

New collider search window for ALPs

Jonathan Machado Rodríguez

Universidad Autónoma de Madrid - Instituto de Física Teórica

Based on: J. Bonilla , I. Brivio, J. M.R., J. F. de Trocóniz

J. High Energ. Phys. 2022, 113 (2022) [2202.03450]

Axion-Like Particles

- Axion-Like Particles (or ALPs) are **neutral pseudo scalar** pseudo Goldston Bosons
- **Effective Field Theory** (EFT) consistent with SM gauge and CP symmetries
- Either **shift-invariant** and/or **anomalous couplings** interactions
- ALP interactions with SM particles have a **derivative character**: they grow with momentum

$$\left. \begin{array}{l} \frac{\partial_\mu a}{f_a} \bar{\psi} \gamma^\mu \psi \\ \frac{a}{f_a} X_{\mu\nu} \tilde{X}^{\mu\nu} \end{array} \right\}$$

Axion-Like Particles

$$\mathcal{L}_{ALP} \supset -c_{\tilde{B}} \frac{a}{f_a} B_{\mu\nu} \tilde{B}^{\mu\nu} - c_{\tilde{W}} \frac{a}{f_a} W_{\mu\nu}^i \tilde{W}^{i\mu\nu} - c_{\tilde{G}} \frac{a}{f_a} G_{\mu\nu}^A \tilde{G}^{A\mu\nu}$$

- Classical searches: ALP couplings to gluons and photons

f_a : New physics scale

- ALP couplings to **EWK bosons**: WW , ZZ , and $Z\gamma$

- **Depend on 2 parameters**

- ALP-gauge interactions at ATLAS and CMS:
 - Mono-X
 - Resonant
 - **New idea: nonresonant ALP searches**

$$\left[\begin{array}{l} g_{a\gamma\gamma} = \frac{4}{f_a} (s_\theta^2 c_{\widetilde{W}} + c_\theta^2 c_{\widetilde{B}}) \\ g_{a\gamma Z} = \frac{4}{f_a} s_{2\theta} (c_{\widetilde{W}} - c_{\widetilde{B}}) \\ g_{aZZ} = \frac{4}{f_a} (c_\theta^2 c_{\widetilde{W}} + s_\theta^2 c_{\widetilde{B}}) \\ g_{aWW} = \frac{4}{f_a} c_{\widetilde{W}} \end{array} \right] \quad \theta: \text{Weinberg angle}$$

↑
Imposed by gauge invariance

Axion-Like Particles

$$\mathcal{L}_{ALP} \supset -c_{\tilde{B}} \frac{a}{f_a} B_{\mu\nu} \tilde{B}^{\mu\nu} - c_{\tilde{W}} \frac{a}{f_a} W_{\mu\nu}^i \tilde{W}^{i\mu\nu} - c_{\tilde{G}} \frac{a}{f_a} G_{\mu\nu}^A \tilde{G}^{A\mu\nu}$$

- Classical searches: ALP couplings to gluons and photons

f_a : New physics scale

- ALP couplings to **EWK bosons**: WW , ZZ , and $Z\gamma$

- Depend on 2 parameters**



- ALP-gauge interactions at ATLAS and CMS:
 - Mono-X
 - Resonant
 - New idea: nonresonant ALP searches**

$$\left[\begin{array}{l} g_{a\gamma\gamma} = \frac{4}{f_a} (s_\theta^2 c_{\widetilde{W}} + c_\theta^2 c_{\widetilde{B}}) \\ g_{a\gamma Z} = \frac{4}{f_a} s_{2\theta} (c_{\widetilde{W}} - c_{\widetilde{B}}) \\ g_{aZZ} = \frac{4}{f_a} (c_\theta^2 c_{\widetilde{W}} + s_\theta^2 c_{\widetilde{B}}) \\ g_{aWW} = \frac{4}{f_a} c_{\widetilde{W}} \end{array} \right]$$

θ : Weinberg angle

↑
Imposed by gauge invariance

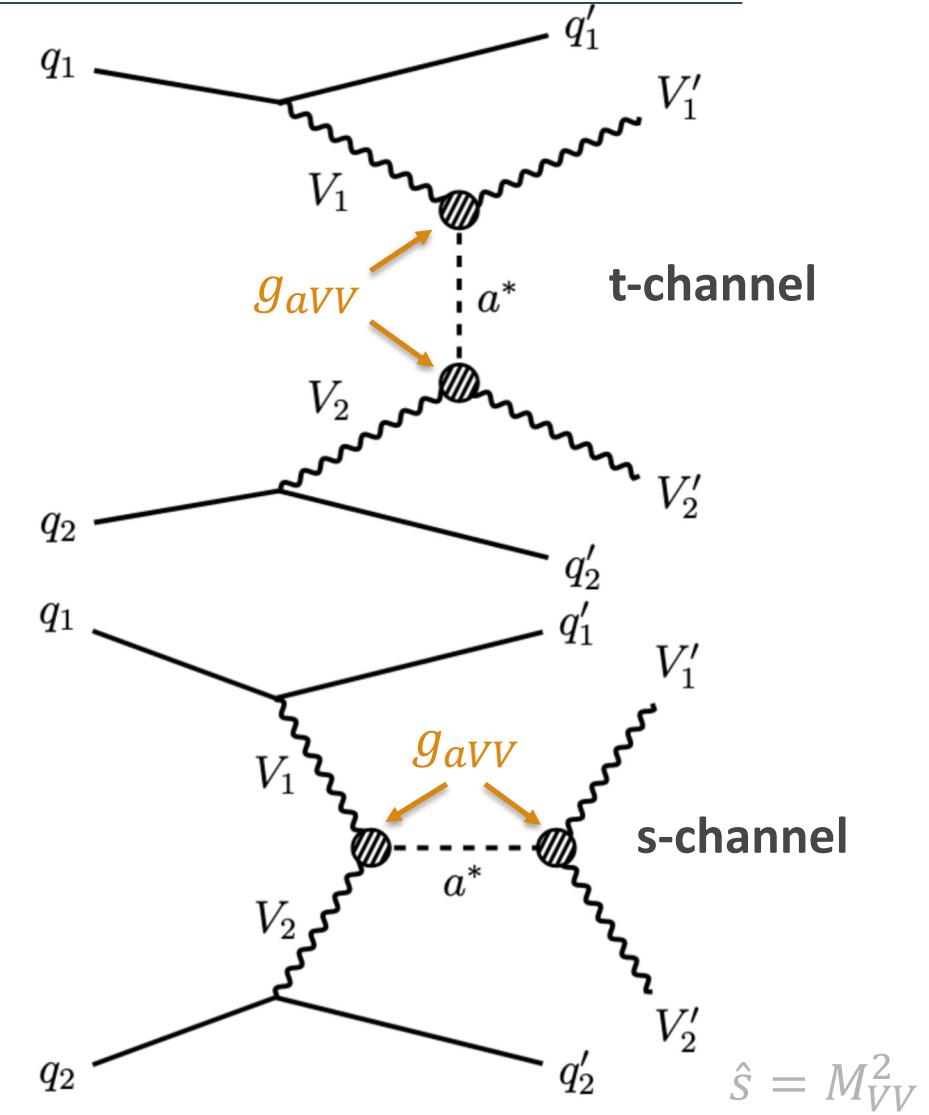
A Novel Approach: Nonresonant ALPs in VBS

- Vector Boson Scattering (VBS)
- Nonresonant ALP searches proposed by M.B. Gavela, J.M. No, V. Sanz and J.F. de Trocóniz [1905.12953]
 - ALP as a very off-shell mediator $m_a^2 \ll \hat{s}$
 - Signals **independent of ALP mass and width up to $m_a \lesssim 100$ GeV**
 - **Harder scaling with $\hat{s} = M_{VV}^2$**

Why VBS?

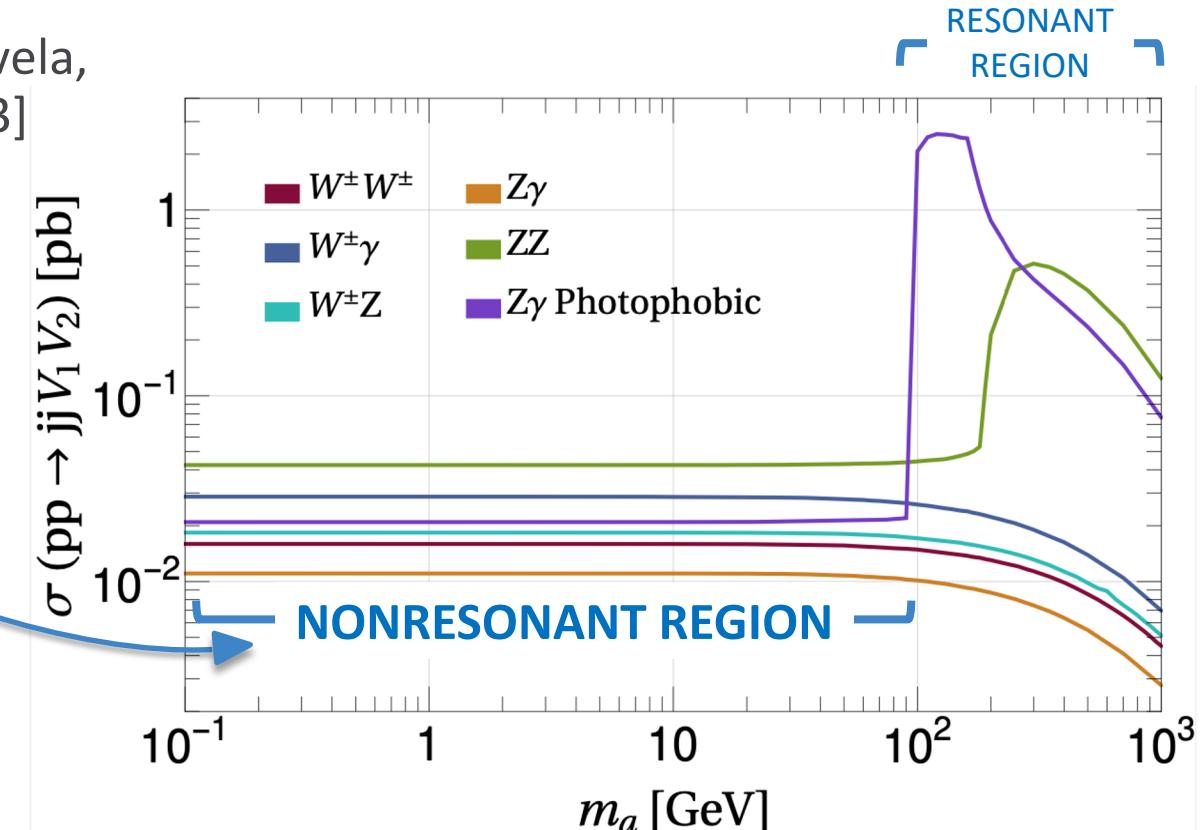
- **Limits on ALP couplings to vector boson independently of the gluon coupling**

→ $c_{\tilde{B}}$ $c_{\tilde{W}}$ ~~$c_{\tilde{G}}$~~ → **2 parameters**



A Novel Approach: Nonresonant ALPs in VBS

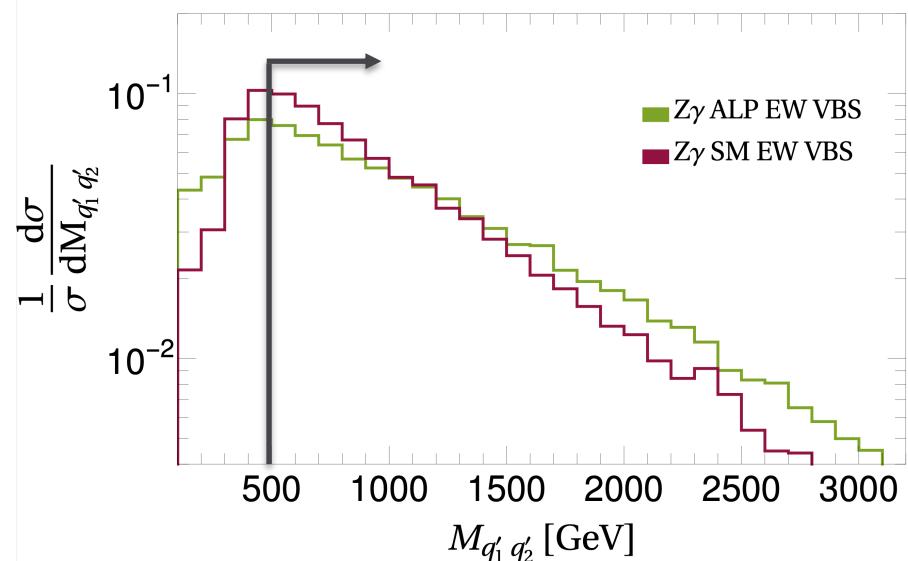
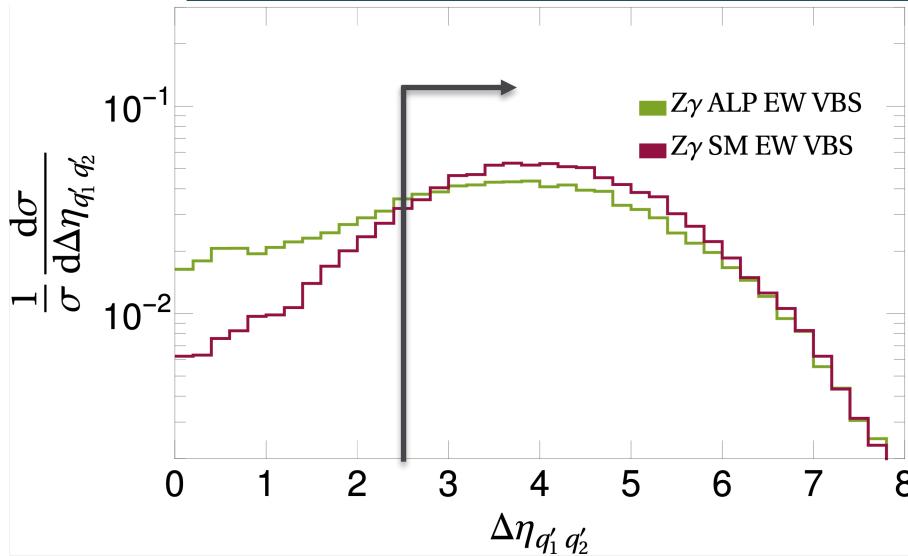
- Vector Boson Scattering (VBS)
- Nonresonant ALP searches proposed by M.B. Gavela, J.M. No, V. Sanz and J.F. de Trocóniz [1905.12953]
 - ALP as a very off-shell mediator $m_a^2 \ll \hat{s}$
 - Signals **independent of ALP mass and width up to $m_a \lesssim 100$ GeV**
 - **Harder scaling with $\hat{s} = M_{VV}^2$**
- **Why VBS?**
- **Limits on ALP couplings to vector boson independently of the gluon coupling**



→ $c_{\tilde{B}}$ $c_{\tilde{W}}$ ~~$c_{\tilde{G}}$~~ → 2 parameters

$$\hat{s} = M_{VV}^2$$

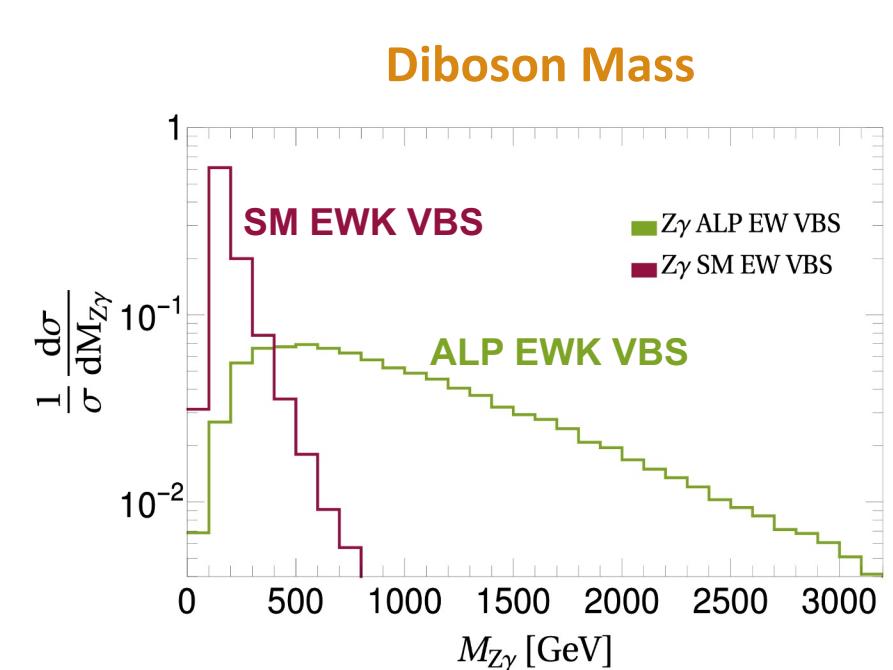
A Novel Approach: Nonresonant ALPs in VBS



Wide
rapidity separation

VBS
CHARACTERISTIC
OBSERVABLES

Large
dijet mass



Diboson Mass
Derivative nature of ALP interactions:
deviations in the tail of the **diboson mass** with respect to the SM

ALP Diboson Mass in CMS Leptonic Analyses

- **ATLAS/CMS Run 2 measurements:** first comparison to data, calibration of simulation tools and calculation of educated predictions for higher luminosities.
- **Reinterpretation of five CMS VBS analyses** with lepton/photon final states:
 - ZZ: CMS-SMP-20-01
 - $W^\pm W^\pm$ and WZ: CMS-SMP-19-012
 - $Z\gamma$: CMS-SMP-20-016
 - $W\gamma$: CMS-SMP-19-008
- Look at **high energy deviations** in the tail of the transverse momentum/mass spectra
- Selections cuts, data and backgrounds in the CMS papers
- **Generation of ALP VBS:** MadGraph_aMC@NLO + Pythia8 + Delphes3
- **Calibrate** our **Delphes detector simulation** using the SM EWK channel

ALP Diboson Mass in CMS Leptonic Analyses

The diagram illustrates a process flow from signal and interference cross sections to expected lepton events. It starts with two boxes: one orange box labeled $c_{\tilde{W}} = c_{\tilde{B}}$ signal / interf. [fb] and one blue box labeled Photophobic signal / interf. [fb]. Arrows point from these boxes to a third box labeled Expected Lepton Events, which then points to a final column labeled Int. lum. [fb⁻¹]. Below the first two boxes are the equations $g_{a\gamma Z} = 0$ and $g_{a\gamma\gamma} = 0$.

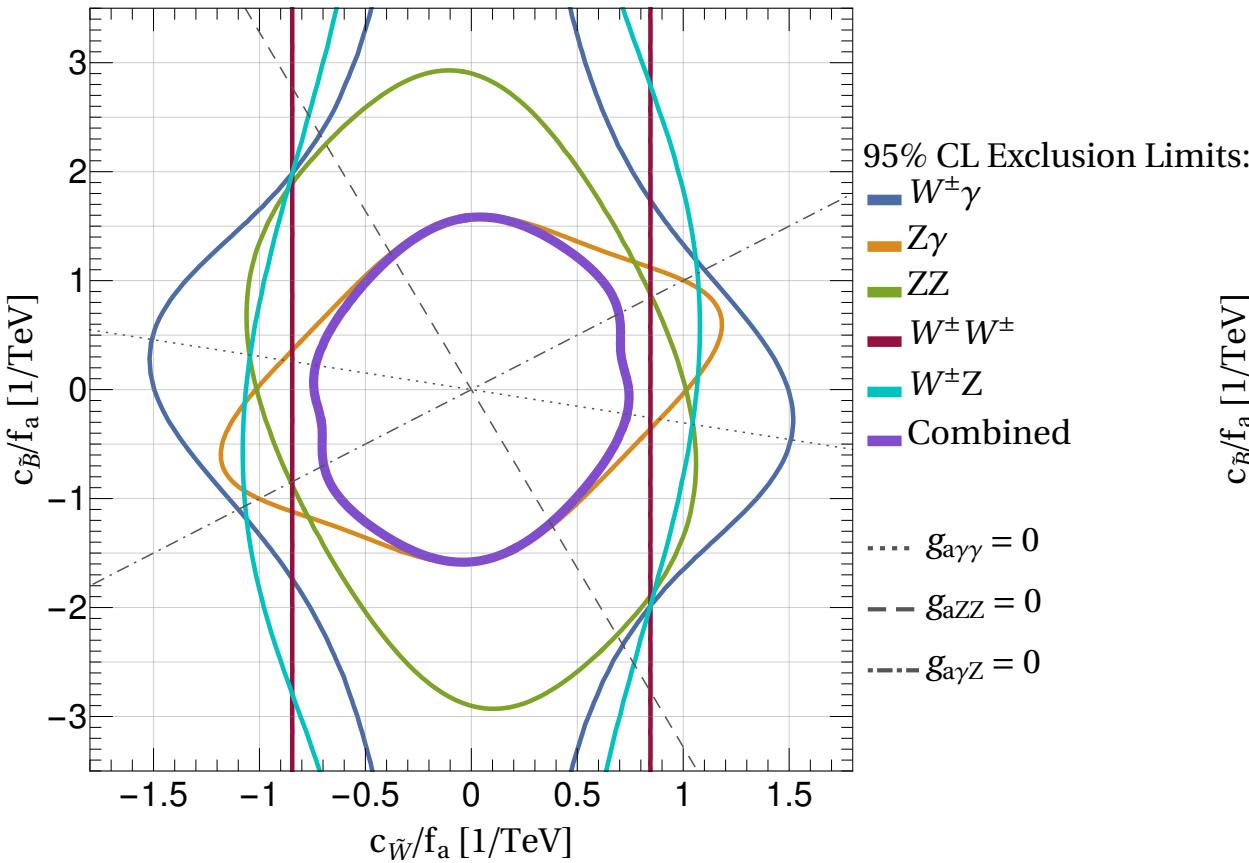
	$c_{\tilde{W}} = c_{\tilde{B}}$ signal / interf. [fb]	Photophobic signal / interf. [fb]	Expected Lepton Events	Int. lum. [fb ⁻¹]
ZZ	42.4 / -13.5	18.5 / -9.3	9.3 / -3.2	137
WZ	18.4 / 1.7	23.9 / -0.14	4.2 / 0.05	137
$W^\pm W^\pm$	16.0 / -4.0	16.0 / -4.0	18 / -5.5	137
$W\gamma$	28.7 / 4.3	5.4 / 1.7	3.6 / -0.04	35.9
$Z\gamma$	11.1 / 0.3	20.9 / -9.1	15.1 / 0.07	137

Results

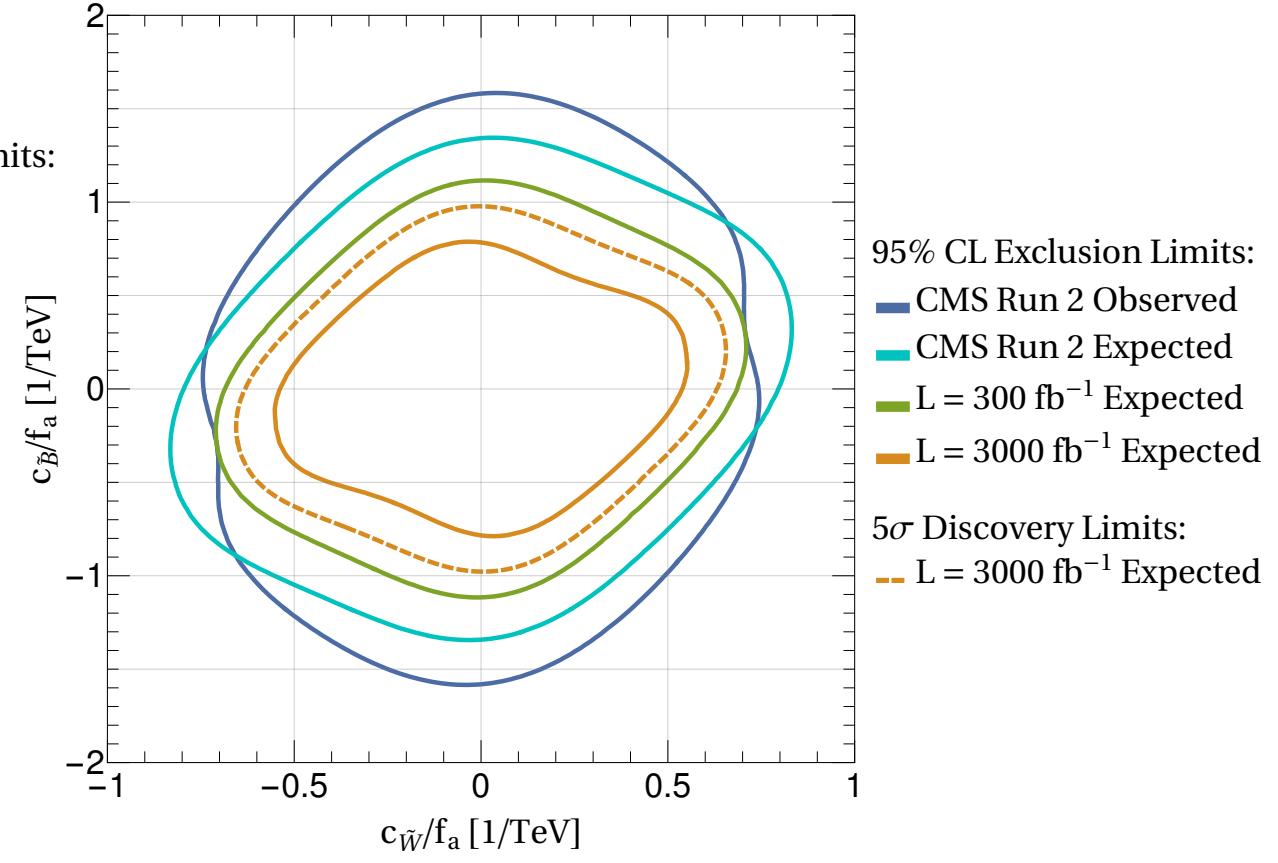
- Maximum likelihood fit of signal and background to the **diboson invariant/transverse masses**
- No excess found with respect to SM expectations
- **Current limits** with CMS Run 2 data and **projected limits** at Run 3 and HL-LHC in the ALP ($c_{\tilde{W}}, c_{\tilde{B}}$) parameter space
- **Expected diff. cross-sections are parameterized** in the $(c_{\tilde{W}}, c_{\tilde{B}})$ plane with quartic / quadratic **polynomials** for pure signal / interference ALP components.

Results

CURRENT LIMITS



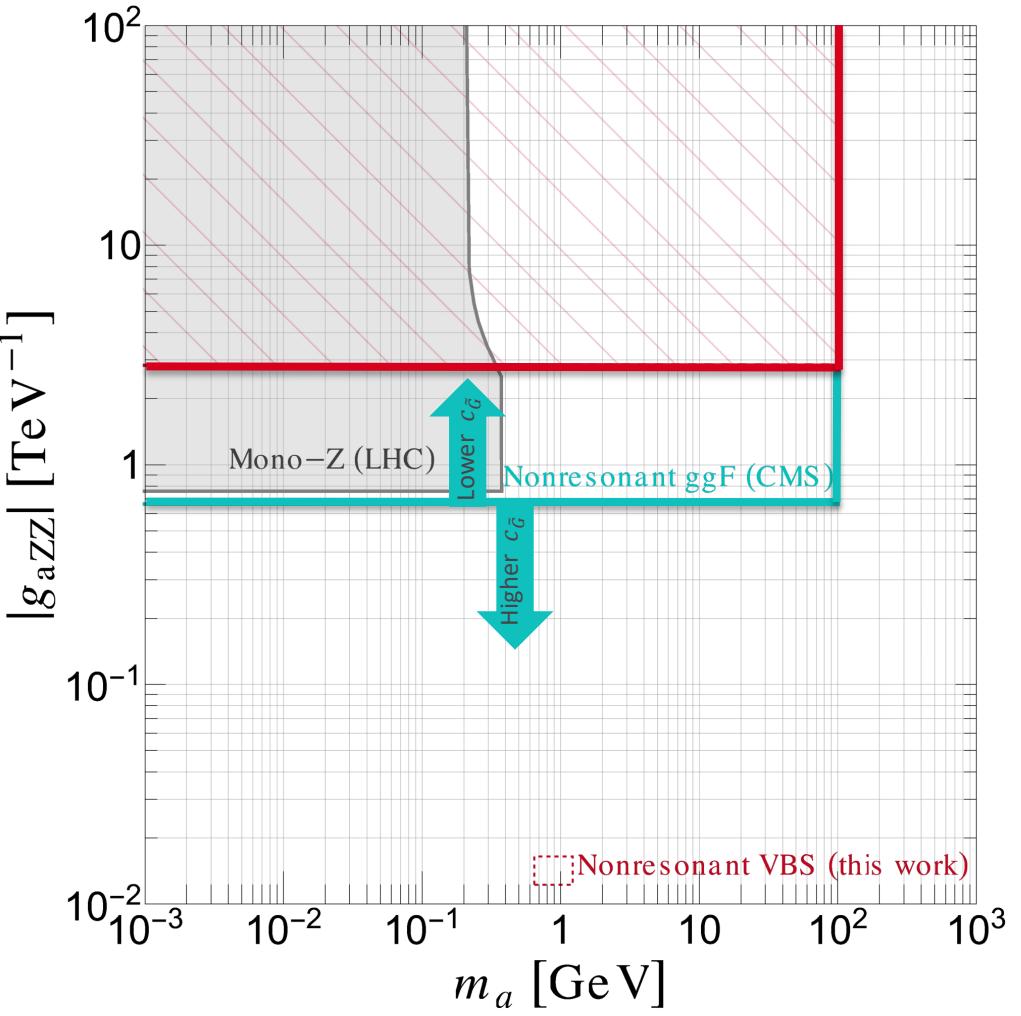
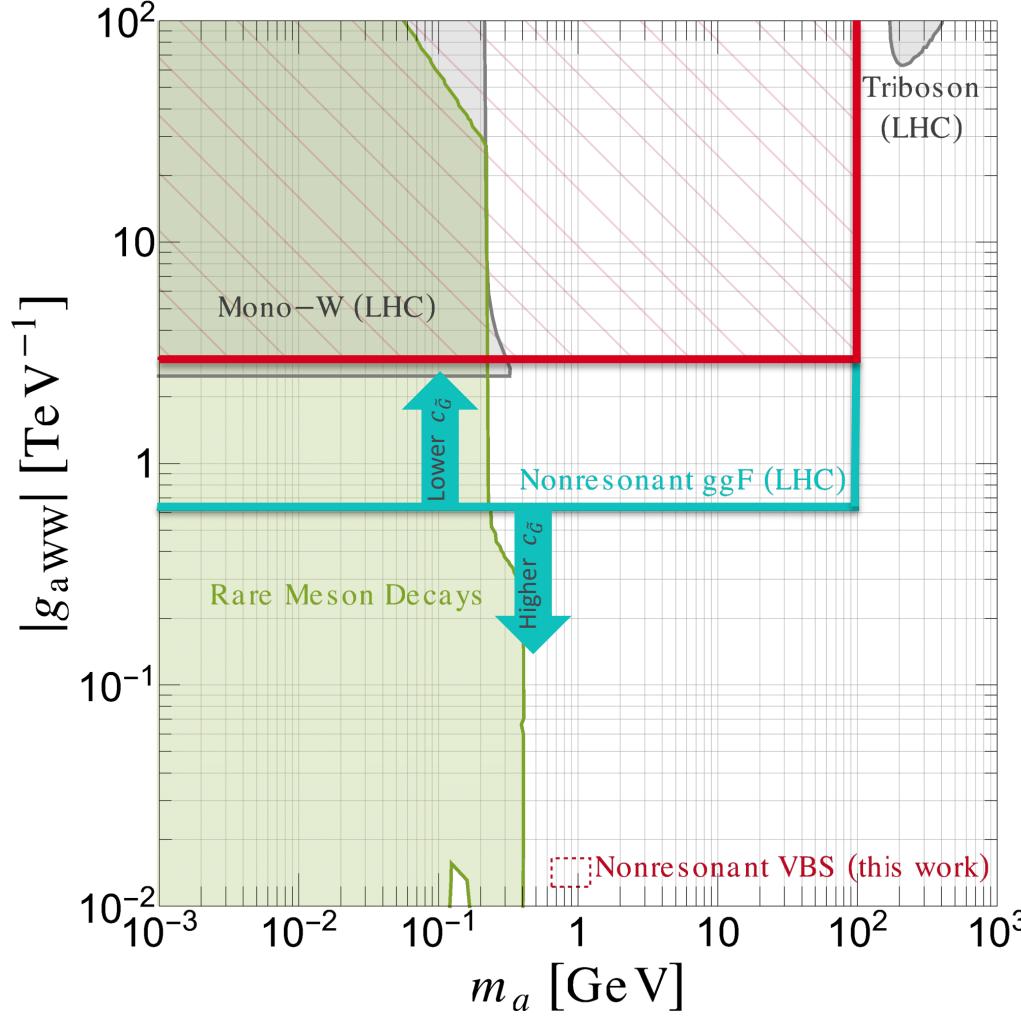
PROJECTED LIMITS



(for $f_a > 4 \text{ TeV}$ and $M_a < 100 \text{ GeV}$)

Results: comparison to existing bounds

- Limits are **very competitive** and probe **previously unexplored regions** of the param. space



Conclusions

- Access to **EWK couplings independently of the gluons**
- Current limits (CMS Run 2 data) and projected limits (Run 3 and HL-LHC)
- Limits **independent** of the **ALP mass and decay width** ($m_a \lesssim 100$ GeV)
- **Limits** are **very competitive** and probe **previously unexplored regions** of the param. space
- Great opportunity for **dedicated ALP searches** at Run 3 and HL-LHC

J. High Energ. Phys. 2022, 113 (2022)
ArXiv: 2202.03450

Conclusions

- Access to **EWK couplings independently of the gluons**
- Current limits (CMS Run 2 data) and projected limits (Run 3 and HL-LHC)
- Limits **independent** of the **ALP mass and decay width** ($m_a \lesssim 100$ GeV)
- **Limits** are **very competitive** and probe **previously unexplored regions** of the param. space
- Great opportunity for **dedicated ALP searches** at Run 3 and HL-LHC

J. High Energ. Phys. 2022, 113 (2022)
ArXiv: 2202.03450

Thank you!

Lagrangian and physical couplings

$$\mathcal{L}_{ALP} = \frac{1}{4}\partial_\mu a\partial^\mu a - \frac{m_a^2}{2}a^2 - c_{\tilde{B}}\frac{a}{f_a}B_{\mu\nu}\tilde{B}^{\mu\nu} - c_{\widetilde{W}}\frac{a}{f_a}W_{\mu\nu}^i\tilde{W}^{i\mu\nu} - c_{\widetilde{G}}\frac{a}{f_a}G_{\mu\nu}^A\tilde{G}^{A\mu\nu}$$

EWSB
↓

$$\mathcal{L}_{ALP,EW} = -\frac{g_{a\gamma\gamma}}{4}aF^{\mu\nu}\tilde{F}_{\mu\nu} - \frac{g_{a\gamma Z}}{4}aZ^{\mu\nu}\tilde{F}_{\mu\nu} - \frac{g_{aZZ}}{4}aZ^{\mu\nu}\tilde{Z}_{\mu\nu} - \frac{g_{aWW}}{2}aW^{+\mu\nu}\tilde{W}_{\mu\nu}^-$$

(Unitary gauge)

$$g_{a\gamma\gamma} = \frac{4}{f_a}(s_\theta^2 c_{\widetilde{W}} + c_\theta^2 c_{\tilde{B}})$$

$$g_{aZZ} = \frac{4}{f_a}(c_\theta^2 c_{\widetilde{W}} + s_\theta^2 c_{\tilde{B}})$$

$$g_{a\gamma Z} = \frac{4}{f_a}s_{2\theta}(c_{\widetilde{W}} - c_{\tilde{B}})$$

$$g_{aWW} = \frac{4}{f_a}c_{\widetilde{W}}$$

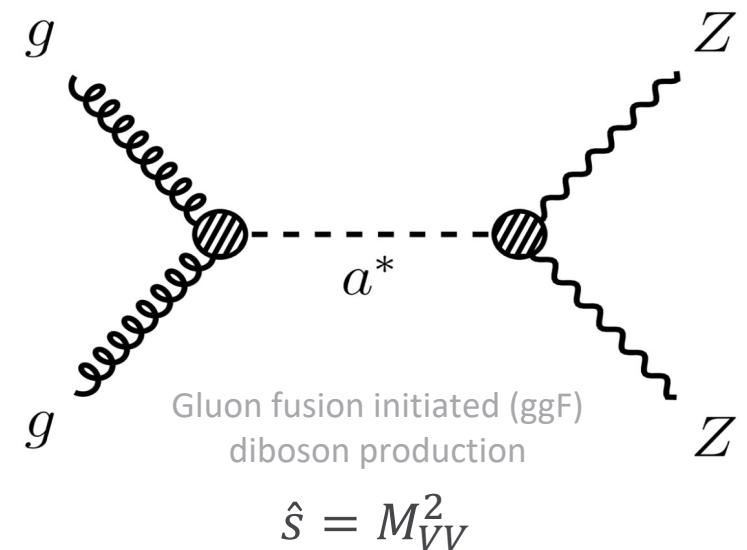
$$g_{agg} = \frac{4}{f_a}c_{\widetilde{G}}$$

Nonresonant ALP-mediated diboson production

- M.B. Gavela, J.M. No, V. Sanz and J.F. de Trocóniz [1905.12953]
- ALP acts as a **very off-shell mediator** $\longrightarrow m_a^2 \ll \hat{s}$
- Signals **independent of ALP mass m_a and its decay width Γ_a** up to $m_a \lesssim 100$ GeV: allows to explore large areas in the parameter space
- Suppression from \hat{s} **compensated by derivative character** of ALP interactions

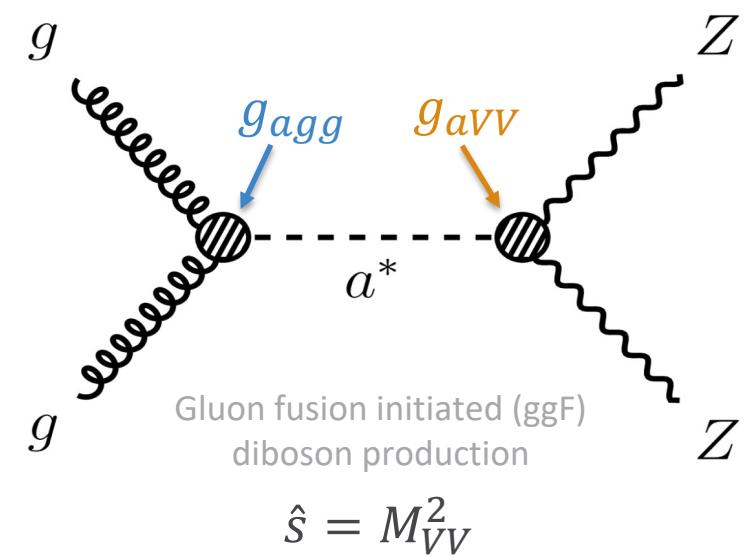
$$\hat{\sigma} \propto \hat{s}/f_a^4$$

Harder scaling



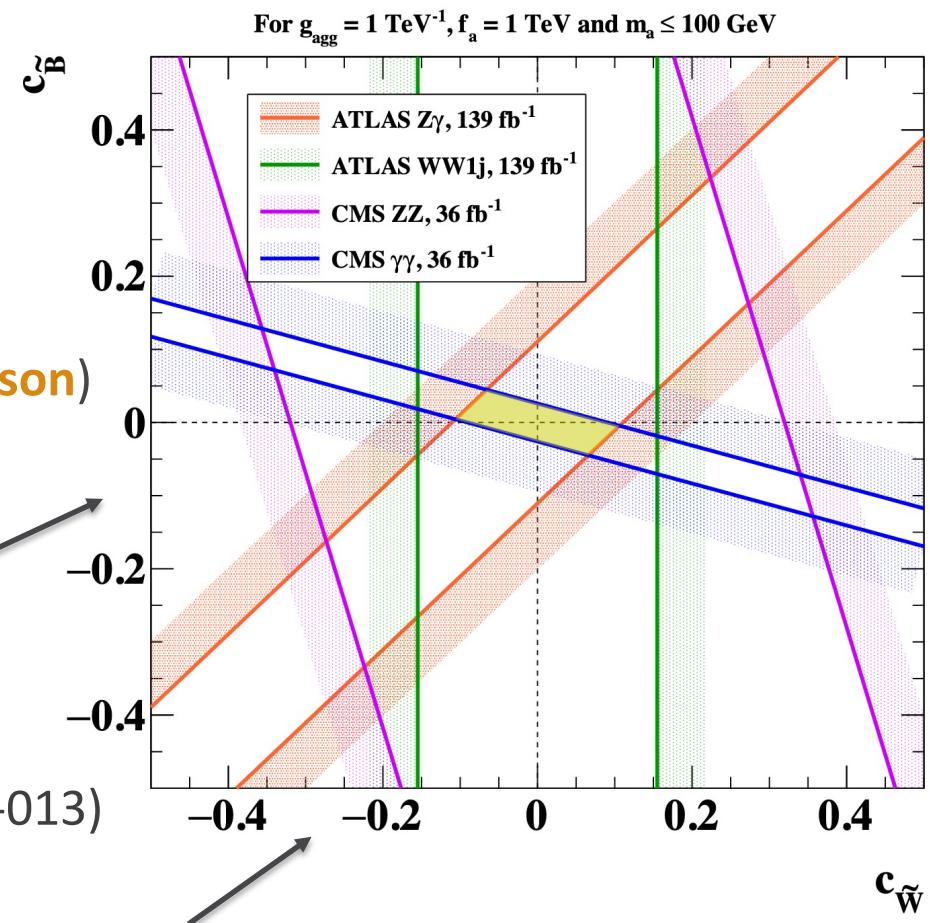
Nonresonant ALP-mediated diboson production

- Reinterpretation of CMS analyses:
 - $gg \rightarrow ZZ$ (CMS-B2G-17-013)
 - $gg \rightarrow \gamma\gamma$ (CMS-EXO-17-017)
- Sensitive to (**ALP coupling to gluons** \times **ALP coupling to EWK diboson**)
 $g_{agg} \times g_{aVV}$
- Cross-sections large enough to constrain significantly the theoretical models using Run 2 data.
- **Dedicated ALP search at CMS:** $gg \rightarrow a^* \rightarrow ZZ/ZH$ (CMS-B2G-20-013)



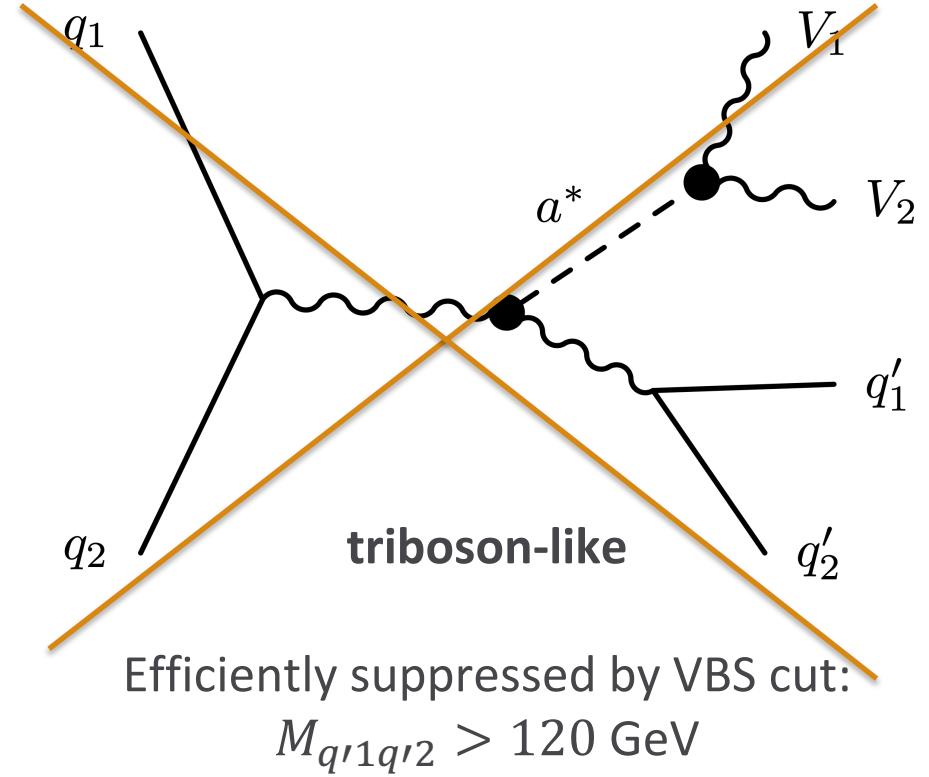
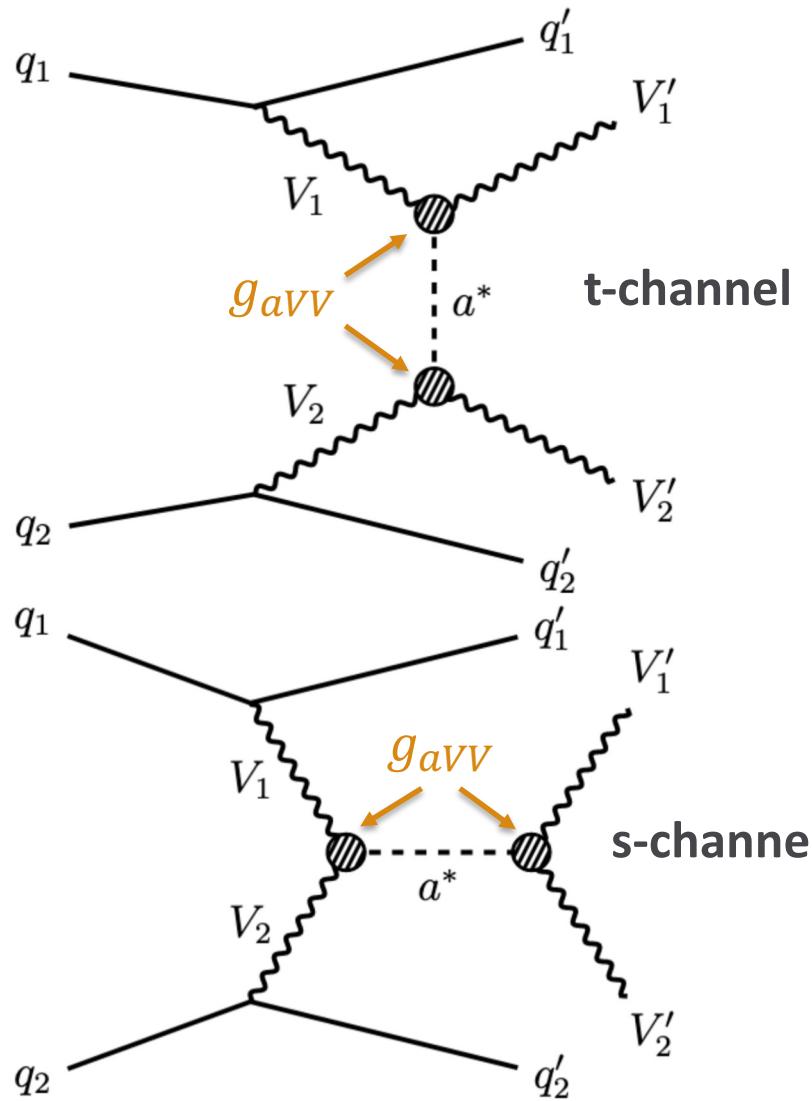
Nonresonant ALP-mediated diboson production

- Reinterpretation of CMS analyses:
 - $gg \rightarrow ZZ$ (CMS-B2G-17-013) 
 - $gg \rightarrow \gamma\gamma$ (CMS-EXO-17-017) 
 - Sensitive to **(ALP coupling to gluons \times ALP coupling to EWK diboson)**
$$g_{agg} \times g_{aVV}$$
 - Cross-sections large enough to constrain significantly the theoretical models using Run 2 data.
 - **Dedicated ALP search at CMS:** $gg \rightarrow a^* \rightarrow ZZ/ZH$ (CMS-B2G-20-013)



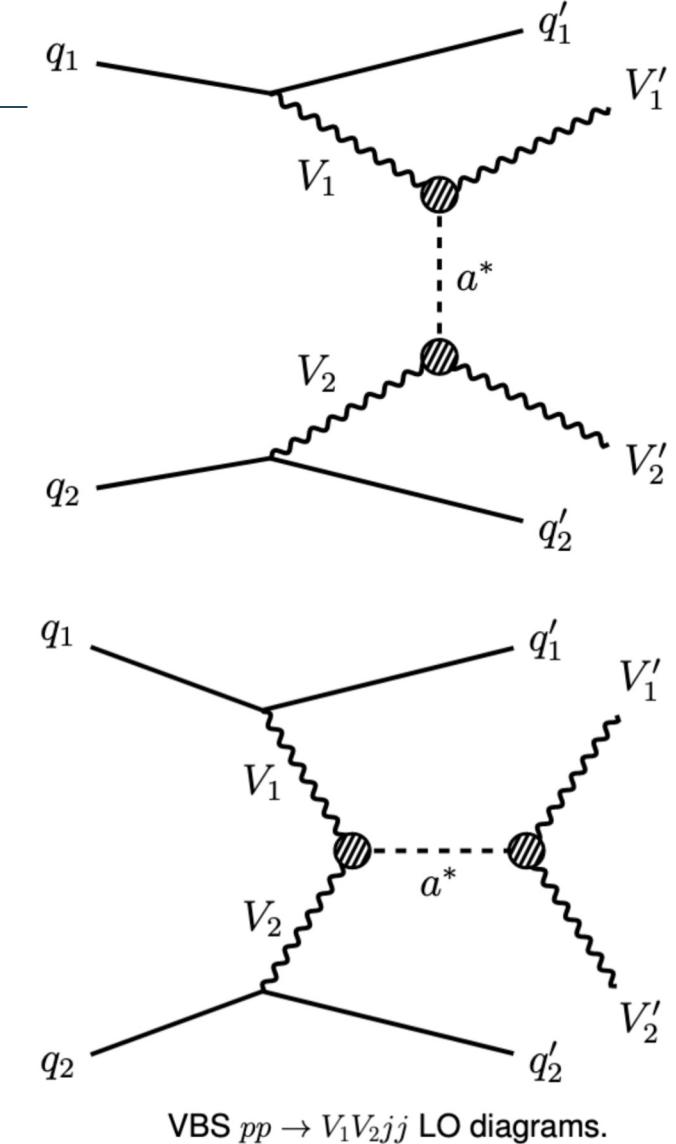
S. Carrá, V. Goumarre, R. Gupta, S. Heim, B. Heinemann, J. Küchler, F. Meloni, P. Quilez and Y.C. Yap
 reinterpretation of ATLAS analyses, [2106.10085]: $Z\gamma$ (CERN- EP-2019-228), WWj (CERN-EP-2021-030)

Nonresonant ALPs in VBS



Contribution from gluons

- Same-sign WW, WZ, W γ : ALP QCD absent at tree level
- ZZ and Z γ : ALP QCD strongly reduced
 - Consistency with previous nonresonant limits [1905.12953], [2106.10085], [2111.13669]
 - VBS selection cuts
 - Large diboson masses
- For the tested region of the ALP parameter space, the theoretical prediction is dominated by ALP VBS
- Conservative: QCD ALP is positive with a subdominant contribution from its interference with EWK ALP



ALP Diboson Mass in CMS Leptonic Analyses

- Calibrate our **Delphes detector simulation** using the SM EWK channel → **EWK scale factor**
- Compare EWK SM VBS expected yields from the CMS simulation and ours →

$$\rho = \frac{N_{SM\ EWK}^{simulated}}{N_{SM\ EWK}^{CMS}}$$

Channel	Obs.	Lum. [fb ⁻¹]	Selection Criteria	ρ
ZZ	M_{ZZ}	137	$M_{jj} > 100$ GeV	0.8 ± 0.1
$Z\gamma$	$M_{Z\gamma}$	137	$M_{jj} > 500$ GeV, $\Delta\eta_{jj} > 2.5$, $p_T^\gamma > 120$ GeV	1.4 ± 0.2
$W^\pm\gamma$	$M_{W\gamma}$	35.9*	$M_{jj} > 800$ GeV, $\Delta\eta_{jj} > 2.5$, $p_T^\gamma > 100$ GeV	3.1 ± 0.5
$W^\pm Z$	M_{WZ}^T	137	$M_{jj} > 500$ GeV, $\Delta\eta_{jj} > 2.5$	1.5 ± 0.4
$W^\pm W^\pm$	M_{WW}^T	137	$M_{jj} > 500$ GeV, $\Delta\eta_{jj} > 2.5$	1.3 ± 0.2

Table 3. Summary of the CMS VBS analyses: the diboson mass observable, the integrated luminosity, the most important selection criteria and the normalization scale factor ρ .

* Recent update at 138 fb⁻¹: CMS PAS SMP-21-011

ALP Diboson Mass in CMS Leptonic Analyses

- **~20 % signal systematics:** PDFs + renormalization and factorization scales + MadGraph@aMC
- **Background uncertainties** from CMS analyses
- Consistency of the ALP EFT and estimation of the impact of the highest-energy bins
 ➡ **upper cut on diboson mass M_{vv}**
- Two benchmarks:
 - $M_{vv} < 2 \text{ TeV}$: ~85 % efficiency
 - $M_{vv} < 4 \text{ TeV}$: >99 % efficiency

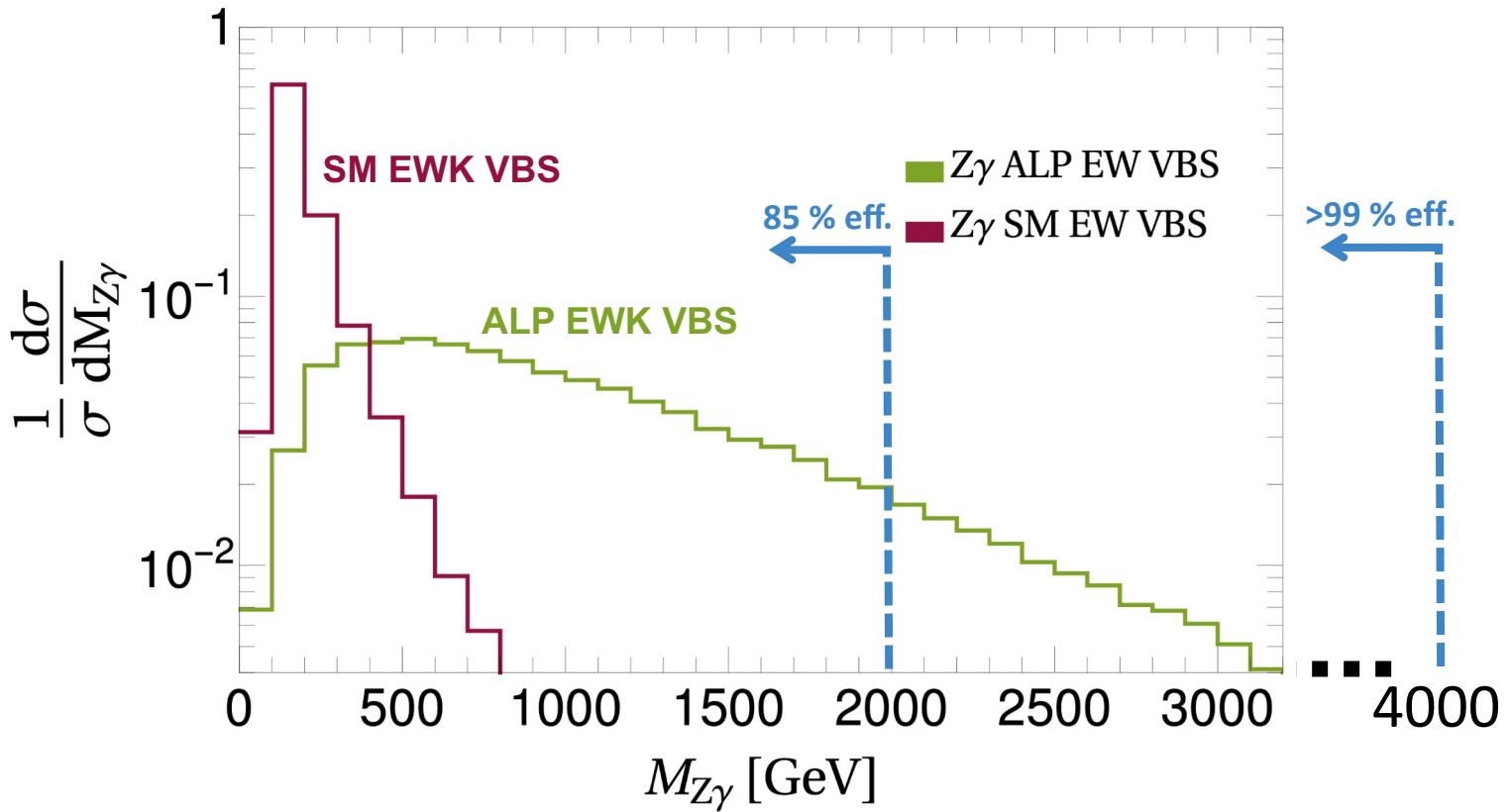
ALP Diboson Mass in CMS Leptonic Analyses

- Consistency of the ALP EFT and estimation of the impact of the highest-energy bins

→ upper cut on diboson mass M_{vv}

- Two benchmarks:

- $M_{vv} < 2 \text{ TeV}$: ~85 % efficiency
- $M_{vv} < 4 \text{ TeV}$: >99 % efficiency



ALP Diboson Mass in CMS Leptonic Analyses

- **Generation cuts:** $p_T(q'_{1,2}) > 20 \text{ GeV}$, $\eta(q'_{1,2}) < 6$, $\Delta R(q'_1 q'_2) > 0.1$, $M_{q'_1 q'_2} > 120 \text{ GeV}$
 $p_T(\gamma) > 10 \text{ GeV}$, $\eta(\gamma) < 2.5$, $\Delta R(\gamma q'_{1,2}) > 0.4$,
- **Selection cuts:**

Channel	Obs.	Lum. [fb ⁻¹]	Selection Criteria	ρ
ZZ	M_{ZZ}	137	$M_{jj} > 100 \text{ GeV}$	0.8 ± 0.1
$Z\gamma$	$M_{Z\gamma}$	137	$M_{jj} > 500 \text{ GeV}$, $\Delta\eta_{jj} > 2.5$, $p_T^\gamma > 120 \text{ GeV}$	1.4 ± 0.2
$W^\pm\gamma$	$M_{W\gamma}$	35.9	$M_{jj} > 800 \text{ GeV}$, $\Delta\eta_{jj} > 2.5$, $p_T^\gamma > 100 \text{ GeV}$	3.1 ± 0.5
$W^\pm Z$	M_{WZ}^T	137	$M_{jj} > 500 \text{ GeV}$, $\Delta\eta_{jj} > 2.5$	1.5 ± 0.4
$W^\pm W^\pm$	M_{WW}^T	137	$M_{jj} > 500 \text{ GeV}$, $\Delta\eta_{jj} > 2.5$	1.3 ± 0.2

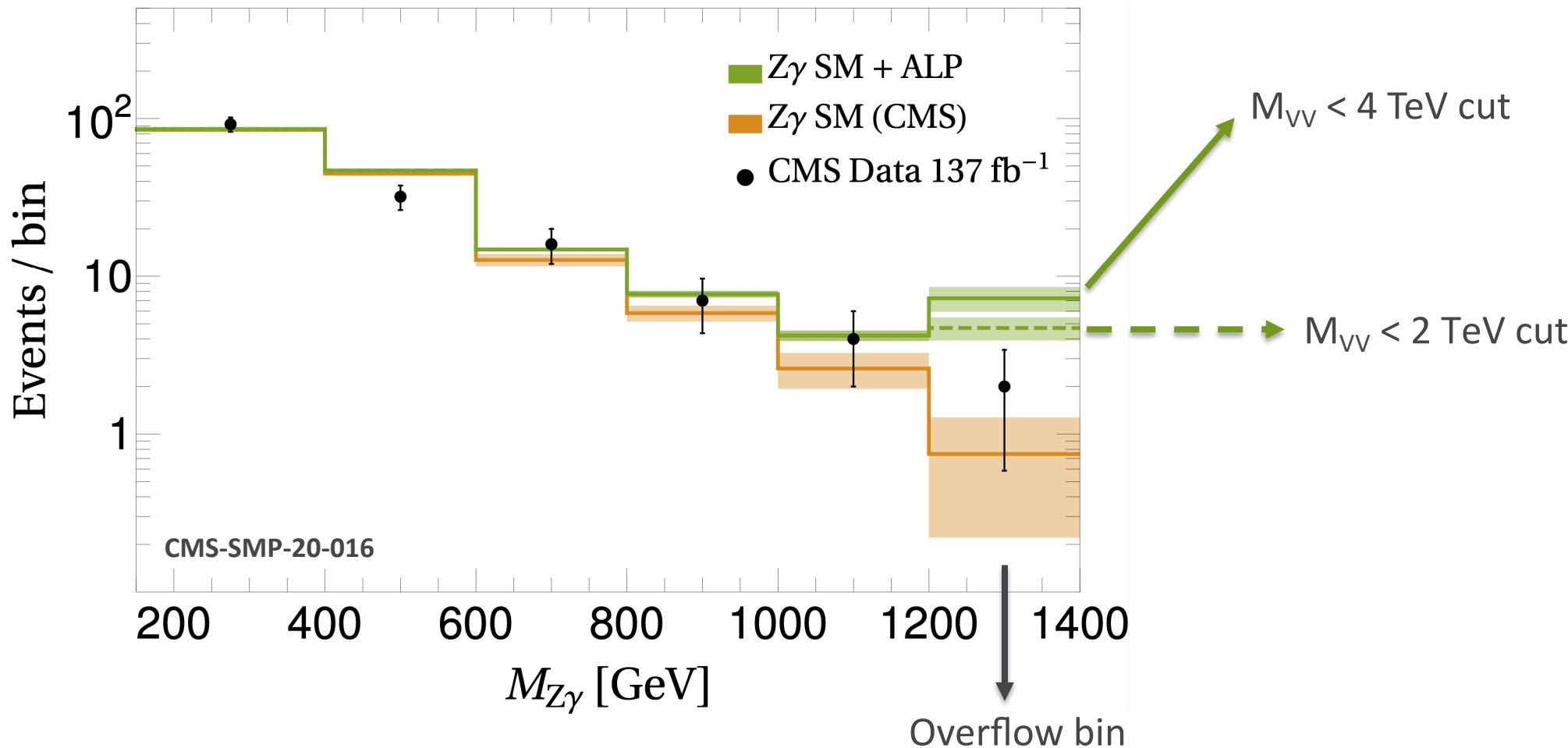
Table 3. Summary of the CMS VBS analyses: the diboson mass observable, the integrated luminosity, the most important selection criteria and the normalization scale factor ρ .

ALP Diboson Mass in CMS Leptonic Analyses

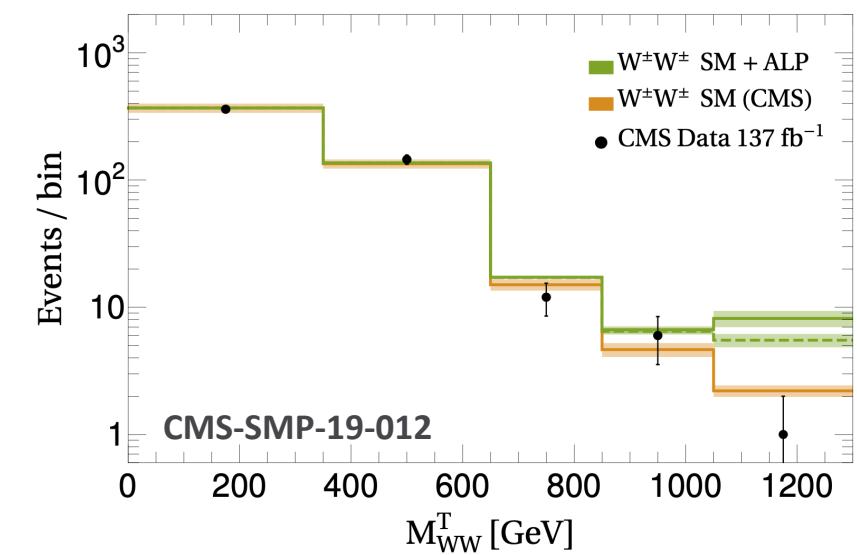
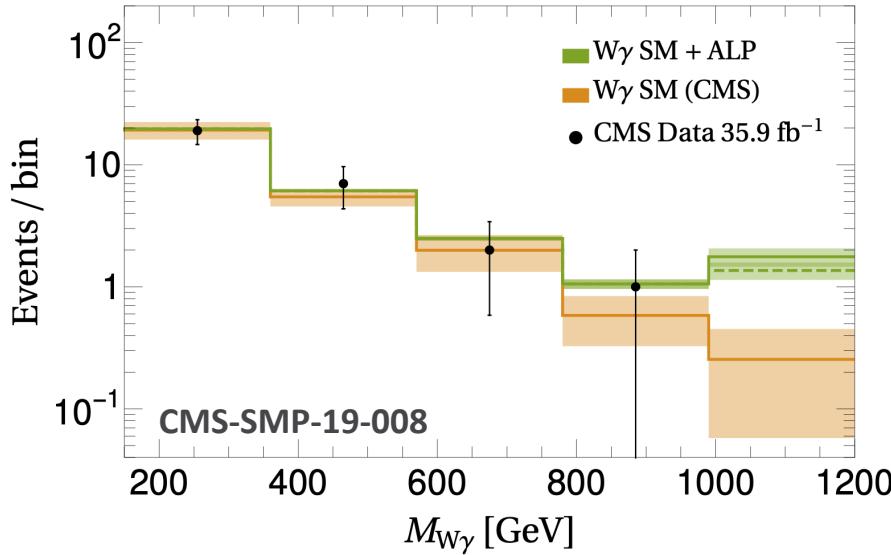
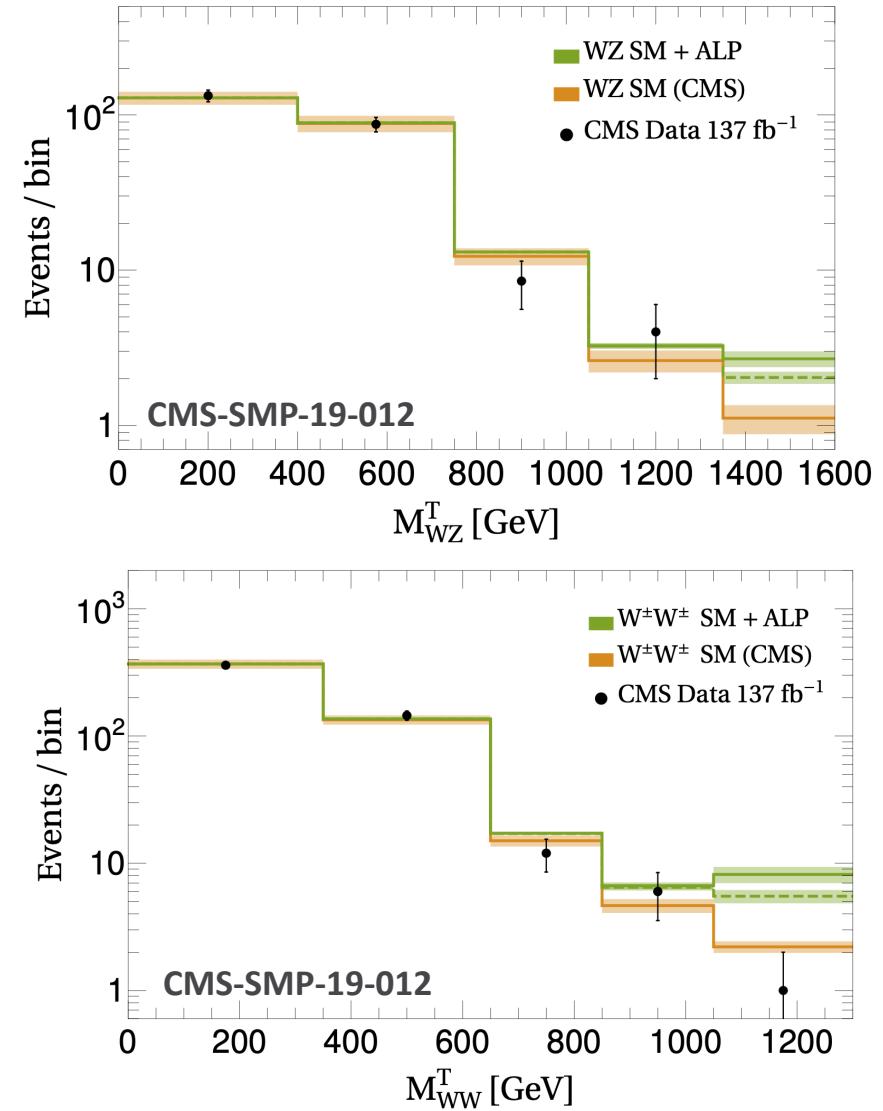
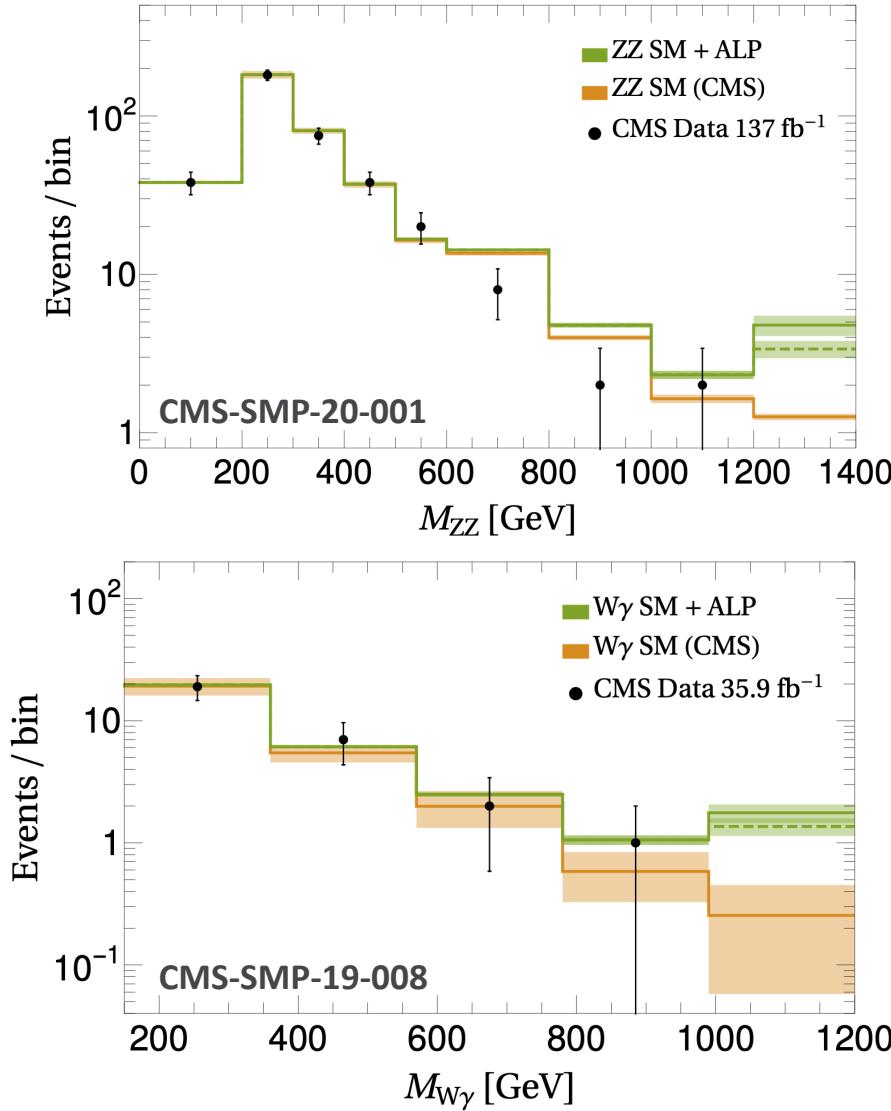
Process	σ_{SM} [fb]	Point	$\sigma_{\text{interf.}}$ [fb]	σ_{signal} [fb]
$pp \rightarrow jjZZ$	98 ± 1	p_0	-13.5 ± 0.1	42.4 ± 0.2
		p_4	-9.3 ± 0.1	18.5 ± 0.1
$pp \rightarrow jjZ\gamma$	393 ± 1	p_0	0.3 ± 0.1	11.1 ± 0.1
		p_4	-9.1 ± 0.1	20.9 ± 0.1
$pp \rightarrow jjW^\pm\gamma$	994 ± 3	p_0	4.3 ± 0.1	28.7 ± 0.1
		p_4	1.7 ± 0.1	5.4 ± 0.1
$pp \rightarrow jjW^\pm Z$	386 ± 1	p_0	1.7 ± 0.1	18.4 ± 0.1
		p_4	0.1 ± 0.1	23.9 ± 0.1
$pp \rightarrow jjW^\pm W^\pm$	256 ± 1	p_0, p_4	-4.0 ± 0.1	16.0 ± 0.1

- **EW VBS SM and ALP signal cross sections**
- Generation cuts
- Two benchmark points:
 - $p_0 \equiv g_{a\gamma Z} = 0$
 - $p_4 \equiv g_{a\gamma\gamma} = 0$

ALP Diboson Mass in CMS Leptonic Analyses



ALP Diboson Mass in CMS Leptonic Analyses



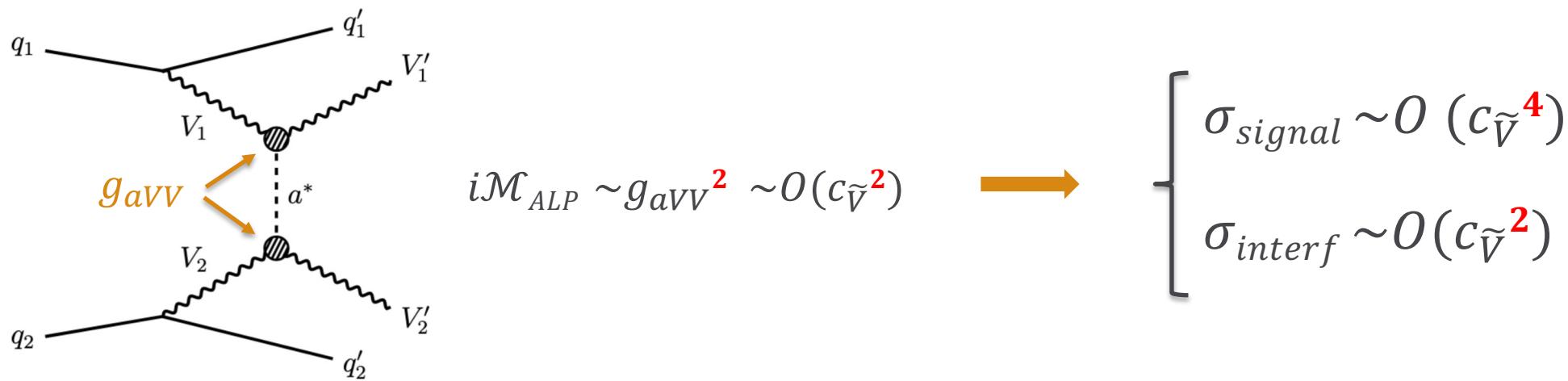
Branching fractions and selection efficiencies

Analysis	ZZ	$Z\gamma$	$W^\pm\gamma$	$W^\pm Z$	$W^\pm W^\pm$
Branching fraction	0.45%	6.7%	22%	1.5%	4.8%
Efficiency	35.7%	14.0%	1.6%	11.3%	17.0%

Table 4. Summary of branching fractions and selection efficiencies for each VBS channel. The efficiencies are relative to the simulated events in which the W and Z bosons decay to electrons or muons.

Differential cross-sections

- Expected diff. cross-sections are parameterized in the $(c_{\tilde{W}}, c_{\tilde{B}})$ plane with quartic / quadratic polynomials for pure signal / interference ALP components.



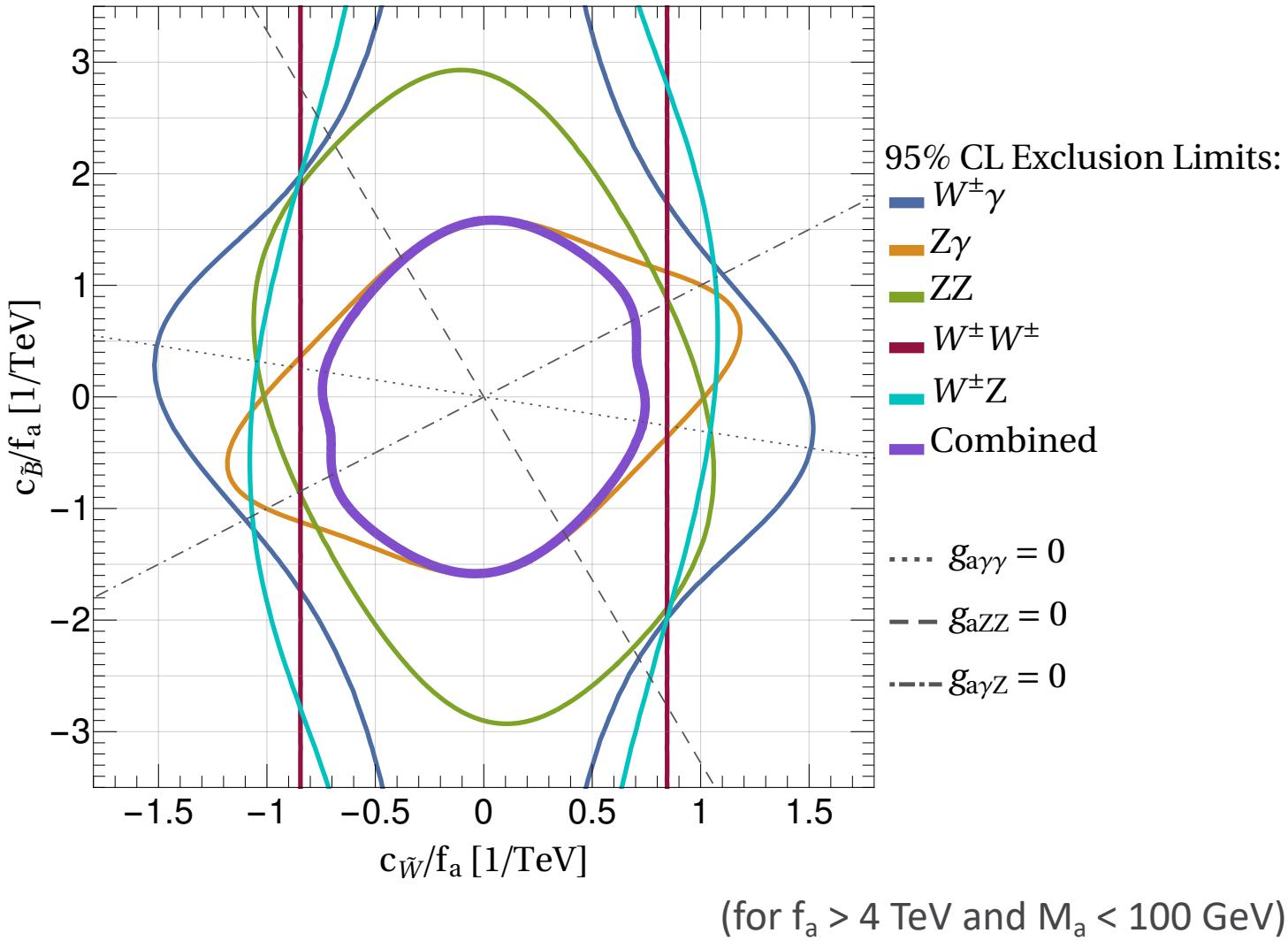
- MadGraph5_aMC@NLO reweighting tool for the generation at different points in the $(c_{\tilde{W}}, c_{\tilde{B}})$ plane:

$$g_{a\gamma Z} = 0 \longrightarrow p_0 = (1, 1), \quad p_1 = (0, 2), \quad p_2 = (1, 0),$$
$$p_3 = (1, -1), \quad p_4 = (1, -0.305), \quad p_5 = (1, -3.279) \quad g_{aZZ} = 0$$
$$g_{a\gamma\gamma} = 0$$

Results

CURRENT LIMITS

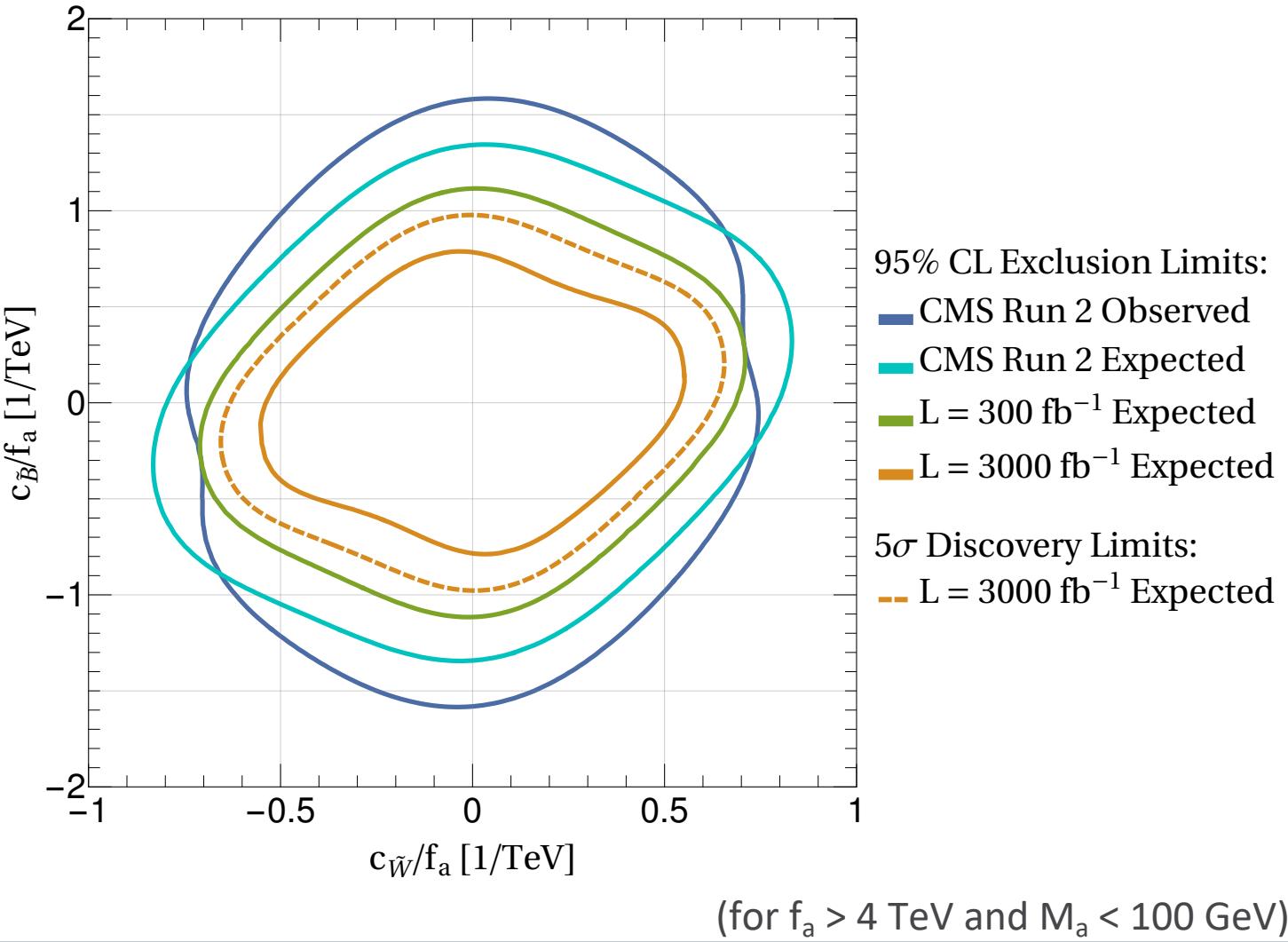
- Only ZZ and $Z\gamma$ can constrain $c_{\tilde{B}}/f_a$
- Most promising channels: $Z\gamma$ and $W^\pm W^\pm$



Results

PROJECTED LIMITS

- HL-LHC limits on ALP couplings decrease by a factor $\sim 1.5 - 1.7$
- HL-LHC limits on ALP cross sections decrease by a factor $\sim 5 - 8$

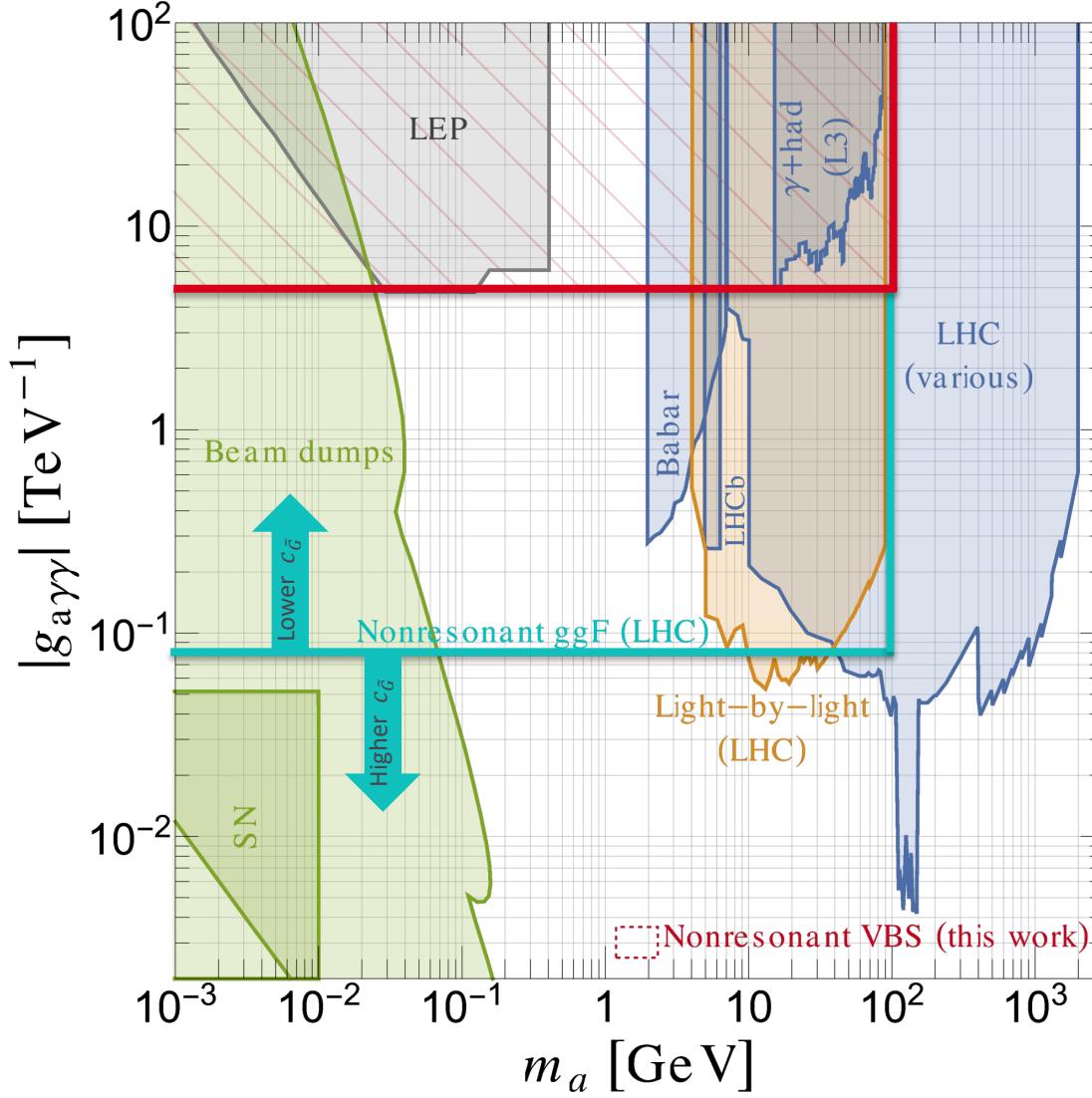


Results

Coupling [TeV ⁻¹]	Run 2 Observed (Expected)		300 fb ⁻¹		3000 fb ⁻¹		• 95% CL upper limits
	$M_{V_1 V_2} < 4 \text{ TeV}$	$< 2 \text{ TeV}$	$< 4 \text{ TeV}$	$< 2 \text{ TeV}$	$< 4 \text{ TeV}$	$< 2 \text{ TeV}$	
$ c_{\widetilde{W}}/f_a $	0.75 (0.83)	0.86 (0.94)	0.71	0.80	0.55	0.62	• CMS Run 2 measurements
$ c_{\widetilde{B}}/f_a $	1.59 (1.35)	1.73 (1.47)	1.12	1.23	0.79	0.87	• CMS Run 2 measurements
$ g_{a\gamma\gamma} $	4.99 (4.24)	5.45 (4.63)	3.50	3.84	2.43	2.68	• Projected sensitivities @ $\sqrt{s} = 14 \text{ TeV}$
$ g_{a\gamma Z} $	5.54 (4.74)	6.15 (5.25)	3.98	4.42	2.94	3.30	• Projected sensitivities @ $\sqrt{s} = 14 \text{ TeV}$
$ g_{aZZ} $	2.84 (3.02)	3.19 (3.38)	2.53	2.81	1.94	2.16	• Projected sensitivities @ $\sqrt{s} = 14 \text{ TeV}$
$ g_{aWW} $	2.98 (3.33)	3.43 (3.74)	2.84	3.18	2.21	2.49	• Projected sensitivities @ $\sqrt{s} = 14 \text{ TeV}$

(for $f_a > 4 \text{ TeV}$ and $M_a < 100 \text{ GeV}$)

Comparison to existing bounds



- **Red:** this work
- **Green:** no assumptions
- **Light blue:** nonresonant ggF. Depend on the coupling to gluons and assume $g_{agg} = 1 \text{ TeV}^{-1}$
- **Dark blue:** gluon dominance $g_{agg} \gg g_{aV_1V_2}$
- **Orange:** $BR(a \rightarrow \gamma\gamma) = 1$
- **Grey:** more elaborate assumptions on the EWK sector

Comparison to existing bounds

