

Stars and Galaxies: a pathway to the Dark Side of Fundamental Physics

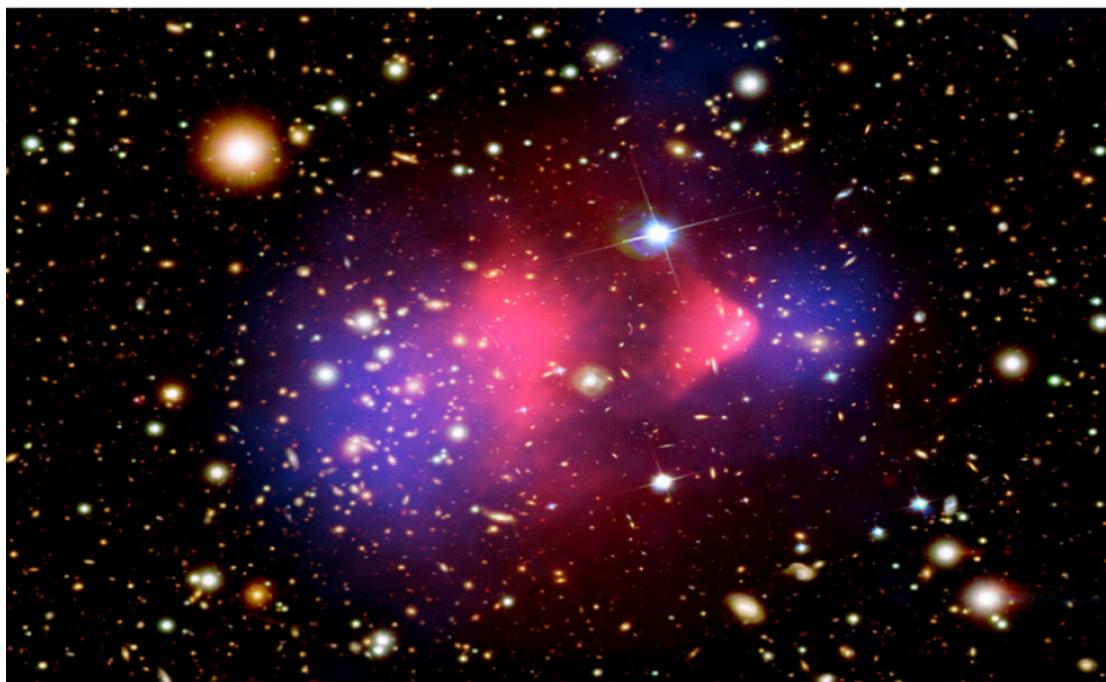
Andrea Caputo

Annecy, 21 March 2023



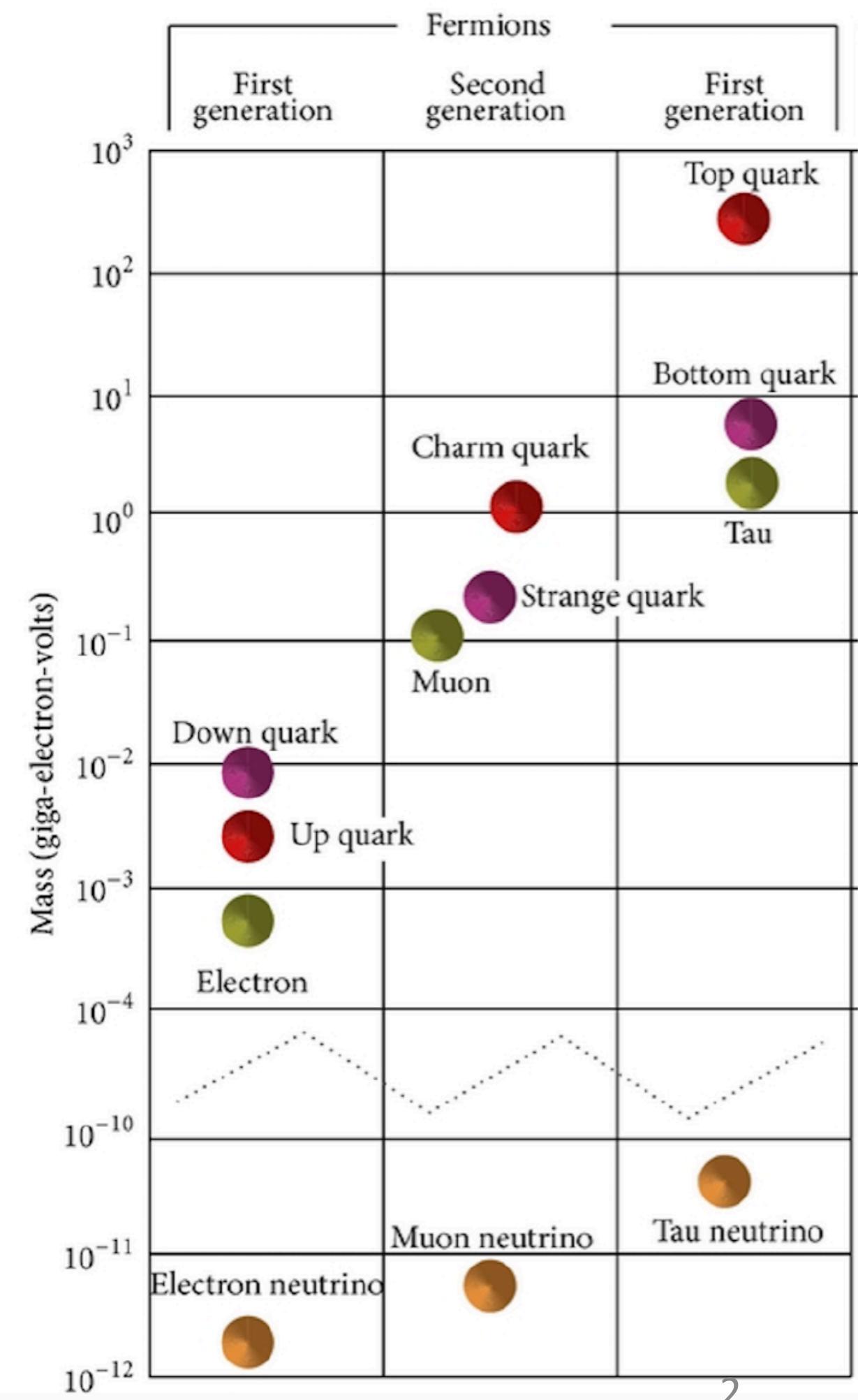
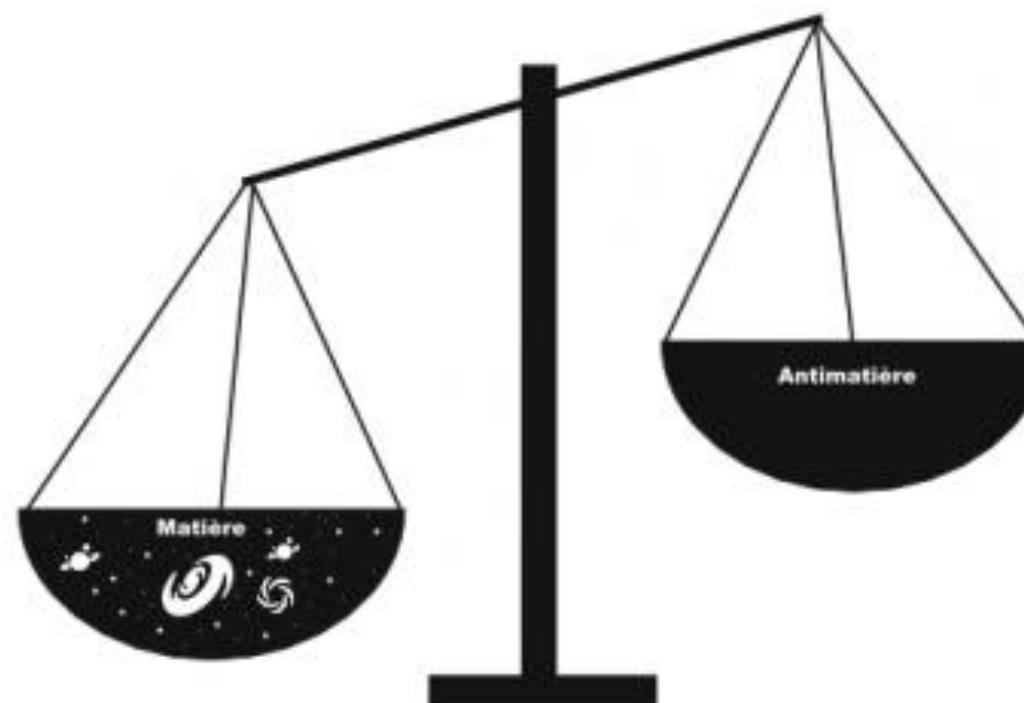
We have a lot of evidences for physics beyond the Standard Model

Dark Matter



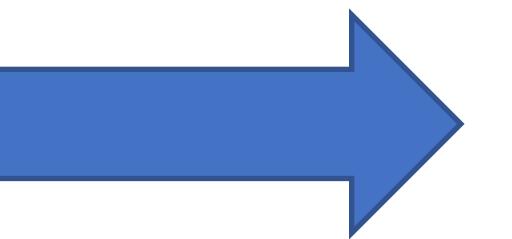
Neutrino Masses

Matter Antimatter asymmetry



The strong CP problem

We know that QCD suffers from the fact
one can add to the Lagrangian the
following term



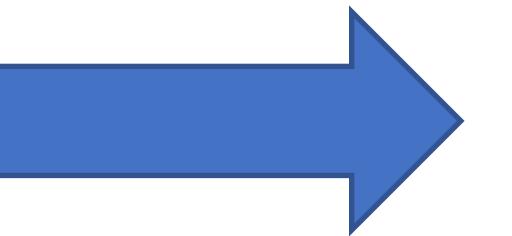
$$\mathcal{L} \supset -\theta \frac{g^2}{16\pi} Tr(G\tilde{G})$$

$$G_{\mu\nu} = G_{\mu\nu}^a T^a$$

$$\tilde{G}_{\mu\nu} = \epsilon_{\mu\nu\rho\sigma} G^{\rho\sigma}$$

The strong CP problem

We know that QCD suffers from the fact one can add to the Lagrangian the following term



$$\mathcal{L} \supset -\theta \frac{g^2}{16\pi} Tr(G\tilde{G})$$

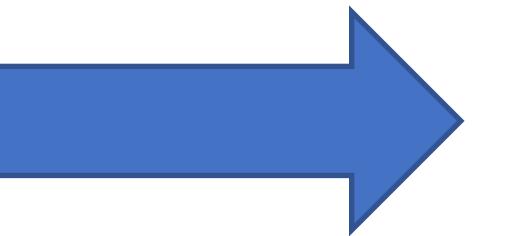
$$G_{\mu\nu} = G_{\mu\nu}^a T^a$$

$$\tilde{G}_{\mu\nu} = \epsilon_{\mu\nu\rho\sigma} G^{\rho\sigma}$$

It can be shown that this term is a *total derivative* and one would naively think we can just forget about it.

The strong CP problem

We know that QCD suffers from the fact one can add to the Lagrangian the following term



$$\mathcal{L} \supset -\theta \frac{g^2}{16\pi} Tr(G\tilde{G})$$

$$G_{\mu\nu} = G_{\mu\nu}^a T^a$$

$$\tilde{G}_{\mu\nu} = \epsilon_{\mu\nu\rho\sigma} G^{\rho\sigma}$$

It can be shown that this term is a *total derivative* and one would naively think we can just forget about it. However, the structure of QCD vacuum makes this term important, with physical consequences.

The strong CP problem

We know that QCD suffers from the fact one can add to the Lagrangian the following term



$$\mathcal{L} \supset -\theta \frac{g^2}{16\pi} \text{Tr}(G\tilde{G})$$

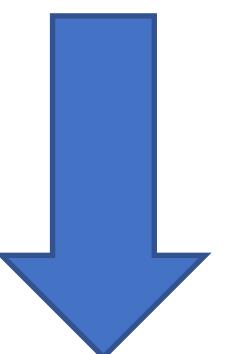
$$G_{\mu\nu} = G_{\mu\nu}^a T^a$$

$$\tilde{G}_{\mu\nu} = \epsilon_{\mu\nu\rho\sigma} G^{\rho\sigma}$$

It can be shown that this term is a *total derivate* and one would naively think we can just forget about it. However, the structure of QCD vacuum makes this term important, with physical consequences



It leads to a dipole moment for the neutron



$$\theta < 10^{-9}$$

The axion

(Peccei, Quinn 1977; Weinberg 1978; Wilczek 1978)

The idea is simple and consists on promoting the dimensionless parameter θ to a *dynamical field*

$$\mathcal{L} \supset -\frac{a(x)}{f_a} \frac{g^2}{16\pi} Tr(G\tilde{G})$$

- (pseudo) Goldstone Boson associated to $U(1)_{PQ}$
- Good dark matter candidate

The axion

(Peccei, Quinn 1977; Weinberg 1978; Wilczek 1978)

The idea is simple and consists on promoting the dimensionless parameter θ to a *dynamical field*

$$\mathcal{L} \supset -\frac{a(x)}{f_a} \frac{g^2}{16\pi} \text{Tr}(G\tilde{G})$$

CP Conservation in the Presence of Instantons

R.D. Peccei (Stanford U., ITP), Helen R. Quinn (Stanford U., ITP)

Mar, 1977

8 pages

Part of *Gauge models of CP violation*, 1440-1443

Published in: *Phys.Rev.Lett.* 38 (1977) 1440-1443

DOI: [10.1103/PhysRevLett.38.1440](https://doi.org/10.1103/PhysRevLett.38.1440)

Report number: ITP-568-STANFORD

View in: [OSTI Information Bridge Server](#), [ADS Abstract Service](#)

links cite claim

reference search

A New Light Boson?

Steven Weinberg (Harvard U.)

Dec, 1977

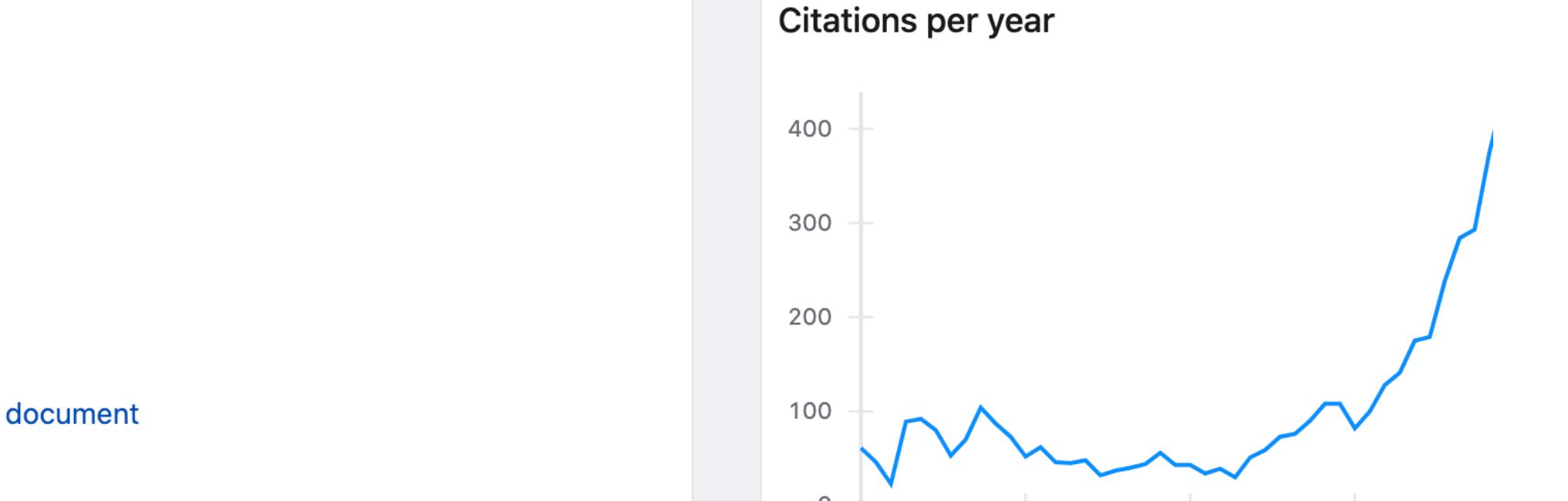
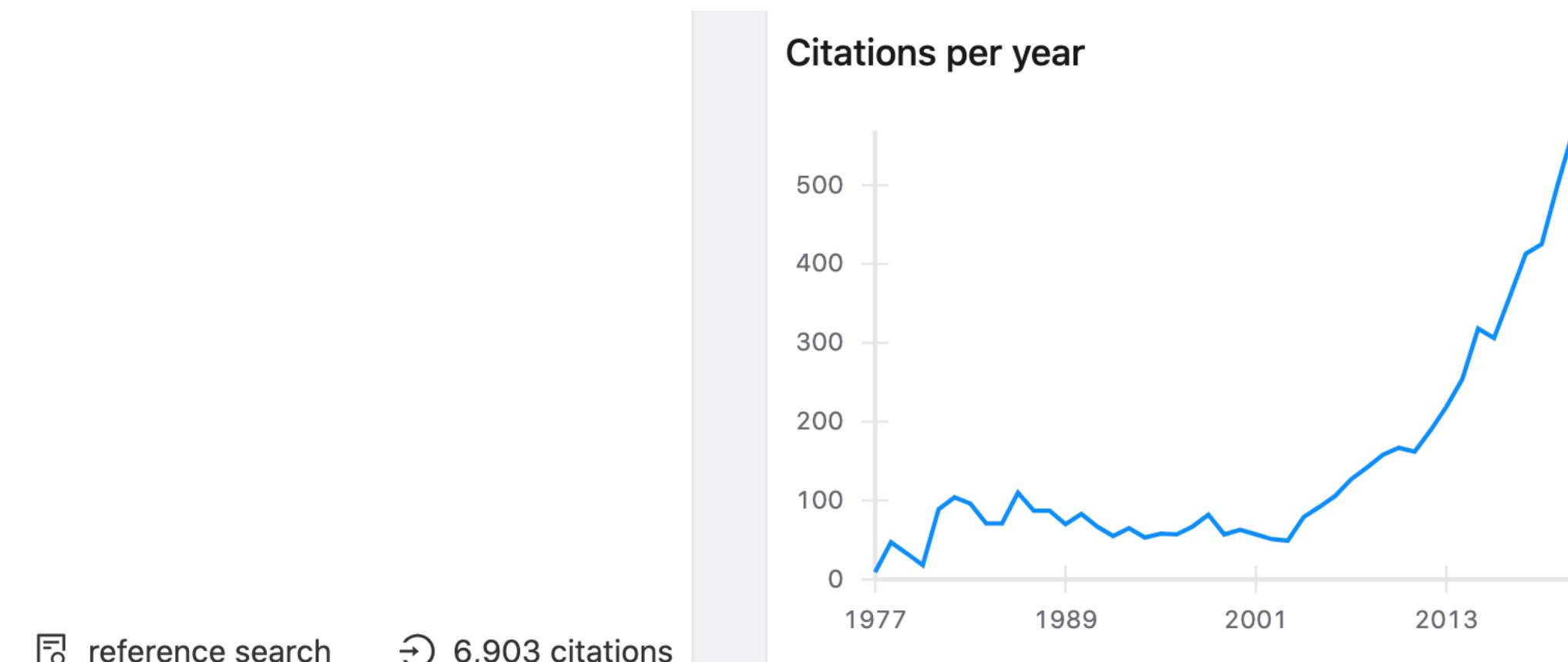
12 pages

Published in: *Phys.Rev.Lett.* 40 (1978) 223-226

DOI: [10.1103/PhysRevLett.40.223](https://doi.org/10.1103/PhysRevLett.40.223)

Report number: HUTP-77/A074

View in: [OSTI Information Bridge Server](#), [ADS Abstract Service](#), [KEK scanned document](#)



The axion

(Peccei, Quinn 1977; Weinberg 1978; Wilczek 1978)

The idea is simple and consists on promoting the dimensionless parameter θ to a *dynamical field*

$$\mathcal{L} \supset -\frac{a(x)}{f_a} \frac{g^2}{16\pi} Tr(G\tilde{G})$$

- (pseudo) Goldstone Boson associated to $U(1)_{PQ}$
- Good dark matter candidate

In a theory with the quarks charged under electromagnetism one gets a coupling to the photon performing a *chiral* (and *anomalous*) transformation



$$\mathcal{L}_{a\gamma\gamma} = -\frac{1}{4} g_{a\gamma} a \tilde{F}^{\mu\nu} F_{\mu\nu}$$

The axion

(Peccei, Quinn 1977; Weinberg 1978; Wilczek 1978)

The idea is simple and consists on promoting the dimensionless parameter θ to a *dynamical field*

$$\mathcal{L} \supset -\frac{a(x)}{f_a} \frac{g^2}{16\pi} \text{Tr}(G\tilde{G})$$

- (pseudo) Goldstone Boson associated to $U(1)_{PQ}$
- Good dark matter candidate

Axions In String Theory

Peter Svrcek, Edward Witten

In a theory with the quarks charged under electromagnetism one gets a coupling to the photon performing a *chiral* (and *anomalous*) transformation



$$\mathcal{L}_{a\gamma\gamma} = -\frac{1}{4} g_{a\gamma} a \tilde{F}^{\mu\nu} F_{\mu\nu}$$

The axion

(Peccei, Quinn 1977; Weinberg 1978; Wilczek 1978)

The idea is simple and consists on promoting the dimensionless parameter θ to a *dynamical field*

$$\mathcal{L} \supset -\frac{a(x)}{f_a} \frac{g^2}{16\pi} \text{Tr}(G\tilde{G})$$

- (pseudo) Goldstone Boson associated to $U(1)_{PQ}$
- Good dark matter candidate

Axions In String Theory

Peter Svrcek, Edward Witten

In a theory with the quarks charged under electromagnetism one gets a coupling to the photon performing a *chiral* (and *anomalous*) transformation

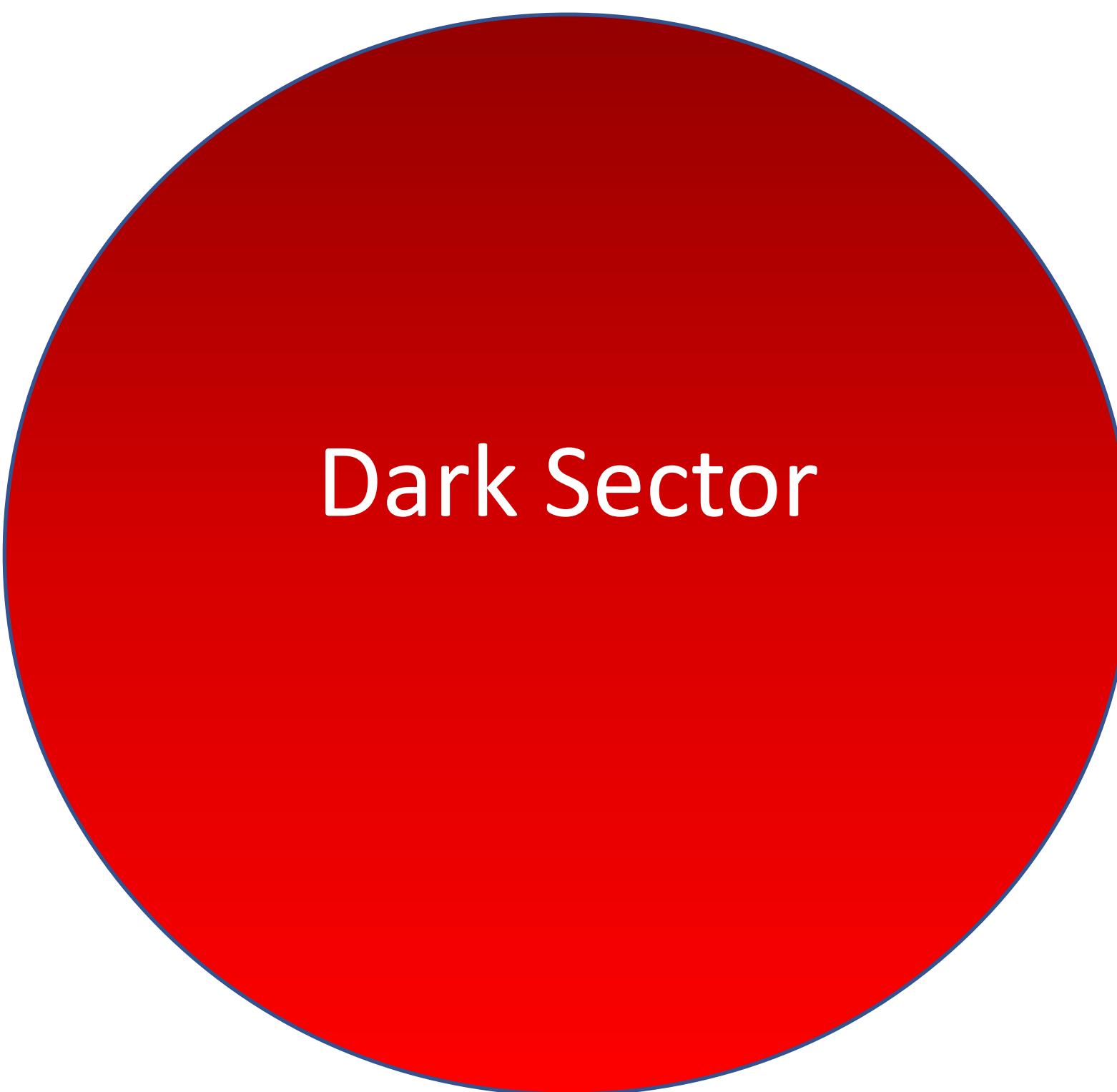


$$\mathcal{L}_{a\gamma\gamma} = -\frac{1}{4} g_{a\gamma} a \tilde{F}^{\mu\nu} F_{\mu\nu}$$

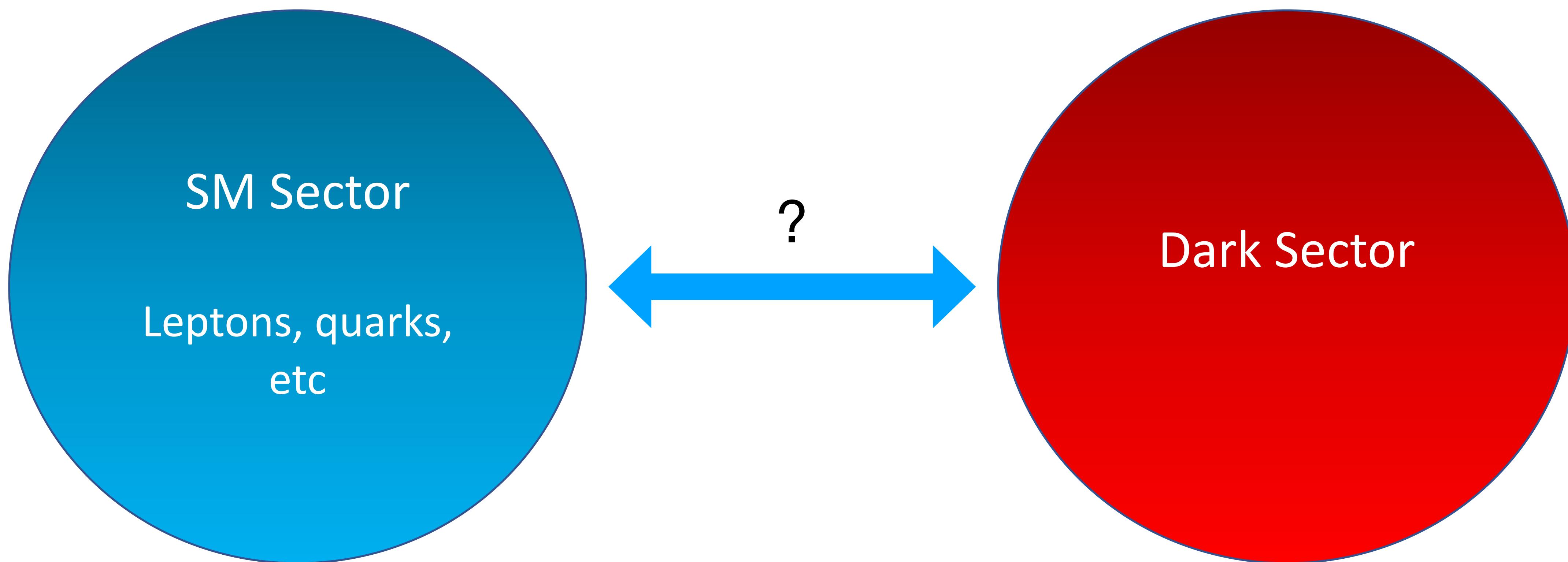
String Axiverse

Asimina Arvanitaki, Savas Dimopoulos, Sergei Dubovsky, Nemanja Kaloper, John March-Russell

More generally we can have a rich Dark Sector, that is to say new particles which can be connected in different ways to the SM particles

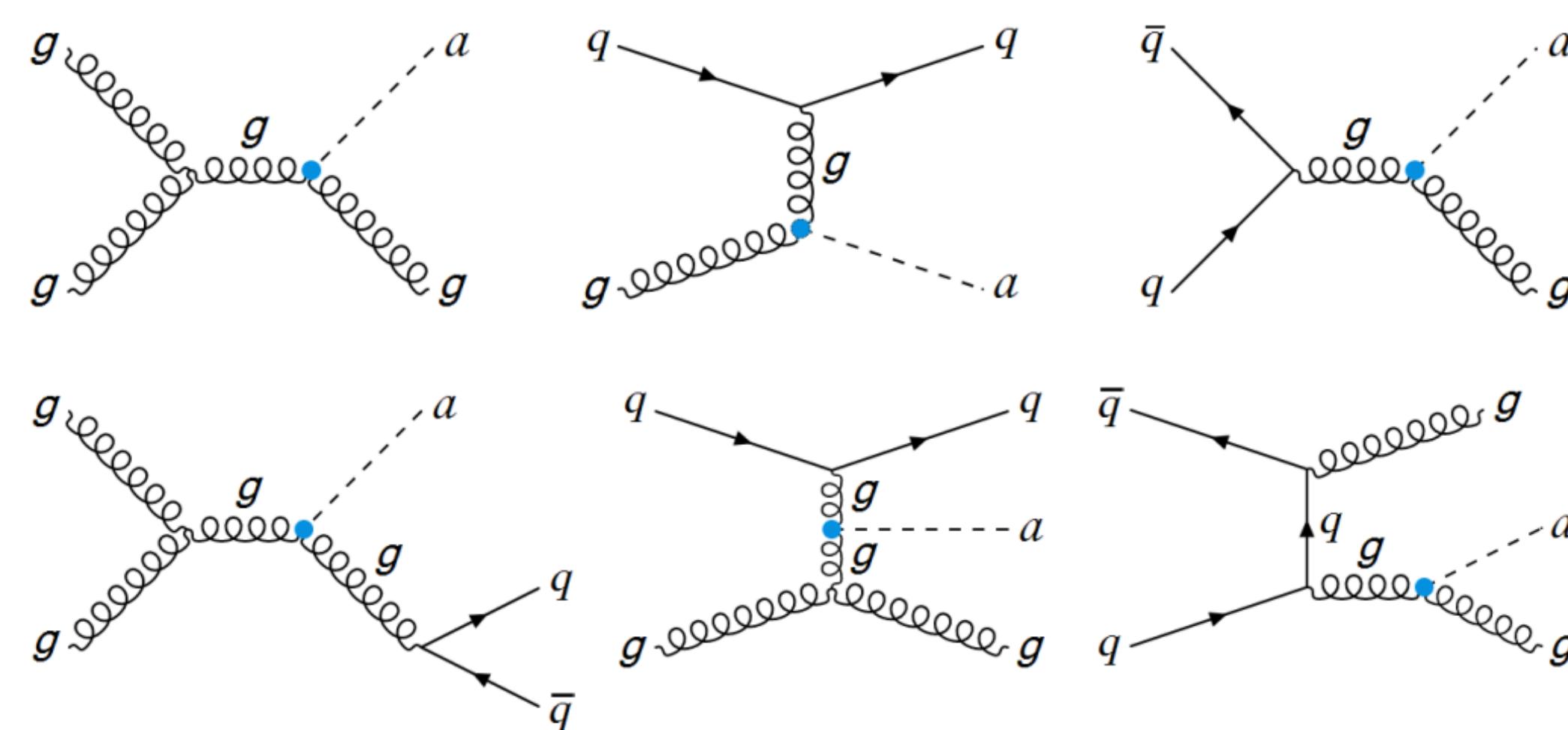


More generally we can have a rich Dark Sector, that is to say new particles which can be connected in different ways to the SM particles



How to look for these new particles?

Collider searches. We have been doing this successfully for decades, however it is more and more difficult to reach high energies



Caputo, P. Hernandez, M. Kekic, J. Lopez-Pavon, J.Salvado, *Eur.Phys.J.C* 77 (2017) 4, 258

Caputo, P. Hernandez, J. Lopez-Pavon, J.Salvado, *JHEP* 06 (2017) 112

D. Barducci, E. Bertuzzo, Caputo, P. Hernandez, *JHEP* 06 (2020) 185

D. Barducci, E. Bertuzzo, Caputo, P. Hernandez, B. Mele, *JHEP* 03 (2021) 117

How to look for these new particles?



“The Universe is a “poor man”
accelerator” Zel’dovich

Zel'dovich

Be creative and use **astrophysical objects and the universe!**

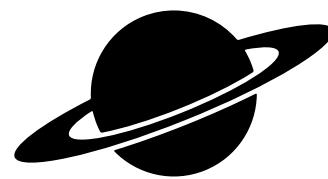
Two parts

- Dark Matter particles
- New particles which are NOT Dark Matter

I. Looking for dark matter



Stars and new physics



- Caputo et al., *JCAP* 08 (2022) 08, 045
Caputo et al., *Phys.Rev.Lett.* 128 (2022) 22, 221103 (CCSNe)
Caputo et al., *Phys.Rev.D* 105 (2022) 3, 035022 (CCSNe)
Caputo et al., *Phys.Rev.Lett.* 127 (2021) 18, 181102 (Hypernovae)
Caputo et al., *Phys.Rev.D* 101 (2020) 12, 123004 (Sun)
O'Hare, Caputo, et al., *Phys.Rev.D* 102 (2020) 4, 043019 (Sun)

Dark Matter and telescopes



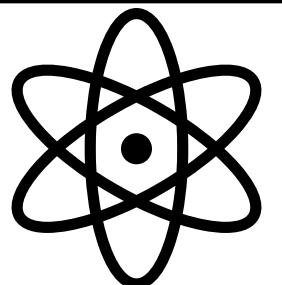
- Caputo et al., *JCAP* 03 (2019) 027 (radio)
Caputo et al., *Phys.Rev.D* 98 (2018) 8, 083024 (radio)
Caputo et al., *JCAP* 03 (2020) 001 (X-rays)
Caputo et al., *JCAP* 05 (2021) 046 (IR)

Gravitational Waves

- Caputo et al., *JAstrophys.J.* 892 (2020) 2, 90
Toubiana, Sberna, Caputo, et al.,
Phys.Rev.Lett. 126 (2021) 10, 101105

My research

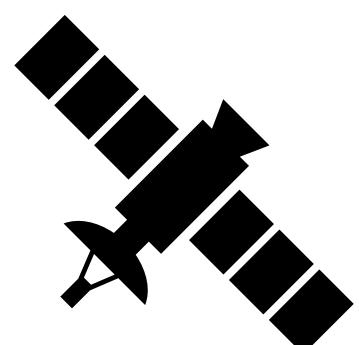
Dark Matter direct detection



- Caputo et al., *Phys.Rev.D* 100 (2019) 11, 116007 (sub-GeV DM)
Caputo et al., *Phys.Lett.B* 802 (2020) 135258 (sub-GeV DM)
Caputo et al., *Phys.Rev.D* 103 (2021) 5, 055017 (sub-GeV DM)
Bloch, Caputo et al., *JHEP* 01 (2021) 178 (sub-GeV DM)

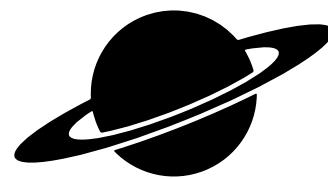
Caputo et al., *Phys.Rev.D* 104 (2021) 9, 095029 (Dark Photon DM)

Dark Matter and other indirect probes



- Caputo et al., *Phys.Rev.Lett.* 127 (2021) 1, 011102 (21 cm)
Bolton, Caputo, et al., *Phys.Rev.Lett.* 129 (2022) 21, 211102 (Ly-alpha)
Caputo et al., *Phys.Rev.Lett.* 125 (2020) 22, 221303 (CMB)
Bernal, Caputo, et al., *arXiv* 2208.13794 (blazars data)
Caputo et al., *Phys.Rev.D* 100 (2019) 6, 063515 (pulsars)
Caputo et al., *Phys.Dark Univ.* 19 (2018) 1-11 (pulsars)

Stars and new physics



Caputo et al., *JCAP* 08 (2022) 08, 045

Caputo et al., *Phys.Rev.Lett.* 128 (2022) 22, 221103 (CCSNe)

Caputo et al., *Phys.Rev.D* 105 (2022) 3, 035022 (CCSNe)

Caputo et al., *Phys.Rev.Lett.* 127 (2021) 18, 181102 (Hypernovae)

Caputo et al., *Phys.Rev.D* 101 (2020) 12, 123004 (Sun)

O'Hare, Caputo, et al., *Phys.Rev.D* 102 (2020) 4, 043019 (Sun)

Dark Matter and telescopes



Caputo et al., *JCAP* 03 (2019) 027 (radio)

Caputo et al., *Phys.Rev.D* 98 (2018) 8, 083024 (radio)

Caputo et al., *JCAP* 03 (2020) 001 (X-rays)

Caputo et al., *JCAP* 05 (2021) 046 (IR)

Bernal, Caputo, et al., *Phys.Rev.D* 103 (2021) 6, 063523 (optic)

Bernal, Caputo, et al., *Phys.Rev.Lett.* 127 (2021) 13, 131102 (optic)

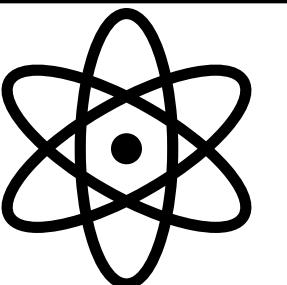
Gravitational Waves

Caputo et al., *J Astrophys.J.* 892 (2020) 2, 90

Toubiana, Sberna, Caputo, et al.,
Phys.Rev.Lett. 126 (2021) 10, 101105

My research

Dark Matter direct detection



Caputo et al., *Phys.Rev.D* 100 (2019) 11, 116007 (sub-GeV DM)

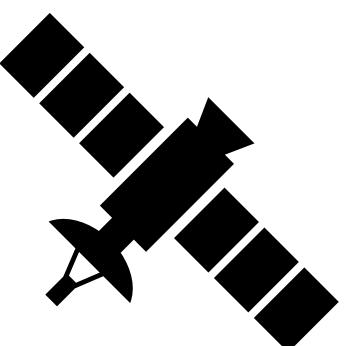
Caputo et al., *Phys.Lett.B* 802 (2020) 135258 (sub-GeV DM)

Caputo et al., *Phys.Rev.D* 103 (2021) 5, 055017 (sub-GeV DM)

Bloch, Caputo et al., *JHEP* 01 (2021) 178 (sub-GeV DM)

Caputo et al., *Phys.Rev.D* 104 (2021) 9, 095029 (Dark Photon DM)

Dark Matter and other indirect probes



Caputo et al., *Phys.Rev.Lett.* 127 (2021) 1, 011102 (21 cm)
Bolton, Caputo, et al., *Phys.Rev.Lett.* 129 (2022) 21, 211102 (Ly-alpha)

Caputo et al., *Phys.Rev.Lett.* 125 (2020) 22, 221303 (CMB)

Bernal, Caputo, et al., *arXiv* 2208.13794 (blazars data)

Caputo et al., *Phys.Rev.D* 100 (2019) 6, 063515 (pulsars)

Caputo et al., *Phys.Dark Univ.* 19 (2018) 1-11 (pulsars)

I will focus on the axion parameter space

$$\mathcal{L}_{a\gamma\gamma} = -\frac{1}{4}g_{a\gamma}a\tilde{F}^{\mu\nu}F_{\mu\nu}$$

I will focus on the axion parameter space

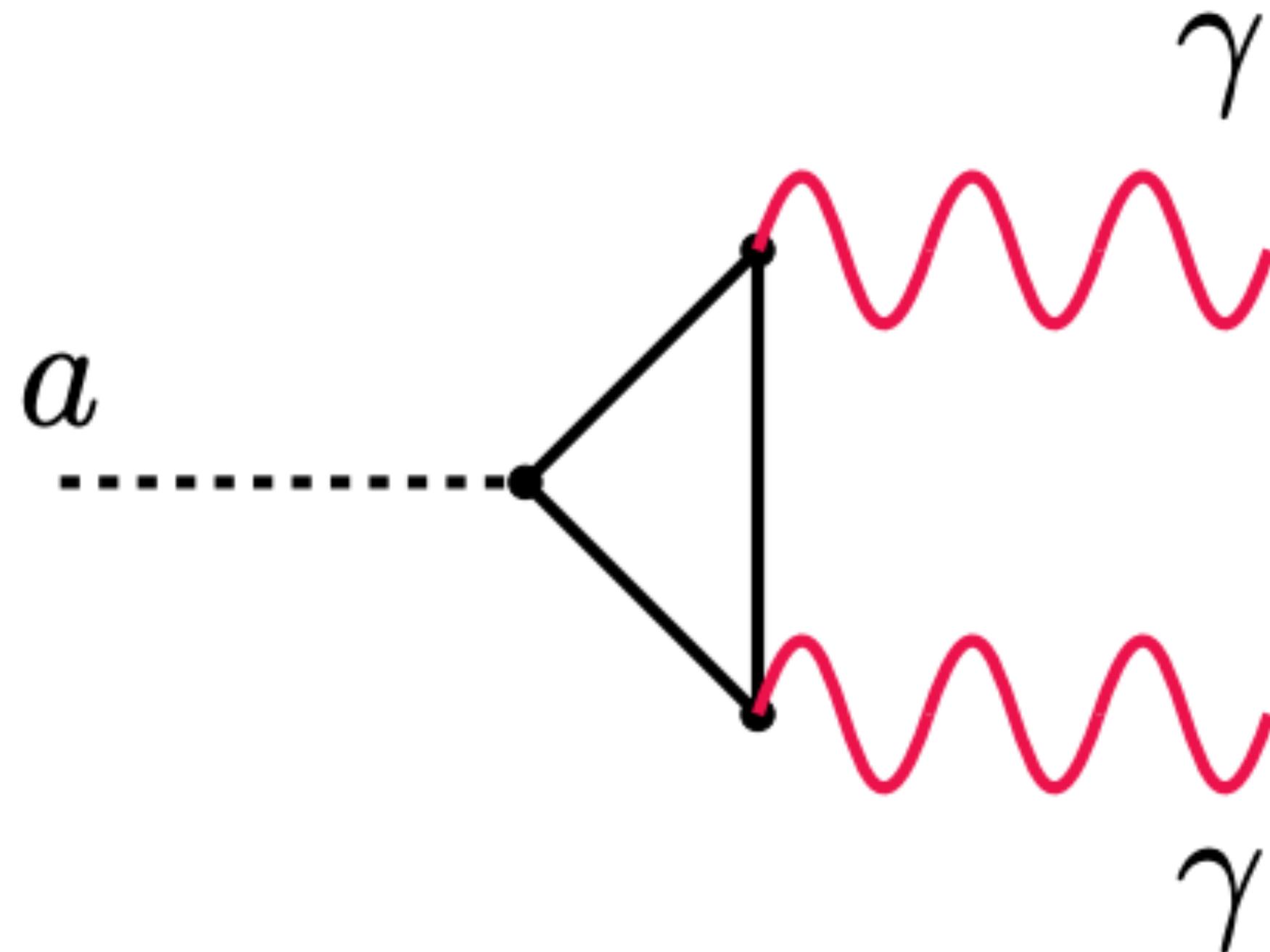
$$\mathcal{L}_{a\gamma\gamma} = -\frac{1}{4}g_{a\gamma}a\tilde{F}^{\mu\nu}F_{\mu\nu}$$

However the ideas will be further more generic, applying to sterile neutrinos DM, dipole DM, scalar DM, hidden photino DM, etc..

Constraining light dark matter with diffuse X-ray and gamma-ray observations

[Rouven Essig](#), [Eric Kuflik](#), [Samuel D. McDermott](#)✉, [Tomer Volansky](#) & [Kathryn M. Zurek](#)

Axion dark matter decaying into photons



$$\tau_a = \frac{64 \pi}{m_a^3 g_{a\gamma\gamma}^2}$$

General strategy

A good story should answer few basic questions

General strategy

A good story should answer few basic questions

What?

General strategy

A good story should answer few basic questions

What?

Looking for a photon monochromatic emission at $E_\gamma = m_a/2$
given by ALP decay from regions with high dark-matter density

General strategy

A good story should answer few basic questions

Where?

General strategy

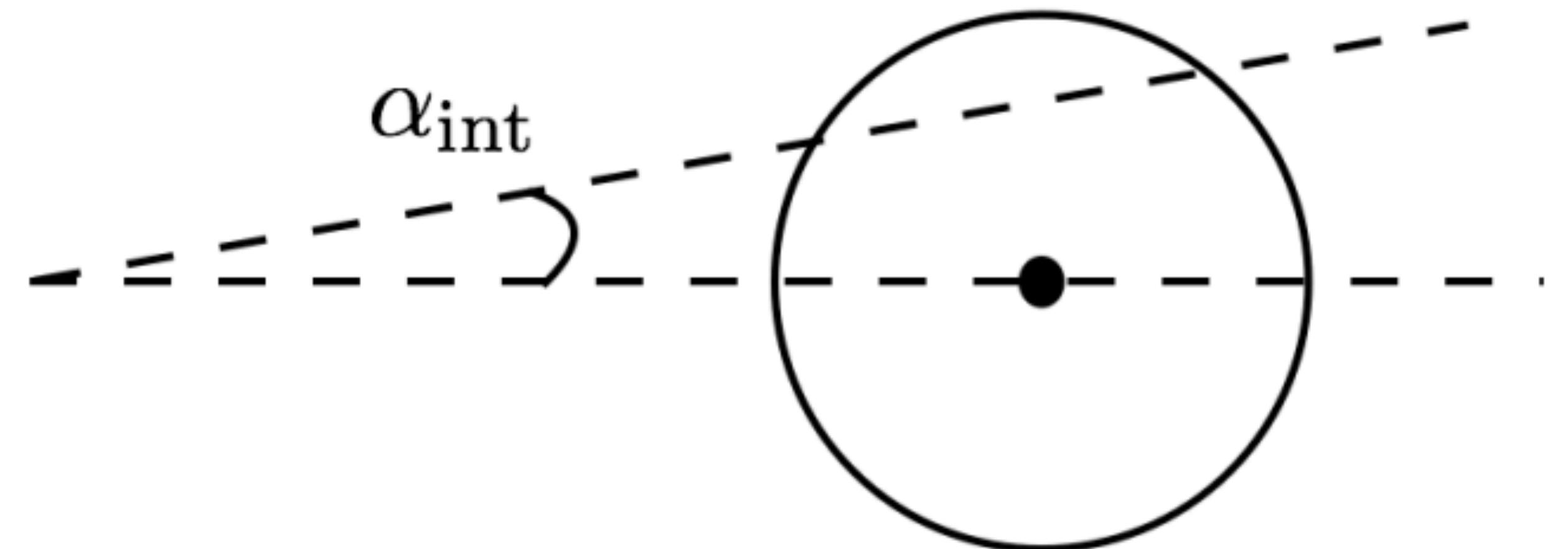
A good story should answer few basic questions

Where?

A) Which astrophysical targets?

$$S_{sd} \propto \int d\Omega d\ell \frac{\rho_a(\Omega, \ell)}{\tau_a}$$

Galactic center, clusters of galaxies, Dwarf galaxies



General strategy

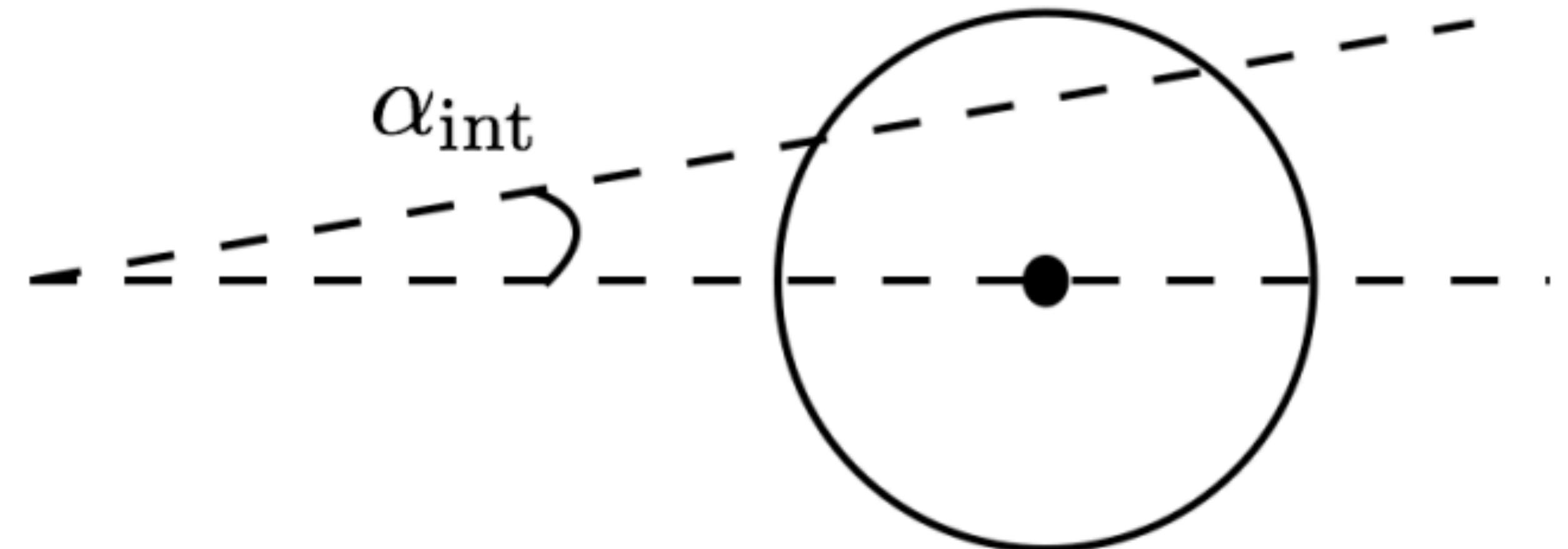
A good story should answer few basic questions

Where?

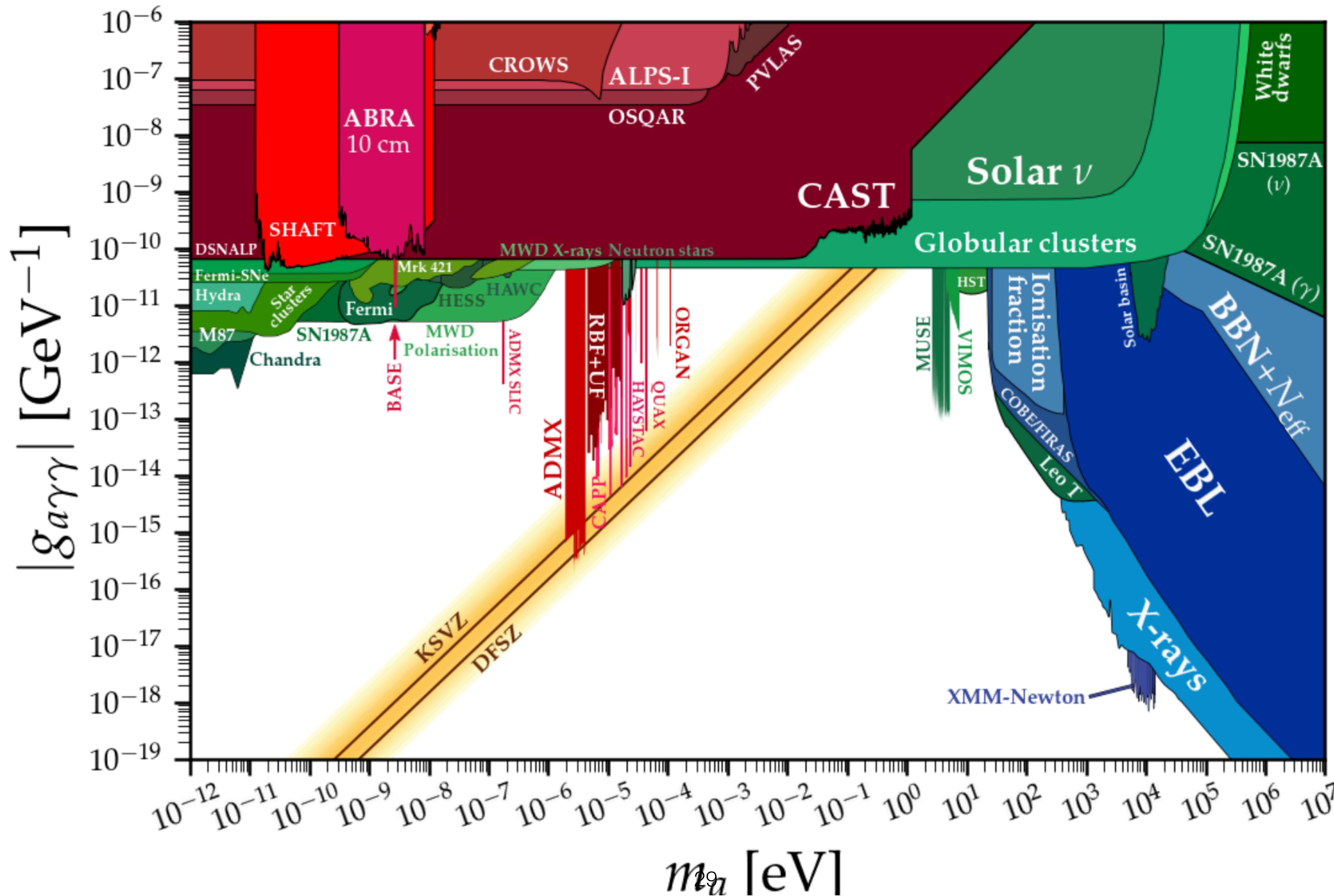
A) Which astrophysical targets?

$$S_{sd} \propto \int d\Omega d\ell \frac{\rho_a(\Omega, \ell)}{\tau_a}$$

Galactic center, clusters of galaxies, Dwarf galaxies

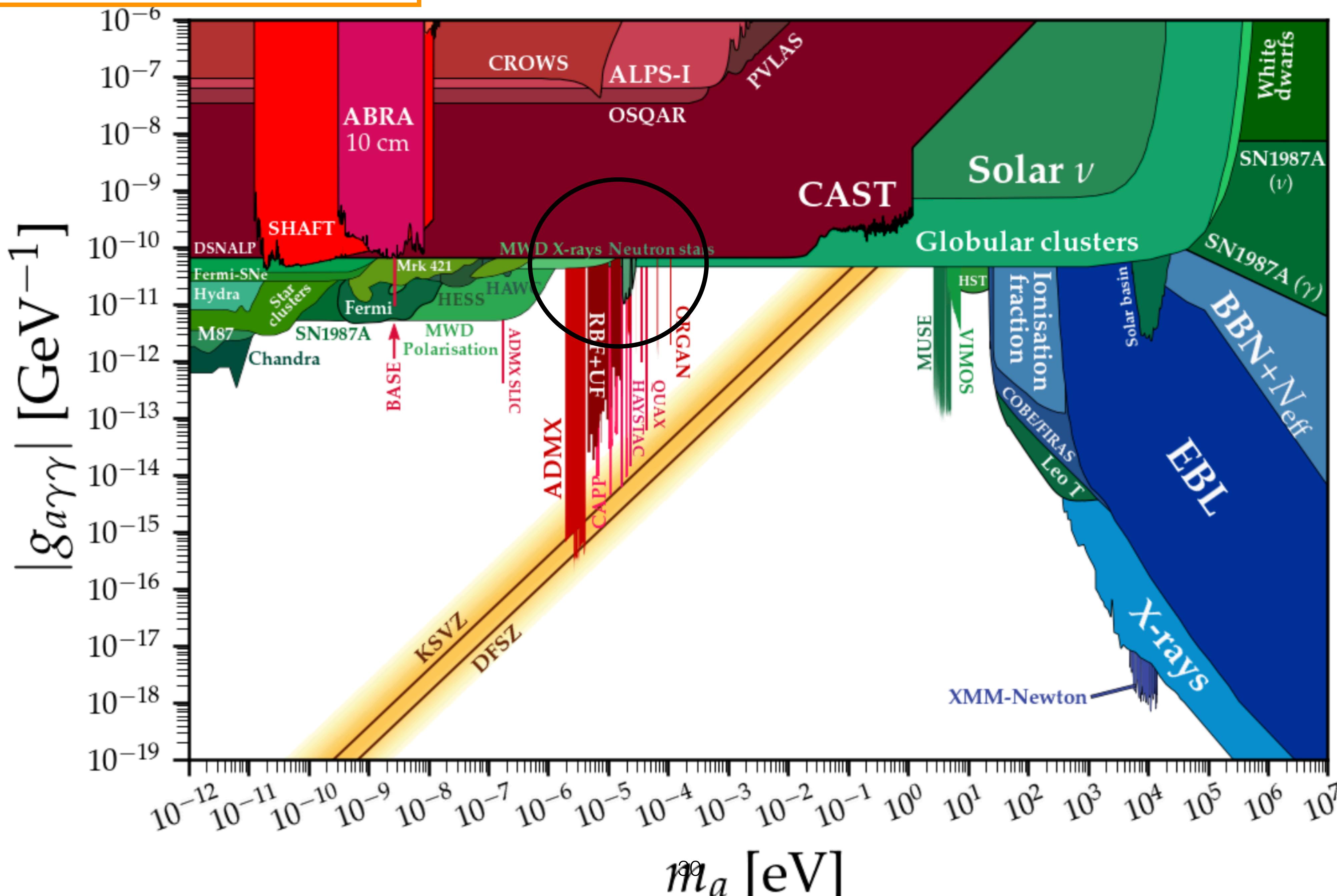


B) Which energy range?



micro-eV (radio)

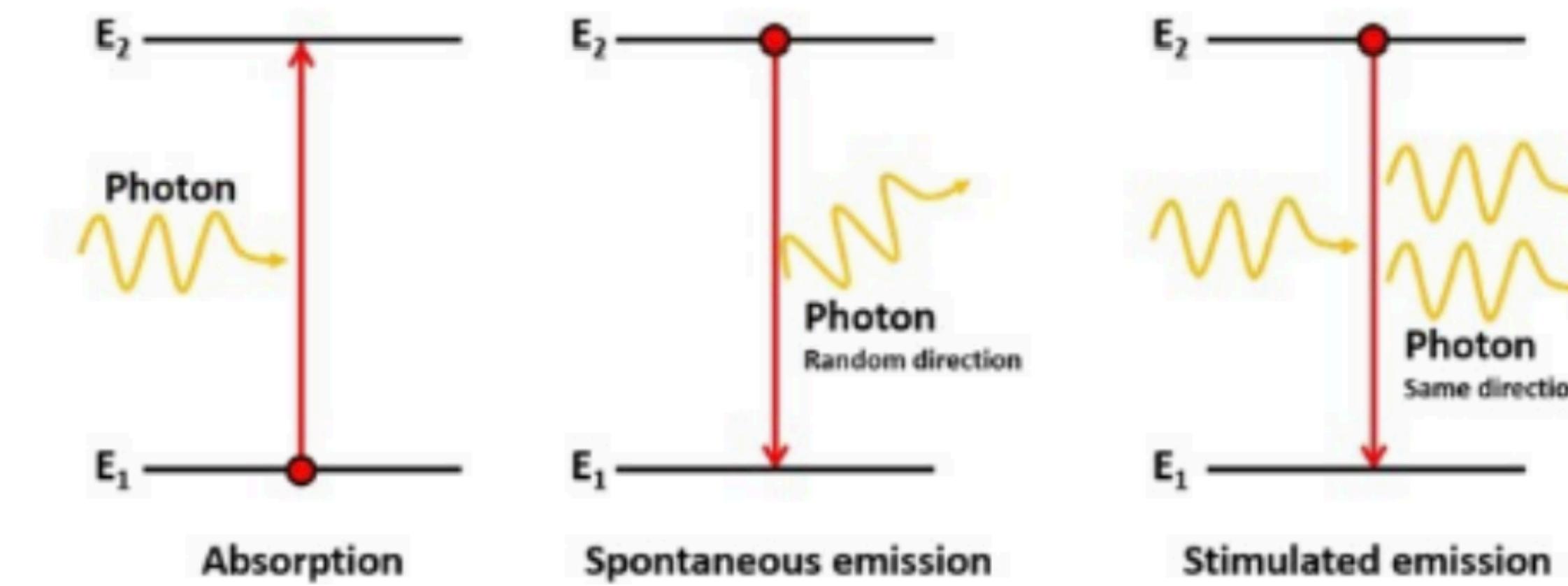
<https://github.com/cajohare/AxionLimits>



In this energy range we realised that **stimulated emission** is important,
the ALP decay is not happening in vacuum!

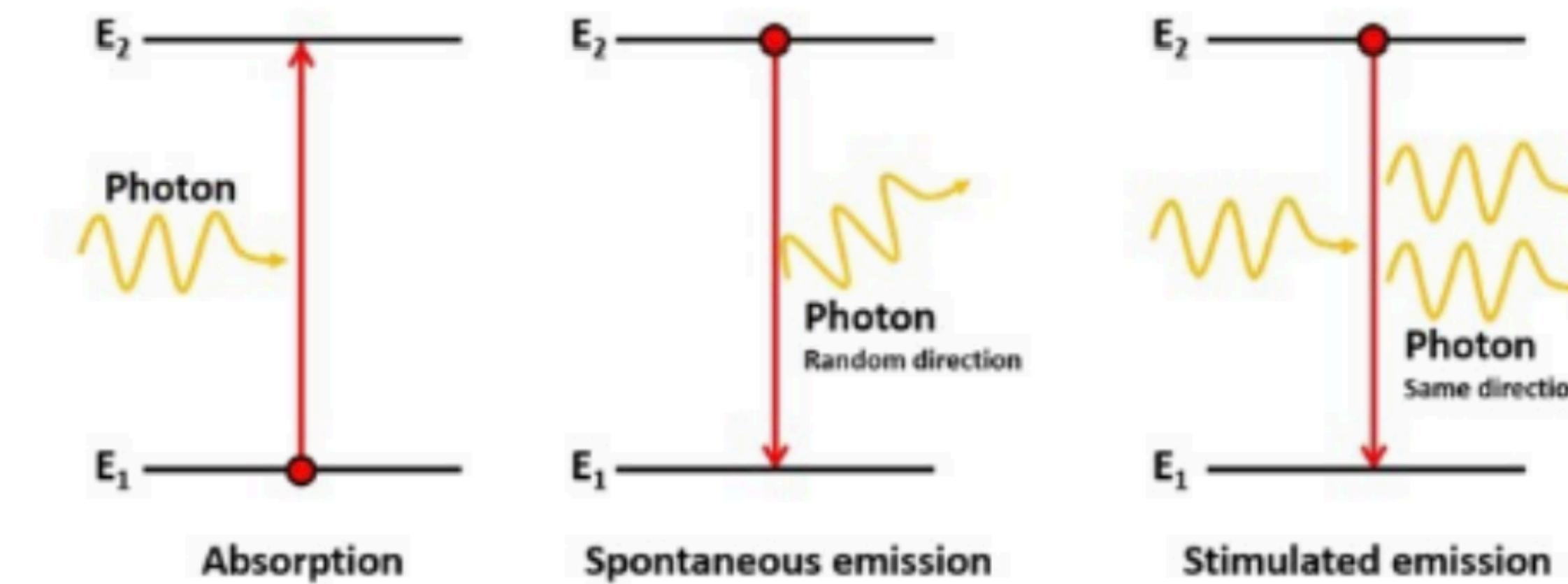
In this energy range we realised that **stimulated emission** is important,
the ALP decay is not happening in vacuum!

Stimulated decay



In this energy range we realised that **stimulated emission** is important,
the ALP decay is not happening in vacuum!

Stimulated decay



$$S_{\text{decay}} = \frac{\Gamma_a}{4\pi\Delta\nu} \int d\Omega d\ell \rho_a(\ell, \Omega) [1 + 2f_\gamma(\ell, \Omega, m_a)]$$

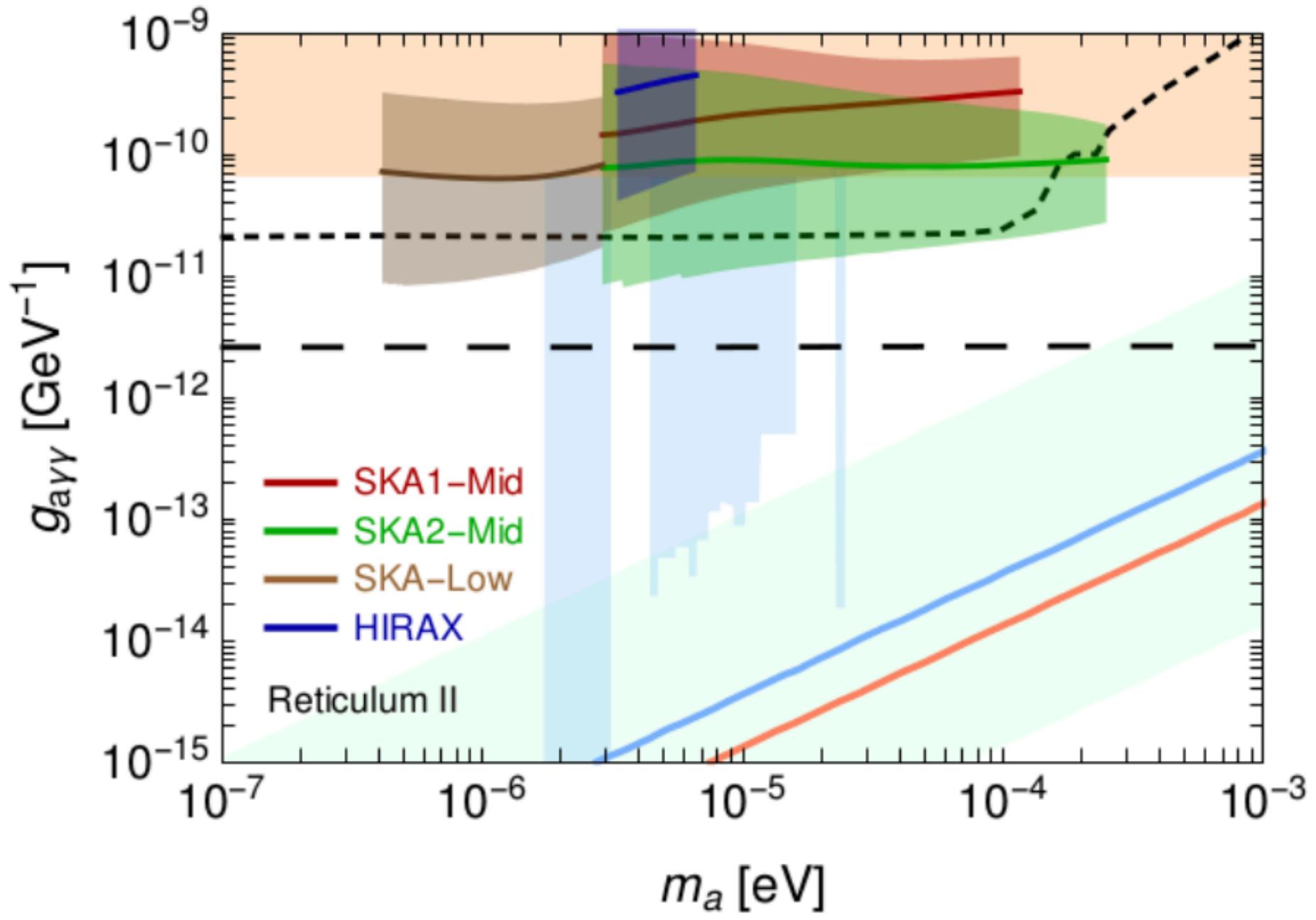
$$f_\gamma = \frac{\pi^2 \rho_\gamma}{E_\gamma^3}$$

A golden era has started for radio astronomy
with SKA and its precursors.

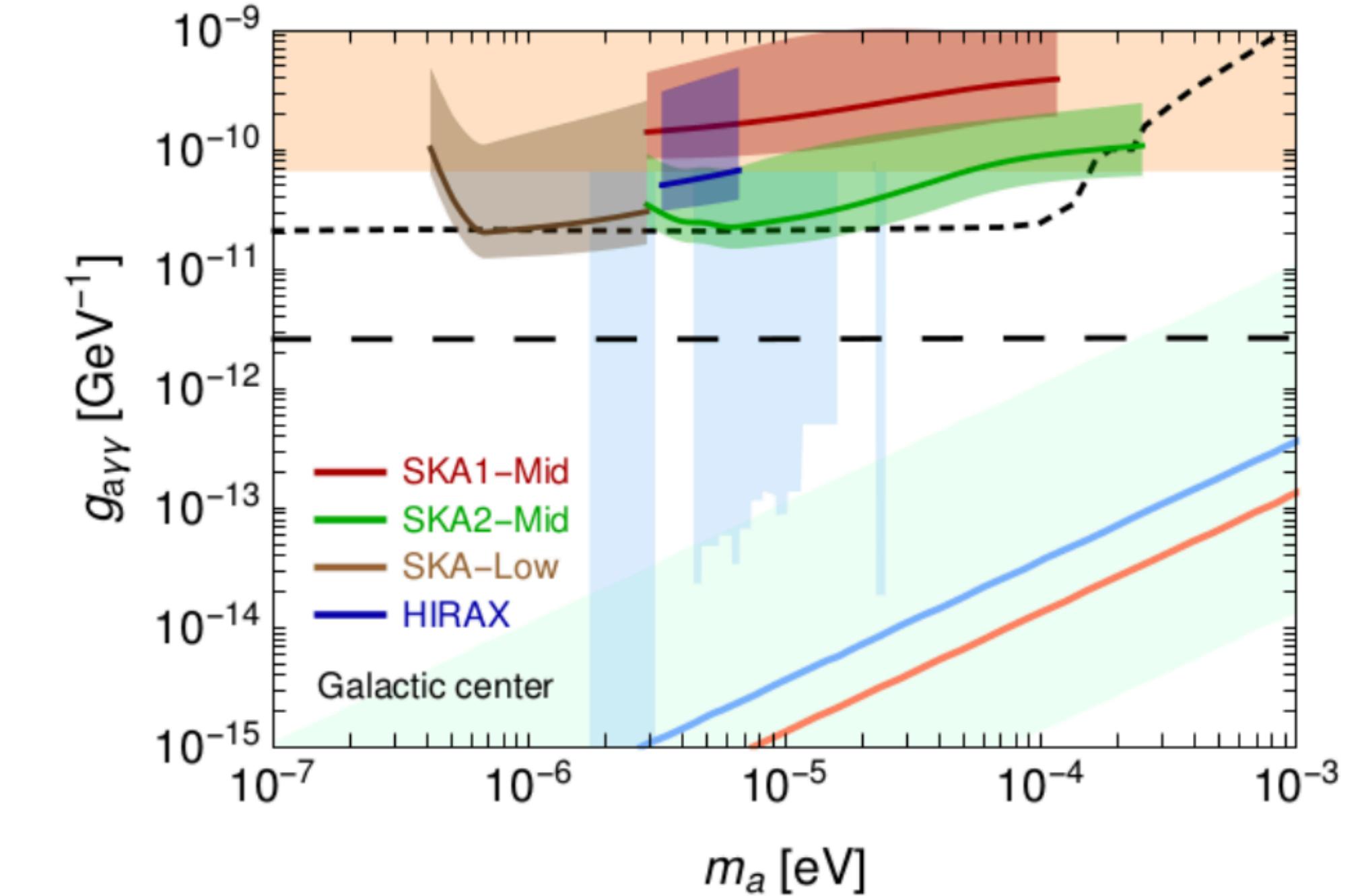
**SKA1-Low: 100 hours of observation;
sensitivity of $180\mu\text{Jy}/\text{beam}$**
 $\Delta\nu/\nu = 10^{-4}$



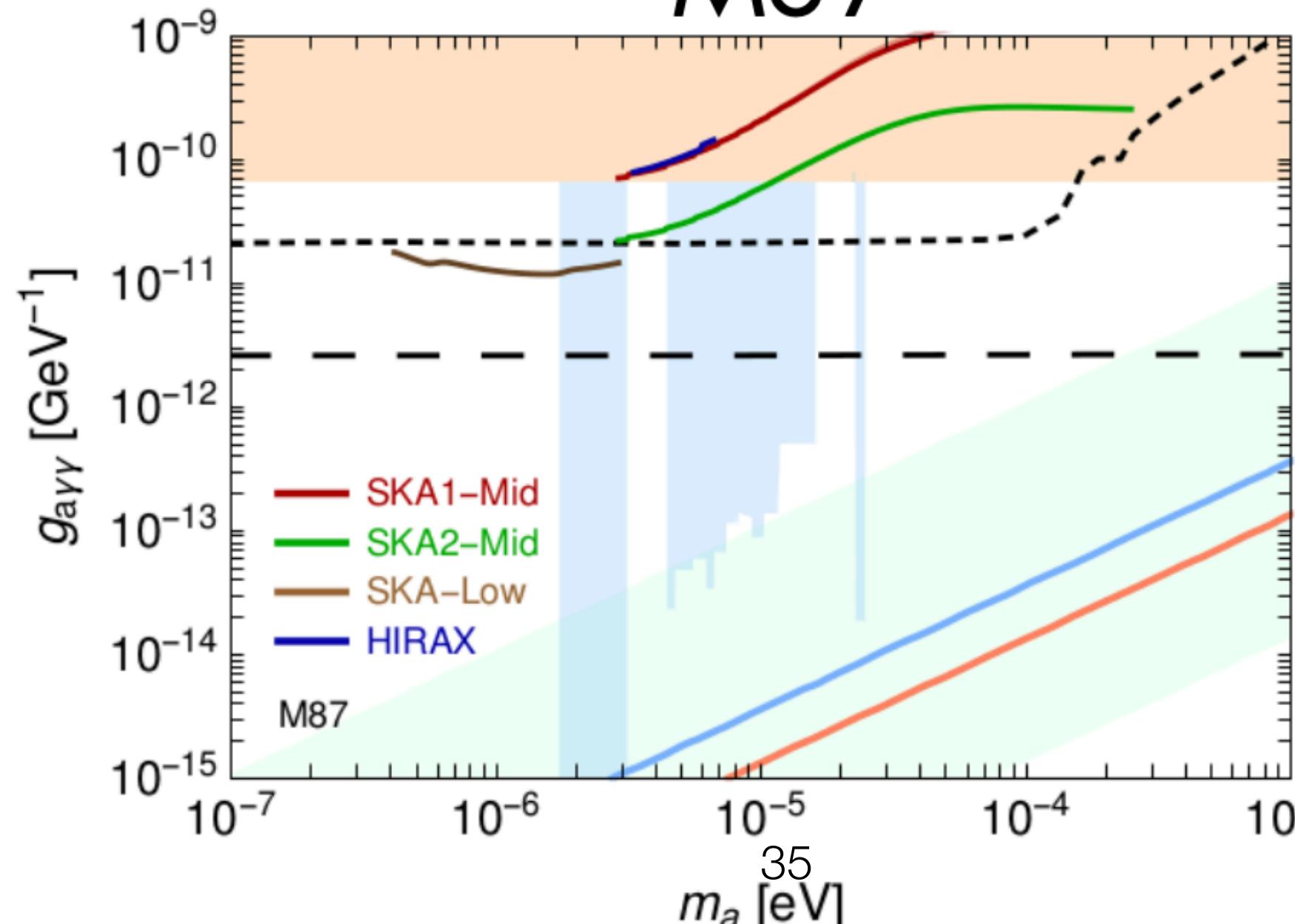
Dwarf spheroidal galaxy



Galactic Center



M87



Caputo, M.Regis, M. Taoso, S. Witte,
JCAP 03 (2019) 027



andrea0292 Create README.md

93e589e on Jan 3, 2021 3 commits

README.md

Create README.md

3 years ago

Try_DMmap.ipynb

Add files via upload

3 years ago

dm_map.py

Add files via upload

3 years ago

README.md

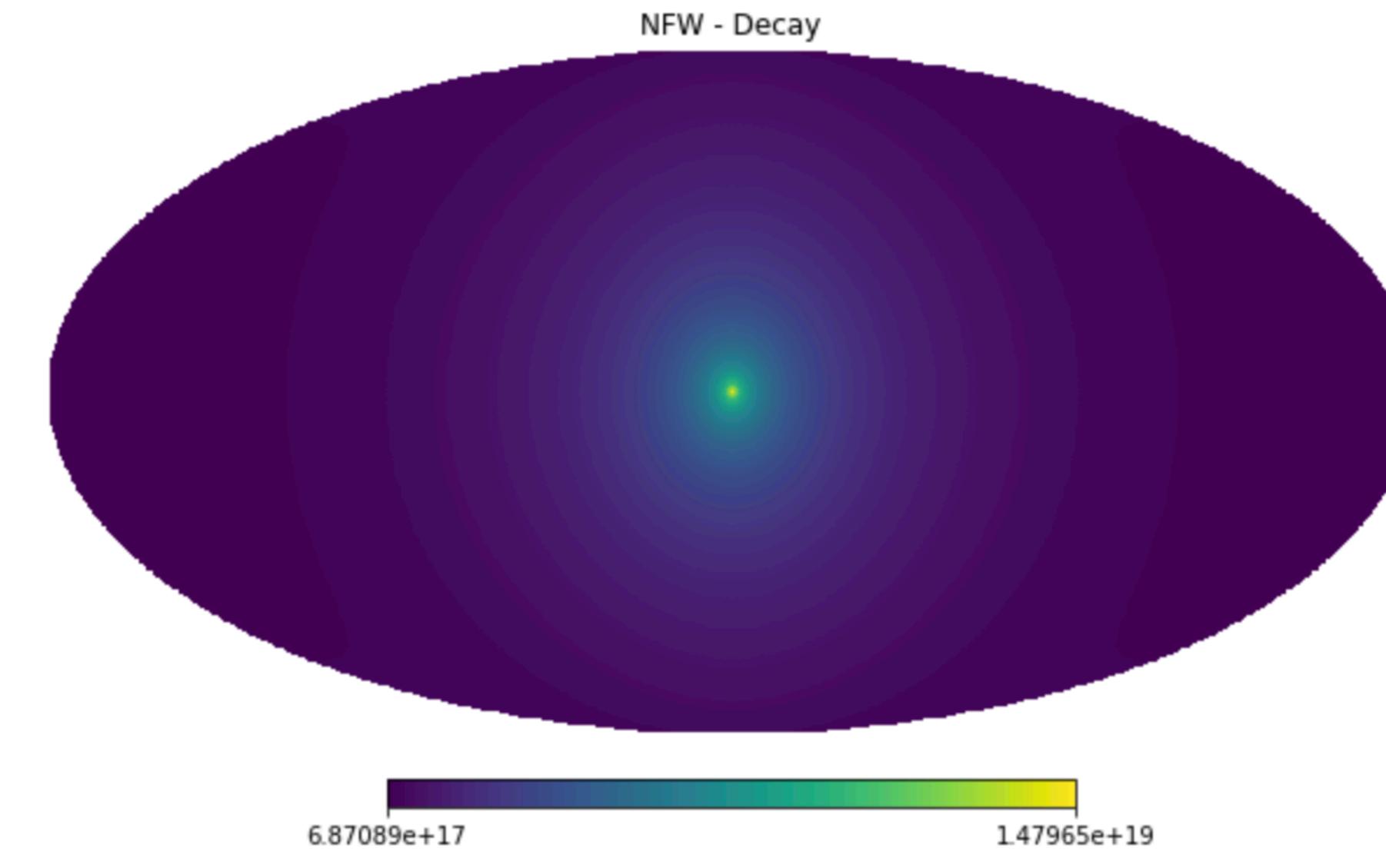


Dark-Matter-Map

I write down a simple Python code to plot with healpy the dark matter map for NFW and Burkert profiles. The .py file contains the class which does the job, while the Jupyter notebook gives an example on how to use it for the NFW generalized profile.

100% | 12582912/12582912 [5:15:47<00:00, 664.08it/s]

...done!



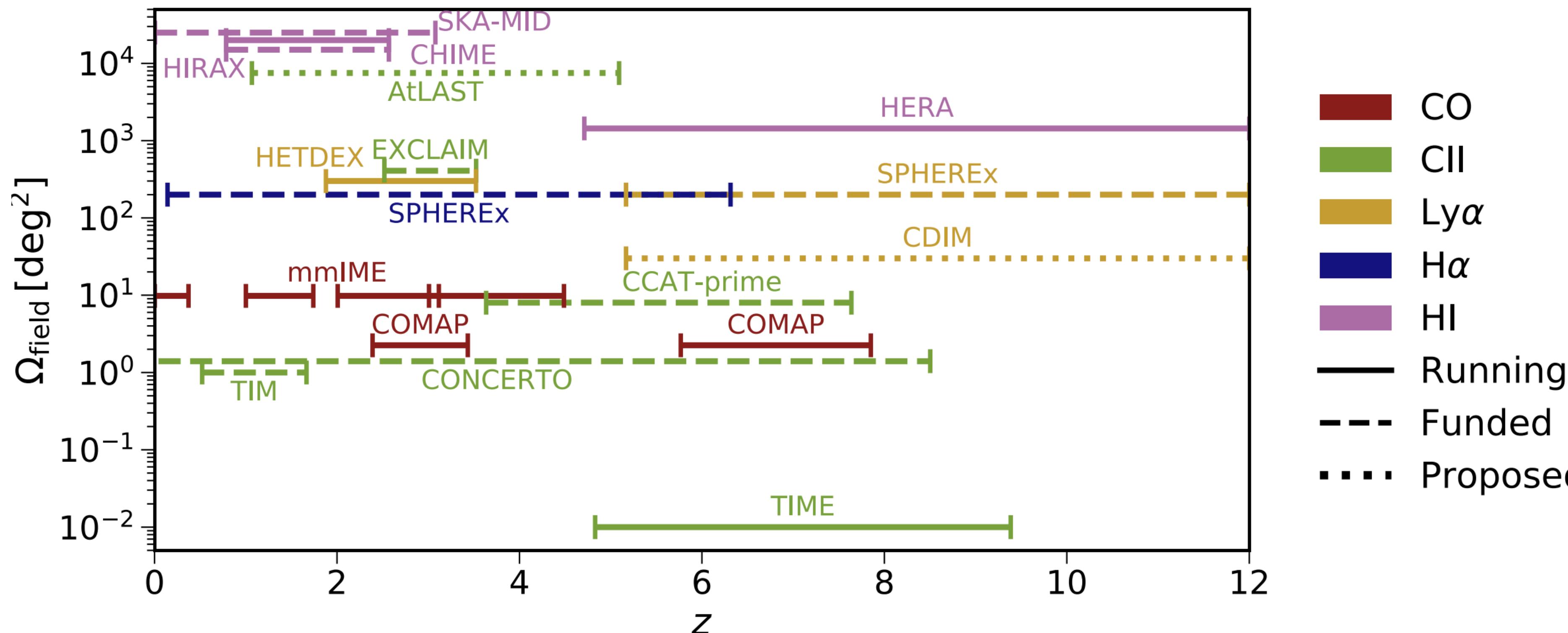
In this case we have targeted specific objects..however there is another possibility

Using the **integrated** signal from all galaxies in the past (**line intensity mapping**) with

$$E_{\text{obs}} = m_a/2/(1 + z_e)$$

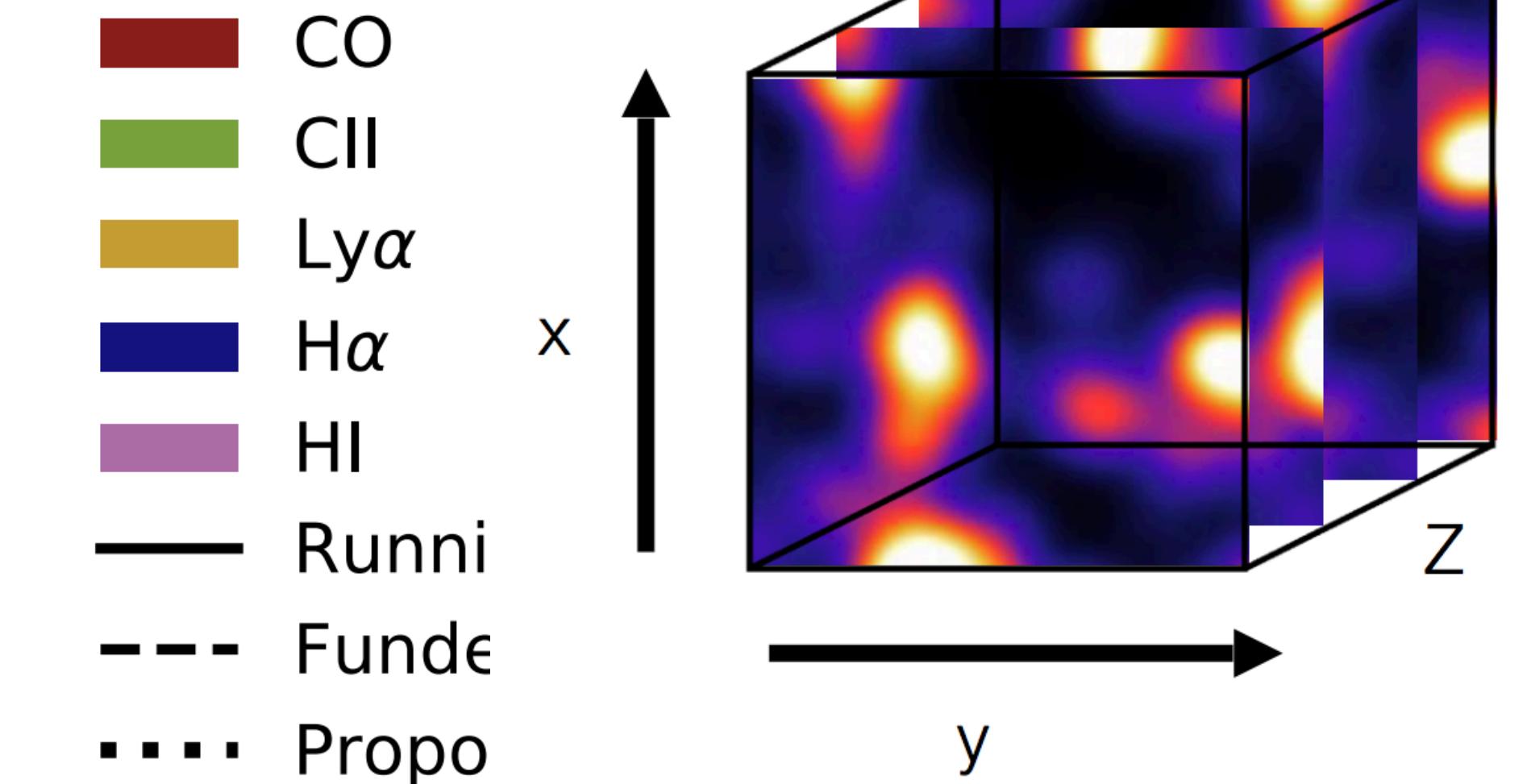
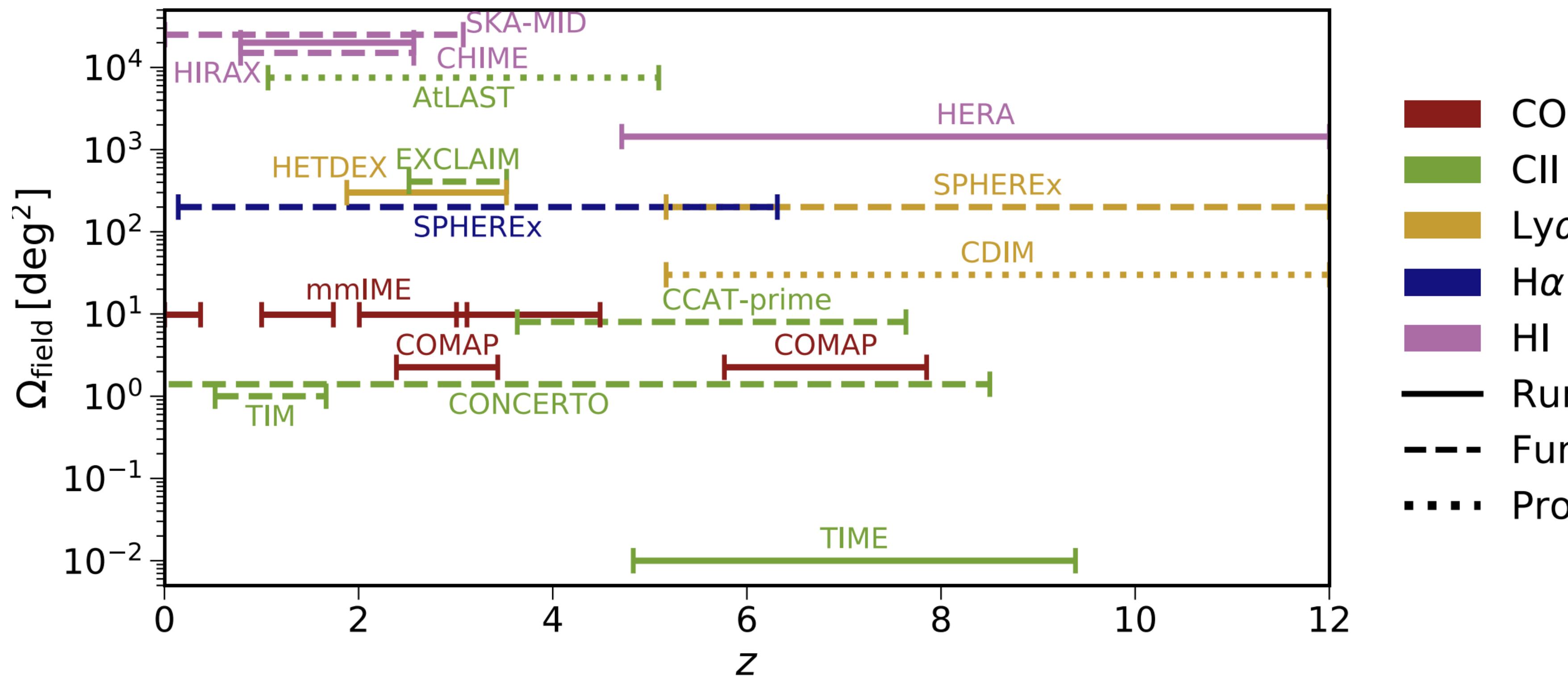
People already do this for atomic and molecular spectral lines!

$$\nu_{obs} = \nu_0 / (1 + z)$$



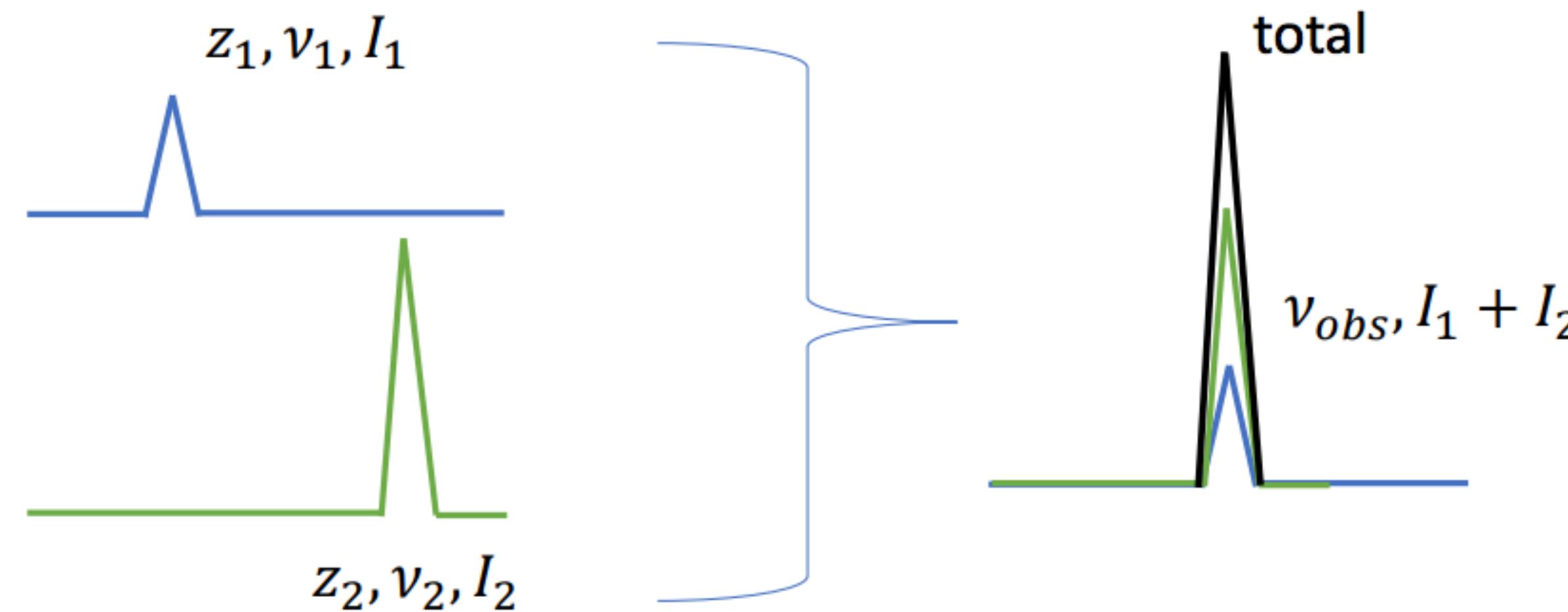
People already do this for atomic and molecular spectral lines!

$$\nu_{obs} = \nu_0 / (1 + z)$$



Contamination of intensity maps

- Continuous foregrounds: problem for HI surveys, less severe at higher frequencies
- **Line interlopers:** Main problem for higher freq. LIM surveys
 - $\nu_{obs} = \nu/(1+z) = \nu'/(1+z') \rightarrow$ other lines redshifted to same ν_{obs}



Contamination of intensity maps

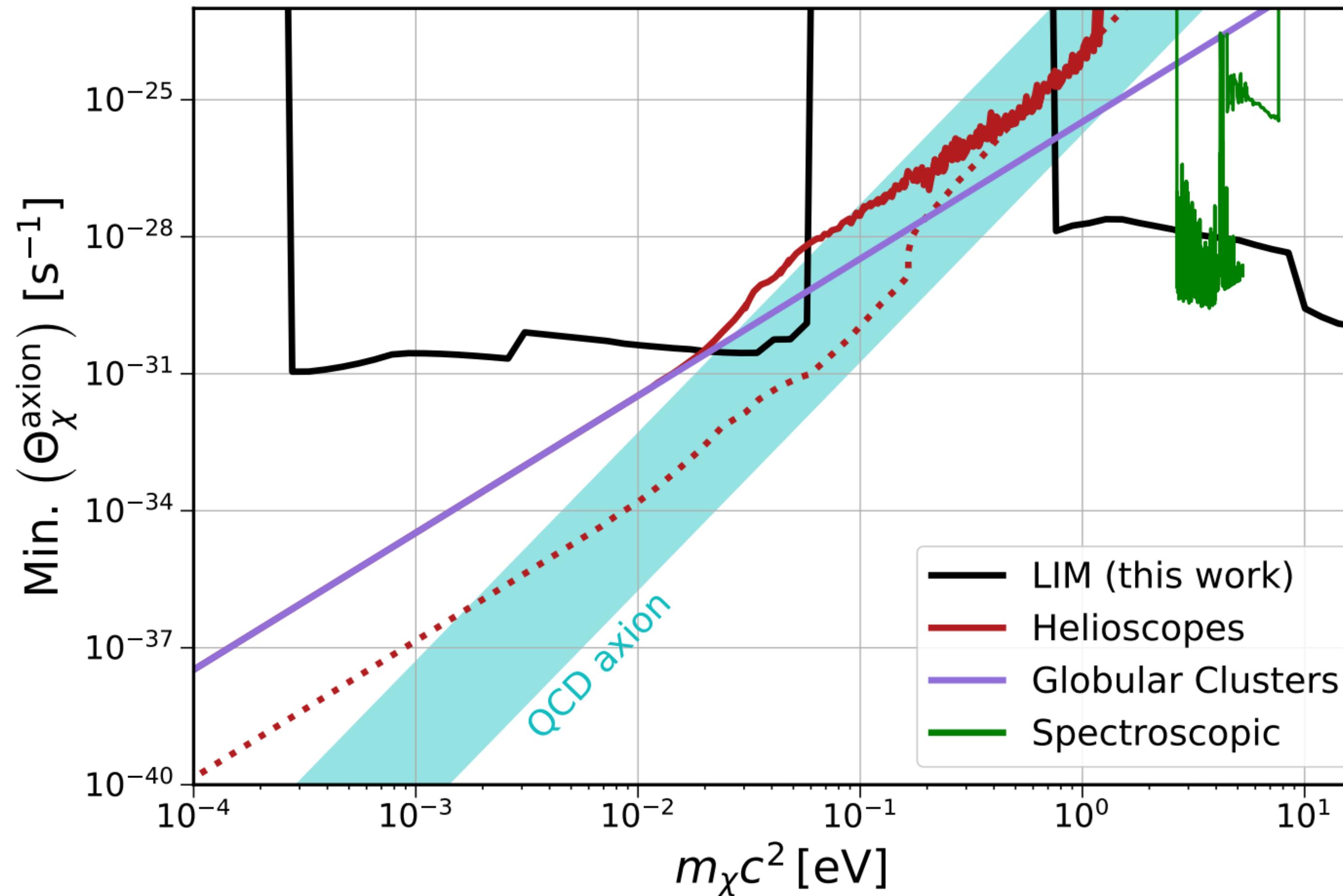
- Continuous foregrounds: problem for HI surveys, less severe at higher frequencies
- **Line interlopers:** Main problem for higher freq. LIM surveys
 - $\nu_{obs} = \nu/(1 + z) = \nu'/(1 + z') \rightarrow$ other lines redshifted to same ν_{obs}
 - Two approaches:
 - Masking: targeted (external data) and blind (contaminated voxels are expected to be brighter)
 - Model the effect of known interlopers in the likelihood and analyses

Exotic radiative decays would be inadvertently detected as a line interloper!!

Observables

- Clustering anisotropy parametrized by monopole, dipole, quadrupole, hexadecapole in angle wrt LOS
 - Clustering along line of sight
 - Angular clustering
- Voxel-intensity distribution (VID) (one-point PDF)

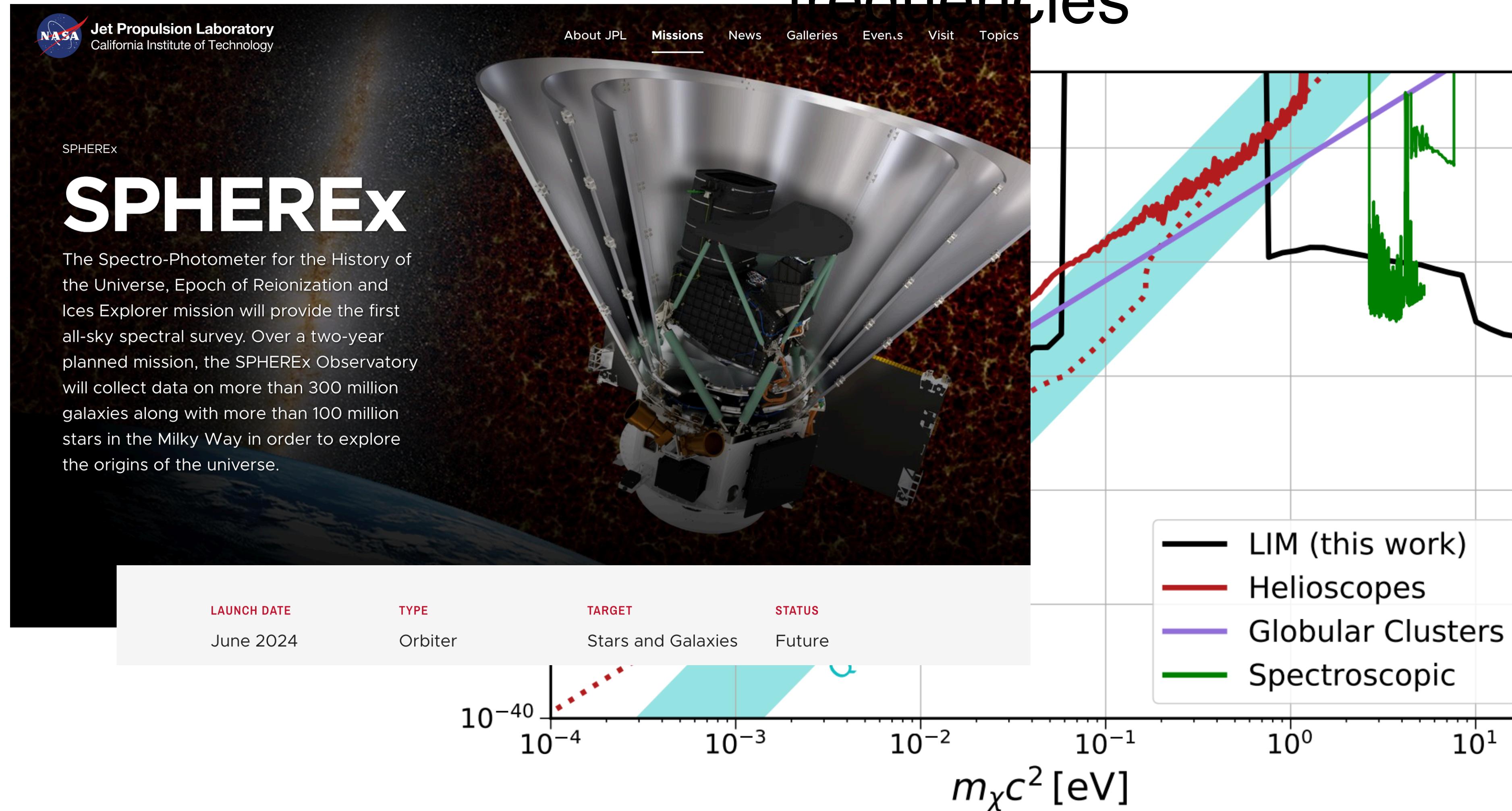
Line-intensity mapping is a promising avenue at different frequencies



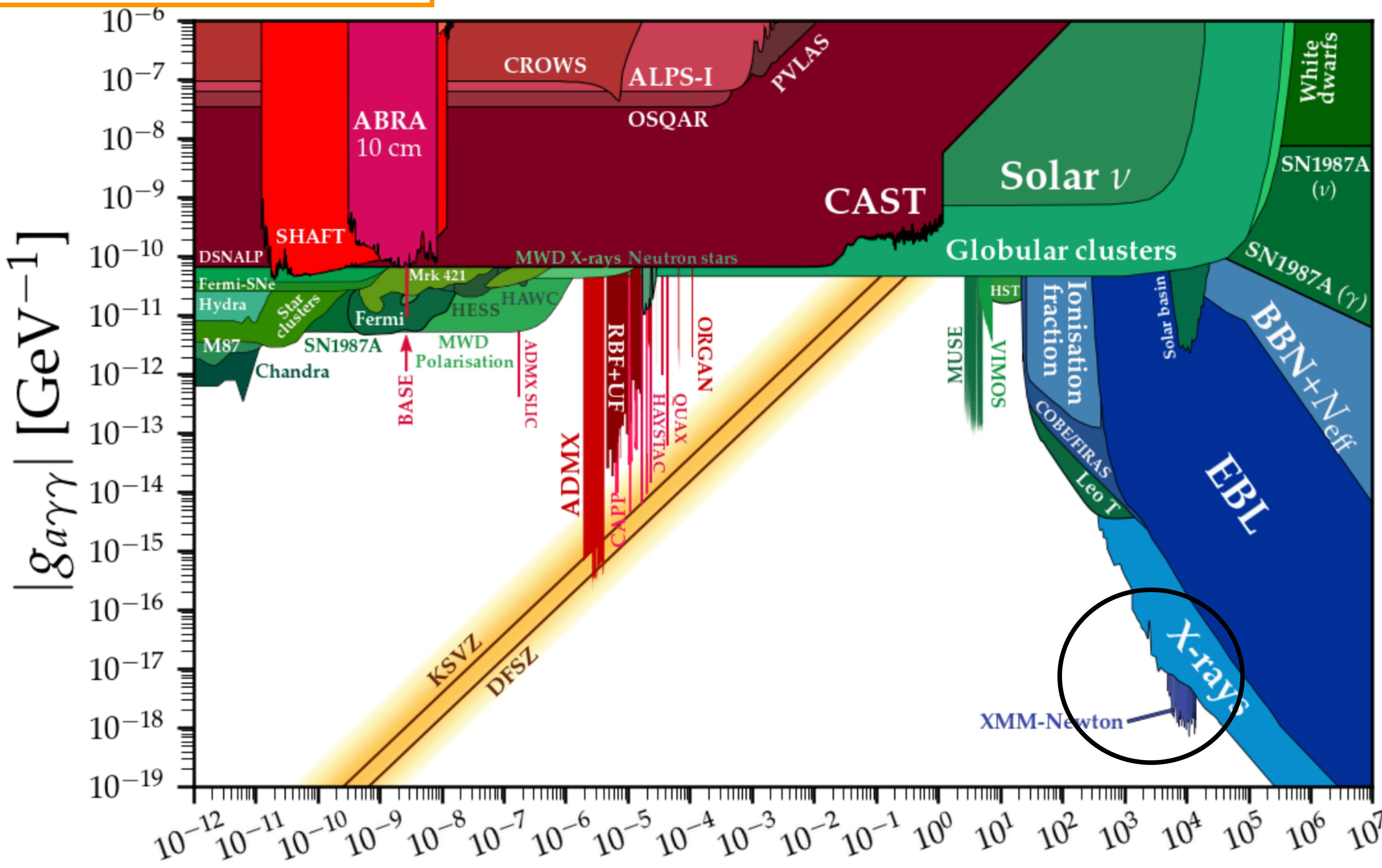
J. Bernal, Caputo, M. Kamionkowski, *Phys.Rev.D* 103 (2021) 6, 063523

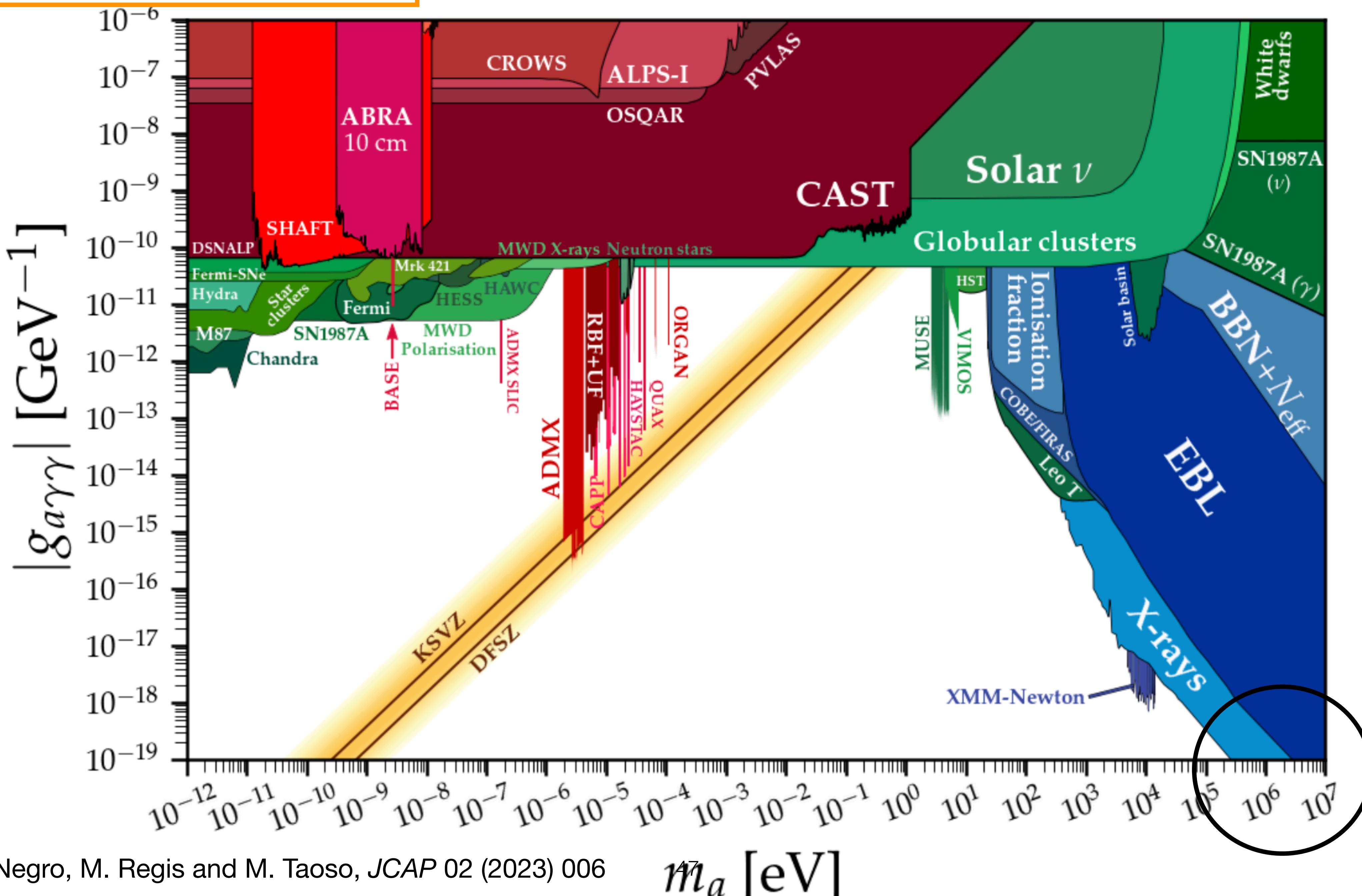
J. Bernal, Caputo, P. Villaescusa-Navarro, M. Kamionkowski, *Phys.Rev.Lett.* 127 (2021) 13, 131102

Line-intensity mapping is a promising avenue at different frequencies



J. Bernal, Caputo, M. Kamionkowski, *Phys.Rev.D* 103 (2021) 6, 063523
J. Bernal, Caputo, P. Villaescusa-Navarro, M⁴⁵Kamionkowski, *Phys.Rev.Lett.* 127 (2021) 13, 131102





Oct 18, 2021
RELEASE 21-134

NASA Selects Gamma-ray Telescope to Chart Milky Way Evolution

COSI-Telescope (approved by NASA)

Compton Spectrometer and Imager (COSI) wide-FOV telescope designed to survey the gamma-ray sky at 0.2-5 MeV → Imaging with high-resolution spectroscopy

[Caputo](#), M. Negro, M. Regis and M. Taoso, *JCAP* 02 (2023) 006

General strategy

A good story should answer few basic questions.

We considered **what** and **where**.

General strategy

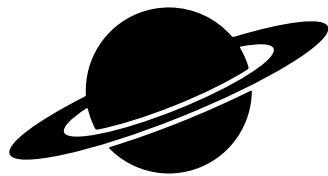
A good story should answer few basic questions.

We considered **what** and **where**. **When?** Very soon! All the missions I talked about will operate in the next 5-6 years.

II. Looking for new particles (not to be dark matter)



Stars and new physics



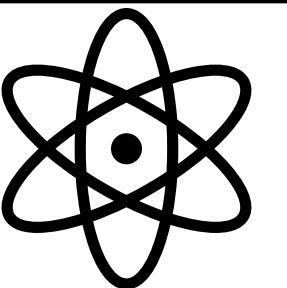
- Caputo et al., *JCAP* 08 (2022) 08, 045
Caputo et al., *Phys.Rev.Lett.* 128 (2022) 22, 221103 (CCSNe)
Caputo et al., *Phys.Rev.D* 105 (2022) 3, 035022 (CCSNe)
Caputo et al., *Phys.Rev.Lett.* 127 (2021) 18, 181102 (Hypernovae)
Caputo et al., *Phys.Rev.D* 101 (2020) 12, 123004 (Sun)
O'Hare, Caputo, et al., *Phys.Rev.D* 102 (2020) 4, 043019 (Sun)

Gravitational Waves

- Caputo et al., *J Astrophys.J.* 892 (2020) 2, 90
Toubiana, Sberna, Caputo, et al.,
Phys.Rev.Lett. 126 (2021) 10, 101105

My research

Dark Matter direct detection



- Caputo et al., *Phys.Rev.D* 100 (2019) 11, 116007 (sub-GeV DM)
Caputo et al., *Phys.Lett.B* 802 (2020) 135258 (sub-GeV DM)
Caputo et al., *Phys.Rev.D* 103 (2021) 5, 055017 (sub-GeV DM)
Bloch, Caputo et al., *JHEP* 01 (2021) 178 (sub-GeV DM)

Caputo et al., *Phys.Rev.D* 104 (2021) 9, 095029 (Dark Photon DM)

Dark Matter and telescopes

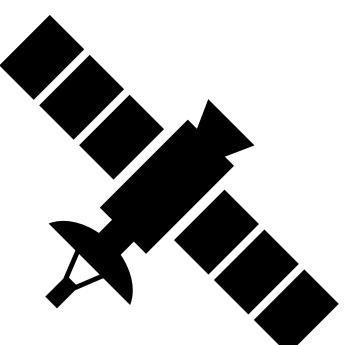


- Caputo et al., *JCAP* 03 (2019) 027 (radio)
Caputo et al., *Phys.Rev.D* 98 (2018) 8, 083024 (radio)
Caputo et al., *JCAP* 03 (2020) 001 (X-rays)
Caputo et al., *JCAP* 05 (2021) 046 (IR)
Bernal, Caputo, et al., *Phys.Rev.D* 103 (2021) 6, 063523 (optic)
Bernal, Caputo, et al., *Phys.Rev.Lett.* 127 (2021) 13, 131102 (optic)

Colliders phenomenology

- Caputo, P. Hernandez, M. Kekic, J. Lopez-Pavon, J. Salvado,
Eur.Phys.J.C 77 (2017) 4, 258
Caputo, P. Hernandez, J. Lopez-Pavon, J. Salvado, *JHEP* 06 (2017) 112
D. Barducci, E. Bertuzzo, Caputo, P. Hernandez, *JHEP* 06 (2020) 185
D. Barducci, E. Bertuzzo, Caputo, P. Hernandez, B. Mele,
JHEP 03 (2021) 117

Dark Matter and other indirect probes



- Caputo et al., *Phys.Rev.Lett.* 127 (2021) 1, 011102 (21 cm)
Bolton, Caputo, et al., *Phys.Rev.Lett.* 129 (2022) 21, 211102 (Ly-alpha)
Caputo et al., *Phys.Rev.Lett.* 125 (2020) 22, 221303 (CMB)
Bernal, Caputo, et al., *arXiv* 2208.13794 (blazars data)
Caputo et al., *Phys.Rev.D* 100 (2019) 6, 063515 (pulsars)
Caputo et al., *Phys.Dark Univ.* 19 (2018) 1-11 (pulsars)

Stars and new physics



- Caputo et al., *JCAP* 08 (2022) 08, 045
Caputo et al., *Phys.Rev.Lett.* 128 (2022) 22, 221103 (CCSNe)
Caputo et al., *Phys.Rev.D* 105 (2022) 3, 035022 (CCSNe)
Caputo et al., *Phys.Rev.Lett.* 127 (2021) 18, 181102 (Hypernovae)
Caputo et al., *Phys.Rev.D* 101 (2020) 12, 123004 (Sun)
O'Hare, Caputo, et al., *Phys.Rev.D* 102 (2020) 4, 043019 (Sun)

Dark Matter and telescopes



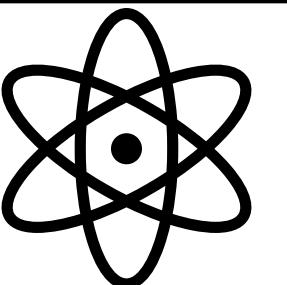
- Caputo et al., *JCAP* 03 (2019) 027 (radio)
Caputo et al., *Phys.Rev.D* 98 (2018) 8, 083024 (radio)
Caputo et al., *JCAP* 03 (2020) 001 (X-rays)
Caputo et al., *JCAP* 05 (2021) 046 (IR)

Gravitational Waves

- Caputo et al., *JAstrophys.J.* 892 (2020) 2, 90
Toubiana, Sberna, Caputo, et al.,
Phys.Rev.Lett. 126 (2021) 10, 101105

My research

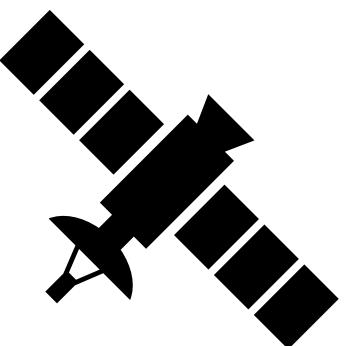
Dark Matter direct detection



- Caputo et al., *Phys.Rev.D* 100 (2019) 11, 116007 (sub-GeV DM)
Caputo et al., *Phys.Lett.B* 802 (2020) 135258 (sub-GeV DM)
Caputo et al., *Phys.Rev.D* 103 (2021) 5, 055017 (sub-GeV DM)
Bloch, Caputo et al., *JHEP* 01 (2021) 178 (sub-GeV DM)

Caputo et al., *Phys.Rev.D* 104 (2021) 9, 095029 (Dark Photon DM)

Dark Matter and other indirect probes



- Caputo et al., *Phys.Rev.Lett.* 127 (2021) 1, 011102 (21 cm)
Bolton, Caputo, et al., *Phys.Rev.Lett.* 129 (2022) 21, 211102 (Ly-alpha)
Caputo et al., *Phys.Rev.Lett.* 125 (2020) 22, 221303 (CMB)
Bernal, Caputo, et al., *arXiv* 2208.13794 (blazars data)
Caputo et al., *Phys.Rev.D* 100 (2019) 6, 063515 (pulsars)
Caputo et al., *Phys.Dark Univ.* 19 (2018) 1-11 (pulsars)

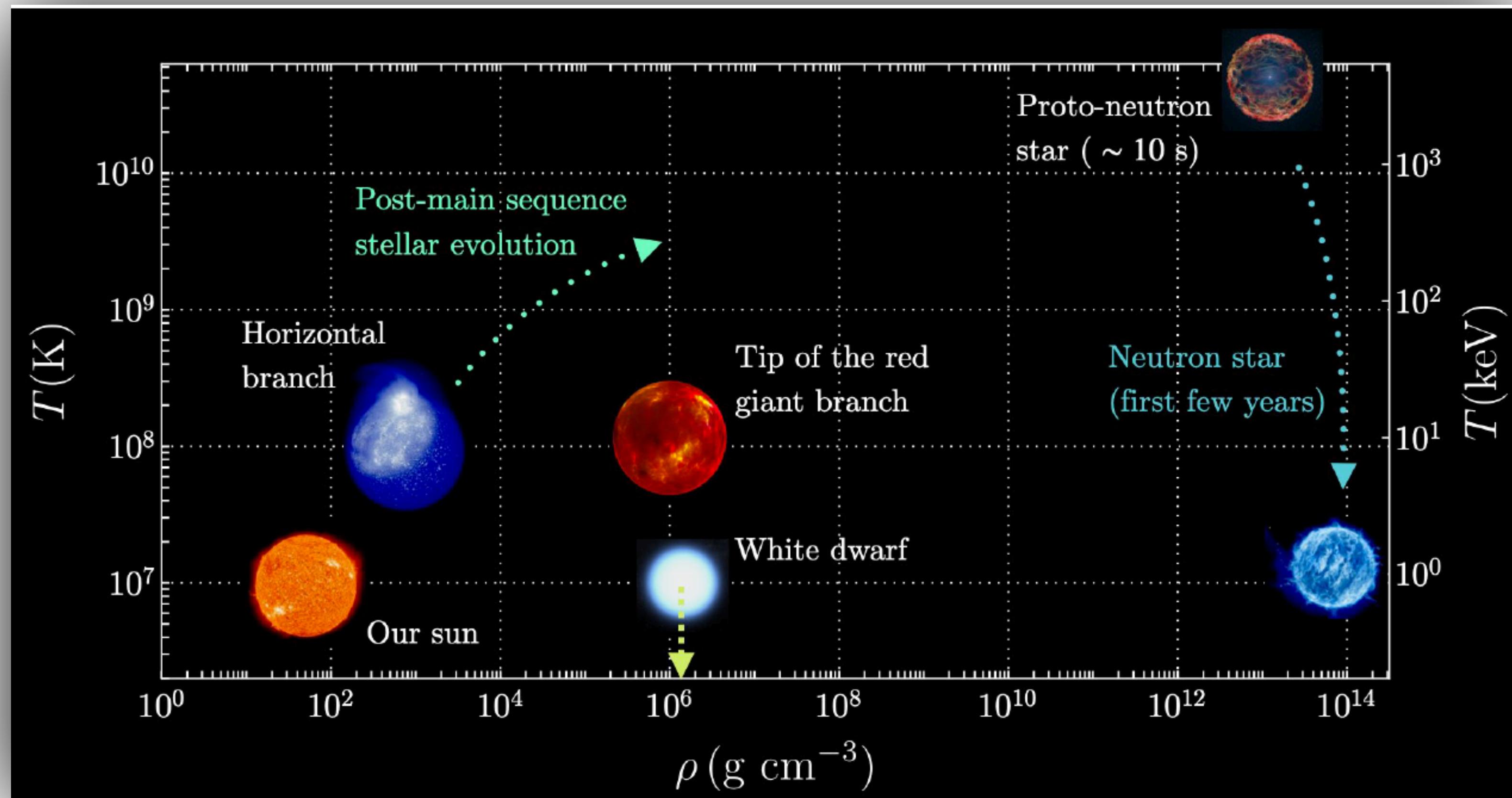
Now we will focus on the same axion coupling, but ask the question:
what can we do if ALPs are **not** dark matter?

Now we will focus on the same axion coupling, but ask the question:
what can we do if ALPs are **not** dark matter?

ALPs can modify stellar evolution!

Caputo, G. Raffel, E.Vitagliano, *Phys.Rev.D* 105 (2022) 3, 035022

Caputo, T.H. Janka, G. Raffel, E.Vitagliano, *Phys.Rev.Lett.* 128 (2022) 22, 221103

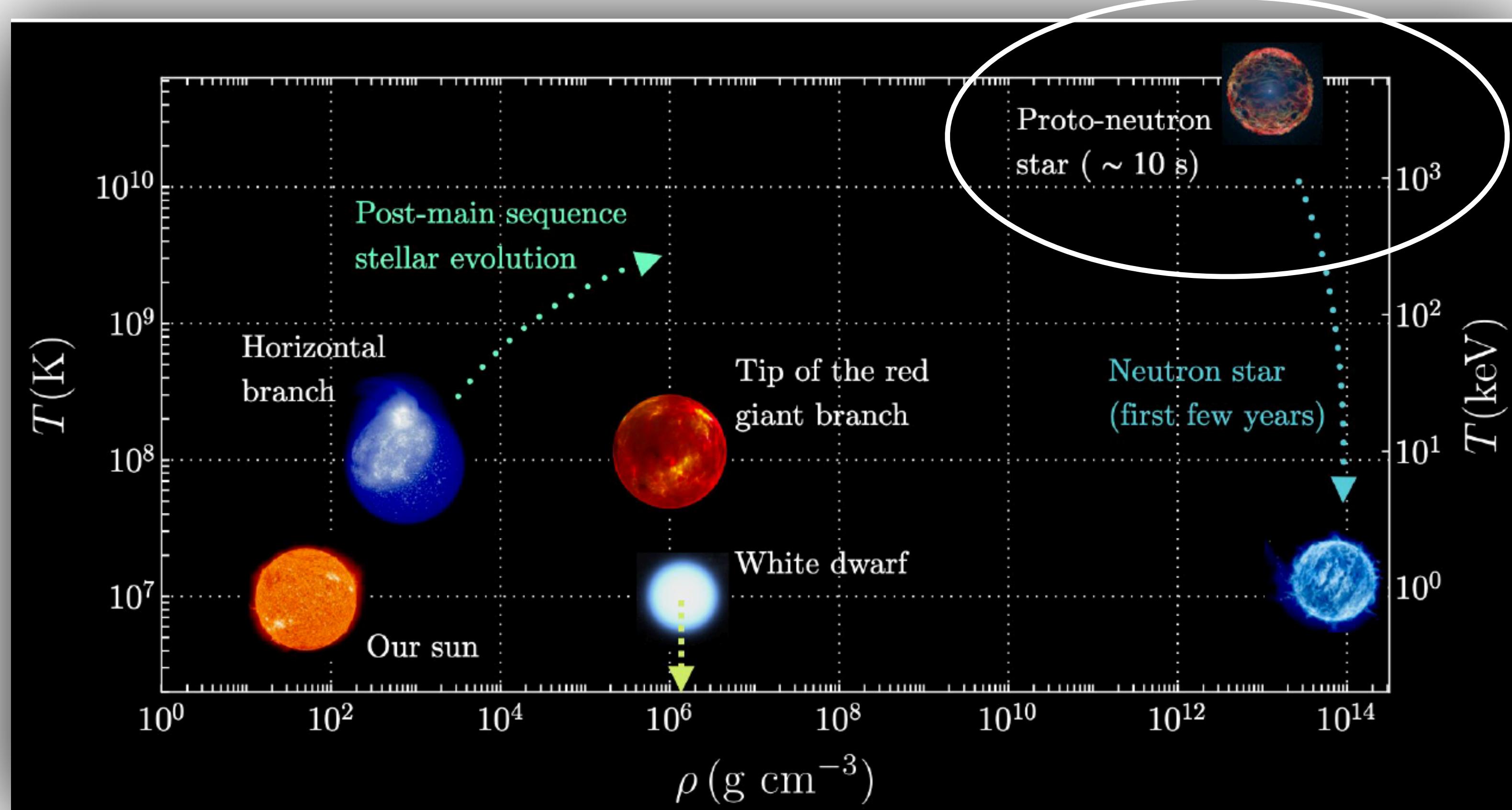


Now we will focus on the same axion coupling, but ask the question:
what can we do if ALPs are not dark matter?

ALPs can modify stellar evolution!

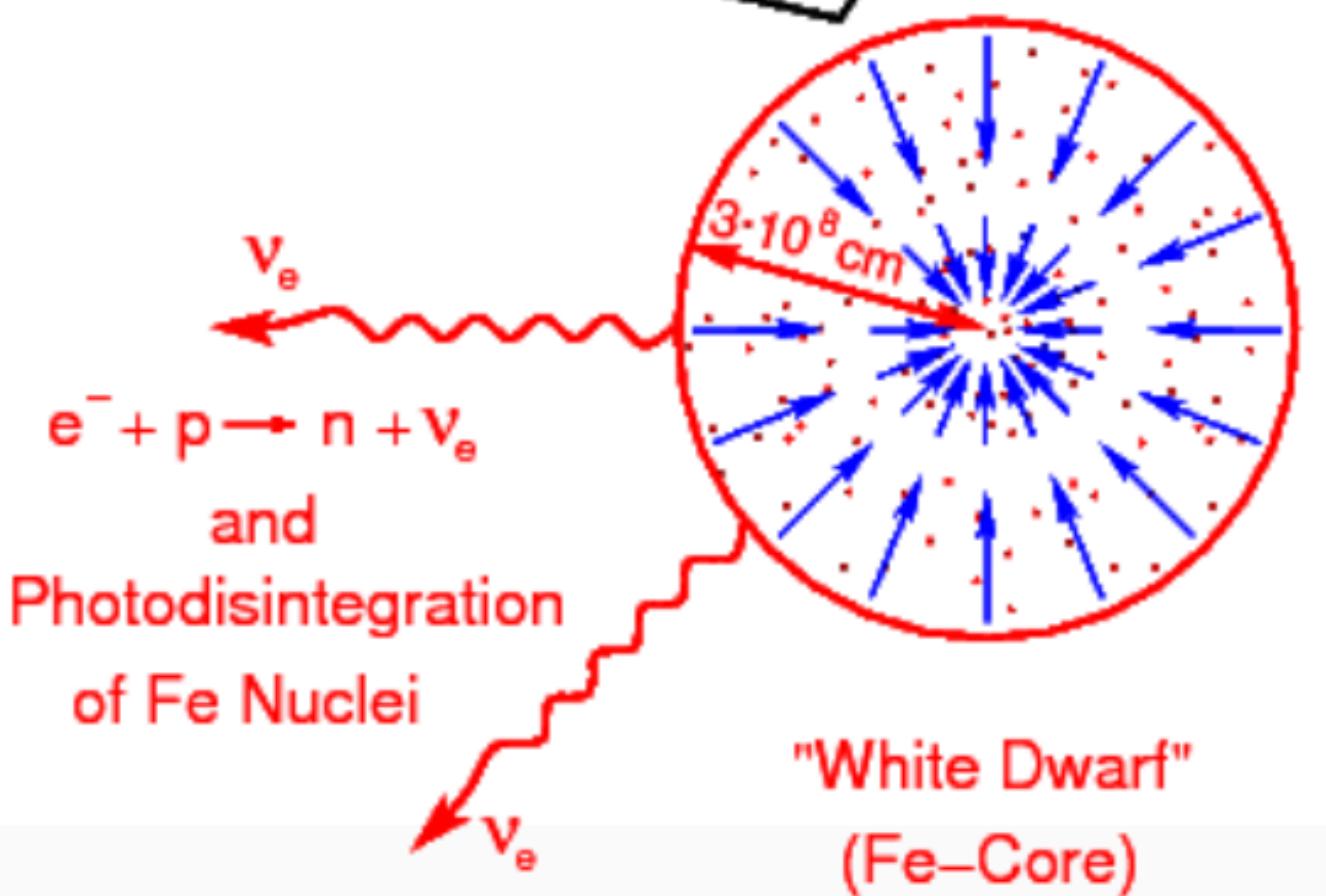
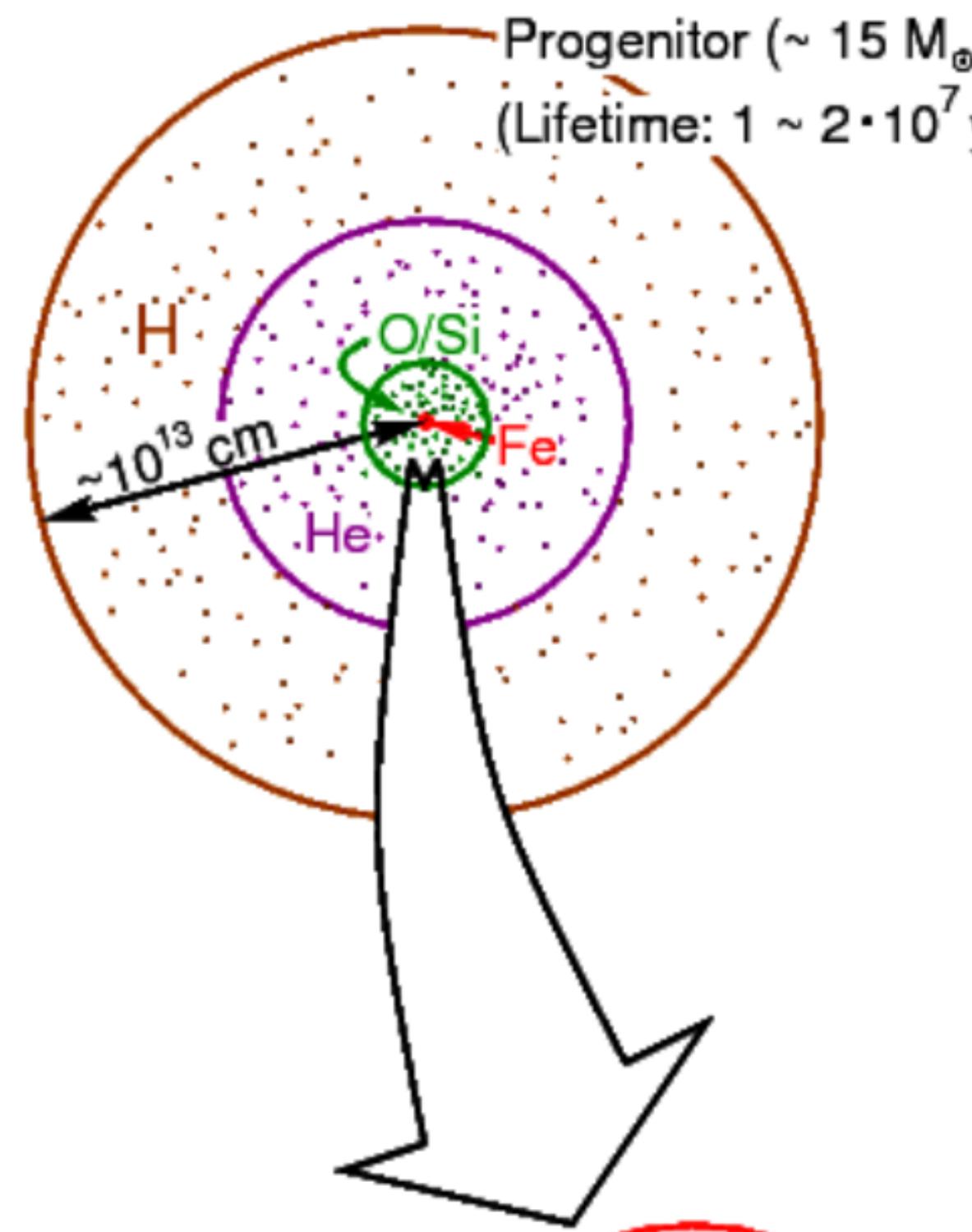
Caputo, G. Raffel, E.Vitagliano, *Phys.Rev.D* 105 (2022) 3, 035022

Caputo, T.H. Janka, G. Raffel, E.Vitagliano, *Phys.Rev.Lett.* 128 (2022) 22, 221103



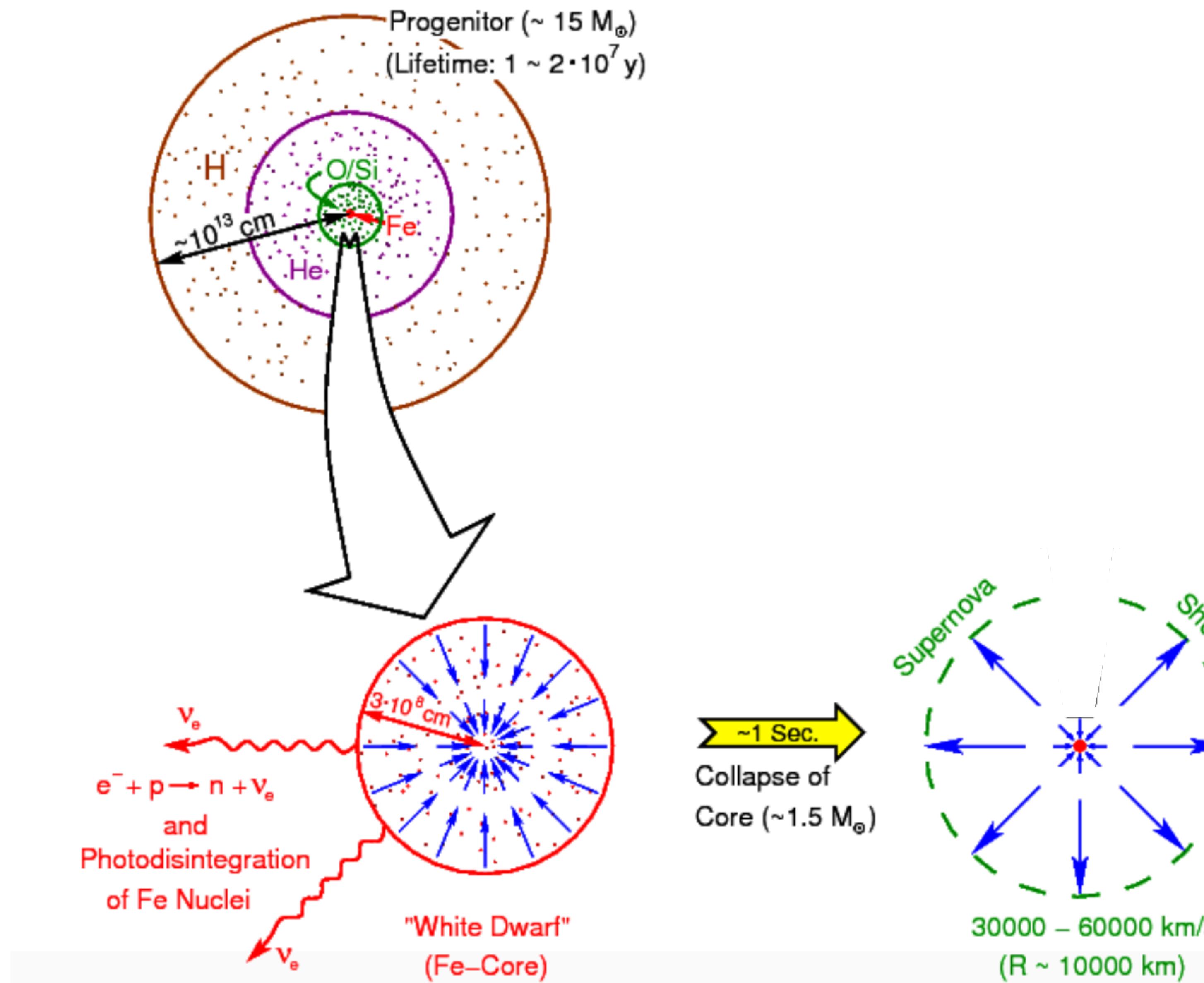
Crash-course on Supernovae (SNe)

A little bit of SN Physics



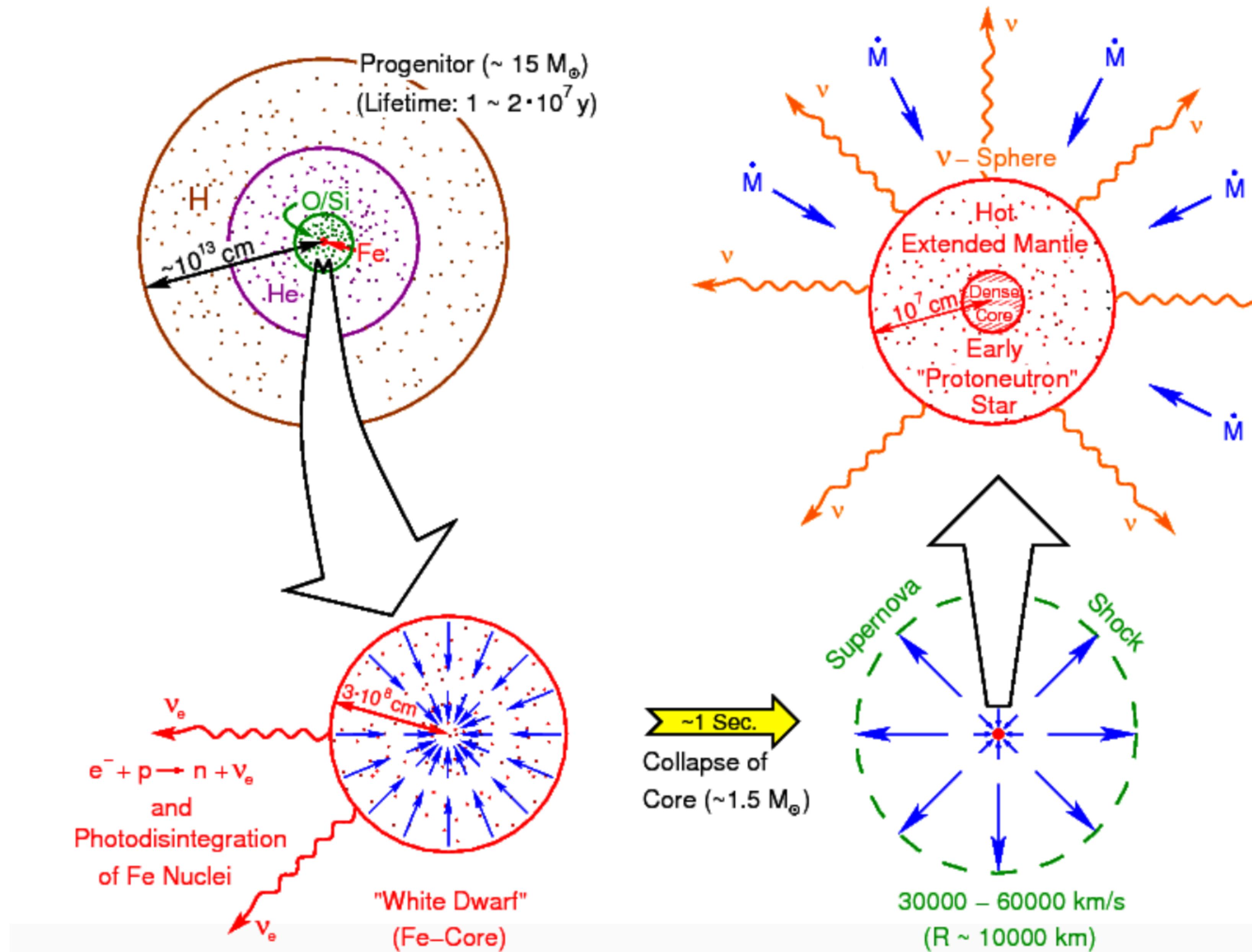
See T. Janka Cern Colloquium
<https://indico.cern.ch/event/1037035/>

A little bit of SN Physics



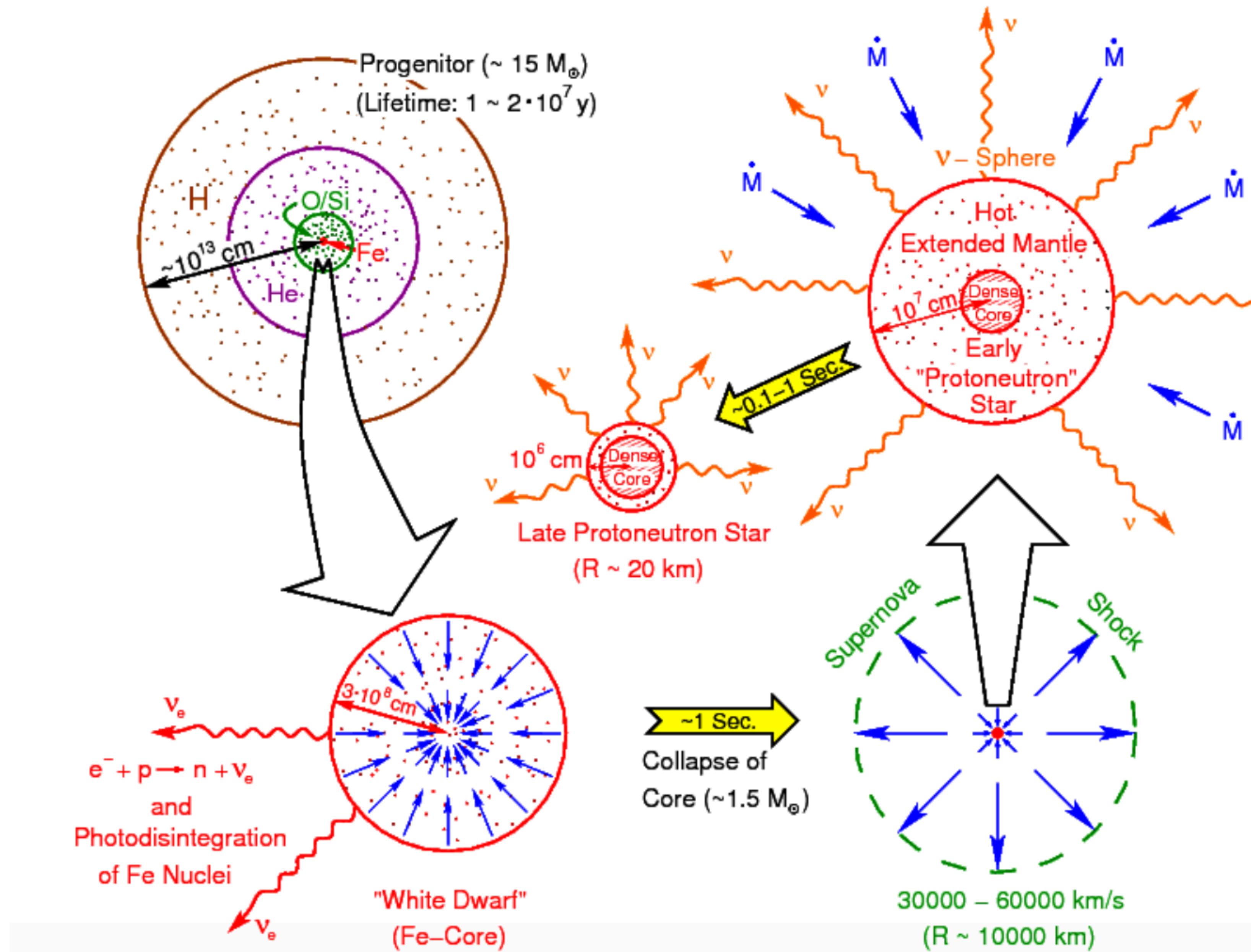
See T. Janka Cern Colloquium
<https://indico.cern.ch/event/1037035/>

A little bit of SN Physics



See T. Janka Cern Colloquium
<https://indico.cern.ch/event/1037035/>

A little bit of SN Physics



See T. Janka CERN Colloquium
<https://indico.cern.ch/event/1037035/>

PNS: protoneutron star

- $r \approx 10 \text{ km}$
- $T \approx 40 \text{ MeV}$
- Nuclear density

Neutrino emitted from the “neutrino sphere”, cooling the PNS in $\sim 10\text{s}$



PNS: protoneutron star

- $r \approx 10 \text{ km}$
- $T \approx 40 \text{ MeV}$
- Nuclear density

Neutrino emitted from the “neutrino sphere”, cooling the PNS in $\sim 10\text{s}$

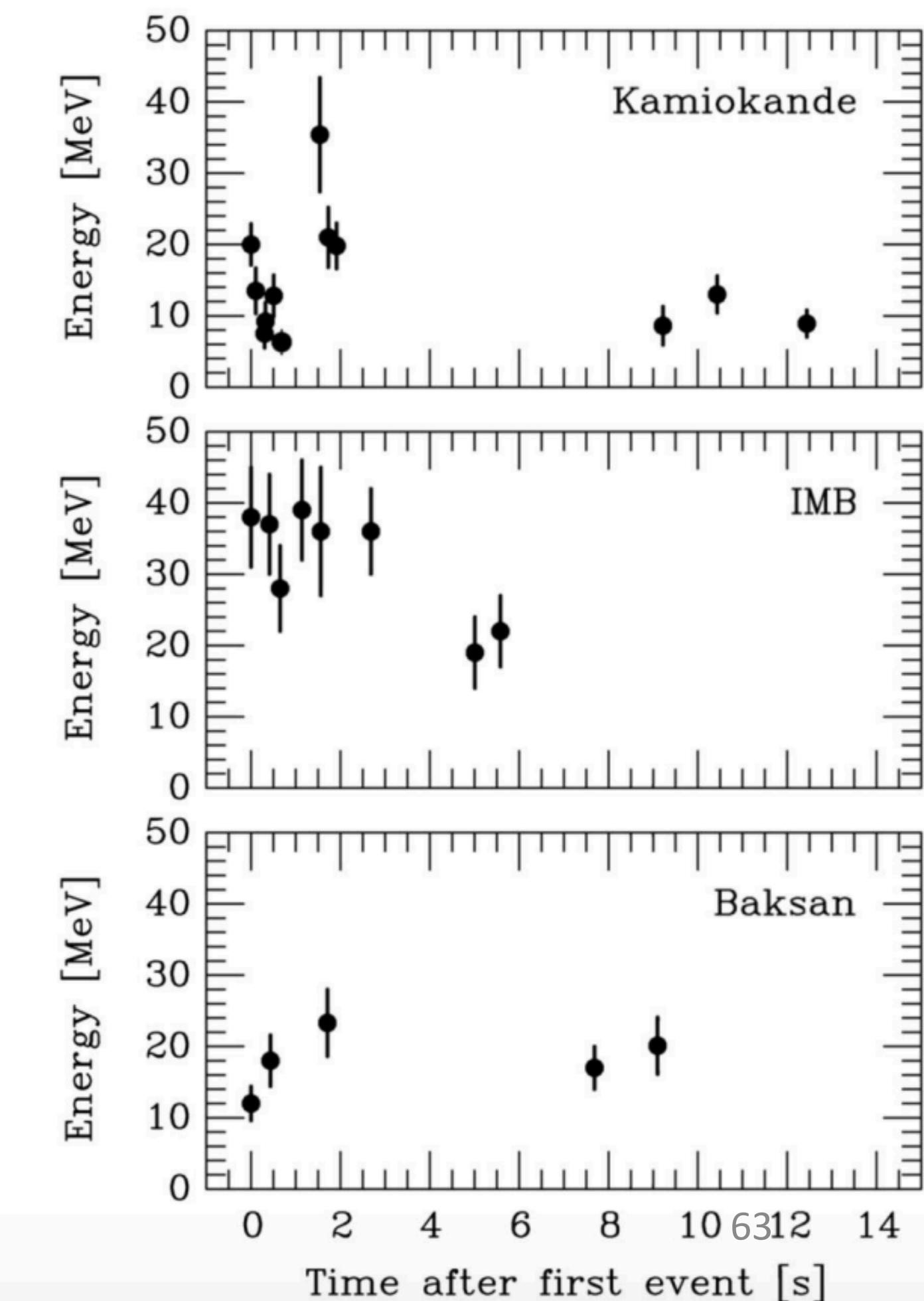
$$E_b \approx \frac{3}{5} \frac{G_N M^2}{R} = 1.60 \times 10^{53} \text{ erg} \left(\frac{M}{M_\odot} \right)^2 \left(\frac{10 \text{ km}}{R} \right)$$

$$T = \frac{2}{3} \langle E_{\text{kin}} \rangle \approx 17 \text{ MeV}$$

$$t_{\text{diff}} \approx R^2 / \lambda$$



SN 1987A neutrino signal



PNS: protoneutron star

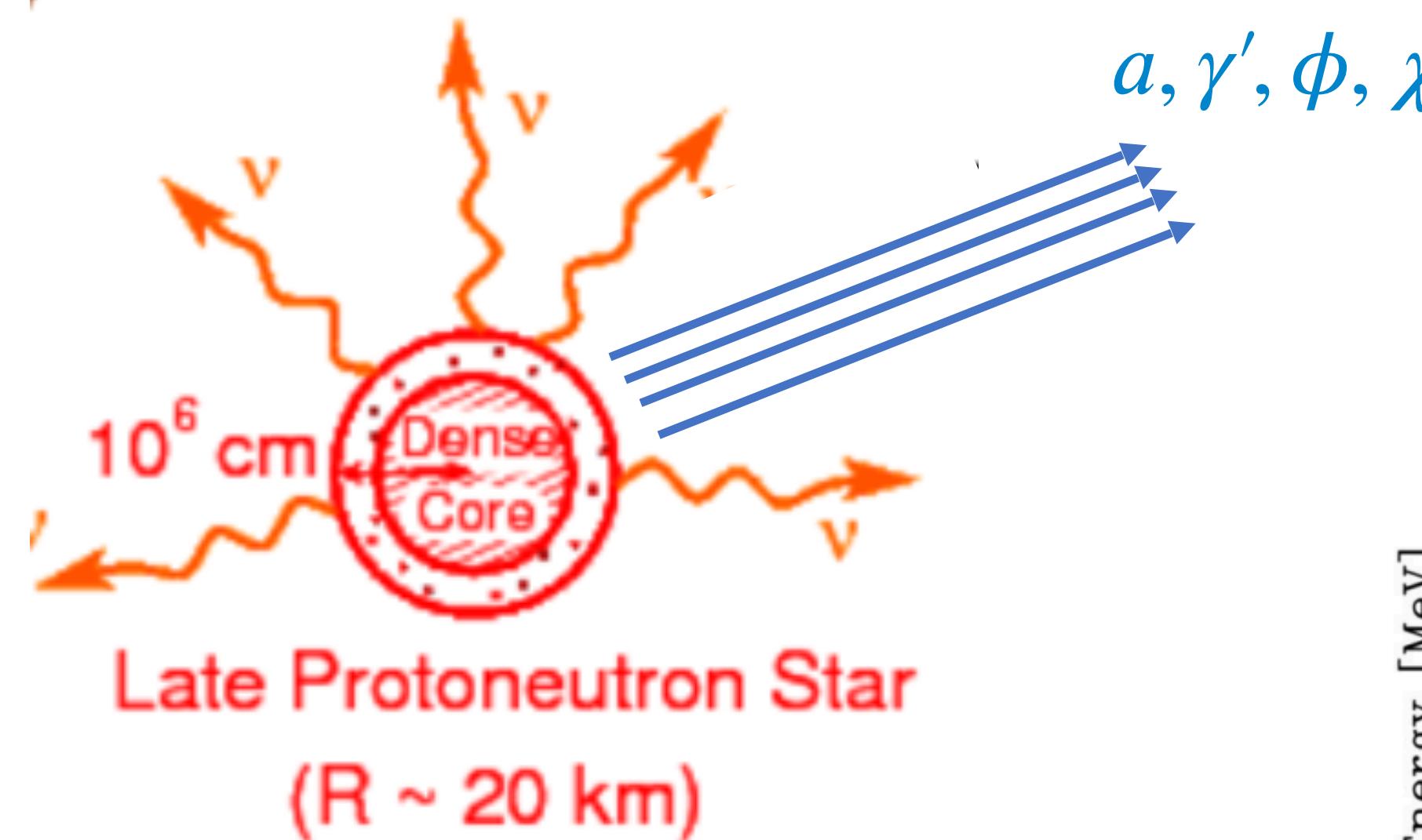
- $r \approx 10 \text{ km}$
- $T \approx 40 \text{ MeV}$
- Nuclear density

Neutrino emitted from the “neutrino sphere”, cooling the PNS in $\sim 10\text{s}$

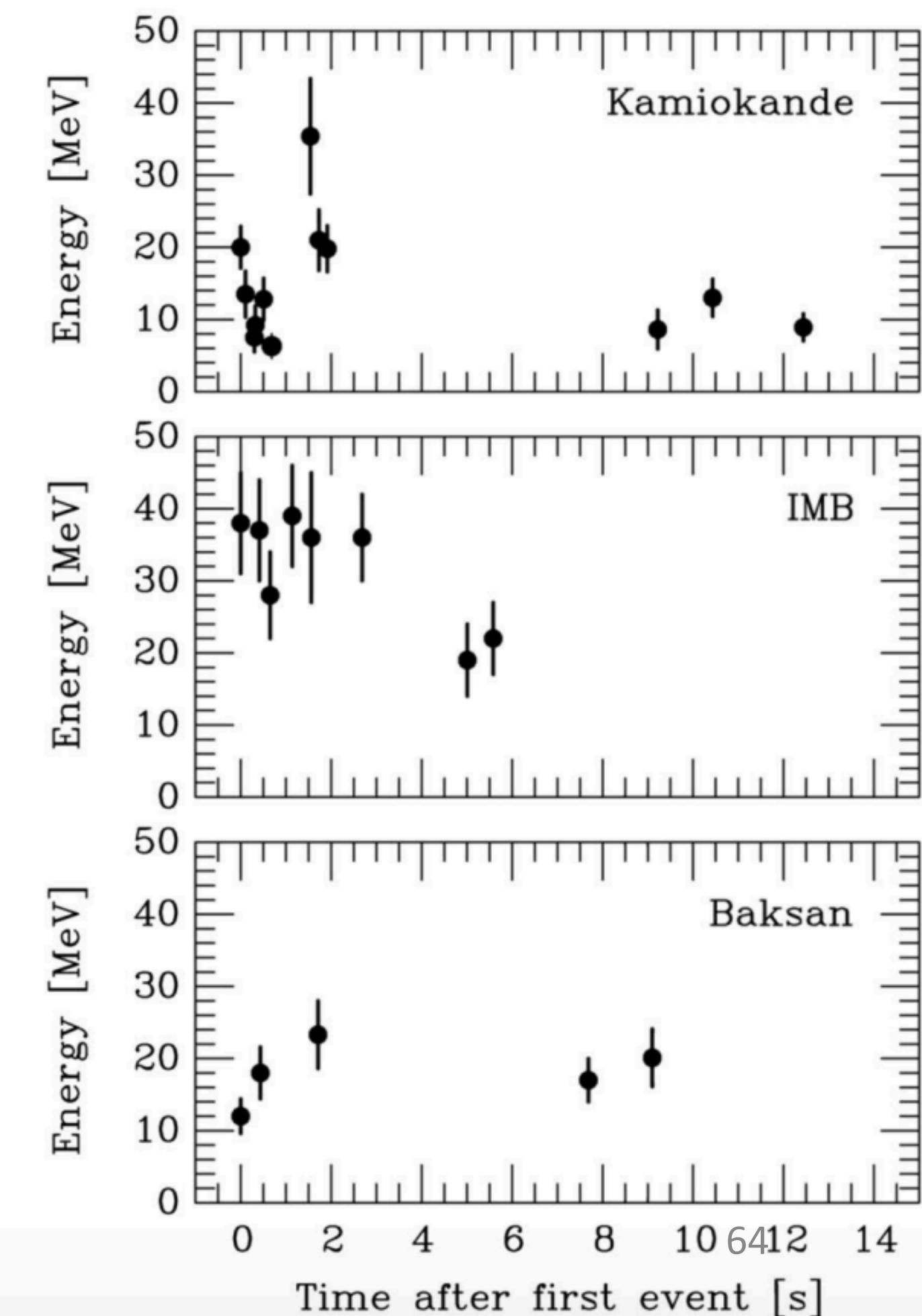
$$E_b \approx \frac{3}{5} \frac{G_N M^2}{R} = 1.60 \times 10^{53} \text{ erg} \left(\frac{M}{M_\odot} \right)^2 \left(\frac{10 \text{ km}}{R} \right)$$

$$T = \frac{2}{3} \langle E_{\text{kin}} \rangle \approx 17 \text{ MeV}$$

$$t_{\text{diff}} \approx R^2 / \lambda$$



SN 1987A neutrino signal



PNS: protoneutron star

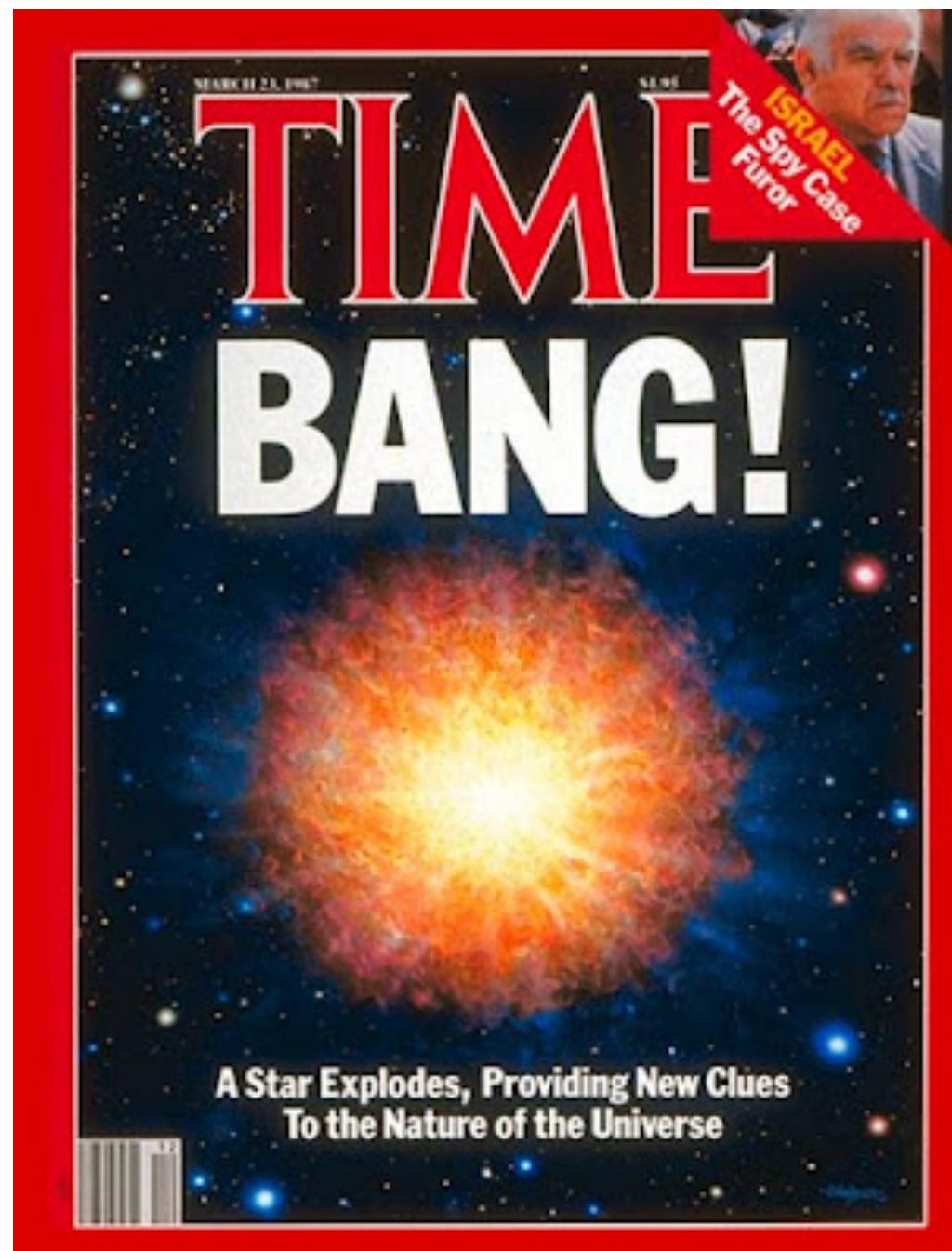
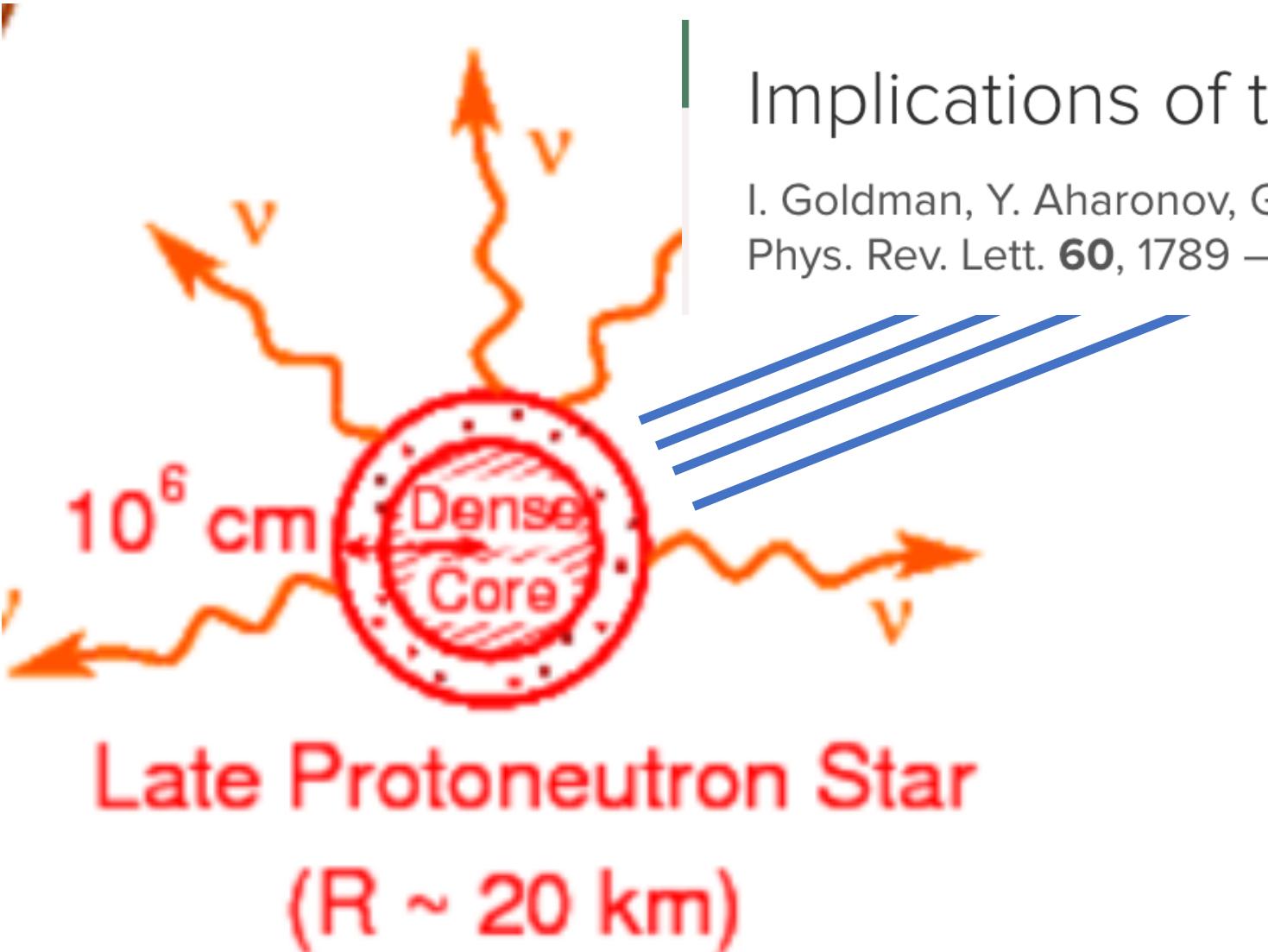
- $r \approx 10 \text{ km}$
- $T \approx 40 \text{ MeV}$
- Nuclear density

Neutrino emitted from the “neutrino sphere”, cooling the PNS in $\sim 10\text{s}$

$$E_b \approx \frac{3}{5} \frac{G_N M^2}{R} = 1.60 \times 10^{53} \text{ erg} \left(\frac{M}{M_\odot} \right)^2 \left(\frac{10 \text{ km}}{R} \right)$$

$$T = \frac{2}{3} \langle E_{\text{kin}} \rangle \approx 17 \text{ MeV}$$

$$t_{\text{diff}} \approx R^2 / \lambda$$

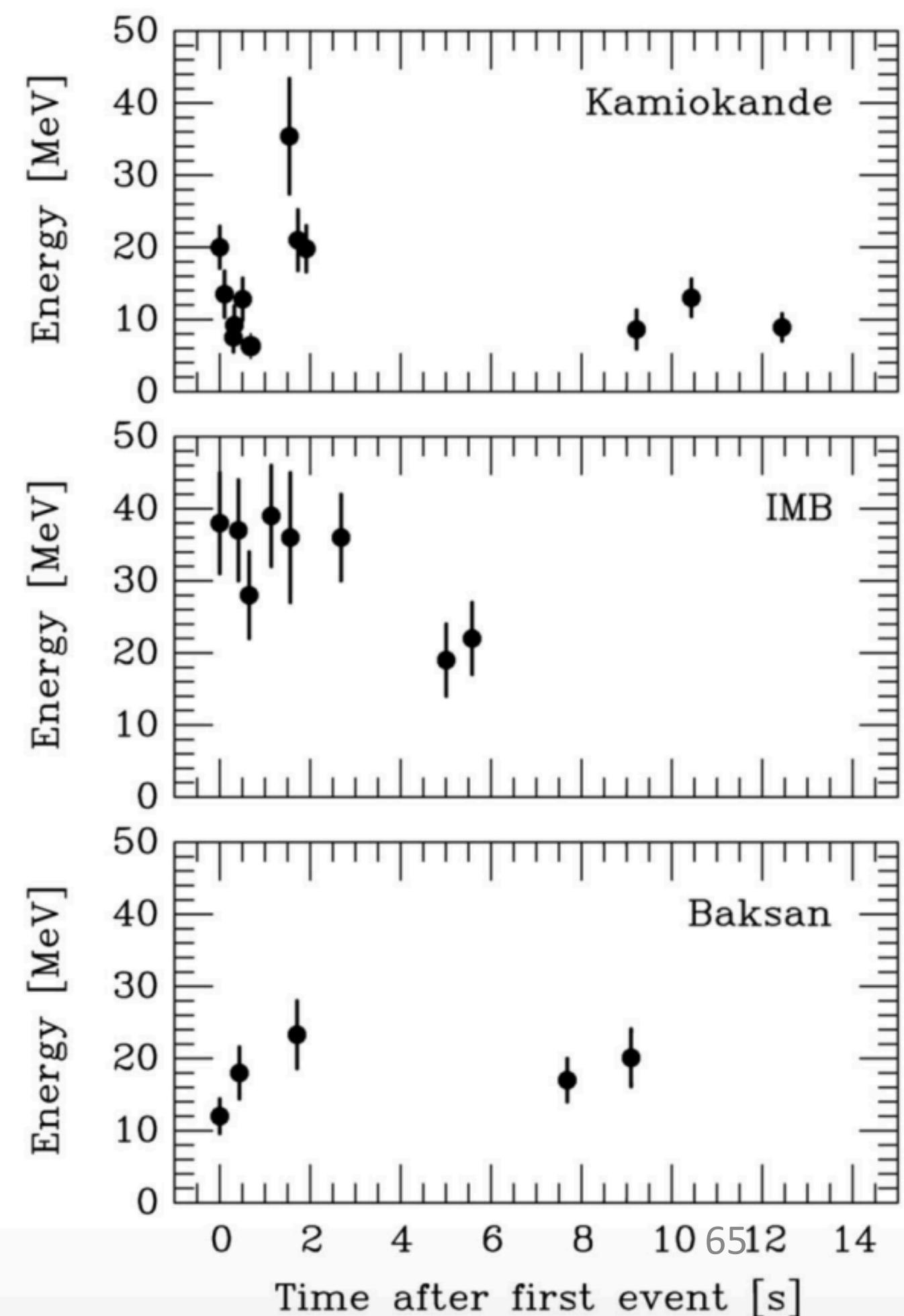


Riccardo Barbieri and Rabindra N. Mohapatra
Phys. Rev. D **39**, 1229 – Published 15 February 1989

Implications of the supernova SN1987A neutrino signals

I. Goldman, Y. Aharonov, G. Alexander, and S. Nussinov
Phys. Rev. Lett. **60**, 1789 – Published 2 May 1988

SN 1987A neutrino signal



Example model: axion coupling to photons

$$\mathcal{L}_{a\gamma\gamma} = -\frac{1}{4}g_{a\gamma}a\tilde{F}^{\mu\nu}F_{\mu\nu}$$

Example model: axion coupling to photons

$$\mathcal{L}_{a\gamma\gamma} = -\frac{1}{4}g_{a\gamma}a\tilde{F}^{\mu\nu}F_{\mu\nu}$$

$$Q_a = \int \frac{2d^3k}{(2\pi)^3} \frac{\omega}{e^{\omega/T} - 1} \hat{n} \sigma_P$$

$$= \hat{n} \alpha G_{a\gamma\gamma}^2 \frac{\pi^2 T^4}{30} \langle f_P \rangle, \quad \langle f_P \rangle = 20 \frac{m_a^2 + 3m_a T + 3T^2}{\pi^4 T^2} e^{-m_a/T}$$

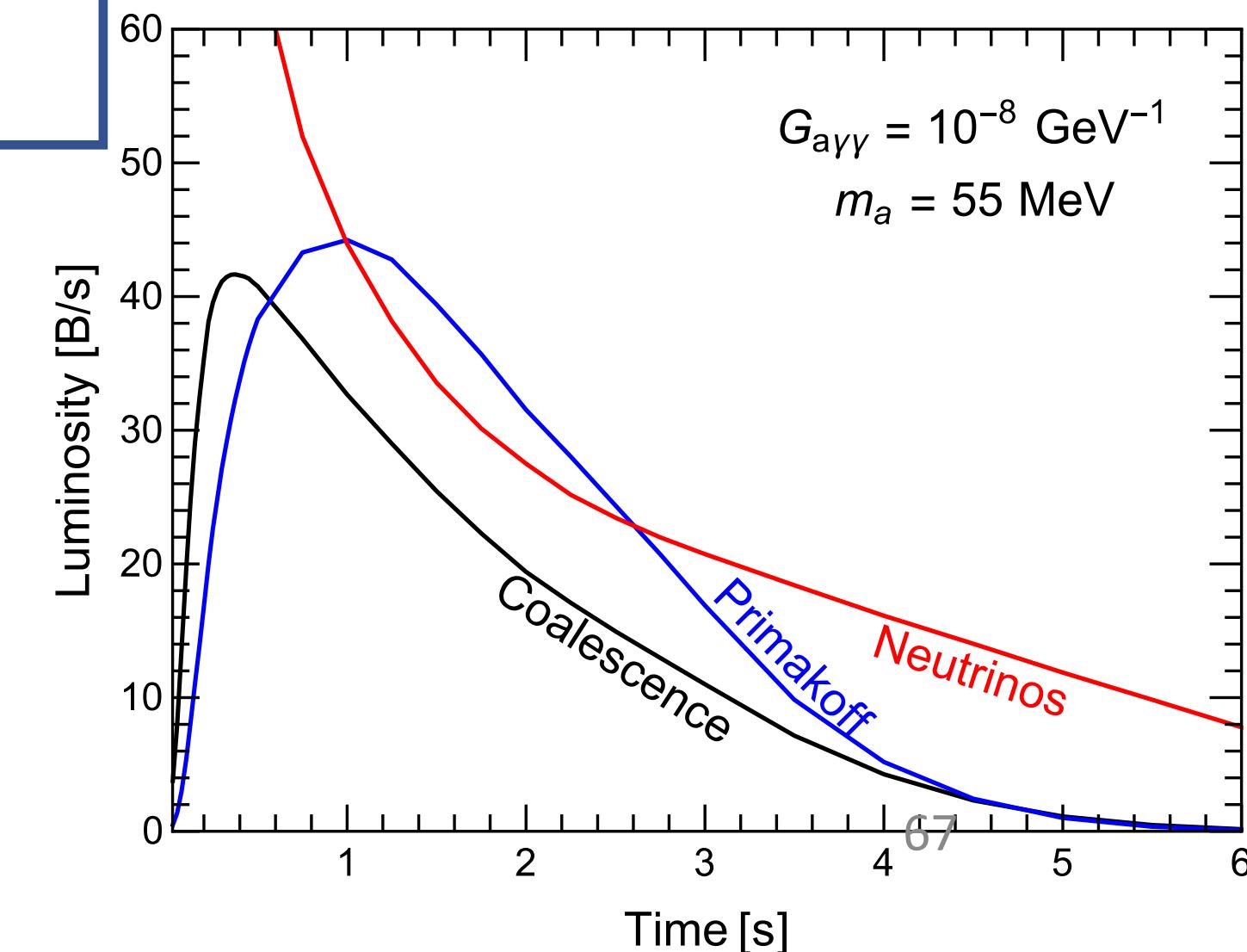
$$\sigma_P = \frac{Z^2 \alpha G_{a\gamma\gamma}^2}{2} f_P,$$

$$Q_a = \int \frac{d^3p}{(2\pi)^3} \omega e^{-\omega/T} \Gamma_A = \frac{G_{a\gamma\gamma}^2 T^3 m_a^4}{128\pi^3} F(m_a/T)$$

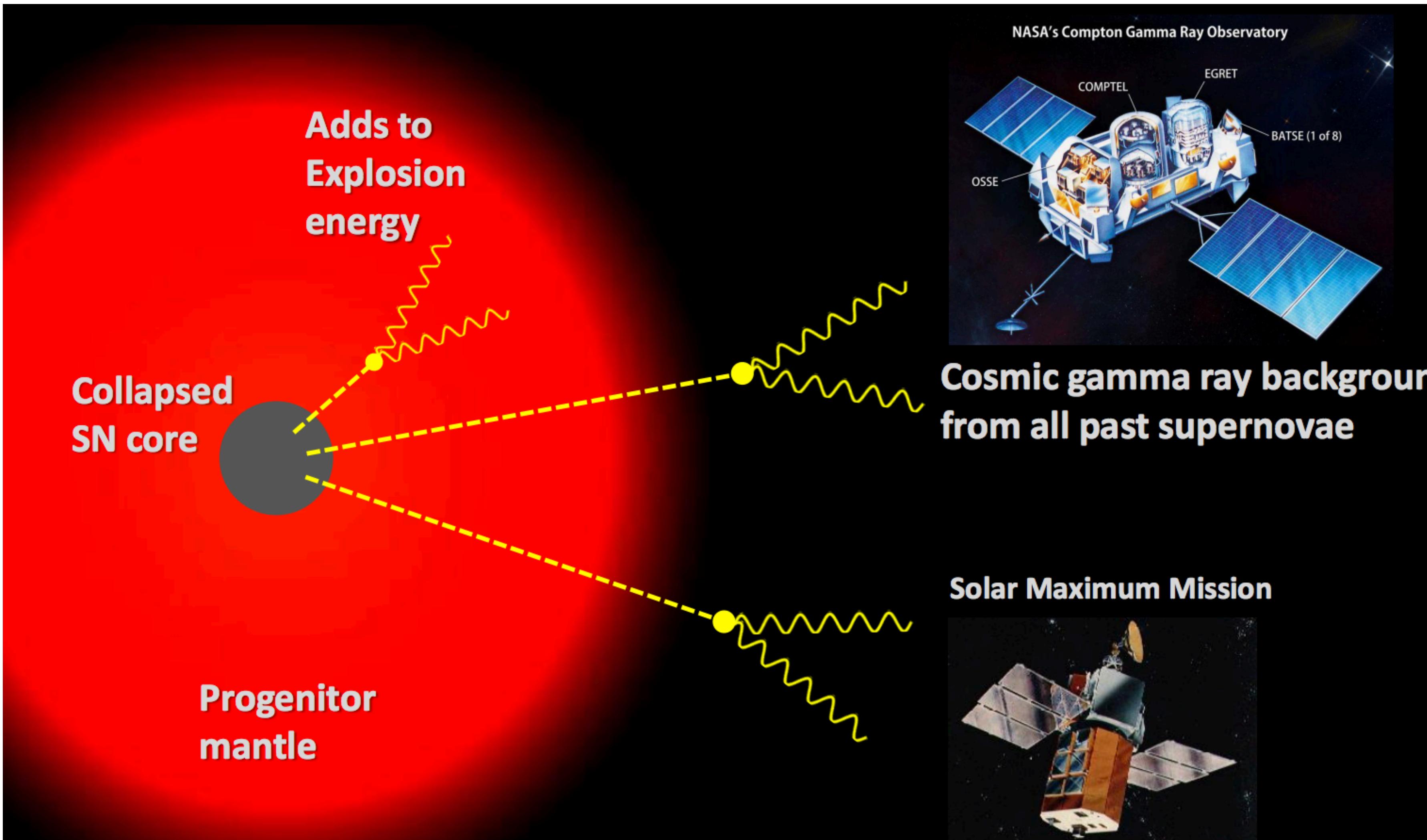
$$F(\mu) = \int_{\mu}^{\infty} dx x \sqrt{x^2 - \mu^2} e^{-x} f_B$$

Coalescence

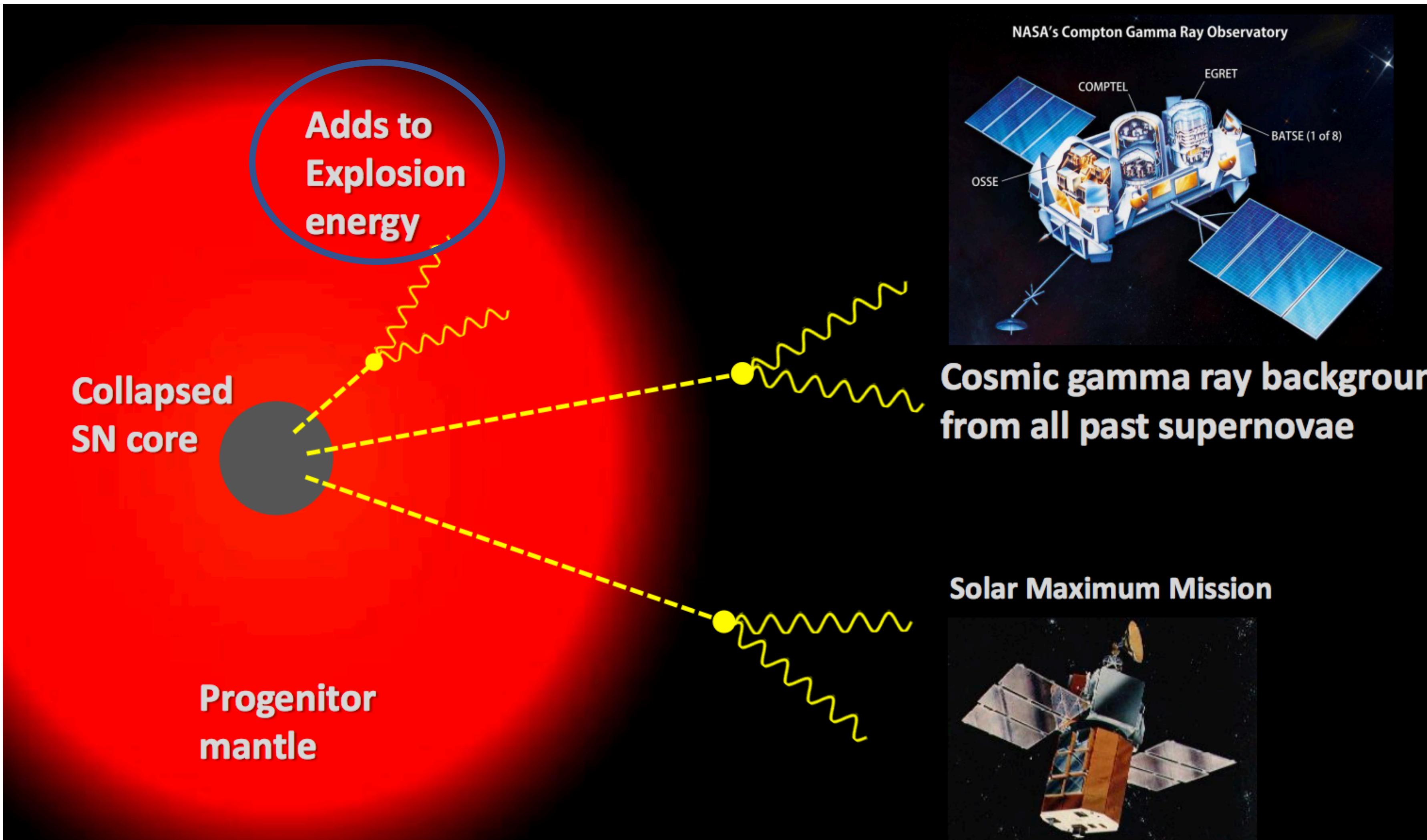
Primakoff



Use these SNe to get a bound on exotic particles which decay into SM relics



Use these SNe to get a bound on exotic particles which decay into SM relics



The axions get produced and then can decay back into photons!

$$\lambda_{a \rightarrow 2\gamma} = \frac{64\pi}{G_{a\gamma\gamma}^2} \frac{\sqrt{E_a^2 - m_a^2}}{m_a^4} \simeq \frac{4.0 \times 10^{13} \text{ cm}}{G_9^2} \frac{E_{100}}{m_{10}^4}$$

Gamma ray photons **quickly absorbed** by pair production on nuclei and electrons

The axions get produced and then can decay back into photons!

$$\lambda_{a \rightarrow 2\gamma} = \frac{64\pi}{G_{a\gamma\gamma}^2} \frac{\sqrt{E_a^2 - m_a^2}}{m_a^4} \simeq \frac{4.0 \times 10^{13} \text{ cm}}{G_9^2} \frac{E_{100}}{m_{10}^4}$$

Gamma ray photons **quickly absorbed** by pair production on nuclei and electrons

So axions are produced in the SN core and they then **decay in the SN mantle**, they can **dump energy** which should then show up either in **kinetic energy or radiation energy**.

$$E_{\text{mantle}} = \int dt \int_0^{R_{\text{NS}}} dR \int_{m'_a(R)}^{\infty} dE_a \frac{dL_a(R, E_a, t)}{dR dE_a} \\ \times \{ \exp[-(R_{\text{NS}} - R)/\lambda_a] - \exp[-(R_* - R)/\lambda_a] \},$$

The axions get produced and then can decay back into photons!

$$\lambda_{a \rightarrow 2\gamma} = \frac{64\pi}{G_{a\gamma\gamma}^2} \frac{\sqrt{E_a^2 - m_a^2}}{m_a^4} \simeq \frac{4.0 \times 10^{13} \text{ cm}}{G_9^2} \frac{E_{100}}{m_{10}^4}$$

Gamma ray photons **quickly absorbed** by pair production on nuclei and electrons

So axions are produced in the SN core and they then **decay in the SN mantle**, they can **dump energy** which should then show up either in **kinetic energy or radiation energy**.

Axion Luminosity

$$E_{\text{mantle}} = \int dt \int_0^{R_{\text{NS}}} dR \int_{m'_a(R)}^{\infty} dE_a \frac{dL_a(R, E_a, t)}{dR dE_a} \\ \times \left\{ \exp[-(R_{\text{NS}} - R)/\lambda_a] - \exp[-(R_* - R)/\lambda_a] \right\},$$

The axions get produced and then can decay back into photons!

$$\lambda_{a \rightarrow 2\gamma} = \frac{64\pi}{G_{a\gamma\gamma}^2} \frac{\sqrt{E_a^2 - m_a^2}}{m_a^4} \simeq \frac{4.0 \times 10^{13} \text{ cm}}{G_9^2} \frac{E_{100}}{m_{10}^4}$$

Gamma ray photons **quickly absorbed** by pair production on nuclei and electrons

So axions are produced in the SN core and they then **decay in the SN mantle**, they can **dump energy** which should then show up either in **kinetic energy or radiation energy**.

Axion Luminosity

$$E_{\text{mantle}} = \int dt \int_0^{R_{\text{NS}}} dR \int_{m'_a(R)}^{\infty} dE_a \frac{dL_a(R, E_a, t)}{dR dE_a} \\ \times \{\exp[-(R_{\text{NS}} - R)/\lambda_a] - \exp[-(R_* - R)/\lambda_a]\},$$

Neutron star radius

The axions get produced and then can decay
back into photons!

$$\lambda_{a \rightarrow 2\gamma} = \frac{64\pi}{G_{a\gamma\gamma}^2} \frac{\sqrt{E_a^2 - m_a^2}}{m_a^4} \simeq \frac{4.0 \times 10^{13} \text{ cm}}{G_9^2} \frac{E_{100}}{m_{10}^4}$$

Gamma ray photons quickly absorbed by pair production on nuclei and electrons

So axions are produced in the SN core and they then decay in the SN mantle, they can dump energy which should then show up either in kinetic energy or radiation energy.

$$E_{\text{mantle}} = \int dt \int_0^{R_{\text{NS}}} dR \int_{m'_a(R)}^{\infty} dE_a \frac{dL_a(R, E_a, t)}{dR dE_a} \times \{\exp[-(R_{\text{NS}} - R)/\lambda_a] - \exp[-(R_* - R)/\lambda_a]\},$$

The axions get produced and then can decay back into photons!

$$\lambda_{a \rightarrow 2\gamma} = \frac{64\pi}{G_{a\gamma\gamma}^2} \frac{\sqrt{E_a^2 - m_a^2}}{m_a^4} \simeq \frac{4.0 \times 10^{13} \text{ cm}}{G_9^2} \frac{E_{100}}{m_{10}^4}$$

Gamma ray photons **quickly absorbed** by pair production on nuclei and electrons

So axions are produced in the SN core and they then **decay in the SN mantle**, they can **dump energy** which should then show up either in **kinetic energy or radiation energy**.

$$E_{\text{mantle}} = \int dt \int_0^{R_{\text{NS}}} dR \int_{m'_a(R)}^{\infty} dE_a \frac{dL_a(R, E_a, t)}{dR dE_a} \times \{\exp[-(R_{\text{NS}} - R)/\lambda_a] - \exp[-(R_* - R)/\lambda_a]\},$$

Axion Luminosity
Neutron star radius
Progenitor radius
Axion MFP

Low-energy Supernovae (LESNe)

- 10–100 times dimmer than normal core-collapse SNe (CCSNe)
- 2–3 times lower photospheric expansion velocities
- Observations point to **0.1 B** (or smaller) explosion energies

Published: 04 June 2009

A low-energy core-collapse supernova without a hydrogen envelope

S. Valenti✉, A. Pastorello, E. Cappellaro, S. Benetti, P. A. Mazzali, J. Manteca, S. Taubenberger, N.

Elias-Rosa, R. Ferrando, A. Harutyunyan, V. P. Hentunen, M. Nissinen, E. Pian, M. Turatto, L. Zampieri &

S. J. Smartt

Monthly Notices
of the
ROYAL ASTRONOMICAL SOCIETYMNRAS 439, 2873–2892 (2014)
Advance Access publication 2014 February 21

doi:10.1093/mnras/stu156

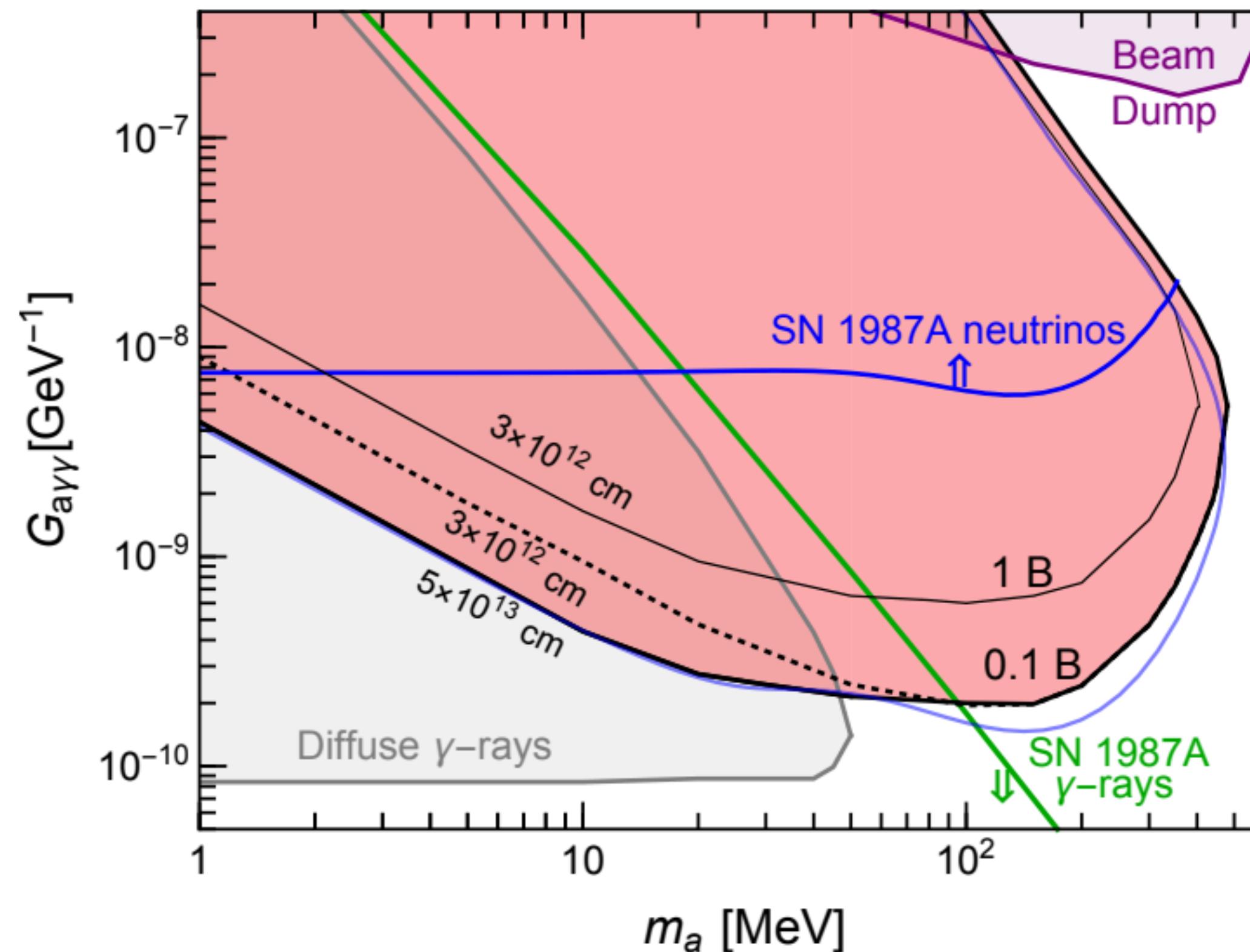
Low-luminosity Type II supernovae: spectroscopic and photometric evolution

A. Pastorello,^{1,2,5}★ L. Zampieri,² M. Turatto,² E. Cappellaro,³ S. Benetti,² D. Branch,⁵ E. Baron,⁵ F. Patat,⁶ M. Armstrong,⁷ C. M. Salvo⁸ and M. Riello^{2,1}¹Dipartimento di Astronomia, Università di Padova, Vicolo dell' Osservatorio 2, I-35122 Padova, Italy²INAF – Osservatorio Astronomico di Padova, Vicolo dell' Osservatorio 5, I-35122 Padova, Italy³INAF – Osservatorio Astronomico di Capodimonte, Via Moiariello 16, I-80131 Napoli, Italy⁴Astrophysics Group, Blackett Laboratory, Imperial College, Prince Consort Road, London SW7 2BZ⁵Department of Physics and Astronomy, University of Oklahoma, 440 W. Brooke St., Norman, OK 73019, USA⁶European Southern Observatory, Karl-Schwarzschild-Strasse 2, D-85748, Garching bei München, Germany⁷UK Supernova Patrol, British Astronomical Association, Rolvenden, Kent⁸Australian National University, Mount Stromlo Observatory, Cotter Road, Weston ACT 2611, Australia

Low luminosity Type II supernovae – II. Pointing towards moderate mass precursors

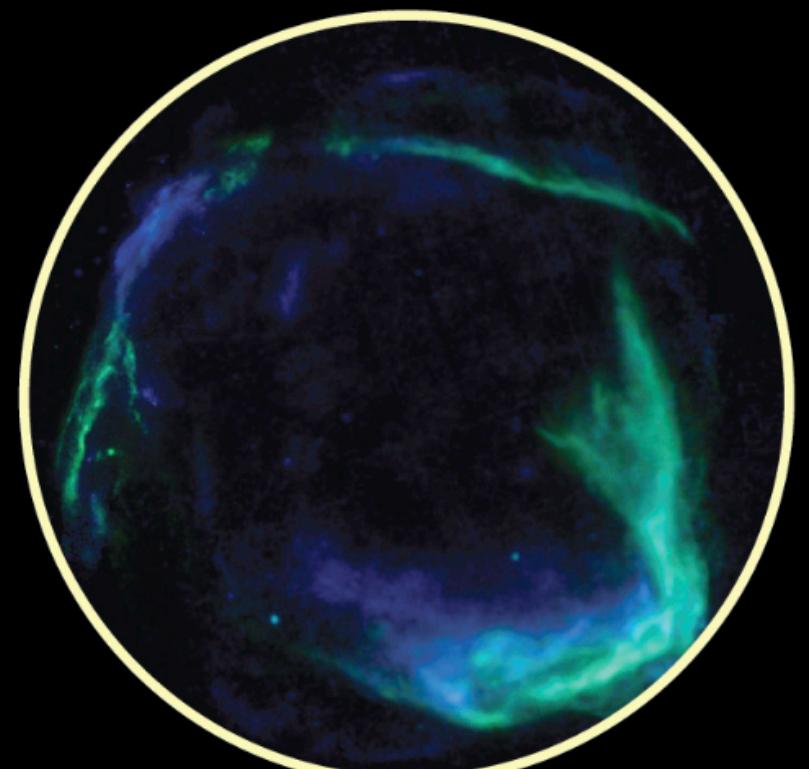
S. Spiro,¹★ A. Pastorello,¹ M. L. Pumo,^{1,2} L. Zampieri,¹ M. Turatto,¹ S. J. Smartt,³ S. Benetti,¹ E. Cappellaro,¹ S. Valenti,^{4,5} I. Agnoletto,¹ G. Altavilla,⁶ T. Aoki,⁷ E. Brocato,⁸ E. M. Corsini,^{1,9} A. Di Cianno,¹⁰ N. Elias-Rosa,¹¹ M. Hamuy,¹² K. Enya,¹³ M. Fiaschi,⁹ G. Folatelli,¹⁴ S. Desidera,¹ A. Harutyunyan,¹⁵ D. A. Howell,^{4,5} A. Kawka,¹⁶ Y. Kobayashi,¹⁷ B. Leibundgut,¹⁸ T. Minezaki,⁷ H. Navasardyan,¹ K. Nomoto,^{19,20} S. Mattila,²¹ A. Pietrinferni,¹⁰ G. Pignata,²² G. Raimondo,¹⁰ M. Salvo,²³ B. P. Schmidt,²³ J. Sollerman,²⁴ J. Spyromilio,¹⁸ S. Taubenberger,²⁵ G. Valentini,¹⁰ S. Vennes¹⁶ and Y. Yoshii⁷

Imposing the energy deposit in the mantle is not too big gives very strong bounds.

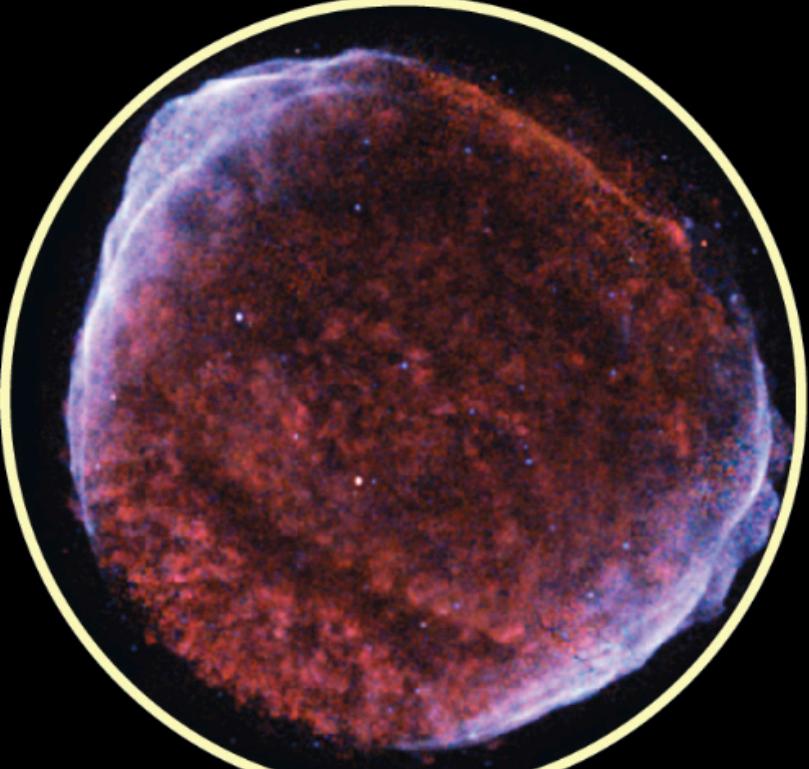


Caputo, H.T. Janka, G. Raffelt, E. Vitagliano, *Phys.Rev.Lett.* 128 (2022) 22, 221103

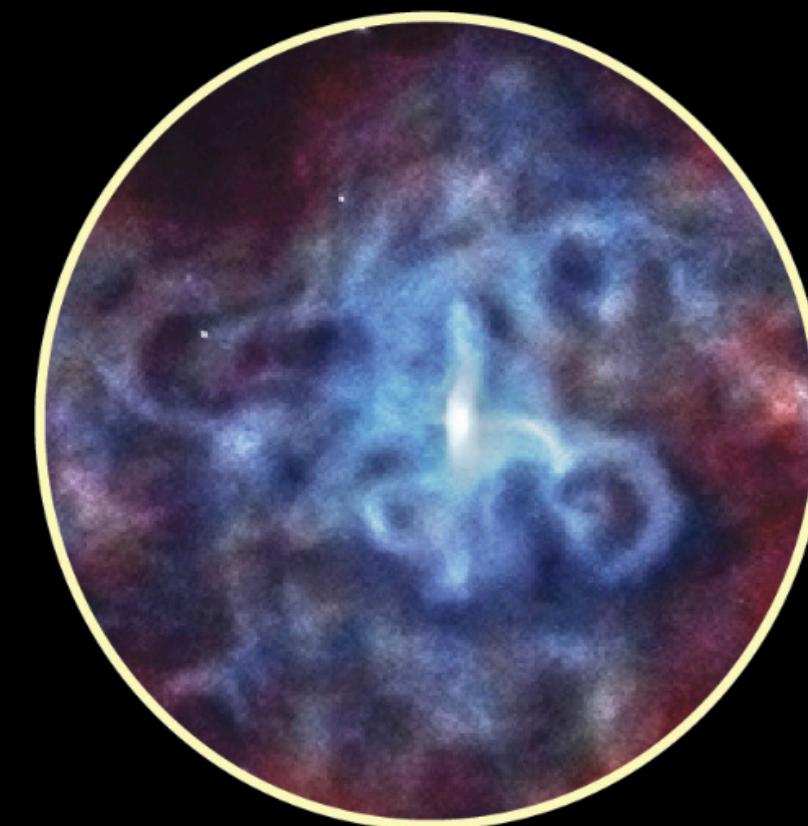
Why this is exciting to me?



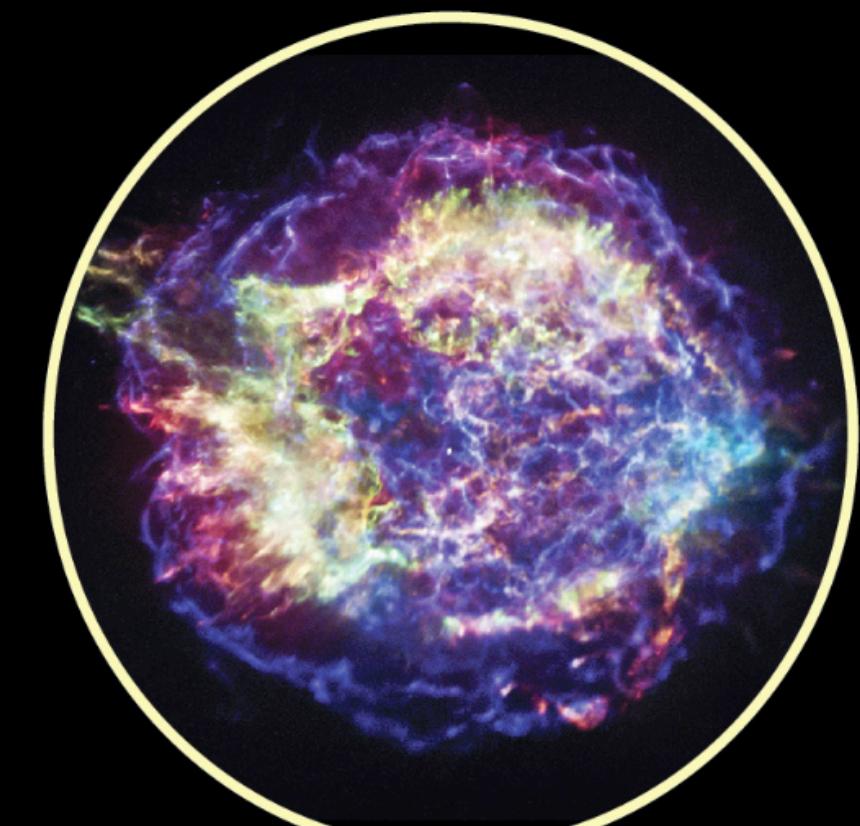
A.D. 185
RCW 86
Historical Observers: Chinese
Likelihood of Identification: Possible
Distance Estimate: 8,200 light years
Type: Core collapse of massive star



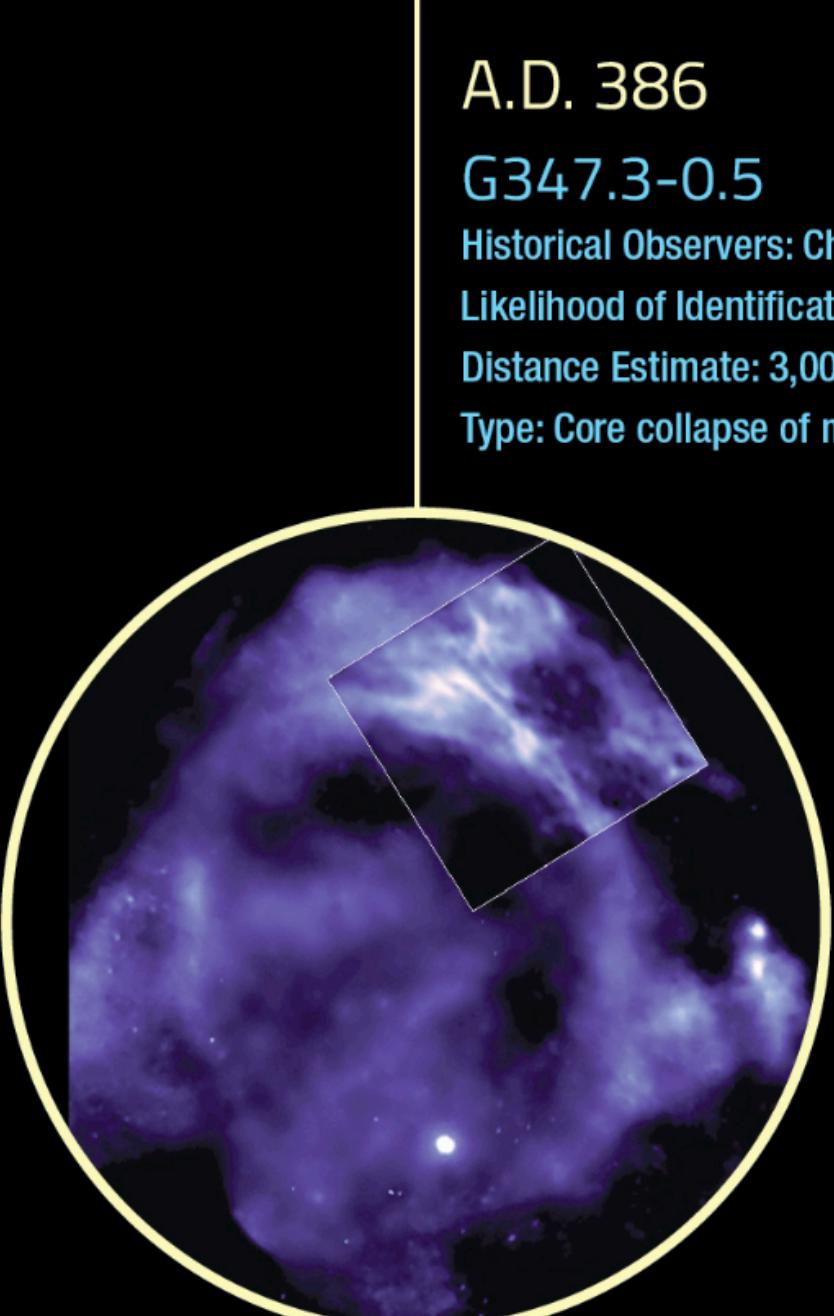
A.D. 393
SN 1006
Historical Observers: Chinese,
Japanese, Arabic, European
Likelihood of Identification: Definite
Distance Estimate: 7,000 light years
Type: Thermonuclear explosion of white dwarf



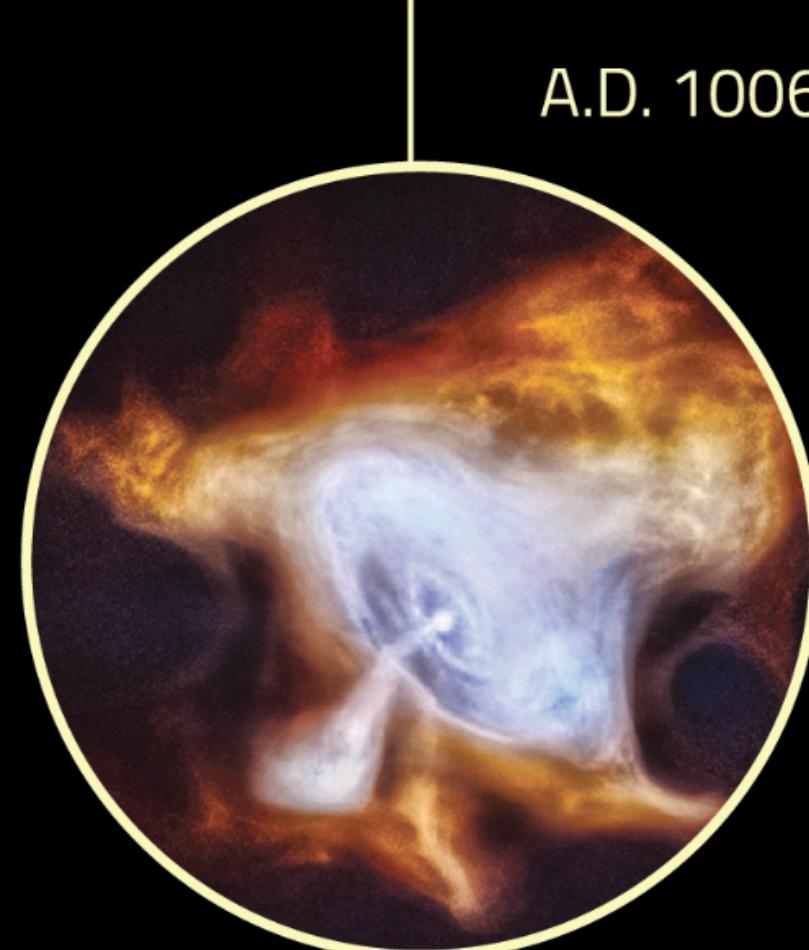
A.D. 1054
3C58
Historical Observers: Chinese, Japanese
Likelihood of Identification: Possible
Distance Estimate: 10,000 light years
Type: Core collapse of massive star



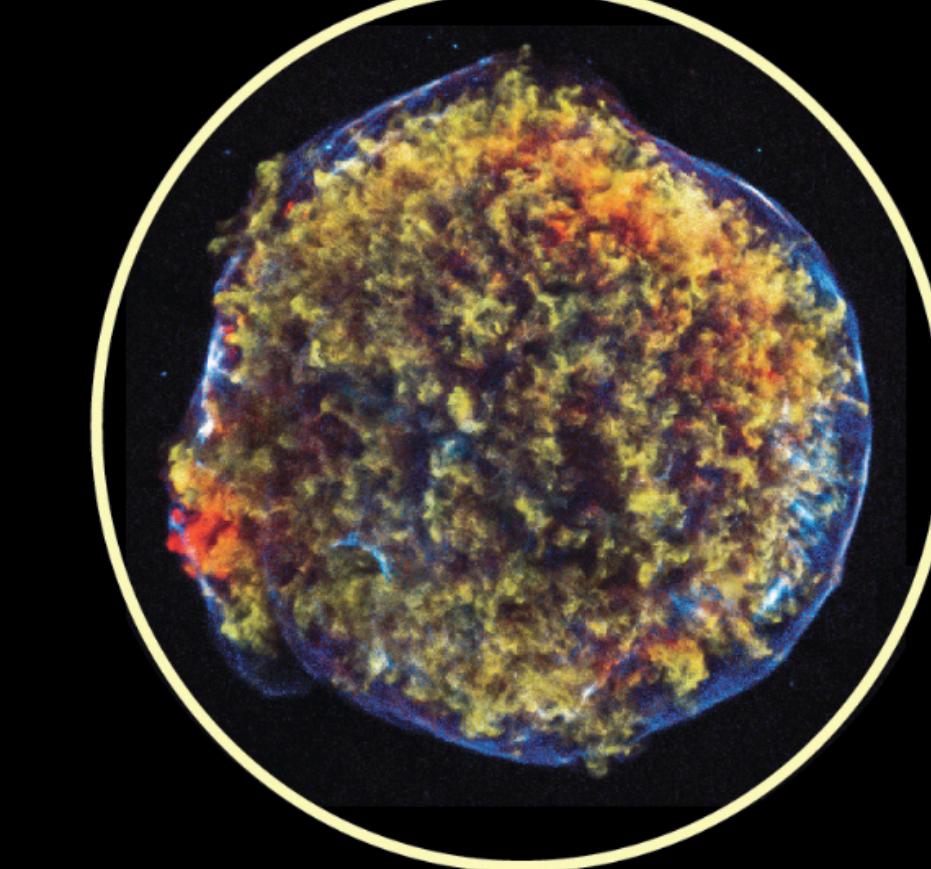
A.D. 1680
Cassiopeia A
Historical Observers: European?
Likelihood of Identification: Unlikely
Distance Estimate: 10,000 light years
Type: Core collapse of massive star



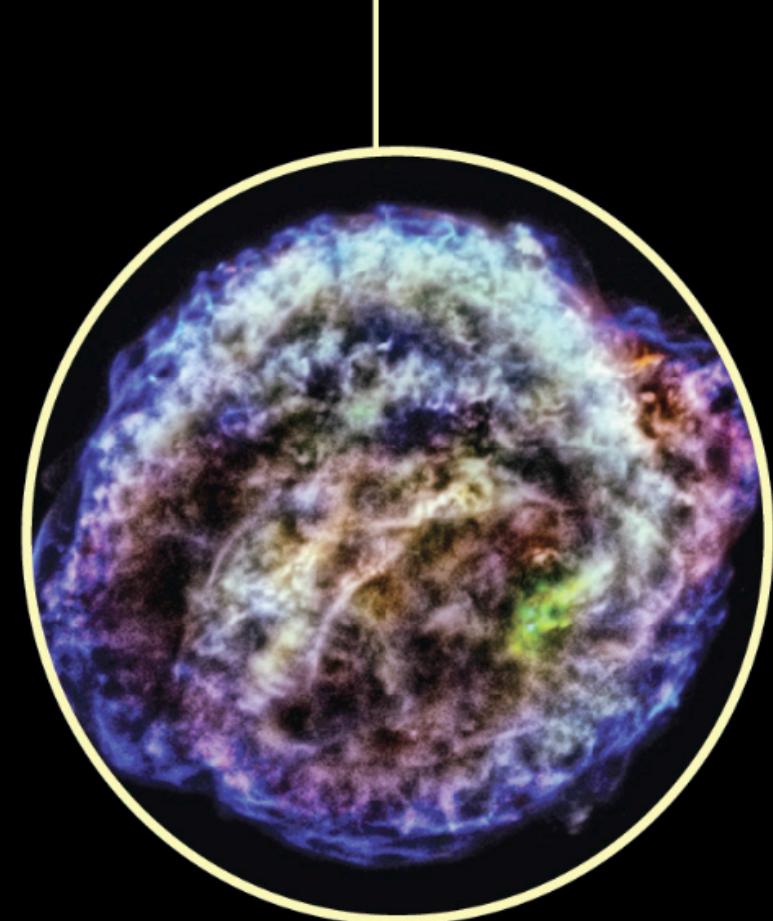
LIGHT YEAR: the distance that light, moving at a constant speed of 300,000 m/s, travels in one year. One light year is just under 0 trillion kilometers.



Crab Nebula
Historical Observers: Chinese, Japanese,
Arabic, Native American
Likelihood of Identification: Definite
Distance Estimate: 6,000 light years
Type: Core collapse of massive star



A.D. 1572
Tycho's SNR
Historical Observers: European, Chinese, Korean
Likelihood of Identification: Definite
Distance Estimate: 7,500 light years
Type: Thermonuclear explosion of white dwarf



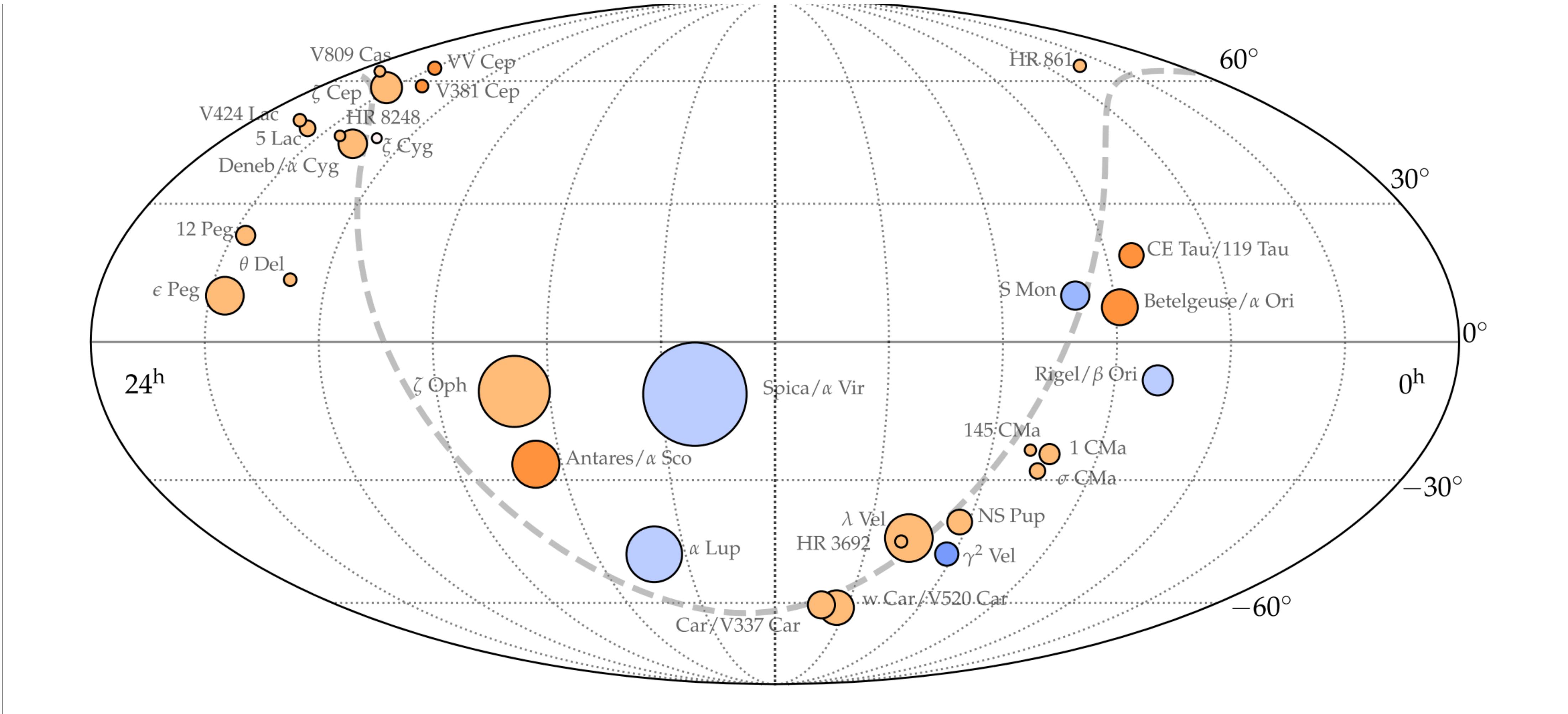
A.D. 1604
Kepler's SNR
Historical Observers: European, Chinese, Korean
Likelihood of Identification: Definite
Distance Estimate: 13,000 light years
Type: Thermonuclear explosion of white dwarf?

[nature](#) > [news feature](#) > [article](#)

NEWS FEATURE | 21 February 2022

A supernova could light up the Milky Way at any time. Astronomers will be watching

Next SN candidates within 1 kpc from us



Plans for the future

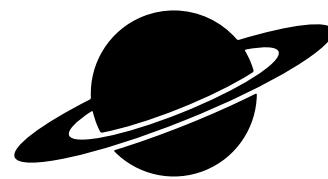
- Detailed simulations and output from radiative transport simulations in supernovae;
- Extension of the LESN criterion to other relevant models;
- Cross-correlation studies for line intensity mapping + foreground analysis + extension to other dark matter models;
- Study dark matter impact on star formation (JWST is taking data, amazing opportunity).

Thanks for the attention!



Back-up slides

Stars and new physics



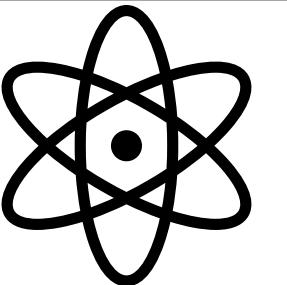
- Caputo et al., *JCAP* 08 (2022) 08, 045
Caputo et al., *Phys.Rev.Lett.* 128 (2022) 22, 221103 (CCSNe)
Caputo et al., *Phys.Rev.D* 105 (2022) 3, 035022 (CCSNe)
Caputo et al., *Phys.Rev.Lett.* 127 (2021) 18, 181102 (Hypernovae)
Caputo et al., *Phys.Rev.D* 101 (2020) 12, 123004 (Sun)
O'Hare, Caputo, et al., *Phys.Rev.D* 102 (2020) 4, 043019 (Sun)

Gravitational Waves

- Caputo et al., *JAstrophys.J.* 892 (2020) 2, 90
Toubiana, Sberna, Caputo, et al.,
Phys.Rev.Lett. 126 (2021) 10, 101105

My research

Dark Matter direct detection



- Caputo et al., *Phys.Rev.D* 100 (2019) 11, 116007 (sub-GeV DM)
Caputo et al., *Phys.Lett.B* 802 (2020) 135258 (sub-GeV DM)
Caputo et al., *Phys.Rev.D* 103 (2021) 5, 055017 (sub-GeV DM)
Bloch, Caputo et al., *JHEP* 01 (2021) 178 (sub-GeV DM)

Caputo et al., *Phys.Rev.D* 104 (2021) 9, 095029 (Dark Photon DM)

Dark Matter and telescopes

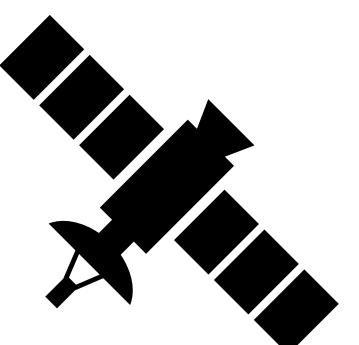


- Caputo et al., *JCAP* 03 (2019) 027 (radio)
Caputo et al., *Phys.Rev.D* 98 (2018) 8, 083024 (radio)
Caputo et al., *JCAP* 03 (2020) 001 (X-rays)
Caputo et al., *JCAP* 05 (2021) 046 (IR)
Bernal, Caputo, et al., *Phys.Rev.D* 103 (2021) 6, 063523 (optic)
Bernal, Caputo, et al., *Phys.Rev.Lett.* 127 (2021) 13, 131102 (optic)

Colliders phenomenology

- Caputo, P. Hernandez, M. Kekic, J. Lopez-Pavon, J.Salvado,
Eur.Phys.J.C 77 (2017) 4, 258
Caputo, P. Hernandez, J. Lopez-Pavon, J.Salvado, *JHEP* 06 (2017) 112
D. Barducci, E. Bertuzzo, Caputo, P. Hernandez, *JHEP* 06 (2020) 185
D. Barducci, E. Bertuzzo, Caputo, P. Hernandez, B. Mele,
JHEP 03 (2021) 117

Dark Matter and other indirect probes



- Caputo et al., *Phys.Rev.Lett.* 127 (2021) 1, 011102 (21 cm)
Bolton, Caputo, et al., *Phys.Rev.Lett.* 129 (2022) 21, 211102 (Ly-alpha)
Caputo et al., *Phys.Rev.Lett.* 125 (2020) 22, 221303 (CMB)
Bernal, Caputo, et al., *arXiv* 2208.13794 (blazars data)
Caputo et al., *Phys.Rev.D* 100 (2019) 6, 063515 (pulsars)
Caputo et al., *Phys.Dark Univ.* 19 (2018) 1-11 (pulsars)

Stars and new physics



- Caputo et al., *JCAP* 08 (2022) 08, 045
Caputo et al., *Phys.Rev.Lett.* 128 (2022) 22, 221103 (CCSNe)
Caputo et al., *Phys.Rev.D* 105 (2022) 3, 035022 (CCSNe)
Caputo et al., *Phys.Rev.Lett.* 127 (2021) 18, 181102 (Hypernovae)
Caputo et al., *Phys.Rev.D* 101 (2020) 12, 123004 (Sun)
O'Hare, Caputo, et al., *Phys.Rev.D* 102 (2020) 4, 043019 (Sun)

Dark Matter and telescopes



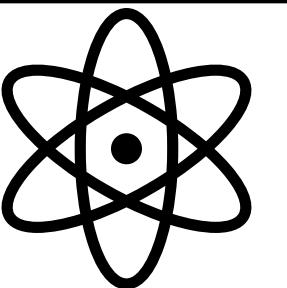
- Caputo et al., *JCAP* 03 (2019) 027 (radio)
Caputo et al., *Phys.Rev.D* 98 (2018) 8, 083024 (radio)
Caputo et al., *JCAP* 03 (2020) 001 (X-rays)
Caputo et al., *JCAP* 05 (2021) 046 (IR)

Gravitational Waves

- Caputo et al., *JAstrophys.J.* 892 (2020) 2, 90
Toubiana, Sberna, Caputo, et al.,
Phys.Rev.Lett. 126 (2021) 10, 101105

My research

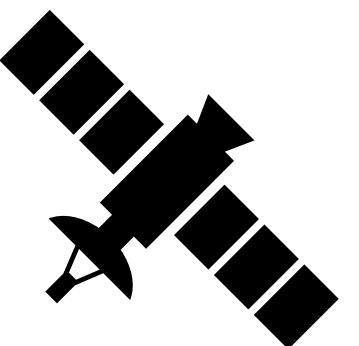
Dark Matter direct detection



- Caputo et al., *Phys.Rev.D* 100 (2019) 11, 116007 (sub-GeV DM)
Caputo et al., *Phys.Lett.B* 802 (2020) 135258 (sub-GeV DM)
Caputo et al., *Phys.Rev.D* 103 (2021) 5, 055017 (sub-GeV DM)
Bloch, Caputo et al., *JHEP* 01 (2021) 178 (sub-GeV DM)

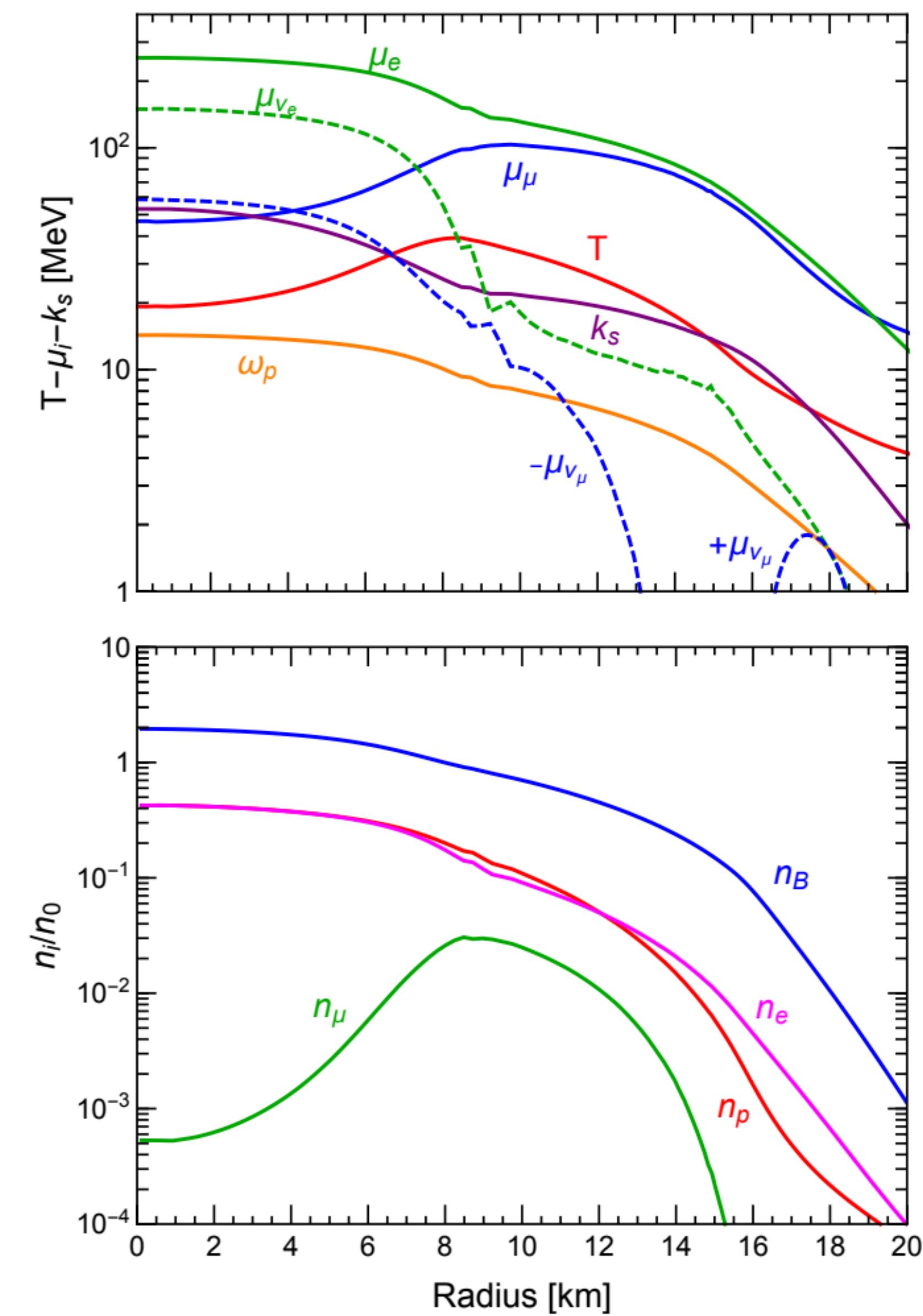
Caputo et al., *Phys.Rev.D* 104 (2021) 9, 095029 (Dark Photon DM)

Dark Matter and other indirect probes



- Caputo et al., *Phys.Rev.Lett.* 127 (2021) 1, 011102 (21 cm)
Bolton, Caputo, et al., *Phys.Rev.Lett.* 129 (2022) 21, 211102 (Ly-alpha)
Caputo et al., *Phys.Rev.Lett.* 125 (2020) 22, 221303 (CMB)
Bernal, Caputo, et al., *arXiv* 2208.13794 (blazars data)
Caputo et al., *Phys.Rev.D* 100 (2019) 6, 063515 (pulsars)
Caputo et al., *Phys.Dark Univ.* 19 (2018) 1-11 (pulsars)

Garching SN model



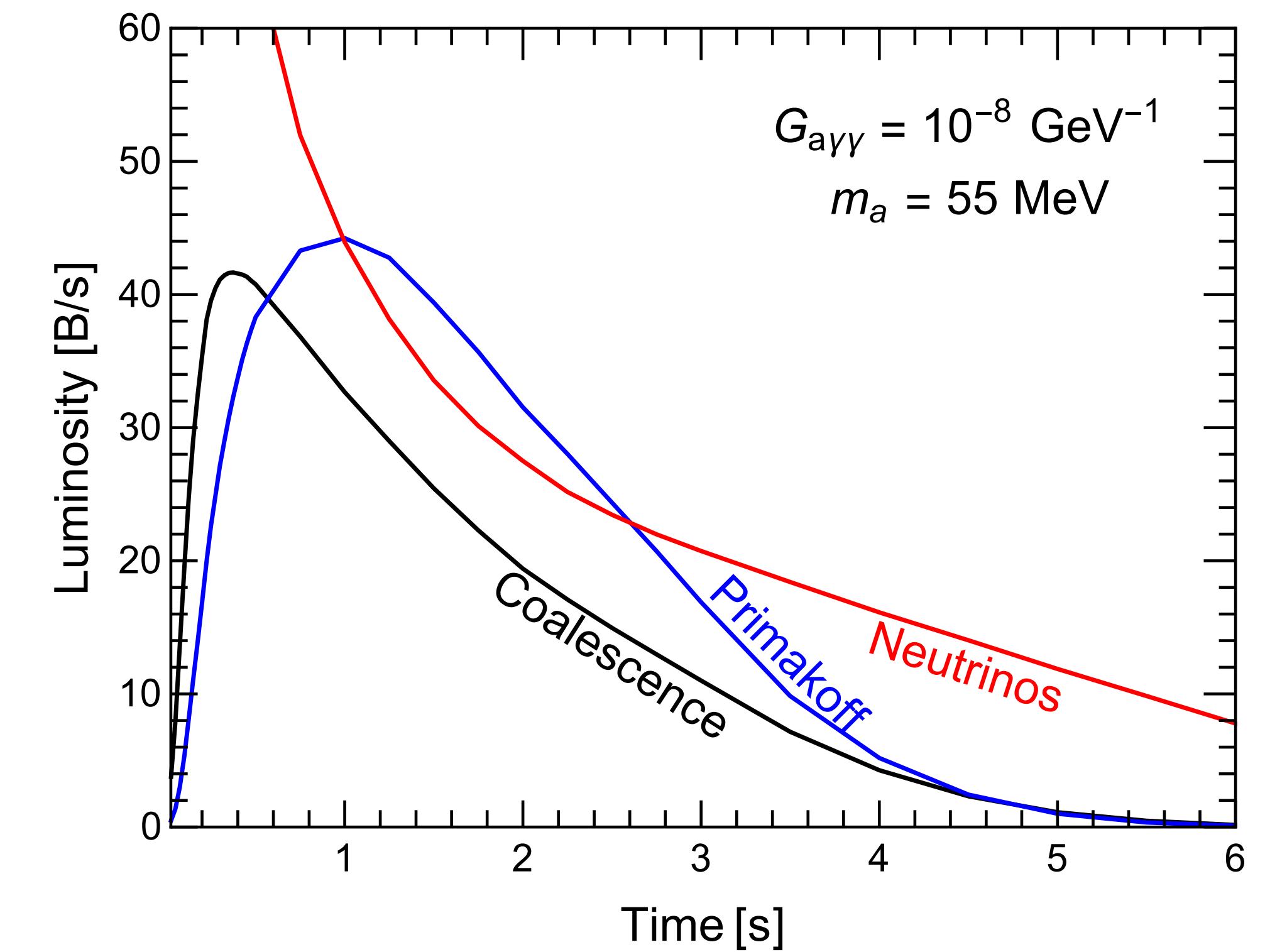
Primakoff and Coalescence calculation

$$Q_a = \int \frac{2d^3\mathbf{k}}{(2\pi)^3} \frac{\omega}{e^{\omega/T} - 1} \hat{n} \sigma_P \quad \sigma_P = \frac{Z^2 \alpha G_{a\gamma\gamma}^2}{2} f_P;$$

$$= \hat{n} \alpha G_{a\gamma\gamma}^2 \frac{\pi^2 T^4}{30} \langle f_P \rangle, \quad \langle f_P \rangle = 20 \frac{m_a^2 + 3m_a T + 3T^2}{\pi^4 T^2} e^{-m_a/T}$$

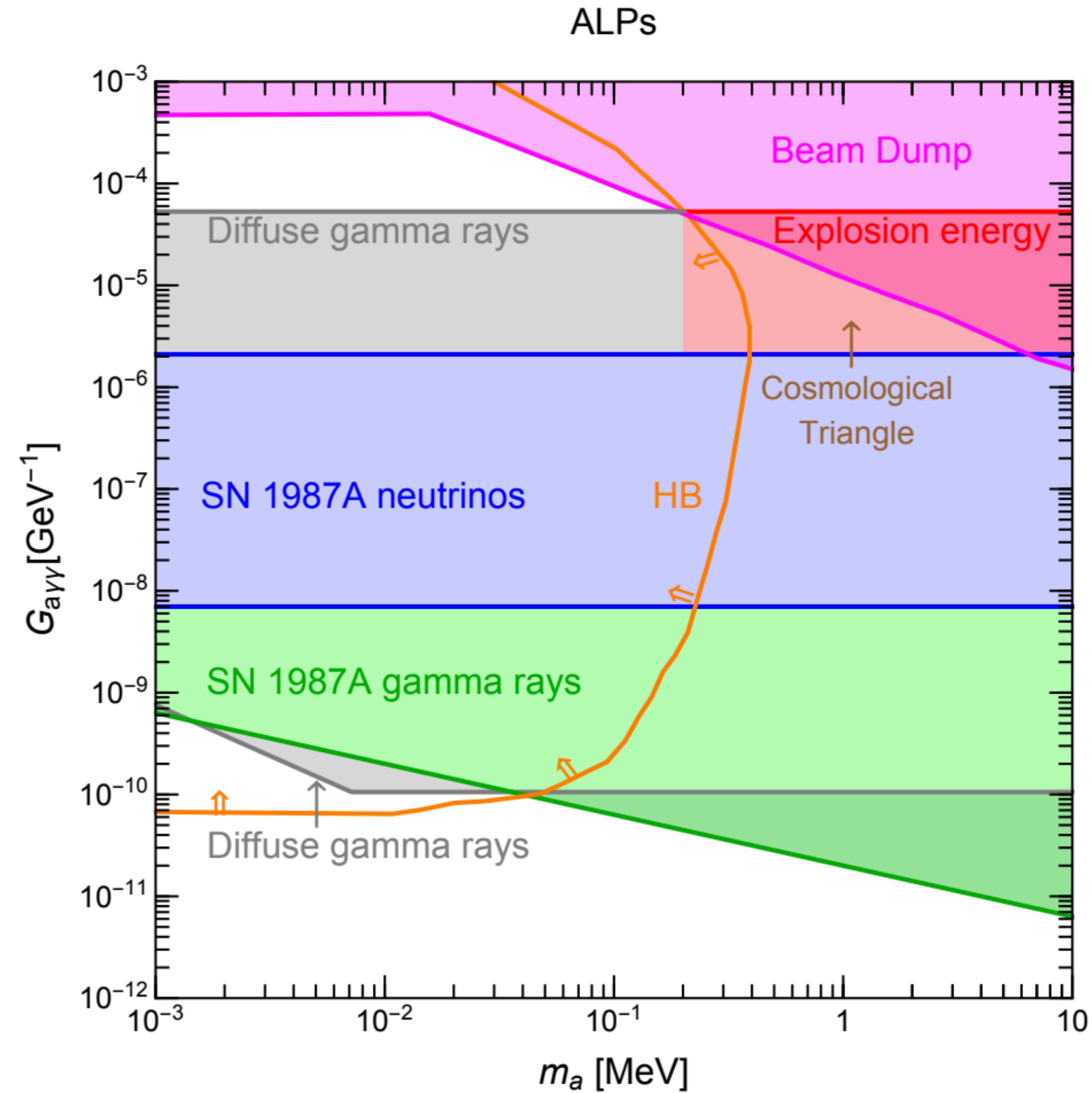
$$Q_a = \int \frac{d^3\mathbf{p}}{(2\pi)^3} \omega e^{-\omega/T} \Gamma_A = \frac{G_{a\gamma\gamma}^2 T^3 m_a^4}{128\pi^3} F(m_a/T)$$

$$F(\mu) = \int_{\mu}^{\infty} dx x \sqrt{x^2 - \mu^2} e^{-x} f_B$$



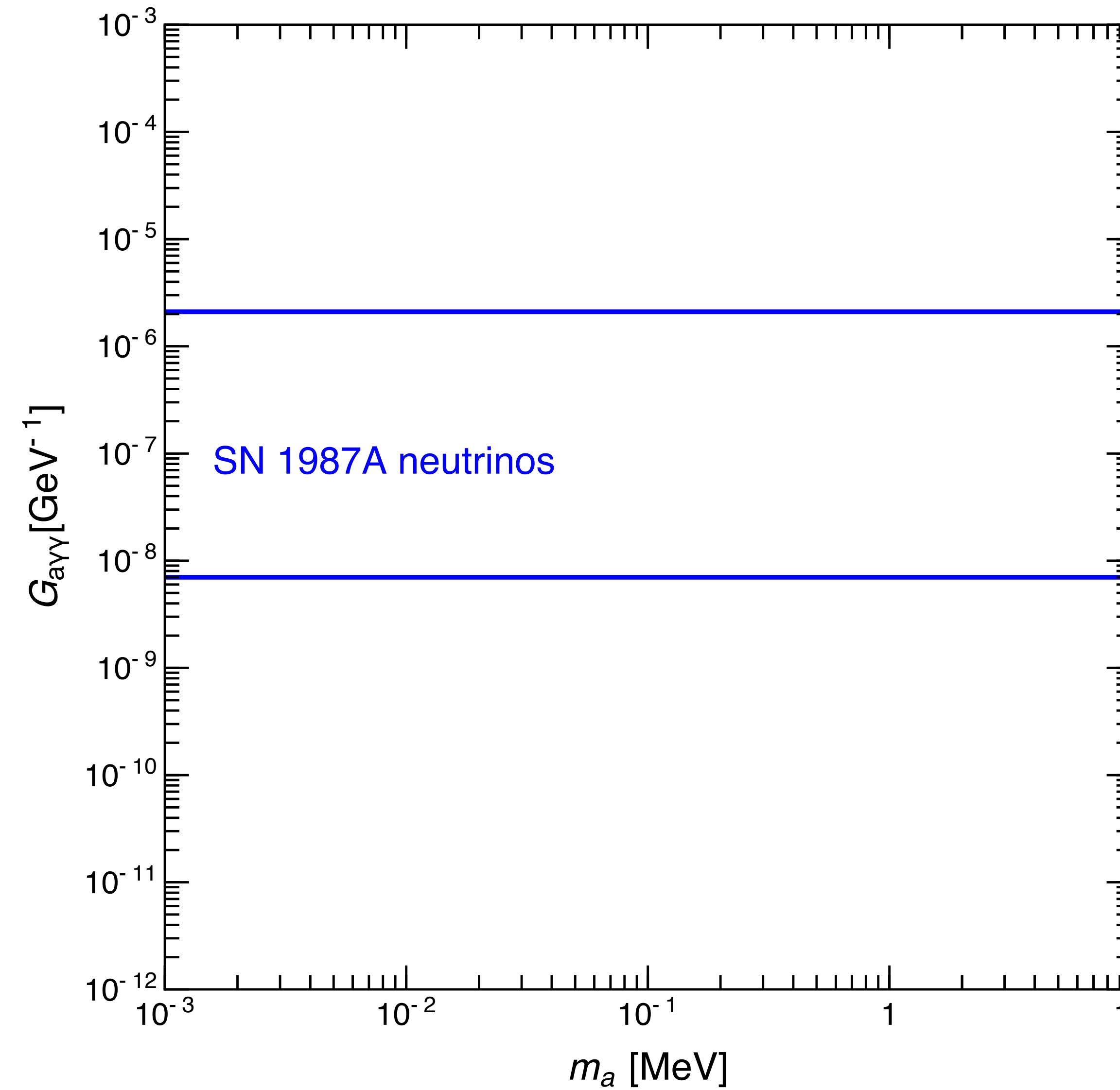
Other bounds

For masses above 1–10 MeV, muonphilic particles can also be efficiently probed at colliders. In particular, [electron beam-dump experiments](#), such as the SLAC E137 experiment or the planned Jefferson Lab BDX experiment provide an excellent source of secondary muons, which can then be used to look for muonic (pseudo)scalars.

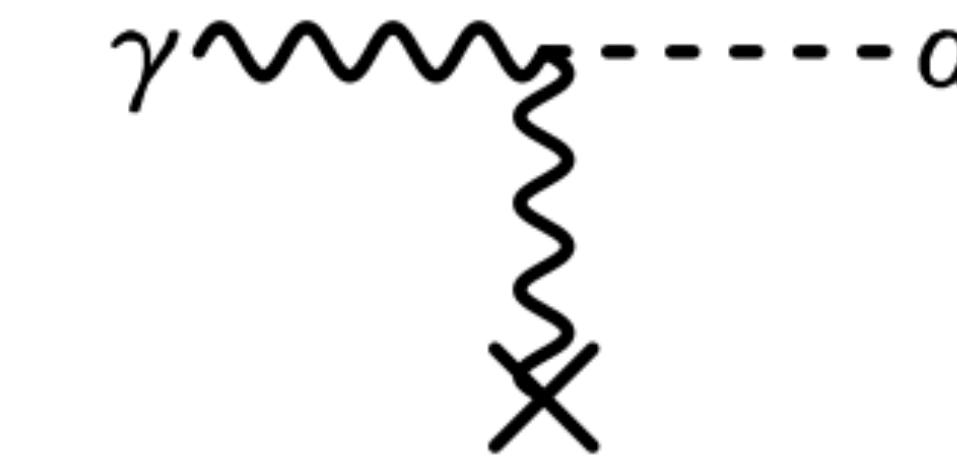


Cooling bounds from SN1987A (free streaming)

ALPs



The duration of several seconds of the SN 1987A neutrino signal is incompatible with excessive energy loss in hypothetical new forms of radiation. The main emission process is photo production on charged particles (aka Primakoff)

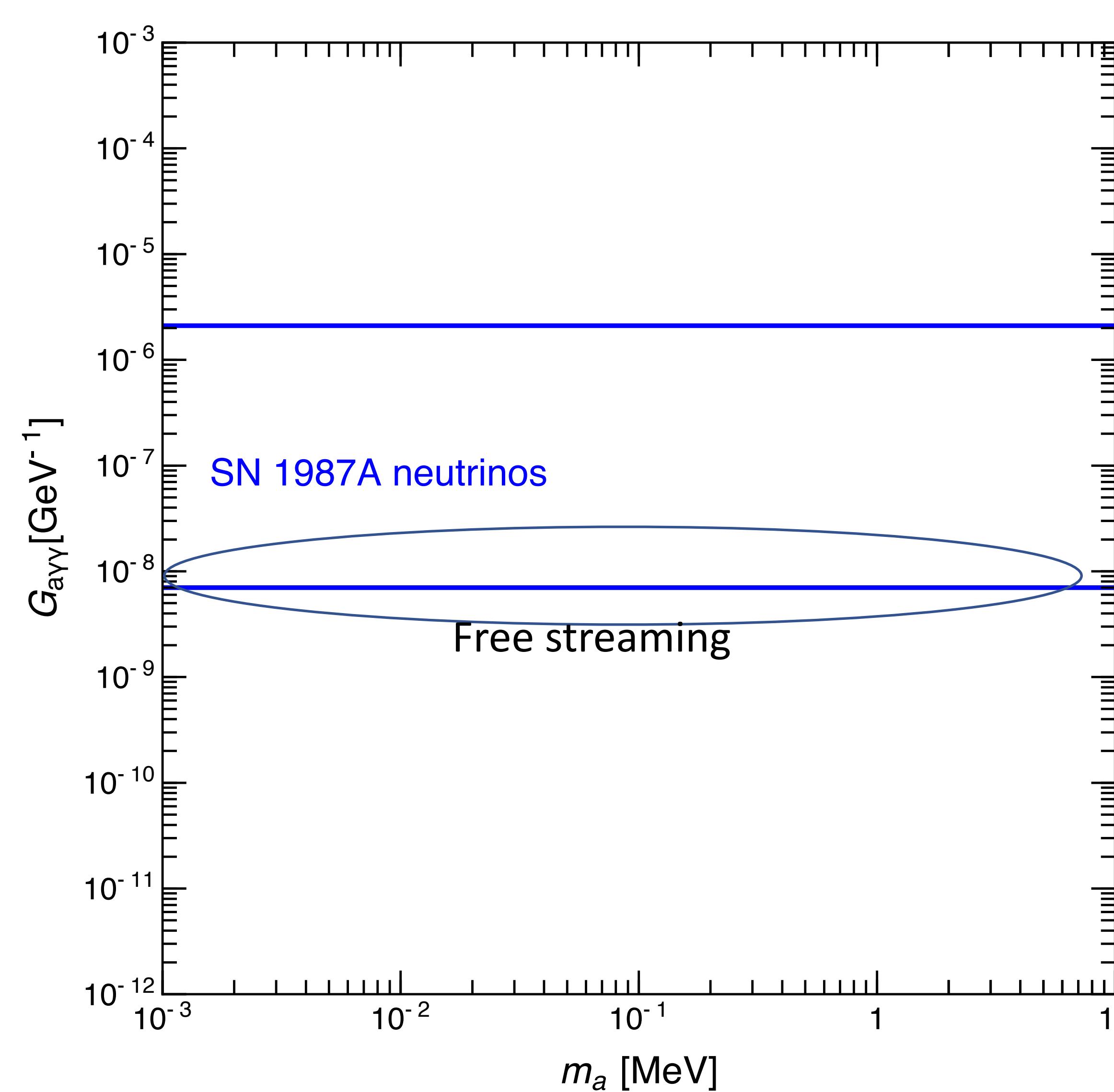


$$Q_P \simeq \hat{n} \frac{2\alpha G_{a\gamma\gamma}^2}{3\pi^2} (m_a^2 + 3m_a T + 3T^2) T^2 e^{-m_a/T}$$

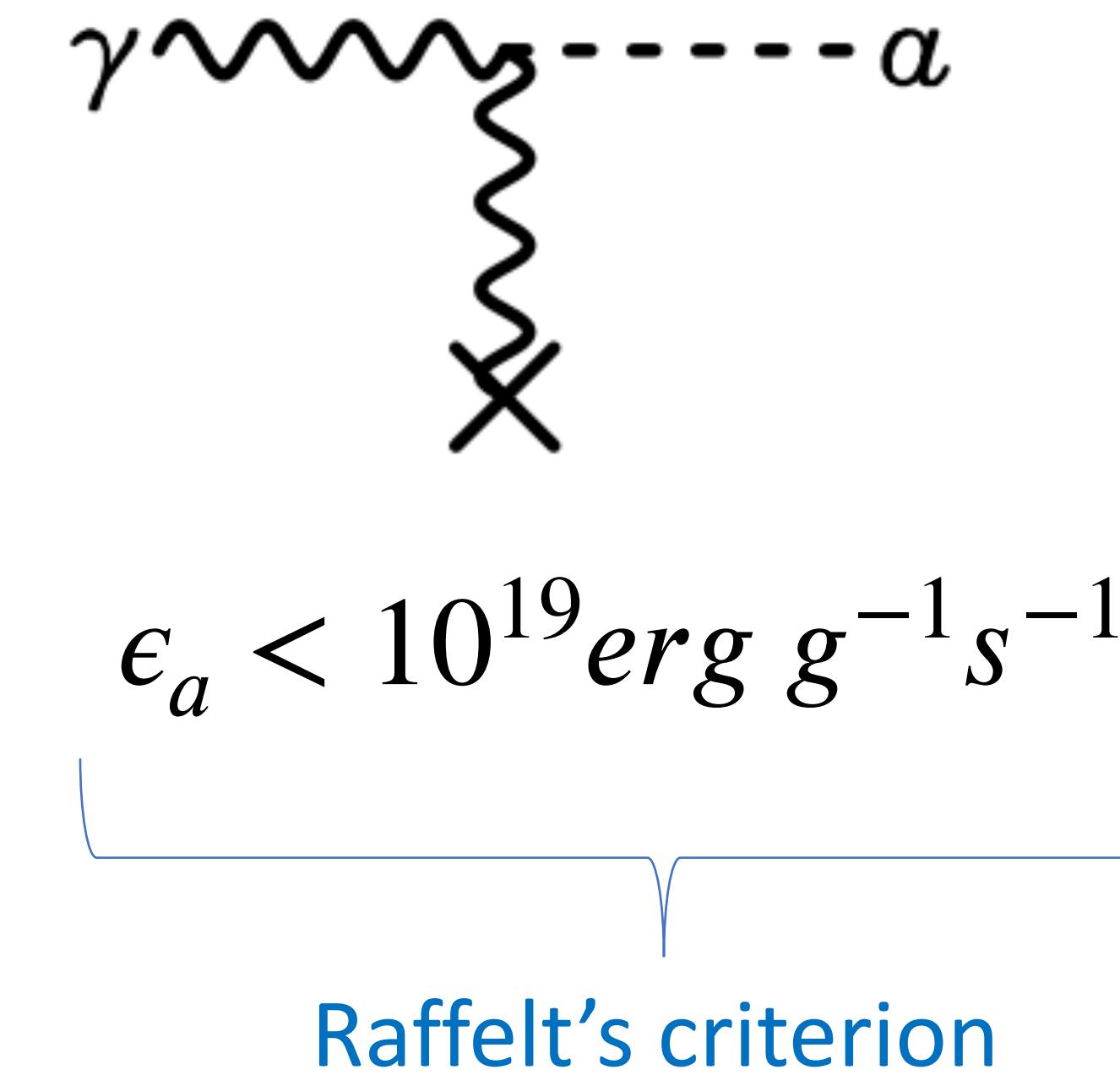
A.C, T. Janka, G. Raffelt and E. Vitagliano, in preparation

Cooling bounds from SN1987A

ALPs

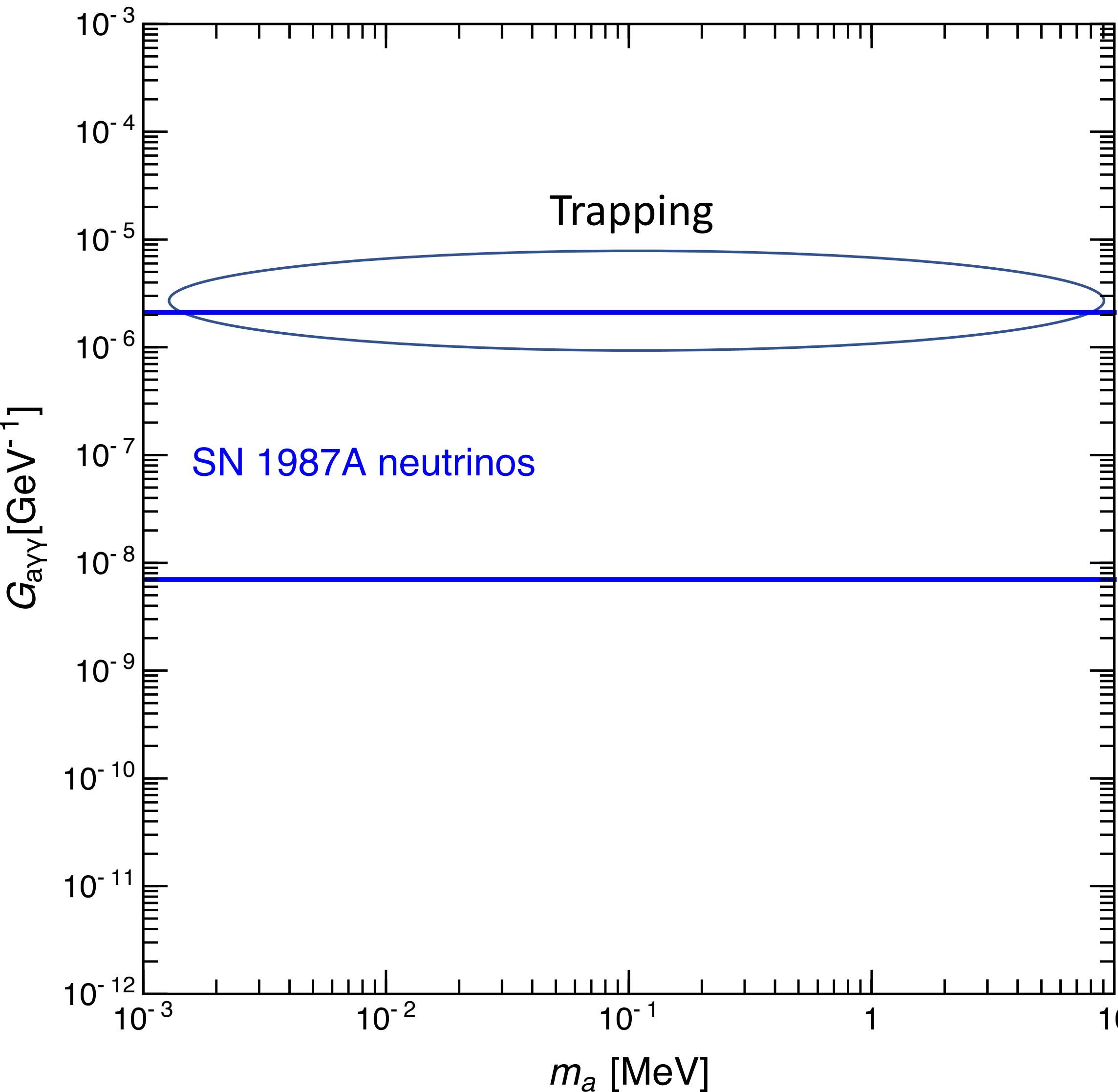


The duration of several seconds of the SN 1987A neutrino signal is incompatible with excessive energy loss in hypothetical new forms of radiation. The main emission process in this case is photo-production on charged particles (aka Primakoff)



Cooling bounds from SN1987A (trapping)

ALPs



In the trapping limit, our bosons emerge from a region near the PNS surface whence they escape without being reabsorbed on their way out, in analogy to the neutrino sphere

$$L_\phi = 4\pi R_\phi^2 \frac{\pi^2}{120} T^4(R_\phi)$$

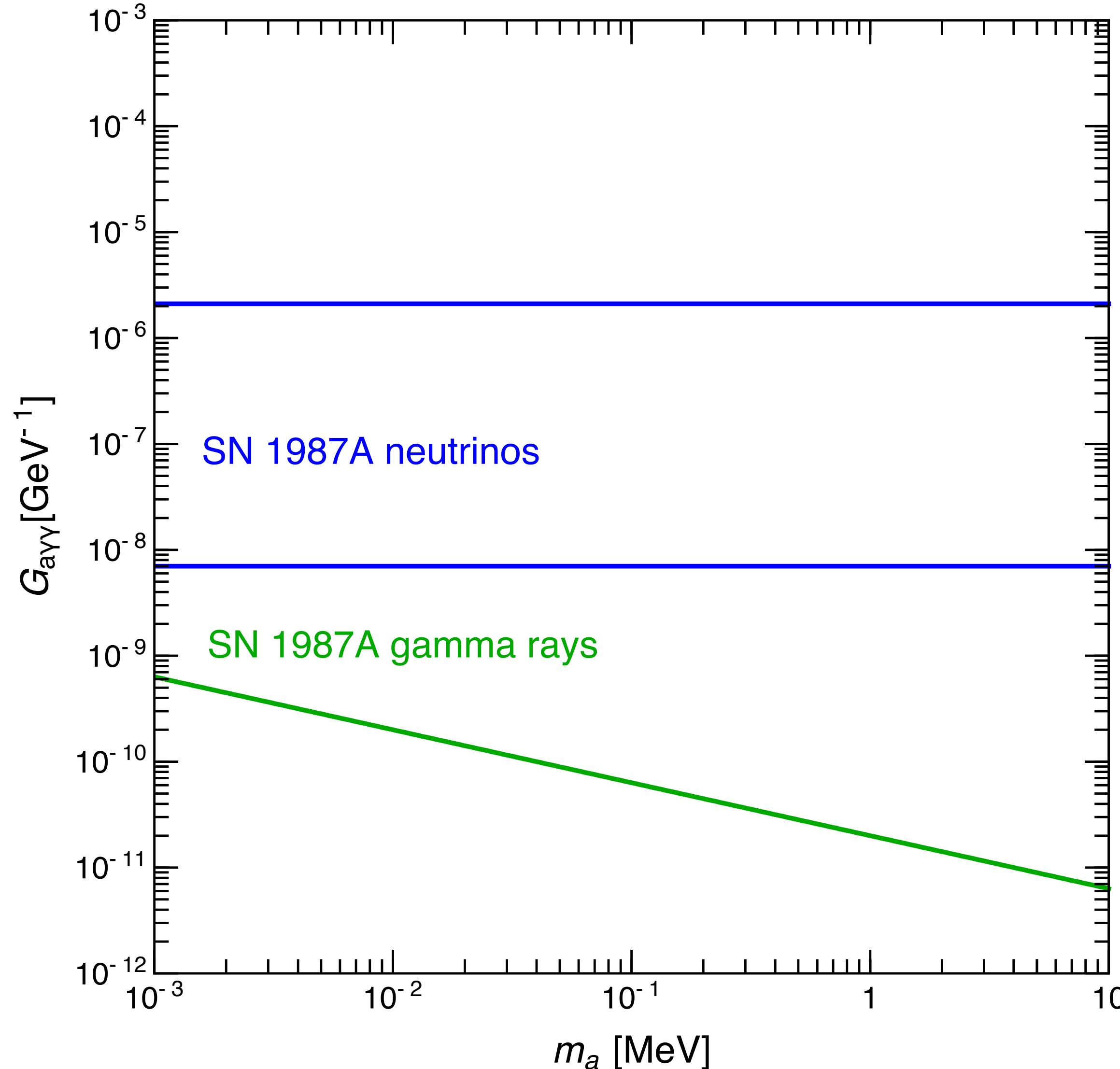
The radius of emission is determined by the optical depth becoming of order one

$$\tau(r) = \int_r^\infty dr' \Gamma(r')$$

where in natural units the interaction rate is the same as the inverse mfp.

Bounds from Gamma-ray signals from SN1987A

ALPs

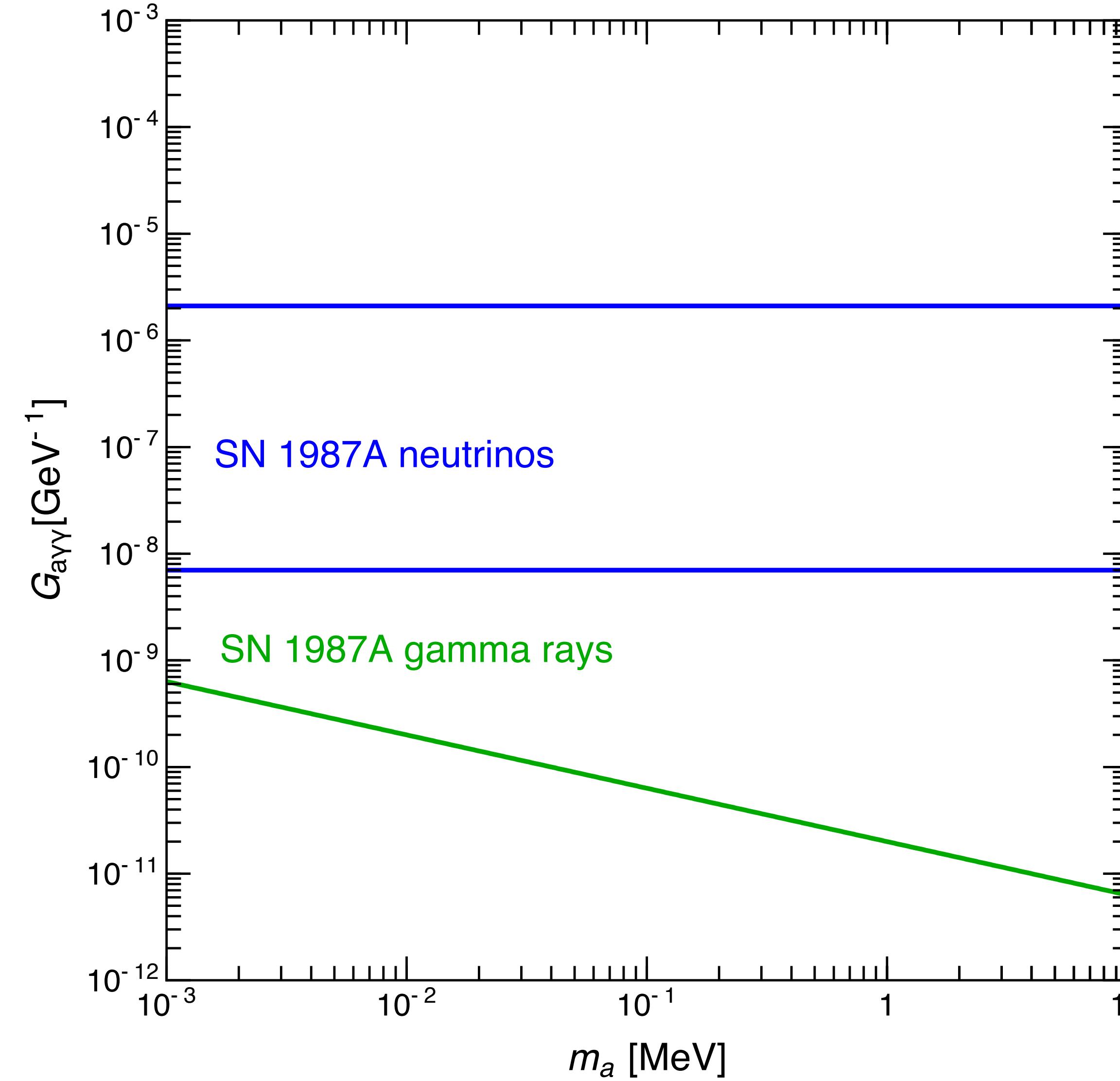


The (putative) produced bosons from 1987A could have [decayed into photons](#), which would have been picked up by the Gamma-Ray Spectrometer (GRS) on board the Solar Maximum Mission (SMM) satellite that operated 02/1980–12/1989 .

$$\frac{dF_\gamma}{dE_\gamma dt} = 2 \frac{2E_\gamma}{m_a \tau_a} e^{-2E_\gamma t/m_a \tau_a} \int_{E_\gamma}^{\infty} dE_a \frac{\Phi_a(E_a)}{E_a}$$

Bounds from Gamma-ray signals from SN1987A

ALPs



The (putative) produced bosons from 1987A could have decayed into photons, which would have been picked up by the Gamma-Ray Spectrometer (GRS) on board the Solar Maximum Mission (SMM) satellite that operated 02/1980–12/1989 .

$$\frac{dF_\gamma}{dE_\gamma dt} = 2 \frac{2E_\gamma}{m_a \tau_a} e^{-2E_\gamma t/m_a \tau_a} \int_{E_\gamma}^{\infty} dE_a \frac{\Phi_a(E_a)}{E_a}$$

fluence

\downarrow

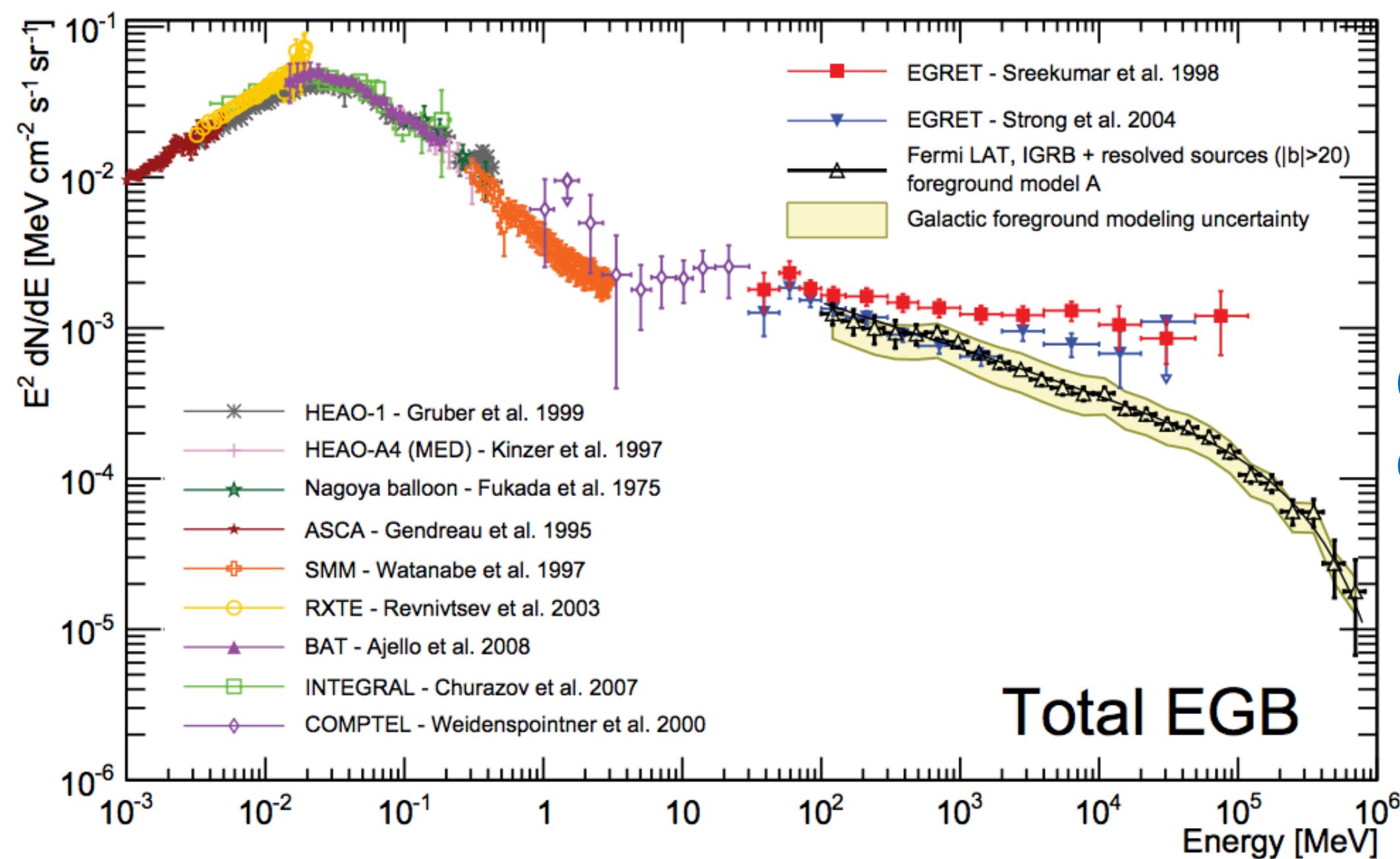
rest-frame boson lifetime

Good agreement with J. Jaeckel, P. C. Malta and J. Redondo, Phys. Rev. D 98 (2018) 055032, although we have been able to provide simple analytical expressions.

Diffuse SNe Gamma-rays

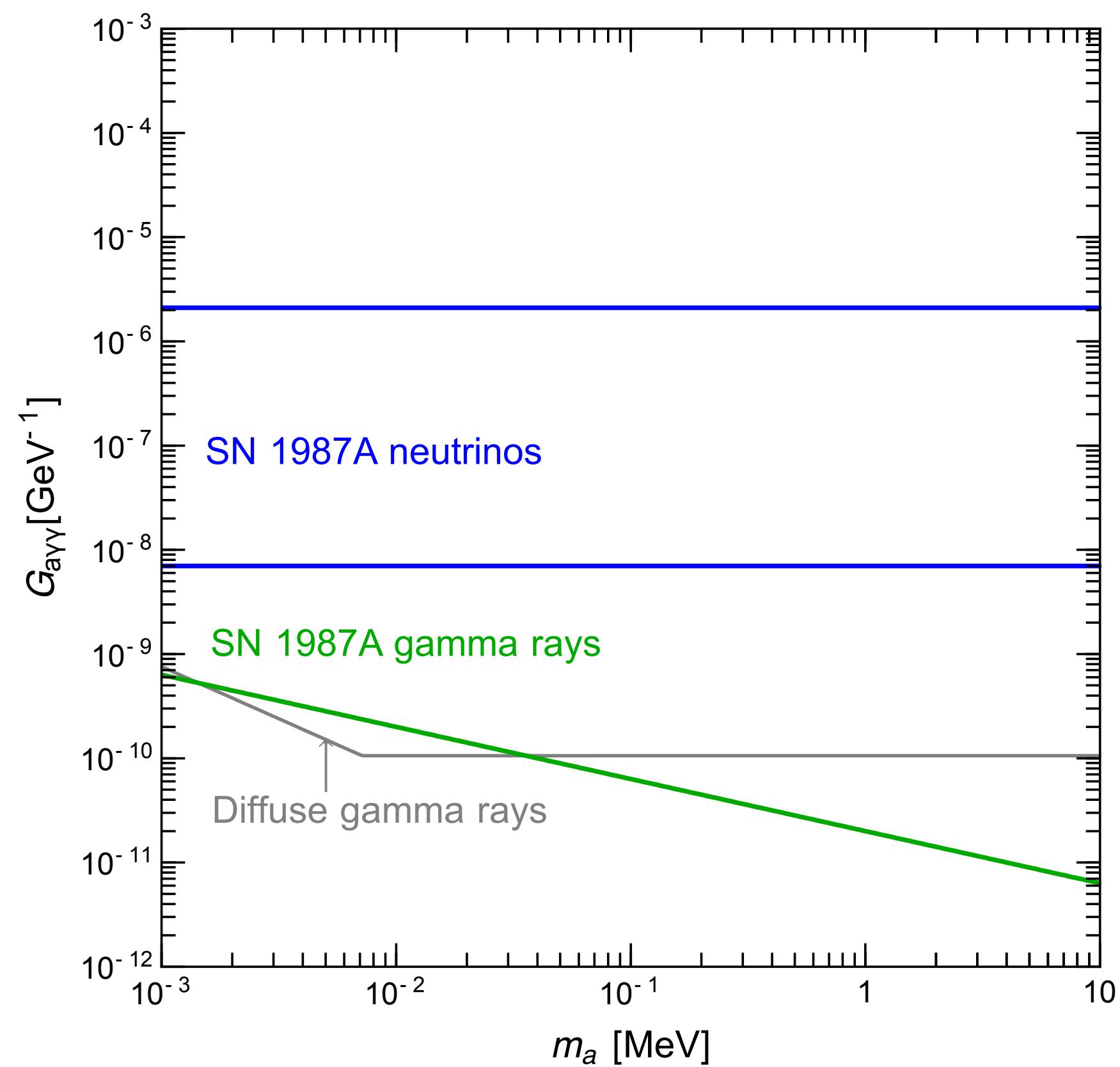
The scalar boson emission from **all past SNe** creates a cosmic background density in analogy to the diffuse SN neutrino background (DSNB). The radiative decays of these particles contribute to the **diffuse cosmic γ -ray background** and thus can be constrained.

$$\frac{dn_\gamma}{d\omega} = \int_0^\infty dz(1+z)n'_{cc}(z) \int_{\omega_z}^\infty dE_z f_D(E_z) \frac{2}{E_z} F_a(E_z)$$



Compare with
data!

Fermi-LAT Collaboration, M. Ackermann et al,
Astrophys. J. 799 (2015) 86



Diffuse SNe Gamma-rays

The scalar boson emission from **all past SNe** creates a cosmic background density in analogy to the diffuse SN neutrino background (DSNB). The radiative decays of these particles contribute to the **diffuse cosmic γ -ray background** and thus can be constrained.

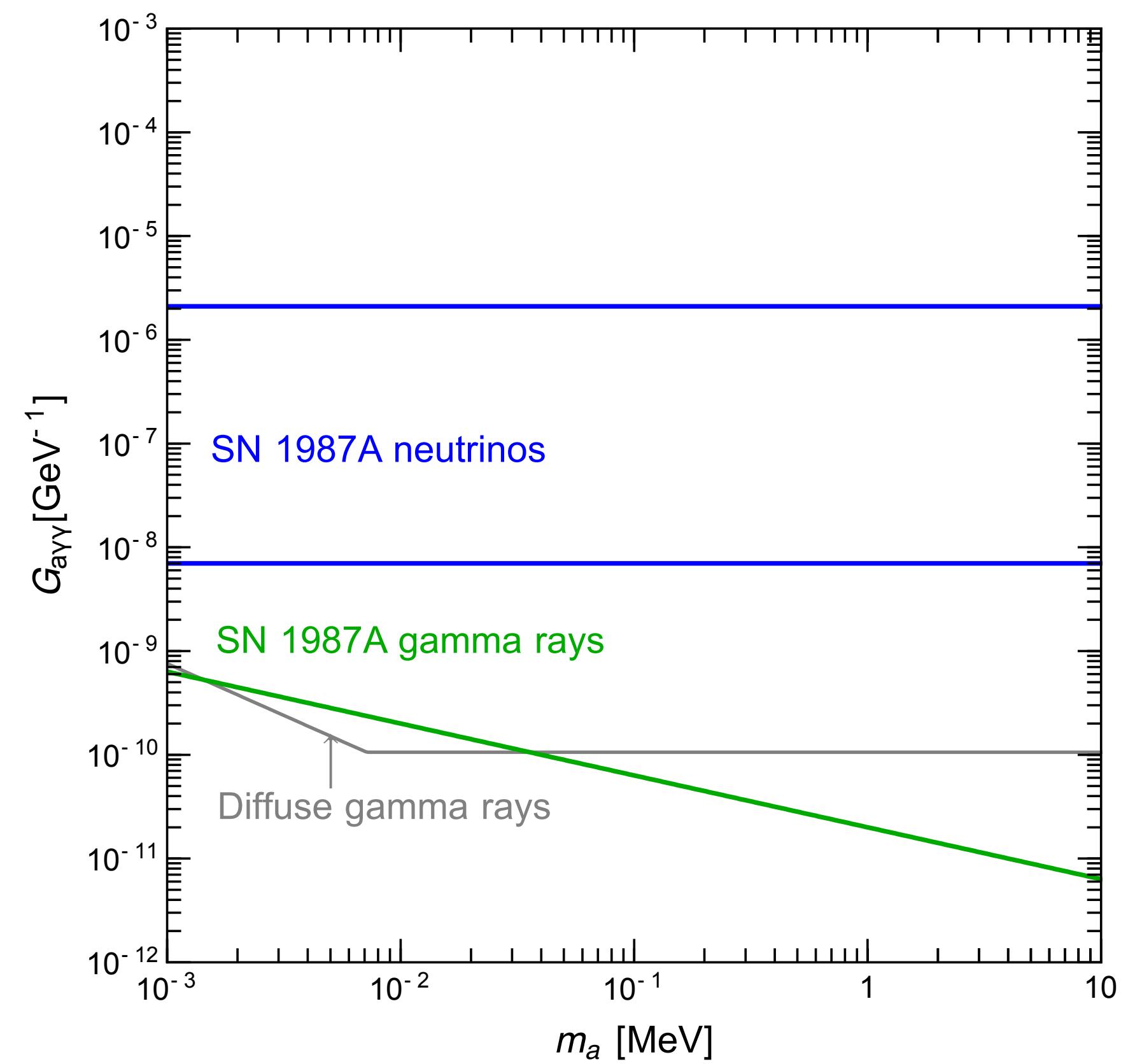
$$\frac{dn_\gamma}{d\omega} = \int_0^\infty dz(1+z)n'_{cc}(z) \int_{\omega_z}^\infty dE_z f_D(E_z) \frac{2}{E_z} F_a(E_z)$$

ALPs

Future is exciting! Gamma-ray astronomy!



See also *JCAP* 02 (2015) 006 for a small axion masses and conversion in the MW magnetic field



X-rays cross-correlation

X-ray telescopes

	eROSITA	Athena WFI	Athena X-IFU
Energy range [keV]	0.3-10	0.1-12	0.3-12
A_{eff} at 3 keV [m^2]	0.03	0.79	0.68
Ω_{Fov} [deg 2]	0.66	0.69	0.014
HEW [arcsec]	28	5	5
Spectral resolution (FWHM) at 7 keV [eV]	138	138	2.5
F_{sens} [erg cm $^{-2}$ s $^{-1}$]	1.1×10^{-14}	2.4×10^{-17}	2.4×10^{-17}
Particle bkg [counts keV $^{-1}$ s $^{-1}$ sr $^{-1}$]	1.2×10^3	1.2×10^3	5.8×10^3

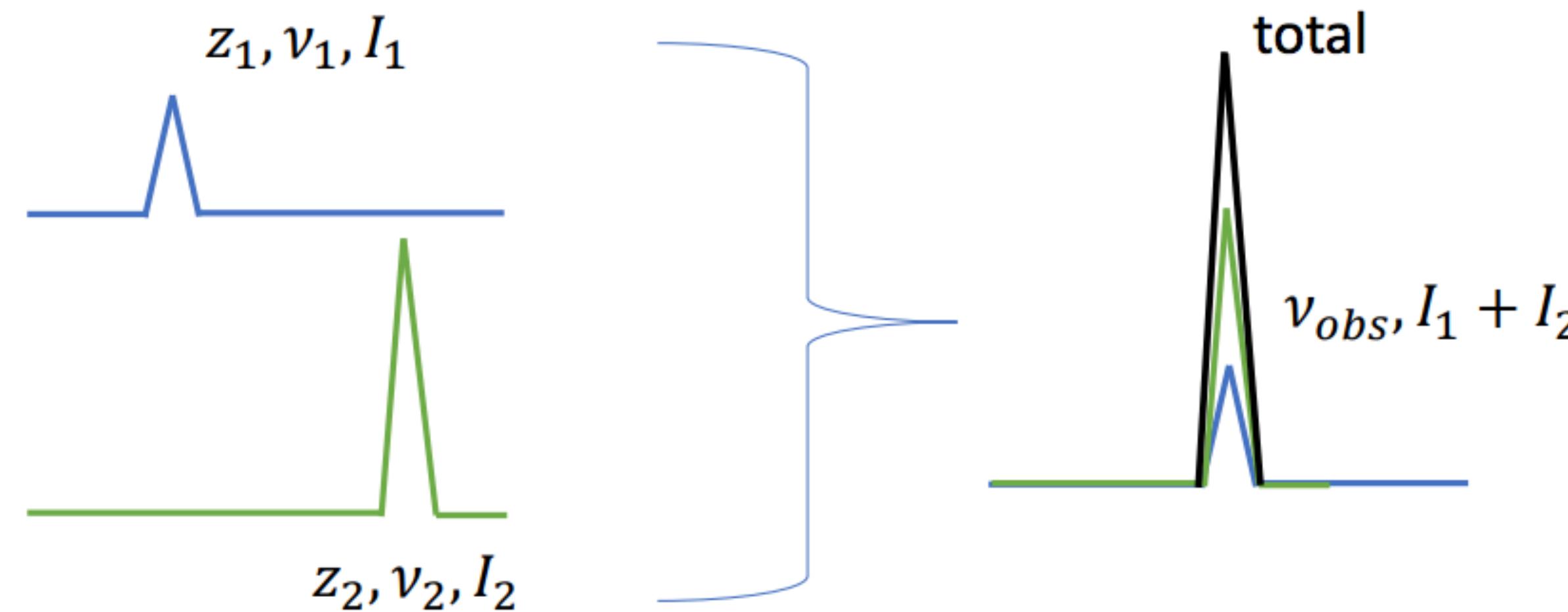
Galaxy surveys

	2MPZ	SDSS	DES	DESI	Euclid	LSST
# of galaxies	6.7×10^5	1.5×10^7	2×10^8	3.5×10^7	2×10^9	3.6×10^9
Sky coverage	0.66	0.26	0.12	0.34	0.36	0.49
z -range	0-0.08	0.08-0.8	0.08-2	0-2	0-2	0-3
# of z -bins	1	4	6	132*	7	8

Line intensity mapping

Contamination of intensity maps

- Continuous foregrounds: problem for HI surveys, less severe at higher frequencies
- **Line interlopers:** Main problem for higher freq. LIM surveys
 - $\nu_{obs} = \nu/(1+z) = \nu'/(1+z') \rightarrow$ other lines redshifted to same ν_{obs}



Contamination of intensity maps

- Continuous foregrounds: problem for HI surveys, less severe at higher frequencies
- **Line interlopers:** Main problem for higher freq. LIM surveys
 - $\nu_{obs} = \nu/(1 + z) = \nu'/(1 + z') \rightarrow$ other lines redshifted to same ν_{obs}
 - Two approaches:
 - Masking: targeted (external data) and blind (contaminated voxels are expected to be brighter)
 - Model the effect of known interlopers in the likelihood and analyses

Exotic radiative decays would be inadvertently detected as a line interloper!!

Observables

- Clustering anisotropy parametrized by monopole, dipole, quadrupole, hexadecapole in angle wrt LOS
 - Clustering along line of sight
 - Angular clustering
- Voxel-intensity distribution (VID) (one-point PDF)

Exotic radiative decays

- Decaying dark matter: $\chi \rightarrow \gamma + \gamma$

$$\nu_\gamma = m_\chi c^2 / 2 h_P$$

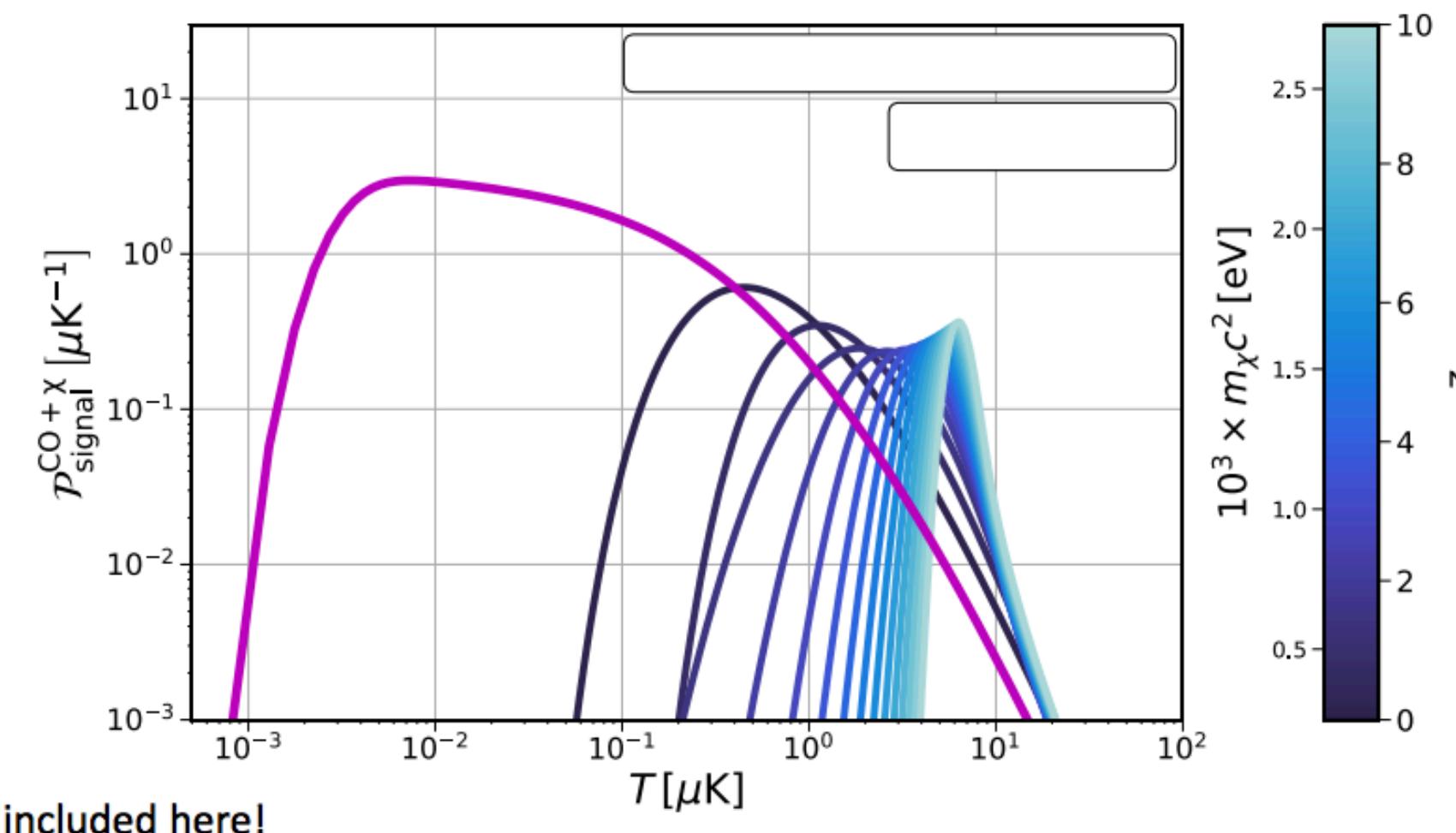
$$\rho_L^\chi(\mathbf{x}, z) = \rho_\chi(\mathbf{x}, z) c^2 \Theta_\chi \boxed{\Gamma_\chi f_\chi f_{\gamma\gamma} f_{esc}} (1 + 2\mathcal{F}_\gamma)$$

- Traces directly the DM density field

Effect in VID

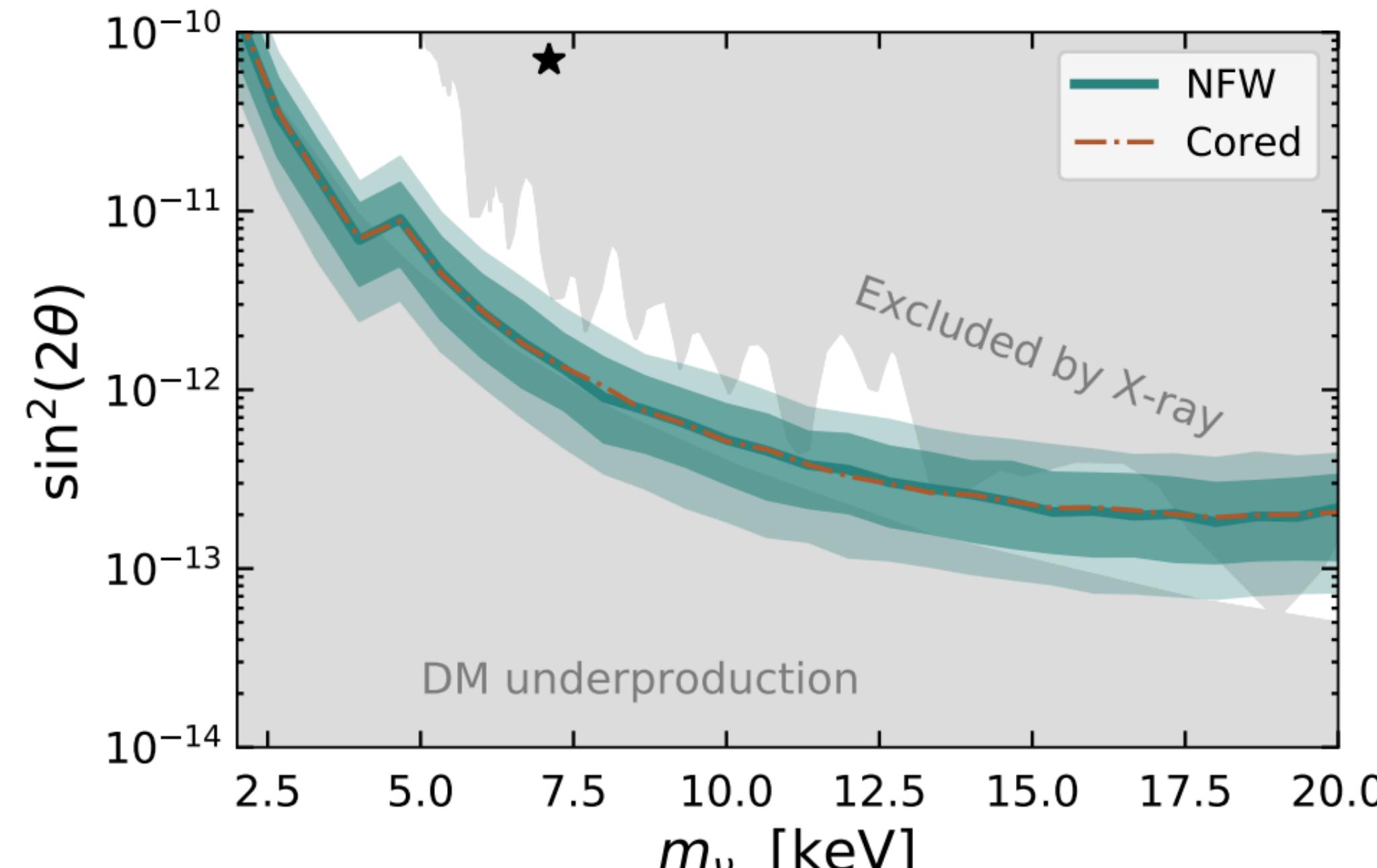
- Each voxel receives contributions from both emissions:

$$\mathcal{P}_{tot+\chi}(T) = ((\mathcal{P}_l * \mathcal{P}_\chi) * \mathcal{P}_{noise})(T); \quad \mathcal{P}_\chi = \mathcal{P}_{\tilde{\rho}} / \langle T_\chi \rangle$$



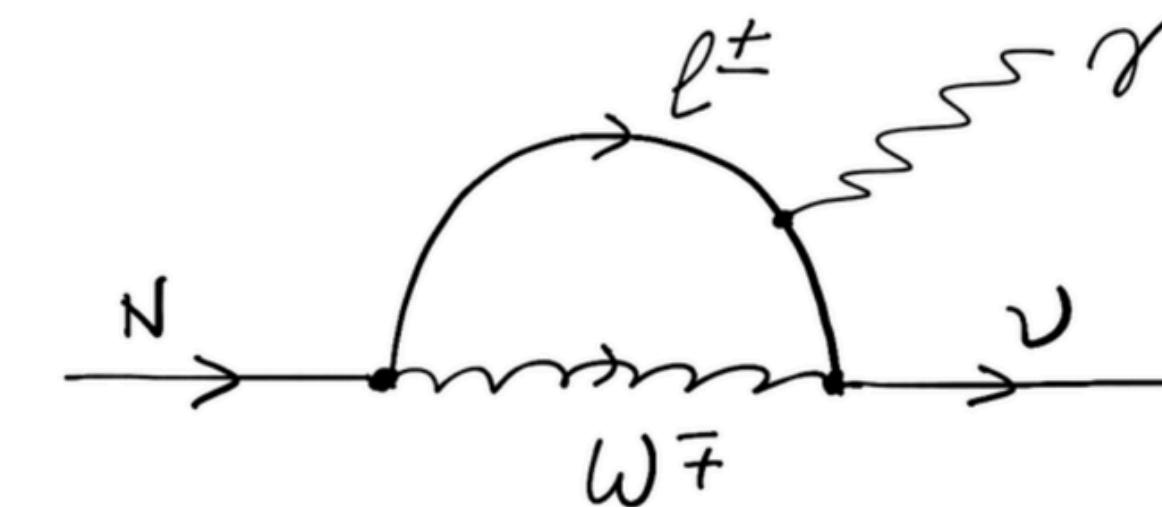
Sterile neutrinos X-ray

Point towards the galactic center and use



A. Dekker, et al (2021)

$$\mathcal{L} = \mathcal{L}_{SM} - \sum_{\alpha,i} \bar{L}^\alpha Y^{\alpha i} \tilde{\Phi} N_R^i - \sum_{i,j} \frac{1}{2} \bar{N}_R^{ic} M^{ij} N_R^j + h.c$$

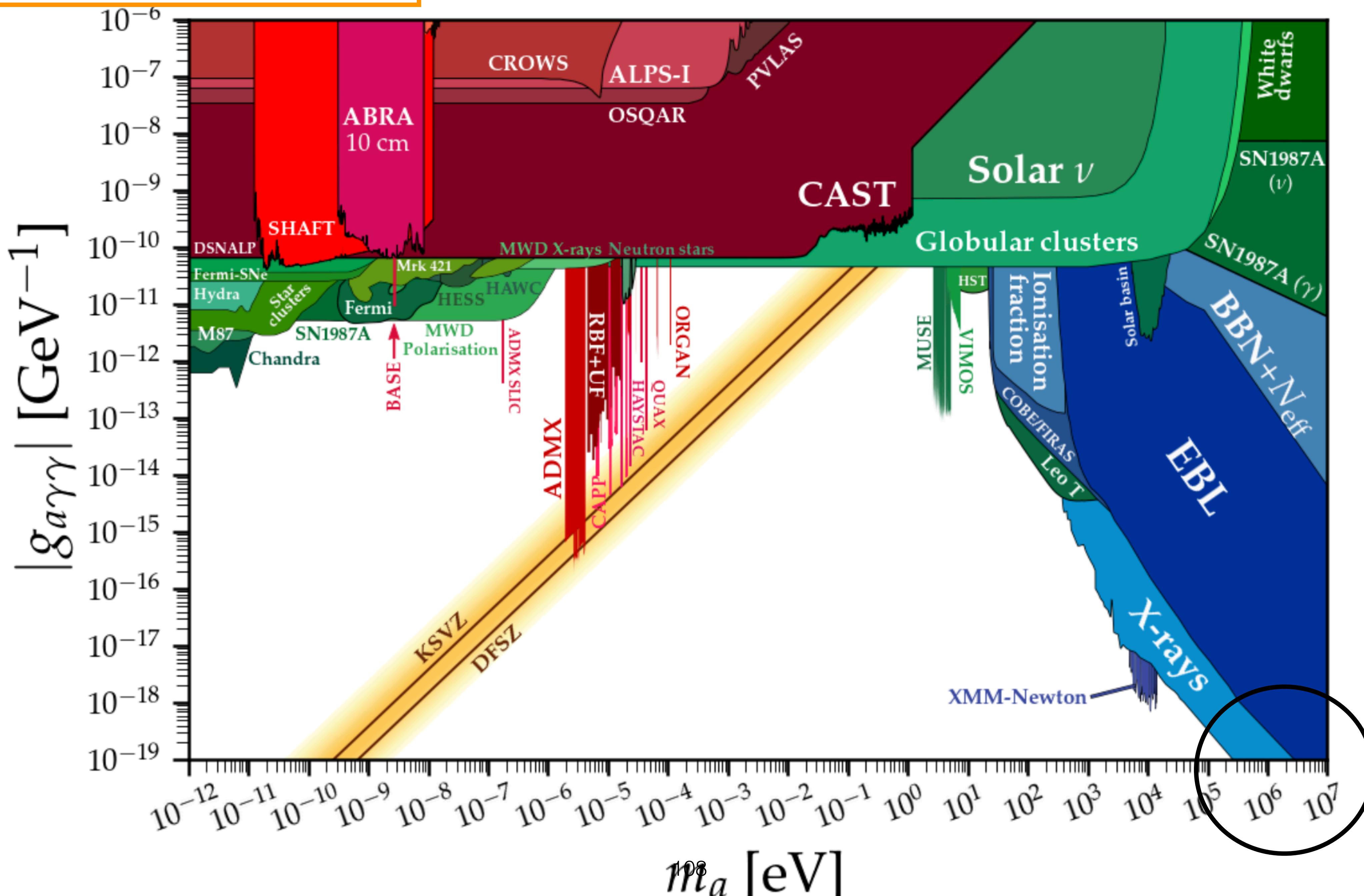


eROSITA [0.2 - 8 keV]

Launched in 2019, Russian-German mission

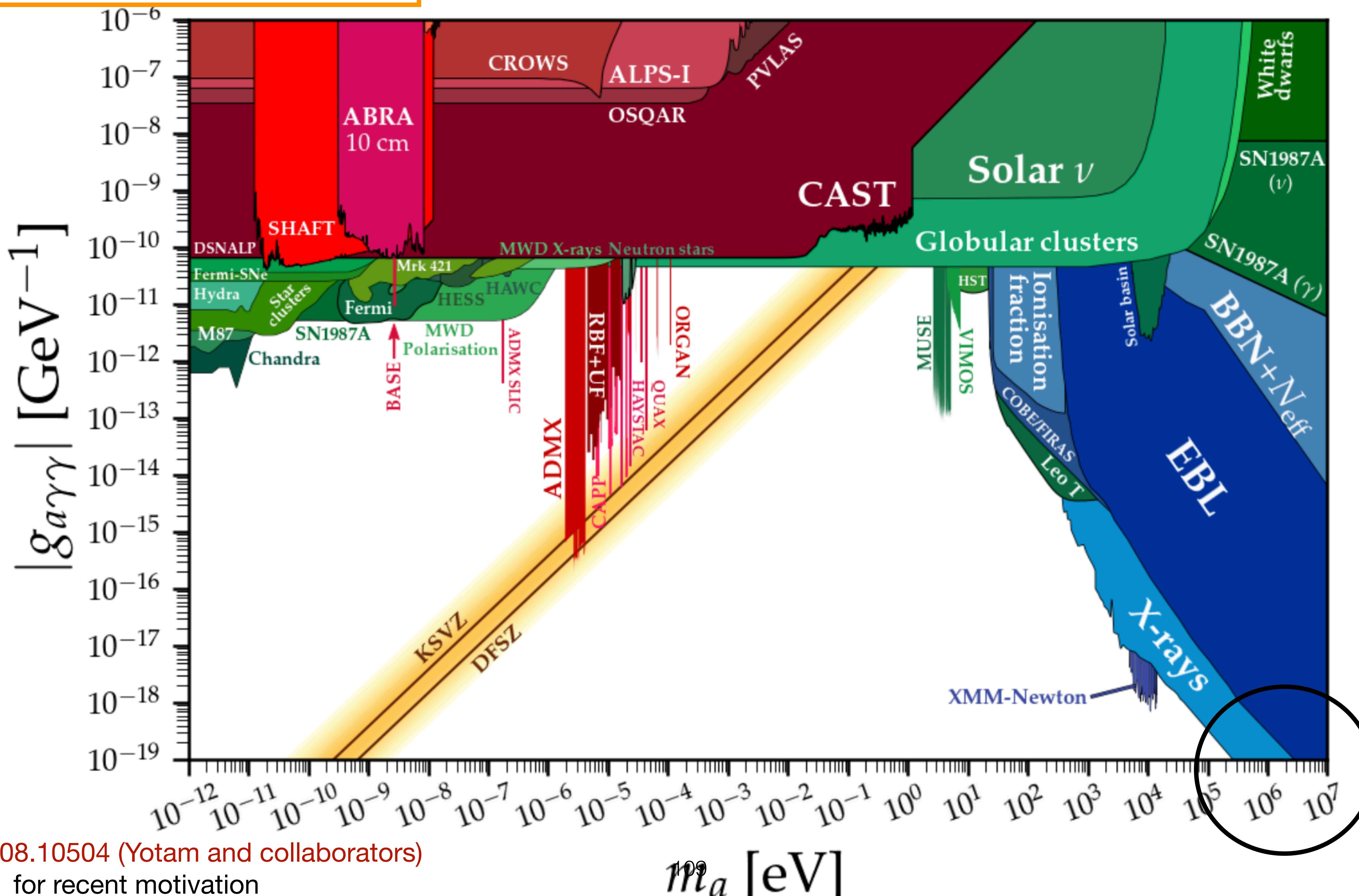
MeV (gamma-rays)

<https://github.com/cajohare/AxionLimits>



MeV (gamma-rays)

<https://github.com/cajohare/AxionLimits>



See arXiv 2208.10504 (Yotam and collaborators)
for recent motivation

NASA Selects Gamma-ray Telescope to Chart Milky Way Evolution

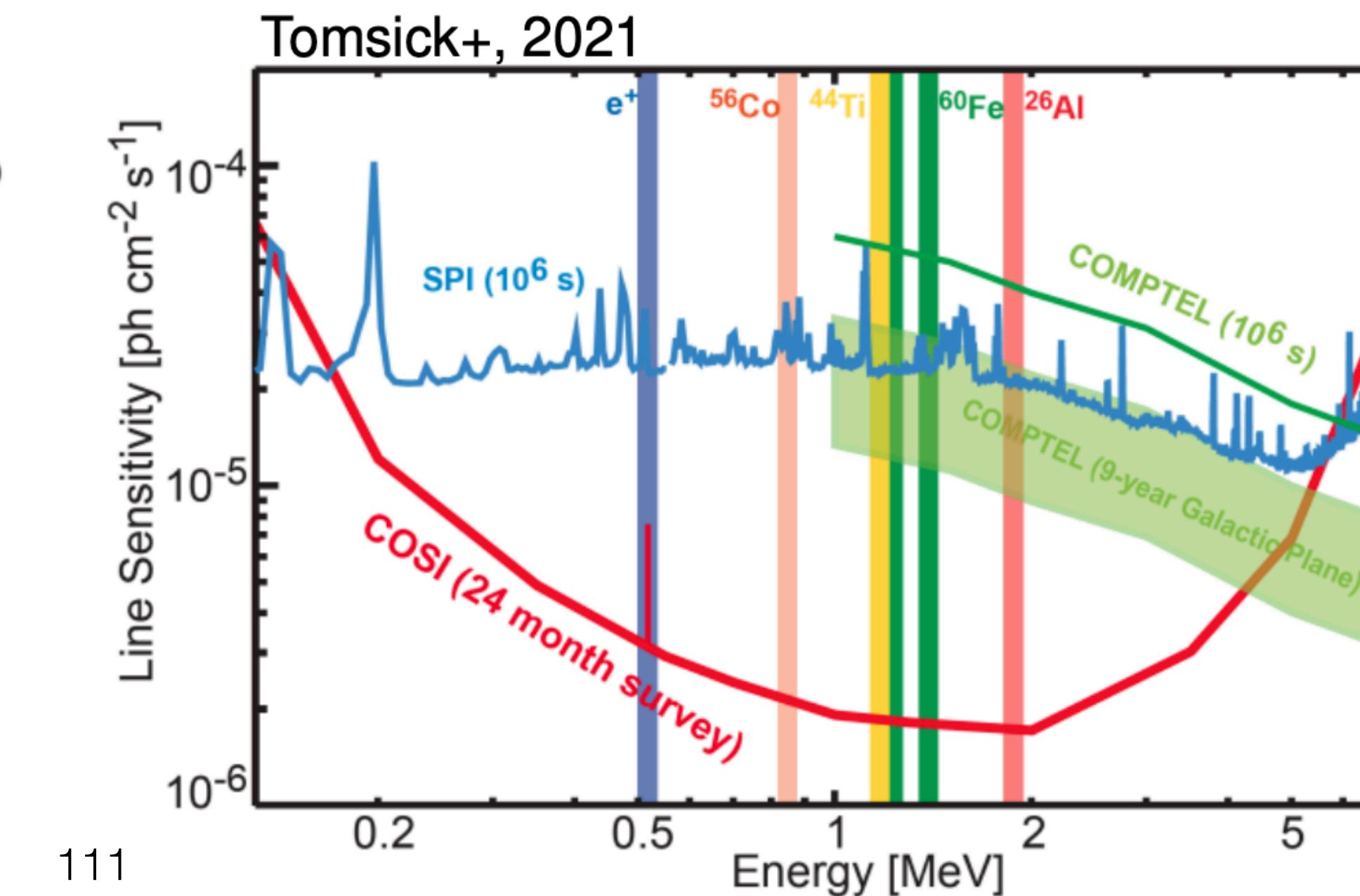
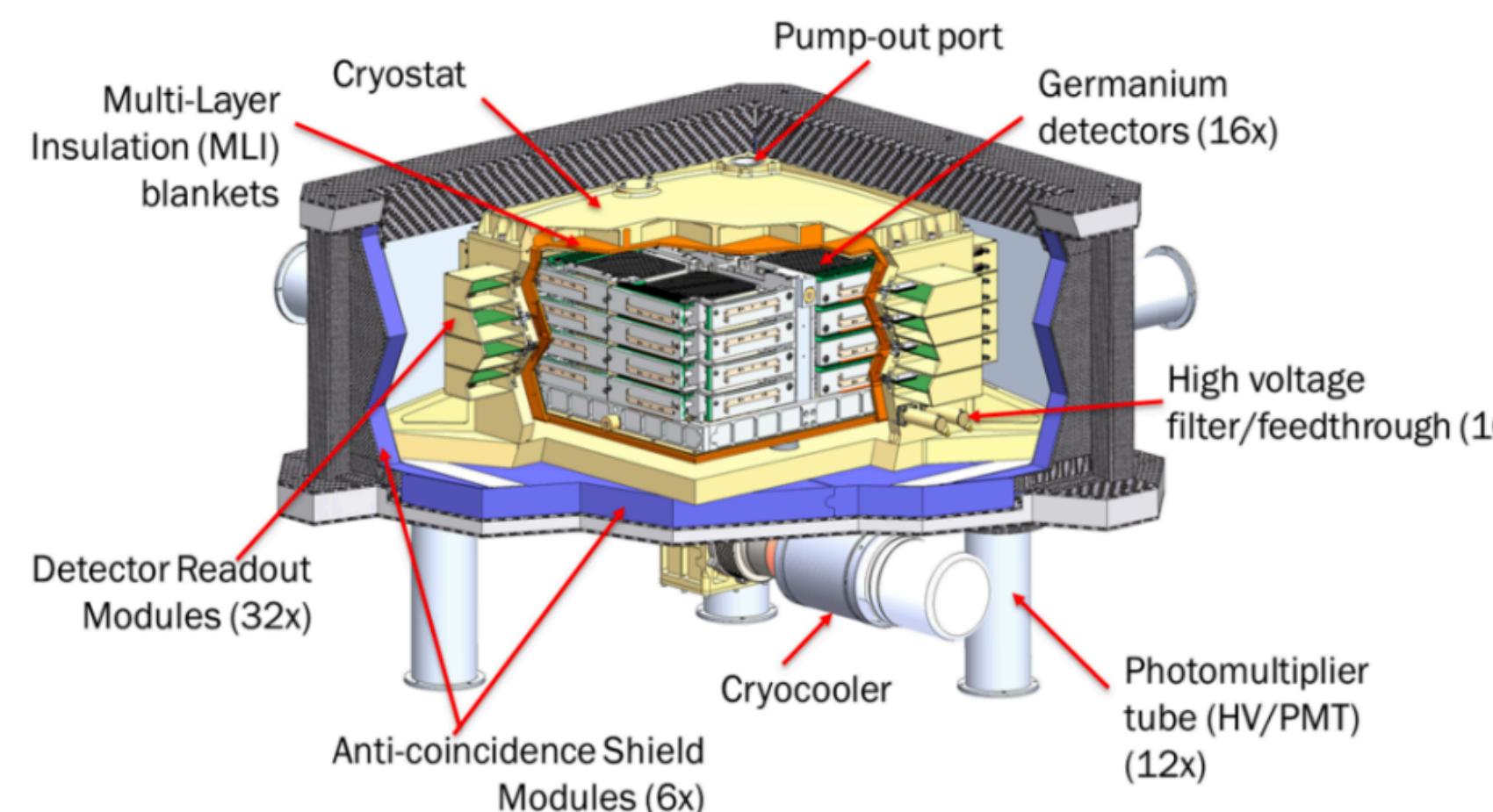
COSI-Telescope (approved by NASA)

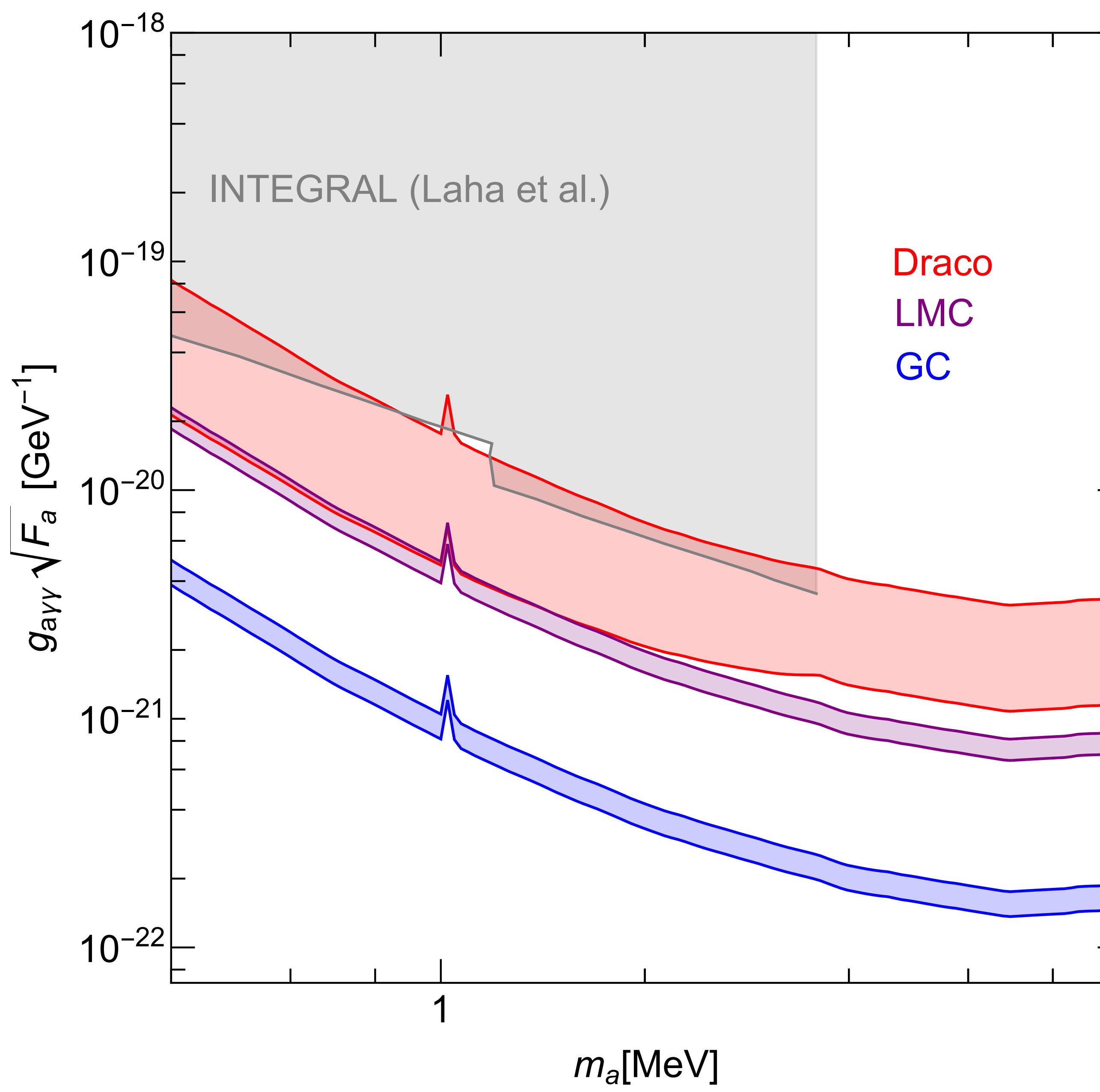
Compton Spectrometer and Imager (COSI) wide-FOV telescope designed to survey the gamma-ray sky at 0.2-5 MeV → Imaging with high-resolution spectroscopy

NASA Selects Gamma-ray Telescope to Chart Milky Way Evolution

COSI-Telescope (approved by NASA)

Compton Spectrometer and Imager (COSI) wide-FOV telescope designed to survey the gamma-ray sky at 0.2-5 MeV → Imaging with high-resolution spectroscopy





Caputo, M. Negro, M. Regis and M. Taoso, arXiv 2210.09310 (accepted JCAP)