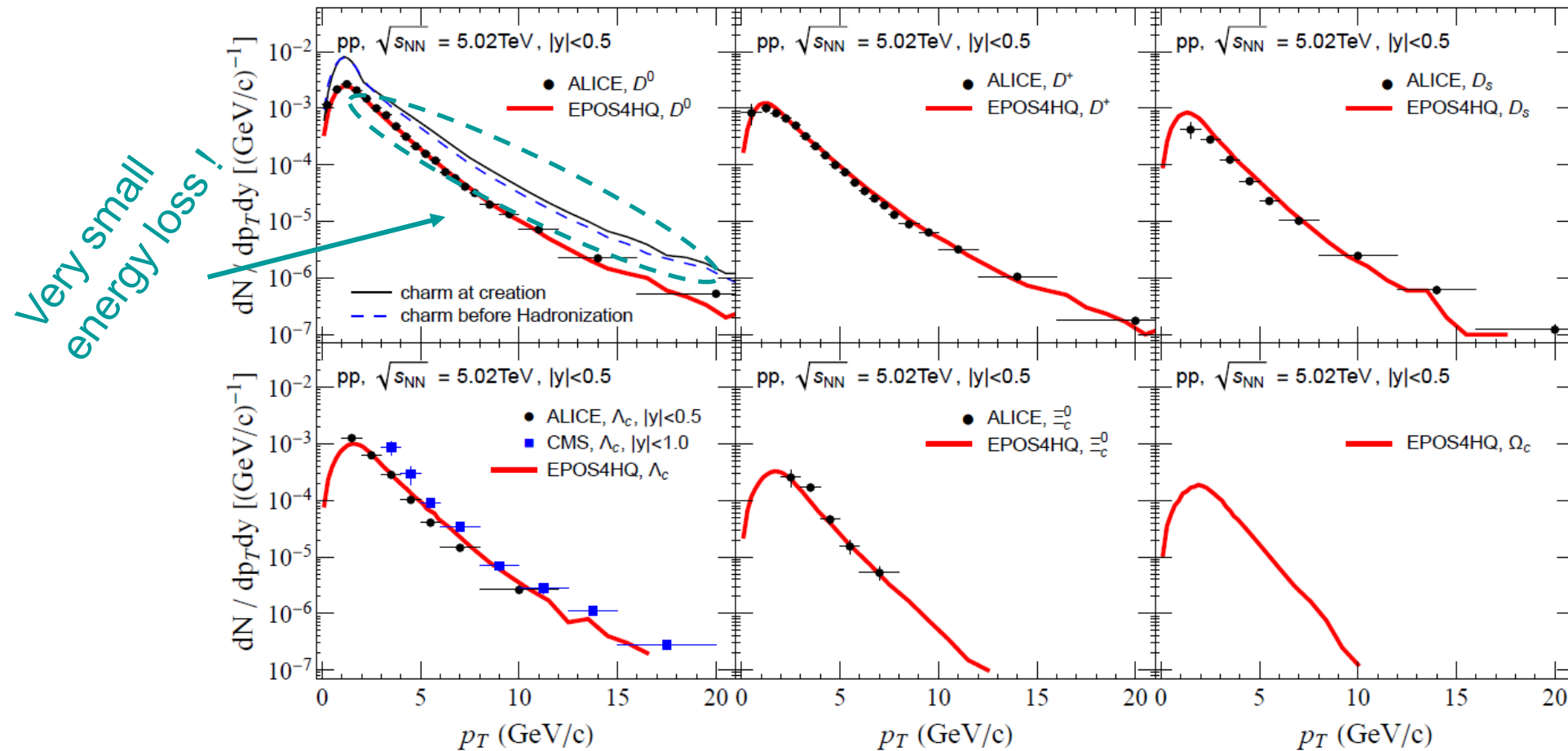


Some overshooting of the FONLL uncertainty band below 1 GeV, but very good agreement at large p_T where FONLL is best justified

Hadrons yield in pp

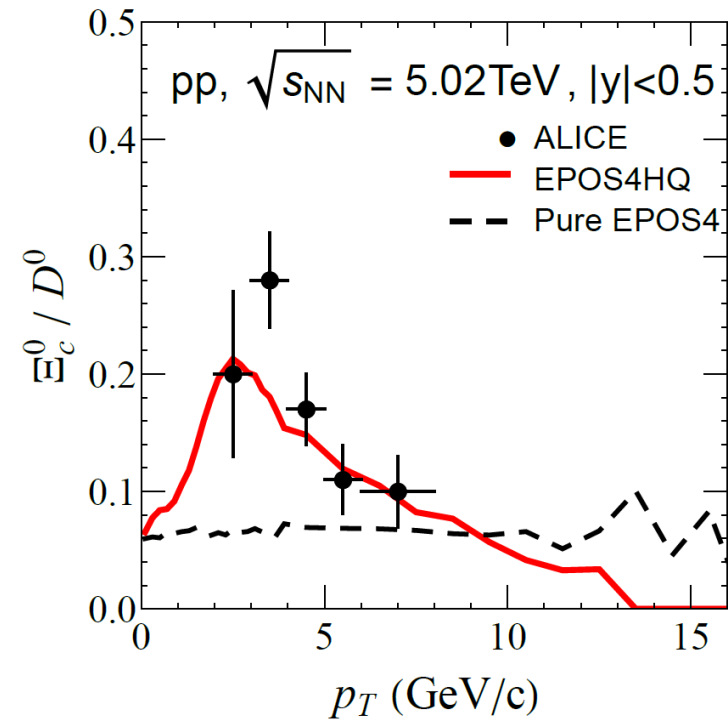
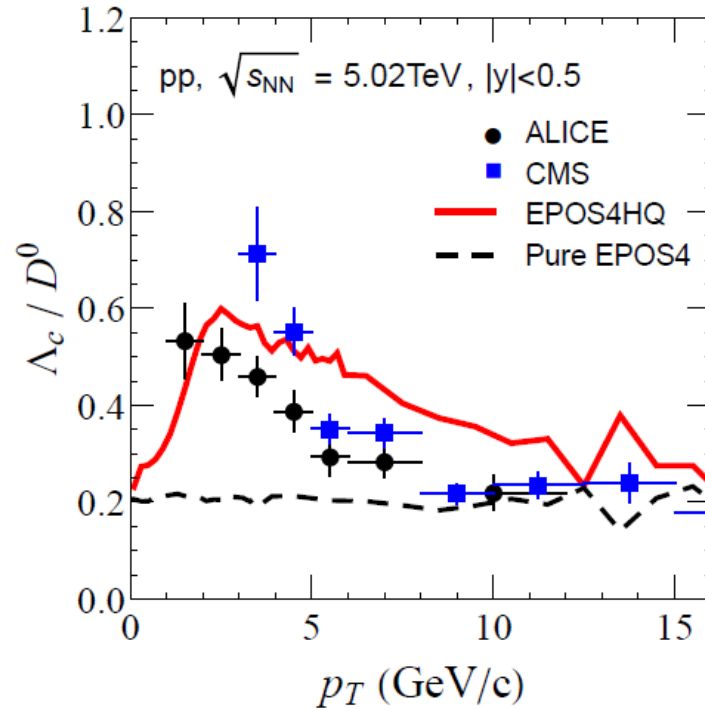
Heavy flavor as a probe of hot QCD matter produced in proton-proton collisions

Jiaying Zhao, Joerg Aichelin, Pol Bernard Gossiaux, and Klaus Werner
Phys. Rev. D **109**, 054011 – Published 6 March 2024



Good agreement in the pp sector, essentially due to the **coalescence + fragmentation hadronization**

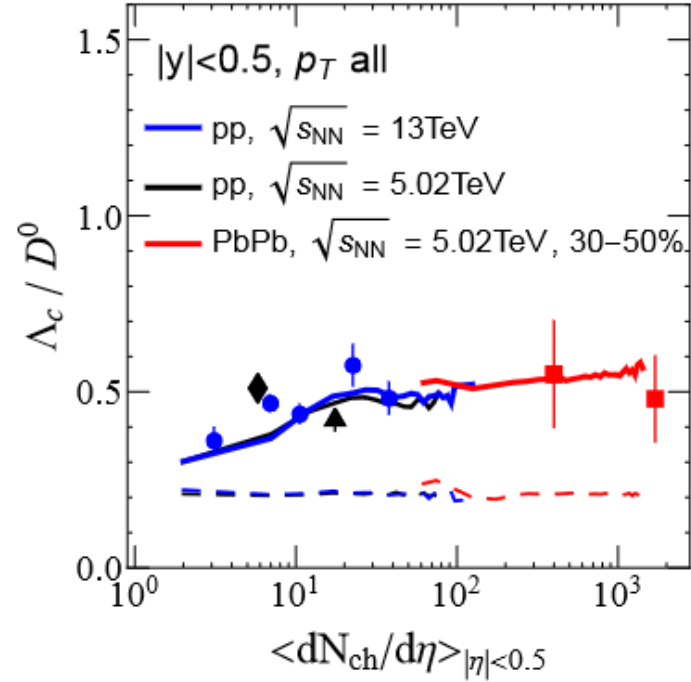
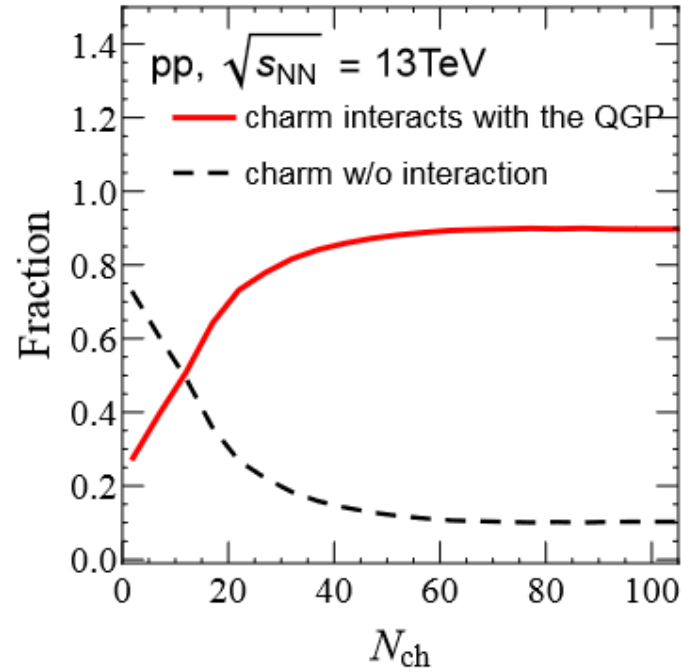
Yield ratios in pp



The coalescence + fragmentation hadronization is also successful in describing the yield ratio between charmed baryon to meson !

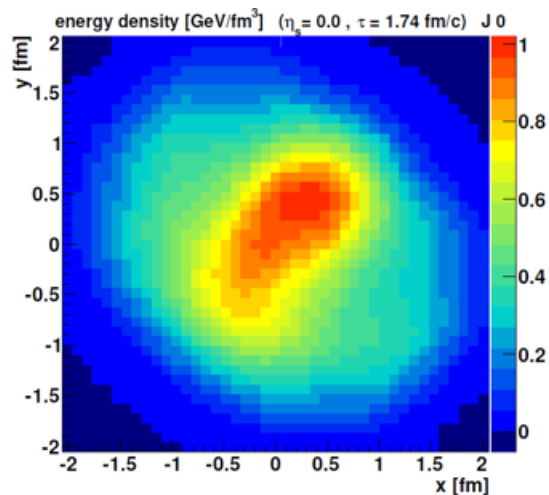
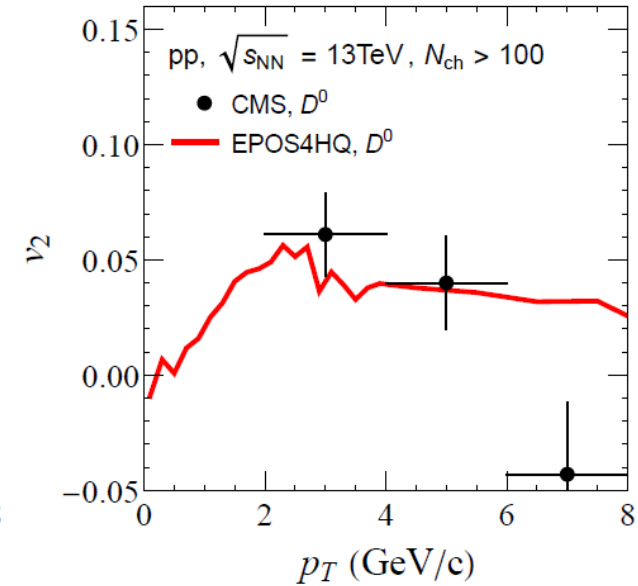
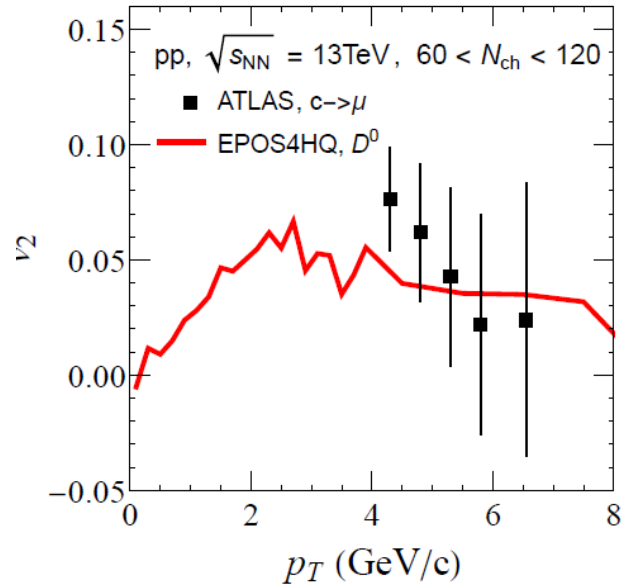
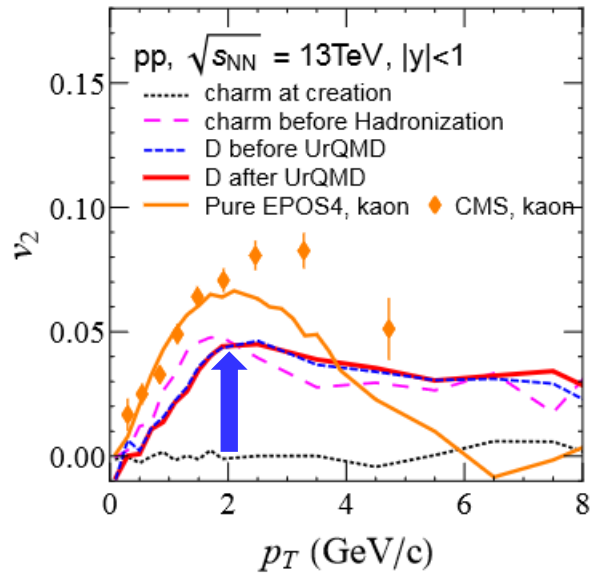
See as well : M. He and R. Rapp, Phys. Lett. B 795, 117 (2019), V. Minissale, S. Plumari, and V. Greco, Phys. Lett. B 821, 136622 (2021), H.-h. Li, F.-l. Shao, and J. Song, Chin. Phys. C 45, 113105 (2021), A. Beraudo, A. De Pace, D. Pablos, F. Prino, M. Monteno, and M. Nardi, arXiv:2306.02152

QGP droplet in pp ?



The yield ratio increase with centrality is correlated with the fraction of c quarks which interact with the QGP droplet

Azimuthal distributions in pp: v_2



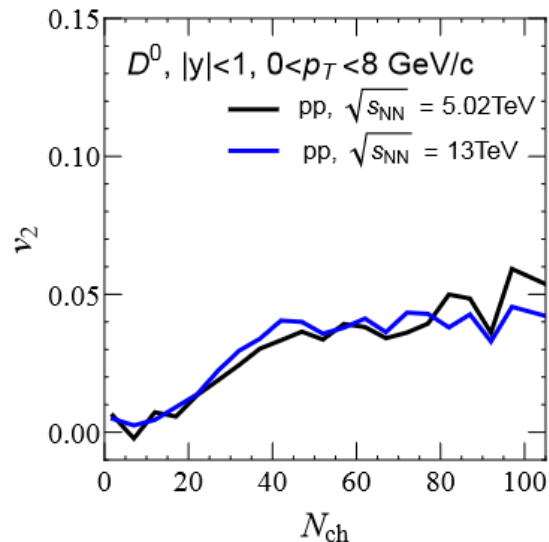
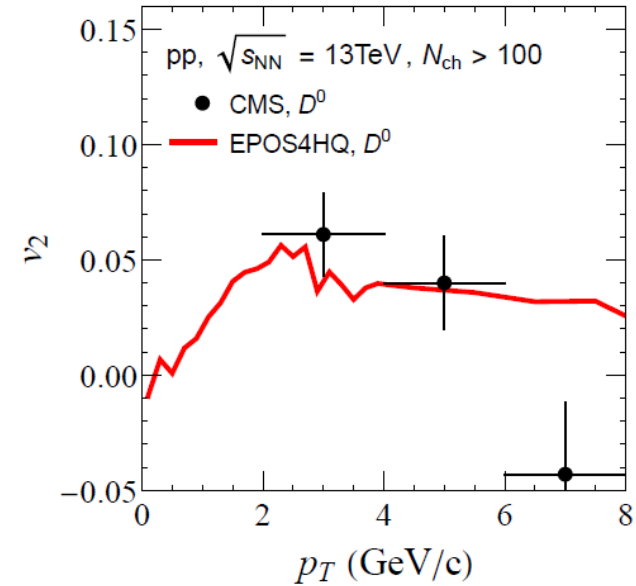
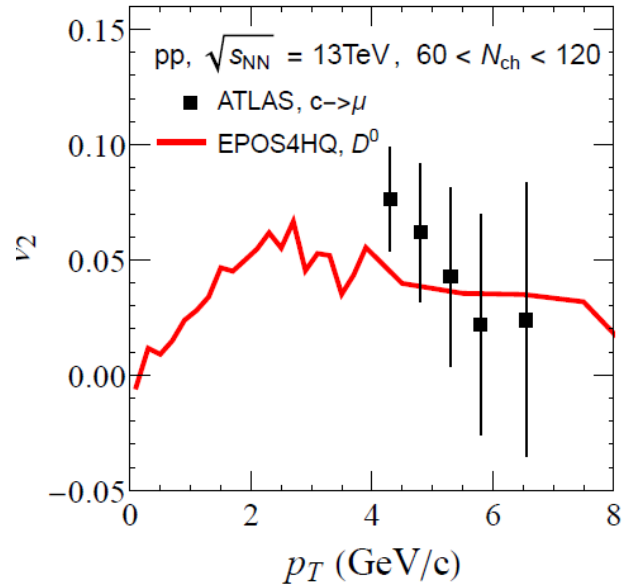
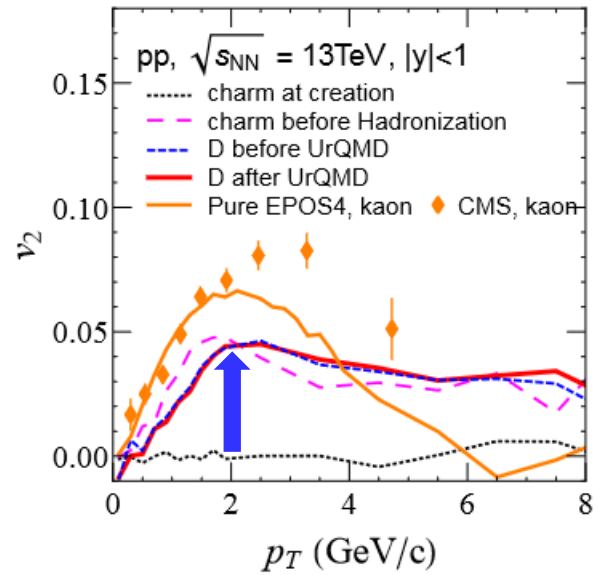
EPOS4HQ describes well the elliptic flow of D meson !

Clear sign of momentum redistribution during the short evolution (only comes through the v_2 in pp)

Little effect of the hadronic stage

See as well : A. Beraudo, A. De Pace, D. Pablos, F. Prino, M. Monteno, and M. Nardi, arXiv:2306.02152

Azimuthal distributions in pp: v_2



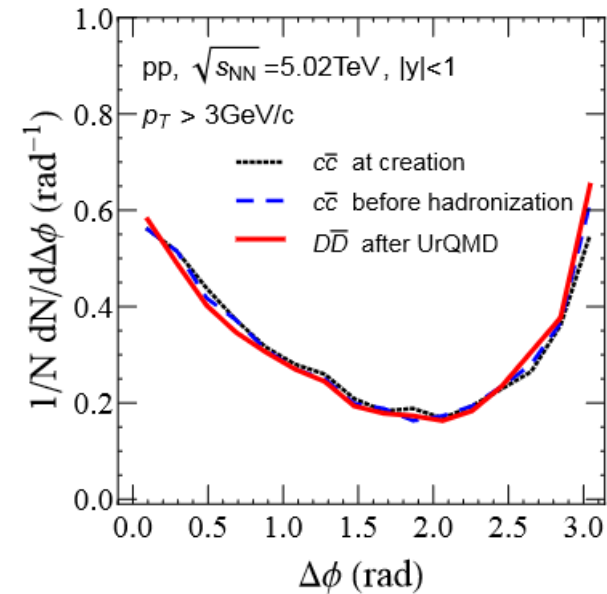
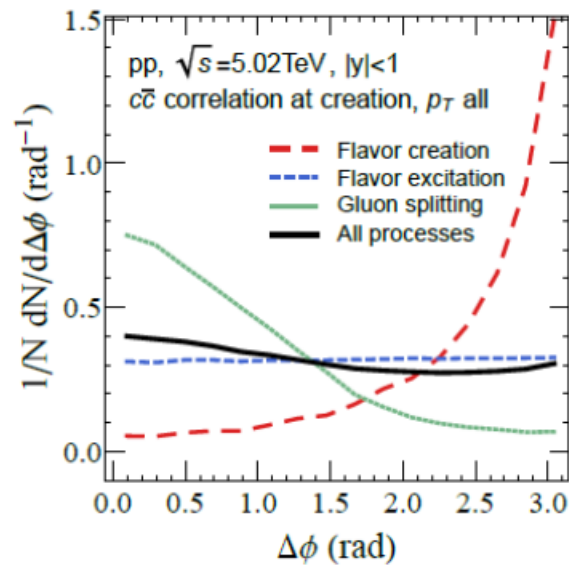
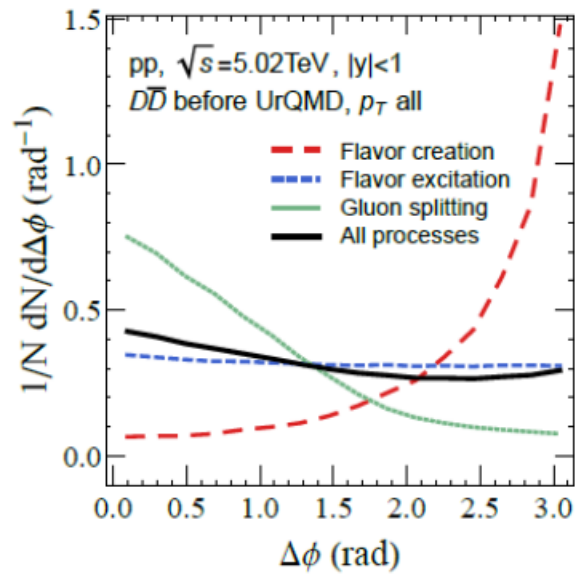
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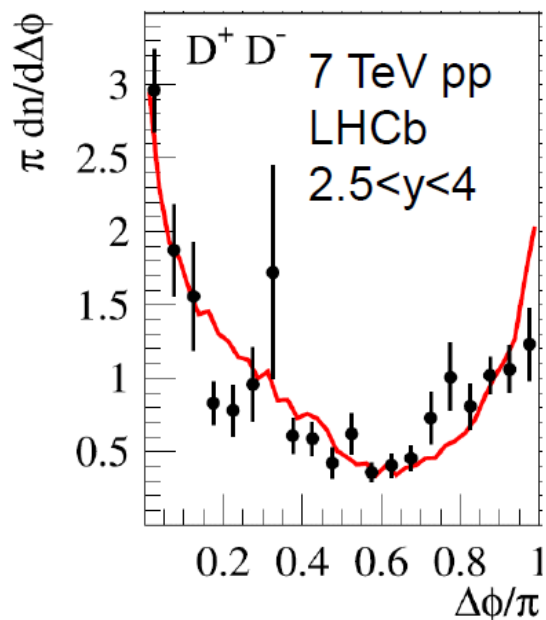
Little effect of the hadronic stage

See as well : A. Beraudo, A. De Pace, D. Pablos, F. Prino, M. Monteno, and M. Nardi, arXiv:2306.02152

Azimuthal distributions in pp



JHEP06(2012)141



**Small modification to the azimuthal correlations
 => constrain the HQ production process via final
 $D\bar{D}$ correlations!**

Good agreement with the experiment (also for other correlations)

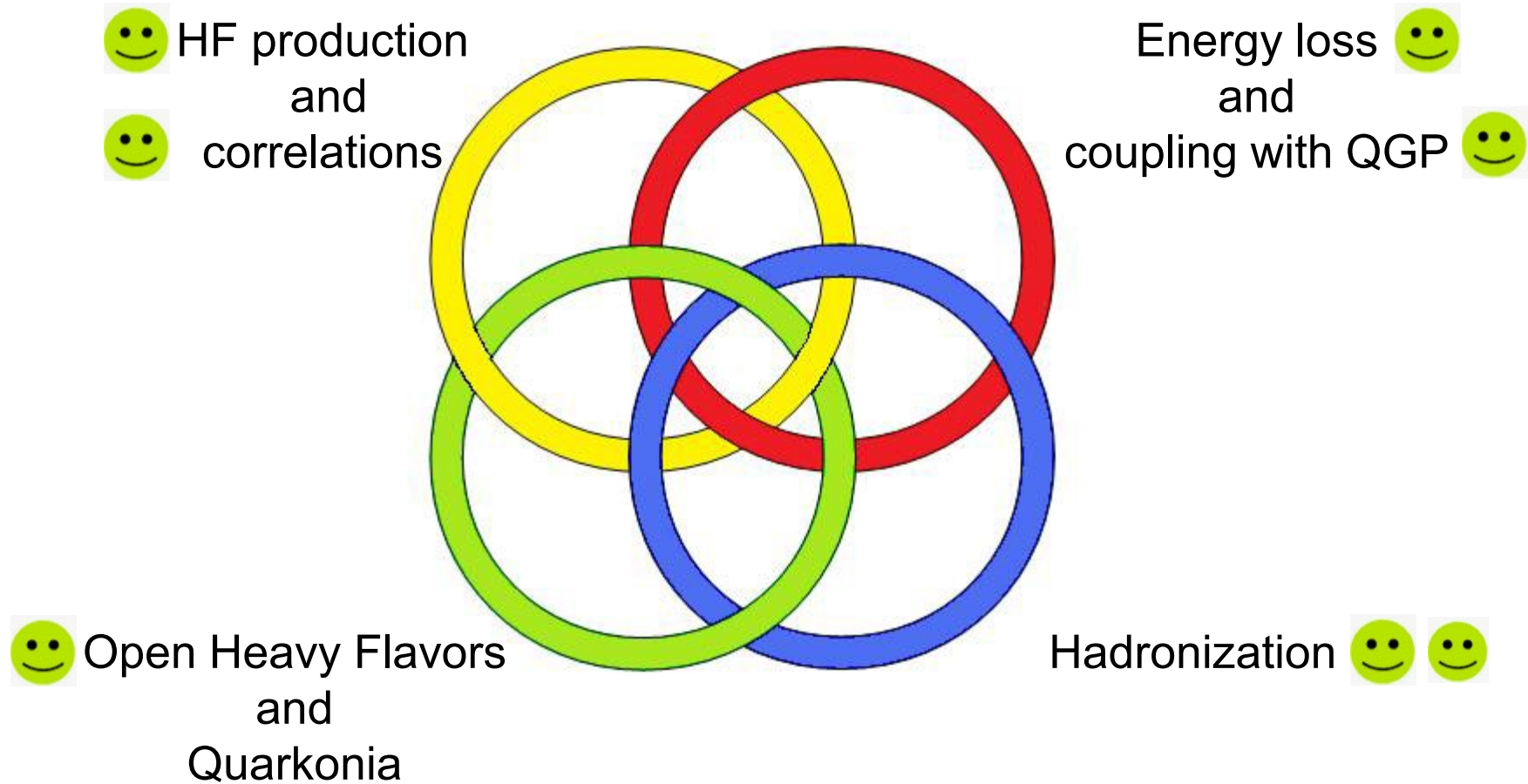
Recap : HF in pp

| observable | HQ energy loss in QGP | Coalescence in the presence of a QGP droplet |
|------------------------|-----------------------|--|
| Hadron pt spectra | Little effect | LARGE effect |
| Hadron yield ratios | Little effect | LARGE effect |
| v_2 | LARGE effect | Little effect |
| Azimuthal correlations | Little effect | Little effect |

According to EPOS4-HQ: Everything consistent with the production of a short-lived QGP in most active pp collisions at LHC

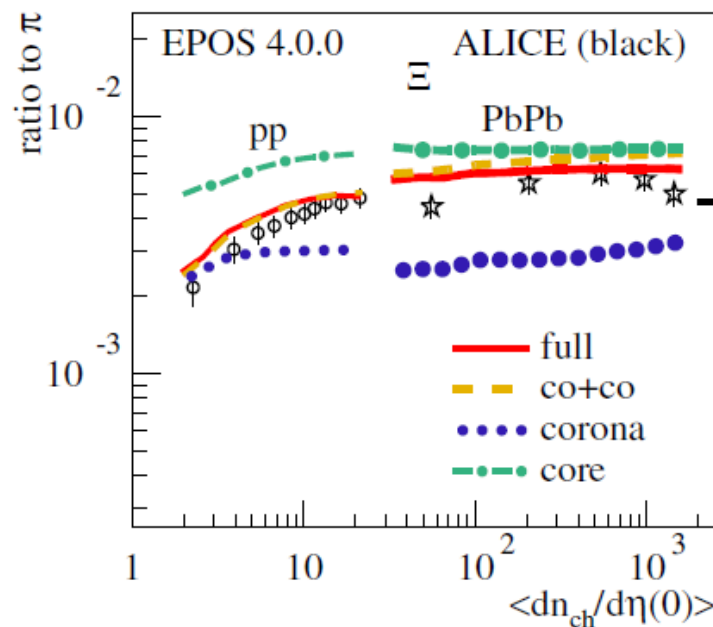
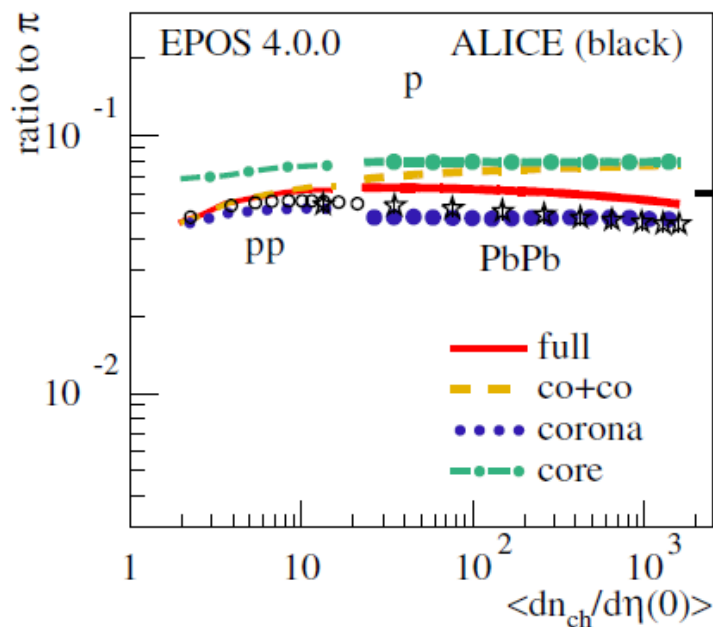
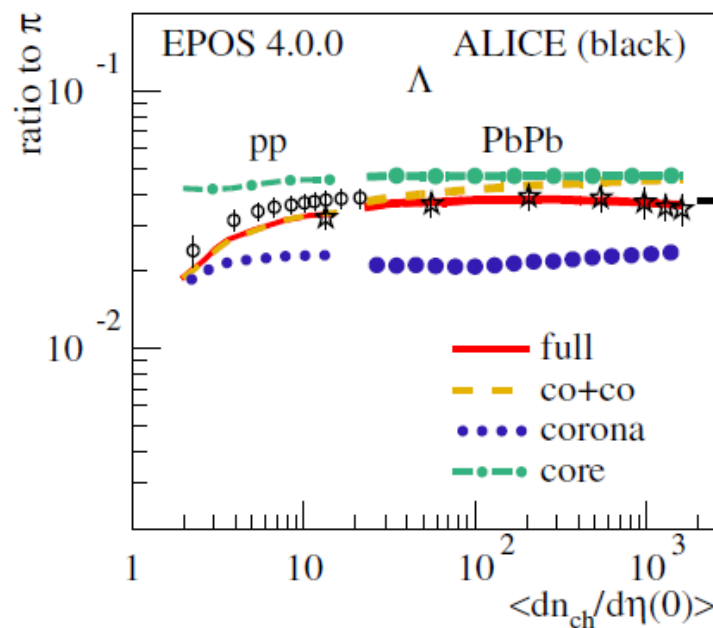
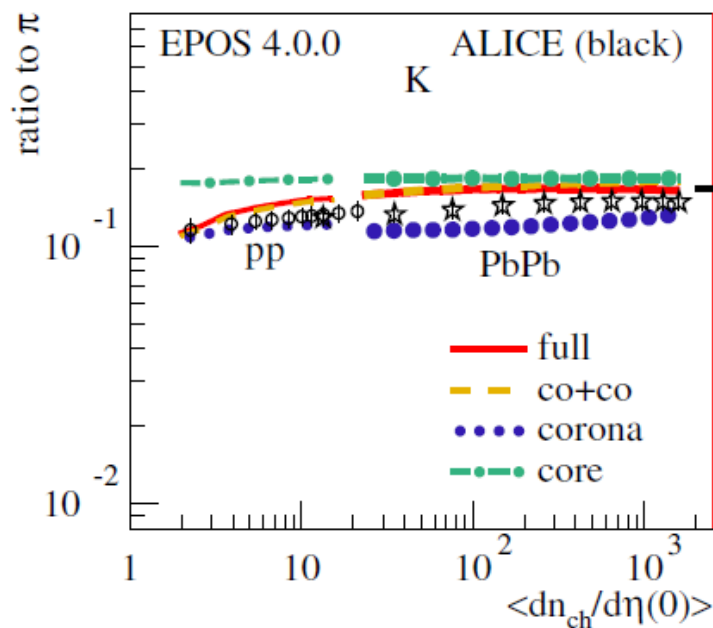
To come : adding the quarkonia component...

HF Borromeo's rings



Crucial to consider all facets of the problem, as in the new EPOS4-HQ framework

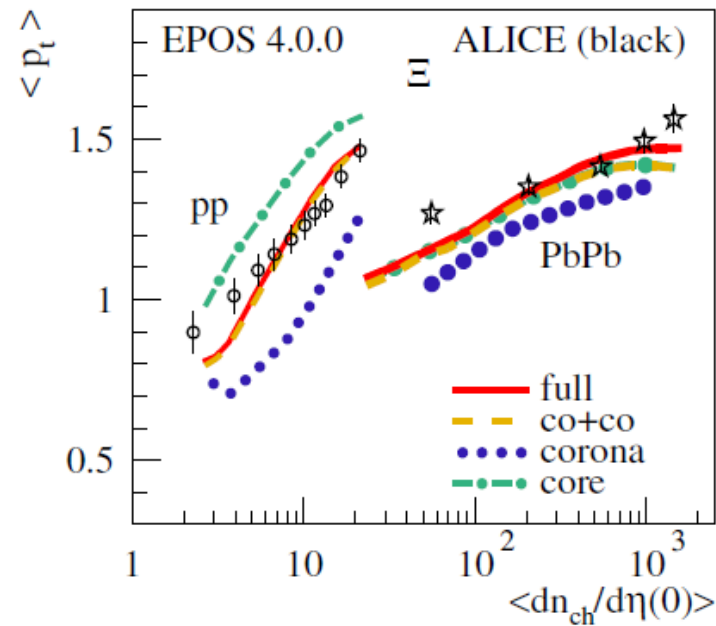
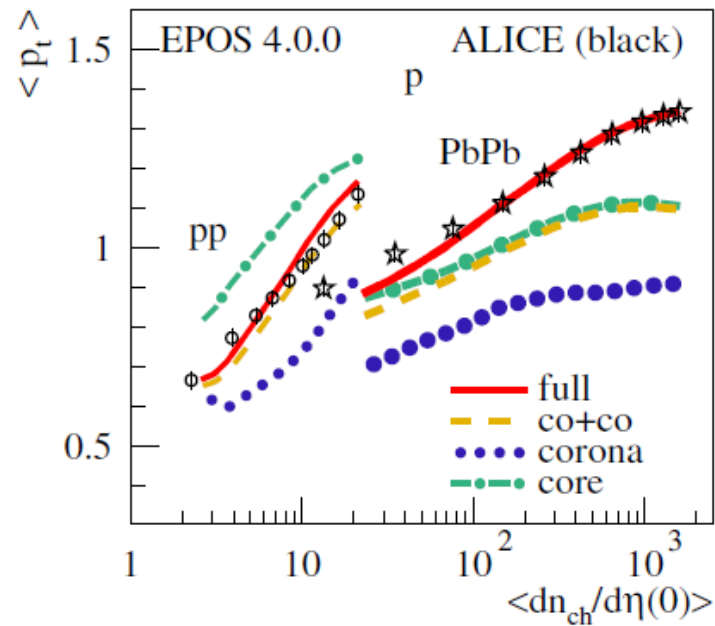
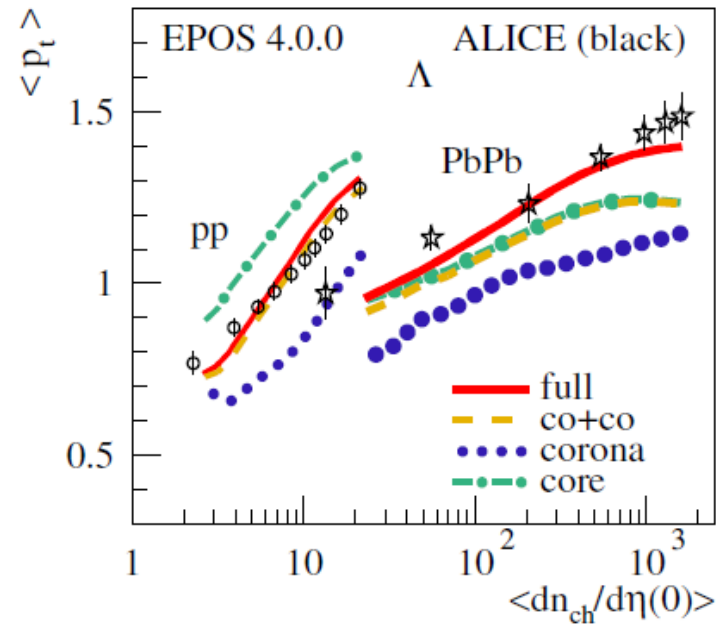
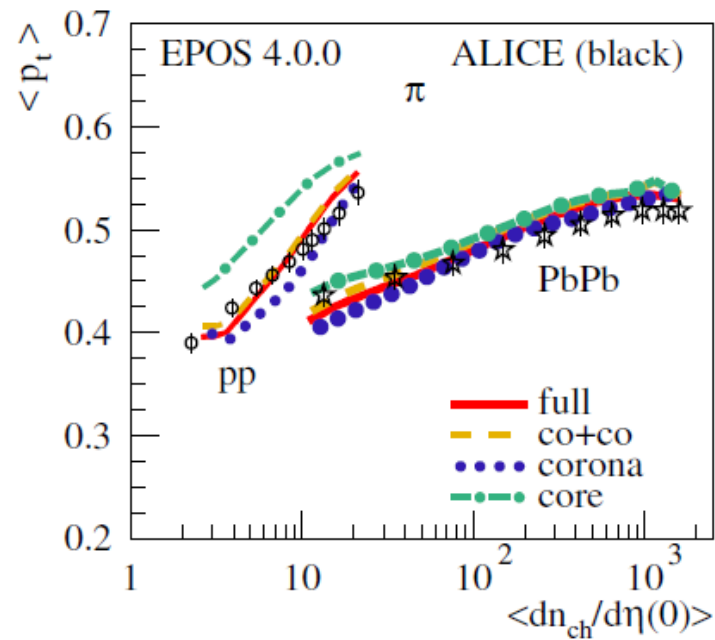
Back up



continuous curves

Affected by:

- core-corona
- microcanonica
- hadronic cascade



discontinuities

curves
affected by:

saturation

flow

core-corona

hadronic
cascade

Charmed/bottom mesons

D 17 states

| Charmed Mesons ($C = + -1$) |
|-------------------------------|
| $D^+ -$ |
| D^0 |
| $D^*(2007)^0$ |
| $D^*(2010)^+ -$ |
| $D^*(0)(2400)^0$ |
| $D^*(0)(2400)^+ -$ |
| $D(1)(2420)^0$ |
| $D(1)(2420)^+ -$ |
| $D(1)(2430)^0$ |
| $D^*(2)(2460)^0$ |
| $D^*(2)(2460)^+ -$ |
| $D(2550)^0$ |
| $D^*(J)(2600)$ |
| $D^*(2640)^+ -$ |
| $D(2740)^0$ |
| $D(2750)$ |
| $D(3000)^0$ |

D_s 10 states

| Charmed, Strange Mesons ($C = S = +-1$) |
|---|
| $D(s)^+-$ |
| $D^*(s)^+-$ |
| $D^*(s_0)(2317)^+-$ |
| $D(s_1)(2460)^+-$ |
| $D(s_1)(2536)^+-$ |
| $D(s_2)(2573)^+-$ |
| $D^*(s_1)(2700)^{+-}$ |
| $D^*(s_1)(2860)^+-$ |
| $D^*(s_3)(2860)^+-$ |
| $D(s_J)(3040)^+-$ |

Charmed/bottom baryons

 Σ_c

54 states

38 states

 Λ_c

TABLE II: Masses of the Λ_Q ($Q = c, b$) heavy baryons (in MeV).

| $I(J^P)$ | Qd state | Q = c | | Q = b | |
|--------------------|----------|-------|---------------------------------------|-------|----------------------|
| | | M | M^{exp} [1] | M | M^{exp} [1] |
| $0(\frac{1}{2}^+)$ | 1S | 2286 | 2286.46(14) | 5620 | 5620.2(1.6) |
| $0(\frac{3}{2}^+)$ | 2S | 2769 | 2766.6(2.4)? | 6089 | |
| $0(\frac{1}{2}^+)$ | 3S | 3130 | | 6455 | |
| $0(\frac{3}{2}^+)$ | 4S | 3437 | | 6756 | |
| $0(\frac{1}{2}^+)$ | 5S | 3715 | | 7015 | |
| $0(\frac{3}{2}^+)$ | 6S | 3973 | | 7256 | |
| $0(\frac{1}{2}^-)$ | 1P | 2598 | 2595.4(6) | 5930 | |
| $0(\frac{3}{2}^-)$ | 2P | 2983 | 2939.3($\frac{1}{3}, \frac{4}{3}$)? | 6326 | |
| $0(\frac{1}{2}^-)$ | 3P | 3303 | | 6645 | |
| $0(\frac{3}{2}^-)$ | 4P | 3588 | | 6917 | |
| $0(\frac{1}{2}^-)$ | 5P | 3852 | | 7157 | |
| $0(\frac{3}{2}^-)$ | 1P | 2627 | 2628.1(6) | 5942 | |
| $0(\frac{1}{2}^-)$ | 2P | 3005 | | 6333 | |
| $0(\frac{3}{2}^-)$ | 3P | 3322 | | 6651 | |
| $0(\frac{1}{2}^-)$ | 4P | 3606 | | 6922 | |
| $0(\frac{3}{2}^-)$ | 5P | 3869 | | 7171 | |
| $0(\frac{1}{2}^+)$ | 1D | 2874 | | 6190 | |
| $0(\frac{3}{2}^+)$ | 2D | 3189 | | 6526 | |
| $0(\frac{1}{2}^+)$ | 3D | 3480 | | 6811 | |
| $0(\frac{3}{2}^+)$ | 4D | 3747 | | 7060 | |
| $0(\frac{1}{2}^+)$ | 1D | 2880 | 2881.53(35) | 6196 | |
| $0(\frac{3}{2}^+)$ | 2D | 3209 | | 6531 | |
| $0(\frac{1}{2}^+)$ | 3D | 3500 | | 6814 | |
| $0(\frac{3}{2}^+)$ | 4D | 3767 | | 7063 | |
| $0(\frac{1}{2}^-)$ | 1F | 3097 | | 6408 | |
| $0(\frac{3}{2}^-)$ | 2F | 3375 | | 6705 | |
| $0(\frac{1}{2}^-)$ | 3F | 3646 | | 6964 | |
| $0(\frac{3}{2}^-)$ | 4F | 3900 | | 7196 | |
| $0(\frac{1}{2}^-)$ | 1F | 3078 | | 6411 | |
| $0(\frac{3}{2}^-)$ | 2F | 3393 | | 6708 | |
| $0(\frac{1}{2}^-)$ | 3F | 3667 | | 6966 | |
| $0(\frac{3}{2}^-)$ | 4F | 3922 | | 7197 | |
| $0(\frac{1}{2}^+)$ | 1G | 3270 | | 6598 | |
| $0(\frac{3}{2}^+)$ | 2G | 3546 | | 6867 | |
| $0(\frac{1}{2}^+)$ | 1G | 3284 | | 6509 | |
| $0(\frac{3}{2}^+)$ | 2G | 3564 | | 6868 | |
| $0(\frac{1}{2}^-)$ | 1H | 3444 | | 6767 | |
| $0(\frac{3}{2}^-)$ | 1H | 3460 | | 6766 | |

TABLE III: Masses of the Σ_Q ($Q = c, b$) heavy baryons (in MeV).

| $I(J^P)$ | Qd state | Q = c | | Q = b | |
|--------------------|----------|-------|---------------------------------------|-------|----------------------|
| | | M | M^{exp} [1] | M | M^{exp} [1] |
| $1(\frac{1}{2}^+)$ | 1S | 2443 | 2453.76(18) | 5808 | 5807.8(2.7) |
| $1(\frac{3}{2}^+)$ | 2S | 2901 | | 6213 | |
| $1(\frac{1}{2}^+)$ | 3S | 3271 | | 6575 | |
| $1(\frac{3}{2}^+)$ | 4S | 3581 | | 6869 | |
| $1(\frac{1}{2}^+)$ | 5S | 3861 | | 7124 | |
| $1(\frac{3}{2}^+)$ | 1S | 2519 | 2518.0(5) | 5834 | 5829.0(3.4) |
| $1(\frac{1}{2}^+)$ | 2S | 2936 | 2939.3($\frac{1}{3}, \frac{4}{3}$)? | 6226 | |
| $1(\frac{3}{2}^+)$ | 3S | 3293 | | 6583 | |
| $1(\frac{1}{2}^+)$ | 4S | 3598 | | 6876 | |
| $1(\frac{3}{2}^+)$ | 5S | 3873 | | 7129 | |
| $1(\frac{1}{2}^-)$ | 1P | 2799 | 2802($\frac{1}{3}$) | 6101 | |
| $1(\frac{3}{2}^-)$ | 2P | 3172 | | 6440 | |
| $1(\frac{1}{2}^-)$ | 3P | 3488 | | 6756 | |
| $1(\frac{3}{2}^-)$ | 4P | 3770 | | 7024 | |
| $1(\frac{1}{2}^-)$ | 1P | 2713 | | 6095 | |
| $1(\frac{3}{2}^-)$ | 2P | 3125 | | 6430 | |
| $1(\frac{1}{2}^-)$ | 3P | 3455 | | 6742 | |
| $1(\frac{3}{2}^-)$ | 4P | 3743 | | 7008 | |
| $1(\frac{1}{2}^-)$ | 1P | 2798 | 2802($\frac{1}{3}$) | 6096 | |
| $1(\frac{3}{2}^-)$ | 2P | 3172 | | 6430 | |
| $1(\frac{1}{2}^-)$ | 3P | 3486 | | 6742 | |
| $1(\frac{3}{2}^-)$ | 4P | 3768 | | 7009 | |
| $1(\frac{1}{2}^-)$ | 1P | 2773 | 2766.6(2.4)? | 6087 | |
| $1(\frac{3}{2}^-)$ | 2P | 3151 | | 6423 | |
| $1(\frac{1}{2}^-)$ | 3P | 3469 | | 6736 | |
| $1(\frac{3}{2}^-)$ | 4P | 3753 | | 7003 | |
| $1(\frac{1}{2}^-)$ | 1P | 2789 | | 6084 | |
| $1(\frac{3}{2}^-)$ | 2P | 3161 | | 6421 | |
| $1(\frac{1}{2}^-)$ | 3P | 3475 | | 6732 | |
| $1(\frac{3}{2}^-)$ | 4P | 3757 | | 6999 | |
| $1(\frac{1}{2}^+)$ | 1D | 3041 | | 6311 | |
| $1(\frac{3}{2}^+)$ | 2D | 3370 | | 6636 | |
| $1(\frac{1}{2}^+)$ | 1D | 3043 | | 6326 | |
| $1(\frac{3}{2}^+)$ | 2D | 3366 | | 6647 | |
| $1(\frac{1}{2}^+)$ | 1D | 3040 | | 6285 | |
| $1(\frac{3}{2}^+)$ | 2D | 3364 | | 6612 | |
| $1(\frac{1}{2}^+)$ | 1D | 3038 | | 6284 | |
| $1(\frac{3}{2}^+)$ | 2D | 3365 | | 6612 | |
| $1(\frac{1}{2}^+)$ | 1D | 3023 | | 6270 | |
| $1(\frac{3}{2}^+)$ | 2D | 3349 | | 6598 | |
| $1(\frac{1}{2}^+)$ | 1D | 3013 | | 6260 | |
| $1(\frac{3}{2}^+)$ | 2D | 3342 | | 6590 | |
| $1(\frac{1}{2}^-)$ | 1F | 3288 | | 6550 | |
| $1(\frac{3}{2}^-)$ | 1F | 3283 | | 6564 | |
| $1(\frac{1}{2}^-)$ | 1F | 3254 | | 6501 | |
| $1(\frac{3}{2}^-)$ | 1F | 3253 | | 6500 | |
| $1(\frac{1}{2}^-)$ | 1F | 3227 | | 6472 | |
| $1(\frac{3}{2}^-)$ | 1F | 3209 | | 6459 | |
| $1(\frac{1}{2}^+)$ | 1G | 3495 | | 6749 | |
| $1(\frac{3}{2}^+)$ | 1G | 3483 | | 6761 | |
| $1(\frac{1}{2}^+)$ | 1G | 3444 | | 6688 | |
| $1(\frac{3}{2}^+)$ | 1G | 3442 | | 6687 | |
| $1(\frac{1}{2}^+)$ | 1G | 3410 | | 6648 | |
| $1(\frac{3}{2}^+)$ | 1G | 3386 | | 6635 | |

Charmed/bottom baryons

Ξ_c

54 states

38 states

Ξ_c

TABLE IV: Masses of the Ξ_Q ($Q = c, b$) heavy baryons with the scalar diquark (in MeV).

| $I(J^P)$ | Qd state | $Q = c$ | | $Q = b$ | |
|-------------------------------|------------|---------|--------------------------|---------|----------------------|
| | | M | $M^{\text{exp}} [1]$ | M | $M^{\text{exp}} [1]$ |
| $\frac{1}{2}(\frac{1}{2}^+)$ | 1S | 2476 | 2470.88($\frac{3}{2}$) | 5803 | 5790.5(2.7) |
| $\frac{1}{2}(\frac{3}{2}^+)$ | 2S | 2959 | | 6266 | |
| $\frac{1}{2}(\frac{5}{2}^+)$ | 3S | 3323 | | 6601 | |
| $\frac{1}{2}(\frac{7}{2}^+)$ | 4S | 3632 | | 6913 | |
| $\frac{1}{2}(\frac{9}{2}^+)$ | 5S | 3909 | | 7165 | |
| $\frac{1}{2}(\frac{11}{2}^+)$ | 6S | 4166 | | 7415 | |
| $\frac{1}{2}(\frac{1}{2}^-)$ | 1P | 2792 | 2791.8(3.3) | 6120 | |
| $\frac{1}{2}(\frac{3}{2}^-)$ | 2P | 3179 | | 6496 | |
| $\frac{1}{2}(\frac{5}{2}^-)$ | 3P | 3500 | | 6805 | |
| $\frac{1}{2}(\frac{7}{2}^-)$ | 4P | 3785 | | 7068 | |
| $\frac{1}{2}(\frac{9}{2}^-)$ | 5P | 4048 | | 7302 | |
| $\frac{1}{2}(\frac{11}{2}^-)$ | 1P | 2819 | 2819.6(1.2) | 6130 | |
| $\frac{1}{2}(\frac{3}{2}^-)$ | 2P | 3201 | | 6502 | |
| $\frac{1}{2}(\frac{5}{2}^-)$ | 3P | 3519 | | 6810 | |
| $\frac{1}{2}(\frac{7}{2}^-)$ | 4P | 3804 | | 7073 | |
| $\frac{1}{2}(\frac{9}{2}^-)$ | 5P | 4066 | | 7306 | |
| $\frac{1}{2}(\frac{11}{2}^-)$ | 1D | 3059 | 3054.2(1.3) | 6366 | |
| $\frac{1}{2}(\frac{13}{2}^-)$ | 2D | 3388 | | 6690 | |
| $\frac{1}{2}(\frac{15}{2}^-)$ | 3D | 3678 | | 6966 | |
| $\frac{1}{2}(\frac{17}{2}^-)$ | 4D | 3945 | | 7208 | |
| $\frac{1}{2}(\frac{19}{2}^-)$ | 1D | 3076 | 3079.9(1.4) | 6373 | |
| $\frac{1}{2}(\frac{21}{2}^-)$ | 2D | 3407 | | 6696 | |
| $\frac{1}{2}(\frac{23}{2}^-)$ | 3D | 3699 | | 6970 | |
| $\frac{1}{2}(\frac{25}{2}^-)$ | 4D | 3965 | | 7212 | |
| $\frac{1}{2}(\frac{1}{2}^-)$ | 1F | 3278 | | 6577 | |
| $\frac{1}{2}(\frac{3}{2}^-)$ | 2F | 3575 | | 6863 | |
| $\frac{1}{2}(\frac{5}{2}^-)$ | 3F | 3845 | | 7114 | |
| $\frac{1}{2}(\frac{7}{2}^-)$ | 4F | 4098 | | 7339 | |
| $\frac{1}{2}(\frac{9}{2}^-)$ | 1F | 3292 | | 6581 | |
| $\frac{1}{2}(\frac{11}{2}^-)$ | 2F | 3592 | | 6867 | |
| $\frac{1}{2}(\frac{13}{2}^-)$ | 3F | 3865 | | 7117 | |
| $\frac{1}{2}(\frac{15}{2}^-)$ | 4F | 4120 | | 7342 | |
| $\frac{1}{2}(\frac{1}{2}^-)$ | 1G | 3469 | | 6760 | |
| $\frac{1}{2}(\frac{3}{2}^-)$ | 2G | 3745 | | 7020 | |
| $\frac{1}{2}(\frac{5}{2}^-)$ | 1G | 3483 | | 6762 | |
| $\frac{1}{2}(\frac{7}{2}^-)$ | 2G | 3763 | | 7032 | |
| $\frac{1}{2}(\frac{9}{2}^-)$ | 1H | 3643 | | 6933 | |
| $\frac{1}{2}(\frac{11}{2}^-)$ | 1H | 3658 | | 6934 | |

TABLE V: Masses of the Ξ_Q ($Q = c, b$) heavy baryons with the axial vector diquark (in MeV).

| $I(J^P)$ | Qd state | M | $Q = c$ | | $Q = b$ | |
|-------------------------------|------------|------|----------------------|------|----------------------|-----|
| | | | $M^{\text{exp}} [1]$ | M | $M^{\text{exp}} [1]$ | M |
| $\frac{1}{2}(\frac{1}{2}^+)$ | 1S | 2579 | 2577.9(2.9) | 5936 | | |
| $\frac{1}{2}(\frac{3}{2}^+)$ | 2S | 2983 | 2971.4(3.3) | 6329 | | |
| $\frac{1}{2}(\frac{5}{2}^+)$ | 3S | 3377 | | 6687 | | |
| $\frac{1}{2}(\frac{7}{2}^+)$ | 4S | 3695 | | 6978 | | |
| $\frac{1}{2}(\frac{9}{2}^+)$ | 5S | 3978 | | 7229 | | |
| $\frac{1}{2}(\frac{11}{2}^+)$ | 1S | 2649 | 2645.9(0.5) | 5963 | | |
| $\frac{1}{2}(\frac{3}{2}^+)$ | 2S | 3026 | | 6342 | | |
| $\frac{1}{2}(\frac{5}{2}^+)$ | 3S | 3396 | | 6695 | | |
| $\frac{1}{2}(\frac{7}{2}^+)$ | 4S | 3709 | | 6984 | | |
| $\frac{1}{2}(\frac{9}{2}^+)$ | 5S | 3989 | | 7234 | | |
| $\frac{1}{2}(\frac{11}{2}^+)$ | 1P | 2936 | 2931(6) | 6233 | | |
| $\frac{1}{2}(\frac{13}{2}^+)$ | 2P | 3313 | | 6611 | | |
| $\frac{1}{2}(\frac{15}{2}^+)$ | 3P | 3630 | | 6915 | | |
| $\frac{1}{2}(\frac{17}{2}^+)$ | 4P | 3912 | | 7174 | | |
| $\frac{1}{2}(\frac{19}{2}^+)$ | 1P | 2854 | | 6227 | | |
| $\frac{1}{2}(\frac{21}{2}^+)$ | 2P | 3267 | | 6604 | | |
| $\frac{1}{2}(\frac{23}{2}^+)$ | 3P | 3598 | | 6906 | | |
| $\frac{1}{2}(\frac{25}{2}^+)$ | 4P | 3887 | | 7164 | | |
| $\frac{1}{2}(\frac{1}{2}^-)$ | 1P | 2935 | 2931(6) | 6234 | | |
| $\frac{1}{2}(\frac{3}{2}^-)$ | 2P | 3311 | | 6605 | | |
| $\frac{1}{2}(\frac{5}{2}^-)$ | 3P | 3628 | | 6905 | | |
| $\frac{1}{2}(\frac{7}{2}^-)$ | 4P | 3911 | | 7163 | | |
| $\frac{1}{2}(\frac{9}{2}^-)$ | 1P | 2912 | | 6224 | | |
| $\frac{1}{2}(\frac{11}{2}^-)$ | 2P | 3293 | | 6598 | | |
| $\frac{1}{2}(\frac{13}{2}^-)$ | 3P | 3613 | | 6900 | | |
| $\frac{1}{2}(\frac{15}{2}^-)$ | 4P | 3898 | | 7159 | | |
| $\frac{1}{2}(\frac{17}{2}^-)$ | 1P | 2929 | 2931(6) | 6226 | | |
| $\frac{1}{2}(\frac{19}{2}^-)$ | 2P | 3303 | | 6596 | | |
| $\frac{1}{2}(\frac{21}{2}^-)$ | 3P | 3619 | | 6897 | | |
| $\frac{1}{2}(\frac{23}{2}^-)$ | 4P | 3902 | | 7156 | | |
| $\frac{1}{2}(\frac{25}{2}^-)$ | 1D | 3163 | | 6447 | | |
| $\frac{1}{2}(\frac{27}{2}^-)$ | 2D | 3505 | | 6767 | | |
| $\frac{1}{2}(\frac{29}{2}^-)$ | 1D | 3167 | | 6459 | | |
| $\frac{1}{2}(\frac{31}{2}^-)$ | 2D | 3506 | | 6775 | | |
| $\frac{1}{2}(\frac{33}{2}^-)$ | 1D | 3160 | | 6431 | | |
| $\frac{1}{2}(\frac{35}{2}^-)$ | 2D | 3497 | | 6751 | | |
| $\frac{1}{2}(\frac{37}{2}^-)$ | 1D | 3166 | | 6432 | | |
| $\frac{1}{2}(\frac{39}{2}^-)$ | 2D | 3504 | | 6751 | | |
| $\frac{1}{2}(\frac{41}{2}^-)$ | 1D | 3153 | | 6420 | | |
| $\frac{1}{2}(\frac{43}{2}^-)$ | 2D | 3493 | | 6740 | | |
| $\frac{1}{2}(\frac{1}{2}^-)$ | 1D | 3147 | 3122.9(1.3) | 6414 | | |
| $\frac{1}{2}(\frac{3}{2}^-)$ | 2D | 3486 | | 6736 | | |
| $\frac{1}{2}(\frac{5}{2}^-)$ | 1F | 3418 | | 6675 | | |
| $\frac{1}{2}(\frac{7}{2}^-)$ | 1F | 3408 | | 6686 | | |
| $\frac{1}{2}(\frac{9}{2}^-)$ | 1F | 3394 | | 6640 | | |
| $\frac{1}{2}(\frac{11}{2}^-)$ | 1F | 3393 | | 6641 | | |
| $\frac{1}{2}(\frac{13}{2}^-)$ | 1F | 3373 | | 6619 | | |
| $\frac{1}{2}(\frac{15}{2}^-)$ | 1F | 3357 | | 6610 | | |
| $\frac{1}{2}(\frac{17}{2}^-)$ | 1G | 3623 | | 6867 | | |
| $\frac{1}{2}(\frac{19}{2}^-)$ | 1G | 3608 | | 6876 | | |
| $\frac{1}{2}(\frac{21}{2}^-)$ | 1G | 3584 | | 6822 | | |
| $\frac{1}{2}(\frac{23}{2}^-)$ | 1G | 3582 | | 6821 | | |
| $\frac{1}{2}(\frac{25}{2}^-)$ | 1G | 3558 | | 6792 | | |
| $\frac{1}{2}(\frac{27}{2}^-)$ | 1G | 3536 | | 6782 | | |

Charmed/bottom baryons

54 states

Ω_c

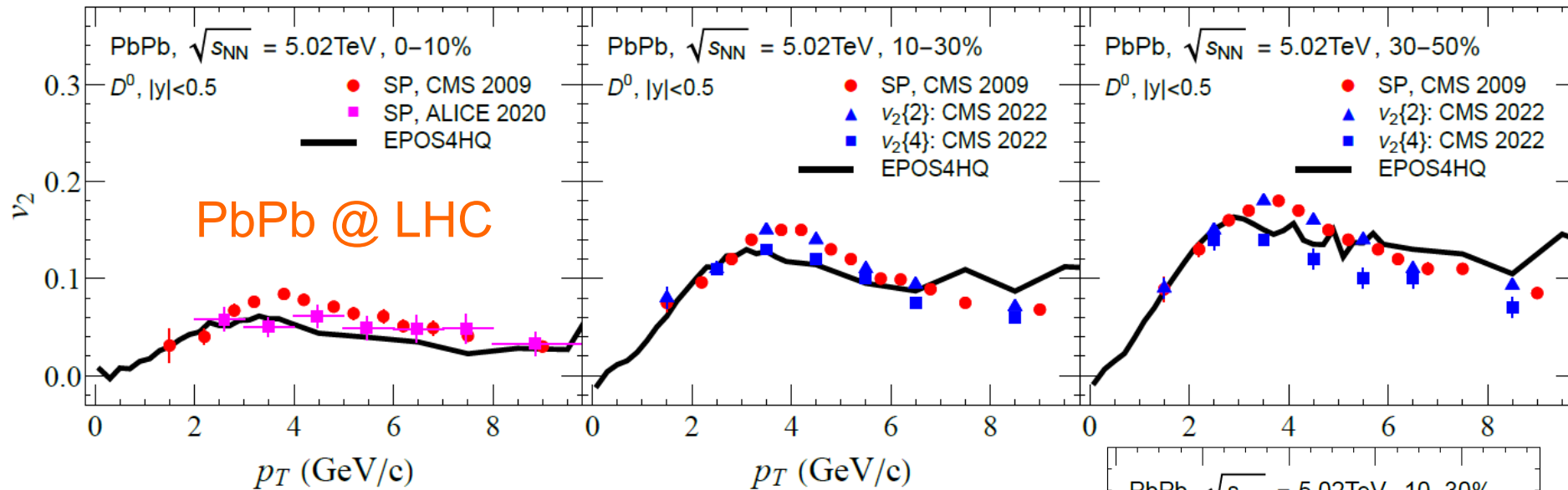
TABLE VI: Masses of the Ω_Q ($Q = c, b$) heavy baryons (in MeV).

| $I(J^P)$ | Qd state | $Q = c$ | | $Q = b$ | |
|--------------------|------------|---------|----------------------|---------|----------------------|
| | | M | M^{exp} [1] | M | M^{exp} [1] |
| $0(\frac{1}{2}^+)$ | 1S | 2698 | 2695.2(1.7) | 6064 | 6071(40) |
| $0(\frac{1}{2}^+)$ | 2S | 3088 | | 6450 | |
| $0(\frac{1}{2}^+)$ | 3S | 3489 | | 6804 | |
| $0(\frac{1}{2}^+)$ | 4S | 3814 | | 7091 | |
| $0(\frac{1}{2}^+)$ | 5S | 4102 | | 7338 | |
| $0(\frac{1}{2}^+)$ | 1S | 2768 | 2765.9(2.0) | 6088 | |
| $0(\frac{1}{2}^+)$ | 2S | 3123 | | 6461 | |
| $0(\frac{1}{2}^+)$ | 3S | 3510 | | 6811 | |
| $0(\frac{1}{2}^+)$ | 4S | 3830 | | 7096 | |
| $0(\frac{1}{2}^+)$ | 5S | 4114 | | 7343 | |
| $0(\frac{1}{2}^-)$ | 1P | 3055 | | 6339 | |
| $0(\frac{1}{2}^-)$ | 2P | 3435 | | 6710 | |
| $0(\frac{1}{2}^-)$ | 3P | 3754 | | 7009 | |
| $0(\frac{1}{2}^-)$ | 4P | 4037 | | 7265 | |
| $0(\frac{1}{2}^-)$ | 1P | 2966 | | 6330 | |
| $0(\frac{1}{2}^-)$ | 2P | 3384 | | 6706 | |
| $0(\frac{1}{2}^-)$ | 3P | 3717 | | 7003 | |
| $0(\frac{1}{2}^-)$ | 2P | 4009 | | 7257 | |
| $0(\frac{1}{2}^-)$ | 1P | 3054 | | 6340 | |
| $0(\frac{1}{2}^-)$ | 2P | 3433 | | 6705 | |
| $0(\frac{1}{2}^-)$ | 3P | 3752 | | 7002 | |

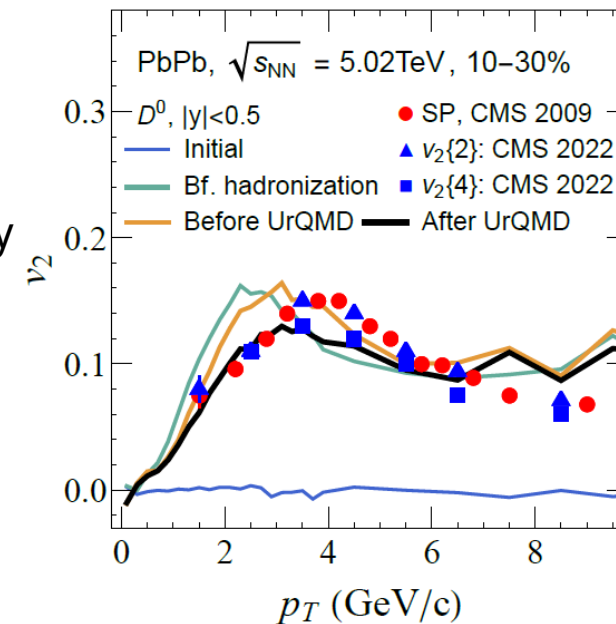
TABLE VI: (continued)

| $I(J^P)$ | Qd state | $Q = c$ | | $Q = b$ | |
|--------------------|------------|---------|----------------------|---------|----------------------|
| | | M | M^{exp} [1] | M | M^{exp} [1] |
| $0(\frac{3}{2}^-)$ | 4P | 4036 | | 7258 | |
| $0(\frac{3}{2}^-)$ | 1P | 3029 | | 6331 | |
| $0(\frac{3}{2}^-)$ | 2P | 3415 | | 6699 | |
| $0(\frac{3}{2}^-)$ | 3P | 3737 | | 6998 | |
| $0(\frac{3}{2}^-)$ | 4P | 4023 | | 7250 | |
| $0(\frac{3}{2}^-)$ | 1P | 3051 | | 6334 | |
| $0(\frac{3}{2}^-)$ | 2P | 3427 | | 6700 | |
| $0(\frac{3}{2}^-)$ | 3P | 3744 | | 6996 | |
| $0(\frac{3}{2}^-)$ | 4P | 4028 | | 7251 | |
| $0(\frac{1}{2}^+)$ | 1D | 3287 | | 6540 | |
| $0(\frac{1}{2}^+)$ | 2D | 3623 | | 6857 | |
| $0(\frac{1}{2}^+)$ | 1D | 3298 | | 6549 | |
| $0(\frac{1}{2}^+)$ | 2D | 3627 | | 6863 | |
| $0(\frac{1}{2}^+)$ | 1D | 3282 | | 6530 | |
| $0(\frac{1}{2}^+)$ | 2D | 3613 | | 6846 | |
| $0(\frac{1}{2}^+)$ | 1D | 3297 | | 6529 | |
| $0(\frac{1}{2}^+)$ | 2D | 3626 | | 6846 | |
| $0(\frac{1}{2}^+)$ | 1D | 3286 | | 6520 | |
| $0(\frac{1}{2}^+)$ | 2D | 3614 | | 6837 | |
| $0(\frac{1}{2}^+)$ | 1D | 3283 | | 6517 | |
| $0(\frac{1}{2}^+)$ | 2D | 3611 | | 6834 | |
| $0(\frac{1}{2}^-)$ | 1F | 3533 | | 6763 | |
| $0(\frac{1}{2}^-)$ | 1F | 3522 | | 6771 | |
| $0(\frac{1}{2}^-)$ | 1F | 3515 | | 6737 | |
| $0(\frac{1}{2}^-)$ | 1F | 3514 | | 6736 | |
| $0(\frac{1}{2}^-)$ | 1F | 3498 | | 6719 | |
| $0(\frac{1}{2}^-)$ | 1F | 3485 | | 6713 | |
| $0(\frac{1}{2}^+)$ | 1G | 3739 | | 6952 | |
| $0(\frac{1}{2}^+)$ | 1G | 3721 | | 6959 | |
| $0(\frac{1}{2}^+)$ | 1G | 3707 | | 6916 | |
| $0(\frac{1}{2}^+)$ | 1G | 3705 | | 6915 | |
| $0(\frac{1}{2}^+)$ | 1G | 3685 | | 6892 | |
| $0(\frac{1}{2}^+)$ | 1G | 3665 | | 6884 | |

Elliptic flow in EPOS4-HQ



- Good agreement with data, maybe some excess at large p_T (one should investigate path length anisotropy in the bulk)
- Most of the v_2 generated by the interaction of the c-quarks with the QGP

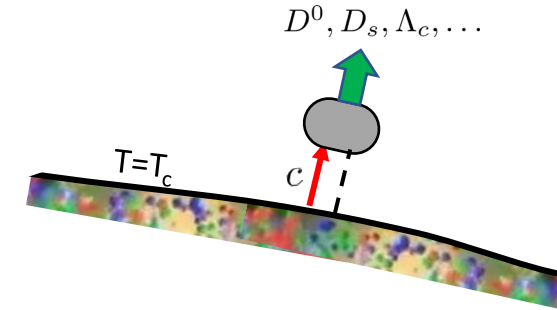
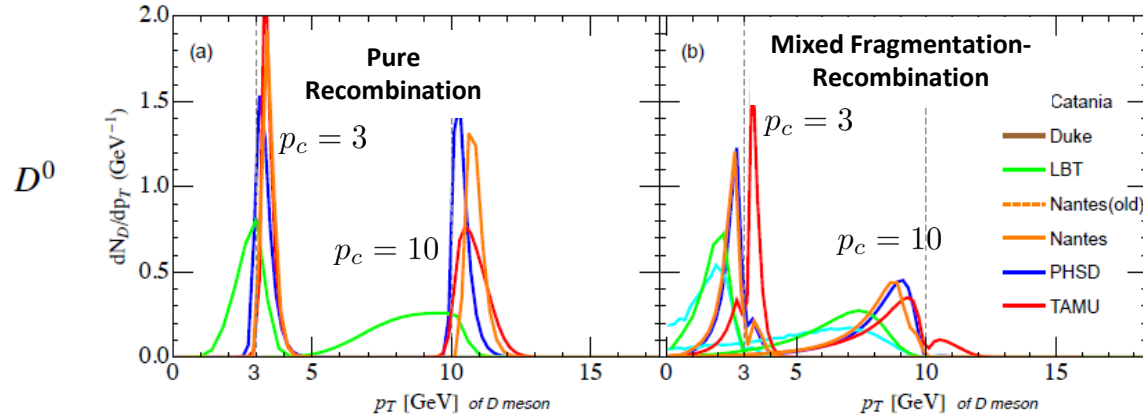


Hadronization of heavy quarks

- Recent effort of theorists to compare their hadronization schemes at the end of the QGP

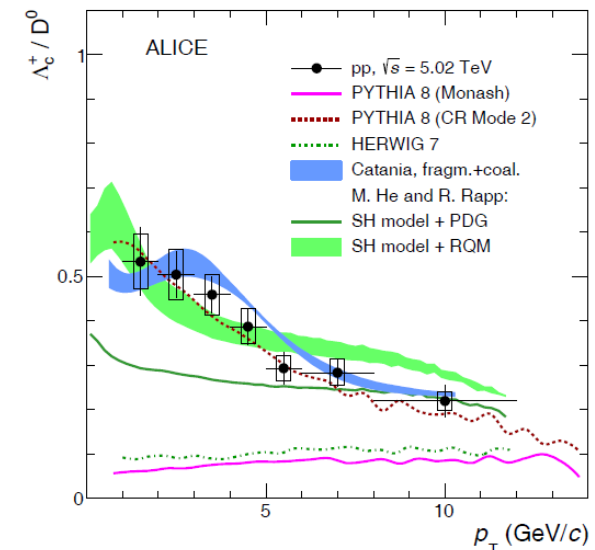
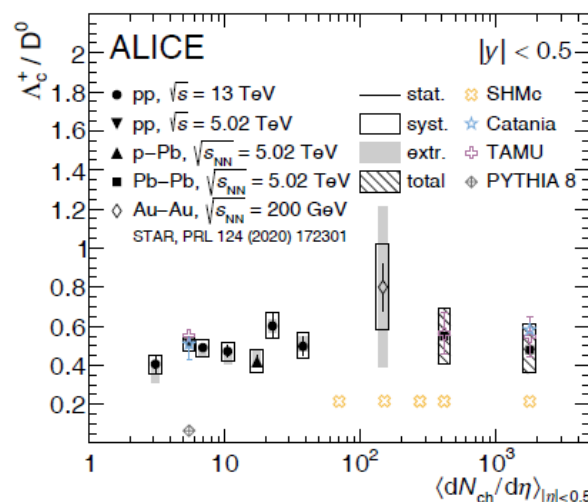
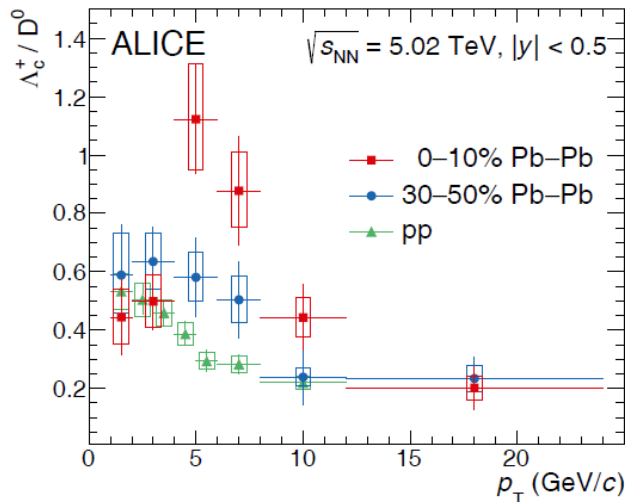
dN_D/dp_T of the direct D^0 meson produced by a c -quark with $p_T = 3$ GeV and 10 GeV

Jiaxing Zhao et al., 2311.10621



- Diversity => things to learn ! ... Hadrochemistry of Heavy Flavor will be a major subject of investigation for ALICE 3.

- But also in small systems like pp (many signs of **collectivity in small systems => QGP ?**)



2023: Bundle 4: EPOS4-HQ

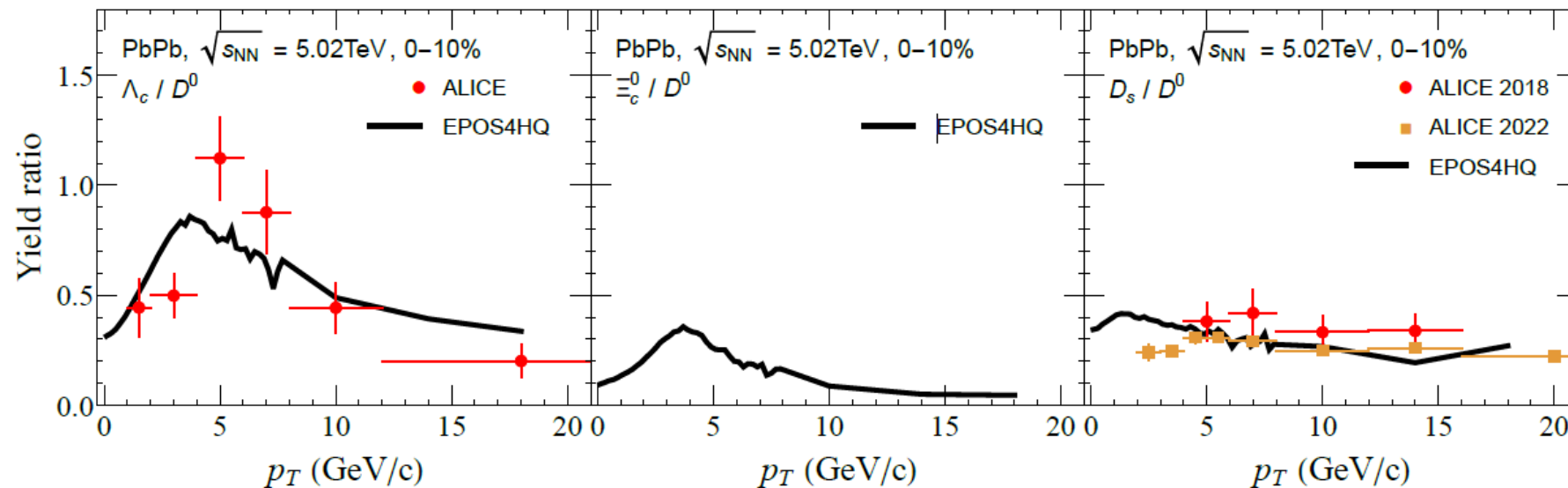
| Ingredient | B1 | B2 (MC@ _s HQ+EPOS2) | B4 (EPOS4-HQ) |
|-------------------------|------------------------------------|---|---------------|
| hydro | Kolb Heinz | vHLLE (0 viscosity) | Viscous vHLLE |
| Init cond (soft) | Glauber | EPOS | EPOS4 |
| Init state fluctuations | No | Yes | Yes |
| hadronization | Covar. Inst. Coal + frag | Same | New scheme |
| HQ production | FONLL (p) + Glauber (space) | FONLL (p) + EPOS (space): position of NN interactions | EPOS4 |
| CNM | No shadowing, initial k_T broad. | EPS09 | EPOS4 |
| Hadronic interaction | None | None | URQMD |

Still no modification of the ELOSS model... maybe as an outcome of the comparison

In the following results : elastic + radiative

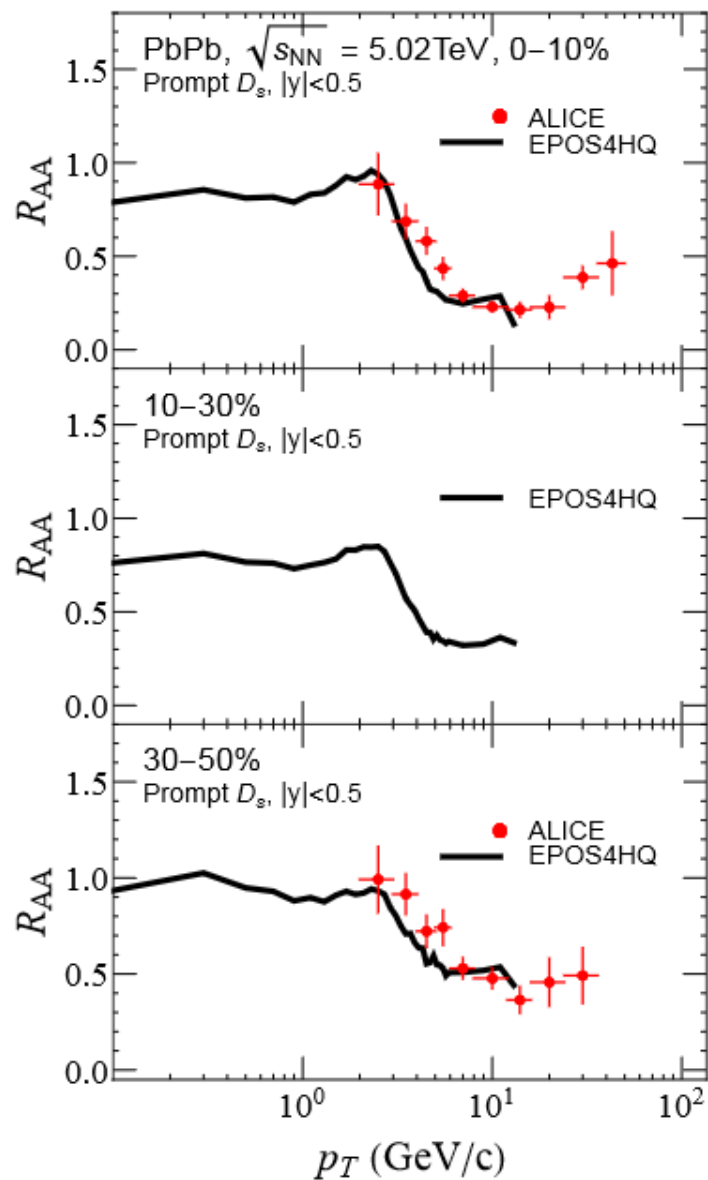
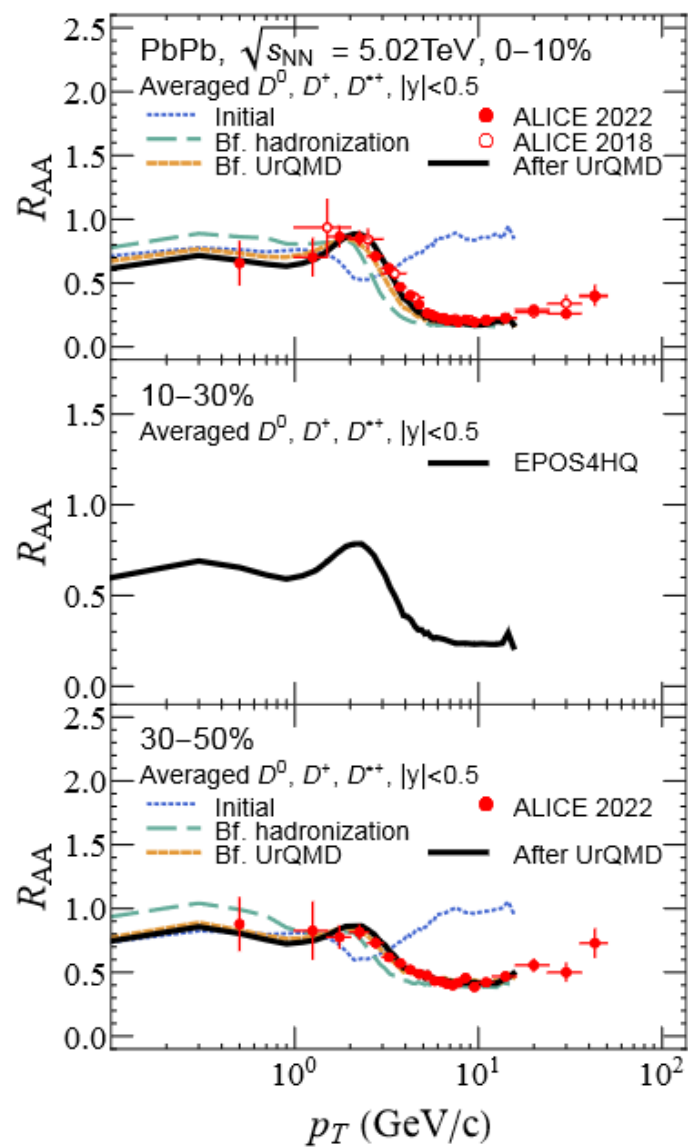
OHF hadrons p_T distributions in EPOS4-HQ

Yield ratios in PbPb @ LHC



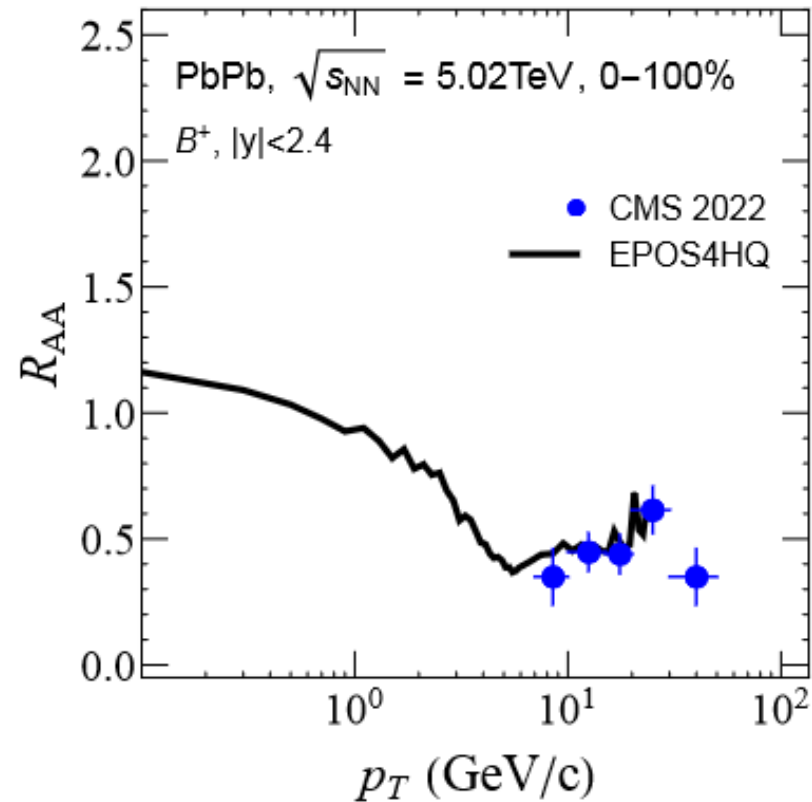
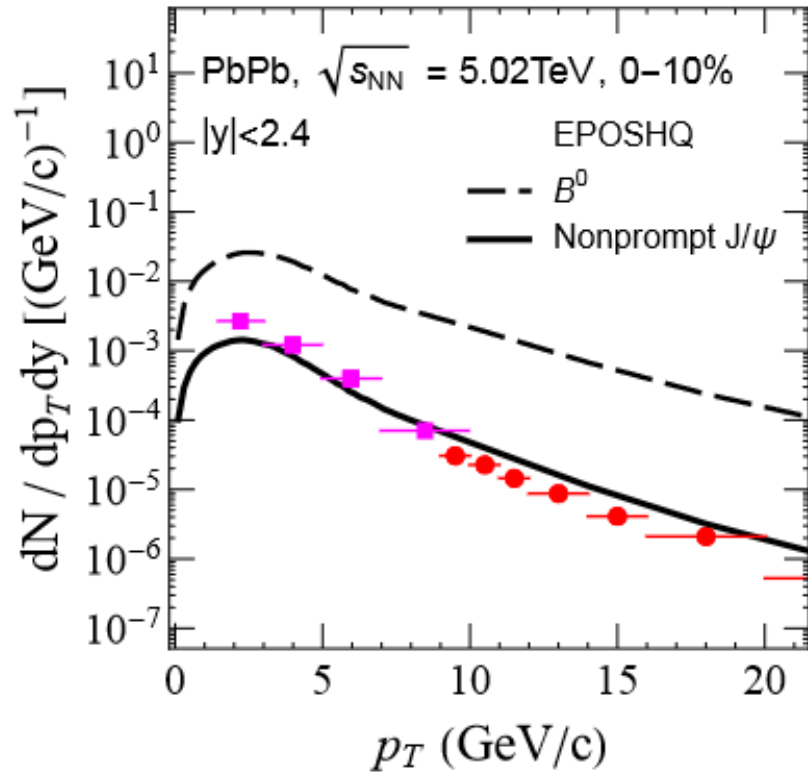
- Experimental trends well reproduced
- Need for more precise data to improve the hadronization “chemistry”

hadrons p_T distributions in EPOS4-HQ



hadrons p_T distributions in EPOS4-HQ

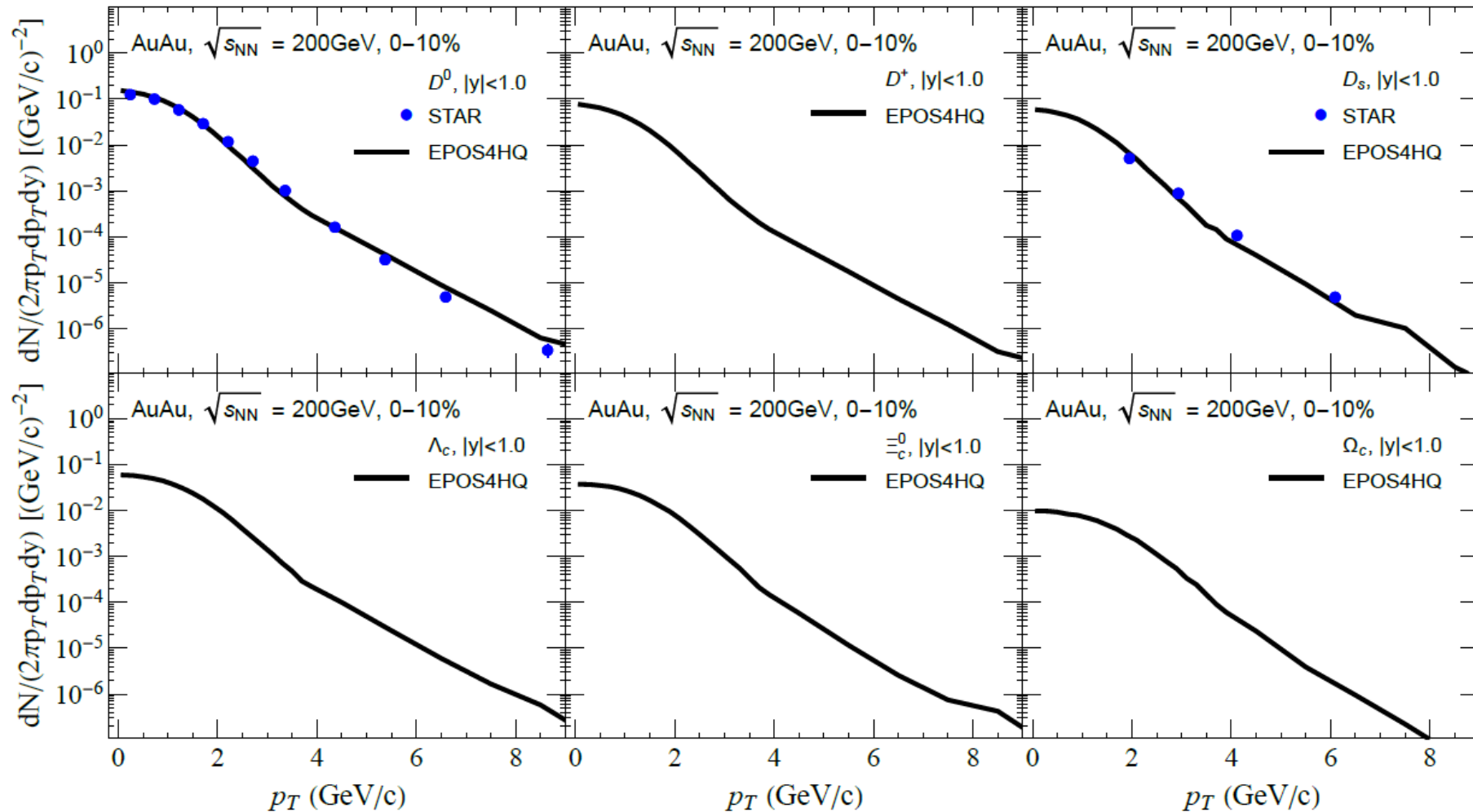
PbPb semi-central @ LHC



Very good agreement for all resonances

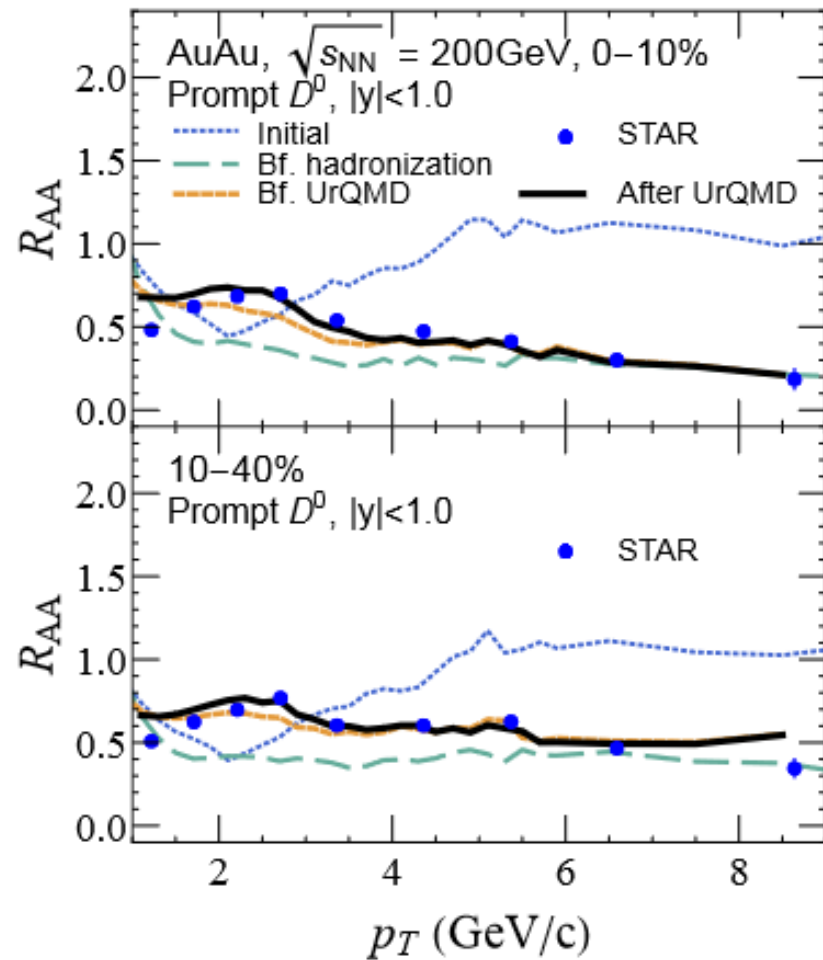
hadrons p_T distributions in EPOS4-HQ

AuAu central @ RHIC



Very good agreement for all resonances

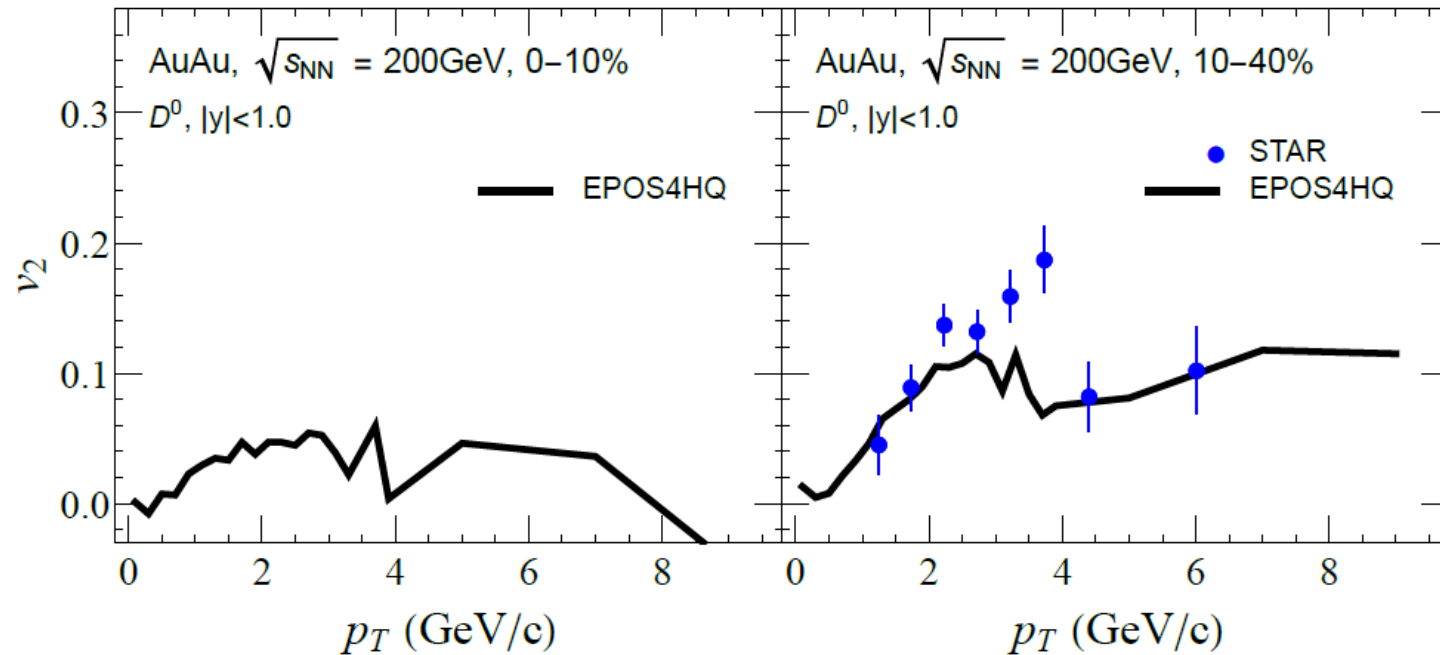
hadrons p_T distributions in EPOS4-HQ



Very good agreement for all resonances

Elliptic flow in EPOS4-HQ

AuAu @ RHIC



- Slight underestimation around 3 GeV/c... probably need sPHENIX data before concluding