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# Tadpole & Moduli in Type lIB on the 1<sup>9</sup> LG Model

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#### **Outline**

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# Flux Compactification of IIB with  $h^{1,1}(M) = 0$

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	- Goal: Study the mirror dual setup flux compactifications of type IIB on spaces with  $h^{1,1}(M) = 0 \rightsquigarrow$  non-geometric.
	- Internal SCFT (with appropriate  $c = 9$ ) given by Landau-Ginzburg models no Kähler moduli to sabilize, and *all* complex structure moduli stabilized by a flux-induced superpotential.
- Becker, Gonzalo, Walcher, Wrase (2022): All supersymmetric flux vacua found in the paper above leave a large number of moduli massless!

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- Absence of Kähler moduli, and the non-renormalization of the flux induced superpotential make these models attractive. One can test various swampland conjectures in this strongly coupled region of the landscape.
- Our main interest is in studying the competition between the flux tadpole and the number of stabilized moduli in these models. An exhaustive test of the tadpole conjecture may be possible.



#### <span id="page-10-0"></span>2D,  $N = 2$  SCFT with  $c = 9$

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- Take *M* to be an  $N = 2$  LG QFT of chiral superfields:

 $S = \int d^2z d^4\theta K(\Phi_1, \bar{\Phi}_1, \dots, \Phi_n \bar{\Phi}_n) + (\int d^2z d^2\theta W(\Phi_1, \dots, \Phi_n) + \text{ h.c.})$ 



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- Tensor *N* of these together  $\rightsquigarrow$   $c = \frac{3Nk}{k+2}$ , which equals 9 if  $(k, N) = (1, 9), (2, 6), (3, 5), (6, 4).$  We will focus on the first model.

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#### The Model

• 1 <sup>⊗</sup>9/**Z**3: Tensor nine *k* = 1 minimal models with worlsheet superpotential  $W=\sum_{i=1}^9 \Phi_i^3$ , orbifold by the diagonal **Z**<sub>3</sub>,  $g: \Phi_i \mapsto \omega \Phi_i$ ,  $\omega=e^{2\pi i/3}$ .



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- Complex structure moduli: a basis given by the Z<sub>3</sub>-invariant monomials of the  $(c, c)$  chiral ring  $\mathbf{C}[\Phi_1, \dots, \Phi_9]/(\Phi_1^2, \dots, \Phi_9^2) \rightsquigarrow h^{2,1}(M) = 84.$



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### Branes and Fluxes in the LG Language I

• Supersymmetric cyles: A- and B-branes distinguished by boundary conditions of  $N = (2, 2)$  supercurrents: A: *G* ±  $L^{\pm}(z) = G_R^{\mp}$  $G^{\pm}_R(\bar{z})$  at  $z = \bar{z}$ ; B:  $G^{\pm}_L$  $g_L^{\pm}(z) = G_R^{\pm}$  $B^{\pm}_R(\bar{z})$  at  $z=\bar{z}$ .



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- Cohomology basis of the space of A-brane charges are spanned by the RR ground states  $|\ell\rangle$ ,  $\ell = 1, 2$ , which are equivalently represented by the chiral ring **C**[Φ]/Φ<sup>2</sup>. The correspondence is  $\ket{\ell} \leftrightarrow \#\Phi^{\ell-1}.$  To compute RR charges:  $\langle V_n | \ell \rangle \sim \int_{V_n} \Phi^{\ell-1} e^{-\Phi^3}$ .

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- 1 <sup>⊗</sup>9/**Z**3: The orbifold projects RR ground states to  $H^{(p,q)}(M) \ni \Omega_{\vec{\ell}} \leftrightarrow |\ell^1, \ldots, \ell^9\rangle$  with  $p+q=3$ ,  $\ell^i=1, 2, \sum_{i=1}^9 \ell^i=0 \mod 3$ .

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### Branes and Fluxes in the LG Language II

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The 1<sup>9</sup> Landau-Ginzburg model

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• The 3-form fluxes  $H_{NS}$  and  $F_{RR}$  are supported on A-branes. Define:  $G = F_{RR} - \tau H_{NS}$ ,  $\tau = C_0 + i e^{-\varphi}$ . For SUSY Minkowski vacua:  $G \in H^{(2,1)}(M)$ .



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- The orbifold of the tensor model also identifies brane states by summing over orbits. Suffice it to say, ∃ 128 independent 3-cycles after orientifolding, and we impose Dirac quantization condition on fluxes:  $\int_{\Gamma}$  *G* = *N* − *τM*, with *N*, *M* ∈ **Z**, with Γ's forming a basis of *H*<sub>3</sub>(*M*, **Z**).



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- $\bullet$  Tadpole cancellation:  $N_{\rm flux}+N_{\rm D3}=$  12, where  $N_{\rm flux}=\int_M F_{RR}\wedge H_{NS}.$

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### Properties of the Flux Tadpole

 $\bullet$  Take  $G = \sum_{l=1}^{63} B_l \Omega_l$ , and impose Dirac quantization on a basis of 3-cycles  $\rightarrow$  real and imaginary parts of  $B$ <sub>*I*</sub> get related to 126 integer variables, the flux quantum numbers – denote them by  $y_i, i = 1(1)$ 126.



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- It generates  $N_{\text{flux}} = \frac{81\sqrt{3}}{2 \text{ Im } \tau} \sum_{l=1}^{63} |B_l|^2 \xrightarrow{\tau=\omega} 81 \sum_{l=1}^{63} |B_l|^2 = Q_{ij} y_i y_j.$

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- To ensure *N*<sub>flux</sub> < 12, sufficient to restrict to up to 12-Ω solutions.
- Already in BBVW and BGWW, solutions with *N*<sub>flux</sub> ≤ 12 were found. The minimum value of  $N_{\text{flux}}$  among these solutions is 8. On: Is there a "shorter" solution?

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# Flux-induced Superpotential and the Rank of the Mass Matrix



• Gukov, Vafa, Witten (2000): For M/F/IIA compactifications on CY 4-folds, fluxes induce a spacetime superpotential  $W = \int_M G \wedge \Omega$ . This continues to hold for type IIB on 3-folds.



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• Burgess, Escoda, Quevedo (2006): The GVW superpotential in type IIB flux compactifications is non-renormalized in all orders of perturbation theory.

[Motivation](#page-2-0)

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- The Hessian *DDW* has the same rank *r* as the mass matrix (BGWW). *r* = the number of massive moduli. Using a "homogeneous" basis of 3-cycles yields simpler formulas for the matrix elements of *DDW*.

<span id="page-45-0"></span>

# Systematic Exploration of Solutions

- Tadpole cancellation  $\Rightarrow N_{\text{flux}} + N_{\text{D3}} = 12$ .
- It suffices to restrict search for fluxes with  $N_{\text{flux}}$  < 12.
- For each  $\Omega$ <sub>*I*</sub> turned on,  $N_{\text{flux}}$  increases by at least 1.
- Hence, restrict search to  $\leq$  12- $\Omega$ <sub>I</sub>'s turned on.
- Searching through all possible 12- $\Omega$  configurations computationally not affordable – we organize solutions with fewer Ω*<sup>I</sup>* 's turned on.
- *n* is a crude lower bound for *N*<sub>flux</sub> for the set of *n*-Ω flux backgrounds. We find the best lower bound, and the full set of fluxes that saturate it, for  $n = 1, 2, 3, 4.$

<span id="page-46-0"></span>

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# 1, 2, 3-Ω Solutions





#### <span id="page-47-0"></span>4-Ω Solutions

Minimum Flux Tadpole Contribution:  $T_{min} = 12$ First example within the physical bound.

$$
G=\frac{1}{3\sqrt{3}}\left(-\Omega_1+\Omega_{\vec{\ell}_{\alpha}}+\Omega_{\vec{\ell}_{\beta}}-\Omega_{\vec{\ell}_{\gamma}}\right)
$$

Family of Solutions:

$$
G=\pm\omega^r\frac{1}{3\sqrt{3}}\left(-\Omega_1+\omega^{\rho}\Omega_{\alpha}+\omega^{q-\rho}\Omega_{\beta}-\omega^q\Omega_{\gamma}\right)
$$

Quantization Condition:

$$
\frac{1}{3}(-\omega^a + \omega^b + \omega^c - \omega^d) = N - \omega M, \quad \text{iff} \quad (-a - d + b + c) \text{ mod } 3 = 0
$$

<span id="page-48-0"></span>

### 4-Ω Solutions



<span id="page-49-0"></span>

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## 8-Ω Solutions

$$
G = \frac{1}{9} \left( -\Omega_1 + \Omega_{\vec{\ell}_{a_2}} - \Omega_{\vec{\ell}_{a_3}} + \Omega_{\vec{\ell}_{a_4}} - \Omega_{\vec{\ell}_{a_5}} + \Omega_{\vec{\ell}_{a_6}} - \Omega_{\vec{\ell}_{a_7}} + \Omega_{\vec{\ell}_{a_8}} \right)
$$
  
Quantization Condition:  $(-\vec{\ell}_1 + \vec{\ell}_{a_2} - \vec{\ell}_{a_3} + \vec{\ell}_{a_4} - \vec{\ell}_{a_5} + \vec{\ell}_{a_6} - \vec{\ell}_{a_7} + \vec{\ell}_{a_8}) \text{ mod } 3 = 0$ 

Tadpole Contribution:

 $T_{\min}$  = 8

 $#$  Massive Scalars:

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<span id="page-50-0"></span>

#### The Shortest Vector Problem

- ∃ solutions with *N*flux = 8. Are their "shorter" fluxes? Answer: No.
- To rule out  $N_{\text{flux}} < T$ , we proceed as follows. (In our code we implement this for  $T = 7$ ).
- $\bullet$  Choose  $T$  complex  $B_i$ 's to be (possibly) non-zero and set the remaining to zero  $\rightsquigarrow$  a set of (126  $-$  2 $T)$  **R**-linear equations in the integer flux quanta  $y_i$ .
- $\bullet$  The set  $\{\boldsymbol{B}_{i_1},\ldots,\boldsymbol{B}_{i_T}\}\neq\boldsymbol{0}$  gives a reduced expression for  $\mathrm{N}_{\text{flux}}$

$$
N_{flux}^{red}=Q_{ij}\,c_i\,c_j\quad\text{where}\quad c_i\in\mathbb{Z},\,i=1,\ldots,2T
$$

and Q is symmetric, positive definite.

 $\bullet$  Is  $\mathrm{N}^{\mathrm{red}}_{\mathrm{flux}}(\vec{c})$   $\leq$  T for a choice of integers  $c_i$ ?

# <span id="page-51-0"></span>The Eigensieve Algorithm

- Define the level set in  $\mathbb{R}^{2T}$ :  $L_T = \{(x_1, \ldots, x_{2T}) : \text{N}^{\text{red}}_{\text{flux}}(\vec{x}) = T\} \subset \mathbb{R}^{2T}$ .
- $\bullet$  Points  $\vec{x}$  outside of the ellipsoid  $L_{\mathcal{T}}$  will have values,  $\mathrm{N}^{\mathrm{red}}_{\mathrm{flux}} > \mathcal{T}$ .
- Are there integer lattice points on  $L<sub>T</sub>$  or in its interior?
- Eigensieve: A combination of enumeration and sieving using the eigenvalues of  $Q$ ,  $\{\lambda_1, \ldots, \lambda_{2\tau}\}\$  where  $0 < \lambda_1 \leq \lambda_2 \leq \ldots \leq \lambda_{2\tau}\$  and normalized eigenvectors,  $\{\vec{v}_1, \ldots, \vec{v}_{2T}\}.$ First carve out a finite region of  $\mathbb{Z}^{27}$  using the smallest eigenvalue of *Q*. Lattice points  $\vec{\rho}$  in this region satisfying  $|\vec{\rho}.\vec{v}_i| > \sqrt{\frac{7}{\lambda}}$  $\frac{1}{\lambda_i}$ ,  $i = 1(1)$ 2*T* , all lie outside of  $L_{\mathcal{T}}$  are sieved off. Explicit evaluation of  $\text{N}^{\text{red}}_{\text{flux}}$  done on the remaining points.

<span id="page-52-0"></span>

<sup>9</sup> [Landau-Ginzburg model](#page-10-0) [Flux Tadpole, Massive Moduli](#page-33-0) [A Large Set of Solutions](#page-45-0) [The Shortest Vector](#page-50-0) [Conclusion](#page-52-0)



# Summary & Open Questions

• A systematic search of solutions with the lowest value of  $N_{\text{flux}}$ , organized by number of non-zero components, has been launched. All *N*flux-minimizing solutions up to four Ω*<sup>I</sup>* 's turned on found.

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- $\bullet$  The SVP for the lattice of supersymmetric flux vacua in 1 $^9/Z_3$  solved.  $(N_{\text{flux}})_{\text{min}} = 8$ . A large class of 8- $\Omega$  solutions found to saturate this.

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- Finally, we aim to extend all our analyses to other Gepner models.

#### *Thank You!*

 $\bullet$