



Search for S and T currents with ^{32}Ar

M.Versteegen

P.Alfaurt, D.Aтанасов, P.Ascher, B.Blank, F.Cresto, L.Daudin, X.Fléchard, A.Husson, A.Garcia, M.Gerbaux, J.Giovinazzo, S.Grévy, J.Ha, R.Lica, E.Liénard, D.Melconian, C.Mihai, M.Nasser, C.Neacsu, A.Ortega-Moral, M.Pomorski, M.Roche, N.Severijns, S.Vanlangendonck, D.Zakoucky

PhyNuBE

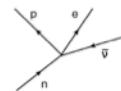
26-31 March 2023



Weak exotic couplings

■ Lee-Yang Lagrangian : n beta decay

$$\begin{aligned}-\mathcal{L}_{LY} = & \textcolor{blue}{C_V} \left(\bar{p} \gamma^\mu n + \frac{\textcolor{blue}{C_A}}{\textcolor{blue}{C_V}} \bar{p} \gamma^\mu \gamma_5 n \right) \times \bar{e} \gamma_\mu (1 - \gamma_5) \nu_e \\ & + \textcolor{magenta}{C_S} \bar{p} n \times \bar{e} (1 - \gamma_5) \nu_e + \frac{1}{2} \textcolor{magenta}{C_T} \bar{p} \sigma^{\mu\nu} n \times \bar{e} \sigma_{\mu\nu} (1 - \gamma_5) \nu_e + hc \\ & + \textit{right-handed neutrinos}\end{aligned}$$



SM "V-A" structure

Exotic couplings : S and T
P omitted

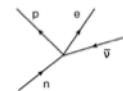
M. González-Alonso, Colloque GANIL (2019)
T.Lee, C-N Yang Phys. Rev. 104 (1956)



Weak exotic couplings

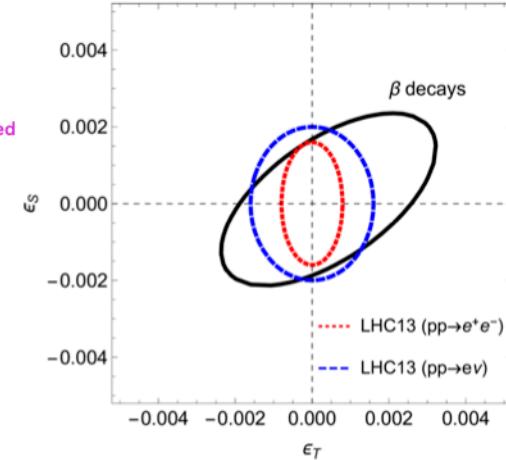
■ Lee-Yang Lagrangian : n beta decay

$$\begin{aligned}
 -\mathcal{L}_{LY} = & C_V \left(\bar{p} \gamma^\mu n + \frac{C_A}{C_V} \bar{p} \gamma^\mu \gamma_5 n \right) \times \bar{e} \gamma_\mu (1 - \gamma_5) \nu_e \\
 & + C_S \bar{p} n \times \bar{e} (1 - \gamma_5) \nu_e + \frac{1}{2} C_T \bar{p} \sigma^{\mu\nu} n \times \bar{e} \sigma_{\mu\nu} (1 - \gamma_5) \nu_e + hc \\
 & + \text{right-handed neutrinos}
 \end{aligned}$$



SM "V-A" structure

Exotic couplings : S and T
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M. González-Alonso, Colloque GANIL (2019)
T.Lee, C-N Yang Phys. Rev. 104 (1956)

M. González-Alonso, O. Naviliat-Cuncic, N. Severijns Prog. Part. Nucl. Phys. (2019)
A. Falkowski, M. González-Alonso, O. Naviliat-Cuncic JHEP04 (2021)

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■ Effective Field Theory

Model independent approach : no assumption on NP origin

Wilson coefficients at quark level :

$$\epsilon_i \propto \left(\frac{m_W}{\Lambda} \right)^2 \sim 10^{-3}$$

TeV NP scale

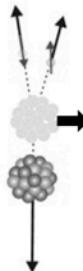
EFT

$$\begin{aligned}
 \bar{C}_V + \bar{C}'_V &= 2g_V(1 + \epsilon_L + \epsilon_R) \\
 \bar{C}_A + \bar{C}'_A &= -2g_A(1 + \epsilon_L - \epsilon_R) \\
 \bar{C}_S + \bar{C}'_S &= 2g_S\epsilon_S \\
 \bar{C}_P + \bar{C}'_P &= 2g_P\epsilon_P \\
 \bar{C}_T + \bar{C}'_T &= 8g_T\epsilon_T
 \end{aligned}$$



Nuclear beta decay

■ Decay rate distribution for polarized nuclei



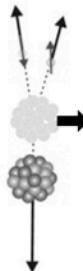
$$\frac{dW(\mathbf{J})}{dE_e d\Omega_e d\Omega_\nu} = dW_0 \times \xi \left\{ 1 + a \frac{\mathbf{p}_e \cdot \mathbf{p}_\nu}{E_e E_\nu} + b \frac{m_e}{E_e} + \frac{<\mathbf{J}>}{J} \cdot \left(A \frac{\mathbf{p}_e}{E_e} + B \frac{\mathbf{p}_\nu}{E_\nu} + D \frac{\mathbf{p}_e \times \mathbf{p}_\nu}{E_e E_\nu} \right) \right\}$$

J.D Jackson, S.B Treiman, H.W Wyld Nuclear Phys 4 (1957)



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$\beta-\nu$ correlation coefficient

CP conserving

Access to C_s and C_T quadratically

Fierz interference term

CP conserving

Access to C_s and C_T linearly

« D » coefficient

CP violating

Access to C_A , C_A' , C_V , C_V' linearly

Correlation measurements
Beta spectrum shape measurements



J.D Jackson, S.B Treiman, H.W Wyld Nuclear Phys 4 (1957)

Angular correlation measurement

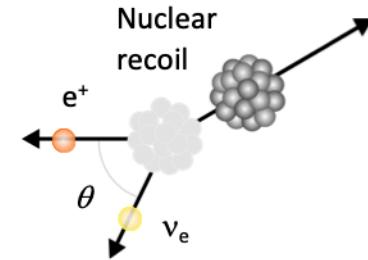


■ Decay rate for non polarized nuclei

$$\frac{dW(\mathbf{J})}{dE_e d\Omega_e d\Omega_\nu} = dW_0 \times \xi \left\{ 1 + a \frac{\mathbf{P}_e \cdot \mathbf{P}_\nu}{E_e E_\nu} + b \frac{m_e}{E_e} + \frac{<\mathbf{J}>}{J} \cdot \left(A \frac{\mathbf{P}_e}{E_e} + B \frac{\mathbf{P}_\nu}{E_\nu} + D \frac{\mathbf{P}_e \times \mathbf{P}_\nu}{E_e E_\nu} \right) \right\}$$

$$dW = dW_0 \times \xi \left(1 + a \frac{\mathbf{P}_e \cdot \mathbf{P}_\nu}{E_e E_\nu} + b \frac{m_e}{E_e} \right)$$

a > 0 : \theta = 0^\circ \text{ favored and large recoil}
a < 0 : \theta = 180^\circ \text{ favored and small recoil}



Angular correlation measurement

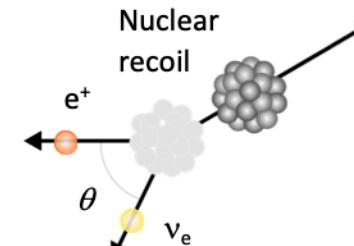


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■ Correlation measurement = recoil measurement

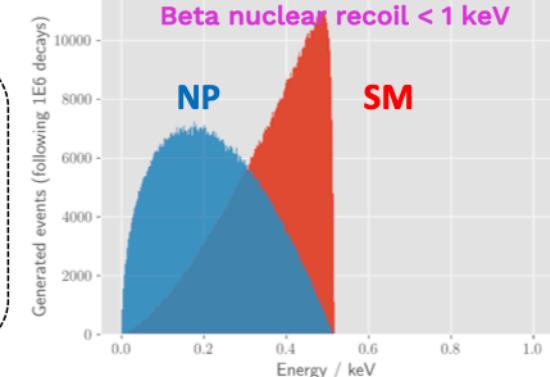
access to :

$$\tilde{a} \sim \frac{a}{1 + b} < \frac{m_e}{E_e} >$$

Pure Fermi transition $\Delta J=0 S=0$
 Vector Coupling

$$a_F \cong 1 - \frac{|C_S|^2 + |C'_S|^2}{|C_V|^2} = 1 \text{ standard model}$$

$$b_F \approx \pm \text{Re} \left(\frac{C_S + C'_S}{C_V} \right) = 0$$



Many Projects



$$\frac{dW(\mathbf{J})}{dE_e d\Omega_e d\Omega_\nu} = dW_0 \times \xi \left\{ 1 + a \frac{\mathbf{p}_e \cdot \mathbf{p}_\nu}{E_e E_\nu} + b \frac{m_e}{E_e} + \frac{<\mathbf{J}>}{J} \cdot \left(A \frac{\mathbf{p}_e}{E_e} + B \frac{\mathbf{p}_\nu}{E_\nu} + D \frac{\mathbf{p}_e \times \mathbf{p}_\nu}{E_e E_\nu} \right) \right\}$$

^6He @ LPC (Paul trap)
 ^8Li @ ANL (Paul trap)
 ^6He @ ANL (MOT)
 ^{32}Ar @ Texas A&M (Penning)
 $^{38\text{m}}\text{K}$ @ TRIUMF (MOT)
 n @ aSPECT
 n @ nab
...

^{114}In @ ISOLDE
 ^6He @ LPC (bSTILED)
 ^6He @ NSCL
...



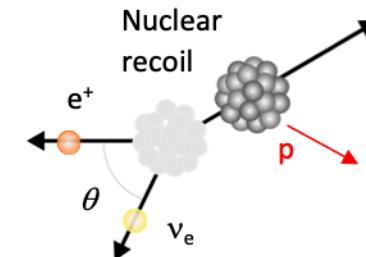
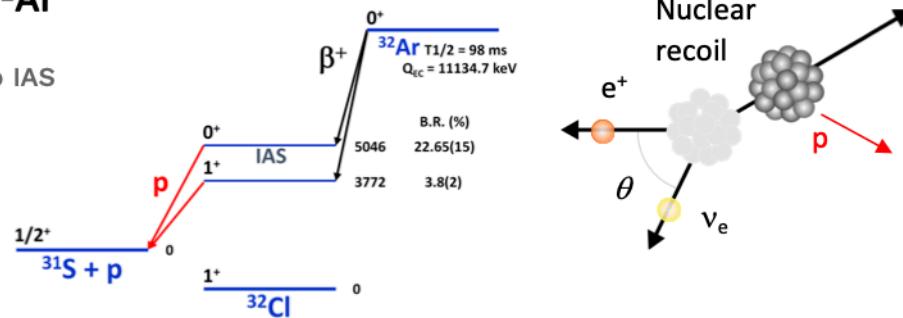
The case of ^{32}Ar



■ β -delayed proton emission in ^{32}Ar

- Fermi $0^+ \rightarrow 0^+$ transition from GS to IAS
- Recoil energy ~ 640 eV
- Beta delayed p emission ~ 3.3 MeV
- IAS : $\Gamma \sim 20$ eV $\Leftrightarrow T_{1/2} \sim 10^{-17}$ s

⇒ p emission in flight from the recoil



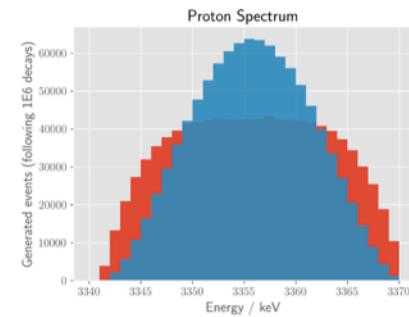
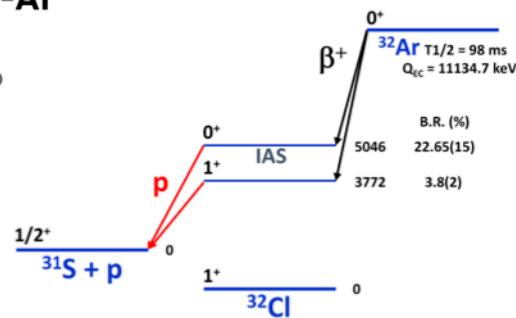
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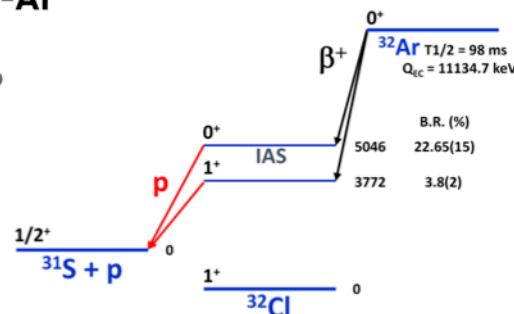
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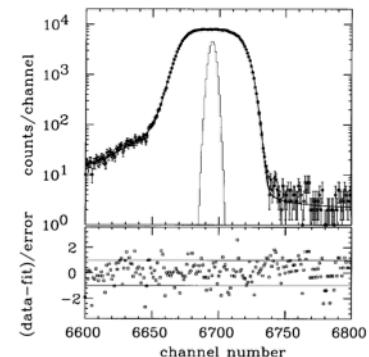
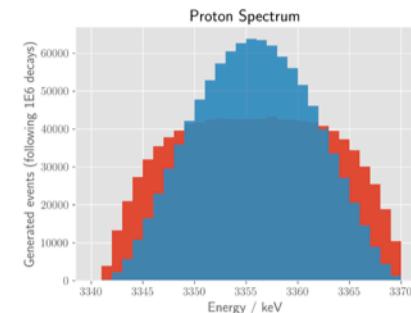
■ Broadening measurement

- ^{32}Ar beam at ISOLDE (60 keV)
- 9x9 mm² cooled p-i-n diodes
high resolution ~ 3 keV FWHM (pulser)
- 3.5 T magnetic field

$$\tilde{a} = 0.9989(52)_{\text{stat}}(39)_{\text{syst}}$$

⇒ precision level : 0.65%

E. G. Adelberger et al. Phys. Rev. Lett. 83 (1999)

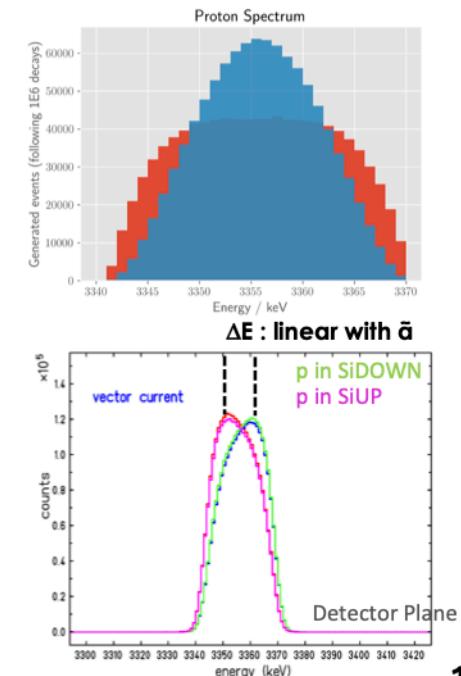
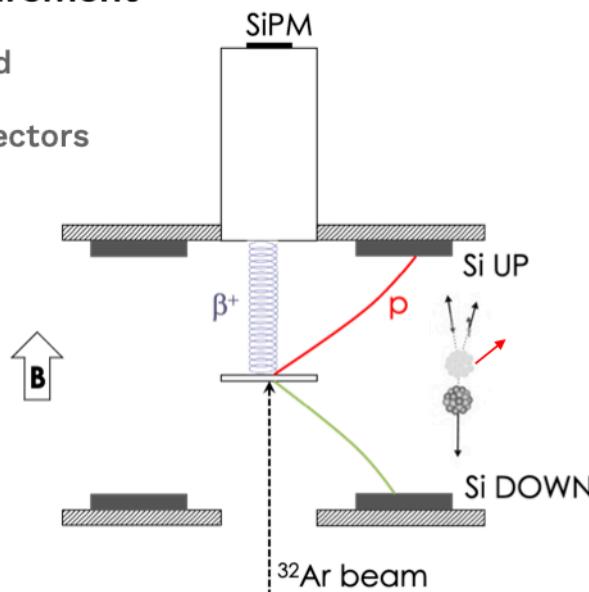




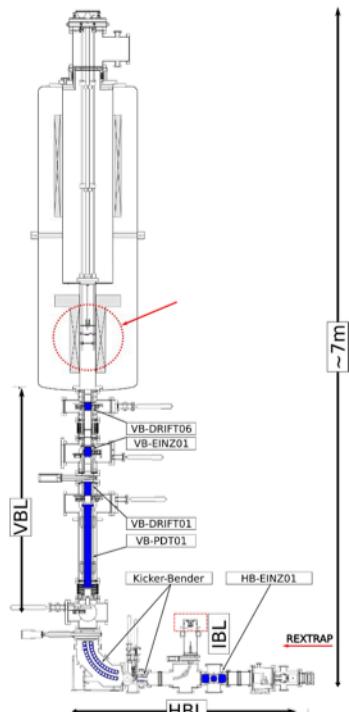
■ β -p coincidence measurement

- Strong magnetic field
- 2 symmetrical p detectors
high resolution
high solid angle
- Beta detector
low detection threshold

⇒ precision level : 0.1%

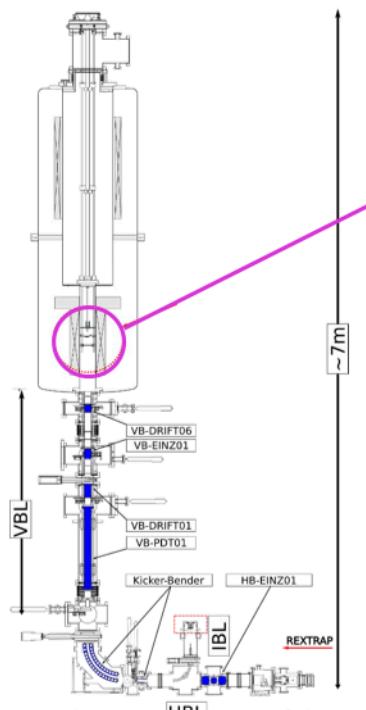


WISArD setup @ISOLDE

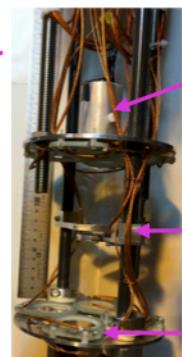


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WISArD setup @ISOLDE



■ Proof-of-Principle Setup 2018



β detector
plastic scintillator + 1 SiPM 6x6 mm²
Hamamatsu

catcher
Al-mylar 6.7(1) μ m

p detectors
2 x 4 Si surface barrier 300 μ m
Dead layer ~ 430 (300) nm
Resolution ~ 35 keV

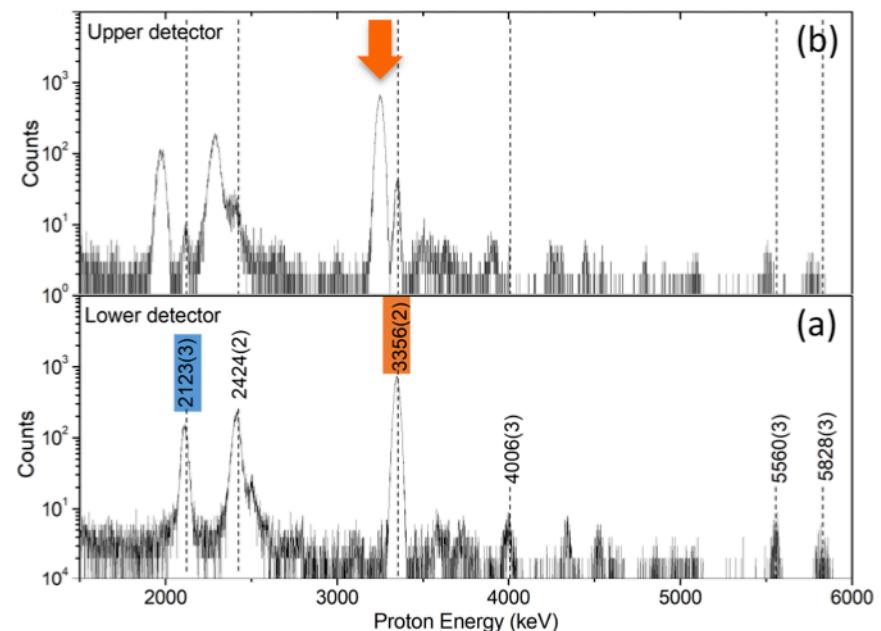
in $B = 4$ T
(WITCH magnet)
+ FASTER DAQ

Proof-of-Principle Experiment - 2018



■ Single p spectra calibration

- Si UP
Catcher thickness
dead layer
- Si DOWN
dead layer

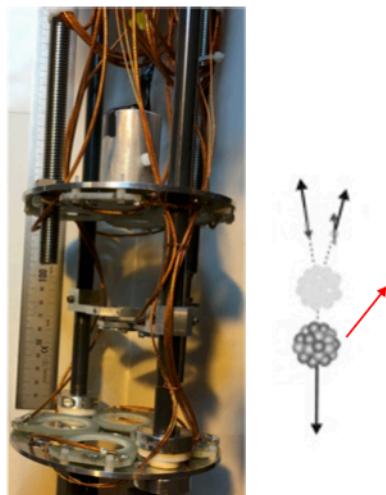


Proof-of-Principle Experiment - 2018



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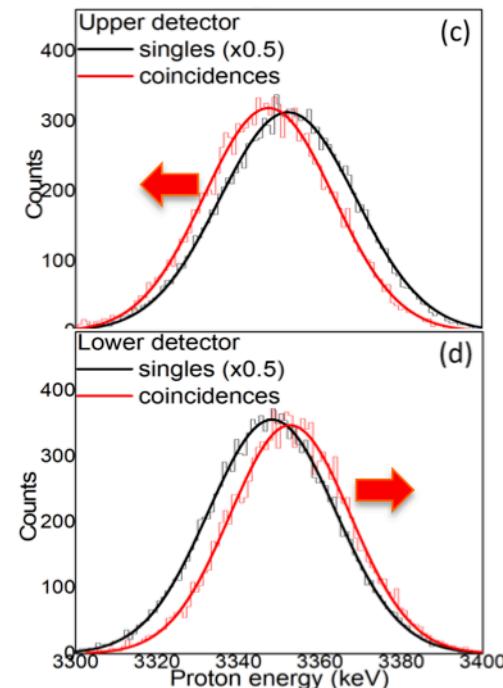


■ β -p coincidence

$$\Delta E = |\bar{E}_{coinc} - \bar{E}_{single}|$$

$$\Delta E_F = 4.49(3) \text{ keV}$$

Mean value over 8 detectors



V. Araujo-Escalona et al. Phys. Rev. C 101 (2020)

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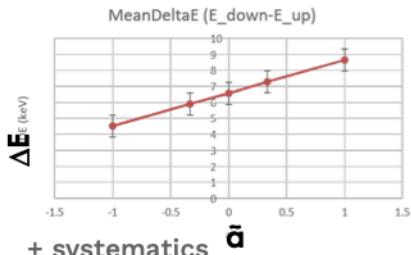
Proof-of-Principle Experiment - 2018



■ Extraction of \tilde{a}

$$\Delta E_F = 4.49(3) \text{ keV}$$

Monte Carlo simulations



35h of beam
Ion transmission 12%
 $N_{\text{coinc}} \sim 10^5$

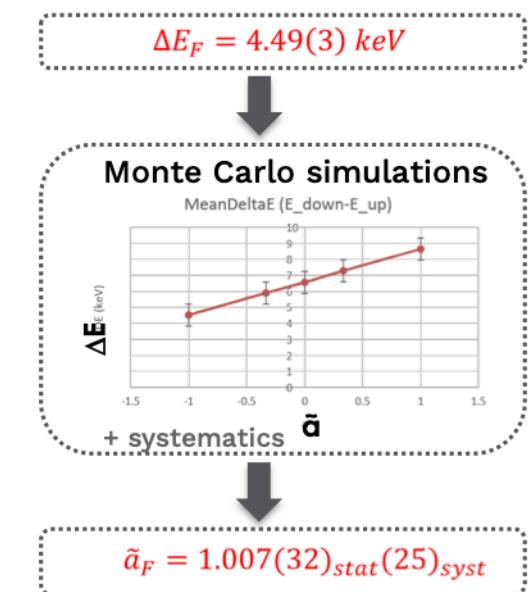
⇒ precision level : 4%

$$\tilde{a}_F = 1.007(32)_{\text{stat}}(25)_{\text{syst}}$$

Proof-of-Principle Experiment - 2018



■ Extraction of \tilde{a}

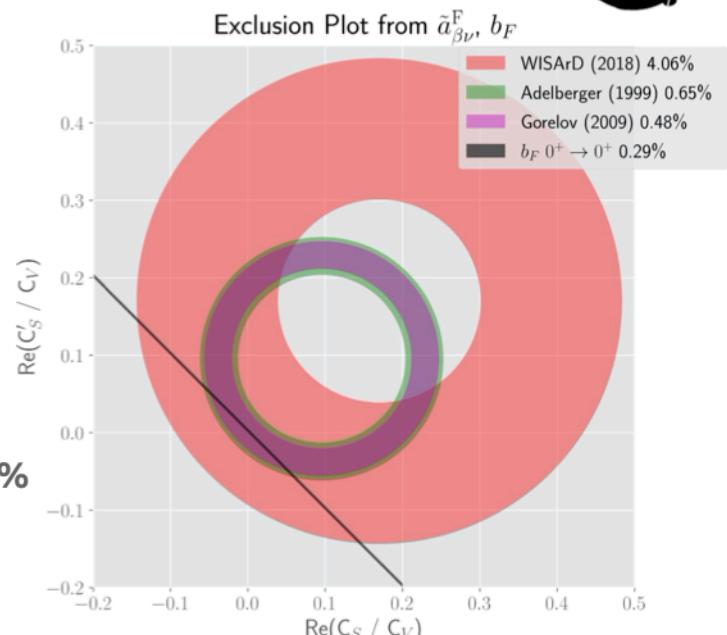


V. Araujo-Escalona et al. Phys. Rev. C 101 (2020)

35h of beam
Ion transmission 12%
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⇒ precision level : 4%

$$a_F \cong 1 - \frac{|C_S|^2 + |C'_S|^2}{|C_V|^2}$$



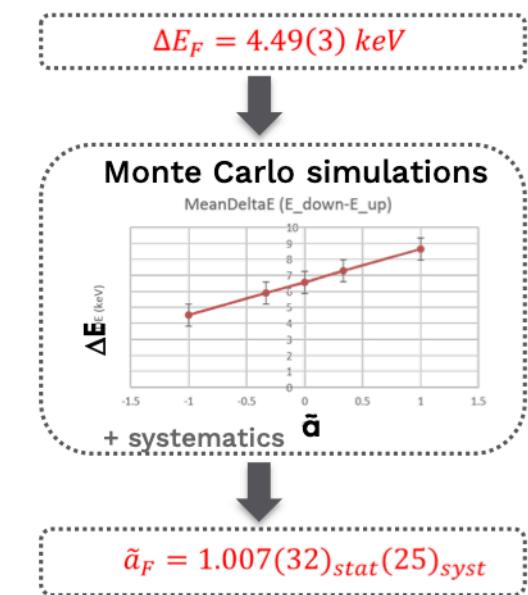
Exclusion plot from D. Atanasov

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Proof-of-Principle Experiment - 2018



■ Extraction of \tilde{a}



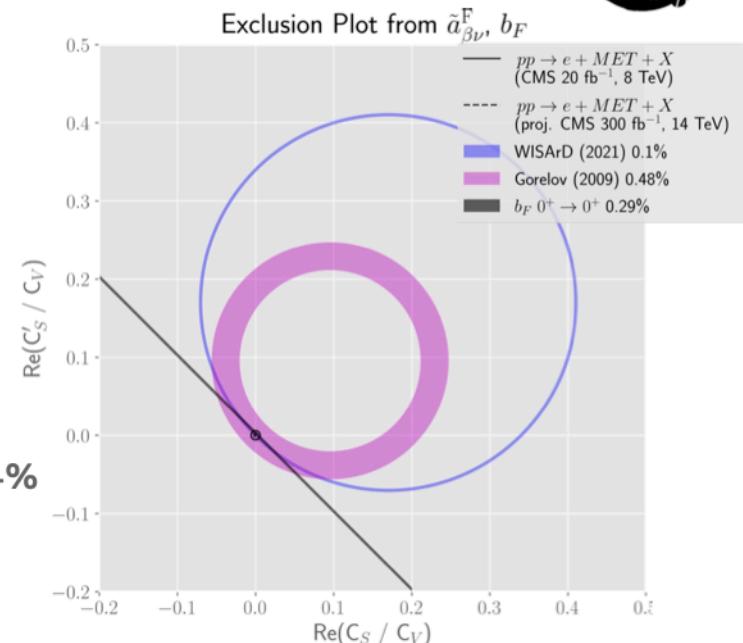
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Exclusion plot from D. Atanasov

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Proof-of-Principle Experiment - 2018

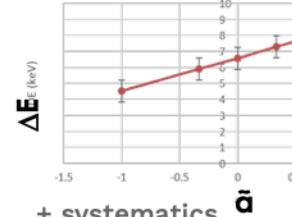


■ Extraction of \tilde{a}

$\Delta E_F = 4.49(3) \text{ keV}$

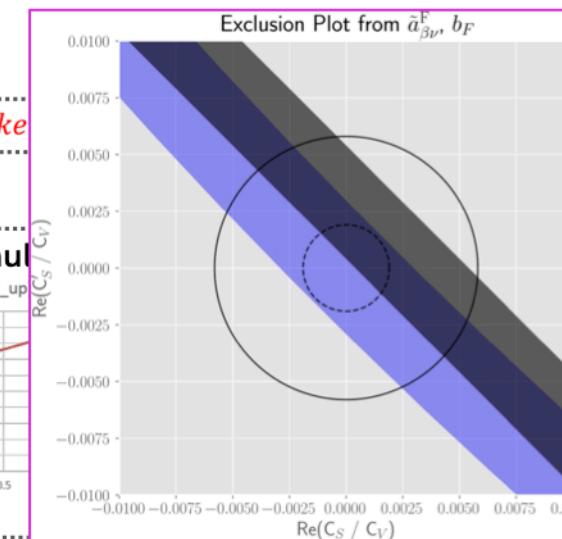
Monte Carlo simulation

MeanDeltaE ($E_{\text{down}} - E_{\text{up}}$)



+ systematics

\tilde{a}

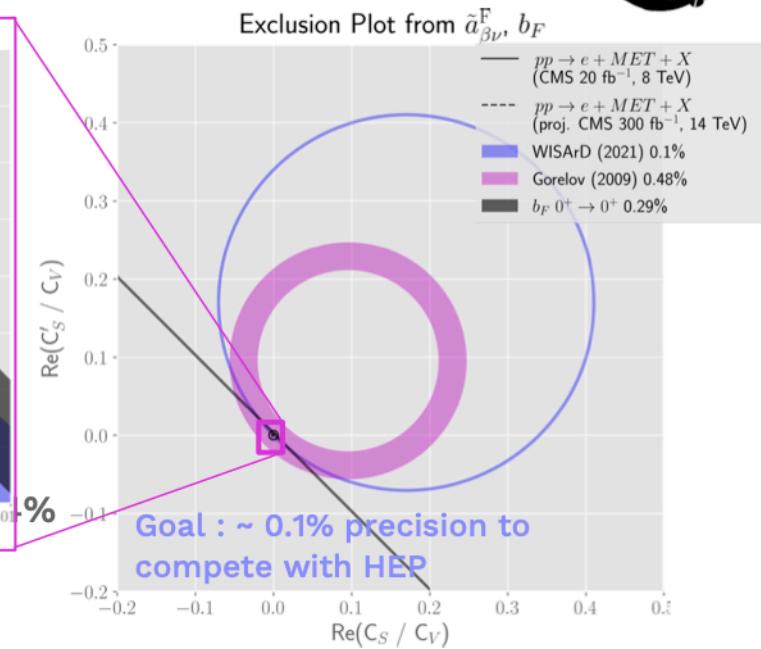


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Exclusion plot from D. Atanasov

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Proof-of-Principle Experiment - 2018



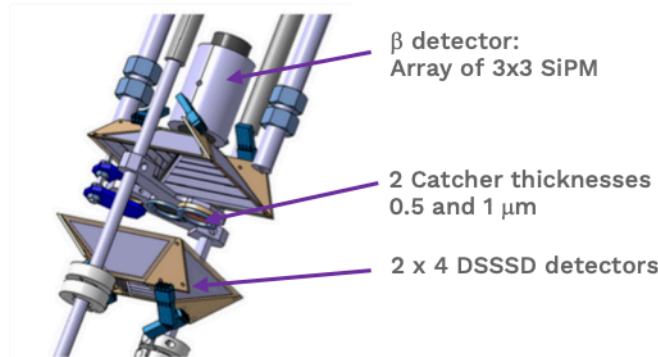
	Source	Uncertainty	$\Delta \tilde{a}_F (\times 10^{-3})$
Background	False coinc.	8%	<1
Proton	Det. calibration	0.2%	9
	Det. position	1 mm	<1
	Source position	3 mm	3
	Source radius	3 mm	1
	<i>B</i> field	1%	<1
Positron	Silicon dead layer	0.3 μm	5
	Mylar thickness	0.15 μm	2
	Detector backscattering	15%	2
	Catcher backscattering	15%	21
	Threshold	12 keV	8
Total			25

Upgrade 2019 - 2022



■ Setup geometry

- 57% maximum solid angle
- 90° p incident angle



Upgrade 2019 - 2022

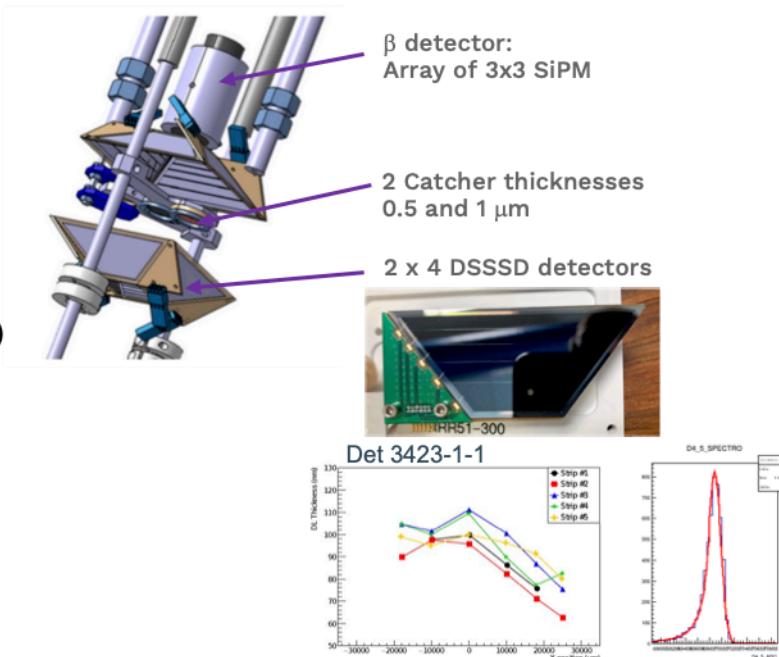


■ Setup geometry

- 57% maximum solid angle
- 90° p incident angle

■ p detectors

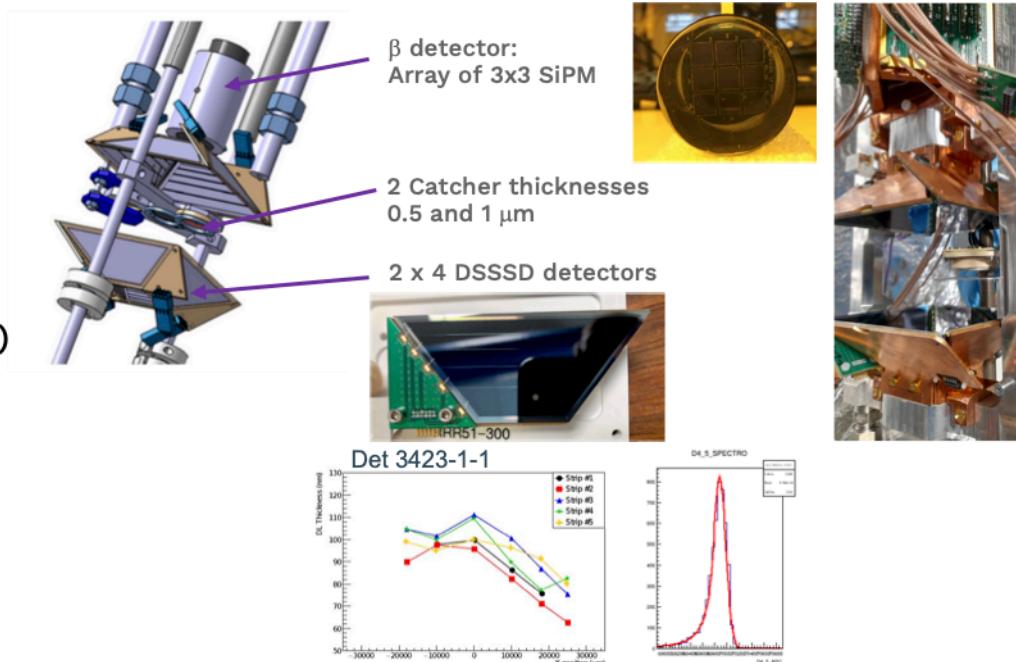
- Cooled -23°
- < 100 nm dead layer
α 700 keV
- 15-20 keV FWHM (α 5.3 MeV)



Upgrade 2019 - 2022



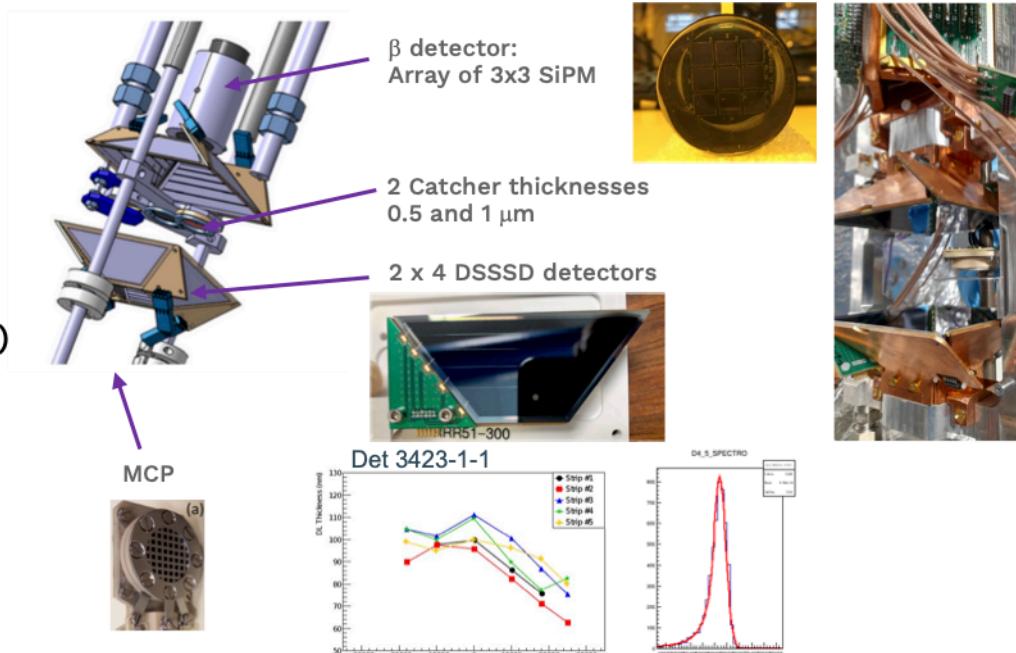
- **Setup geometry**
 - 57% maximum solid angle
 - 90° p incident angle
- **p detectors**
 - Cooled -23°
 - < 100 nm dead layer
α 700 keV
 - 15-20 keV FWHM (α 5.3 MeV)
- **β detector**
 - 130 keV FWHM @ 1 MeV
 - Dual gain output
 - Mult 3 threshold < 23 keV



Upgrade 2019 - 2022



- **Setup geometry**
 - 57% maximum solid angle
 - 90° p incident angle
- **p detectors**
 - Cooled -23°
 - < 100 nm dead layer
α 700 keV
 - 15-20 keV FWHM (α 5.3 MeV)
- **β detector**
 - 130 keV FWHM @ 1 MeV
 - Dual gain output
 - Mult 3 threshold < 23 keV
- **MCP + segmented FC**
 - Resolution < 300 μm FWHM
- **Beam line transmission : 90%**

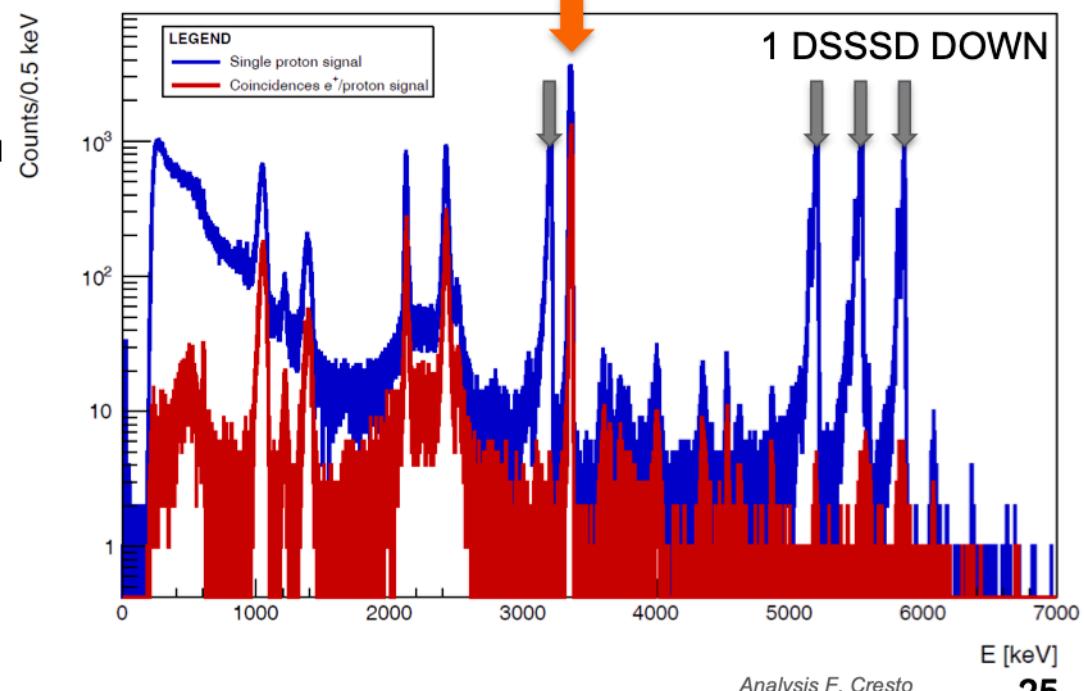


Test run - 2021



■ 43h of ^{32}Ar beam time

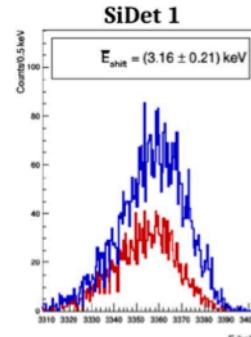
- 4- α source pollution
- Resolution 7-10 keV FWHM
- F + GT transitions



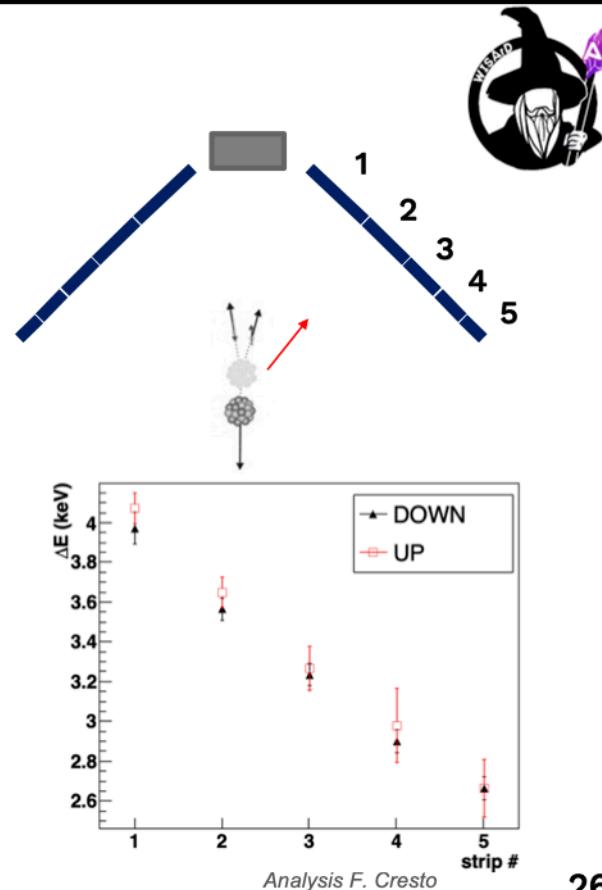
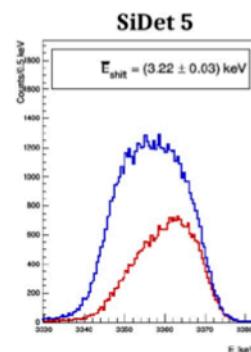
Test run - 2021

- 43h of ^{32}Ar beam time
 - 4- α source pollution
 - F + GT transitions
- Kinematic broadening visible
 - IAS peak 19 keV FWHM
- Kinematic shift (down)
 $\Delta E_F = 3.97(8) - 2.66(6)$
- Preliminary value of $\tilde{\alpha}$
 1.002 ± 0.017 (stat.)
◊ stat. precision level : 2%

Upper detectors



Lower detectors

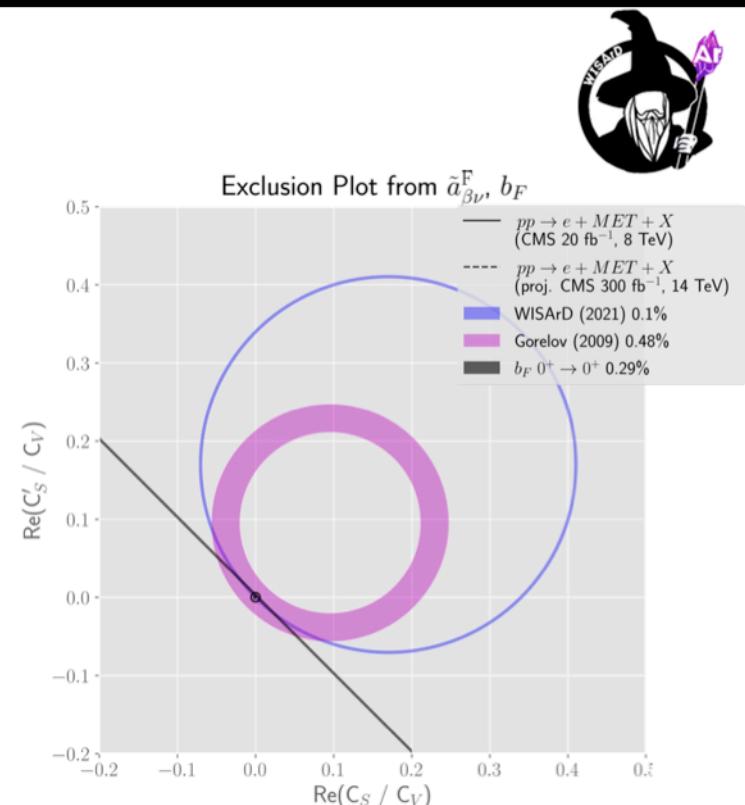


Next data taking : 2023

■ 24 shifts

- ~150h ^{32}Ar : 3×10^7 β -p events
- 2 shifts ^{33}Ar for calibration
- 3 shifts for beam tuning

	Source	Uncertainty	$\Delta \tilde{a}_F$	
Background	False coinc.	8%	<1	
Proton	Det. calibration	0.2%	9	■
	Det. position	1 mm	<1	■
	Source position	3 mm	3	■
	Source radius	3 mm	1	■
	B field	1%	<1	■
	Silicon dead layer	0.3 μm	5	■
	Mylar thickness	0.15 μm	2	■
Positron	Detector backscattering	15%	2	■
	Catcher backscattering	15%	21	■
	Threshold	12 keV	8	■
Total			25	



Exclusion plot from D. Atanasov

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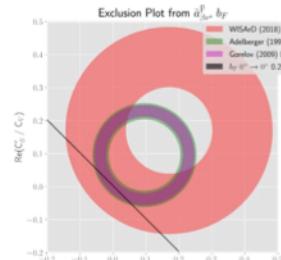
Conclusion



2018 2021 2023



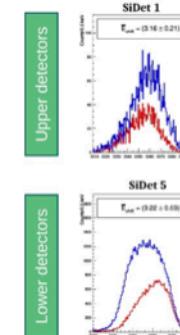
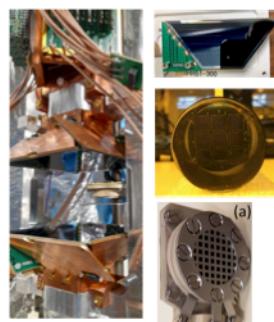
■ Proof-of-Principle



$$\Delta E_F = 4.49(3) \text{ keV}$$
$$\tilde{a}_F = 1.007(32)_{\text{stat}}(25)_{\text{syst}}$$

V. Araujo-Escalona et al. Phys. Rev. C 101 (2020)

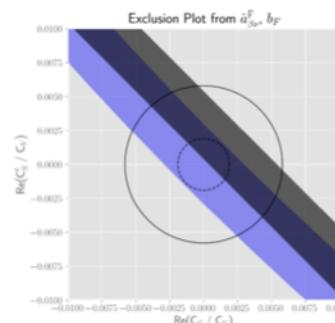
■ Upgrade



$$1.002 \pm 0.017 \text{ (stat.)}$$

D. Atanasov et al. NIM A 1050 168159 (2023)

■ Next data taking ^{32}Ar



outlook



2018

2021

2023

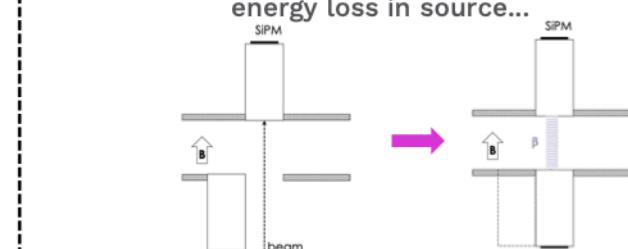
...

■ β spectrum shape measurement

$$dW = dW_0 \times \xi \left(1 + b \frac{m}{E_e} \right)$$

□ Challenges :

backscattering, dead layer,
energy loss in source...



Outlook



2018

2021

2023

...

■ β - 2α emitter : ${}^8\text{Li}$

- GT transition : T current?
- Best limit to date in trap (BPL)

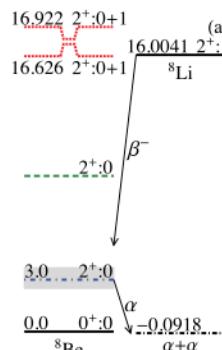
$$a_{\beta\nu} = -0.3325 \pm 0.0013_{\text{stat}} \pm 0.0019_{\text{syst}}$$

$$|C_T/C_A| < 0.087 \text{ at the 95.5\% C.L.}$$

M.T. Burkey et al. Phys. Rev. Lett. 128 (2022)

■ Other β -p candidates

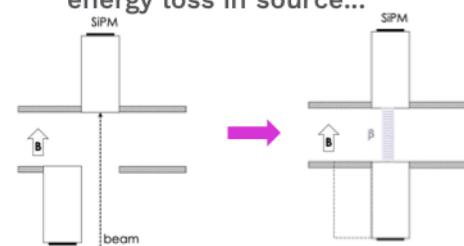
- ${}^{20}\text{Mg}$: sensitivity to $b \times 2$
- ${}^{24}\text{Si}, {}^{28}\text{S}$



■ β spectrum shape measurement

$$dW = dW_0 \times \xi \left(1 + b \frac{m}{E_e} \right)$$

- Challenges : backscattering, dead layer, energy loss in source...



Thank you for your attention

P.Alfaurt, D.Atanasov, P.Ascher, B.Blank, F.Cresto, L.Daudin, X.Fléchard, A.Husson, A.Garcia, M.Gerbaux, J.Giovinazzo,
S.Grévy, J.Ha, R.Lica, E.Liénard, D.Melconian, C.Mihai, M.Nasser, C.Neacsu, A.Ortega-Moral, M.Pomorski, M.Roche,
N.Severijns, S.Vanlangendonck, M.Versteegen, D.Zakoucky



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Back up

Test run 2021 : preliminary systematics



■ Evaluated

- Catcher thickness
- β detection threshold

For mylar (down) : 5 %o

■ To be evaluated

- Backscattering in catcher and at the surface of the β detector
- p spectrum calibration
- Implantation point position

HE vs HP

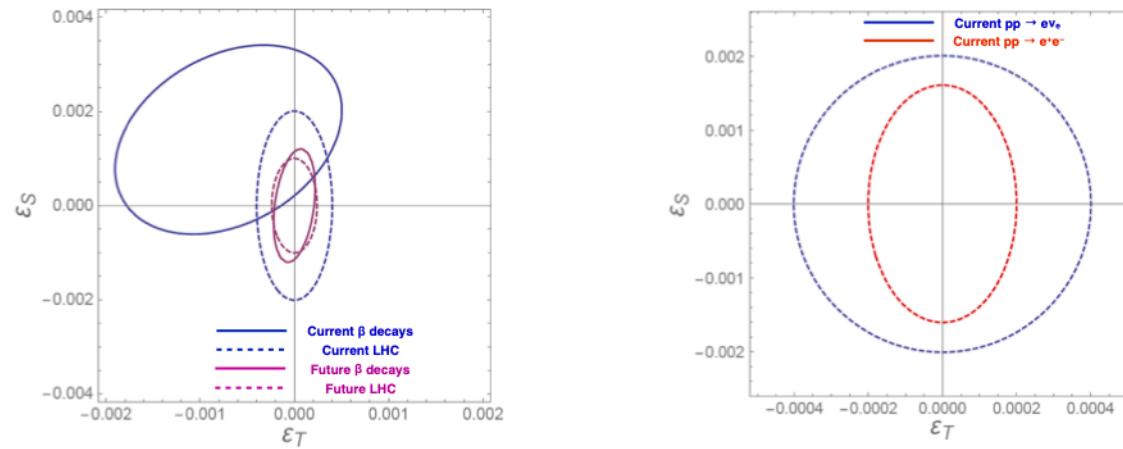


FIG. 8. Current and projected 90% C.L. constraints on ϵ_S and ϵ_T defined at 2 GeV in the $\overline{\text{MS}}$ scheme. (Left) The beta-decay constraints are obtained from the recent review article, Ref. [81]. The current and future LHC bounds are obtained from the analysis of the $pp \rightarrow e + MET + X$. We have used the ATLAS results [82], at $\sqrt{s} = 13$ TeV and integrated luminosity of 36 fb^{-1} . We find that the strongest bound comes from the cumulative distribution with a cut on the transverse mass at 2 TeV. The projected future LHC bounds are obtained by assuming that no events are observed at transverse mass greater than 3 TeV with an integrated luminosity of 300 fb^{-1} . (Right) Comparison of current LHC bounds from $pp \rightarrow e + MET + X$ versus $pp \rightarrow e^+e^- + X$.

R. Gupta et al Phys. Rev. D 98 (2018)

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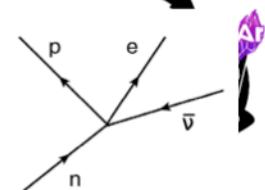
Beta decay

■ n beta decay Lagrangian

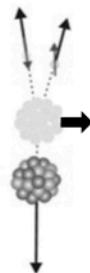
$$-\mathcal{L}_{LY} = \mathcal{C}_V \left(\bar{p} \gamma^\mu n + \frac{\mathcal{C}_A}{\mathcal{C}_V} \bar{p} \gamma^\mu \gamma_5 n \right) \times \bar{e} \gamma_\mu (1 - \gamma_5) \nu_e \\ + \mathcal{C}_S \bar{p} n \times \bar{e} (1 - \gamma_5) \nu_e + \frac{1}{2} \mathcal{C}_T \bar{p} \sigma^{\mu\nu} n \times \bar{e} \sigma_{\mu\nu} (1 - \gamma_5) \nu_e + hc$$

SM “V-A” structure

Exotic currents : S and T
P omitted



■ Decay rate distribution for polarized nuclei



$$\frac{dW(\mathbf{J})}{dE_e d\Omega_e d\Omega_\nu} = dW_0 \times \xi \left\{ 1 + a \frac{\mathbf{p}_e \cdot \mathbf{p}_\nu}{E_e E_\nu} + b \frac{m_e}{E_e} + \frac{<\mathbf{J}>}{J} \cdot \left(A \frac{\mathbf{p}_e}{E_e} + B \frac{\mathbf{p}_\nu}{E_\nu} + D \frac{\mathbf{p}_e \times \mathbf{p}_\nu}{E_e E_\nu} \right) \right\}$$

- Ft values : \mathcal{V}_{ud} , b
- Beta spectrum shape : b
- Correlation measurements : a , b , D

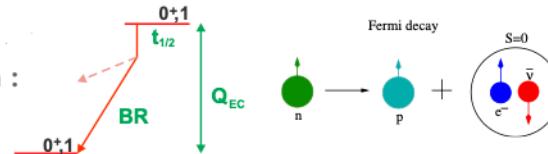
T.Lee, C-N Yang Phys. Rev. 104 (1956)
M. González-Alonso, Colloque GANIL (2019)
J.D Jackson, S.B Treiman, H.W Wyld Nuclear Phys 4 (1957)

35

Ft values : total decay rates

■ Unitarity test of the CKM matrix 1st row : $|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 1 + \text{NP}$

$0^+ \rightarrow 0^+$ superallowed Fermi transition :



Statistical rate function $f \propto \int dW_0$

$$t = \frac{t_{1/2}}{BR} (1 + P_{EC})$$

Corrections:
Radiative < 1%
Structure < 1%

$$\mathcal{F}t = \frac{K}{2G_F^2 V_{ud}^2 (1 + \Delta_R^V)}$$

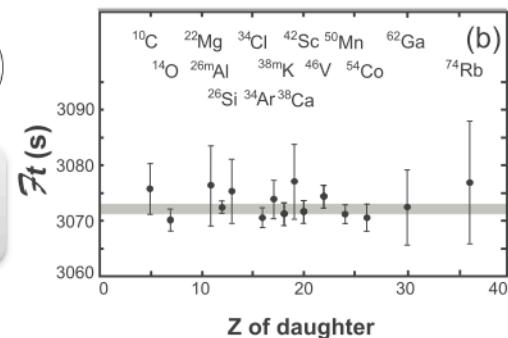
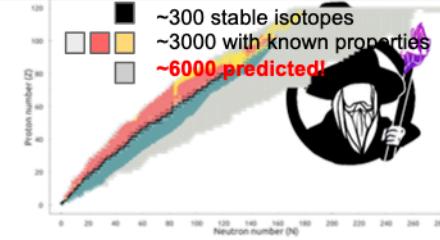
Trap : Q_{EC}
 $\Delta m \sim 10^{-8}$

Beta counting and Ge
with calibrated ε :
 $t_{1/2}$ and BR
 $\Delta\varepsilon \sim 0.2\%$

Theoretical
Calculations
uncertainties < 0.1%
(except ^{62}Ga & ^{74}Rb)

15 transitions with
uncertainties < 0.3%

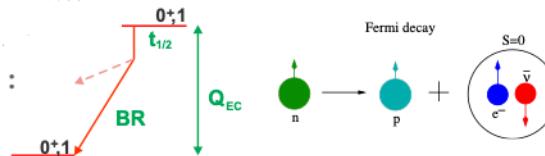
222 individual measurements from 23 decays : $|V_{ud}| = 0.97373 \pm 0.00031$



Ft values : total decay rates

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$0^+ \rightarrow 0^+$ superallowed Fermi transition :



$$\text{Statistical rate function } f \propto \int dW_0$$

$$\text{Partial half-life } t = \frac{t_{1/2}}{BR} (1 + P_{EC})$$

Corrections:
Radiative < 1%
Structure < 1%

$$f t = \frac{K}{2G_F^2 V_{ud}^2 (1 + \Delta_R^V)}$$

Trap : Q_{EC}
 $\Delta m \sim 10^{-8}$

Beta counting and Ge
with calibrated ε :
 $t_{1/2}$ and BR
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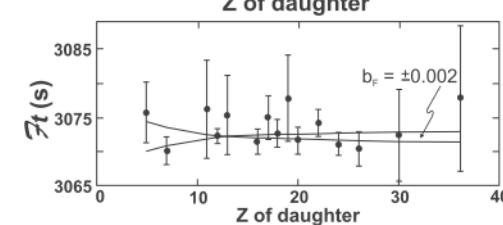
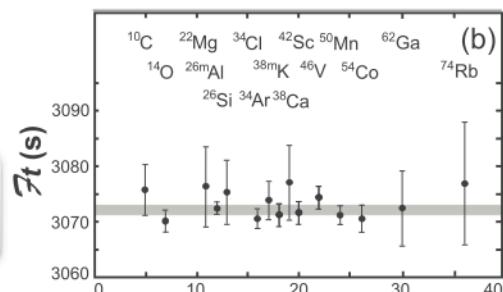
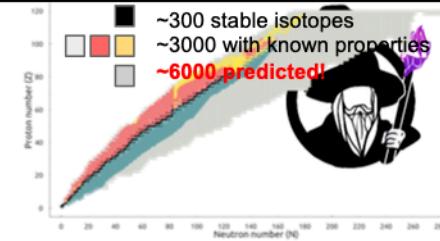
15 transitions with
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222 individual measurements from 23 decays : $|V_{ud}| = 0.97373 \pm 0.00031$

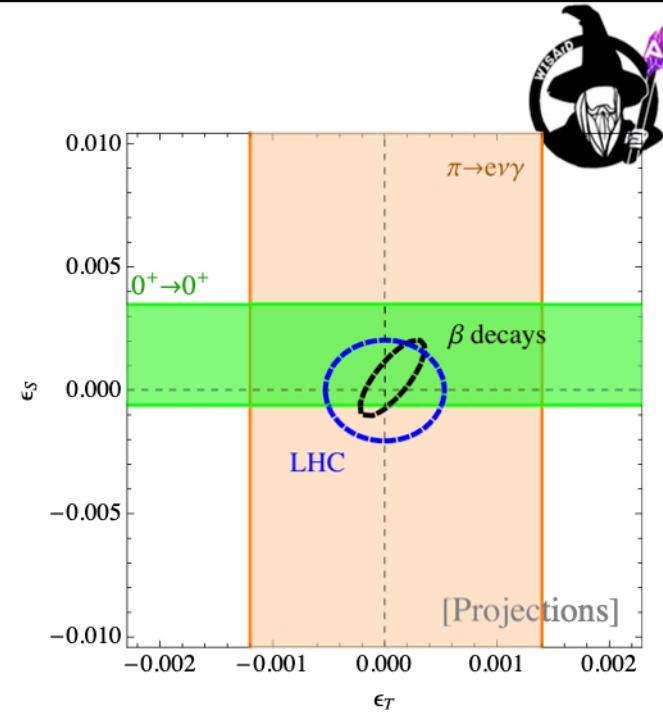
■ Sensitivity to exotic scalar currents : $b = \pm 0.002$

$$\text{If } b \neq 0 \text{ then } f \text{ is affected : } ft' = ft \times \frac{1}{1 + b < \frac{m_e}{E_e} >}$$

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J.C. Hardy & I.S. Towner Phys. Rev. C 102 (2020)



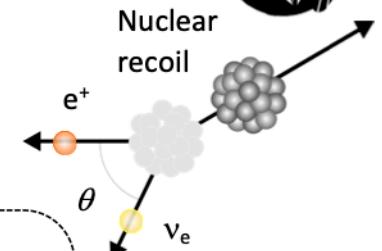
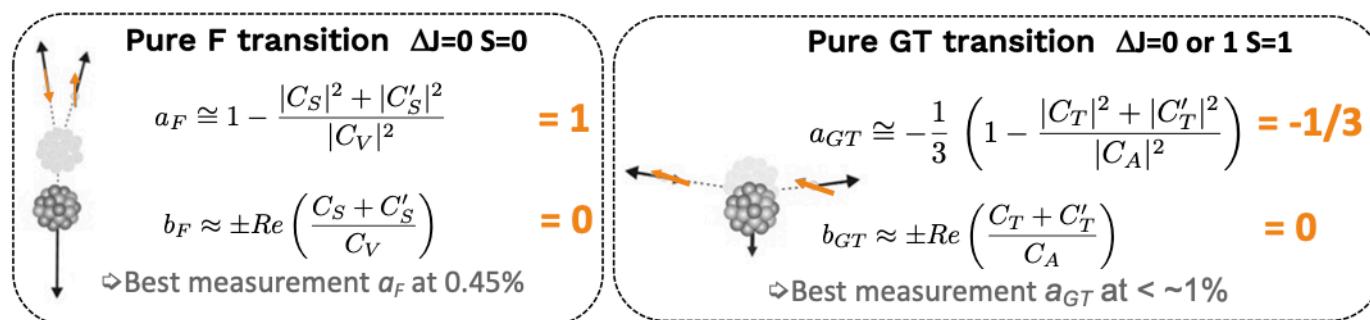
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Correlation measurements : WISArD



■ Decay rate for non polarized nuclei :

$$dW = dW_0 \left(1 + a \frac{\mathbf{p}_e \cdot \mathbf{p}_\nu}{E_e E_\nu} + b \frac{m_e}{E_e} \right)$$



Gorelov, A. et al, Phys. Rev. Lett. 94, 1425 (2005)
 Johnson, C.H. et al, Phys. Rev. 132, 1149 (1963)

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