FINAL ACT (2007-2022) Thibaut Louis

6 meter telescope (4x Planck angular resolution) Located in the Atacama desert (Chile). Observing in 5 frequencies bands : $-$ > 30, 40, 90, 150, 220 GHz With 6000 detectors (transition edges sensor)

Next data release: ACT DR6 (2024)

ACT survey: approx. 40% of the sky

ACT polarisation field

ACT + PLANCK PLANCK

E modes

(video from Sigurd Naess) ⁵

Signal to noise: ACT vs Planck

DR6 data errors

150 GHz

The Atacama Cosmology Telescope: DR6 Power spectra, Likelihood, and constrains on LCDM Louis, La Posta, Li, et al (expected 2024)

ACT Temperature anisotropies field

ACT DR6 + PLANCK PLANCK

9 (video from Sigurd Naess)

Signal to noise: ACT vs Planck

Radio sources: mostly AGN

Galaxy clusters

thermal Sunyaev-Zel'dovich effect

Hot electron gas

Galaxy clusters

tSZ-selected ACT Clusters

ACT has published around 4000 galaxy clusters, expect 7000 with the final data release

150 GHz

150 GHz

DR6 data errors 150 GHz

The Atacama Cosmology Telescope: DR6 Power spectra, Likelihood, and constrains on LCDM Louis, La Posta, Li, et al (expected 2024)

Strong tests on LCDM

Discriminate $f_{EDE} \sim 0.1$ model from
 $\triangle CDM$ at \sim 10-20 σ

ACT DR6 lensing result

The Atacama Cosmology Telescope: A Measurement of the DR6 CMB Lensing Power Spectrum and its Implications for Structure Growth: Qu et al. (April 2023)

The Atacama Cosmology Telescope: DR6 Gravitational Lensing Map and Cosmological Parameters: [Madhavacheril](https://arxiv.org/search/astro-ph?searchtype=author&query=Madhavacheril,+M+S) et al. (April 2023)

The Atacama Cosmology Telescope: Mitigating the impact of extragalactic foregrounds for the DR6 CMB lensing analysis: MacCrann et al. (April 2023)

Image ESA: Planck

Unlensed

Lensed

What does CMB lensing tells us ?

Lensing power spectrum is a projected matter power spectrum

Redshift kernel

How do we measure this effect ?

We assume that the cosmological principle is correct The key assumption here is that, in an isotropic universe, the angular covariance matrix describing the CMB statistical properties is only a function of the angular separation of the different line of sights $\xi^{TT}(\hat{n}_1, \hat{n}_2) = \langle T(\hat{n}_1)T(\hat{n}_2) \rangle = \xi(\hat{n}_1, \hat{n}_2) = \xi(\cos \theta)$

How do we measure this effect ?

We assume that the cosmological principle is correct The key assumption here is that, in an isotropic universe, the angular covariance matrix describing the CMB statistical properties is only a function of the angular separation of the different line of sights $\xi^{TT}(\hat{n}_1, \hat{n}_2) = \langle T(\hat{n}_1)T(\hat{n}_2) \rangle = \xi(\hat{n}_1, \hat{n}_2) = \xi(\cos \theta)$

If this is true, then it implies that in harmonics space The different $a_{\ell m}^T$ are uncorrelated

$$
\langle a^T_{\ell m} a^{T,*}_{\ell' m'} \rangle = C_\ell \delta_{\ell,\ell'} \delta_{m,m'}
$$

How do we measure this effect ?

We assume that the cosmological principle is correct The key assumption here is that, in an isotropic universe, the angular covariance matrix describing the CMB statistical properties is only a function of the angular separation of the different line of sights

$$
\xi^{TT}(\hat{n}_1, \hat{n}_2) = \langle T(\hat{n}_1) T(\hat{n}_2) \rangle = \xi(\hat{n}_1, \hat{n}_2) = \xi(\cos \theta)
$$

If this is true, then it implies that in harmonics space The different $a_{\ell m}^T$ are uncorrelated

$$
\langle a^T_{\ell m} a^{T,*}_{\ell' m'} \rangle = C_\ell \delta_{\ell,\ell'} \delta_{m,m'}
$$

Lensing break the isotropy, part of the sky are magnified while *other are not. The way we reconstruct the lensing field is by measuring the correlation between different* $a_{\ell m}^T$

ACT DR6 lensing map

-Covers a quarter of the sky

- -You can see the projected dark matter distribution
- Few degree-scale structure corresponding to the $P(k)$ peak at $z=1-2$

The Atacama Cosmology Telescope: DR6 Gravitational Lensing Map and Cosmological Parameters: Madhavacheril et al. (April 2023)

ZOOM IN: ACT DR6 Gravitational potential map

ZOOM IN: ACT DR6 Gravitational potential map + Planck CIB

ZOOM IN: ACT DR6 Gravitational potential map + Planck CIB

2x SNR per mode compared to Planck. Reconstruction on mostly linear scales.

CMB lensing power spectrum

The Atacama Cosmology Telescope: A Measurement of the DR6 CMB Lensing Power Spectrum and its Implications for Structure Growth: Qu et al. (April 2023)

CMB lensing power spectrum

- Excellent agreement of our measurement (with no free parameters) with the LCDM theory predictions based on Planck 2018 CMB power spectra. A PTE of 0.17
- Amplitude of lensing (relative to theory amplitude) determined to 2.3%

 $A_{\rm lens} = 1.013 \pm 0.023$

SNR of 43

The Atacama Cosmology Telescope: A Measurement of the DR6 CMB Lensing Power Spectrum and its Implications for Structure Growth: Qu et al. (April 2023)

CMB lensing power spectra

Cosmological constraint

$$
S_8^{\text{CMBL}} \equiv \sigma_8 \left(\frac{\Omega_m}{0.3}\right)^{0.25}
$$
\n
$$
S_8^{\text{CMBL}} = 0.818 \pm 0.022
$$
\nEarly time CMB predictions

\n
$$
S_8^{\text{CMBL}} = 0.828 \pm 0.020
$$
\n
$$
S_8^{\text{CMBL}} = 0.828 \pm 0.020
$$
\nPlanck 2018 CMB aniso.

\n
$$
S_8^{\text{CMBL}} = 0.823 \pm 0.011
$$

ACT lensing is not low !!

ACT lensing is not low!!

Planck CMB aniso. Planck CMB aniso. $(+A_{lens}$ marg.) Planck CMB lensing + BAO SPT CMB lensing + BAO **ACT CMB lensing + BAO ACT+Planck CMB lensing + BAO** DES-Y3 galaxy lensing + BAO KiDS-1000 galaxy lensing + BAO HSC-Y3 galaxy lensing (Fourier) + BAO HSC-Y3 galaxy lensing (Real) + BAO

 0.5 instead of 0.25

Conclusion:

- ACT DR6 lensing papers are out, lensing maps are going to be public on Lambda soon
- ACT lensing is not low
- Lot of cross correlation papers coming
- ACT power spectra/likelihood and parameters are coming !

The Atacama Cosmology Telescope: DR6 Power spectra, Likelihood, and constrains on LCDM Louis, La Posta, Li, et al (expected 2024)

Rosette Nebula

WISE

Planck

ACT + Planck

 $\left| \right|$ act dr6 + planck R:f090 G:f150

BACK UP

Wide-field Infrared Survey Explorer

BAO likelihoods

3.1. BAO likelihoods

Weak lensing measurements depend primarily on the amplitude of matter fluctuations σ_8 , the matter density $\Omega_{\rm m}$, and the Hubble constant H_0 . In order to reduce degeneracies of our σ_8 constraint with the latter parameters and allow for more powerful comparisons of lensing probes with different degeneracy directions, we include information from the 6dF and SDSS surveys. The data we include measures the BAO signature in the clustering of galaxies with samples spanning redshifts up to $z \approx 1$, including 6dFGS (Beutler et al. 2011), SDSS DR7 Main Galaxy Sample (MGS; Ross et al. 2015), BOSS DR12 luminous red galaxies (LRGs; Alam et al. 2017), and eBOSS DR16 LRGs (Alam et al. 2021). We do not use the higher-redshift Emission Line Galaxy (ELG; Comparat et al. 2016), Lyman- α (du Mas des Bourboux et al. 2020), and quasar samples (Hou et al. 2021), though we hope to include these in future analyses. We only include the BAO information from these surveys (which provides constraints in the $\Omega_{\rm m}$ -H₀ plane) and do not include the structure growth information in the redshiftspace distortion (RSD) component of galaxy clustering. We make this choice so as to isolate information on structure formation purely from lensing alone.

$$
\xi^{TT}(\hat{n}_1, \hat{n}_2) = \langle T(\hat{n}_1)T(\hat{n}_2) \rangle = \xi(\hat{n}_1, \hat{n}_2) = \xi(\cos \theta)
$$

$$
\langle a_{\ell m}^T a_{\ell' m'}^{T,*} \rangle = \langle \int d\hat{n} T(\hat{n}) Y_{\ell m}^*(\hat{n}) \int d\hat{n}' T(\hat{n}') Y_{\ell' m'}(\hat{n}') \rangle
$$

$$
\langle a_{\ell m}^T a_{\ell' m'}^{T,*} \rangle = \int d\hat{n} d\hat{n}' \langle T(\hat{n})T(\hat{n}') \rangle Y_{\ell m}^*(\hat{n}) Y_{\ell' m'}(\hat{n}')
$$

$$
= \int d\hat{n} d\hat{n}' \xi(\hat{n}.\hat{n}') Y_{\ell m}^*(\hat{n}) Y_{\ell' m'}(\hat{n}')
$$

We can expand a function of $\cos \theta$ in Legendre polynomials, and expand the legendre polynomial in ${\rm spherical}$ harmonics

$$
\xi(\hat{n}.\hat{n}') = \sum_{\ell_0=0}^{\infty} \frac{2\ell_0 + 1}{4\pi} C_{\ell_0} P_{\ell_0}(\hat{n}.\hat{n}'))
$$

$$
= \sum_{\ell_0=0}^{\infty} C_{\ell_0} \sum_{m=-\ell_0}^{\ell_0} Y_{\ell_0 m_0}(\hat{n}) Y_{\ell_0 m_0}^*(\hat{n}')
$$

$$
\begin{array}{lcl} \langle a_{\ell m}^T a_{\ell' m'}^{T,*} \rangle & = & \displaystyle \int d\hat{n} d\hat{n}' \sum_{\ell_0 = 0}^{\infty} C_{\ell_0} \sum_{m = -\ell_0}^{\ell_0} Y_{\ell_0 m_0}(\hat{n}) Y_{\ell m_0}^*(\hat{n}') Y_{\ell m'}^*(\hat{n}') \\ \\ & = & \displaystyle \sum_{\ell_0 = 0}^{\infty} C_{\ell_0} \sum_{m = -\ell_0}^{\ell_0} \int d\hat{n} Y_{\ell_0 m_0}(\hat{n}) Y_{\ell m}^*(\hat{n}) \int d\hat{n}' Y_{\ell m_0}^*(\hat{n}') Y_{\ell' m'}(\hat{n}') \\ \\ & = & \displaystyle \sum_{\ell_0 = 0}^{\infty} C_{\ell_0} \sum_{m = -\ell_0}^{\ell_0} \delta_{\ell_0, \ell} \delta_{m_0, m} \delta_{\ell_0, \ell'} \delta_{m_0, m'} = C_{\ell} \delta_{\ell, \ell'} \delta_{m, m'} \end{array}
$$

Where we use the orthonormality of spherical harmonics

$$
\int d\hat{n} Y_{\ell_0 m_0}(\hat{n}) Y_{\ell m}^*(\hat{n}) = \delta_{\ell_0,\ell} \delta_{m_0,m}
$$

$$
\langle a^T_{\ell m} a^{T,*}_{\ell' m'} \rangle = C_\ell \delta_{\ell,\ell'} \delta_{m,m'}
$$

