

A Large Set of Solutions

The Shortest Vector

Conclusior

Tadpole & Moduli in Type IIB on the 1⁹ LG Model

Nathan Brady, Anindya Sengupta (Based on 2310.00770, with Katrin Becker)

Department of Physics & Astronomy Texas A&M University

Brady, Sengupta

Deconstructing the String Landscape, IPhT, Nov 29, 2023



A Large Set of Solutions

The Shortest Vector

Conclusion

Outline

Motivation

- 2 The 1⁹ Landau-Ginzburg model
- 📵 Flux Tadpole, Massive Moduli
- 4 A Large Set of Solutions
- 6 The Shortest Vector

6 Conclusion



A Large Set of Solutions

The Shortest Vector

Conclusion o

Flux Compactification of IIB with $h^{1,1}(M) = 0$

• Fluxes in string theory compactifications are important in myriad ways: partial SUSY breaking, generating large hierarchies, stabilizing moduli, ...



The Shortest Vector

Conclusion o

- Fluxes in string theory compactifications are important in myriad ways: partial SUSY breaking, generating large hierarchies, stabilizing moduli, ...
- DeWolfe, Giryavets, Kachru, Taylor (2005): Massive type IIA on a rigid CY $(h^{2,1}(M) = 0)$ with fluxes and O6 planes \rightsquigarrow all moduli stabilized.



The Shortest Vector

Conclusion o

- Fluxes in string theory compactifications are important in myriad ways: partial SUSY breaking, generating large hierarchies, stabilizing moduli, ...
- DeWolfe, Giryavets, Kachru, Taylor (2005): Massive type IIA on a rigid CY $(h^{2,1}(M) = 0)$ with fluxes and O6 planes \rightsquigarrow all moduli stabilized.
- Becker, Becker, Vafa, Walcher (2006):
 - Goal: Study the mirror dual setup flux compactifications of type IIB on spaces with h^{1,1}(M) = 0 → non-geometric.



Conclusion

- Fluxes in string theory compactifications are important in myriad ways: partial SUSY breaking, generating large hierarchies, stabilizing moduli, ...
- DeWolfe, Giryavets, Kachru, Taylor (2005): Massive type IIA on a rigid CY $(h^{2,1}(M) = 0)$ with fluxes and O6 planes \rightsquigarrow all moduli stabilized.
- Becker, Becker, Vafa, Walcher (2006):
 - Goal: Study the mirror dual setup flux compactifications of type IIB on spaces with $h^{1,1}(M) = 0 \rightsquigarrow$ non-geometric.
 - Internal SCFT (with appropriate c = 9) given by Landau-Ginzburg models no Kähler moduli to sabilize, and *all* complex structure moduli stabilized by a flux-induced superpotential.
- Becker, Gonzalo, Walcher, Wrase (2022): All supersymmetric flux vacua found in the paper above leave a large number of moduli massless!



A Large Set of Solutions

The Shortest Vector

Conclusion o

Flux Compactification of IIB with $h^{1,1}(M) = 0$

 BBVW did not make a systematic search for solutions. (e.g., they reported only Minkowski solutions in the 1⁹ model; BGWW found AdS soltions as well).



A Large Set of Solutions

The Shortest Vector

Conclusion

- BBVW did not make a systematic search for solutions. (e.g., they reported only Minkowski solutions in the 1⁹ model; BGWW found AdS soltions as well).
- Lattice of supersymmetric, Dirac quantized, flux backgrounds in the non-geometric LG models have large rank brute force search for "short" vectors is expensive.



A Large Set of Solutions

The Shortest Vector

Conclusion

- BBVW did not make a systematic search for solutions. (e.g., they reported only Minkowski solutions in the 1⁹ model; BGWW found AdS soltions as well).
- Lattice of supersymmetric, Dirac quantized, flux backgrounds in the non-geometric LG models have large rank brute force search for "short" vectors is expensive.
- Absence of Kähler moduli, and the non-renormalization of the flux induced superpotential make these models attractive. One can test various swampland conjectures in this strongly coupled region of the landscape.



Conclusion

- BBVW did not make a systematic search for solutions. (e.g., they reported only Minkowski solutions in the 1⁹ model; BGWW found AdS soltions as well).
- Lattice of supersymmetric, Dirac quantized, flux backgrounds in the non-geometric LG models have large rank brute force search for "short" vectors is expensive.
- Absence of Kähler moduli, and the non-renormalization of the flux induced superpotential make these models attractive. One can test various swampland conjectures in this strongly coupled region of the landscape.
- Our main interest is in studying the competition between the flux tadpole and the number of stabilized moduli in these models. An exhaustive test of the tadpole conjecture may be possible.



The Shortest Vector

Conclusion

2D, *N* = 2 SCFT with *c* = 9

• Setting: Realize total c = 15 via $M_4 \times M$, where M is a c = 9, N = 2 SCFT.



A Large Set of Solutions

The Shortest Vector

Conclusion

2D, N = 2 SCFT with c = 9

- Setting: Realize total c = 15 via $M_4 \times M$, where M is a c = 9, N = 2 SCFT.
- Take *M* to be an N = 2 LG QFT of chiral superfields:

 $S = \int d^2z d^4\theta K(\Phi_1, \bar{\Phi}_1, \dots, \Phi_n \bar{\Phi}_n) + (\int d^2z d^2\theta \mathcal{W}(\Phi_1, \dots, \Phi_n) + \text{ h.c.})$



A Large Set of Solutions

The Shortest Vector

Conclusion

2D, N = 2 SCFT with c = 9

- Setting: Realize total c = 15 via $M_4 \times M$, where M is a c = 9, N = 2 SCFT.
- Take *M* to be an *N* = 2 LG QFT of chiral superfields: $S = \int d^2z d^4\theta K(\Phi_1, \bar{\Phi}_1, \dots, \Phi_n \bar{\Phi}_n) + (\int d^2z d^2\theta W(\Phi_1, \dots, \Phi_n) + \text{ h.c.})$
- It flows in the infrared to a CFT provided K is chosen suitably. W only
 receives wavefunction renormalization. Choose W to be
 quasi-homogeneous polynomial, and absorb this in an overall rescaling.



A Large Set of Solutions

The Shortest Vector

Conclusion

2D, N = 2 SCFT with c = 9

- Setting: Realize total c = 15 via $M_4 \times M$, where M is a c = 9, N = 2 SCFT.
- Take *M* to be an *N* = 2 LG QFT of chiral superfields: $S = \int d^2z d^4\theta K(\Phi_1, \bar{\Phi}_1, \dots, \Phi_n \bar{\Phi}_n) + (\int d^2z d^2\theta W(\Phi_1, \dots, \Phi_n) + \text{ h.c.})$
- It flows in the infrared to a CFT provided *K* is chosen suitably. *W* only receives wavefunction renormalization. Choose *W* to be quasi-homogeneous polynomial, and absorb this in an overall rescaling.
- Lerche, Vafa, Warner (1989): For one Φ with $\mathcal{W} = \Phi^{k+2}$, central charge of the infrared fixed point is $c_k = \frac{3k}{k+2} \rightsquigarrow$ it is the k^{th} minimal model of the N = 2 superconformal algebra.



The Shortest Vector

Conclusion

2D, N = 2 SCFT with c = 9

- Setting: Realize total c = 15 via $M_4 \times M$, where M is a c = 9, N = 2 SCFT.
- Take *M* to be an *N* = 2 LG QFT of chiral superfields: $S = \int d^2z d^4\theta K(\Phi_1, \bar{\Phi}_1, \dots, \Phi_n \bar{\Phi}_n) + (\int d^2z d^2\theta W(\Phi_1, \dots, \Phi_n) + \text{ h.c.})$
- It flows in the infrared to a CFT provided *K* is chosen suitably. *W* only receives wavefunction renormalization. Choose *W* to be quasi-homogeneous polynomial, and absorb this in an overall rescaling.
- Lerche, Vafa, Warner (1989): For one Φ with $\mathcal{W} = \Phi^{k+2}$, central charge of the infrared fixed point is $c_k = \frac{3k}{k+2} \rightsquigarrow$ it is the k^{th} minimal model of the N = 2 superconformal algebra.
- Tensor *N* of these together $\rightsquigarrow c = \frac{3Nk}{k+2}$, which equals 9 if (k, N) = (1, 9), (2, 6), (3, 5), (6, 4). We will focus on the first model.



A Large Set of Solutions

The Shortest Vector

Conclusion o

The Model

Brady, Sengupta



A Large Set of Solutions

The Model

• $1^{\otimes 9}/\mathbb{Z}_3$: Tensor nine k = 1 minimal models with worlsheet superpotential $\mathcal{W} = \sum_{i=1}^{9} \Phi_i^3$, orbifold by the diagonal \mathbb{Z}_3 , $g : \Phi_i \mapsto \omega \Phi_i$, $\omega = e^{2\pi i/3}$.



- $1^{\otimes 9}/\mathbb{Z}_3$: Tensor nine k = 1 minimal models with worlsheet superpotential $\mathcal{W} = \sum_{i=1}^{9} \Phi_{i}^{3}$, orbifold by the diagonal **Z**₃, $\boldsymbol{g} : \Phi_{i} \mapsto \omega \Phi_{i}, \omega = \boldsymbol{e}^{2\pi i/3}$.
- Complex structure moduli: a basis given by the Z₃-invariant monomials of the (c, c) chiral ring $\mathbf{C}[\Phi_1, ..., \Phi_9]/(\Phi_1^2, ..., \Phi_9^2) \rightsquigarrow h^{2,1}(M) = 84$.



- $1^{\otimes 9}/\mathbb{Z}_3$: Tensor nine k = 1 minimal models with worlsheet superpotential $\mathcal{W} = \sum_{i=1}^{9} \Phi_i^3$, orbifold by the diagonal \mathbb{Z}_3 , $g : \Phi_i \mapsto \omega \Phi_i$, $\omega = e^{2\pi i/3}$.
- Complex structure moduli: a basis given by the Z₃-invariant monomials of the (*c*, *c*) chiral ring C[Φ₁,...,Φ₉]/(Φ²₁,...,Φ²₉) → h^{2,1}(M) = 84.
- Kähler moduli: no non-trivial (a, c) primaries in the **Z**₃ orbifold \rightsquigarrow $h^{1,1}(M) = 0$.



A Large Set of Solutions

The Shortest Vector

Conclusion

- $1^{\otimes 9}/\mathbb{Z}_3$: Tensor nine k = 1 minimal models with worlsheet superpotential $\mathcal{W} = \sum_{i=1}^{9} \Phi_i^3$, orbifold by the diagonal \mathbb{Z}_3 , $g : \Phi_i \mapsto \omega \Phi_i$, $\omega = e^{2\pi i/3}$.
- Complex structure moduli: a basis given by the Z₃-invariant monomials of the (*c*, *c*) chiral ring C[Φ₁,...,Φ₉]/(Φ²₁,...,Φ²₉) → h^{2,1}(M) = 84.
- Kähler moduli: no non-trivial (a, c) primaries in the **Z**₃ orbifold $\rightarrow h^{1,1}(M) = 0$.
- Orientifold: To get N = 1 SUSY in 4D, take an orientifold. We choose: Worldsheet parity operator dressed with g_1 : $g_1 : (\Phi_1, \Phi_2, \Phi_3, \Phi_4, \Phi_5, \Phi_6, \Phi_7, \Phi_8, \Phi_9) \mapsto -(\Phi_2, \Phi_1, \Phi_3, \Phi_4, \Phi_5, \Phi_6, \Phi_7, \Phi_8, \Phi_9).$



- $1^{\otimes 9}/\mathbb{Z}_3$: Tensor nine k = 1 minimal models with worlsheet superpotential $\mathcal{W} = \sum_{i=1}^{9} \Phi_i^3$, orbifold by the diagonal \mathbb{Z}_3 , $g : \Phi_i \mapsto \omega \Phi_i$, $\omega = e^{2\pi i/3}$.
- Complex structure moduli: a basis given by the Z₃-invariant monomials of the (*c*, *c*) chiral ring C[Φ₁,...,Φ₉]/(Φ²₁,...,Φ²₉) → h^{2,1}(M) = 84.
- Kähler moduli: no non-trivial (a, c) primaries in the **Z**₃ orbifold \rightsquigarrow $h^{1,1}(M) = 0$.
- Orientifold: To get N = 1 SUSY in 4D, take an orientifold. We choose: Worldsheet parity operator dressed with g₁: g₁ : (Φ₁, Φ₂, Φ₃, Φ₄, Φ₅, Φ₆, Φ₇, Φ₈, Φ₉) → -(Φ₂, Φ₁, Φ₃, Φ₄, Φ₅, Φ₆, Φ₇, Φ₈, Φ₉).
 ⇒ h^{2,1}(M)⁺ = 63.



A Large Set of Solutions

The Shortest Vector

Conclusior



A Large Set of Solutions

The Shortest Vector

Conclusion o

Branes and Fluxes in the LG Language I

Supersymmetric cyles: A- and B-branes distinguished by boundary conditions of N = (2,2) supercurrents:
 A: G[±]_L(z) = G[±]_R(z̄) at z = z̄; B: G[±]_L(z) = G[±]_R(z̄) at z = z̄.



A Large Set of Solutions

The Shortest Vector

Conclusion o

- Supersymmetric cyles: A- and B-branes distinguished by boundary conditions of N = (2, 2) supercurrents:
 A: G[±]_L(z) = G[∓]_R(z̄) at z = z̄; B: G[±]_L(z) = G[±]_R(z̄) at z = z̄.
- A-branes in the Φ -space are preimages of $\text{Im}\mathcal{W} = 0$.



A Large Set of Solutions

The Shortest Vector

Conclusion

- Supersymmetric cyles: A- and B-branes distinguished by boundary conditions of N = (2, 2) supercurrents:
 A: G[±]_L(z) = G[∓]_R(z̄) at z = z̄; B: G[±]_L(z) = G[±]_R(z̄) at z = z̄.
- A-branes in the Φ -space are preimages of $\mathrm{Im}\mathcal{W} = 0$.
- The building block minimal model at hand has three A-branes V_0 , V_1 , V_2 :





urg model Flux Tadpole

A Large Set of Solut

The Shortest Vector

Conclusion

- Supersymmetric cyles: A- and B-branes distinguished by boundary conditions of N = (2, 2) supercurrents:
 A: G[±]_L(z) = G[±]_R(z̄) at z = z̄;
 B: G[±]_L(z) = G[±]_R(z̄) at z = z̄.
- A-branes in the Φ -space are preimages of $\text{Im}\mathcal{W} = 0$.
- The building block minimal model at hand has three A-branes V_0 , V_1 , V_2 : $V_0 + V_1 + V_2 = 0$.



A Large Set of Solutions

The Shortest Vector

Conclusion

Branes and Fluxes in the LG Language I

• Supersymmetric cyles: A- and B-branes distinguished by boundary conditions of N = (2, 2) supercurrents:

A: $G_L^{\pm}(z) = G_R^{\mp}(\bar{z})$ at $z = \bar{z}$; B: $G_L^{\pm}(z) = G_R^{\pm}(\bar{z})$ at $z = \bar{z}$.

- A-branes in the $\Phi\text{-space}$ are preimages of $\text{Im}\mathcal{W}=0.$
- The building block minimal model at hand has three A-branes V_0 , V_1 , V_2 : $V_0 + V_1 + V_2 = 0$.
- Cohomology basis of the space of A-brane charges are spanned by the RR ground states $|\ell\rangle$, $\ell = 1, 2$, which are equivalently represented by the chiral ring $\mathbf{C}[\Phi]/\Phi^2$. The correspondence is $|\ell\rangle \leftrightarrow \#\Phi^{\ell-1}$. To compute RR charges: $\langle V_n | \ell \rangle \sim \int_{V_n} \Phi^{\ell-1} e^{-\Phi^3}$.



A Large Set of Solutions

The Shortest Vector

Conclusion o

Branes and Fluxes in the LG Language I

• Supersymmetric cyles: A- and B-branes distinguished by boundary conditions of N = (2, 2) supercurrents:

A: $G_L^{\pm}(z) = G_R^{\mp}(\bar{z})$ at $z = \bar{z}$; B: $G_L^{\pm}(z) = G_R^{\pm}(\bar{z})$ at $z = \bar{z}$.

- A-branes in the $\Phi\text{-space}$ are preimages of $\text{Im}\mathcal{W}=0.$
- The building block minimal model at hand has three A-branes V_0 , V_1 , V_2 : $V_0 + V_1 + V_2 = 0$.
- Cohomology basis of the space of A-brane charges are spanned by the RR ground states $|\ell\rangle$, $\ell = 1, 2$, which are equivalently represented by the chiral ring $\mathbf{C}[\Phi]/\Phi^2$. The correspondence is $|\ell\rangle \leftrightarrow \#\Phi^{\ell-1}$. To compute RR charges: $\langle V_n | \ell \rangle \sim \int_{V_n} \Phi^{\ell-1} e^{-\Phi^3}$.
- $1^{\otimes 9}/\mathbb{Z}_3$: The orbifold projects RR ground states to $H^{(p,q)}(M) \ni \Omega_{\vec{\ell}} \leftrightarrow |\ell^1, \dots, \ell^9\rangle$ with p+q=3, $\ell^i=1,2$, $\sum_{i=1}^9 \ell^i=0 \mod 3$.



The 1⁹ Landau-Ginzburg model

Flux Tadpole, Massive Moduli

A Large Set of Solutions

The Shortest Vector

Conclusion

Branes and Fluxes in the LG Language II

| $\sum_{i} \ell^{i}$ | 9 | 12 | 15 | 18 |
|---------------------|-------------|-------------|-------------|-------------|
| $H^{(p,q)}$ | $H^{(3,0)}$ | $H^{(2,1)}$ | $H^{(1,2)}$ | $H^{(0,3)}$ |



Branes and Fluxes in the LG Language II

• The "3-forms" are classified as:

| $\sum_{i} \ell^{i}$ | 9 | 12 | 15 | 18 |
|---------------------------|-------------|-------------|-------------|-------------|
| H ^(p,q) | $H^{(3,0)}$ | $H^{(2,1)}$ | $H^{(1,2)}$ | $H^{(0,3)}$ |

• The 3-form fluxes H_{NS} and F_{BB} are supported on A-branes. Define: $G = F_{BB} - \tau H_{NS}$, $\tau = C_0 + ie^{-\varphi}$. For SUSY Minkowski vacua: $G \in H^{(2,1)}(M)$.



A Large Set of Solutions

The Shortest Vector

Conclusion

Branes and Fluxes in the LG Language II

| $\sum_{i} \ell^{i}$ | 9 | 12 | 15 | 18 |
|---------------------|-------------|-------------|--------------------|-------------|
| $H^{(p,q)}$ | $H^{(3,0)}$ | $H^{(2,1)}$ | H ^(1,2) | $H^{(0,3)}$ |

- The 3-form fluxes H_{NS} and F_{RR} are supported on A-branes. Define: $G = F_{RR} - \tau H_{NS}$, $\tau = C_0 + ie^{-\varphi}$. For SUSY Minkowski vacua: $G \in H^{(2,1)}(M)$.
- The orbifold of the tensor model also identifies brane states by summing over orbits. Suffice it to say, \exists 128 independent 3-cycles after orientifolding, and we impose Dirac quantization condition on fluxes: $\int_{\Gamma} G = N - \tau M$, with $N, M \in \mathbb{Z}$, with Γ 's forming a basis of $H_3(M, \mathbb{Z})$.



The 1⁹ Landau-Ginzburg model

A Large Set of Solutions

The Shortest Vector

Conclusion o

Branes and Fluxes in the LG Language II

| $\sum_{i} \ell^{i}$ | 9 | 12 | 15 | 18 |
|---------------------|-------------|-------------|-------------|-------------|
| $H^{(p,q)}$ | $H^{(3,0)}$ | $H^{(2,1)}$ | $H^{(1,2)}$ | $H^{(0,3)}$ |

- The 3-form fluxes H_{NS} and F_{RR} are supported on A-branes. Define: $G = F_{RR} - \tau H_{NS}$, $\tau = C_0 + ie^{-\varphi}$. For SUSY Minkowski vacua: $G \in H^{(2,1)}(M)$.
- The orbifold of the tensor model also identifies brane states by summing over orbits. Suffice it to say, \exists 128 independent 3-cycles after orientifolding, and we impose Dirac quantization condition on fluxes: $\int_{\Gamma} G = N - \tau M$, with $N, M \in \mathbb{Z}$, with Γ 's forming a basis of $H_3(M, \mathbb{Z})$.
- Can construct B-branes using matrix factorization [Kapustin, Li; Brunner, Herbst, Lerche, Scheuner (2003)] \rightsquigarrow useful to compute D3-brane charge of O-planes. For the orientifold we consider, Q_3 (O-plane) = 12.



The 1⁹ Landau-Ginzburg model

Flux Tadpole, Massive Moduli

A Large Set of Solutions

The Shortest Vector

Conclusion

Branes and Fluxes in the LG Language II

| $\sum_{i} \ell^{i}$ | 9 | 12 | 15 | 18 |
|---------------------|-------------|-------------|-------------|-------------|
| $H^{(p,q)}$ | $H^{(3,0)}$ | $H^{(2,1)}$ | $H^{(1,2)}$ | $H^{(0,3)}$ |

- The 3-form fluxes H_{NS} and F_{RR} are supported on A-branes. Define: $G = F_{RR} - \tau H_{NS}$, $\tau = C_0 + ie^{-\varphi}$. For SUSY Minkowski vacua: $G \in H^{(2,1)}(M)$.
- The orbifold of the tensor model also identifies brane states by summing over orbits. Suffice it to say, \exists 128 independent 3-cycles after orientifolding, and we impose Dirac quantization condition on fluxes: $\int_{\Gamma} G = N - \tau M$, with $N, M \in \mathbb{Z}$, with Γ 's forming a basis of $H_3(M, \mathbb{Z})$.
- Can construct B-branes using matrix factorization [Kapustin, Li; Brunner, Herbst, Lerche, Scheuner (2003)] \rightsquigarrow useful to compute D3-brane charge of O-planes. For the orientifold we consider, Q_3 (O-plane) = 12.
- Tadpole cancellation: $N_{\rm flux} + N_{\rm D3} =$ 12, where $N_{\rm flux} = \int_M F_{RR} \wedge H_{NS}$.



A Large Set of Solutions

The Shortest Vector

Conclusion o

Properties of the Flux Tadpole

Brady, Sengupta



The Shortest Vector

Conclusion

Properties of the Flux Tadpole

• Take $G = \sum_{l=1}^{63} B_l \Omega_l$, and impose Dirac quantization on a basis of 3-cycles \rightsquigarrow real and imaginary parts of B_l get related to 126 integer variables , the flux quantum numbers – denote them by y_i , i = 1(1)126.



A Large Set of Solutions

The Shortest Vector

Conclusion o

Properties of the Flux Tadpole

• Take $G = \sum_{l=1}^{63} B_l \Omega_l$, and impose Dirac quantization on a basis of 3-cycles \rightsquigarrow real and imaginary parts of B_l get related to 126 integer variables , the flux quantum numbers – denote them by y_i , i = 1(1)126.

• It generates
$$N_{\text{flux}} = \frac{81\sqrt{3}}{2 \operatorname{Im}\tau} \sum_{l=1}^{63} |B_l|^2 \stackrel{\tau=\omega}{==} 81 \sum_{l=1}^{63} |B_l|^2 = Q_{ij} y_i y_j.$$



The Shortest Vector

Conclusion

Properties of the Flux Tadpole

- Take $G = \sum_{l=1}^{63} B_l \Omega_l$, and impose Dirac quantization on a basis of 3-cycles \rightarrow real and imaginary parts of B_l get related to 126 integer variables , the flux quantum numbers denote them by y_i , i = 1(1)126.
- It generates $N_{\text{flux}} = \frac{81\sqrt{3}}{2 \text{ Im}\tau} \sum_{l=1}^{63} |B_l|^2 \xrightarrow{\tau=\omega} 81 \sum_{l=1}^{63} |B_l|^2 = Q_{ij} y_i y_j.$
- The coefficients Q are **Z**-valued. Moreover, for each I, $81|B_I|^2$ is a homogenous quadratic in y_i with coefficients in **Z** \Rightarrow turning on an Ω_I must contribute at least 1 to N_{flux} .



The Shortest Vector

Conclusion

Properties of the Flux Tadpole

- Take $G = \sum_{l=1}^{63} B_l \Omega_l$, and impose Dirac quantization on a basis of 3-cycles \rightarrow real and imaginary parts of B_l get related to 126 integer variables , the flux quantum numbers denote them by y_i , i = 1(1)126.
- It generates $N_{\text{flux}} = \frac{81\sqrt{3}}{2 \text{ Im}\tau} \sum_{l=1}^{63} |B_l|^2 \xrightarrow{\tau=\omega} 81 \sum_{l=1}^{63} |B_l|^2 = Q_{ij} y_i y_j.$
- The coefficients Q are **Z**-valued. Moreover, for each I, $81|B_I|^2$ is a homogenous quadratic in y_i with coefficients in **Z** \Rightarrow turning on an Ω_I must contribute at least 1 to N_{flux} .
- To ensure $N_{\rm flux} \leq$ 12, sufficient to restrict to up to 12- Ω solutions.

The Shortest Vector

Conclusion

Properties of the Flux Tadpole

- Take $G = \sum_{l=1}^{63} B_l \Omega_l$, and impose Dirac quantization on a basis of 3-cycles \rightarrow real and imaginary parts of B_l get related to 126 integer variables , the flux quantum numbers denote them by y_i , i = 1(1)126.
- It generates $N_{\text{flux}} = \frac{81\sqrt{3}}{2 \text{ Im}\tau} \sum_{l=1}^{63} |B_l|^2 \xrightarrow{\tau=\omega} 81 \sum_{l=1}^{63} |B_l|^2 = Q_{ij} y_i y_j.$
- The coefficients Q are **Z**-valued. Moreover, for each I, $81|B_I|^2$ is a homogenous quadratic in y_i with coefficients in **Z** \Rightarrow turning on an Ω_I must contribute at least 1 to N_{flux} .
- To ensure $N_{\rm flux} \leq$ 12, sufficient to restrict to up to 12- Ω solutions.
- Already in BBVW and BGWW, solutions with $N_{\rm flux} \leq 12$ were found. The minimum value of $N_{\rm flux}$ among these solutions is 8. Qn: Is there a "shorter" solution?



A Large Set of Solutions

The Shortest Vector

Conclusion o

Flux-induced Superpotential and the Rank of the Mass Matrix



A Large Set of Solutions

The Shortest Vector

Conclusion

Flux-induced Superpotential and the Rank of the Mass Matrix

• Gukov, Vafa, Witten (2000): For M/F/IIA compactifications on CY 4-folds, fluxes induce a spacetime superpotential $W = \int_M G \wedge \Omega$. This continues to hold for type IIB on 3-folds.



- Gukov, Vafa, Witten (2000): For M/F/IIA compactifications on CY 4-folds, fluxes induce a spacetime superpotential $W = \int_M G \wedge \Omega$. This continues to hold for type IIB on 3-folds.
- Burgess, Escoda, Quevedo (2006): The GVW superpotential in type IIB flux compactifications is non-renormalized in all orders of perturbation theory.

Flux-induced Superpotential and the Rank of the Mass Matrix

- Gukov, Vafa, Witten (2000): For M/F/IIA compactifications on CY 4-folds, fluxes induce a spacetime superpotential $W = \int_M G \wedge \Omega$. This continues to hold for type IIB on 3-folds.
- Burgess, Escoda, Quevedo (2006): The GVW superpotential in type IIB flux compactifications is non-renormalized in all orders of perturbation theory.
- Considering a type IIB D5-brane which wraps a 3-cycle in internal *M*, and is a domain wall in spacetime, one can derive the superpotential, and also show that it receives no non-perturbative corrections beyond tree level.

Flux-induced Superpotential and the Rank of the Mass Matrix

- Gukov, Vafa, Witten (2000): For M/F/IIA compactifications on CY 4-folds, fluxes induce a spacetime superpotential $W = \int_M G \wedge \Omega$. This continues to hold for type IIB on 3-folds.
- Burgess, Escoda, Quevedo (2006): The GVW superpotential in type IIB flux compactifications is non-renormalized in all orders of perturbation theory.
- Considering a type IIB D5-brane which wraps a 3-cycle in internal *M*, and is a domain wall in spacetime, one can derive the superpotential, and also show that it receives no non-perturbative corrections beyond tree level.
- Dependence of *W* on complex structure moduli $t^{\vec{\ell}}$: $\int_{\Gamma} \Omega = \int_{\Gamma} \exp -W(t^{\vec{\ell}})$, where $W(t^{\vec{\ell}}) = \sum_{i=1}^{9} \Phi_i^3 \sum_{\vec{\ell}} t^{\vec{\ell}} \vec{\Phi}^{\vec{\ell}-\vec{1}}$.

Flux-induced Superpotential and the Rank of the Mass Matrix

- Gukov, Vafa, Witten (2000): For M/F/IIA compactifications on CY 4-folds, fluxes induce a spacetime superpotential $W = \int_M G \wedge \Omega$. This continues to hold for type IIB on 3-folds.
- Burgess, Escoda, Quevedo (2006): The GVW superpotential in type IIB flux compactifications is non-renormalized in all orders of perturbation theory.
- Considering a type IIB D5-brane which wraps a 3-cycle in internal *M*, and is a domain wall in spacetime, one can derive the superpotential, and also show that it receives no non-perturbative corrections beyond tree level.
- Dependence of *W* on complex structure moduli $t^{\vec{\ell}}$: $\int_{\Gamma} \Omega = \int_{\Gamma} \exp -\mathcal{W}(t^{\vec{\ell}})$, where $\mathcal{W}(t^{\vec{\ell}}) = \sum_{i=1}^{9} \Phi_i^3 \sum_{\vec{\ell}} t^{\vec{\ell}} \vec{\Phi}^{\vec{\ell}-\vec{1}}$.
- The Hessian *DDW* has the same rank r as the mass matrix (BGWW). r = the number of massive moduli. Using a "homogeneous" basis of 3-cycles yields simpler formulas for the matrix elements of *DDW*.



The Shortest Vector

Conclusion o

Systematic Exploration of Solutions

- Tadpole cancellation $\Rightarrow N_{\text{flux}} + N_{\text{D3}} = 12$.
- It suffices to restrict search for fluxes with $N_{\rm flux} \leq$ 12.
- For each Ω_l turned on, N_{flux} increases by at least 1.
- Hence, restrict search to \leq 12- Ω_{l} 's turned on.
- Searching through all possible $12-\Omega$ configurations computationally not affordable we organize solutions with fewer Ω_{I} 's turned on.
- *n* is a crude lower bound for N_{flux} for the set of $n-\Omega$ flux backgrounds. We find the best lower bound, and the full set of fluxes that saturate it, for n = 1, 2, 3, 4.



1, 2, $3-\Omega$ Solutions

| $oldsymbol{G} = \sum \mathrm{B} \Omega$ | T _{min} | Index Range | # Massive Scalars |
|--|------------------|---|----------------------------|
| $G = \frac{1}{\sqrt{2}}\Omega_{\mathrm{I}}$ | 27 | $\begin{array}{l} \alpha = 1, \dots, 35\\ \alpha = 57, \dots, 63 \end{array}$ | 16 22 |
| V S | 54 | $A = 36, \dots, 63$ | 22 |
| $G = \frac{i}{2} (\Omega_1 - \Omega_2)$ | 18 | $\alpha = 2, \ldots, 35$ | 16, 24, or 26 |
| $\mathbf{G} = \frac{1}{3} (\mathbf{M}_1 + \mathbf{M}_{\alpha})$ | 10 | $\alpha = 57, \ldots, 63$ | 28 or 32 |
| $oldsymbol{G} = rac{\mathrm{i}}{3} \left(\Omega_1 + \Omega_lpha + \Omega_eta ight)$ | | $\alpha,\beta=2,\ldots,35,$ | ∫16 20 22 24 28 29 30 |
| or | 27 | $57, \ldots, 63$ | |
| $G=rac{1}{3}\left(\Omega_{57}+\Omega_{lpha}+\Omega_{eta} ight)$ | | $\alpha, \beta = 58, \dots, 63$ | 32, 34, 36, 38, 40, 42, 46 |



au-Ginzburg model Fl

Flux Tadpole, Massive Moduli

A Large Set of Solutions

The Shortest Vector

Conclusion

$4-\Omega$ Solutions

 $\begin{array}{ll} \mbox{Minimum Flux Tadpole Contribution:} & T_{min} = 12 \\ \mbox{First example within the physical bound.} \end{array}$

$$G = rac{1}{3\sqrt{3}} \left(-\Omega_1 + \Omega_{ec{\ell}_lpha} + \Omega_{ec{\ell}_eta} - \Omega_{ec{\ell}_\gamma}
ight)$$

Family of Solutions:

$$G=\pm \omega^{r}rac{1}{3\sqrt{3}}\left(-\Omega_{1}+\omega^{
ho}\Omega_{lpha}+\omega^{q-
ho}\Omega_{eta}-\omega^{q}\Omega_{\gamma}
ight)$$

Quantization Condition:

$$\frac{1}{3}(-\omega^a+\omega^b+\omega^c-\omega^d)=N-\omega M,\qquad\text{iff}\qquad(-a-d+b+c)\,\,\,\text{mod}\,\,3=0$$



A Large Set of Solutions

The Shortest Vector

Conclusion

$4-\Omega$ Solutions

| # Massive Scalars | (α, β, γ) | | | |
|-------------------|--|---|------------------------------|---|
| | (2,6,8) (2,7,9) | (3,5,8) (3,7,10) | (4,5,9) (5,12,17) | |
| 12 | (2, 12, 14) (2, 13, 15) | (3, 11, 14) (3, 13, 16) | (5, 13, 18) (6, 13, 19) | |
| | (2,23,25) | (3, 21, 24) (3, 23, 26) | (5,22,27) (7,22,29) | |
| | (2,33,34) (2,29,30) | (3,32,34) (3,28,30) | (4, 31, 34) (4, 27, 30) | (5,26,30) |
| 22 | (2, 19, 20) (7, 31, 35) (11, 29, 35) | (3, 18, 20) (13, 27, 35) (17, 23, 35) | (13, 24, 34) | (5, 16, 20) (6, 15, 20) (7, 14, 20) |
| | (11, 29, 55) | (17,25,55) | | (8, 13, 20) |
| 26 | (2, 59, 60) | (4, 57, 60) (5, 59, 61) | (13, 57, 62) (21, 59, 63) | |



A Large Set of Solutions

The Shortest Vector

Conclusion

8- Ω Solutions

$$G = \frac{1}{9} \left(-\Omega_1 + \Omega_{\vec{\ell}_{a_2}} - \Omega_{\vec{\ell}_{a_3}} + \Omega_{\vec{\ell}_{a_4}} - \Omega_{\vec{\ell}_{a_5}} + \Omega_{\vec{\ell}_{a_6}} - \Omega_{\vec{\ell}_{a_7}} + \Omega_{\vec{\ell}_{a_8}} \right)$$

Quantization Condition: $(-\vec{\ell}_1 + \vec{\ell}_{a_2} - \vec{\ell}_{a_3} + \vec{\ell}_{a_4} - \vec{\ell}_{a_5} + \vec{\ell}_{a_6} - \vec{\ell}_{a_7} + \vec{\ell}_{a_8}) \mod 3 = 0$

Tadpole Contribution:

 $T_{min} = 8$

Massive Scalars:

14

| $(a_2, a_3, a_4, a_5, a_6, a_7, a_8)$ | | | | |
|---------------------------------------|-------------------------------|--|--|--|
| (2, 8, 6, 15, 13, 19, 20) | (2, 8, 6, 25, 23, 29, 30) | | | |
| (2, 9, 7, 14, 12, 19, 20) | (2, 9, 7, 24, 22, 29, 30) | | | |
| (3, 8, 5, 16, 13, 18, 20) | (3, 8, 5, 26, 23, 28, 30) | | | |
| (3, 10, 7, 14, 11, 18, 20) | (3, 10, 7, 24, 21, 28, 30) | | | |
| (4, 9, 5, 16, 12, 17, 20) | (4, 9, 5, 26, 22, 27, 30) | | | |
| (2, 14, 12, 25, 23, 33, 34) | (3, 14, 11, 26, 23, 32, 34) | | | |
| (4, 15, 11, 26, 22, 31, 34) | (5, 17, 12, 28, 23, 33, 35) | | | |



The Shortest Vector

Conclusion

The Shortest Vector Problem

- \exists solutions with $N_{\text{flux}} = 8$. Are their "shorter" fluxes? Answer: No.
- To rule out $N_{\text{flux}} \leq T$, we proceed as follows. (In our code we implement this for T = 7).
- Choose *T* complex *B_i*'s to be (possibly) non-zero and set the remaining to zero → a set of (126 – 2*T*) **R**-linear equations in the integer flux quanta *y_i*.
- The set $\{\textit{B}_{i_1},\ldots,\textit{B}_{i_T}\}\neq 0$ gives a reduced expression for $N_{flux},$

$$N^{red}_{flux} = Q_{ij}\,c_i\,c_j \quad \text{where} \quad c_i \in \mathbb{Z}, \, i=1,\ldots,2T$$

and Q is symmetric, positive definite.

• Is $N_{flux}^{red}(\vec{c}) \leq T$ for a choice of integers c_i ?

A Large Set of Solutions

The Shortest Vector

Conclusion

The Eigensieve Algorithm

- Define the level set in \mathbb{R}^{2T} : $L_T = \{(x_1, \dots, x_{2T}) : N_{\text{flux}}^{\text{red}}(\vec{x}) = T\} \subset \mathbb{R}^{2T}$.
- Points \vec{x} outside of the ellipsoid L_T will have values, $N_{flux}^{red} > T$.
- Are there integer lattice points on L_T or in its interior?
- Eigensieve: A combination of enumeration and sieving using the eigenvalues of Q, $\{\lambda_1, \ldots, \lambda_{2T}\}$ where $0 < \lambda_1 \leq \lambda_2 \leq \ldots \leq \lambda_{2T}$ and normalized eigenvectors, $\{\vec{v}_1, \ldots, \vec{v}_{2T}\}$. First carve out a finite region of \mathbb{Z}^{2T} using the smallest eigenvalue of Q. Lattice points \vec{p} in this region satisfying $|\vec{p}.\vec{v}_i| > \sqrt{\frac{T}{\lambda_i}}$, i = 1(1)2T, all lie outside of L_T are sieved off. Explicit evaluation of $N_{\text{flux}}^{\text{red}}$ done on the remaining points.



The 1⁹ Landau-Ginzburg model

Flux Tadpole, Massive Modul

A Large Set of Solutions

The Shortest Vector

Conclusion

Summary & Open Questions

Brady, Sengupta



A Large Set of Solutions

The Shortest Vector

Conclusion

Summary & Open Questions

• A systematic search of solutions with the lowest value of $N_{\rm flux}$, organized by number of non-zero components, has been launched. All $N_{\rm flux}$ -minimizing solutions up to four Ω_1 's turned on found.



The Shortest Vector

Conclusion

Summary & Open Questions

- A systematic search of solutions with the lowest value of $N_{\rm flux}$, organized by number of non-zero components, has been launched. All $N_{\rm flux}$ -minimizing solutions up to four Ω_{I} 's turned on found.
- The SVP for the lattice of supersymmetric flux vacua in $1^9/Z_3$ solved. $(N_{\rm flux})_{\rm min} = 8$. A large class of 8- Ω solutions found to saturate this.

The Shortest Vector

Conclusion

Summary & Open Questions

- A systematic search of solutions with the lowest value of $N_{\rm flux}$, organized by number of non-zero components, has been launched. All $N_{\rm flux}$ -minimizing solutions up to four Ω_{l} 's turned on found.
- The SVP for the lattice of supersymmetric flux vacua in $1^9/Z_3$ solved. $(N_{\rm flux})_{\rm min} = 8$. A large class of 8- Ω solutions found to saturate this.
- All solutions found so far have number of massive moduli below the tadpole conjecture bound.
- Need to push exhaustive search up to 12-Ω solutions moving away from symbolic computation in Mathematica necessary.

The Shortest Vector

Conclusion

Summary & Open Questions

- A systematic search of solutions with the lowest value of $N_{\rm flux}$, organized by number of non-zero components, has been launched. All $N_{\rm flux}$ -minimizing solutions up to four Ω_{l} 's turned on found.
- The SVP for the lattice of supersymmetric flux vacua in $1^9/Z_3$ solved. $(N_{\rm flux})_{\rm min} = 8$. A large class of 8- Ω solutions found to saturate this.
- All solutions found so far have number of massive moduli below the tadpole conjecture bound.
- Need to push exhaustive search up to 12-Ω solutions moving away from symbolic computation in Mathematica necessary.
- Finally, we aim to extend all our analyses to other Gepner models.

Thank You!

.