

**REVIEW TALK:  
REALISTIC VACUA IN THE STRING THEORY LANDSCAPE**

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# REALISTIC STRING VACUA

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*Any realistic string theory vacuum should have (at least):*

- four macroscopic spacetime dimensions (obviously)
- broken / no supersymmetry
- dark energy / positive cosmological constant
- Standard Model matter (gauge groups, chiral fermions, ...)

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# KNOWN STRING VACUA

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*String vacua that we understand well have:*

- extended ( $\mathcal{N} \geq 2$ ) supersymmetry
- negative or vanishing cosmological constant (AdS or Mink.)

*side note:*

*SUSY breaking and positive vacuum energy (e.g. de Sitter) are related  
(no SUSY algebra with unitary representations in de Sitter)*

Unknown whether string theory has stable  $\left\{ \begin{array}{l} \text{non-SUSY vacua!} \\ \text{de Sitter vacua!} \end{array} \right.$

# NON-SUPERSYMMETRIC STRING THEORY

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➤ Bosonic string

- Target space **tachyon!**

➤ Type 0 string

- Target space **tachyon!**

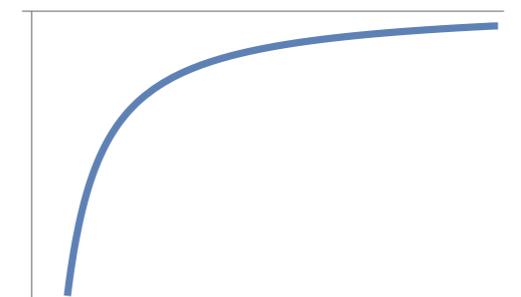
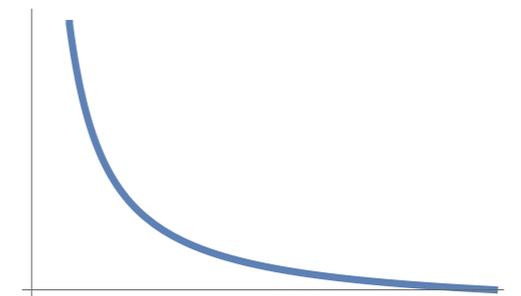
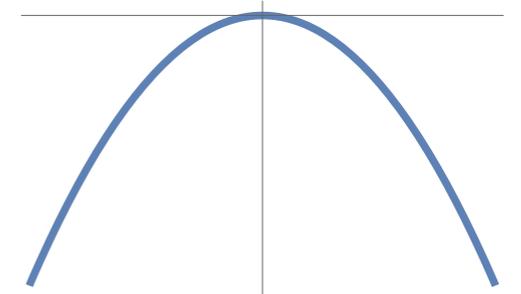
➤  $O(16) \times O(16)$  Heterotic string

- String frame: positive cosmological constant
- Einstein frame:  $V \sim e^{-5\phi/2}$  (**run-away!**)

➤ Scherk-Schwarz supersymmetry breaking

- anti-periodic fermion boundary conditions on circle

- Potential for radius:  $V \sim -\frac{1}{R^\alpha}$  (**run-away!**)



# DINE-SEIBERG PROBLEM

[Dine, Seiberg '85]

- .....
- Fundamental problem of string compactifications:

Moduli!  
(e.g. dilaton, comp. volume, ...)

=

*massless scalar fields  
at tree (classical) level*

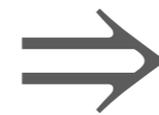
Broken Supersymmetry:

- Quantum effect: generate a potential for moduli!

*assume:*

$\phi \rightarrow \infty$ :

*weakly coupled regime, SUSY restored,  
effective tree-level description valid*



$$\lim_{\phi \rightarrow \infty} V = 0$$

# DINE-SEIBERG PROBLEM

[Dine, Seiberg '85]

.....

potential from first order quantum corrections:

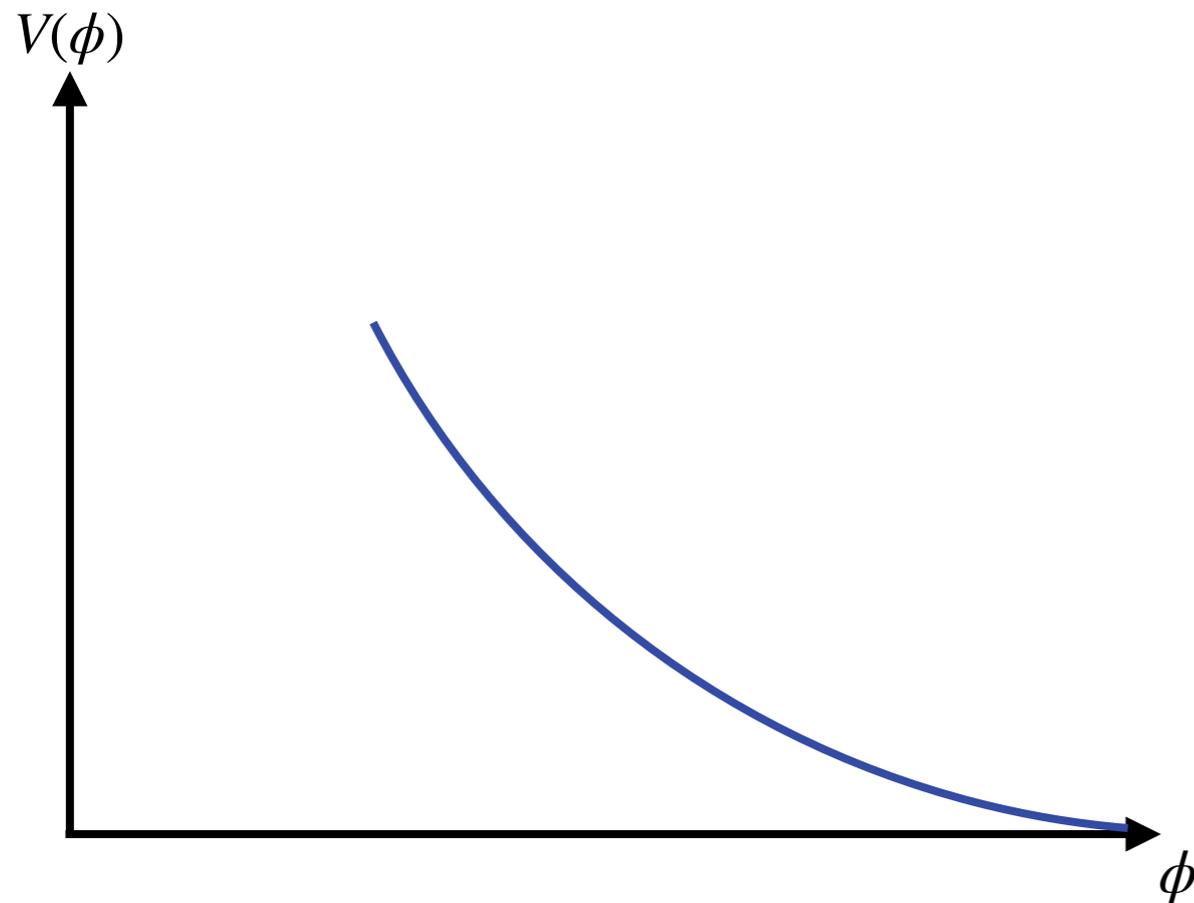
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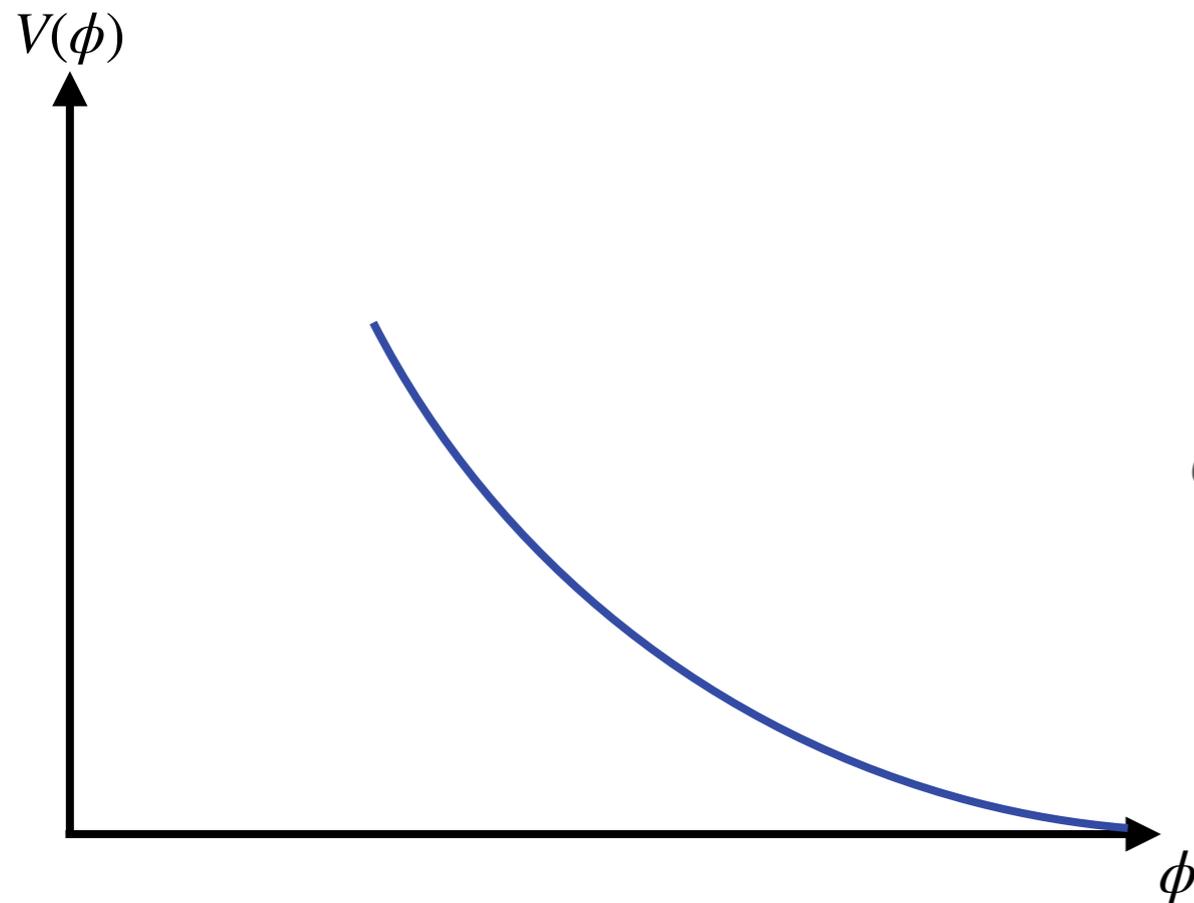


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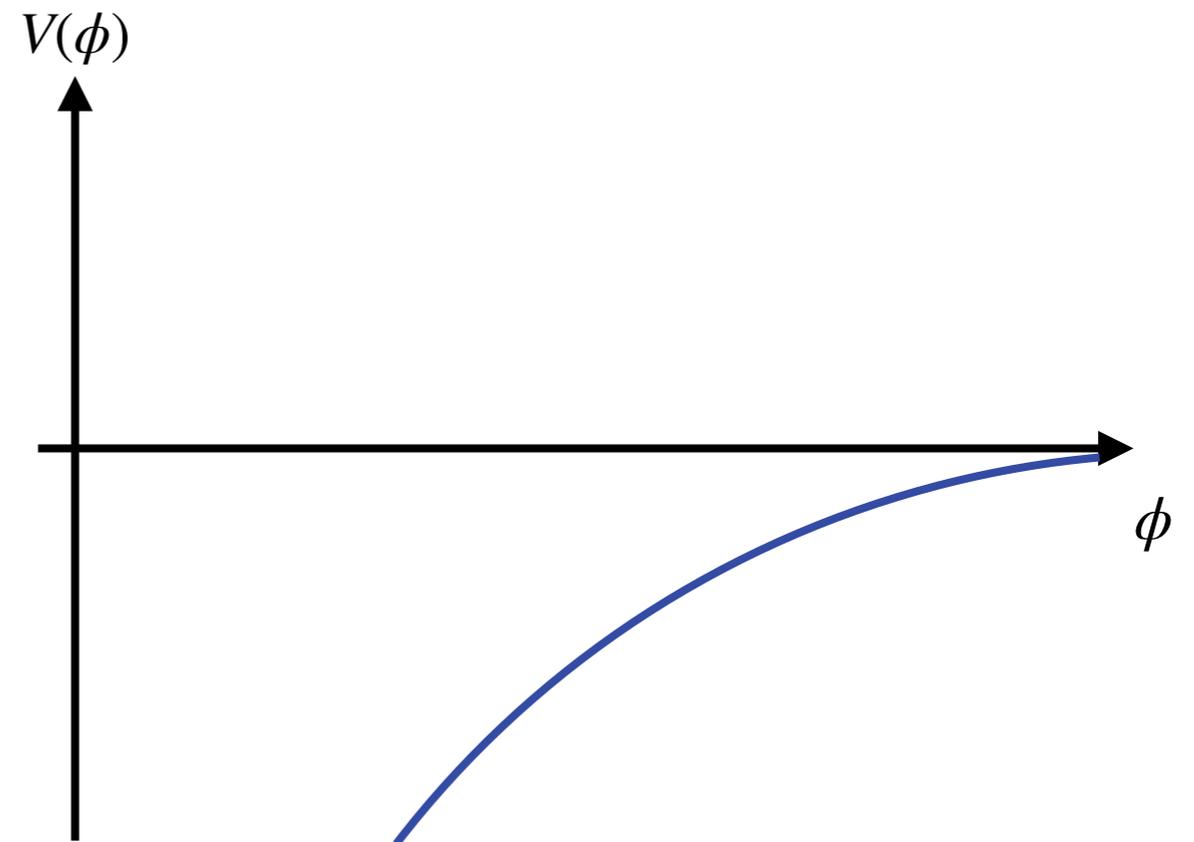
[Dine, Seiberg '85]

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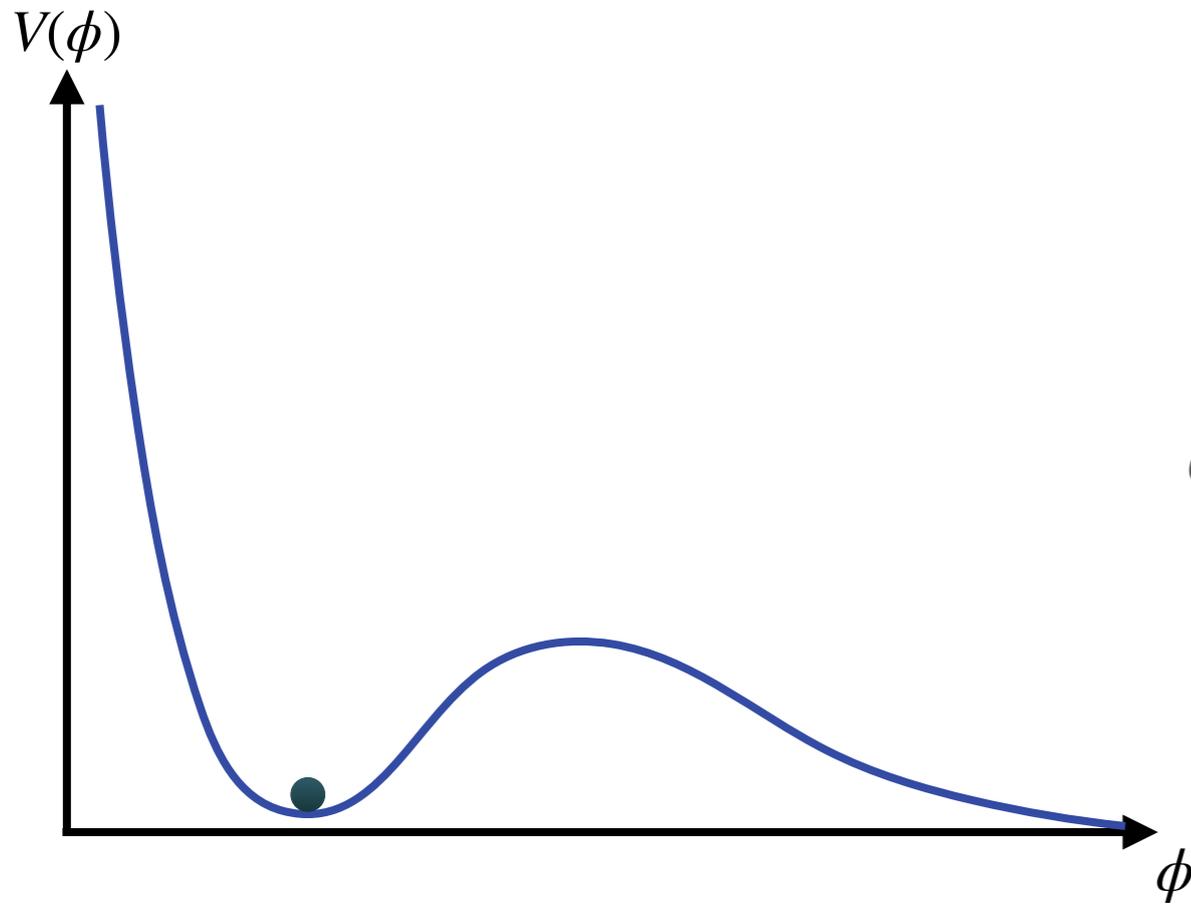


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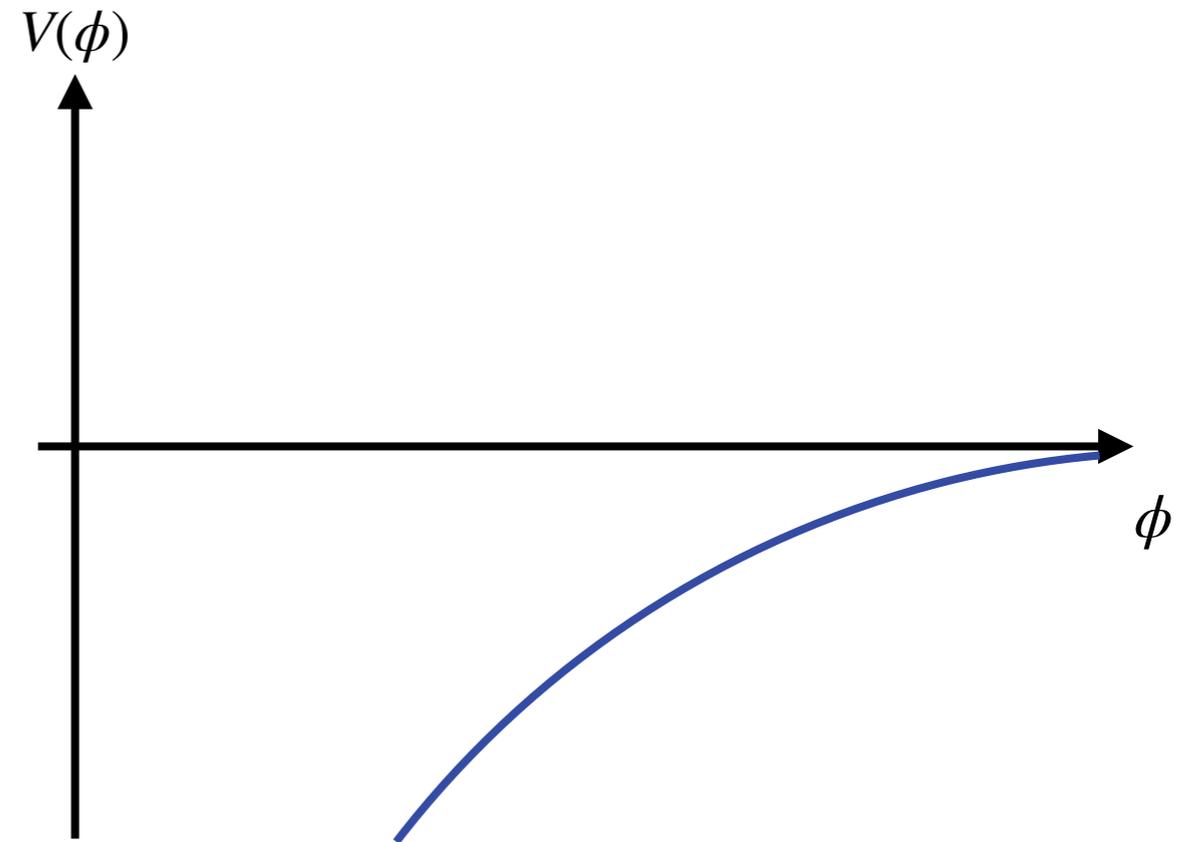
[Dine, Seiberg '85]

take higher order corrections into account:

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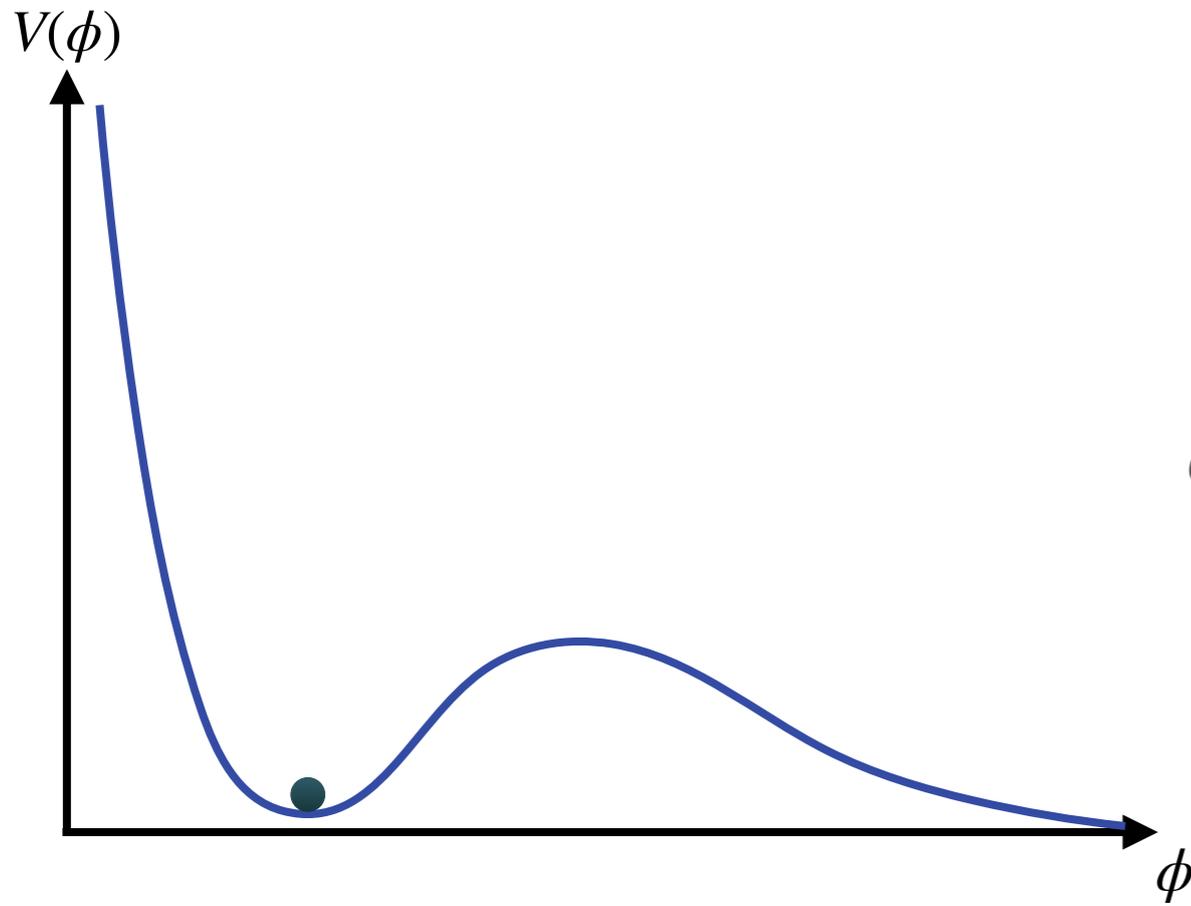


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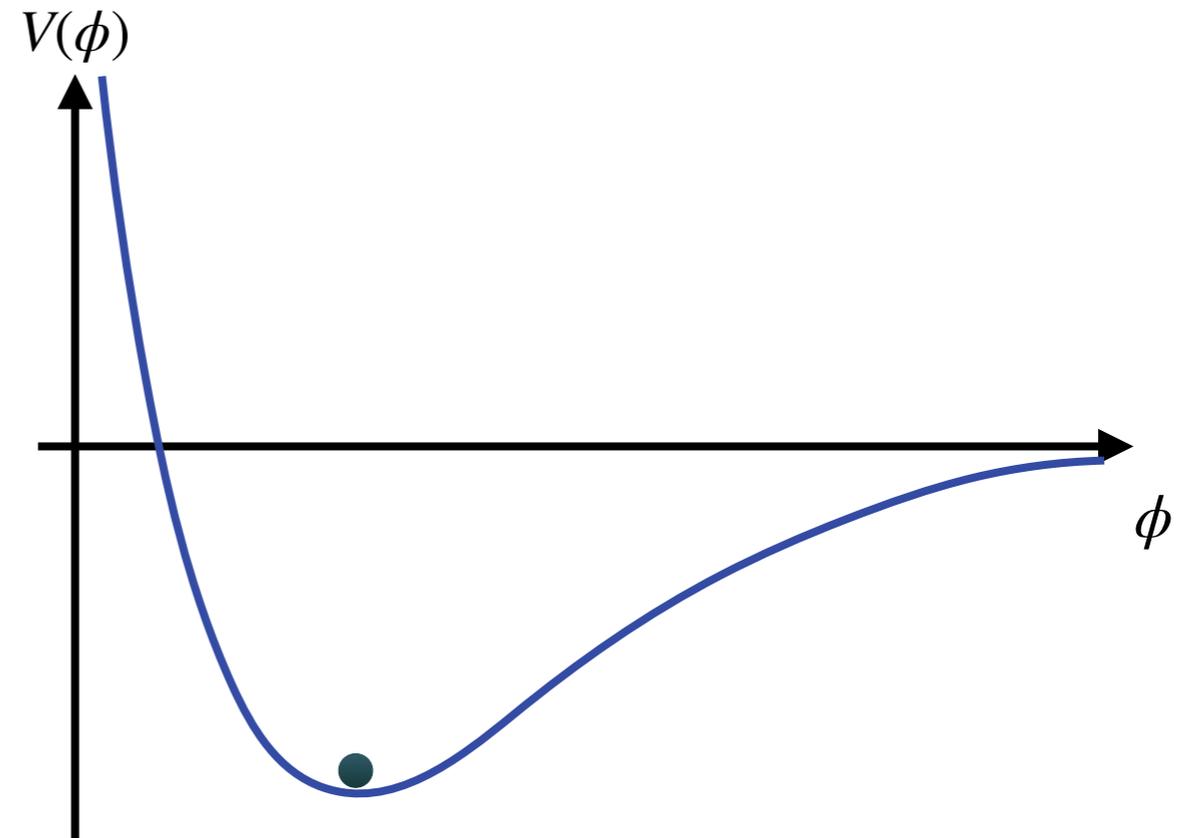
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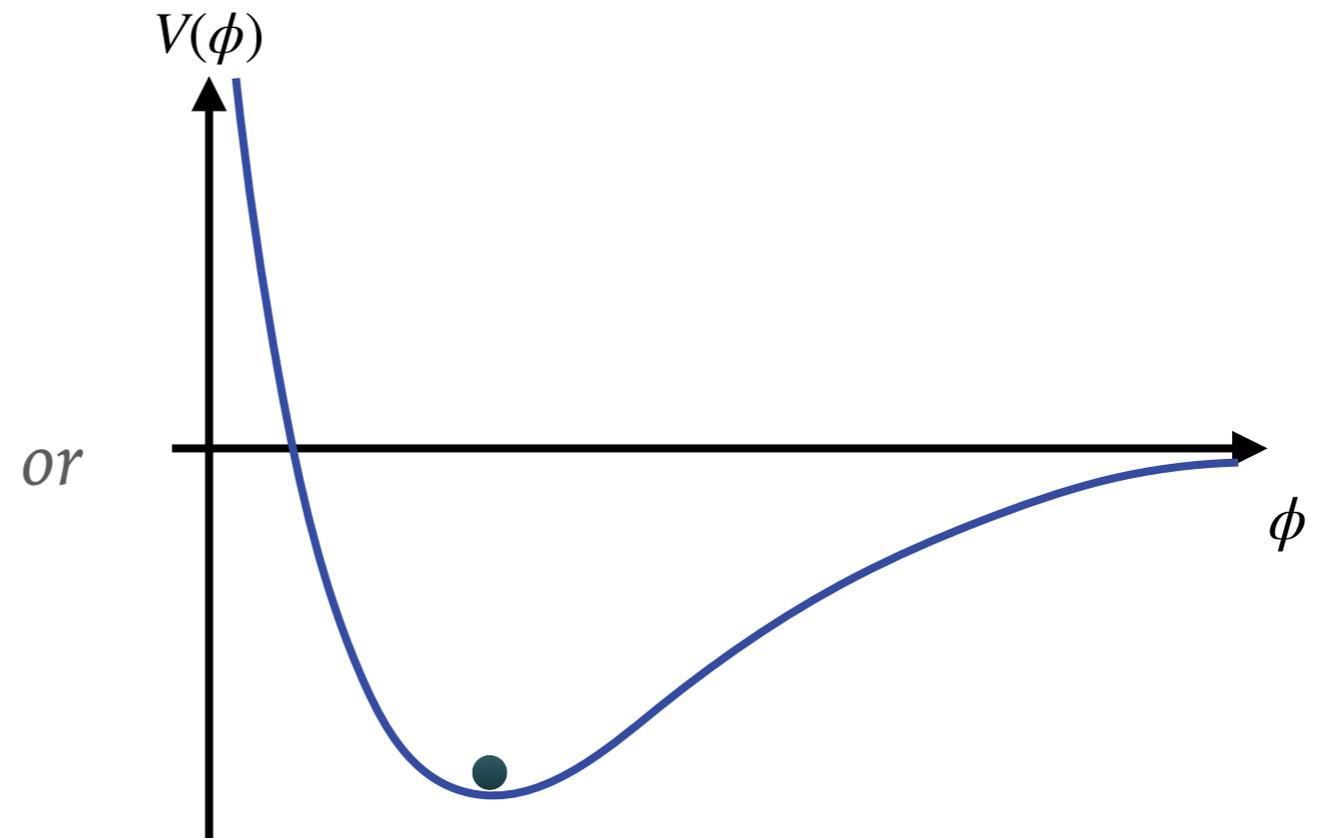
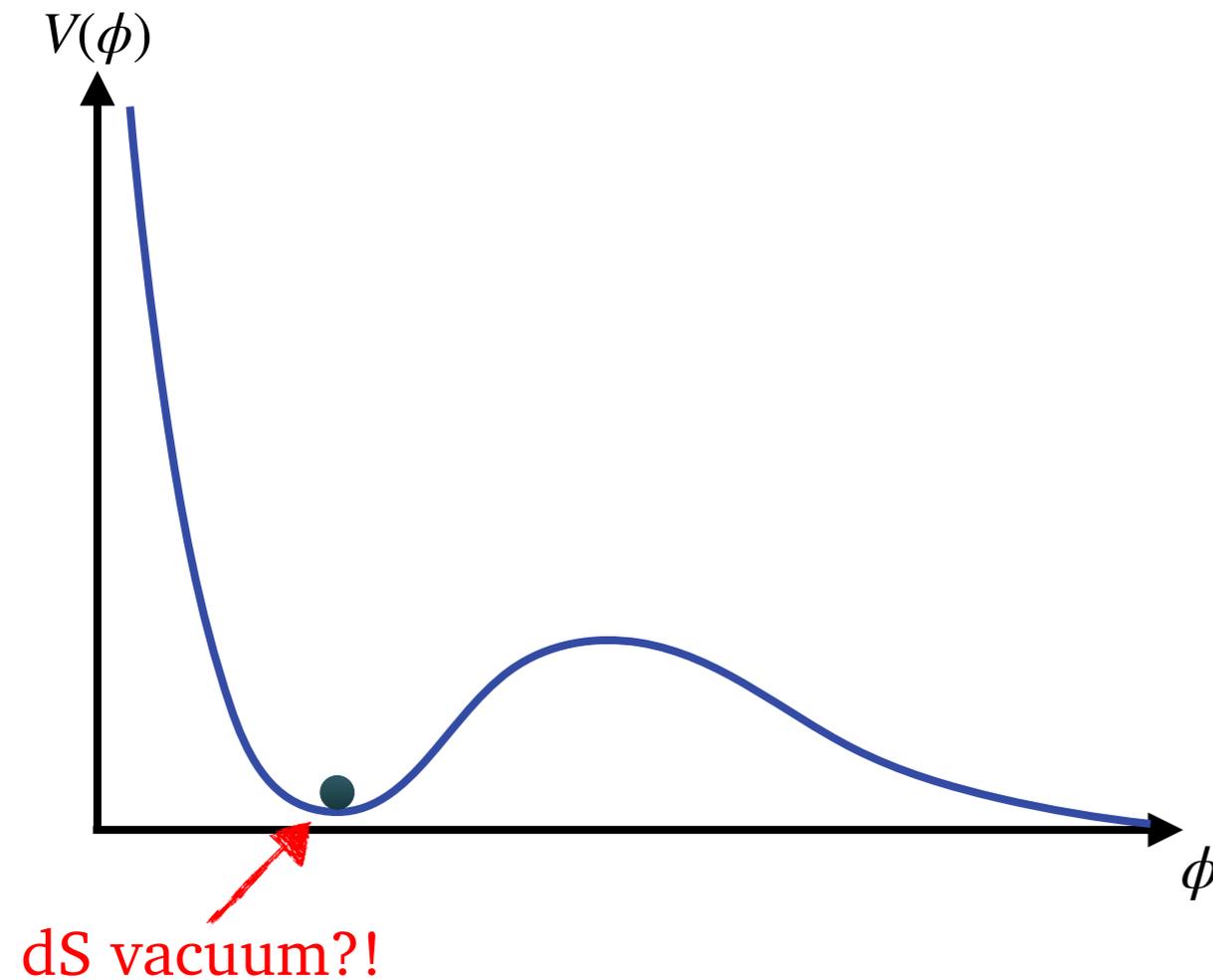


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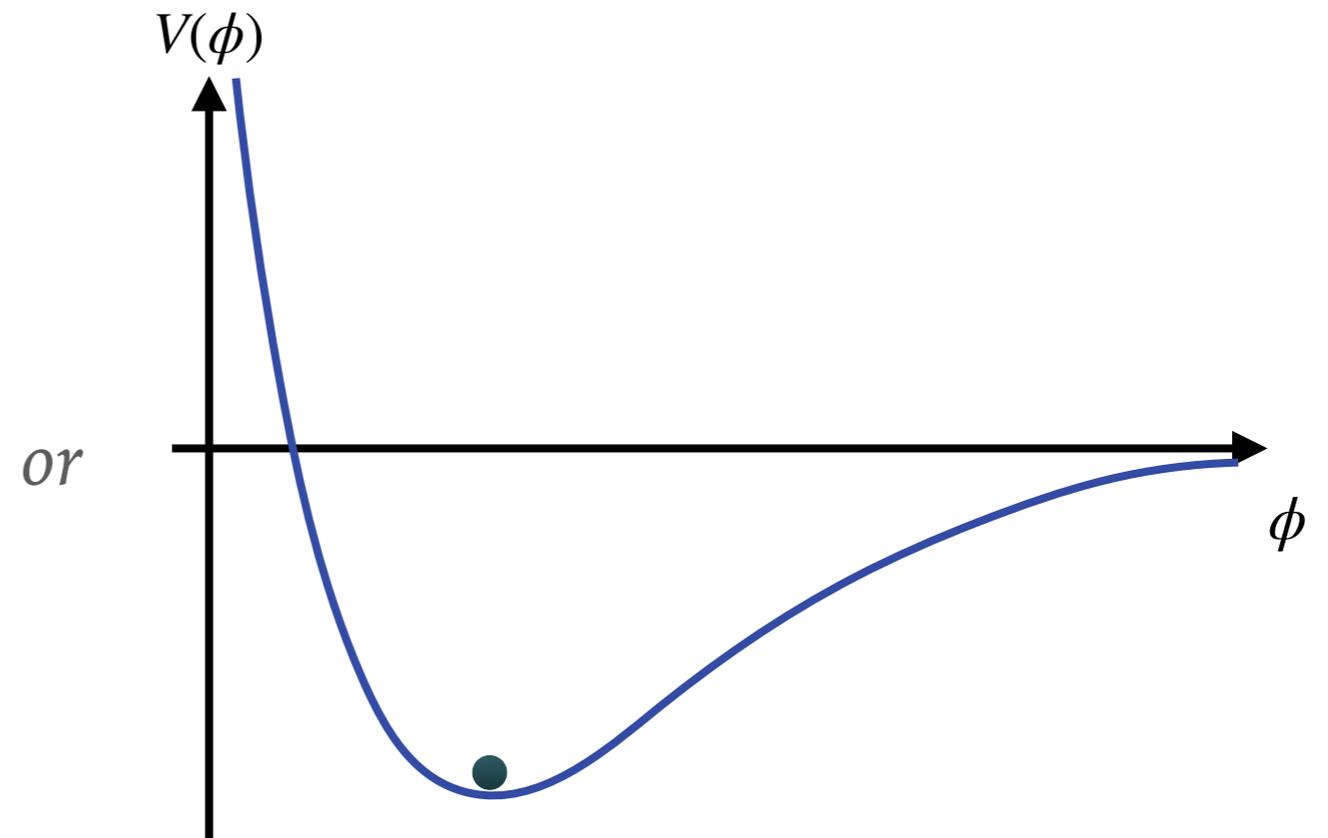
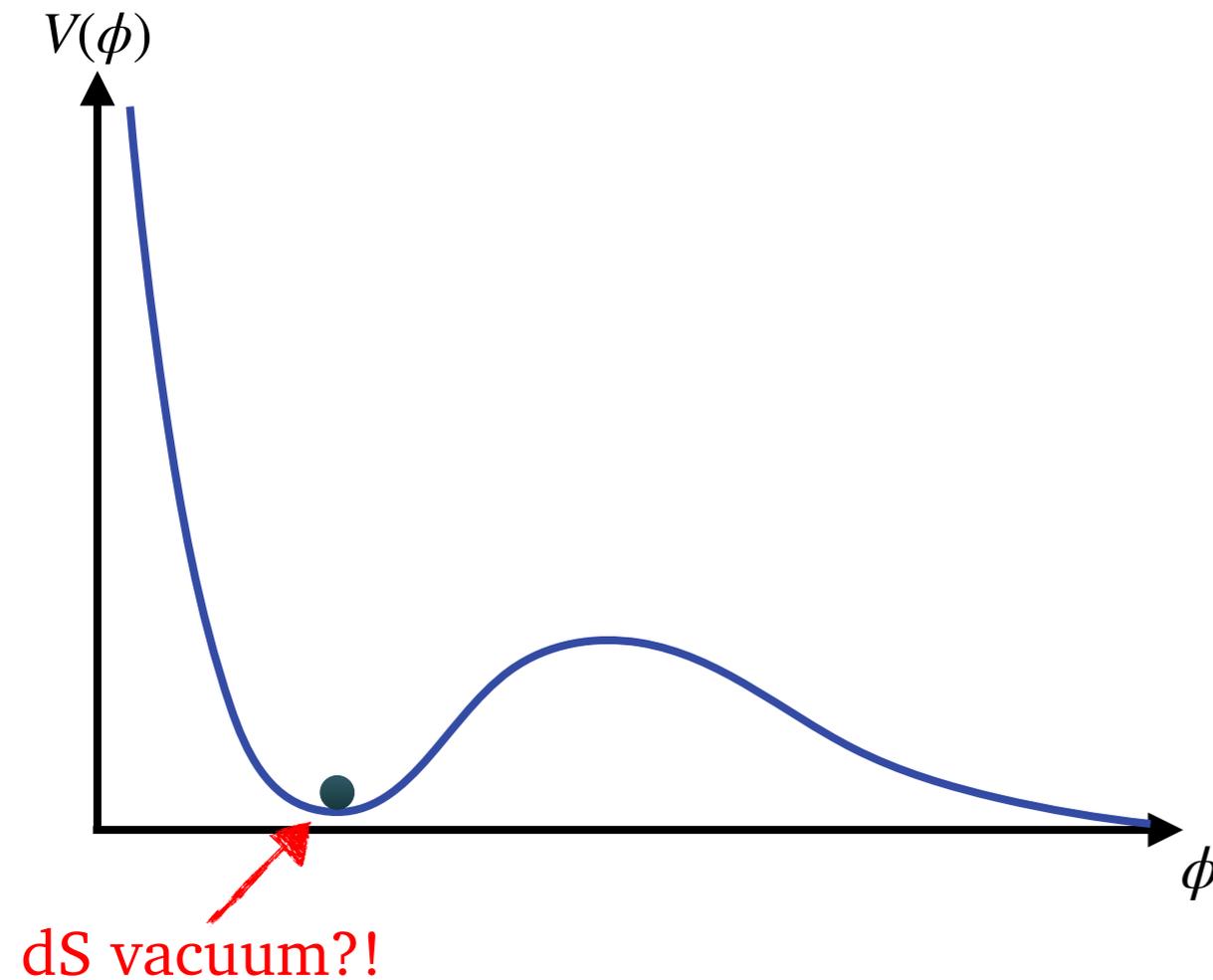


# DINE-SEIBERG PROBLEM

[Dine, Seiberg '85]

take higher order corrections into account:

$$\lim_{\phi \rightarrow \infty} V = 0$$



at minimum of  $V$ :

higher order  
corrections

$\approx$

first order  
corrections

$\rightarrow$

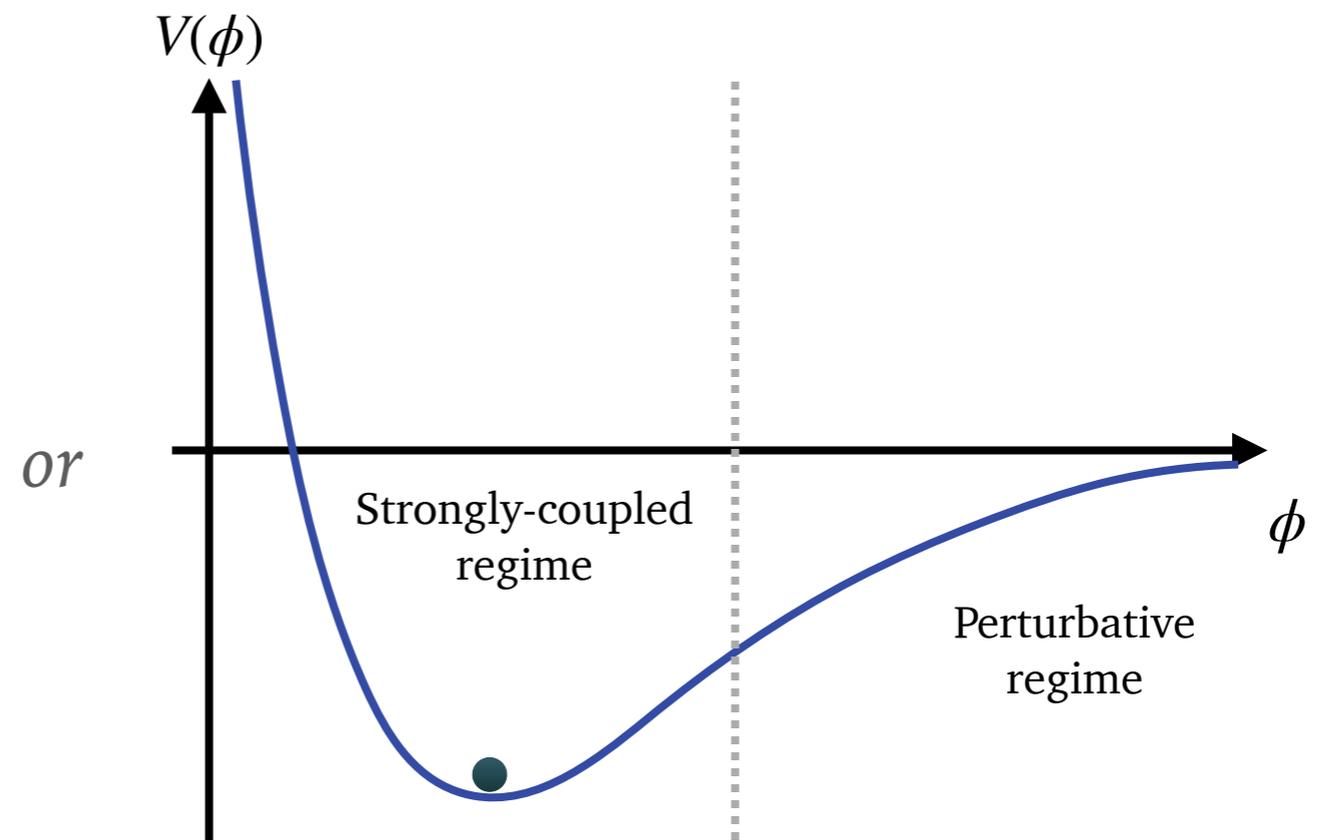
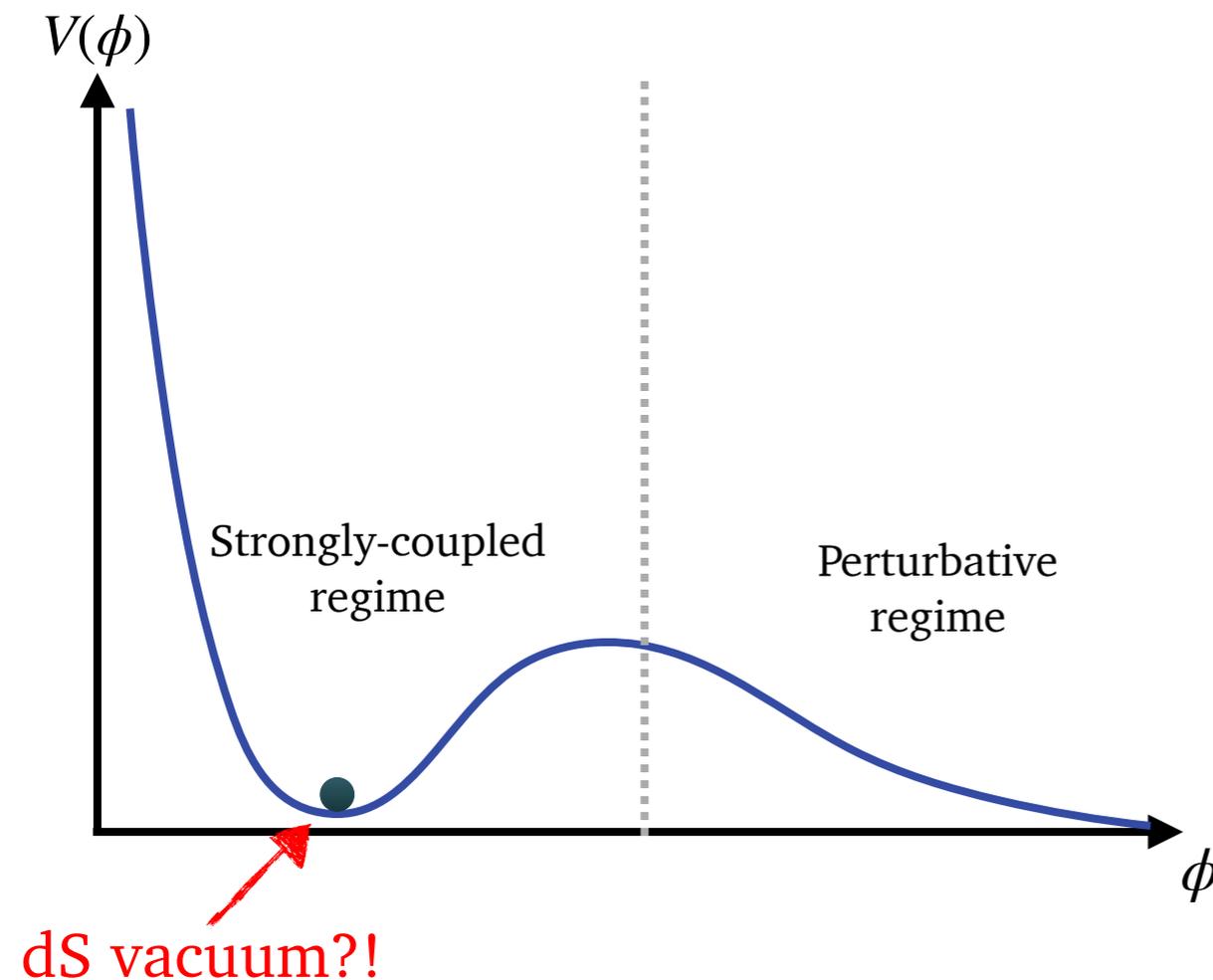
*strong coupling!*

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[Dine, Seiberg '85]

take higher order corrections into account:

$$\lim_{\phi \rightarrow \infty} V = 0$$



at minimum of  $V$ :

higher order corrections  $\approx$  first order corrections  $\rightarrow$  *strong coupling!*

# DINE-SEIBERG PROBLEM

---

de Sitter vacua from quantum corrections only at strong coupling!

higher order corrections:  
generally only known for extended SUSY ( $\mathcal{N} \geq 2$ )!

*“when corrections can be computed, they are not important,  
and when they are important, they cannot be computed”*

*F. Denef, Les Houches Lecture, 2008*

# FLUX COMPACTIFICATION:

---

- Alternative strategy:

Stabilize moduli at the classical level!

➔ *Fluxes!*

*non-vanishing p-form field strengths  $F_{m_1\dots m_p} \neq 0$   
along cycles of the internal geometry*

- Fluxes generate a **potential**:

$$V_F \sim \int \sqrt{g} g^{m_1 n_1} \dots g^{m_p n_p} F_{m_1 \dots m_p} F_{n_1 \dots n_p}$$

- Dependence on volume  $V \sim r^d$ :

$$V_F \sim r^{-d-2p} \int F^2 \quad \rightarrow \quad \textit{runaway towards decompactification!}$$

# FLUX COMPACTIFICATION AND DE SITTER NO-GO

---

- Balance against potential from internal curvature:

$$V_R \sim r^{-2-d} \int R$$

- Schematic form of the overall potential (fluxes + curvature):

$$V = \sum_p r^{-2p-d} \int F_p^2 - r^{-2-d} \int R$$

- For  $V > 0$  (and  $p \geq 1$ ) this potential satisfies

$$\frac{|V'|}{V} \geq \frac{d+2}{\phi} \quad \rightarrow \text{no de Sitter minima!}$$

(AdS minima are easily possible, e.g. Freund-Rubin type  $AdS_{D-d} \times S^d$ )

# DE SITTER NO-GO

---

- [Maldacena, Nuñez '00]:

*From any two-derivative supergravity there is no smooth compactification to de Sitter!*

- de Sitter vacua from String Theory must involve:

a) *quantum effects*

*or*

b) *stringy ingredients* (higher-derivative terms, O-planes, ...)

➔ Danger of Dine-Seiberg like control issues!

- Swampland de Sitter conjecture [Obied, Ooguri, Spodyneiko, Vafa '18]

$$\frac{|\nabla V|}{V} \geq \mathcal{O}(1)$$



*Likely true in asymptotic limits  
but not necessarily everywhere in field space*

# DE SITTER CONSTRUCTIONS

---

➤ new strategy:

*combine different effects (classical + corrections)  
to avoid Dine-Seiberg!*

➤ two main competitors (both in IIB or F-theory):

- *KKLT* [Kachru, Kallosh, Linde, Trivedi '03]

- *Large Volume Scenario (LVS)*

[Balasubramanian, Berglund, Conlon, Quevedo '05]

➤ many other ideas (not in this talk)

- *classical* [Danielson et al. '11], [Andriot '19] for reviews

- *non-geometric* e.g. [de Carlos, Guarino, Moreno '09]

**DE SITTER FROM IIB  
KKLT (AND LVS)**

# IIB DE SITTER VACUA

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Three step procedure [KKLT '03]

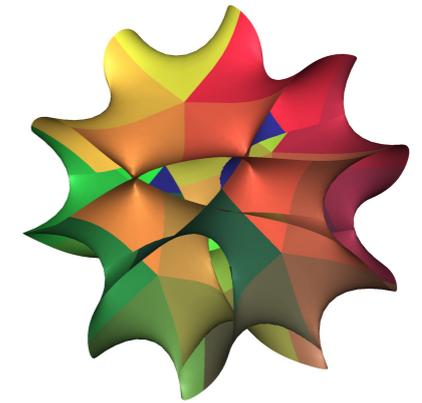
1. Calabi-Yau orientifold with **complex structure-moduli** stabilized by three-form **fluxes**
2. Stabilize **Kähler moduli** by
  - a) **non-perturbative** quantum effects (KKLT)
  - b)  $\alpha'$  corrections (LVS)  
→ (supersymmetric) AdS-vacuum
3. Supersymmetry breaking by an **anti-D3-brane** at the bottom of a warped throat  
→ exp. suppressed **uplift to dS** due to strong warping

# COMPACTIFICATION ON CALABI-YAU MANIFOLDS

---

► IIB on Calabi-Yau:

$$M_{10} = M_4 \times CY_3$$



→ 4d N=2 supergravity

→ many massless scalar fields (moduli):

- $h^{1,1}$  Kähler moduli (volumes of 2 or 4-cycles)
- $h^{2,1}$  complex structure moduli (volumes of 3-cycles)

→ *very well understood!*

*(moduli space geometry, quantum corrections, Mirror symmetry, BPS states, ...)*

# SUPERSYMMETRY BREAKING FROM ORIENTIFOLDING

---

- ▶ divide  $CY_3$  by a discrete involution (i.e.  $\mathbb{Z}_2$  action) combined with world sheet parity  $\Omega_p$  (and  $(-1)^F$ )
- effect A: **break SUSY** from  $\mathcal{N} = 2$  to  $\mathcal{N} = 1$ !
- effect B: fixed points of  $\mathbb{Z}_2$  action:
  - **orientifold planes** ( $O3$  and  $O7$ )

O-planes carry charge!

*Tadpole cancellation condition:* 
$$N_{D3} - \frac{1}{4}N_{O3} + \frac{1}{2} \int H_3 \wedge F_3 = 0$$

(for O3-planes; generally includes also O7 and D7)

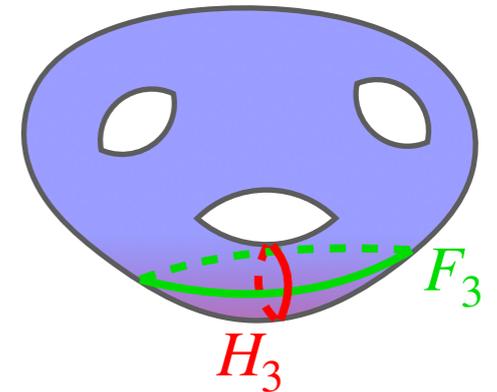
Need D3-branes and/or **fluxes!**

# FLUX COMPACTIFICATION

---

- ▶ 3-form fluxes:  $\langle F_3 \rangle \neq 0, \langle H_3 \rangle \neq 0$

fix volume of 3-cycles  
(i.e. complex structure moduli!)



- ▶ Kinetic term for  $G_3 = F_3 - \tau H_3$ :

$$\mathcal{L}_{kin} \sim \int G_3 \wedge \star \bar{G}_3 \sim \int d^6 y \sqrt{g} g^{il} g^{jm} g^{kn} G_{ijk} \bar{G}_{lmn}$$

depends on  
CY-metric  $g_{ij}$



Classical potential  
for complex structure moduli!

# THE NO-SCALE POTENTIAL

- **classical** superpotential from fluxes (GVW): [Gukov, Vafa, Witten '99]

$$W = \int G_3 \wedge \Omega \quad (G_3 = F_3 - \tau H_3)$$

- $W$  depends on complex structure moduli but not on Kähler moduli!

$$\partial_\rho W = 0 \quad \Rightarrow \quad D_\rho W \sim \frac{1}{\rho - \bar{\rho}} W$$

$\rho$ : volume modulus

*no-scale  
structure:*

$$K(\rho, \bar{\rho}) = -2 \log(\mathcal{V}) = -3 \log[-i(\rho - \bar{\rho})]$$

- $\mathcal{N} = 1$  supergravity potential:

$$V = e^K \left( g^{a\bar{b}} D_a W \bar{D}_{\bar{b}} \bar{W} - 3 |W|^2 \right) = e^K g^{i\bar{j}} D_i W \bar{D}_{\bar{j}} \bar{W}$$

*all moduli*

*complex structure  
moduli*

# KÄHLER MODULI: THE NO-SCALE POTENTIAL (2)

- No-scale potential from fluxes:

$$V = e^K g^{i\bar{j}} D_i W \bar{D}_{\bar{j}} \bar{W} \longrightarrow \text{depends only on complex structure moduli!}$$

*potential for Kähler moduli only from quantum or stringy corrections!*

- KKLT: non-perturbative corrections to superpotential

$$W = \int G_3 \wedge \Omega + \sum_{\mathbf{k}} \mathcal{A}_{\mathbf{k}}(z^i, G_3) e^{-2\pi k^\alpha T_\alpha} \xleftarrow{\text{Kähler moduli}}$$

- LVS: perturbative corrections to Kähler potential

$$K(\rho, \bar{\rho}) = -2 \log \left[ \mathcal{V} + g_s^{-3/2} \xi \right] \xleftarrow{\xi \sim \chi(\text{CY}_3)}$$



*avoid control problems à la Dine-Seiberg by balancing corrections against classical terms?*

# KKLT: KÄHLER MODULI STABILIZATION [Kachru, Kallosh, Linde, Trivedi '03]

- Example: one Kähler modulus  $\rho = i\sigma$

$$W_0 = \langle W_{\text{flux}} \rangle = \left\langle \int G_3 \wedge \Omega \right\rangle$$

$$W = W_0 + Ae^{iap}$$

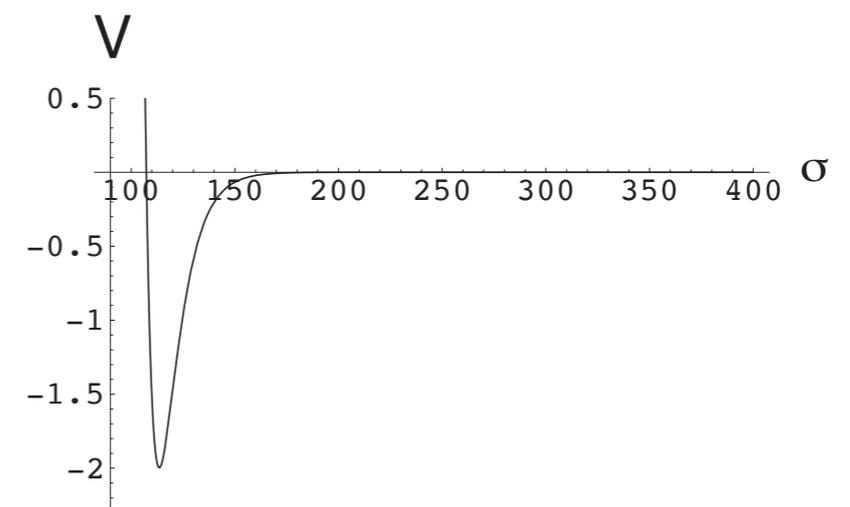
- F-term condition:  $D_\rho W = 0$

$$W_0 = -Ae^{-a\sigma} \left( 1 + \frac{2}{3}a\sigma \right)$$

Balance classical  $W_0$   
against non-pert.  $e^{-a\sigma}$   
→ avoid Dine-Seiberg

- AdS minimum:

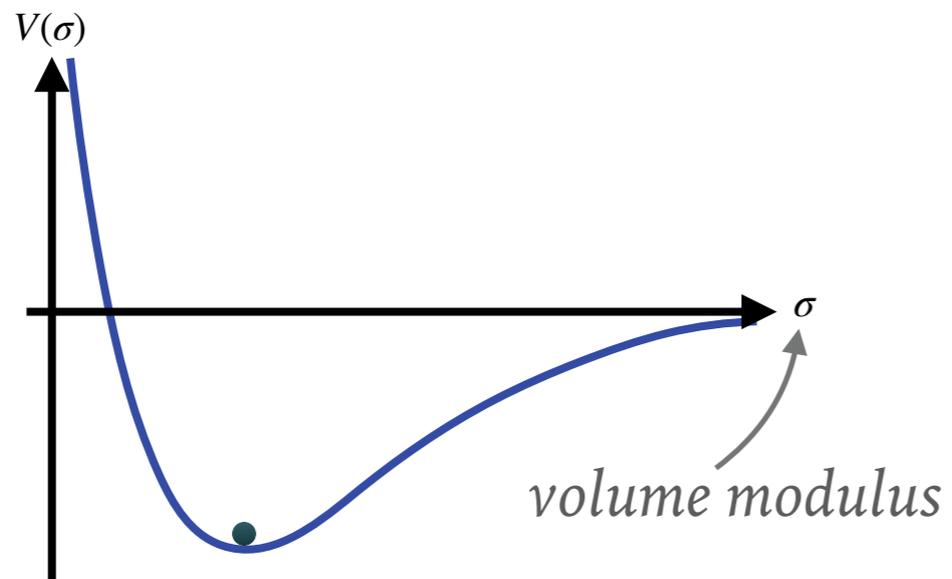
$$V_{AdS} = -3e^K W^2 = -\frac{a^2 A^2 e^{-2a\sigma}}{6\sigma}$$



# ANTI-BRANE UPLIFT

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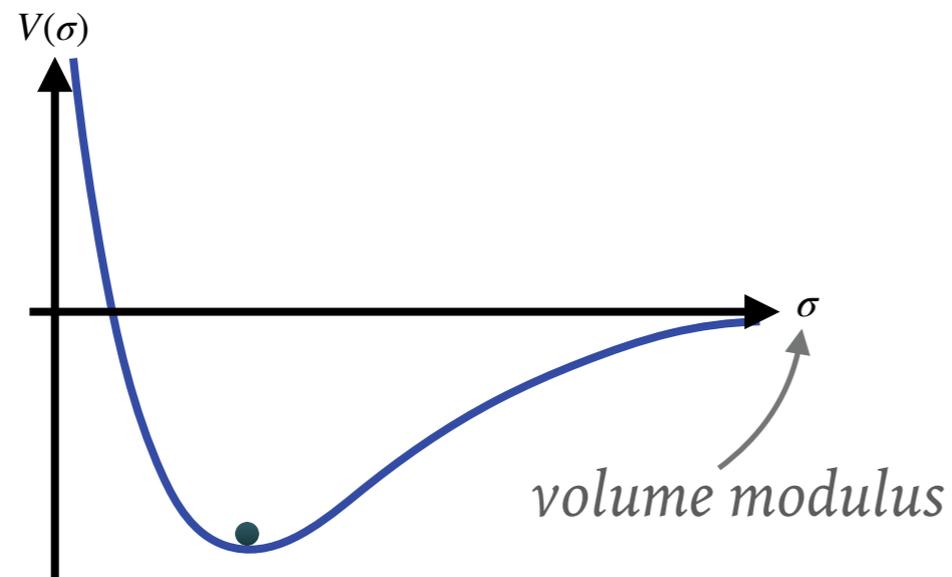
- so far: AdS vacuum ( $\mathcal{N} = 1$  supersymmetric for KKLT)



# ANTI-BRANE UPLIFT

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- so far: AdS vacuum ( $\mathcal{N} = 1$  supersymmetric for KKLT)



- next: raise vacuum energy (and break SUSY)

*add an anti-D3-brane:*

(D3-brane with negative charge)

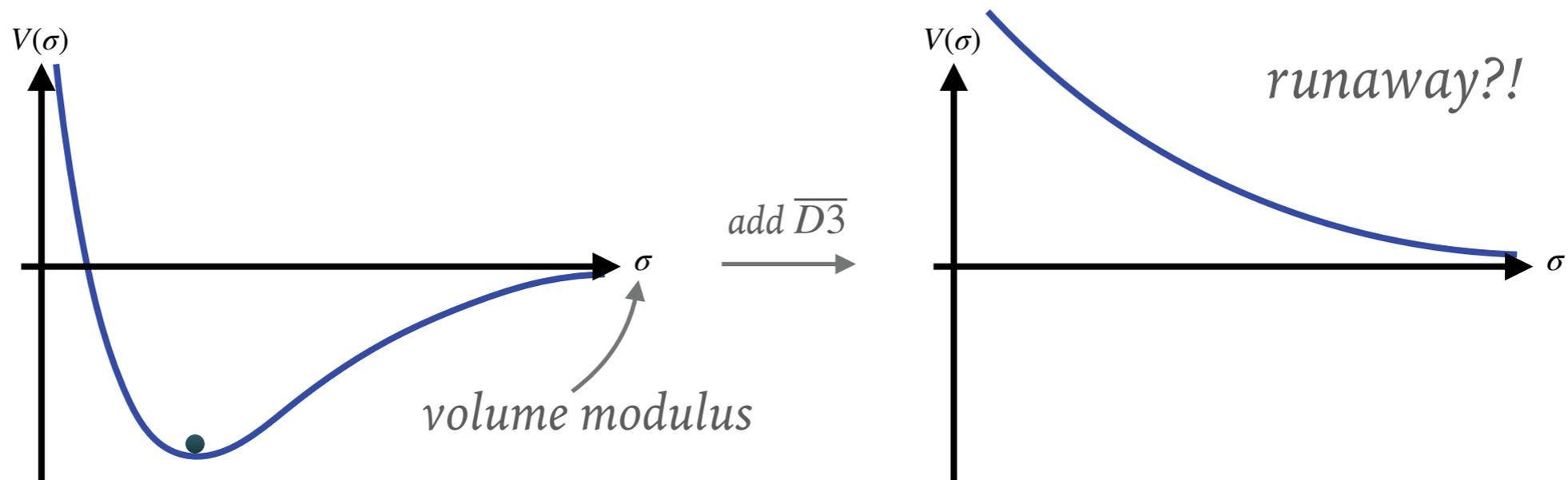


*contribution to potential:*

$$V_{\overline{D3}} \sim \frac{1}{\sigma^3} T_{D3}$$

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contribution to potential:

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- Problem: potential for  $\sigma$  too shallow,  $\overline{D3}$  too heavy!

# BACKREACTION OF FLUXES

[Graña, Polchinski '00, '01]  
 [Giddings, Kachru, Polchinski '02]

- Three-form fluxes in IIB

$$\int_{(A_I, B^I)} F_3 = (M^I, M_I) \in \mathbb{Z} \qquad \int_{(A_I, B^I)} H_3 = (K^I, K_I) \in \mathbb{Z}$$

$A_I, B^I$ : symplectic basis of 3-cycles on the CY  
 $(I = 0, \dots, h^{2,1})$

- back-reaction: warped background:

A: warp factor Calabi-Yau  
(orbifold) metric

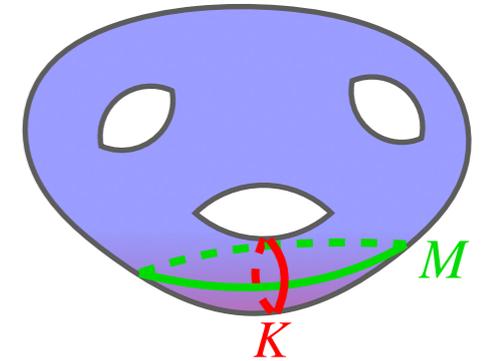
$$ds_{10}^2 = e^{2A} ds_4^2 + e^{-2A} ds_6^2$$

$$F_5 = (1 + \star) vol_4 \wedge de^{4A}$$

# ANTI-BRANE UPLIFT AND WARPED THROATS

---

- Solution: use back-reaction of fluxes to create region with large redshift in Calabi-Yau



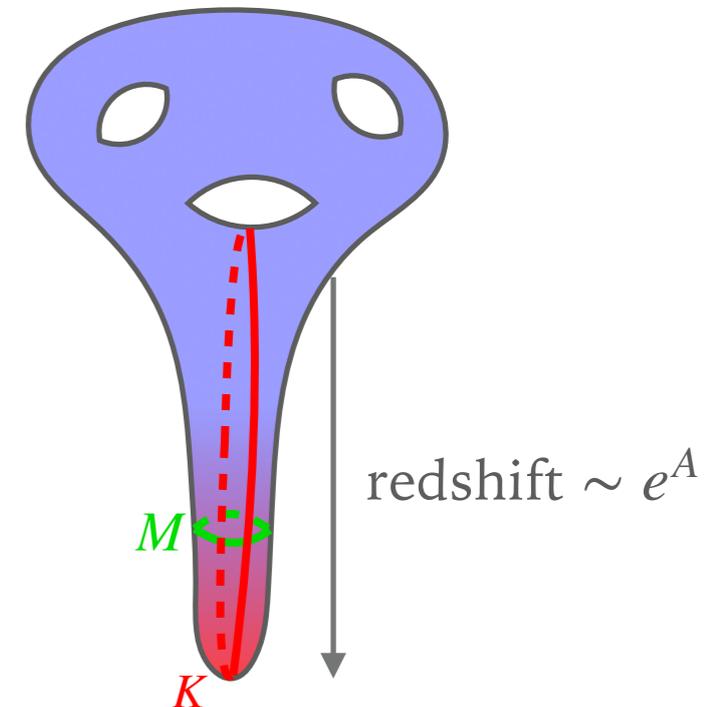
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- Solution: use back-reaction of fluxes to create region with **large redshift** in Calabi-Yau

(local) geometry described by *Klebanov-Strassler* throat:

$$e^{4A} \sim \exp\left(-\frac{8\pi K}{g_s M}\right)$$



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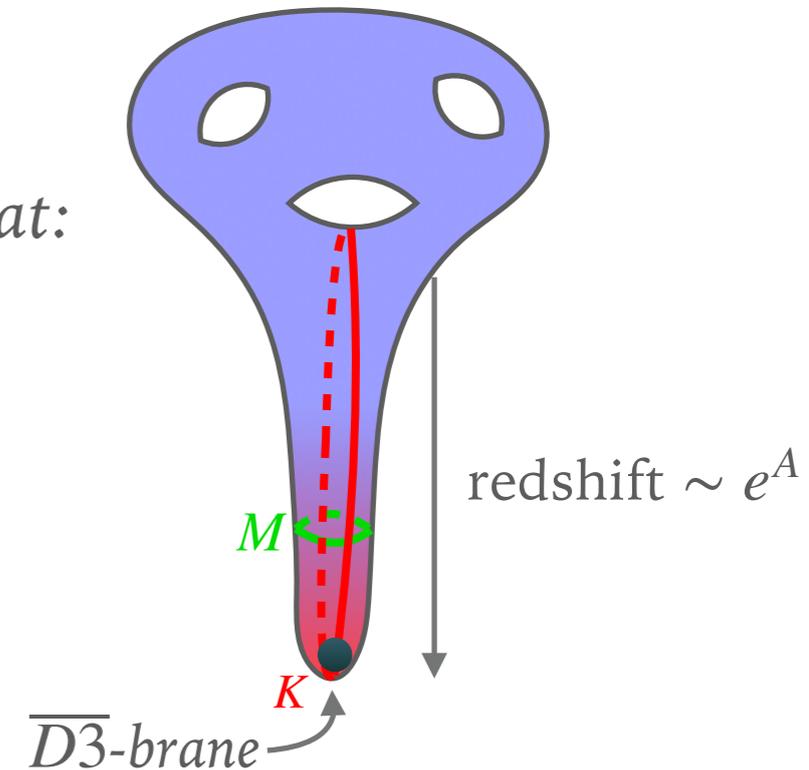
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$\overline{D3}$  potential in warped throat:

$$V_{\overline{D3}} \sim \frac{1}{\sigma^3} e^{4A} T_{D3}$$



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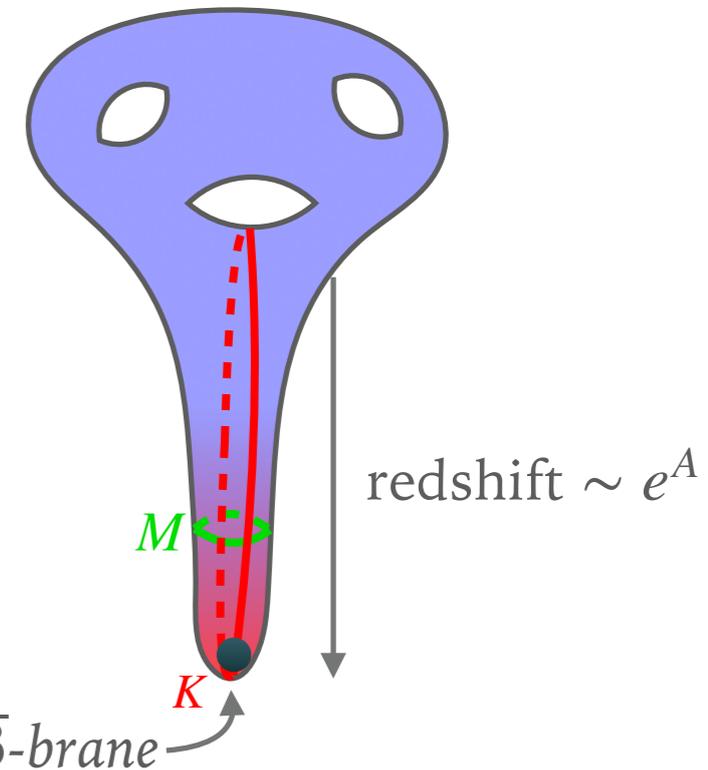
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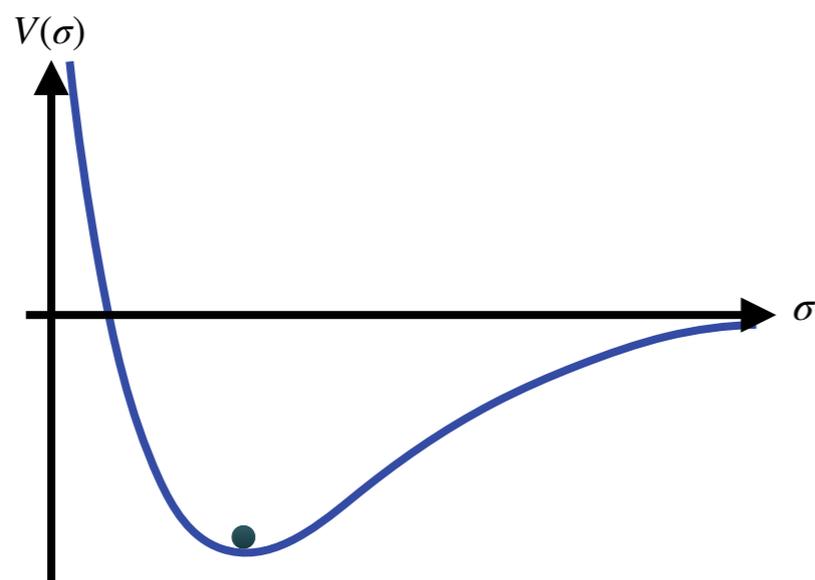


$\overline{D3}$  potential in warped throat:

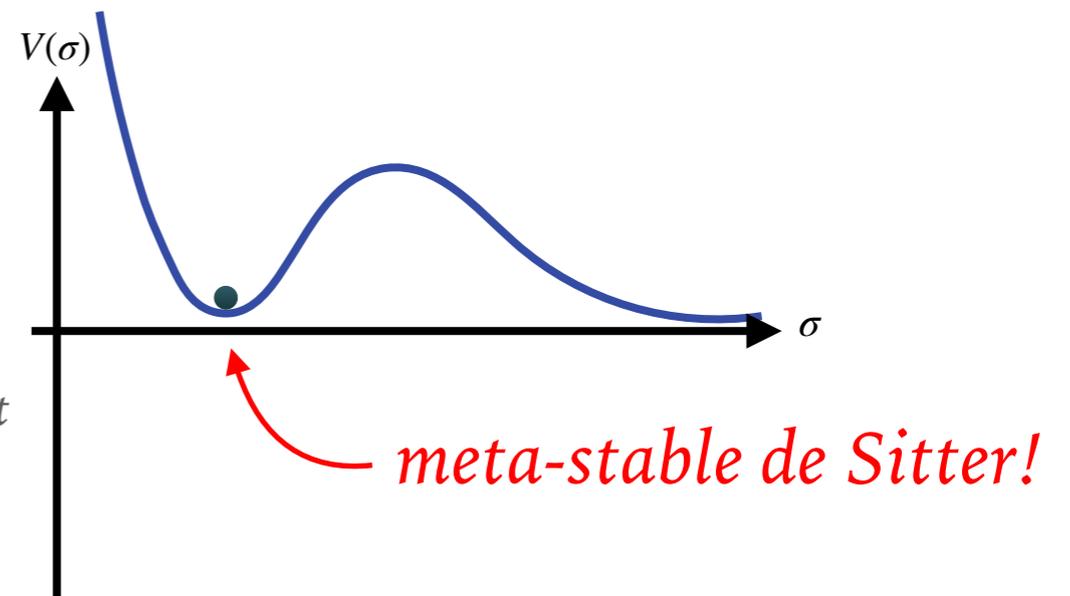
$$V_{\overline{D3}} \sim \frac{1}{\sigma^3} e^{4A} T_{D3}$$



- Effect on volume modulus potential:



add  $\overline{D3}$   
in warped throat



# **DISCUSSION**

**RECENT DEVELOPMENTS AND ISSUES**

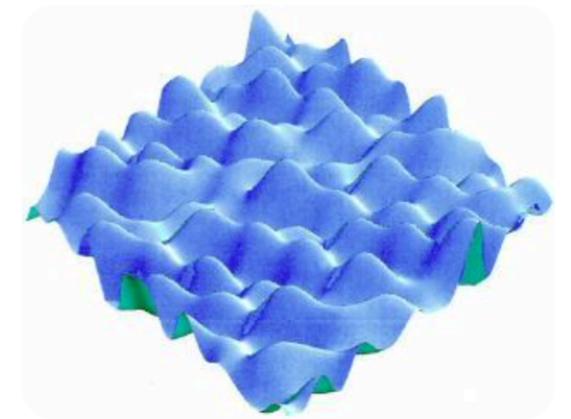
# CONCRETE REALISATIONS IN THE LANDSCAPE

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- Flux Landscape:

Huge combinatorial number of possible Calabi-Yau and flux choices

$$> 10^{500}$$



- Promises high statistical probability to find meta-stable KKLT / LVS de Sitter vacuum [Ashok, Douglas '03], [Denef, Douglas '04]

- Only few concrete realisations (difficult for models with many moduli)

*KKLT/LVS are rather scenarios than concrete models*

- Impressive recent efforts for KKLT AdS vacua

[Demirtas, Kim, McAllister, Moritz, Rios-Tascon '21]

# TADPOLE BOUNDS

---

- Fluxes are constraint by tadpole cancellation condition:

$$\frac{1}{2} \int H_3 \wedge F_3 + Q_{loc} = 0$$

- Tadpole conjecture: [Bena, Blåbäck, Graña, SL '20]

$$\frac{1}{2} \int H_3 \wedge F_3 \text{ grows faster than } Q_{loc} \text{ with the number of moduli}$$

- If, true: Landscape much smaller than anticipated!

➔ *more difficult to obtain required fine-tuning*

[Graña, Grimm, van de Heisteeg, Herraez, Plauschinn '22]

- Confirmed in asymptotic limits in moduli space

- also in deep interior? [SL, Wiesner '22]

[Marchesano, Prieto, Wiesner '21]  
[Plauschinn '21] [SL '21]

# CONTROL ISSUES

---

Careful estimates of the size of relevant corrections:

- KKLT: “Fit” warped throat into Calabi-Yau [Carta, Moritz, Westphal ’18]  
[Gao, Hebecker, Junghans ’20]

*“Singular bulk problem”*

Control over supergravity approximation only if  $Q_{\text{loc}} \gg 1$

(possible resolution: [Carta, Moritz ’20])

- LVS: [Junghans ’22 (2x)]  
[Gao, Hebecker, Schreyer, Venken ’22]

Simultaneous control over all correction requires large  $Q_{\text{loc}}$

- KKLT with fluxes along the lines of [Demirtas, Kim, McAllister, Moritz ’20]:

Controlled mass-hierarchy only if  $Q_{\text{loc}} \gg 1$ ?

# MODULAR CONSTRUCTION

---

- **Bottom-up** (EFT) construction that combines different **top-down** (string theory) ingredients
  - (fluxes, quantum effects, warped throats, anti-branes, ...)*
- Individual ingredients well understood but **interaction between ingredients** often neglected
- incomplete list of possible issues:
  - warped throats and Kähler moduli stabilisation  
[Carta, Moritz, Westphal '18] [Gao, Hebecker, Junghans '20]
  - anti-brane uplift and complex structure moduli stabilisation  
[Bena, Dudas, Graña, SL '18] [Blumenhagen, Kläwer, Schlechter '18]
  - backreaction of fluxes  
[Randall, SL '22]
- related: no genuine  $\mathcal{N} = 1$  formulation, instead treatment as approximative  $\mathcal{N} = 2$

# SCALE SEPARATION

---

- When can we trust a **lower-dimensional EFT** description?

[Gautason, Schillo, Van Riet, Williams '15]

necessary condition:  $|\Lambda| \ll m_{KK}$  “Scale separation”

cosmological constant  $\curvearrowright$  Kaluza-Klein scala = size of the extra-dimension

→ violated for many AdS vacua with extended SUSY (e.g.  $AdS \times S$ )!

- **Swampland AdS conjecture:** [D. Lüst, Palti, Vafa '19]

AdS vacua: tower of states with mass

$$m_{tower} \sim |\Lambda|^\alpha$$

➔ same for **dS**?!

➔ *dark dimension scenario*

[Montero, Vafa, Valenzuela '19]

- KKLT / LVS: appear to satisfy scale separation

*but:* maybe presence of similar tower from warped throat?

[Blumenhagen, Gligovic, Kaddachi '22] 29

# HOLOGRAPHY

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- KKLT: Before de Sitter-uplift: SUSY AdS vacuum

➔ *holographically dual CFT?!*

[Minasian, Tsimpis '99]

[Kounnas, Lüst, Petropoulos, Tsimpis '07]

- CFT dual to flux vacuum (+ non pert. effects) as system of D5/NS5-branes

[SL, Vafa, Wiesner, Xu '22]

count degrees of freedom:

$$c_{CFT} \lesssim Q_{loc}$$

↙ central charge
↘ localised D3-charge in tadpole cancellation condition

- AdS/CFT holography:  $|\Lambda_{AdS}| \sim \frac{1}{c_{CFT}^2} \gtrsim \frac{1}{Q_{loc}^2}$  bounded from below!

➔ *no weakly coupled, scale-separated AdS vacua from KKLT?!*

# DE SITTER VACUA UNDER THE LAMPPOST?

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- weakly coupled, geometric, supersymmetric vacua  
→ only a small fraction of the Landscape
- but: no phenomenological reason to focus on these vacua!



*We study supersymmetric Calabi-Yau vacua  
not because we should but because we can...*

**THANK YOU!**