

# (experimental) LHC physics

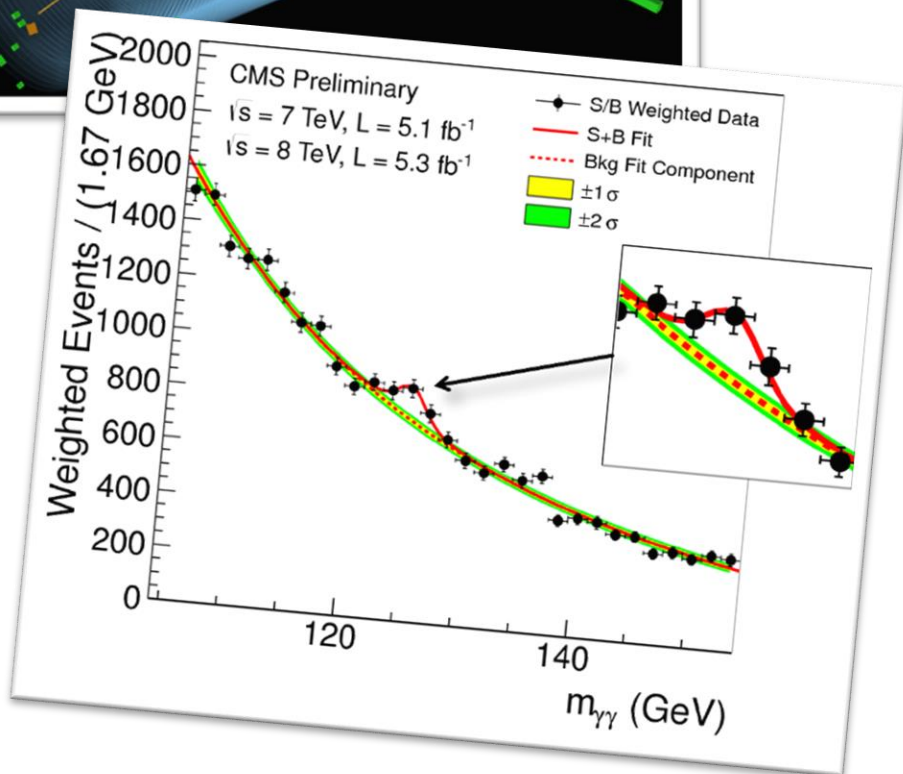
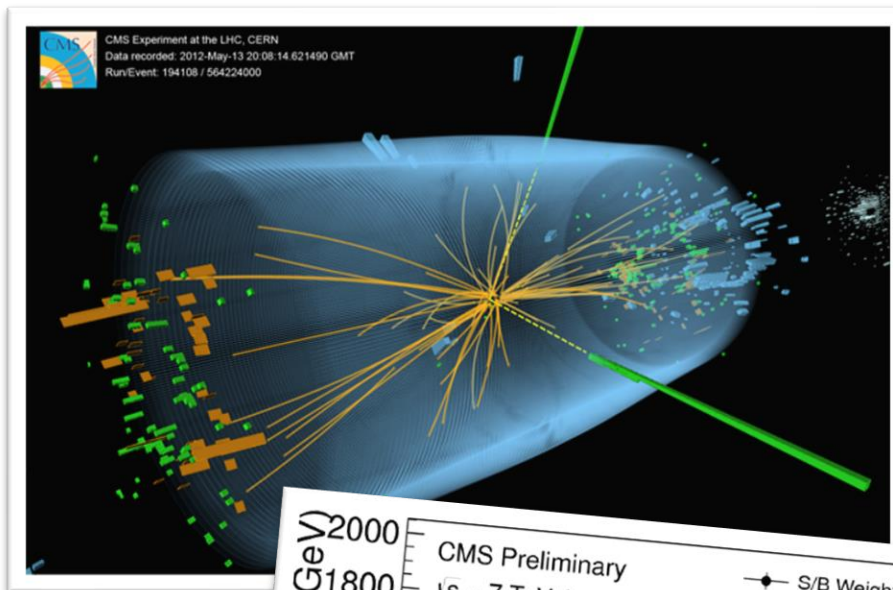


- how (which) particles are produced and measured?

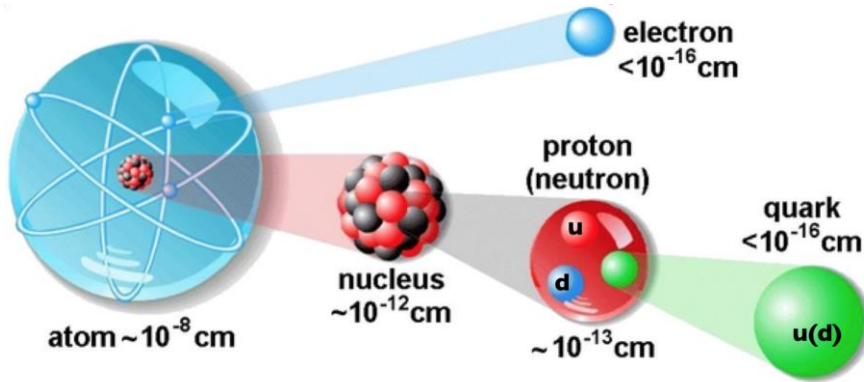
Roberto Covarelli

# Experiment = probing/building theories with data!

$$\begin{aligned}
 & -\frac{1}{2}\partial_\nu g_\mu^\nu \partial_\nu g_\mu^\nu - g_s f^{abc} \partial_\mu g_\nu^a g_\mu^b g_\nu^c - \frac{1}{4}g_s^2 f^{abc} f^{ade} g_\mu^a g_\nu^b g_\mu^c g_\nu^d + \\
 & \frac{1}{2}g_s^2 (\bar{q}_i^\mu \gamma^\nu q_j^\mu) g_\mu^\nu + \bar{C}^a \partial^2 C^a + g_s f^{abc} \partial_\mu \bar{C}^a C^b g_\mu^c - \partial_\nu W_\mu^+ \partial_\nu W_\mu^- - \\
 & M^2 W_\mu^+ W_\mu^- - \frac{1}{2}\partial_\nu Z_\mu^0 \partial_\nu Z_\mu^0 - \frac{1}{2c_w^2} M^2 Z_\mu^0 Z_\mu^0 - \frac{1}{2}\partial_\mu A_\nu \partial_\mu A_\nu - \frac{1}{2}\partial_\mu H \partial_\mu H - \\
 & \frac{1}{2}m_h^2 H^2 - \partial_\mu \phi^+ \partial_\mu \phi^- - M^2 \phi^+ \phi^- - \frac{1}{2}\partial_\mu \phi^0 \partial_\mu \phi^0 - \frac{1}{2c_w} M \phi^0 \phi^0 - \beta_h \left[ \frac{2M^2}{g^2} + \right. \\
 & \left. \frac{2M}{g} H + \frac{1}{2}(H^2 + \phi^0 \phi^0 + 2\phi^+ \phi^-) \right] + \frac{2M^4}{g^2} \alpha_h - ig_{c_w} [\partial_\nu Z_\mu^0 (W_\mu^+ W_\nu^- - \\
 & W_\nu^+ W_\mu^-) - Z_\mu^0 (W_\mu^+ \partial_\nu W_\nu^- - W_\nu^- \partial_\mu W_\mu^+) + Z_\mu^0 (W_\nu^+ \partial_\nu W_\mu^- - \\
 & W_\nu^- \partial_\mu W_\mu^+)] - ig_{s_w} [\partial_\nu A_\mu (W_\mu^+ W_\nu^- - W_\nu^+ W_\mu^-) - A_\nu (W_\mu^+ \partial_\nu W_\mu^- - \\
 & W_\nu^- \partial_\mu W_\mu^+) + A_\mu (W_\nu^+ \partial_\nu W_\mu^- - W_\nu^- \partial_\mu W_\mu^+)] - \frac{1}{2}g^2 W_\mu^+ W_\nu^- W_\nu^+ W_\mu^- + \\
 & \frac{1}{2}g^2 W_\mu^+ W_\nu^- W_\nu^+ W_\mu^- + g^2 c_w^2 (Z_\mu^0 W_\mu^+ Z_\nu^0 W_\nu^- - Z_\mu^0 Z_\nu^0 W_\mu^+ W_\nu^-) + \\
 & g^2 s_w^2 (A_\mu W_\mu^+ A_\nu W_\nu^- - A_\mu A_\nu W_\mu^+ W_\nu^-) + g^2 s_w c_w [A_\mu Z_\nu^0 (W_\mu^+ W_\nu^- - \\
 & W_\nu^+ W_\mu^-) - 2A_\mu Z_\mu^0 W_\nu^+ W_\nu^-] - g\alpha [H^3 + H\phi^0 \phi^0 + 2H\phi^+ \phi^-] - \\
 & \frac{1}{8}g^2 \alpha_h [H^4 + (\phi^0)^4 + 4(\phi^+ \phi^-)^2 + 4(\phi^0)^2 \phi^+ \phi^- + 4H^2 \phi^+ \phi^- + 2(\phi^0)^2 H^2] - \\
 & gM W_\mu^+ W_\nu^- H - \frac{1}{2}g \frac{M}{c_w} Z_\mu^0 Z_\nu^0 H - \frac{1}{2}ig (W_\mu^+ \partial_\nu \phi^0 \phi^- - \phi^- \partial_\mu \phi^0) - \\
 & W_\mu^- (\phi^0 \partial_\mu \phi^+ - \phi^+ \partial_\mu \phi^0) + \frac{1}{2}g [W_\mu^+ (H \partial_\mu \phi^- - \phi^- \partial_\mu H) - W_\mu^- (H \partial_\mu \phi^+ - \\
 & \phi^+ \partial_\mu H)] + \frac{1}{2}g \frac{1}{c_w} (Z_\mu^0 (H \partial_\mu \phi^0 - \phi^0 \partial_\mu H) - ig \frac{s_w^2}{c_w} M Z_\mu^0 (W_\mu^+ \phi^- - W_\mu^- \phi^+) + \\
 & ig_{s_w} M A_\mu (W_\mu^+ \phi^- - W_\mu^- \phi^+) - ig \frac{1-2c_w^2}{2c_w} Z_\mu^0 (\phi^+ \partial_\mu \phi^- - \phi^- \partial_\mu \phi^+) + \\
 & ig_{s_w} A_\mu (\phi^+ \partial_\mu \phi^- - \phi^- \partial_\mu \phi^+) - \frac{1}{4}g^2 W_\mu^+ W_\nu^- [H^2 + (\phi^0)^2 + 2\phi^+ \phi^-] - \\
 & \frac{1}{4}g^2 \frac{1}{c_w} Z_\mu^0 Z_\nu^0 [H^2 + (\phi^0)^2 + 2(2s_w^2 - 1)^2 \phi^+ \phi^-] - \frac{1}{2}g^2 \frac{s_w^2}{c_w} Z_\mu^0 \phi^0 (W_\mu^+ \phi^- + \\
 & W_\mu^- \phi^+) - \frac{1}{2}ig \frac{2s_w^2}{c_w} Z_\mu^0 H (W_\mu^+ \phi^- - W_\mu^- \phi^+) + \frac{1}{2}g^2 s_w A_\mu \phi^0 (W_\mu^+ \phi^- + \\
 & W_\mu^- \phi^+) + \frac{1}{2}ig^2 s_w A_\mu H (W_\mu^+ \phi^- - W_\mu^- \phi^+) - g^2 s_w (2c_w^2 - 1) Z_\mu^0 A_\mu \phi^+ \phi^- - \\
 & g^1 s_w^2 A_\mu A_\nu \phi^+ \phi^- - e^\lambda (\gamma \partial + m_\lambda^2) e^\lambda + \frac{2}{3}(\bar{u}_j^\lambda \gamma^\mu u_j^\lambda) - \frac{1}{3}(\bar{d}_j^\lambda \gamma^\mu d_j^\lambda) + \\
 & \bar{d}_j^\lambda (\gamma \partial + m_d^2) d_j^\lambda + ig_{s_w} A_\mu [-(e^\lambda \gamma^\mu e^\lambda) + \frac{2}{3}(\bar{u}_j^\lambda \gamma^\mu u_j^\lambda) - \frac{1}{3}(\bar{d}_j^\lambda \gamma^\mu d_j^\lambda)] + \\
 & \frac{ig}{4c_w} Z_\mu^0 [(\bar{\nu}^\lambda \gamma^\mu (1 + \gamma^5) \nu^\lambda) + (e^\lambda \gamma^\mu (4s_w^2 - 1 - \gamma^5) e^\lambda) + (\bar{u}_j^\lambda \gamma^\mu (\frac{2}{3}s_w^2 - \\
 & 1 - \gamma^5) u_j^\lambda) + (\bar{d}_j^\lambda \gamma^\mu (1 - \frac{8}{3}s_w^2 - \gamma^5) d_j^\lambda)] + \frac{ig}{2\sqrt{2}} W_\mu^+ [(\bar{\nu}^\lambda \gamma^\mu (1 + \gamma^5) \nu^\lambda) + \\
 & (\bar{u}_j^\lambda \gamma^\mu (1 + \gamma^5) C_{\lambda\kappa} d_j^\kappa)] + \frac{ig}{2\sqrt{2}} W_\mu^- [(\bar{e}^\lambda \gamma^\mu (1 + \gamma^5) \nu^\lambda) + (\bar{d}_j^\kappa C_{\lambda\kappa}^\dagger \gamma^\mu (1 + \\
 & \gamma^5) u_j^\lambda)] + \frac{ig}{2\sqrt{2}} \frac{m_\lambda^2}{M} [-\phi^+ (\bar{\nu}^\lambda (1 - \gamma^5) e^\lambda) + \phi^- (\bar{e}^\lambda (1 + \gamma^5) \nu^\lambda)] - \\
 & \frac{g}{2} \frac{m_\lambda^2}{M} [H (\bar{e}^\lambda e^\lambda) + i\phi^0 (\bar{e}^\lambda \gamma^5 e^\lambda)] + \frac{ig}{2M\sqrt{2}} \phi^+ [-m_d^2 (\bar{u}_j^\lambda C_{\lambda\kappa} (1 - \gamma^5) d_j^\kappa) + \\
 & m_u^2 (\bar{u}_j^\lambda C_{\lambda\kappa} (1 + \gamma^5) d_j^\kappa) + \frac{ig}{2M\sqrt{2}} \phi^- [m_d^2 (\bar{d}_j^\kappa C_{\lambda\kappa}^\dagger (1 + \gamma^5) u_j^\kappa) - m_u^2 (\bar{d}_j^\kappa C_{\lambda\kappa}^\dagger (1 - \\
 & \gamma^5) u_j^\kappa) - \frac{g}{2} \frac{m_\lambda^2}{M} H (\bar{u}_j^\lambda u_j^\lambda) - \frac{g}{2} \frac{m_\lambda^2}{M} H (\bar{d}_j^\lambda d_j^\lambda) + \frac{ig}{2} \frac{m_\lambda^2}{M} \phi^0 (\bar{u}_j^\lambda \gamma^5 u_j^\lambda) - \\
 & \frac{ig}{2} \frac{m_\lambda^2}{M} \phi^0 (\bar{d}_j^\lambda \gamma^5 d_j^\lambda) + \bar{X}^+ (\partial^2 - M^2) X^+ + \bar{X}^- (\partial^2 - M^2) X^- + X^0 (\partial^2 - \\
 & \frac{M^2}{c_w^2}) X^0 + \bar{Y} \partial^2 Y + ig_{c_w} W_\mu^+ (\partial_\mu \bar{X}^0 X^- - \partial_\mu \bar{X}^+ X^0) + ig_{s_w} W_\mu^+ (\partial_\mu \bar{X}^- X^0 - \\
 & \partial_\mu \bar{X}^+ X^0) + ig_{c_w} W_\mu^- (\partial_\mu \bar{X}^- X^0 - \partial_\mu \bar{X}^0 X^+) + ig_{s_w} W_\mu^- (\partial_\mu \bar{X}^- X^0 - \\
 & \partial_\mu \bar{X}^+ X^0) + ig_{c_w} Z_\mu^0 (\partial_\mu \bar{X}^+ X^+ - \partial_\mu \bar{X}^- X^-) + ig_{s_w} A_\mu (\partial_\mu \bar{X}^+ X^+ - \\
 & \partial_\mu \bar{X}^- X^-) - \frac{1}{2}gM [\bar{X}^+ X^+ H + \bar{X}^- X^- H + \frac{1}{c_w} \bar{X}^0 X^0 H] + \\
 & \frac{1-2c_w^2}{2c_w} igM [\bar{X}^+ X^0 \phi^+ - \bar{X}^- X^0 \phi^-] + \frac{1}{2c_w} igM [\bar{X}^0 X^- \phi^+ - \bar{X}^0 X^+ \phi^-] + \\
 & igMs_w [\bar{X}^0 X^- \phi^+ - \bar{X}^0 X^+ \phi^-] + \frac{1}{2}igM [\bar{X}^+ X^+ \phi^0 - \bar{X}^- X^- \phi^0]
 \end{aligned}$$



# The Standard Model of particle physics...



	three generations of matter (fermions)			interactions / force carriers (bosons)	
	I	II	III		
mass	$= 2.2 \text{ MeV}/c^2$	$= 1.28 \text{ GeV}/c^2$	$= 173.1 \text{ GeV}/c^2$	0	$= 125.09 \text{ GeV}/c^2$
charge	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	0	0
spin	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	0	0
	<b>u</b> up	<b>c</b> charm	<b>t</b> top	<b>g</b> gluon	<b>H</b> higgs
	<b>d</b> down	<b>s</b> strange	<b>b</b> bottom	<b><math>\gamma</math></b> photon	
	<b>e</b> electron	<b><math>\mu</math></b> muon	<b><math>\tau</math></b> tau	<b>Z</b> Z boson	
	<b><math>\nu_e</math></b> electron neutrino	<b><math>\nu_\mu</math></b> muon neutrino	<b><math>\nu_\tau</math></b> tau neutrino	<b>W</b> W boson	

**QUARKS** (left side)

**LEPTONS** (left side)

**GAUGE BOSONS VECTOR BOSONS** (right side)

**SCALAR BOSONS** (right side)

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i\bar{\Psi} \not{D} \psi + D_\mu \Phi^\dagger D^\mu \Phi - V(\Phi) + \bar{\Psi}_L \hat{Y} \Phi \Psi_R + h.c.$$

Gauge bosons

Gauge boson coupling to fermions (EW, QCD)

Higgs coupling to fermions (fermion masses)

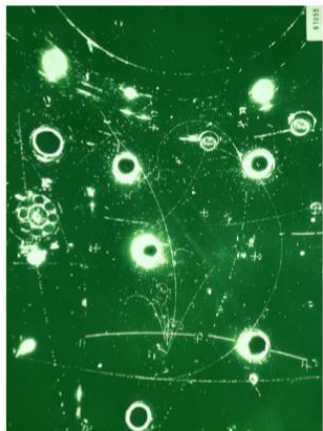
Higgs coupling to bosons (boson masses)

Higgs self-coupling (Higgs potential)

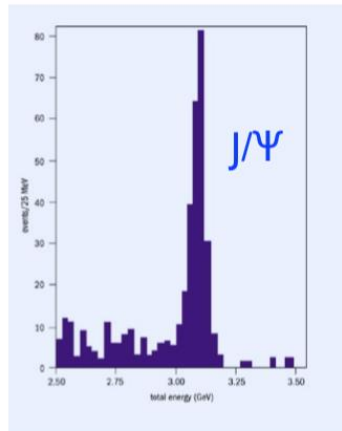


# A theory built (and probed) over time...

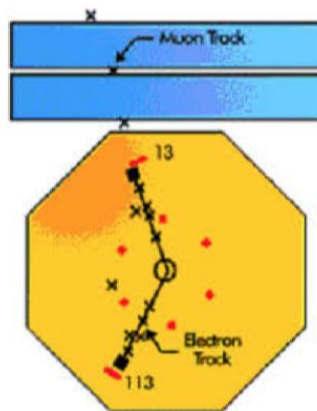
1972 – CERN  
Neutral currents



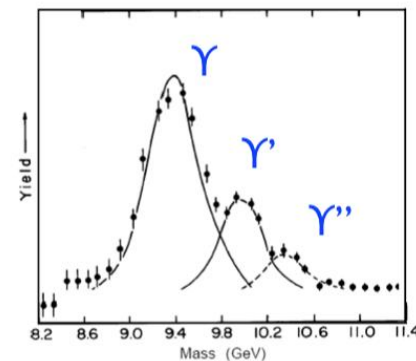
1974 – BNL, SLAC  
Charm



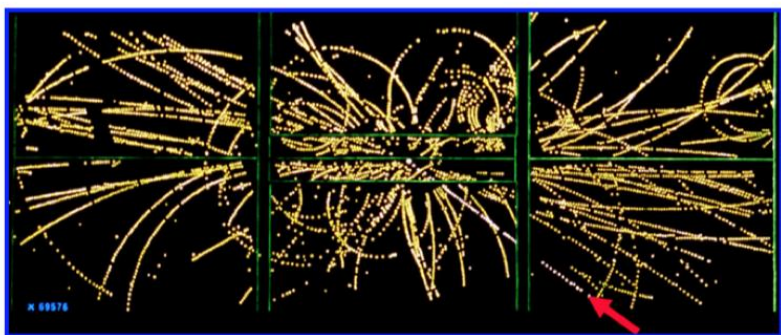
1976 – SLAC  
Tau lepton



1979 – Fermilab  
Beauty

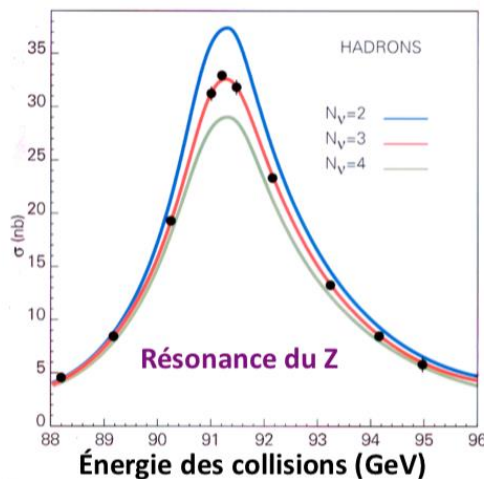


1983 – CERN/SppS  
W and Z bosons



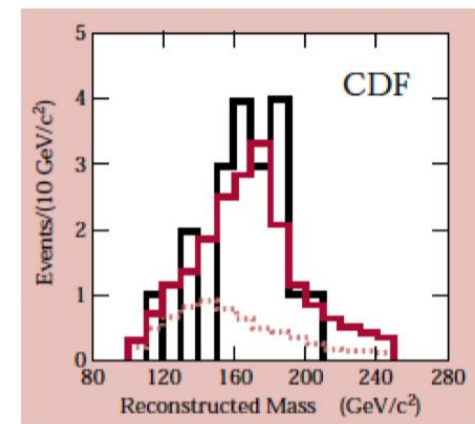
UA1, UA2

1990 – CERN/LEP  
Three families of neutrinos



ALEPH, DEPHI, L3, OPAL  
(experimental) LHC physics

1994 – Fermilab/TeVatron  
Top quark



CDF, D0

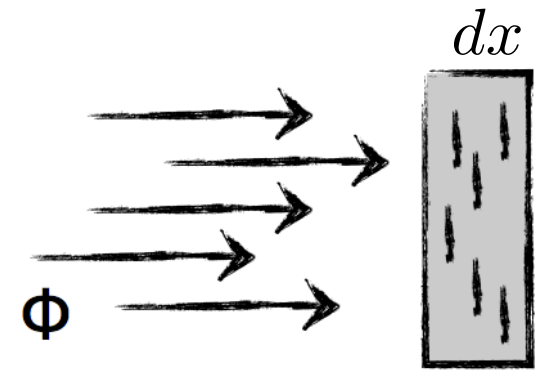


# How do we compare experiment and predictions in a **quantum** field theory?

- Through two fundamental quantities:
- $\sigma$  (**cross section**): **probability** of a particle of **being produced** in collisions **at a given energy** (es. 13 TeV at LHC)
  - ✓ May be *differential*, that is, as a function of the energy of the particle, the angles of its trajectory, etc.
- $\Gamma$  (decay rate): probability of a particle of decaying into certain specific final particles
  - ✓ The sum of all  $\Gamma$ 's is the total decay rate, and because of resonance theory it is the inverse of its decay time:  $\tau = 1/\Gamma$

# Interaction cross section

Flux  $\Phi = \frac{1}{S} \frac{dN_i}{dt}$   $[L^{-2} t^{-1}]$



area obscured by target particle

Reactions  
per unit of  
time

$$d\dot{N}_{\text{reac}} = \Phi \underbrace{\sigma}_{[?]} N_{\text{target}} dx \quad [t^{-1}]$$

$[L^{-2} t^{-1}]$      $[L^{-1}]$      $[L]$

Reaction rate  
per target particle

$$W_{if} = \underbrace{\Phi \sigma}_{[t^{-1}]}$$

Cross section  
per target  
particle

$$\sigma = \frac{W_{if}}{\Phi} \quad [L^2] \quad = \text{reaction rate per unit of flux}$$

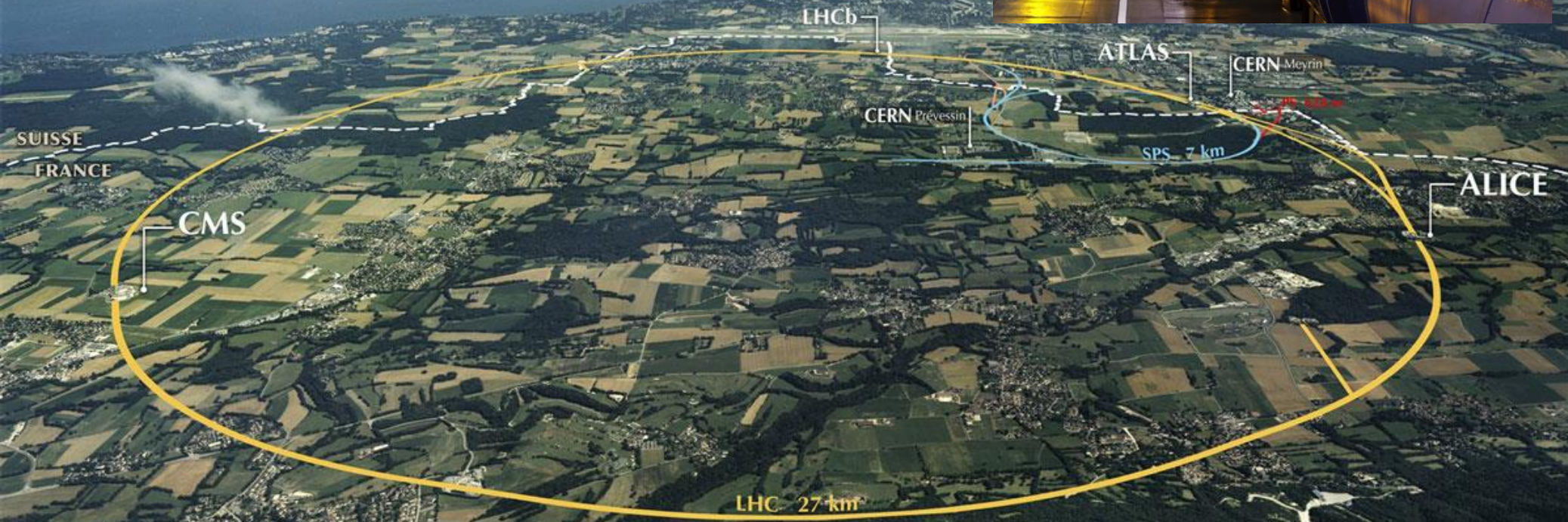
1 b =  $10^{-28}$  m<sup>2</sup> (roughly the area of a nucleus with A = 100)



# LHC

$pp$  collider (2008-present)

$\sqrt{s} = 7-8-13$  TeV



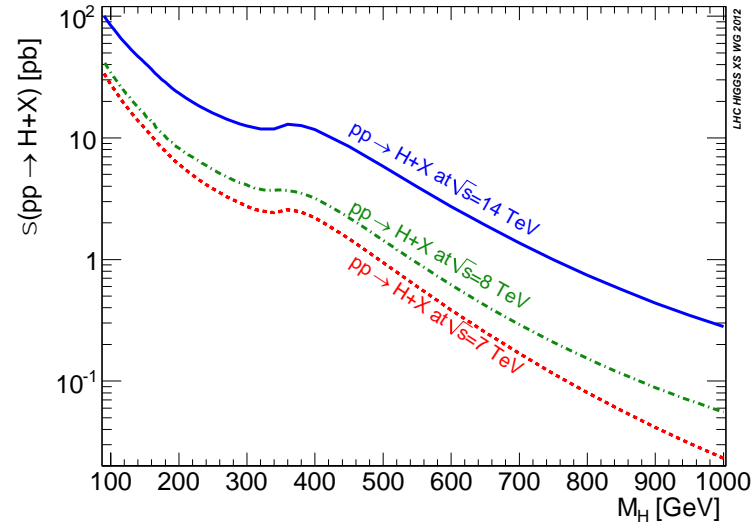


# Luminosity

Number of events  
in unit of time

$$\dot{N} = \mathcal{L} \cdot \sigma$$

$[\text{t}^{-1}]$ 
 $[\text{L}^{-2} \text{t}^{-1}]$ 
 $[\text{L}^2]$



In a collider ring...

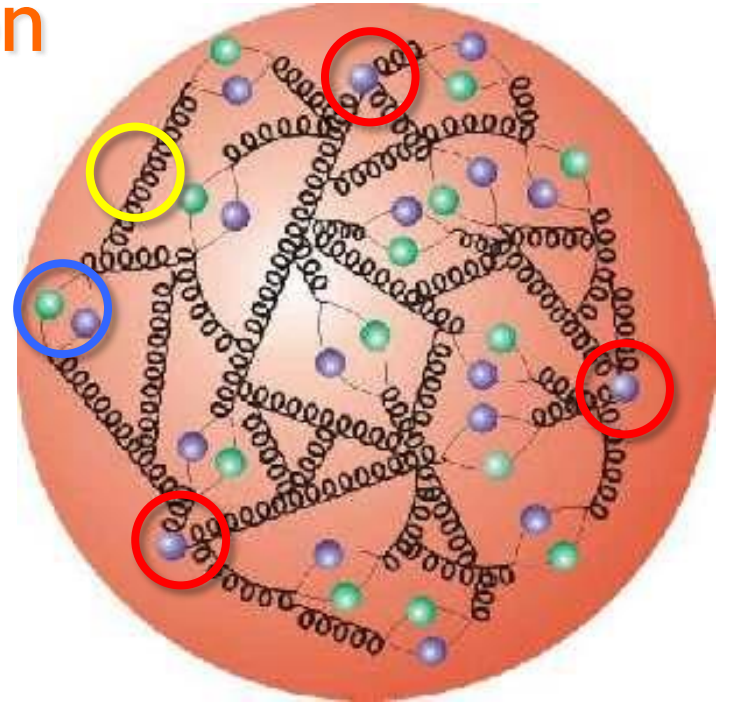
$$\mathcal{L} = \frac{1}{4\pi} \frac{fkN_1N_2}{\sigma_x\sigma_y}$$

Current  
  
 Beam sizes (RMS)

# About the inner life of a proton

- **Protons have substructure!**

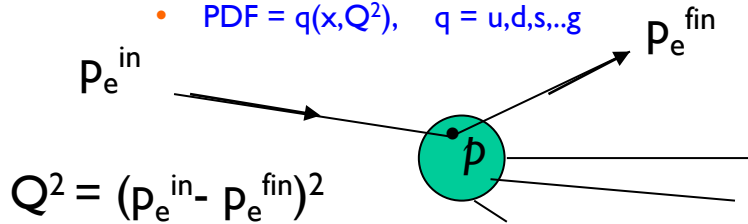
- ✓ partons = quarks & gluons
- ✓ 3 valence (colored) quarks bound by gluons
- ✓ Gluons (colored) have self-interactions
- ✓ Virtual quark pairs can pop-up (sea-quark)
- ✓  $p$  momentum shared among constituents
  - described by  $p$  structure functions



- **Parton energy not 'monochromatic'**

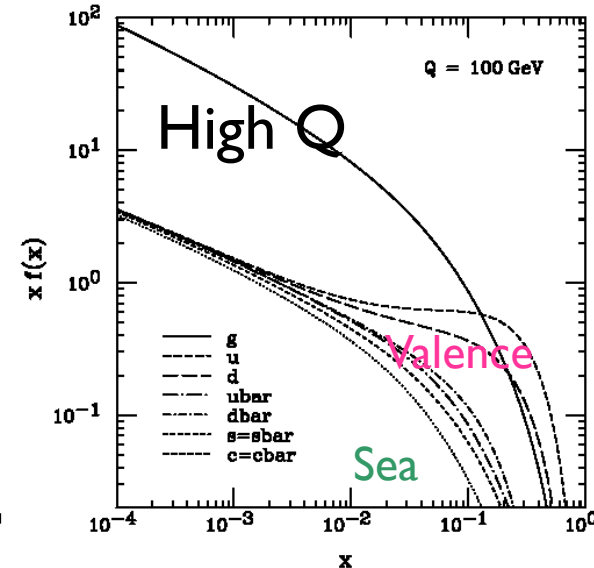
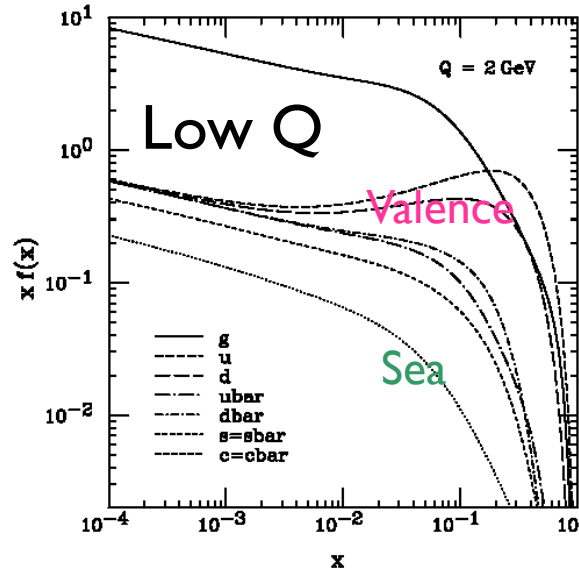
- ✓ Parton Distribution Function

- PDF =  $q(x, Q^2)$ ,  $q = u, d, s, \dots, g$

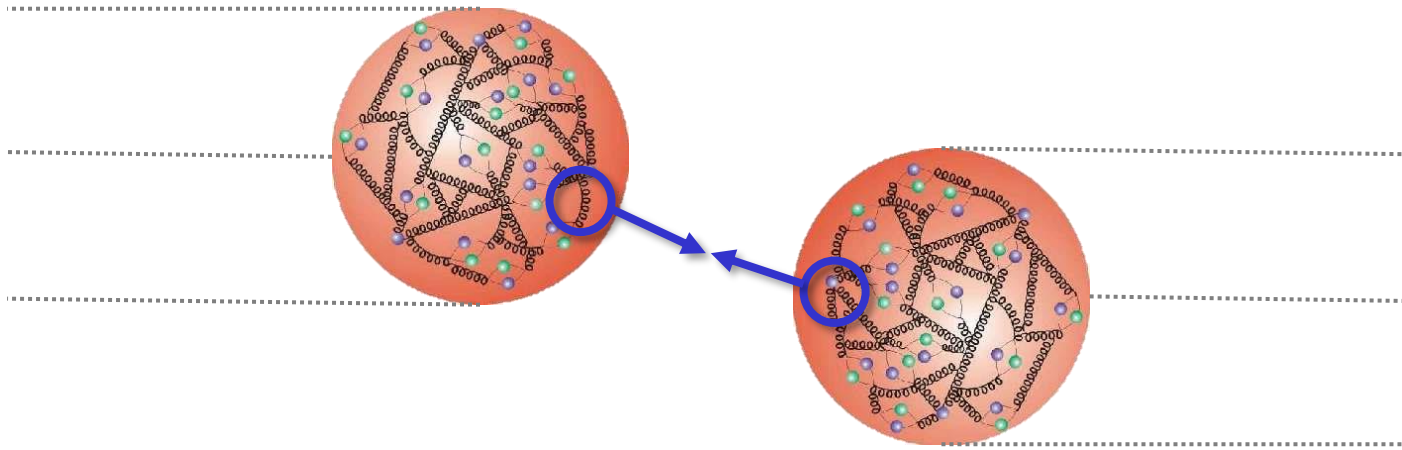


- **Kinematic variables**

- ✓ Bjorken- $x$ : fraction of the proton momentum carried by struck parton
  - $x = p_{\text{parton}}/p_{\text{proton}}$
- ✓  $Q^2$ : 4-momentum<sup>2</sup> transfer

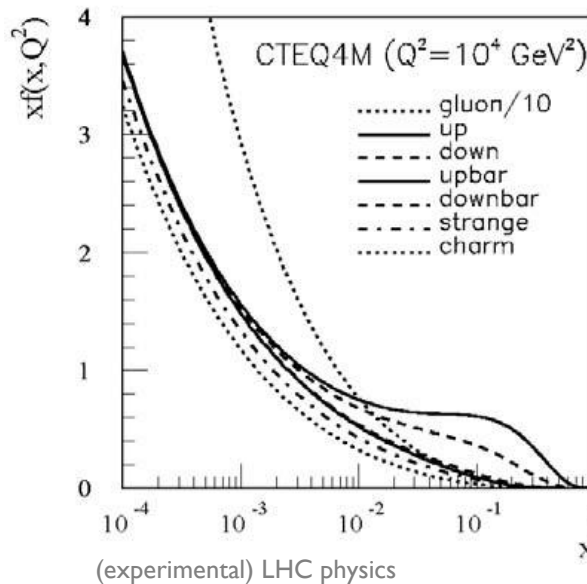
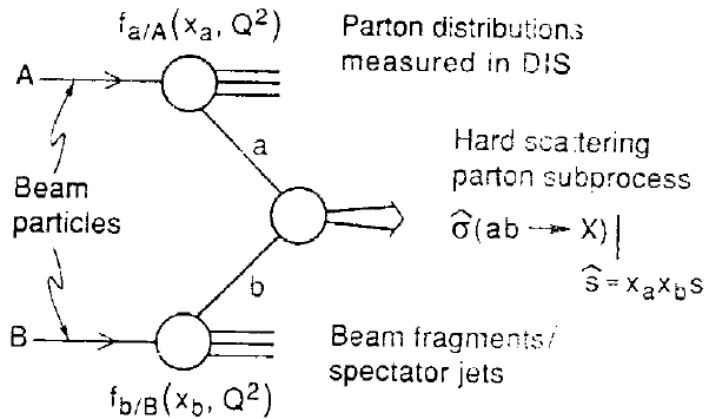


# Cross sections at a proton-proton collider



$$\sqrt{\hat{s}} = \sqrt{x_a x_b S}$$

$$\sigma = \sum_{a,b} \int dx_a dx_b f_a(x, Q^2) f_b(x, Q^2) \hat{\sigma}_{ab}(x_a, x_b)$$



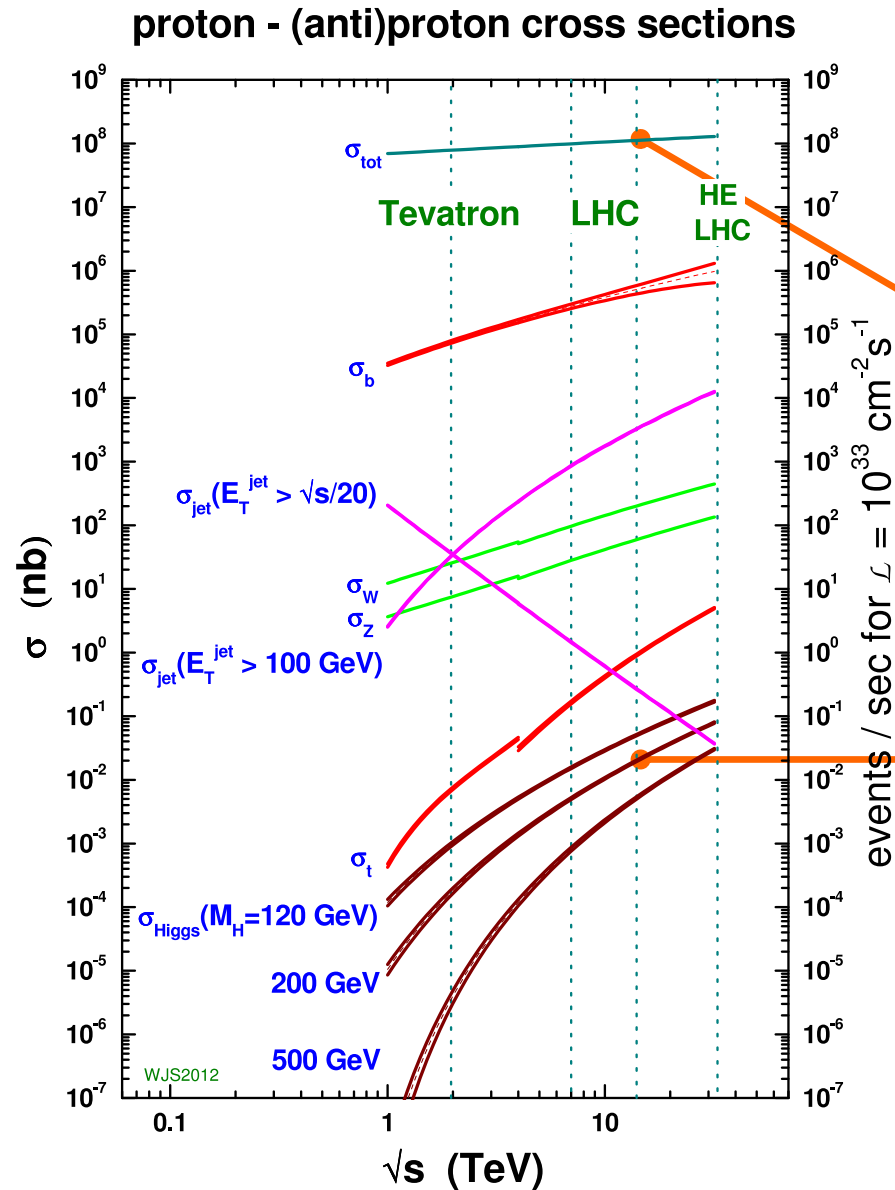
Example: to produce a particle with mass  $m = 100 \text{ GeV}$

$$\sqrt{\hat{s}} = 100 \text{ GeV}$$

$$\sqrt{s} = 14 \text{ TeV} \rightarrow x_a x_b = 0.007$$



# Cross-sections at LHC



$$1 \text{ nb} = 10^{-33} \text{ cm}^2$$

$$\sigma_{\text{tot}}(13 \text{ TeV}) = 10^8 \text{ nb}$$

$$\sigma_H(13 \text{ TeV}) = 0.05 \text{ nb}$$

$$\text{LHC instantaneous luminosity } \mathcal{L} = 1 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$$

inelastic  $pp$  collisions

$10^9$  events/s

$\sim 10^{10}$

$10^{-1}$  events/s

$\sim 1$  Higgs boson  
every 2 seconds

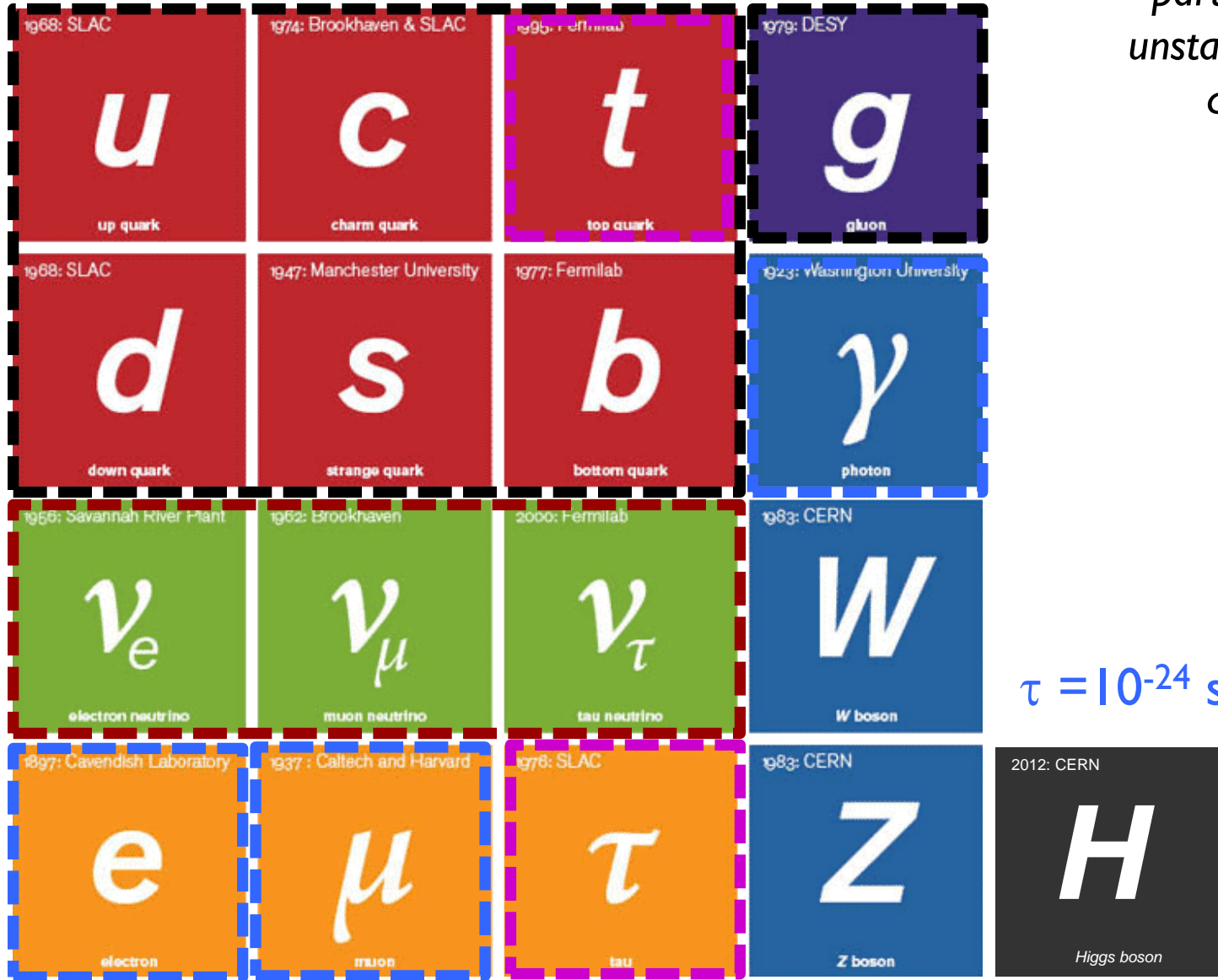
$[m_H \sim 125 \text{ GeV}]$

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  - ✓ May be *differential*, that is, as a function of the energy of the particle, the angles of its trajectory, etc.
- $\Gamma$  (decay rate): **probability** of a particle of **decaying into certain specific final particles**
  - ✓ The sum of all  $\Gamma$ 's is the **total decay rate**, and because of **resonance theory** it is the inverse of its **decay time**:  $\tau = 1/\Gamma$

# What do we want to measure?

... “stable”  
particles from  
unstable particle  
decays!



$\tau = \infty$

$\tau = 10^{-24} \text{ s}$

$\tau = 2.2 \mu\text{s}$

(experimental) LHC physics



# What do we want to measure?

decays?

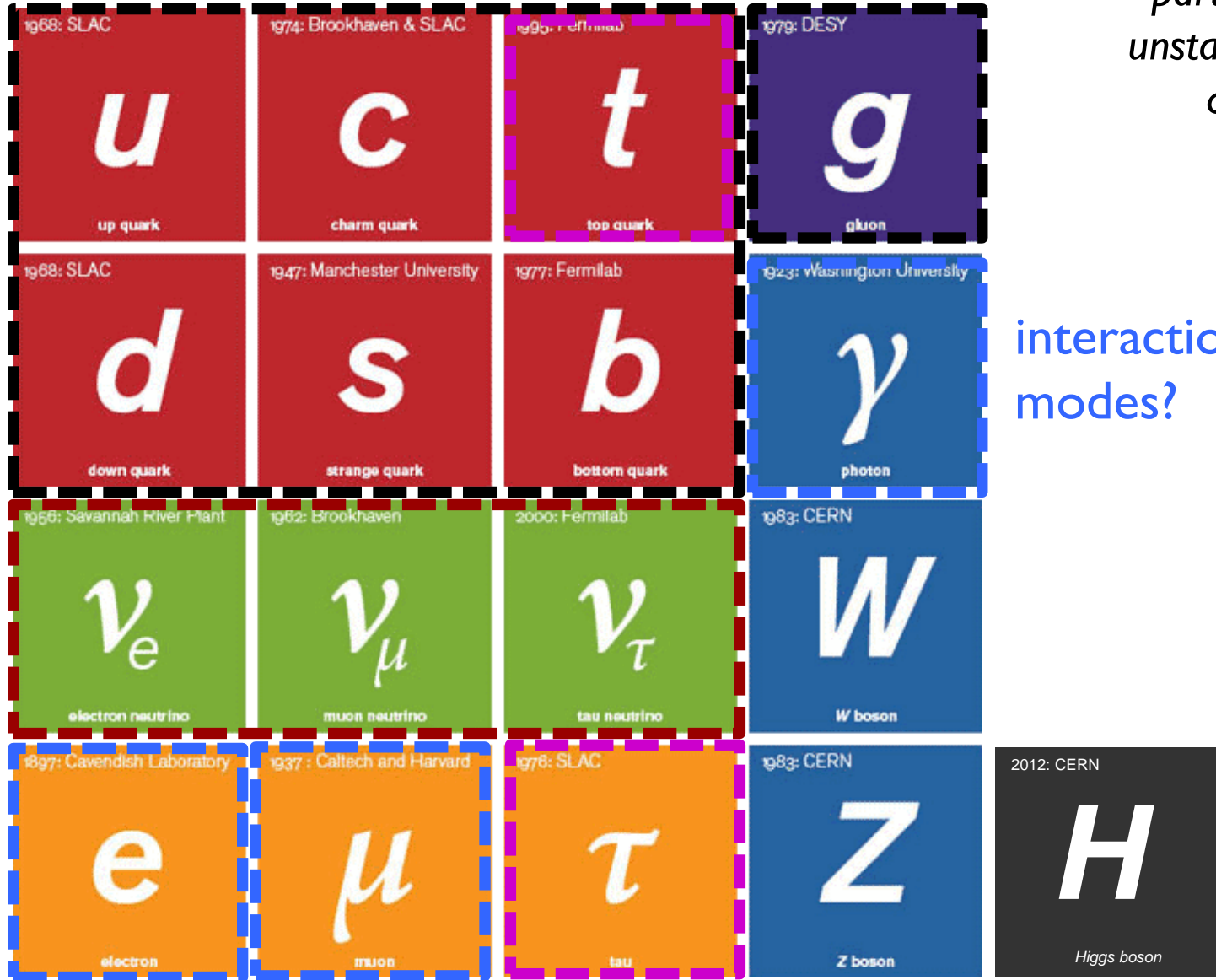
... “stable”  
particles from  
unstable particle  
decays!

hadron  
jets

interaction  
modes?

invisible  
*in particle  
detectors at  
accelerators*

interaction  
modes?



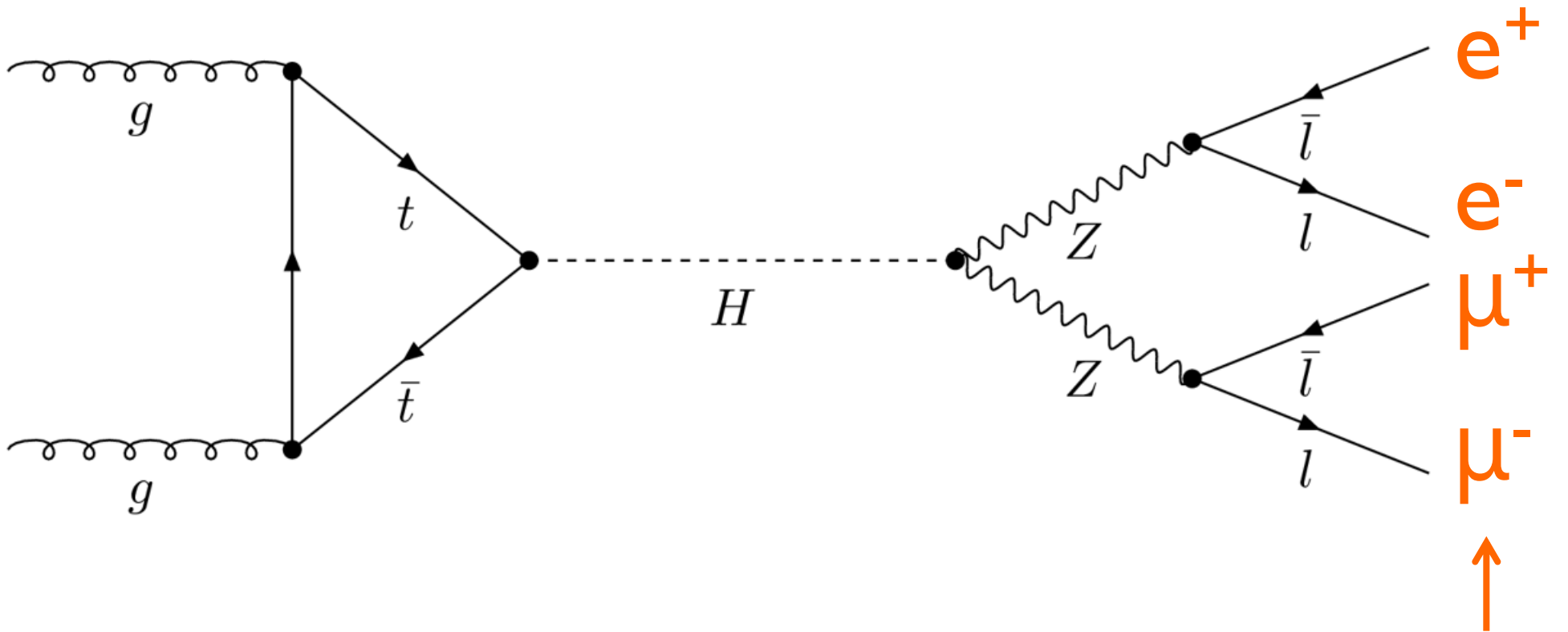
decays?

# What do we want to measure?

Example: let's assume a Higgs boson is produced at the LHC ...

It is a **SM particle**, we **can predict** how and how frequently

... we look for “stable” particles from an unstable particle decays



this is what we are looking for...

# Identifying and measuring “stable” particles

- Particles are characterized by
  - ✓ **Mass** [Unit: eV/c<sup>2</sup> or eV]
  - ✓ **Charge** [Unit: e]
  - ✓ **Energy** [Unit: eV]
  - ✓ **Momentum** [Unit: eV/c or eV]
  - ✓ (+ spin, lifetime, ...)

Particle identification via measurement of:

e.g. (E, p, Q) or (p, β, Q)  
(p, m, Q) ...

- ... and move at **relativistic speed** (here in “natural” unit:  $\hbar = c = 1$ )

$$\beta = \frac{v}{c} \quad \gamma = \frac{1}{\sqrt{1 - \beta^2}}$$

$$l = \frac{l_0}{\gamma} \quad \text{length contraction}$$

$$t = t_0 \gamma \quad \text{time dilation}$$

$$E^2 = \vec{p}^2 + m^2$$
$$E = m\gamma \quad \vec{p} = m\gamma\vec{\beta}$$
$$\vec{\beta} = \frac{\vec{p}}{E}$$



# Center of mass energy

- In the **center-of-mass frame** the total momentum is 0
- In **laboratory frame**, the center of mass energy can be computed as:

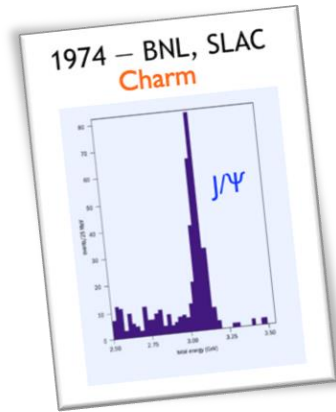
$$E_{\text{cm}} = \sqrt{s} = \sqrt{\left(\sum E_i\right)^2 - \left(\sum \vec{p}_i\right)^2}$$

Hint: it can be computed as the “length” of the total four-momentum, that is invariant:

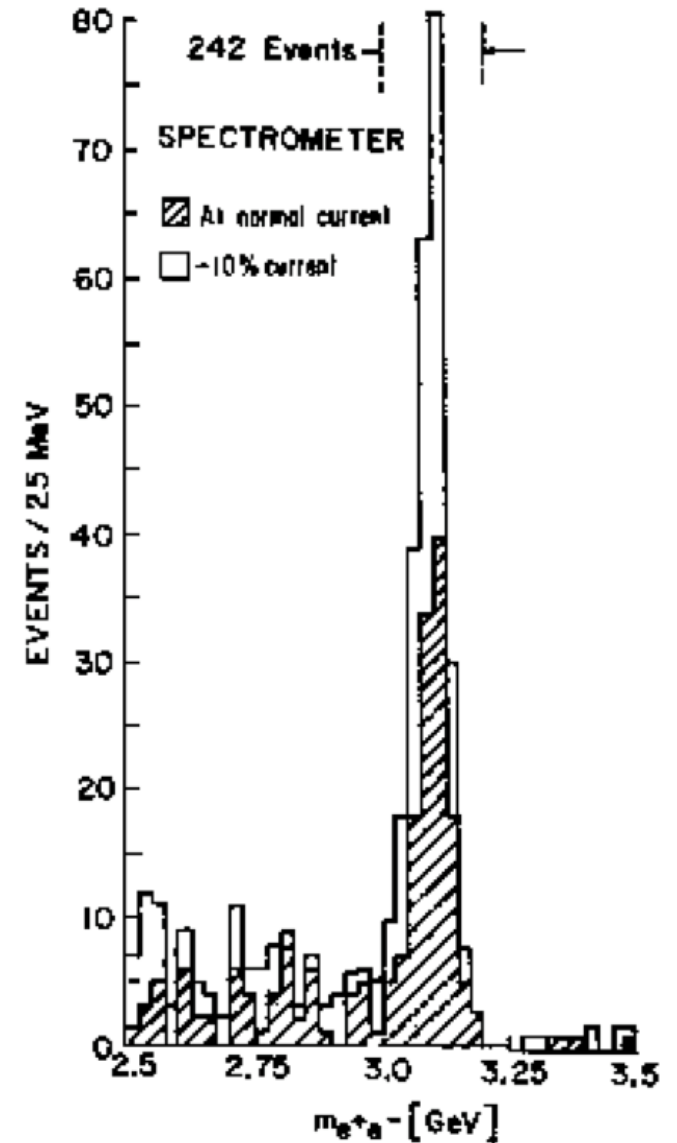
$$p = (E, \vec{p}) \quad \sqrt{p \cdot p}$$

What is the “length” of a the four-momentum of a particle?

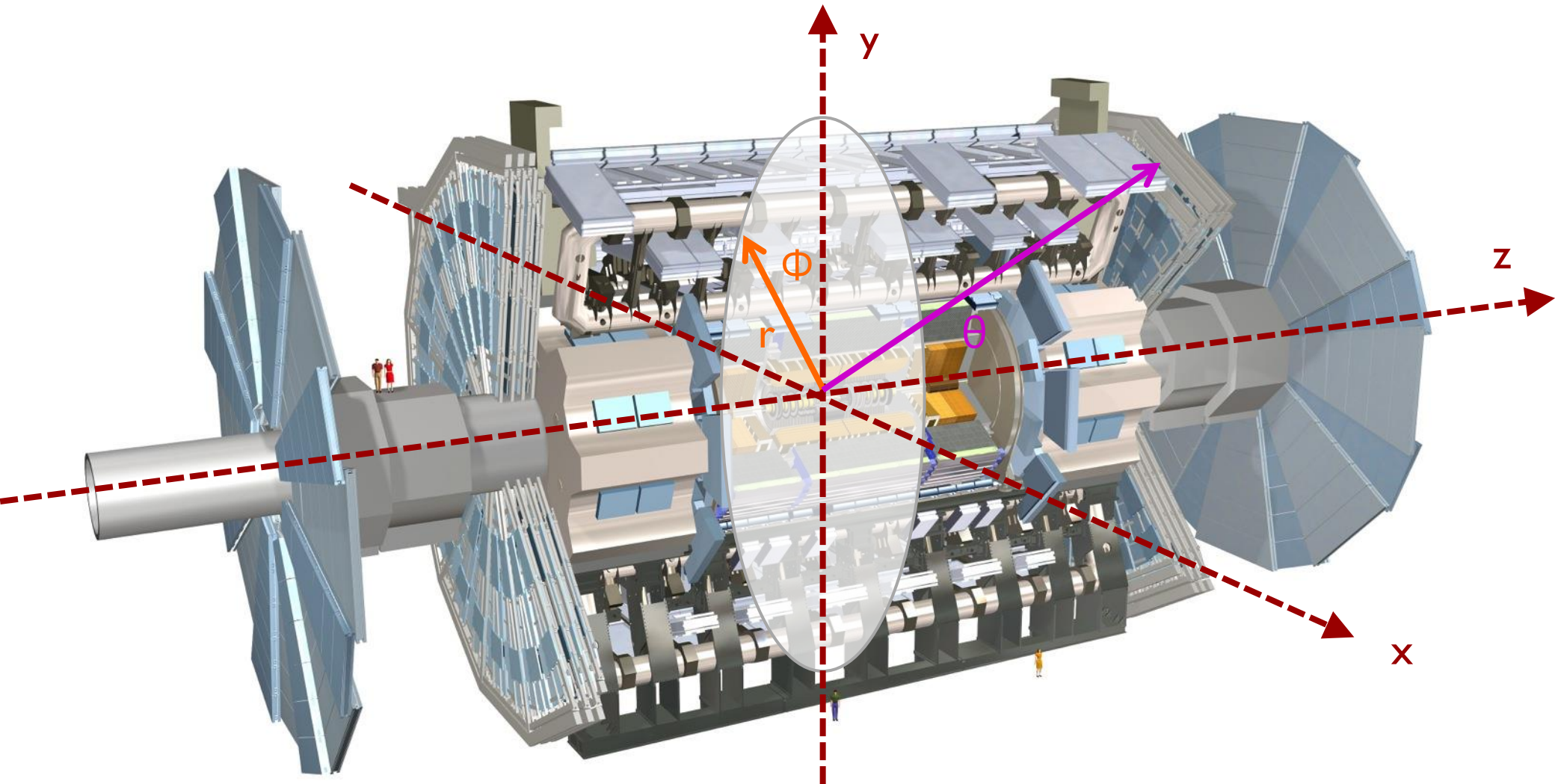
# Invariant mass



$$M = \sqrt{\left(\sum E_i\right)^2 - \left(\sum \vec{p}_i\right)^2}$$



# A collider experiment

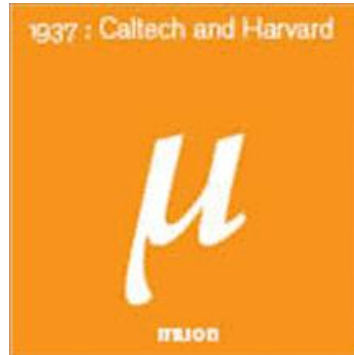




# Interaction mode cheat sheet (“light” particles)



- electrically charged
- ionization ( $dE/dx$ )
- *electromagnetic shower...*



- electrically charged
- ionization ( $dE/dx$ )
- can emit photons
  - ✓ electromagnetic shower induced by emitted photon...
  - but it's rare...



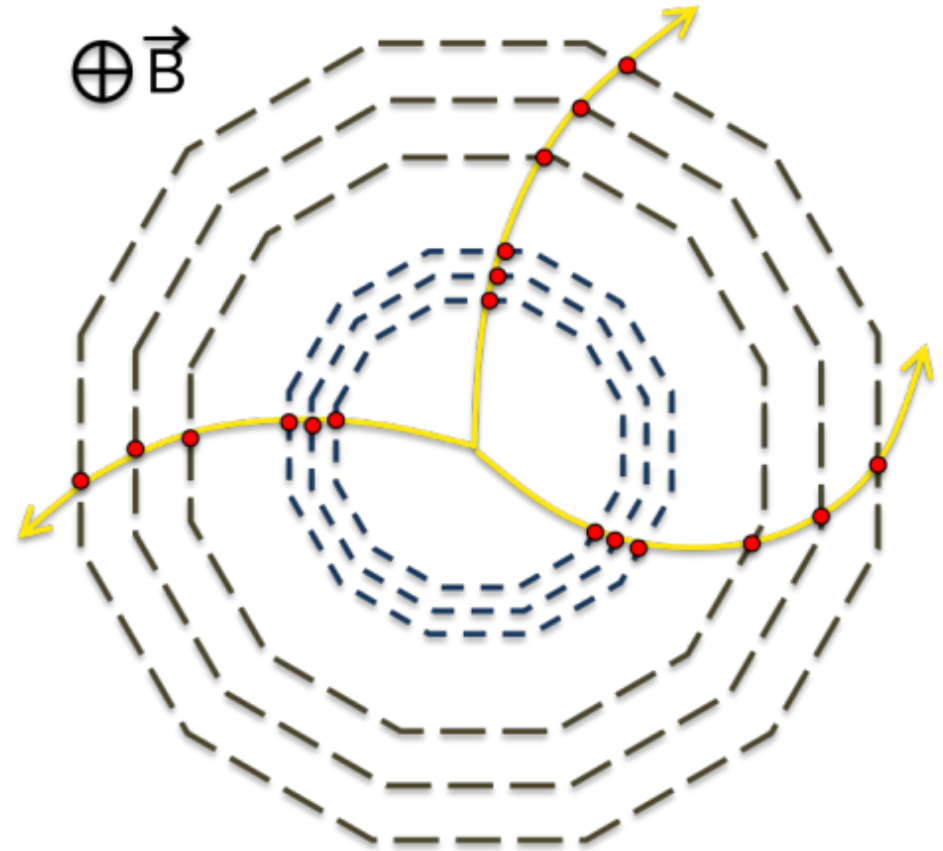
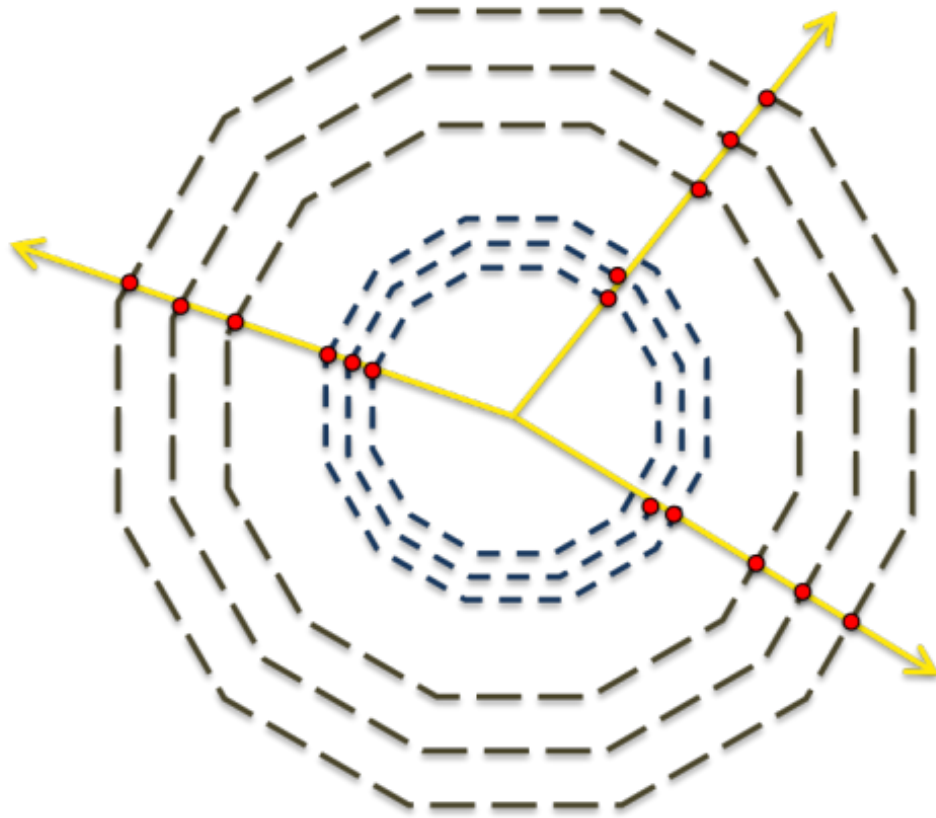
- electrically neutral
- pair production
  - ✓  $E > 1 \text{ MeV}$
- *electromagnetic shower...*



- produce *hadron(s)* jets via QCD hadronization process

# Magnetic spectrometer for ionizing particles

- A system to measure (charged) particle momentum
- Tracking device + magnetic field



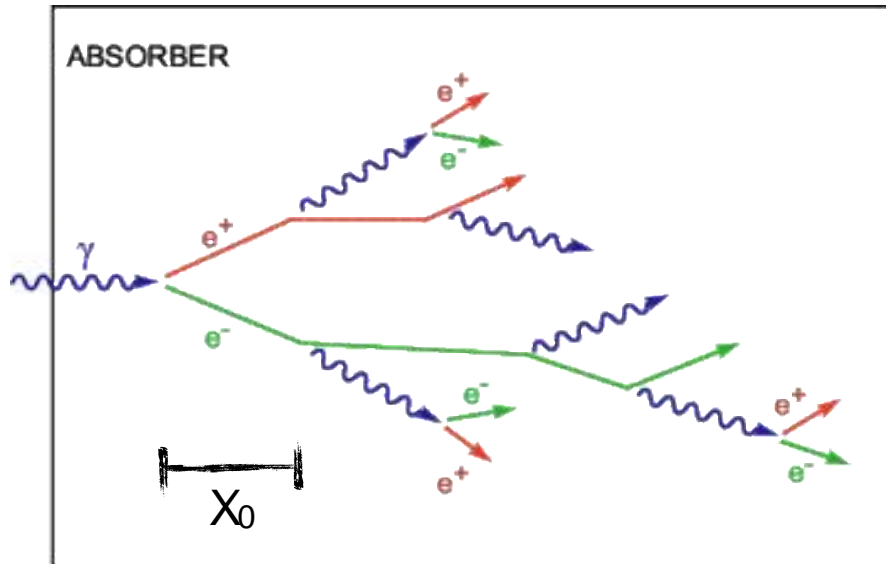
$$p[\text{GeV}] = 0.3B[\text{T}]\rho[\text{m}]$$

# Calorimeters for showering particles

- Electromagnetic shower

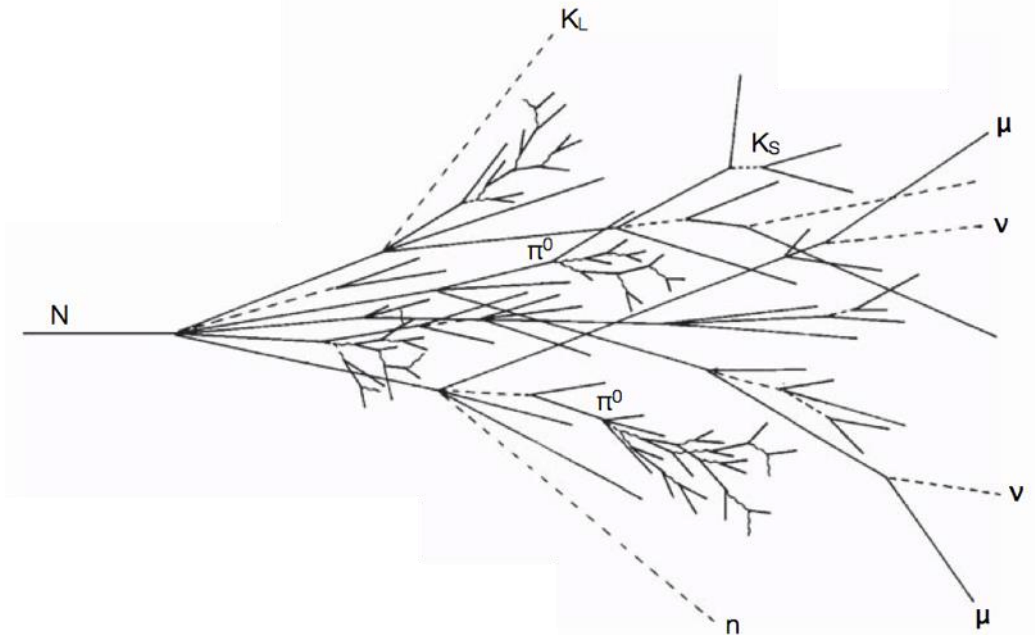
- ✓ Photons: pair production
  - Until below  $e^+e^-$  threshold
- ✓ Electrons: bremsstrahlung
  - Until brem cross-section smaller than ionization

$$\left. \frac{dE}{dx}(E_c) \right|_{\text{Brems}} = \left. \frac{dE}{dx}(E_c) \right|_{\text{Ion}}$$

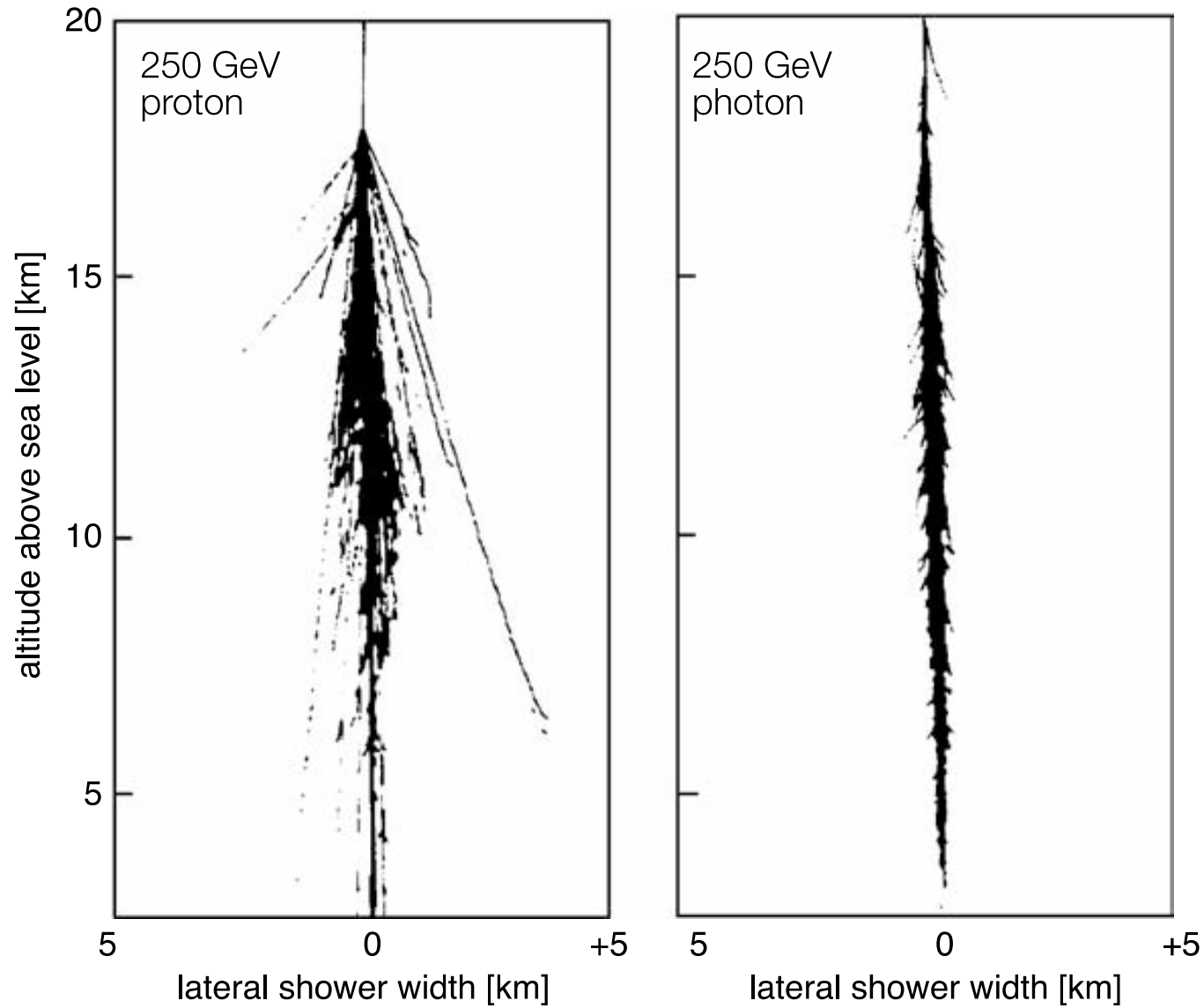


- Hadronic showers

- ✓ Inelastic scattering w/ nuclei
  - Further inelastic scattering until below pion production threshold
- ✓ Sequential decays
  - $\pi^0 \rightarrow \gamma\gamma$
  - Fission fragment:  $\beta$ -decay,  $\gamma$ -decay
  - Neutron capture, spallation, ...

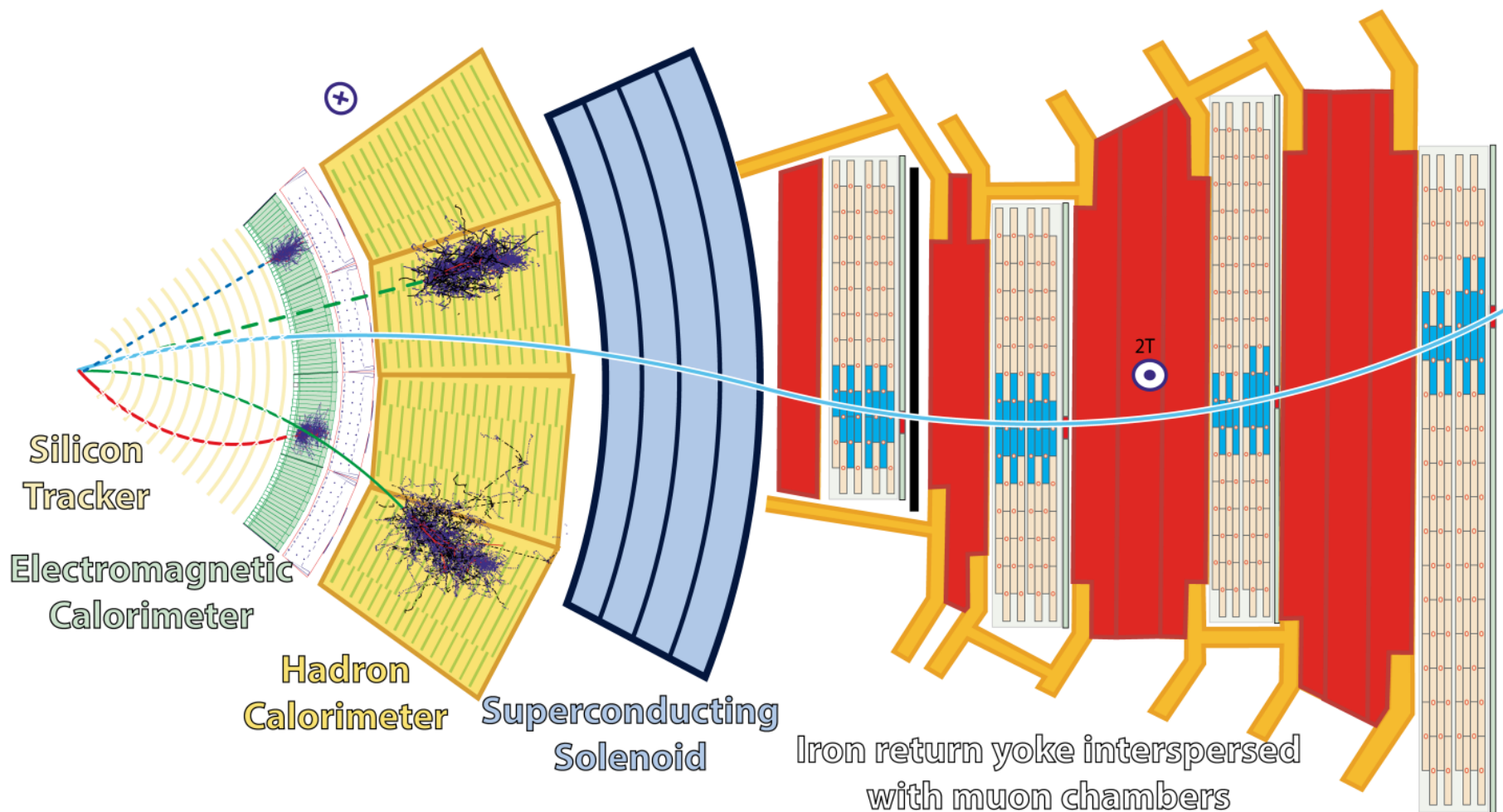


# Hadronic vs. EM showers





# Particle identification with CMS@LHC



— Muon

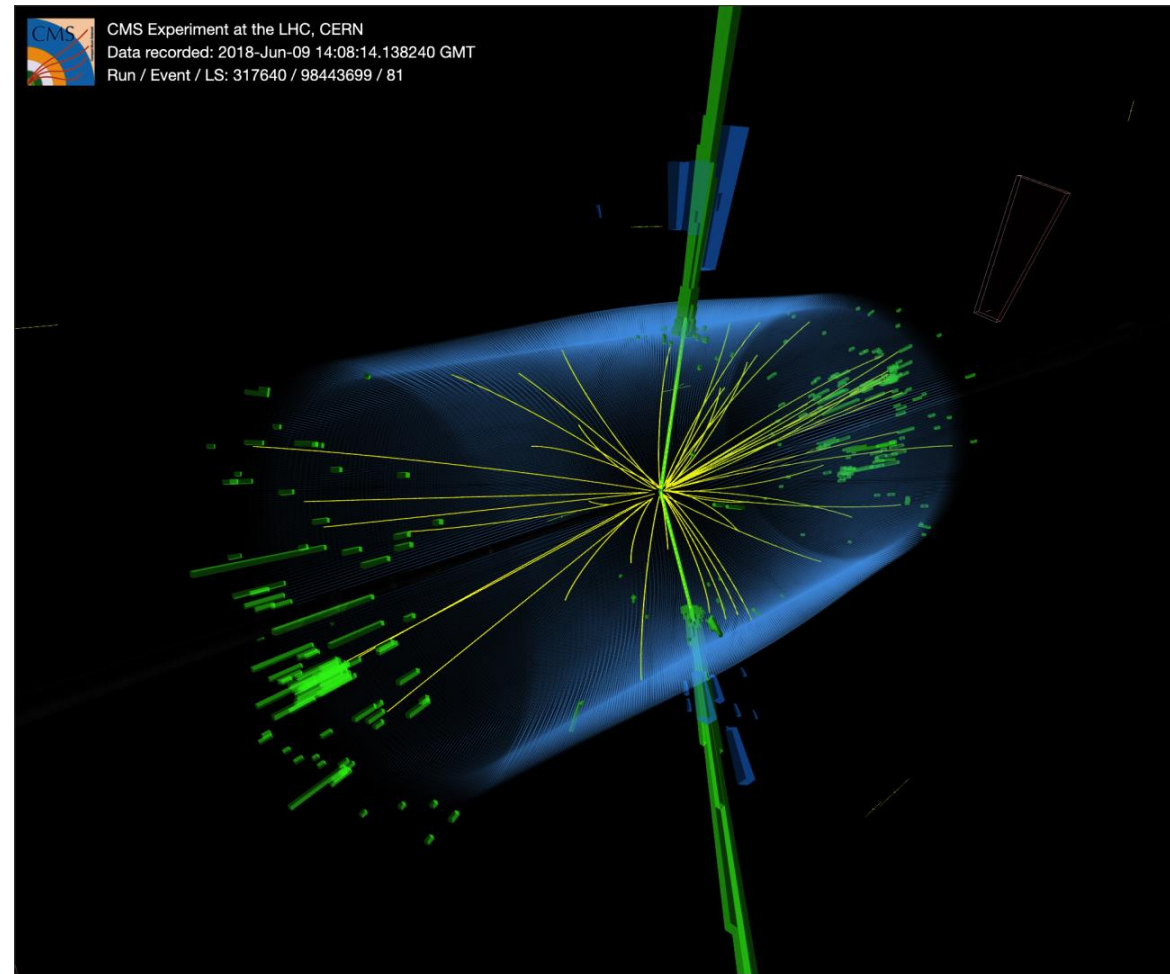
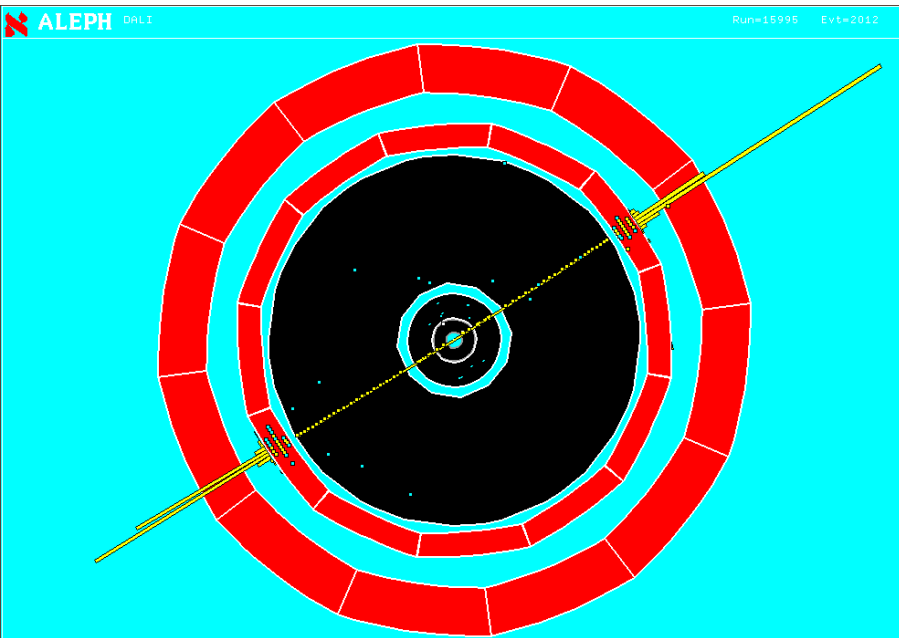
— Electron

— Charged hadron (e.g. pion)

- - - Neutral hadron (e.g. neutron)

- - - Photon

# A $Z \rightarrow e^+e^-$ event at LEP and ad LHC



ALEPH @ LEP

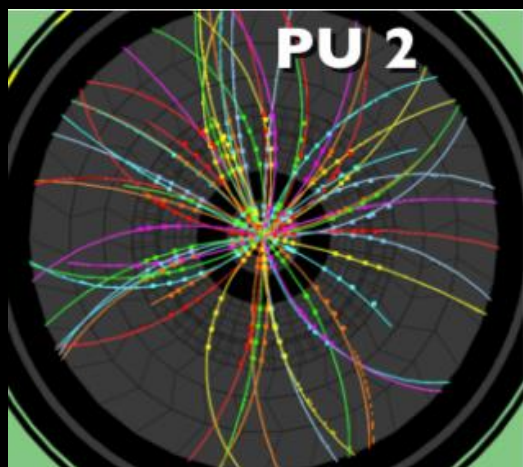
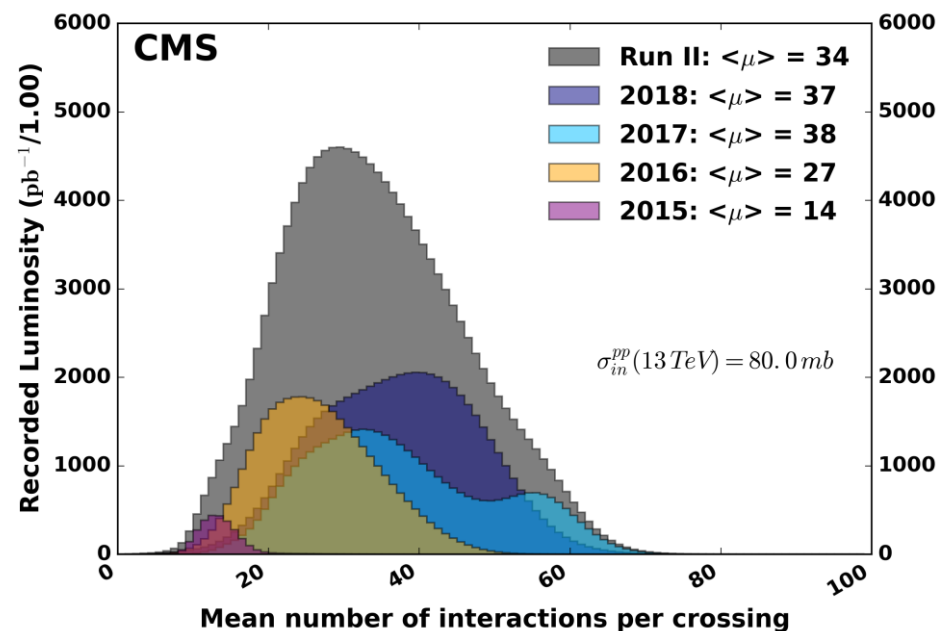
CMS @ LHC

# Pile-Up

$$\mathcal{L} = \frac{1}{4\pi} \frac{fk N_1 N_2}{\sigma_x \sigma_y}$$

PU = number of inelastic interactions per beam bunch crossing

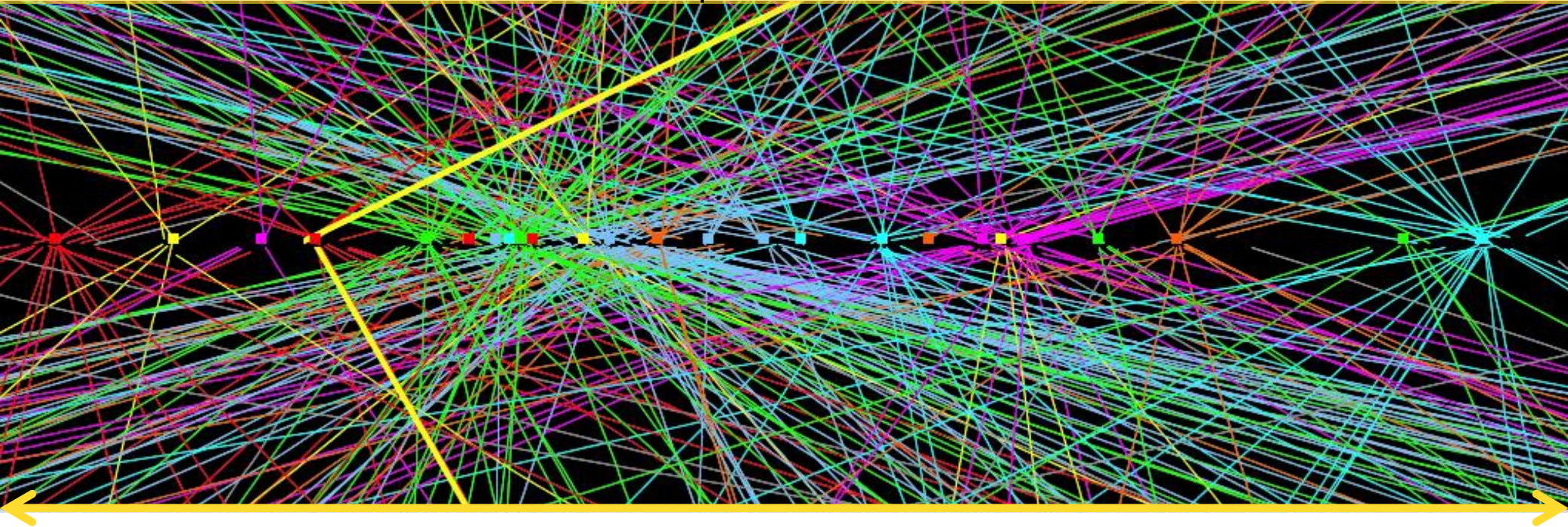
CMS Average Pileup (pp,  $\sqrt{s}=13$  TeV)





# $Z \rightarrow \mu\mu$ event with 25 reconstructed vertices

April 15<sup>th</sup>, 2012



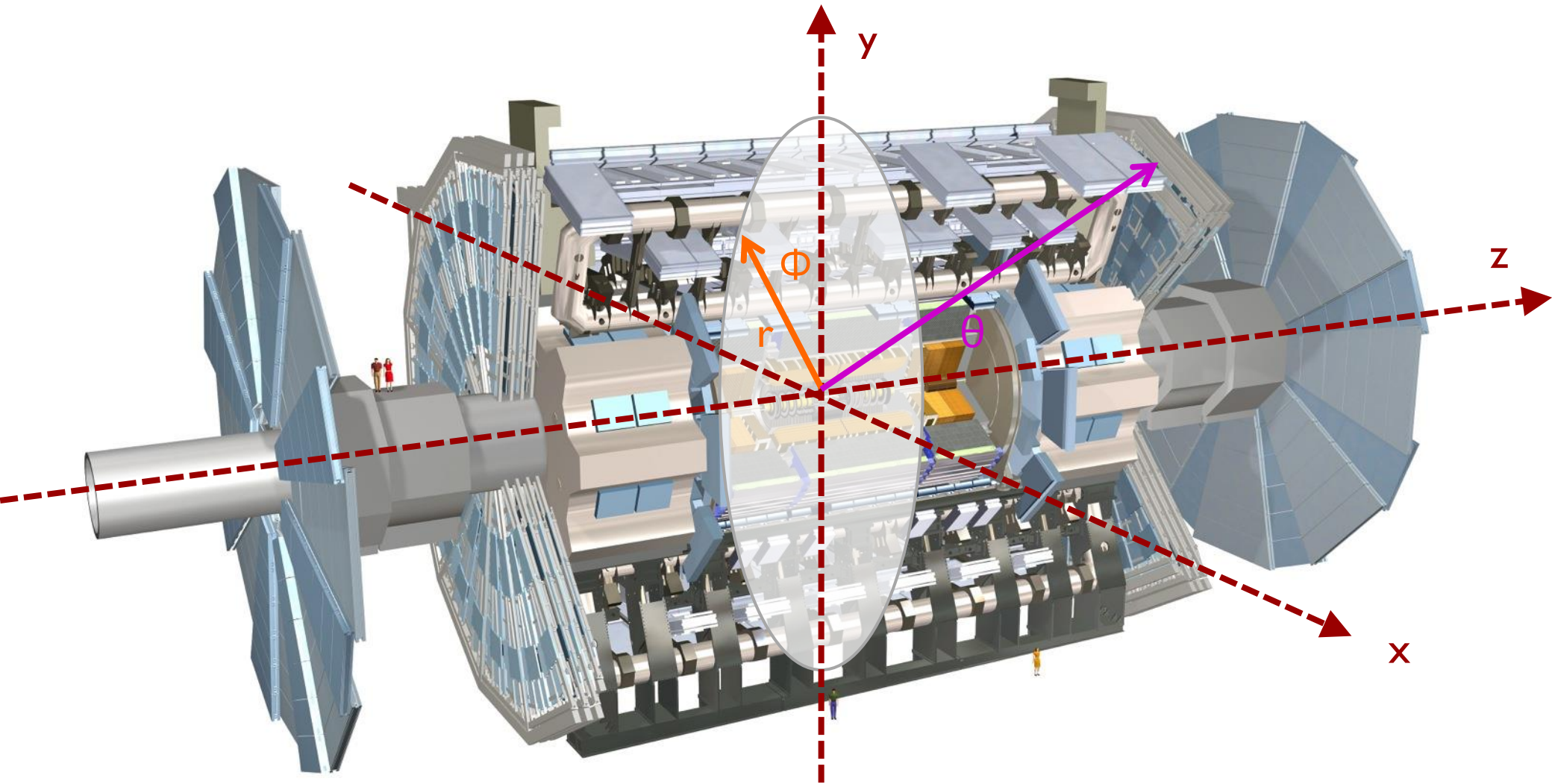
~5 cm



# Additional information

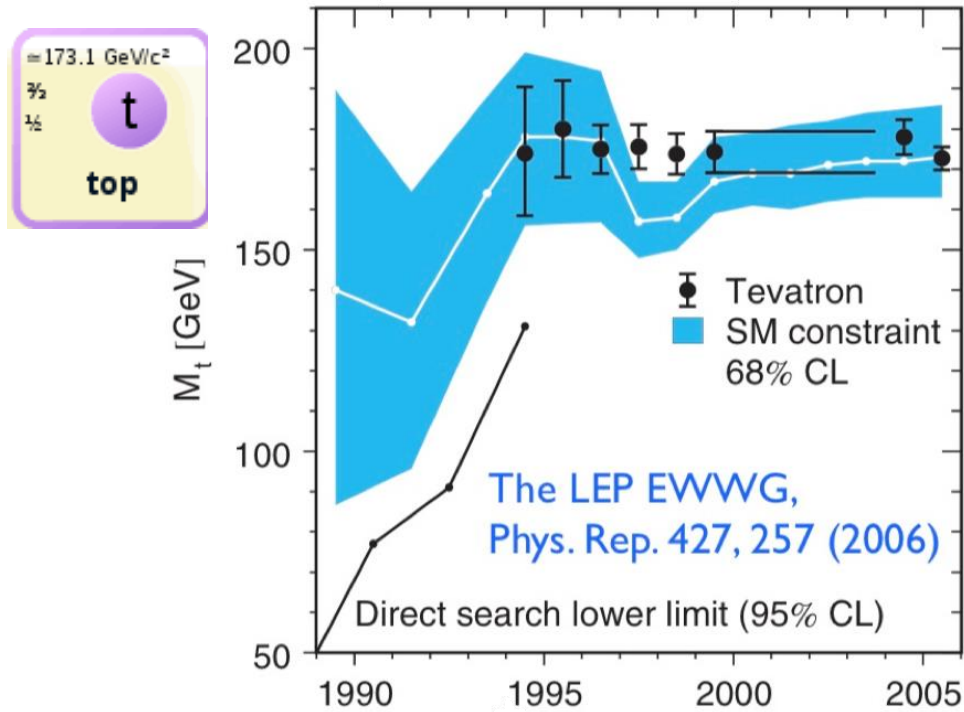
(I find you lack of faith disturbing)

# Collider experiment coordinates





# Before the LHC startup



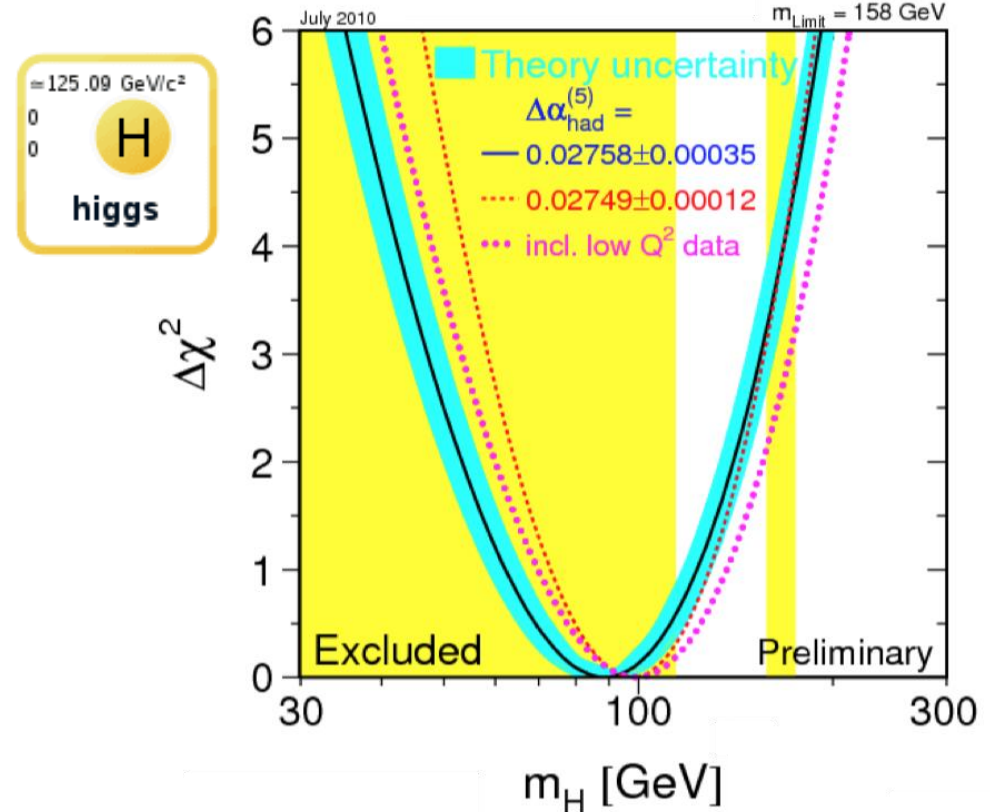
$m_W$  measurement  
at SppS and LEP-I  
precision  
measurements

top quark  
discovery  
(1994)

$m_W$  measurement  
at LEP-2

electroweak fit  
and indirect limit on  $m_H$

Direct limits on Higgs  
production from LEP-2  
and Tevatron



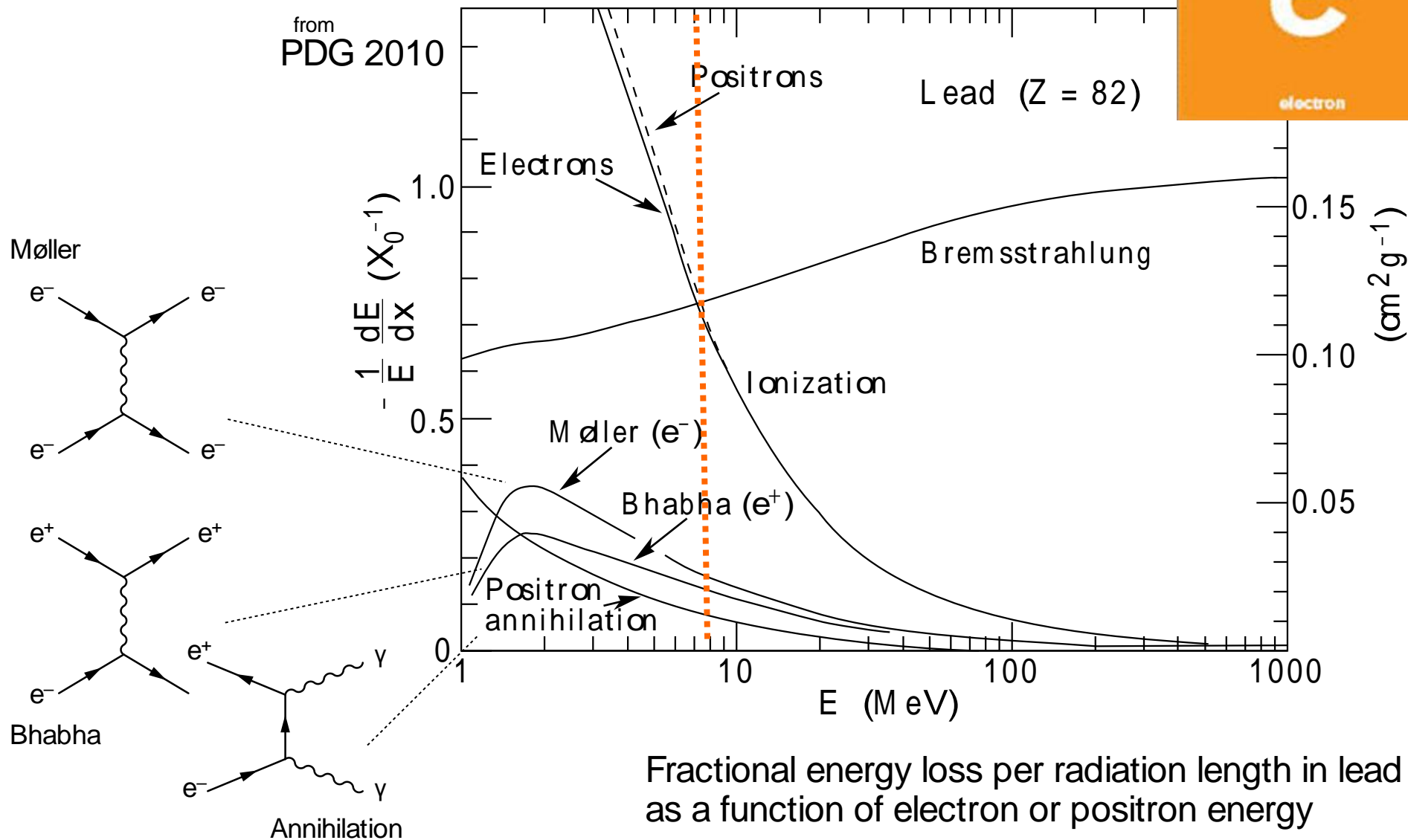
## LHC “no lose theorem”

*Either the Higgs boson is discovered,  
or New Physics should manifest to avoid unitarity violation in WW scattering at TeV scale*

# Electron energy loss

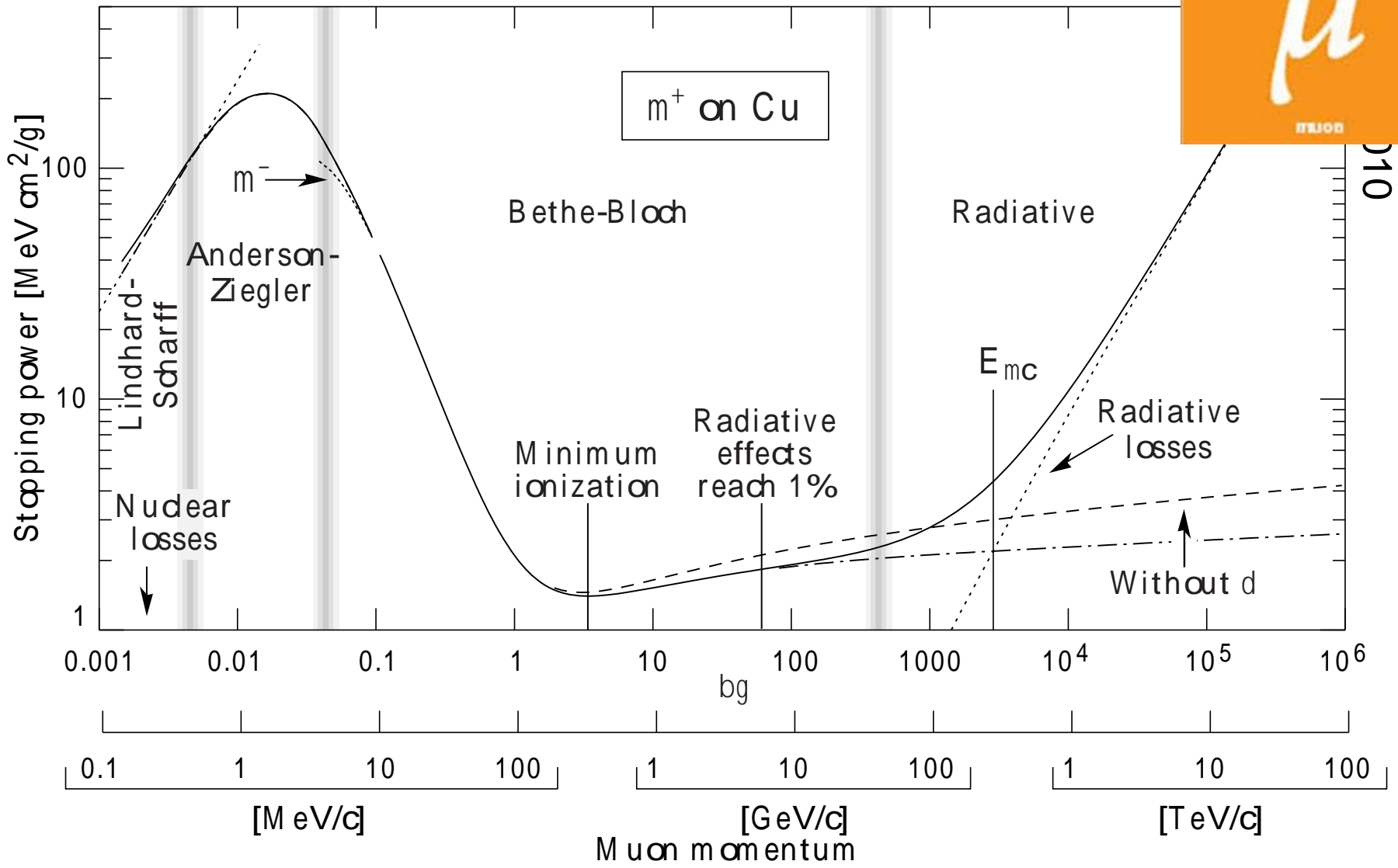


electron

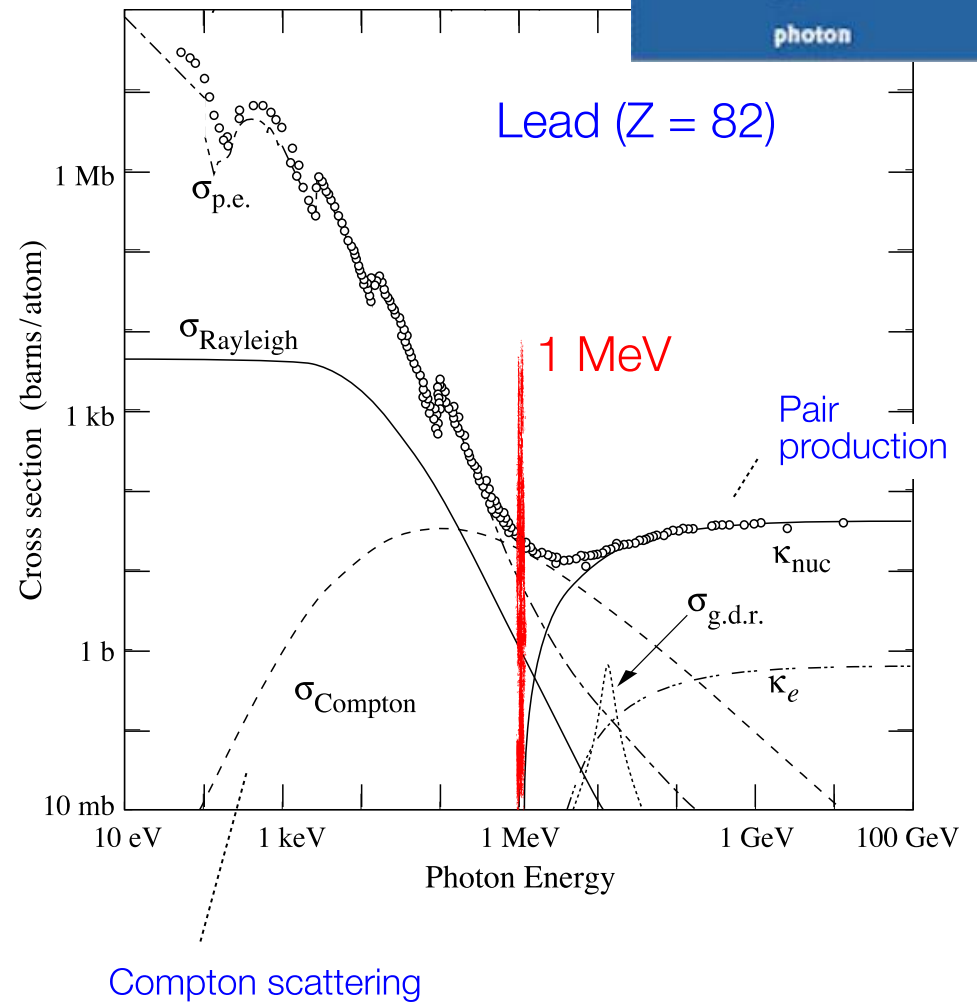
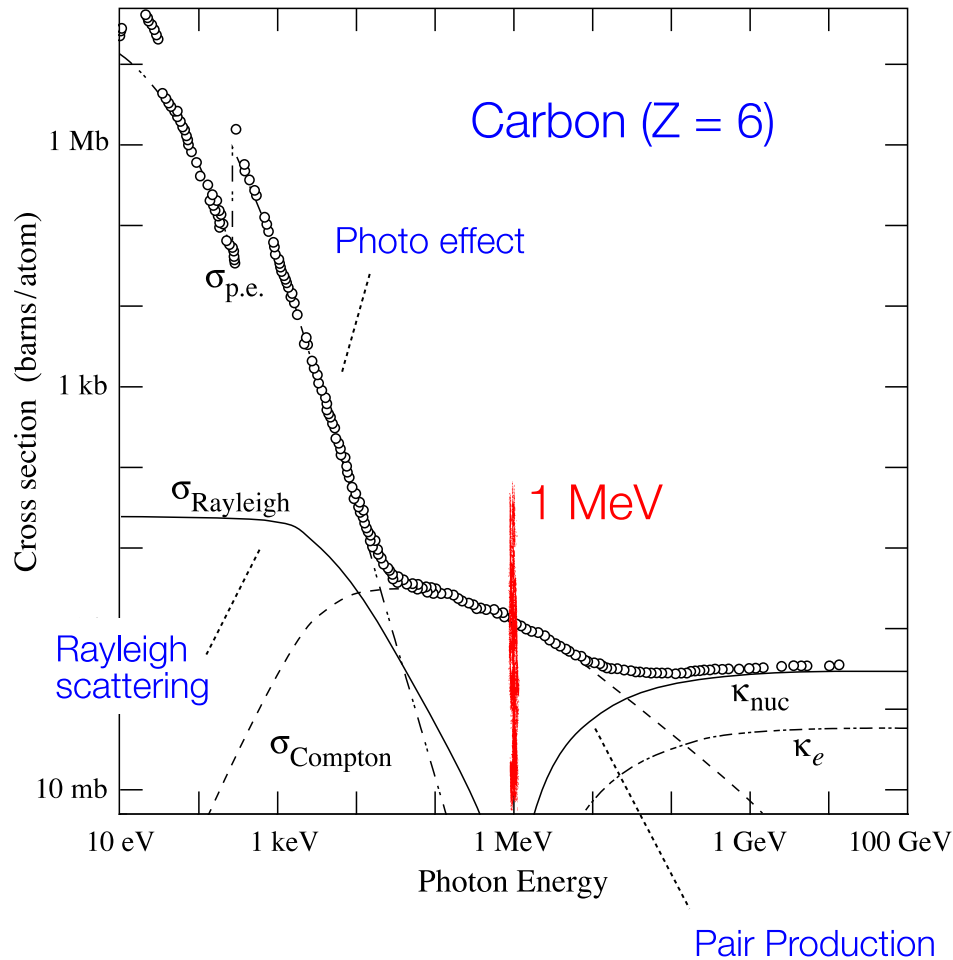




# Muon energy loss



# Interaction of photons with matter



# HEP, SI and “natural” units

Quantity	HEP units	SI units
length	1 fm	$10^{-15}$ m
charge	e	$1.602 \cdot 10^{-19}$ C
energy	1 GeV	$1.602 \times 10^{-10}$ J
mass	1 GeV/c <sup>2</sup>	$1.78 \times 10^{-27}$ kg
$\hbar = h/2\pi$	$6.588 \times 10^{-25}$ GeV s	$1.055 \times 10^{-34}$ Js
c	$2.988 \times 10^{23}$ fm/s	$2.988 \times 10^8$ m/s
$\hbar c$	197 MeV fm	...

## “natural” units ( $\hbar = c = 1$ )

mass	1 GeV
length	1 GeV <sup>-1</sup> = 0.1973 fm
time	1 GeV <sup>-1</sup> = $6.59 \times 10^{-25}$ s

# Relativistic kinematics in a nutshell

$$E^2 = \vec{p}^2 + m^2$$

$$l = \frac{l_0}{\gamma}$$

$$E = m\gamma$$

$$t = t_0\gamma$$

$$\vec{p} = m\gamma\vec{\beta}$$

$$\vec{\beta} = \frac{\vec{p}}{E}$$



# Cross section: magnitude and units

Standard

cross section unit:

$$[\sigma] = \text{mb}$$

with  $1 \text{ mb} = 10^{-27} \text{ cm}^2$

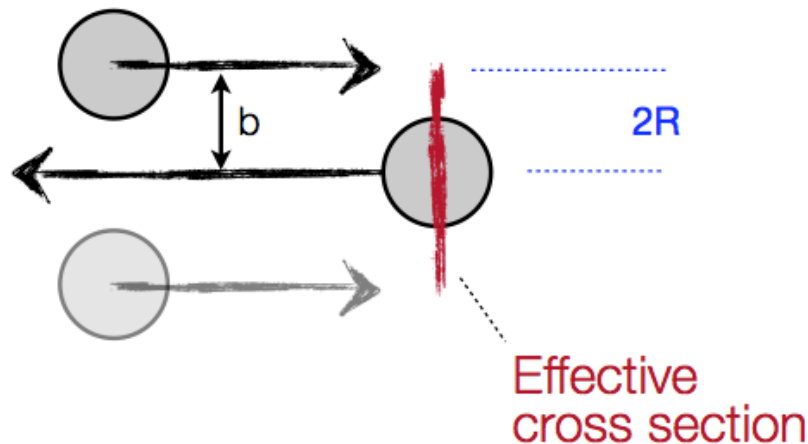
or in

natural units:

$$[\sigma] = \text{GeV}^{-2}$$

with  $1 \text{ GeV}^{-2} = 0.389 \text{ mb}$   
 $1 \text{ mb} = 2.57 \text{ GeV}^{-2}$

Estimating the  
proton-proton cross section:



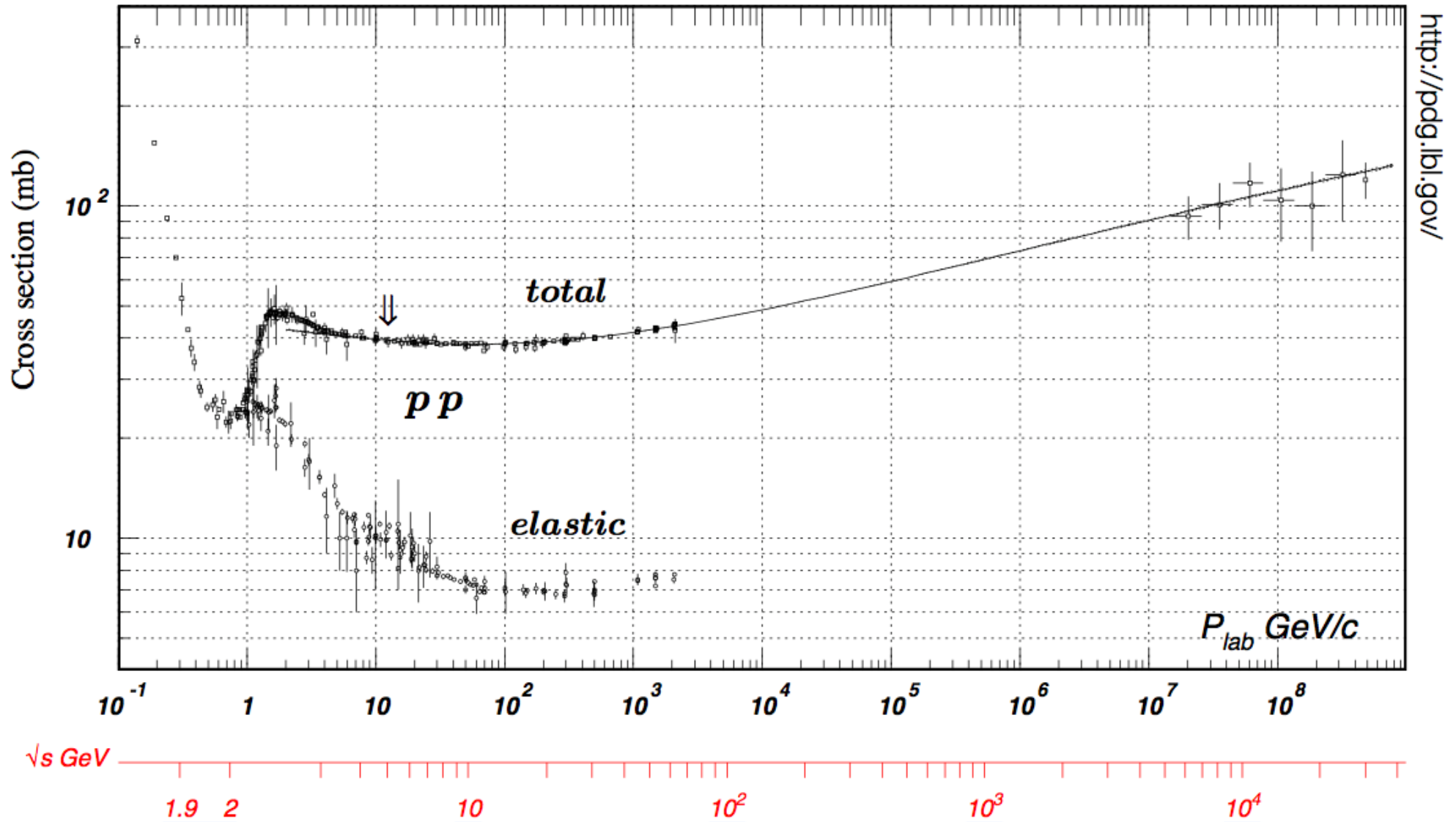
---

using:  $\hbar c = 0.1973 \text{ GeV fm}$   
 $(\hbar c)^2 = 0.389 \text{ GeV}^2 \text{ mb}$

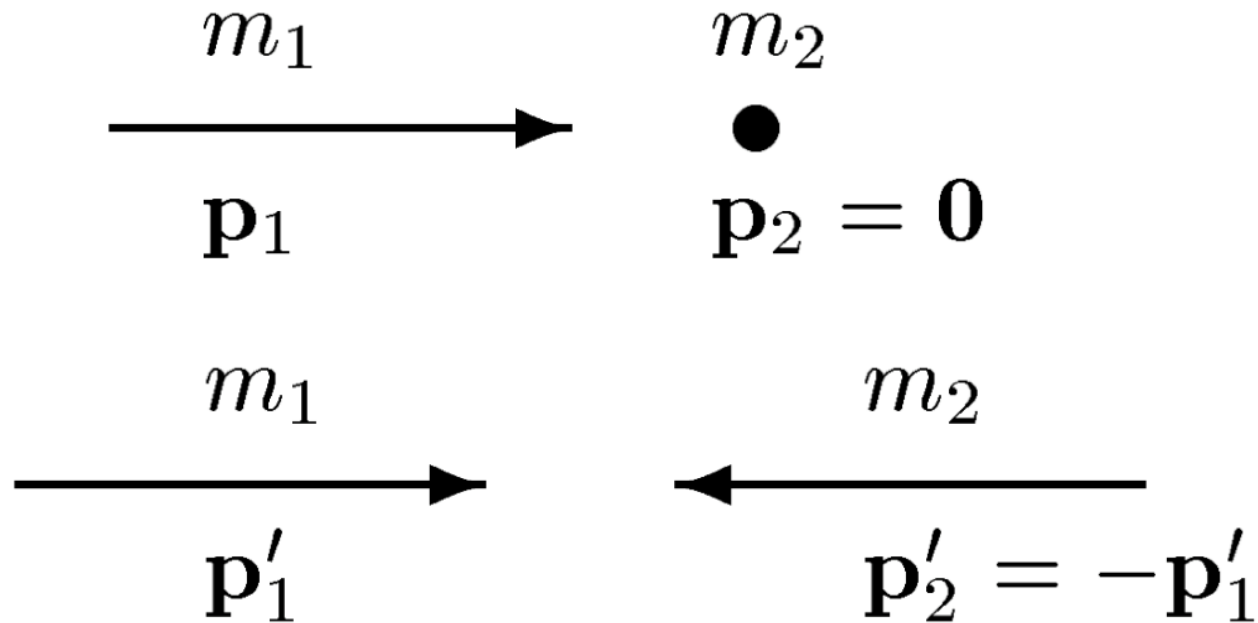
Proton radius:  $R = 0.8 \text{ fm}$   
Strong interactions happens up to  $b = 2R$

$$\begin{aligned}\sigma &= \pi (2R)^2 = \pi \cdot 1.6^2 \text{ fm}^2 \\ &= \pi \cdot 1.6^2 \cdot 10^{-26} \text{ cm}^2 \\ &= \pi \cdot 1.6^2 \cdot 10 \text{ mb} \\ &= 80 \text{ mb}\end{aligned}$$

# Proton-proton scattering cross-section



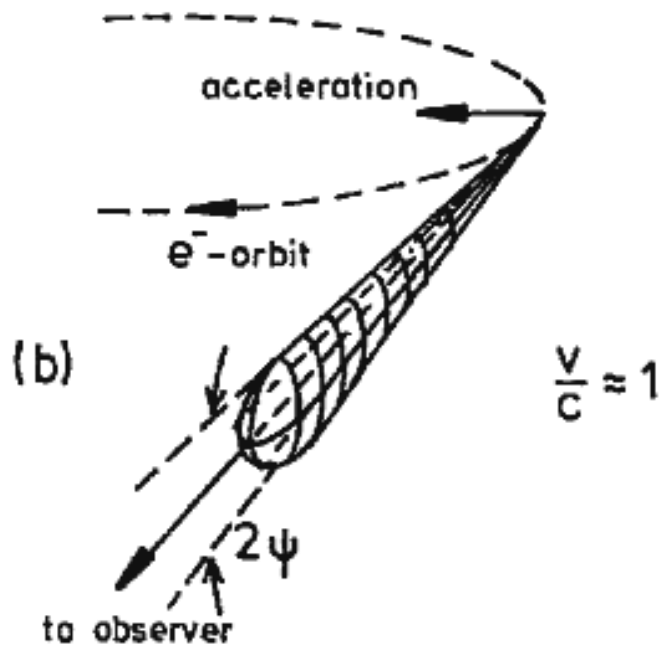
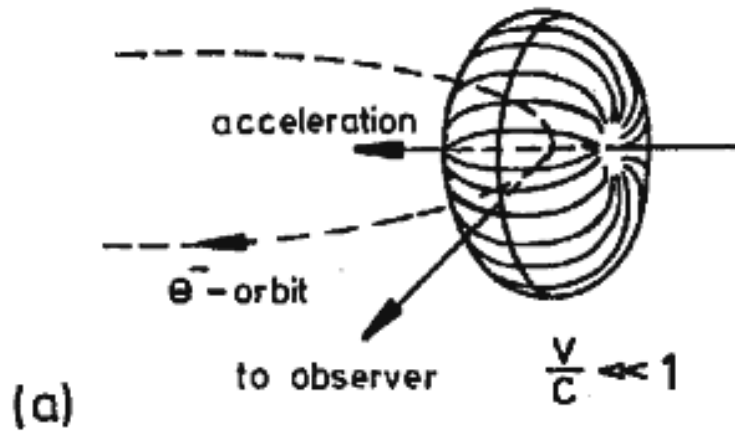
# Fixed target vs. collider



How much energy should a fixed target experiment have to equal the center of mass energy of two colliding beam?

$$E_{\text{fix}} = 2 \frac{E_{\text{col}}^2}{m} - m$$

# Synchrotron radiation



energy lost per revolution

$$\Delta E = \frac{4\pi}{3} \frac{1}{4\pi\epsilon_0} \left( \frac{e^3 \beta^3 \gamma^4}{R} \right)$$

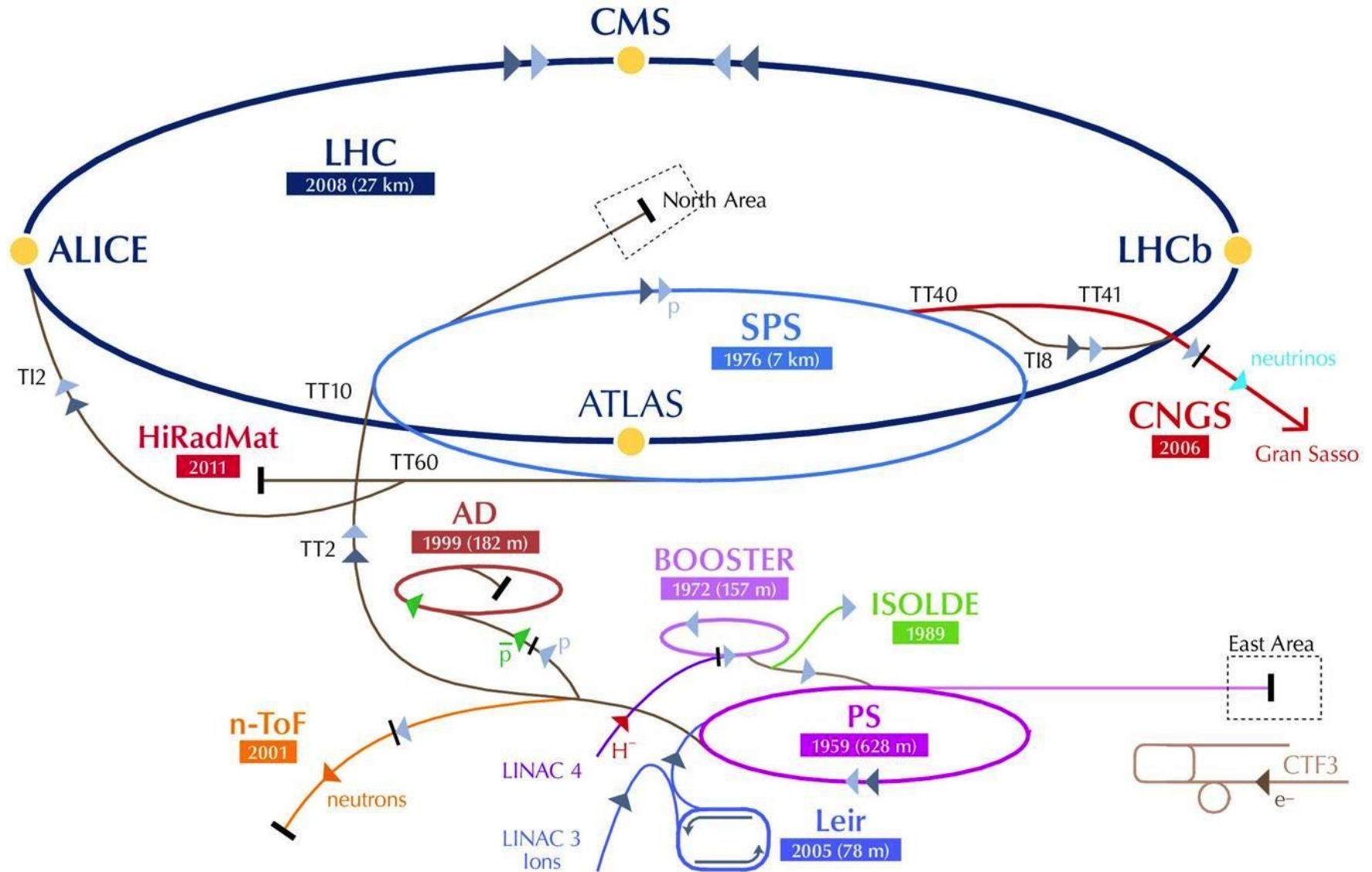
electrons vs. protons

$$\frac{\Delta E_e}{\Delta E_p} \simeq \left( \frac{m_p}{m_e} \right)^4$$

It's easier to accelerate protons to higher energies, but protons are fundamentals...



# CERN accelerator complex



# Magnetic spectrometer

Charged particle in  
magnetic field

$$\frac{d\vec{p}}{dt} = q\vec{\beta} \times \vec{B}$$

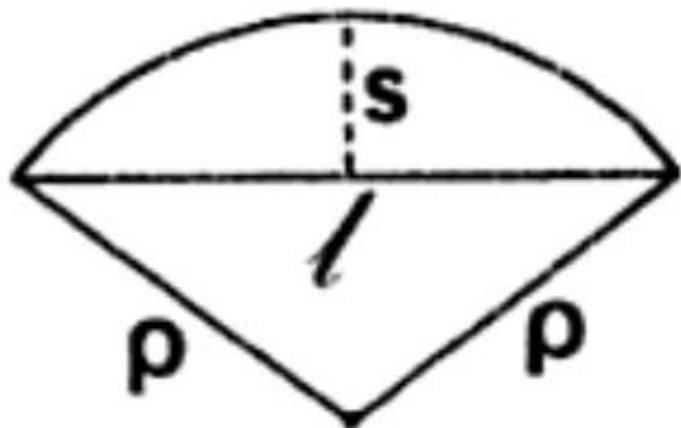
If the field is constant and we neglect presence of matter, **momentum magnitude is constant** with time, **trajectory is helical**

$$p[\text{GeV}] = 0.3B[\text{T}]\rho[\text{m}]$$

Actual trajectory differ from exact helix because of:

- **magnetic field inhomogeneity**
- **particle energy loss** (ionization, multiple scattering)

# Momentum measurement



$s$  = sagitta

$l$  = chord

$\rho$  = radius

$$\rho \simeq \frac{l^2}{8s} \quad p = 0.3 \frac{Bl^2}{8s}$$

$$\left| \frac{\delta p}{p} \right| = \left| \frac{\delta s}{s} \right|$$

*smaller for larger number of points*      *measurement error (RMS)*

Momentum resolution due to measurement error

$$\left| \frac{\delta p}{p} \right| = A_N \underbrace{\frac{\epsilon}{L^2}}_{\text{projected track length in magnetic field}} \underbrace{\frac{p}{0.3B}}_{\text{resolution is improved faster by increasing } L \text{ then } B}$$

*Momentum resolution gets worse for larger momenta*

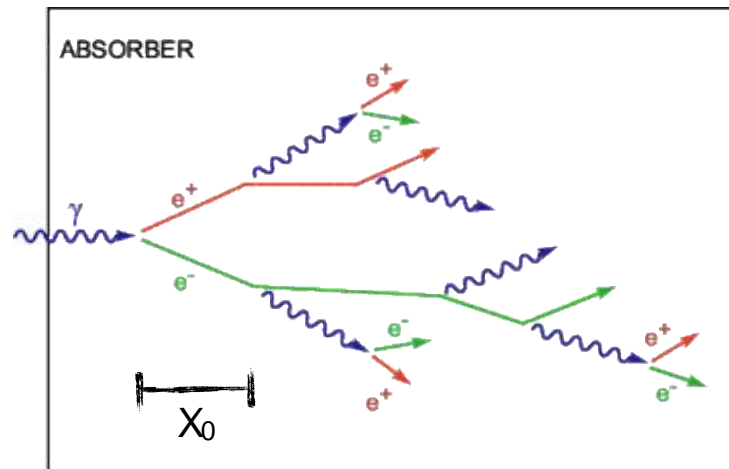
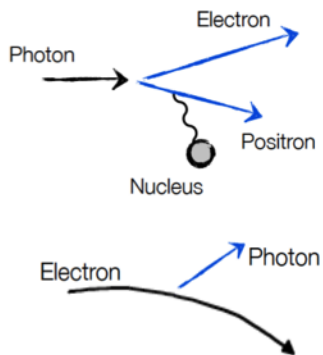
*resolution is improved faster by increasing  $L$  then  $B$*

# Electromagnetic showers

Dominant processes  
at high energies ...

Photons : Pair production

Electrons : Bremsstrahlung



Pair production:

$$\sigma_{\text{pair}} \approx \frac{7}{9} \left( 4\alpha r_e^2 Z^2 \ln \frac{183}{Z^{1/3}} \right)$$

$$= \frac{7}{9} \frac{A}{N_A X_0} \quad [X_0: \text{radiation length}]$$

[in cm or g/cm<sup>2</sup>]

Bremsstrahlung:

$$\frac{dE}{dx} = 4\alpha N_A \frac{Z^2}{A} r_e^2 \cdot E \ln \frac{183}{Z^{1/3}} = \frac{E}{X_0}$$

Absorption  
coefficient:

$$\mu = n\sigma = \rho \frac{N_A}{A} \cdot \sigma_{\text{pair}} = \frac{7}{9} \frac{\rho}{X_0}$$

$$E = E_0 e^{-x/X_0}$$

After passage of one  $X_0$  electron  
has only  $(1/e)^{\text{th}}$  of its primary energy ...  
[i.e. 37%]

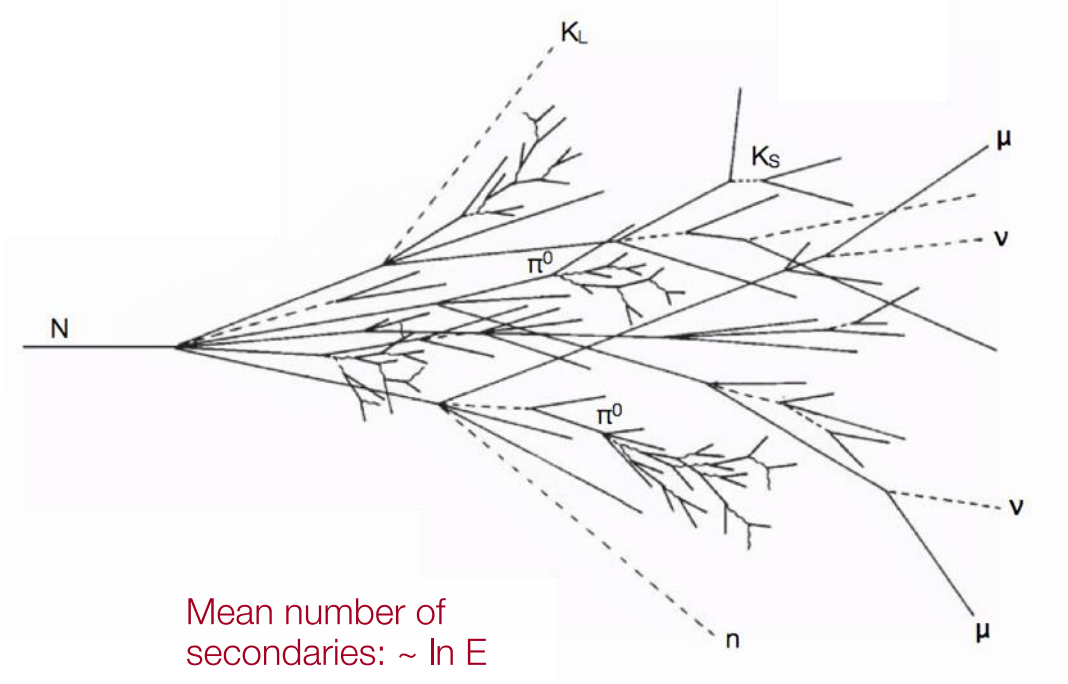
Critical energy:  $\left. \frac{dE}{dx}(E_c) \right|_{\text{Brems}} = \left. \frac{dE}{dx}(E_c) \right|_{\text{Ion}}$



# Hadronic showers

## Shower development:

1.  $p + \text{Nucleus} \rightarrow \text{Pions} + N^* + \dots$
2. Secondary particles ...  
undergo further inelastic collisions until they fall below pion production threshold
3. Sequential decays ...  
 $\pi_0 \rightarrow \gamma\gamma$ : yields electromagnetic shower  
 Fission fragments  $\rightarrow \beta$ -decay,  $\gamma$ -decay  
 Neutron capture  $\rightarrow$  fission  
 Spallation ...



Mean number of secondaries:  $\sim \ln E$

Typical transverse momentum:  $p_t \sim 350 \text{ MeV}/c$

Substantial electromagnetic fraction

$$f_{em} \sim \ln E$$

[variations significant]

### Cascade energy distribution:

[Example: 5 GeV proton in lead-scintillator calorimeter]

Ionization energy of charged particles ( $p, \pi, \mu$ )	1980 MeV [40%]
Electromagnetic shower ( $\pi^0, \eta^0, e$ )	760 MeV [15%]
Neutrons	520 MeV [10%]
Photons from nuclear de-excitation	310 MeV [ 6%]
Non-detectable energy (nuclear binding, neutrinos)	1430 MeV [29%]
	5000 MeV [29%]

# Homogeneous calorimeters

- ★ In a homogeneous calorimeter the whole detector volume is filled by a high-density material which simultaneously serves as absorber as well as as active medium ...

Signal	Material
Scintillation light	BGO, BaF <sub>2</sub> , CeF <sub>3</sub> , ...
Cherenkov light	Lead Glass
Ionization signal	Liquid noble gases (Ar, Kr, Xe)

- ★ Advantage: homogenous calorimeters provide optimal energy resolution
- ★ Disadvantage: very expensive
- ★ Homogenous calorimeters are exclusively used for electromagnetic calorimeter, i.e. energy measurement of electrons and photons

# Sampling calorimeters

## Scheme of a sandwich calorimeter

Principle:

Alternating layers of absorber and active material [sandwich calorimeter]

Absorber materials:  
[high density]

Iron (Fe)

Lead (Pb)

Uranium (U)  
[For compensation ...]

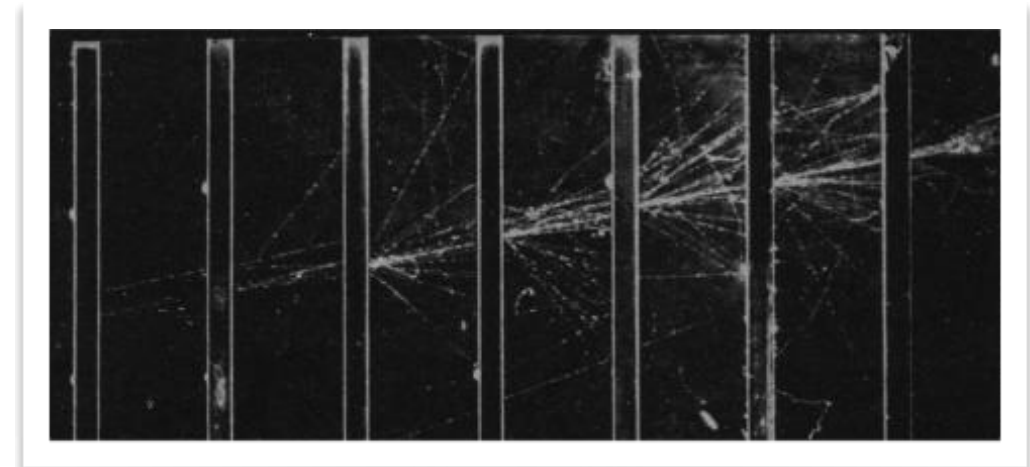
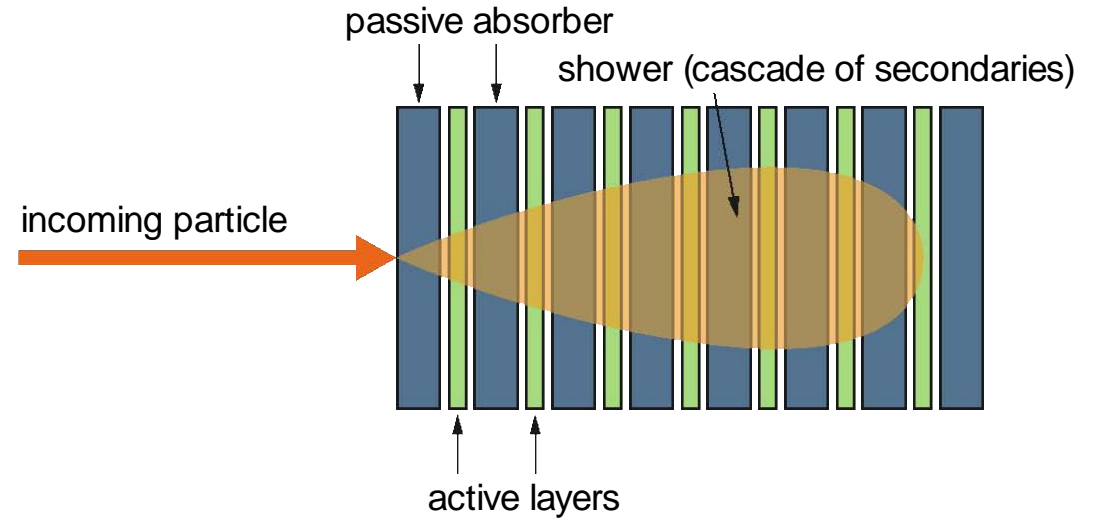
Active materials:

Plastic scintillator

Silicon detectors

Liquid ionization chamber

Gas detectors



Electromagnetic shower

# A typical HEP calorimetry system

Typical Calorimeter: two components ...

Electromagnetic (EM) +  
Hadronic section (Had) ...

Different setups chosen for  
optimal energy resolution ...

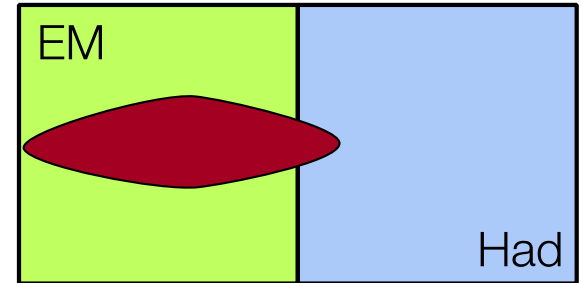
But:

Hadronic energy measured in  
both parts of calorimeter ...

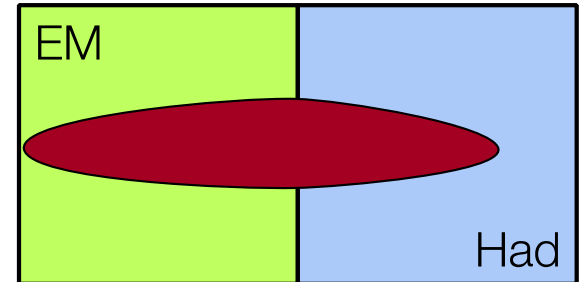
Needs careful consideration of  
different response ...

Schematic of a  
typical HEP calorimeter

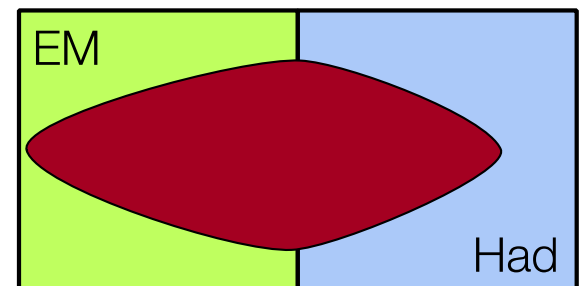
Electrons  
Photons



Taus  
Hadrons



Jets





# Energy resolution in calorimeters

Energy resolution:

e.g. inhomogeneities  
shower leakage

e.g. electronic noise  
sampling fraction variations

$$\frac{\sigma_E}{E} = \sqrt{\frac{A}{E} \oplus B \oplus \frac{C}{E}}$$

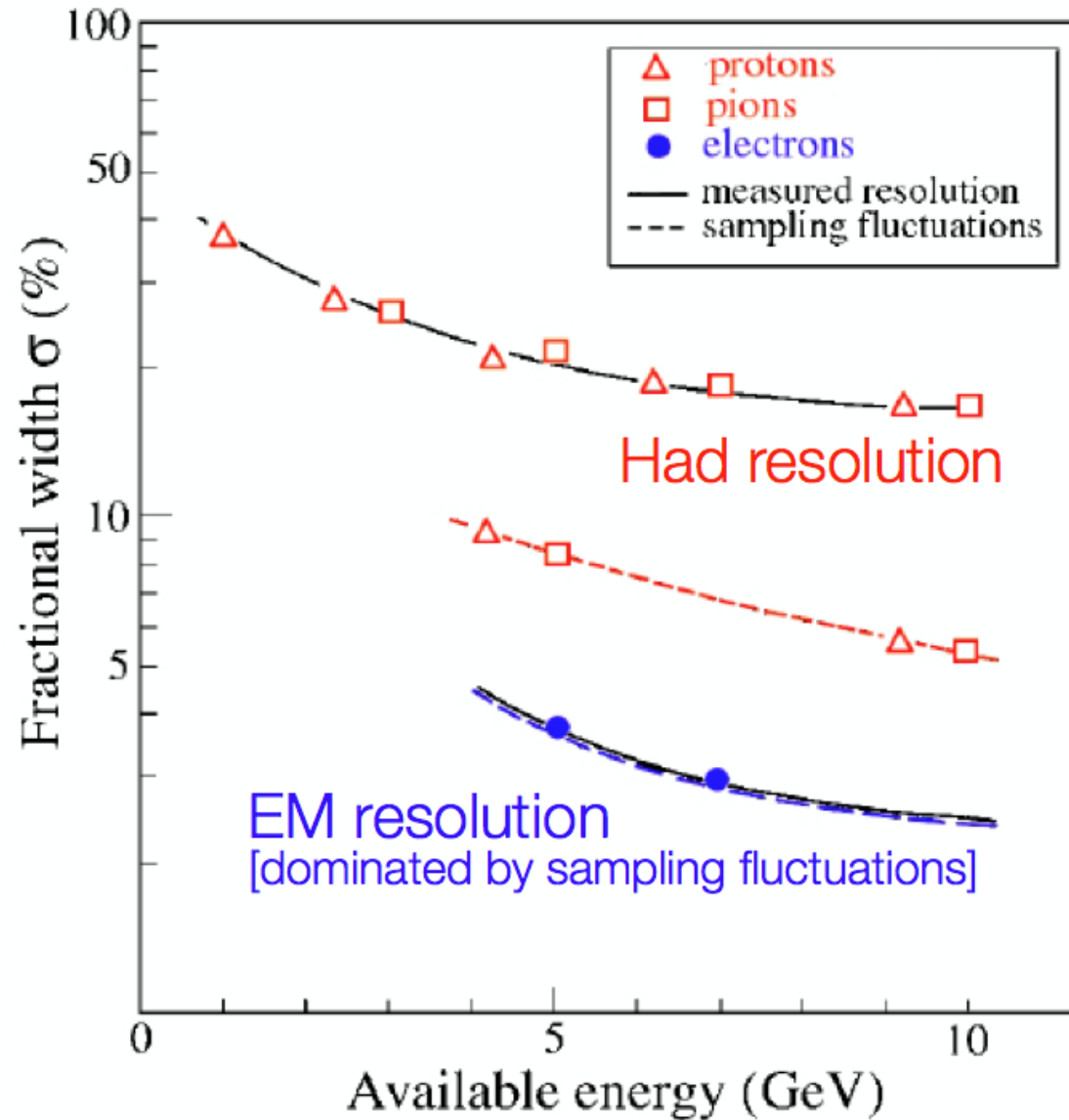
Fluctuations:

- Sampling fluctuations
- Leakage fluctuations
- Fluctuations of electromagnetic fraction
- Nuclear excitations, fission, binding energy fluctuations ...
- Heavily ionizing particles

Typical:

- A: 0.5 – 1.0 [Record:0.35]
- B: 0.03 – 0.05
- C: few %

# Resolution: EM vs. HAD



Sampling fluctuations only minor contribution to hadronic energy resolution

[AFM Collaboration]