

Parameter space for testable leptogenesis

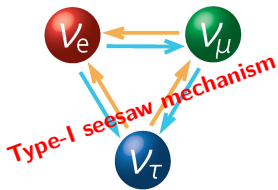
Yannis Georis

based on work in collaboration with M. Drewes and J. Klarić
[arXiv:2106.16226]

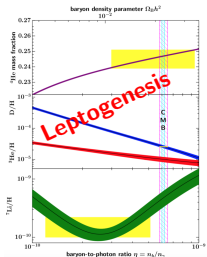
57th Rencontres de Moriond: EW session
March 20, 2023



Right-handed neutrinos (RHN)



Neutrino masses

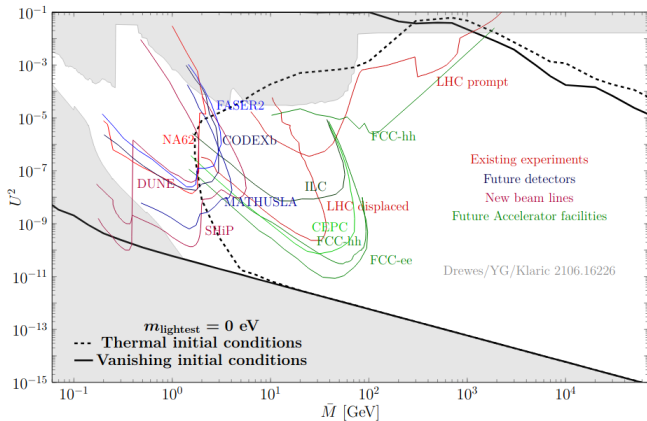


Baryon asymmetry

Spin-1/2 fermions						Spin-1 bosons		Spin-0 Higgs boson					
Quarks	Left	u	Right	Left	c	Right	Left		t	Right	Force carriers	g	Z ⁰
	Left	d	Right	Left	s	Right	Left	b	Right	γ			
	Left	ν_1	Right	Left	ν_2	Right	Left	ν_3	Right	W^\pm			
Leptons	Left	e	Right	Left	μ	Right	Left	τ	Right				

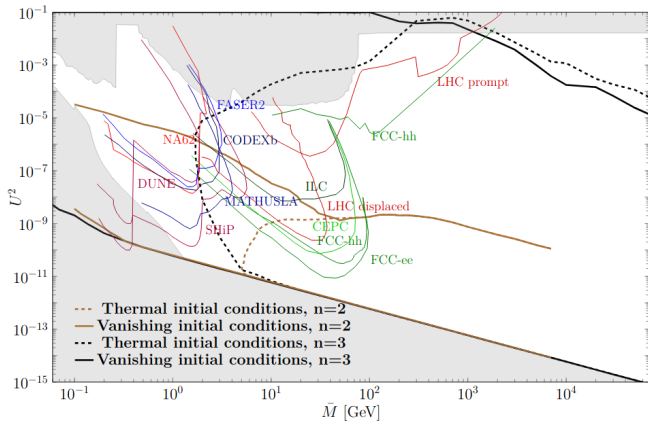
- Minimal scenario: 2 RHN generations (neutrino masses)
- In this work: 3 RHN generations (LRSM, flavour symmetries,...)
- Focus on low scale models: RHN mass at MeV-TeV scale [see talk by J. Klarić]

Leptogenesis parameter space ($m_{\text{lightest}} = 0 \text{ eV}$)



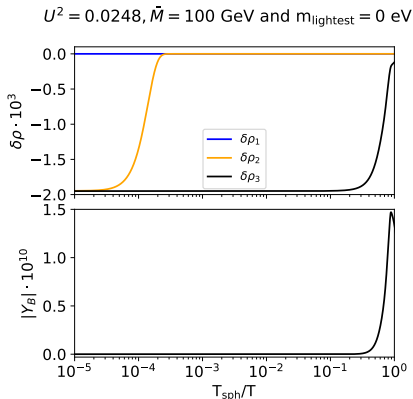
- ▶ Experiments will cut deep into $n = 3$ parameter space.
- ▶ Most optimistic case: Produce thousands of displaced vertices at HL-LHC. Testability !
- ▶ Thermal initial conditions within reach of e.g. NA62.

Leptogenesis parameter space ($m_{\text{lightest}} = 0 \text{ eV}$)



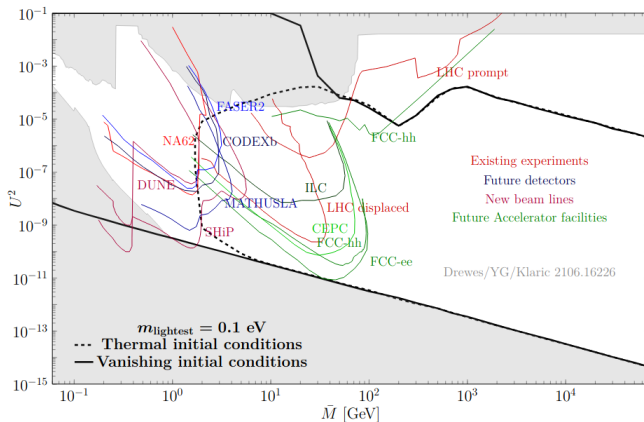
- ▶ Experiments will cut deep into $n = 3$ parameter space.
- ▶ Most optimistic case: Observe thousands of displaced vertices at HL-LHC. Testability !
- ▶ Thermal initial conditions within reach of e.g. NA62.

How can we understand the parameter space ?



- ▶ Large mixing angles allow late equilibration of one RHN
 $U_i^2 \ll 1$.
 \hookrightarrow Late BAU production, less time for washout.
- ▶ For thermal initial conditions, resonant and flavour effects partly compensate for $\frac{M^2}{T^2}$ suppression.
- ▶ At large \bar{M} , parameter space for thermal I.C. is larger because asymmetry produced during freeze-in and freeze-out have opposite signs.

Leptogenesis parameter space ($m_{\text{lightest}} = 0.1 \text{ eV}$)



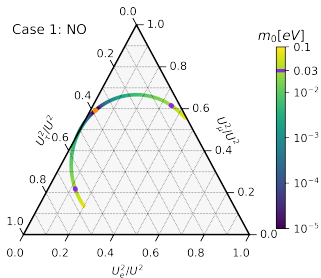
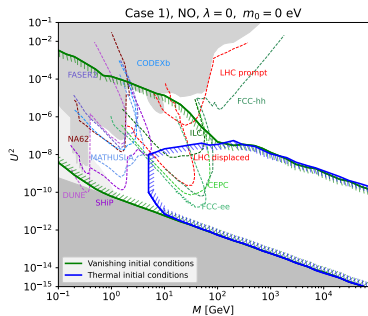
- ▶ Parameter space slightly smaller for $m_{\text{lightest}} = 0.1 \text{ eV}$.
 \hookrightarrow Weaker Yukawa hierarchy: $F = \frac{i}{v} U_\nu \sqrt{m_\nu^{\text{diag}}} R \sqrt{M_M}$.
- ▶ Dip at $\bar{M} \sim 100 \text{ GeV}$ due to transition between freeze-in and freeze-out regime.

- ▶ Parameter space much larger for models with 3 RHN due to late equilibration of the additional RHN.
- ▶ Testability prospects largely improved: can hope to test RHN as solution to neutrinos masses and baryon asymmetry problem !
- ▶ Leptogenesis with thermal initial conditions work as low as $\mathcal{O}(1)$ GeV, within reach of current experiments !
- ▶ Smaller parameter space for non-zero m_{lightest} but remains large.

Thanks for your attention !

Backup slides

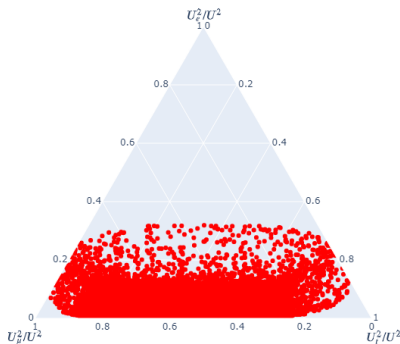
Parameter space in presence of flavour symmetries



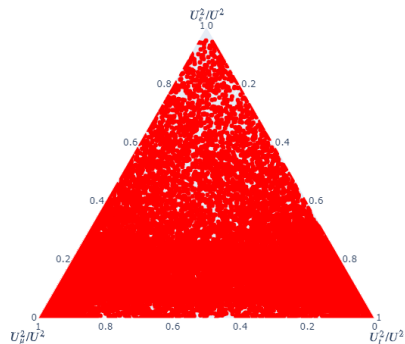
[Drewes/YG/Hagedorn/Klarič, 23xx.xxxx]

Enhanced predictivity in presence of flavour symmetries (e.g. $\Delta(6n^2)$)

Flavour branching ratios



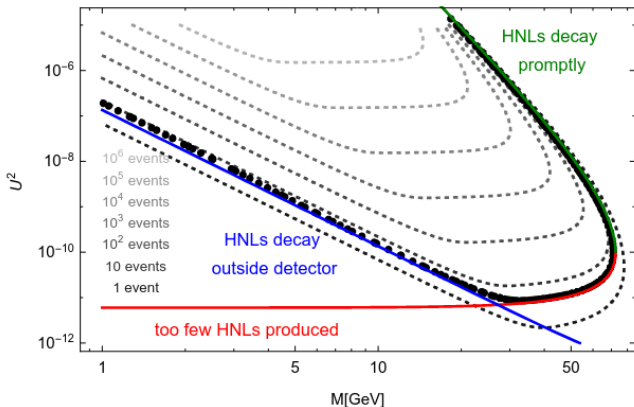
$$m_{\text{lightest}} = 0 \text{ eV}$$



$$m_{\text{lightest}} = 0.1 \text{ eV}$$

[Drewes/YG/Klarić 23xx.xxxx]

Expected number of events at FCC-ee



Expected number of events for
 $U_e^2 : U_\mu^2 : U_\tau^2 = 0 : 1 : 0$, $N_Z = 2.5 \cdot 10^{12}$, $l_0 = 400 \mu\text{m}$, $d_{\text{cyl}} = 10\text{m}$, $l_{\text{cyl}} = 8.6\text{m}$.

[Drewes, 2210.17110]

$$i\frac{d\rho}{dt} = [\mathbf{H}, \delta\rho] - \frac{i}{2}\{\Gamma, \delta\rho\} - i \sum_{a \in \{e, \mu, \tau\}} \tilde{\Gamma}_a \frac{\mu_a}{T} f_F(1 - f_F),$$

$$i\frac{d\bar{\rho}}{dt} = -[\mathbf{H}, \delta\bar{\rho}] - \frac{i}{2}\{\Gamma, \delta\bar{\rho}\} + i \sum_{a \in \{e, \mu, \tau\}} \tilde{\Gamma}_a \frac{\mu_a}{T} f_F(1 - f_F),$$

$$\frac{d}{dt} n_{\Delta_a} = -\frac{2i\mu_a}{T} \int \frac{d^3\vec{k}}{(2\pi)^3} \text{Tr}[\Gamma_a] f_F(1 - f_F) + i \int \frac{d^3\vec{k}}{(2\pi)^3} \text{Tr}[\tilde{\Gamma}_a(\delta\bar{\rho} - \delta\rho)].$$

Density matrix/Matter-antimatter asymmetry
Effective Hamiltonian/Interaction rates

- ▶ Momentum-averaged rates from Klaric/Shaposhnikov/Timiryasov [2103.165451]
- ▶ Cover a mass range from 50 MeV to 70 TeV.

Naive seesaw bound

$$U_i^2 \sim \frac{\sqrt{\Delta m_{atm}^2 + m_{light}^2}}{M} \lesssim 10^{-10} \frac{\text{GeV}}{M_i}$$

B-L approximate symmetry

Majorana mass

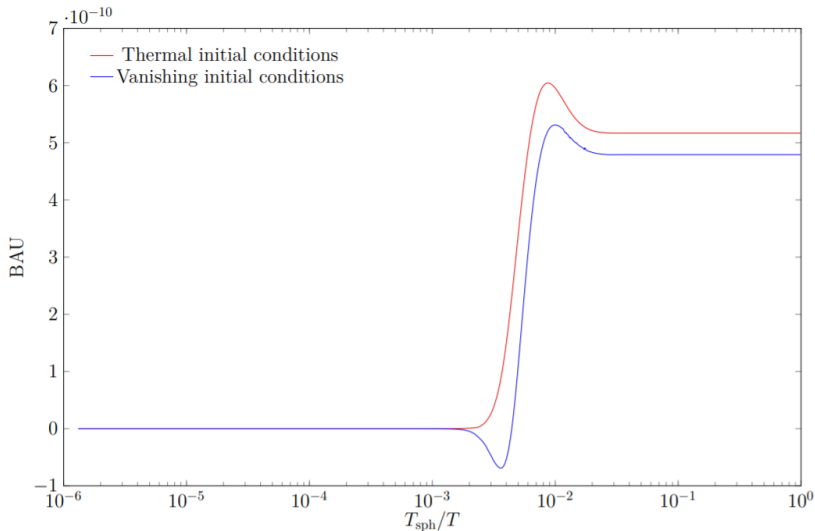
$$\bar{M} \cdot \begin{pmatrix} 1 - \mu & 0 & 0 \\ 0 & 1 + \mu & 0 \\ 0 & 0 & \mu' \end{pmatrix}$$

Yukawa coupling

$$\begin{pmatrix} f_e(1 + \epsilon_e) & if_e(1 - \epsilon_e) & f_e \epsilon'_e \\ f_\mu(1 + \epsilon_\mu) & if_\mu(1 - \epsilon_\mu) & f_\mu \epsilon'_\mu \\ f_\tau(1 + \epsilon_\tau) & if_\tau(1 - \epsilon_\tau) & f_\tau \epsilon'_\tau \end{pmatrix}$$

Smallness of light neutrino masses from the smallness of the symmetry breaking parameters $\mu, \epsilon, \epsilon' \ll 1$.

Thermal vs vanishing initial conditions



- ▶ Asymmetries generated during freeze-in and freeze-out have opposite signs.

Consistency with ν -oscillation data induced by Casas-Ibarra parametrisation

$$F = \frac{i}{v} U_\nu \sqrt{m_\nu^{diag}} \mathbf{R} \sqrt{M_M}$$

n=2

- 2 CP-violating phases**
- 3 PMNS angles** (fixed)
- 2 light neutrino masses** (fixed)
- 1 complex Euler angle**
- 2 Majorana masses**

6 free parameters

n=3

- 3 CP-violating phases**
- 3 PMNS angles** (fixed)
- 3 light neutrino masses** (2 fixed)
- 3 complex Euler angles**
- 3 Majorana masses**

13 free parameters