

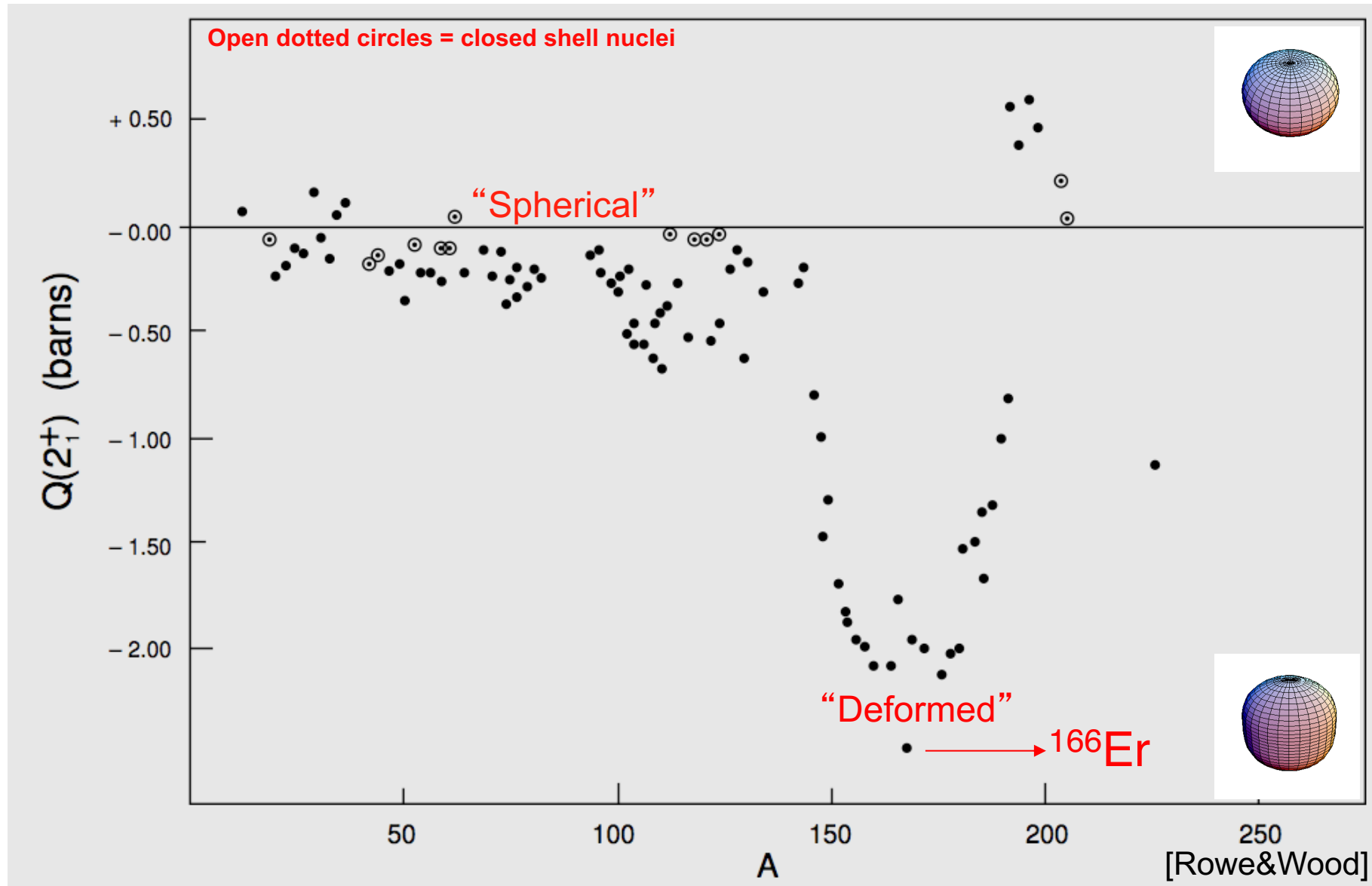
Updated View on Empirical Moments of Inertia of Axially Asymmetric Nuclei

J.M. Allmond – ORNL

ORNL is managed by UT-Battelle, LLC for the US Department of Energy

First 2^+ Quadrupole Moment

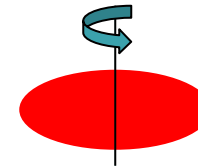
Nuclei are predominantly between spherical and prolate deformed



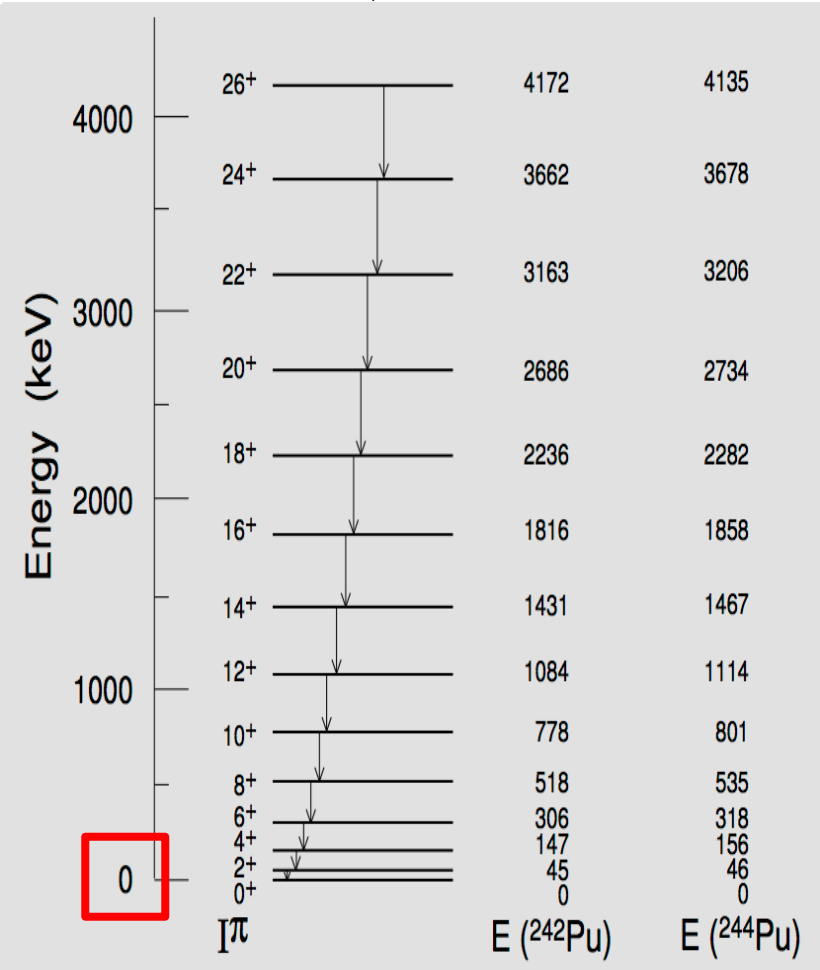
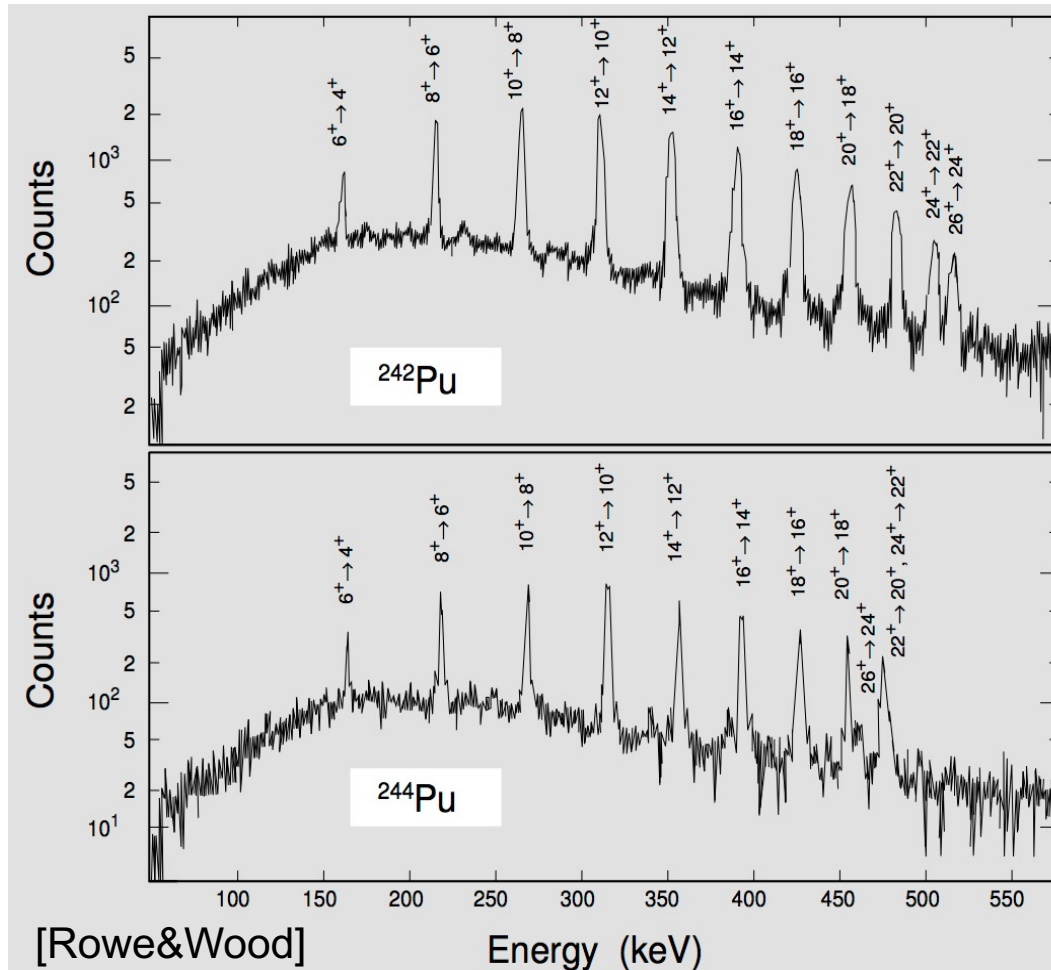
Rotor Example

Sequence of γ -rays following Coulomb excitation of Pu isotopes reveals rotor-like pattern

condition= γ -ray singles



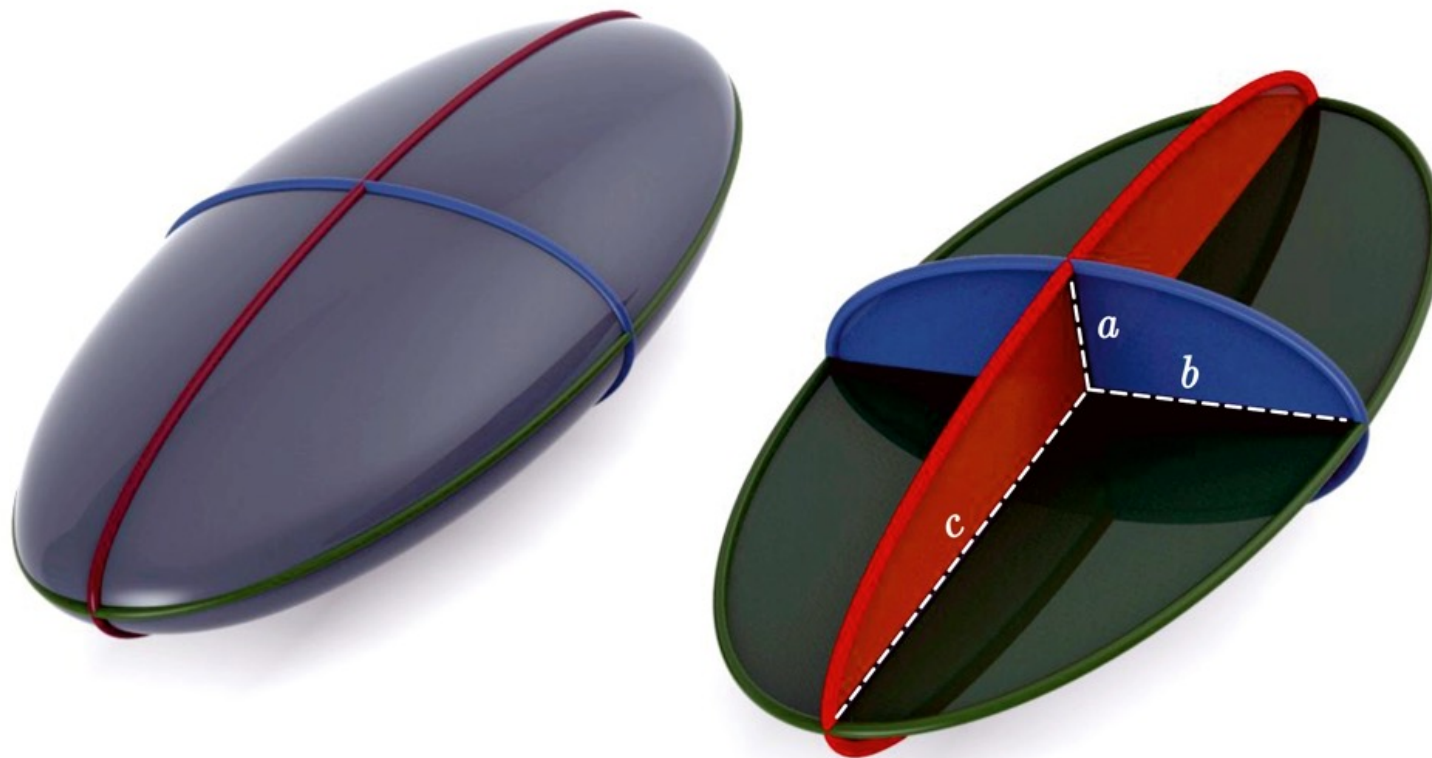
$$E = A I(I+1)$$



[Rowe&Wood]

Triaxial (ellipsoid) Shapes?

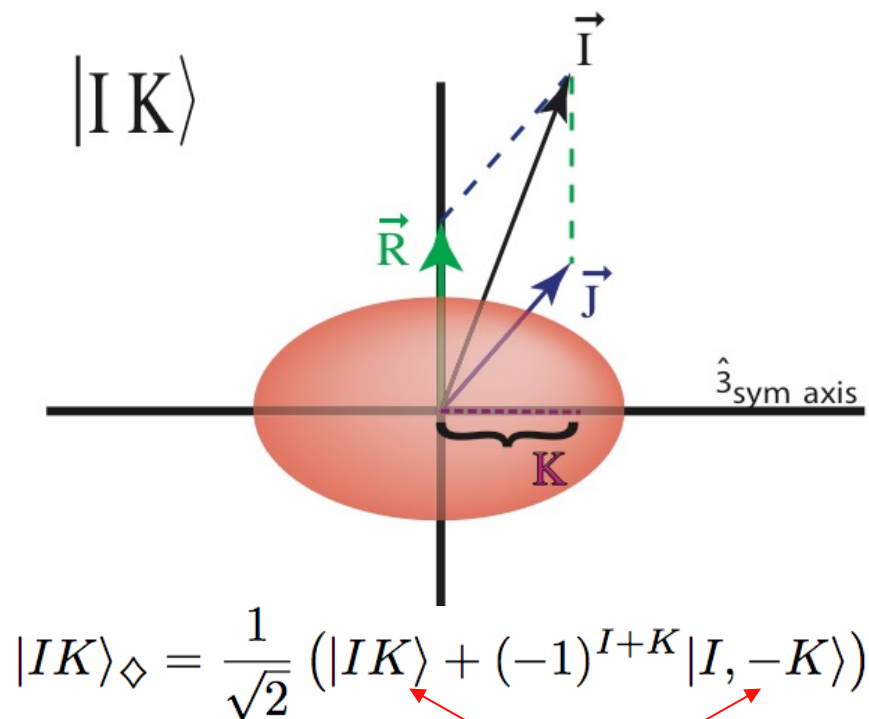
Three unique axes



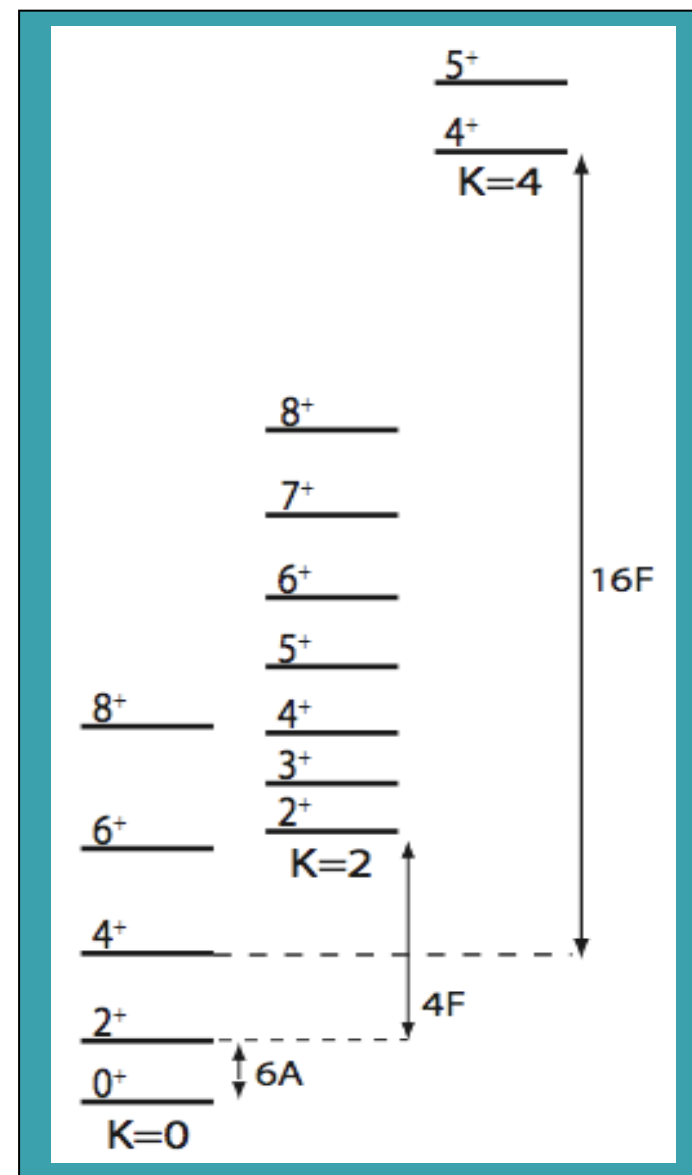
$$c > b > a$$

Rotor Basis

Body-frame projection (K) is used instead of lab-frame projection (M)

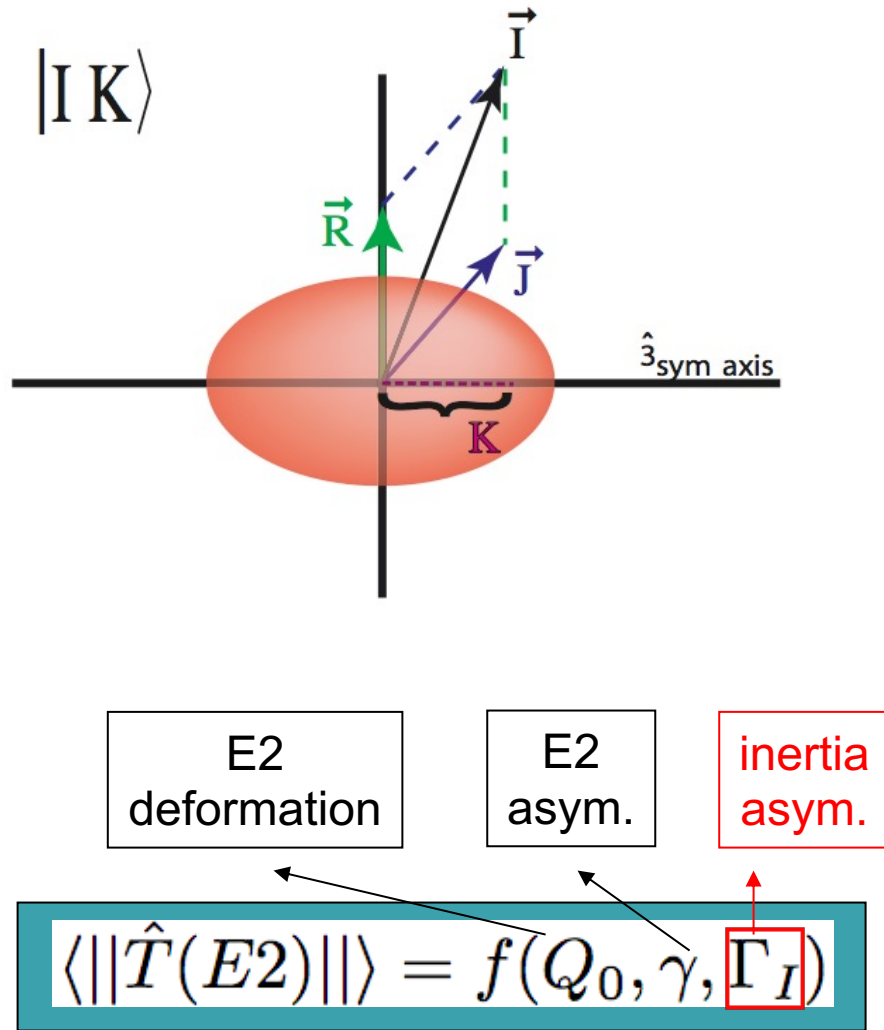


*Because nucleus has a plane of reflection symmetry



$$\diamond \langle IK | \hat{H} | IK \rangle_{\diamond} = AI(I+1) + FK^2$$

Generalized Triaxial Rotor Model



Γ Consequence of making no assumption on rigid or irrotational flow MOI in Hamiltonian. Induces $\Delta K=2$ mixing from differences in 1- and 2-axis MOI.

Hamiltonian

$$H = A_1 \hat{I}_1^2 + A_2 \hat{I}_2^2 + A_3 \hat{I}_3^2$$

$$= A \hat{I}^2 + F \hat{I}_3^2 + G(\hat{I}_+^2 + \hat{I}_-^2),$$

where

$$A = \frac{1}{2}(A_1 + A_2), \quad F = A_3 - A, \quad G = \frac{1}{4}(A_1 - A_2),$$

$$H(2^+) = \begin{pmatrix} 6A & 4\sqrt{3}G \\ 4\sqrt{3}G & 6A + 4F \end{pmatrix}$$

$$|I_g\rangle = \cos \Gamma_I |I0\rangle_{\diamond} - \sin \Gamma_I |I2\rangle_{\diamond}$$

$$|I_{\gamma}\rangle = \sin \Gamma_I |I0\rangle_{\diamond} + \cos \Gamma_I |I2\rangle_{\diamond}$$

$$\tan(2\Gamma_2) = 2\sqrt{3} \frac{G}{F}$$

E2 Operator

$$\hat{T}(E2) = \cos \gamma \hat{T}_0^{(2)} + \frac{\sin \gamma}{\sqrt{2}} (\hat{T}_{+2}^{(2)} + \hat{T}_{-2}^{(2)})$$

Rigid and Irrotational Couplings of Inertia and E2 Tensors

Generic model with use of Γ can recapture the traditional assumptions, e.g., of Davydov-Filippov Irrotational Model

$$\mathcal{J}_{rigid, k} = B_{rigid} \left[1 - \sqrt{\frac{5}{4\pi}} \beta \cos \left(\gamma - k \frac{2\pi}{3} \right) \right]$$

and

$$\mathcal{J}_{irrot., k} = 4B_{irrot.} \beta^2 \sin^2 \left(\gamma - k \frac{2\pi}{3} \right),$$

Becomes equivalent to Davydov-Filippov model

$$\boxed{\Gamma} = \frac{1}{2} \tan^{-1} \left(\sqrt{3} \frac{\mathcal{J}_2 - \mathcal{J}_1}{\frac{2\mathcal{J}_1\mathcal{J}_2}{\mathcal{J}_3} - \mathcal{J}_2 - \mathcal{J}_1} \right) \xrightarrow{\text{Apply Irrotational MOI Assumption}} \Gamma_{irrot.} = -\frac{1}{2} \cos^{-1} \left(\frac{\cos 4\gamma + 2 \cos 2\gamma}{\sqrt{9 - 8 \sin^2 3\gamma}} \right)$$

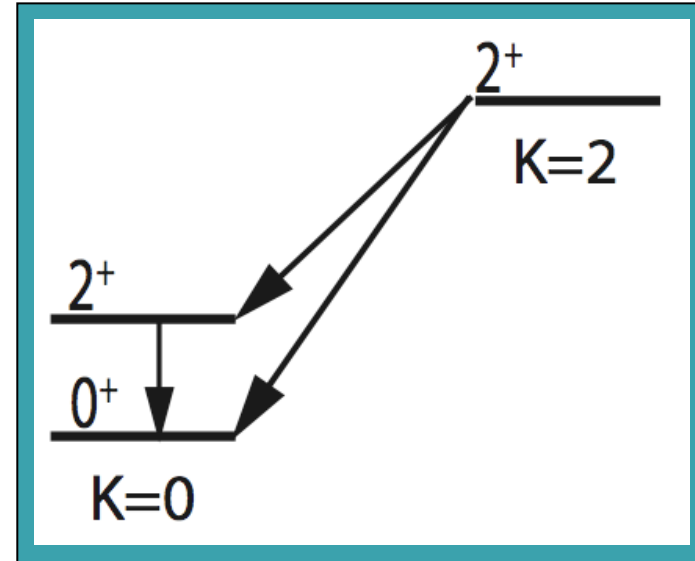
$\boxed{\Gamma}$ Consequence of making no assumption on rigid or irrotational flow MOI in Hamiltonian. Induces $\Delta K=2$ mixing from differences in 1- and 2-axis MOI.

$$\begin{aligned} \gamma = 0 &\rightarrow \Gamma = 0 \\ \gamma = 30 &\rightarrow \Gamma = -30 \\ \gamma = 60 &\rightarrow \Gamma = -60 \end{aligned} \quad \text{Independent of } \beta$$

Triaxial I=0,2 Model Space and Shape Parameters

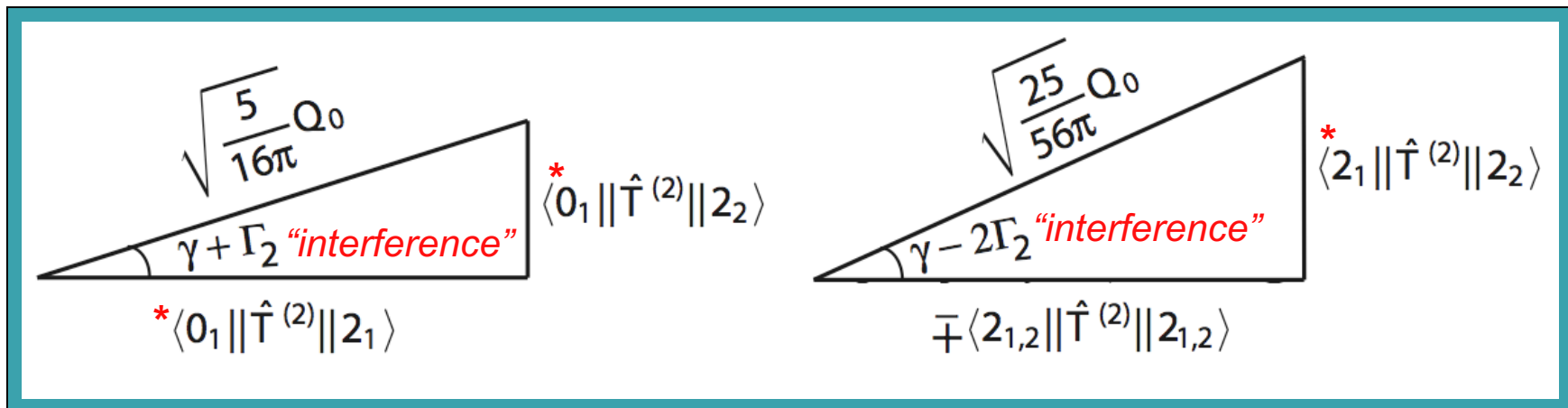
*Only need 3 of 5 M.E.s to determine Q_0, γ, Γ ; Typically use the 3 transition M.E.s due to precision

$$\begin{aligned}
 * \langle 0_1 || \hat{T}^{(2)} || 2_1 \rangle &= \sqrt{\frac{5}{16\pi}} Q_0 \cos(\gamma + \Gamma_2) \\
 * \langle 0_1 || \hat{T}^{(2)} || 2_2 \rangle &= \sqrt{\frac{5}{16\pi}} Q_0 \sin(\gamma + \Gamma_2) \\
 * \langle 2_1 || \hat{T}^{(2)} || 2_2 \rangle &= \sqrt{\frac{25}{56\pi}} Q_0 \sin(\gamma - 2\Gamma_2) \\
 \langle 2_1 || \hat{T}^{(2)} || 2_1 \rangle &= -\sqrt{\frac{25}{56\pi}} Q_0 \cos(\gamma - 2\Gamma_2) \\
 &= -\langle 2_2 || \hat{T}^{(2)} || 2_2 \rangle
 \end{aligned}$$



2^+ mixing,
 Γ_2 (deg)

*generated
from difference
in 1- and 2-axis
MOI



Parametrization reveals a destructive interference effect between E2 and inertia asymmetry

Empirical Moments of Inertia

Summary of Procedure:

- β, γ, Γ from E2 Matrix Elements
- A, F, G from 2+ energies and Γ
- MOI from A, F, G

$$\mathcal{J}_1 = \frac{1}{2} \frac{\hbar^2}{A + 2G},$$

$$\mathcal{J}_2 = \frac{1}{2} \frac{\hbar^2}{A - 2G},$$

$$\mathcal{J}_3 = \frac{1}{2} \frac{\hbar^2}{A + F}.$$

where,

$$F = \frac{E(2_{\gamma}^+) - E(2_g^+)}{4\sqrt{1 + \tan^2(2\Gamma)}},$$

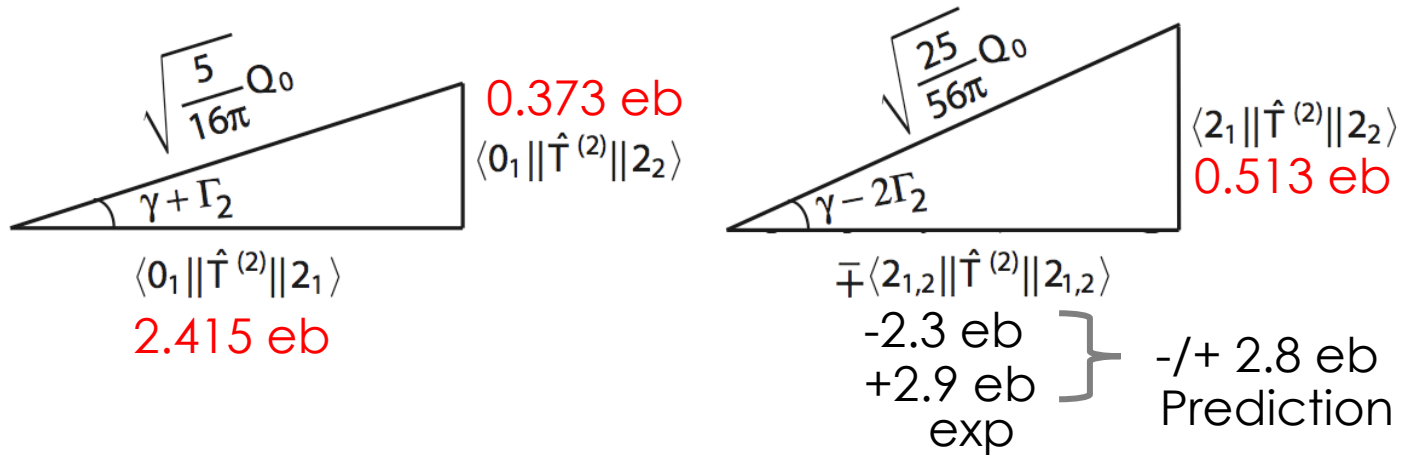
$$A = \frac{E(2_g^+) + E(2_{\gamma}^+) - 4F}{12},$$

$$G = \frac{F}{2\sqrt{3}} \tan 2\Gamma,$$

Experimental 2+ energies

Mixing angle (inertia asymmetry)
from experimental E2 Matrix
Elements (*see previous slide*)

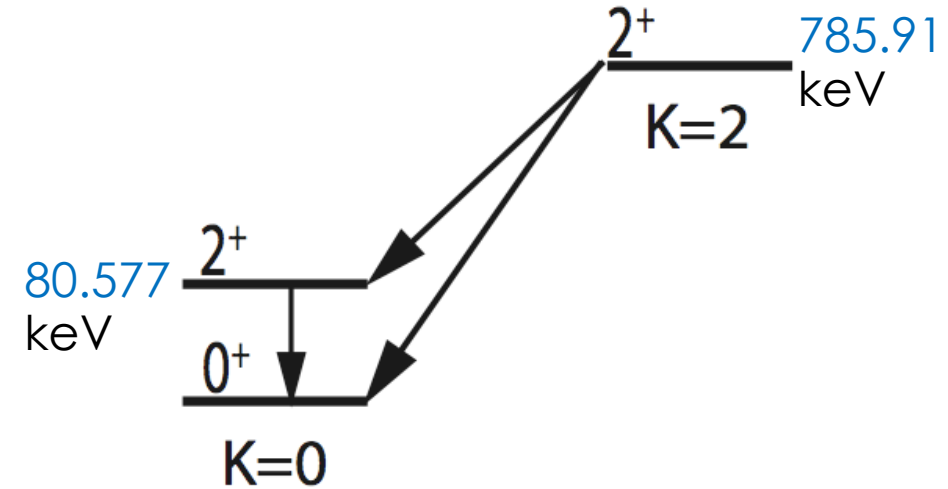
Example Fit to ^{166}Er



$$Q_0 = 7.75(3) \text{ eb}$$

$$\gamma = 9.23(17)^\circ$$

$$\Gamma_2 = -0.45(12)^\circ$$



$$F = \frac{E(2_\gamma^+) - E(2_g^+)}{4\sqrt{1 + \tan^2(2\Gamma)}}, \quad F = 176.31 \text{ keV}$$

$$A = \frac{E(2_\gamma^+) + E(2_g^+) - 4F}{12}, \quad A = 13.44 \text{ keV}$$

$$G = \frac{F}{2\sqrt{3}} \tan 2\Gamma, \quad G = -0.79 \text{ keV}$$

MOI from combining **E2** and **Energy** values

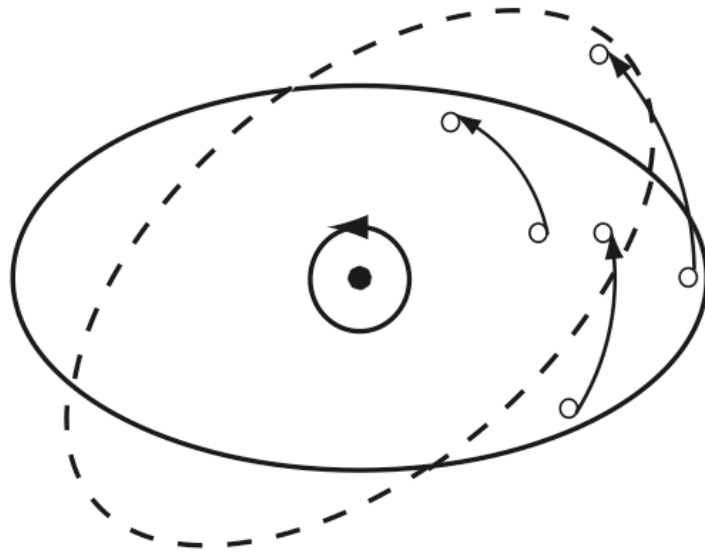
$$\mathcal{J}_1 = \frac{1}{2} \frac{\hbar^2}{A + 2G}, \quad = 42.2 \text{ h-bar}^2/\text{MeV}$$

$$\mathcal{J}_2 = \frac{1}{2} \frac{\hbar^2}{A - 2G}, \quad = 33.3 \text{ h-bar}^2/\text{MeV}$$

$$\mathcal{J}_3 = \frac{1}{2} \frac{\hbar^2}{A + F}, \quad = 2.6 \text{ h-bar}^2/\text{MeV}$$

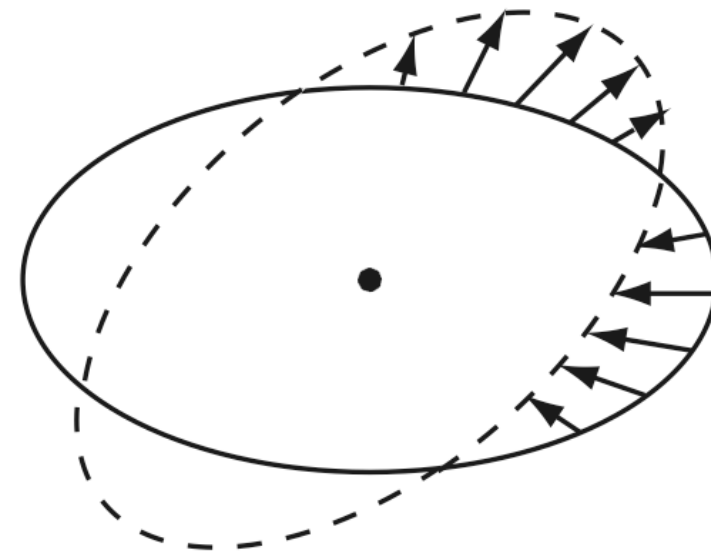
How Do Nuclei Rotate: Rigid or Irrotational Flow?

Rigid Flow



e.g., like a spinning football

Irrotational Flow



e.g., like a surface wave

View in 2017 Study

E(4)/E(2) > 2.7 cases

Physics Letters B 767 (2017) 226–231



ELSEVIER

Contents lists available at [ScienceDirect](#)

Physics Letters B

www.elsevier.com/locate/physletb



Empirical moments of inertia of axially asymmetric nuclei

J.M. Allmond^{a,*}, J.L. Wood^b

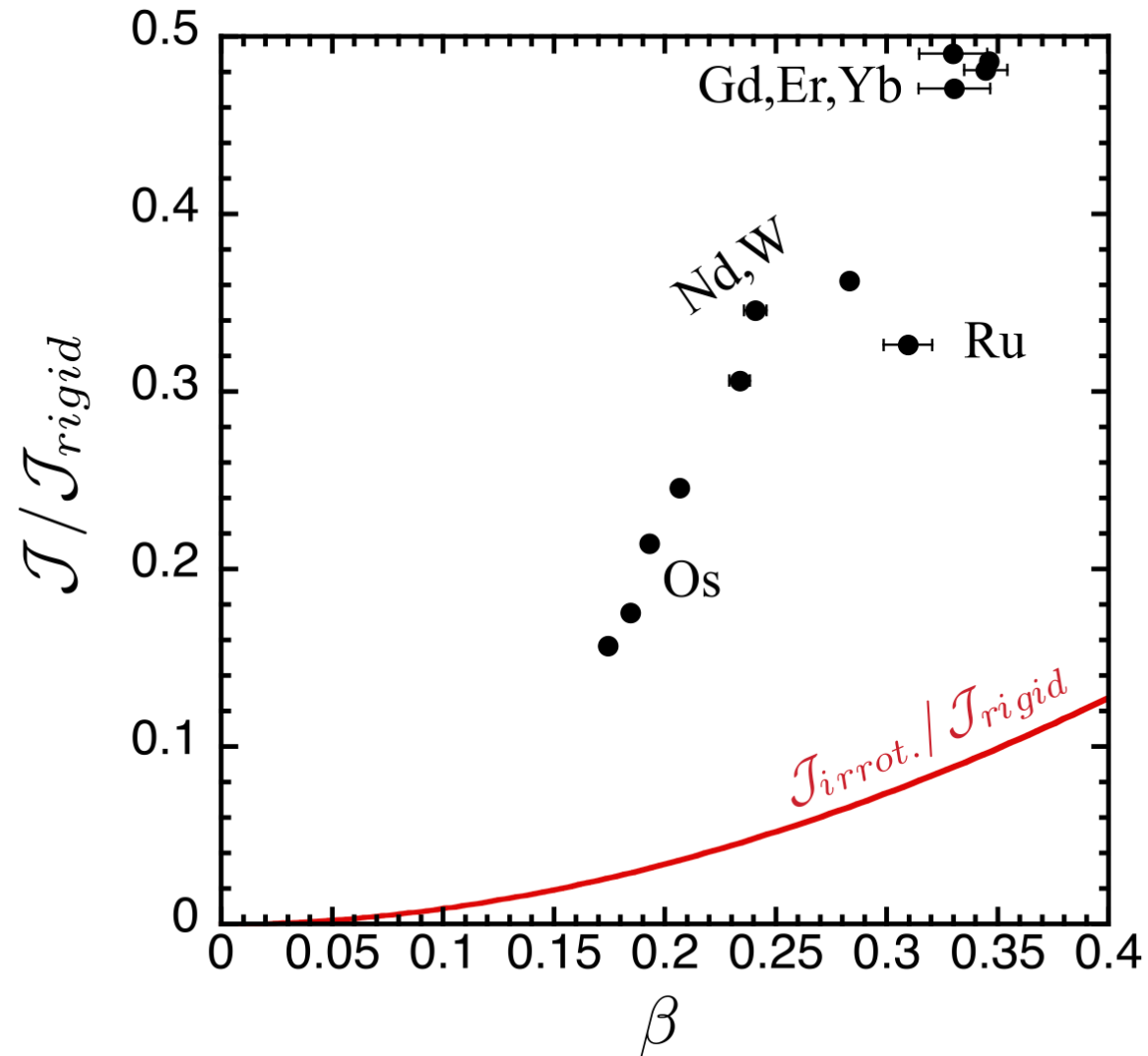
^a Physics Division, Oak Ridge National Laboratory, Oak Ridge, TN 37831, USA

^b School of Physics, Georgia Institute of Technology, Atlanta, GA 30332, USA



Empirical 1D View

$E(4)/E(2) > 2.7$

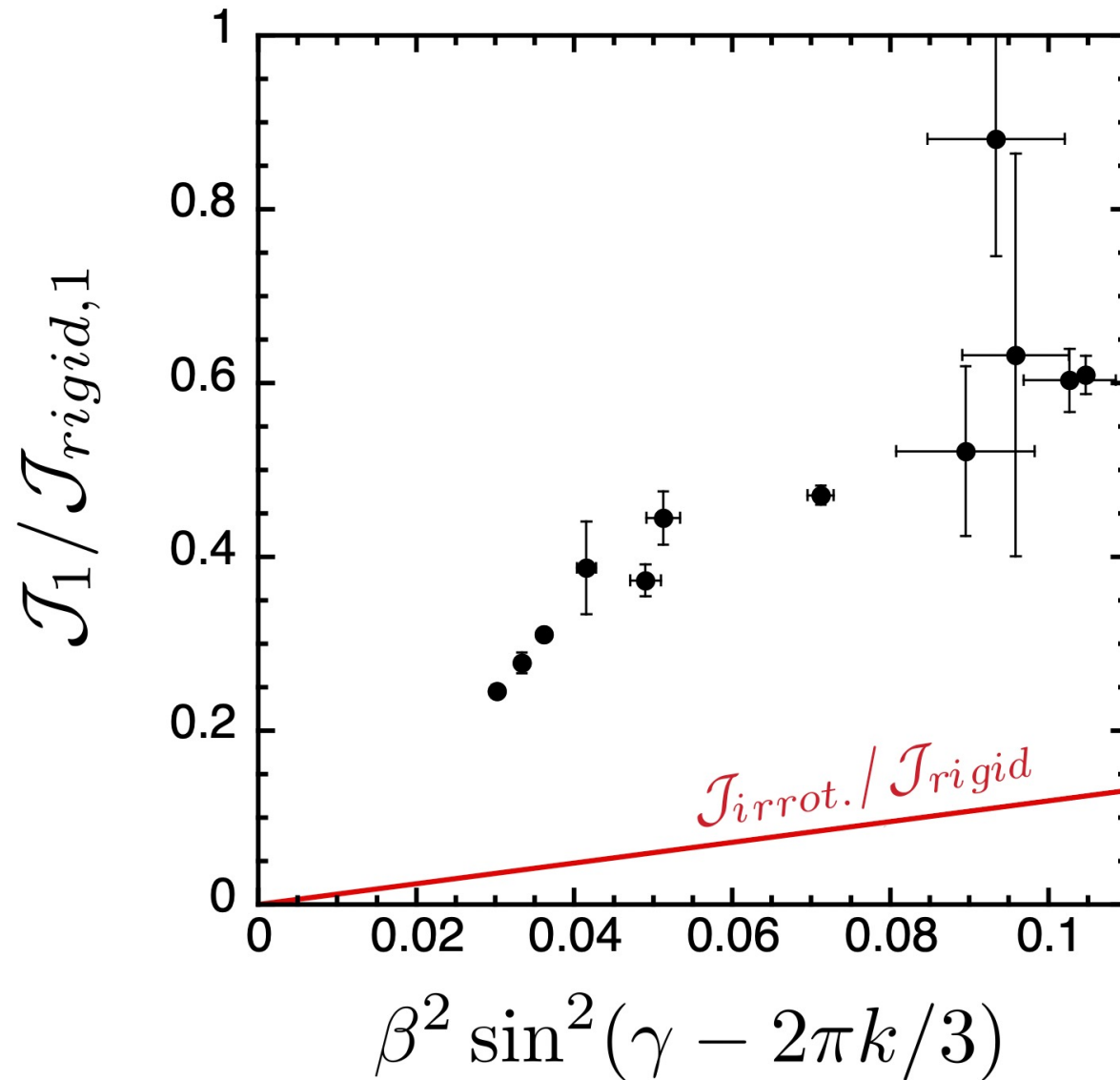


Presumes Axial Symmetry, $\gamma=0$

1-axis = 2-axis MOI = $3/E(2_1)$
3-axis = 0

Empirical 3D View: the 1-Axis

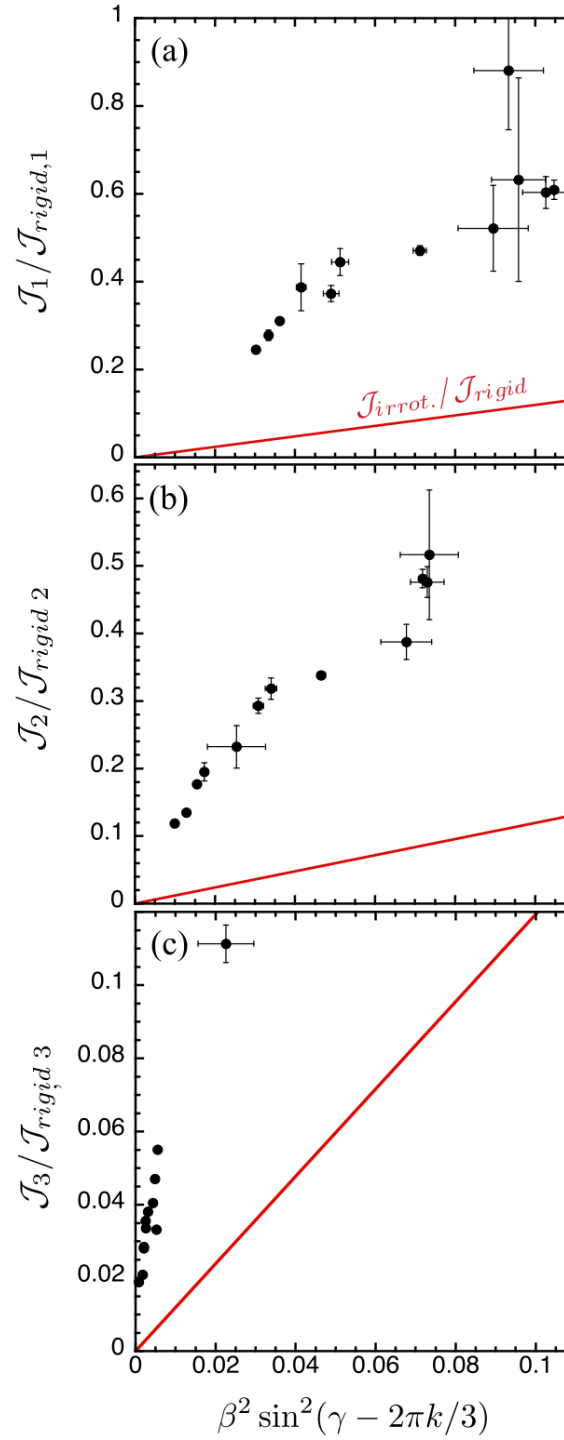
$E(4)/E(2) > 2.7$



Presumes Axial Asymmetry
1-axis \neq 2-axis \neq 3-axis

Empirical 3D View

$E(4)/E(2) > 2.7$



1-axis

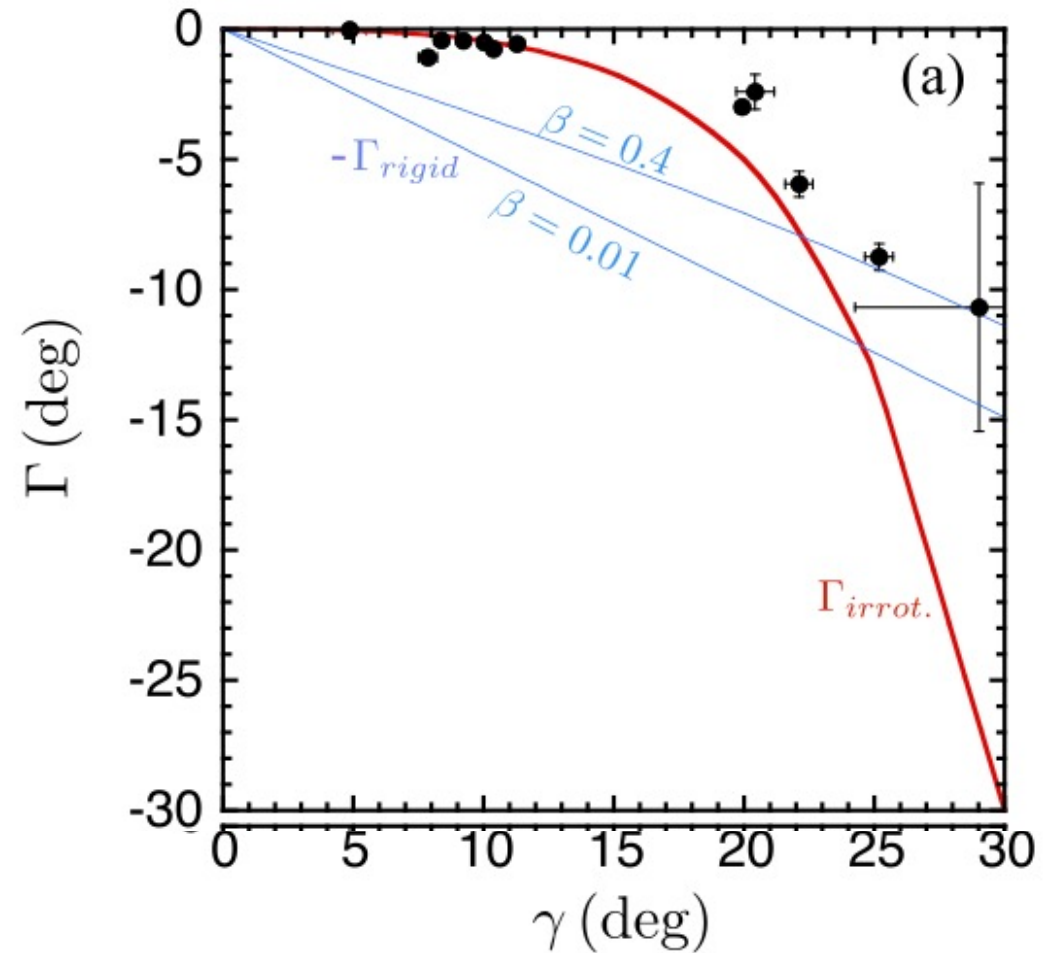
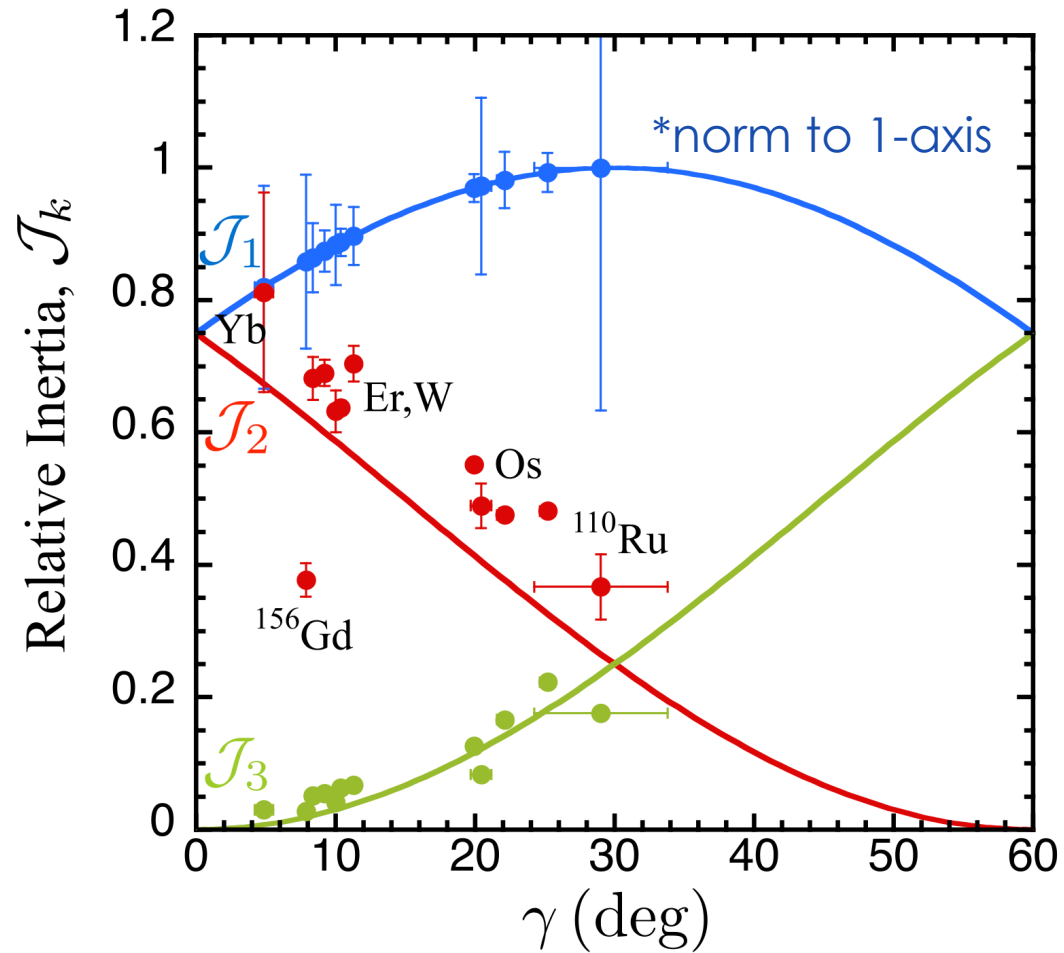
Presumes Axial Asymmetry
1-axis \neq 2-axis \neq 3-axis

2-axis

3-axis

Empirical 3D View

$E(4)/E(2) > 2.7$



Absolute values: Between Rigid and Irrotational Flow

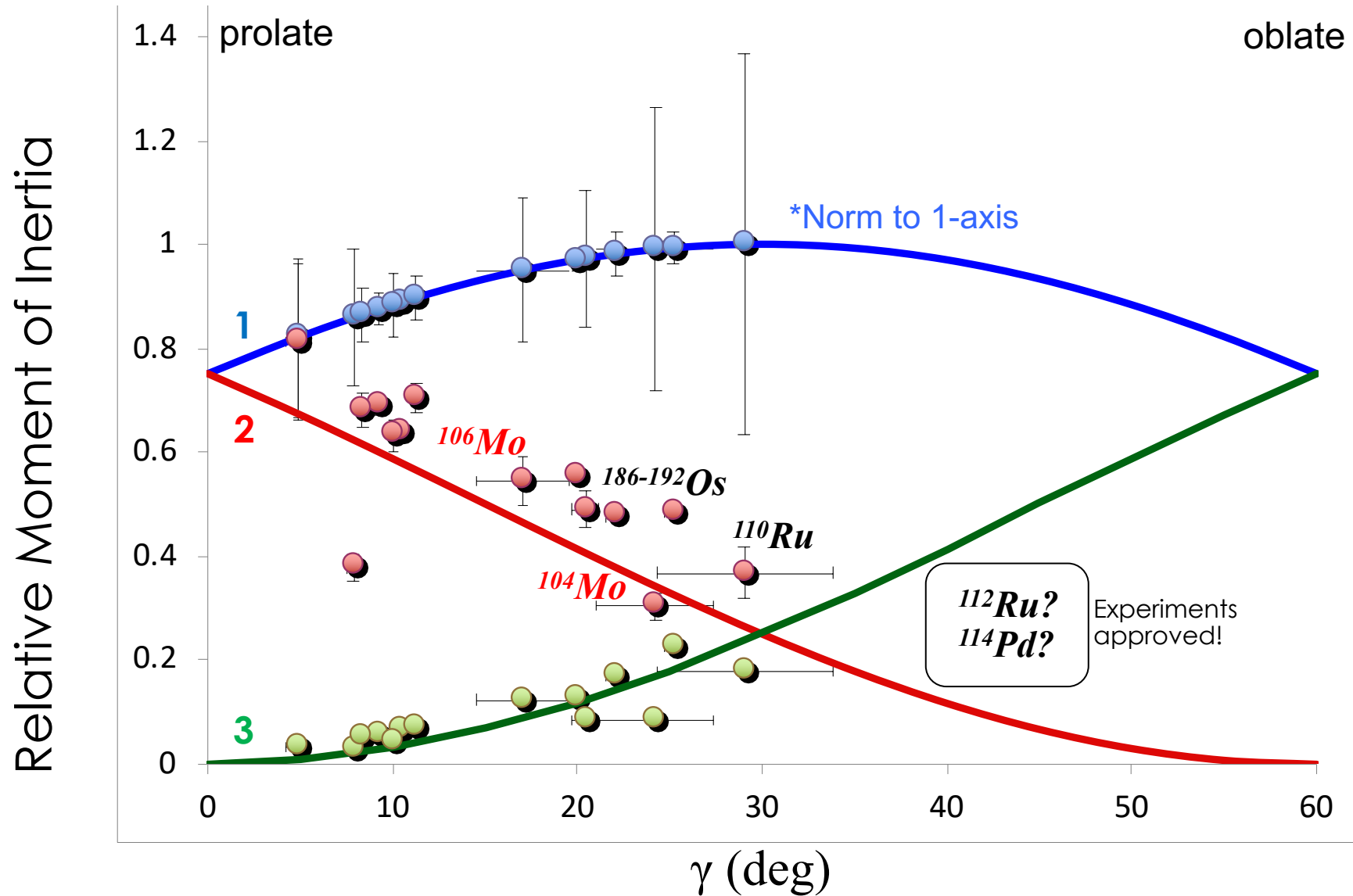
Relative values: Qualitatively consistent with SO(5) symmetry of Bohr Hamiltonian

**irrotational flow happens to be SO(5) symmetric

**E2 matrix elements only dependent on relative MOI

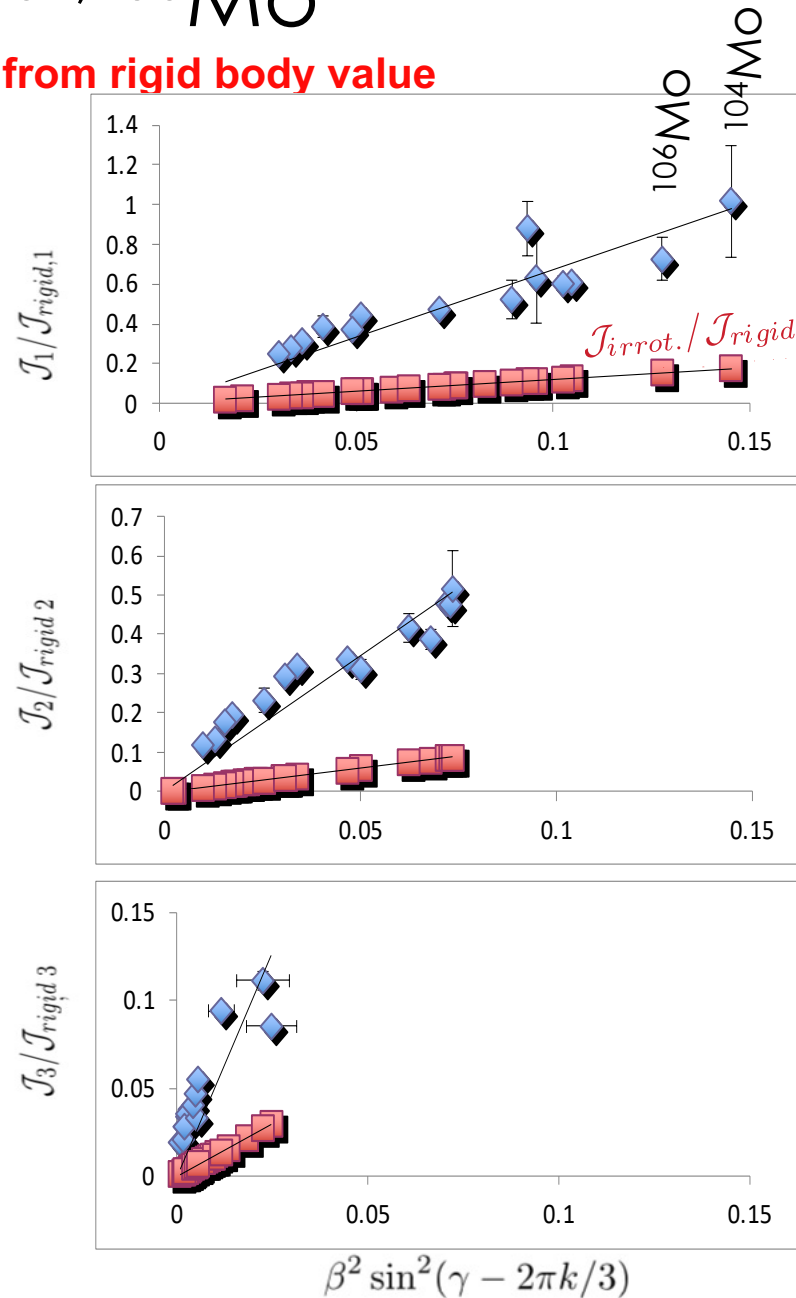
Update to Include $^{104,106}\text{Mo}$

Add $E(4)/E(2) = 2.92$ and 3.03 cases; $^{104,106}\text{Mo}$ are similar to the stable Os isotopes



Update to Include $^{104,106}\text{Mo}$

1-axis MOI of ^{104}Mo nearly unquenched from rigid body value

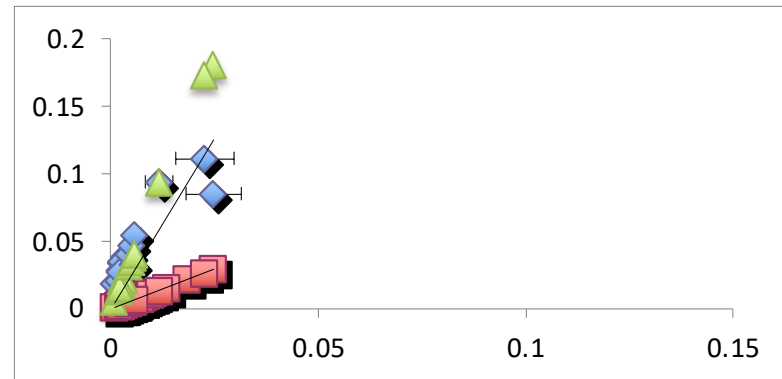
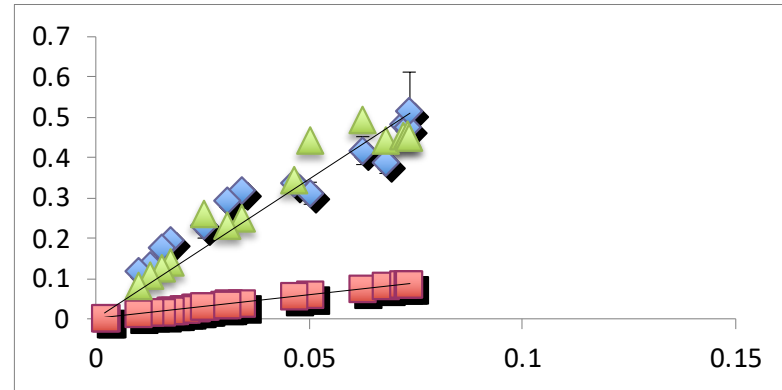
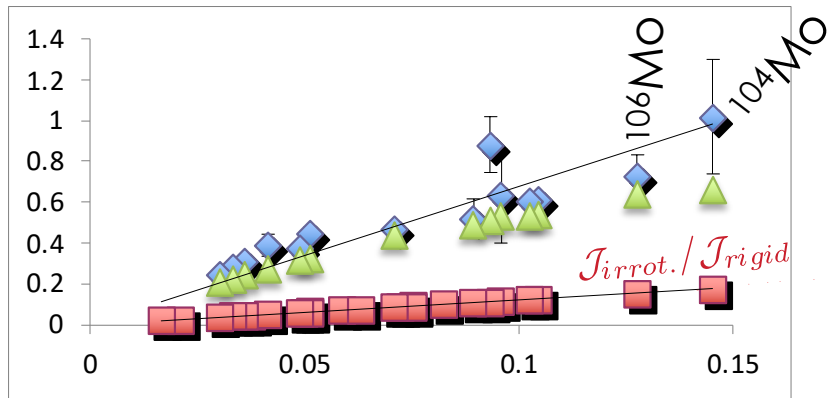


1-axis

2-axis

3-axis

3D MOI Values ~ Similar to a Trapped Cold Fermi Gas



$$\beta^2 \sin^2(\gamma - 2\pi k/3)$$

PHYSICAL REVIEW A **67**, 033611 (2003)

Slow rotation of a superfluid trapped Fermi gas

Michael Urban¹ and Peter Schuck^{1,2}

Green Points (Cold Fermi Gas MOI)

- 41-MeV “trapping potential”
- 1.5-MeV “interaction strength”
- β and γ from Coulex data

PHYSICAL REVIEW C **100**, 031301(R) (2019)

Rapid Communications

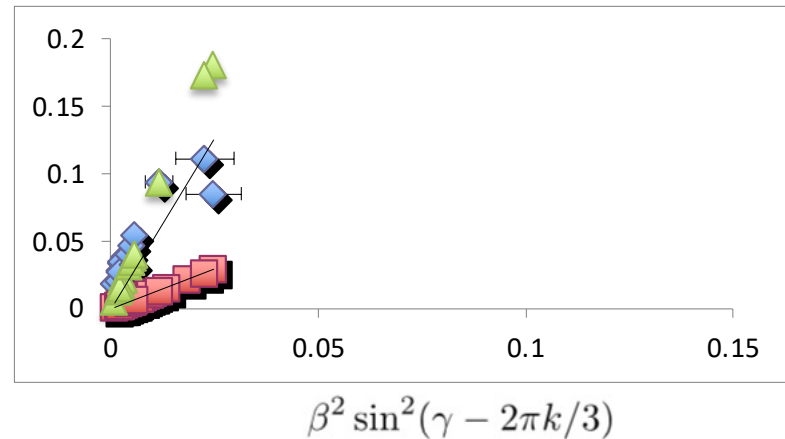
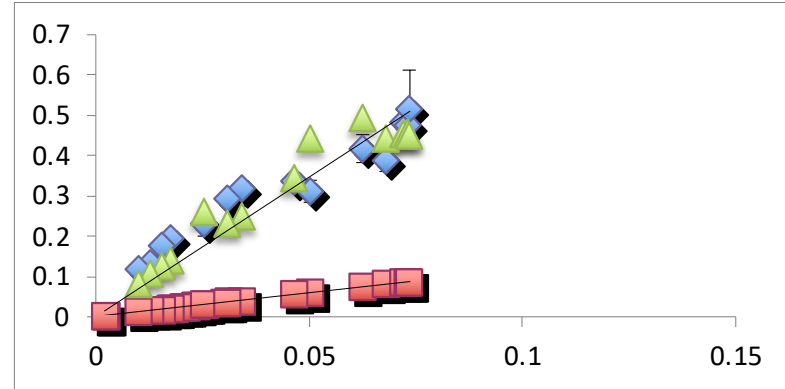
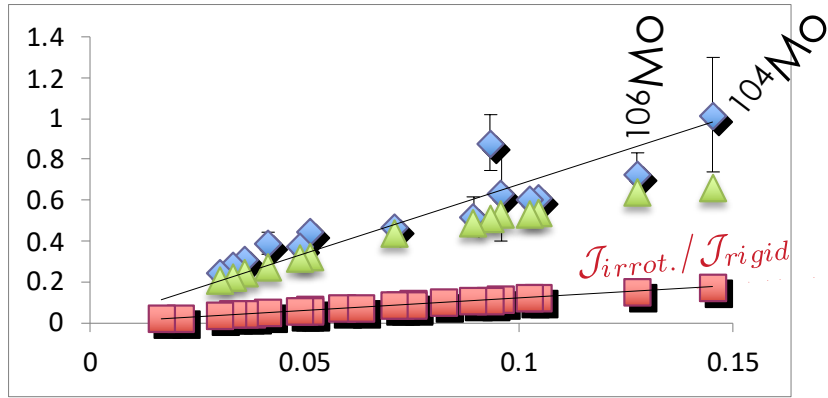
Macroscopic manifestations of rotating triaxial superfluid nuclei

P. Schuck^{1,2,*} and M. Urban^{1,†}

¹Institut de Physique Nucléaire, CNRS-IN2P3, Université Paris-Sud, Université Paris-Saclay, F-91406 Orsay Cedex, France

²Université Grenoble Alpes, CNRS, LPMMC, F-38000 Grenoble, France

3D MOI Values ~ Similar to a Trapped Cold Fermi Gas



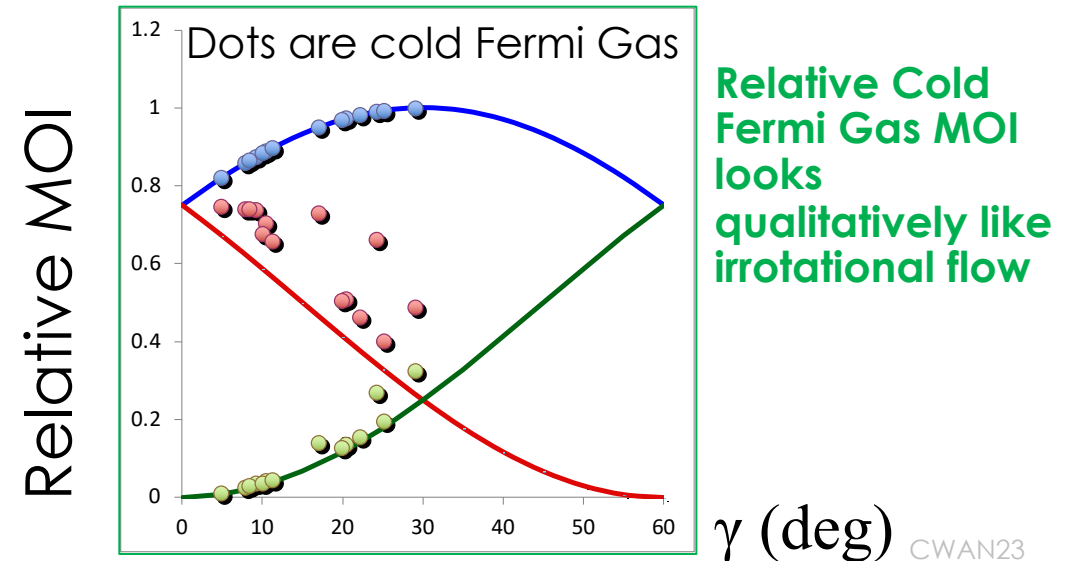
PHYSICAL REVIEW A **67**, 033611 (2003)

Slow rotation of a superfluid trapped Fermi gas

Michael Urban¹ and Peter Schuck^{1,2}

Green Points (Cold Fermi Gas MOI)

- 41-MeV “trapping potential”
- 1.5-MeV “interaction strength”
- β and γ from Coulex data



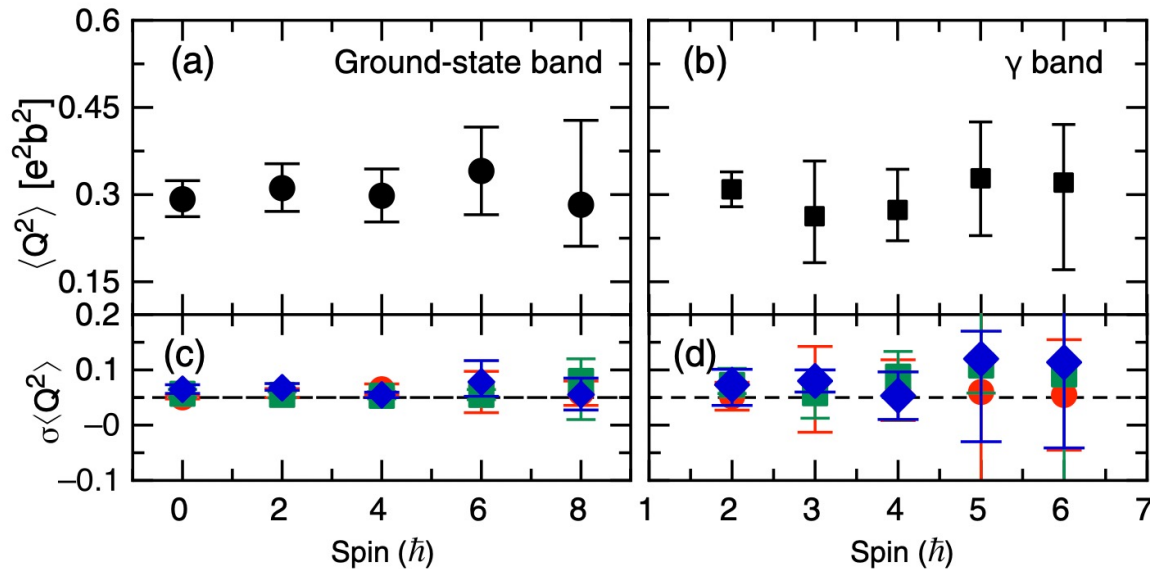
Recent Surprises: ^{76}Ge is β and γ Rigid

Energies imply shape softness, $E(4)/E(2) = 2.5$, but E2 matrix elements imply shape rigidity

PHYSICAL REVIEW LETTERS 123, 102501 (2019)

Evidence for Rigid Triaxial Deformation in ^{76}Ge from a Model-Independent Analysis

A. D. Ayangeakaa^{1,*}, R. V. F. Janssens^{2,3,†}, S. Zhu^{4,‡}, D. Little^{2,3}, J. Henderson⁵, C. Y. Wu⁵, D. J. Hartley¹, M. Albers⁴, K. Auranen⁴, B. Bucher^{5,§}, M. P. Carpenter⁴, P. Chowdhury⁶, D. Cline⁷, H. L. Crawford⁸, P. Fallon⁸, A. M. Forney⁹, A. Gade^{10,11}, A. B. Hayes⁷, F. G. Kondev⁴, Krishichayan^{3,12}, T. Lauritsen⁴, J. Li⁴, A. O. Macchiavelli⁸, D. Rhodes^{10,11}, D. Seweryniak⁴, S. M. Stolze⁴, W. B. Walters⁹ and J. Wu⁴

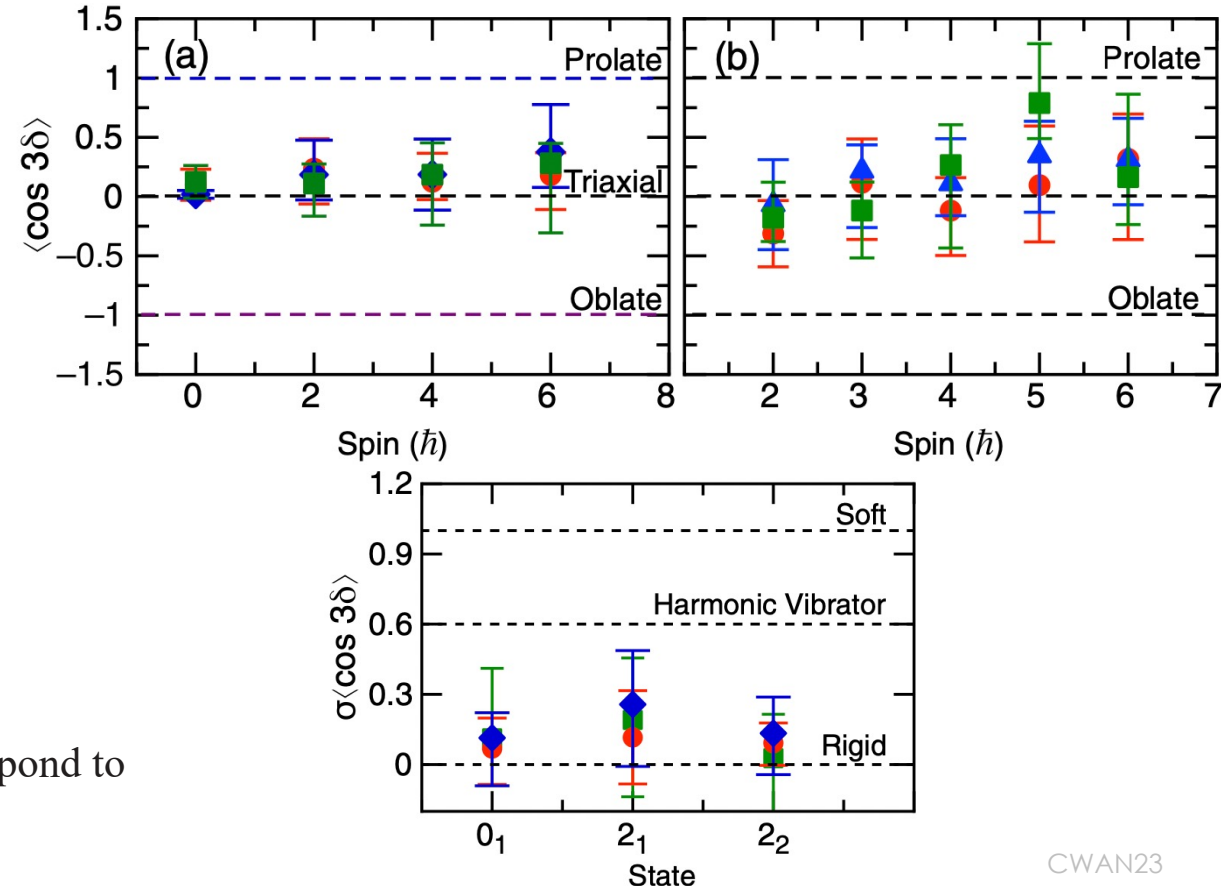


red (circle), blue (diamond), and green (square) correspond to $J=0, 2, 4$ E2 vector coupling solutions, giving three independent measures of variance

PHYSICAL REVIEW C 107, 044314 (2023)

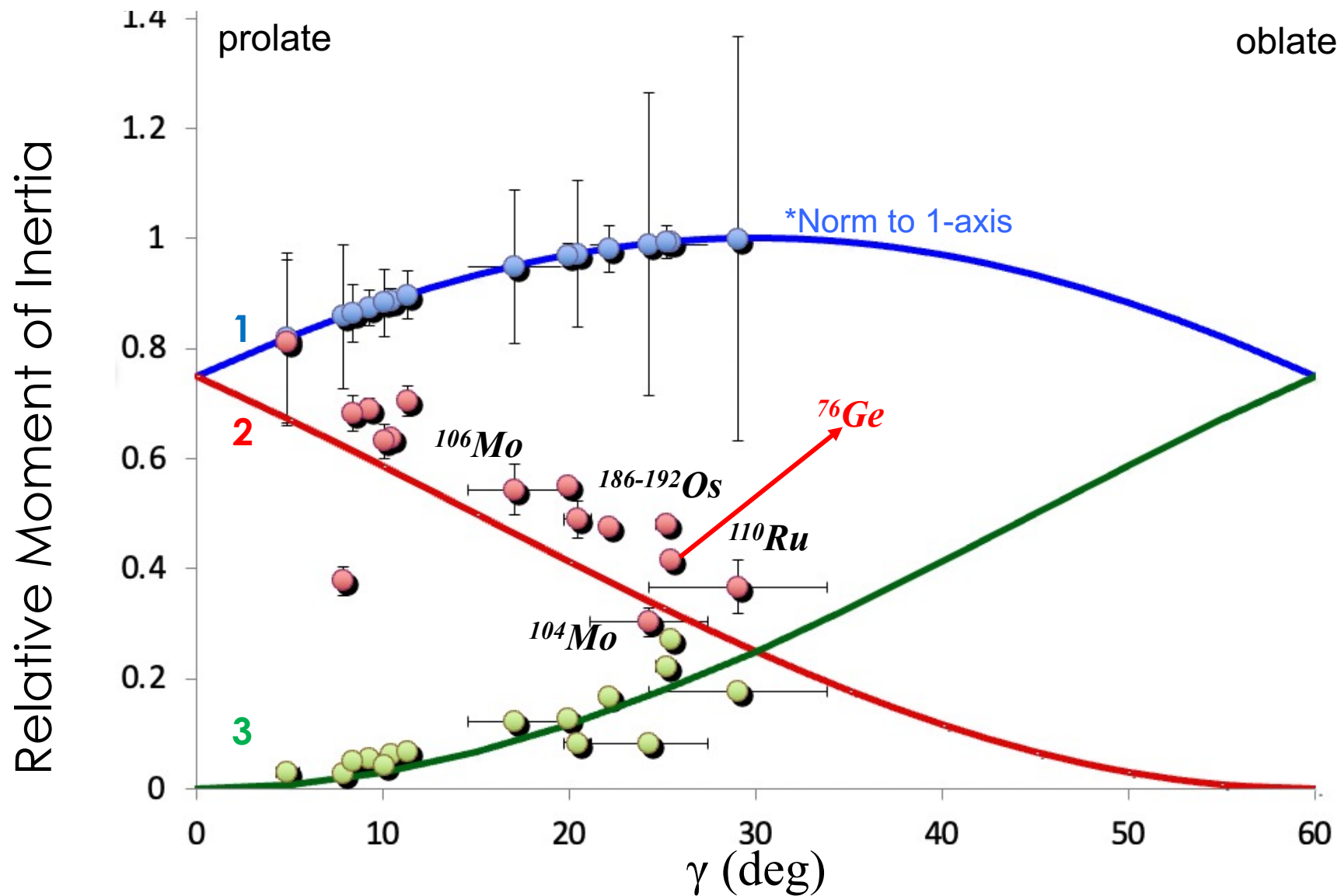
Triaxiality and the nature of low-energy excitations in ^{76}Ge

A. D. Ayangeakaa^{1,2,3,*}, R. V. F. Janssens^{1,2,†}, S. Zhu^{4,‡}, J. M. Allmond⁵, B. A. Brown^{6,7}, C. Y. Wu⁸, M. Albers^{9,§}, K. Auranen⁹, B. Bucher^{8,||}, M. P. Carpenter⁹, P. Chowdhury¹⁰, D. Cline¹¹, H. L. Crawford¹², P. Fallon¹², A. M. Forney¹³, A. Gade^{6,7}, D. J. Hartley³, A. B. Hayes¹¹, J. Henderson⁸, F. G. Kondev⁹, Krishichayan^{2,14}, T. Lauritsen⁹, J. Li⁹, D. Little^{1,2}, A. O. Macchiavelli^{12,¶}, D. Rhodes^{6,7,#}, D. Seweryniak⁹, S. M. Stolze⁹, W. B. Walters¹³ and J. Wu^{9,**}



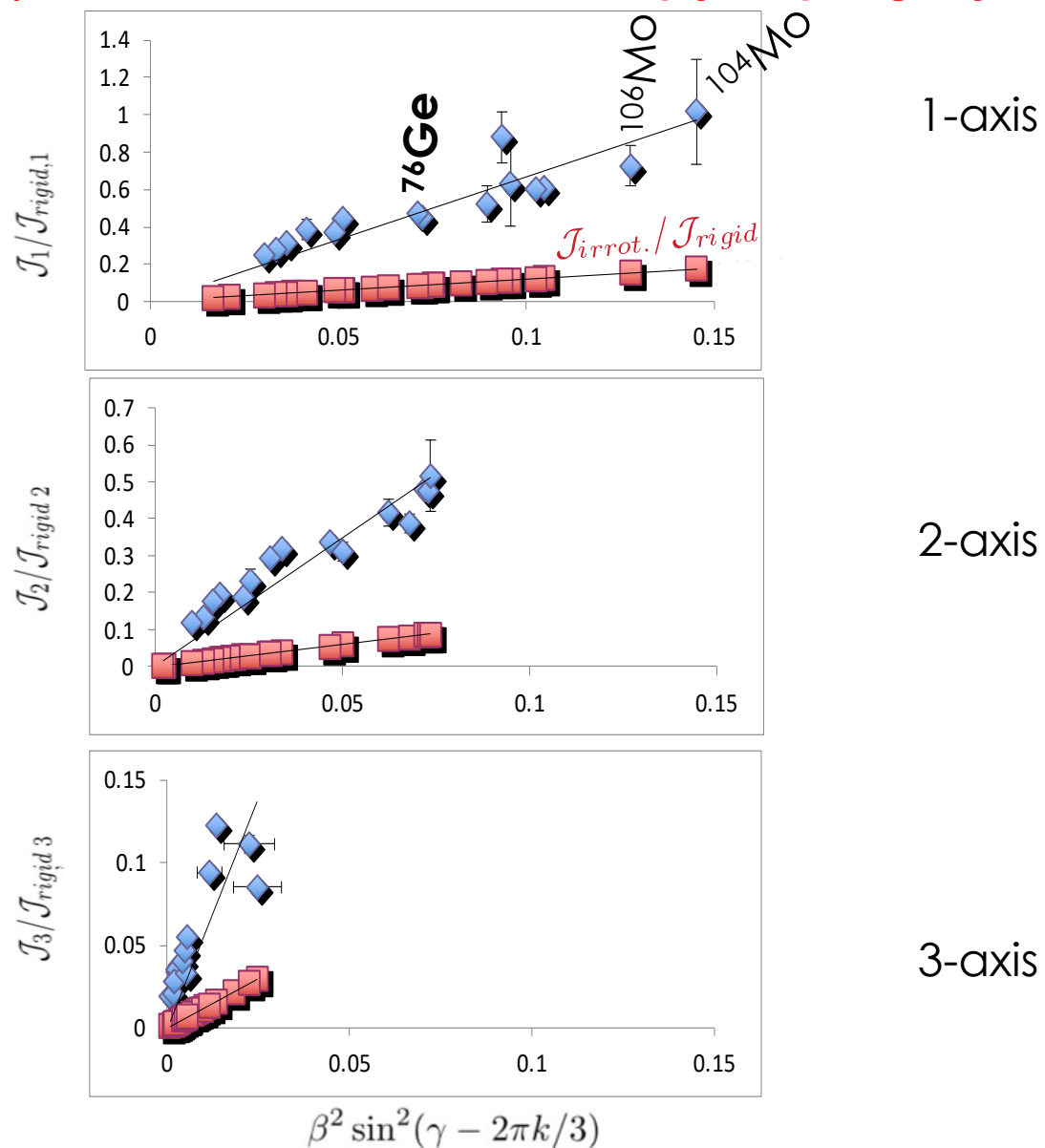
Recent Surprises: ^{76}Ge is β and γ Rigid

Energies imply shape softness, $E(4)/E(2) = 2.5$, but E2 matrix elements imply shape rigidity



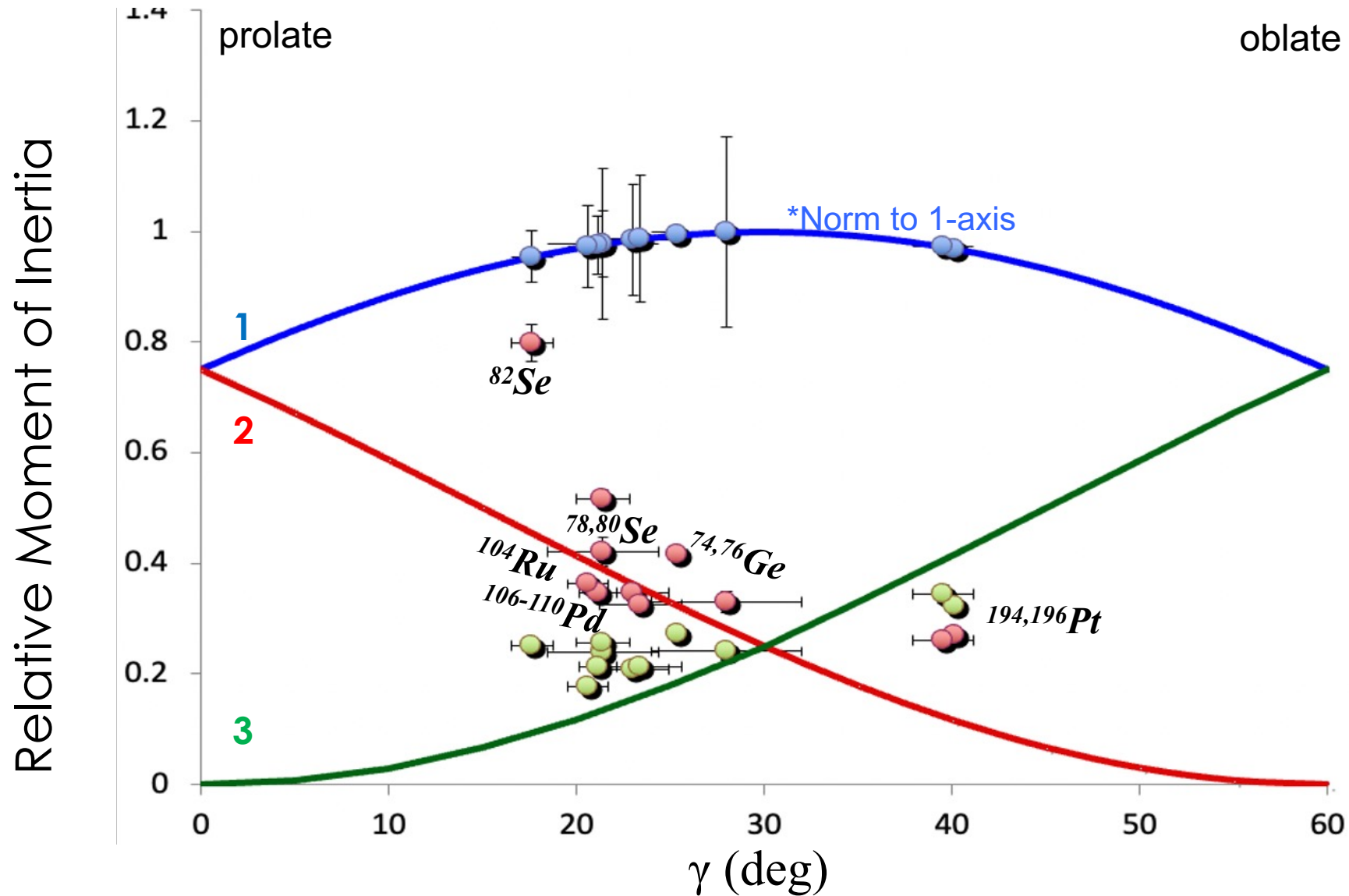
Recent Surprises: ^{76}Ge is β and γ Rigid

Energies imply shape softness, $E(4)/E(2) = 2.5$, but E2 matrix elements imply shape rigidity



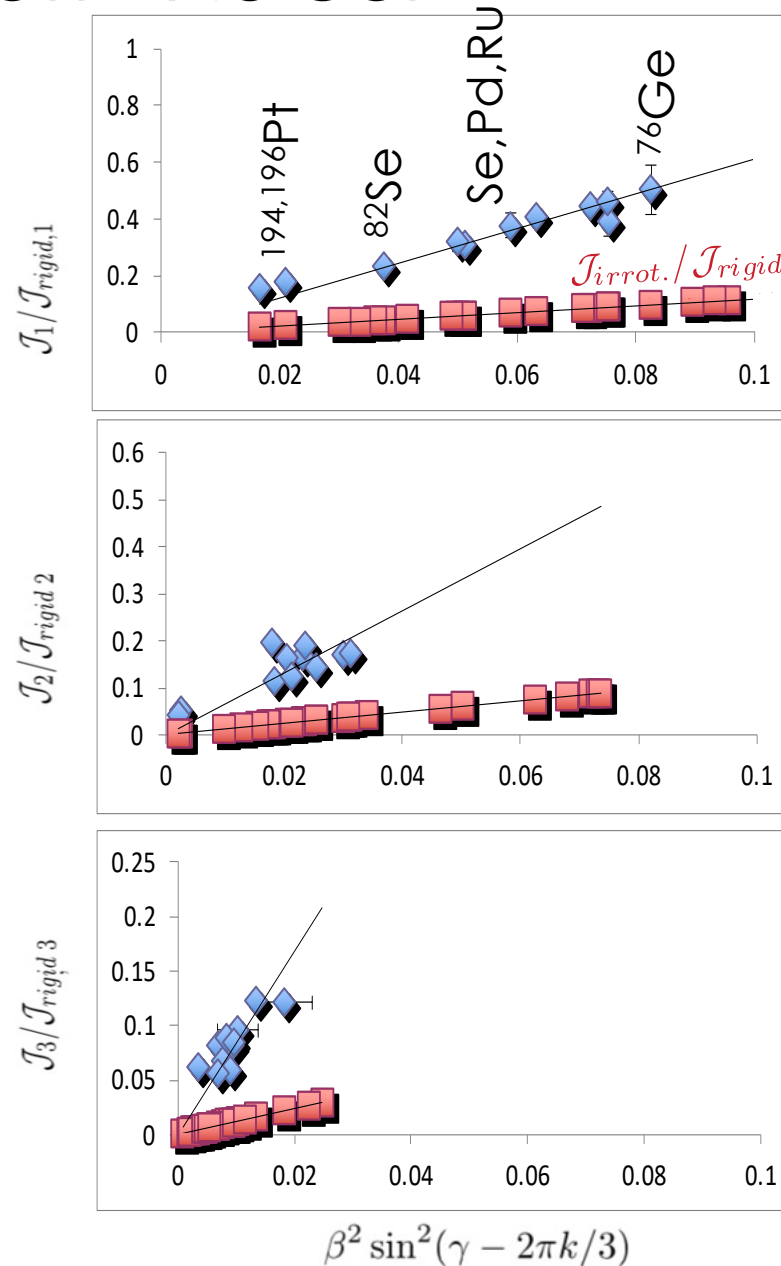
Empirical MOI of "Soft" Nuclei

Cases with $E(4)/E(2) = 2.4 - 2.7$



Empirical MOI of “Soft” Nuclei

Cases with $E(4)/E(2) = 2.4 - 2.7$



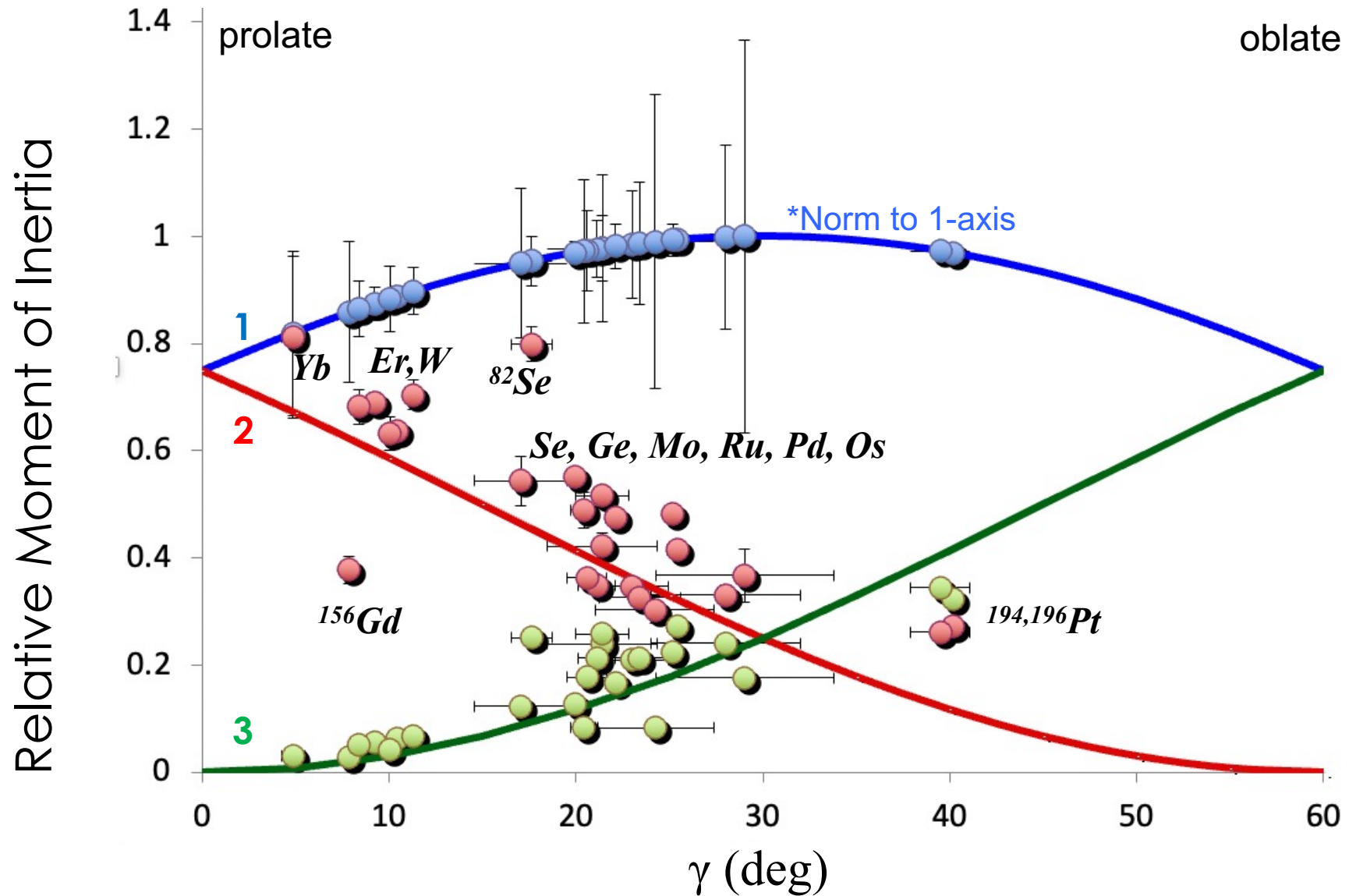
1-axis

2-axis

3-axis

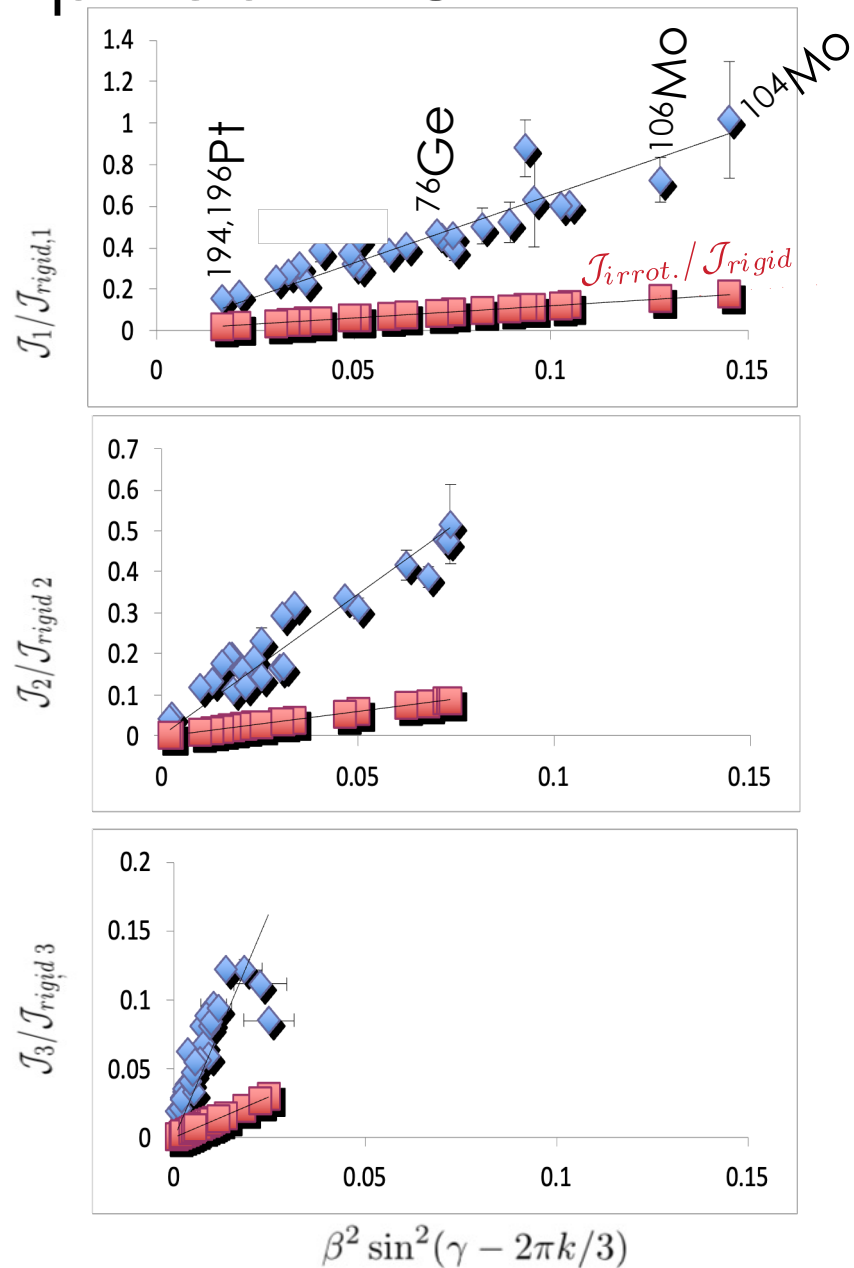
Updated View of Empirical MOI

Cases with $E(4)/E(2) = 2.4 - 3.3$



Updated View of Empirical MOI

Cases with $E(4)/E(2) = 2.4 - 3.3$



1-axis

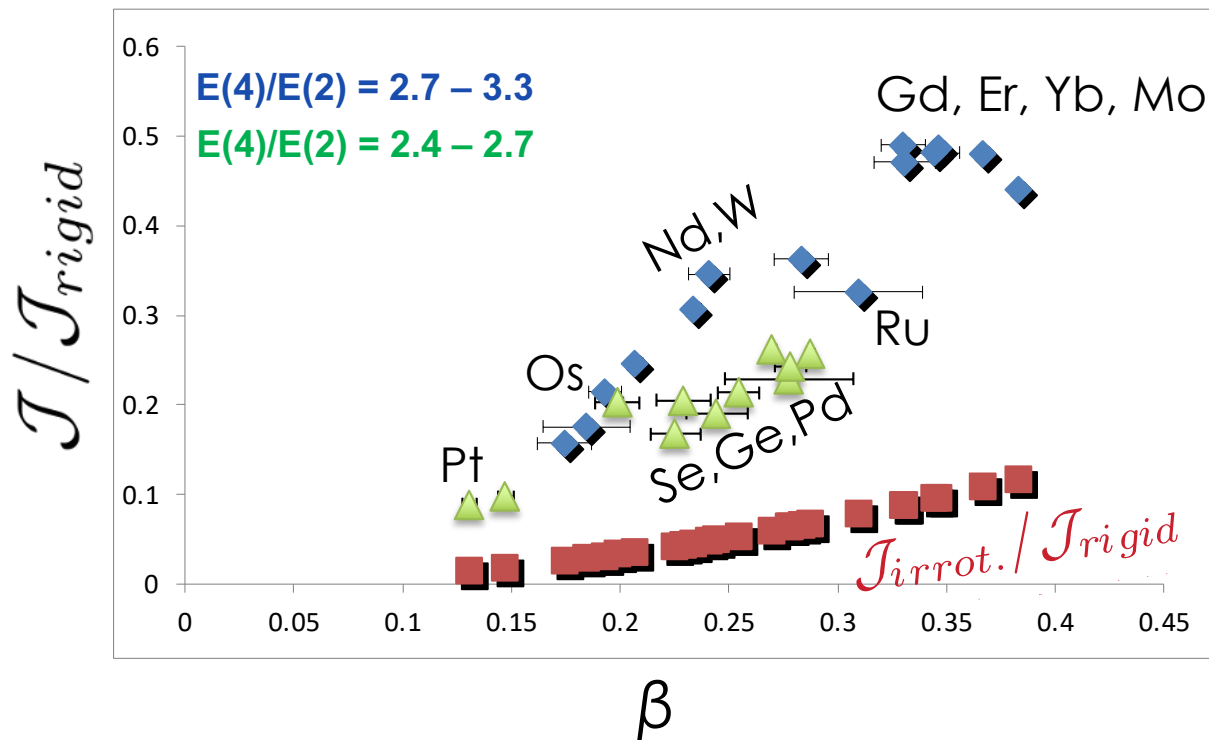
2-axis

3-axis

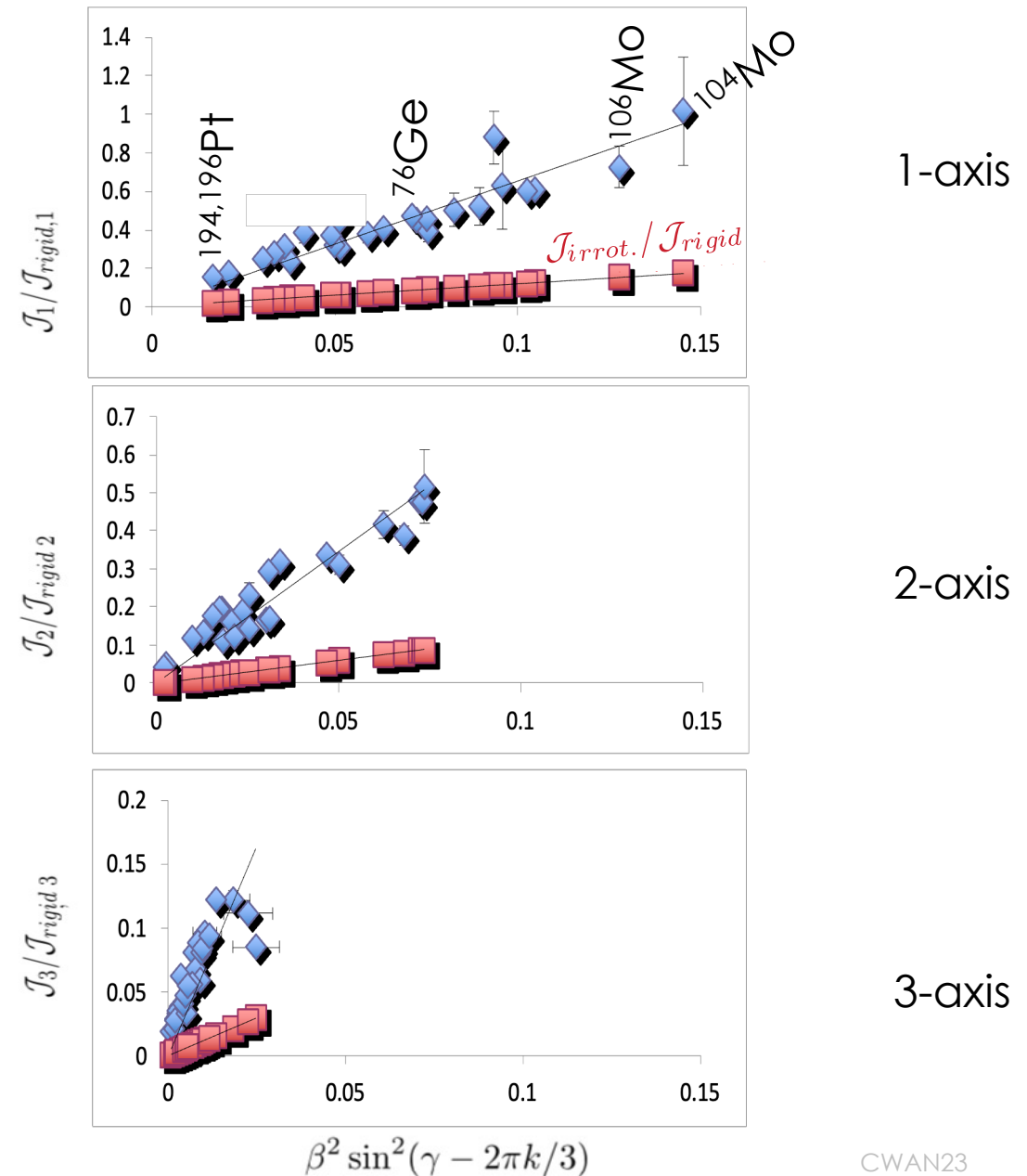
Updated View of Empirical MOI

Cases with $E(4)/E(2) = 2.4 - 3.3$

1D View shows bifurcation in MOI from "soft" nuclei



3D view removes strong bifurcation

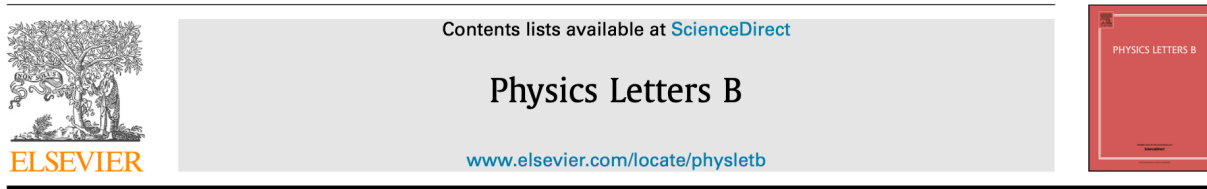


Summary

- Absolute MOI are between rigid and irrotational inertial flow values
- Relative MOI are qualitatively (not quantitatively) consistent with irrotational flow
- Bohr Hamiltonian is $SO(5)$ symmetric and Irrotational flow is also fortuitously $SO(5)$ symmetric, explaining limited success of Davydov-Filippov Model
- Abs. and Rel. MOI are qualitatively consistent with the slow rotation of a trapped cold Fermi Gas
- Similar MOI behavior for both $E(4)/E(2) > 2.7$ and $E(4)/E(2) = 2.4 - 2.7$
- 2- and 3-axis crossing happens after $\gamma = 30$ degrees
- Inertial dynamics may not be understood or constrained for more oblate nuclei, $\gamma > 40$ degrees
- Independent E2 and Inertia tensor required empirically and must begin thinking in terms of inertia or energy “softness” versus E2 “softness”; E2 character could be significantly more “rigid” than expected from energies.

Recent Case in Point for $E(4)/E(2) < 2.4$ and Future Teaser

- **E2 character could be significantly more “rigid” than expected from energies, cf. results on ^{106}Cd**
- Analysis of other Cd isotopes underway
- Preliminary results hint at constant $\langle Q^2 \rangle$ values like ^{76}Ge
- Do low-energy quadrupole “vibrations” even exist?



E2 rotational invariants of 0_1^+ and 2_1^+ states for ^{106}Cd : The emergence of collective rotation

T.J. Gray^{a,*}, J.M. Allmond^a, R.V.F. Janssens^{b,c}, W. Korten^d, A.E. Stuchbery^e, J.L. Wood^f, A.D. Ayangeakaa^{b,c}, S. Bottoni^{g,h}, B.M. Bucher^{i,j}, C.M. Campbell^k, M.P. Carpenter^g, H.L. Crawford^k, H. David^g, D.T. Doherty^{d,l}, P. Fallon^k, M.T. Febraro^a, A. Galindo-Uribarri^a, C.J. Gross^a, M. Komorowska^m, F.G. Kondev^g, T. Lauritsen^g, A.O. Macchiavelli^{k,a}, P. Napiorkowski^m, E. Padilla-Rodalⁿ, S.D. Pain^a, W. Reviol^{o,g}, D.G. Sarantites^o, G. Savard^g, D. Seweryniak^g, C.Y. Wuⁱ, C.-H. Yu^a, S. Zhu^{g,p,1}

