

International Conference on Chirality and Wobbling in Atomic Nuclei
Huizhou, July 10-14, 2023

Microscopic study of nuclear shapes and excitations

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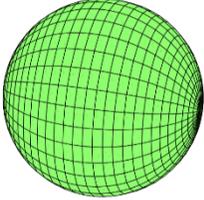
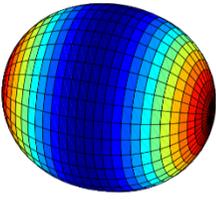
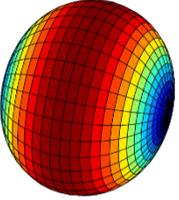
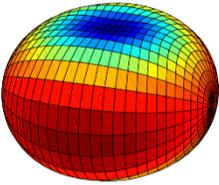
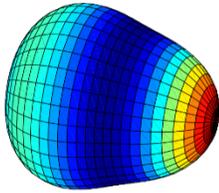
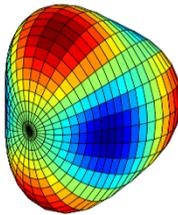
Outline

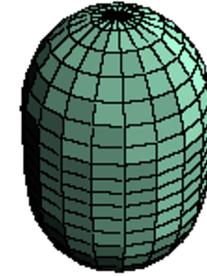


- 1. Introduction**
- 2. Microscopic collective Hamiltonian**
- 3. Results and discussion**
 - a. Quadrupole shape coexistence**
 - b. Octupole shapes and parity doublets**
- 4. Summary and outlook**

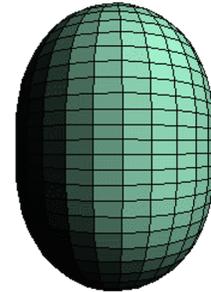
Nuclear shapes and excitations

$$R = R_0 \left[1 + \beta_{00} + \sum_{\lambda=1}^{\infty} \sum_{\mu=-\lambda}^{\lambda} \beta_{\lambda\mu}^* Y_{\lambda\mu}(\theta, \varphi) \right]$$

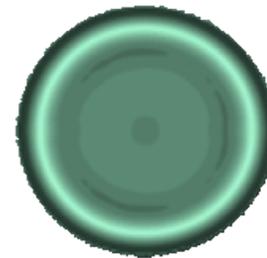
$\beta_{\lambda\mu} = 0$	$\beta_{20} > 0$	$\beta_{20} < 0$
		
$\beta_{22} \neq 0$	$\beta_{30} \neq 0$	$\beta_{32} \neq 0$
		



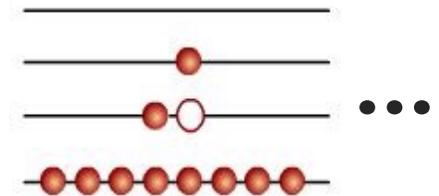
vibration



rotation



fission

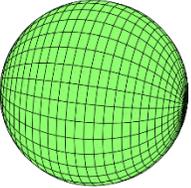
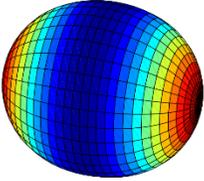
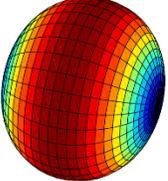
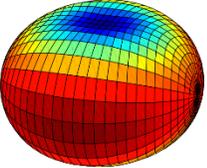
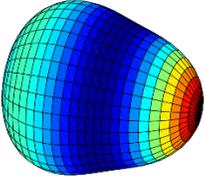
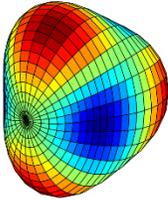


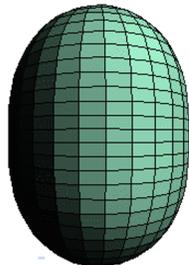
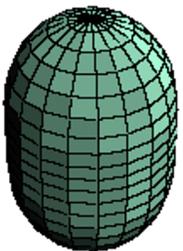
1p-1h excitation

s.p. excitation

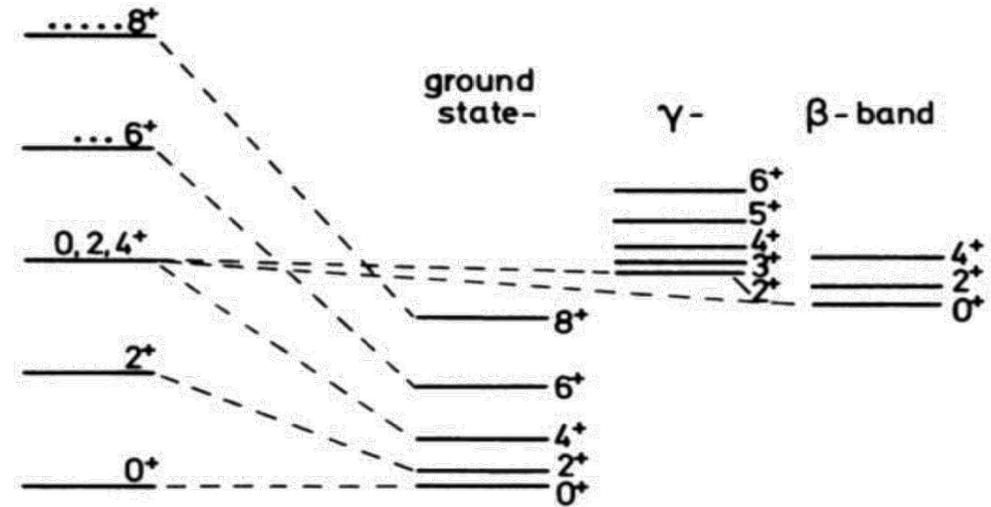
B.N. Lu, Ph.D. thesis, Academy of Sciences (2012)

Nuclear quadrupole shapes

$\beta_{\lambda\mu} = 0$	$\beta_{20} > 0$	$\beta_{20} < 0$
		
$\beta_{22} \neq 0$	$\beta_{30} \neq 0$	$\beta_{32} \neq 0$
		



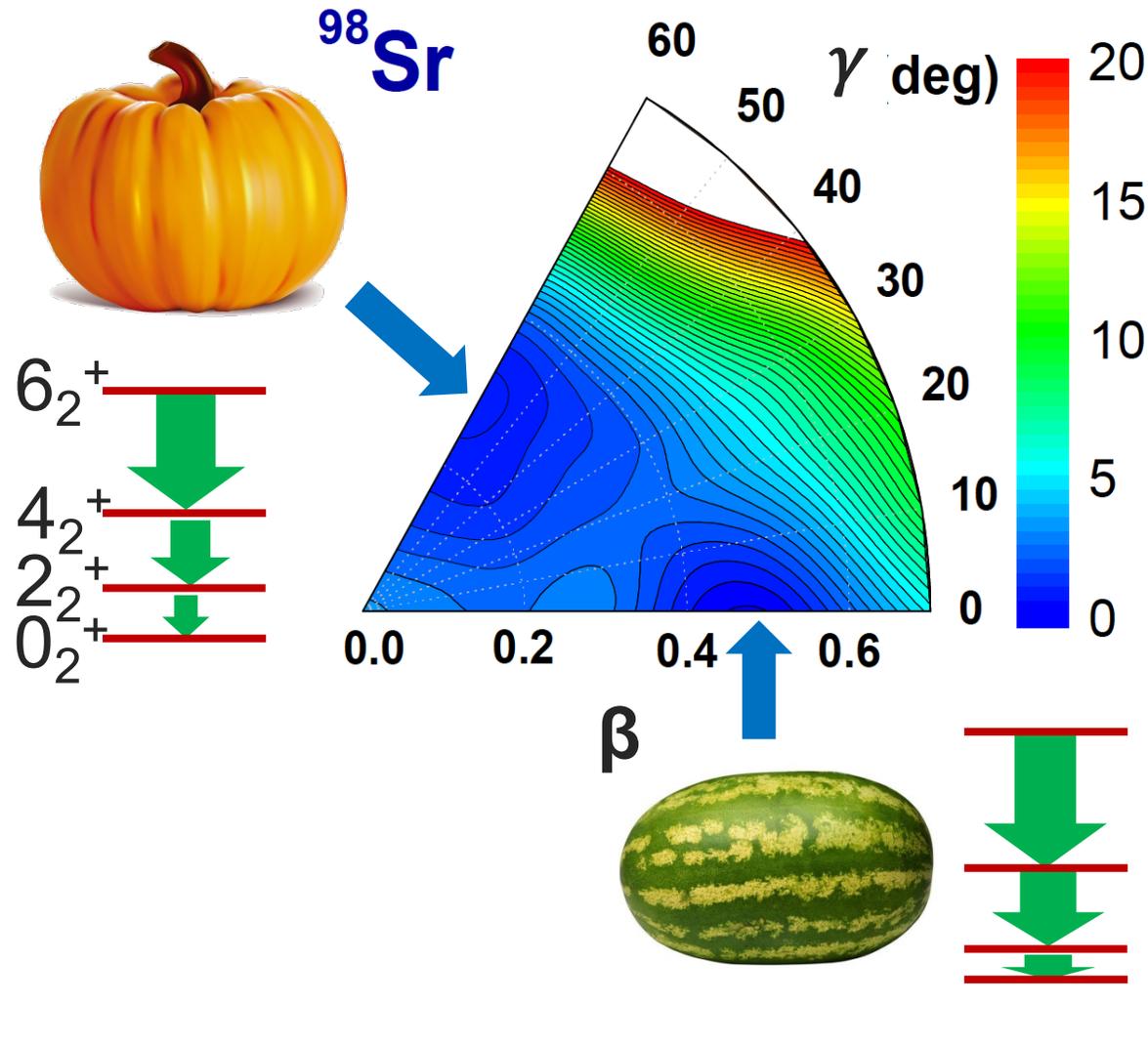
$$\hat{H} = -\frac{\hbar^2}{2B} \left[\frac{1}{\beta^4} \frac{\partial}{\partial \beta} \beta^4 \frac{\partial}{\partial \beta} + \frac{1}{\beta^2} \frac{1}{\sin(3\gamma)} \frac{\partial}{\partial \gamma} \sin(3\gamma) \frac{\partial}{\partial \gamma} \right] + \sum_{k=1}^3 \frac{\hat{j}_k^2}{2\mathcal{J}_k} + V(\beta, \gamma)$$



A. Bohr, Mat. Fys. Medd. K. Dan. Vidensk Selsk. 26, 14 (1952).

A. Bohr & Mottelson, Mat. Fys. Medd. K. Dan. Vidensk Selsk. 27, 16 (1953).

Nuclear shape coexistence

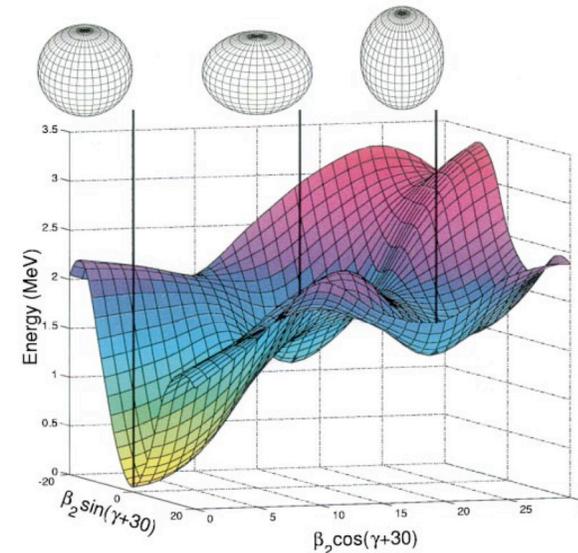


□ There may be deformation in the 0^+ excited state of ^{16}O .

H. Morinaga, Phys. Rev. 101, 254 (1956).

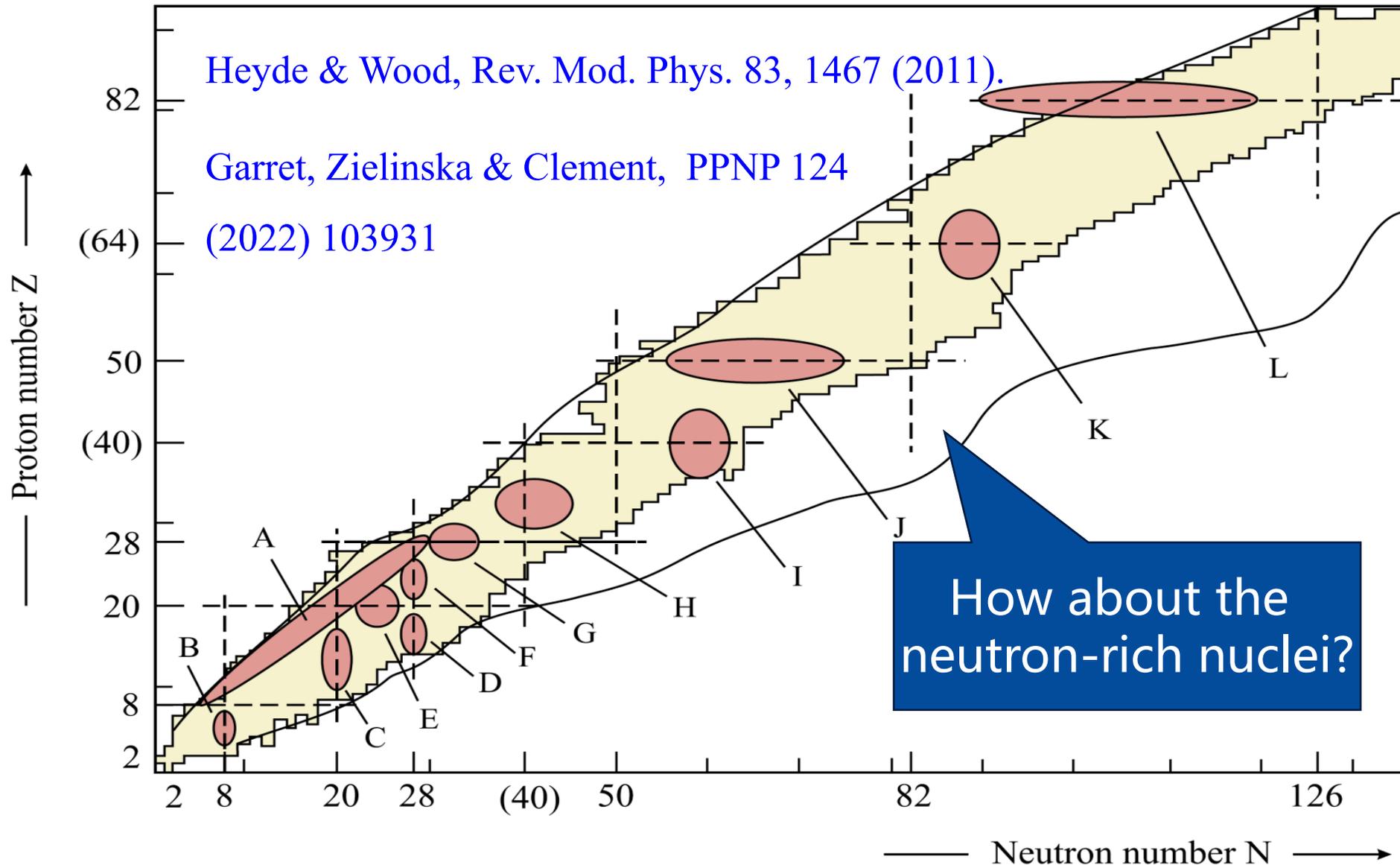
□ Late 1960s, near $Z = 50, 82$

□ In 2000, three shapes coexisted in ^{186}Pb .

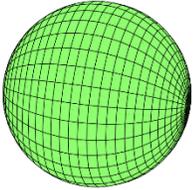
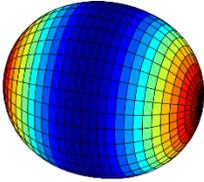
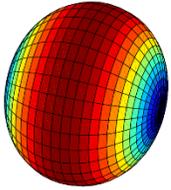
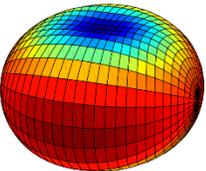
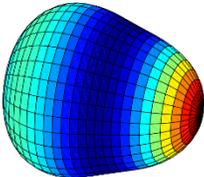
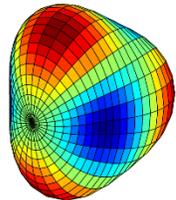


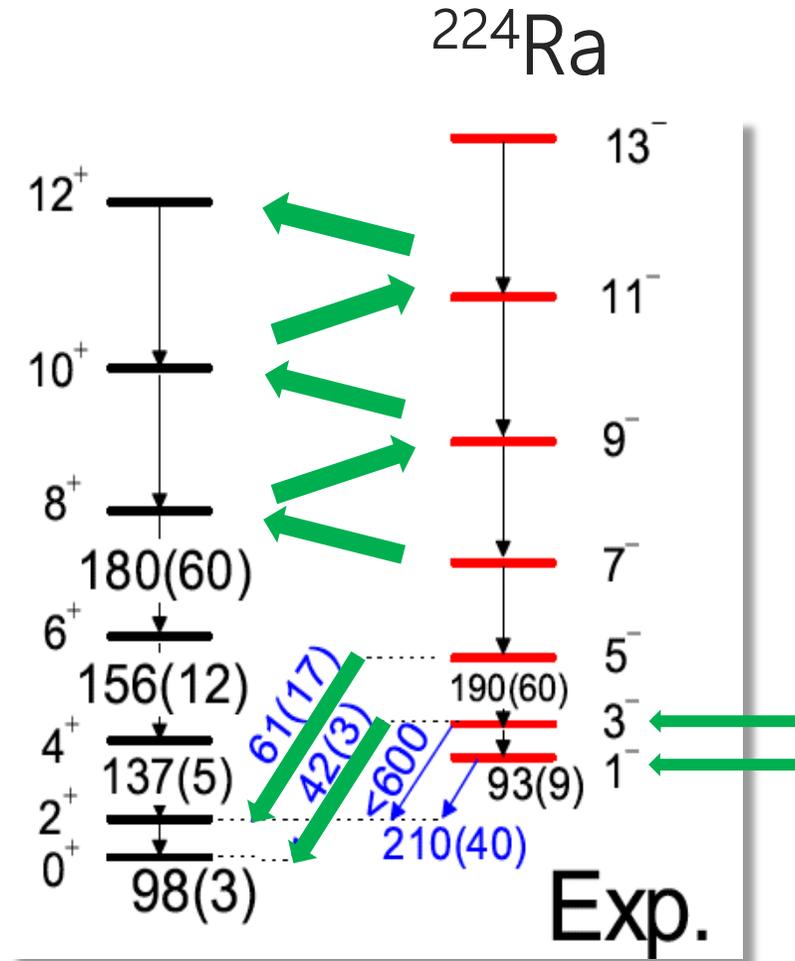
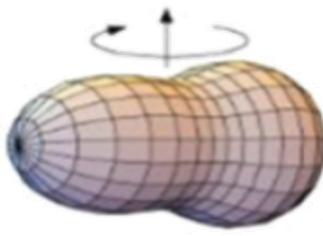
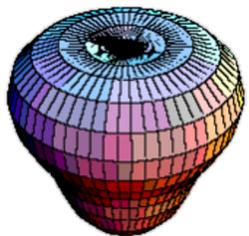
Andreyev et al., Nature 405, 430 (2000).

Nuclear shape coexistence



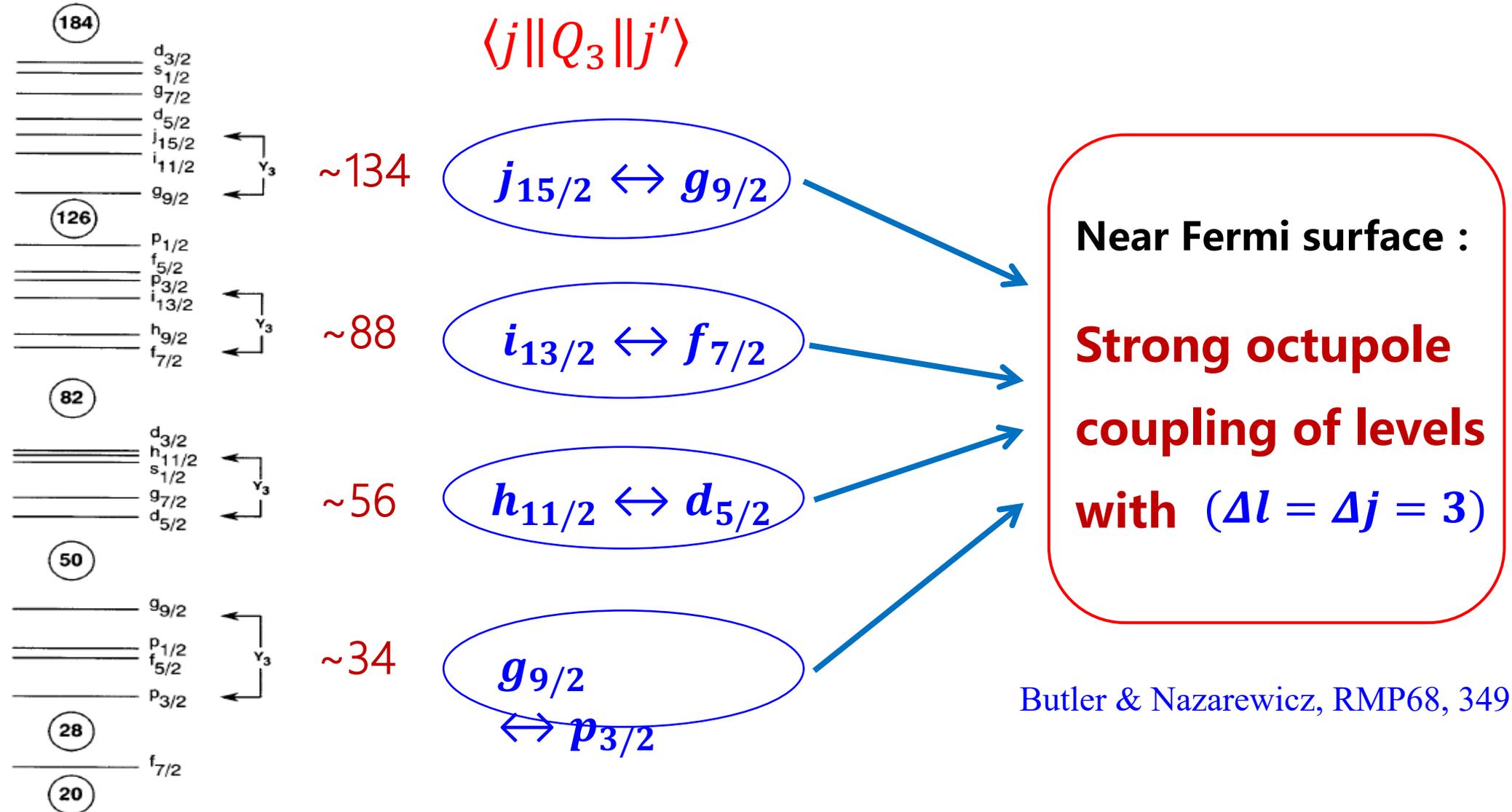
Axially symmetric octupole shapes

$\beta_{\lambda\mu} = 0$	$\beta_{20} > 0$	$\beta_{20} < 0$
		
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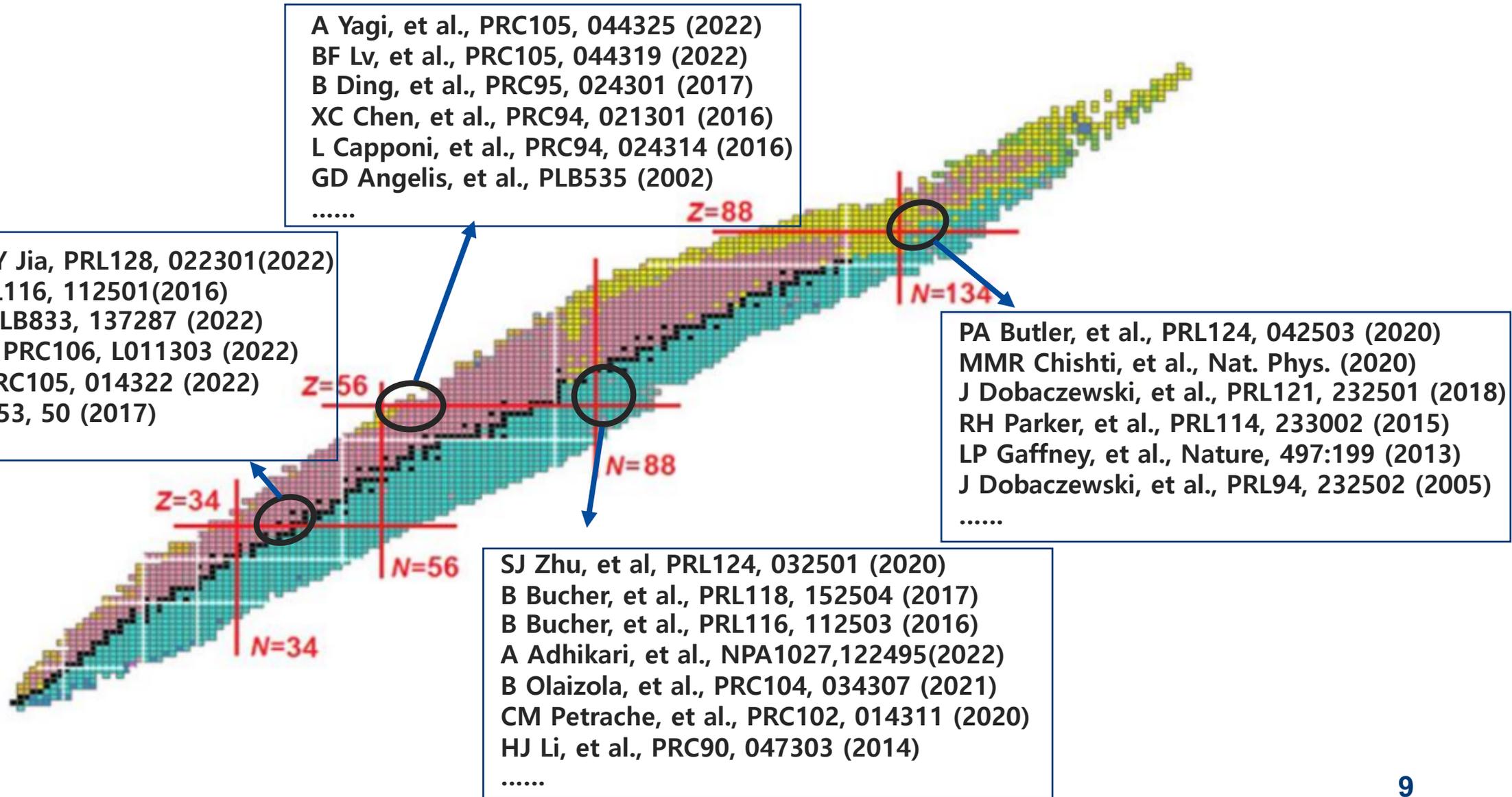


Butler & Nazarewicz, RMP68, 349 (1996).

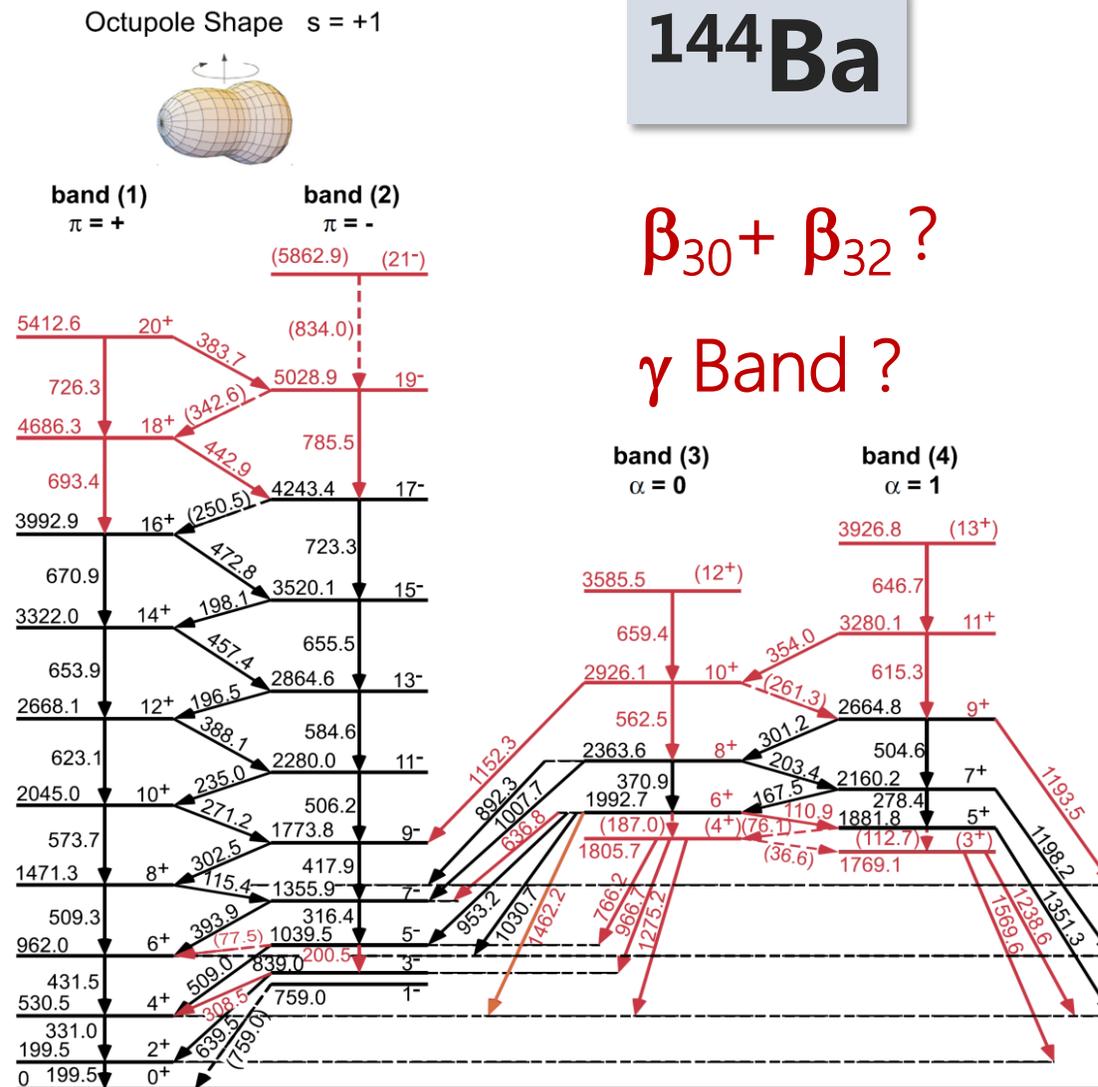
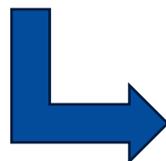
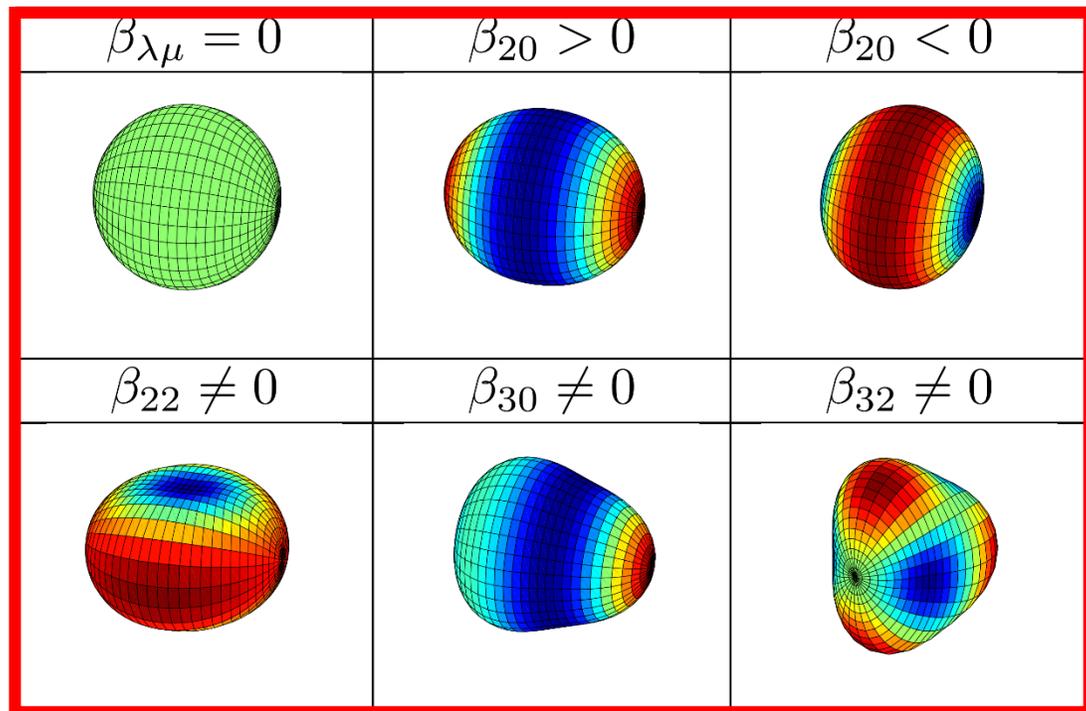
Axially symmetric octupole shapes



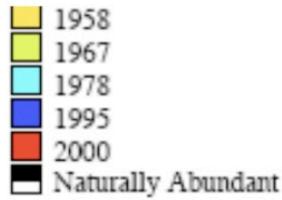
Experimental progress for octupole correlations



Triaxial octupole deformation ?



Microscopic models for nuclear excitation

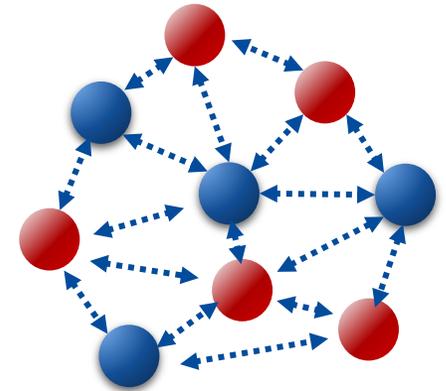


Ab initio:
Quantum Monte Carlo,
No-core Shell Model,
Coupled-Cluster, ...

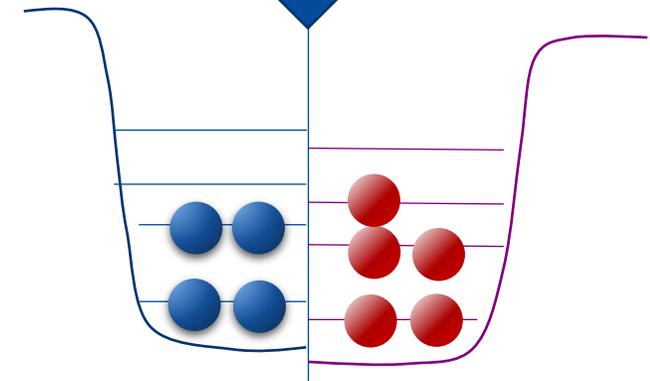
Density Functionals

Configuration interaction
(Interacting Shell-Model)

**Global study on
nuclear ground and
excitation properties**



MF appr.



- ◆ Simpler wave function
- ◆ More correlations: ph, pp

Limitation of MF

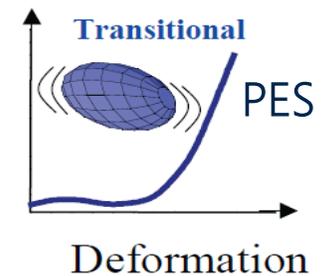
- intrinsic ground state properties: mass, radius, shape...; PES



“Real” ground state (good quantum numbers) and excitations:

beyond the MF approximation

- Restoration of broken symmetry, e.g. rotational, parity
- Mixing of different shapes
- Inclusion of quasiparticles



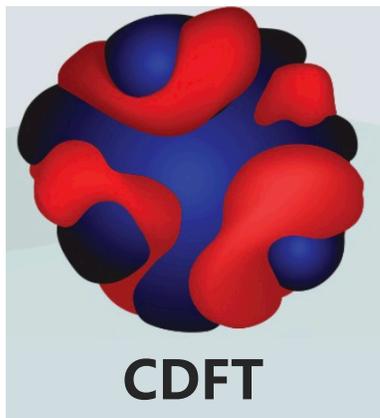
Collective Hamiltonian (DFT)

Outline



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Collective Hamiltonian based on CDFT



Construct



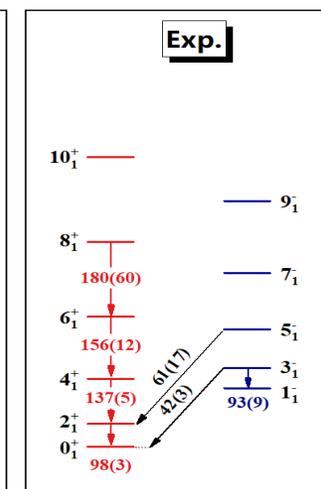
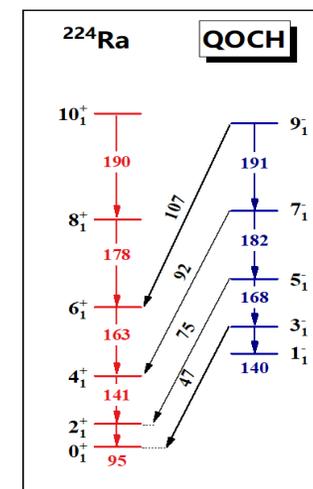
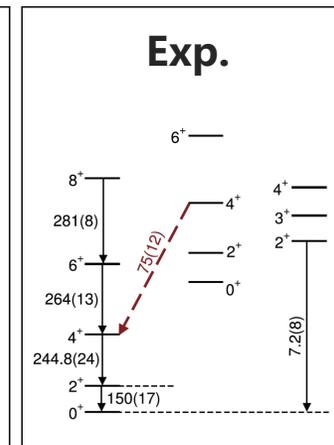
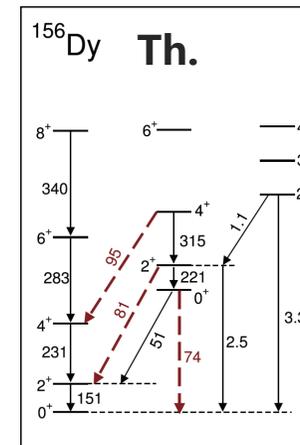
Collective parameters

Collective Hamiltonian

$$\hat{H}_{\text{coll}} = \hat{T}_{\text{vib}} + \hat{T}_{\text{rot}} + V_{\text{coll}}$$

$(\beta_{20}, \beta_{22}, \beta_{30}, \beta_{32}, \Omega \dots)$

Solve



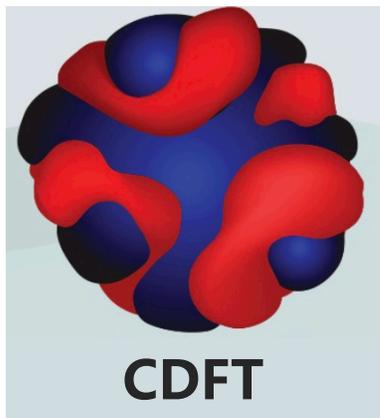
ZPLi, Niksic, Vretenar, Meng, Lalazissis, & Ring, PRC (2009)

ZPLi, Song, Yao, Vretenar & Meng. PLB (2013)

Xiang, ZPLi, Niksic, Vretenar & Long, PRC (2020)

ZPLi & Vretenar, *Model for Collective Motion. In: Handbook of Nuclear Physics*. Springer, Singapore.

Collective Hamiltonian based on CDFT



Construct



Collective parameters

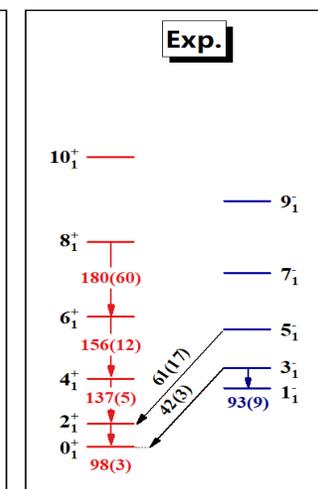
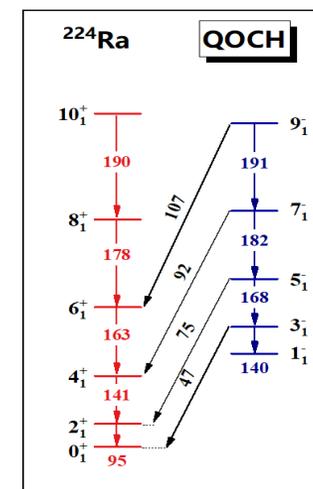
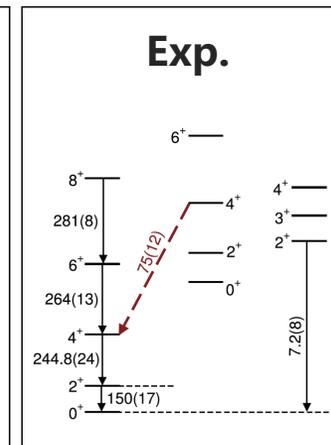
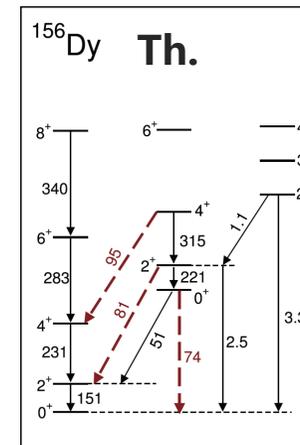
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$$(\beta_{20}, \beta_{22}, \beta_{30}, \beta_{32}, \Omega \dots)$$

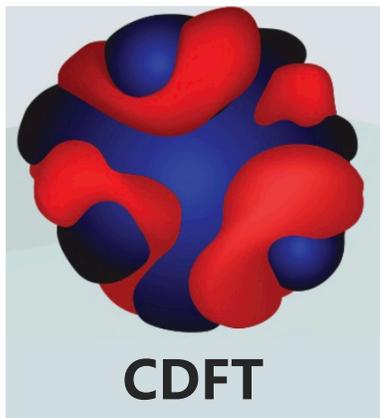
Solve

$$\begin{aligned}
 E_{\text{CDF}} &= \int d^3(\mathbf{r}) \varepsilon_{\text{CDF}}(\mathbf{r}) \\
 &= \sum_k \int d\mathbf{r} v_k^2 \bar{\psi}_k(\mathbf{r}) (-i\gamma\nabla + m)\psi_k(\mathbf{r}) \\
 &+ \int d\mathbf{r} \left(\frac{\alpha_S}{2} \rho_S^2 + \frac{\beta_S}{3} \rho^3 + \frac{\gamma_S}{3} \rho_S^4 + \frac{\delta_S}{2} \rho_S \Delta \rho_S + \frac{\alpha_V}{2} j_\mu j^\mu + \frac{\gamma_V}{4} (j_\mu j^\mu)^2 \right. \\
 &\left. + \frac{\delta_V}{2} j_\mu \Delta j^\mu + \frac{\alpha_{TV}}{2} \vec{j}_{TV}^\mu \cdot (\vec{j}_{TV})_\mu + \frac{\delta_{TV}}{2} \vec{j}_{TV}^\mu \cdot \Delta (\vec{j}_{TV})_\mu + \frac{e^2}{2} \rho_p A^0 \right)
 \end{aligned}$$



PC-PK1 functional

Collective Hamiltonian based on CDFT



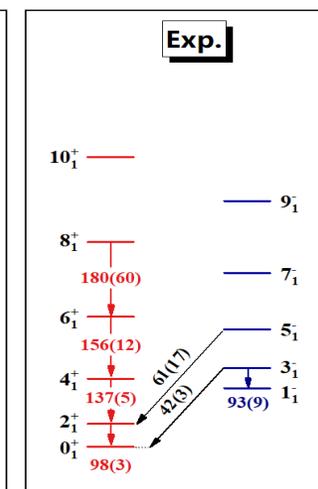
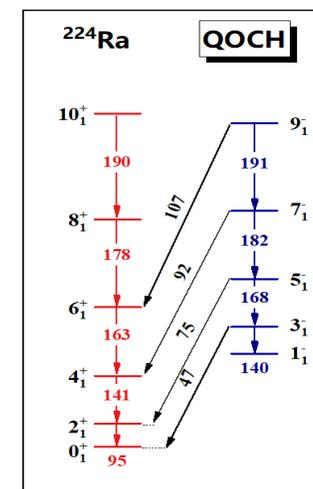
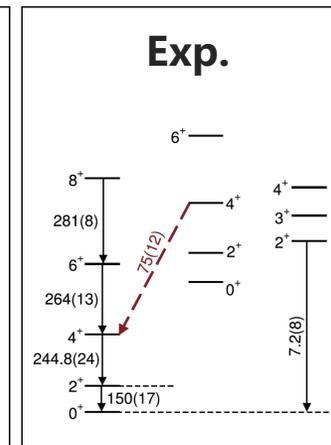
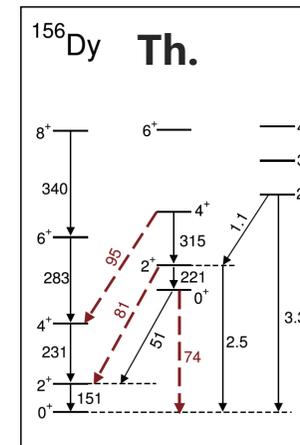
Construct
Collective parameters

Collective Hamiltonian

$$\hat{H}_{\text{coll}} = \hat{T}_{\text{vib}} + \hat{T}_{\text{rot}} + V_{\text{coll}}$$

$$(\beta_{20}, \beta_{22}, \beta_{30}, \beta_{32}, \Omega \dots)$$

Solve

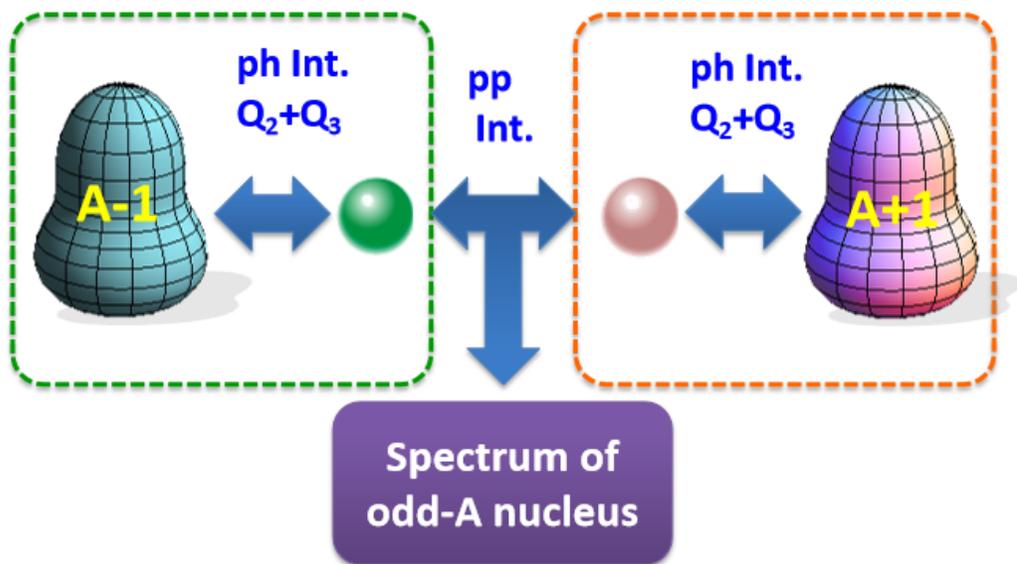


$$\hat{H}_{\text{coll}} = -\frac{\hbar^2}{2\sqrt{G}} \sum_{ij} \frac{\partial}{\partial \beta_i} \sqrt{G} B_{vib}^{-1}{}_{ij} \frac{\partial}{\partial \beta_j} + \sum_{k=1}^3 \frac{\hat{I}_k^2}{2\mathcal{J}_k} + V_{\text{coll}}(\beta)$$

Mass tensor:

$$B_{vib} = \begin{pmatrix} B_{\beta_{20}\beta_{20}} & B_{\beta_{20}\beta_{22}} & B_{\beta_{20}\beta_{30}} & B_{\beta_{20}\beta_{32}} \\ B_{\beta_{22}\beta_{20}} & B_{\beta_{22}\beta_{22}} & B_{\beta_{22}\beta_{30}} & B_{\beta_{22}\beta_{32}} \\ B_{\beta_{30}\beta_{20}} & B_{\beta_{30}\beta_{22}} & B_{\beta_{30}\beta_{30}} & B_{\beta_{30}\beta_{32}} \\ B_{\beta_{32}\beta_{20}} & B_{\beta_{32}\beta_{22}} & B_{\beta_{32}\beta_{30}} & B_{\beta_{32}\beta_{32}} \end{pmatrix}$$

➤ Microscopic core-quasiparticle coupling model (CQC)



Sun, Quan, ZPLi, et al, PRC100, 044319 (2019)

$$|\alpha J M_J \pi\rangle^A = \sum_{\mu\Omega} \left\{ U_{\mu\Omega} [a_{\mu}^{\dagger} | \Omega \rangle]_{J M_J \pi}^{A-1} + V_{\mu\Omega} [a_{\mu} | \Omega \rangle]_{J M_J \pi}^{A+1} \right\}$$

$$\begin{aligned}
 H &= H_{\text{qp}} + H_c \quad \text{s.p. levels} \\
 &= \begin{pmatrix} (\varepsilon^{A-1} & -\varepsilon_f) + \Gamma^{A-1} \\ & \Delta^{\dagger A-1} \end{pmatrix} \quad \text{Multipole int.} \\
 &\quad + \begin{pmatrix} & \Delta^{A+1} \\ -(\varepsilon^{A+1} & -\varepsilon_f) - \Gamma^{A+1} \end{pmatrix} \quad \text{Pairing gap} \\
 &\quad + \begin{pmatrix} E^{A-1} & 0 \\ 0 & E^{A+1} \end{pmatrix} \quad \text{Collective levels}
 \end{aligned}$$

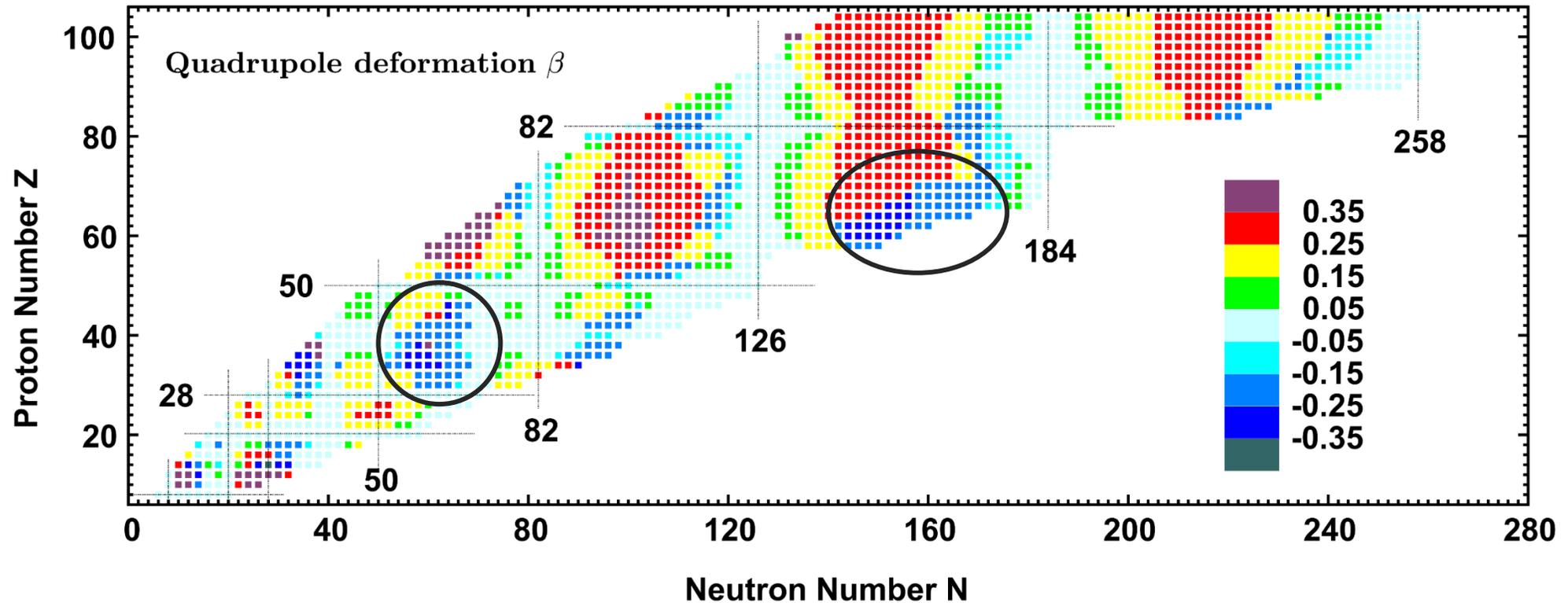
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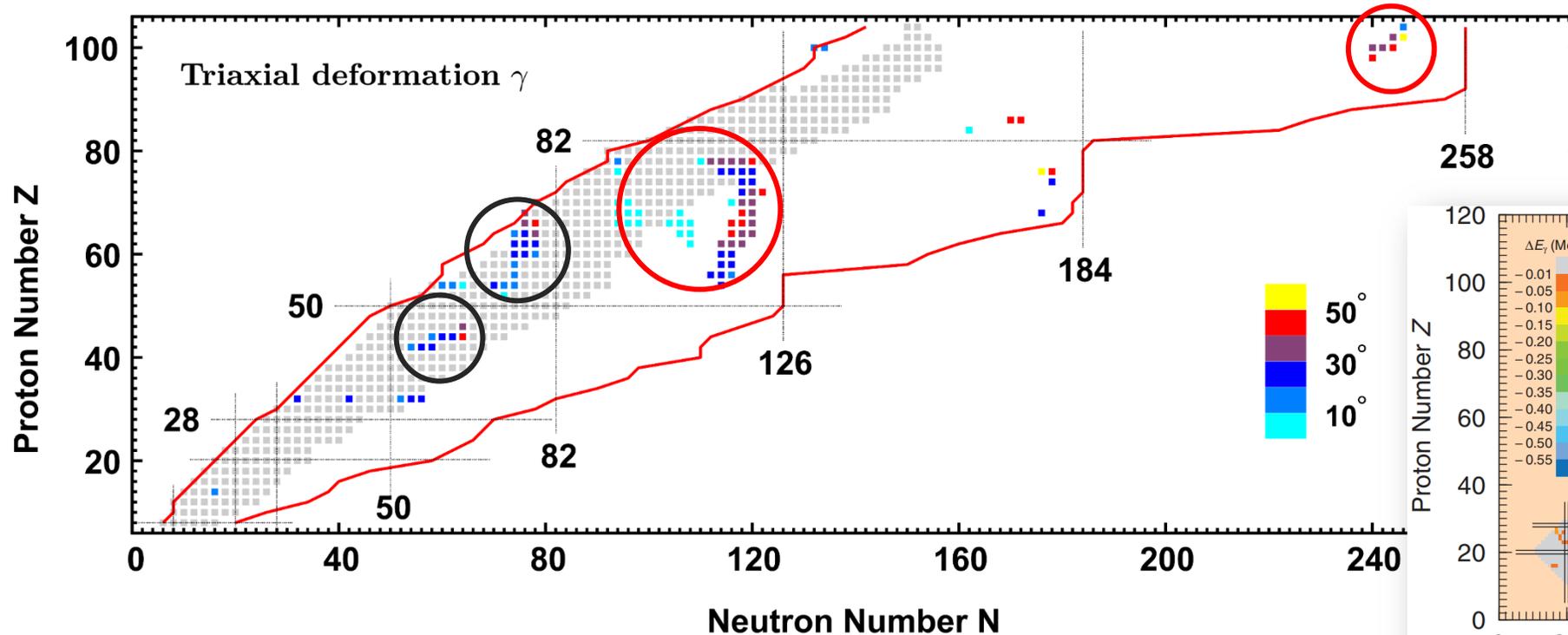


Static quadrupole deformation β

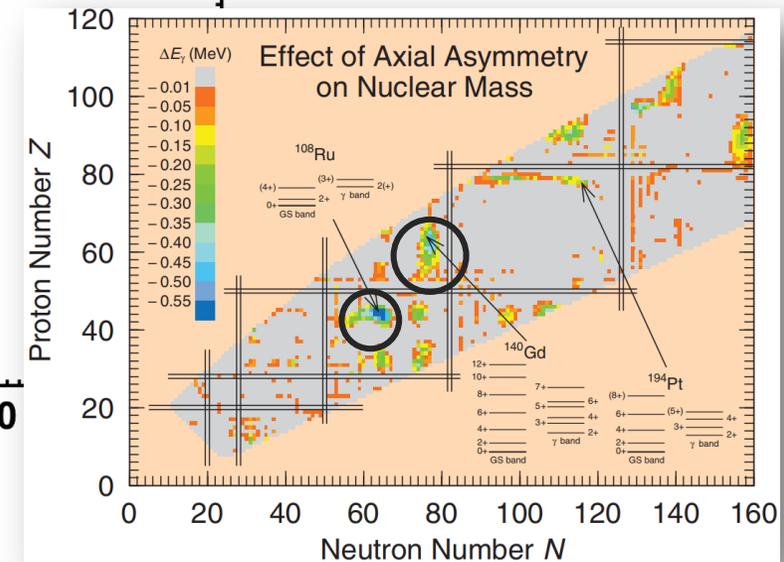


- Most of the deformed nuclei are prolate ($\beta > 0$). The nuclei with oblate deformation are mainly distributed in $(N, Z) \sim (60, 40), (150, 60)$.
- The quadrupole deformation of most nuclei is in the range of $-0.35 < \beta < 0.35$.

Static triaxial deformation γ



P. Möller et al., *Phys. Rev. Lett.*
97, 162502 (2006)



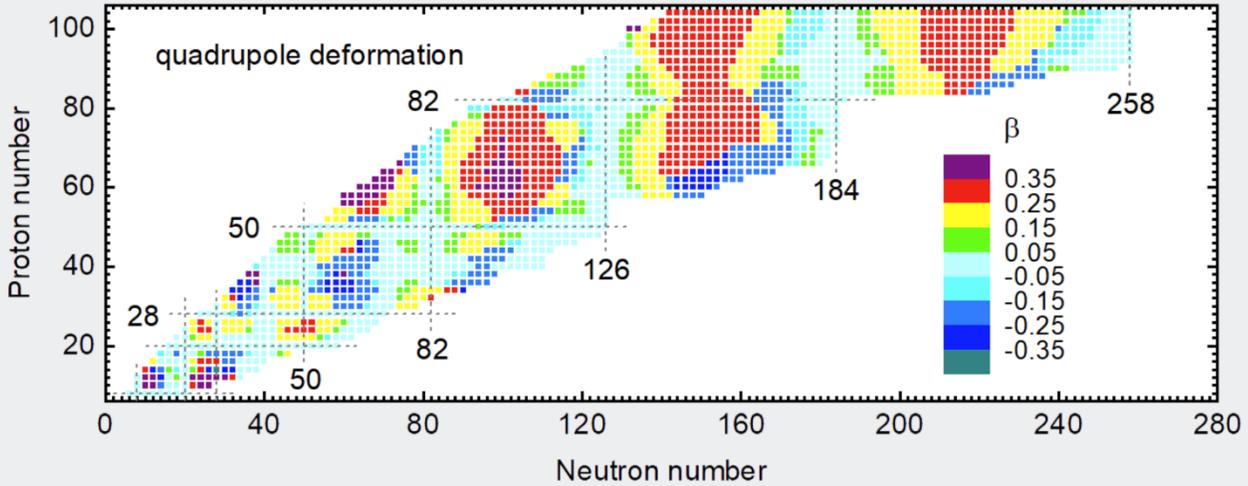
➤ 108 triaxially deformed even-even nuclei; Triaxial deformation energy: 0~1.89 MeV

➤ Finite-range liquid-drop model:

>700 triaxially deformed nuclei; Triaxial deformation energy : 0~0.63 MeV

Triaxial Relativistic Hartree-Bogoliubov Results with the PC-PK1 Density Functional

Search a nuclide:



quadrupole deformation

Proton number

Neutron number

β

- 0.35
- 0.25
- 0.15
- 0.05
- 0.05
- 0.15
- 0.25
- 0.35

Results for Krypton 74 (Z=36, N=38)

Ground-state properties

Spectroscopy (coming soon)

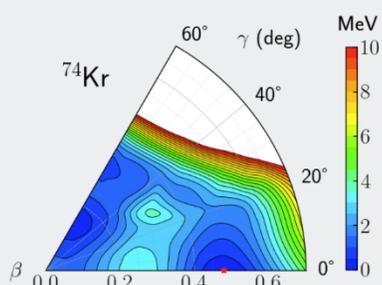
$E_{\text{RHB}} = -626.72$ MeV
 $E_{\text{SDCH}} = -630.61$ MeV
 $E_{\text{exp}} = -631.44$ MeV
 $\beta = 0.48$
 $\gamma = 0^\circ$

To know the meaning of a quantity, please hold the mouse still on it.

Select a nuclide

Proton number

Neutron number



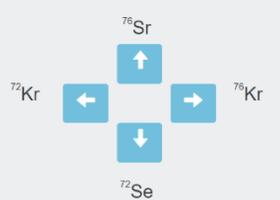
^{74}Kr

γ (deg)

MeV

β

Potential energy surface calculated by RHB theory with PC-PK1 density functional. All energies are normalized with respect to the binding energy of the absolute minimum. The contours join points on the surface with the same energy, and the energy difference between adjacent contours is 0.5 MeV.

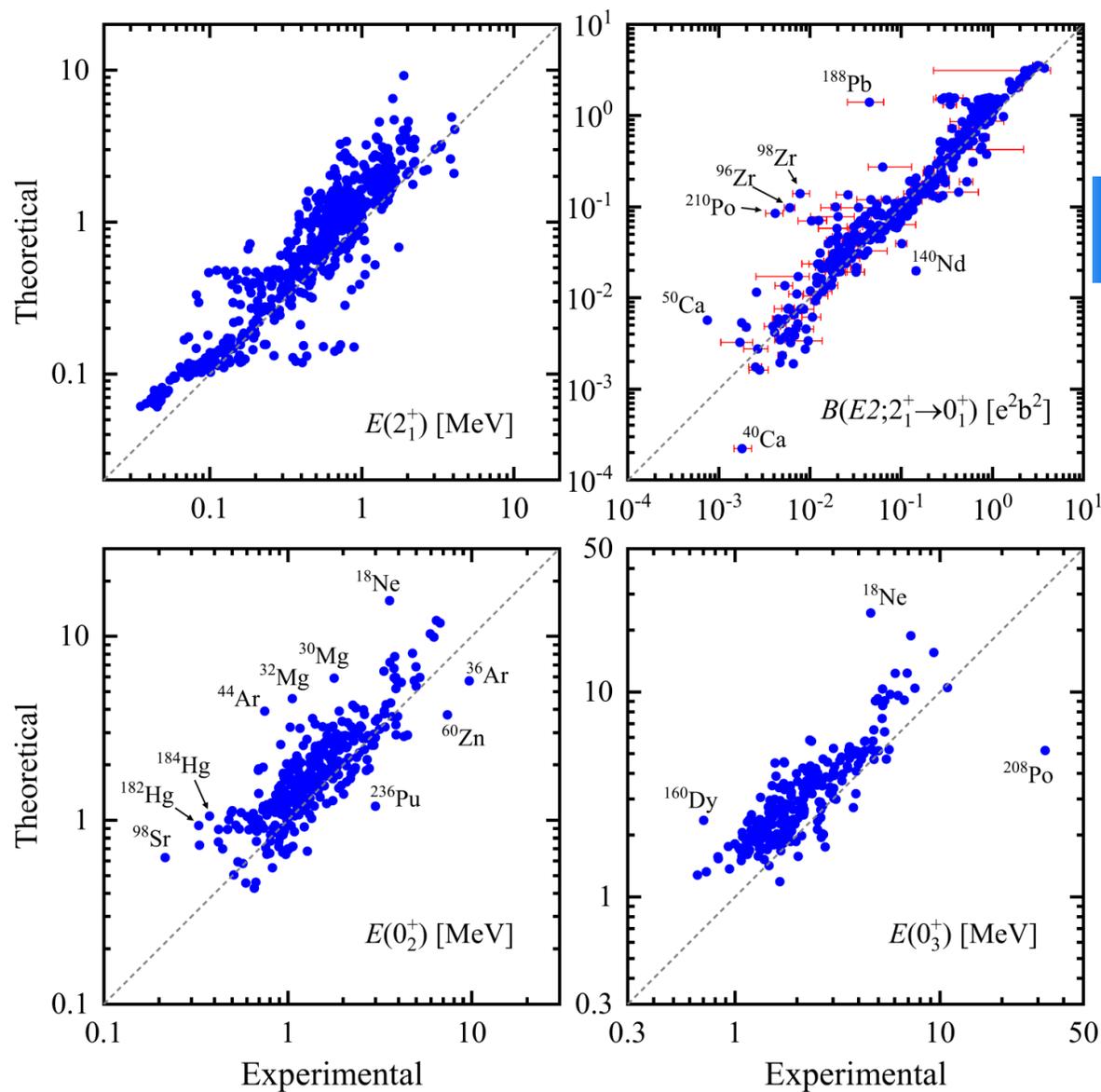


^{76}Sr

^{72}Kr ^{78}Kr

^{72}Se

Quadrupole excitations

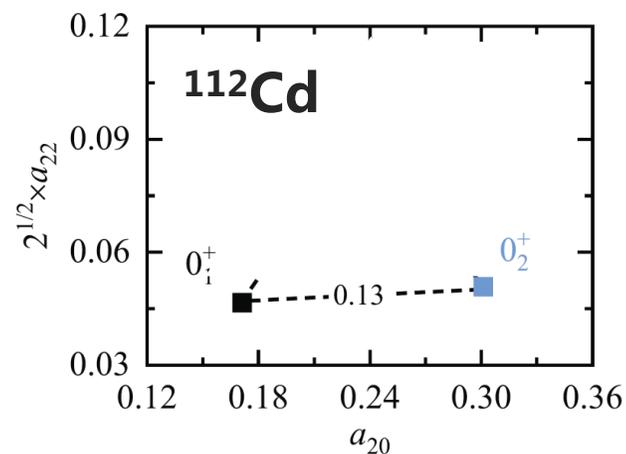


$$q_2(0_i^+) = \sum_j \langle 0_i^+ || \hat{Q}_2 || 2_j^+ \rangle \langle 2_j^+ || \hat{Q}_2 || 0_i^+ \rangle$$

$$= \left(\frac{3ZeR^2}{4\pi} \right)^2 [(a_{20}^{\text{eff}})^2 + 2(a_{22}^{\text{eff}})^2]$$

$$q_3(0_i^+) = -\sqrt{\frac{7}{10}} \sum_{ik} \langle 0_i^+ || \hat{Q}_2 || 2_j^+ \rangle \langle 2_j^+ || \hat{Q}_2 || 2_k^+ \rangle \langle 2_k^+ || \hat{Q}_2 || 0_i^+ \rangle$$

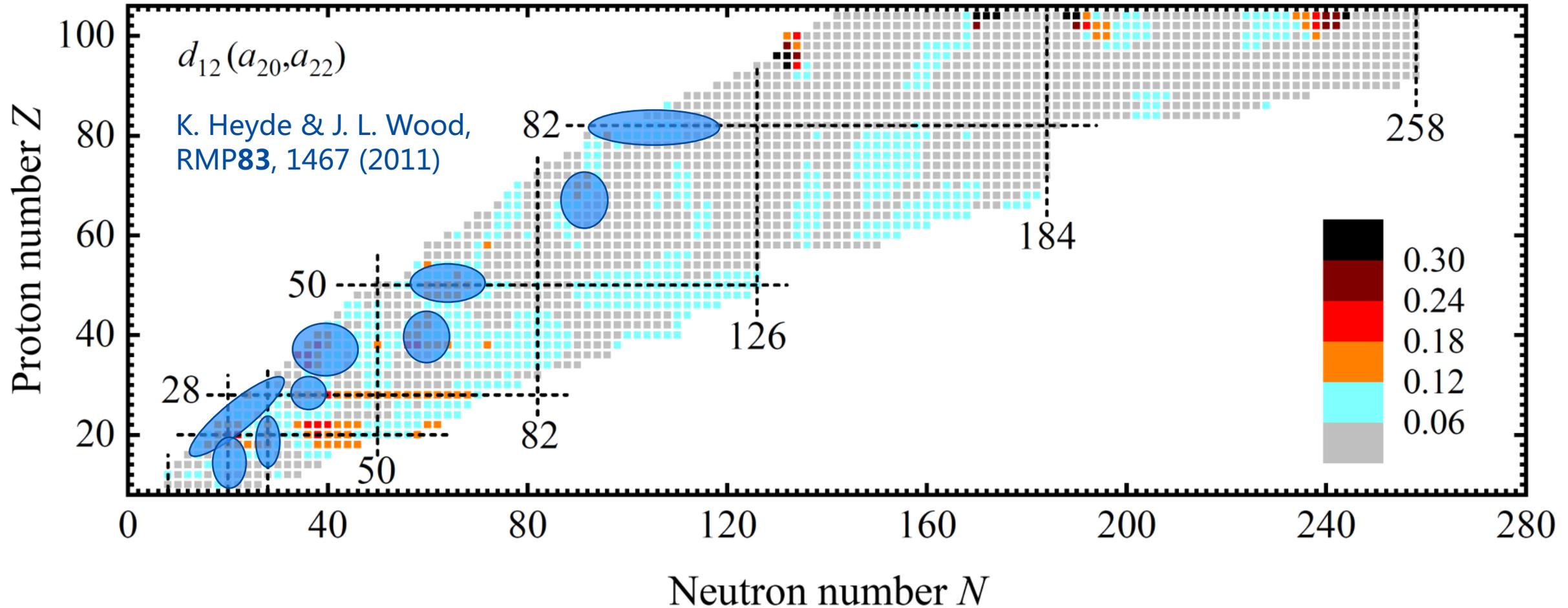
$$= \left(\frac{3ZeR^2}{4\pi} \right)^3 [(a_{20}^{\text{eff}})^3 - 6a_{20}^{\text{eff}}(a_{22}^{\text{eff}})^2]$$



➤ **0^+ states have different shapes**

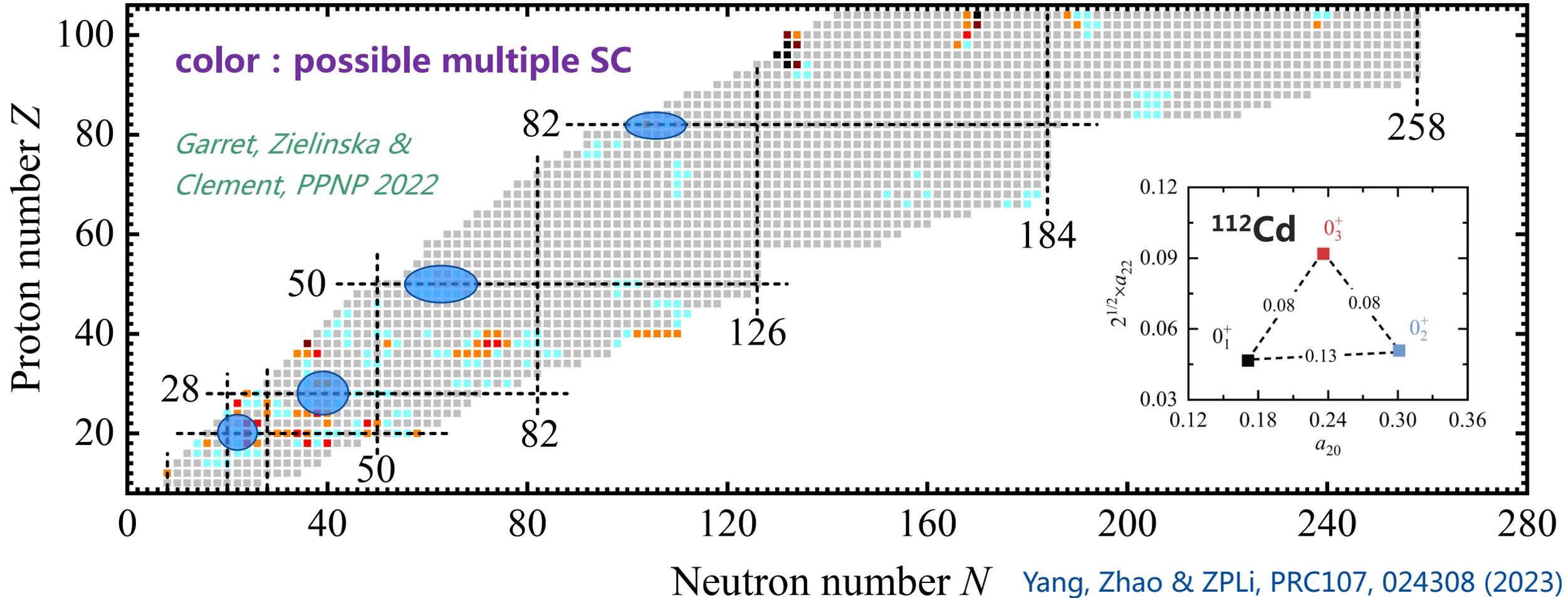
Garrett et al., PRL 2019

Nuclear shape coexistence



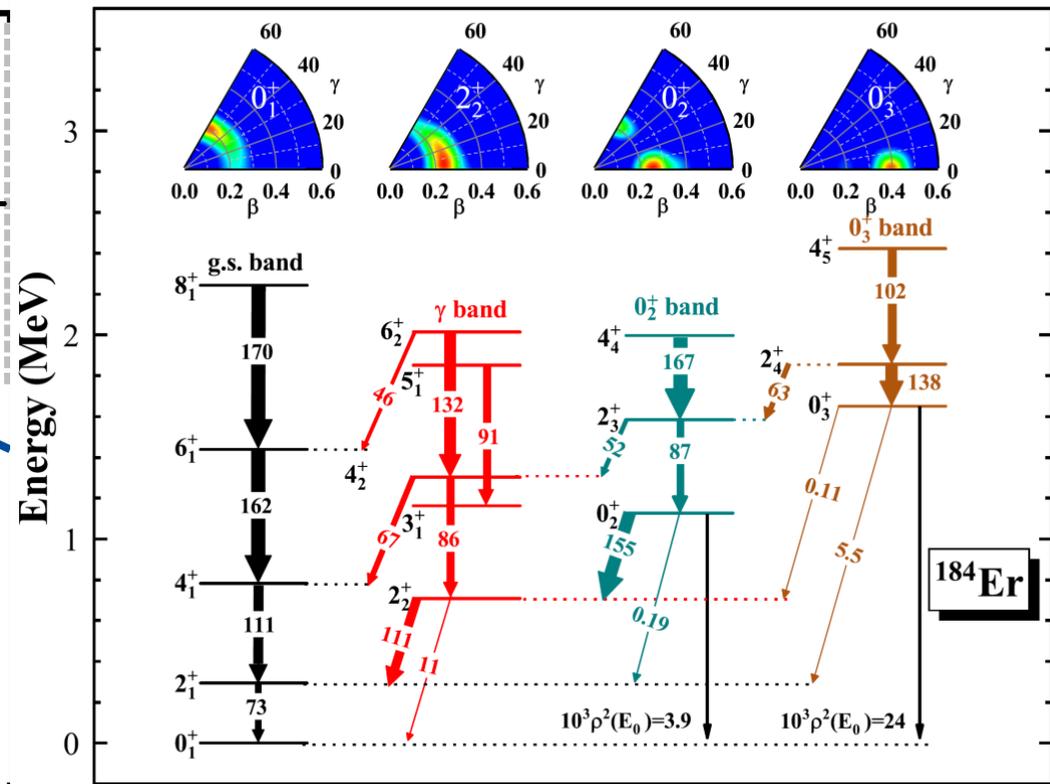
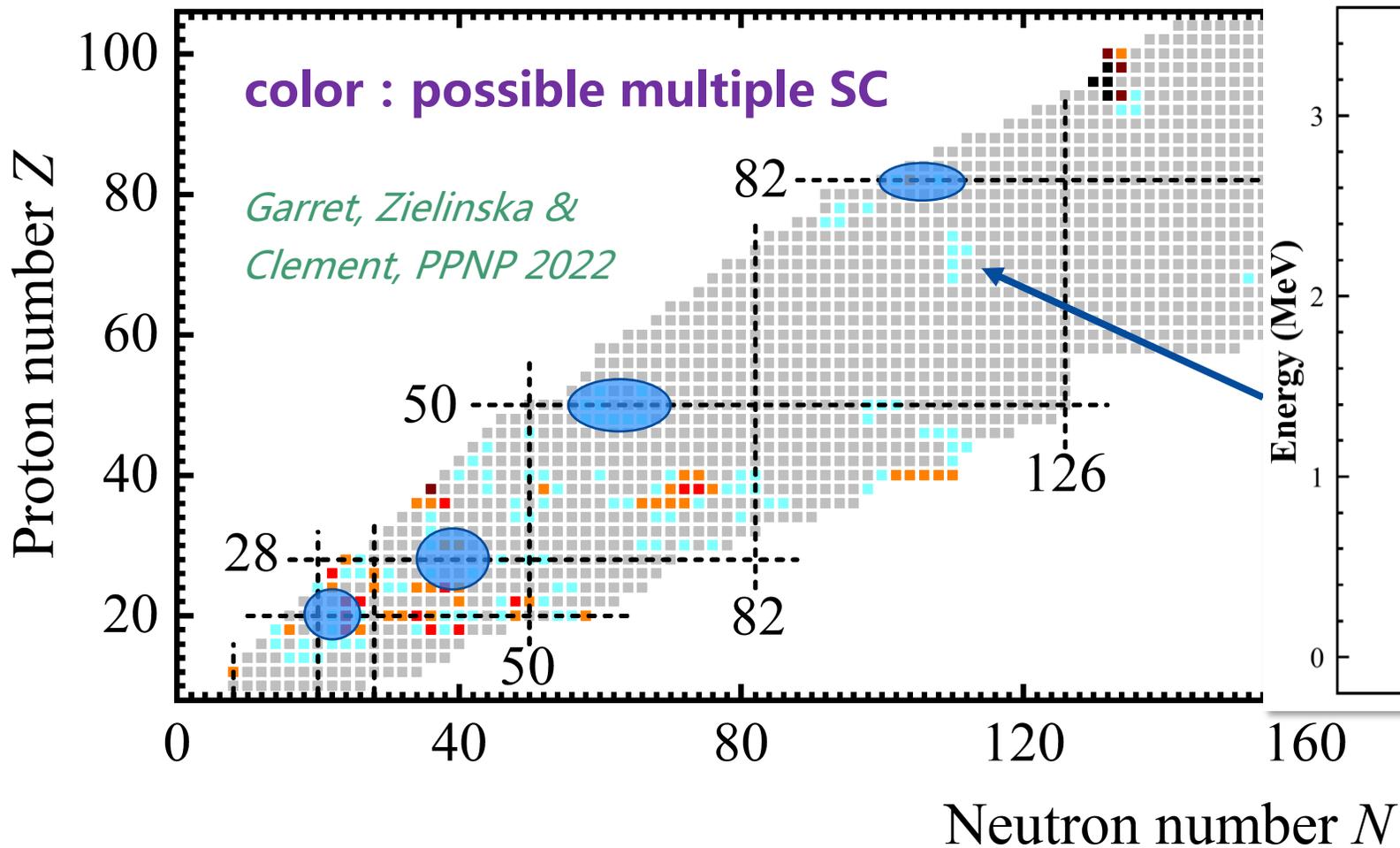
- **$Z \sim$ shell or sub-shell; transitional regions**
- **New regions in neutron-rich side and heavy nuclei**

Multiple shape coexistence



➤ $Z \sim 20, 28, 40, 50, 82$; transitional regions

Multiple shape coexistence



➤ $Z \sim 20, 28, 40, 50, 82$; transitional regions

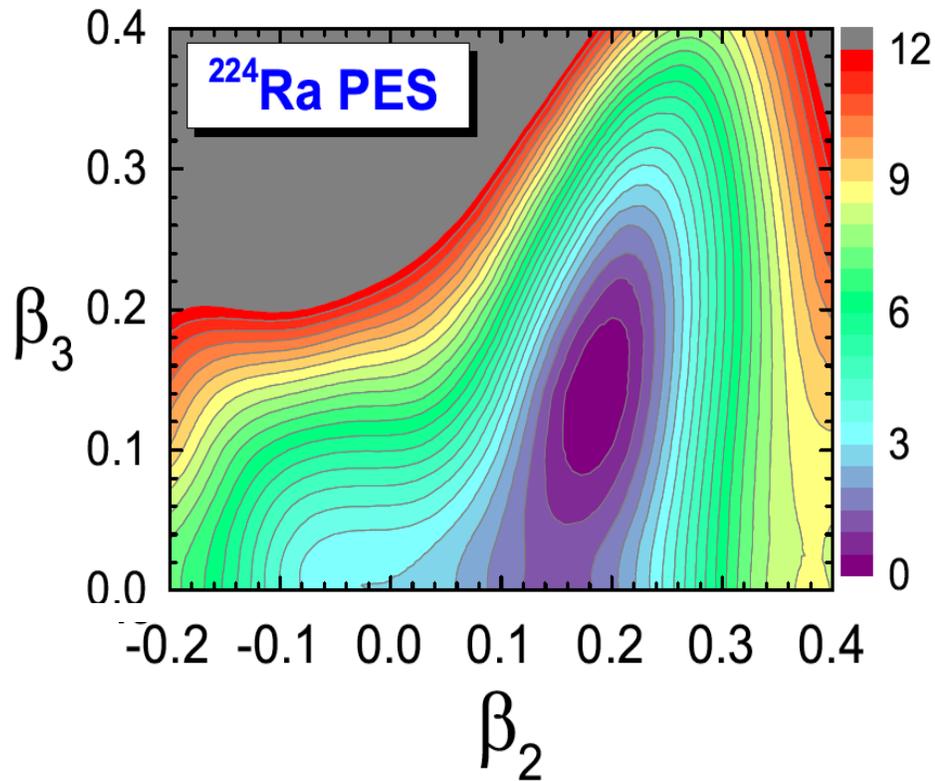
Outline



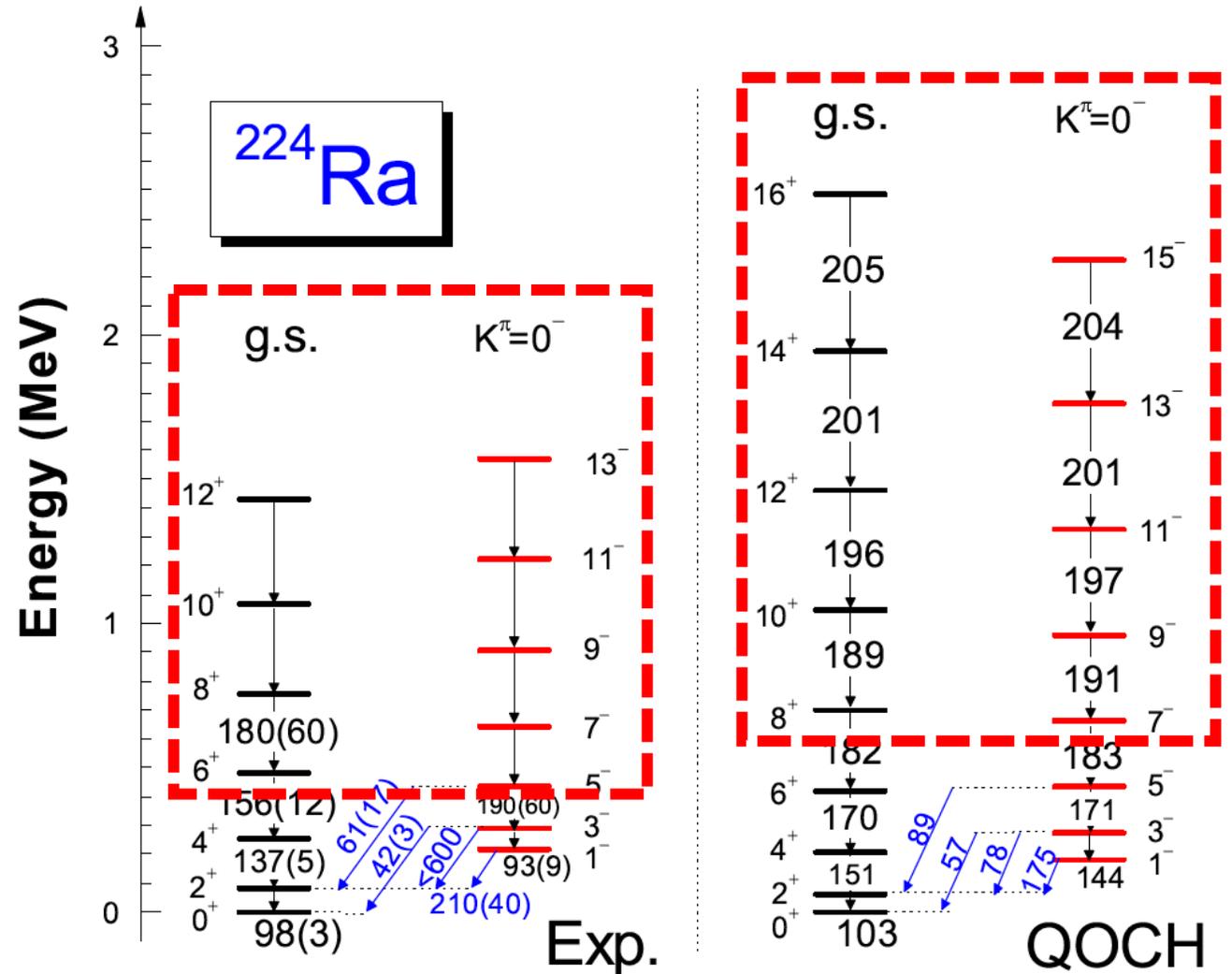
1. Introduction
2. Microscopic collective Hamiltonian
- 3. Results and discussion**
 - a. Quadrupole shape coexistence
 - b. Octupole shapes and parity doublets**
4. Summary and outlook

Octupole deformation in ^{224}Ra

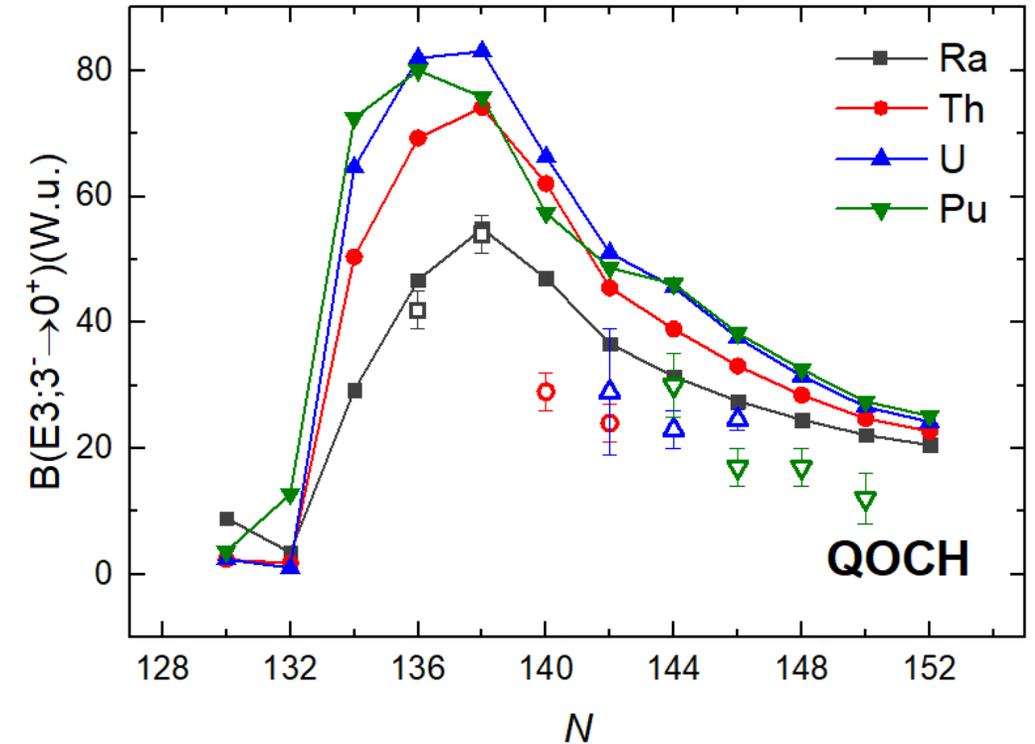
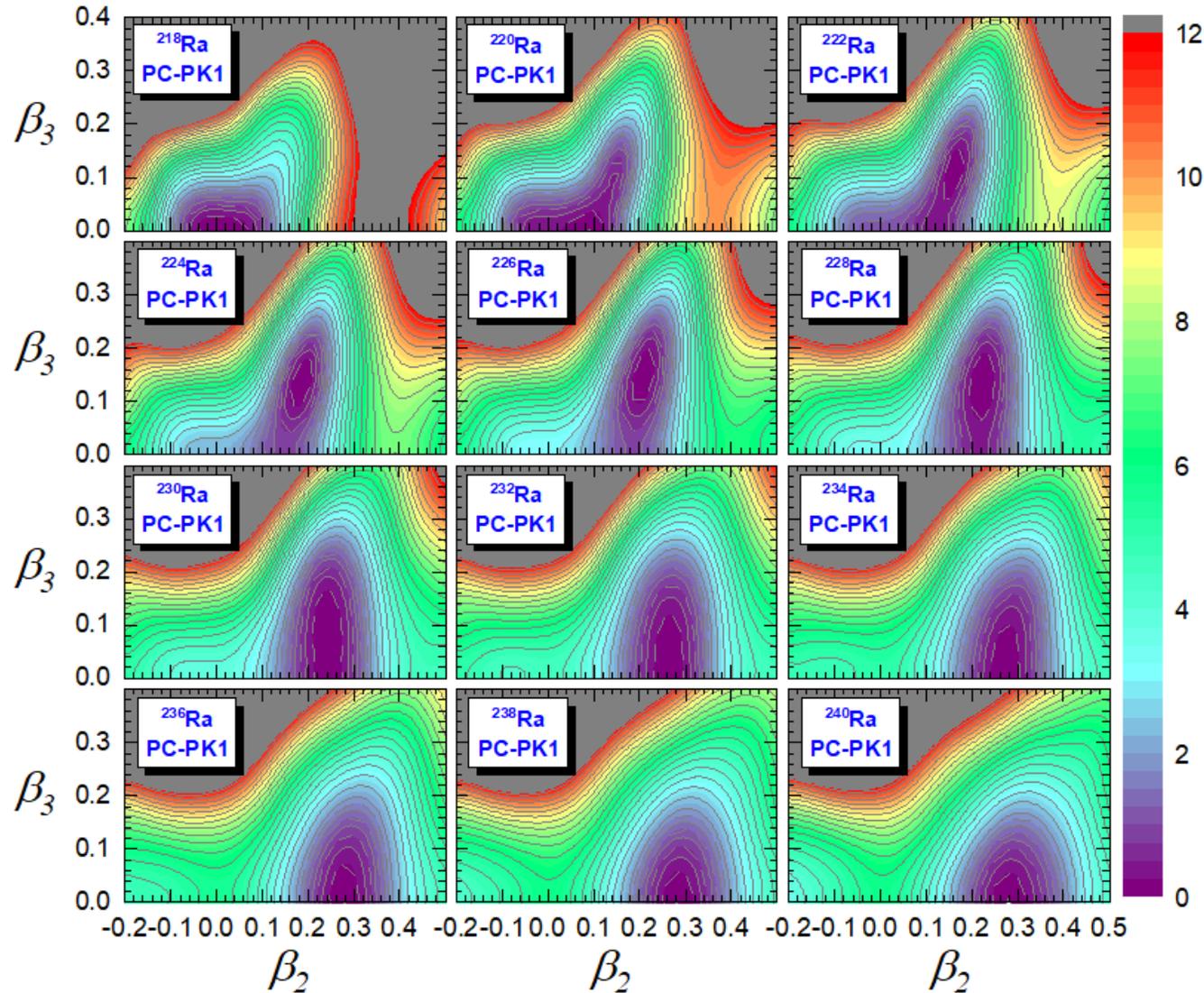
➤ PES of ^{224}Ra



➤ Low-lying spectrum of ^{224}Ra



Octupole shape transition

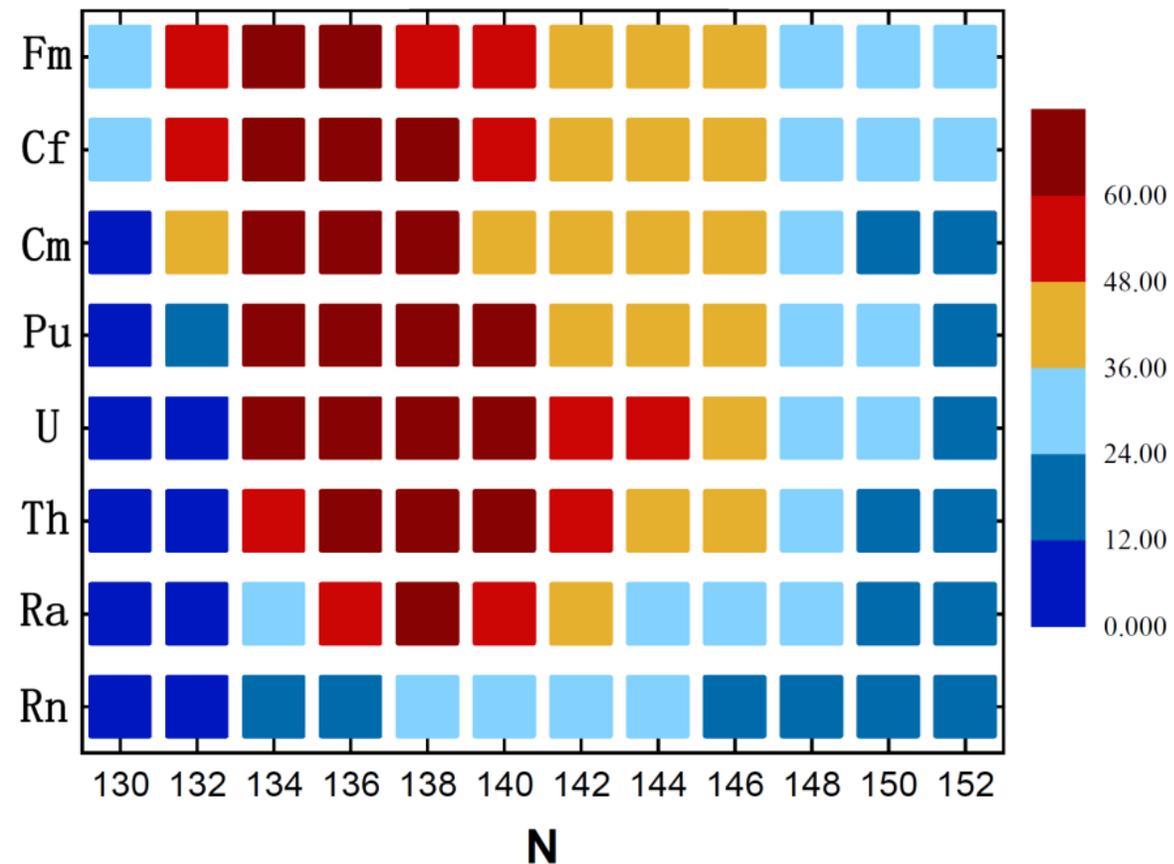
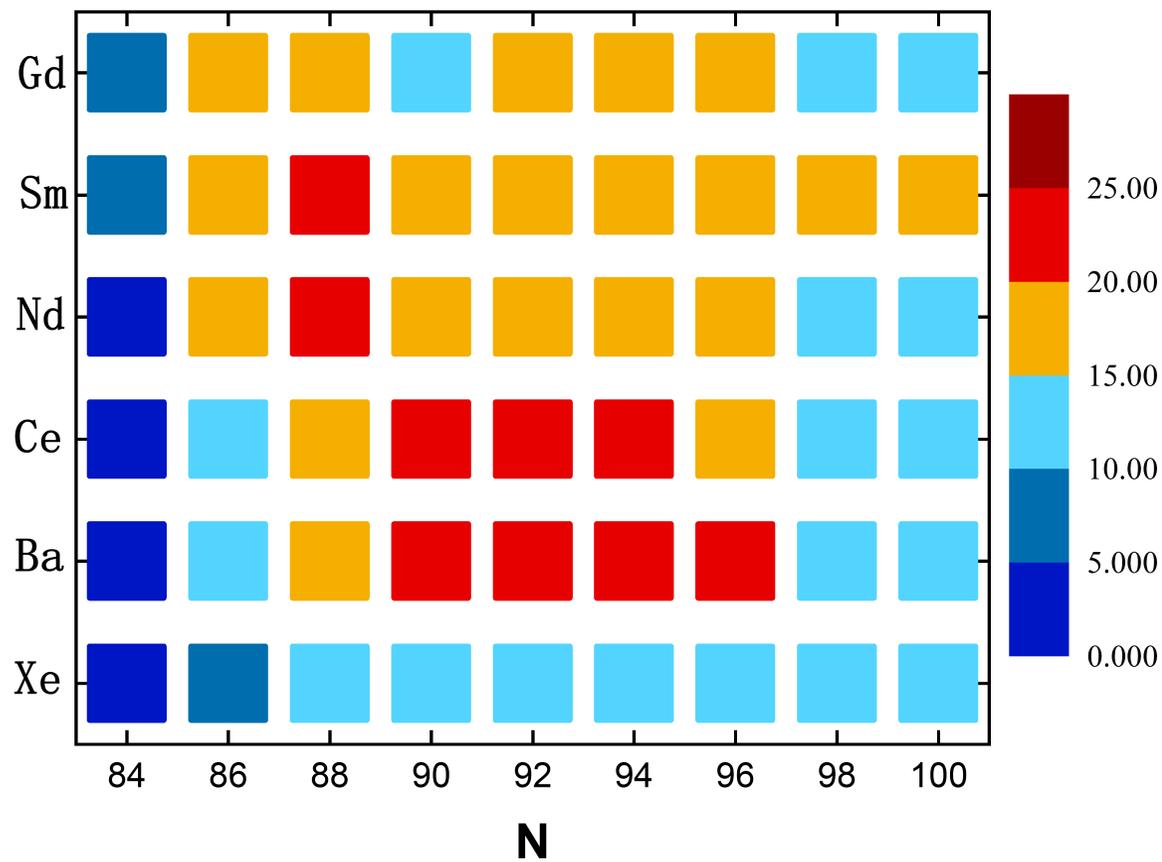


- **Ra, Th, U, Pu (N=130~152)**
near sph. → octupole def.
→ octupole soft → quad. def.

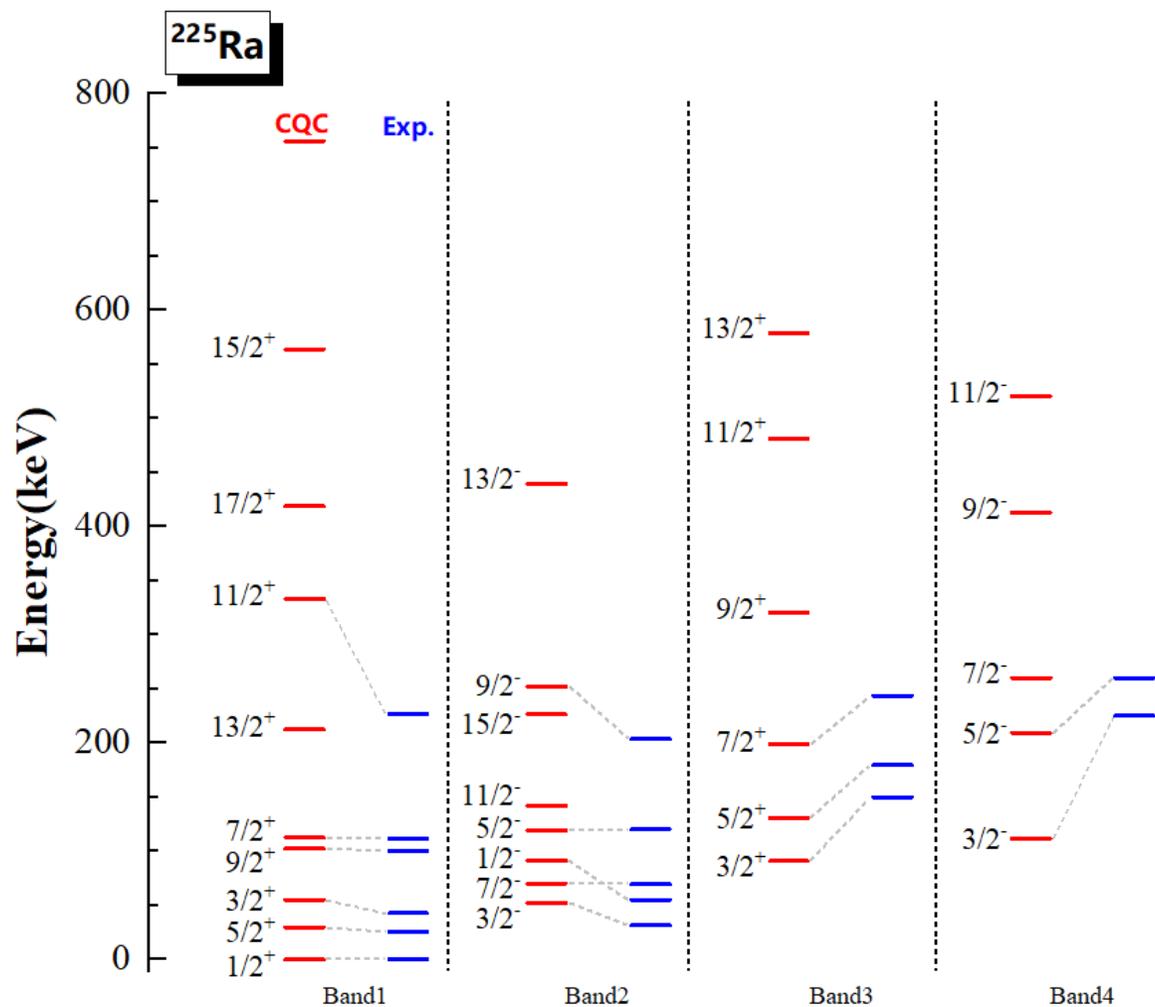
➤ **Largest B(E3) at N~136**

Octupole shape transition

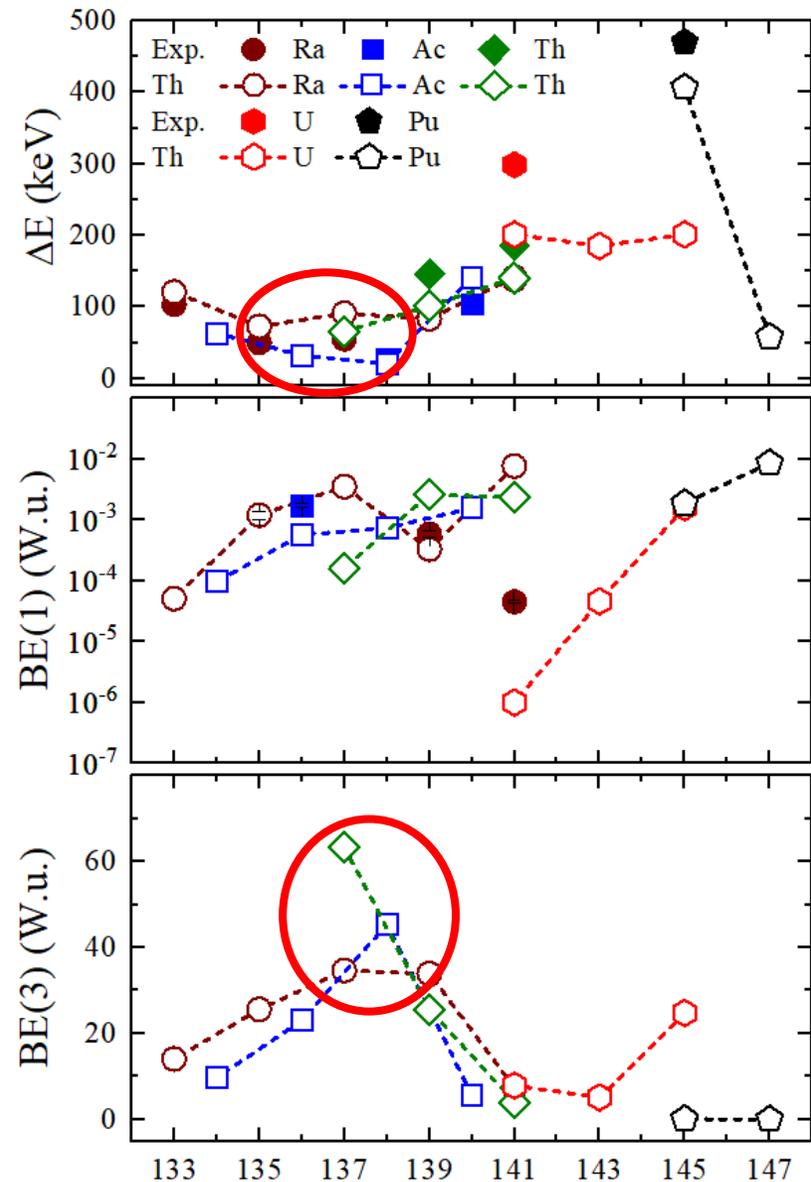
➤ $B(E3; 3^- \rightarrow 0^+)$



Parity doublets in odd-A Ra and actinides

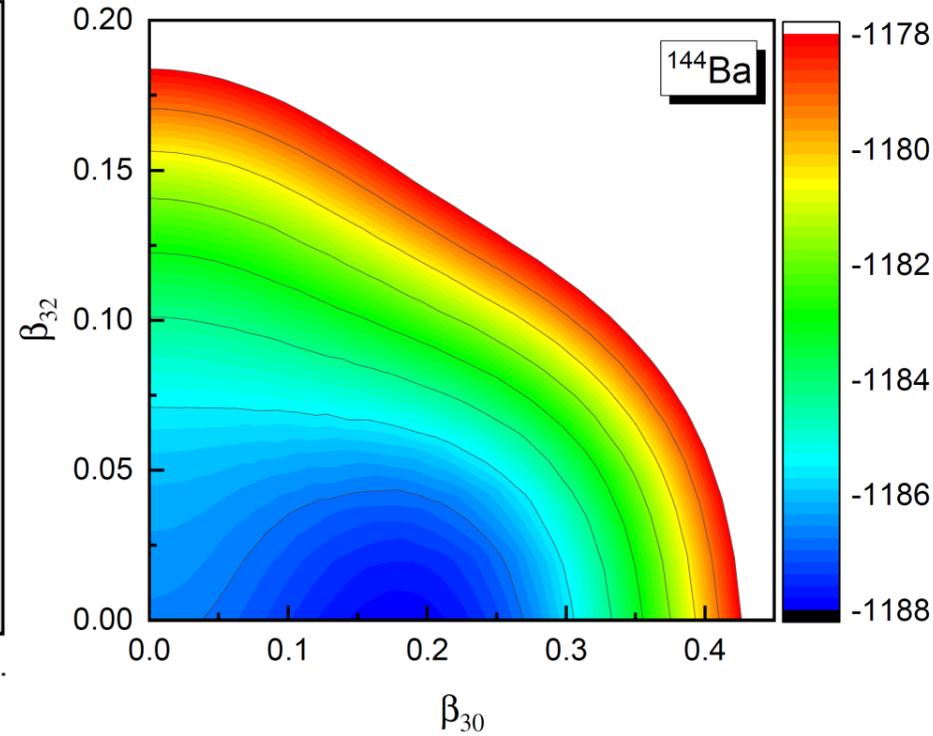
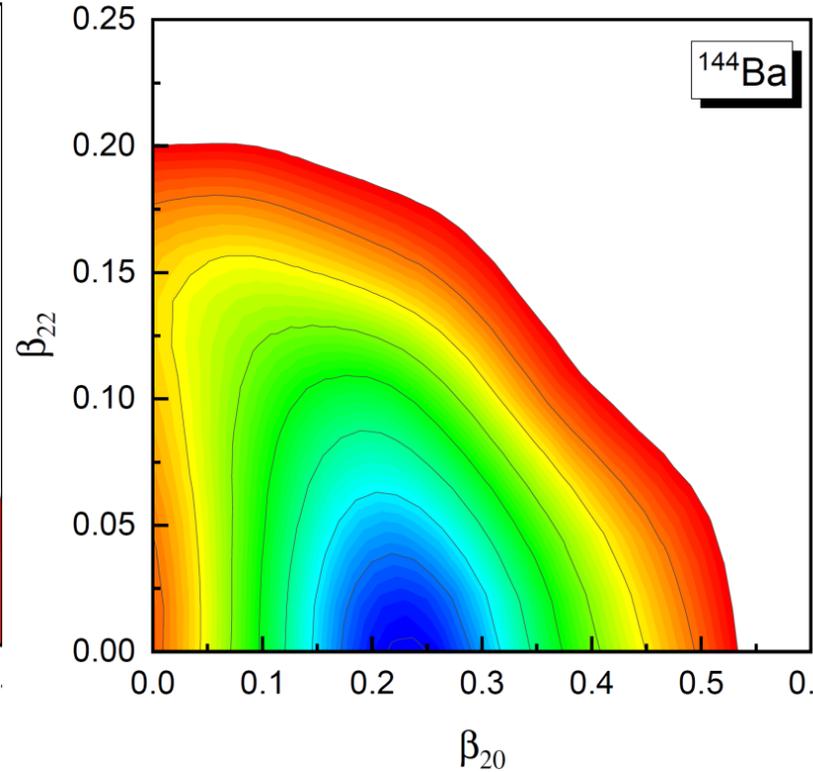
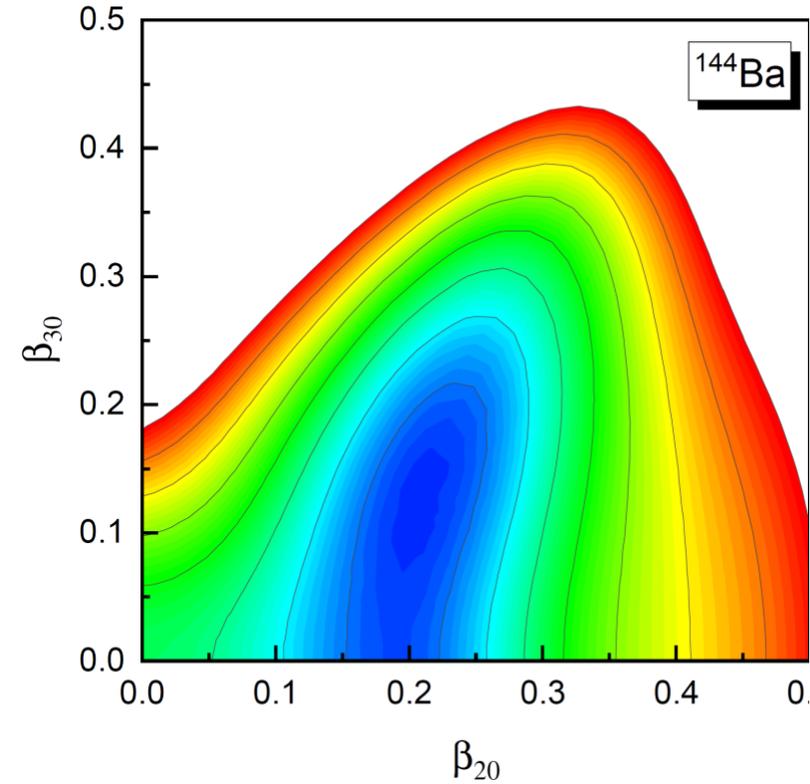


➤ Parity doublets in ^{225}Ra is reproduced.



PRL 118,
152504
(2017)

Triaxial octupole d.o.f.



➤ β_{30} minimum;

β_{22} : quite stiff

β_{32} : quite soft

➤ 7D collective Hamiltonian (7DCH) has been constructed, checking ...

◆ Summary

- Global description of nuclear deformations, low-lying states and $B(E2)$
- Using quadrupole shape invariants to analyze shape coexistence: reproduce the known regions and **predict new regions**
- Octupole shapes around **$N \sim 90$** and **$N \sim 136$** ; parity doublets in odd- A actinides.

◆ Outlook: triaxial octupole, chiral bands, ...

Thank you for your attention!

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Nuclear Physics Group

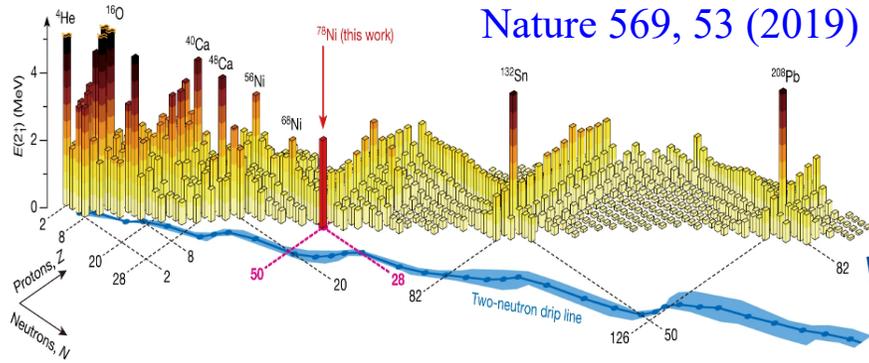
Peking University

Chongqing Normal Uni.

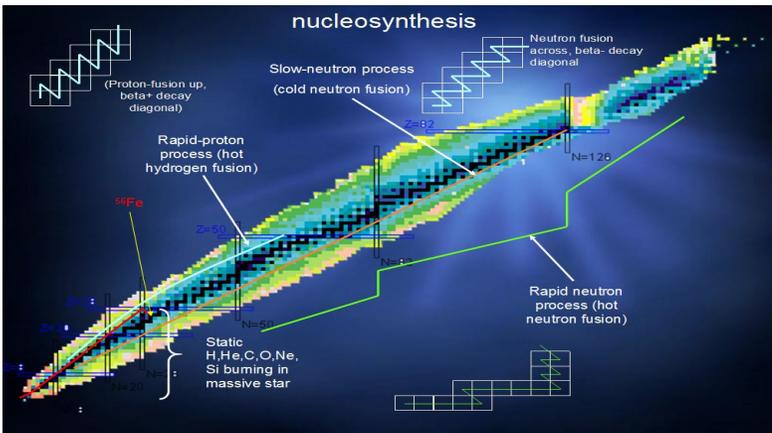
University of Zagreb

Southwest University

Nuclear shapes and excitations

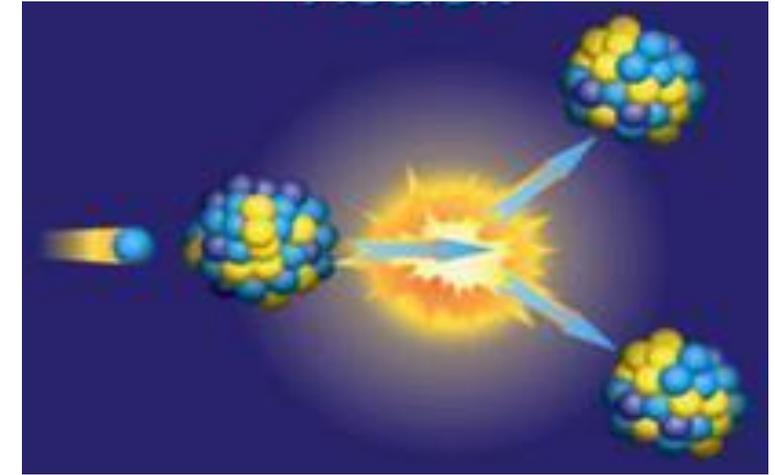


Structure: shell, shell evolution, effective interactions ...

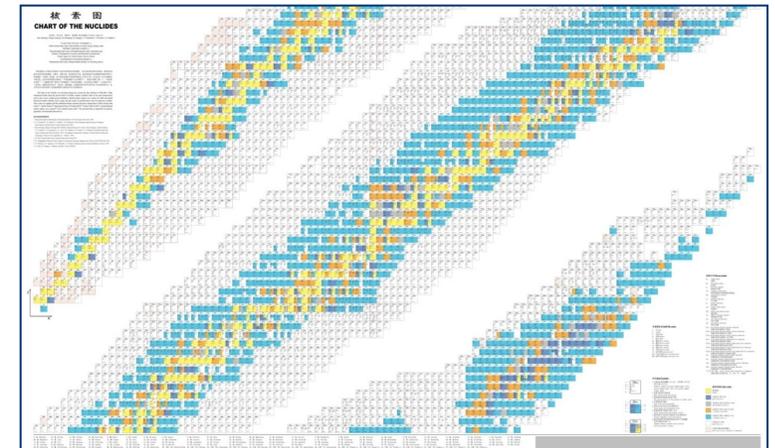


Astrophysics: r process, rp process, supernova ...

Nuclear shapes and excitations



Reaction: fission, reaction dynamics ...



Nuclear data, engineering 34