

Rare B decays at LHCb (excluding b->sll)





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Highly **suppressed** in SM:

- Higher orders diagrams
- FCNC box or penguin diagrams
- \implies b \rightarrow sll, dedicated presentation later

signal different from SM? Possible new physics!

Rare B decays





LHCb detector for rare B decays

The LHC has a large cross section of b and c hadrons:

$$\sigma(b\bar{b})_{7\ TeV} = 295\ \mu b$$

•
$$\sigma(b\bar{b})_{13\ TeV} = 590\ \mu b$$

- LHCb designed as forward spectrometer to focus on bb production
- LHCb uses luminosity levelling:
 - Proton beams are defocused
 - Keeps run conditions more stable during fills
 - Reduces interactions per bunch crossing to 1-2
- All results presented today make used of the full available proton-proton collisions collected by the LHCb experiment:

Run 1:
• 1 fb⁻¹ at
$$\sqrt{s} = 7$$

• 2 fb⁻¹ at $\sqrt{s} = 8$



Run 2: $^{\circ}$ 6 fb⁻¹ at $\sqrt{s} = 13$ TeV TeV TeV



LHCb detector for rare B decays













More B decays in <u>presentation</u> by Jike Wang.

Rare B decays

- In <u>next presentation</u> by Sara Celani all about $b \rightarrow sll$.



LHCb-PAPER-2022-022 arXiv:2210.10412

Motivation:

- Lepton and baryon number violation (LNV and BNV).
- Forbidden in the SM with a prediction of

 $\mathscr{B}(\bar{b} \to uul^{-}) < 2.4 \times 10^{-27}$ [1].

Highly sensitive to new physics: GUT models. Never searched before.

[1] Phys. Rev. D 72 (2005) 095001

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First time!







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$$R_{\text{norm}} \equiv \frac{\mathcal{B}\left(B^0 \to K^+ \pi^-\right)}{\mathcal{B}(B^+ \to J/\psi\left(\mu^+ \mu^-\right)K^+)} = \frac{N_{B^0 \to K^+ \pi^-} \times \varepsilon_{B^+ \to J/\psi}}{N_{B^+ \to J/\psi\left(\mu^+ \mu^-\right)K^+} \times \varepsilon_{B^+ \to J/\psi}}$$

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Results:

- Signal extraction using fit on $p\mu^-$ mass.
- Not evidence found, limits at 90% (95%) CL:

$$\mathcal{B}(B^{0} \to p\mu^{-}) < 2.6 \ (3.1) \times 10^{-9}$$

$$\mathcal{B}(B^{0}_{s} \to p\mu^{-}) < 12.1 \ (14.0) \times 10^{-9}$$
First ULs on these decays!

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ρμ

S

/ (30 MeV/c²) 700 400 400

300

200

100

Candidates /

Mass distribution example for two MLP bins and one dataset:





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Motivation:

The two-body baryonic decays are suppressed [1,2]

Theoretical estimations disagree about the \mathscr{B} [3,4]. Waiting for experimental results to clarify.

Previous search of $B_s^0 \rightarrow p\bar{p}$ at LHCb with Run 1 [5]:

 $\mathscr{B}(B_s^0 \to p\bar{p}) < 1.5 \times 10^{-8}$ at 90% CL

[1] High Energy Physics **2022** (2022) 4343824 [3] JHEP **04** (2020) 035 [2] Phys. Rev. **D91** (2015) 036003

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 B_s^0

LHCb-PAPER-2022-004 arXiv:2206.06673

Strategy:

- Tight PID selection to reduce background from mis-identified hadrons.
- Multivariate technique: boosted decision tree (BDT), to reduce combinatorial background.
- Mass signal blinded until analysis procedure was finalised.
- Solution Normalisation channel: $B^0 \to K^+ \pi^-$

$$\mathcal{B}(B^0_{(s)} \to p\overline{p}) = \frac{N(B^0_{(s)} \to p\overline{p})}{N(B^0 \to K^+\pi^-)} \times \frac{\varepsilon_{B^0 \to K^+\pi^-}}{\varepsilon_{B^0_{(s)} \to p\overline{p}}} \times \overline{P}$$

 $\rightarrow p\bar{p}$







LHCb-PAPER-2022-004 arXiv:2206.06673

Results:

No evidence of the decay is found.

Upper limit on the branching fraction at 90 (95)% CL:

$$\mathcal{B}(B^0_s \rightarrow p\bar{p}) < 4.4~(5.1) \times 10^{-9}$$

Most stringent UL of this decay to date!









Motivation:

FCNC transition, no tree level ⇒ highly suppressed.
Theoretically predicted [1]:
\$\mathcal{B}(B^0 \rightarrow \mu^+ \mu^- \mu^+ \mu^-) = (0.4 - 4.0) × 10^{-12}\$
\$\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^- \mu^+ \mu^-) = (0.9 - 1.0) × 10^{-10}\$
Sensitive to new physics: MSSM sgoldstino [2].
Previous LHCb limits with Run 1 data at 95% CL [3]:
\$\mathcal{B}(B^0 \rightarrow \mu^+ \mu^- \mu^+ \mu^-) < 6.9 × 10^{-10}\$
\$\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^- \mu^+ \mu^-) < 2.5 × 10^{-9}\$

[1] Phys. Atom. Nucl. **81** (2018) 347 [2] Phys. Rev. **D85** (2012) 077701 [3] J. High Energ. Phys. **2017** (2017) 001

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 $\mu^{+}\mu^{-}\mu^{+}\mu^{-}$







Strategy:

- Three considerations, with specific selection requirements: Direct decay
 - Decay via BSM light scalar a [1,2]
 - Decay via J/ψ [3,4]: $\mathscr{B}(B^0 \to J/\psi \;(\mu^+\mu^-)\;\mu^+\mu^-) \sim 10^{-13}$ $\mathscr{B}(B_s^0 \to J/\psi \;(\mu^+\mu^-)\;\mu^+\mu^-) \sim 10^{-11}$
- Normalisation channel: $B_s^0 \to J/\psi(\mu^+\mu^-)\Phi(\mu^+\mu^-)$.
- Combinatorial bkg = muons that do not all originate from the same b-hadron decay.
- Multivariate classifier to reduce combinatorial bkg.

[1] Phys. Rev. Lett. **119** (2017) 031802

[2] JHEP **03** (2019) 008

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 $\mu^+\mu^-\mu^+\mu^-$



[3] Nucl. Phys. **B577** (2000) 240 [4] Prog. Theor. Exp. Phys. **2020** (2020) 083C01





Results:

- Applied Φ and J/ψ vetoes.
- So evidence found of the $B_{(s)}^0 \to \mu^+ \mu^- \mu^+ \mu^-$ decay.
- Limits at 95% CL:

$$\mathscr{B}(B^0 \to \mu^+ \mu^- \mu^+ \mu^-) < 1.8 \times 10^{-10}$$

$$\mathscr{B}(B_s^0 \to \mu^+ \mu^- \mu^+ \mu^-) < 8.6 \times 10^{-10}$$



 $\rightarrow \mu^+ \mu^- \mu^+ \mu^-$

Candidates / (50 MeV/ c^2)

2Ē

CLs

0.8

0.4

02

LHCb-PAPER-2021-039 JHEP 03 (2022) 109

a = BSM light narrow scalar resonance. In SM extension involving a new strongly interacting sector.

Two mutually exclusive pairs of opposite-sign muons satisfying:

$$|q_{ij}^2 - q_{kl}^2| < 2\sqrt{\sigma^2(q_{ij}^2) + \sigma^2(q_{kl}^2)}$$

Limits at 95% CL: 43

$$\mathscr{B}(B^0 \to a(\mu^+\mu^-)a(\mu^+\mu^-)) < 2.3 \times 10^{-10}$$

 $\mathscr{B}(B_s^0 \to a(\mu^+\mu^-)a(\mu^+\mu^-)) < 5.8 \times 10^{-10}$

 $\mu^+\mu^-\mu^+\mu^-$





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LHCb-PAPER-2023-004 arXiv:2210.10412

Motivation:

- Search for $D * (2007)^0 \rightarrow \mu^+ \mu^-$ [1].
- D^{*0} production ($b \rightarrow cd\bar{u}$) + $D^{*0} \rightarrow \mu^+\mu^-$ decay.
- The B decay allows better bkg. discrimination.
- SM prediction at $\mathcal{O}(10^{-19})$ level [2].
- Potential probe of physics beyond the SM.
- Limit by CMD-3 collaboration at 90% CL [3]: $\mathscr{B}(D^{*0} \to e^+e^-) < 1.7 \times 10^{-6}.$

[1] Eur. Phys. J. **C82** (2022) 459 [2] JHEP **11** (2015) 142

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[3] Phys. Atom. Nucl. 83 (2020) 954

LHCb-PAPER-2023-004 arXiv:2210.10412

Strategy:

- Displace vertex and exclusive final state provide powerful bkg. rejection.
- Use a BDT against combinatorial bkg.
- Solution Use PID info to suppress $K \to \pi$ and $hh \to \mu\mu$ misID bkgs.
- Signal signature = simultaneous peaks in $m(\mu^+\mu^-)$ and $m(\pi^-\mu^+\mu^-)$.
- Solution Normalisation channel $B^- \to J/\psi (\to \mu^+ \mu^-) K^-$.

$$\mathcal{B}\left(D^{*0} \to \mu^{+}\mu^{-}\right) = \frac{N_{D^{*0}\pi^{-}}}{N_{J/\psi K^{-}}} \cdot \frac{\varepsilon_{J/\psi K^{-}}}{\varepsilon_{D^{*0}\pi^{-}}} \cdot \frac{\mathcal{B}\left(B^{-} \to J/\psi K^{-}\right)}{\mathcal{B}\left(B^{-} \to D^{*0}\pi^{-}\right)}$$

 $B^- \rightarrow D^{*0}$ (

 $\rightarrow \mu^+\mu^-)\pi^-$





LHCb-PAPER-2023-004 arXiv:2210.10412

Results:

- First search for a rare charm-meson decay exploiting its production in beauty-meson decays.
- No excess with respect to the background-only hypothesis is observed.
- Upper limits at 90 (95)% CL:

$$\mathscr{B}(D^{*0} \to \mu^+ \mu^-) < 2.6 \ (3.4) \times 10^{-8}$$

Most stringent limit on
$$D^{*0}$$
 decays to leptons!
And first limit on D^{*0} to muons!

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$B^- \rightarrow D^{*0} (\rightarrow \mu^+ \mu^-) \pi^-$







In this presentation:

•
$$B_{(s)}^{0} \rightarrow p\mu^{-}$$
 First time!
• $B_{s}^{0} \rightarrow p\bar{p}$ Best limit!
• $B_{(s)}^{0} \rightarrow \mu^{+}\mu^{-}\mu^{+}\mu^{-}$ Most
• $B^{-} \rightarrow D^{*0}(\rightarrow \mu^{+}\mu^{-})\pi^{-}$

Rare B decays are dominated by statistical uncertainties, Run 3 is providing more statistics, and the LHCb upgrade (see presentation) will help us to push the limits of the SM.

stringent limits on each of the six decays! Most stringent limit on D^{*0} decays to leptons!

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New physics can be still hidden in rare B decays!
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Backup



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 $B^0_{(s)} \rightarrow p\mu^-$

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All mass distributions in MLP bins in Run 1:







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All mass distributions in MLP bins in Run 2:



Mass distributions for all BDT bins:



Figure 9: Distribution of the $\mu^+\mu^-\mu^+\mu^-$ invariant mass of candidates passing the $B_{(s)}^0 \to \mu^+ \mu^- \mu^+ \mu^-$ selection in (top left) the lowest BDT interval, (top right) the second lowest BDT interval and (bottom) the second highest BDT interval, with the fit models used to determine the branching fraction of $B^0 \rightarrow \mu^+ \mu^- \mu^+ \mu^-$ overlaid.

 $\rightarrow \mu^+ \mu^- \mu^+ \mu^-$



Figure 8: Distribution of the $\mu^+\mu^-\mu^+\mu^-$ invariant mass of candidates passing the $B^0_{(s)} \to \mu^+ \mu^- \mu^+ \mu^-$ selection in (top left) the lowest BDT interval, (top right) the second lowest BDT interval and (bottom) the second highest BDT interval, with the fit models used to determine the branching fraction of $B_s^0 \rightarrow \mu^+ \mu^- \mu^+ \mu^-$ overlaid.





$B^- \rightarrow D^{*0}$ (

LHCb-PAPER-2023-004 Submitted to Eur. Phys. J. C.

Differential branching fraction of the $B^+ \rightarrow \pi^+ \mu^+ \mu^-$ decay as a function of q2. The hashed regions correspond to theoretical predictions.

[1] Eur. Phys. J. **C82** (2022) 459

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 $\rightarrow \mu^+ \mu^-)\pi^-$





$B^- \rightarrow I$

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$$\mathcal{B}(D^{*0} \to \mu^+ \mu^-) = (-1.06 \pm 1.85) \times 10^{-8}$$

$$\mathcal{B}(D^{*0} \to \mu^+ \mu^-) < 2.6 (3.4) \times 10^{-8} \text{ at } 90 (95)\% \text{ CL}$$

Component	Yield
$B^- \to D^{*0}(\mu^+\mu^-) \pi^-$	-2 ± 3
$B^- \rightarrow \pi^- \mu^+ \mu^-$	17 ± 7
$B^- \rightarrow K^- \mu^+ \mu^-$	17 ± 8
Combinatorial bkg.	90 ± 13

Figure 2: Two-dimensional distribution of $\mu^+\mu^-$ invariant mass versus $\pi^-\mu^+\mu^-$ invariant mass for the selected $B^- \to D^{*0}(\mu^+\mu^-)\pi^-$ candidates. The red box corresponds to a range of about $\pm 3\sigma$ around the expected signal peak position in each dimension.

Table 2: Input parameters used in the estimation of the $D^{*0} \rightarrow \mu^+ \mu^-$ branching fraction. The uncertainties correspond to the statistical and systematic uncertainties added in quadrature.

Parameter	Value
$\mathcal{B}(B^- \to J/\psi K^-)$	$(10.20 \pm 0.19) \times 10^{-4} [25]$
$\mathcal{B}\left(B^{-} \rightarrow D^{*0}\pi^{-} ight)$	$(4.90 \pm 0.17) \times 10^{-3} [25]$
$\mathcal{B}\left(J\!/\psi\! ightarrow\mu^+\mu^- ight)$	$(59.61 \pm 0.33) \times 10^{-3} [25]$
$arepsilon_{J\!/\psi K^-}/arepsilon_{D^{*0}\pi^-}$	1.21 ± 0.03
$N_{J\!/\psi K^-}$	$(2316\pm8) imes10^3$







 D^*



0 vs. D^{0}

D^0

spin-0 \Rightarrow no γ mediator

Receives two contributions within SM:

Short distance: $\mathcal{B}(D^0 \to \mu^+ \mu^-) < 10^{-18}$ (*Z*-penguins, *W*-boxes) Long distance : $\mathcal{B}(D^0 \to \mu^+ \mu^-) < 10^{-11}$ ($D^0 \to \gamma \gamma$ transitions)



 $\mathscr{B}(D^0 \to \mu^+ \mu^-) < 2.94 \ (3.25) \times 10^{-9} \text{ at } 90 \ (95) \% \text{ CL}$

LHCb-PAPER-2022-029

