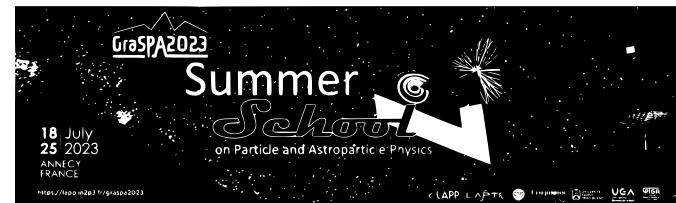




Neutrino Physics: theory

GraSPA Summer School 2023



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LPT - IJCLAB

two lectures

- ▶ Basics: brief history and basic concepts
- ▶ Oscillation phenomena and searches from many fronts
- ▶ Properties and Nature
- ▶ Theoretical frameworks and (Minimal) New Physics Models

Some references

- ▶ C. Giunti, C.W. Kim, “Fundamentals of Neutrino Physics and Astrophysics, Oxford University Press.
- ▶ R. N. Mohapatra and P. Pal, “Massive Neutrinos in Physics and Astrophysics, World Scientific
- ▶ M. Fukugita, T. Yanagida, “Physics of Neutrinos: and Application to Astrophysics (Theoretical and Mathematical Physics) ", Springer

Part 1

- ▶ Neutrino Problem: brief chronology
- ▶ Oscillation phenomena
- ▶ Searches from many fronts: present situation
- ▶ Why neutrino physics is challenging

Neutrinos are the most elusive particles of the Standard Model

$$Q_{em} = 0, \quad Q_{color} = 0.$$

Provide informations on the essential features of the SM:



“left” nature of the weak interaction and family structure

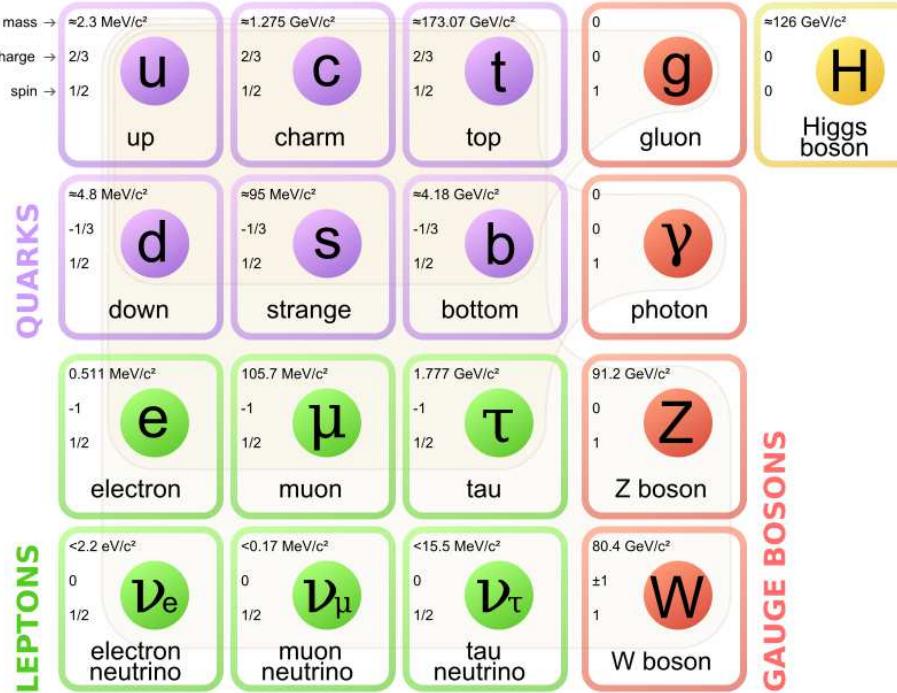


more importantly, they call for physics beyond the Standard Model

The Standard Model of particle physics

► **Standard Model:** renormalisable QFT formulation based on $SU(3)_c \times SU(2)_L \times U(1)_Y$

⇒ successful description of (most) elementary particles and their interactions



► **Gauge bosons**

strong, weak, electromagnetic interactions

► **Quarks** (strong, weak, electric);

charged leptons (weak, electric);

neutral leptons (weak)

► **Higgs boson:** EW symmetry breaking;
elementary particle masses

👉 Despite its *remarkable success*, is the SM the ultimate description of Nature?

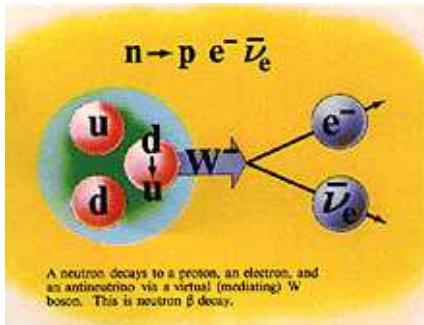
Theoretical caveats (hierarchy problem, choice of gauge group, family/flavour puzzle, ...)

Observational problems: dark matter candidate, baryon asymmetry of the Universe,
massive neutrinos!



Brief history of the neutrinos ν

- ν birth: “Rescue” conservation of energy in nucleus beta decay $n \rightarrow p + e^- + \bar{\nu}_e$



“Dear Radioactive Ladies and Gentlemen,

*... because of the wrong statistics of the N and Li^6 nuclei...and the continuous beta spectrum, I have hit upon a desperate remedy to save the "exchange theorem" of statistics and the law of conservation of energy. ... electrically neutral particles, that that I wish to call neutrons, which have spin 1/2 and obey the exclusion principle ...
...The continuous beta spectrum would then become understandable...”*

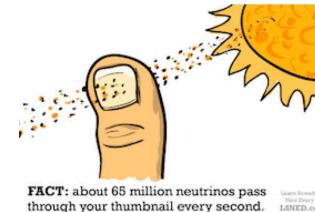
Pauli, 1930

- Enter the “neutrino”: following the discovery of the “neutron” in 1933 by Chadwick, Pauli postulates the existence of a “**massless neutrino**”
- Electron neutrino: detected in **1956** by Cowan and Reines
- Muon neutrino: discovery in **1962** by Lederman, Schwartz and Steinberger
- 3 neutrino families: Z boson decay width, CERN **1989**
- Tau neutrino: direct evidence in **2000** by DONUT team
- Neutrinos in the SM: 3 massless states! ν_e , ν_μ and ν_τ

Studying neutrinos: sources & detectors

- Neutrino sources have been experimentally and observationally explored, huge impact for particle & astroparticle physics and astronomy!

- At every second, **70×10^9 neutrinos** cross a cm^2 !



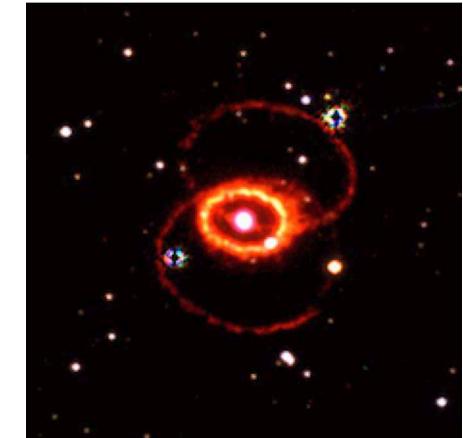
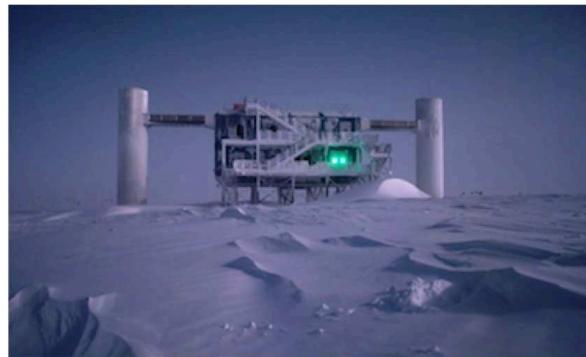
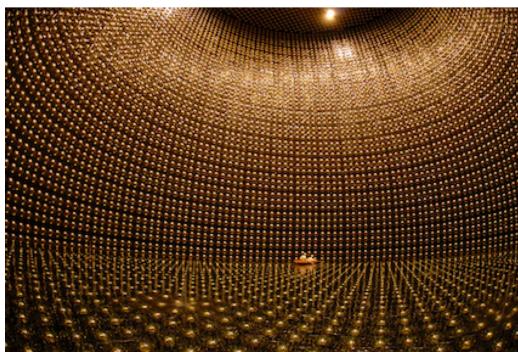
FACT: about 65 million neutrinos pass through your thumbnail every second. Learn something new at LINED.com

- A *world-wide effort* to detect and study ν 's from different sources, using distinct methods...

Laboratory: reactors, accelerators

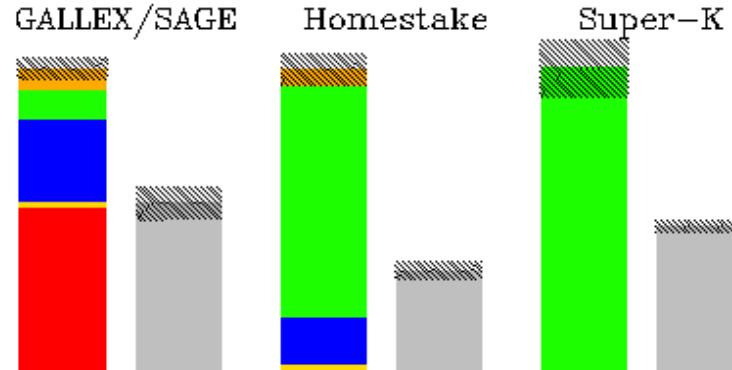
Cosmic rays: atmospheric neutrinos (ν_{\oplus}), ultra-high energy neutrinos

Astrophysical: solar neutrinos (ν_{\odot}), supernovae



Studying neutrinos: unexpected news

- A puzzling and surprising discovery: the solar ν_e and atmospheric ν_μ fluxes...



Results of Solar Neutrino experiments			
experiment	method	flux	Data/SSM (BP95)
^{37}Cl	$\nu_e {}^{37}\text{Cl}$	$2.54 \pm 0.14 \pm 0.14 \text{ SNU}$	0.27 ± 0.02
GALLEX	$\nu_e {}^{73}\text{Ga}$	$69.7 \pm 6.7 \pm 3.9 / 4.5 \text{ SNU}$	0.51 ± 0.06
SAGE	$\nu_e {}^{73}\text{Ga}$	$73 \pm 10 / 11 \text{ SNU}$	$0.53 \pm 0.07 / -0.08$
Kamiokande	ν_e scat.	$(2.80 \pm 0.19 \pm 0.83) \times 10^6 \text{ /cm}^2\text{/sec}$	0.42 ± 0.06
Super-K	ν_e scat.	$(2.44 \pm 0.06 \pm 0.25 / 0.09) \times 10^6 \text{ /cm}^2\text{/sec}$	$0.37 \pm 0.04 / -0.02$

BP95: J.N.Bahcall and M.H.Pinsonneault Rev.Mod.Phys.67(1995)781.

Solar neutrino problem: detection of only 1/3 of expected flux of solar ν_e 's

Atmospheric neutrino problem: detection of $\nu_e \sim \nu_\mu$ (expected $\nu_e \sim 2\nu_\mu$)

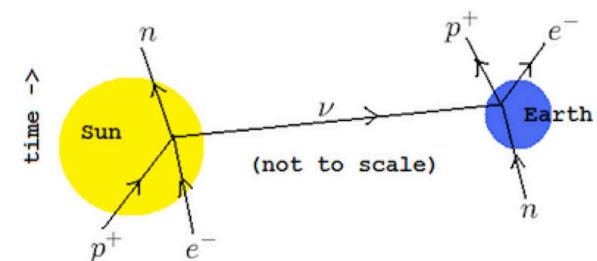


- Hypotheses:

“Unexpected” production of ν_α : *do charged currents violate lepton flavours?*

“Disappearance” of propagating ν_α : *do neutrinos oscillate?*

Standard Solar Model predictions: *to be challenged?*



👉 Facts: ν change flavours after propagating a finite distance

Solar	$\Delta m_{\text{sol}}^2 \simeq 7.4 \times 10^{-5} \text{ eV}^2$	SNO, BOREXino, Super-Kamiokande,
$\nu_e \rightarrow \nu_{\mu,\tau}$	$\sin^2 \theta_{\text{sol}} \simeq 0.30$	GALLEX/GNO, SAGE, Homestake, Kamiokande
Atmospheric		IMB, MACRO, Soudan-2,
$\nu_\mu \rightarrow \nu_\tau$		Kamiokande, Super-Kamiokande
LBL Accelerator	$\Delta m_{\text{atm}}^2 \simeq 2.5 \times 10^{-3} \text{ eV}^2$	
ν_μ disappearance	$\sin^2 \theta_{\text{atm}} \simeq 0.58$	K2K, T2K, MINOS
LBL Accelerator		
$\nu_\mu \rightarrow \nu_\tau$		Opera
LBL Accelerator		
$\nu_\mu \rightarrow \nu_e$	Δm_{atm}^2	T2K, MINOS
LBL Reactor	$\sin^2 \theta_{\text{Chooz}} \simeq 0.022$	Daya Bay, RENO
$\bar{\nu}_e$ disappearance		Double Chooz
SBL Accelerator		
$\nu_\mu (\bar{\nu}_\mu) \rightarrow \nu_e (\bar{\nu}_e)$	$\Delta m^2 \simeq 1 \text{ eV}^2 (?)$	LSND, MiniBooNE
SBL Reactor	$\sin^2 \theta \simeq 0.1 (?)$	++ Solar: GALLEX, SAGE++
$\bar{\nu}_e$ disappearance		Bugey, ILL, Rovno,...

☞ **Indisputable:** ν s are massive and mix

→ The minimal SM is incomplete!

Neutrino oscillations: massive states, leptonic mixing!

- A simple solution to both problems! Illustrative 2-family example

Two **massive states** ($\Delta m_\nu \neq 0$) related to flavour eigenstates as $\nu_\alpha = U_{\alpha i} \nu_i$

$$\begin{pmatrix} \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} \nu_2 \\ \nu_3 \end{pmatrix}$$

- What happens to a relativistic neutrino, e.g. a ν_μ produced in the atmosphere?

Production of weak eigenstate at $t = 0$:

$$|\nu_{t=0}\rangle = |\nu_\mu\rangle = \cos \theta |\nu_2\rangle + \sin \theta |\nu_3\rangle$$

Travel distance L to the **detector**, during which it **oscillates**

$$|\nu(t)\rangle = \cos \theta e^{-iE_2 t} |\nu_2\rangle + \sin \theta e^{-iE_3 t} |\nu_3\rangle$$

At the detector, it produces μ in **charged current** scattering, with probability

$$\mathcal{P}_{\mu \rightarrow \mu}^{2\nu}(L, t) = |\langle \nu_\mu | \nu(t) \rangle|^2 = 1 - \sin^2 2\theta \sin^2 \left(\frac{\Delta m_\nu^2 L}{4E} \right) \neq 1$$

- **Oscillations are possible if and only if** neutrinos are **massive** and **mix!**

👉 this is not accounted for by the SM!

Parametrisation with 3 flavours

The charged current interaction is not diagonal in flavour space:

$$\mathcal{L}_{int} = -\frac{g}{\sqrt{2}} \bar{\ell}_L^i \gamma^\mu \nu_L^j U_{ij} W_\mu^+ + h.c. ,$$

n=3 → Pontecorvo-Maki-Nakagawa-Sakata (PMNS) mixing matrix

$$U = \begin{pmatrix} c_{12} & c_{13} & & s_{12} & c_{13} & & s_{13} e^{-i\delta} \\ -s_{12} & c_{23} - c_{12} s_{23} & s_{13} e^{i\delta} & c_{12} & c_{23} - s_{12} s_{23} & s_{13} e^{i\delta} & s_{23} & c_{13} \\ s_{12} & s_{23} - c_{12} c_{23} & s_{13} e^{i\delta} & -c_{12} & s_{23} - s_{12} c_{23} & s_{13} e^{i\delta} & c_{23} & c_{13} \end{pmatrix} \text{Diag} \left\{ e^{i\alpha_1}, e^{i\alpha_2}, 1 \right\}$$

[Chau-Keung parametrisation]

δ Dirac phase, $\alpha_{1,2}$ Majorana phases, $\theta_{12}, \theta_{23}, \theta_{13}$ ($\nu = \bar{\nu}$)

m_1, m_2, m_3 mass eigenvalues, if $m_3 > 0$, $m_{1,2} = |m_{1,2}| e^{i\alpha_{1,2}}$

Transition Probabilities

$$\begin{aligned}
 P(\nu_\alpha \rightarrow \nu_\beta; L) &= \delta_{\alpha\beta} - 4 \sum_{j<k} \operatorname{Re} (U_{\alpha j} U_{\beta j}^* U_{\alpha k}^* U_{\beta k}) \sin^2 \left(\frac{\Delta m_{jk}^2 L}{4E} \right) \\
 &\quad + 2 \sum_{j<k} \operatorname{Im} (U_{\alpha j} U_{\beta j}^* U_{\alpha k}^* U_{\beta k}) \sin \left(\frac{\Delta m_{jk}^2 L}{2E} \right), \quad \Delta m_{jk}^2 = m_j^2 - m_k^2
 \end{aligned}$$

- 👉 Oscillations are possible if ν are **massive** ($\Delta m_{jk}^2 \neq 0$) and **mix** ($U_{\alpha j} U_{\beta j} \neq 0$)
- 👉 Oscillation experiments do not give the nature : **Dirac** or **Majorana** : $\bar{\nu} \equiv \nu$!

👉 $n = 2$:

$$P(\nu_\alpha \rightarrow \nu_\beta) = \sin^2 2\theta \sin^2 \left(\frac{L}{L_{vac}} \pi \right), \quad L_{vac} = \frac{4\pi E}{\Delta m^2} \simeq 2.48 \text{km} \left(\frac{E(\text{GeV})}{\Delta m^2(\text{eV}^2)} \right)$$



- 👉 oscillations arise when $L \sim L_{vac} \Rightarrow \frac{\Delta m^2 L}{4\pi E} \sim 1 \Leftrightarrow \Delta m^2(\text{eV}^2) \sim \frac{E(\text{GeV})}{L(\text{km})}$

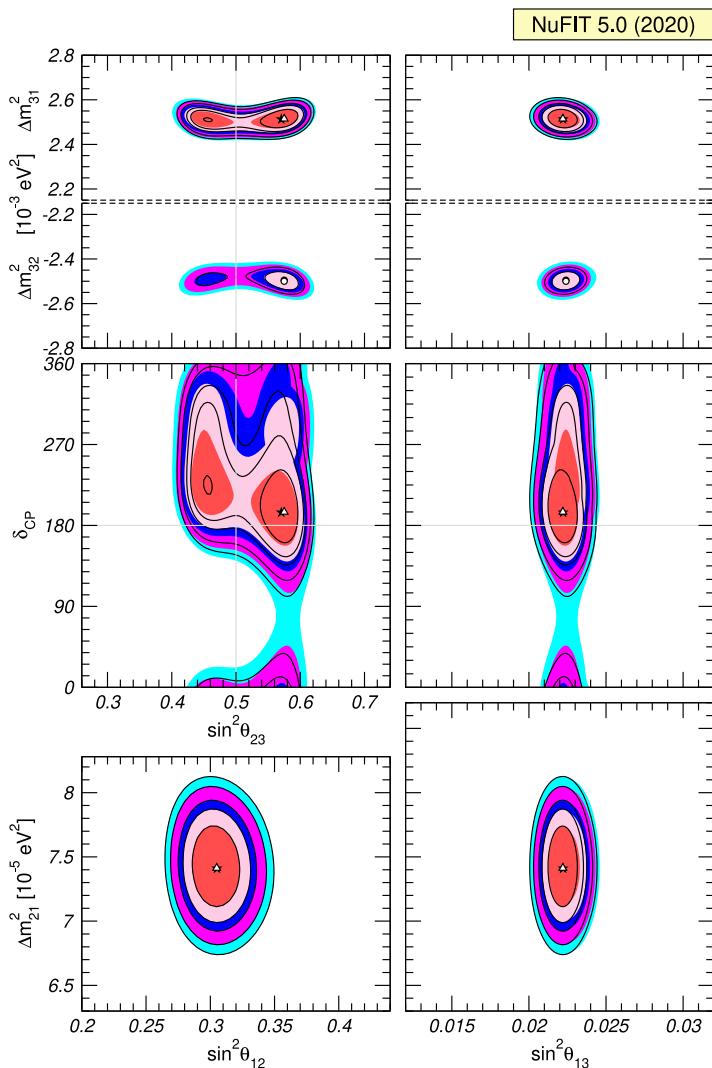
Accessible Δm^2

Depending on L and E neutrino sources:

i.e. ν source and position of the detector

$L(\text{km})$	$E(\text{GeV})$	$\Delta m^2(\text{eV}^2)$	Source
10^8	10^{-3}	10^{-11}	solar ν
10^4	1	10^{-4}	atmospheric ν
10^3	10	10^{-2}	ν from accelerators (long distance)
0.1	1	10	ν from accelerators (short distance)
0.1	10^{-3}	10^{-2}	ν from reactors

Lepton mixing & neutrino data: current status



	Normal Ordering (best fit)		Inverted Ordering ($\Delta\chi^2 = 7.1$)		
	bfp $\pm 1\sigma$	3σ range	bfp $\pm 1\sigma$	3σ range	
with SNO atmospheric data	$\sin^2 \theta_{12}$	$0.304^{+0.012}_{-0.012}$	$0.269 \rightarrow 0.343$	$0.304^{+0.013}_{-0.012}$	$0.269 \rightarrow 0.343$
	$\theta_{12}/^\circ$	$33.44^{+0.77}_{-0.74}$	$31.27 \rightarrow 35.86$	$33.45^{+0.78}_{-0.75}$	$31.27 \rightarrow 35.87$
	$\sin^2 \theta_{23}$	$0.573^{+0.016}_{-0.020}$	$0.415 \rightarrow 0.616$	$0.575^{+0.016}_{-0.019}$	$0.419 \rightarrow 0.617$
	$\theta_{23}/^\circ$	$49.2^{+0.9}_{-1.2}$	$40.1 \rightarrow 51.7$	$49.3^{+0.9}_{-1.1}$	$40.3 \rightarrow 51.8$
	$\sin^2 \theta_{13}$	$0.02219^{+0.00062}_{-0.00063}$	$0.02032 \rightarrow 0.02410$	$0.02238^{+0.00063}_{-0.00062}$	$0.02052 \rightarrow 0.02428$
	$\theta_{13}/^\circ$	$8.57^{+0.12}_{-0.12}$	$8.20 \rightarrow 8.93$	$8.60^{+0.12}_{-0.12}$	$8.24 \rightarrow 8.96$
	$\delta_{CP}/^\circ$	197^{+27}_{-24}	$120 \rightarrow 369$	282^{+26}_{-30}	$193 \rightarrow 352$
	$\frac{\Delta m_{21}^2}{10^{-5} \text{ eV}^2}$	$7.42^{+0.21}_{-0.20}$	$6.82 \rightarrow 8.04$	$7.42^{+0.21}_{-0.20}$	$6.82 \rightarrow 8.04$
	$\frac{\Delta m_{3\ell}^2}{10^{-3} \text{ eV}^2}$	$+2.517^{+0.026}_{-0.028}$	$+2.435 \rightarrow +2.598$	$-2.498^{+0.028}_{-0.028}$	$-2.581 \rightarrow -2.414$

- “Precision era” for neutrino physics
- Only three oscillation parameters unknown...
- θ_{23} octant; δ_{CP} ; ν -mass ordering
- Exciting and rich experimental program ahead!



So far, only Δm_{ij}^2 , butwhat about absolute neutrino masses?

Lepton mixing & neutrino data: current status

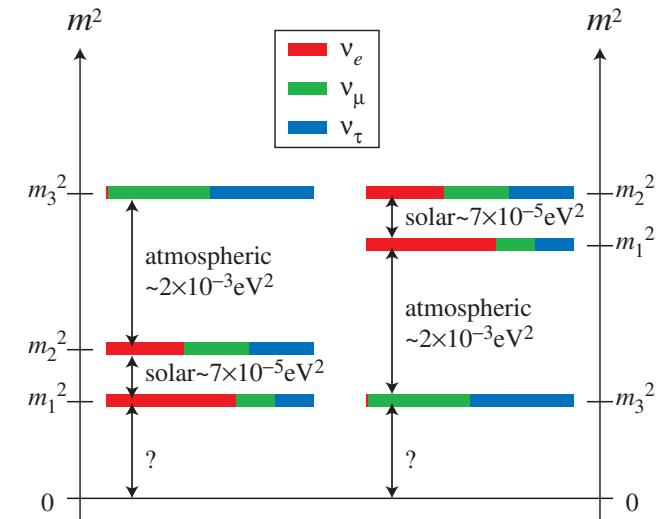
- Oscillation data: only two squared-mass differences

Undetermined mass ordering:

normal [$m_{\nu_1} < m_{\nu_2} \ll m_{\nu_3}$]

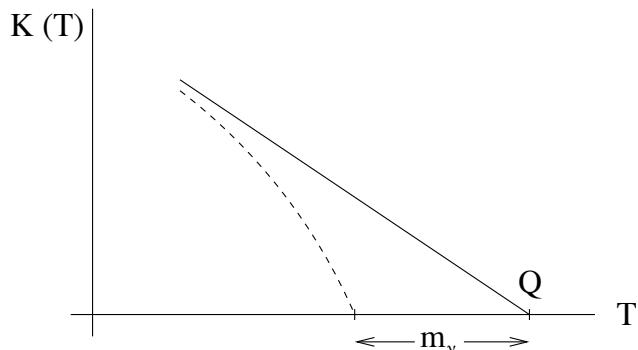
inverted [$m_{\nu_3} \ll m_{\nu_1} \lesssim m_{\nu_2}$]

Unknown absolute mass scale

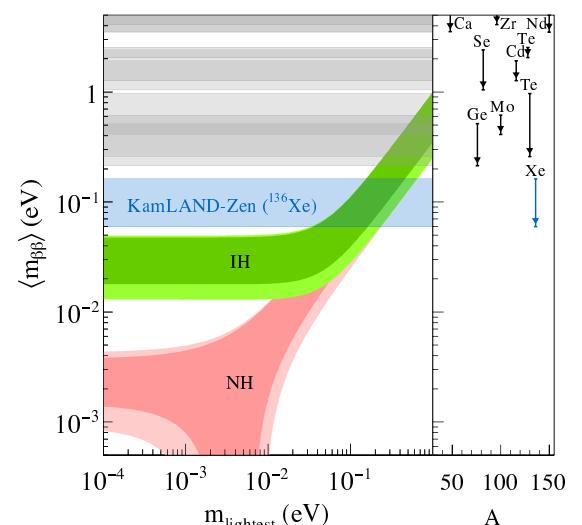


- Resolving the absolute mass scale m_{lightest}

- Tritium decays (${}^3\text{H} \rightarrow {}^3\text{He} + \bar{\nu}_e + e^-$): $m_{\nu_e} \lesssim 0.8 \text{ eV}$ [KATRIN'22] the only direct mass determination
- $0\nu 2\beta$ decays - if Majorana nature: $|m_{ee}| \lesssim 0.06 - 0.16 \text{ eV}$ [KamLAND-Zen]
- Cosmology (CMB, LSS, Lyα): $\sum_i m_{\nu_i} \lesssim 0.26 \rightarrow 0.39 \text{ eV}$



[KamLAND-Zen Coll., '16]



Leptonic CP Asymmetry

$$\begin{aligned}\Delta_{CP}(\alpha\beta) &\equiv P(\nu_\alpha \rightarrow \nu_\beta) - P(\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta) \\ &= 4 \sum_{j>k} \text{Im} \left(U_{\alpha j} U_{\beta j}^* U_{\alpha j}^* U_{\beta k} \right)^* \sin \left(\Delta m_{jk}^2 \frac{L}{2E} \right)\end{aligned}$$

☞ Cannot be observed in **appearance experiments**

$$CPT \rightarrow \Delta_{CP}(e\mu) = \Delta_{CP}(\mu\tau) = \Delta_{CP}(\tau e) \equiv 16 \mathcal{J} \ell_{12} \ell_{23} \ell_{31}$$

☞ $\mathcal{J} \equiv \text{Im} \left(U_{e3} U_{e1}^* U_{\mu 3}^* U_{\mu 1} \right) \simeq \sin 2\theta_{23} \sin 2\theta_{12} \sin \theta_{13} \sin \delta$
(Jarlskog Invariant)

☞ $\ell_{ij} \equiv \sin \left(1.27 \Delta m_{ij}^2 (\text{eV})^2 \frac{L(\text{km})}{E(\text{GeV})} \right)$

☞ θ_{23} large (OK) and also Δm_{13}^2 .

☞ θ_{12} large (OK) and also Δm_{12}^2 .

☞ θ_{13} conditions the measurement of CPV phase: $\theta_{13} \sim 8.5^\circ$

☞ **Indisputable:** ν s are massive and mix

→ **The minimal SM is incomplete!**



An observational Caveat that is also theoretical one!



► ν mixings "add fuel to the fire": add to the fermion flavour puzzle!

$$U_{CKM} = \begin{pmatrix} 1 - \lambda^2/2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix}, \lambda \sim 0.2, A \simeq 0.8, \rho \simeq 0.1, \eta \simeq 0.4$$

→ Quarks: small mixing angles, 1 Dirac CPV phase

$$U_{PMNS} = \begin{pmatrix} c_{13}c_{12} & c_{13}s_{12} & s_{13}e^{-i\delta} \\ -c_{23}s_{12} - s_{23}s_{13}c_{12}e^{i\delta} & c_{23}c_{12} - s_{23}s_{13}s_{12}e^{i\delta} & -s_{23}c_{13} \\ s_{23}s_{12} - c_{23}s_{13}c_{12}e^{i\delta} & -s_{23}c_{12} - c_{23}s_{13}s_{12}e^{i\delta} & c_{23}c_{13} \end{pmatrix} \times \text{diag}(e^{i\alpha_1}, e^{i\alpha_2}, 1)$$

Leptons: 2 large mixing angles, 1 Dirac + 2 Majorana CPV phases

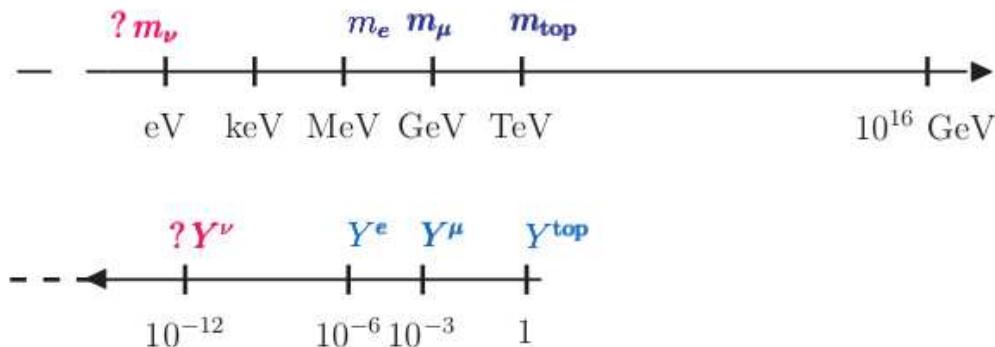
⇒ Very different mixing pattern for Leptons and Quarks



→ Is this related to different mass generation mechanisms?

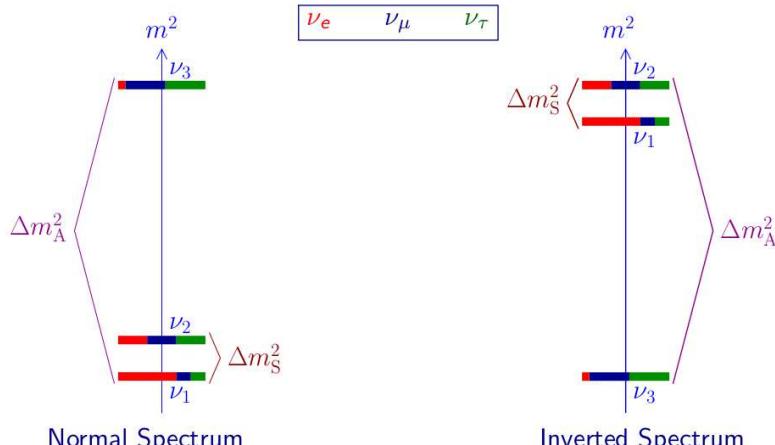
- ▶ ν data worsens fermion hierarchy problem!

→ Why are ν so light?

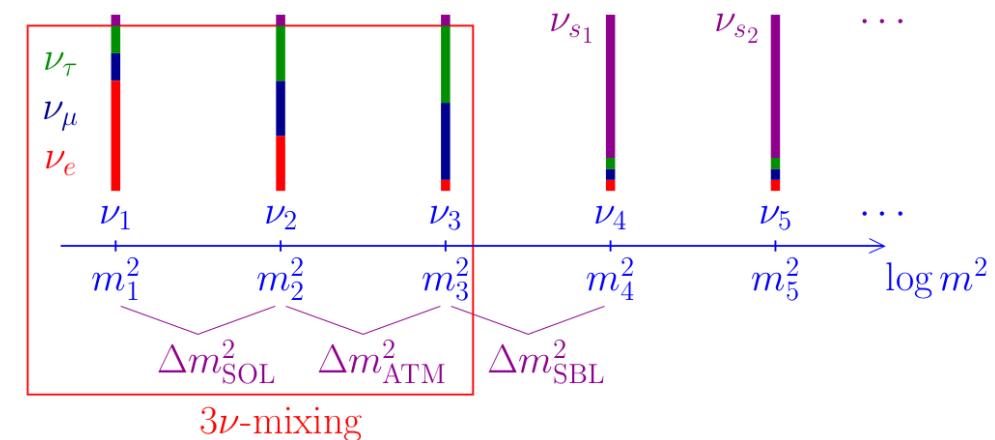


→ What is the absolute neutrino mass scale?

► Are there some extra fermionic gauge singlets (steriles)?



3- ν mixing scheme



3+? ν mixing schemes

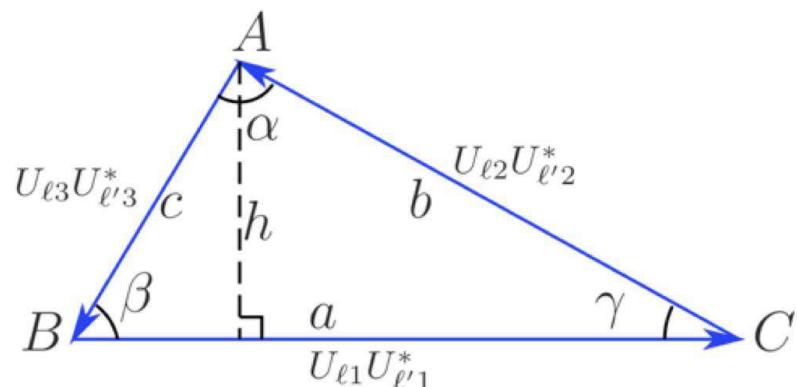
→ Does this mean that U_{PMNS} is incomplete? Non-Unitary ?



► Strong Potential for CP violation

☞ Unitarity triangle surface $\propto J_{\text{CP}}^{\text{lepton}}$: $J_{\text{CP},\text{max}} \simeq 1000 \times J_{\text{CP}}^{\text{quark}}$

$$J_{\text{CP}}^{\text{quark}} = 2.96 \times 10^{-5}, \quad J_{\text{CP},\text{max}} \simeq 3.29 \times 10^{-2}$$



Unitarity Triangle (in e, μ)

$$\mathcal{J} \equiv \text{Im} (U_{e3}U_{e1}^*U_{\mu 3}^*U_{\mu 1})$$

$$\mathcal{J} = \sin 2\theta_{23} \sin 2\theta_{12} \sin \theta_{13} \sin \delta$$

Jarlskog Invariant

→ New possibility for having Baryogenesis from Leptogenesis ?

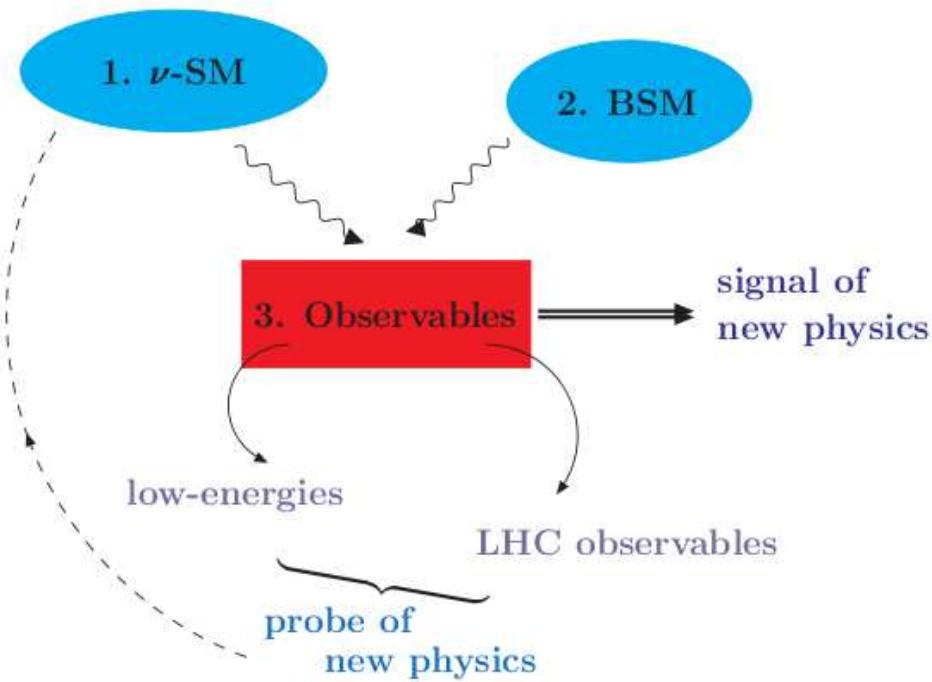
Lepton mixing & massive neutrinos: a gateway to NP

- ☞ ν -SM = New Physics just to explain ν masses and mixings
- ▶ Need other d.o.f, for instance Right-Handed Neutrinos, $m_\nu \leftarrow H Y^\nu \nu_L \nu_R + \dots$
- ▶ What is the neutrino mass generation mechanism?
 - ▶ $\nu \leftrightarrow \bar{\nu}$ the **only particle** that can have *both* Dirac or Majorana descriptions
 - ▶ If ν has Majorana nature → **New physics scale (\neq EW scale)**
- ☞ ν -SM will allow for many new phenomena
 - ▶ Lepton flavour violation in neutral sector, not in the charged one? $\ell_i \rightarrow \ell_j \ell_k \ell_l, \ell_i \rightarrow \ell_j \gamma, \dots$
 - ▶ Contributions to other observables like $g - 2$, Lepton EDMs
 - ▶ Collider searches for new heavy states ...
- ☞ SM has other issues that call for larger BSM frameworks
 - ▶ observational problems (ν masses & mixings): BAU and Dark Matter
 - ▶ theoretical caveats: fine-tuning, hierarchy and flavour problems

→ Determination of ν -SM/BSM model requires combinations of \neq observables



👉 How to proceed?



► Ingredients:

1. mass generation mechanism (**seesaw**, radiative corrections, extra dim, ...)
2. extension of SM: SM + new d.o.f, or BSM (extend Higgs and/or gauge sectors, e.g. SUSY, ...)

► Observables (peculiar to these extensions):

- Produce directly new d.o.f at LHC (if accessible)
- or study impact of 1. (and 1. + 2.) on **observables** at low-energy/high-intensity experiments (MEG, ...) and high-energy (LHC)

► Probe New Physics: 🐢 interplay between low- and high-energy observables