

# Introduction and motivation for Effective Field Theories in high energy physics Experimental view

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# Outline

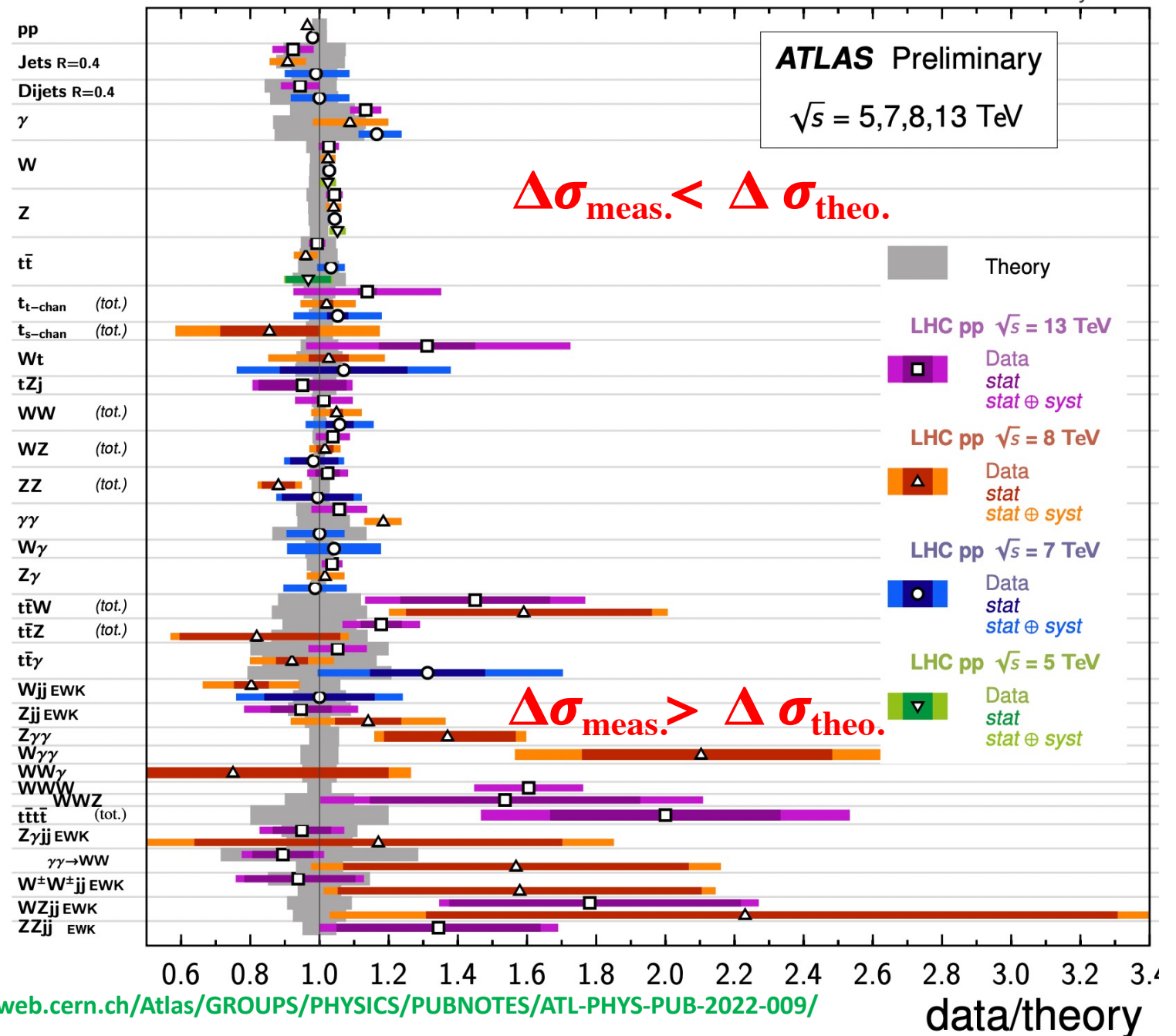
- **Quest for Beyond Standard Model: where do we stand**  
(gauge bosons, top, Higgs sector and new particles)
- **Why we have not (yet) found New Physics** (“anomalies” exist)
- **How can we use LHC data in the best way**
- **What is an Effective Field Theory (EFT)**
- **Example of EFT: Fermi theory of weak interactions**  
**Multipole expansion**  
**(SM) EFT fit @ LHC**  
(important to know well the SM)
- **Summary**

# Quest for Beyond Standard Model: where do we stand ?

## Gauge bosons and top quark

- Comparison data/theory over many orders of magnitudes ( $\sim 10^{11}\text{pb} \rightarrow \sim 10^{-3}\text{pb}$ )

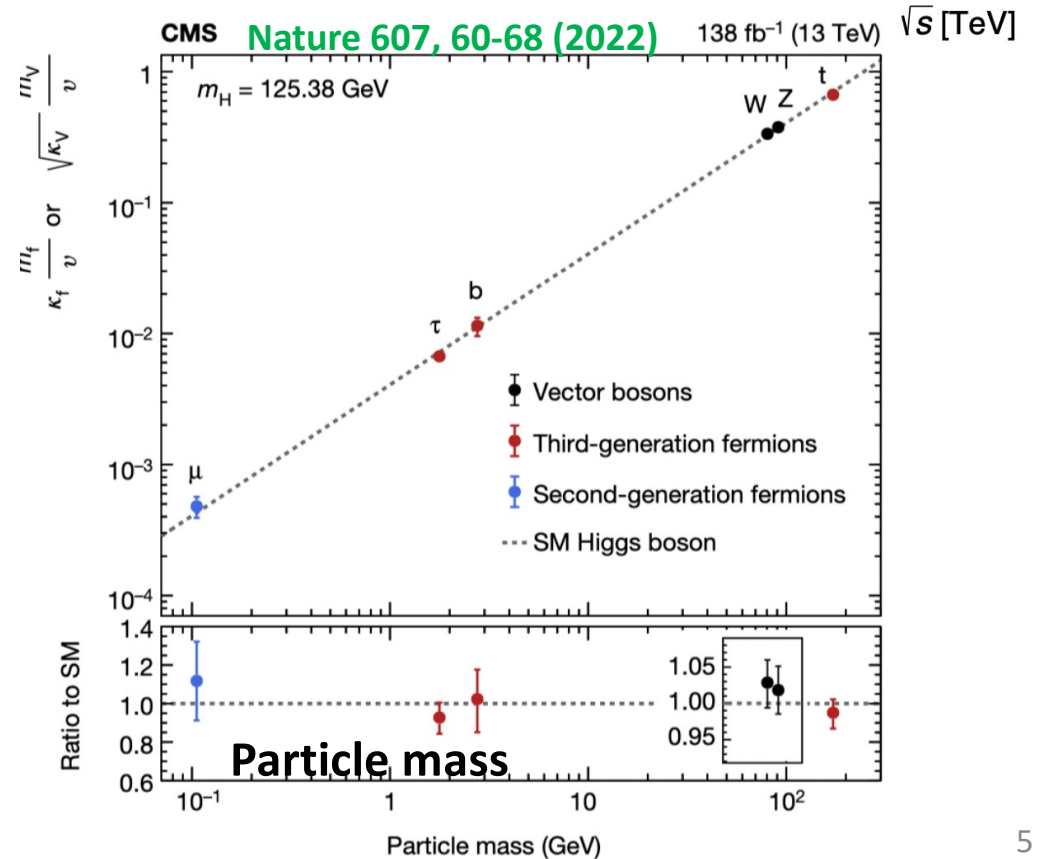
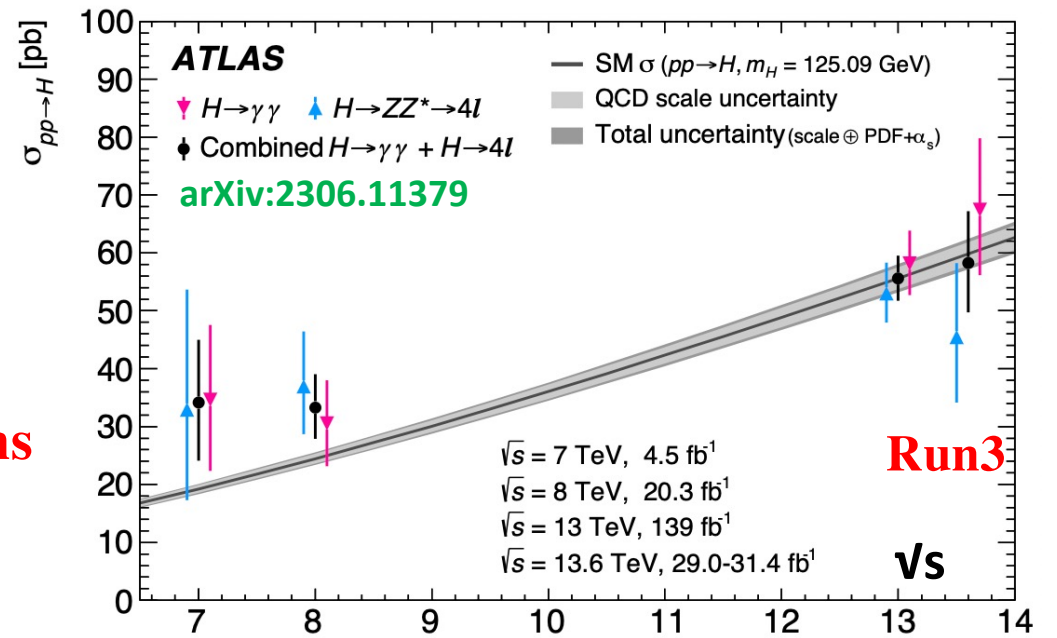
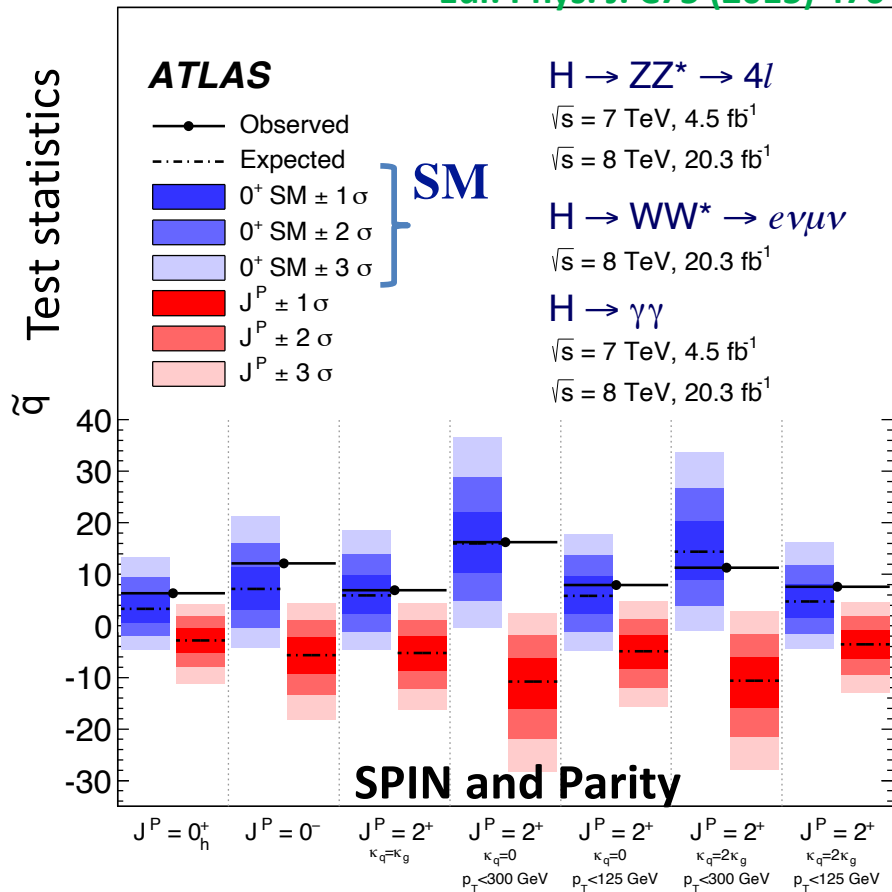
- No significant deviations so far



# Higgs sector

- No significant deviations
- Spin and Parity: Results consistent with SM predictions (SM Higgs  $J^{PC} = 0^{++}$ )

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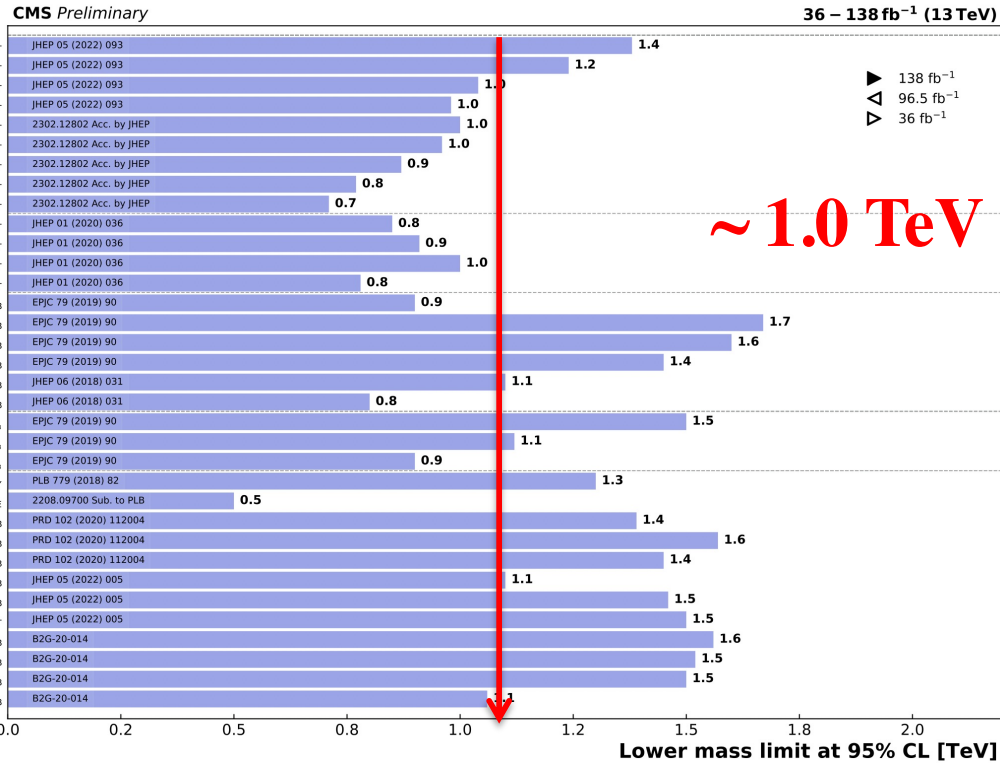
# New particles

- Many direct searches (too many to display in one page)
- No new particles
- Set lower mass limit @ 95% CL up to  $\approx 6.6$  TeV ( $Z' \rightarrow tt$ )

## Very heavy fermions

- Pair prod.**
- (qb)T
    - $b Z (Z \rightarrow \nu\nu)$  ( $\Gamma/m=0.3$ , Singlet)  $M_T$
    - $b Z (Z \rightarrow \nu\nu)$  ( $\Gamma/m=0.2$ , Singlet)  $M_T$
    - $b Z (Z \rightarrow \nu\nu)$  ( $\Gamma/m=0.1$ , Singlet)  $M_T$
    - $b Z (Z \rightarrow \nu\nu)$  ( $\Gamma/m=0.05$ , Singlet)  $M_T$
    - $b tH (H \rightarrow \gamma\gamma)$  ( $\Gamma/m=0.05$ , Singlet)  $M_T$
    - $b tH (H \rightarrow \gamma\gamma)$  ( $\Gamma/m=0.04$ , Singlet)  $M_T$
    - $b tH (H \rightarrow \gamma\gamma)$  ( $\Gamma/m=0.03$ , Singlet)  $M_T$
    - $b tH (H \rightarrow \gamma\gamma)$  ( $\Gamma/m=0.02$ , Singlet)  $M_T$
    - $b tH (H \rightarrow \gamma\gamma)$  ( $\Gamma/m=0.01$ , Singlet)  $M_T$
  - (qt)T
    - $t H (H \rightarrow b\bar{b})$  ( $\Gamma/m=0.3$ , Doublet)  $M_T$
    - $t H (H \rightarrow b\bar{b})$  ( $\Gamma/m=0.3$ , Singlet)  $M_T$
    - $t H (H \rightarrow b\bar{b})$  ( $\Gamma/m=0.1$ , Singlet)  $M_T$
    - $t H (H \rightarrow b\bar{b})$  ( $\Gamma/m=0.05$ , Singlet)  $M_T$
  - (qt)(qb)B
    - $b W t \rightarrow lep. + jets$  ( $\Gamma/m=0.1$ , LH)  $M_B$
    - $b W t \rightarrow lep. + jets$  ( $\Gamma/m=0.3$ , LH)  $M_B$
    - $b W t \rightarrow lep. + jets$  ( $\Gamma/m=0.2$ , LH)  $M_B$
    - $b W t \rightarrow lep. + jets$  ( $\Gamma/m=0.1$ , LH)  $M_B$
  - (qt)X
    - $b H b (H \rightarrow b\bar{b})$  ( $\Gamma/m=0.3$ , Doublet)  $M_B$
    - $b H b (H \rightarrow b\bar{b})$  ( $\Gamma/m=0.2$ , Doublet)  $M_B$
    - $t W t \rightarrow lep. + jets$  ( $\Gamma/m=0.3$ , LH)  $M_{X,50}$
    - $t W t \rightarrow lep. + jets$  ( $\Gamma/m=0.2$ , LH)  $M_{X,50}$
    - $t W t \rightarrow lep. + jets$  ( $\Gamma/m=0.1$ , LH)  $M_{X,50}$
- Other**
- $Y_{-4/3} Y_{-4/3} \rightarrow bW bW \rightarrow l\nu q\bar{q}q\bar{q}$   $M_Y$
  - $EE \rightarrow 4b \tau\tau/\nu\nu$   $M_E$
  - $BB \rightarrow bq\bar{q} bq\bar{q}$  ( $B(tZ) = 1$ )  $M_B$
  - $BB \rightarrow bq\bar{q} bq\bar{q}$  ( $B(tH) = 1$ )  $M_B$
  - $BB \rightarrow bq\bar{q} bq\bar{q}$  (Singlet)  $M_B$
  - $BB \rightarrow lep. + jets$  (Doublet)  $M_B$
  - $BB \rightarrow lep. + jets$  (Singlet)  $M_B$
  - $TT \rightarrow lep. + jets$  (Singlet and Doublet)  $M_T$
  - $BB \rightarrow lep. + jets$  ( $B(tH) = 1$ )  $M_B$
  - $BB \rightarrow lep. + jets$  ( $B(tZ) = 1$ )  $M_B$
  - $BB \rightarrow lep. + jets$  (Doublet)  $M_B$
  - $BB \rightarrow lep. + jets$  (Singlet)  $M_B$

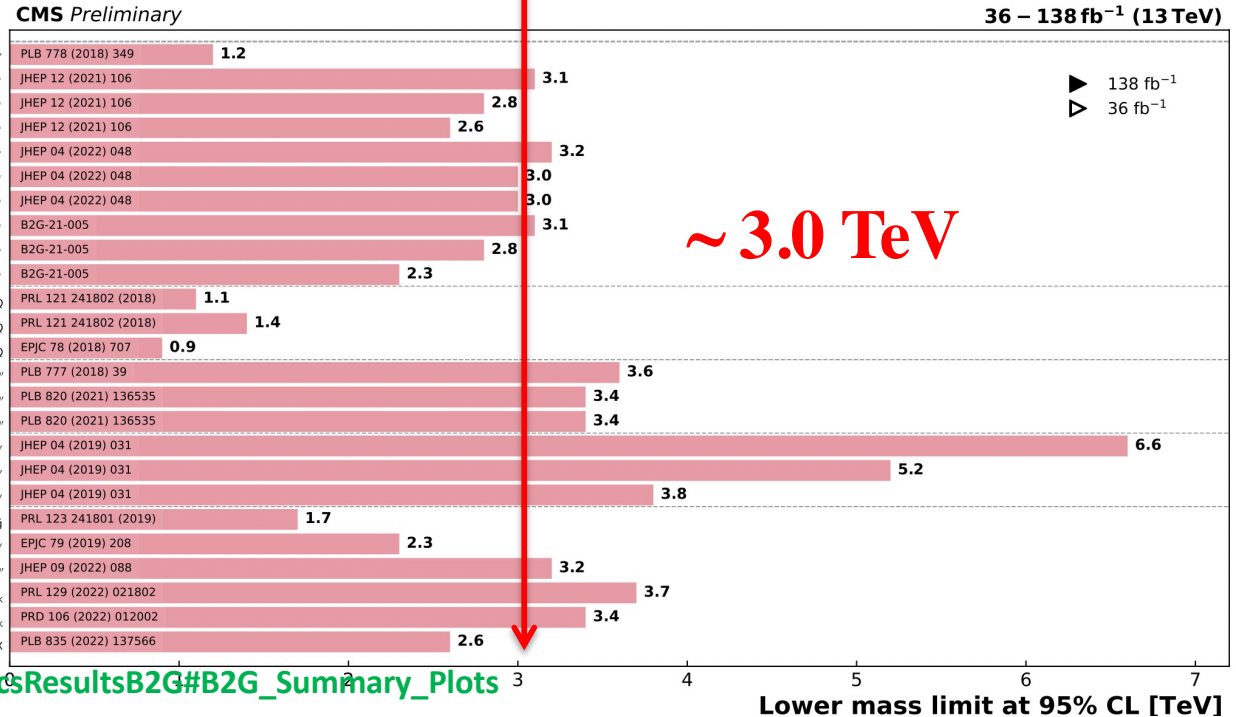
## Overview of CMS B2G Results



$\sim 1.0$  TeV

## Resonances

- Excited quarks**
- $t^* \bar{t}^* \rightarrow l\nu b\bar{b} + jets$  (R-S model,  $B = 1$ )
  - $b^* \rightarrow tW \rightarrow bq\bar{q} q\bar{q}$  (LH+RH)
  - $b^* \rightarrow tW \rightarrow bq\bar{q} q\bar{q}$  (RH)
  - $b^* \rightarrow tW \rightarrow bq\bar{q} q\bar{q}$  (LH)
  - $b^* \rightarrow tW \rightarrow bq\bar{q} l\nu$  (LH+RH)
  - $b^* \rightarrow tW \rightarrow bq\bar{q} l\nu$  (RH)
  - $b^* \rightarrow tW \rightarrow bq\bar{q} l\nu$  (LH)
  - $b^* \rightarrow tW \rightarrow bl\nu q\bar{q}$  (LH+RH)
  - $b^* \rightarrow tW \rightarrow bl\nu q\bar{q}$  (RH)
  - $b^* \rightarrow tW \rightarrow bl\nu q\bar{q}$  (LH)
- LQ**
- $LQ\bar{LQ} \rightarrow b\nu b\nu$
  - $LQ\bar{LQ} \rightarrow t\mu t\mu$
  - $LQ\bar{LQ} \rightarrow t\tau t\tau$
- $W' \rightarrow tb$**
- $W' \rightarrow tb, 1l$  (RH)  $M_{W'} > M_W$
  - $W' \rightarrow tb, 0l$ , (LH)
  - $W' \rightarrow tb, 0l$ , (RH)
- $Z' \rightarrow tt$**
- $Z' \rightarrow t\bar{t}$  ( $\Gamma/M_{Z'} = 30\%$ )
  - $Z' \rightarrow t\bar{t}$  ( $\Gamma/M_{Z'} = 10\%$ )
  - $Z' \rightarrow t\bar{t}$  ( $\Gamma/M_{Z'} = 1\%$ )
- Other**
- Stealth  $\tilde{g} \rightarrow \tilde{\chi}_1^0 q\bar{q}$  ( $\gamma + 2jets$ ,  $M_{\tilde{\chi}_1^0} = 0.2\text{TeV}$ )
  - $Z' \rightarrow t\bar{t} \rightarrow tZt/tHt \rightarrow l\nu + jets$  ( $M_{Z'} = 1.5$  TeV)
  - $W' \rightarrow Tb/Bt$  ( $M_{W'} = 2/3 M_W$ )
  - $W_{KK} \rightarrow RW \rightarrow WWW$  ( $0l + 1l$ )
  - $W_{KK} \rightarrow RW \rightarrow WWW$  ( $1l$ )
  - $X \rightarrow aa \rightarrow b\bar{b}b\bar{b}$  ( $M_a = 0.1$  TeV,  $M_X/N_f = 8$ )



$\sim 3.0$  TeV

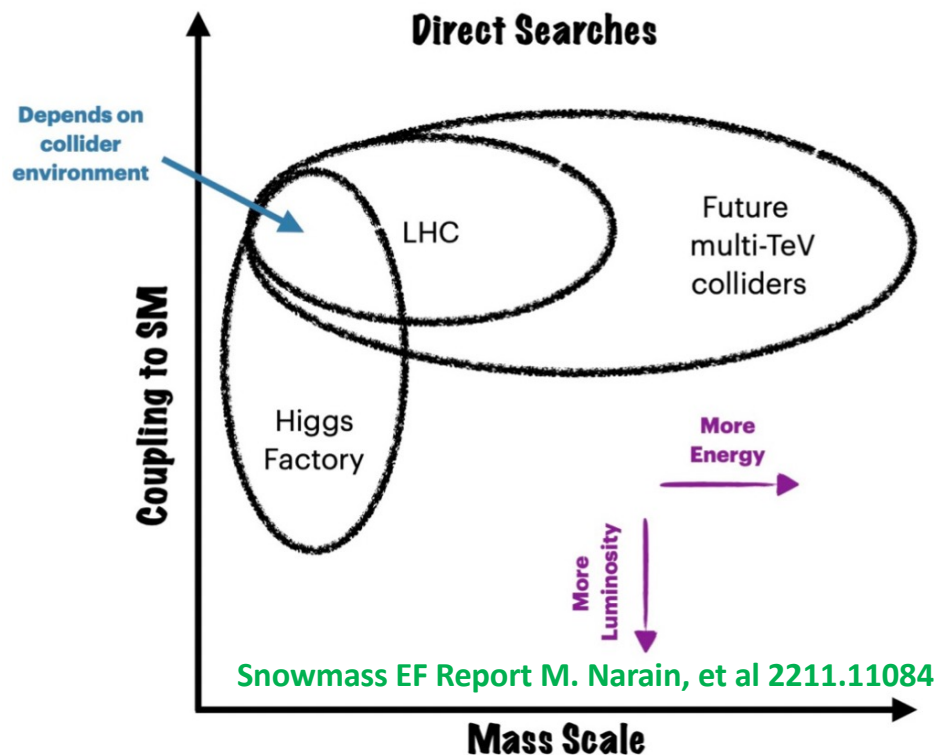


# Why we have not (yet) found new physics (wrt SM) (\*)?

(\*) Few anomalies exist

1. New physics is buried in the backgrounds
2. New physics is weakly coupled

Improve: trigger (TLA) analysis methods (ML), **theory prediction of bkg** and/or more luminosity



- Direct searches of new states:

$$\sqrt{s'} > Mx$$

Need a machine with enough energy

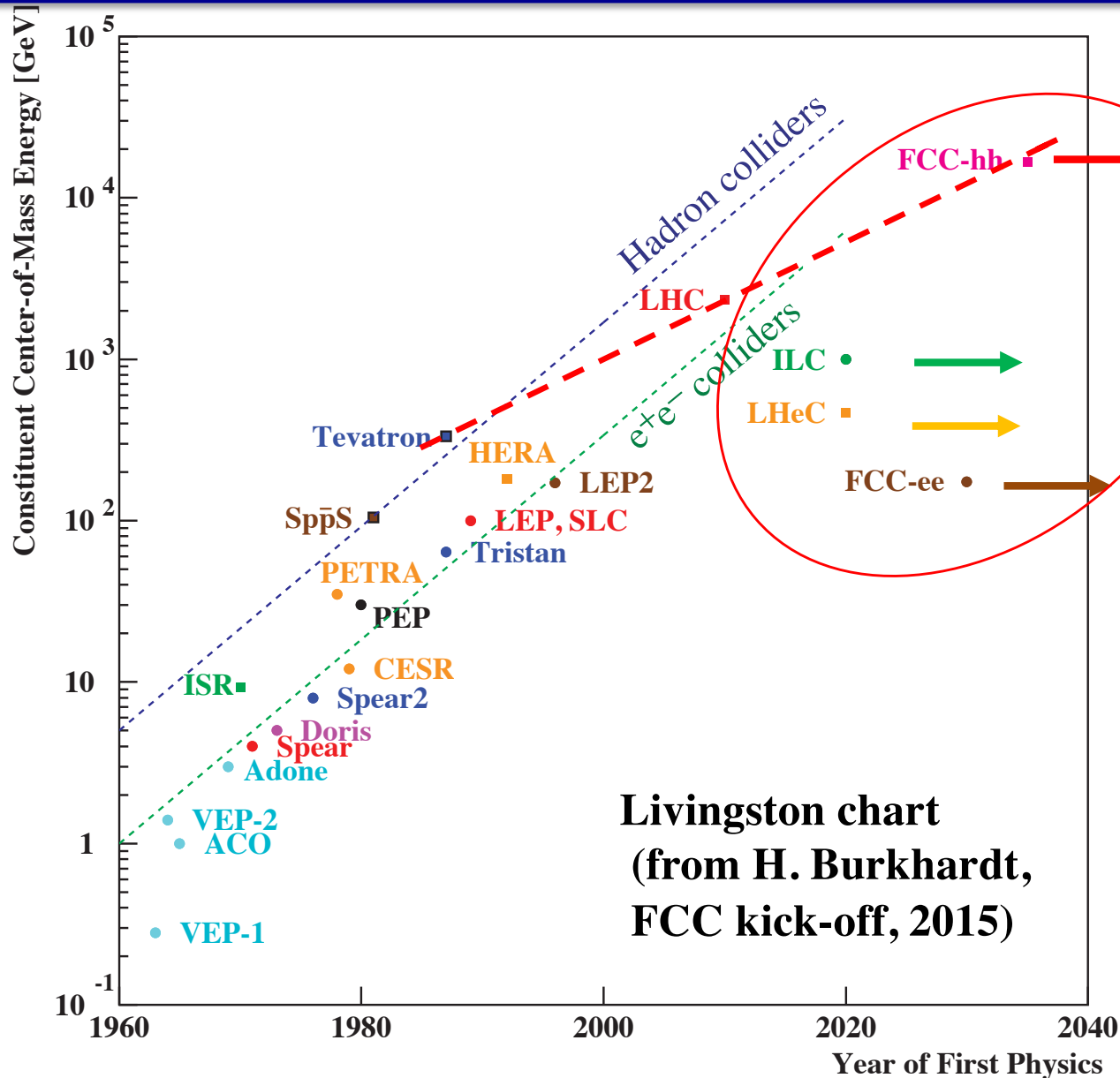
- \* Search for a peak (often) in the invariant mass of the decay products or in the production cross section (ex. LEP)

- \* In most cases, direct searches motivated by a 'descriptive' theory (ex. SUSY)

3. A gap exists between the reached energy and new physics

# A gap between the reached energy and new physics

→ Increase the energy in the center of mass ( $\sqrt{s'} > M_X$ ) ?



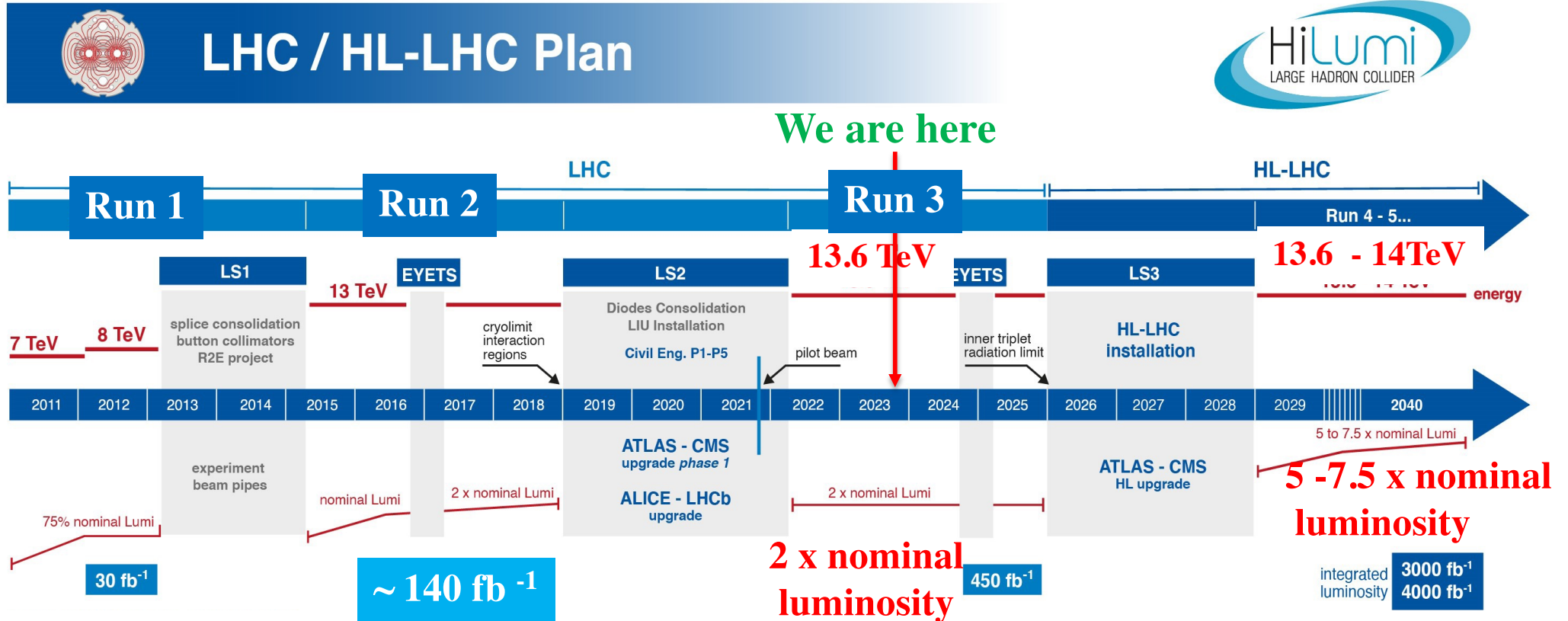
Future accelerators

- Accelerator Physics offers the unique opportunity to explore many fundamental questions at high intensity and energy with known, tunable initial conditions
- Change in slope: (it takes time and money)



# (HL-)LHC: a powerful tool already NOW

## How can we use it in the best way ?



- Energy ~ constant (since 2015)
- Significant luminosity increase (@HL-LHC ~ 15 times the present data sample)

# Complementary search scenarios wrt direct searches

- Search for new physics by measuring **very precisely** key SM observables (EWPO)

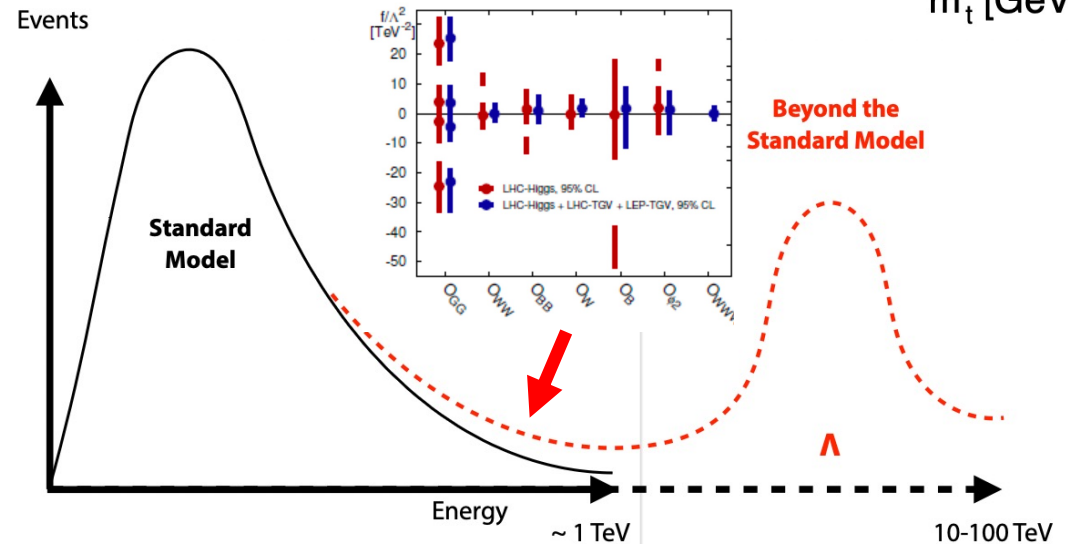
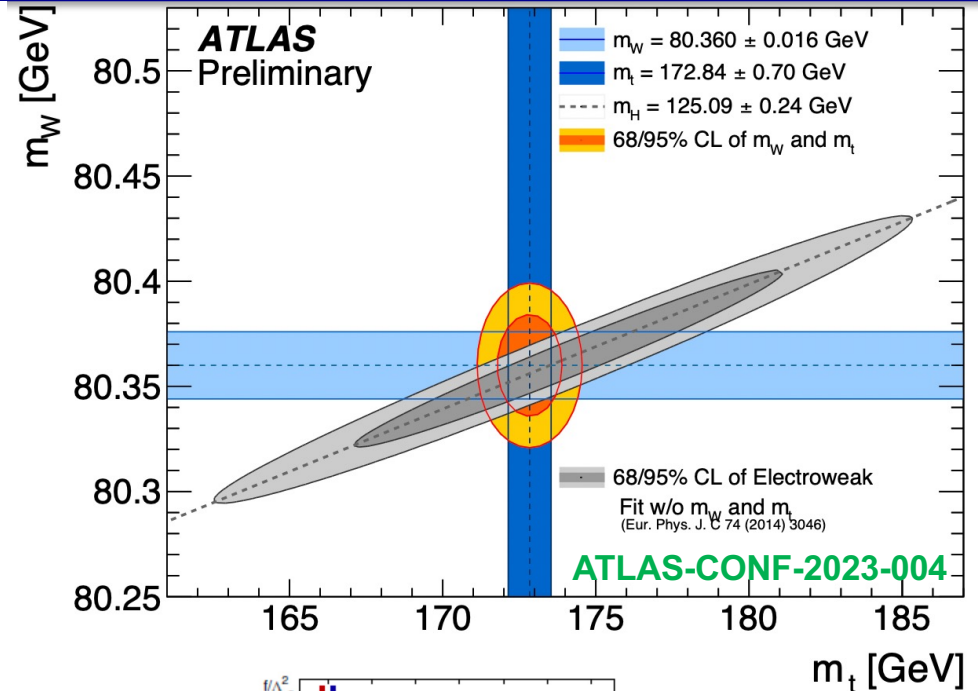
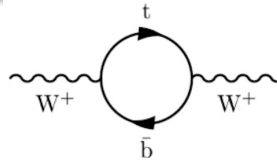
Expansion in coupling strength

Any significant inconsistency wrt SM

→ existence of New Physics
  - Search for deviations from SM in high energy regions (often tails)

Use an **Effective Field Theory** to interpret the data

Expansion in scale ratios
- In both cases probe energy scales beyond the direct kinematic reach



# What is an EFT ?

- It's a quantum field theory describing the behaviour of **an underlying unknown physical theory** (“UV complete”) in some limited regime (“IR limit”).
- In an EFT, the details of (new) Physics are irrelevant **@  $E \ll \Lambda$**   
→ EFT use degrees of freedom relevant to the IR regime (simplification)

**$\Lambda$**  is a scale (**energy or inverse of a distance**):

- **@  $E \ll \Lambda$**  (new) Physics effects are described by local, analytic operators with  **$1/\Lambda^n$**  suppressions
- possible divergences **@  $E \sim \Lambda$**  and above, can be neglected

In practice: Taylor expansion of the Lagrangian in powers of  **$E/\Lambda$**

$$\mathcal{L}_{\text{EFT}} = \sum_{\mathcal{D} > 0, i} \frac{c_i^{(\mathcal{D})} O_i^{(\mathcal{D})}}{\Lambda^{\mathcal{D}-d}}$$

$d$  = space-time dimensions

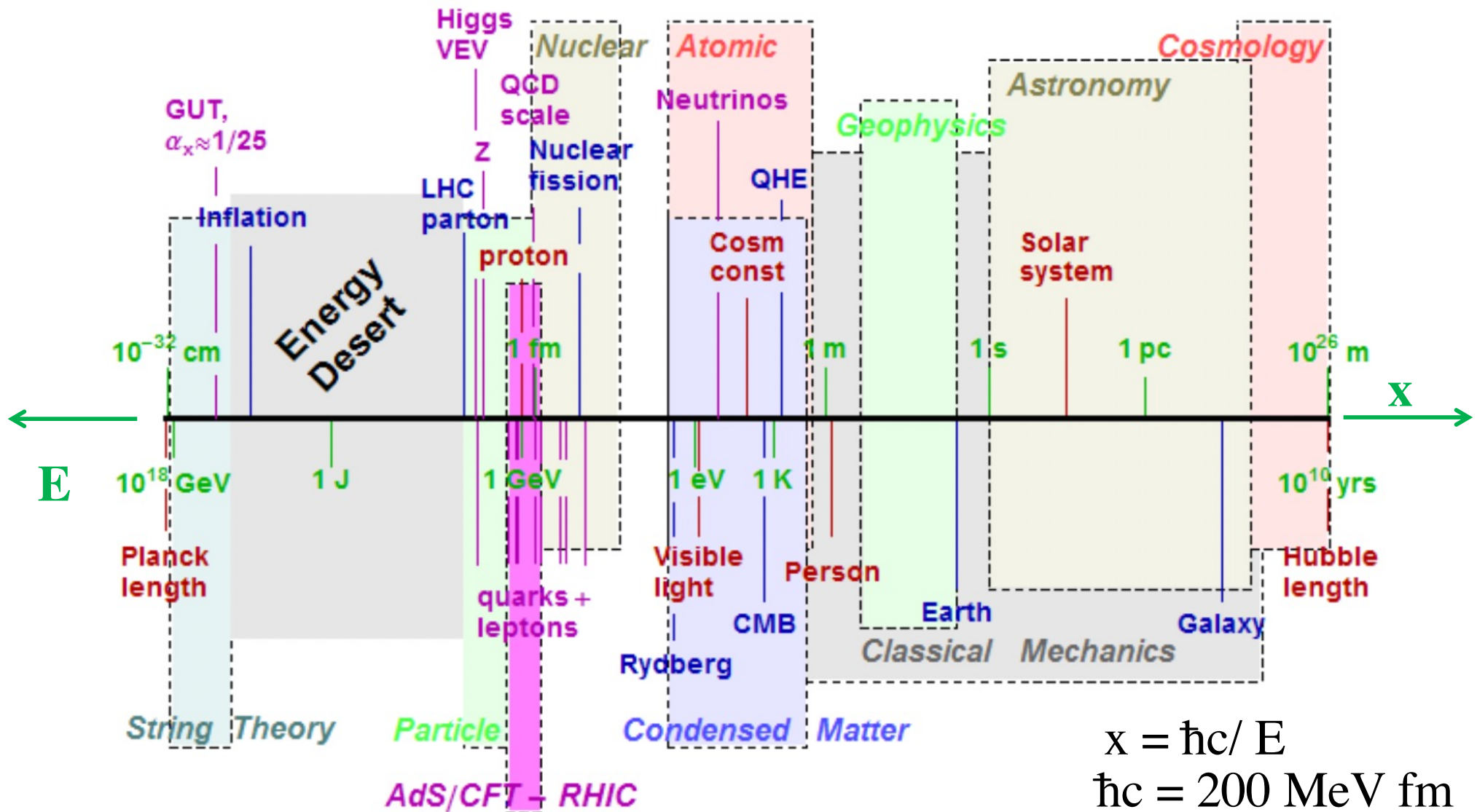
$c_i$  = Wilson coefficients

$O_i$  = allowed operators

$\mathcal{D}$  = operator dimension

- New physics expressed in terms of coefficients of higher dimension operators including constraints of locality, gauge and Lorentz invariance

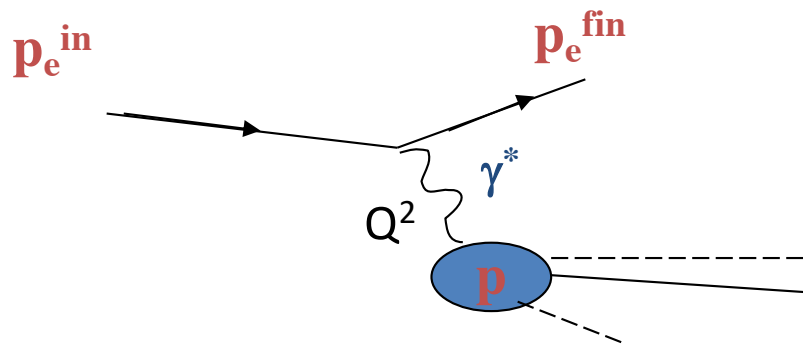
# Nature has naturally many scales (“divide et impera”)



<https://phy.princeton.edu/research/high-energy-theory/gubser-group/outreach/energy-scales-in-physics>

# Nature has naturally many scales (“divide et impera”)

## Electron scattering on a proton

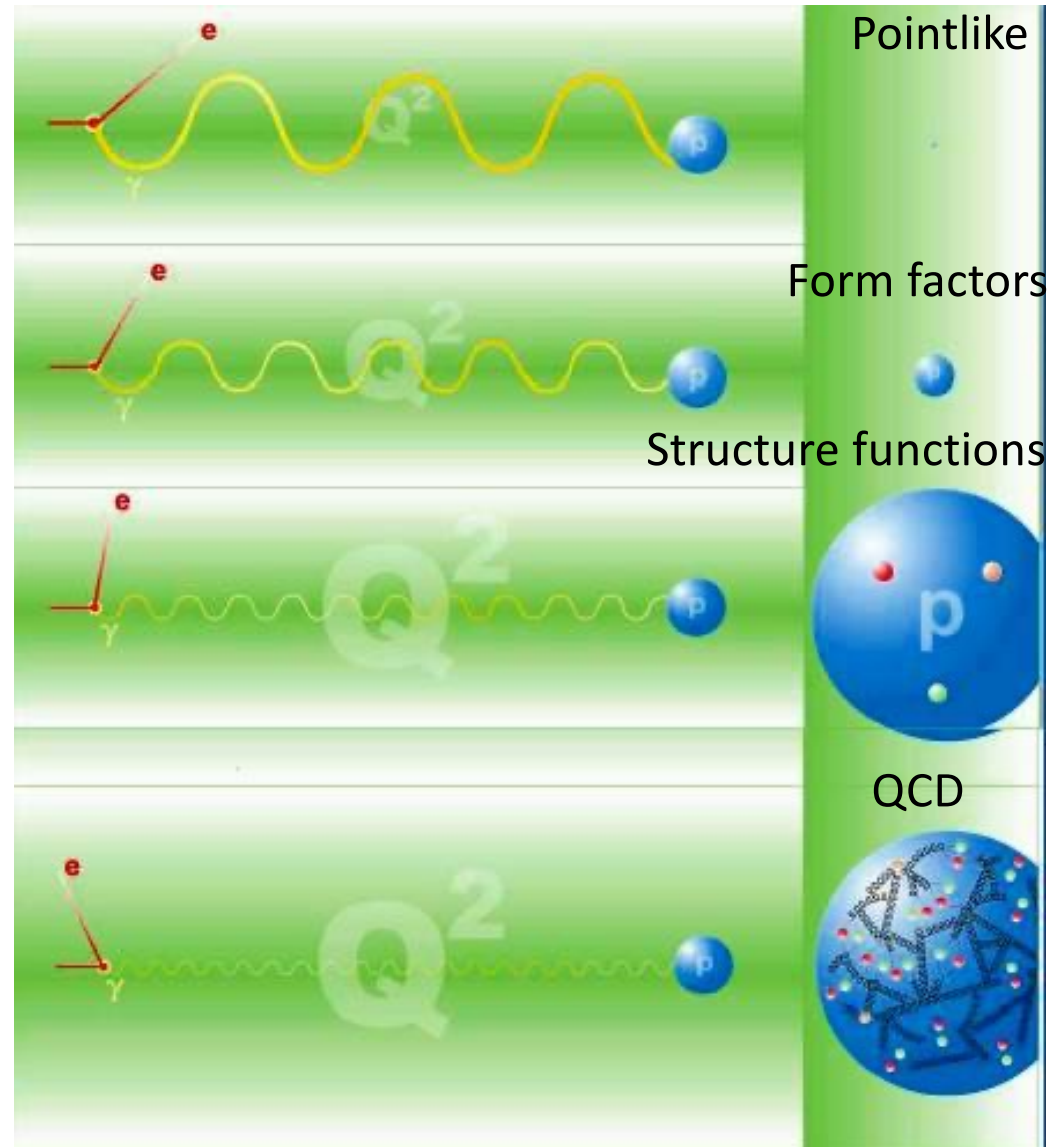


$$Q^2 = (p_e^{in} - p_e^{fin})^2$$

@ Higher and higher  $Q^2$   
 (= smaller and smaller  $\lambda$ )  
 new physics regime

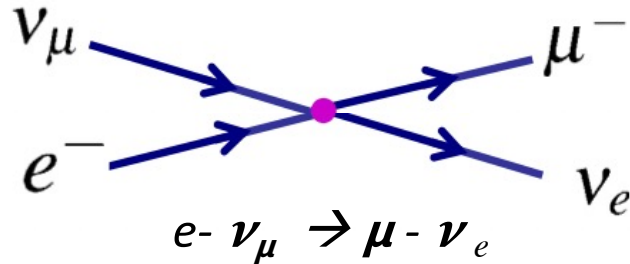
$$\Lambda \sim \sqrt{Q^2} \quad \text{or} \quad \Lambda \sim 1/\lambda$$

$Q^2$





# Successful example of EFT: Fermi theory of weak interactions



Current-current ‘universal’\* interaction

‘Contact interaction’

\* Same coupling constant for all weak interactions

$$M_{fi} = \frac{G_F}{\sqrt{2}} g_{\mu\nu} [\bar{\psi} \gamma^\mu (1 - \gamma^5) \psi] [\bar{\psi} \gamma^\nu (1 - \gamma^5) \psi]$$

- Neglecting the lepton masses ( $m_e \sim m_\mu \ll E_\nu$ ):

$$\sigma_{\nu_\mu e^-} = \frac{G_F^2 s}{\pi}$$

$s = \text{energy}^2$  in the center of mass =  $2 m_e E_\nu$

$G_F =$  Fermi constant measured in the  $\mu$  decay

(@ a scale  $\Lambda \sim m_\mu$ ) =  $1.16 * 10^{-5} \text{ GeV}^{-2}$

- The cross section cannot be higher than the value corresponding to the max of the scattering probability ( ‘unitarity violation’)  $\sigma_{max}^{elastic}$

$$\sigma_{max}^{elastic} = (2l + 1) 16 \pi / s \quad \text{From partial wave analysis neglecting spins}$$

$l = \text{angular momentum}$

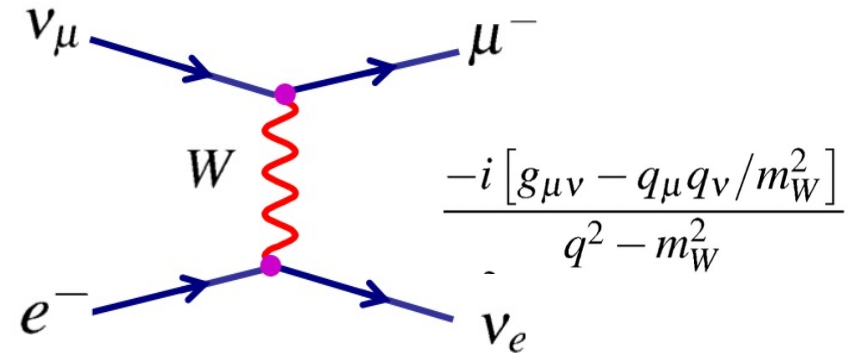
- The Fermi cross section violates unitarity  $\sqrt{s} \approx \sqrt{G_F} \approx 300 \text{ GeV}$



# Successful example of EFT: Fermi theory of weak interactions

$$e - \nu_\mu \rightarrow \mu - \nu_e$$

- The divergent behaviour is 'cured' by introducing an 'intermediate' massive boson of Spin 1 as propagator



$$M_{fi} = \left[ \frac{g_W}{\sqrt{2}} \bar{\psi} \frac{1}{2} \gamma^\mu (1 - \gamma^5) \psi \right] \frac{g_{\mu\nu} - q_\mu q_\nu / m_W^2}{q^2 - m_W^2} \left[ \frac{g_W}{\sqrt{2}} \bar{\psi} \frac{1}{2} \gamma^\nu (1 - \gamma^5) \psi \right]$$

$q^2 = (p_e^{in} - p_{\nu_e}^{fin})^2$

For  $q^2 \ll m_W^2$ : one obtains back the contact interaction ( $\Delta \sim m_W$ )

Expansion in  $\delta = q/M_W$

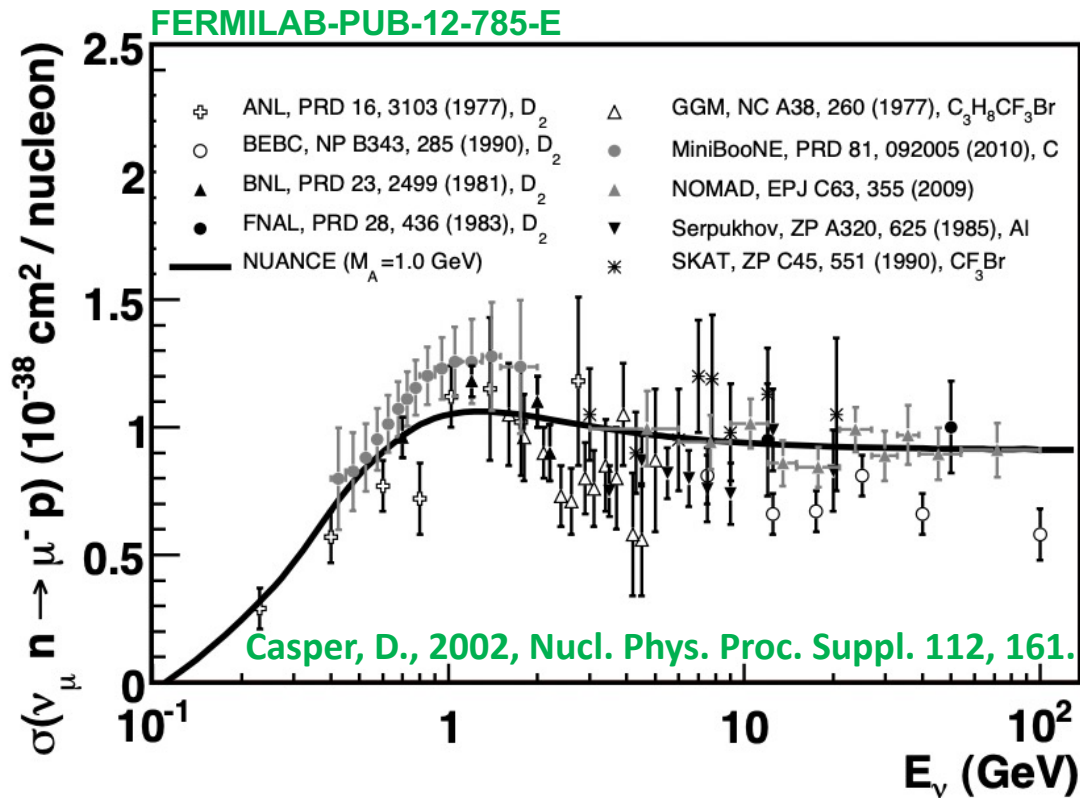
$$\frac{1}{q^2 - M_W^2} = -\frac{1}{M_W^2} \left( 1 + \frac{q^2}{M_W^2} + \frac{q^4}{M_W^4} + \dots \right)$$

$$M_{fi} = \frac{i}{M_W^2} \left( \frac{-ig_W}{2\sqrt{2}} \right)^2 (\bar{\nu}_\mu \gamma^\mu P_L \mu) (\bar{e} \gamma^\mu P_L \nu_e, \nu_e) + \mathcal{O} \left( \frac{1}{M_W^4} \right) \quad \frac{G_F}{\sqrt{2}} \equiv \frac{g_W^2}{8M_W^2}$$

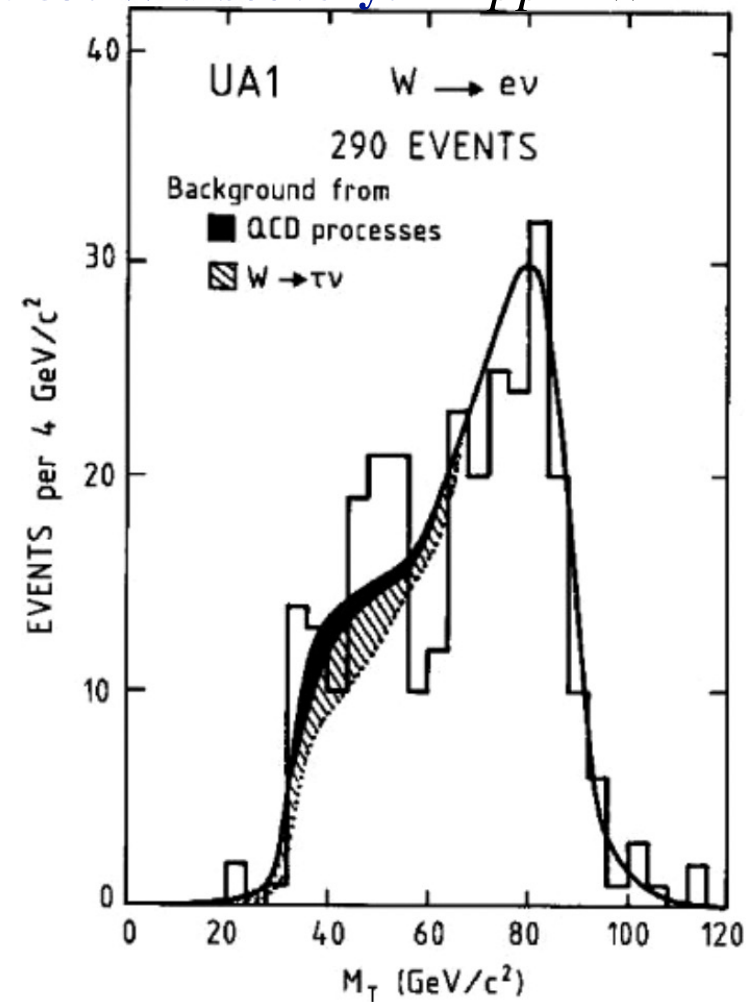
# Great success

Intermediate Energy Cross Sections:  
 $E_\nu \sim 0.1 - 100 \text{ GeV}$

Quasi-elastic scattering:  $\nu_\mu n \rightarrow \mu^- p$

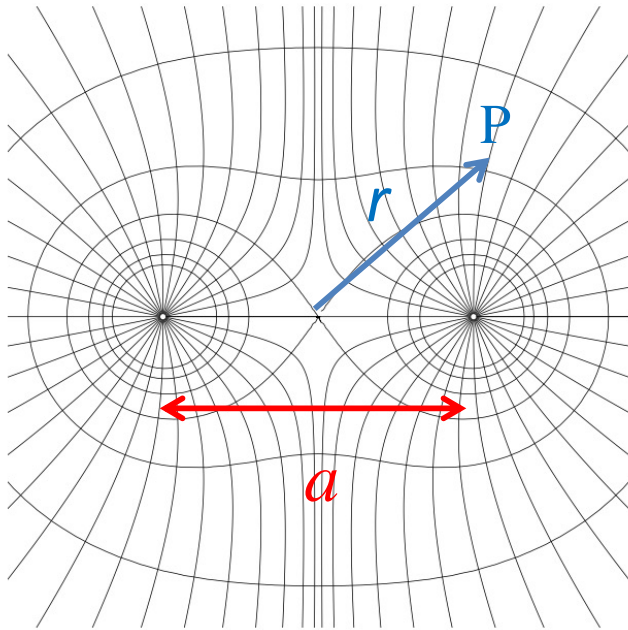


1983: W discovery:  $pp \rightarrow W X$



C. Albajar et al., (UA1 Collaboration),  
 Z. Phys. C 44, 15 (1989)

# Another example of EFT : multipole expansion



- Potential from point-like charges (sources)

$$V(\mathbf{r}) = \frac{1}{r} \sum_{l,m} b_{lm} \frac{1}{r^l} Y_{lm}(\Omega) \quad b_{lm} \equiv c_{lm} a^l$$

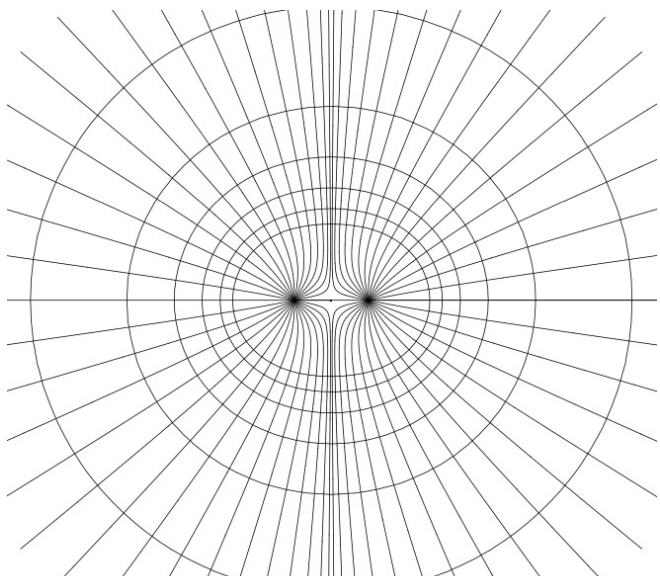
$c_{lm}$  dimensionless coefficient

$a$  short distance (high energy scale)  $\Lambda \sim 1/a$ ,

$$V(\mathbf{r}) = \frac{1}{r} \sum_{l,m} c_{lm} \left(\frac{a}{r}\right)^l Y_{lm}(\Omega) \quad \text{Expansion in } l$$

$$\delta = a/r$$

- Simplification if  $a \ll r$  (far)
- The field far away from the sources looks like a point-like charge (truncation to small  $l$ )



# Other examples of EFT

- **HQET Heavy quark effective theories**

Low-energy dynamics of hadrons containing a heavy quark.

Expansion parameter:  $\Lambda_{\text{QCD}}/m_Q$

- **Chiral Perturbation Theory**

Interactions of pions and nucleons at low momentum transfer. Expansion parameter:  $p/\Lambda_\chi$

$\Lambda_\chi \sim 1 \text{ GeV}$

- **Soft-collinear effective theory SCET**

Energetic QCD processes, the final states (jets) with small invariant mass  $M_J$  wrt to the center-of-mass energy of the collision ( $Q$ ). Expansion parameter:  $M_J/Q$

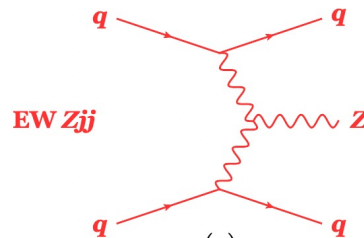
- **SMEFT** describes deviations from the SM. Expansion parameter:  $1/\Lambda$

- **SM is itself an effective theory valid at our accessible energies**

**EFT methods allow us to separate scales in a multi-scale problem, and organize the calculation in a systematic way**

# Example of (SM)EFT fit @ LHC

## Vector Boson Fusion (VBF) Differential cross section

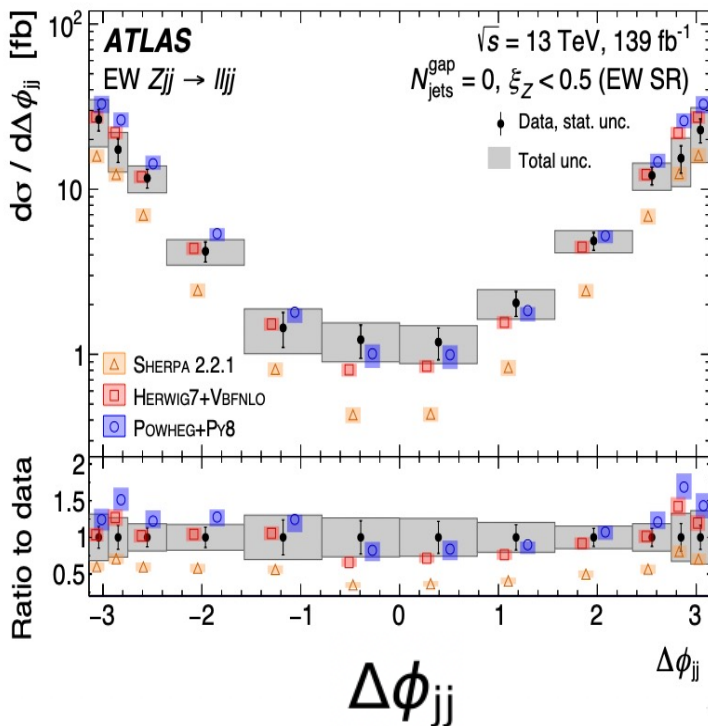


$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i$$

$$c_i = c_W$$

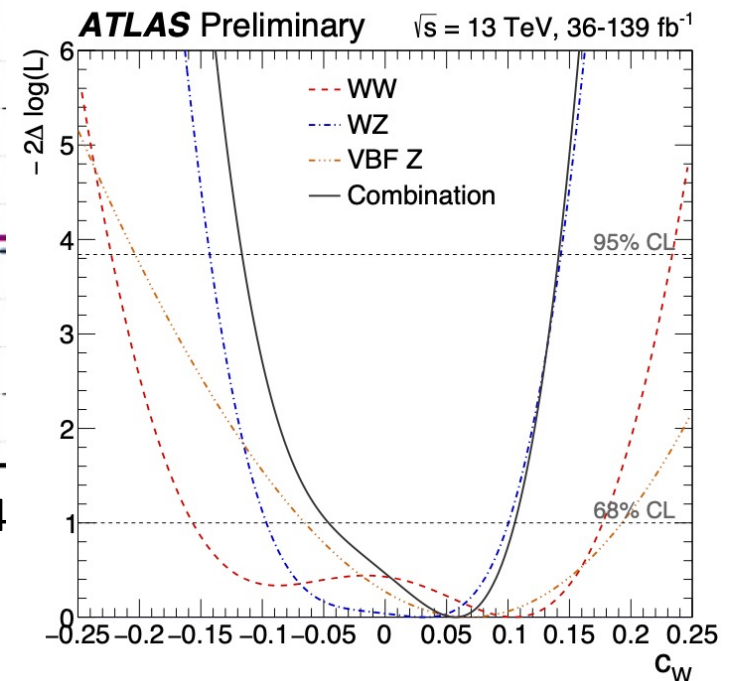
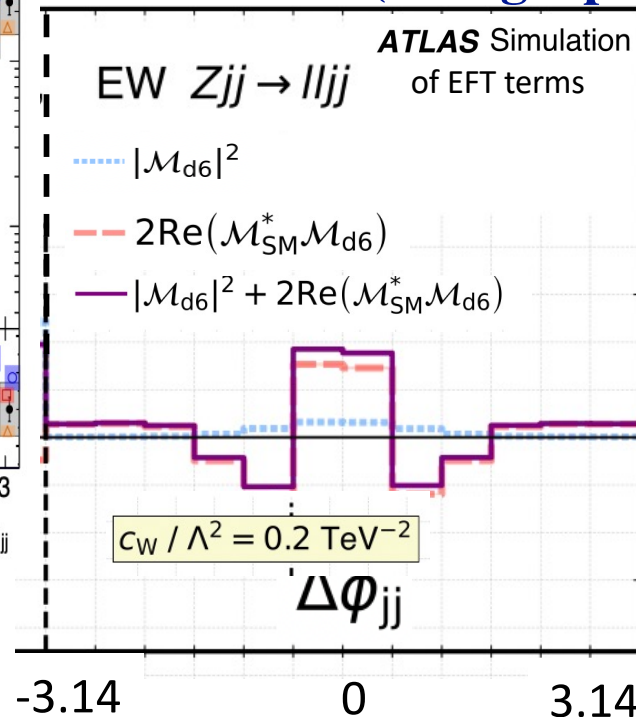
$$\mathcal{O}_i = \varepsilon^{IJK} W_{\mu}^{I\nu} W_{\nu}^{J\rho} W_{\rho}^{K\mu}$$

## One-dimensional confidence level intervals



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## Ratio to SM (Madgraph)

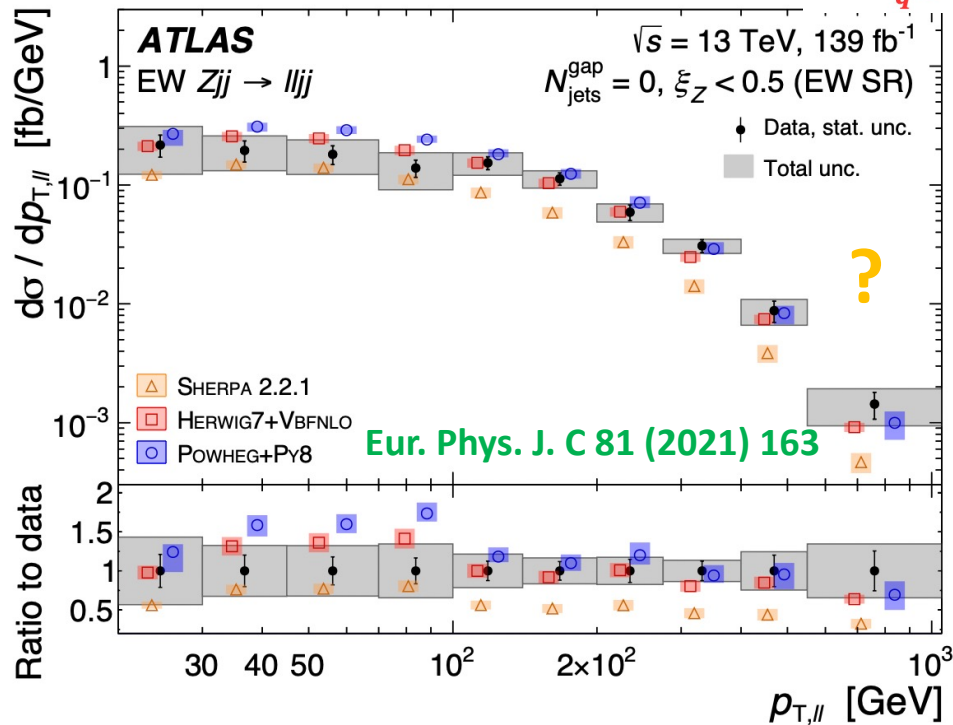
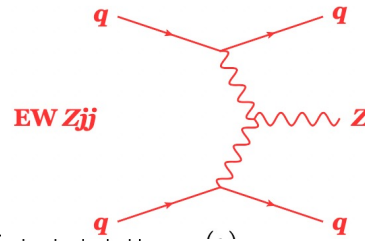


- EFT quantify degree of compatibility wrt SM in a theoretically and statistically well sounded way in many distributions, comparison and combinations among experiments

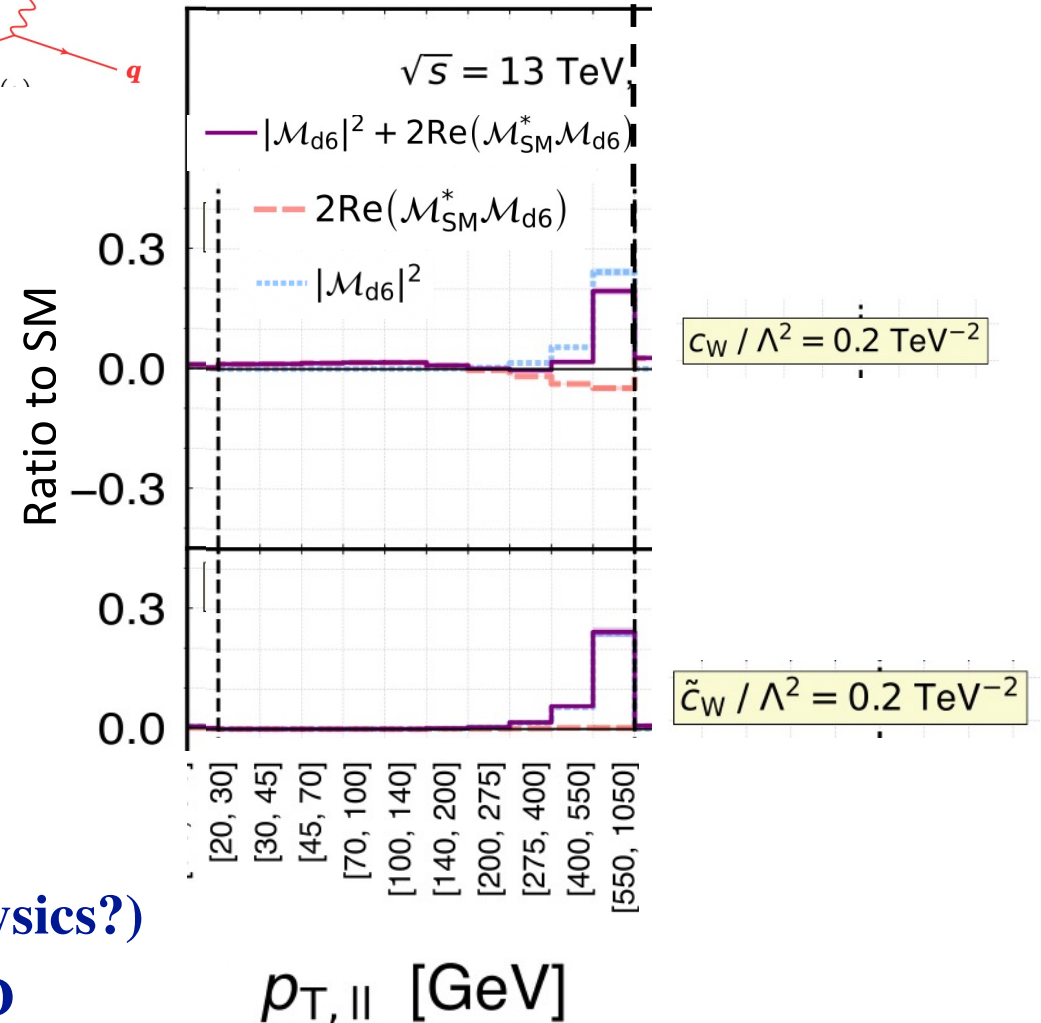


# Important to know well the SM

## Vector Boson Fusion (BVF) Differential cross section



## ATLAS Simulation of EFT terms EW $Zjj \rightarrow lljj$



- Understand the predictions  
(Sherpa has less tail or it is New physics?)
- Know higher orders in QCD & QED



# Summary

## Motivations for EFT mainly from an experimental point of view

- **It takes time to build higher energy machines**
- **Nature has naturally many scales (“divide et impera”)**
- **Few examples from the past show that the EFT approach was successful**
- **EFT allow to quantify small deviations or the degree of compatibility wrt SM in a theoretically and statistically sounded way in many distributions, comparisons and combinations among experiments**
- **Reduced model dependence wrt direct searches**

# Introduction and experimental motivation: End



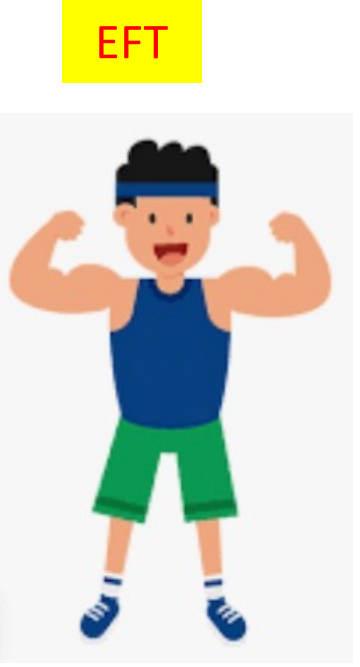
Supersymmetry



Composite Higgs



Extra dimensions



EFT

(.. “but it is not as easy as it seems” ..)

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HEP-EPS 2023 : Patrick Meade

# Additional references

Aneesh V. Manohar, Introduction to Effective Field Theories : arXiv:1804.05863v1

[https://www.hep.phy.cam.ac.uk/~thomson/lectures/partIIIparticles/Handout10\\_2009.pdf](https://www.hep.phy.cam.ac.uk/~thomson/lectures/partIIIparticles/Handout10_2009.pdf)

Mark Thomson lectures