Recent Developments in Relativistic Embeddings of Modified Newtonian Dynamics

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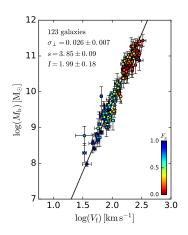


Figure: Power-law relation between v_{∞} and M_b . Slope consistent with 4. (Lelli et al., 2019)

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- Squaring again find that $v_{\infty}^4 = (a_0 G) M_b$, the baryonic Tully-Fisher relation, with the constant now identified!



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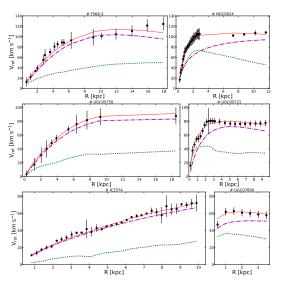
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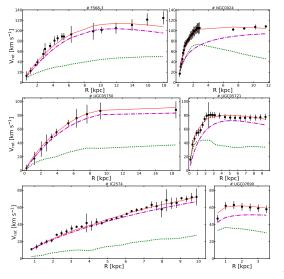
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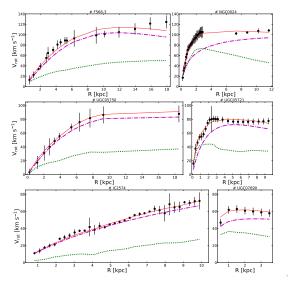
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- Another approach to getting the equations is to have a hierarchy such that Φ_N satisfies the Poisson equation and enforce an algebraic relation between $\nabla\Phi$ and $\nabla\Phi_N$ (QUMOND).



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- Rotation curve tracks trend of baryonic contribution.

Additional galactic regularity: the Radial Acceleration Relation (RAR)

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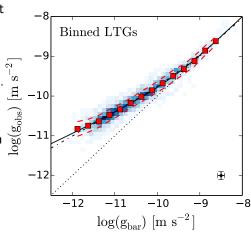


Figure: The baryons alone predict the dynamics (153 LTG) (Lelli et al., 2019).

Recent news.

Additional regularity: the radial acceleration relation (RAR) extended

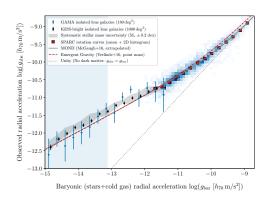


Figure: RAR extended by weak lensing agrees with deep MOND behaviour (slope 1/2) (Brouwer et al., 2019).

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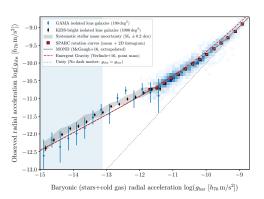


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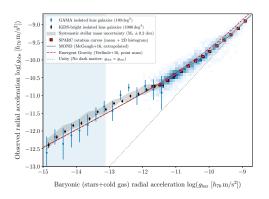


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- Signal of $\sim 10^5$ lenses (KiDS and GAMA) of isolated late-type and early-type galaxies, stacked.
- Consistent with MOND behaviour persisting!

Not-so-old news. To larger scales: galaxy clusters

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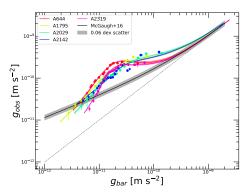


Figure: The RAR of galaxy clusters (tSZ and X-ray obs.) departs from the galactic RAR (Eckert et al., 2022).

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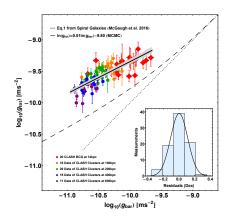


Figure: RAR of galaxy clusters inferred from lensing also departs from galactic RAR. (Tian et al., 2020)

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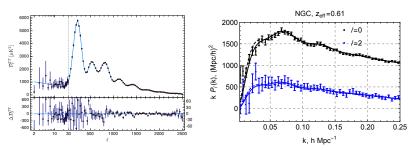


Figure: Left: Planck angular power spectrum. (Aghanim et al., 2018). Right: Matter power spectrum from BOSS (SDSS), monopole and quadrupole. (Ivanov et al., 2021)

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- A natural starting point, to not spoil all the successes of general relativity, is to have a metric theory, involving $g_{\mu\nu}$.
- Just as general relativity reduces to Newtonian gravity in the weak-field, slow-motion $(v \ll c)$ regime, so we need to find a theory whose weak-field, slow motion and low acceleration regime $a \ll a_0$ is governed by MOND.

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- In general relativity, the potential Φ is only the diagonal part of the larger metric tensor $g_{\mu\nu}$: $g_{00} = -1 + 2\Phi$, $g_{ii} = 1 + 2\Phi$.

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- Function \mathcal{F} of $Q = A^{\mu} \nabla_{\mu} \phi$ (time, cosmology) and $\mathcal{Y} = \nabla_{\mu} \phi \nabla_{\nu} \phi (g^{\mu\nu} + A^{\mu} A^{\nu})$ (space) is undetermined.
- Different choices lead to different behaviour.
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AeST vacuum solutions

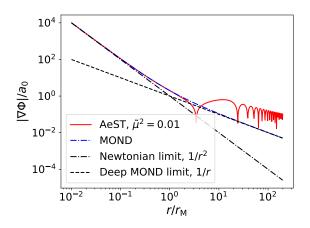


Figure: The magnitude of the force away from a point source in AeST (red). Newtonian regime $1/r^2$, MOND regime 1/r and oscillatory regime (with power-law enveloped). Distance in units of $r_{\rm M} = \sqrt{GM/a_0}$.

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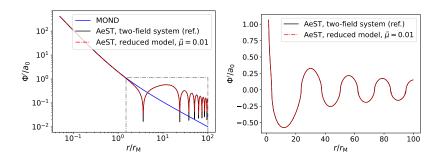


Figure: Oscillations in force. Left: Log-log plot of abs. value of force (AeST red, MOND blue). Right: Linear plot of force.

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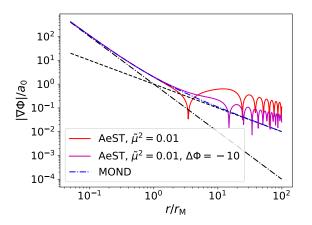


Figure: Onset of oscillatory regime depends on the boundary value of potential. Can be delayed by a lowered potential (magenta line).

Static spherically symmetric weak-field solutions of AeST: isothermal case

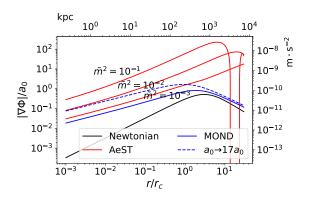


Figure: PRELIMINARY RESULT. Isothermal gas in hydrostatic equilibrium: $\rho = \exp(-\Phi)$. The force in AeST (red line) appears to be stronger than MOND (blue line) and also MOND with $a_0 \rightarrow 17a_0$ (dashed blue line).

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- Though further quantitative investigation is needed there are new weak-field effects in the theory that enhance the force compared to MOND.
- AeST has a rich phenomenology with other elements (non-static effects and vector field) which may address other issues, such as the apparent dark matter distribution around colliding clusters.

Role of vector fields

 Away from spherical symmetry the vector field cannot a priori be neglected. The weak field AeST system has been derived and reads

$$\begin{split} \nabla^{2}\tilde{\Phi} + \tilde{\mathcal{K}}_{2}\mathcal{Q}_{0}^{2}\left(\tilde{\Phi} + \chi\right) &= 4\pi G \rho_{b} \\ \nabla^{2}\tilde{\Phi} &= \nabla \cdot \left(\mathcal{J}'\left(\nabla \chi + \mathcal{Q}_{0}\beta\right)\right) \\ - \frac{\mathcal{K}_{B}}{2\left(2 - \mathcal{K}_{B}\right)}\nabla^{2}\beta + \mathcal{Q}_{0}^{2}\beta + \mathcal{J}'\left(\nabla \chi + \mathcal{Q}_{0}\beta\right) + \nabla \Phi &= 0 \end{split}$$

where β is a divergenceless vector.

- Whether the new terms can be important for colliding galaxy clusters is currently under investigation.
- It appears that they have minor effect on the rotation curve of a MW-like galaxy (Mistele).

Addressing the displaced halo in galaxy clusters

- By construction the static equations have $\dot{\Phi}, \dot{\chi} = 0$.
- When expanding around Minkowski space it turns out that scalar, vector and tensor perturbations travel at different speeds.
- Could the scalar field lag behind baryons set in motion, only to join again when caught up? That would give a displacement of the baryonic distribution from the scalar field and potentially a displaced lensing signal.
- Another possibility: Though the galaxy clusters themselves are spherically symmetric the configuration that has both of them in collision is not, so there may be effects from the vector field.