HS Symmetry at Infinite Distance and its Stringy Origin José Calderón Infante

Based on 2305.05693 with Florent Baume and ongoing work with Irene Valenzuela Landscapia, IPhT, CEA/Saclay, 29/11/2023







Swampland Distance Conjecture (SDC) [Ooguri, Vafa '06]

There is an infinite tower of states becoming light at infinitedistance points in moduli space:

 $M_{tower} \sim e^{-\alpha \Delta \phi}$ as $\Delta \phi \to \infty$ $(M_{Pl} = 1)$



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Precise order one bounds on the exponential rates
Lightest tower: [Etheredge, Heidenreich, Kaya, Qiu, Rudelius '22]
Species scale: [van de Heisteeg, Vafa, Wiesner, Wu '23] [JCI, Castellano, Herráez, Ibáñez '23]
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Progress:





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+ Bottom-up motivations [Hamada, Montero, Vafa, Valenzuela '21] [Stout '21+'22] [JCI, Castellano, Herráez, Ibáñez '23]





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Conformal manifold of local CFT in d>2



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Local CFT: Posses energy-momentum tensor

Dynamical gravity in the bulk!



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[Baume, JCI '20] [Perlmutter, Rastelli, Vafa, Valenzuela '20] How does the SDC look like in AdS/CFT **?**

[Perlmutter, Rastelli, Vafa, Valenzuela '20] **CFT Distance Conjecture:** Conformal manifold of local CFT in d>2 I. HS point → Infinite distance II. Infinite distance → HS point $\lim \gamma_{\ell} = \Delta_{\ell} - (\ell + d - 2) \sim e^{-\alpha_{\ell} \ell}$

Zamolodchikov distance

Local CFT: Posses energy-momentum tensor Dynamical gravity in the bulk!



Can we prove this using CFT techniques **?**





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[Baume, JCI '23]

Towards a CFT Distance Theorem:

Conformal perturbation theory + HS symmetry



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I. Finite vs infinite distance criterion!

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$$\blacksquare \leftrightarrow \alpha_{\ell} \sim \left\langle K_{\ell-1} K_{\ell-1} \mathcal{O} \right\rangle_{HS} \neq 0$$

Evaluated at HS point! -

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Evaluated at HS point! -



No extra assumption, e.g., no supersymmetry + existence of energy-momentum is crucial!

HS

 \mathcal{M}_{CFT}



 \mathcal{M}_{CFT}



 $\delta \left\langle J_{\ell} J_{\ell} \right\rangle_{t} = \delta t \int d^{d} y \left\langle J_{\ell} J_{\ell} \mathcal{O}(y) \right\rangle_{t}$



$$\delta \left\langle J_{\ell} J_{\ell} \right\rangle_{t} = \delta t \int d^{d} y \left\langle J_{\ell} J_{\ell} \mathcal{O}(y) \right\rangle_{t}$$

$$\downarrow$$

$$\delta \gamma_{\ell} = -C_{JJ\mathcal{O}}(t) \,\delta t$$



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$$\left\{ \frac{d\gamma_{\ell}}{dt} = -C_{JJ\mathcal{O}} \right\}$$

in the second second



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 \mathcal{M}_{CFT}

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$$\downarrow$$

$$\left(\frac{d\gamma_{\ell}}{dt} = -C_{JJ\mathcal{O}}\right)$$
Trajectory!
(not necessarily a g

Sketch of the Proof



geodesic)






(Weakly-broken) HS symmetry



(Weakly-broken) HS symmetry

$$\partial \cdot J_{\ell} = g(t) K_{\ell-1}$$



(Weakly-broken) HS symmetry

$$\partial \cdot J_{\ell} = \underbrace{g(t)}_{\ell - 1} K_{\ell - 1}$$

$$g(t) \to 0 \text{ as } t \to H^{2}$$



(Weakly-broken) HS symmetry

$$\partial \cdot J_{\ell} = \underbrace{g(t)}_{\ell - 1} K_{\ell - 1}$$

$$C_{JJ\mathcal{O}} \simeq C_{JJ\mathcal{O}}^{HS} + C_{JK\mathcal{O}}^{HS} g + C_{KK\mathcal{O}}^{HS} g^2 + \cdots$$



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Essentially, as $g \rightarrow 0$:

$$\partial \cdot \left\langle J_{\ell} J_{\ell} \mathcal{O} \right\rangle = g \left\langle J_{\ell} K_{\ell-1} \mathcal{O} \right\rangle$$



(Weakly-broken) HS symmetry

$$\partial \cdot J_{\ell} = \underbrace{g(t)}_{\ell-1} K_{\ell-1}$$

 $g(t) \to 0 \text{ as } t \to HS$

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$$C_{JJ\mathcal{O}} \qquad C_{JK\mathcal{O}}$$



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Essentially, as $g \rightarrow 0$:



A bit more complicated than this... → More details in [Baume, JCI '23]



(Weakly-broken) HS symmetry

$$\partial \cdot J_{\ell} = \underbrace{g(t)}_{\ell-1} K_{\ell-1}$$

$$\downarrow g(t) \to 0 \text{ as } t \to \mathsf{H}$$

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(Weakly-broken) HS symmetry



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 $\partial \cdot J_{\ell} = g(t) K_{\ell-1}$ $g(t) \to 0 \text{ as } t \to \text{HS}$ $C_{JJO} \simeq C_{JJO}^{HS} + C_{JKO}^{HS} g + C_{KKO}^{HS} g^2 + \cdots$

 $|C_{JJO}| \lesssim \gamma_{\ell} \operatorname{as} \gamma_{\ell} \to 0$

HS symmetry constraints

 $\gamma_{\ell} \sim g^2$



$$\left|\frac{d\gamma_{\ell}}{dt}\right| = \left|C_{JJO}\right| \lesssim \gamma_{\ell} \operatorname{as} \gamma_{\ell} \to 0$$



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Asymptotic version of bound in [van de Heisteeg, Vafa, Wiesner '20]



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 \mathcal{M}_{CFT}

$$\left|\frac{d\gamma_{\ell}}{dt}\right| = \left|C_{JJ\mathcal{O}'}\right| \lesssim \gamma_{\ell} \operatorname{as} \gamma_{\ell} \to 0$$



Sketch of the Proof



O'



Conformal perturbation theory + (Weakly-broken) HS symmetry

$$\left|\frac{d\gamma_{\ell}}{dt}\right| = \left|C_{JJ\mathcal{O}'}\right| \lesssim \gamma_{\ell} \operatorname{as} \gamma_{\ell} \to 0$$





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Conformal perturbation theory + (Weakly-broken) HS symmetry

$$\left|\frac{d\gamma_{\ell}}{dt}\right| = \left|C_{JJ\mathcal{O}''}\right| \lesssim \gamma_{\ell} \operatorname{as} \gamma_{\ell} \to 0$$







$$\left|\frac{d\gamma_{\ell}}{dt}\right| = \left|C_{JJO}\right|$$



$$\lesssim \gamma_{\ell} \operatorname{as} \gamma_{\ell} o 0 \quad \forall \mathcal{O}$$



$$\left| \left| \frac{d\gamma_{\ell}}{dt} \right| = \left| C_{JJO} \right| \lesssim \gamma_{\ell} \text{ as } \gamma_{\ell} \to 0 \quad \forall J_{\ell}, \mathcal{O}$$

$$\gamma_{\ell} \to 0 \text{ as } t \to \infty \quad \forall J_{\ell}, \mathcal{O} : \text{All HS points are at infinit}$$



CFT Distance Conjecture





CFT Distance Conjecture

CFT Distance Conjecture



Conformal perturbation theory

$$\frac{d\gamma_{\ell}}{dt} = -C_{JJO}$$

$$\bigoplus$$

$$C_{JJO} \lesssim \gamma_{\ell} \text{ as } \gamma_{\ell} \to 0$$

(Weakly-broken) HS symmetry

II. Infinite distance → HS point

I. HS point → Infinite distance

Conformal perturbation theory

$$\frac{d\gamma_{\ell}}{dt} = -C_{JJO}$$

$$C_{JJ\mathcal{O}} \lesssim \gamma_{\ell} \text{ as } \gamma_{\ell}
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(+)

(Weakly-broken) HS symmetry

CFT Distance Conjecture



CFT Distance Conjecture

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II. Infinite distance → HS point
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 $\exists \mathcal{O} : C_{JJ\mathcal{O}} \neq 0$ **CFT Distance** Sufficient but **Criterion** not necessary

Finite distance



CFT Distance Conjecture

- II. Infinite distance → HS point
 - No HS symmetry
 - ↓ ?
 - $\exists \mathscr{O}: C_{JJ\mathscr{O}} \neq 0$
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 $\blacksquare. \ \gamma_{\ell} \sim e^{-\alpha_{\ell} t}$



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$$\blacksquare. \ \gamma_{\ell} \sim e^{-\alpha_{\ell} t}$$

$$\frac{d\gamma_{\ell}}{dt} \simeq -C_{KKO}\gamma_{\ell} \text{ as } \gamma_{\ell} \rightarrow$$
Conformal perturbation theorem

(+)(Weakly-broken) HS symmetry







CFT Distance Conjecture

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↓ ?

 $\exists \mathcal{O} : C_{JJ\mathcal{O}} \neq 0$ **CFT Distance** Sufficient but

n v not necessary

Finite distance

III.
$$\gamma_{\ell} \sim e^{-\alpha_{\ell} t}$$

$$\exists \mathcal{O} : C_{KKO}^{HS} \neq 0 ?$$

$$\downarrow^{d} \gamma_{\ell}$$

$$\frac{d\gamma_{\ell}}{dt} \simeq -C_{KKO} \gamma_{\ell} \text{ as } \gamma_{\ell} \rightarrow$$
Conformal perturbation theor
$$\bigoplus$$
(Weakly-broken) HS symmetric







Interesting questions for the future!

... But let me ask something else

Stringy nature of HS points **?**

CFT Distance Conjecture

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 - No HS symmetry

 $\exists \mathcal{O}: C_{II\mathcal{O}} \neq 0$

not necessary

Finite distance

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$$\gamma_{\ell} \sim e^{-\alpha_{\ell} t}$$

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Strings in the Conformal Manifold



- Inspiration: Emergent String Conjecture [Lee, Lerche, Weigand '19]
 - KK modes \rightarrow Decompactification
 - Excitations of weakly-coupled string

Strings in the Conformal Manifold



KK tower \rightarrow No HS fields





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Strings in the Conformal Manifold



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Problem: $M_s \lesssim R_{AdS}^{-1} \rightarrow$ String in a highly-curved background... hard to study!



- Inspiration: Emergent String Conjecture [Lee, Lerche, Weigand '19]

 - KK modes → Decompactification
 Excitations of weakly-coupled string



- Rely on CFT results and extract clues

In flat space: Value of $\alpha \rightarrow$ Nature of the tower

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$$\alpha = \sqrt{\frac{d-2+n}{n(d-2)}} \longrightarrow \begin{array}{l} \text{Decompactific} \\ n \text{ extra dime} \end{array}$$



In flat space: Value of $\alpha \rightarrow$ Nature of the tower

$$\alpha = \sqrt{\frac{d-2+n}{n(d-2)}} \longrightarrow \frac{\text{Decc}}{n \text{ ex}}$$



Caveat: Different values found for decompactification to running solution [Etheredge, Heidenreich, McNamara, Rudelius, Ruiz, Valenzuela '23]

In flat space: Value of $\alpha \rightarrow$ Nature of the tower

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Decompactification of *n* extra dimensions

- $\alpha = \frac{1}{\sqrt{d-2}} \longrightarrow$ Emergent string limit
- **Caveat:** Different values found for decompactification to running solution [Etheredge, Heidenreich, McNamara, Rudelius, Ruiz, Valenzuela '23]
- From the CFT: Restrict to zoo of 4d SCFTs with simple gauge group (Lagrangian) admitting large N
In flat space: Value of $\alpha \rightarrow$ Nature of the tower

$$\alpha = \sqrt{\frac{d-2+n}{n(d-2)}} \quad \blacksquare$$

From the CFT: Restrict to zoo of 4d SCFTs with simple gauge group (Lagrangian) admitting large N

- Decompactification of $\alpha = \frac{1}{\sqrt{d-2}} \rightarrow \frac{1}{\text{string limit}}$
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Three different values: $\alpha = \left\{ \sqrt{\frac{2}{3}}, \sqrt{\frac{7}{12}}, \frac{1}{\sqrt{2}} \right\}$ [Perlmutter, Rastelli, Vafa, Valenzuela '20]

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From the CFT: Restrict to zoo of 4d SCFTs with simple gauge group (Lagrangian) admitting large N

But...
$$\alpha \neq \frac{1}{\sqrt{3}}$$
 for all of them?

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Decompactificant *n* extra dimer

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From the CFT: Restrict to zoo of 4d SCFTs with simple gauge group (Lagrangian) admitting large N

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 $\left\{\frac{7}{12}, \frac{1}{\sqrt{2}}\right\}$ [Perlmutter, Rastelli, Vafa, Valenzuela '20] Suggests three different strings in AdS **y...** Match $n = \{3,4,6\}$ ctification to $D = \{8,9,11\}$?

In flat space: Value of $\alpha \rightarrow$ Nature of the tower

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From the CFT: Restrict to zoo of 4d SCFTs with simple gauge group (Lagrangian) admitting large N $\left\{\frac{7}{12}, \frac{1}{\sqrt{2}}\right\}$ [Perlmutter, Rastelli, Vafa, Valenzuela '20] Suggests three different strings in AdS **y...** Match $n = \{3,4,6\}$ **So...** What is going on?! ctification to $D = \{8,9,11\}$? Irene's puzzle

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From the CFT: Restrict to zoo of 4d SCFTs with simple gauge group (Lagrangian) admitting large N Three different values: $\alpha = \left\{ \left(\sqrt{\frac{2}{3}}, \sqrt{\frac{7}{12}}, \frac{1}{\sqrt{2}} \right\} \right\}$ [Perlmutter, Rastelli, Vafa, Valenzuela '20] E.g. $\mathcal{N} = 2 \text{ SCQCD}$

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From the CFT: Restrict to zoo of 4d SCFTs with simple gauge group (Lagrangian) admitting large N $\left(\frac{\overline{2}}{3}, \sqrt{\frac{7}{12}}, \frac{1}{\sqrt{2}}\right)$ [Perlmutter, Rastelli, Vafa, Valenzuela '20] Three different values: $\alpha = \langle \alpha \rangle$ E.g. $\mathcal{N} = 2$ SCQC E.g. $\mathcal{N} = 1$ SCQCD

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Type IIB on an 5-sphere

 $S = \frac{M_{Pl}^3}{2} \left[d^5 x \sqrt{-g} \left(R - (\partial \hat{\Phi})^2 - (\partial \hat{R})^2 - V(\hat{\Phi}, \hat{R}) \right) \right]$

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► Â











 $\mathcal{N} = 4$ SU(N) gauge theory in 4d

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Notice:

Convex hulls for AdS and CFT glue nicely together! (see later)



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Summary





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Relax condition

 $\Delta_{BPS} \sim \mathcal{O}(1) \longleftrightarrow M_{KK} R_{AdS} \sim \mathcal{O}(1)$

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ADC parameter [Lust, Palti, Vafa '19]

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Recap: 4d SCFTs with simple gauge group (Lagrangian) admitting large N

$$\alpha = \left\{ \sqrt{\frac{2}{3}}, \sqrt{\frac{7}{12}}, \frac{1}{\sqrt{2}} \right\} \quad [Pe]$$

erlmutter, Rastelli, Vafa, Valenzuela '20]

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EFT!

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Setup: $AdS_5 \times S^5/Z_k \leftrightarrow \mathcal{N} = 2$ necklace quivers

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 $s \leftrightarrow \mathcal{N} = 2$ necklace quivers

 \rightarrow S¹ of orbifold singularities



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 S^1 of orbifold singularities

A very peculiar limit:

Driven by only axions \rightarrow Typically finite distance



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String propagating in $AdS_5 \times S^1$!

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String propagating in AdS₅ × S¹! Candidate for new emergent string in AdS \mathbf{Z} [Baume, JCI '20]

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Thank you for your attention!

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