Black Holes as Transducers for Ultralight Bosons

Yifan Chen Niels Bohr International Academy yifan.chen@nbi.ku.dk

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[Introduction to Ultralight Bosons and Superradiance](#page-2-0)

[Black Holes as Neutrino Factories](#page-5-0)

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$$
-\frac{1}{2}\nabla^{\mu}a\nabla_{\mu}a-\frac{1}{4}B^{\mu\nu}B_{\mu\nu}+\mathcal{L}_{\text{EH}}(H)-V(\Psi), \quad \Psi=a,\phi,B^{\mu}\text{ and }H^{\mu\nu}.
$$

 \triangleright Axion: hypothetical pseudoscalar motivated by strong CP problem.

 \blacktriangleright Prediction from fundamental theories with extra dimensions:

e.g.
$$
g^{MN}(5D) \rightarrow g^{\mu\nu}(4D) + B^{\mu}(4D)
$$
, $B^M(5D) \rightarrow B^{\mu}(4D) + a(4D)$.

String axiverse/photiverse: logarithmic mass window, $m_\Psi \propto e^{-\mathcal{V}_{6D}}$.

► Coherent wave dark matter candidates when $m_\Psi < 1$ eV:

$$
\Psi(x^{\mu}) \simeq \Psi_0(\mathbf{x}) \cos \omega t; \qquad \Psi_0 \simeq \frac{\sqrt{\rho}}{m_{\Psi}}; \qquad \omega \simeq m_{\Psi}.
$$

Superradiant Gravitational Atoms

Gravitational Atom between BH and axion cloud:

BL coordinate : $\Psi^{\text{GA}}(x^{\mu}) = e^{-i\omega t} e^{im\phi} S_{lm}(\theta) R_{lm}(r), \qquad \omega \simeq m_{\Psi} + i\Gamma.$

$$
\bigotimes_{\omega \simeq m_{\Psi} + i\Gamma}
$$

IF Superradiance [Penrose, Zeldovichi, Starobinsky, Damour et al, Brito et al review]: boson cloud exponentially extracting BH rotation energy when

Compton wavelength
$$
\lambda_c \simeq
$$
 gravitational radius r_g .
\n $m_\Psi \sim 10^{-21} \text{ eV} \leftrightarrow M_{\text{BH}} \sim 10^9 M_\odot$.

 $\blacktriangleright \psi_{\text{max}}^{\text{GA}} \equiv \Psi_0$ approaches M_{pl} when $M_{\text{cloud}} \leq 10\%$ M_{BH} :

$$
\frac{\textit{M}_{\text{cloud}}}{\textit{M}_{\text{BH}}} \approx \left\{ \begin{array}{c} 0.5\% \; \left(\frac{\Psi_{0}}{10^{16} \, \text{GeV}}\right)^{2} \; \left(\frac{0.4}{\alpha}\right)^{4} \; \text{for scalar,} \\ 0.8\% \; \left(\frac{\Psi_{0}}{10^{17} \, \text{GeV}}\right)^{2} \; \left(\frac{0.4}{\alpha}\right)^{4} \; \text{for vector,} \end{array} \right.
$$

 $\alpha \equiv G_N M_{\rm BH} m_{\rm W}$ gravitational fine-structure constant

 \blacktriangleright Black holes are powerful transducers for ultralight bosons.

Superradiant Saturating Cloud

 \triangleright Self interaction or matter interaction triggers cloud energy leakage, balancing SR, invalidating spin constraints.

- Two examples for axion:
	- Ionized axion waves for $\Psi_0 \sim f_a < 10^{16}$ GeV [Yoshino et al 12', Baryakht et al 20'].

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- Parametric γ production for $g_{a\gamma}\Psi_0 \sim 1$ [Rosa et al 17', Spieksma et al 23'].
- **St[ro](#page-5-0)[n](#page-3-0)[g fi](#page-4-0)[e](#page-5-0)[ld](#page-1-0) frontier**: similar to preheating an[d s](#page-3-0)trong field [Q](#page-4-0)[E](#page-5-0)[D](#page-1-0)[.](#page-2-0)

Black Holes as

Neutrino Factories

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based on arXiv:2308.00741

YC, Xiao Xue, Vitor Cardoso.

Neutrino Production and Acceleration from Scalar Cloud

Neutrino coupled to majoron: $\omega_{\nu}^2 = k^2 + m_{\text{eff}}^2$, $m_{\text{eff}} = m_{\nu} + g_{\phi \nu} \phi_0 \cos \mu t$.

► Fermi sphere $k_* = \sqrt{g_{\phi\nu}\phi_0\mu}/2$ is pumped as $m_{\rm eff} \sim 0$ [Greene Kofman 98' 00']. Production rate: $\Gamma_{\phi\nu} \approx (g_{\phi\nu} \phi_0)^{3/2} \mu^{5/2}/(48\pi^3)$.

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Further neutrino acceleration under majoron cloud background:

$$
\tfrac{d\rho_\nu^\alpha}{dt}=-\tfrac{1}{\rho_\nu^0}\Gamma_{\kappa\beta}^\alpha\rho_\nu^\kappa\, \rho_\nu^\beta -\tfrac{1}{2\rho_\nu^0}\nabla^\alpha m_{\rm eff}^2. \leftarrow \text{scalar force [Uzan et al 20']}
$$

Spin Measurement and Neutrino/Boosted DM Flux

- **IDED** Neutrino emission from saturation phase $\Gamma_{\phi\nu} = \Gamma_{\text{SR}}$. Point-like sources surpass atmospheric neutrino at ∼ TeV.
- High spin excludes region $\Gamma_{\phi\nu} \ll \Gamma_{\text{SR}}$.
- Multi-messenger observation:
	-
	- GW and EM searches for BHs. Neutrin[o a](#page-6-0)[nd](#page-8-0) [b](#page-6-0)[oo](#page-7-0)[st](#page-8-0)[e](#page-4-0)[d](#page-5-0) [d](#page-7-0)[a](#page-8-0)[rk](#page-4-0) [m](#page-7-0)a[tte](#page-0-0)[r.](#page-41-0)

Probing Ultralight Bosons

with Event Horizon Telescope

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EHT and ngEHT for new physics

Event Horizon Telescope: best-ever spatial resolution from VLBI.

geodesics: photons propagating multiple times around BH enhance intensity on the image plane.

- \rightarrow Precise test of general relativity.
	- \blacktriangleright Astrometry for new physics?

Stokes Q, U EVPA χ \equiv $arg(Q + i U)/2$ [EHT 21']

Linear polarization from synchrotron radiation reveals magnetic field structure.

Four days' observations show slight difference.

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 \blacktriangleright New interactions?

Photon Ring Astrometry

for Superradiant Clouds

based on arXiv:2211.03794, Phys. Rev. Lett. 130 (2023) no.11, 111401

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YC, Xiao Xue, Richard Brito, Vitor Cardoso.

Gravitational Atom-induced Geodesics Deflections

In Superradiant clouds generate local oscillatory metric perturbations $g_{\mu\nu}\simeq g_{\mu\nu}^{\rm K}+\epsilon h_{\mu\nu}$ that deflect geodesics $x^\mu\simeq x^\mu_{(0)}+\epsilon x^\mu_{(1)}$:

- Scalar cloud mainly causes time delay [Khmelnitsky, Rubakov 13'].
- Polarized vector or tensor cloud contribute to both time delay and spatial deflection.

Astrometrical Photon Ring Autocorrelations

A photon pair executing different half orbits number N:

 \triangleright Intensity fluctuation correlation: $\langle \Delta I(t, \varphi) \Delta I(t+T, \varphi+\Phi) \rangle$, peaks at $T \approx N \tau_0$ and $\Phi \approx N \delta_0$ [Hadar, Johnson, Lupsasca, Wong 20'].

Observables: $\Delta \Phi^N = \Phi_0^N \cos (\omega t + \delta)$ for $N = 1$ and 2.

► Probe M_{cloud}/M_{BH} to 10^{-3} for vector and 10^{-7} for tensor.

Hunting Axions with Event Horizon Telescope

Polarimetric Measurements

based on arxiv: 1905.02213, Phys. Rev. Lett. 124 (2020) no.6, 061102, arxiv: 2105.04572, Nature Astron. 6 (2022) no.5, 592-598, arxiv: 2208.05724, JCAP 09 (2022), 073.

YC, Chunlong Li, Yuxin Liu, Ru-Sen Lu, Yosuke Mizuno, Jing Shu, Xiao Xue, Qiang Yuan, Yue Zhao, Zihan Zhou.

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Axion Cloud Induced Birefringence (Xiao Xue's talk)

Axion-induced Birefringence: rotation of *linear polarization*:

 $\mathbf{g}_{a\gamma}a\mathbf{F}_{\mu\nu}\tilde{\mathbf{F}}^{\mu\nu}/2 \rightarrow \Delta\chi = g_{a\gamma}[a(t_{\rm obs}, \mathbf{x}_{\rm obs}) - a(t_{\rm emit}, \mathbf{x}_{\rm emit})].$

 \triangleright Extended sources, plasma and curved space-time effects?

Covariant radiative transfer [IPOLE simulation]

with an accretion flow model outside SMBH: [Strominger 19']

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Summary

- \blacktriangleright Rotating black holes are powerful transducers for ultralight bosons due to superradiance.
- Strong field frontier:
	- Parametric particle production and acceleration.
- Multi-messenger correlation:

neutrino/dark matter detection \leftrightarrow GW/EM observation.

- Event Horizon Telescope:
	- Photon geodesics deflection.
	-

Thank you!

Appendix

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Superradiance for Boson with Negligible Interaction

For bosons with negligible interaction, superradiance stops after BH spins down and M_{cloud} takes up to $10\%M_{BH}$.

- **High spin** excludes boson mass in SR range with reasonable τ_{BH} . [Arvanitaki, Brito, Davoudiasl, Denton, Stott, Unal, Saha et al]
- GW from boson annihilation and transition slowly decreases M_{cloud} .

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[Yoshino, Brito, Isi, Siemonsen, Sun, Palomba, Zhu, Tsukada, Yuan, LVK et al]

Weakly Saturating Axion Cloud

Strong self-interaction region $a^{GA} \simeq f_a$ happens when $f_a < 10^{16}$ GeV:

$$
V(a) = m_a^2 f_a^2 \left(1 - \cos \frac{a}{f_a}\right) = \frac{m_a^2 a^2}{2} - \frac{m_a^2 a^4}{24f_a^2} + \dots;
$$

 \triangleright A quasi-equilibruim phase where superradiance and non-linear interaction induced emission balance each other with $a_{\rm max}^{\rm GA}\simeq {\cal O}(1)\,f_a.$

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[Yoshino, Kodama 12' 15', Baryakht et al 20']

Neutrino Acceleration from Boson Cloud

In Neutrino propagation under majoron cloud background:

 $\frac{{\rm d} p_{\nu}^{\alpha}}{{\rm d} t} = -\frac{1}{p_{\nu}^{0}}\Gamma_{\kappa\beta}^{\alpha}\rho_{\nu}^{\kappa}\,\rho_{\nu}^{\beta} \!-\! \frac{1}{2p_{\nu}^{0}}\nabla^{\alpha}m_{\rm eff}^{2}.$ \leftarrow scalar force

- \blacktriangleright Two parts of scalar force:
	- $-\vec{\nabla} m_{\text{eff}}^2 \propto \alpha^2 \hat{r} \frac{2 r_{\text{g}}}{r \cos(\alpha t \phi) \sin \theta} \hat{n}_{\perp} + \cdots$
	- Outer region: pure radial acceleration.
	- Inner region: polar trapping.

Find momentum:
$$
\bar{\omega}_{\text{acc}}^{\nu} \sim g_{\phi\nu} \Psi_0
$$
.

 \blacktriangleright Both spatial and temporal variation are necessary for acceleration.

Gravitational Atom-induced Geodesics Deflections

Backward ray-tracing:

Two phases of evolution:

- Perturbative generation of oscillatory deviations;
- Photon ring instability leads to exponential growth of the oscillatory deviations between two sequential crossing the equa[tor](#page-20-0)i[al](#page-22-0) [pl](#page-20-0)[an](#page-21-0)[e.](#page-22-0)

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Photon Ring Autocorrelations as Astrometry

► Photon ring autocorrelation exclusion criteria: $\Delta \Phi^N > \ell_{\phi} \approx 4.3^{\circ}$ or ngEHT's smearing kernel for φ : 10°.

- \triangleright A tensor with linear coupling to stress tensors is more sensitive than a vector with quadratic couplings.
- $N = 2$ correlation peak can probe large unexplored parameter space of cloud mass.

 $\mathbf{A} \equiv \mathbf{A} + \mathbf{B} + \math$

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 \triangleright Sources with shorter correlation time, e.g., hotspots or pulsars can significantly increase the sensitivity.

Stringent Constraints on Axion-Photon Coupling

Next-generation EHT is expected to significan[tly](#page-22-0)i[nc](#page-24-0)[re](#page-22-0)[as](#page-23-0)[e](#page-24-0) [se](#page-7-0)[n](#page-8-0)[sit](#page-41-0)[iv](#page-7-0)[it](#page-8-0)[y.](#page-41-0)

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Event Horizon Telescope: an Earth-sized Telescope

- For single telescope with diameter D , the angular resolution for photon of wavelength λ is around $\frac{\lambda}{D}$;
- \triangleright VLBI: for multiple radio telescopes, the effective D becomes the maximum separation between the telescopes.

As good as being able to see \leq on the moon from the Earth.

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Supermassive Black Hole (SMBH) M87* [EHT 19' 21']

Event Horizon Telescope: best-ever spatial resolution from VLBI.

Total intensity I

Linear polarization Q, U EVPA $\chi \equiv$ $arg(Q + i U)/2$

- First-time: shadow and the ring;
- Ring size determines $6.5 \times 10^9 M_{\odot}$;
- Polarization map reveals magnetic field structure.
- Four days' observations show slight difference.

From other observations:

- Nearly extreme Kerr black hole: $a_1 > 0.8$;
- Almost face-on disk with a 17° inclination angle;
- Rich information under strong gravity, what el[se](#page-24-0) [ca](#page-26-0)[n](#page-24-0) [w](#page-25-0)[e](#page-26-0) [le](#page-7-0)[a](#page-8-0)[rn?](#page-41-0)

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Axion Cloud and Birefringence

 $\blacktriangleright \Delta\langle \chi(\varphi) \rangle$: propagating wave along φ on the sky plane BL[c](#page-27-0)[o](#page-25-0)ordin[a](#page-7-0)te: $\;$ a $^{\rm GA}$ $^{\rm GA}$ $^{\rm GA}$ \propto cos $[m_at-\phi]\rightarrow \Delta\langle \chi(\varphi)\rangle \propto {\cal A}(\varphi)$ cos $[m_at+\varphi+\delta(\varphi)].$ $[m_at+\varphi+\delta(\varphi)].$ $[m_at+\varphi+\delta(\varphi)].$ $[m_at+\varphi+\delta(\varphi)].$ $[m_at+\varphi+\delta(\varphi)].$

Axion Birefringence for RIAF around M87* (IPOLE simulation)

 $\Delta\langle \chi(\varphi) \rangle = \mathcal{A}(\varphi) \cos[m_a t + \varphi + \delta(\varphi)].$

Scan axion mass: $\alpha \equiv r_{\rm g} m_{\rm a} \in [0.10, 0.44]$ with **period [5, 20] days**.

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- \triangleright $\delta(\varphi) \approx -5 \alpha \sin 17^\circ \cos \varphi$: phase delay at different φ .
- Asymmetry of $A(\varphi) = O(1)g_{\alpha\gamma}f_{\alpha}$: washout from lensed photon with $\delta_{12} = \omega \delta t - \delta \phi$!

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Lensed Photon Washout

 \triangleright The ratio between linear polarization from lensed photon and direct emissions vary from RIAF models, giving different washout effects.

 \triangleright Universal birefringence signals for direct emission only:

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Prospect for next-generation EHT

 \triangleright Next-generation EHT is expected to significantly increase sensitivity.

Recent updates:

- Constraints from EVPAs on the whole image.
- Closure traces for EVPA variations with specifi[c p](#page-29-0)a[tt](#page-31-0)[er](#page-29-0)[ns](#page-30-0) [\[](#page-31-0)[Br](#page-7-0)[o](#page-8-0)[de](#page-41-0)[ri](#page-7-0)[ck](#page-8-0) [e](#page-41-0)[t a](#page-0-0)[l\].](#page-41-0) 4 ロ) 4 \overline{r}) 4 \overline{z}) 4 \overline{z})

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Prospect for next-generation EHT

 \triangleright Correlation between $\Delta\chi$ at different radius and frequency.

At 86 GHz, lensed photon is suppressed due to higher optical thickness.

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Longer and sequential observations.

- Better resolution of EVPA.
- Better understanding of accretion flow and jet. Intrinsic variations of EVPA from GRMHD simulation?

Birefringence from Soliton Core Dark Matter

 \triangleright Ultralight axion dark matter forms soliton core in the galaxy center. Quantum pressure balences gravitational interactions $a \sim 10^{10}$ GeV.

 $2Q$

- Linearly polarized photon from pulsar. [Liu et al 19' Caputo et al 19']
- Polarized radiation from Sgr A^* . [Yuan, Xia, YC, Yuan et al 20']
- \triangleright Coherent signals at each pixel increase the sensitivity.

Axion QED: Achromatic Birefringence [Carroll, Field, Jackiw 90']

$$
\mathcal{L}=-\frac{1}{4}F_{\mu\nu}F^{\mu\nu}-\frac{1}{2}g_{a\gamma}aF_{\mu\nu}\tilde{F}^{\mu\nu}+\frac{1}{2}\partial^{\mu}a\partial_{\mu}a-V(a),
$$

 \triangleright Chiral dispersions for photons propragating under axion background:

$$
[\partial_t^2 - \nabla^2] A_{L,R} = \mp 2 g_{a\gamma} n^{\mu} \partial_{\mu} a k A_{L,R}, \qquad \omega_{L,R} \sim k \mp g_{a\gamma} n^{\mu} \partial_{\mu} a.
$$

$$
n^{\mu}: \text{ unit directional vector}
$$

Rotation of electric vector position angle of linear polarization:

$$
\Delta \chi = g_{a\gamma} \int_{\text{emit}}^{\text{obs}} n^{\mu} \partial_{\mu} a \, dl
$$

= $g_{a\gamma} [a(t_{\text{obs}}, \mathbf{x}_{\text{obs}}) - a(t_{\text{emit}}, \mathbf{x}_{\text{emit}})].$

Topological effect for each photon: only $a(x_{\text{emit}}^{\mu})$ and $a(x_{\text{obs}}^{\mu})$ dependent.

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Accretion Flow around M87?

- ▶ EHT polarimetric measurements prefer Magnetically Arrested Disk with vertical \vec{B} around M87*.
- Analytic model: sub-Kep radiatively inefficient accretion flow:

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Dimensionless thickness parameter $H = 0.05$ and 0.3 as benchmark.

EHT Polarization Data Characterization

 \triangleright Four days' polarization map with slight difference on sequential days:

Uncertainty of the azimuthal bin EVPA from polsolve:

 $\mathbf{E} = \mathbf{A} \oplus \mathbf{A} + \mathbf{A$

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ranging from $\pm 3^{\circ}$ to $\pm 15^{\circ}$ for the bins used.

Landscape of SMBH and Accretion Flow (IPOLE simulation)

Horizon scale SMBH landscape with nnngEHT (space, L2):

Universal birefringence signals for direct emission only:

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Axion cloud can't keep growing exponentially. What's the fate of it?

- **If Self interaction** of axion becomes important for $f_a < 10^{16}$ GeV. [Yoshino, Kodama 12', Baryakht et al 20']
- \triangleright Black hole spins down until the superradiance condition is violated for $f_a > 10^{16}$ GeV. [Arvanitakia, Dubovsky 10']
- \triangleright Formation of a **binary system** leads to the decay/transition of the bound state. [Chia et al 18']
- \triangleright Electromagnetic blast for strong (large field value) axion-photon coupling. [Boskovic et al 18']

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Superradiant evolution of

the shadow and photon ring of Sgr A^*

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based on arxiv: 2205.06238, Phys. Rev. D 106 (2022) no.4, 043021.

YC, Rittick Roy, Sunny Vagnozzi, and Luca Visinelli.

Superradiant Evolution for Bosons ۱t

▶ Superradiant evolution for scalar, vector or tensor \rightarrow spin decreases:

10-2 $\sim \mathcal{O}(10)$ yrs for vector or tensor outside SgrA * . $l = 0$ and $j = m = 1$ or 2 from intrinsic spin. ▶ Superradiant timescale \propto M_{BH}, and is shorter for vector or tensor due to

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Large Inclination Angle: Shadow Drift

- ► Center of the shadow contour drifts $\sim \mathcal{O}(1) r_g$ once the spin decreases. The drift is more manifest at large inclination angles.
- \triangleright Resolution to the shadow center benefits from long observation time $\sim \mathcal{O}(1)$ yr.

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Low Inclination Angle: Azimuthal Lapse

At low inclination angles,

photon ring autocorrelation for intensity fluctuations: $\mathcal{C}(\mathcal{T}, \varphi) \equiv \iint dr dr' r r' \langle \Delta I(t, r, \phi) \Delta I(t+T, r', \phi + \varphi) \rangle$ peaks at $T = \tau_0$ and $\varphi = \delta_0$, where δ_0 is the azimuthal lapse.

 δ_0 is sensitive to spin evolution due to frame dragging.

[Chael Palumbo]

 $\mathbf{A} \equiv \mathbf{A} + \mathbf{B} + \math$ 2990