

Naturalness, Gegenbauer polynomials, and the Higgs self-coupling

Gauthier Durieux
(CERN)

Gegenbauer Goldstones, JHEP 01 (2022) 076, [2110.06941]

Gegenbauer's Twin, JHEP 05 (2022) 140, [2202.01228]

Charting the Higgs self-coupling boundaries, JHEP 12 (2022) 148, [2209.00666]

with Matthew McCullough and Ennio Salvioni



The Higgs at 10

Ten years of LHC measurements
exacerbated the naturalness puzzle.

Small mass?

SM-like couplings?

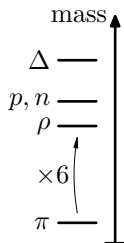
Composite Higgs

realise the Higgs as the
pseudo-Nambu-Goldstone boson (pNGB)
of a new strong sector

e.g. global $SO(5) \rightarrow SO(4)$ spontaneous breaking
at scale f

small mass obtained from the
explicit breaking of $SO(5)$
by e.g. the SM

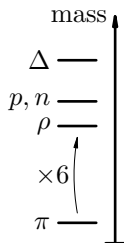
SM-like couplings require $v^2/f^2 \ll 1$
which requires tuning in minimal models



Composite Higgs

realise the Higgs as the
pseudo-Nambu-Goldstone boson (pNGB)
of a new strong sector

e.g. global $SO(5) \rightarrow SO(4)$ spontaneous breaking
at scale f



small mass in the
exp (5)
SM

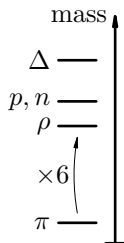
small-ish mass

SM-like couplings require $v^2/f^2 \ll 1$
which requires tuning in minimal models

Composite Higgs

realise the Higgs as the
pseudo-Nambu-Goldstone boson (pNGB)
of a new strong sector

e.g. global $SO(5) \rightarrow SO(4)$ spontaneous breaking
at scale f



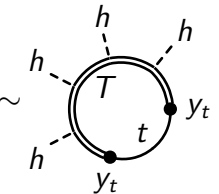
small mass in the
exp (5)

small-ish mass

SM-like couplings $v^2/f^2 \ll 1$
which require small models

small vev?

Minimal composite Higgs

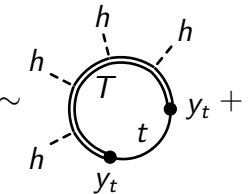


$$V(h) \sim \text{loop} + \dots \sim \kappa \frac{y_t^2 N_c}{16\pi^2} f^2 M_T^2 \left(-\sin^2 \frac{h}{f} + \delta \sin^4 \frac{h}{f} \right)$$

$$\rightarrow \frac{v^2}{f^2} = \sin^2 \frac{\langle h \rangle}{f} = \frac{1}{2\delta} \quad \text{vs.} \quad \frac{C_{hVV}}{C_{hVV}^{\text{SM}}} = \sqrt{1 - \frac{v^2}{f^2}} \gtrsim 0.94$$

$$\rightarrow m_h^2 = \kappa \frac{y_t^2 N_c}{4\pi^2} M_T^2 \left(1 - \frac{1}{2\delta} \right) \quad \text{vs.} \quad M_T \gtrsim 1.3 \text{ TeV}$$

Minimal composite Higgs



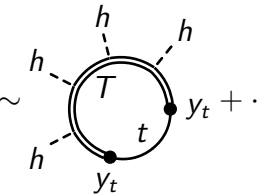
$$V(h) \sim \dots \sim \kappa \frac{y_t^2 N_c}{16\pi^2} f^2 M_T^2 \left(-\sin^2 \frac{h}{f} + \delta \sin^4 \frac{h}{f} \right)$$

$$\rightarrow \frac{v^2}{f^2} = \sin^2 \frac{\langle h \rangle}{f} = \frac{1}{2\delta} \quad \boxed{1/\delta \lesssim 0.23} \quad \frac{v}{M_T} = \sqrt{1 - \frac{v^2}{f^2}} \gtrsim 0.94$$

$$\rightarrow m_h^2 = \kappa \frac{y_t^2 N_c}{4\pi^2} M_T^2 \left(1 - \frac{1}{2\delta} \right) \quad \boxed{\kappa \lesssim 0.14} \quad M_T \gtrsim 1.3 \text{ TeV}$$

Few percent tuning wrt $1/\delta \gtrsim 1$, $\kappa \simeq 1$ expectation

Minimal composite Higgs



The diagram shows a top quark loop with two top quark vertices labeled y_t and a top quark line labeled t . Four Higgs boson lines labeled h are attached to the loop. The loop is labeled T . The potential is given by:

$$V(h) \sim \dots \sim \kappa \frac{y_t^2 N_c}{16\pi^2} f^2 M_T^2 \left(-\sin^2 \frac{h}{f} + \delta \sin^4 \frac{h}{f} \right)$$

Gegenbauer fix

$$\rightarrow \frac{v^2}{f^2} = \sin^2 \frac{\langle h \rangle}{f} = \frac{1}{2\delta} \quad 1/\delta \lesssim 0.23 \quad \frac{v}{M_T} = \sqrt{1 - \frac{v^2}{f^2}} \gtrsim 0.94$$

$$\rightarrow m_h^2 = \kappa \frac{y_t^2 N_c}{4\pi^2} M_T^2 \left(1 - \frac{1}{2\delta} \right) \quad \kappa \lesssim 0.14 \quad M_T \gtrsim 1.3 \text{ TeV}$$

Twin-Higgs fix

Few percent tuning wrt $1/\delta \gtrsim 1$, $\kappa \simeq 1$ expectation

Naturally small vev

=
radiatively stable potential
with deepest minimum close to the origin

agnostic of the UV
considering the $SO(5) \rightarrow SO(4)$ pNGB EFT
with $\vec{\phi} = \left(\frac{\vec{h}}{h} \sin \frac{h}{f}, \cos \frac{h}{f} \right)$ where $h \equiv |\vec{h}|$

Radiatively stable potentials

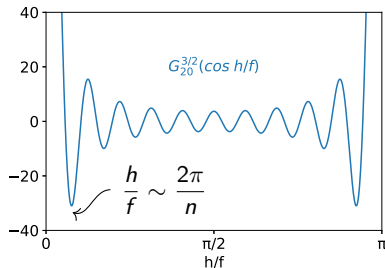
Explicit $SO(5) \rightarrow SO(4)$ breaking by an irrep spurion
(symmetric traceless)

$$K^{i_1 \dots i_n} \phi_{i_1} \dots \phi_{i_n}$$

No other invariant, linear in K , can be constructed,
so all-loop linear renormalisation can only be multiplicative.

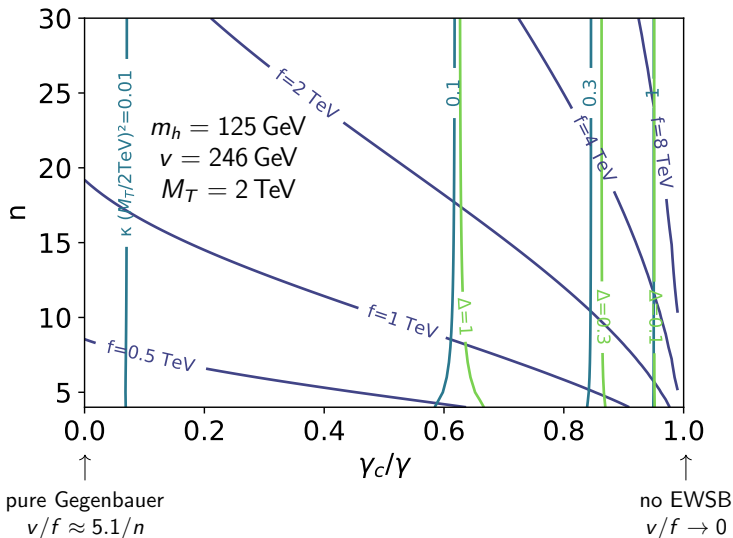
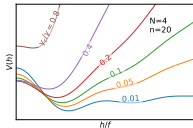
Obtain Gegenbauer polynomials:

$$K^{i_1 \dots i_n} \phi_{i_1} \dots \phi_{i_n} = G_n^{3/2}(\cos \frac{h}{f})$$



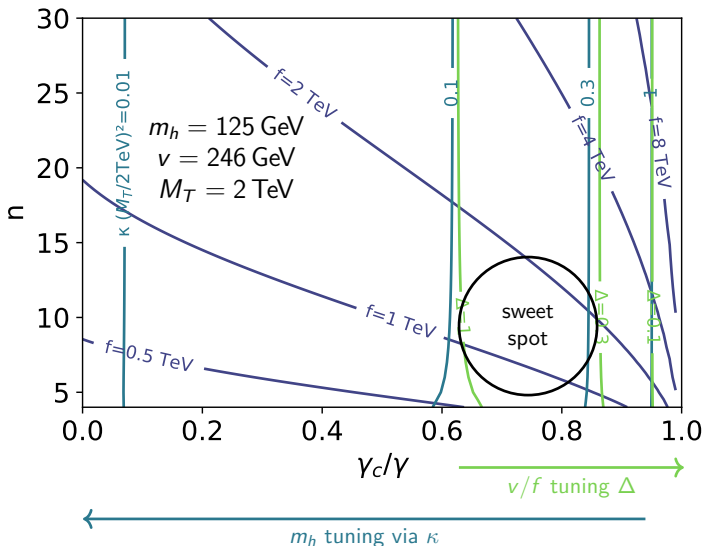
The Gegenbauer Higgs

$$V(h) = \kappa \frac{N_c y_t^2}{16\pi^2} f^2 M_T^2 \left[\sin^2 \frac{h}{f} + \gamma G_n^{3/2} \left(\cos \frac{h}{f} \right) \right]$$



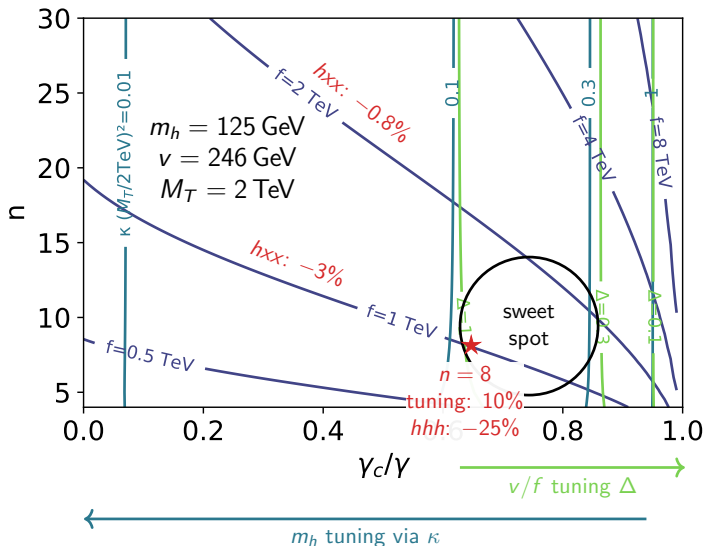
The Gegenbauer Higgs

$$V(h) = \kappa \frac{N_c y_t^2}{16\pi^2} f^2 M_T^2 \left[\sin^2 \frac{h}{f} + \gamma G_n^{3/2} \left(\cos \frac{h}{f} \right) \right]$$



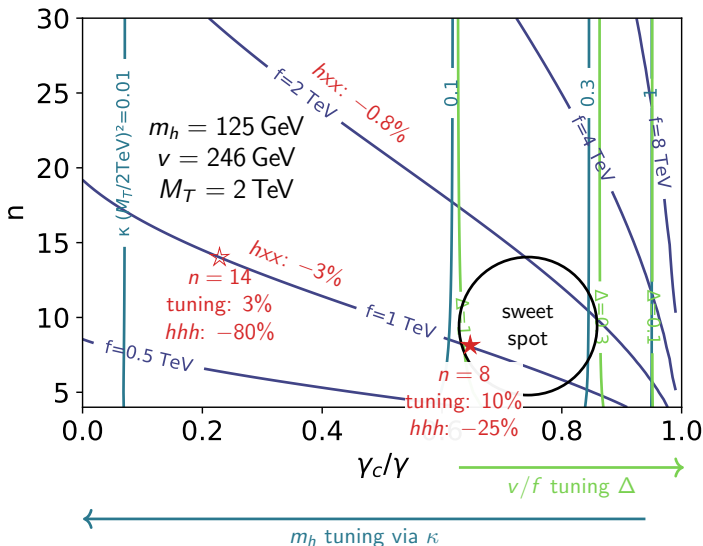
The Gegenbauer Higgs

$$V(h) = \kappa \frac{N_c y_t^2}{16\pi^2} f^2 M_T^2 \left[\sin^2 \frac{h}{f} + \gamma G_n^{3/2} \left(\cos \frac{h}{f} \right) \right]$$

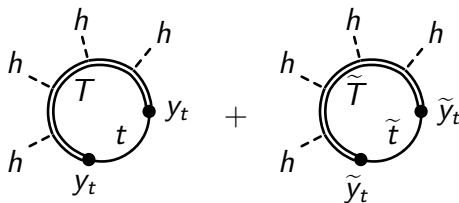


The Gegenbauer Higgs

$$V(h) = \kappa \frac{N_c y_t^2}{16\pi^2} f^2 M_T^2 \left[\sin^2 \frac{h}{f} + \gamma G_n^{3/2} \left(\cos \frac{h}{f} \right) \right]$$



Adding twins

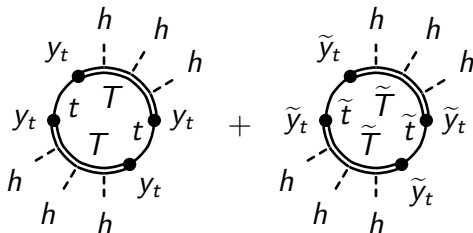


$$\frac{N_c y_t^2}{16\pi^2} f^2 M_T^2 \sin^2 \frac{h}{f} + \frac{N_{\tilde{c}} \tilde{y}_t^2}{16\pi^2} f^2 M_{\tilde{T}}^2 \cos^2 \frac{h}{f}$$

no M_T^2 sensitivity

if twin parity enforces $y_t = \tilde{y}_t$ and $M_T = M_{\tilde{T}}$

Adding twins



$$\frac{N_c y_t^4}{16\pi^2} f^4 \sin^4 \frac{h}{f} \log M_T \quad + \quad \frac{N_c \tilde{y}_t^4}{16\pi^2} f^4 \cos^4 \frac{h}{f} \log M_{\tilde{T}}$$

retaining $\log M_T$ sensitivity only

Gegenbauer's Twin

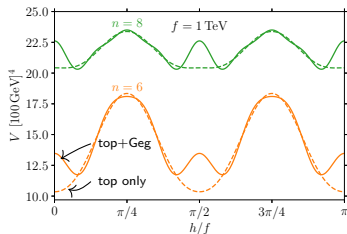
- global $SO(8) \supset SO(4) \times \widetilde{SO(4)}$
- spontaneous $SO(8) \rightarrow SO(7)$
 - 7 NGBs
 - 6 eaten by W^\pm, Z and $\widetilde{W}^\pm, \widetilde{Z}$
 - 1 Higgs: $\vec{\phi} = (\vec{0}_3, \sin \frac{h}{f}; \vec{0}_3, \cos \frac{h}{f})^T$ in unitary gauge



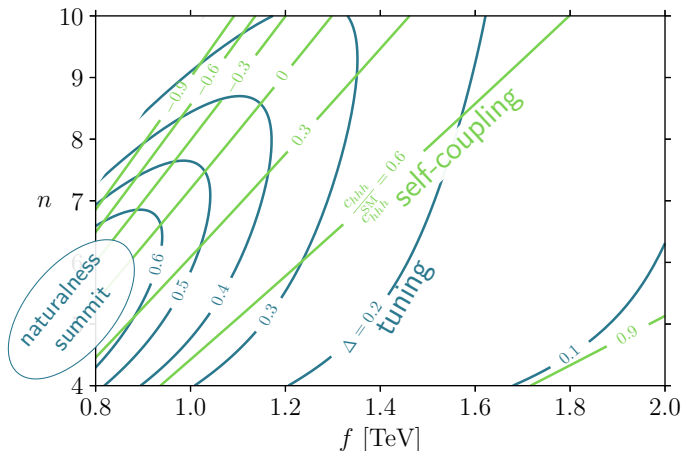
L.B. Gegenbauer
1849–1903

- minimal explicit breaking is insufficient

- explicit $SO(8) \rightarrow SO(4) \times \widetilde{SO(4)}$
 - radiative stability from irrep spurion
 - $G_n^{3/2}(\cos \frac{2h}{f})$ potential



Gegenbauer's Twin

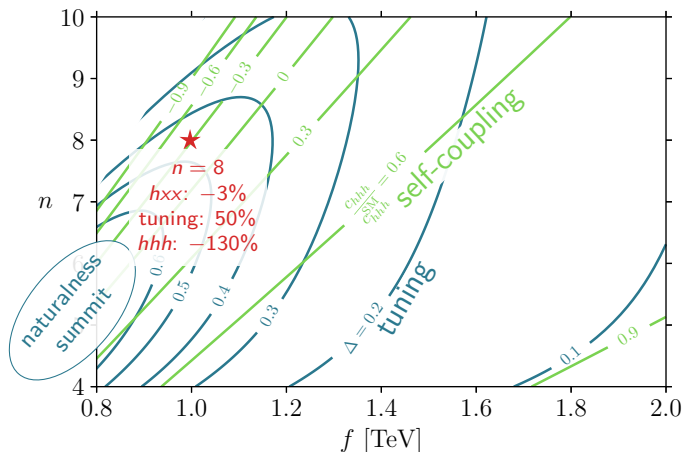


HL-LHC 2σ reach: hVV : 2.6%
 hhh : 100%

[ECFA report '19]

(and possibly large M_T , with unitarity violating H scattering towards 6 TeV)

Gegenbauer's Twin



HL-LHC 2σ reach: hVV : 2.6%
 hhh : 100%

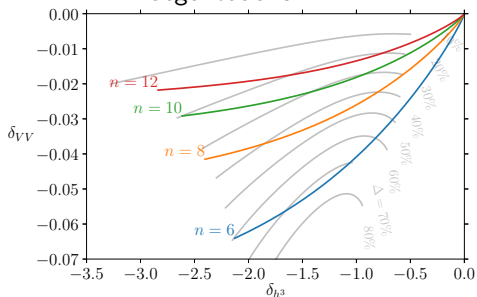
[ECFA report '19]

(and possibly large M_T , with unitarity violating H scattering towards 6 TeV)

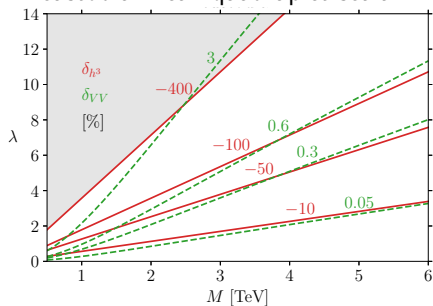
Structurally large hhh/hxx

see also: [Di Luzio, Gröber, Spannowsky '17]
 [Gupta, Rzehak, Wells '13] [Falkowski, Rattazzi '19]
 [Logan, Rentalá '15] [Chala, Krause, Nardini '18]
 [etc.]

Gegenbauer's Twin



custodial weak-quadruplet scalar



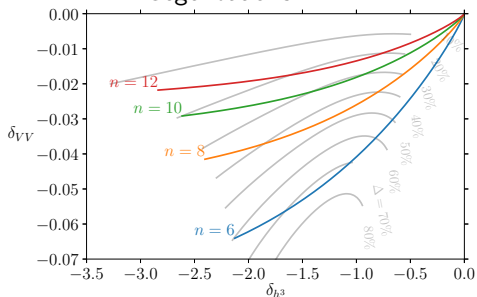
$$\lambda H^* H^* (\epsilon H) \Phi + \lambda \frac{1}{\sqrt{3}} H^* H^* H^* \tilde{\Phi}$$

- a loop factor allowed dimensionally (or v^2/M^2 if dim-6/dim-8)
- $\text{dim} \gg 6$ operators may be very relevant
- vacuum stability constrains large self-couplings

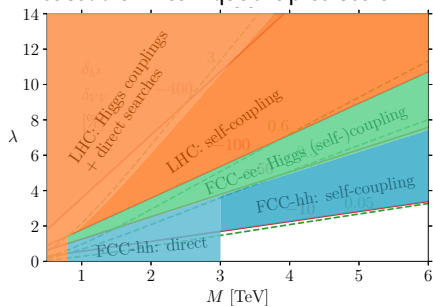
Structurally large hhh/hxx

see also: [Di Luzio, Gröber, Spannowsky '17]
 [Gupta, Rzehak, Wells '13] [Falkowski, Rattazzi '19]
 [Logan, Rental '15] [Chala, Krause, Nardini '18]
 [etc.]

Gegenbauer's Twin



custodial weak-quadruplet scalar



$$\lambda H^* H^* (\epsilon H) \Phi + \lambda \frac{1}{\sqrt{3}} H^* H^* H^* \tilde{\Phi}$$

- a loop factor allowed dimensionally (or v^2/M^2 if dim-6/dim-8)
- $\text{dim} \gg 6$ operators may be very relevant
- vacuum stability constrains large self-couplings

Naturalness, Gegenbauer polynomials, and the Higgs self-coupling

Gegenbauer potentials are eigenfunctions of linear renorm.
for $SO(N + 1) \rightarrow SO(N)$ pNGBs
(on spherical manifolds, more generally).

Explicit global symmetry breakings in non-minimal irreps
lead to naturally small pNGB Higgs vev's.

Structurally large hhh/hxx ratios in coupling deviations occur,
making self-coupling measurements essential probes.

Extras

Radiatively stable potentials (II)

Linear one-loop correction to $V(\frac{h}{f})$:

$$\frac{\Lambda^2}{32\pi^2 f^2} \left(V''' + 3 \cot \frac{h}{f} V' \right)$$

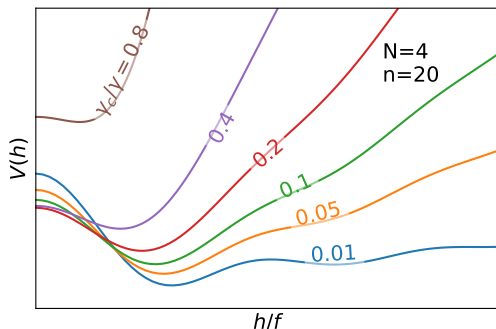
Radiative stability at one-loop and linear order order if $\propto V$

Differential equation of Gegenbauer polynomials

$$V\left(\frac{h}{f}\right) \propto G_n^{3/2}\left(\cos \frac{h}{f}\right)$$

Gegenbauer Higgs potential

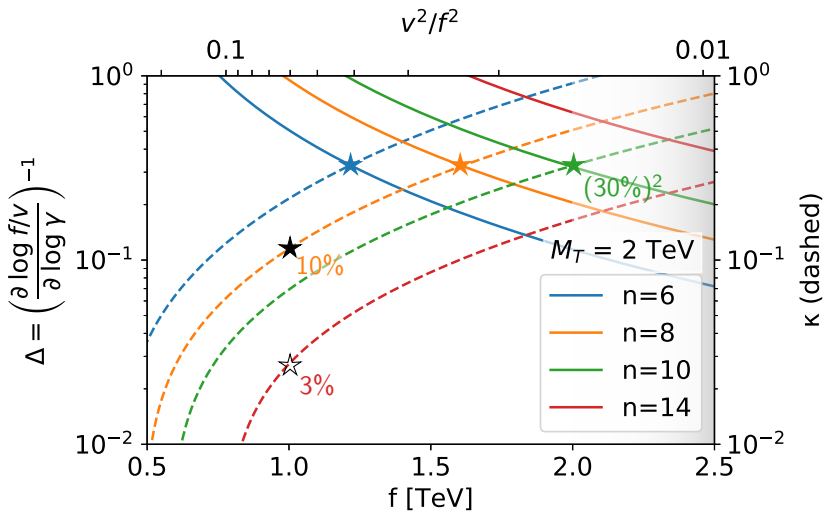
$$V(h) = \kappa \frac{N_c y_t^2}{16\pi^2} f^2 M_T^2 \left[\sin^2 \frac{h}{f} + \gamma G_n^{3/2} \left(\cos \frac{h}{f} \right) \right]$$



$$v/f \rightarrow 0 \quad \text{as} \quad \gamma \rightarrow \gamma_c$$

$$\frac{m_h^2}{\kappa \frac{N_c y_t^2}{16\pi^2} M_T^2} \rightarrow \quad \text{as} \quad \gamma \rightarrow \gamma_c \quad \text{relaxing} \quad \kappa \rightarrow 1$$

Gegenbauer Higgs tunings

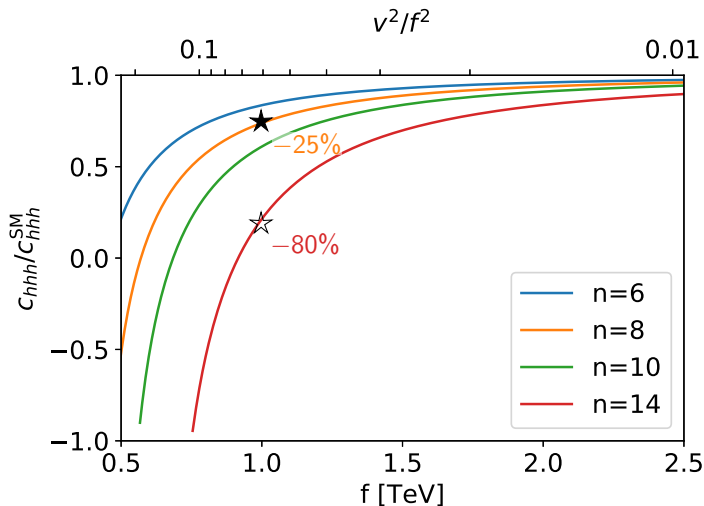


$$\Delta \approx 30\% \left(\frac{f}{4v} \frac{5.1}{n}\right)^{-2.1}$$

$$\kappa \approx 30\% \left(\frac{f}{4v} \frac{5.1}{n} \frac{2 \text{ TeV}}{M_T}\right)^2$$

$(M_T \gtrsim y_t f / \sqrt{2} \text{ since } m_t \sim M_T v / f \text{ for } y_t f \gg M_T)$

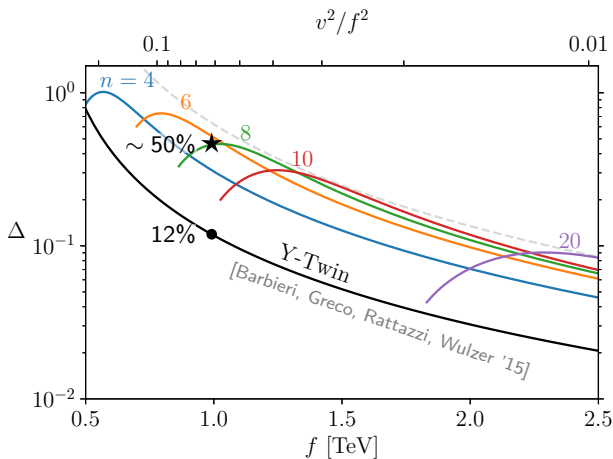
Gegenbauer Higgs self-coupling



$$\frac{C_{hhh}}{C_{hhh}^{SM}} \approx 1 - 1.2 \left(\frac{f}{v} \frac{5.1}{n + \lambda} \right)^{-2} \quad \text{when close to 1}$$

Gegenbauer's Twin tunings

- conservative definition RMS(eig. log-derivative matrix)
- dominated by top-sector dependence of v/f
- about 4 times better than usual $\Delta \approx 2v^2/f^2$ minimum



Gegenbauer's Twin self-coupling

