

In collaboration with: M. Aghaie, G. Armando, P. Panci and R. Ziegler

Motivations for axions and ALPs

• Axion: solution to the **strong CP problem**

$$
\mathcal{L}_\theta = \theta \tfrac{\alpha_s}{8\pi} G^a_{\mu\nu} \tilde{G}^{a, \mu\nu}
$$

- Pseudo Nambu-Goldstone Bosons of a spontaneously broken symmetry
- Plausible **Dark Matter candidates** with several production mechanisms

$$
\ddot{a} + 3H\dot{a} + m_a^2 a = 0 \qquad \qquad \dot{n}_X + 3n_X H \approx g_{B_1} \int \frac{d^3 p_{B_1}}{(2\pi)^3} \frac{f_{B_1} \Gamma_{B_1}}{\gamma_{B_1}}
$$
\nMisalignment

\nFreeze-In

ALP with Flavor-Violating Couplings

Consider an ALP model with flavor-violating (FV) couplings to SM fermions f

$$
\mathcal{L} = \frac{1}{2} (\partial_{\mu} a)^2 - \frac{m_a^2}{2} a^2 + \frac{\partial_{\mu} a}{2f_a} \bar{f}_i \gamma^{\mu} (C_{f_i, f_j}^V + C_{f_i, f_j}^A \gamma^5) f_j
$$

Features:

- Free parameters: ALP mass m_a , the scale f_a and the FV couplings $C_{q_i,q_j}^{V,A}$ V, A
g_i.g_i.
- Anomaly-free: No couplings to gluons or EW gauge bosons \longrightarrow Crucial for DM Pheno.

- Leptons: DM scenario considered in arxiv:2209.03371, Panci et al.
- Quarks \longrightarrow This talk.

ALP with Flavor-Violating Couplings: why?

Theory Misalignment between the U(1) PQ charges $X_{L,R}$ and the Yukawa matrix

$$
V_{\text{CKM}} = U_{u_L}^{\dagger} U_{d_L} \quad \begin{array}{c} C_u^{V, A} = U_{u_R}^{\dagger} X_{u_R} U_{u_R} \pm U_{u_L}^{\dagger} X_{Q_L} U_{u_L}, \\ C_d^{V, A} = U_{d_R}^{\dagger} X_{d_R} U_{d_R} \pm U_{d_L}^{\dagger} X_{Q_L} U_{d_L}, \end{array}
$$

Phenomenology Potentially interesting experimental signatures, for example…

$$
K\to \pi a
$$

Colliders Supernovae

 $\Lambda \rightarrow na$

arxiv: 2012.11632, Camalich et al.

Indirect probes of high-energy scales f_a (up to 10^{12} GeV)

ALP with Flavor-Violating Couplings: how?

What about the $U_{L,R}$ matrices? Who knows...

A Benchmark case: two-flavor scenario $S_{d_R} = \text{diag}(0,1,-1) \, \boxed{\text{Switchle rotation in a plane}} \hspace{1cm} C_{V,A}^d = \begin{pmatrix} 0 & 0 & 0 \ 0 & \sin\alpha & \cos\alpha \ 0 & \cos\alpha & -\sin\alpha \end{pmatrix} \, .$ In this way, 3 parameters only m_a α f_a $b-s$

ALP Dark Matter - Stability

Extremely light axion, $m_a < \Lambda_{\rm OCD}$ \sim \sim \times PT needed Decay into pions: $a \rightarrow \pi \pi \pi$ Decay into photons: $a \rightarrow \gamma \gamma$ $m_a \gtrsim 3m_{\pi}$ $\tau_{\gamma\gamma}\gtrsim 10^{30}~{\rm sec}~~|~~~\Gamma_{\gamma\gamma}=\frac{4\pi\alpha_{em}^2m_a^3}{f_a^2}|C_{\gamma\gamma}|^2$ Excluded by X-ray searches $C_{\gamma\gamma} \approx -\frac{1}{10^{-2}} \sum \frac{C_{ii}^A Q_i^2}{4} \frac{m_a^2}{m_a^2} - \frac{C_{ss}^A}{48 \pi^2} \frac{m_a^2}{m_a^2}$, $(\pi^8 - \text{mixing})$ $(i = c, b, t)$

$$
C_{\gamma\gamma} \approx -\frac{1}{16\pi^2} \sum_{i} \frac{C_{ii}^A Q_i^2 m_a^2}{4 m_i^2} - \frac{C_{uu}^A - C_{dd}^A}{32\pi^2} \frac{m_a^2}{m_\pi^2 - m_a^2} , \quad (\pi^3 - \text{mixing}).
$$

ALP Dark Matter – Production I

Large $f_a \succeq$ Small couplings \searrow DM produced out of equilibrium from the early Universe thermal bath

DM freeze-in

$$
\Omega h^2|_{\text{dec}} \approx 0.12 \left(\frac{m_a}{0.1 \,\text{MeV}}\right) \left(\frac{6 \times 10^9 \text{GeV}}{f_a/C_{q_i q_j}}\right)^2 \left(\frac{m_{q_i}}{\text{GeV}}\right) \left(\frac{79}{g^*(m_q)}\right)^{\frac{3}{2}} \qquad q_i \to q_j a
$$
\n
$$
\Omega h^2|_{\text{scatt}} \approx 0.12 \left(\frac{m_a}{0.1 \,\text{MeV}}\right) \left(\frac{2 \times 10^9 \text{GeV}}{f_a/C_{qq}}\right)^2 \left(\frac{m_q}{\text{GeV}}\right) \left(\frac{79}{g^*(m_q)}\right)^{\frac{3}{2}} \quad q_i \overline{q_i} \to g a, \qquad q_i \text{ g} \to q_i a
$$

ALP Dark Matter – Production II

Freeze-In under control? The non-ren. operator

$$
\mathcal{L}_{eff} = -C_{q_i q_j}^A \frac{ia}{f_a} \frac{m_{q_i}}{v} h \bar{q}_i P_R q_j \qquad q_i h \to q_j a
$$

$$
\sum_{n=1}^{\infty} \sum_{j=1}^{\infty} \sum_{q_i \to q_j a} \mathcal{L}^{(1)}(q_i \to q_j a) \qquad \text{Introduces a dependence on } T_R \text{ (unknown)}
$$

Dominant IR contribution for

$$
T_R < \frac{3\pi^3 v^2}{m_{q_i}}
$$

Scenarios involving the top quark have

$$
T_R \approx O(10)\,\text{TeV} \quad \overbrace{\text{---}}\rangle \quad \text{UV Dominated}
$$

Astro Bounds – Lyman-α

DM produced with large free-streaming length

$$
m_a \gtrsim 10\,{\rm keV}\left(\frac{m_{\rm WDM}}{3.5\,{\rm keV}}\right)^{\frac{4}{3}} \left(\frac{79}{g^*(m_q)}\right)^{\frac{1}{3}} \qquad \qquad \ \ {\rm a} \quad \ \ \, \frac{1}{\rm b}
$$

arxiv:2012.01446, D'Eramo, Lenoci

For lower DM masses the large free streaming length suppresses the matter power spectrum

Conflicts with structure formation: constraints from Ly- α data

Astro Bounds - Supernova

Axion emission through nucleon scattering $N' = n, p$ $N + N' \rightarrow N + N' + a$

Plus emittivity constraint from SN 1987A:

 $0.61g_{ap}^2 + g_{an}^2 + 0.53g_{an}g_{ap} < 8.26 \times 10^{-19}$ arxiv:1906.11844, Carenza et al. $g_{ai} \equiv C_i m_i/f_a.$

Astro Bounds – X-ray Searches + CMB

Collider Bounds

from arxiv:2303.13353, Ziegler

Conclusions

- Axions and ALPs are interesting candidates from a phenomenological point of view.
- An ALP that is also a DM candidate meets specific requirements , especially from Cosmology.
- Flavor Violating couplings to fermions: more parameters but also more pheno.
- For ALPs coupled to quarks, QCD dynamics improves DM stability: weaker constraints.
- Less testable than ALP coupled to leptons. Can future collider and X-ray searches cover more parameter space?

Backup Slides

A chiral Lagrangian for ALP and Mesons

$$
\mathcal{L}_{\text{light}} = \frac{1}{2} (\partial_{\mu} a)^{2} - \frac{m_{a}^{2}}{2} a^{2} + \bar{\Psi} (i \not\!\!D - M_{q}) \Psi
$$
\n
$$
+ \frac{\partial_{\mu} a}{2 f_{a}} \bar{\Psi} \gamma^{\mu} \tilde{k}_{L} (1 - \gamma^{5}) \Psi + \frac{\partial_{\mu} a}{2 f_{a}} \bar{\Psi} \gamma^{\mu} \tilde{k}_{R} (1 + \gamma^{5}) \Psi
$$
\n
$$
\Psi \equiv (u, d, s)^{T}.
$$
\n
$$
\mathbf{W} = (u, d, s)^{T}.
$$
\nEXECUTE: The system of the system is given by

\n
$$
\mathbf{S} \mathbf{U} (\mathbf{3}) \mathbf{L} \otimes \mathbf{S} \mathbf{U} (\mathbf{3}) \mathbf{R}
$$

$$
\Sigma = \exp\left(i\sqrt{2\lambda_a} \cdot \pi^a / f_\pi\right)
$$

with

$$
D_\mu \Sigma = \partial_\mu \Sigma + ieA_\mu [Q, \Sigma] + i\frac{\partial_\mu a}{f_a}(\widetilde{k}_L \Sigma - \Sigma \widetilde{k}_R)
$$

$$
\mathcal{L}_{\chi PT} = \frac{1}{2}\partial_\mu a \partial^\mu a - \frac{m_a^2}{2}a^2 + \frac{f_\pi^2}{8}\text{Tr}\left[D_\mu \Sigma D^\mu \Sigma^\dagger\right] + \frac{f_\pi^2}{4}B_0 \text{Tr}\left[M_q \Sigma^\dagger + h.c.\right]
$$

ALP Dark Matter – Production III

Non thermal mechanisms are also allowed

Misalignment $m_a \approx H$

$$
\Omega h^2|_{\text{mis}} \approx 0.015 \left(\frac{H_R}{0.64 \text{keV}}\right)^{\frac{1}{2}} \left(\frac{f_a \theta_0}{4 \times 10^{10} \text{ GeV}}\right)^2 \qquad \text{Negligible}
$$

In our mass range, onset of oscilations prior to reheating: diluted mis. contribution.

Additional Bounds? Too weak!

• Collider Bound for the Top Scenario:

$$
\Delta c_{sd}^{V}(\mu) = \frac{y_t^2}{64\pi^2} \log \frac{f_a}{\mu} \left[2V_{ts}^* V_{td}(c_{tt}^{V} + c_{tt}^A) - \sum_k V_{ts}^* V_{td}(c_{u_{kt}}^{V} - c_{u_{kt}}^A) \right]
$$
\n
$$
- \sum_k V_{ts}^* V_{kd}(c_{u_{kt}}^{V} - c_{u_{kt}}^A) \left[\text{Camalich et al.} \right]
$$
\n
$$
Q_{Lj}
$$
\n
$$
Q_{Lj}
$$
\n
$$
Q_{Lk}
$$
\n
$$
U_{R,k}, D_{R,k}
$$
\n
$$
U_{R,k'}, D_{R,k'}
$$
\n
$$
U_{R,k'}, D_{R,k'}
$$

Combined with NA62 is only able to probe $f_a \approx 10^7$ GeV

• Supernova Bound: $\Lambda \to n \, a$ is only able to probe $f_a \approx 10^3$ GeV

 $F_{sd}^V \gtrsim 7.1 \times 10^9$ GeV

 \tilde{H},H

arxiv: 2012.11632,

Camalich et al.