



Recent advances in precision holography

Valentin Reys

IPhT CEA Paris-Saclay

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Motivation

AdS/CFT provides a gauge theory description of string or M-theory on asymptotically locally AdS backgrounds. Immediate consequence:

$$Z_{\mathsf{CFT}}[J] = Z_{\mathsf{string}/\mathsf{M}}[\phi]$$

► Tested thoroughly in the large *N* limit.

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- ▶ The correspondence is meant to be valid at finite *N*. Schematically,

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- At strong coupling, CFT observables encode information about supergravity. Corrections teach us about string/M-theory beyond the low-energy limit.
- Studying and understanding corrections on both sides of the duality:

Precision Holography

Precision holography

► Need concrete realizations where we can test and make predictions → use supersymmetry to gain computational mileage.

SCFT

Localization gives exact results Solving the matrix models yields 1/N expansion to any order Various techniques available

Analytic & numeric

Supergravity

Exact results out of reach*

Work order-by-order in the derivative expansion

Can also study one-loop effects LO, NLO, NNLO tests

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- ▶ New handle on AdS vacua of string/M-theory with non-trivial fluxes.

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- This talk: recent progress in precision holography for AdS₄/CFT₃.

AdS_4/CFT_3 from M2 branes

- ▶ Study 3d $N \ge 2$ SCFTs arising from low-energy limit of N M2 branes.
- \blacktriangleright No λ in M-theory \rightarrow easier to study the correspondence.

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- M-theory engineers many dual pairs:

Take
$$X_7$$
 to be Sasaki-Einstein \rightarrow CSm

$$egin{aligned} X_7 &= S^7/\mathbb{Z}_k \ (ext{free}) & o ext{ABJM} \ X_7 &= S^7/\mathbb{Z}_r \ (ext{f.p.}) & o ext{ADHM} \ X_7 &= N^{010}, \ V^{52}, \ Q^{111}, \dots \end{aligned}$$



Put the gauge theory on compact M₃ and study susy partition functions.
 Use localization to compute them exactly for various (X₇, M₃).

ABJM on the squashed sphere: $M_3 = S_b^3$

▶ Let us start with simple $X_7 = S^7 / \mathbb{Z}_k$ that gives rise to ABJM theory $\rightarrow U(N)_k \times U(N)_{-k}$ CSm with bifundamental chirals and a superpotential.

[Aharony,Bergman,Jafferis,Maldacena'08]

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- The partition function on S_b^3 localizes to a matrix model.

For
$$\mathcal{F} = \begin{cases} k \ge 1 \& b = 1 \\ k = 1 \& b = \sqrt{3} \end{cases}$$
 matrix model reduces to a free Fermi gas.
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- Fully solved using refined topological string on toric non-compact CYs.
 [Grassi,Hatsuda,Mariño'15]
- ► All perturbative terms in the large *N* limit resum to an Airy function:

$$Z_{\mathcal{F}}(N) = e^{A}C^{-\frac{1}{3}}\operatorname{Ai}[C^{-\frac{1}{3}}(N-B)] + \mathcal{O}(e^{-\sqrt{N/k}}, e^{-\sqrt{Nk}})$$

with (k, b)-dependent coefficients (A, B, C) explicitly known.

Systematic expansion of the free energy at large N:

$$-\log Z_{\mathcal{F}}(N) = \frac{2}{3\sqrt{C}}N^{\frac{3}{2}} - \frac{B}{\sqrt{C}}N^{\frac{1}{2}} + \frac{1}{4}\log N + \mathcal{O}(N^{0})$$

▶ Dual to 4d minimal supergravity solutions (g_4, A) with $\Lambda < 0$ and $\partial M = S_b^3$

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- ▶ Dual to 4d minimal supergravity solutions (g_4, A) with $\Lambda < 0$ and $\partial M = S_b^3$
- LO N^{3/2} term matches the two-derivative regularized on-shell actions.
 [Emparan, Johnson, Myers '99; Martelli, Passias, Sparks '11]
- NLO N^{1/2} term reproduced by including bulk four-derivative terms. [Bobev, Charles, Hristov, VR²[21]
- NNLO log N term is a one-loop effect from summing over the KK modes around the 11d backgrounds. [Bhattacharyya,Grassi,Mariño,Sen'12]

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- NNLO log N term is a one-loop effect from summing over the KK modes around the 11d backgrounds. [Bhattacharyya,Grassi,Mariño,Sen'12]
- ► Precision holography → prediction for all higher-derivative and higher-loop effects in the bulk!

Beyond the ABJM Fermi gas?

- For general (k, b) there is no known Fermi gas.
- Conjecture: we again have an Airy function

[Bobev,Hong,VR'22; Hristov'22]

$$Z_{k,b}(N) = e^{A} C^{-\frac{1}{3}} \operatorname{Ai} \left[C^{-\frac{1}{3}}(N-B) \right] + \mathcal{O}(e^{-\sqrt{N/k}}, e^{-\sqrt{Nk}})$$
$$C = \frac{32}{\pi^{2}k} Q^{-4} \quad \text{and} \quad B = \frac{k}{24} - \frac{2}{k} \left(\frac{1}{3} - 2Q^{-2} \right)$$

with $Q = b + b^{-1}$.

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- Consistent with dual supergravity results up to and including log terms.
- Consistent with ongoing detailed numerical studies of the matrix model.
 Can even access the large-k expansion of the prefactor:

$$A = ig[(b-b^{-1})^2 - 4ig] rac{\zeta(3)}{32\pi^2} \, k^2 + \mathcal{O}(k^0)$$

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$$A = \left[(b - b^{-1})^2 - 4 \right] \frac{\zeta(3)}{32\pi^2} \, k^2 + \mathcal{O}(k^0)$$

Question: what structure produces this Airy function in M-theory? Is there a Fermi gas or topological string picture? Something new?

A word on correlators

The conjecture has implications for ABJM dynamics since the squashing parameter b couples to the stress tensor.

•
$$\mathcal{N} = 2$$
 Ward identities imply [Closset, Dumitrescu, Festuccia, Komargodski'12]
 $\langle T_{\mu\nu}(x) T_{\rho\sigma}(0) \rangle = C_T T_{\mu\nu\rho\sigma}(x) \text{ and } C_T = \frac{\partial^2 \log Z_{k,b}}{\partial b^2} \Big|_{b=1}.$

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- Taking more derivatives gives access to integrated correlators of stress tensors over the round 3-sphere.
- ► Flat space limit gives a way to study scattering amplitudes in M-theory.

[Chester, Pufu, Yin'18]

ABJM on $M_3 = S^1 \times \Sigma_{\mathfrak{g}}$

- ▶ Turn on flux for exact R-sym $\int_{\Sigma_{\mathfrak{g}}} dA_R = 2\pi(1-\mathfrak{g}) \rightarrow \text{topological twist.}$
- Use localization to compute the Topologically Twisted Index (TTI).

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- > The resulting matrix model can be written in the form

[Benini, Hristov, Zaffaroni'16; Closset, Kim, Willett'17]

$$Z_{S^1 imes \Sigma_{\mathfrak{g}}}(N) = \sum_{\{x_i, \tilde{x}_j\} \in \mathsf{BAE}} \mathcal{B}(x_i, \tilde{x}_j)$$

Sum over solutions to transcendental eqs called Bethe Ansatz Equations.

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Sum over solutions to transcendental eqs called Bethe Ansatz Equations.

- Amenable to numerical evaluation:
 - Use known large N solution as init
 - Numerically solve $BAE \rightarrow x_i = x_i^*$ and $\tilde{x}_j = \tilde{x}_j^*$
 - Evaluate $\mathcal{B}(x_i^{\star}, \tilde{x}_j^{\star})$ numerically to high precision
 - Fit with respect to (N, k) and (hopefully) infer analytic expressions

The ABJM twisted index

▶ Detailed numerics \rightarrow we can propose an analytic formula: [Bobev, Hong, VR 22]

$$-\log Z_{S^1 \times \Sigma_{\mathfrak{g}}}(N) = (1-\mathfrak{g}) \frac{\pi \sqrt{2k}}{3} \left(\widehat{N}^{\frac{3}{2}} - \frac{3}{k} \widehat{N}^{\frac{1}{2}} \right) + \frac{1-\mathfrak{g}}{2} \log \widehat{N} \\ + (1-\mathfrak{g}) f(k) + \mathcal{O}(e^{-\sqrt{N/k}}, e^{-\sqrt{Nk}})$$

with $\widehat{N} = N - \frac{k}{24} + \frac{2}{3k}$ acting as a "shifted" N.

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All perturbative corrections in 1/N resum to a (very!) simple form.
 Can also access the large-k expansion of the constant term:

$$f(k) = \frac{3\zeta(3)}{8\pi^2}k^2 - \frac{7}{6}\log k + \mathcal{O}(k^0)$$

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- <u>Question</u>: what is the interpretation of \widehat{N} in M-theory?
- Expand at large $N \to N^{\frac{3}{2}}$, $N^{\frac{1}{2}}$, log N terms... and all perturbative corrections.
- Useful for precision holography.

> The index captures the path integral of M-theory on the 11d background

$$ds_{11}^2 = \frac{L^2}{4} ds_4^2 + L^2 ds_{\mathbb{CP}_3}^2 + L^2 \left(d\psi + \sigma + \frac{1}{4} A \right)^2$$
$$G_4 = \frac{3L^3}{8} \operatorname{vol}_4 - \frac{1}{4} \star_4 F \wedge J$$

with $(g_4, F = dA)$ the Euclidean Romans solution of 4d $\mathcal{N} = 2$ minimal gauged supergravity. [Romans'92; Genolini, Ipiña, Sparks'19; Bobev, Charles, Min'20]

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- When g > 1, can Wick rotate to obtain a Lorentzian magnetic black hole interpolating between AdS₄ and AdS₂ × Σ_g near-horizon.
- LO, NLO and NNLO for BH entropy match the corresponding terms in TTI. [Benini,Hristov,Zaffaroni'16; Bobev,Hong,VR'22; Liu,Pando Zayas,Rathee,Zhao'17]

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- ▶ Precision holography \rightarrow prediction for the corrected Bekenstein-Hawking entropy of this BPS black hole to all orders in 1/N.

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$$\mathcal{N} = 4$$
 with $X_7 = S^7/\mathbb{Z}_r \to \widehat{N} = N + \frac{7r}{24} + \frac{1}{3r}$ and

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•
$$\mathcal{N} = 3$$
 with $X_7 = N^{010}/\mathbb{Z}_k \rightarrow \widehat{N} = N + \frac{5k}{48} + \frac{1}{3k}$ and

$$\frac{\log Z_{S^1 \times \Sigma_{\mathfrak{g}}}}{\mathfrak{g}-1} = \frac{4\pi\sqrt{k}}{3\sqrt{3}} \Big(\widehat{N}^{\frac{3}{2}} - \Big(\frac{k}{4} + \frac{5}{4k}\Big)\widehat{N}^{\frac{1}{2}}\Big) + \frac{1}{2}\log\widehat{N} + h(k)$$

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• $\mathcal{N}=2$ with $X_7=Q^{111}/\mathbb{Z}_k o \widehat{N}=N+rac{k}{6}$ and

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▶ In all cases, black hole solution exists \rightarrow prediction for entropy at finite N.

Valentin Reys (IPhT Saclay)

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 - rotating black holes
 - solutions coupled to matter multiplets

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- ▶ Reorganize the large N expansion into a genus expansion at fixed $\lambda = N/k$ → predictions for type IIA string theory on $M_4 \times KE_6$ to all orders in α'
- Develop bulk/string computations to match SCFT results to all orders? cf. [Hristov,Lodato,VR'17-18; Hristov'22]
- Logarithmic terms currently under investigation.

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- Tests and predictions for AdS vacua, including AdS black holes.
- Can also study
 - rotating black holes [Bobev, Choi, Hong, VR'22]
 - solutions coupled to matter multiplets

▶ Reorganize the large N expansion into a genus expansion at fixed λ = N/k → predictions for type IIA string theory on M₄ × KE₆ to all orders in α'

- Develop bulk/string computations to match SCFT results to all orders? cf. [Hristov,Lodato,VR'17-18; Hristov'22]
- Logarithmic terms currently under investigation.

Thank you for your attention

[Bobev.Hong.VR'22]