

UNIVERSAL ACCELERATING COSMOLOGIES FROM 10D SUPERGRAVITY

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based on [\[2210.10813\]](#) with Dimitrios Tsimpis



Motivations

- Account for the observed **accelerated expansion** of our (homogeneous + isotropic) universe within **string theory**
- **Aim:** construct solutions of 10d SUGRA where the 4d spacetime is **FLRW**,

$$ds_4^2 = -dt^2 + a(t)^2 \left(\frac{1}{1 - kr^2} dr^2 + r^2 d\Omega_2^2 \right) \Rightarrow R_{00} = -\ddot{a}$$

with good phenomenological properties so to model **dark energy/inflation**, and good control on the SUGRA solution as a low-energy limit of string theory

- **Problems:** even the “simplest” instances = de Sitter seem to be ruled out
Early 21st century: accelerating cosmologies from compactifications were thought to be as difficult as de Sitter

Strong Energy Condition and a no-go theorem

- **Strong Energy Condition:** $R_{00} = T_{00} + g^{ij}T_{ij} \geq 0$
 \Rightarrow needs to be violated for accelerated expansion

- Can be done with e.g. a **positive scalar potential**,

$$R_{00} = 2 \left(\dot{\phi}^2 - V(\phi) \right) \Rightarrow \text{acceleration whenever } E_p > 2E_k$$

\Rightarrow *All we need is a compactification for which the 4d EFT has a positive scalar potential?*

- **Problem:** (1) all 10d/11d pure supergravities have their bosonic EMT satisfying the SEC, (2) **the SEC is hereditary** [Gibbons '84, Maldacena, Nuñez '00]

Strong Energy Condition and a no-go theorem

- Cosmic acceleration is forbidden in models descending from 10d/11d SUGRAs, with **time-indep.** compactifications on **non-singular** manifolds **without boundary**

$$ds_{10}^2 = H(y)^2 ds_4^2 + ds_6^2 \quad \Rightarrow \quad R_{00}^{(10)} = R_{00}^{(4)} - \frac{1}{4} H^{-2} \nabla_{\mathcal{M}}^2 H^4$$

$$\left(\int_{\mathcal{M}} H^2 \right) R_{00}^{(4)} = \int_{\mathcal{M}} H^2 R_{00}^{(10)} \quad \Rightarrow \quad \boxed{R_{00}^{(10)} \geq 0 \Rightarrow R_{00}^{(4)} \geq 0}$$

- **Evade the no-go:** (1) SUGRA is not the relevant starting point, (2) \mathcal{M} is not compact, (3) Include brane/orientifold singularities, (4) \mathcal{M} **is time-dependent**
- (4) \Rightarrow **Transient acceleration** is in fact generic in flux compactifications, [Townsend, Wohlfarth '03] although de Sitter space is still ruled out
Late-time acceleration is *not* ruled out [Russo, Townsend '03]
- Still have to engineer a **positive scalar potential**

Positive scalar potential, recipe

- Fluxes: $V_{\text{fluxes}} \sim \sum_p \frac{1}{2^p p!} \int \sqrt{g_6} e^{a_p \phi} |F_p|^2$
- Internal curvature: $V_{\text{curvature}} \sim - \int \sqrt{g_6} R^{(6)}$
- Brane/orientifold sources: $V_{\text{source}} \sim -\text{tension} \times \int \sqrt{g_n}$

$$V_{\text{scalar}} = V_{\text{fluxes}} + V_{\text{curvature}} + V_{\text{source}}$$

The set-up

- **10d massive Type IIA SUGRA with fluxes,**

$$S = \frac{1}{4\kappa^2} \int d^{10}x \sqrt{g} \left(-2R^{(10)} + |\partial\phi|^2 + \frac{1}{2!} e^{3\phi/2} |F_2|^2 + \frac{1}{3!} e^{-\phi} |H|^2 + \frac{1}{4!} e^{\phi/2} |F_4|^2 + m^2 e^{5\phi/2} \right) + \text{Chern-Simons}$$

- **Cosmological ansatz** = fields only depend on time
- Ansatz which “solves the fluxes” [Terrisse, Tsimpis ‘19]
- Ansatz for the 10d metric, two time-dep. **warp factors,**

$$ds_{10}^2 = e^{2A(t)} \left(e^{2B(t)} ds_{\text{FLRW}}^2 + \underbrace{g_{mn} dy^m dy^n}_{\text{Calabi-Yau, Einstein, Einstein-Kähler, ...}} \right)$$

Calabi-Yau, Einstein, Einstein-Kähler, ...

- In 4d Einstein frame,

$$ds_{4E}^2 = -dt^2 + a(t)^2 d\Omega_k^2, \quad a(t) = e^{4A(t)+B(t)}$$

Equations of motion

$$d_\tau^2 A = -\frac{1}{48} (\partial_A U - 4\partial_B U)$$

$$d_\tau^2 B = \frac{1}{12} (\partial_A U - 3\partial_B U)$$

$$d_\tau^2 \phi = -\partial_\phi U$$

$$U(A, B, \phi) = 72(d_\tau A)^2 + 6(d_\tau B)^2 + 48d_\tau A d_\tau B - \frac{1}{2}(d_\tau \phi)^2$$

- The “potential” U encodes all the info about the **fluxes** and internal **geometry**,

$$U(A, B, \phi) = \sum_i c_i e^{\alpha_i A + \beta_i B + \gamma_i \phi}$$

- Changing flux configuration/internal geometry \Leftrightarrow tuning constants $c_i, \alpha_i, \beta_i, \gamma_i$
- Equations cannot be solved analytically except for a few particular cases

Consistent truncations

Equations can be integrated in:

- In all cases, a **1d action**,

$$S_{1d} = \int d\tau \left\{ \frac{1}{\mathcal{N}} \left(-72(d_\tau A)^2 - 6(d_\tau B)^2 - 48d_\tau A d_\tau B + \frac{1}{2}(d_\tau \phi)^2 \right) - \mathcal{N}U(A, B, \phi) \right\}$$

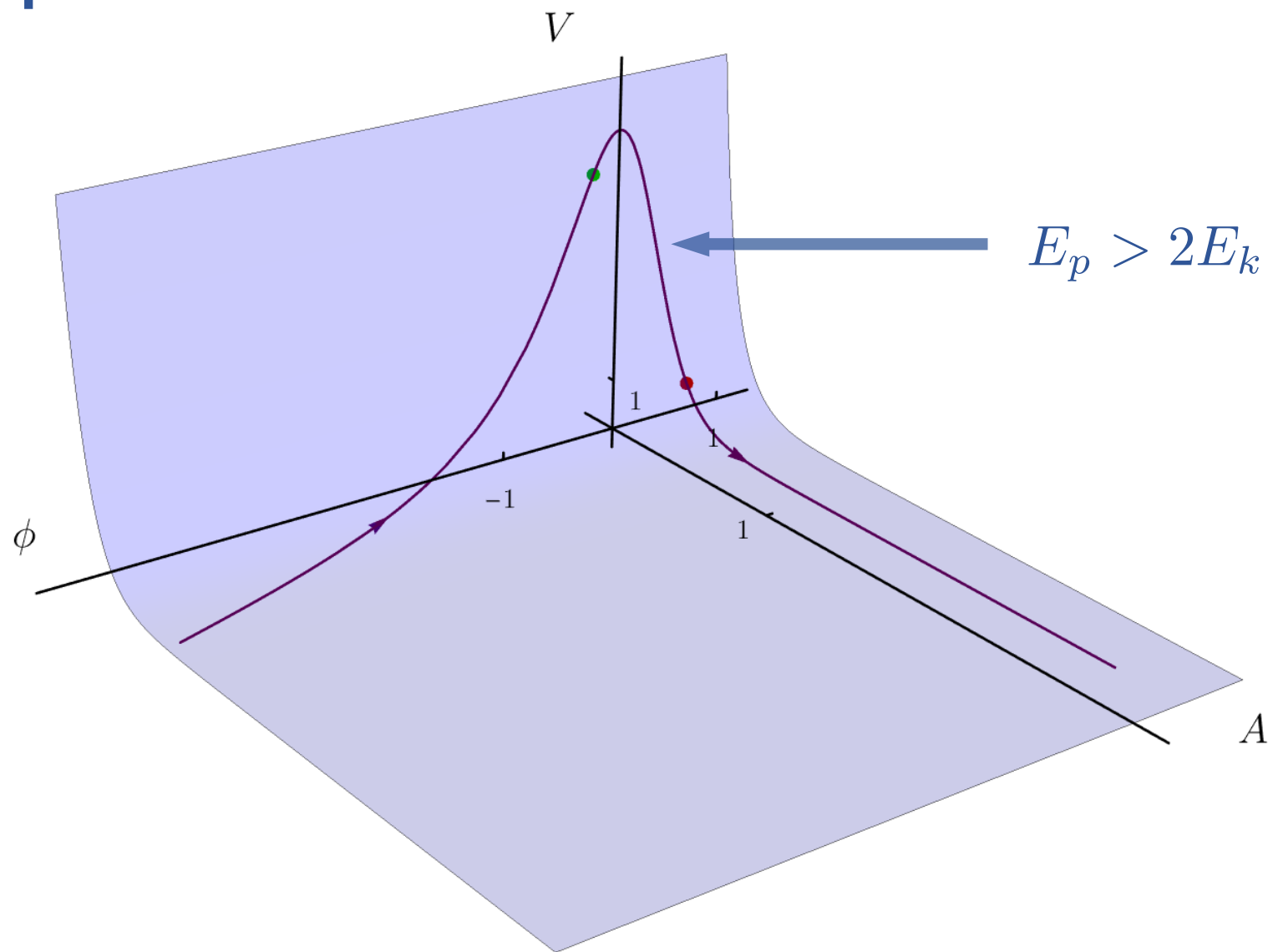
- In most cases, a **4d action**,

$$S_{4d} = \int d^4x \sqrt{g} \left(R^{(4)} - \frac{1}{2}g^{\mu\nu} \partial_\mu A \partial_\nu A - \frac{1}{2}g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - V(A, \phi) \right)$$

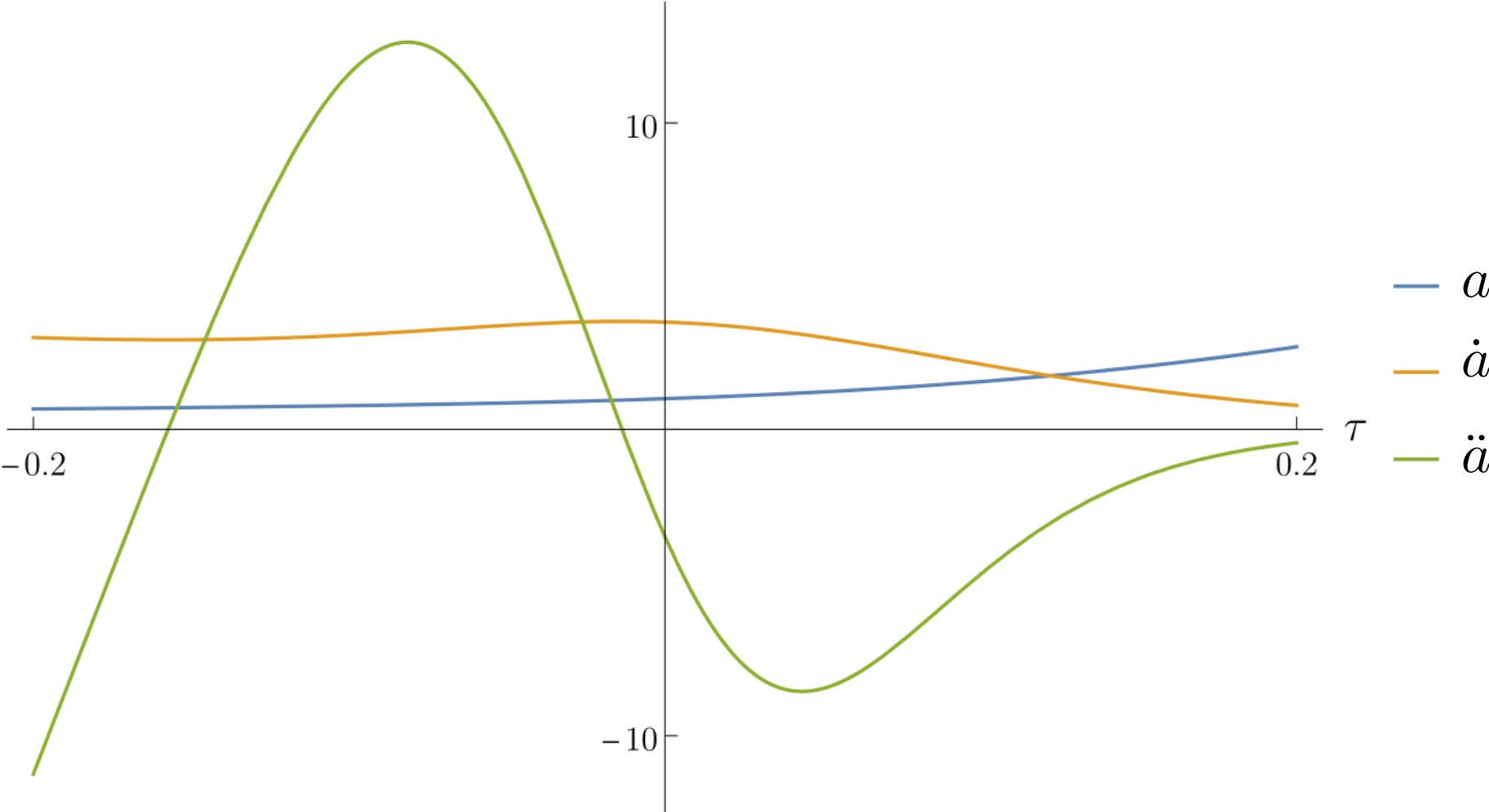
- **Two-scalar cosmology with multi-exponential potential**

The B -field dependency is hidden in the scale factor, in the metric

Field space picture



Transient acceleration



Dynamical system analysis

For a two-exponential potential, $U(A, B, \phi) = c_1 e^{\alpha_1 A + \beta_1 B + \gamma_1 \phi} + c_2 e^{\alpha_2 A + \beta_2 B + \gamma_2 \phi}$

E.O.M.'s can be recast into an **autonomous dynamical system**,

$$x' = \frac{1}{4} \left([\alpha_2 + 2\beta_2(-2 + x)](-1 + x^2 + y^2 + z^2) + [-\alpha_1 - 2\beta_1(-2 + x)]z^2 \right)$$

$$y' = \frac{1}{2} \left((2\sqrt{3}\gamma_2 + \beta_2 y)(-1 + x^2 + y^2 + z^2) - (2\sqrt{3}\gamma_1 + \beta_1 y)z^2 \right)$$

$$z' = \frac{1}{4} z \left(\alpha_1 x + 4\sqrt{3}\gamma_1 y - 2\beta_1(-1 + 2x + z^2) + 2\beta_2(-1 + x^2 + y^2 + z^2) \right)$$

with the constraint, $c_1(1 - x^2 - y^2 - z^2) = c_2 z^2 e^{A(\alpha_2 - \alpha_1) + B(\beta_2 - \beta_1) + \phi(\gamma_2 - \gamma_1)}$

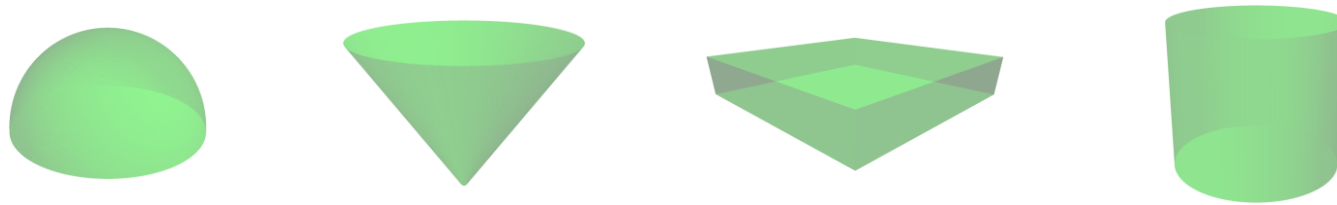
$c_1, c_2 > 0 \Rightarrow$ **the phase space is the unit-ball**

Dynamical system analysis, interpretation

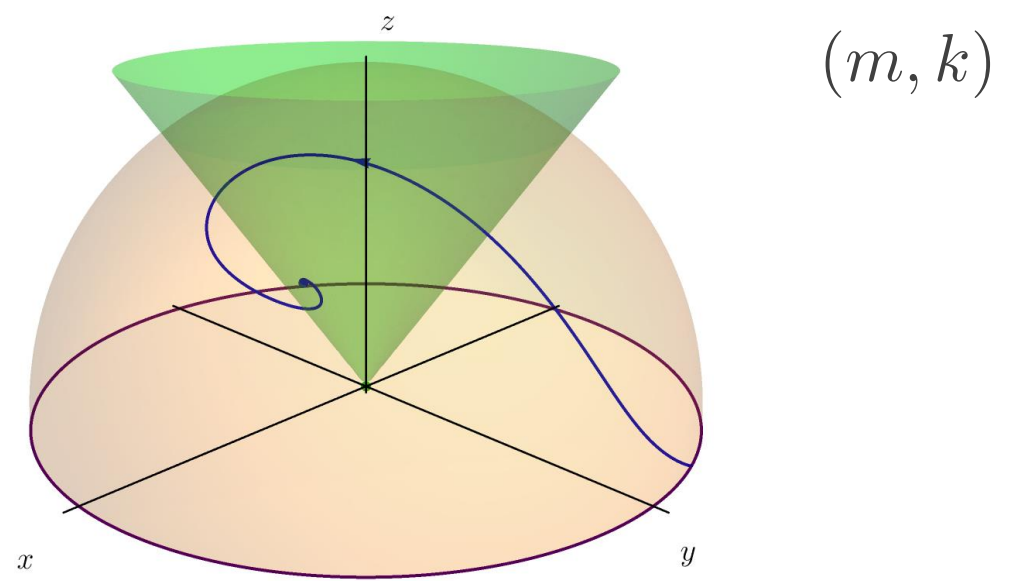
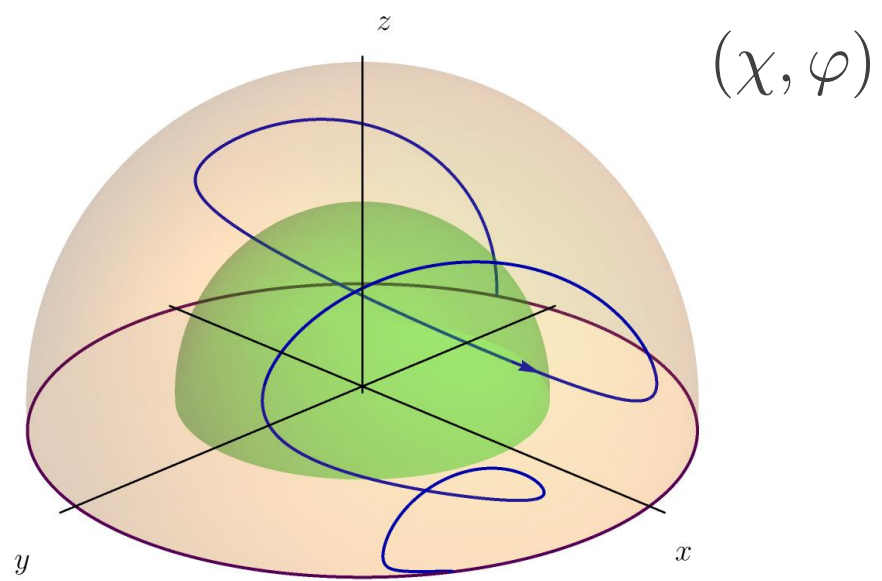
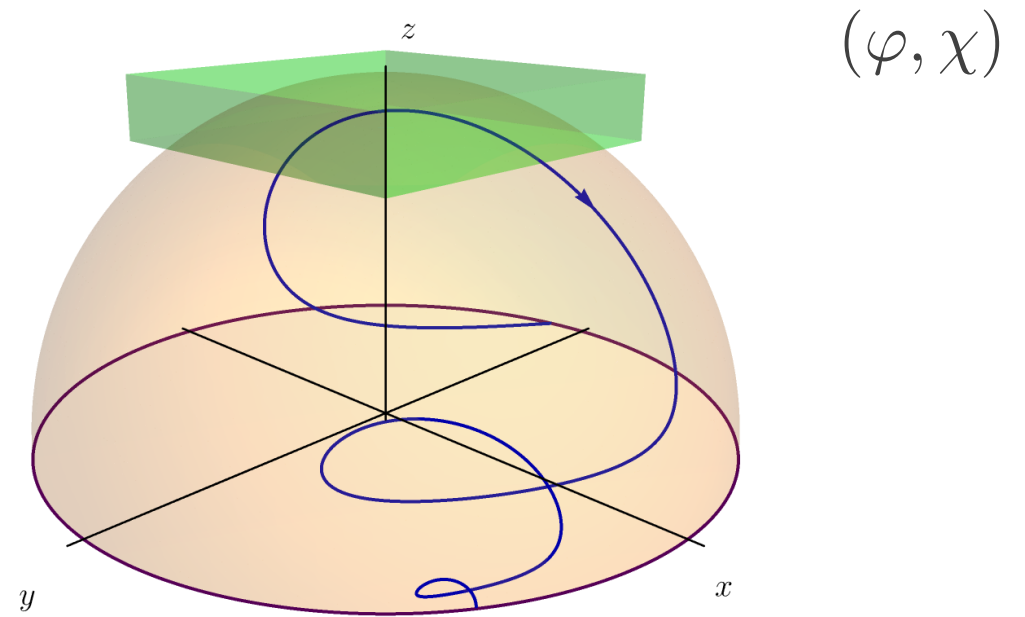
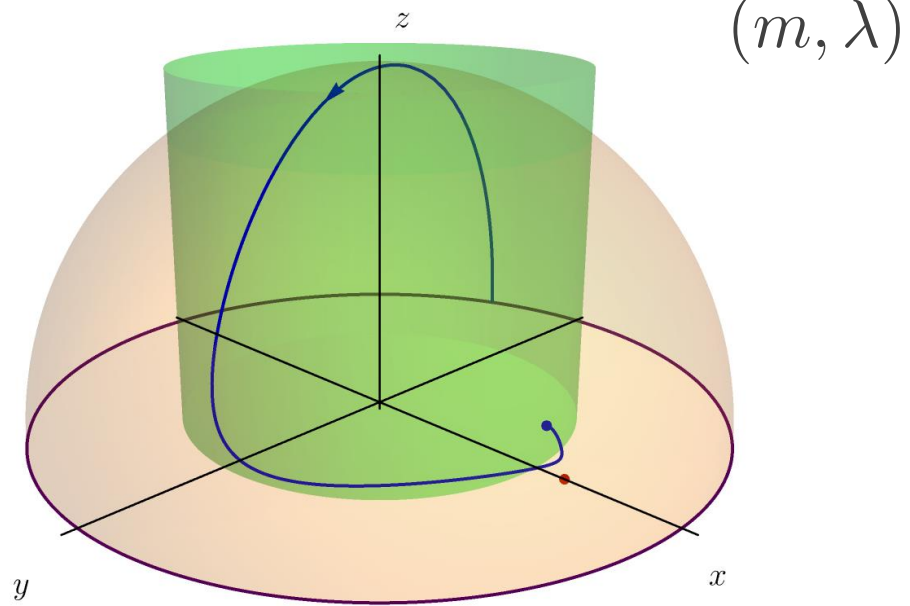
- Condition for **expansion**, $z > 0 \Leftrightarrow$ North hemisphere
- Condition for **acceleration**,

$$\ddot{a} > 0 \quad \Leftrightarrow \quad (\beta_1 - \beta_2)z^2 - \beta_2(x^2 + y^2) + \beta_2 - 4 > 0$$

defines a geometrical **acceleration region** in phase space,



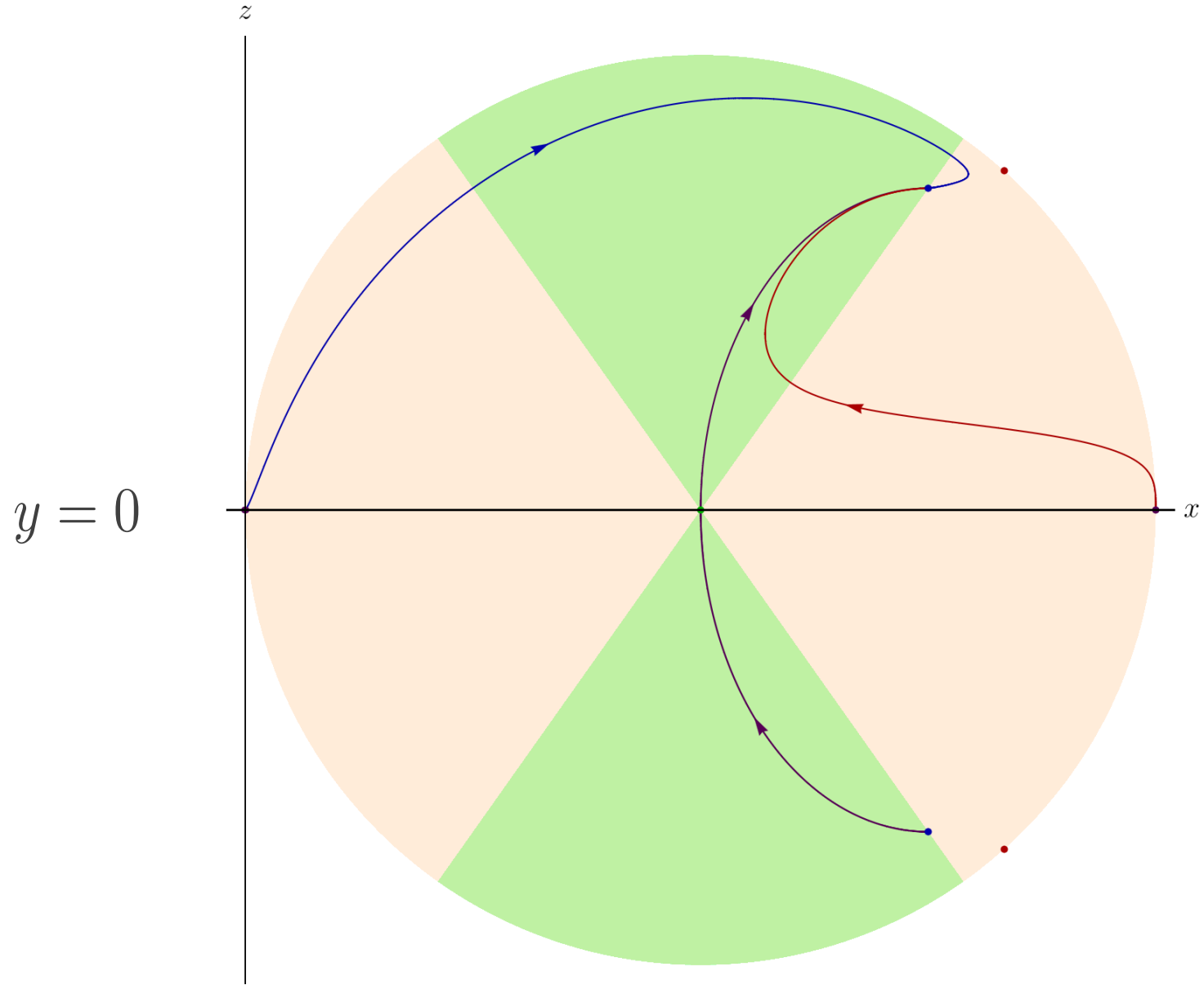
- **Number of e folds**, $N = \int_{\text{in}}^{\text{out}} d\omega = \omega_{\text{out}} - \omega_{\text{in}}$
- **Fixed points = power law expansions**, $a(t) = t^p$, $p \leq 1$
- Generic trajectories interpolate between fixed points

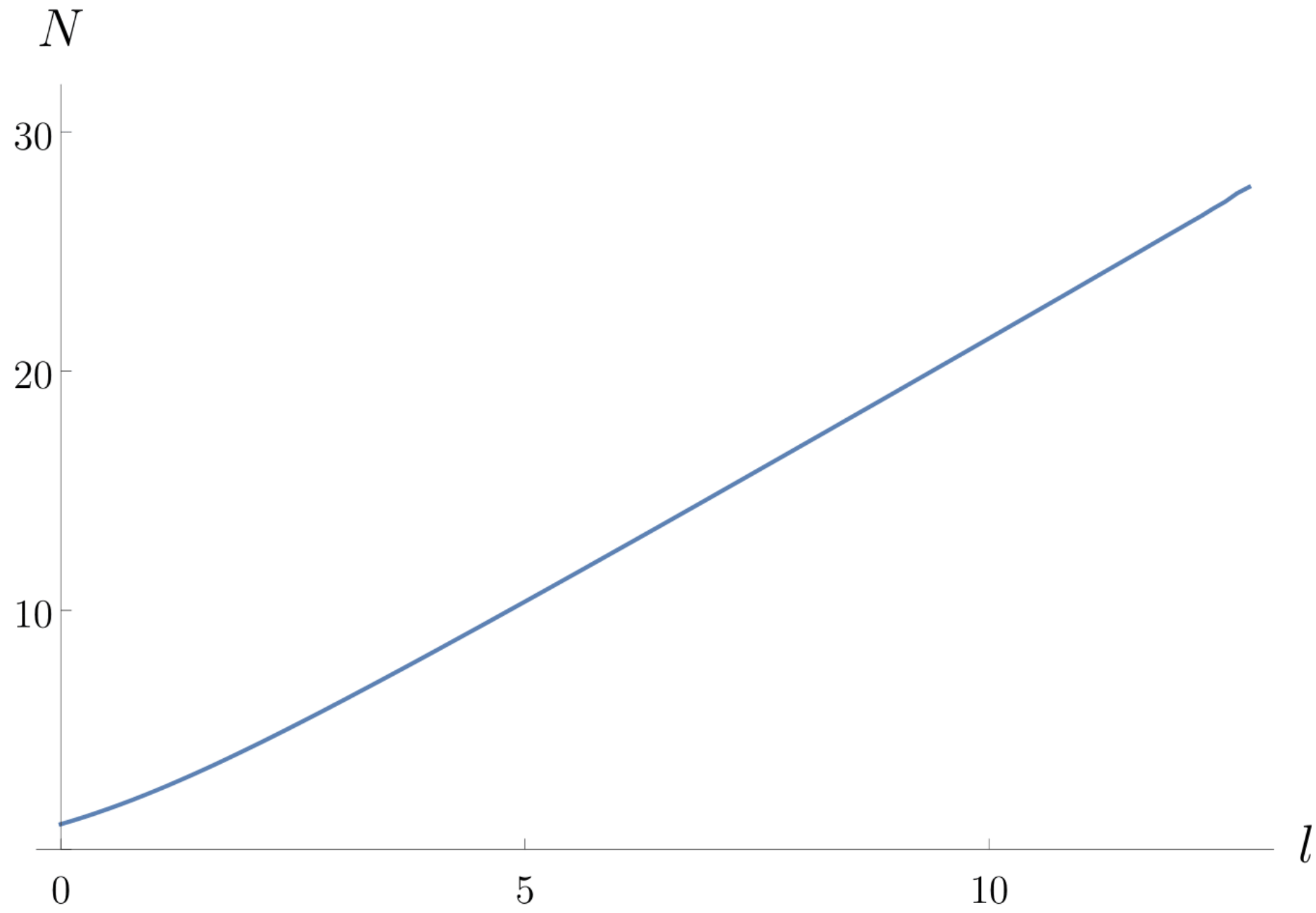


Results

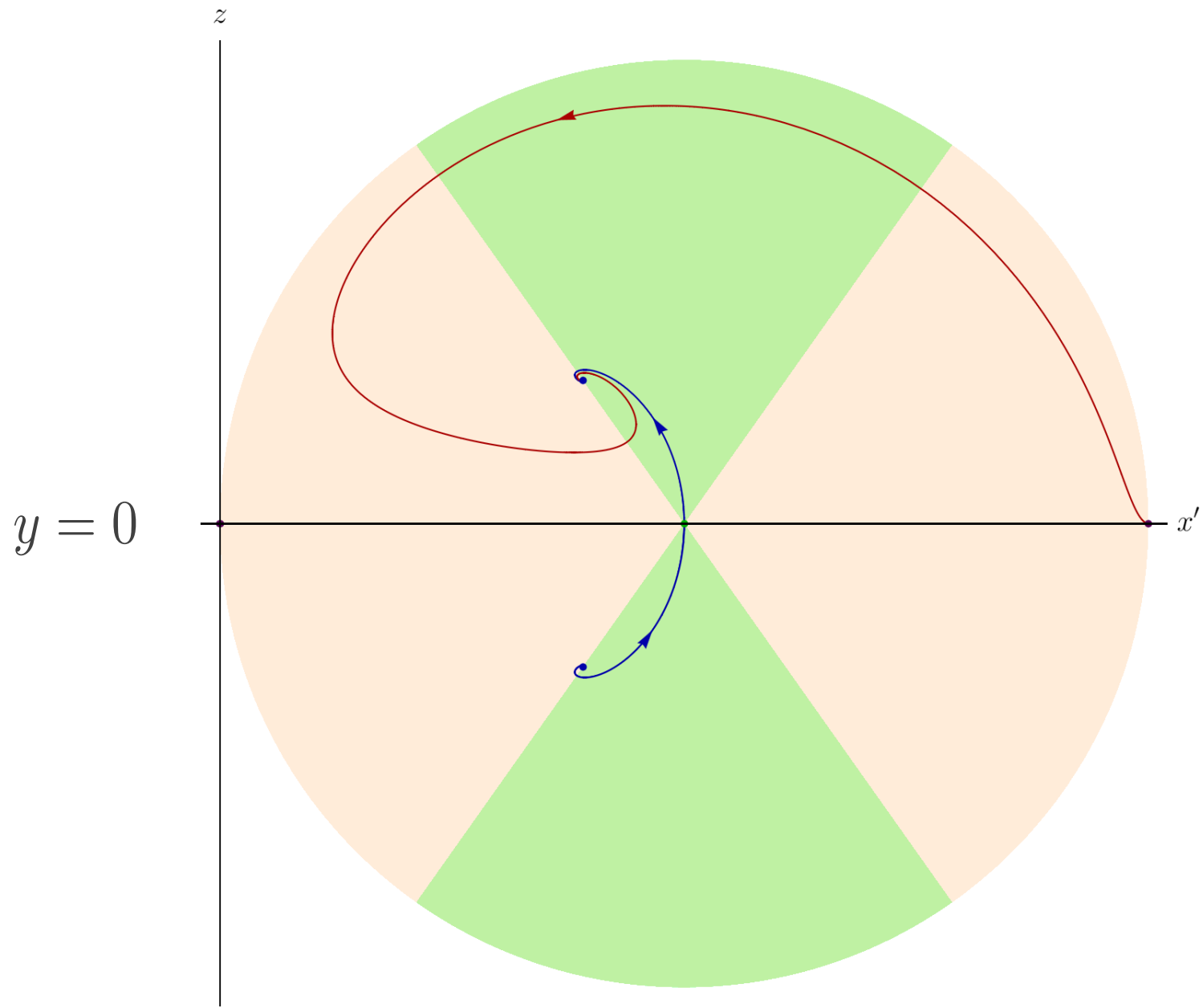
- Continuum of **transient accelerations** with **parametric control on the number of e-folds**, by tuning initial conditions
 - **Semi-eternal inflation** whenever there is a fixed point lying on the boundary of the acceleration region
 - **Eternal inflation** (unique trajectory) whenever there is also a fixed point at the origin of phase space: becomes **de Sitter** around the origin, can be geodesically completed in the past by gluing its mirror in the southern hemisphere
 - Alternating phases of acceleration/deccelaration with **spiraling trajectories** around a fixed point on the boundary
- ⇒ Can be used for “rollercoaster” models of inflation [D’Amico, Kaloper ‘20]

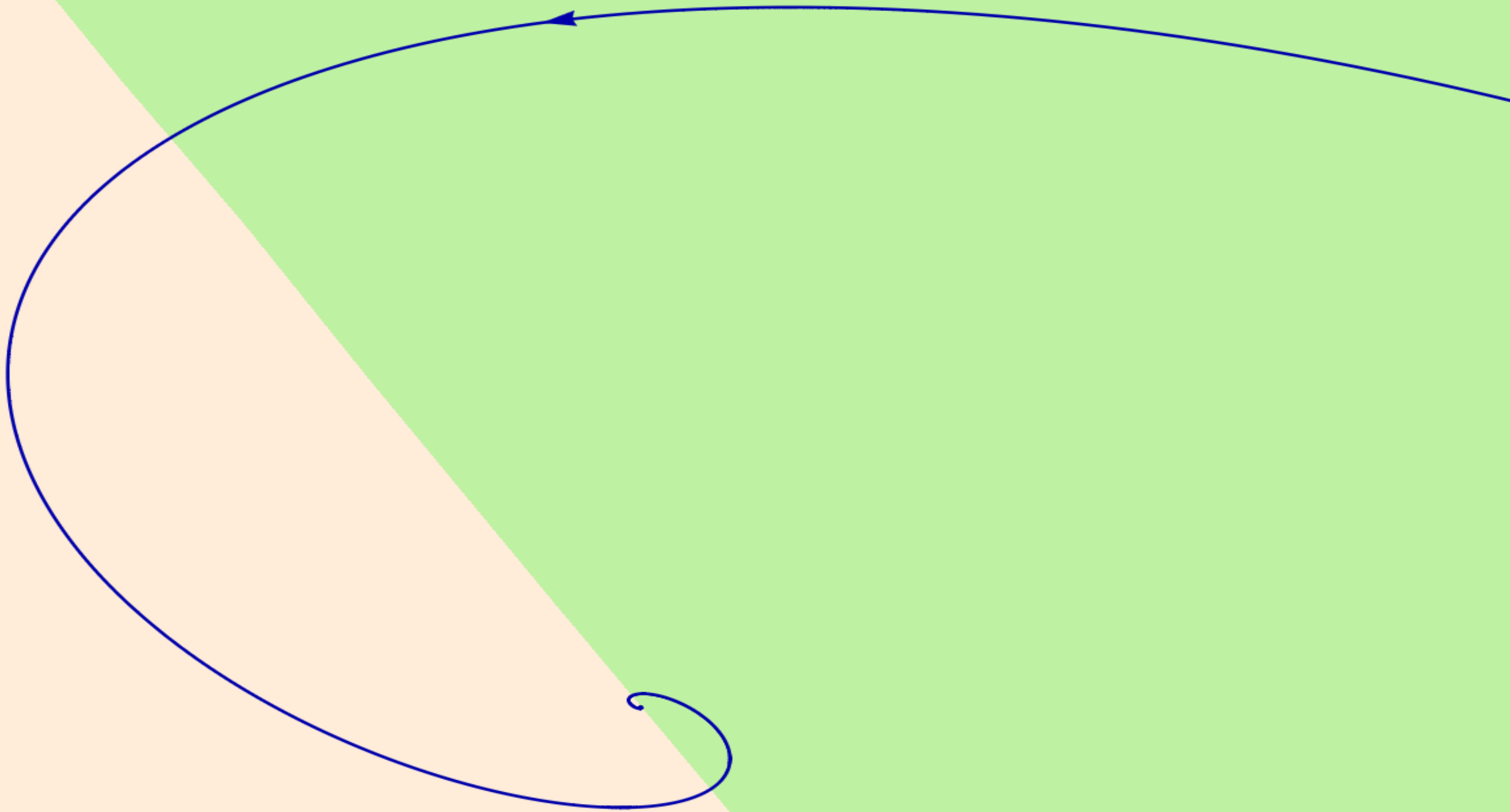
(λ, k)



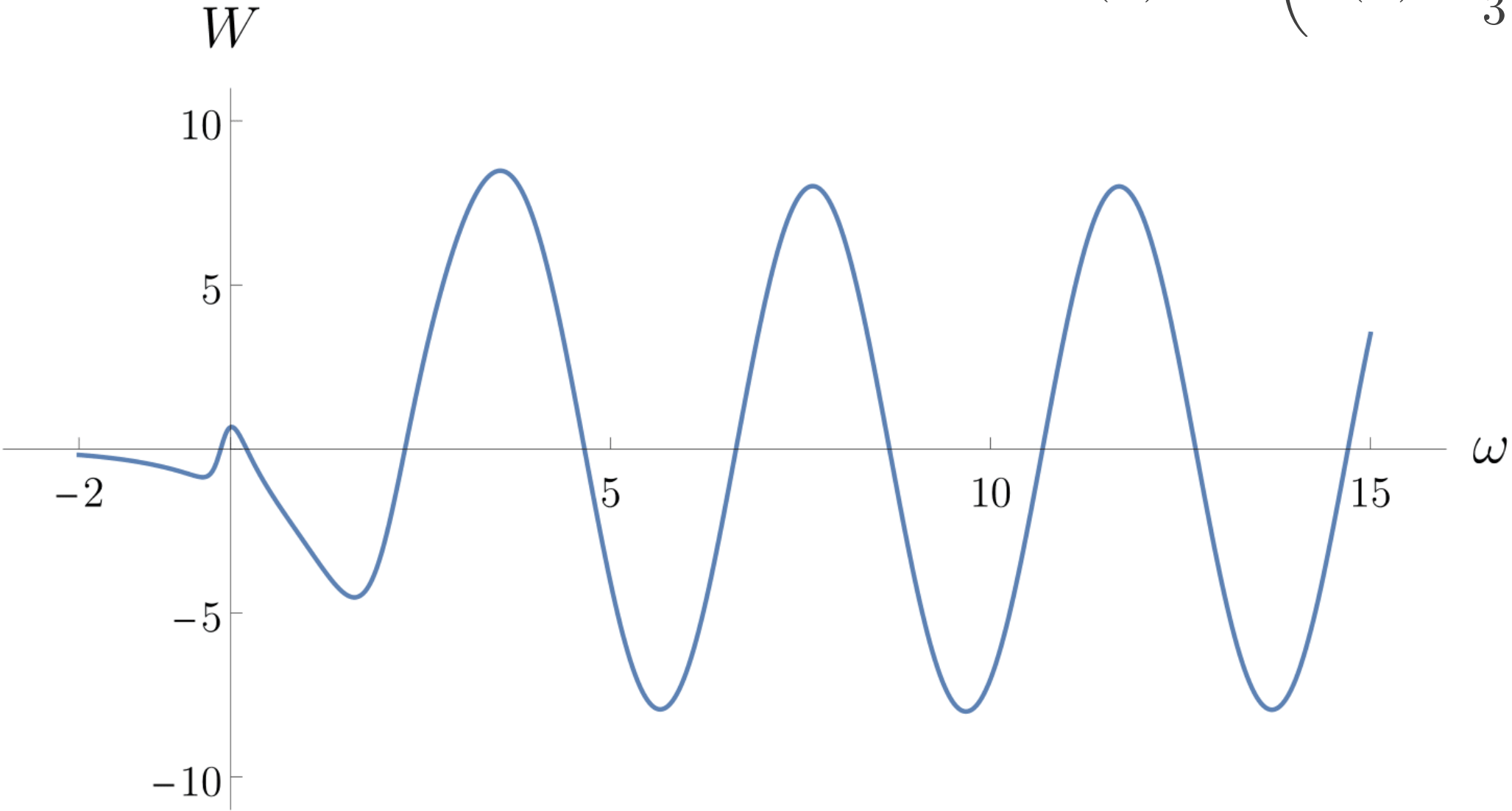


(m, k)





$$W(\omega) = -\left(w(\omega) + \frac{1}{3}\right) e^{\omega}$$



Conclusion

- **Transient acceleration** is generic in flux compactifications
- Cosmologies featuring **(semi-) eternal acceleration**, or a **parametric control** on the number of e-folds also seem **generic!**
They have $k = -1$ and asymptotically vanishing acceleration
- Existence of “**rollercoaster**” **cosmologies** with infinite number of cycles alternating between accelerated and decelerated expansion

Outlook

- Extend the dynamical system analysis **beyond 2 fluxes**
- Inclusion of **orientifolds**, dynamical system analysis of **dS fixed points**
- Study the **phenomenological properties** of these models beyond the mere number of e-folds. Can one hope for **realistic cosmologies?**

THANK YOU!