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#### **Motivations for axions and ALPs**

• Axion: solution to the strong CP problem

$$\mathcal{L}_{\theta} = \theta \frac{\alpha_s}{8\pi} G^a_{\mu\nu} \tilde{G}^{a,\mu\nu}$$

- Pseudo Nambu-Goldstone Bosons of a spontaneously broken symmetry
- Plausible Dark Matter candidates with several production mechanisms

$$\ddot{a} + 3H\dot{a} + m_a^2 a = 0 \qquad \dot{n}_X + 3n_X H \approx g_{B_1} \int \frac{d^3 p_{B_1}}{(2\pi)^3} \frac{f_{B_1} \Gamma_{B_1}}{\gamma_{B_1}}$$
Misalignment
Freeze- In

# **ALP with Flavor-Violating Couplings**

Consider an ALP model with flavor-violating (FV) couplings to SM fermions *f* 

$$\mathcal{L} = \frac{1}{2} (\partial_{\mu} a)^2 - \frac{m_a^2}{2} a^2 + \frac{\partial_{\mu} a}{2f_a} \bar{f}_i \gamma^{\mu} (C_{f_i, f_j}^V + C_{f_i, f_j}^A \gamma^5) f_j$$

Features:

- Free parameters: ALP mass  $m_a$ , the scale  $f_a$  and the FV couplings  $C_{q_i,q_i}^{V,A}$ .
- Anomaly-free: No couplings to gluons or EW gauge bosons \_\_\_\_\_\_\_ Crucial for DM Pheno.

- Leptons: DM scenario considered in arxiv:2209.03371, Panci et al.
- Quarks This talk.

# ALP with Flavor-Violating Couplings: why?

**Theory** Misalignment between the U(1) PQ charges X<sub>L,R</sub> and the Yukawa matrix

$$V_{\text{CKM}} = U_{u_L}^{\dagger} U_{d_L} \qquad C_u^{V,A} = U_{u_R}^{\dagger} X_{u_R} U_{u_R} \pm U_{u_L}^{\dagger} X_{Q_L} U_{u_L},$$
$$C_d^{V,A} = U_{d_R}^{\dagger} X_{d_R} U_{d_R} \pm U_{d_L}^{\dagger} X_{Q_L} U_{d_L},$$

Phenomenology

Potentially interesting experimental signatures, for example...

Colliders

$$K \to \pi a$$

Supernovae

 $\Lambda \to na$ 

arxiv: 2012.11632, Camalich et al.

Indirect probes of high-energy scales  $f_a$  (up to  $10^{12}$  GeV)

**ALP with Flavor-Violating Couplings: how?** 

What about the  $U_{L,R}$  matrices? Who knows...



#### **ALP Dark Matter - Stability**

Extremely light axion,  $m_a < \Lambda_{\text{QCD}}$   $\chi$ PT needed • Decay into pions:  $a \to \pi\pi\pi$ • Decay into photons:  $a \to \gamma\gamma$ • Decay into photons:  $a \to \gamma\gamma$  $\Gamma_{\gamma\gamma} \gtrsim 10^{30} \text{ sec}$   $\Gamma_{\gamma\gamma} = \frac{4\pi\alpha_{em}^2 m_a^3}{f_a^2} |C_{\gamma\gamma}|^2$  Excluded by X-ray searches

$$\begin{split} C_{\gamma\gamma} &\approx -\frac{1}{16\pi^2} \sum_{i} \frac{C_{ii}^A Q_i^2}{4} \frac{m_a^2}{m_i^2} - \frac{C_{ss}^A}{48\pi^2} \frac{m_a^2}{m_\eta^2 - m_a^2} , \quad (\pi^8 - \text{mixing}) \\ C_{\gamma\gamma} &\approx -\frac{1}{16\pi^2} \sum_{i} \frac{C_{ii}^A Q_i^2}{4} \frac{m_a^2}{m_i^2} - \frac{C_{uu}^A - C_{dd}^A}{32\pi^2} \frac{m_a^2}{m_\pi^2 - m_a^2} , \quad (\pi^3 - \text{mixing}). \end{split}$$

#### ALP Dark Matter – Production I

Large  $f_a$  Small couplings DM produced out of equilibrium from the early Universe thermal bath

#### DM freeze-in

$$\begin{split} \Omega h^2|_{\rm dec} &\approx 0.12 \left(\frac{m_a}{0.1\,{\rm MeV}}\right) \left(\frac{6\times 10^9 {\rm GeV}}{f_a/C_{q_iq_j}}\right)^2 \left(\frac{m_{q_i}}{{\rm GeV}}\right) \left(\frac{79}{g^*(m_q)}\right)^{\frac{3}{2}} \qquad q_i \to q_j a \\ \Omega h^2|_{\rm scatt} &\approx 0.12 \left(\frac{m_a}{0.1\,{\rm MeV}}\right) \left(\frac{2\times 10^9 {\rm GeV}}{f_a/C_{qq}}\right)^2 \left(\frac{m_q}{{\rm GeV}}\right) \left(\frac{79}{g^*(m_q)}\right)^{\frac{3}{2}} q_i \overline{q_i} \to ga, \qquad q_i \ g \to q_i a \end{split}$$

### **ALP Dark Matter – Production II**

Freeze-In under control? The non-ren. operator

$$\mathcal{L}_{eff} = -C_{q_i q_j}^A \frac{ia}{f_a} \frac{m_{q_i}}{v} h \bar{q}_i P_R q_j \qquad q_i h \to q_j a$$

$$\sum_{i=1}^{N} \sum_{\substack{i=1\\ i \neq i \neq i}} \Omega h^2|_{q_i \to q_j a} \qquad \text{Introduces a dependence on } T_R \text{ (unknown)}$$

$$\sum_{\substack{i=1\\ i \neq i \neq i \neq i}} \Omega h^2|_{q_i \to q_j a} \qquad \text{Introduces a dependence on } T_R \text{ (unknown)}$$

Dominant IR contribution for

$$T_R < \frac{3\pi^3 v^2}{m_{q_i}}$$

Scenarios involving the top quark have

$$T_R pprox O(10) \, {
m TeV}$$
 UV Dominated

# Astro Bounds – Lyman- $\alpha$

DM produced with large free-streaming length

$$m_a \gtrsim 10 \,\mathrm{keV} \left(\frac{m_{\mathrm{WDM}}}{3.5 \,\mathrm{keV}}\right)^{\frac{4}{3}} \left(\frac{79}{g^*(m_q)}\right)^{\frac{1}{3}}$$

arxiv:2012.01446, D'Eramo, Lenoci

For lower DM masses the large free streaming length suppresses the matter power spectrum

Conflicts with structure formation: constraints from Ly-  $\alpha\,$  data

# Astro Bounds - Supernova

Axion emission through nucleon scattering N' = n, p $N + N' \rightarrow N + N' + a$ 



Plus emittivity constraint from SN 1987A:

 $\begin{array}{ll} 0.61g_{ap}^2 + g_{an}^2 + 0.53g_{an}g_{ap} < 8.26 \times 10^{-19} & \mbox{arxiv:1906.11844,} \\ g_{ai} \equiv C_i m_i/f_a. & \mbox{Carenza et al.} \end{array}$ 

#### Astro Bounds – X-ray Searches + CMB



#### **Collider Bounds**



from arxiv:2303.13353, Ziegler





# Conclusions

- Axions and ALPs are interesting candidates from a phenomenological point of view.
- An ALP that is also a DM candidate meets specific requirements , especially from Cosmology.
- Flavor Violating couplings to fermions: more parameters but also more pheno.
- For ALPs coupled to quarks, QCD dynamics improves DM stability: weaker constraints.
- Less testable than ALP coupled to leptons. Can future collider and X-ray searches cover more parameter space?

# **Backup Slides**

# A chiral Lagrangian for ALP and Mesons

$$\begin{aligned} \mathcal{L}_{\text{light}} &= \frac{1}{2} (\partial_{\mu} a)^2 - \frac{m_a^2}{2} a^2 + \bar{\Psi} (i \not D - M_q) \Psi \\ &+ \frac{\partial_{\mu} a}{2f_a} \bar{\Psi} \gamma^{\mu} \widetilde{k}_L \left( 1 - \gamma^5 \right) \Psi + \frac{\partial_{\mu} a}{2f_a} \bar{\Psi} \gamma^{\mu} \widetilde{k}_R \left( 1 + \gamma^5 \right) \Psi \\ &\Psi \equiv (u, d, s)^T. \end{aligned}$$
 Effective symmetry 
$$\begin{aligned} \text{SU(3)}_L &\otimes \text{SU(3)}_R \\ &\Psi \equiv (u, d, s)^T. \end{aligned}$$

$$\Sigma = \exp\left(i\sqrt{2\lambda_a} \cdot \pi^a/f_{\pi}\right)$$
with
$$D_{\mu}\Sigma = \partial_{\mu}\Sigma + ieA_{\mu}\left[Q,\Sigma\right] + i\frac{\partial_{\mu}a}{f_a}(\tilde{k}_L\Sigma - \Sigma\tilde{k}_R)$$

$$\mathcal{L}_{\chi\rm PT} = \frac{1}{2}\partial_{\mu}a\partial^{\mu}a - \frac{m_a^2}{2}a^2 + \frac{f_{\pi}^2}{8}\mathrm{Tr}\left[D_{\mu}\Sigma D^{\mu}\Sigma^{\dagger}\right] + \frac{f_{\pi}^2}{4}B_0\mathrm{Tr}\left[M_q\Sigma^{\dagger} + h.c.\right]$$

### **ALP Dark Matter – Production III**

Non thermal mechanisms are also allowed

Misalignment  $m_a \approx H_1$ 

$$\Omega h^2|_{\rm mis} \approx 0.015 \left(\frac{H_R}{0.64 \text{keV}}\right)^{\frac{1}{2}} \left(\frac{f_a \theta_0}{4 \times 10^{10} \,\text{GeV}}\right)^2$$
 Negligible

In our mass range, onset of oscilations prior to reheating: diluted mis. contribution.

# Additional Bounds? Too weak!

• Collider Bound for the Top Scenario:

Combined with NA62 is only able to probe  $f_a \approx 10^7$  GeV

• Supernova Bound:  $\Lambda \rightarrow n \ a$  is only able to probe  $f_a \approx 10^3 \ {\rm GeV}$ 

 $F_{sd}^V \gtrsim 7.1 \times 10^9 \text{ GeV}$ 

 $\tilde{H}, H$ 

arxiv: 2012.11632,

Camalich et al.