



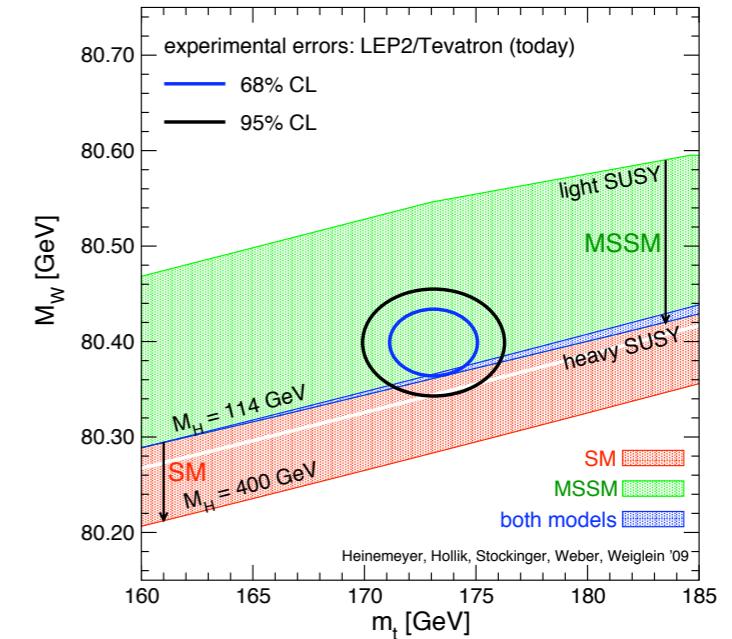
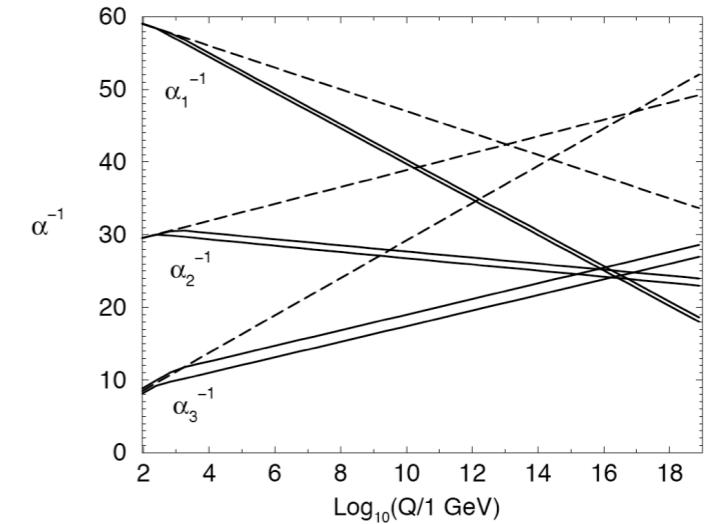
Status and Prospects for SUSY @ 7 TeV

Sabine Kraml
LPSC Grenoble

CMS France
CEA, Saclay, 1 April 2010

Why SUSY?

- Symmetry between fermions and bosons, providing a unified description of matter, gauge and Higgs fields.
- Unique extension of relativistic space-time symmetries.
- Solves the gauge hierarchy problem.
- Predicts a light Higgs.
- Allows for gauge coupling unification → GUT?
- Excellent fit to precision electroweak data.
- Provides cold dark matter candidate
- Arguably the best motivated extension of the SM



MSSM

Minimal supersymmetric standard model

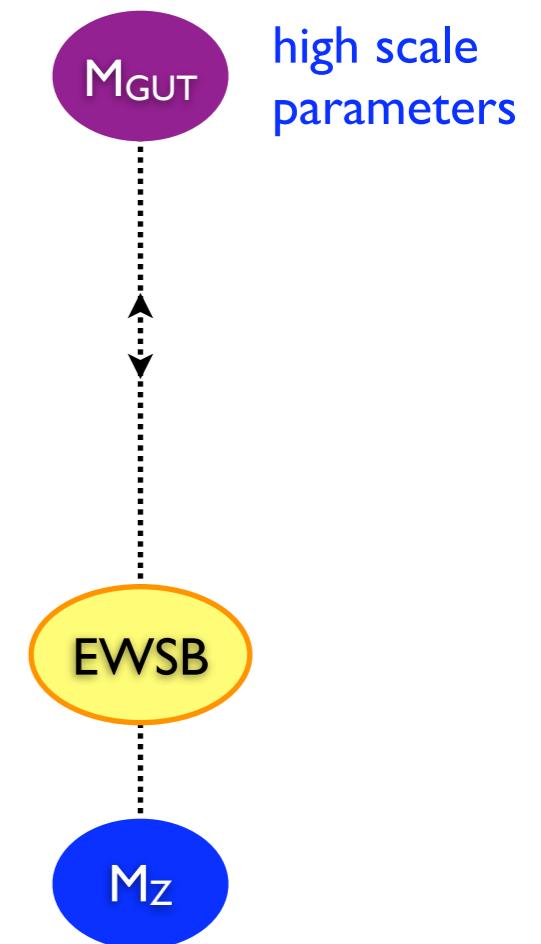
particle	spin	superpartner	spin
quarks	1/2	squarks	0
leptons	1/2	sleptons	0
gauge bosons	1	gauginos	1/2
Higgs bosons	0	higgsinos	1/2
graviton	2	gravitino	3/2

full MSSM: ~100 free parameters (soft breaking terms)

CMSSM (mSUGRA): $m_{1/2}$, m_0 , A_0 , $\tan\beta$, sign μ .

NUHM: $m_{1/2}$, m_0 , $m_{H_1,2}$, A_0 , $\tan\beta$, sign μ .

NUHM-1: $m_{H_1} = m_{H_2}$, NUHM-2: $m_{H_1} \neq m_{H_2}$



gauge & Yukawa
couplings

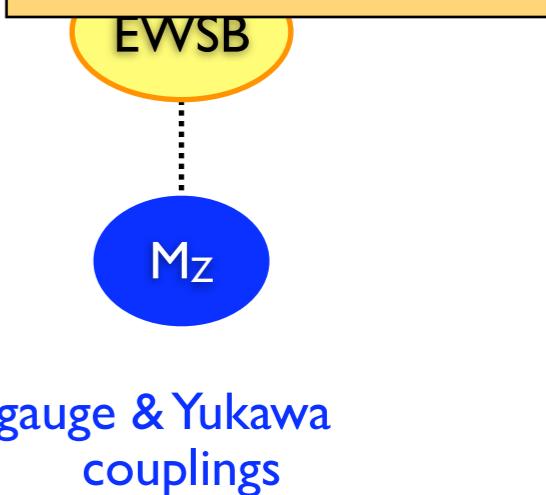
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Disclaimer:
There's been a lot of work recently regarding low-scale SSB*, SSB from extra dimensions, string-inspired models, non-universal / non-minimal models, higgsless SUSY, etc, etc. These will not be discussed here.

* SSB=soft SUSY breaking



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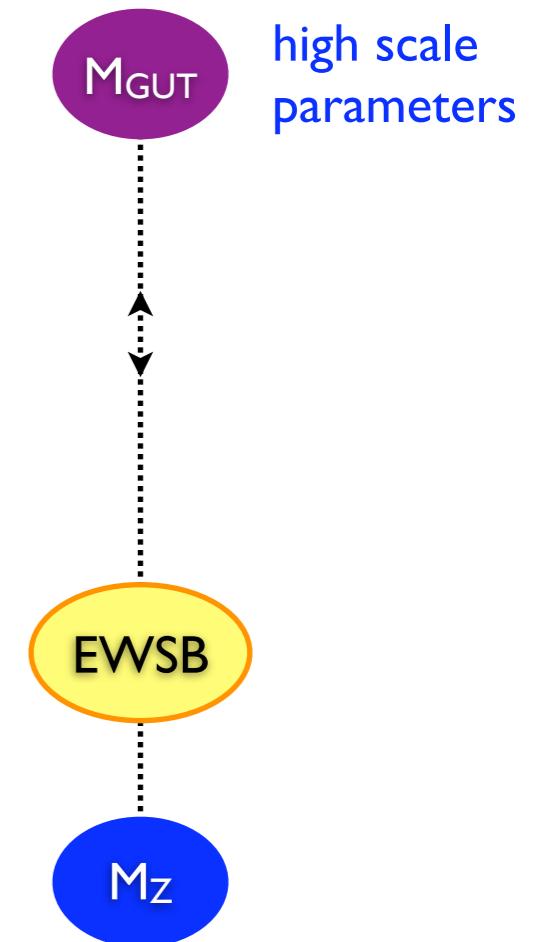
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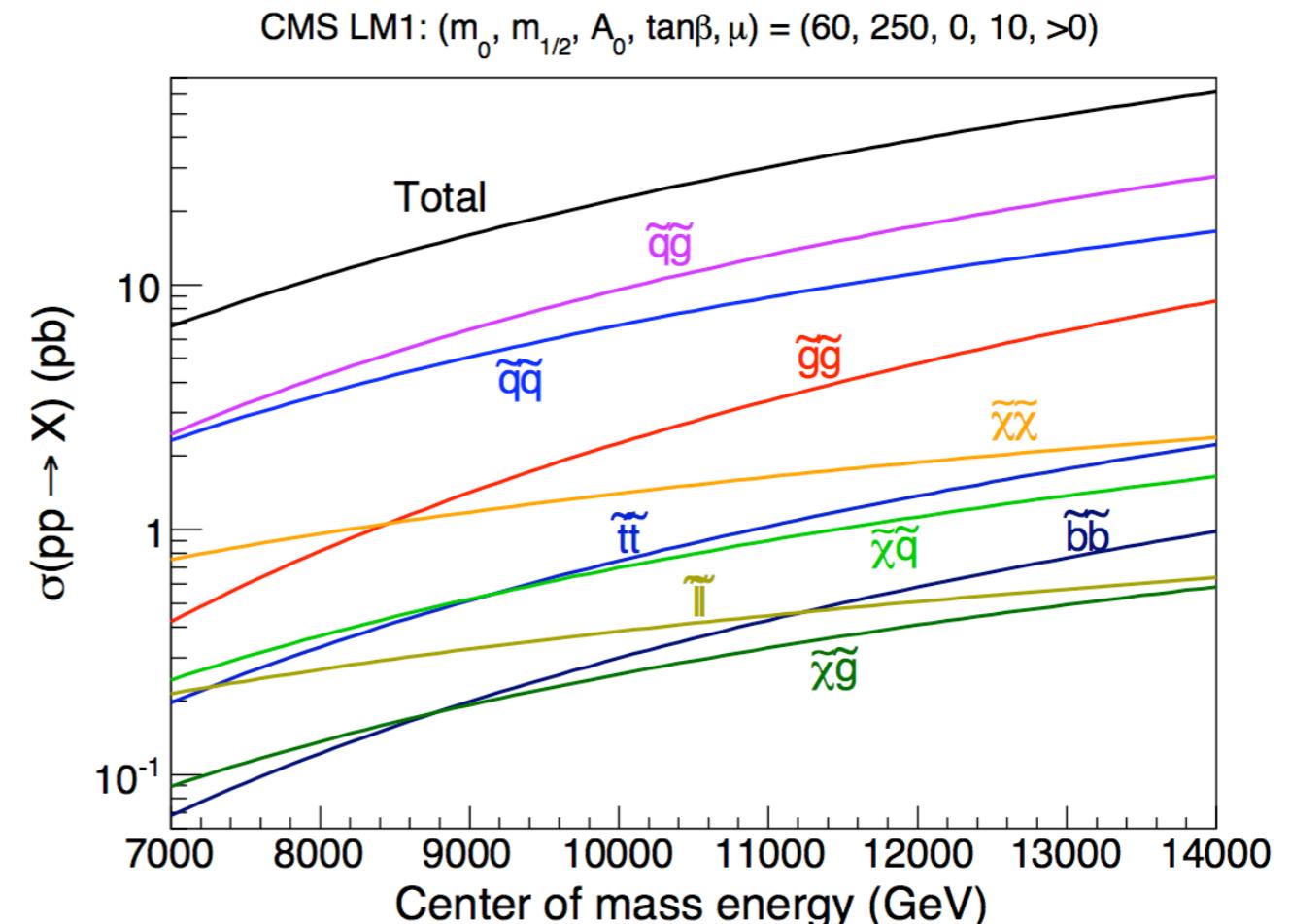
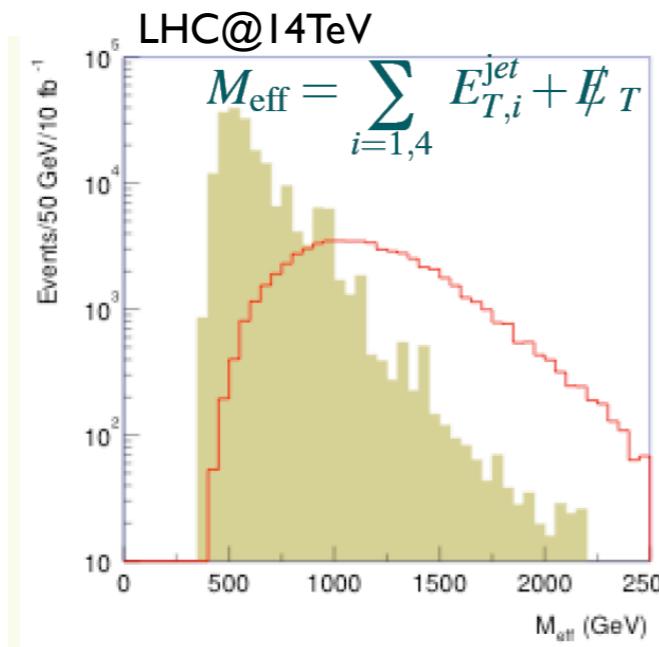
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gauge & Yukawa
couplings

SUSY at LHC

- General expectation at LHC: large Xsections for squarks and/or gluinos.
Strong interaction + the power of phase-space.
- Once produced, squarks/gluinos will decay into lighter sparticles until the LSP* is reached
 - cascade decays
 - high-p_T jets
 - large missing E_T



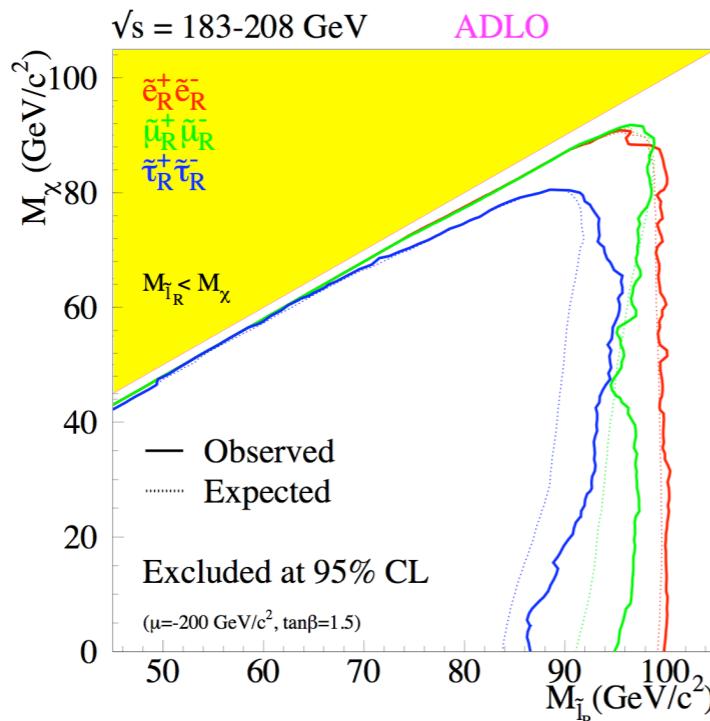
courtesy S. Sekmen

*) LSP = lightest SUSY particle, stable if R-parity is conserved

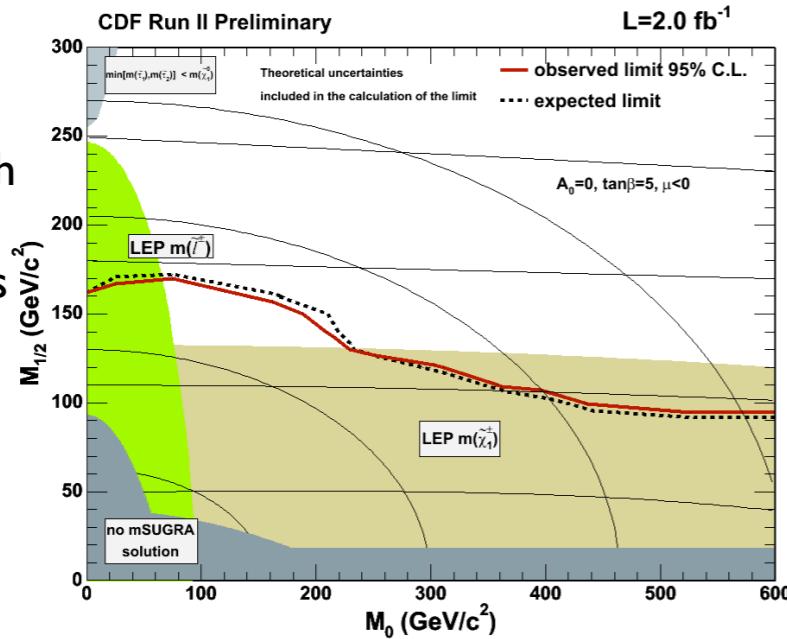
Limits and constraints

So far only lower mass limits and indirect constraints, e.g.,

LEP: charged
sparticle
 $M \gtrsim 100$ GeV

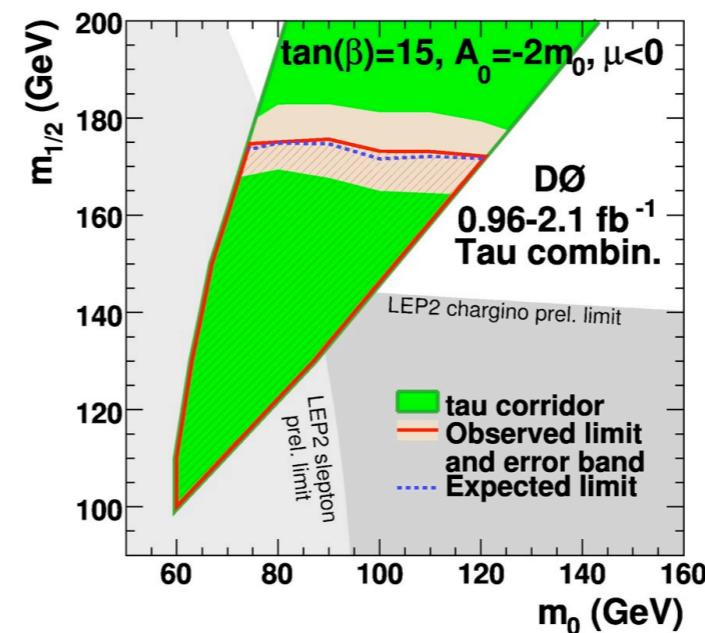


Tevatron:
begins to reach
beyond LEP
but mass limits
quite model
dependent



$\text{BR}_{b \rightarrow s\gamma}^{\text{exp}}/\text{BR}_{b \rightarrow s\gamma}^{\text{SM}}$	$1.117 \pm 0.076_{\text{exp}} \pm 0.082_{\text{th(SM)}}$
$\text{BR}(B_s \rightarrow \mu^+ \mu^-)$	$< 4.7 \times 10^{-8}$
$\text{BR}_{B \rightarrow \tau\nu}^{\text{exp}}/\text{BR}_{B \rightarrow \tau\nu}^{\text{SM}}$	$1.25 \pm 0.40_{[\text{exp+th}]}$
$\text{BR}(B_d \rightarrow \mu^+ \mu^-)$	$< 2.3 \times 10^{-8}$
$\text{BR}_{B \rightarrow X_s \ell\ell}^{\text{exp}}/\text{BR}_{B \rightarrow X_s \ell\ell}^{\text{SM}}$	0.99 ± 0.32
$\text{BR}_{K \rightarrow \mu\nu}^{\text{exp}}/\text{BR}_{K \rightarrow \mu\nu}^{\text{SM}}$	$1.008 \pm 0.014_{[\text{exp+th}]}$
$\text{BR}_{K \rightarrow \pi\nu\bar{\nu}}^{\text{exp}}/\text{BR}_{K \rightarrow \pi\nu\bar{\nu}}^{\text{SM}}$	< 4.5
$\Delta M_{B_s}^{\text{exp}}/\Delta M_{B_s}^{\text{SM}}$	$0.97 \pm 0.01_{\text{exp}} \pm 0.27_{\text{th(SM)}}$
$(\Delta M_{B_s}^{\text{exp}}/\Delta M_{B_s}^{\text{SM}})/(\Delta M_{B_d}^{\text{exp}}/\Delta M_{B_d}^{\text{SM}})$	$1.00 \pm 0.01_{\text{exp}} \pm 0.13_{\text{th(SM)}}$
$\Delta \epsilon_K^{\text{exp}}/\Delta \epsilon_K^{\text{SM}}$	$1.08 \pm 0.14_{[\text{exp+th}]}$
$a_\mu^{\text{exp}} - a_\mu^{\text{SM}}$	$(30.2 \pm 8.8) \times 10^{-10}$
$M_h [\text{GeV}]$	> 114.4 (see text)
$\Omega_{\text{CDM}} h^2$	0.1099 ± 0.0062

B-physics!



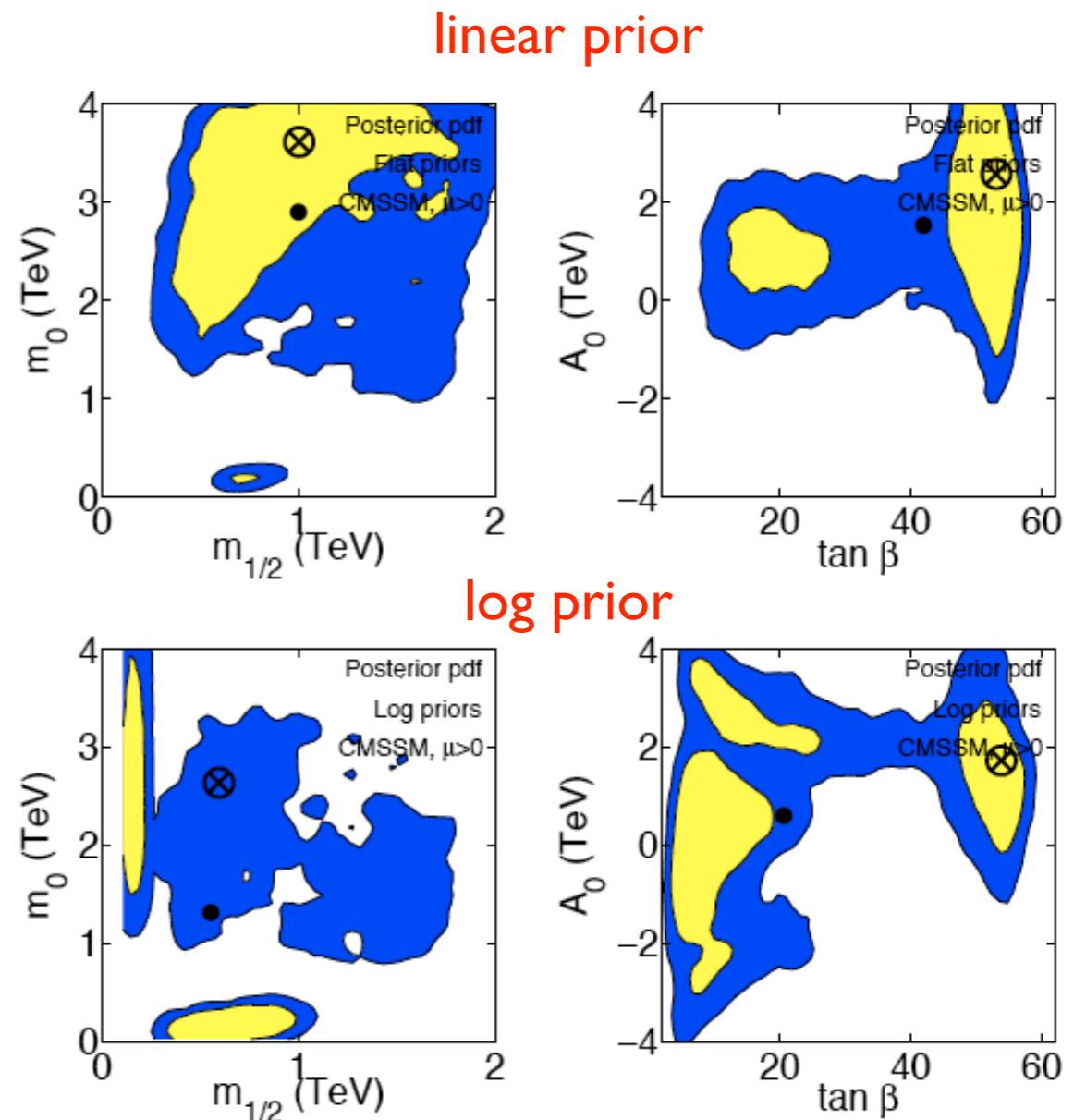
→ severe constraints
on parameter space

Fits: Bayesian

- Markov Chain Monte Carlo (MCMC) or MultiNest sampling.
 - Marginalized posterior PDFs (probability density functions) for parameter inference.

$$p(m|d) = \frac{p(d|\xi)\pi(m)}{p(d)}$$

- Alternative: profile likelihoods, but more difficult to evaluate.
 - Prior dependence!!



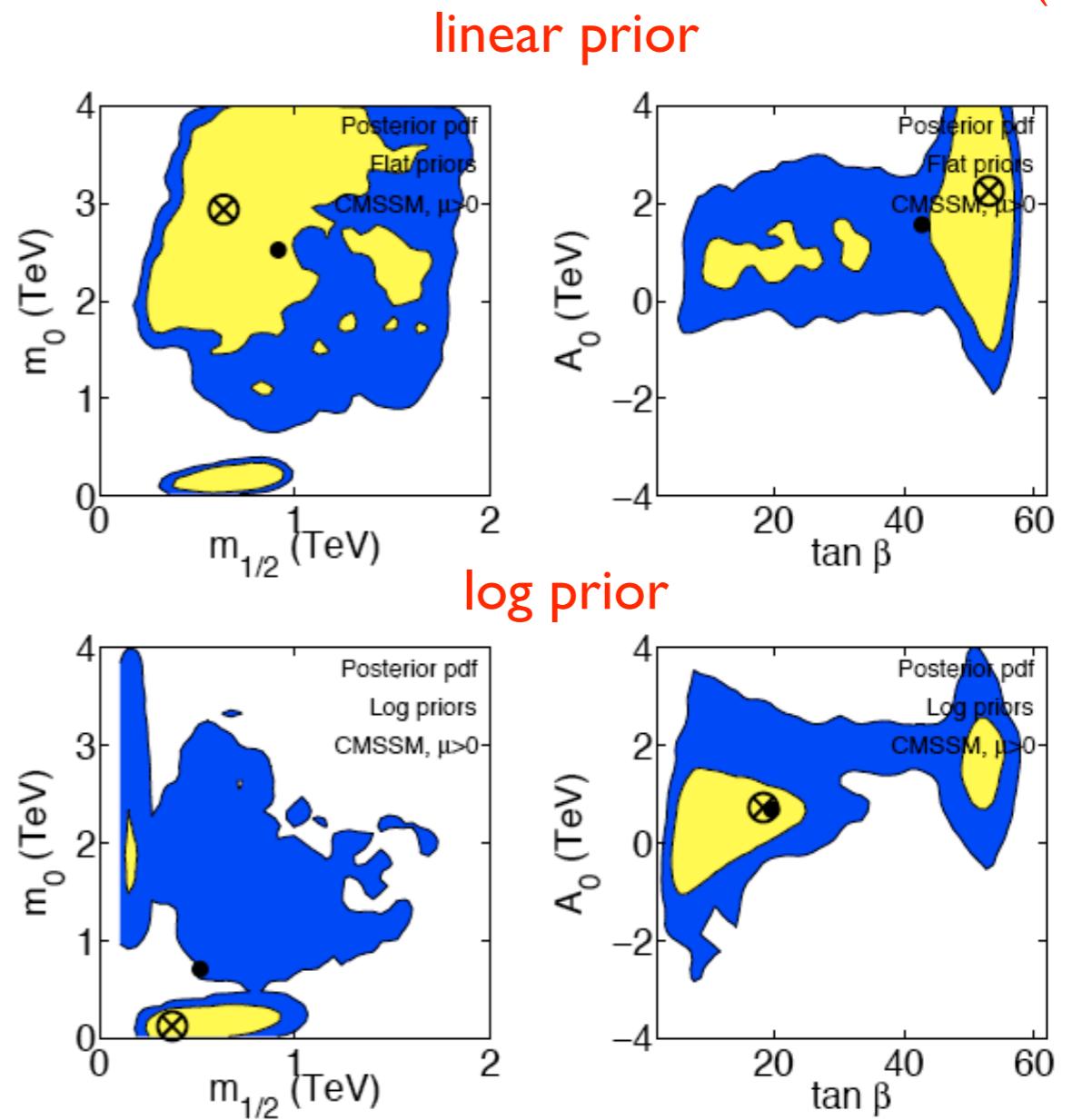
Trotta et al, arXiv:0809.3792

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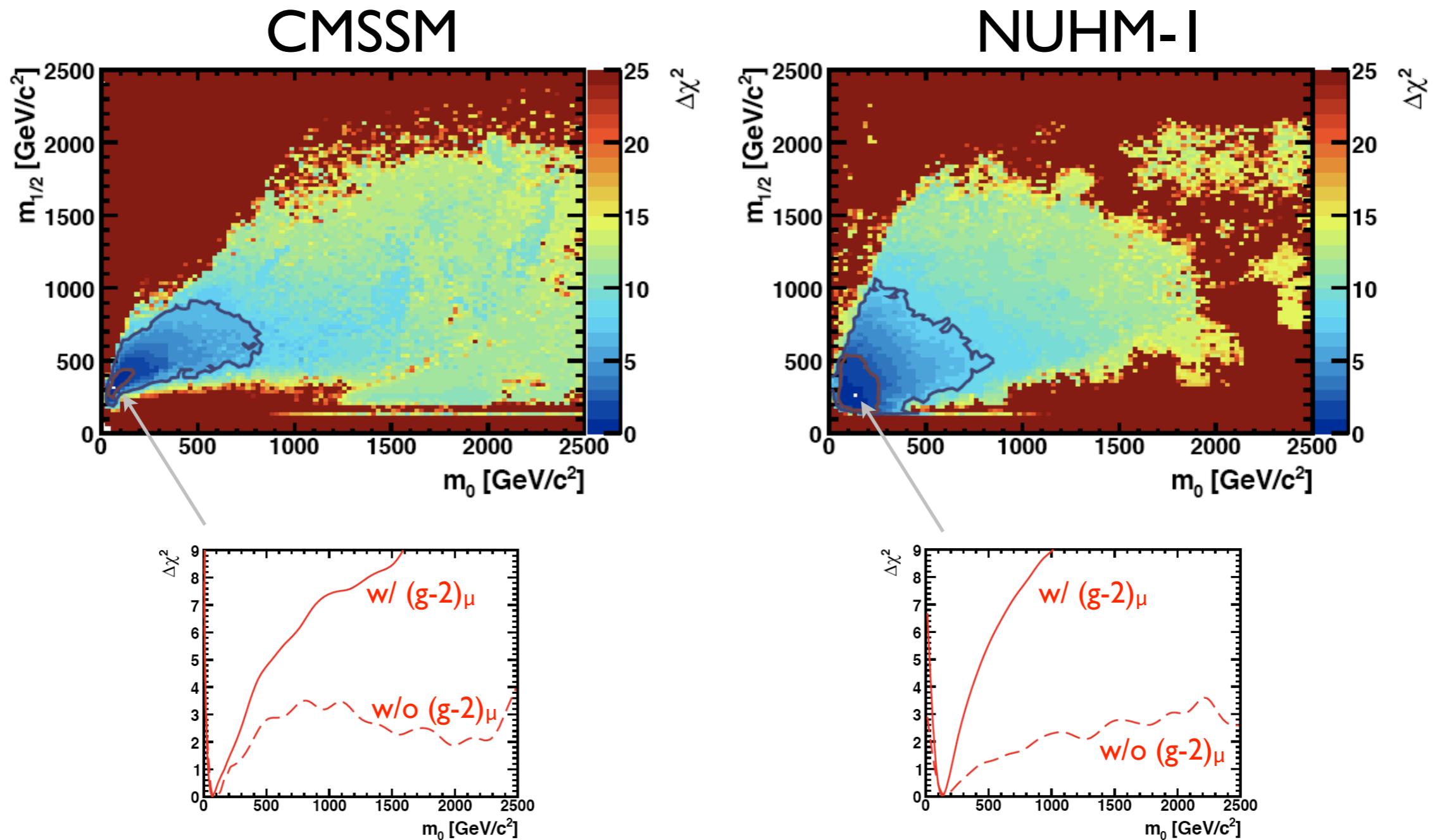
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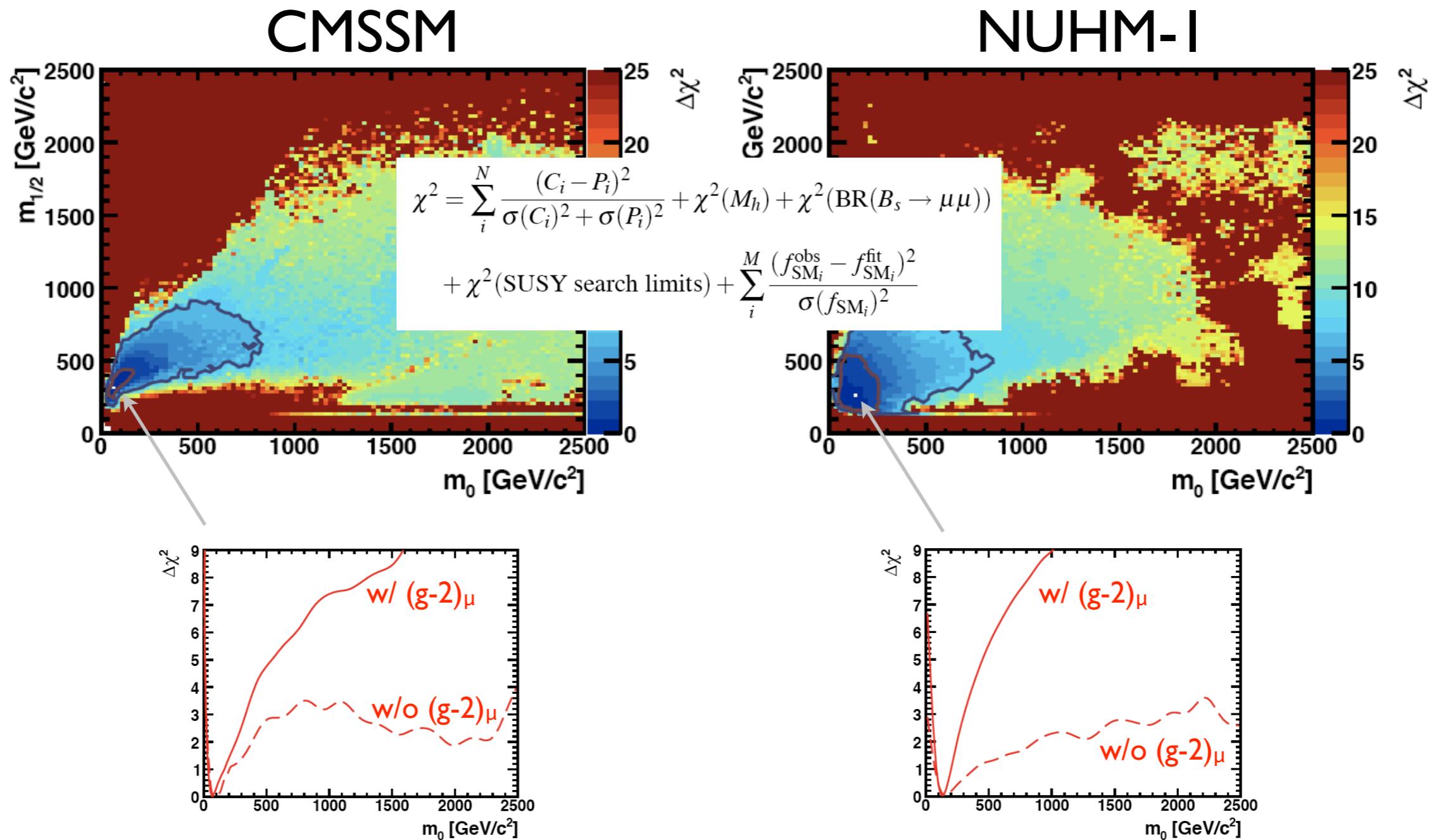
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Fits: Frequentist



Best fit at $m_{\tilde{g}} \approx 750/600$ GeV, $\tan \beta \approx 11$.

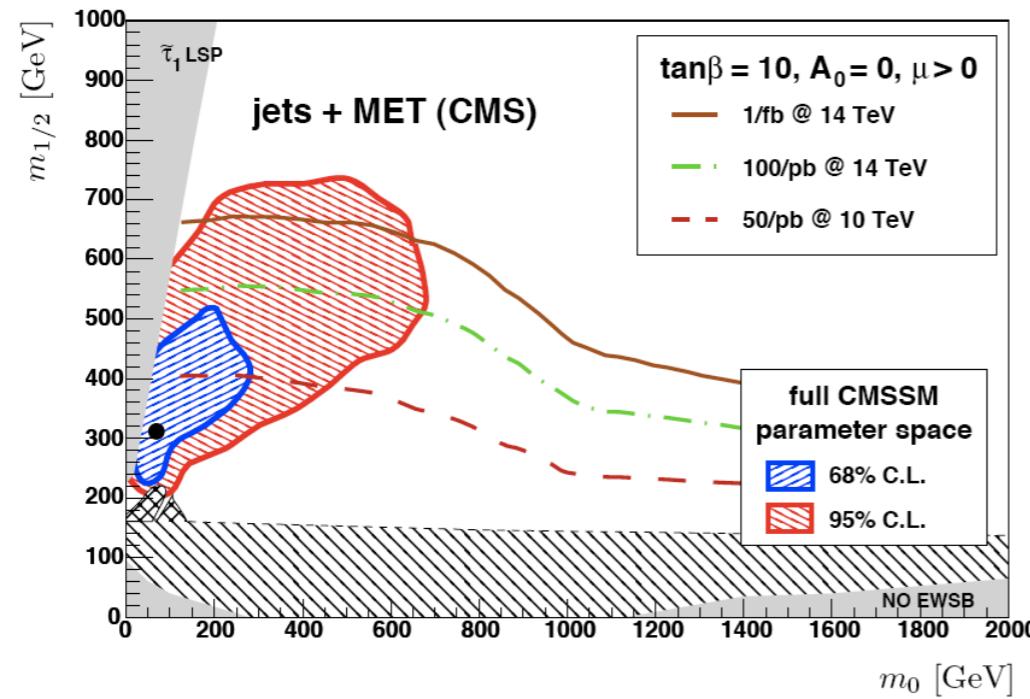
Fits: Frequentist



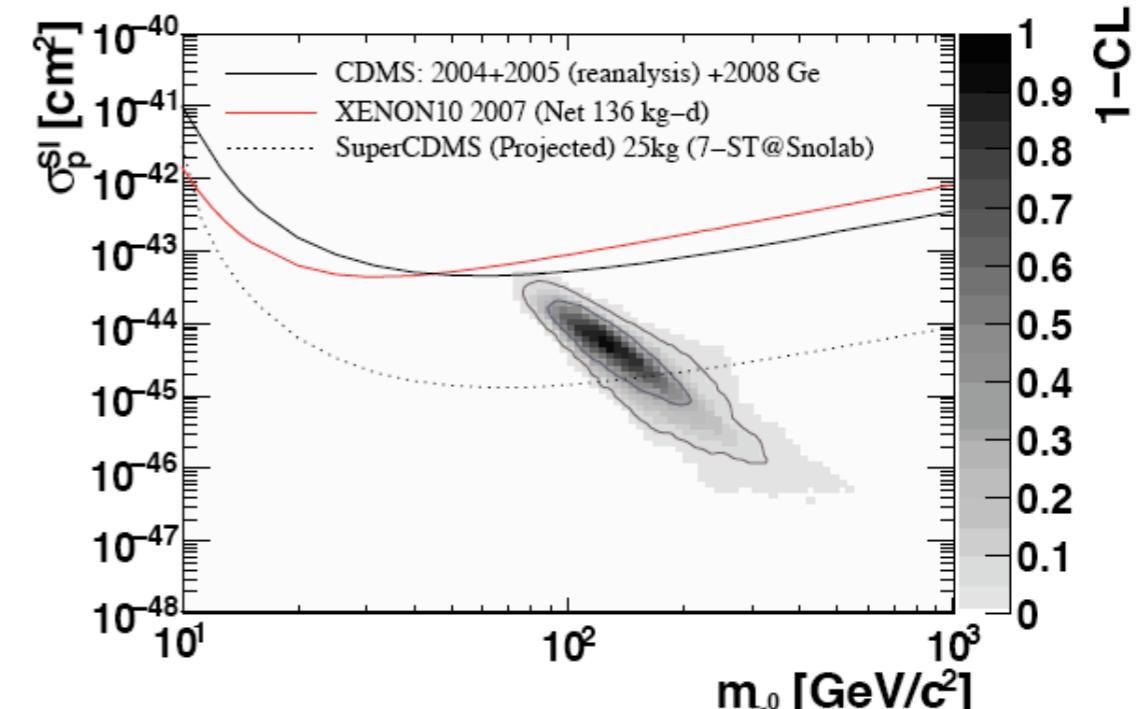
Fits: Frequentist

CMSSM

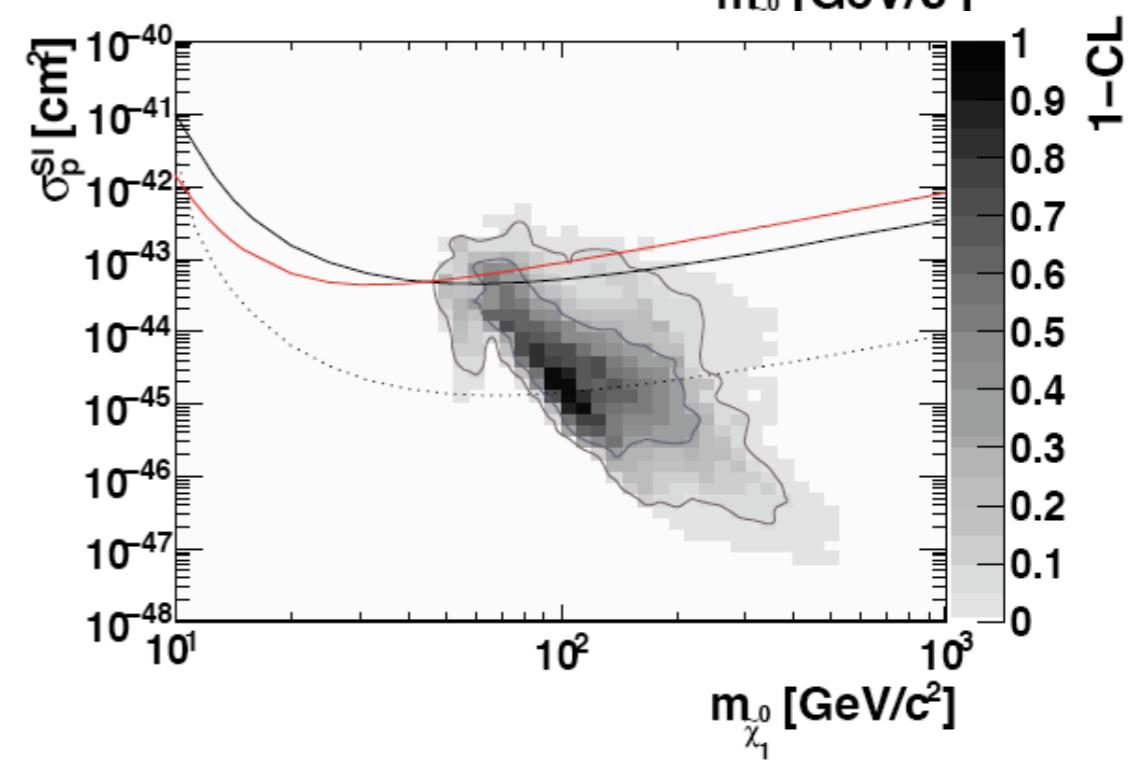
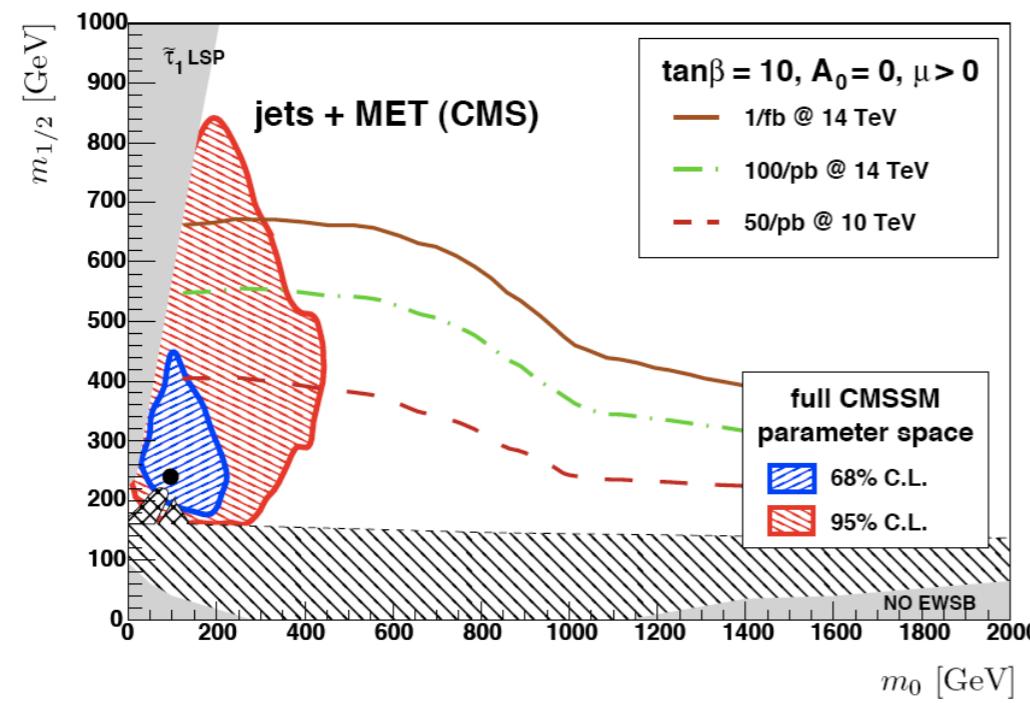
LHC potential



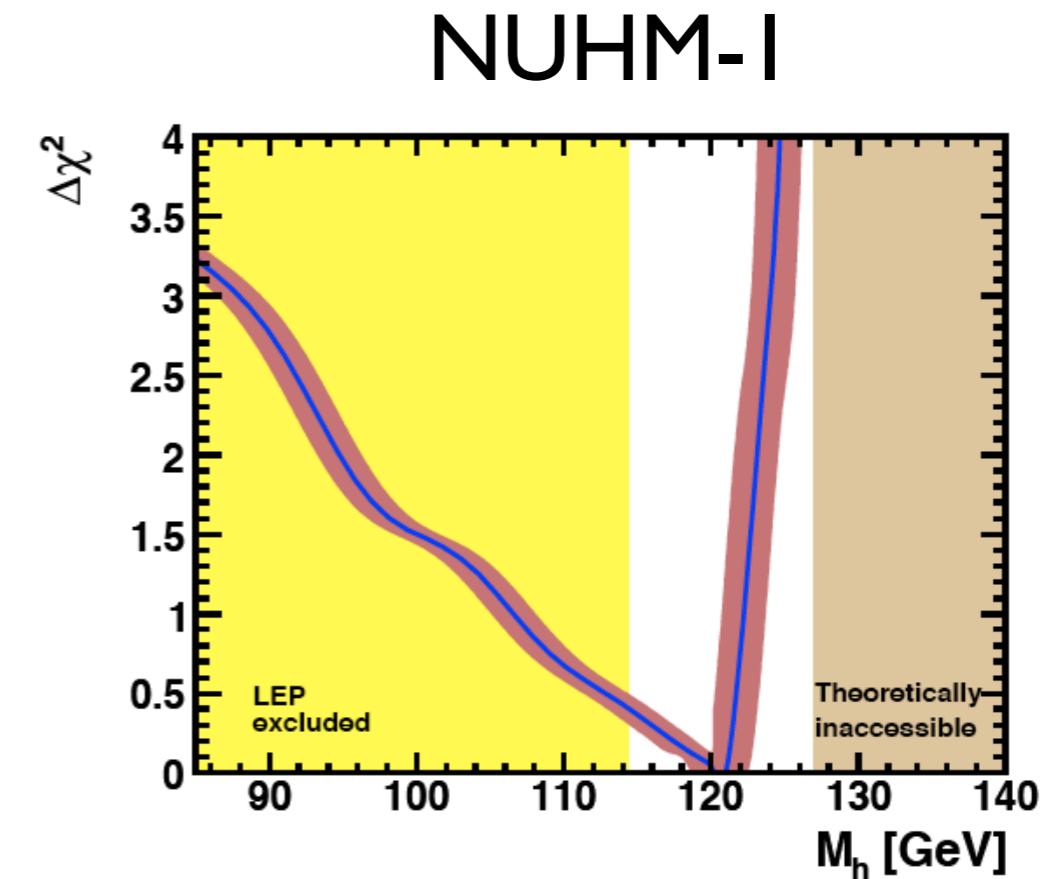
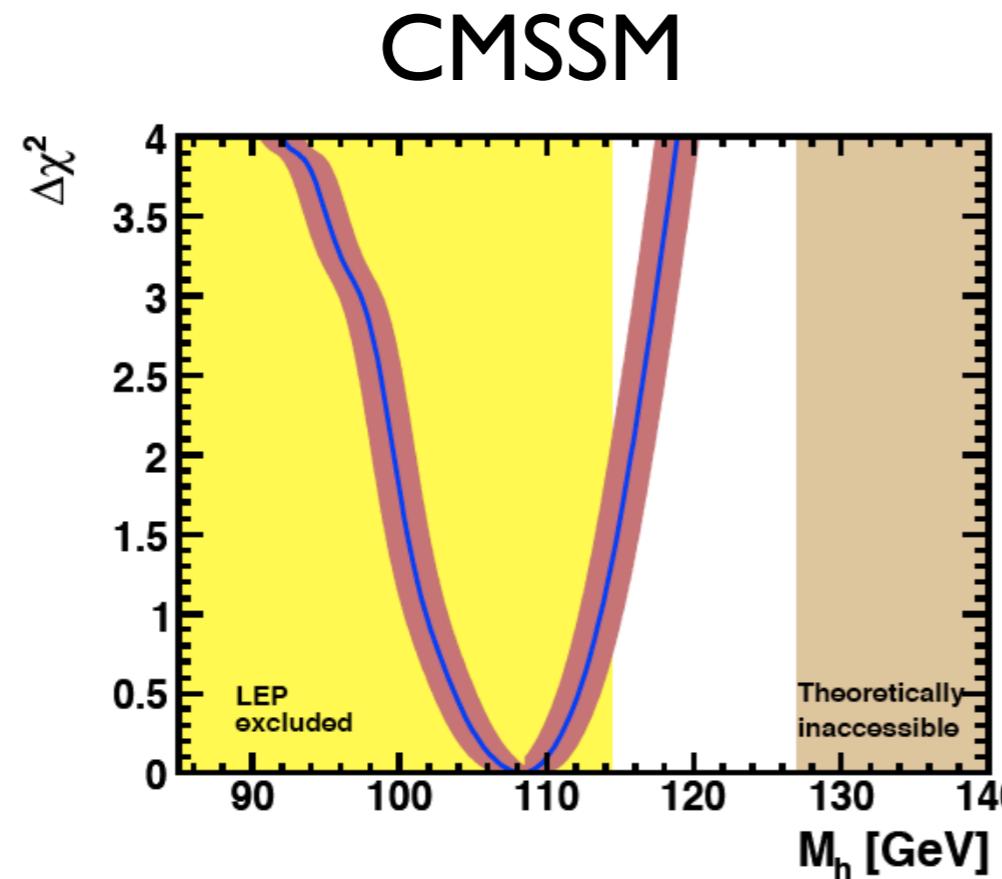
Direct DM search



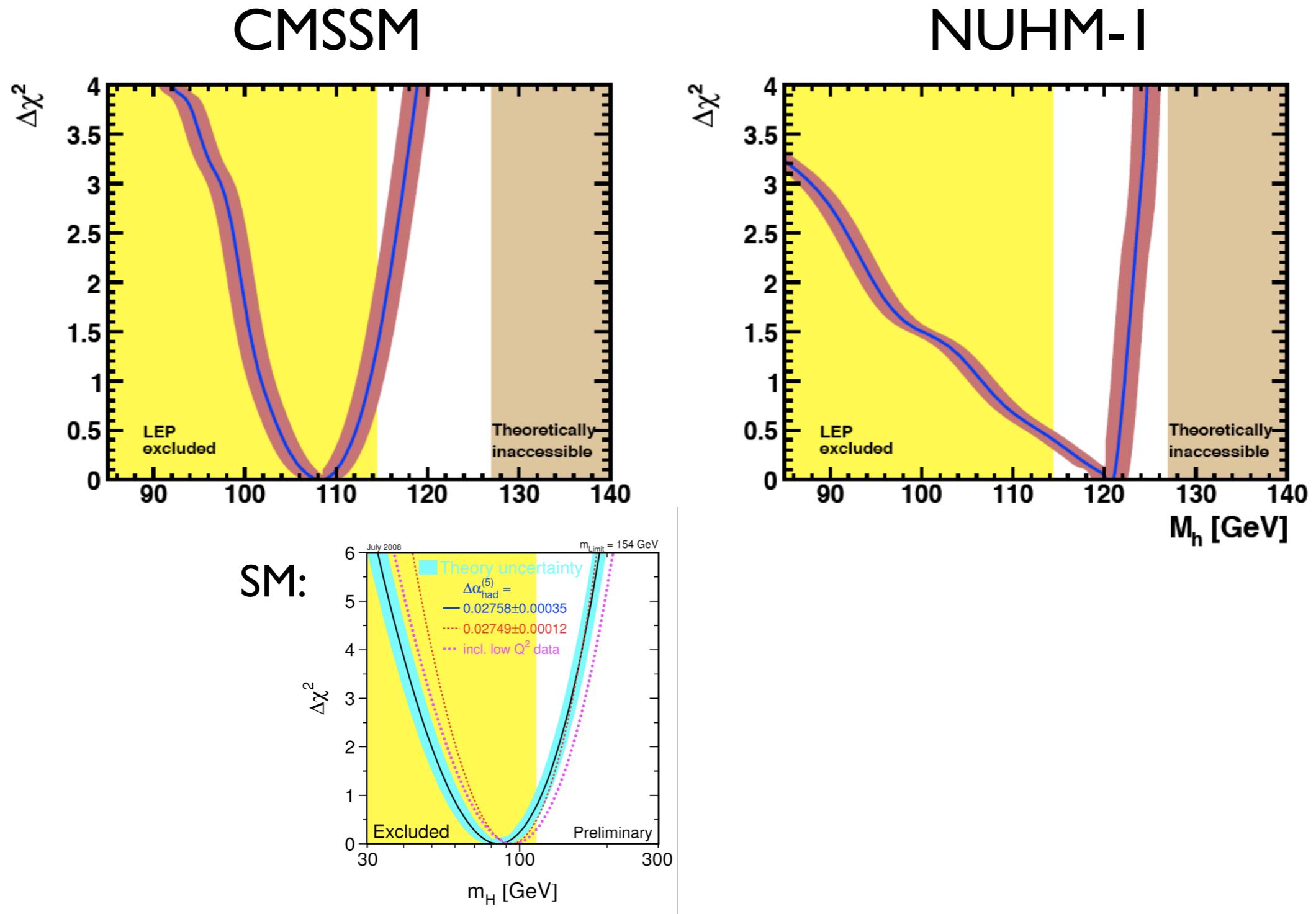
NUHM-I



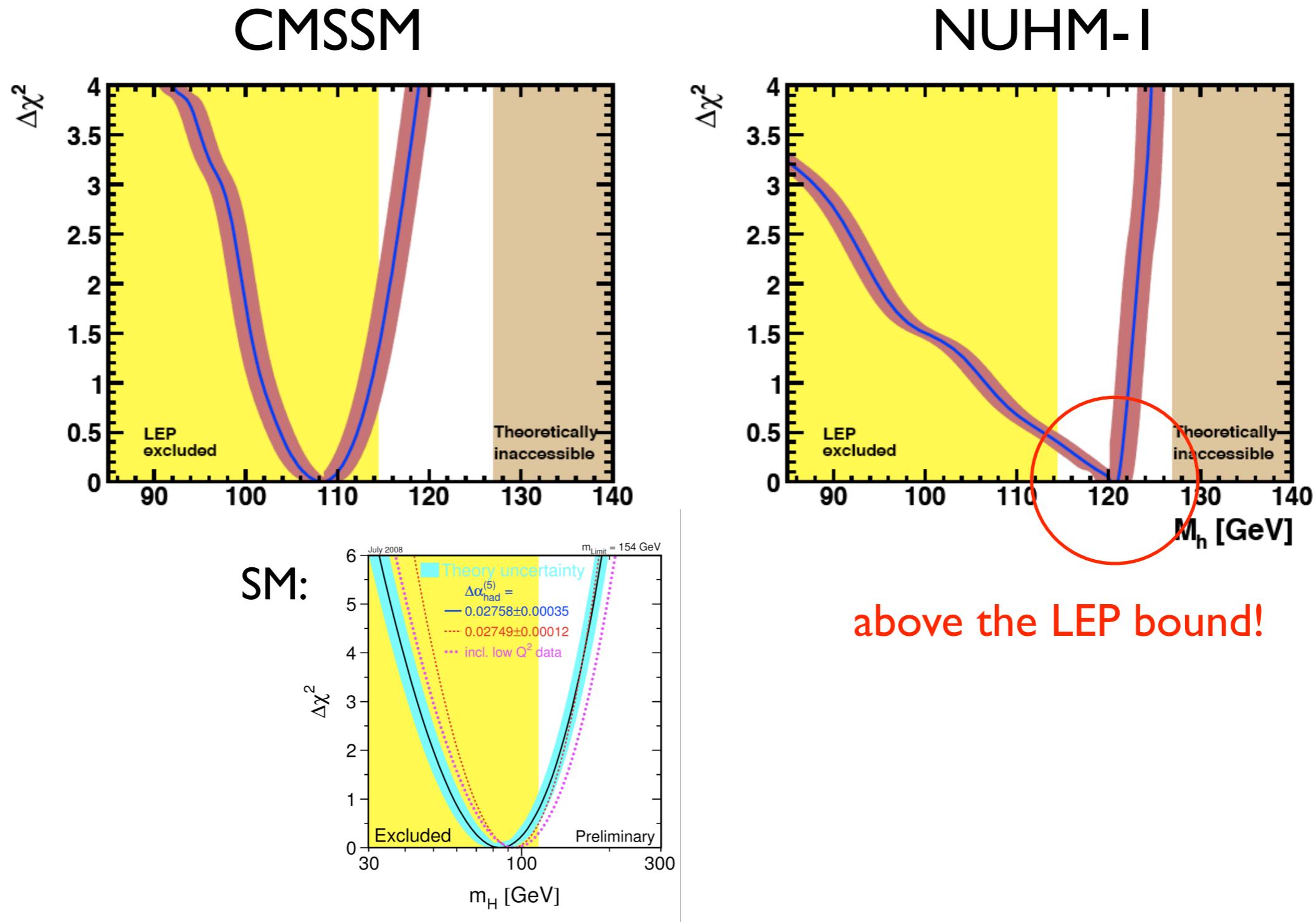
Fit of Higgs mass



Fit of Higgs mass



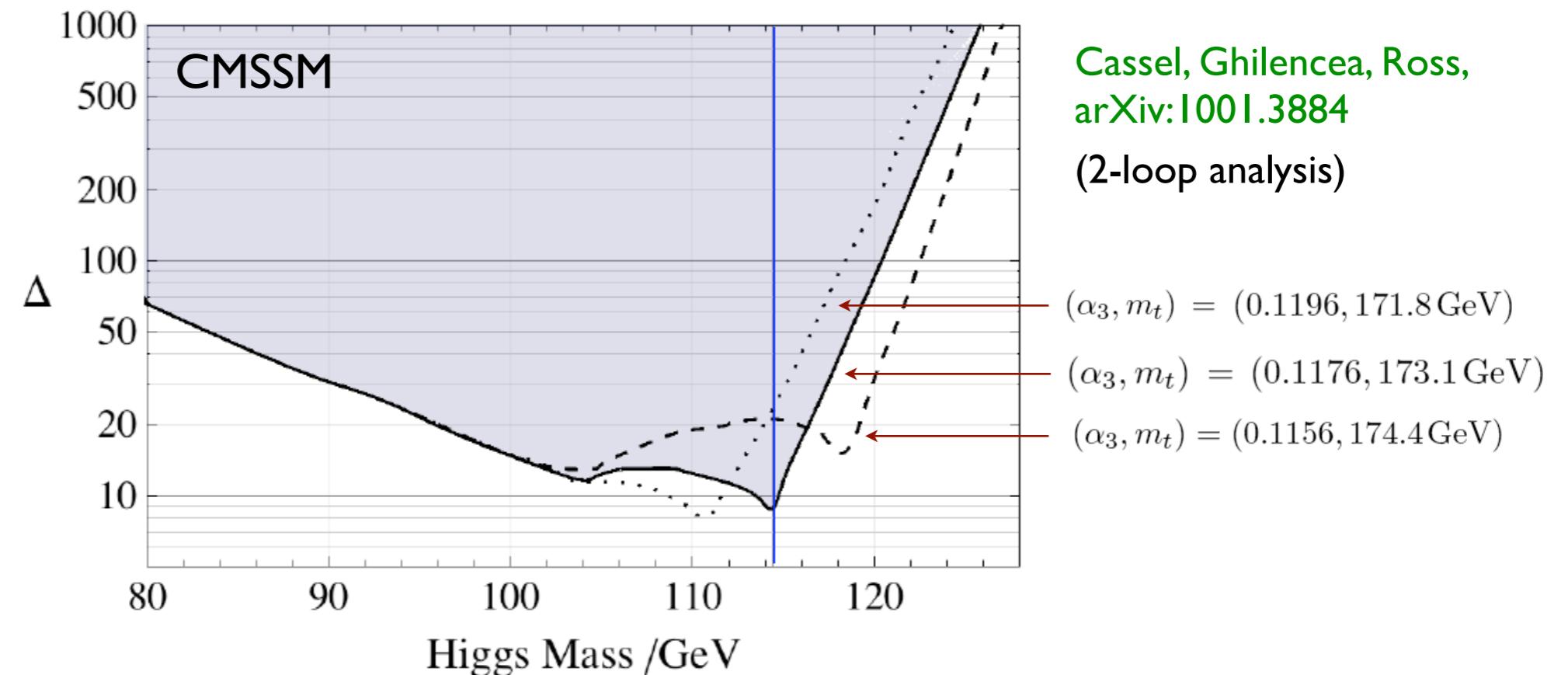
Fit of Higgs mass



The finetuning price

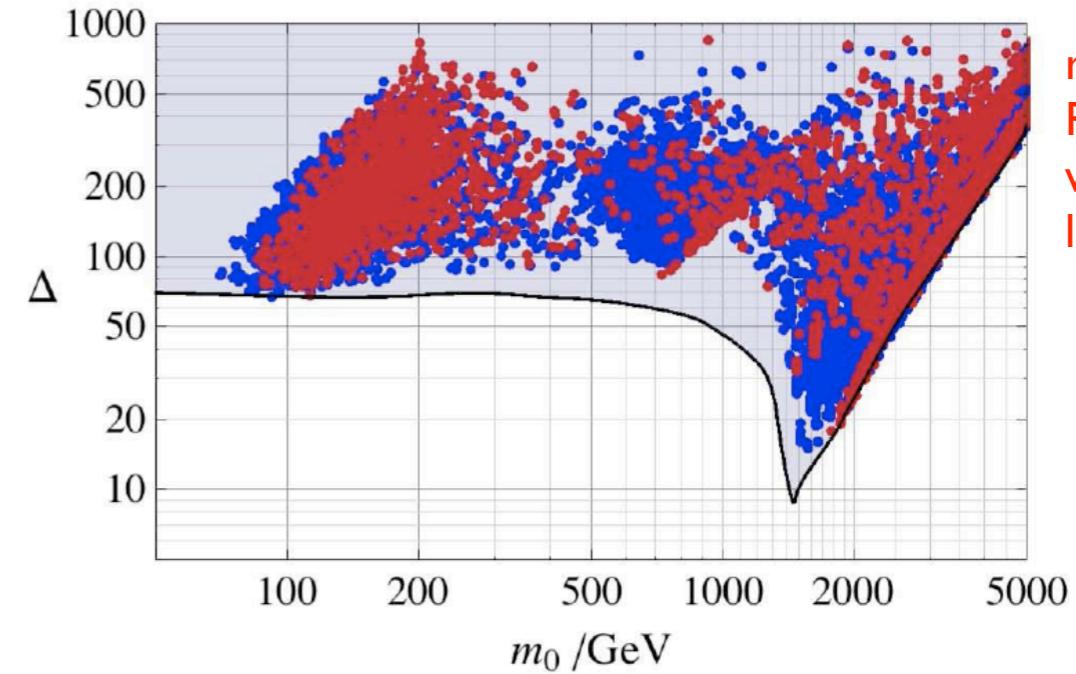
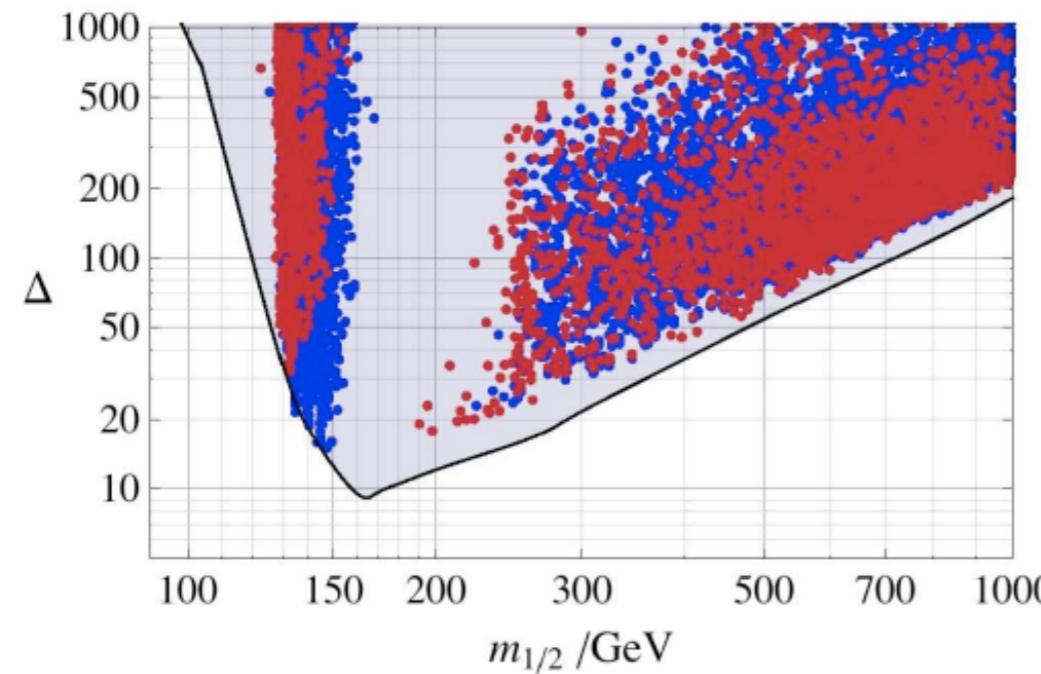
- Finetuning = sensitivity of EW scale to input parameters

$$\frac{M_Z^2}{2} \approx -m_{H_2}^2 - |\mu|^2$$



$$\Delta \equiv \max \left| \Delta_p \right|_{p=\{\mu_0^2, m_0^2, m_{1/2}^2, A_0^2, B_0^2\}}, \quad \Delta_p \equiv \frac{\partial \ln v^2}{\partial \ln p}$$

CMSSM low finetuning



NB: points with lowest finetuning lie in the focus point region

- gaugino-higgsino mixing
- light gluino
- Xsections a few pb at 7TeV LHC

$BR(\tilde{g} \rightarrow \tilde{\chi}_i^0 g) \sim 10 - 20\%$

$BR(\tilde{g} \rightarrow \tilde{\chi}_i^0 b\bar{b}, \tilde{\chi}_i^\pm tb) \sim 20\%$

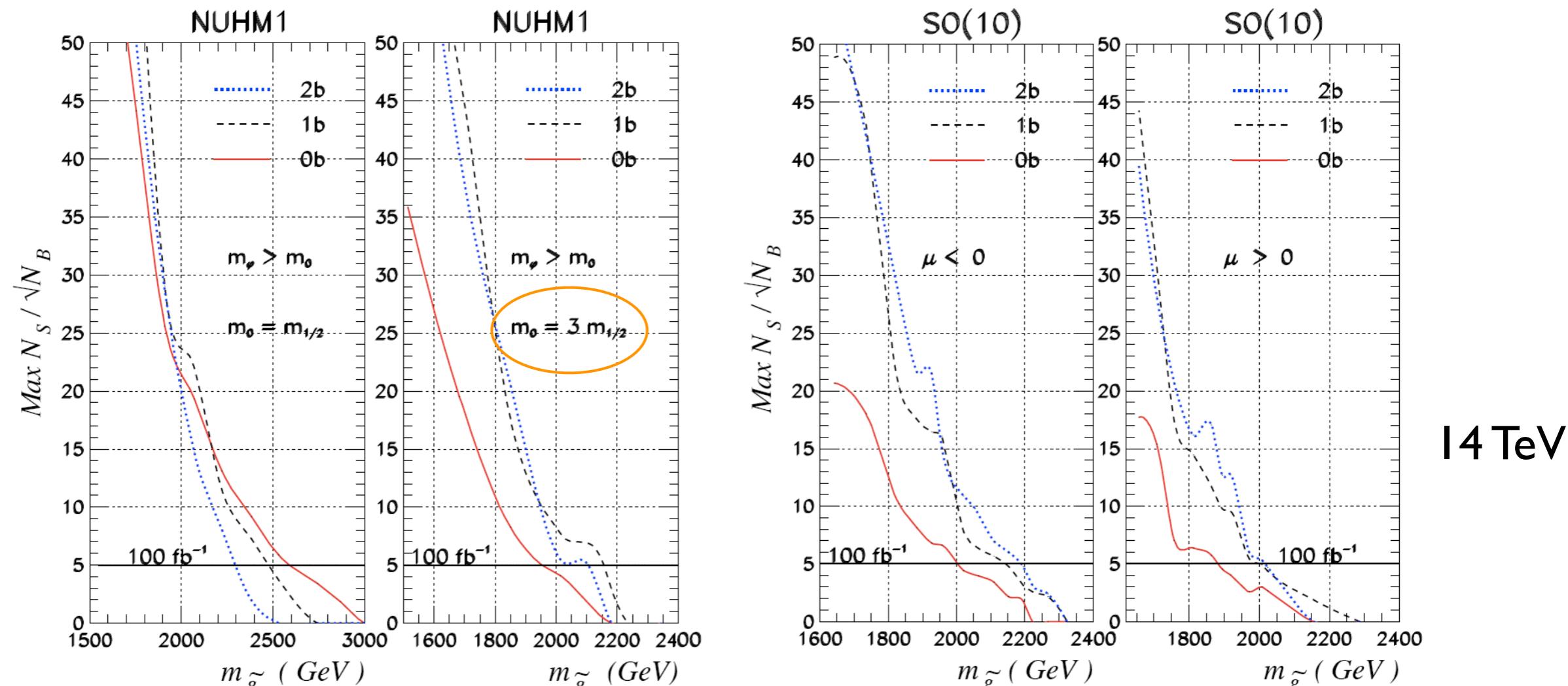
h^0	114.5	$\tilde{\chi}_1^0$	79	\tilde{b}_1	1147	\tilde{u}_L	1444
H^0	1264	$\tilde{\chi}_2^0$	142	\tilde{b}_2	1369	\tilde{u}_R	1446
H^\pm	1267	$\tilde{\chi}_3^0$	255	$\tilde{\tau}_1$	1328	\tilde{d}_L	1448
A^0	1264	$\tilde{\chi}_4^0$	280	$\tilde{\tau}_2$	1368	\tilde{d}_R	1446
\tilde{g}	549	$\tilde{\chi}_1^\pm$	142	$\tilde{\mu}_L$	1406	\tilde{s}_L	1448
$\tilde{\nu}_\tau$	1366	$\tilde{\chi}_2^\pm$	280	$\tilde{\mu}_R$	1406	\tilde{s}_R	1446
$\tilde{\nu}_\mu$	1404	\tilde{t}_1	873	\tilde{e}_L	1406	\tilde{c}_L	1444
$\tilde{\nu}_e$	1404	\tilde{t}_2	1158	\tilde{e}_R	1406	\tilde{c}_R	1446

upper limits
for $\Delta < 100$

\tilde{g}	χ_1^0	χ_2^0	χ_3^0	χ_4^0	χ_1^\pm	χ_2^\pm	\tilde{t}_1	\tilde{t}_2	\tilde{b}_1	\tilde{b}_2
1720	305	550	660	665	550	670	2080	2660	2660	3140

Importance of b-tagging

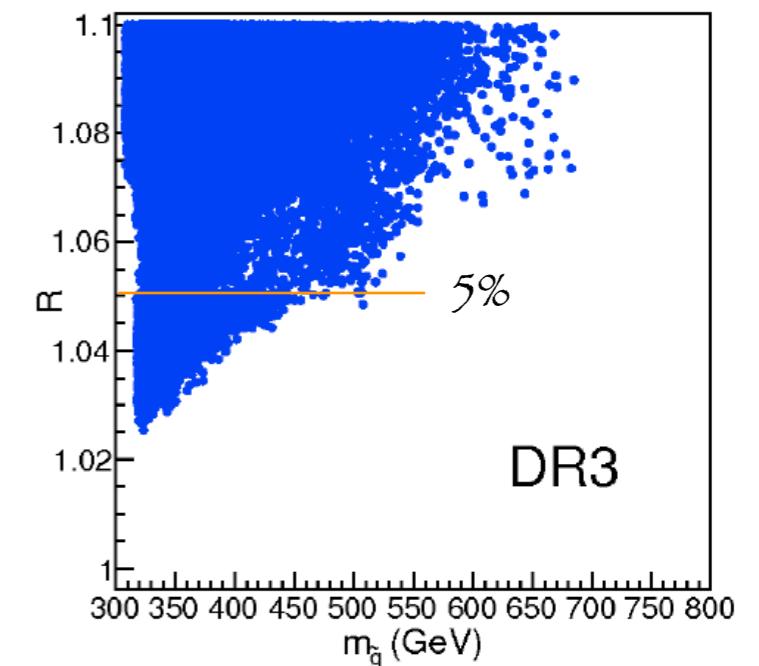
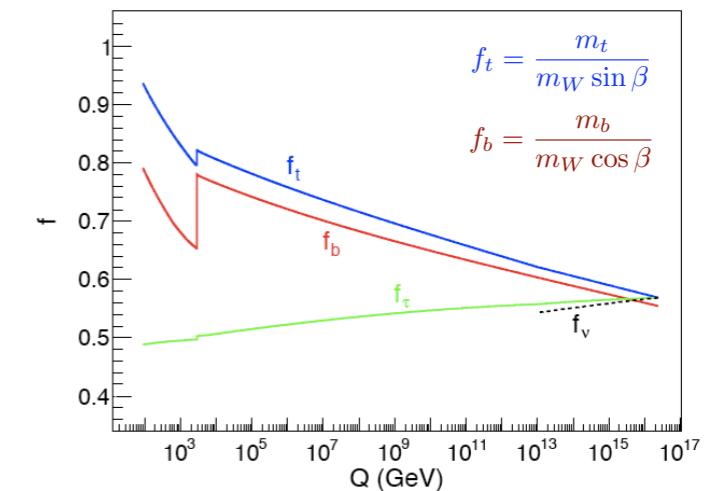
- Requiring 1, 2 or more b-jets can significantly enhance the signal/bg ratio in certain scenarios, e.g., 15-20% in the CMSSM focus point region.



- Typical if 3rd generation is lighter than 1st/2nd generation and $m_{\tilde{g}} \ll m_{\tilde{q}}$; enhances gluino decays into t or b via on- or off-shell stop/sbottom

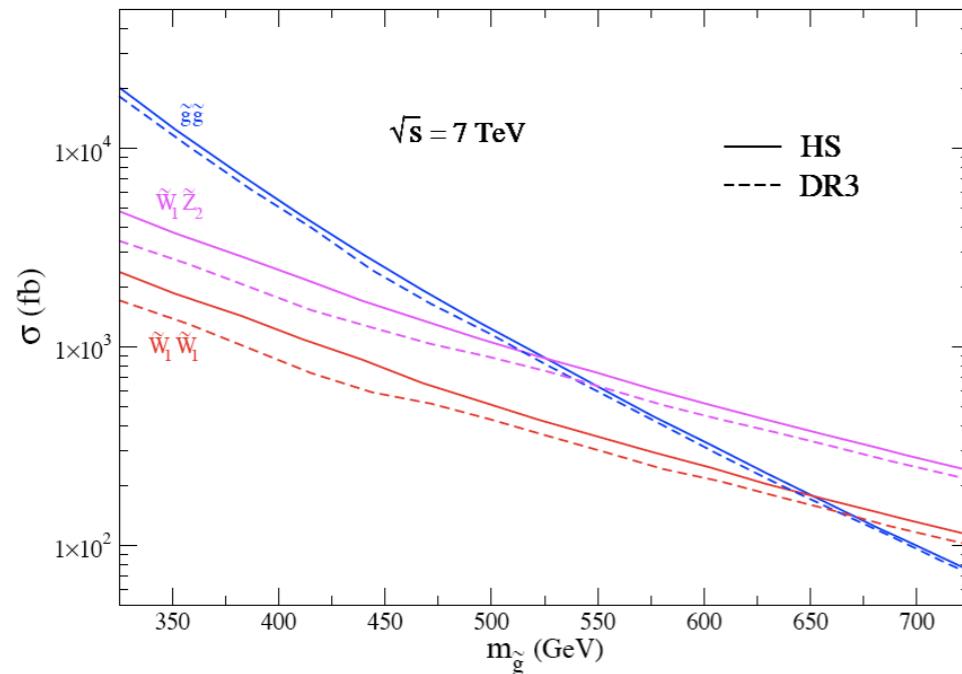
SO(10): Yukawa-unified SUSY

- SUSY GUTs based on SO(10) are particularly compelling
 - unify all matter of one generation in a 16-dim spinor (incl. r.h. neutrino!)
 - automatic anomaly cancellation
- The simplest realizations (Higgs in a 10-plet) require t-b-tau Yukawa coupling unification in addition to gauge coupling unification at MgUT.
- Typical mass spectrum:
 - 1st/2nd generation scalars in the 5-15 TeV range,
 - 3rd gen. scalars, Higgses and higgsinos ca. 1-3 TeV,
 - gauginos: LSP \sim 50-80 GeV, gluino \sim 300-500 GeV!
 - c.f. “effective SUSY” by Cohen, Kaplan, Nelson ’1996



See Baer et al, 0801.1831, 0809.0710, 0812.2693,
0908.0134, 0910.2988, 0911.4739.

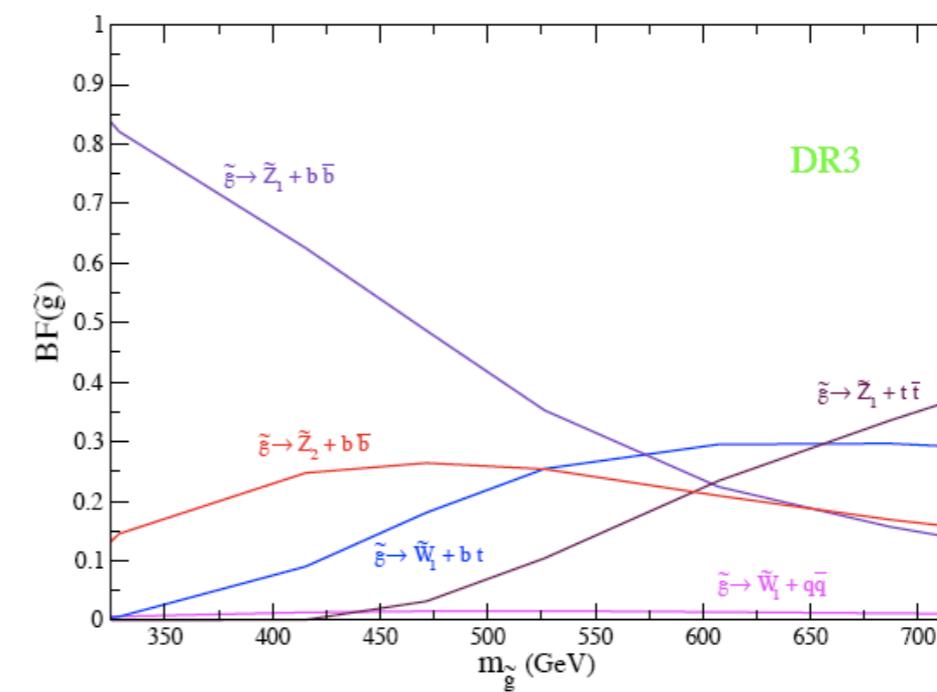
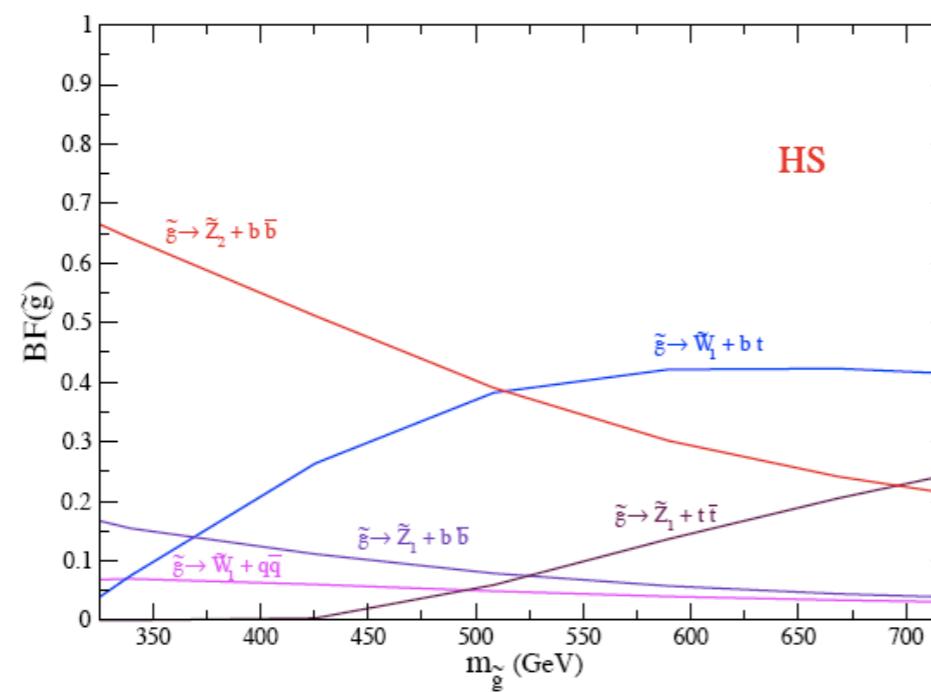
LHC reach at 7 TeV



Gluino-pair prod. dominated by gg fusion.
 $\sigma(\text{LO}) \sim 1 \text{ pb}$ at $m(\text{gluino}) \sim 525 \text{ GeV}$

We consider model lines for HS and DR3 *)
 cases as funct. of $m(\text{gluino})$ up to 700 GeV.

Gluino decays are dominated by heavy
 flavours: $\tilde{g} \rightarrow \tilde{\chi}_{1,2}^0 b\bar{b}$, $\tilde{\chi}_1^\pm tb$



*) HS: only Higgs masses split, DR3: full D-term splitting, see 0908.0134

LHC reach at 7 TeV

(theorists' simulation)

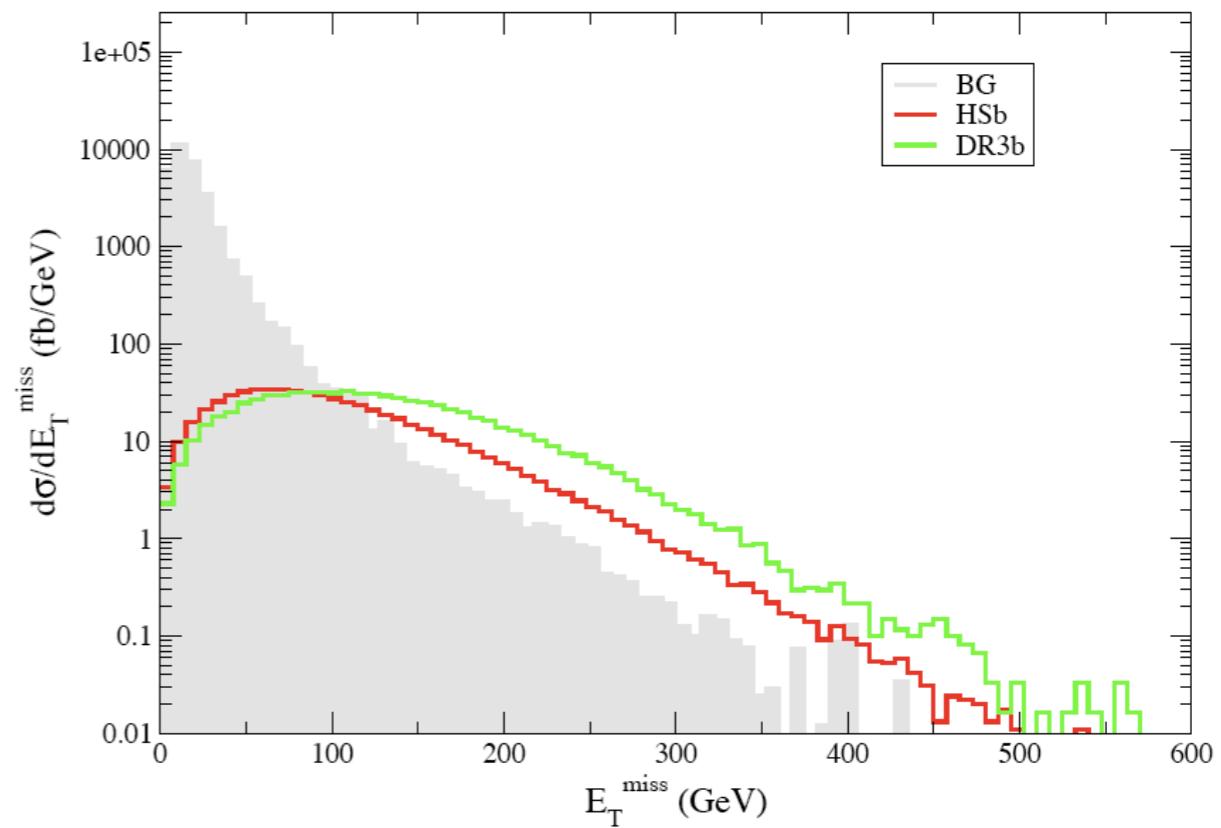
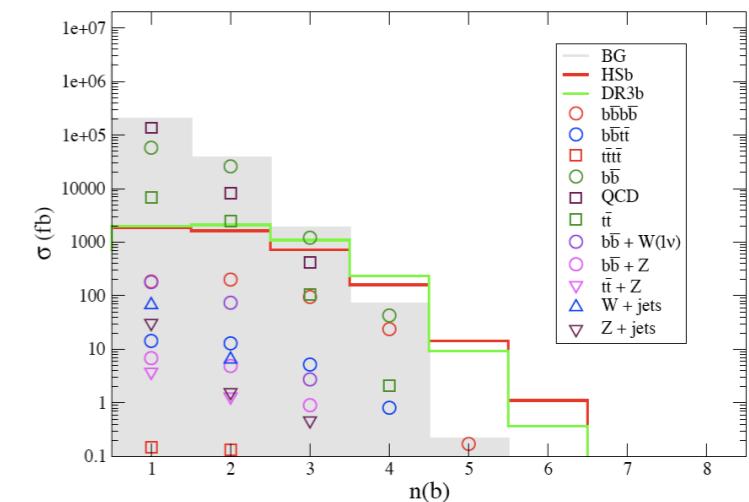
Event simulation:

- Isajet 7.79 for the signal
- QCD, 2- and 3-bdy BGs with Alpgen
- 4t, 4b, 2t2b BGs with Madgraph
- Phythia for showering and hadronization
- Generic LHC detector simulation

Basic Cuts “C0”:

- $n(\text{jets}) \geq 4$ with $p_T > 50\text{GeV}$
- hardest jet $p_T > 100\text{ GeV}$
- $S_T \geq 0.2$ (transv sphericity)
- $n(b) \geq 1$ (b-eff. 60%)

Results after C0-based selection			
	$\sigma(n(b) \geq 3)$	$\sigma(n(b) \geq 4)$	$\sigma(\text{OS})$
HSb	899 fb	176 fb	99 fb
DR3b	1334 fb	243 fb	22 fb
BG	1911 fb	70 fb	11 fb



LHC reach at 7 TeV

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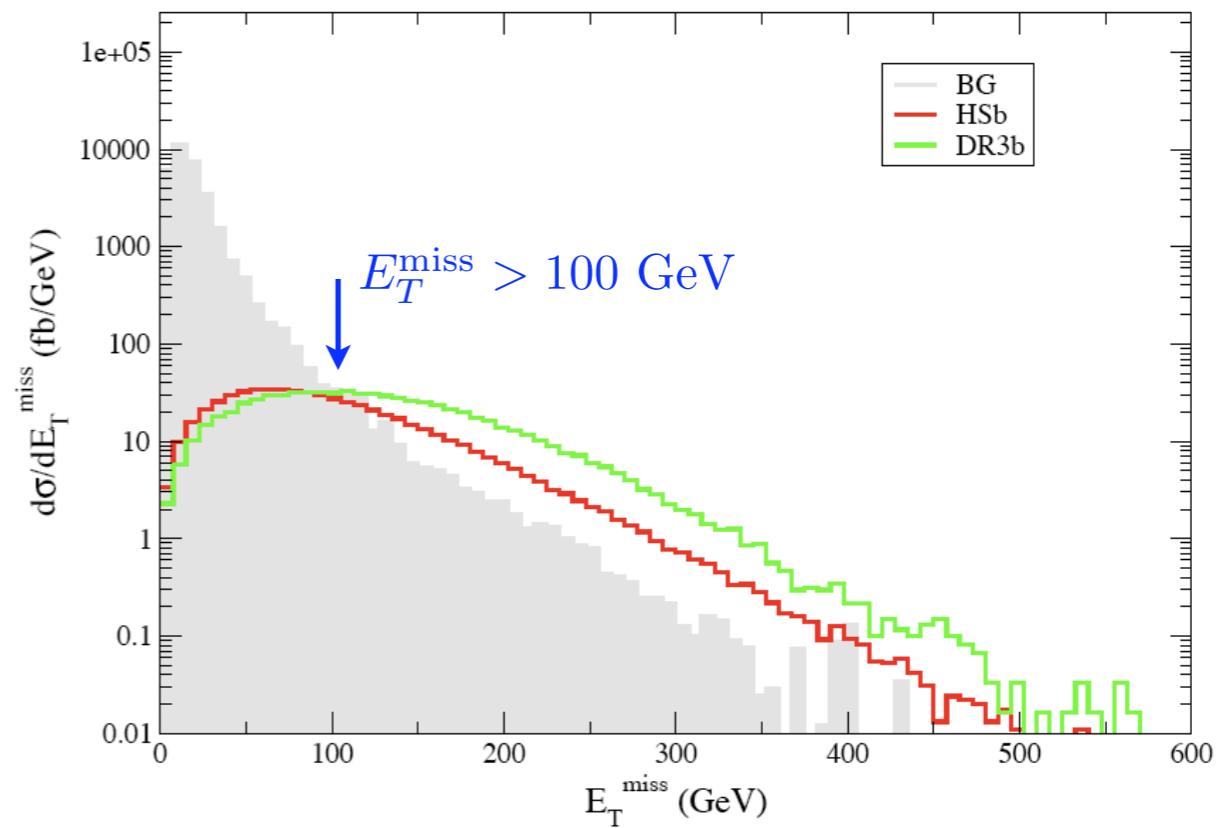
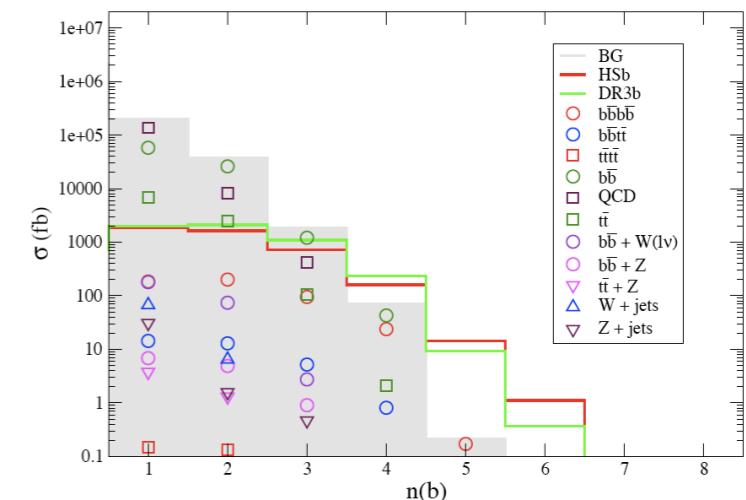
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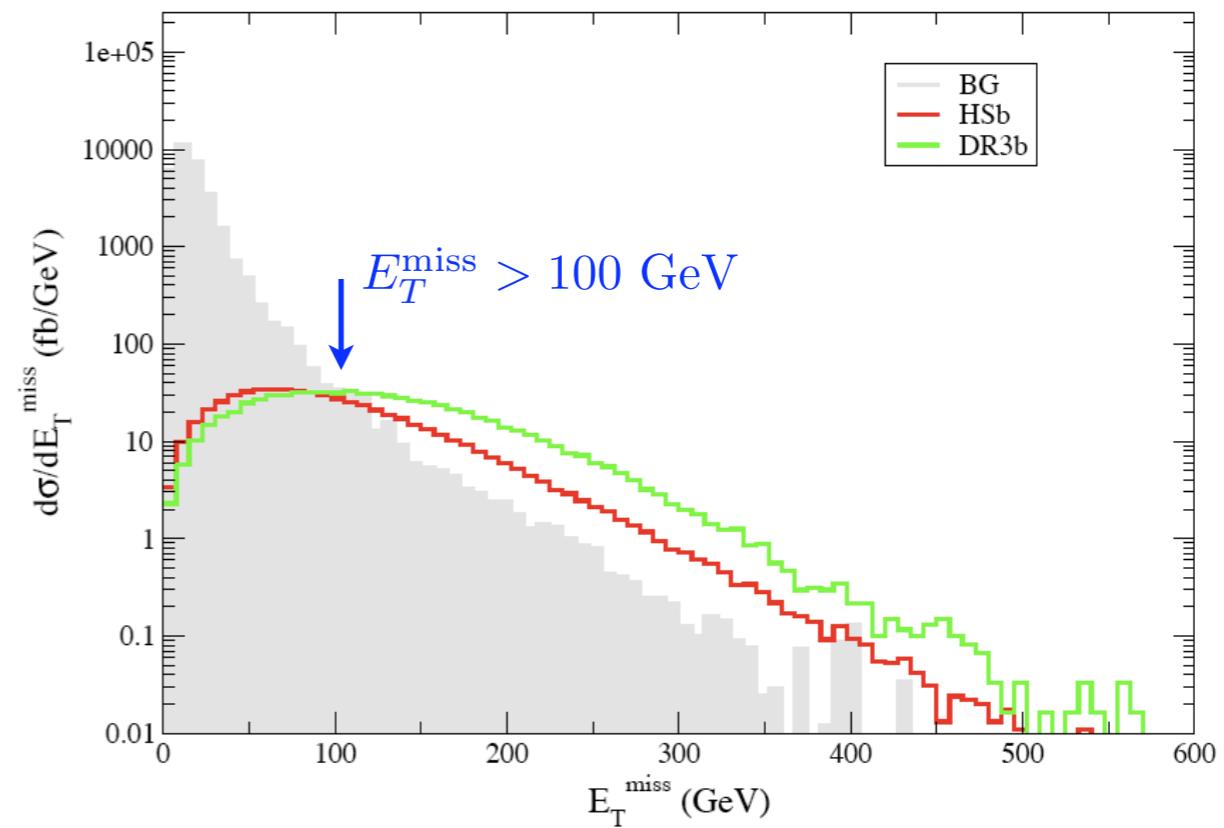
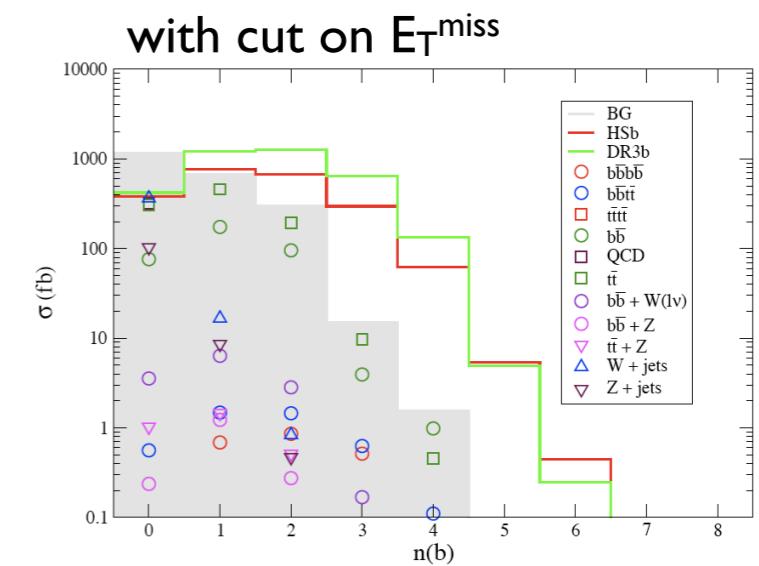
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Basic Cuts “C0”:

- $n(\text{jets}) \geq 4$ with $p_T > 50\text{GeV}$
- hardest jet $p_T > 100\text{ GeV}$
- $S_T \geq 0.2$ (transv sphericity)
- $n(b) \geq 1$ (b-eff. 60%)

Results after C1-based selection			
	$\sigma(n(b) \geq 3)$	$\sigma(n(b) \geq 4)$	$\sigma(\text{OS})$
HSb	364 fb	68 fb	81 fb
DR3b	782 fb	139 fb	23 fb
BG	16 fb	2 fb	9 fb

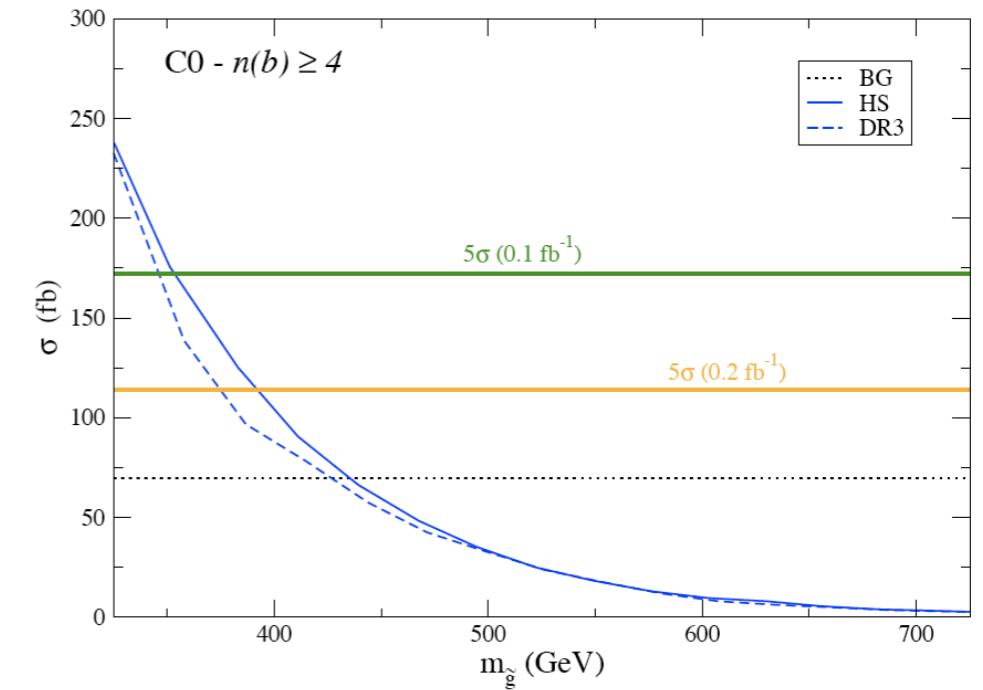
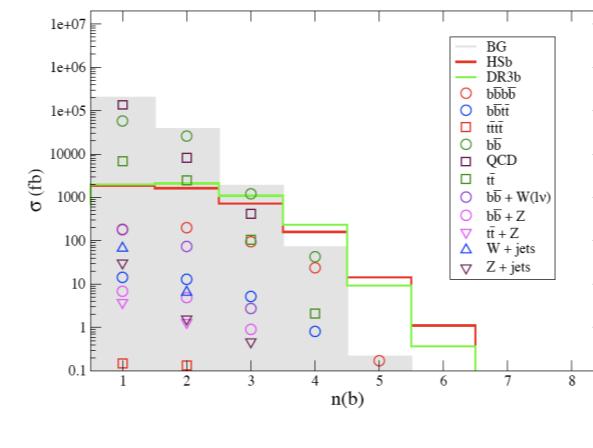


LHC reach at 7 TeV

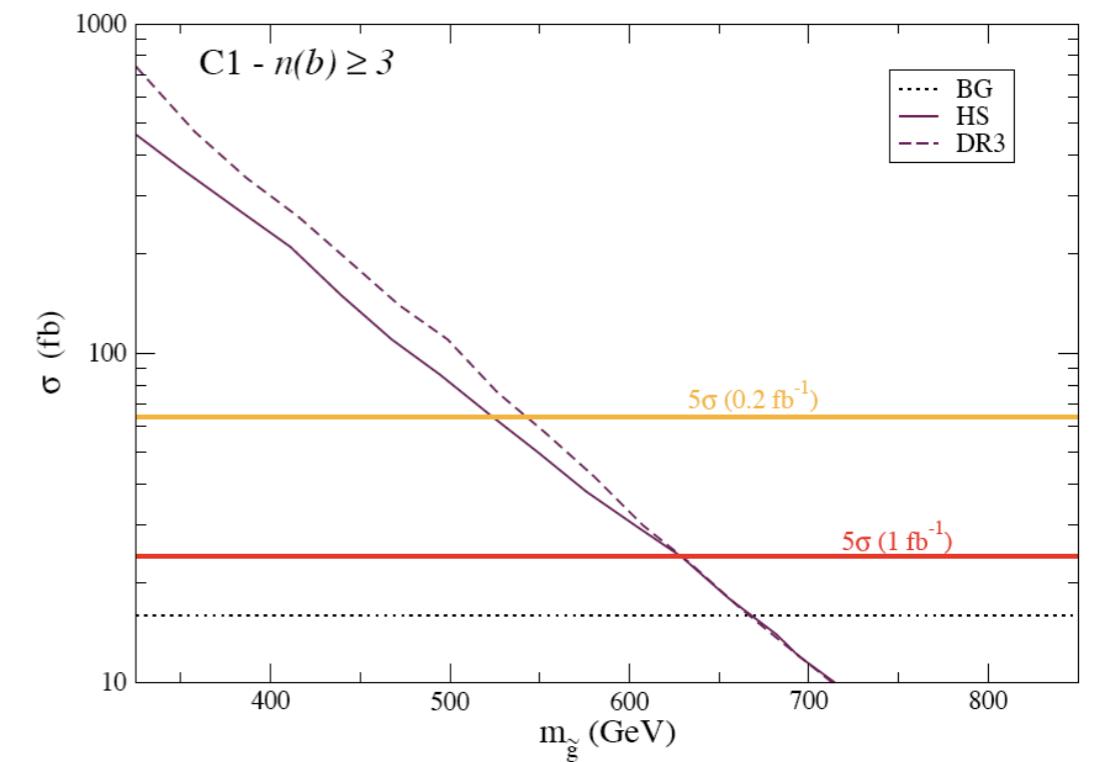
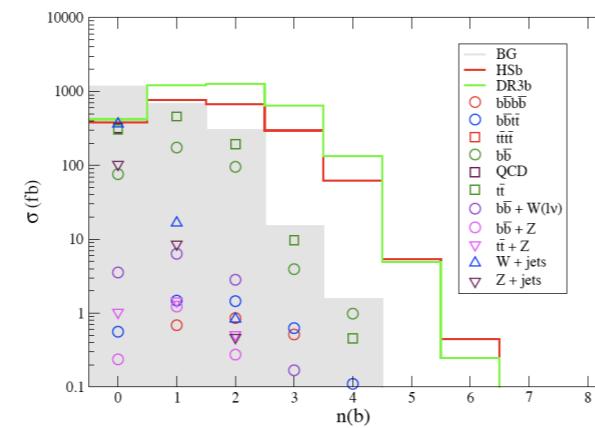
(theorists' simulation)

Without missing energy measurement:
up to $m(\text{gluino})=400$ GeV with 0.2 fb^{-1} of data
requiring 4 b-jets

\lesssim Tevatron with 10 fb^{-1}

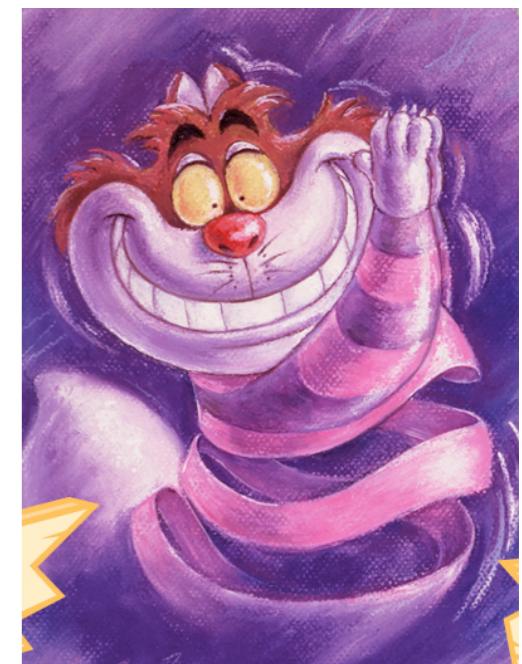


With reliable MET measurement:
reach up to $m(\text{gluino})=540-630$ GeV
with $0.2-1 \text{ fb}^{-1}$ of data,
 $n(b) \geq 3$



Conclusions

- MSSM and constrained versions thereof are in a good shape
 - large regions of parameter space compatible with exp. constraints
 - tension with LEP Higgs limit can be evaded
 - finetuning price may not be too high (CMSSM preferred)
- Important development of fitting tools
(Sfitter, Fittino, SuperBayes, Mastercode, many private codes)
- Much of recent work has gone into Bayesian inference;
fits show strong prior dependence → need more data
- Tevatron/LHC have good potential for “effective SUSY”
(light gauginos [gluino], TeV-scale 3rd generation, multi-TeV 1st/2nd gen)
- Example: Yukawa-unified SUSY SO(10)
- b-tagging may be essential for early SUSY discovery

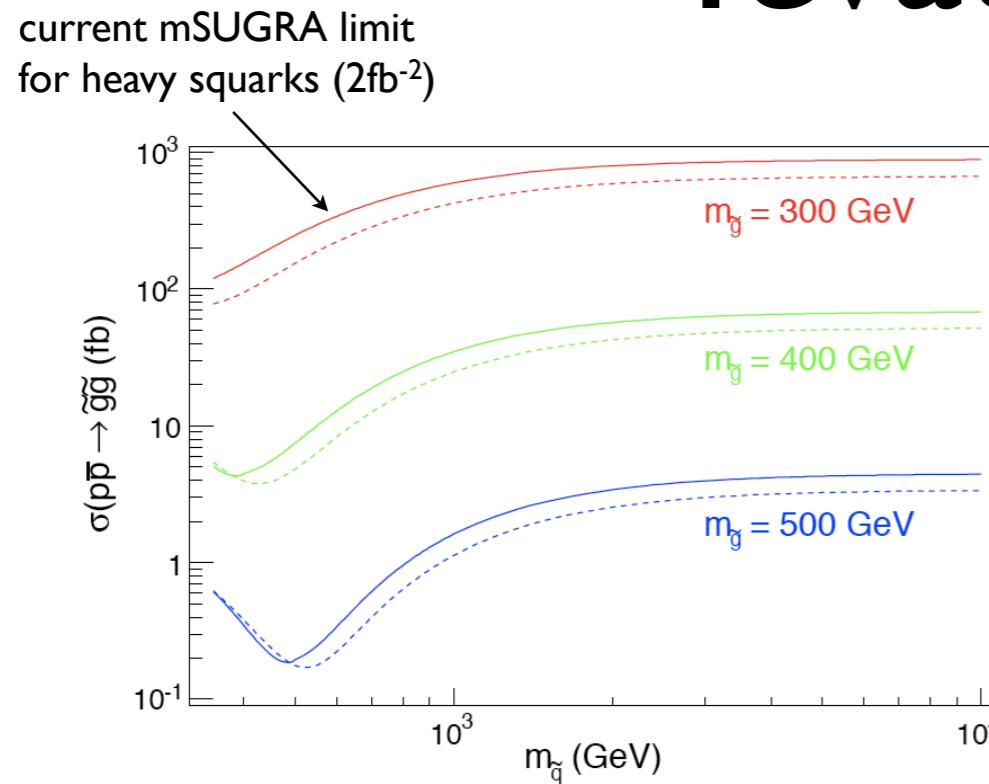




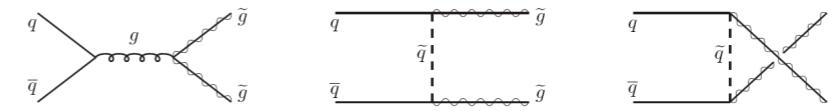
Exciting times !

Backup

Tevatron reach



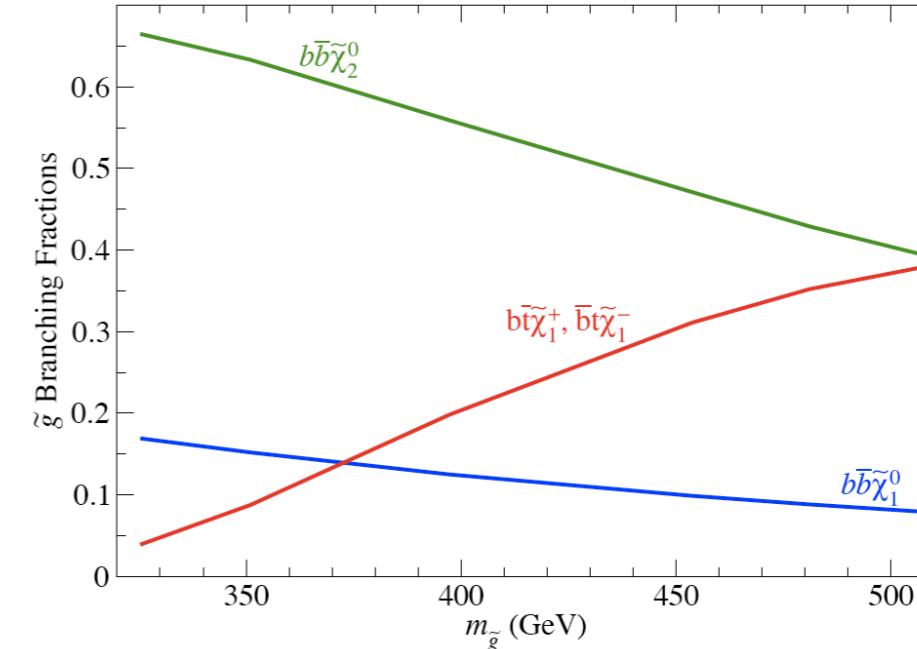
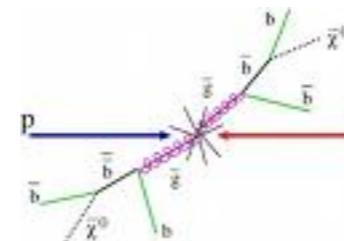
Gluino-pair prod. dominated by $q\bar{q}$ fusion.
Negative interference of s-, t-, u-channels
for $m(\text{squark}) \sim m(\text{gluino})$!



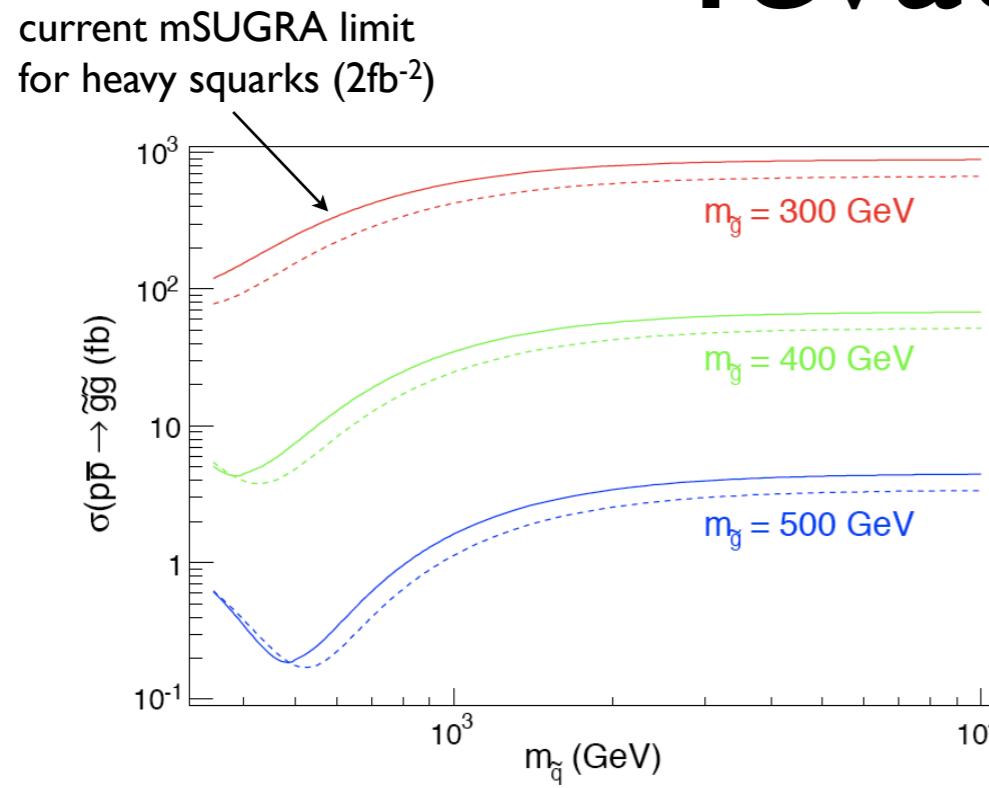
Xsection grows with increasing squark mass!

Gluino decays dominated by $\tilde{\chi}_2^0 b\bar{b}$ channel.
We adopt a YU model line by starting from
a HS point with $m_{16}=10\text{ TeV}$ and $R \sim 1.02$
and varying $m_{1/2}$.

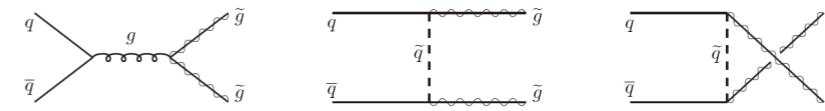
$$\begin{aligned} m_{1/2} &= 35 - 100 \text{ GeV}, \\ m_{\tilde{g}} &= 325 - 508 \text{ GeV}, \\ R &\rightarrow 1.07 \end{aligned}$$



Tevatron reach



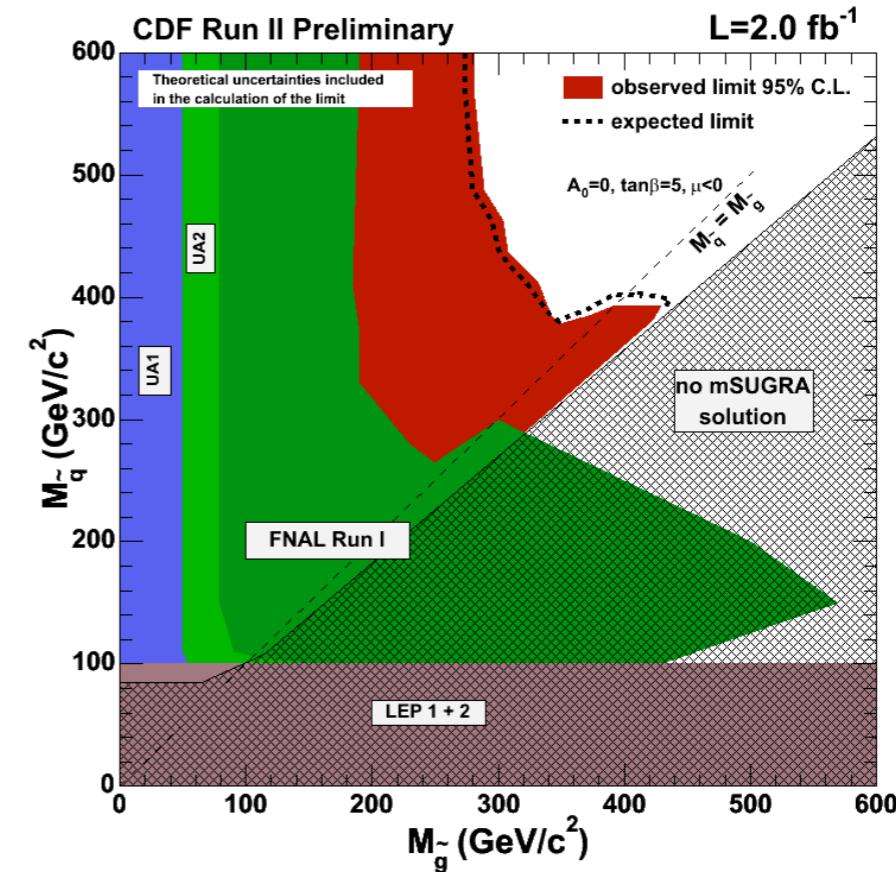
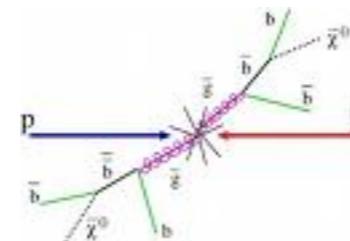
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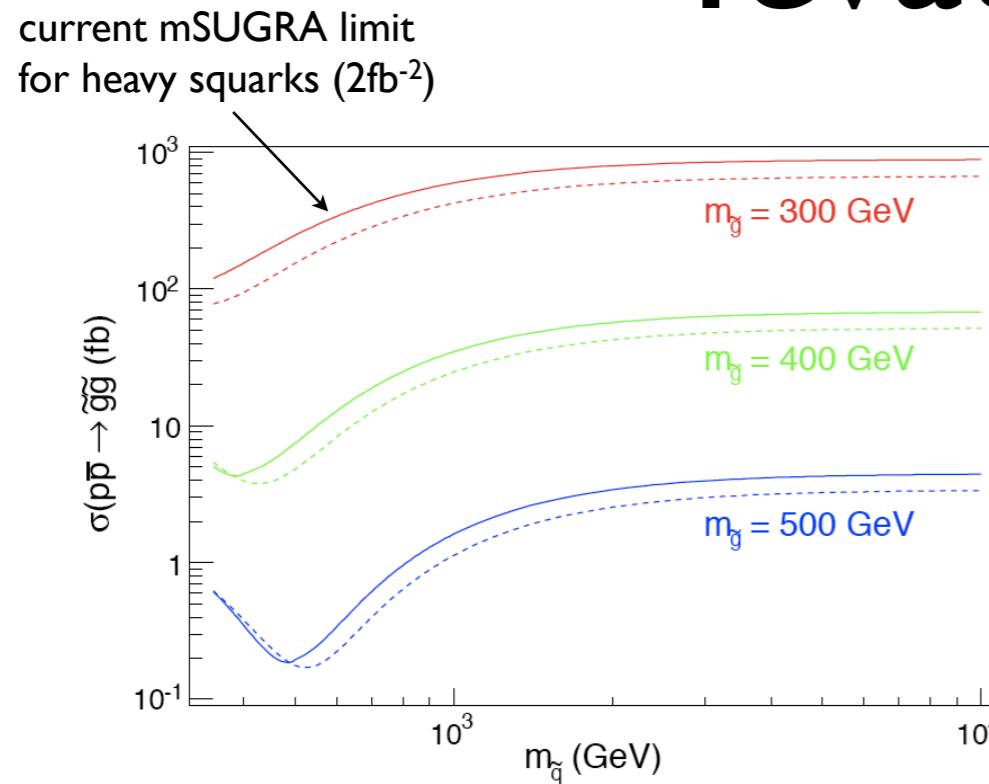
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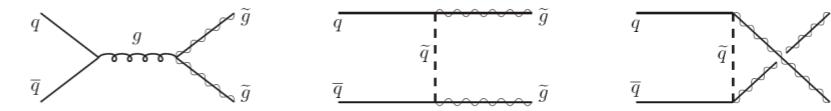
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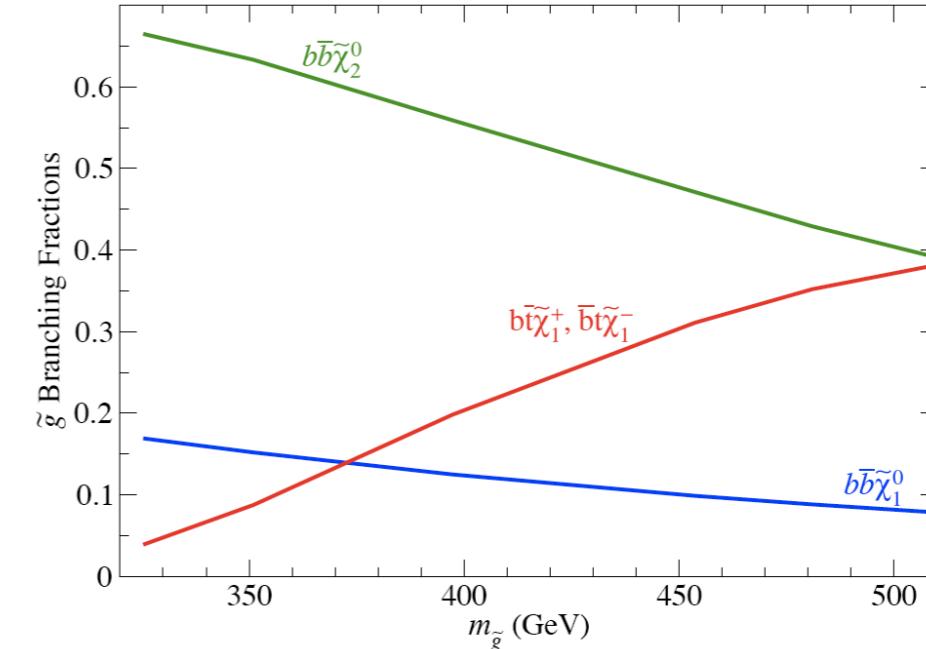
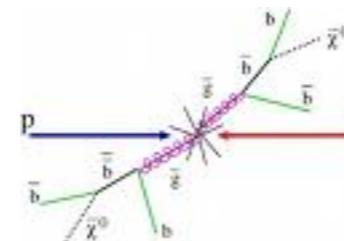
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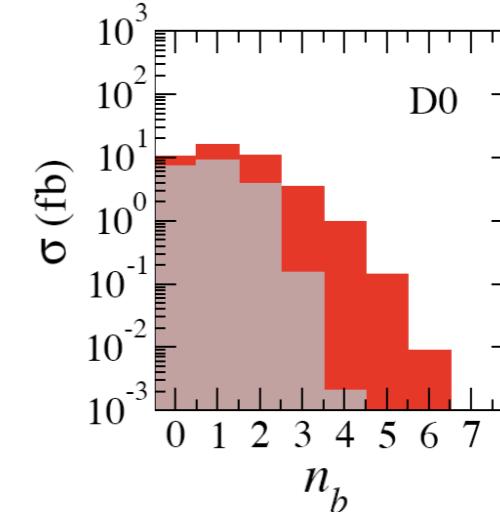
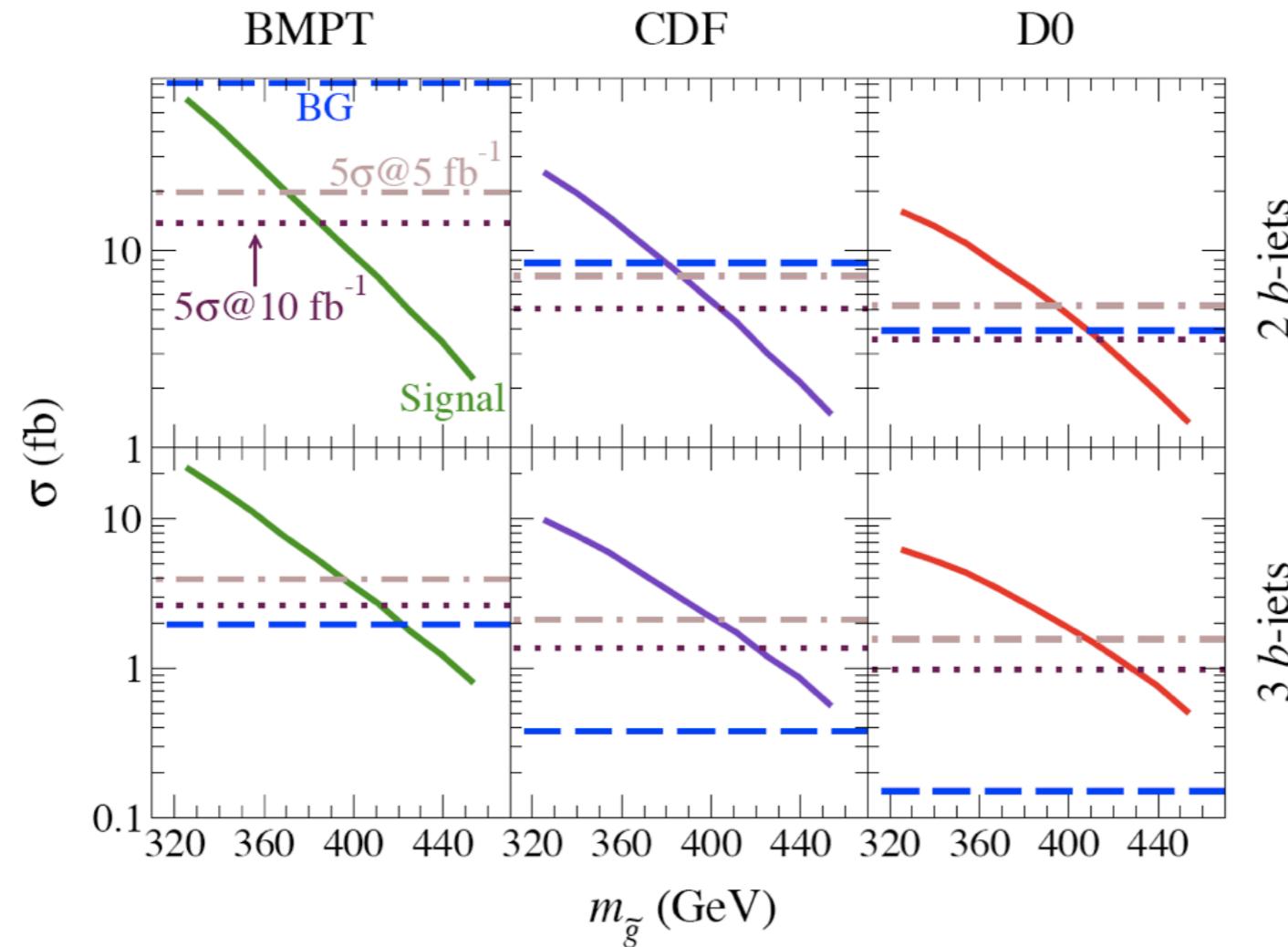
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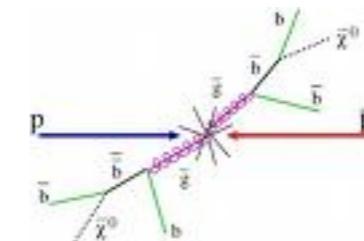
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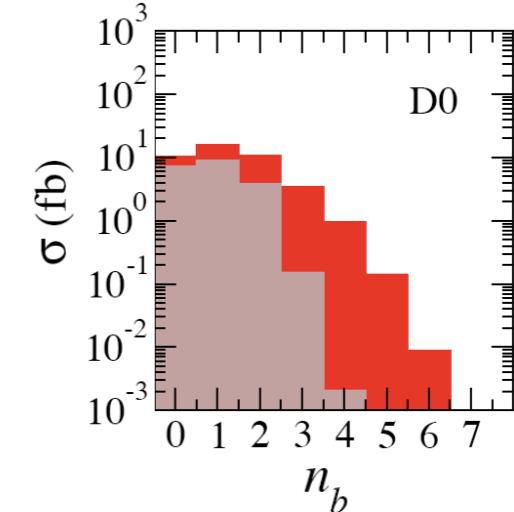
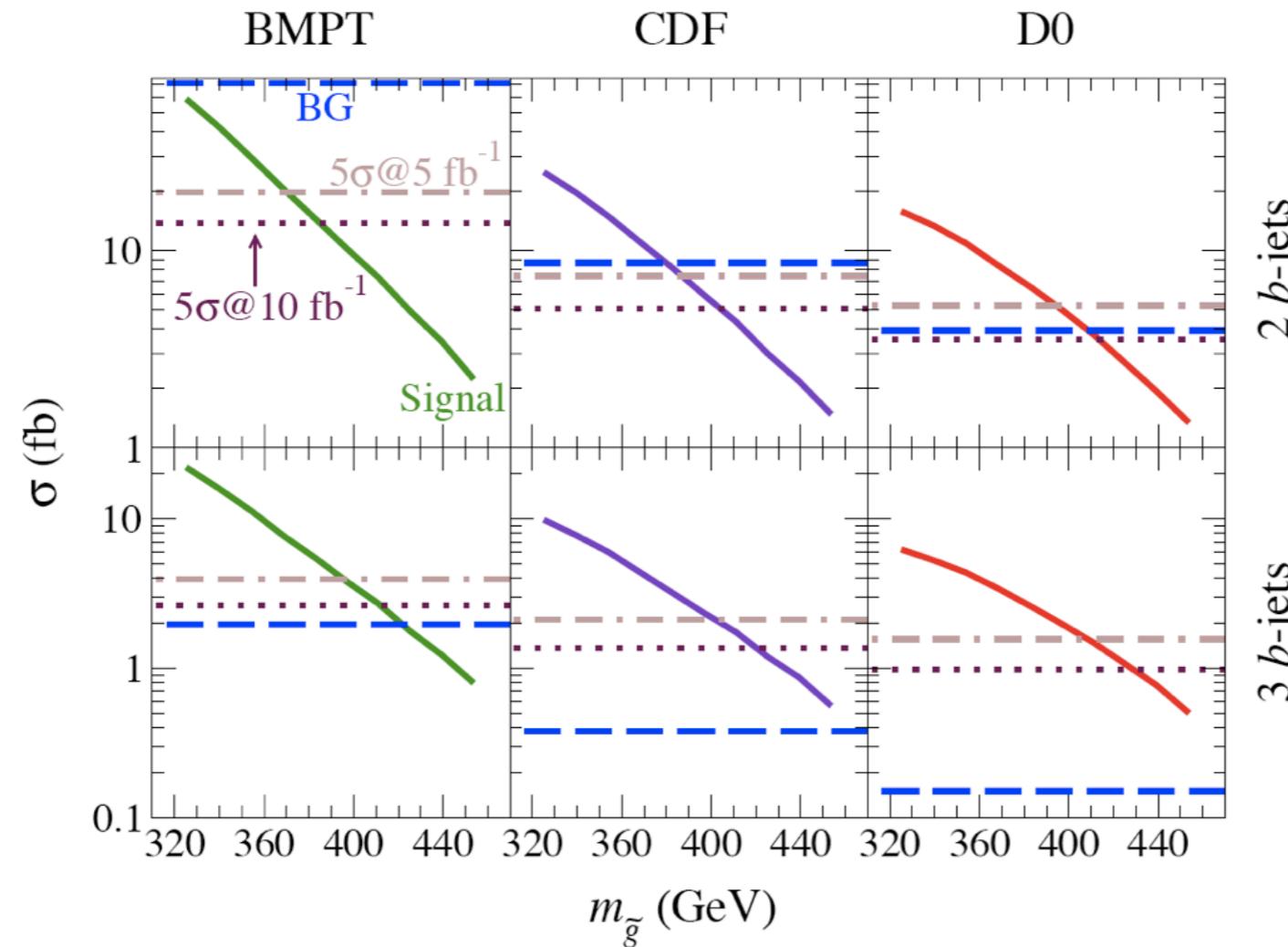
Tevatron reach



cuts	E_T^{miss}	H_T	$E_T(j1)$	$E_T(j2)$	$E_T(j3)$	$E_T(j4)$
BMPT	≥ 75 GeV	–	15	15	15	15
CDF	≥ 90 GeV	280	95	55	55	25
D0	≥ 100 GeV	400	35	35	35	20

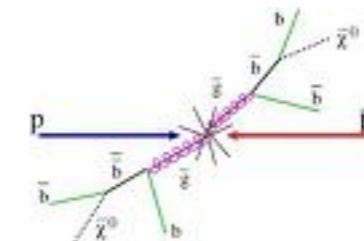


Tevatron reach



With “D0” cuts and requiring ≥ 3 b-jets, 5 σ discovery reach with 10 fb^{-1} about $m(\text{gluino})=430 \text{ GeV}!$

cuts	E_T^{miss}	H_T	$E_T(j1)$	$E_T(j2)$	$E_T(j3)$	$E_T(j4)$
BMPT	$\geq 75 \text{ GeV}$	–	15	15	15	15
CDF	$\geq 90 \text{ GeV}$	280	95	55	55	25
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Conditions for Yukawa unification (YU)

★ For $\mu > 0$, as preferred by $b \rightarrow s\gamma$, Yukawa unification (YU) can only be realized for very particular parameter relations

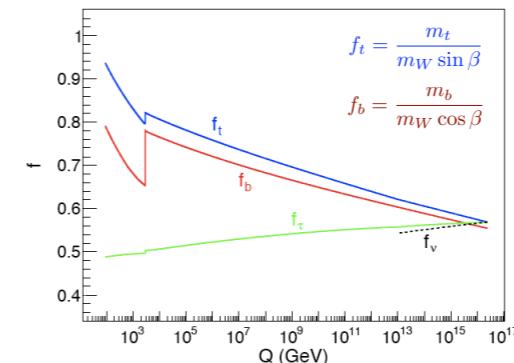
- $m_{16} \sim 5 - 15$ TeV,
- $A_0^2 \simeq 2m_{10}^2 \simeq 4m_{16}^2$, $(A_0 < 0)$
- $m_{1/2} \ll m_{16}$,
- $\tan \beta \sim 50$.

★ D-term splitting

$$\begin{aligned} m_Q^2 = m_E^2 = m_U^2 &= m_{16}^2 + M_D^2 \\ m_D^2 = m_L^2 &= m_{16}^2 - 3M_D^2 \\ m_{\tilde{\nu}_R}^2 &= m_{16}^2 + 5M_D^2 \\ m_{H_{u,d}}^2 &= m_{10}^2 \mp 2M_D^2. \end{aligned}$$

“just-so” Higgs splitting (HS) case

NB: we need $m_{H_u}^2 < m_{H_d}^2$ at M_{GUT} , so $M_D^2 > 0$.



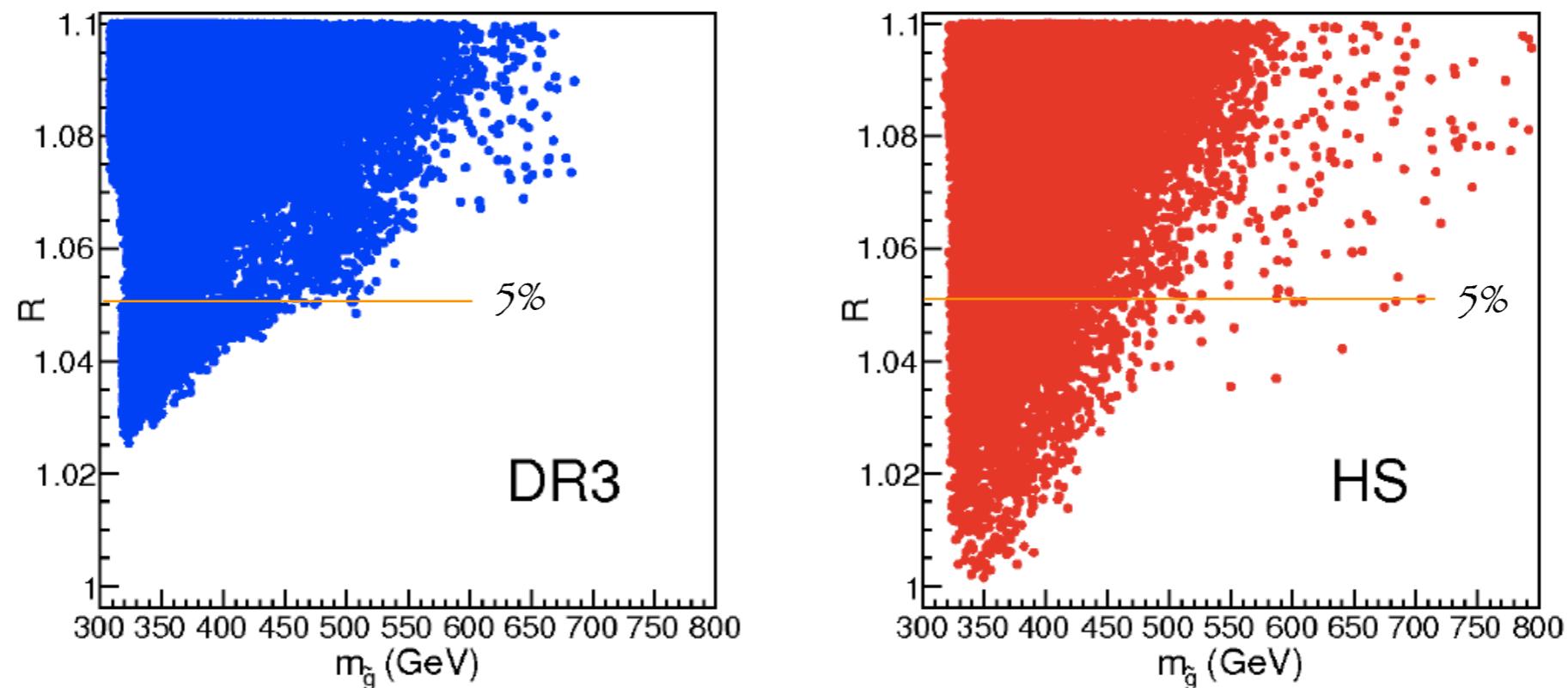
$$R = \frac{\max(f_t, f_b, f_\tau)}{\min(f_t, f_b, f_\tau)}$$

- D-term splitting w/o RHN gives $R \sim 1.08$ (i.e. 8% unification)
- Splitting of only m_H 's (“just-so HS”) allows for $R \sim 1.01$
- D-term splitting with RHN gives $R \sim 1.04, \dots$
- ... but if we allow in addition small non-degeneracy of 3rd vs. 1st/2nd generation, we get $R \sim 1.02$

Typical mass spectra

- 1st/2nd generation scalars in the multi-TeV range (5-15 TeV)
- 3rd gen. scalars, heavy Higgses and higgsinos in the 1-3 TeV range
- light gauginos: LSP \sim 50-80 GeV, gluino \sim 300-500 GeV
- c.f “effective SUSY” by Cohen, Kaplan, Nelson ’1996

$$R = \frac{\max(f_t, f_b, f_\tau)}{\min(f_t, f_b, f_\tau)}$$



Points from a MCMC scan for small R