



### Cabibbo Anomaly and Universality Tests

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#### Outline

- 1. What is the Cabbibo angle anomaly?
- 2. Why this anomaly?
- 3. Prospects
- 4. New Physics Interpretations
- 5. Conclusion and Outlook

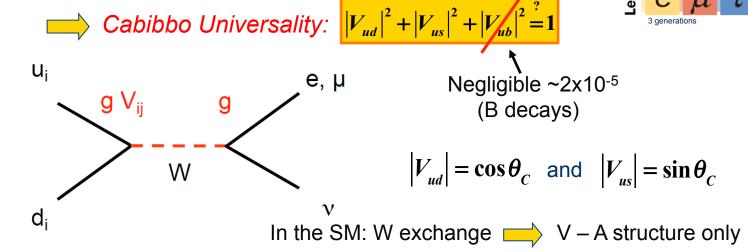
# 1. What is the Cabbibo angle anomaly?

### 1.1 Test of the Standard Model: V<sub>III</sub> and CKM unitarityq

- Extraction of the Cabibbo-Kobayashi-Maskawa matrix element  $V_{us}$ 
  - Fundamental parameter of the Standard Model

Description of the weak interactions:

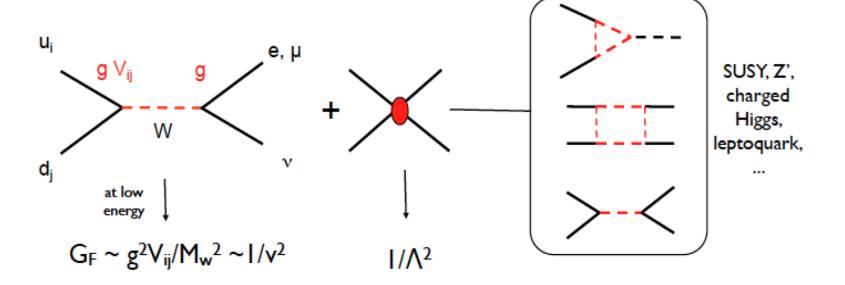
Check unitarity of the first row of the CKM matrix:



#### 1.2 Constraining New Physics

#### semi-reprome decays

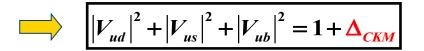
> BSM: sensitive to tree-level and loop effects of a large class of models

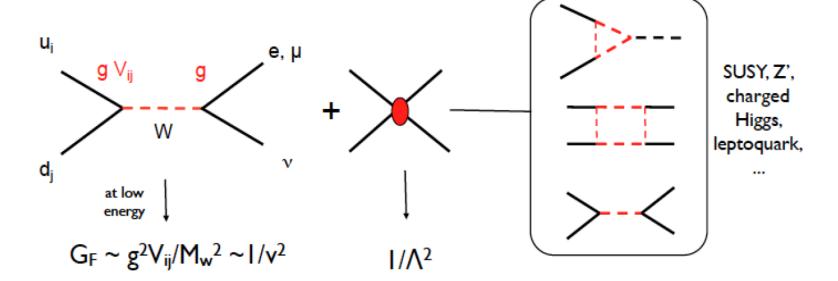


BSM effects : 
$$\Delta \sim \frac{c_n}{g^2} \frac{M_W^2}{\Lambda^2} \le 10^{-2} - 10^{-3} \longleftrightarrow \Lambda \sim 1-10 \text{ TeV}$$

### 1.2 Constraining New Physics

> BSM: sensitive to tree-level and loop effects of a large class of models



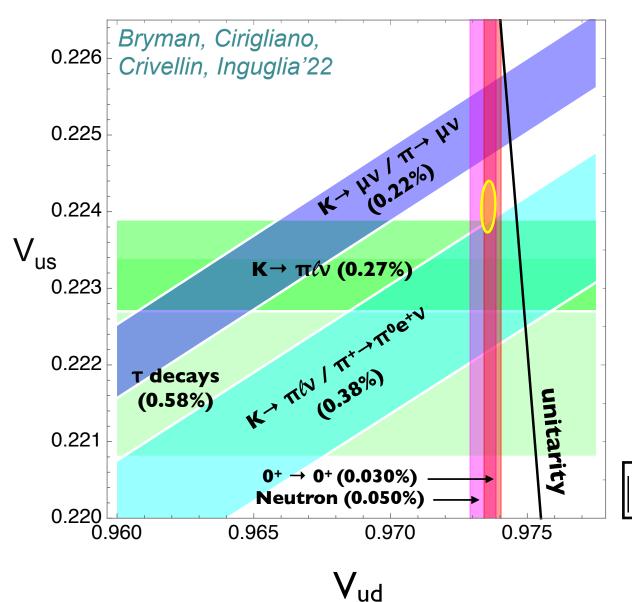


Grossman, E.P., Schacht'20

 $\blacktriangleright$  Look for new physics by comparing the extraction of  $V_{us}$  from different processes: helicity suppressed  $K_{\mu 2}$ , helicity allowed  $K_{l3}$ , hadronic  $\tau$  decays

### 1.2 Cabibbo angle anomaly

Moulson & E.P.@CKM2021



$$|V_{ud}| = 0.97373(31)$$
  
 $|V_{us}| = 0.2231(6)$   
 $|V_{us}|/|V_{ud}| = 0.2311(5)$ 

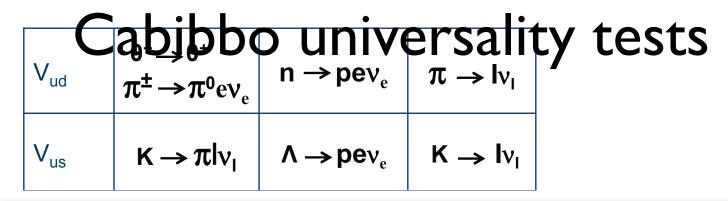
#### Fit results, no constraint

$$V_{ud} = 0.97365(30)$$
  
 $V_{us} = 0.22414(37)$   
 $\chi^2/\text{ndf} = 6.6/1 (1.0\%)$   
 $\Delta_{\text{CKM}} = -0.0018(6)$   
 $-2.7\sigma$ 

$$|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 1 + \Delta_{CKM}$$
Negligible ~2x10<sup>-5</sup>
(B decays)

### Paths to $V_{ud}$ and $V_{us}$

From kaon, pion, baryon and nuclear decays



$$\Gamma_k = (G_F^{(\mu)})^2 \times |V_{ij}|^2 \times |M_{\text{had}}|^2 \times (1 + \delta_{RC}) \times F_{\text{kin}}$$

Channel-dependent effective CKM element

Hadronic matrix element

Radiative corrections

Recent progress on 1) Hadronic matrix elements from lattice QCD

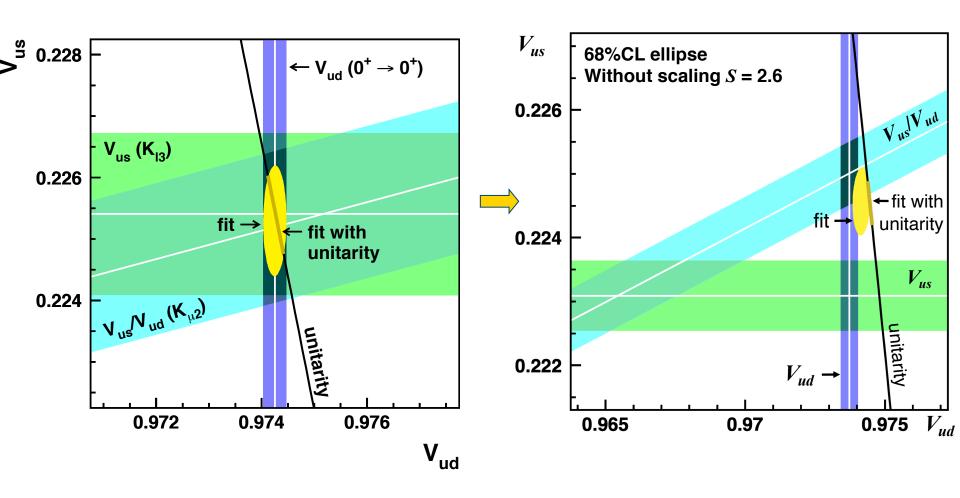
2) Radiative corrections from dispersive methods + Lattice QCD

### 2. Why this anomaly?

### 2.1 Changes on $V_{us}$ and $V_{ud}$ since 2011

Flavianet Kaon WG: Antonelli et al'11

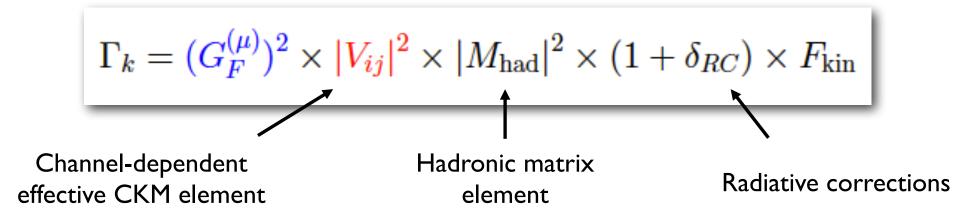
Moulson & E.P.@CKM2021



### 2.1 Changabibboardriversality tests

Almost no change on the experimental side since 2011

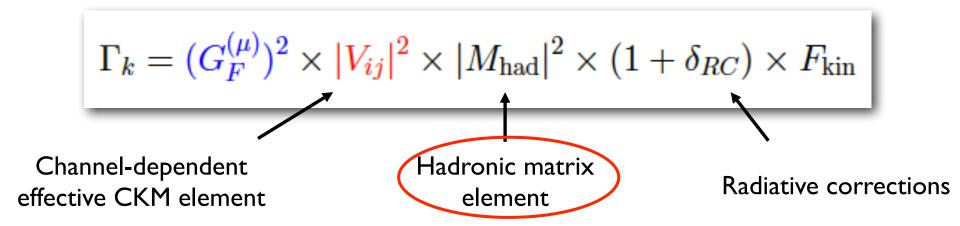
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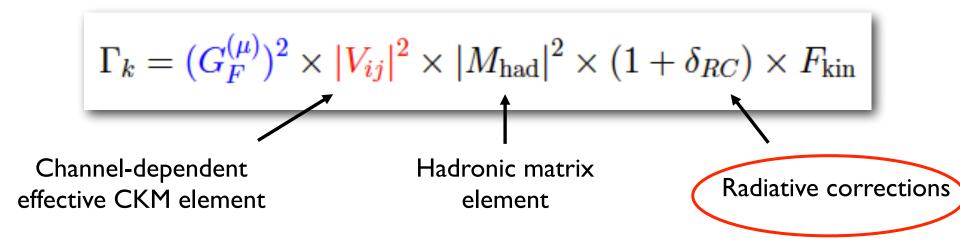


- Changes in theoretical inputs:
  - Impressive progress on hadronic matrix element computations from lattice QCD for V<sub>us</sub> and V<sub>us</sub>/V<sub>ud</sub> extraction from Kaon decays

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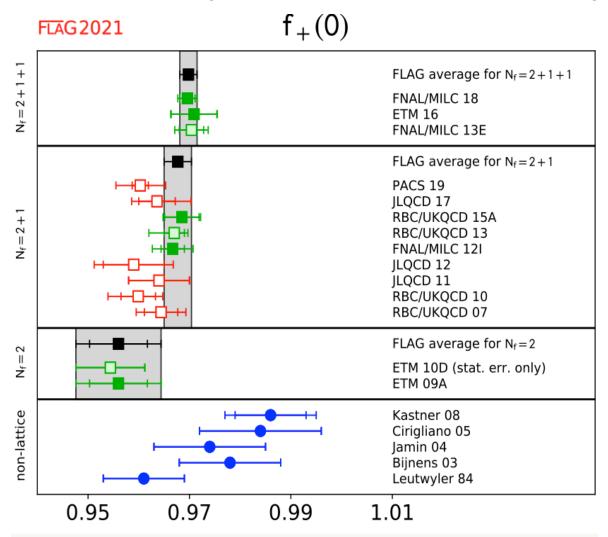


- Changes in theoretical inputs:
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Radiative corrections from dispersive methods for V<sub>ud</sub> extraction

### 2.2 $f_{+}(0)$ from lattice QCD

Recent progress on Lattice QCD for determining f<sub>+</sub>(0)



$$f_{+}(0)_{N_f=2+1+1}^{FLAG21} = 0.9698(17)$$

0.18% uncertainty

to be compared to

$$f_{+}(0)_{N_f=2+1+1}^{FLAG16} = 0.9704(32)$$

$$f_{+}(0)_{N_f=2+1}^{2010} = 0.959(50)$$

Uncertainty divided by ~2 w/ 2016 and by 25 w/ 2011!



Lattice uncertainties at the same level as exp.

 $-3.2\sigma$  away from unitarity!

2011: 
$$V_{us} = 0.2254(5)_{exp}(11)_{lat} \rightarrow V_{us} = 0.2231(4)_{exp}(4)_{lat}$$

### $V_{us}/V_{ud}$ from $K_{l2}/\pi_{l2}$

$$\frac{|V_{us}|}{|V_{ud}|} \frac{f_K}{f_{\pi}} = \left(\frac{\Gamma_{K_{\mu 2(\gamma)}} m_{\pi^{\pm}}}{\Gamma_{\pi_{\mu 2(\gamma)}} m_{K^{\pm}}}\right)^{1/2} \frac{1 - m_{\mu}^2 / m_{\pi^{\pm}}^2}{1 - m_{\mu}^2 / m_{K^{\pm}}^2} \left(1 - \frac{1}{2} \delta_{\text{EM}} - \frac{1}{2} \delta_{\text{SU}(2)}\right)$$

Recent progress on radiative corrections computed on lattice:

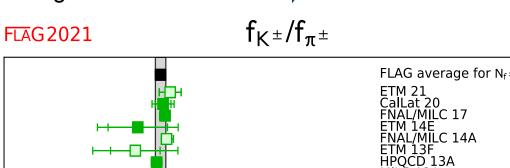
Di Carlo et al.'19

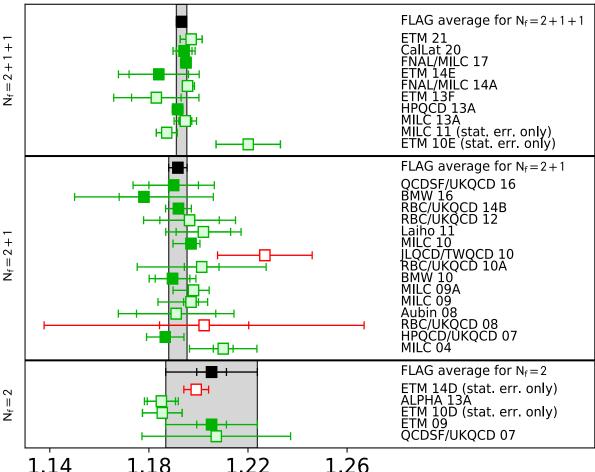
- Main input hadronic input: f<sub>K</sub>/f<sub>π</sub>
- In 2011:  $V_{us}/V_{ud} = 0.2312(4)_{exp}(12)_{lat}$
- In 2021:  $V_{us}/V_{ud} = 0.2311(3)_{exp}(4)_{lat}$  the lattice error is reducing by a factor of 3 compared to 2011! It is now of the same order as the experimental uncertainty.

-1.8σ away from unitarity

### 2.2 $f_K/f_{\pi}$ from lattice QCD

Progress since 2018: new results from *ETM'21* and *CalLat'20* 





Now Lattice collaborations include SU(2) IB corr. For  $N_f$ =2+1+1, FLAG2021

$$f_{K^+}/f_{\pi^+} = 1.1932(21)$$

0.18% uncertainty

Results have been stable over the years

For average substract IB corr.

$$f_{K}/f_{\pi} = 1.1967(18)$$

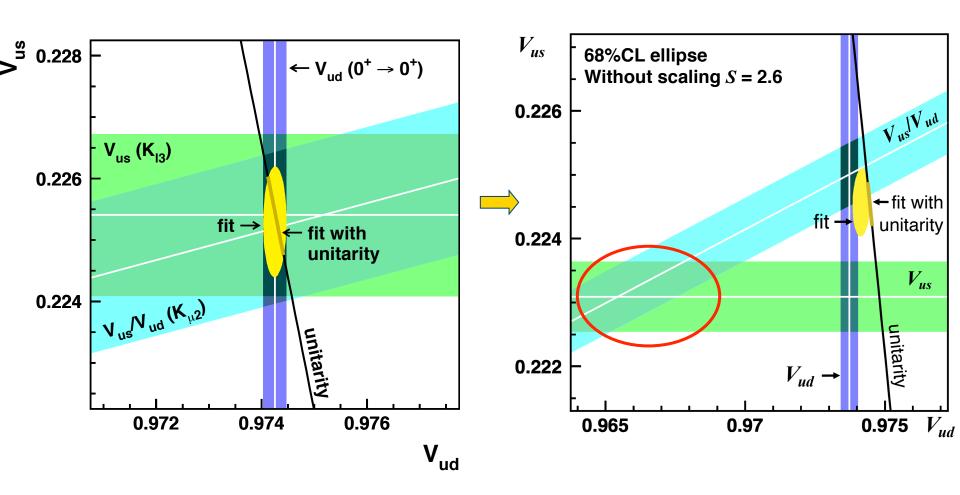
In 2011: 
$$f_K/f_\pi = 1.193(6)$$

$$V_{us}/V_{ud} = 0.23108(29)_{exp}(42)_{lat}$$

### Changes on V<sub>us</sub> and V<sub>ud</sub> since 2011

Flavianet Kaon WG: Antonelli et al'11

Moulson & E.P.@CKM2021

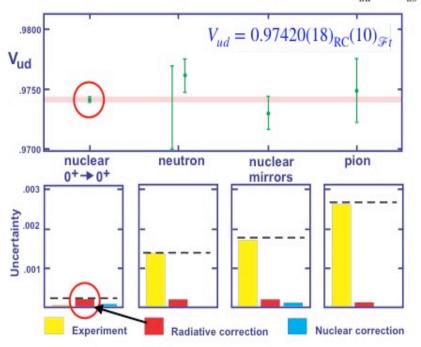


### 2.3 $|V_{ud}|$ from $0^+ \rightarrow 0^+$ superallowed $\beta$ decays

See Talk by Misha Gorshteyn @CKM2021

PDG 2018:

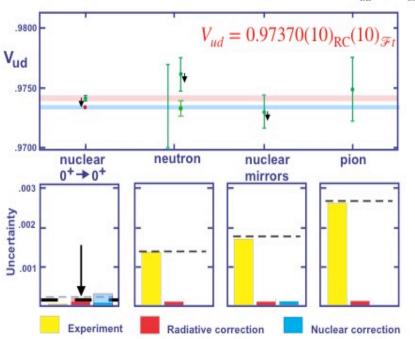
 $|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 0.9994(4)_{V_{ud}}(2)_{V_{us}}$ 



PDG 2020:

Figure adapted from J. Hardy

$$|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 0.9985(3)_{V_{ud}}(4)_{V_{us}}$$



Recent improvement on the theoretical RCs +Nuclear Structure Corrections

Use of a data driven dispersive approach

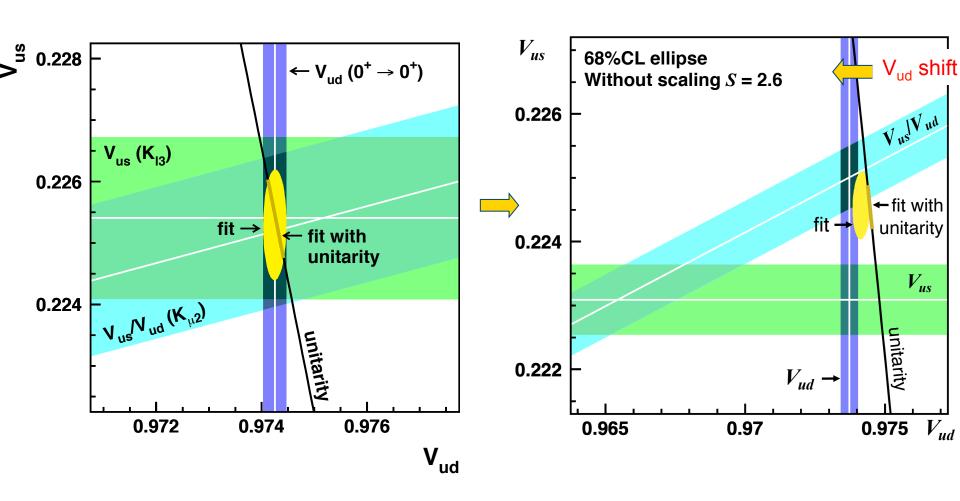
Seng et al.'18'19, Gorshteyn'18

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### Changes on V<sub>us</sub> and V<sub>ud</sub> since 2011

Flavianet Kaon WG: Antonelli et al'11

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### 3. Prospects

### 3.1 Experimental Prospects for V<sub>us</sub>

#### On Kaon side

Cirigliano et al'22

- NA62 could measure several BRs:  $K_{\mu 3}/K_{\mu 2}$ ,  $K \to 3\pi$ ,  $K_{\mu 2}/K \to \pi\pi$
- Note that the high precision measurement of BR( $K_{\mu 2}$ ) (0.3%) comes only from a single experiment: KLOE. It would be good to have another measurement at the same level of accuracy
- *LHCb*: could measure BR( $K_S \rightarrow \pi \mu \nu$ ) at the < 1% level?  $K_S \rightarrow \pi \mu \nu$  measured by KLOE-II but not competitive  $\tau_S$  known to 0.04% (vs 0.41% for  $\tau_L$ , 0.12% for  $\tau_\pm$ )
- V<sub>IIS</sub> from Tau decays at Belle II:

Belle II with 50 ab<sup>-1</sup> and ~4.6 x  $10^{10}~\tau$  pairs will improve  $V_{us}$  extraction from  $\tau$  decays

Inclusive measurement is an opportunity to have a complete independent extraction of  $V_{us}$  not easy as you have to measure many channels



$$|V_{us}| = 0.2184 \pm 0.0018_{\text{exp}} \pm 0.0011_{\text{th}}$$

To be competitive theory error will have to be improved as well

HFI AV'21

### V<sub>us</sub> from Hyperon decays



V<sub>us</sub> can be measured from Hyperon decays:

- Λ → pev<sub>e</sub> Possible measurement at BESIII, Super τ-Charm factory?
- Possibilities at LHCb?

Talk by Dettori@FPCP20

Channel	${\cal R}$	$\epsilon_L$	$\epsilon_D$	$\sigma_L({ m MeV}/c^2)$	$\sigma_D({ m MeV}/c^2$	R = ratio of
$K_{\rm S}^0  o \mu^+\mu^-$	1	1.0 (1.0)	1.8 (1.8)	$\sim 3.0$	$\sim 8.0$	1 4 •
$K_{\scriptscriptstyle \mathrm{S}}^0  o \pi^+\pi^-$	1	$1.1\ (0.30)$	1.9(0.91)	$\sim 2.5$	$\sim 7.0$	production
$K_{\mathrm{S}}^{0}  ightarrow \pi^{0} \mu^{+} \mu^{-}$	1	0.93 (0.93)	1.5 (1.5)	$\sim 35$	$\sim 45$	$\epsilon = \text{ratio of}$
$K_{\rm S}^0 \to \gamma \mu^+ \mu^-$	1	0.85 (0.85)	1.4 (1.4)	$\sim 60$	$\sim 60$	$\epsilon = 1a00 01$
$K_{\rm S}^0 \to \mu^+ \mu^- \mu^+ \mu^-$	1	0.37(0.37)	1.1(1.1)	$\sim 1.0$	$\sim 6.0$	efficiencies
$K_{\rm L}^0  o \mu^+ \mu^-$	$\sim 1$	$2.7 (2.7) \times 10^{-3}$	0.014 (0.014)	$\sim 3.0$	$\sim 7.0$	
$K^+ \to \pi^+ \pi^+ \pi^-$	$\sim 2$	$9.0 (0.75) \times 10^{-3}$	$41 (8.6) \times 10^{-3}$	$\sim 1.0$	$\sim 4.0$	
$K^+ \to \pi^+ \mu^+ \mu^-$	$\sim 2$	$6.3(2.3)\times10^{-3}$	0.030(0.014)	$\sim 1.5$	$\sim 4.5$	
$\Sigma^+ \to p \mu^+ \mu^-$	$\sim 0.13$	0.28 (0.28)	0.64 (0.64)	$\sim 1.0$	$\sim 3.0$	
$\Lambda  o p\pi^-$	$\sim 0.45$	$0.41 \ (0.075)$	1.3(0.39)	$\sim 1.5$	$\sim 5.0$	
$\Lambda  o p \mu^- ar{ u_\mu}$	$\sim 0.45$	0.32(0.31)	0.88 (0.86)	_	_	
$\Xi^-  o \Lambda \mu^- ar{ u_\mu}$	$\sim 0.04$	$39 (5.7) \times 10^{-3}$	0.27 (0.09)	—	_	
$\Xi^-  o \Sigma^0 \mu^- \bar{\nu}$	$\sim 0.03$	$24 (4.9) \times 10^{-3}$	$0.21\ (0.068)$	_	_	
$\Xi \rightarrow p\pi^-\pi^-$	$\sim 0.03$	0.41(0.05)	0.94(0.20)	$\sim 3.0$	$\sim 9.0$	
$\Xi^0 o p\pi^-$	$\sim 0.03$	1.0(0.48)	2.0(1.3)	$\sim 5.0$	$\sim 10$	
$\Omega^- \to \Lambda \pi^-$	$\sim 0.001$	$95(6.7) \times 10^{-3}$	0.32(0.10)	$\sim 7.0$	$\sim 20$	

To be able to extract V<sub>us</sub> one needs to compute form factors precisely

### 3.2 Theoretical Prospects for $V_{us}$

- Lattice Progress on hadronic matrix elements: decay constants,
   FFs
- Full QCD+QED decay rate on the lattice, for Leptonic decays of kaons and pions 

   Inclusion of EM and IB corrections :
  - Perturbative treatment of QED on lattice established
  - Formalism for  $K_{12}$  worked out
- Application of the method for semileptonic Kaon (K<sub>I3</sub>) and Baryon decays

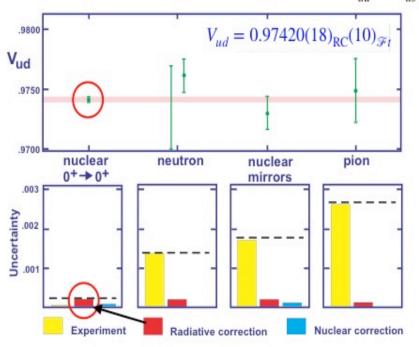
Aim: Per mille level within 10 years

### 3.3 Prospects for $|V_{ud}|$

See Talk by Misha Gorshteyn @CKM2021

PDG 2018:

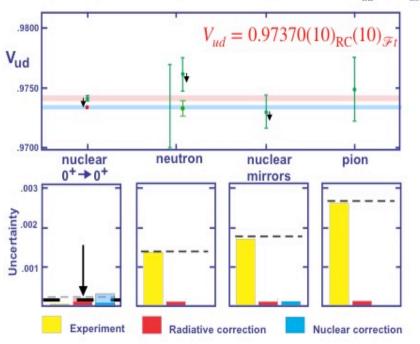
$$|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 0.9994(4)_{V_{ud}}(2)_{V_{us}}$$



PDG 2020:

Figure adapted from J. Hardy

$$|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 0.9985(3)_{V_{ud}}(4)_{V_{us}}$$



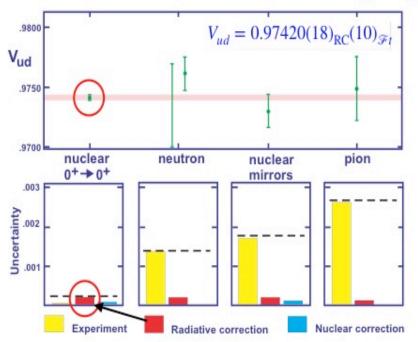
- From neutron decays: very impressive progress recently
- From pion  $\beta$  decay  $\pi^+ \to \pi^0 e^+ v$ : PIONEER experiment

### 3.3 Prospects for $|V_{ud}|$

See Talk by Misha Gorshteyn @CKM2021

PDG 2018:

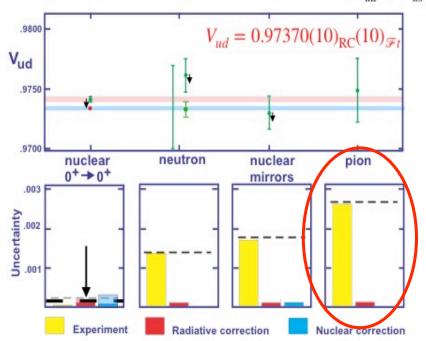
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PDG 2020:

Figure adapted from J. Hardy

$$|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 0.9985(3)_{V_{ud}}(4)_{V_{us}}$$



- From neutron decays
- From pion β decay π<sup>+</sup> → π<sup>0</sup>e<sup>+</sup>v : PIONEER experiment
   (Phase-I) approved at PSI, physics starting in ~2029

### $|V_{ud}|$ from pion $\beta$ decay: $\pi^+ \to \pi^0 e^+ v$

- Theoretically cleanest method to extract V<sub>ud</sub>: corrections computed in SU(2)
   ChPT
   Sirlin'78, Cirigliano et al.'03, Passera et al'11
- Present result: *PIBETA* Experiment (2004) → **Uncertainty: 0.64%**

$$\mathbf{B}(\pi^+ \to \pi^0 e^+ \nu) = (1.036 \pm 0.004_{\text{stat}} \pm 0.004_{\text{syst}} \pm 0.003_{\pi e2}) \times 10^{-8} (\pm 0.6\%)$$

$$|V_{ud}| = 0.9739(28)_{\text{exp}}(1)_{\text{th}}$$
 to be compared to  $|V_{ud}| = 0.97373(31)$ 

- Reduction of the theory error thanks to a new lattice calculation for RC Feng et al'20
- ▶ Wext generation experiment PIONEER Phase II and III measurement at 0.02% level → will be competitive with current 0+ → 0+ extraction
- Would be completely independent check! No nuclear correction and different RCs compared to neutron decay
- Opportunity to extract  $V_{us}/V_{ud}$  from  $\frac{B(K \to \pi l \nu)}{B(\pi^+ \to \pi^0 e^+ \nu)}$  EW Rad. Corr. cancel Improve precision on  $B(\pi^+ \to \pi^0 e^+ \nu)$  by x3  $\longrightarrow$   $V_{us}/V_{ud} < \pm 0.2\%$

### Pion decays and LFU tests

- Lepton Flavor Universality test in
  - Early insight into the V-A structure of weak interactions

$$\pm (R_{e/\mu}({
m SM}) = 1.23524(015) \times 10^{-4}$$

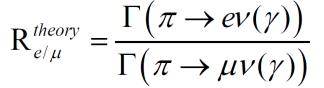
Cirigliano & Rosell'07

G): 
$$R_{e/\mu}^{\text{exp}} = (1.2327 \pm 0.0023) x 10^{-4} \ (\pm 0.19\%)$$

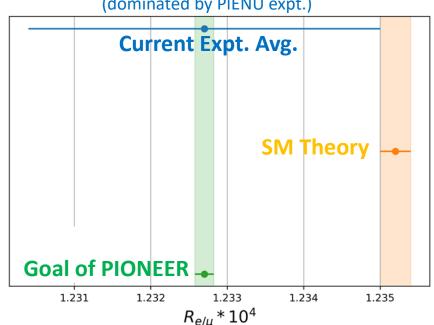
$$R / (Exp) = 1.23270(230) \times 10^{-4}$$

G): 
$$R_{e/\mu}^{\text{exp}} = (1.2327 \pm 0.0023) x 10^{-4} \ (\pm 0.19\%)$$
15 times worse than theory!





(dominated by PIENU expt.)



$$\frac{g_e}{g_{\mu}} = 0.9990 \pm 0.0009 \ (\pm 0.09\%)$$

Goal of PIONEER: reduce unc. by a factor of 10!  $\Longrightarrow$  by far most precise test of LFU

### 4. New Physics Interpretations

### 4.1 Right-handed currents

Bernard, Oertel, E.P., Stern'08

$$\mathcal{L}_{W} = \frac{e(1 - \xi^{2} \rho_{L})}{\sqrt{2}s} \left\{ \bar{N}_{L} V_{MNS} \gamma^{\mu} L_{L} + (1 + \delta) \bar{U}_{L} V_{L} \gamma^{\mu} D_{L} + \epsilon \bar{U}_{R} V_{R} \gamma^{\mu} D_{R} \right\} W_{\mu}^{+} + \text{h.c.}$$

See also Antonelli et al.'09

Alioli, Cirigliano, Dekens, de Vries, Mereghetti'17

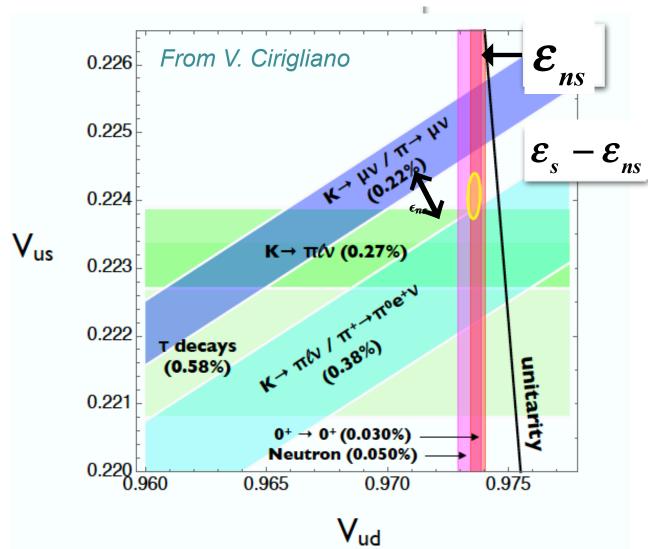
T. Kitahara@HC2NP 2019

### 4.1 Right-handed Currents

$$\begin{split} V_{us}^{K_{l3}} &= \left| \sin \theta_C + \varepsilon_s \right| \;, \qquad \qquad \text{Vector s quark} \\ \left( \frac{V_{us}}{V_{ud}} \right)^{K_{l2}} &= \left| \frac{\sin \theta_C - \varepsilon_s}{\cos \theta_C - \varepsilon_{ns}} \right| \; \longleftarrow \quad \text{Axial} \\ V_{ud}^{\beta} &= \left| \cos \theta_C + \varepsilon_{ns} \right| \; \longleftarrow \quad \text{Vector no s quark} \\ \left( \frac{V_{us}}{V_{ud}} \right)^{K_{\ell3}} &= \left| \frac{\sin \theta_C + \epsilon_s}{\cos \theta_C + \epsilon_{ns}} \right| \; \longleftarrow \quad \text{Vector} \end{split}$$

- The SM is obtained in the limit  $\varepsilon_s = \varepsilon_{ns} = 0$ .
- Perfect fit to data  $\chi^2_{
  m min,RH}=0$
- Not obvious how to define CKM unitarity test in this case

### 4.1 Right-handed Currents

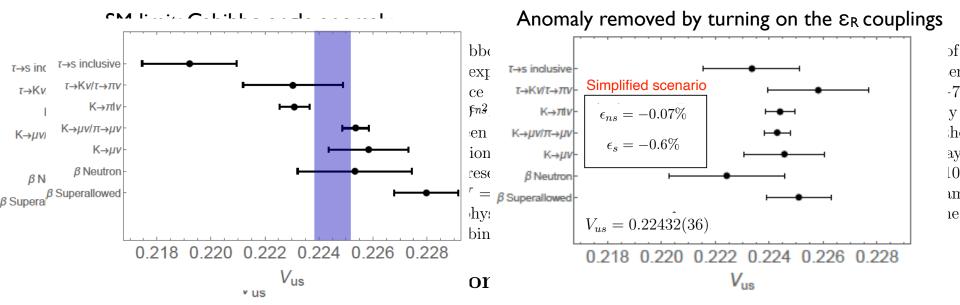


Emilie Passemar

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## 4.1 Right-handed Currents \( \times \times

- Global fit to CC processes involving light quarks and all lepton families
- SM hypothesis ( $\varepsilon_s = \varepsilon_{ns} = 0$ ) disfavored (p-value 0.3%)



In this paper we studied hadronic tau decays in the framework of an EFT for light SM of Ciricular Discrete the for the form of the SM Aloys of the series of hyperone and particles with masses larger than 2 GeV. Focusing on the charged-current into between light quarks and leptons, the leading non-standard effects are parametrized by Wilson coefficients  $\epsilon_X^{q\ell}$ , cf. Eq. (2.1). The main new result of this paper is Eq. (6.1) sum the constraints on  $\epsilon_X^{q\ell}$  from a large set of hadronic tau observables, which include the

Emilie Passemar  $\tau \to \pi(K)\nu$ , 3-body  $\tau \to \pi\pi\nu$ , and inclusive  $\tau \to \nu\bar{u}d(s)$  decays. There we quote per marginalized constraints on six linear combinations of  $\epsilon_X^{D\tau}$ , D=d,s, and we provide the constraints of  $\epsilon_X^{D\tau}$  and  $\epsilon_X^{D\tau}$  are the constraints of  $\epsilon_X^{D\tau}$ .

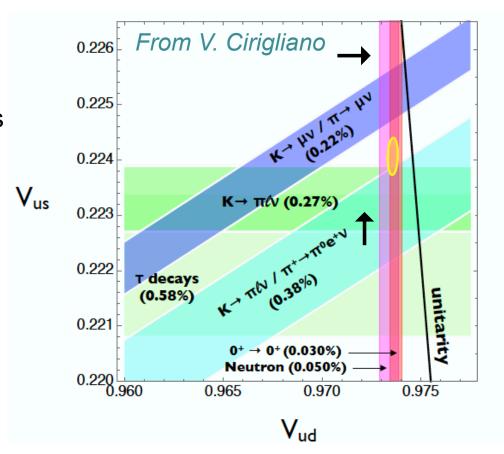
# 4.2 Other New Physics Models V

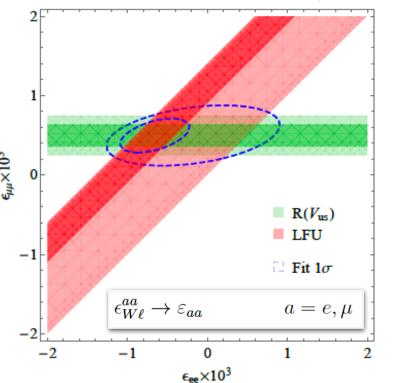
4th quark b'

- Belfatto, Beradze, Berezhiani'19
- Gauge horizontal family symmetry
- Turn on only vertex corrections to leptons Crivellin & Hoferichter'21

Shift the location of the Vud,us bands but do not solve the tension between ratios

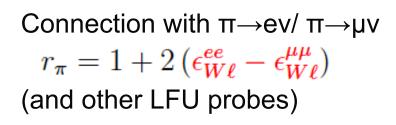
And many more....



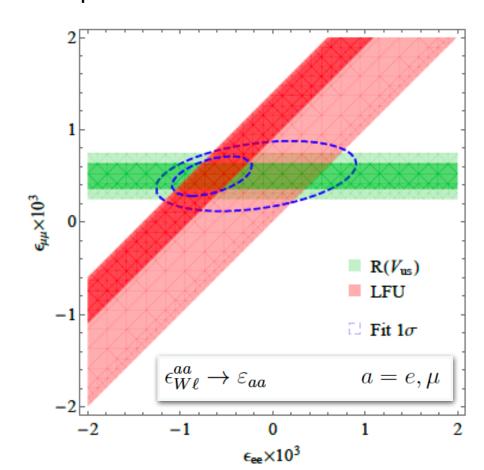


### LFUL

Belfatto, Beradze, Berezhiani'19



And many more....



### 5. Conclusion and Outlook

#### Conclusion and Outlook

- Recent precision determinations of V<sub>us</sub> and V<sub>ud</sub> enable unprecedented tests of the SM and constraints on possible NP models
- Tensions in unitarity of 1<sup>st</sup> row of CKM matrix have reappeared!
- We need to work hard to understand where they come from:
  - On experimental side:
     For V<sub>us</sub>, new measurements in kaons (NA62: K<sub>μ3</sub>/K<sub>μ2</sub>, LHCb?)
     but mainly in tau decays from Belle II
     V<sub>us</sub> from hyperon decays?
     BESSIII, LHCb?
  - For  $V_{ud}$ , understand the situation of the neutron lifetime, beta decay of pion? PIONEER Consider  $R_V = \Gamma\left(K \to \pi l \nu(\gamma)\right)/\Gamma\left(\pi^+ \to \pi^0 e^+ \nu(\gamma)\right)$

On theory side:
 Calculate very precisely radiative corrections, isospin breaking effects and matrix elements
 Be sure the uncertainties are under control

- If these tensions are confirmed what do they tell us?
- Interesting time ahead of us!

# 6. Back-up

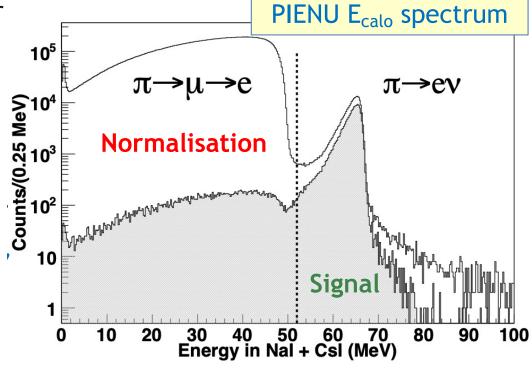
### PIONEER (Phase-I)

#### PIONEER (Phase-I) approved at PSI, physics starting in ~2029

- ➢ Goal: matching the SM precision on R<sub>e/µ</sub>
   ➡ Test of New Physics at 1 PeV scale
- $\triangleright$  Stopped  $\pi^+$  at high rate (300 kHz), focus on reduction of systematics.
- Detectors: highly-segmented LGAD active target, positron tracker, LXe calorimeter
- > Collection of 2×10<sup>8</sup> π<sup>+</sup>→ e<sup>+</sup>ν<sub>e</sub> events in three years.
- > Key point: control of the  $\pi^+ \rightarrow e^+ v_e$  signal tail in the calorimeter to a  $10^{-4}$  precision

#### PIONEER Phase II,III:

 $V_{ud}$  from  $\pi^+ \rightarrow \pi^0 e^+ v_e$ decays to a 0.02% level



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Emilie Passemar

## 1.1 Test of the Standard Model: V<sub>us</sub> and CKM unitarity

- Extraction of the Cabibbo-Kobayashi-Maskawa matrix element  $V_{us}$ 
  - > Fundamental parameter of the Standard Model

Description of the 
$$\frac{g}{\sqrt{2}}W_{\alpha}^{+}$$
 ( $\overline{\mathbf{U}}_{L}\mathbf{V}_{\mathrm{CKM}}\gamma^{\alpha}\mathbf{D}_{L}+\overline{e}_{L}\gamma^{\alpha}\nu_{e\,L}+\overline{\mu}_{L}\gamma^{\alpha}\nu_{\mu\,L}+\overline{\tau}_{L}\gamma^{\alpha}\nu_{\tau}$ )

$$\frac{g}{\sqrt{2}}W_{\alpha}^{+}\left(\overline{\mathbf{U}}_{L}\mathbf{V}_{\mathrm{CKN}}\right)$$

$$|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 1$$

$$|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 1$$

coupling

**Emilie Passen** 

$$G_{\mu}^2 = (g_{\mu}g_e)^2/M_W^4 \quad \stackrel{?}{=} \quad G_{\rm CKM}^2 = (g_qg_\ell)^2 \, (|V_{ud}|^2 + |V_{us}|^2)/M_W^4$$
 u<sub>i</sub> g V<sub>ij</sub> g



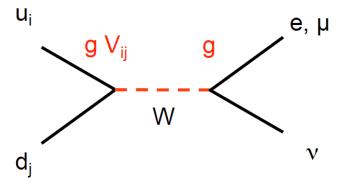
## 1.2 Constraining New Physics

- Extraction of the Cabibbo-Kobayashi-Maskawa matrix element  $V_{us}$ 
  - Fundamental parameter of the Standard Model

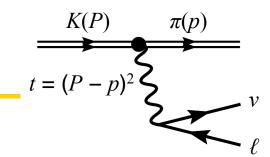
Description of the weak interactions:

$$\mathcal{L}_{EW} = \frac{g}{\sqrt{2}} W_{\alpha}^{+} \left( \overline{D}_{L} V_{CKM} \gamma^{\alpha} U_{L} + \overline{e}_{L} \gamma^{\alpha} v_{eL} + \overline{\mu}_{L} \gamma^{\alpha} v_{\mu L} + \overline{\tau}_{L} \gamma^{\alpha} v_{\tau L} \right) + \text{h.c.}$$

- Look for new physics
  - ➤ In the Standard Model : W exchange only V-A structure



# 2.2 $V_{us}$ from $K_{l3}$ ( $K \rightarrow \pi l \nu_l$ )



• Master formula for  $K \rightarrow \pi lv_l$ :  $K = \{K^+, K^0\}$ ,  $l=\{e, \mu\}$ 

$$\Gamma(K \to \pi l \nu [\gamma]) = Br(K_{l3}) / \tau = C_K^2 \frac{G_F^2 m_K^5}{192\pi^3} S_{EW}^K |V_{us}|^2 |f_+^{K^0 \pi^-}(0)|^2 I_{Kl} (1 + 2\Delta_{EM}^{Kl} + 2\Delta_{SU(2)}^{Kl})$$

Average and work by Flavianet Kaon WG Antonelli et al 11 and then by  $\overline{m_K^2 - m_\pi^2}$  M. Moulson, see e.g. Moulson.@CKM2021

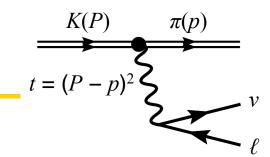
Theoretically 
$$\gamma$$
 =  $\frac{Br(K_{I3})}{mlv}$  =  $C_K^2 \frac{G_F^2 m_K^5}{400} S_{EW}^5 |V_{us}|^2 |f_+^{K^0 \pi^-}(0)|^2 I_K (1 + \delta_{EM}^{KI} + \delta_{SU(2)}^{KI})^2$ . Update on long-distance EM corrections for  $K_{e3}^{EW}$ 

- Improvement on Isospin breaking evaluation due to more precise dominant input: quark mass ratio from  $\eta \to 3\pi$  Colangelo et al.'18
- Progress from lattice QCD on the  $K \rightarrow \pi$  FF

$$\boxed{\left\langle \pi^{-}(p) \middle| \overline{s} \gamma_{\mu} \mathbf{u} \middle| \mathbf{K}^{0}(\mathbf{P}) \right\rangle = \mathbf{f}_{+}^{K^{0}\pi^{-}}(\mathbf{0}) \left[ \left( P + p \right)_{\mu} \overline{f}_{+}^{K^{0}\pi^{-}}(t) + \left( P - p \right)_{\mu} \overline{f}_{-}^{K^{0}\pi^{-}}(t) \right]}$$

Emilie Passemar  $f_+(t)$ 

# 2.2 $V_{us}$ from $K_{l3}$ ( $K \rightarrow \pi l \nu_l$ )



• Master formula for  $K \rightarrow \pi lv_l$ :  $K = \{K^+, K^0\}$ ,  $l=\{e, \mu\}$ 

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Average and work by Flavianet Kaon WG Antonelli et al 11 and then by  $\overline{m_K^2 - m_\pi^2}$  *M. Moulson*, see e.g. *Moulson*.@*CKM2021* 

Theoretically 
$$\gamma$$
 =  $\frac{Br(K_{I3})}{mlv}$  =  $C_K^2 \frac{G_F^2 m_K^5}{400} S_{EW}^5 |V_{us}|^2 |f_{+Seng}^{K^0\pi^-}(0)|^2 I_{EM} + \delta_{SU(2)}^{KI} + \delta_{SU(2)}^{KI}$  Update on long-distance EM corrections for  $K_{e3}^{K}$ 

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- Progress from lattice QCD on the  $K \rightarrow \pi$  FF

$$\left| \left\langle \pi^{-}(p) \right| \overline{s} \gamma_{\mu} \mathbf{u} \left| \mathbf{K}^{0}(\mathbf{P}) \right\rangle = \mathbf{f}_{+}^{K^{0} \pi^{-}}(\mathbf{0}) \left[ \left( P + p \right)_{\mu} \overline{f}_{+}^{K^{0} \pi^{-}}(t) + \left( P - p \right)_{\mu} \overline{f}_{-}^{K^{0} \pi^{-}}(t) \right] \right|$$

 $ilde{f}_+(t)$ 

Emilie Passemar 4

# 2.3 $V_{us}/V_{ud}$ from $K_{l2}/\pi_{l2}$

$$\frac{|V_{us}|}{|V_{ud}|} \frac{f_K}{f_{\pi}} = \left(\frac{\Gamma_{K_{\mu 2(\gamma)}} m_{\pi^{\pm}}}{\Gamma_{\pi_{\mu 2(\gamma)}} m_{K^{\pm}}}\right)^{1/2} \frac{1 - m_{\mu}^2 / m_{\pi^{\pm}}^2}{1 - m_{\mu}^2 / m_{K^{\pm}}^2} \left(1 - \frac{1}{2} \delta_{\text{EM}} - \frac{1}{2} \delta_{\text{SU}(2)}\right)$$

Recent progress on radiative corrections computed on lattice:

#### First lattice calculation of EM corrections to $P_{l2}$ decays

Ensembles from ETM

Giusti et al.'18

•  $N_f = 2+1+1$  Twisted-mass Wilson fermions

$$\delta_{SU(2)} + \delta_{EM} = -0.0122(16)$$

Uncertainty from quenched QED included (0.0006)

Compare to ChPT result from Cirigliano, Neufeld '11:

$$\delta_{SU(2)} + \delta_{EM} = -0.0112(21)$$

Update, extended description, and systematics of Giusti et al.

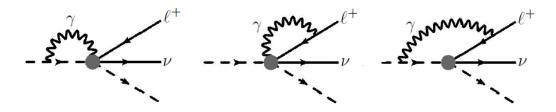
$$\delta_{SU(2)} + \delta_{EM} = -0.0126(14)$$

Di Carlo et al.'19

#### Progress since 2018:

• First experimental measurement of BR of  $K_S \to \pi \mu \nu$ BR( $K_S \to \pi \mu \nu$ ) = (4.56 ± 0.20) ×10<sup>-4</sup> **KLOE-2** PLB 804 (2020)

Theoretically update on long-distance EM corrections:



Up to now computation at fixed order  $e^2p^2$  + model estimate for the LECs

Cirigliano et al. '08

New calculation of complete EW RC using hybrid current algebra and ChPT (Sirlin's representation) with resummation of largest terms to all chiral orders

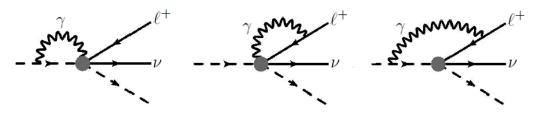
- Reduced uncertainties at  $O(e^2p^4)$
- Lattice evaluation of QCD contributions to γW box diagrams

Seng et al.'21

#### Progress since 2018:

• First experimental measurement of BR of  $K_S \rightarrow \pi \mu \nu$ BR( $K_S \rightarrow \pi \mu \nu$ ) = (4.56 ± 0.20) ×10<sup>-4</sup> **KLOE-2** PLB 804 (2020)

Theoretically update on long-distance EM corrections:



Only  $K_{e3}$  at present For  $K_{\mu3}$  modes continue to use Cirigliano et al. '08

	Cirigliano et al. '08	Seng et al. '21
$\Delta^{EM}(K^0_{e3})$ [%]	$0.50 \pm 0.11$	$0.580 \pm 0.016$
$\Delta^{EM}(K^+_{e3})$ [%]	$0.05 \pm 0.13$	$0.105 \pm 0.024$
ρ	+0.081	-0.039

#### Progress since 2018:

Theoretical progress on isospin breaking correction

$$\begin{split} \Delta^{SU(2)} &\equiv \frac{f_+(0)^{K^+\pi^0}}{f_+(0)^{K^0\pi^-}} - 1 \quad \text{Strong isospin breaking} \\ &= \frac{3}{4} \frac{1}{Q^2} \left[ \frac{\overline{M}_K^2}{\overline{M}_\pi^2} + \frac{\chi_{p^4}}{2} \left( 1 + \frac{m_s}{\hat{m}} \right) \right] \quad Q^2 = \frac{m_s^2 - \hat{m}^2}{m_d^2 - m_u^2} \quad \text{NLO in strong interaction O} \\ &= \frac{3}{4} \frac{1}{Q^2} \left[ \frac{\overline{M}_K^2}{\overline{M}_\pi^2} + \frac{\chi_{p^4}}{2} \left( 1 + \frac{m_s}{\hat{m}} \right) \right] \quad Q^2 = \frac{m_s^2 - \hat{m}^2}{m_d^2 - m_u^2} \quad \text{NLO in strong interaction O} \\ &= \frac{3}{4} \frac{1}{Q^2} \left[ \frac{\overline{M}_K^2}{\overline{M}_\pi^2} + \frac{\chi_{p^4}}{2} \left( 1 + \frac{m_s}{\hat{m}} \right) \right] \quad Q^2 = \frac{m_s^2 - \hat{m}^2}{m_d^2 - m_u^2} \quad \text{NLO in strong interaction O} \\ &= \frac{3}{4} \frac{1}{Q^2} \left[ \frac{\overline{M}_K^2}{\overline{M}_\pi^2} + \frac{\chi_{p^4}}{2} \left( 1 + \frac{m_s}{\hat{m}} \right) \right] \quad Q^2 = \frac{m_s^2 - \hat{m}^2}{m_d^2 - m_u^2} \quad \text{NLO in strong interaction O} \\ &= \frac{3}{4} \frac{1}{Q^2} \left[ \frac{\overline{M}_K^2}{\overline{M}_\pi^2} + \frac{\chi_{p^4}}{2} \left( 1 + \frac{m_s}{\hat{m}} \right) \right] \quad Q^2 = \frac{m_s^2 - \hat{m}^2}{m_d^2 - m_u^2} \quad \text{NLO in strong interaction O} \\ &= \frac{3}{4} \frac{1}{Q^2} \left[ \frac{\overline{M}_K^2}{\overline{M}_\pi^2} + \frac{\chi_{p^4}}{2} \left( 1 + \frac{m_s}{\hat{m}} \right) \right] \quad Q^2 = \frac{m_s^2 - \hat{m}^2}{m_d^2 - m_u^2} \quad \text{NLO in strong interaction O} \\ &= \frac{3}{4} \frac{1}{Q^2} \left[ \frac{\overline{M}_K^2}{\overline{M}_\pi^2} + \frac{\chi_{p^4}}{2} \left( 1 + \frac{m_s}{\hat{m}} \right) \right] \quad Q^2 = \frac{m_s^2 - \hat{m}^2}{m_d^2 - m_u^2} \quad \text{NLO in strong interaction O} \\ &= \frac{3}{4} \frac{1}{Q^2} \left[ \frac{\overline{M}_K^2}{\overline{M}_\pi^2} + \frac{\chi_{p^4}}{2} \left( 1 + \frac{m_s}{\hat{m}} \right) \right] \quad Q^2 = \frac{m_s^2 - \hat{m}^2}{m_d^2 - m_u^2} \quad Q^2 = \frac{m_s^2 - \hat{m}^2}{M_s^2 - m_u^2} \quad Q^2 = \frac{m_s^2 - m_u^2}{M_s^2 - m_u^2} \quad Q^2 = \frac{m_s^2 - m_u^2}{M_s^2 - m_u^2} \quad Q^2 = \frac{m_s^2 - m_u^2}{M_$$

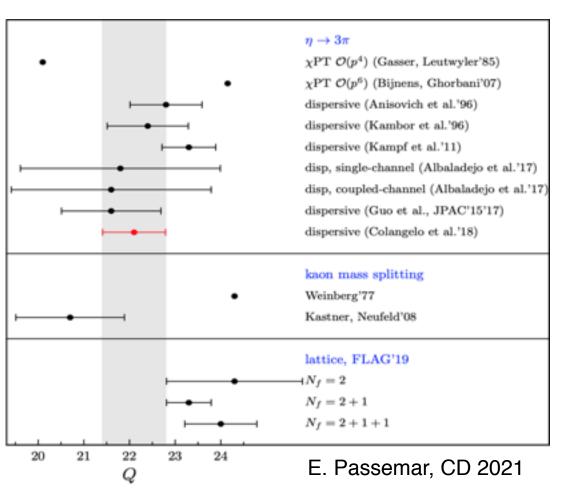
Cirigliano et al., '02; Gasser & Leutwyler, '85

#### = +2.61(17)% Calculated using:

$$Q = 22.1(7)$$
 Colangelo et al. '18, avg. from  $\eta \to 3\pi$   $m_s/\hat{m} = 27.23(10)$  FLAG '20,  $N_f = 2+1+1$  avg.  $M_K = 494.2(3)$  Isospin-limit meson masses from FLAG '17  $M_\pi = 134.8(3)$ 

Test by evaluating  $V_{us}$  from  $K^{\pm}$  and  $K^{0}$  data with **no** corrections: Equality of  $V_{us}$  values would require  $\Delta^{SU(2)} = 2.86(34)\%$ 

Previous to recent results for Q, uncertainty on  $\Delta^{SU(2)}$  was leading contributor to uncertainty on  $V_{us}$  from  $K^{\pm}$  decays



# Reference value of Q from dispersion relation analyses of

 $\eta \to 3\pi$  Dalitz plots

Colangelo et al., '18

$$Q = 22.1 \pm 0.7$$

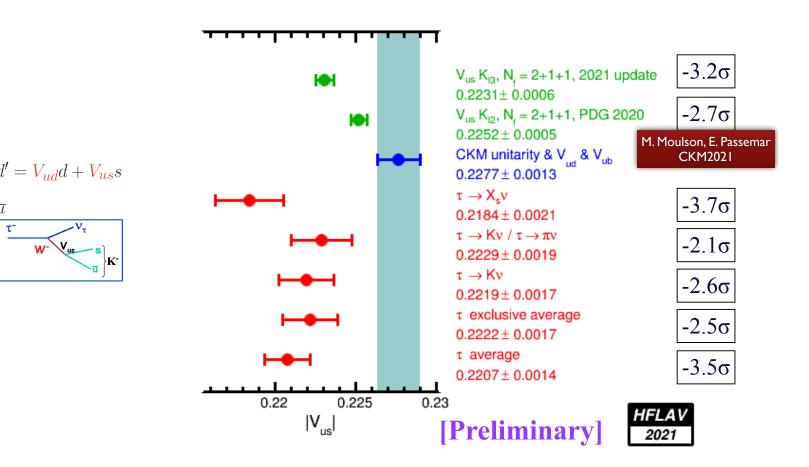
# Lattice results for *Q* somewhat higher than analytical results

But, lattice results have finite correction to LO expectation:

$$Q_M^2 \equiv \frac{\hat{M}_K^2}{\hat{M}_\pi^2} \frac{\hat{M}_K^2 - \hat{M}_\pi^2}{\hat{M}_{K^0}^2 - \hat{M}_{K^+}^2}$$

Low-energy theorem: *Q* has no correction at NLO

# Vus fr Summary of Yays results



- Belle II with 50 ab<sup>-1</sup> and ~4.6 x  $10^{10} \, \tau$  pairs will improve  $V_{us}$  extraction
- Inclusive measurement is an opportunity to have a complete independent measurement of V<sub>us</sub> not easy as you have to measure many channels

#### I vust il otti illetasilleasulleul eseteras

## V<sub>us</sub> from Tau decays

e 13: HFLAV 2021 au branching fractions to strange final states.

Branching fraction	HFLAV 2021 fit (%)
$\mathcal{K}^- u_ au$	$0.6957 \pm 0.0096$
$\mathcal{K}^-\pi^0 u_ au$	$0.4322 \pm 0.0148$
$\mathcal{K}^-2\pi^0 u_ au$ (ex. $\mathcal{K}^0$ )	$0.0634 \pm 0.0219$
$K^-3\pi^0 u_ au$ (ex. $K^0$ , $\eta$ )	$0.0465 \pm 0.0213$
$\pi^-\overline{K}^0 u_ au$	$0.8375 \pm 0.0139$
$\pi^-\overline{K}^0\pi^0 u_ au$	$0.3810 \pm 0.0129$
$\pi^-\overline{K}^0 2\pi^0 u_ au$ (ex. $K^0$ )	$0.0234 \pm 0.0231$
$\overline{\textit{K}}^{0}\textit{h}^{-}\textit{h}^{-}\textit{h}^{+}\nu_{ au}$	$0.0222 \pm 0.0202$
$\mathcal{K}^-\eta u_ au$	$0.0155 \pm 0.0008$
$\mathcal{K}^-\pi^0\eta u_ au$	$0.0048 \pm 0.0012$
$\pi^-\overline{K}^{0}\eta u_{ au}$	$0.0094 \pm 0.0015$
${\sf K}^-\omega u_ au$	$0.0410 \pm 0.0092$
$\mathcal{K}^-\phi(\mathcal{K}^+\mathcal{K}^-) u_ au$	$0.0022 \pm 0.0008$
$\mathcal{K}^-\phi(\mathcal{K}^0_\mathcal{S}\mathcal{K}^0_\mathcal{L}) u_ au$	$0.0015 \pm 0.0006$
$\mathcal{K}^-\pi^-\pi^+ u_ au$ (ex. $\mathcal{K}^{ extsf{0}}$ , $\omega$ )	$0.2924 \pm 0.0068$
$\mathcal{K}^-\pi^-\pi^+\pi^0 u_ au$ (ex. $\mathcal{K}^0$ , $\omega$ , $\eta$ )	$0.0387 \pm 0.0142$
$\mathcal{K}^-2\pi^-2\pi^+ u_ au$ (ex. $\mathcal{K}^0$ )	$0.0001 \pm 0.0001$
$K^- 2\pi^- 2\pi^+ \pi^0  u_{ au}  ext{ (ex.} K^0)$	$0.0001 \pm 0.0001$
$X_s^- u_ au$	$2.9076 \pm 0.0478$

$$R_{\tau} = \frac{\Gamma(\tau^{-} \to v_{\tau} + \text{hadrons})}{\Gamma(\tau^{-} \to v_{\tau}e^{-}\overline{v_{e}})} \approx N_{C}$$
  $\delta R_{\text{theory}}$ 

where  $\delta R_{\rm theory}$  can be defined as scattering data. The literature reports estimate  $\delta R_{\tau} \equiv \frac{R_{\tau,NS}}{|V_{ud}|^2} - \frac{R_{\tau,S}}{|V_{us}|^2}$  ize is in between the value of  $m_s = 95.00 \pm 100$ 

We proceed following the same procedu (17.81543) brogsting (Heartise strongs) to constraint and non-straint hadronic final st

Using the  $\tau$  branching fraction fit results (2.909 ± 0.048)% (see also Table 13) and  $= \frac{R_{\tau,s}}{R}$  branching fraction  $\frac{R_{\tau,s}}{R}$  and  $\frac{R_{\tau,s}}{R}$  branching fraction  $\frac{R_{\tau,s}}{R}$  and  $\frac{R_{\tau,s}}{R}$  branching fraction  $\frac{R_{\tau,s}}{R}$  and  $\frac{R_{\tau,s}}{R}$  branching fraction  $\frac{R_{\tau,s}}{R}$  branching fraction  $\frac{R_{\tau,s}}{R}$  and  $\frac{R_{\tau,s}}{R}$  branching fraction  $\frac{R_{\tau,s}}{R}$  and  $\frac{R_{\tau,s}}{R}$  branching fraction  $\frac{R_{\tau,s}}{R}$  and  $\frac{R_{\tau,s}}{R}$  branching fraction  $\frac{R_{\tau,s}}{R}$  b

$$\frac{1}{\left|V_{ud}\right|^{2}} - \delta R_{\tau,th}$$

$$\frac{1}{\left|V_{ud}\right|^{2}} - \delta R_{\tau,th}$$

$$\frac{1}{\left|V_{ud}\right|^{2}} - \delta R_{\tau,th}$$

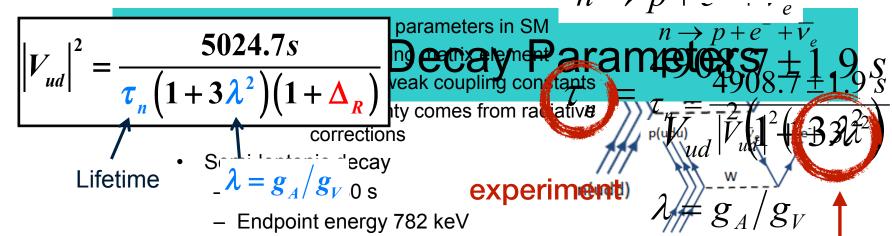
$$\frac{1}{\left|V_{us}\right|_{\text{uni}}} = 1 - \left|V_{ud}\right|^{2}.$$

 $|V_{us}| = 0.2184 \pm 0.0018_{\text{exp}} \pm 0.0011_{\text{th}} \delta R_{\text{theory}}$ 

## V<sub>ud</sub> | from Neutrosemi-leptonic decay

- Lifetime ~880 s
- Endpoint energy 782 keV

Master Formula:



- Needs δλ/λ ≈ 3 × 10<sup>-1</sup>/<sub>20</sub> and δτ<sub>β1</sub>ε 0.3 s to compele with p + 0<sup>44</sup> transitions.
- Theoretically, the radiative corrections are under control (same as for  $0^+$  0+ QC CKM mixing matrix element 1000 7 100 QC
- - New Perkeo III\_result: PERKEO III result into your world-average of beta asymmetry by factor Stibles of it is due to the reduction of the scale payton

$$A = -0.11958(21), S = 1.2 \quad \lambda_A = -1.2757(5)$$

– Tension with aSPECT result:

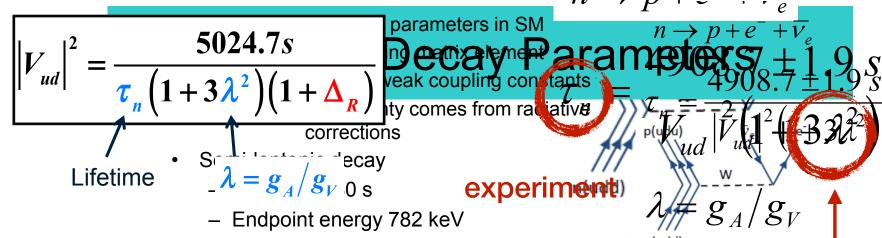
$$\lambda = \frac{\sigma}{\sigma} / \frac{\sigma}{\sigma}$$
 $\lambda_{\text{avg}} = -1.2754(13), S = 2.7$ 

n(udd)

### | V<sub>ud</sub> | from Neutrossemi-leptonic decay

- Lifetime ~880 s
- Endpoint energy 782 keV

Master Formula:



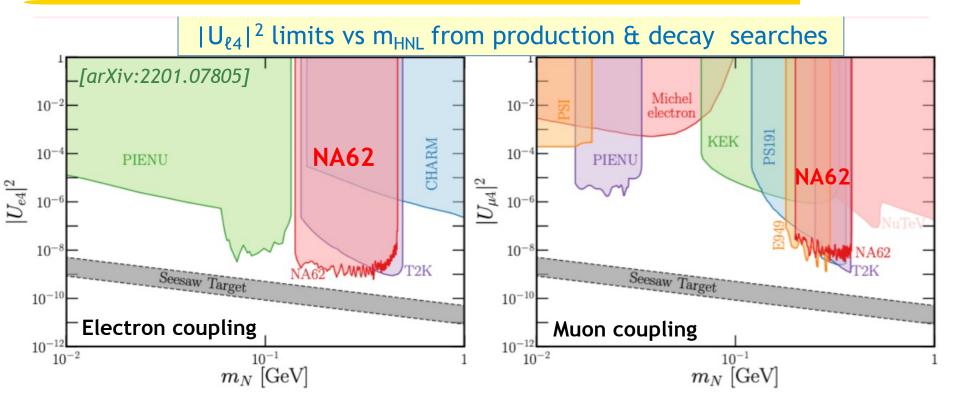
- Needs δλ/λ ≈ 3 × 10<sup>-4</sup> and δτ<sub>β1</sub>ε 0.3 s to compete with p + 0<sup>44</sup> transitions.
- Theoretically, the radiative corrections are under control (same as for 0 = 20)
- Recent progress : 

   Ratio of weak coupling constants
  - New Perkeo III\_result: PERKEO III result inforoves world-average of beta asymmetry by factor of it is due to the reduction of the scale lagger

$$A = -0.11958(21), S = 1.2 \quad \lambda_A = -1.2757(5)$$

- New result for Lifetime from UCNT  $au_n=877.75\pm0.28^{+0.22}_{-0.16}$  s
  - improverse of 2.25 compared to previous resultant

## 3.3 Example: Constraints on Heavy Neutral Leptons



- Strongest |Ue4|2 limits below 400 MeV: K+,  $\pi^+ \rightarrow e^+N$  from NA62 & PIENU.
- Also important limits on |U<sub>u4</sub>|<sup>2</sup> from E949, NA62 and PIENU.
- NA62/E949 limits are complementary to HNL decay searches at T2K.
- Next-generation K<sup>+</sup> and p+ experiments (NA62<sup>++</sup>, PIONEER) to improve by up to factor 10, reaching the seesaw bound.

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