

Resummation in $gg \rightarrow h$ via light quarks

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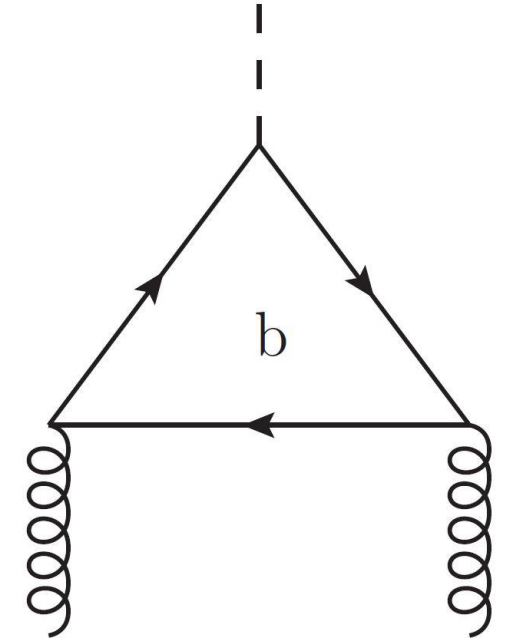
in collaboration with Z.L. Liu, M. Neubert, and X. Wang

arXiv: [2112.00018](https://arxiv.org/abs/2112.00018), [2212.10447](https://arxiv.org/abs/2212.10447)



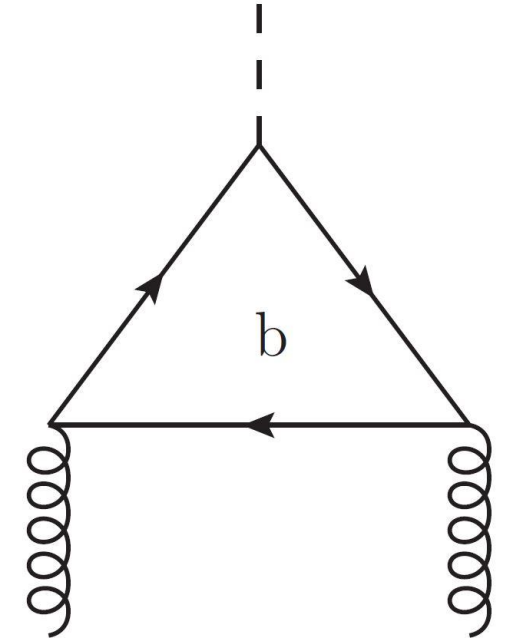
Motivation and Introduction

- Gluon fusion is dominant Higgs production channel
- Light quarks contribute $\sim 15\%$ to amplitude



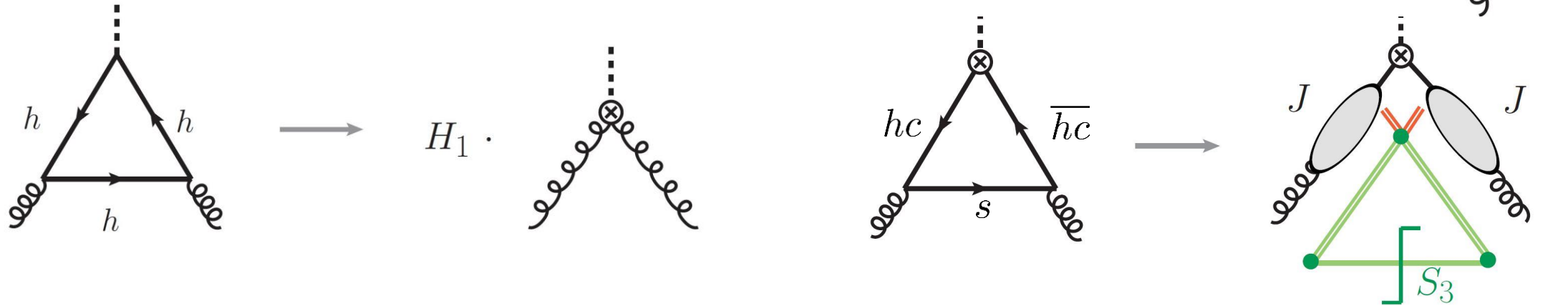
Motivation and Introduction

- Gluon fusion is dominant Higgs production channel
- Light quarks contribute ~15% to amplitude
- N-loop amplitude is $\propto \alpha_s^n L^{2n}$, $L = \ln \left(\frac{-M_h^2}{m_b^2} \right)$
 - $\alpha_s L^2 > 1 \rightarrow$ perturbation series does not converge
 - Need to resum logarithmic corrections to all order of perturbation theory
 - Apply soft-collinear effective theory (SCET)



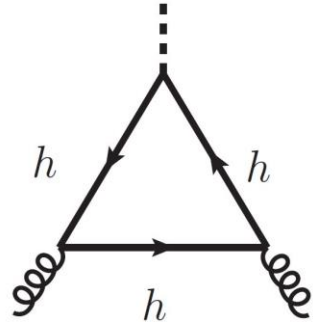
Factorisation theorem

- Amplitude factorises into three terms

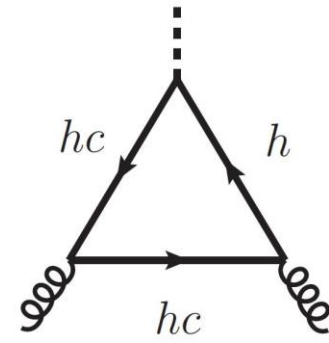
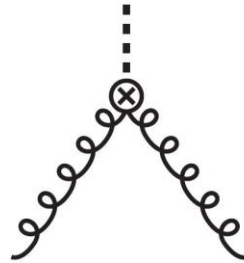


Factorisation theorem

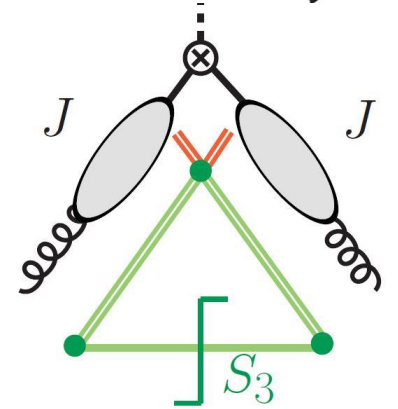
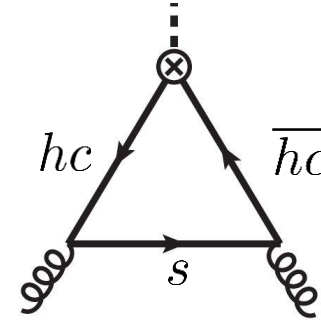
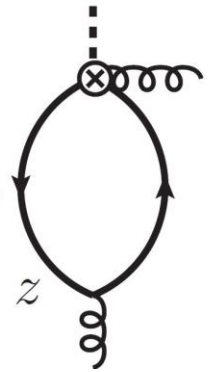
- Amplitude factorises into three terms



$H_1 \cdot$



$H_2(z) \otimes$

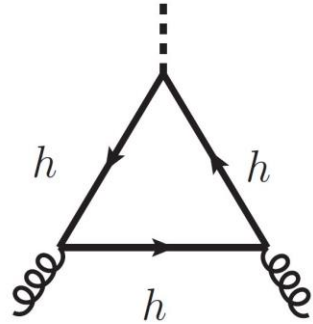


$$F_{gg}^{(0)} = \left(H_1^{(0)} \right) S_1 + 4 \int_0^1 \frac{dz}{z} \left(\bar{H}_2^{(0)}(z) S_2^{(0)}(z) \right)$$

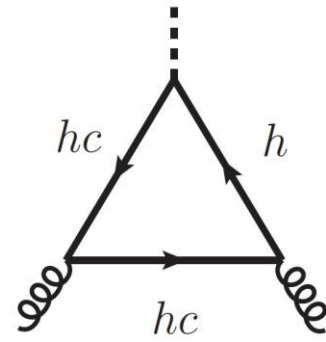
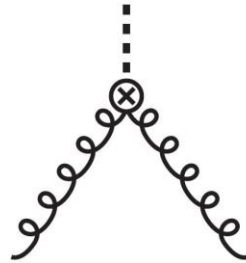
$$+ H_3^{(0)} \int_0^\infty \frac{d\ell_-}{\ell_-} \int_0^\infty \frac{d\ell_+}{\ell_+} J^{(0)}(-M_h \ell_-) J^{(0)}(M_h \ell_+) S_3^{(0)}(\ell_- \ell_+) \Big|_{\text{leading power}}$$

Factorisation theorem

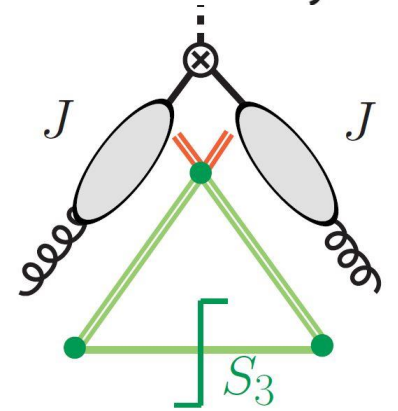
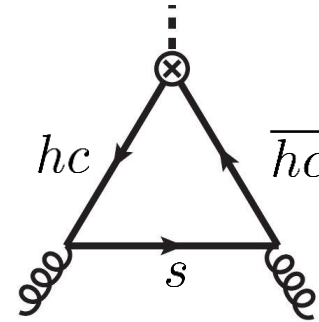
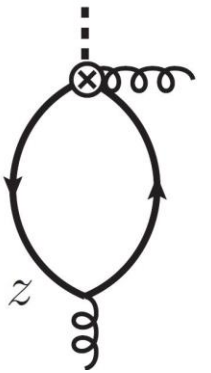
- Amplitude factorises into three terms



H_1



$H_2(z) \otimes$



$$F_{gg}^{(0)} = \left(H_1^{(0)} + \Delta H_1^{(0)} \right) S_1 + 4 \int_0^1 \frac{dz}{z} \left(\bar{H}_2^{(0)}(z) S_2^{(0)}(z) - [\bar{H}_2^{(0)}(z)] [S_2^{(0)}(z)] \right)$$

$$+ \lim_{\sigma \rightarrow -1} H_3^{(0)} \int_0^{M_h} \frac{d\ell_-}{\ell_-} \int_0^{\sigma M_h} \frac{d\ell_+}{\ell_+} J^{(0)}(-M_h \ell_-) J^{(0)}(M_h \ell_+) S_3^{(0)}(\ell_- \ell_+) \Big|_{\text{leading power}}$$

Resummation of large logarithms

- Renormalisation is non-trivial (operator mixing, regularisation and renormalisation do not commute)
- Derive RG equations and solve them to derive resummed formula

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$$\begin{aligned}
 F_{gg}(gg \rightarrow h)^{\text{RGi-LO}} = & \mathcal{M}_0(\mu_h) \frac{L^2}{2} \left\{ F_1(\rho) + \frac{1}{L} \left[\left(\frac{23}{3} \rho - 18 \right) F_3(\rho) - \frac{27}{10} \right] \right. \\
 & + \frac{1}{L^2} \left[\frac{529}{1600} \rho^2 + \frac{25}{16} \rho - \left(\frac{11}{60} \pi^2 + \frac{1577}{300} \right) \rho F_1(\rho) - \frac{327}{40} \rho F_2(\rho) \right. \\
 & \left. \left. + \left(-\frac{529}{1600} \rho^2 + \frac{1909}{480} \rho + \frac{3\pi^2}{5} + \frac{8243}{600} \right) \rho F_3(\rho) \right] \right\}
 \end{aligned}$$

special functions

$$\rho = \frac{\alpha_s}{2\pi} (C_F - C_A) L^2$$

M. Schnubel – Resummation in $gg \rightarrow h$ via light quarks

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special functions

Known from trad. QCD

Infer from $h \rightarrow \gamma\gamma$

New!

$$+ \frac{1}{L^2} \left[\frac{529}{1600} \rho^2 + \frac{25}{16} \rho - \left(\frac{11}{60} \pi^2 + \frac{1577}{300} \right) \rho F_1(\rho) - \frac{327}{40} \rho F_2(\rho) \right.$$

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**Thank you for your
attention!**

Special functions in the resummed amplitude

$$F_1(z) = {}_2F_2 \left(1, 1; \frac{3}{2}, 2; -\frac{z}{4} \right)$$

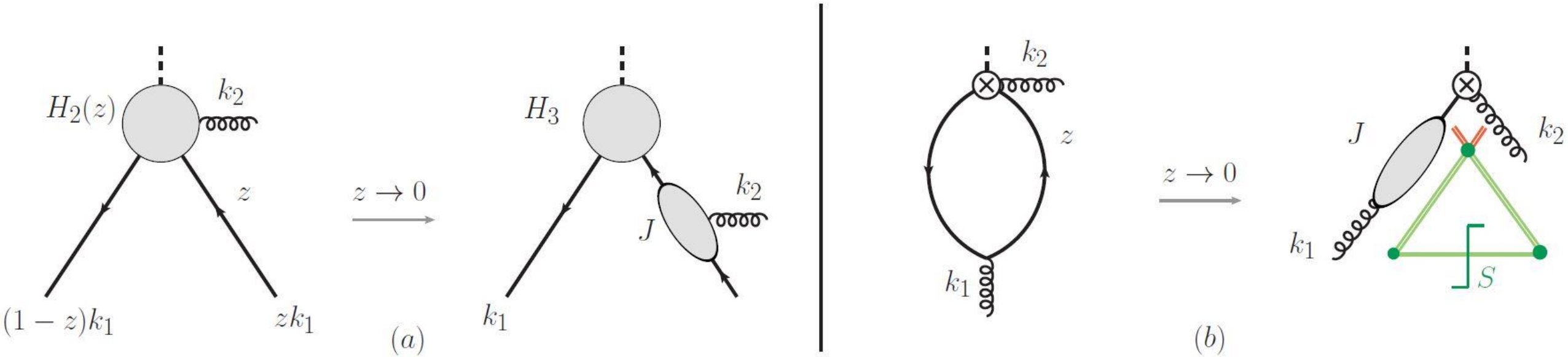
hypergeometric functions

$$F_2(z) = {}_2F_2 \left(1, 1; \frac{1}{2}, 2; -\frac{z}{4} \right)$$

$$F_3(z) = \frac{1}{\sqrt{z}} D(z) = \frac{1}{\sqrt{z}} e^{-z^2} \int_0^z dx e^{x^2}$$

Dawson integral function

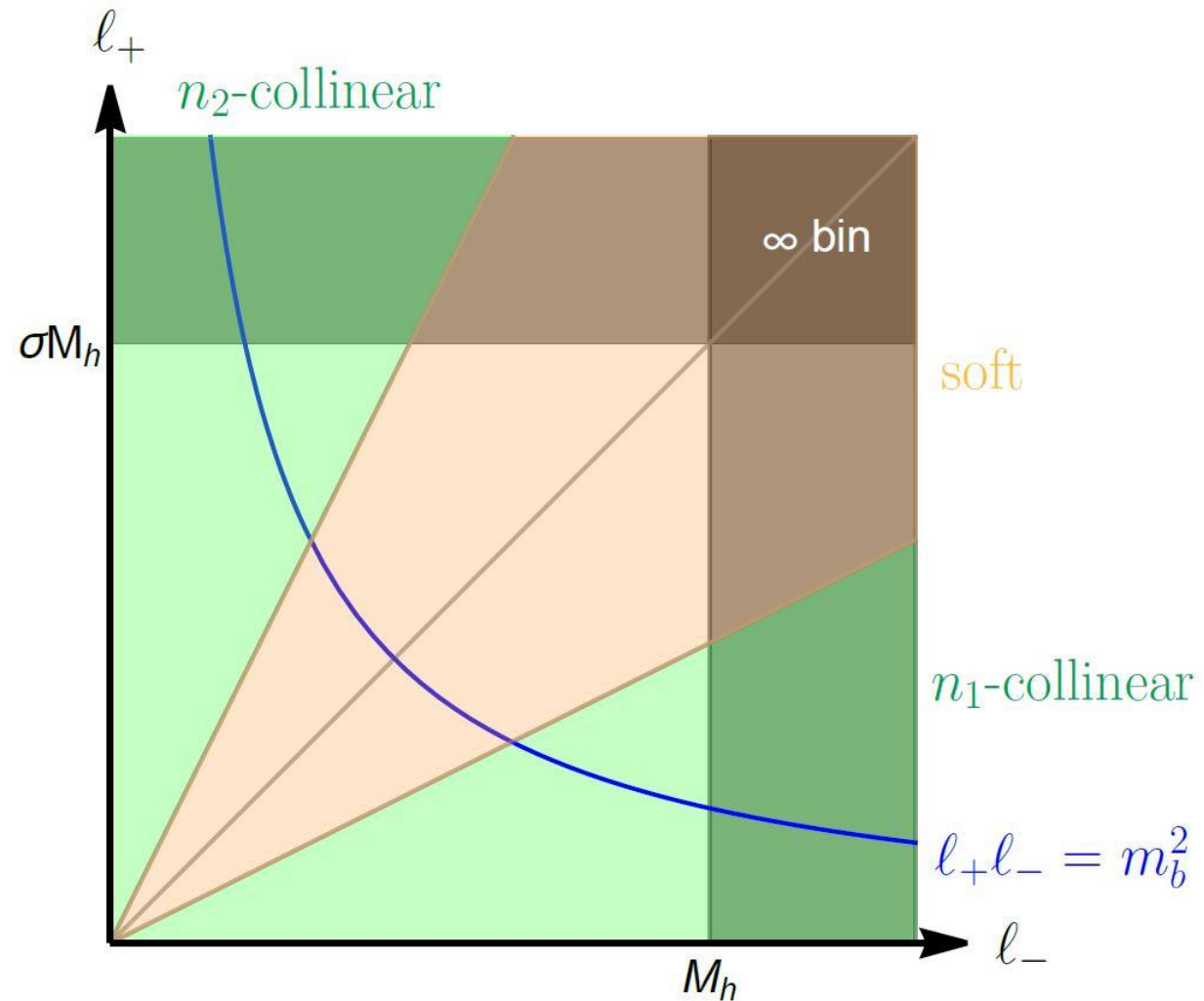
Refactorisation theorems



$$(a) \quad \llbracket \bar{H}_2(z) \rrbracket = H_3^{(0)} J^{(0)} (z M_h^2)$$

$$(b) \quad \llbracket S_2^{(0)}(z) \rrbracket = -\frac{1}{2} \int_0^\infty \frac{d\ell_+}{\ell_+} J^{(0)}(-M_h \ell_+) S_3^{(0)}(z M_h \ell_+)$$

Regularisation of endpoint divergences



Prediction of large logs at higher loops

- Solving the RGEs iteratively, we can predict the leading logs in the 3-loop amplitude:

$$\begin{aligned}
 F_{gg}(\hat{\mu}_h) = T_F \delta_{ab} \frac{\alpha_s(\hat{\mu}_h)}{\pi} \frac{m_b^2}{v} & \left\{ -2 + \frac{L^2}{2} + \frac{\alpha_s(\hat{\mu}_h)}{4\pi} \left[\frac{C_A - C_F}{12} L^4 - C_F L^3 \right. \right. \\
 & + \left(\left(1 + \frac{5\pi^2}{12} \right) C_A - \frac{2\pi^2}{3} C_F \right) L^2 + \left(\left(12 + \frac{2\pi^2}{3} + 16\zeta_3 \right) C_F - 12\zeta_3 C_A \right) L \\
 & + \left(4\zeta_3 - \frac{\pi^4}{5} - 20 \right) C_F + \left(12\zeta_3 + \frac{8\pi^4}{45} - \frac{\pi^2}{3} - 12 \right) C_A \left. \right] \\
 & + \left(\frac{\alpha_s(\hat{\mu}_h)}{4\pi} \right)^2 \left[\frac{(C_A - C_F)^2}{90} L^6 + (C_A - C_F) \left(\frac{\beta_0}{30} - \frac{C_F}{10} \right) L^5 \right. \\
 & \left. \left. + d_4^{\text{OS}} L^4 + d_3^{\text{OS}} L^3 + \dots \right] \right\},
 \end{aligned}$$