

# DUNE IN<sub>2</sub>P<sub>3</sub> workshop

# Ionization

- Charged particles traversing the detector ionize argon
- “Work” or average energy required to dissociate electron to form argon ion – electron pair in liquid:

$$W_{ion} = 23.6_{-0.3}^{+0.5} \text{ eV (} \text{PhysRevA.9.1438)} \text{)}$$

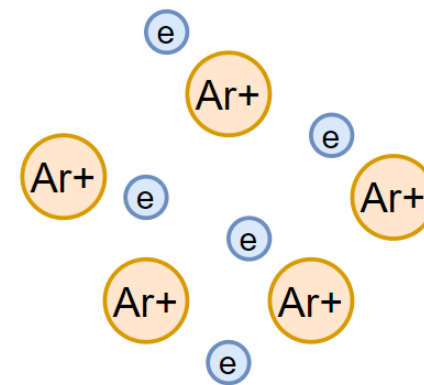
- The ionization electrons are the charge that is detected in LAr TPC by drifting it to sensing electrodes (e.g., wires, strips)
- Trivially one can find a number of electrons produced per 1 MeV of deposited energy: ~42 ke / MeV
- The fluctuation on this number can be estimated from Fano factor  $F = 0.107$  [[NIM. 134 \(1976\)](#)]

$$\frac{\sigma_n}{n} = \sqrt{\frac{F}{n}} \sim 0.2\%$$

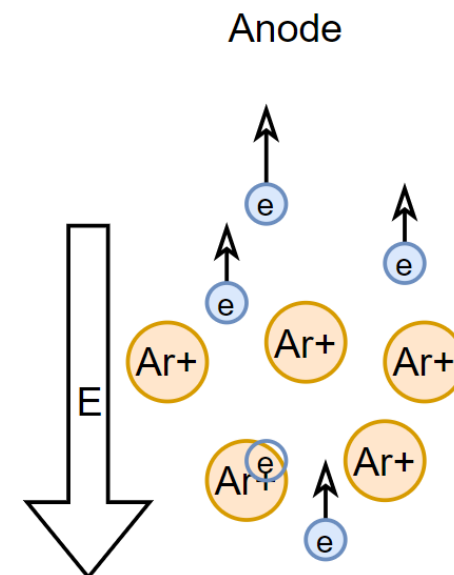
Negligible with respect to fluctuation in  $dE/dx$

# Recombination

- Liquid argon is a dense medium
- Some electrons recombine with ions leading to an intrinsic “loss” of ionization signal
- The fraction of recombined electrons depends on
  - The density of charge depositions (typically taken to be proportional as  $dE/dx$ )
    - The higher the local density of ion-electrons the more likely they will recombine
  - Strength of the electric field in the detector
    - With no field all electrons would recombine
    - With infinitely strong field the recombination would be minimal
- In practice the recombination (model) is fitted as function of  $dE/dx$  and electric field



Jaffe ([Ann. Phys. 42 \(1913\) 303](#))  
columnar model



Electron drift velocity  $O(\text{mm}/\mu\text{s})$   
Ion drift velocity  $O(\text{cm}/\text{s})$

# Recombination pumps light production

- The passage of ionizing particles also leads to formation of excimer states of Ar
- These decay via emission of VUV photons (128 nm)
- Two ways to produce excimer states
  1. Self-trapping:  $\text{Ar}^* + \text{Ar} \rightarrow \text{Ar}_2^* \rightarrow \text{Ar} + \text{Ar} + \gamma$
  2. Recombination:  $\text{Ar}^+ + \text{Ar} + e^-_{\text{therm}} \rightarrow \text{Ar}_2^* \rightarrow \text{Ar} + \text{Ar} + \gamma$
- Conversion of recombined electrons to light
- Max number of photons are produced in zero electric field (maximum recombination)

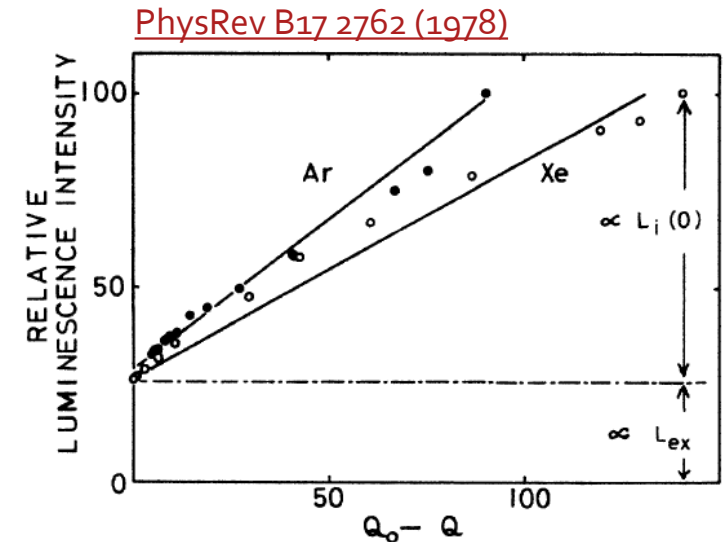


FIG. 4. Relative luminescence intensity against uncollected charge  $Q_0 - Q$ .

$Q_0 - Q = 0$ : very large E-field, so no recombination contribution

# Recombination parametrization: Birks form

- Birks form (ICARUS, [NIMA 523 \(2004\) 275](#)):

$$R = \frac{A}{1 + k/\varepsilon \times dE/dx}$$

$\varepsilon$  – electric field x LAr density,  $dE/dx$  expected energy loss and A, k are constants

- The fitted values (muons) of A and k parameters (NIMA 523) :

$$k = 0.0486 \text{ (kV/cm)(g/MeV cm}^2\text{)}$$

$$A = 0.800$$

# Recombination parametrization: modified Box model

- ArgoNeuT [[JINST 8 P08005 \(2013\)](#)]:

$$R = \frac{\ln(A + \xi)}{\xi}$$
$$\xi = B/\varepsilon \times dE/dx$$

The fit parameters A & B;  $\varepsilon$  – electric field x density

- The fitted parameters (stopping protons) in the paper:

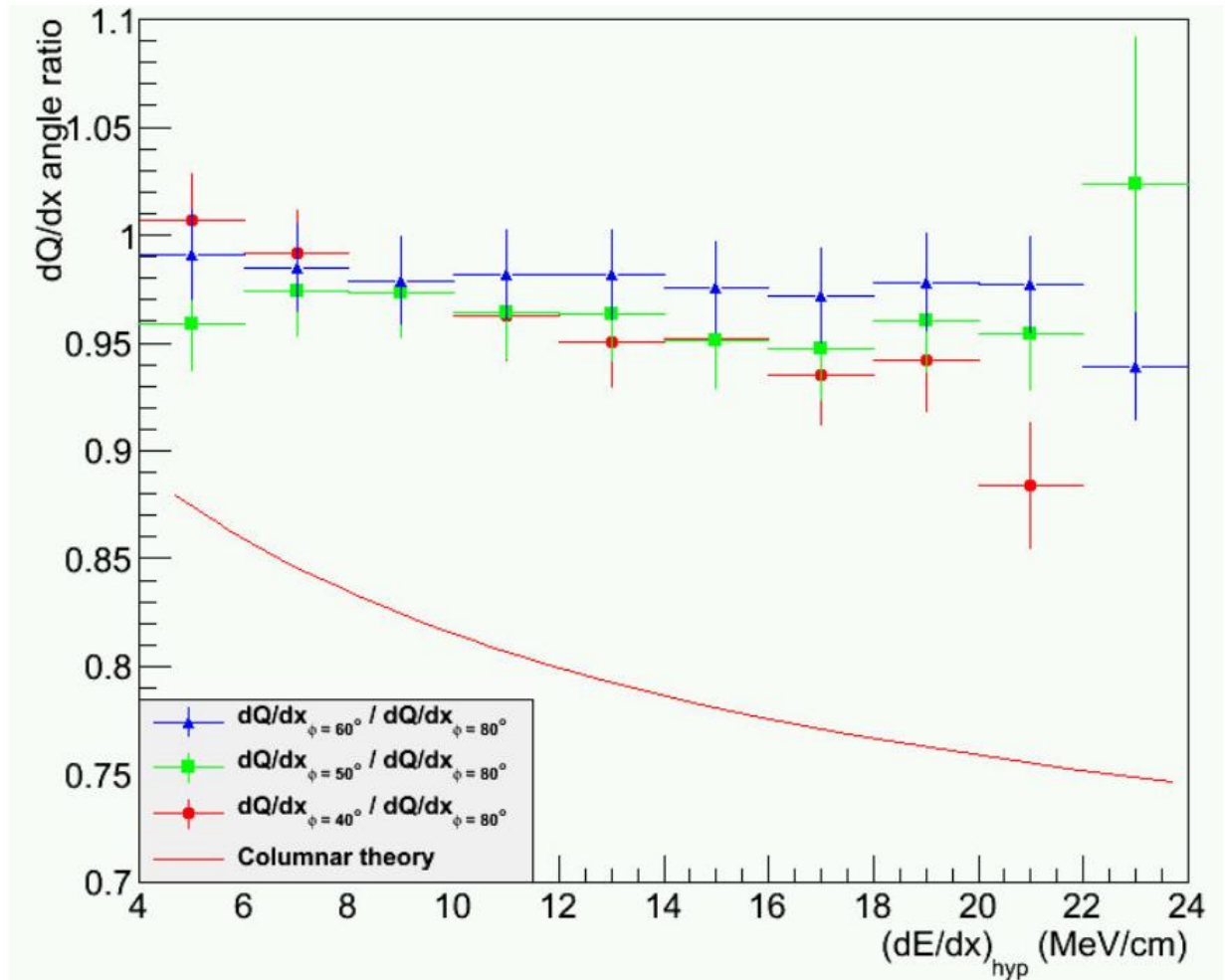
$$B = 0.212 \text{ (kV/cm)(g/MeV cm}^2\text{)}$$

$$A = 0.930$$

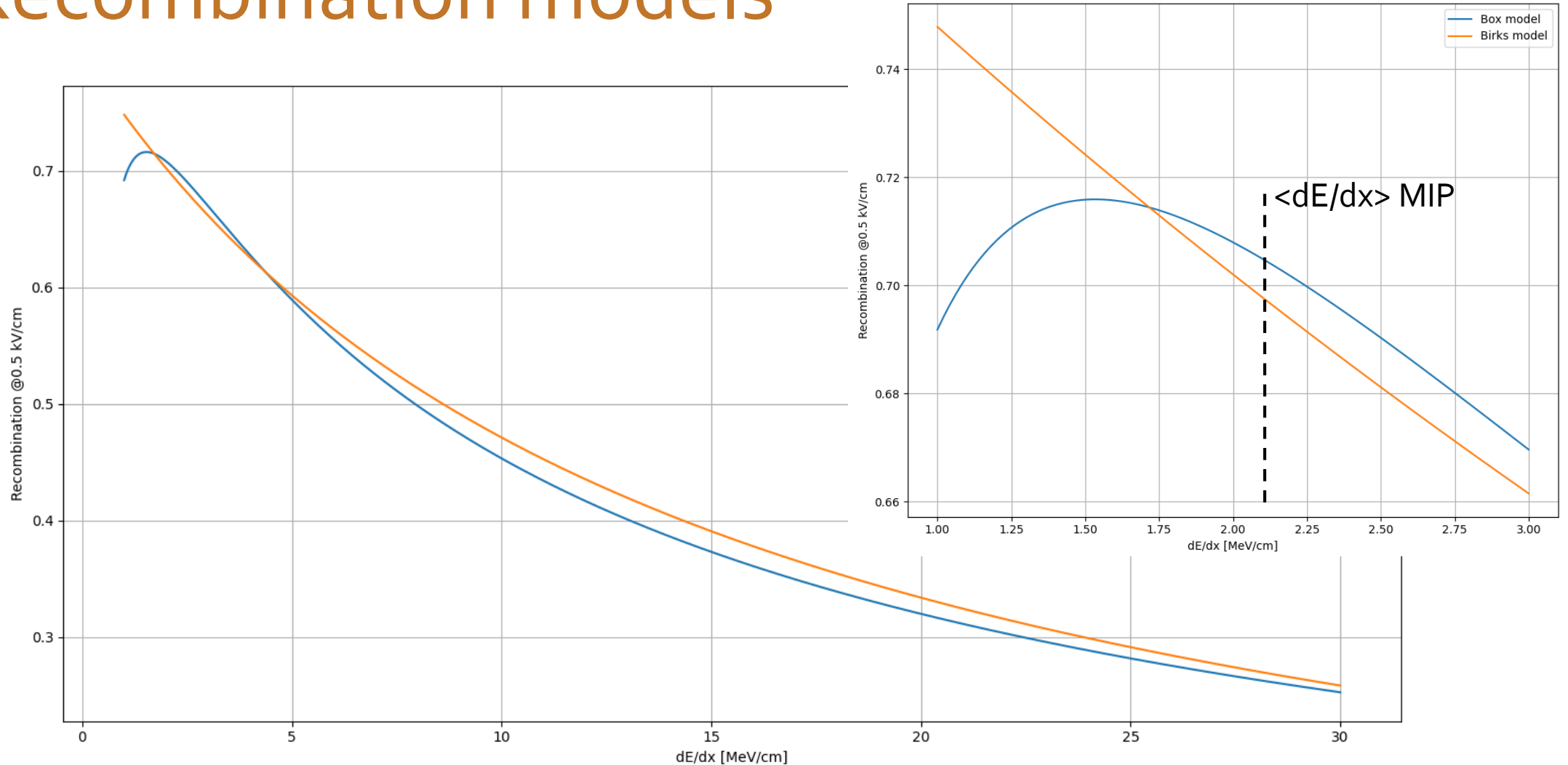
A = 1 in canonical Box model in [[Phys. Rev. A 36 \(1987\) 614](#)] (hence “modified”)

# Angular dependence

- Columnar theory predicts a fairly strong dependence on the angle of “columns” around particle direction relative to the field direction
- No such dependence has been observed in [[JINST 8 P08005 \(2013\)](#)]



# Recombination models



Number of electrons per MIP 2.1MeV/cm @0.5kV/cm: ~62 ke/cm or ~10 fC/cm



# From $dQ/dx$ to $dE/dx$

$$\frac{dQ}{dx} = \frac{1}{W_{ion}} R \left( \frac{dE}{dx}, \varepsilon \right) \frac{dE}{dx}$$

- To obtain  $dE/dx$  from  $dQ/dx$  need to invert recombination model

Birks: 
$$dE/dx = \frac{dQ/dx}{A_B/W_{ion} - k_B \cdot (dQ/dx)/\mathcal{E}}$$

More stable at large  $dQ/dx$   
as in Birks model the  
denominator  $\rightarrow 0$

Box: 
$$dE/dx = (\exp(\beta W_{ion} \cdot (dQ/dx)) - \alpha) / \beta.$$

$$\alpha = A$$

$$\beta = B/\varepsilon$$

# larsoft code

- Look at "[ionization and scintillation](#)" part of larsim for electron/light prod
- Normally the fastest way to find set parameters is to run fhicl-dump of <your\_job>.fcl:

```
fhicl-dump my_job.fcl
```

- **Exception**: if parameters are not specified by default in fcl! (i.e., fhicl-dump will only show what is declared in various fcl files)

## DetectorPropertiesStandard:

```
fhicl::Atom<double> ModBoxAlpha{  
  Name("ModBoxAlpha"),  
  Comment("alpha parameter in the Modified Box recombination model."),  
  util::kModBoxA};  
  
fhicl::Atom<double> ModBoxBeta{  
  Name("ModBoxBeta"),  
  Comment("beta parameter in the Modified Box recombination model."),  
  util::kModBoxB};
```

The default values are used if no fcl parameters are specified

util::kModBoxA(B) come from [PhysicalConstants.h](#)

```
constexpr double kModBoxA = 0.930;  
constexpr double kModBoxB = 0.212;
```

# larsoft code

DetectorPropertiesStandard Birks function can also be found there

```
double DetectorPropertiesStandard::ModBoxCorrection(double dQdx, double E_field) const
{
    // Modified Box model correction has better behavior than the Birks
    // correction at high values of dQ/dx.
    double const rho = Density(); // LAr density in g/cm^3
    constexpr double Wion = 1000. / util::kGeVToElectrons; // 23.6 eV = 1e, Wion in MeV/e
    double const Beta = fModBoxB / (rho * E_field);
    double const Alpha = fModBoxA;
    double const dEdx = (exp(Beta * Wion * dQdx) - Alpha) / Beta;

    return dEdx;
}
```

$$dE/dx = (\exp(\beta W_{ion} \cdot (dQ/dx)) - \alpha) / \beta.$$

This is what is called when one asks  
to calculate dE/dx from dQ/dx in  
CalorimetryAlg

```
double CalorimetryAlg::dEdx_from_dQdx_e(detinfo::DetectorClocksData const& clock_data,
                                         detinfo::DetectorPropertiesData const& det_prop,
                                         double dQdx_e,
                                         double const time,
                                         double const T0,
                                         double const EField) const
{
    if (fDoLifeTimeCorrection) {
        dQdx_e *= LifetimeCorrection(clock_data, det_prop, time, T0); // (dQdx_e in e/cm)
    }

    if (fUseModBox) { return det_prop.ModBoxCorrection(dQdx_e, EField); }

    return det_prop.BirksCorrection(dQdx_e, EField);
}
```

# More FCL parameter complications

## Ex. ISCalcCorrelated.cxx

```
ISCalcCorrelated::ISCalcCorrelated(detinfo::DetectorPropertiesData const& detProp,  
                                   CLHEP::HepRandomEngine& Engine)  
  : fISTPC{*(lar::providerFrom<geo::Geometry>())}  
  , fSCE(lar::providerFrom<spacecharge::SpaceChargeService>())  
  , fBinomialGen{CLHEP::RandBinomial(Engine)}  
{  
  MF_LOG_INFO("ISCalcCorrelated") << "IonizationAndScintillation/ISCalcCorrelated Initialize.";  
  
  fScintPreScale = lar::providerFrom<detinfo::LARPropertiesService>()->ScintPreScale();  
  
  art::ServiceHandle<sim::LArG4Parameters const> LArG4PropHandle;  
  
  // The recombination coefficient is in g/(MeVcm^2), but we report  
  // energy depositions in MeV/cm, need to divide Recombk from the  
  // LArG4Parameters service by the density of the argon we got  
  // above.  
  fRecombA = LArG4PropHandle->RecombA();  
  fRecombk = LArG4PropHandle->Recombk() / detProp.Density(detProp.Temperature());  
  fModBoxA = LArG4PropHandle->ModBoxA();  
  fModBoxB = LArG4PropHandle->ModBoxB() / detProp.Density(detProp.Temperature());  
  fUseModBoxRecomb = (bool)LArG4PropHandle->UseModBoxRecomb();  
  fUseModLarqlRecomb = (bool)LArG4PropHandle->UseModLarqlRecomb();
```

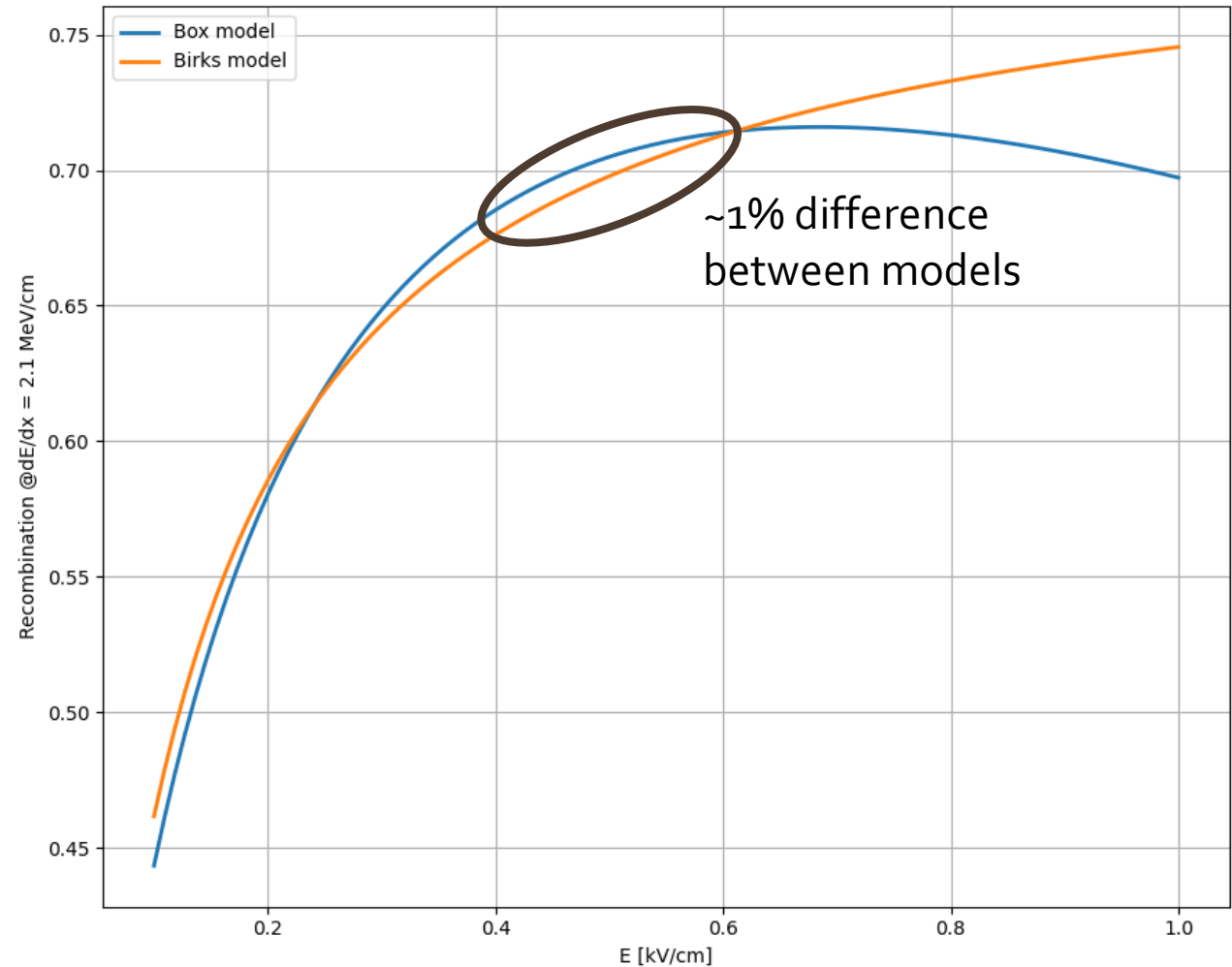
Note that for simulation the recombination parameters are in a LArG4Parameters

Defined in standard\_largeantparameters in [simulationservices.fcl](#)

So as far as I can tell one defines the same parameters in different places.

# From $dQ/dx$ to $dE/dx$

- Converting  $dQ/dx$  to  $dE/dx$  is not a just inversion of recombination
- The charge that is measured needs to be corrected for a number of effects:
  - E.g., attenuation due to attachment to electro-negative impurities
- The drift field non-uniformity needs to be assessed and included, since recombination rate depends on the local field strength



# Drifting charge

- Ionization electrons drift in applied electric field
- Some fraction of the charge is captured by electronegative impurities like O<sub>2</sub>
- The electrons bounce off the atoms as they traverse the medium
  - Diffusion of original cloud both longitudinally and transversely to the direction motion
- The average speed “drift velocity” is a function of field strength and temperature (density)

# Electron attachment to impurities

- With capture rate to a given molecule species  $k_s$  (E-field dependence) and assuming a constant concentration of impurities  $n_s \gg N_e$ , the number of electrons captured per time interval  $dt$ :

$$\frac{dN_e}{dt} = -k_s n_s N_e$$

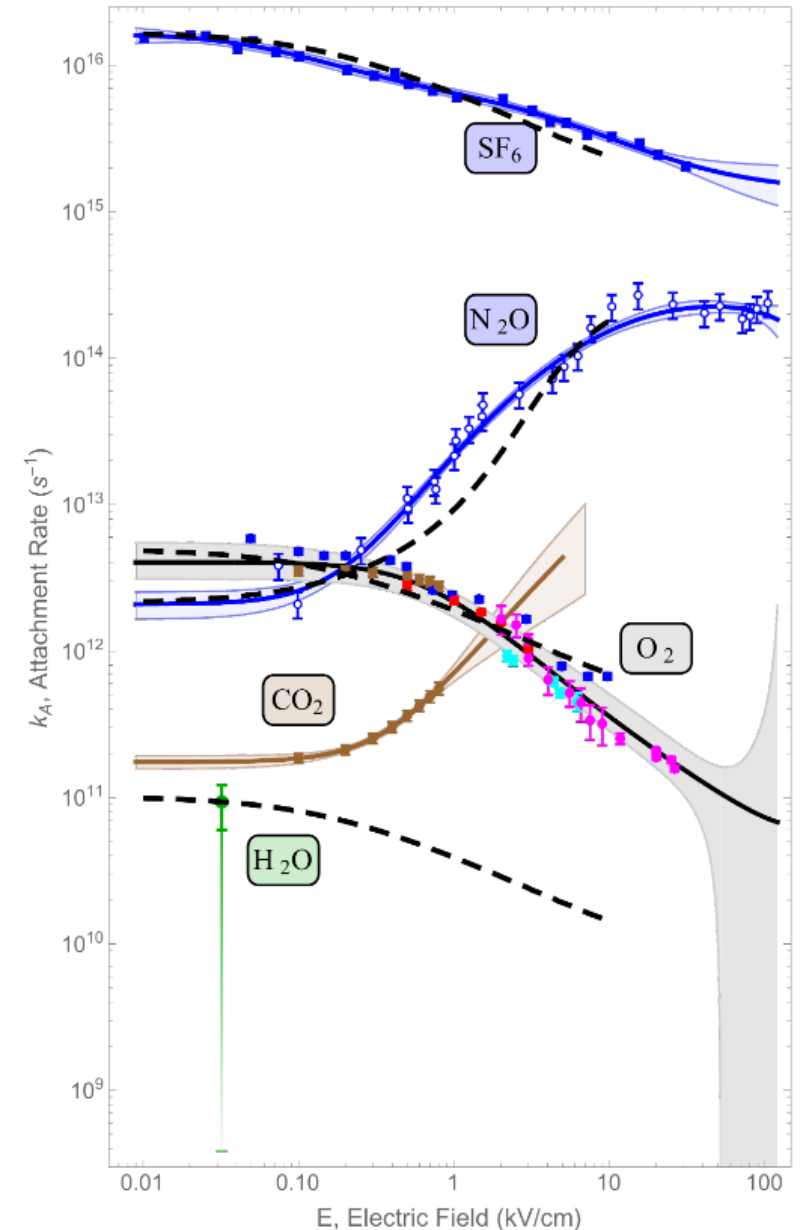
With:

$$\tau_e = \frac{1}{k_s n_s}$$

Fitted from data

$$N_e(t) = N_e(0)e^{-t/\tau_e}$$

From [2022 JINST 17 T11007](#)



# Electron attachment to impurities

$$N_e(t) = N_e(0)e^{-t/\tau_e}$$

- The electron lifetime  $\tau_e$  is measured in s (on a ms scale)
- One expresses concentration of impurities in oxygen equivalents (the dominant)
- If  $\tau_e$  measured in ms, what equivalent concentration of O<sub>2</sub> this corresponds to?

$$\rho_{O_2} [\text{ppb}] = 10^9 \times n_s / n_{Ar}$$

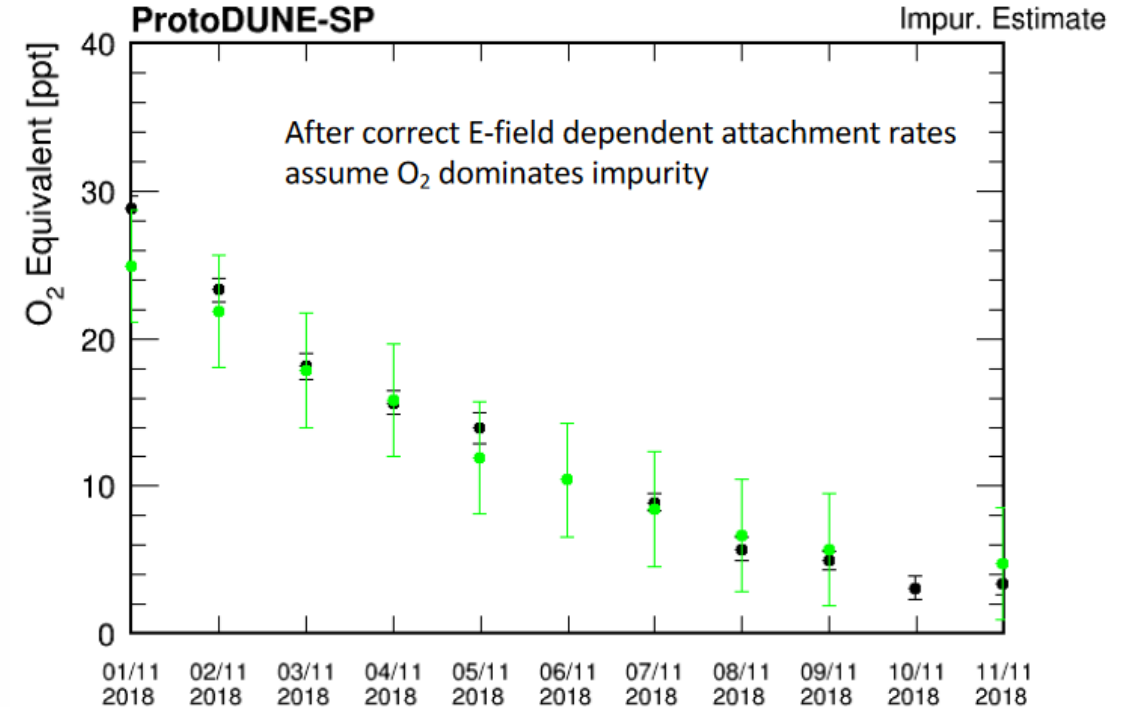
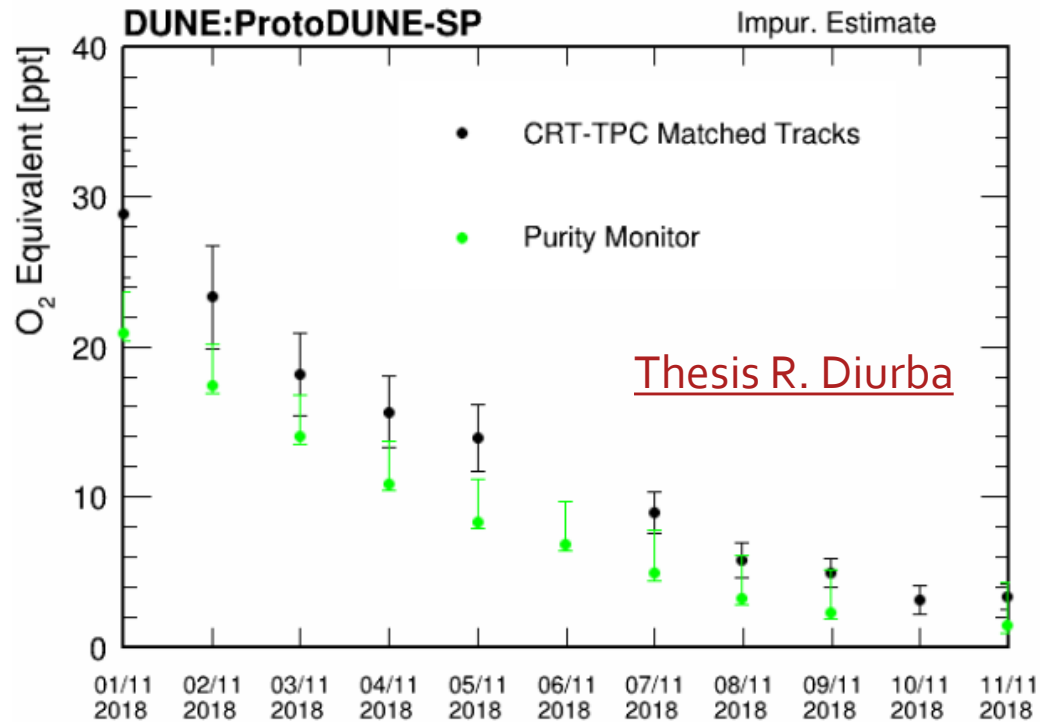
with  $n_{Ar}$  number of moles per L of LAr = 35

With  $k_s \sim 10^{11}$  [L/mol/s] for O<sub>2</sub>

$$\tau_e [\text{ms}] \sim \frac{0.3}{\rho_{O_2} [\text{ppb}]}$$



# Purity monitor and purity in TPC



- Purity monitors (10-100 V/cm) and TPC (500 V/cm) work at different fields: impact on the attachment rate constants
- From parametrization of O<sub>2</sub> data  $k_S(E_{PRM})/k_S(E_{TPC}) \sim 1.3$

# Electron diffusion

- In presence of external electric field (“drift field”) electrons are accelerated between collisions. In each collision:
  - They lose some part of energy
  - They lose the original direction
- On average the electron cloud  $n$  moves or drifts towards anode with “drift” velocity  $v_d$ :

Current conservation of moving and diffusing charge current  $J = v_d n - D \nabla n$        $\frac{\partial n}{\partial t} = D \nabla^2 n - v_d \frac{\partial n}{\partial x}$       Drift direction in X with velocity  $v_d$  & no losses due to impurities

- Electrons have position dependent energy within the cloud
  - Electrons diffused to the leading edge of the swarm tend to be more energetic and therefore less sensitive to drift field effects
  - These fall back to the middle/tail of the bunch when they loose enough energy
  - Semiquantitative treatment in [[Phys. Rev. 181 \(1969\)](#)] solving Boltzmann transport equations
- Leads anisotropic diffusion with two different diffusions coefficients  $D_L$  and  $D_T$  in the direction parallel and perpendicular to the electric field with  $D_L \leq D_T$  :

$$\frac{\partial n}{\partial t} = D_L \frac{\partial^2 n}{\partial x^2} + D_T \left( \frac{\partial^2 n}{\partial y^2} + \frac{\partial^2 n}{\partial z^2} \right) - v_d \frac{\partial n}{\partial x}$$

# Electron diffusion

- Effective solution for anisotropic diffusion:

$$n(r, t) = \frac{n_0}{4\pi D_T t \sqrt{4\pi D_L t}} \exp \left[ -\frac{(x - v_d t)^2}{4D_L t} - \frac{y^2 + z^2}{4D_T t} \right]$$

- The cloud dimensions increase with time:

$$\begin{aligned} \sigma_T^2(t) &= 2D_T t + \rho_0^2 \\ \sigma_L^2(t) &= 2D_L t + \sigma_0^2 \end{aligned}$$

Some initial dimensions of the electron cloud

- For longitudinal diffusion it is more natural to express width in time units (the drift coordinate in TPC):

$$\sigma_{L,t}^2 = \frac{2D_L t}{v_d^2} + \sigma_{0,t}^2 = \frac{2D_L x}{v_d^3} + \sigma_{0,t}^2$$

# Electron diffusion DL/DT values

- The sim values for DL / DT:

$$D_L = 6.2 \text{ cm}^2/\text{s}$$

$$D_T = 16.3 \text{ cm}^2/\text{s}$$

- These come from theoretical calculations and higher than experimental data ([docdb-14407](#))
- The docdb offers: (L) 5.3 cm<sup>2</sup>/s & (T) 12.8 cm<sup>2</sup>/s
- The values are higher than existing data
  - [ICARUS](#):  $D_L = 4.74 \text{ cm}^2/\text{s}$
  - In Protodune-SP data:  $D_L = 3.91 \text{ cm}^2/\text{s}$
- Ideally should try to be coherent at least within DUNE ...

[simulationservices\\_dune.fcl](#)

```
#include "simulationservices.fcl"
```

```
BEGIN_PROLOG
```

```
#FD1-HD
```

```
dunefd_largeantparameters:
```

```
{
```

```
  @table::standard_largeantparameters
```

```
  LongitudinalDiffusion: 6.2e-9 #cm^2/ns
```

```
  TransverseDiffusion: 1.63e-8 #cm^2/ns
```

```
}
```

```
fhicl-dump protodunevd_detsim.fcl | grep Diffusion
```

```
LongitudinalDiffusion: 6.2e-9
```

```
TransverseDiffusion: 1.63e-8
```

# DL measured with gold photocathode

$$D_L = \frac{\mu \epsilon_L}{e} = \left( \frac{a_0 + a_1 E + a_2 E^{3/2} + a_3 E^{5/2}}{1 + (a_1/a_0)E + a_4 E^2 + a_5 E^3} \right) \left( \frac{b_0 + b_1 E + b_2 E^2}{1 + (b_1/b_0)E + b_3 E^2} \right) \left( \frac{T}{T_0} \right)^{-3/2} \left( \frac{T}{T_1} \right)$$

Table 1: Function parameters in Eq. (21)

$a_0$	=	551.6
$a_1$	=	7953.7
$a_2$	=	4440.43
$a_3$	=	4.29
$a_4$	=	43.63
$a_5$	=	0.2053

Table 2: Function parameters in Eq. (22)

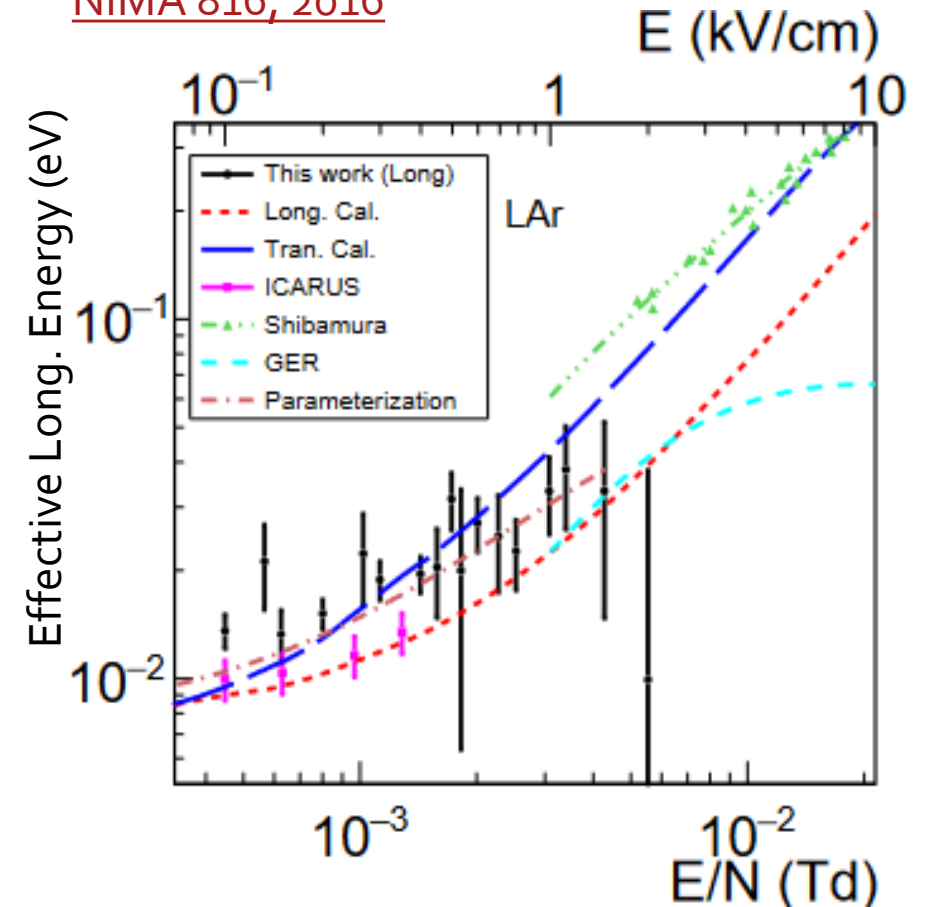
$b_0$	=	0.0075
$b_1$	=	742.9
$b_2$	=	3269.6
$b_3$	=	31678.2

- This calculator from BNL group:

- DL = 6.63 cm<sup>2</sup>/s
- DT = 13.23 cm<sup>2</sup>/s

$$\frac{D_L}{D_T} = 1 + \frac{E}{\mu} \frac{\partial \mu}{\partial E}$$

NIMA 816, 2016



See refs in NIMA 816 for more details

# Drift velocity in liquid argon

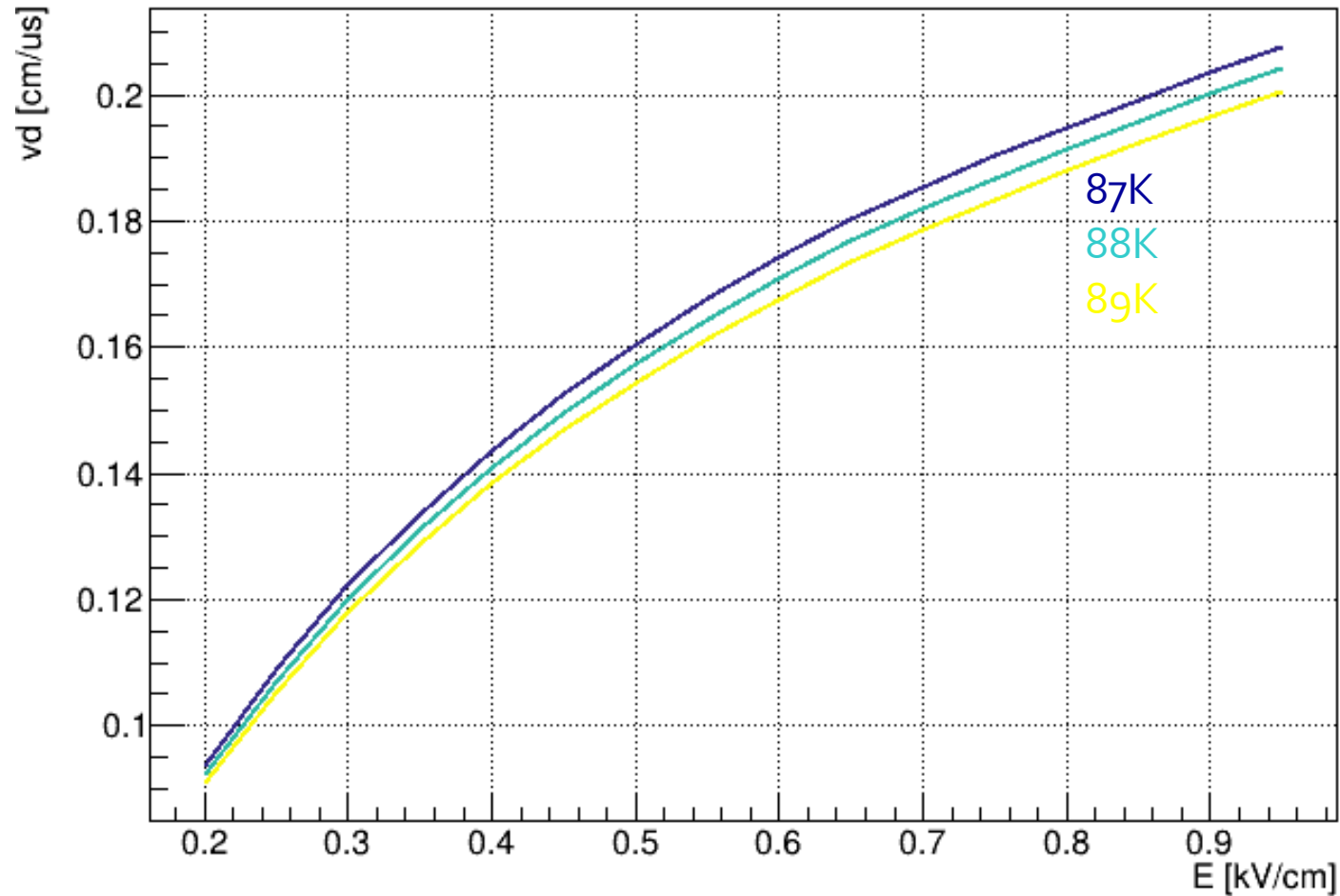
- Drift velocity has dependence on not only field strength but also temperature
- A smooth interpolation of a so-called “Walkowiak” parameterization [NIM A 449 (2000) 288-294] and fitted ICARUS measurements below 0.7 kV/cm is a default option in larsoft
- Not sure about wirecell. In [this](#) paper [2018 *JINST* **13** P07006] it says “the electron drift velocity as function of electric field is taken from recent measurements...”
- At 0.5 kV/cm  $v_d = 1.6$  mm/us
- 3m drift distance in constant field: ~1.9 ms

From [DetectorPropertiesStandard](#)

```
//-----//  
double DetectorPropertiesStandard::DriftVelocity(double efield, double temperature) const  
{  
    // Drift Velocity as a function of Electric Field and LAr Temperature  
    // from : W. Walkowiak, NIM A 449 (2000) 288-294  
    //  
    // Option to use MicroBooNE+ICARUS model (as in arXiv:2008.09765) provided as  
    // well, with temperature dependence as prescribed by Mike Mooney based on  
    // looking at the Walkowiak data.  
    //  
    // Efield should have units of kV/cm  
    // Temperature should have units of Kelvin
```

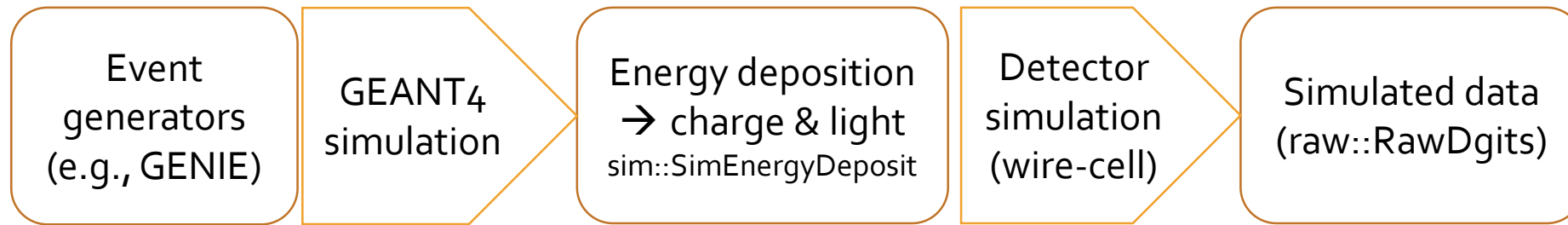
- Need a decent electron lifetime to see charge close to the cathode
- Example:  $\tau_e = 5$  ms, max attenuation ~30%

# Drift velocity in liquid argon



1K diff ~2% effect on drift velocity @0.5kV/cm

# Detector simulation workflow



- Some care must be taken for WCL (Wire-Cell) module to pick up the “transport” parameters from larsoft services:

## In [wirecell\\_dune.fcl](#)

For protodunespmc (set interpedently of larsoft):

```
// Make available parameters via Jsonnet's std.extVar()
params: {
}
structs: {
  // Longitudinal diffusion constant [cm2/s]
  DL: 4.0
  // Transverse diffusion constant [cm2/s]
  DT: 8.8
  // Electron lifetime [ms]
  lifetime: 35.0
  // Electron drift speed from SP measurement
  driftSpeed: 1.565
}
```

Example protodunevd (using values from larsoft services):

```
structs: {
  nticks: @local::protodunevd_services.DetectorPropertiesService.NumberTimeSamples
  lifetime: @local::protodunevd_services.DetectorPropertiesService.Electronlifetime
  DL: @local::dunefd_largeantparameters.LongitudinalDiffusion
  DT: @local::dunefd_largeantparameters.TransverseDiffusion
  efield: @local::protodunevd_services.DetectorPropertiesService.Efield[0] # kV/cm
  temperature: @local::protodunevd_services.DetectorPropertiesService.Temperature # K
}
```

The drift speed is computed by WCL then



# Passing info to WCL

```
wirecell_protodunevd_mc:
{
  module_type : WireCellToolkit
  wcls_main: {
    tool_type: WCLS
    apps: ["Pgrapher"]
    plugins: ["WireCellPgraph", "WireCellGen","WireCellSio","WireCellRoot","WireCellarsoft"]
    // needs to be found via your WIRECELL PATH
    configs: ["pgrapher/experiment/protodunevd/wcls-sim-drift-simchannel.jsonnet"]
    inputers: ["wclsSimDepoSource:electron"]
    outputers: [
      "wclsSimChannelSink:postdrift",
      "wclsFrameSaver:simdigits"
    ]

    // Make available parameters via Jsonnet's std.extVar()
    params: {
    }
    structs: {
      nticks: @local::protodunevd_services.DetectorPropertiesService.NumberTimeSamples
      lifetime: @local::protodunevd_services.DetectorPropertiesService.Electronlifetime
      DL: @local::dunefd_largeantparameters.LongitudinalDiffusion
      DT: @local::dunefd_largeantparameters.TransverseDiffusion
      efield: @local::protodunevd_services.DetectorPropertiesService.Efield[0] # kV/cm
      temperature: @local::protodunevd_services.DetectorPropertiesService.Temperature # K
    }
  }
}
```

## wcls-sim-drift-simchannel.jsonnet

```
local wcls_input = {
  // depos: wcls.input.depos(name="", art_tag="IonAndScint"),
  depos: wcls.input.depos(name='electron', art_tag='IonAndScint'),
};
```

## Need IonAndScint sim::Energy art product

```
local params = base {
  daq: super.daq {
    nticks: std.extVar('nticks'),
  },
  lar: super.lar {
    // Longitudinal diffusion constant
    DL: std.extVar('DL') * wc.cm2 / wc.s,
    // Transverse diffusion constant
    DT: std.extVar('DT') * wc.cm2 / wc.s,
    // Electron lifetime
    lifetime: std.extVar('lifetime') * wc.ms,
    // Electron drift speed
    // drift_speed: std.extVar('driftSpeed') * wc.mm / wc.us,
    drift_speed: util.drift_velocity(std.extVar('efield'), std.extVar('temperature')) * wc.mm / wc.us,
  },
};
```

drift\_speed is computed here from field value and temperature

Q: How does one treat non-uniformities in drift field (e.g., due to SCE)? 25

# Example eventdump.fcl

- To quickly examine art products in a file use eventdump.fcl
- Example: `lar -c eventdump.fcl <my_file>.root -n 1`

```
Begin processing the 1st record. run: 20230221 subRun: 0 event: 1 at 17-Apr-2023 16:57:07 CEST
PRINCIPAL TYPE: Event
PROCESS NAME | MODULE LABEL.. | PRODUCT INSTANCE NAME..... | DATA PRODUCT TYPE..... | SIZE
DummyBetaGen | generator..... | ..... | std::vector<simb::MCTruth>..... | ...1
DummyBetaGen | rns..... | ..... | std::vector<art::RNGsnapshot>..... | ...1
DummyBetaGen | TriggerResults | ..... | art::TriggerResults..... | ...1
G4..... | elecDrift..... | ..... | std::vector<sim::SimChannel>..... | .131
G4..... | rns..... | ..... | std::vector<art::RNGsnapshot>..... | ...3
G4..... | IonAndScint... | ..... | std::vector<sim::SimEnergyDeposit>..... | .158
G4..... | TriggerResults | ..... | art::TriggerResults..... | ...1
G4..... | largeant..... | ..... | std::vector<simb::MCParticle>..... | ..43
G4..... | largeant..... | LArG4DetectorServicevolTPCActive | std::vector<sim::SimEnergyDeposit>..... | .158
G4..... | largeant..... | LArG4DetectorServicevolCryostat. | std::vector<sim::SimEnergyDeposit>..... | ...1
G4..... | largeant..... | ..... | std::map<int,std::set<int> >..... | ...0
G4..... | largeant..... | ..... | art::Assns<simb::MCTruth,simb::MCParticle,sim::GeneratedParticleInfo> | ..43
G4..... | IonAndScint... | priorSCE..... | std::vector<sim::SimEnergyDeposit>..... | .158
```

According to WCL config it should look for IonAndScint product

On to part 2 ...