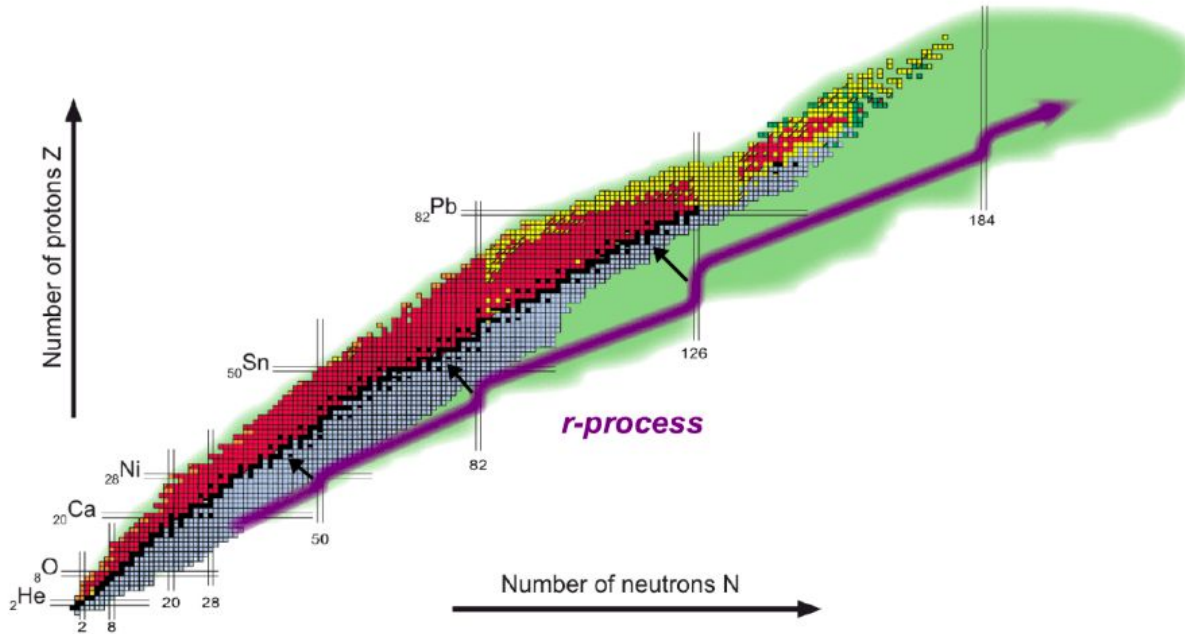


Constructing nuclear functionals for neutron stars and nucleosynthesis applications

Guilherme Grams,
Wouter Ryssens, Guillaume Scamps,
Stephane Goriely and Nicolas Chamel
Astronomy and Astrophysics Institute (IAA) - ULB

Motivation

Nucleosynthesis



Neutron star merger



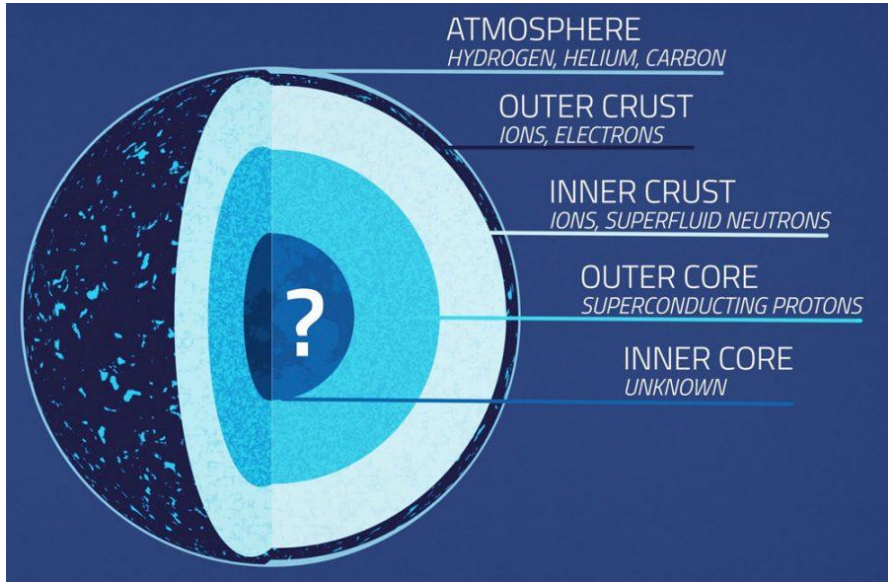
Supernova explosion



Description of **nuclear masses** and **fission** barriers in regions *unknown experimentally*.

Motivation

Neutron star physics



Nuclear physics inputs:

- **nuclear masses** for the NS crust;
- **infinite nuclear matter (INM) properties** for the NS core and neutron fluid on inner crust.

The road to here

Brussels
mass models

AN ENERGY DENSITY NUCLEAR MASS FORMULA

(I). Self-consistent calculation for spherical nuclei

F. TONDEUR

Physique Nucléaire Théorique, Université Libre de Bruxelles, Campus de la Plaine, Cp 229, 1050 Bruxelles

Received 2 December 1977

The road to here

Brussels
mass models

HFB-1 mass model
BSk-1 Skyrme interaction

HFB-32 mass model
BSk-32 Skyrme interaction

AN ENERGY DENSITY NUCLEAR MASS FORMULA

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Received 2 December 1977

A Hartree–Fock–Bogoliubov mass formula

M. Samyn^a, S. Goriely^{a,*}, P.-H. Heenen^b, J.M. Pearson^c, F. Tondeur^d

^a *Institut d'Astronomie et d'Astrophysique, ULB-CP226, 1050 Brussels, Belgium*

^b *Service de Physique Nucléaire Théorique, ULB-CP229, 1050 Brussels, Belgium*

^c *Dépt. de Physique, Université de Montréal, Montréal, PQ, Canada H3C 3J7*

^d *Institut Supérieur Industriel de Bruxelles, 1000 Brussels, Belgium*

Received 28 June 2001; revised 11 September 2001; accepted 25 September 2001

PHYSICAL REVIEW C **93**, 034337 (2016)

Further explorations of Skyrme-Hartree-Fock-Bogoliubov mass formulas. XVI. Inclusion of self-energy effects in pairing

S. Goriely,¹ N. Chamel,¹ and J. M. Pearson²

¹*Institut d'Astronomie et d'Astrophysique, CP-226, Université Libre de Bruxelles, 1050 Brussels, Belgium*

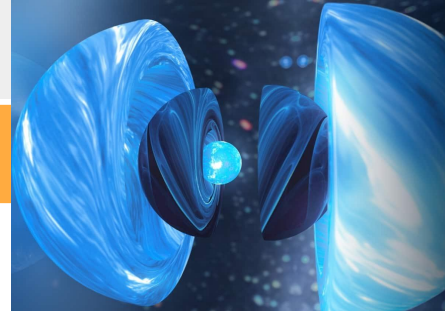
²*Département de Physique, Université de Montréal, Montréal (Québec), H3C 3J7 Canada*

The BSks functionals

From BSk/HFB series

- Skyrme EDF + HFB method;
- fitted to essentially all experimental masses;
- constraints on infinite nuclear matter properties;
- extended Skyrme functional to obtain stiff neutron matter equation of state;
- realistic pairing with self-energy corrections.

The BSk functionals



Astrophysics interest.

From BSk/HFB series

- Skyrme EDF + HFB method;
- fitted to essentially all experimental masses;
- constraints on infinite nuclear matter properties;
- extended Skyrme functional to obtain stiff neutron matter equation of state;
- realistic pairing with self-energy corrections.

- Description of **masses** of nuclei in NS **outer-crust**.
- INM, e.g., the symmetry energy coefficient J , and its slope L , are crucial for many NS properties, such as **crust-core transition and NS radius**.
- Necessary to describe **heavy pulsars**.
- Important for the description of **superfluids in NS**.

See N. Chamel and V. Allard talks.

The road to here

Brussels
mass models

AN ENERGY DENSITY NUCLEAR MASS FORMULA

(I). Self-consistent calculation for spherical nuclei

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HFB-1 mass model
BSk-1 nuclear force

New BSkG family:

Skyrme-Hartree-Fock-Bogoliubov mass models on a 3D mesh: effect of triaxial shape

Guillaume Scamps^{1,a}, Stephane Goriely¹, Erik Olsen¹, Michael Bender², Wouter Ryssens^{1,3}

¹ Institut d'Astronomie et d'Astrophysique, Université Libre de Bruxelles, Campus de la Plaine CP 226, 1050 Brussels, Belgium

² Université de Lyon, Université Claude Bernard Lyon 1, CNRS, IP2I Lyon / IN2P3, UMR 5822, 69622 Villeurbanne, France

³ Center for Theoretical Physics, Sloane Physics Laboratory, Yale University, New Haven, CT 06520, USA

Received: 7 September 2021 / Accepted: 26 November 2021

The BSkG functionals

BSkG1

EPJA 57, 333 (2021)

**Skyrme-Hartree-Fock-Bogoliubov mass models on a 3D mesh:
effect of triaxial shape**

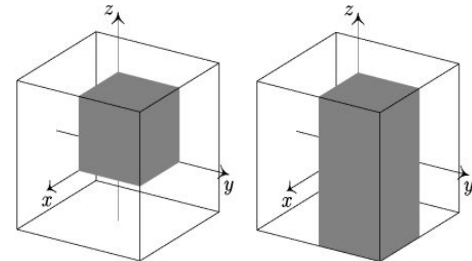
Guillaume Scamps^{1,a}, Stephane Goriely¹, Erik Olsen¹, Michael Bender², Wouter Ryssens^{1,3}

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- Brussels-Skyrme-on-the-grid (BSkG);



The BSkG functionals

BSkG1

EPJA 57, 333 (2021)

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Guillaume Scamps^{1,a}, Stéphane Goriely¹, Erik Olsen¹, Michael Bender², Wouter Ryssens^{1,3}

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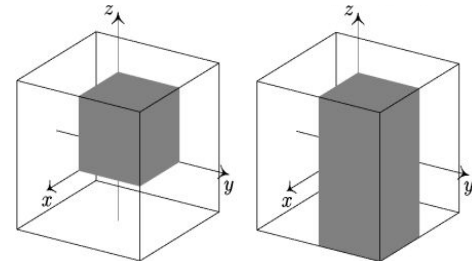
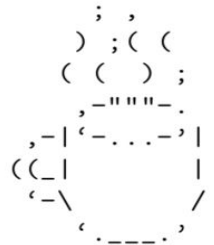
² Université de Lyon, Université Claude Bernard Lyon 1, CNRS, IP2I Lyon / IN2P3, UMR 5822, 69622 Villeurbanne, France

³ Center for Theoretical Physics, Sloane Physics Laboratory, Yale University, New Haven, CT 06520, USA

- Brussels-Skyrme-on-the-grid (BSkG);
- HFB code based coordinate space in 3D mesh -> **MOCCa** code (Ryssens et. al., EPJA 55, 93 (2019));
- better control of numerical accuracy;

Modular Cranking Code = MOCCa

HFB solver at **3D coordinate-space on Lagrangian mesh** from
W. Ryssens et al,
[W. Ryssens PhD Thesis, ULB (2016).]



The BSkG functionals

BSkG1

EPJA 57, 333 (2021)

Skyrme-Hartree-Fock-Bogoliubov mass models on a 3D mesh:
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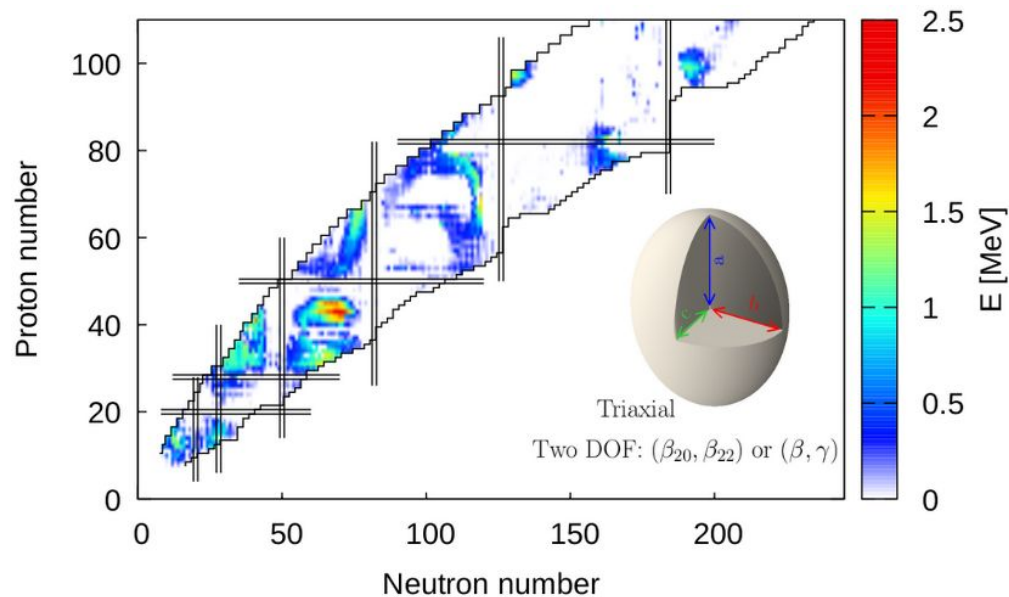
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- Brussels-Skyrme-on-the-grid (BSkG);
- HFB code based coordinate space in 3D mesh -> MOCCa code (Ryssens et. al., EPJA 55, 93 (2019));
- better control of numerical accuracy;
- allow **triaxial deformation**.



The BSkG functionals

BSkG1

EPJA 57, 333 (2021)

**Skyrme-Hartree-Fock-Bogoliubov mass models on a 3D mesh:
effect of triaxial shape**

Guillaume Scamps^{1,a}, Stephane Goriely¹, Erik Olsen¹, Michael Bender², Wouter Ryssens^{1,3}

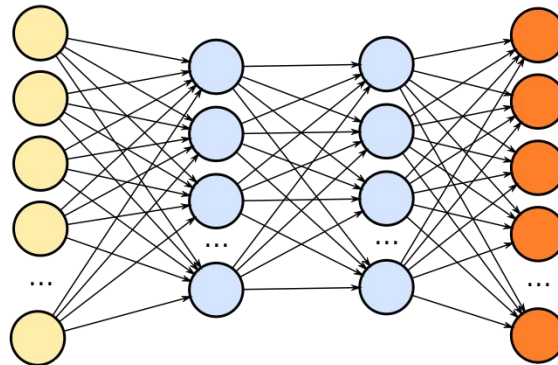
¹ Institut d'Astronomie et d'Astrophysique, Université Libre de Bruxelles, Campus de la Plaine CP 226, 1050 Brussels, Belgium

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- Brussels-Skyrme-on-the-grid (BSkG);
- HFB code based coordinate space in 3D mesh -> MOCCa code (Ryssens et. al., EPJA 55, 93 (2019));
- better control of numerical accuracy;
- allow **triaxial deformation**.
- *machine learning* to accelerate the fit.

machine learning as emulator of MOCCa.
MOCCa predictions for one nucleus ~ 20 minutes.
Machine learning prediction for one nucleus ~ a few seconds.



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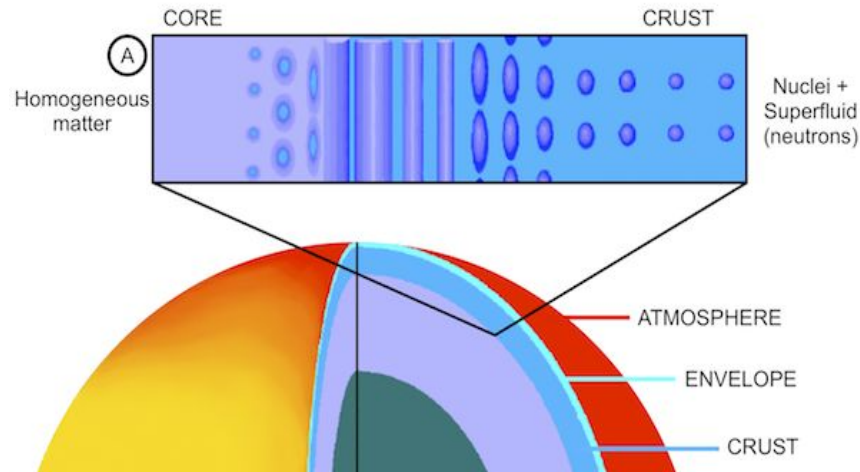
The BSkG3 functional

Astrophysics interest.

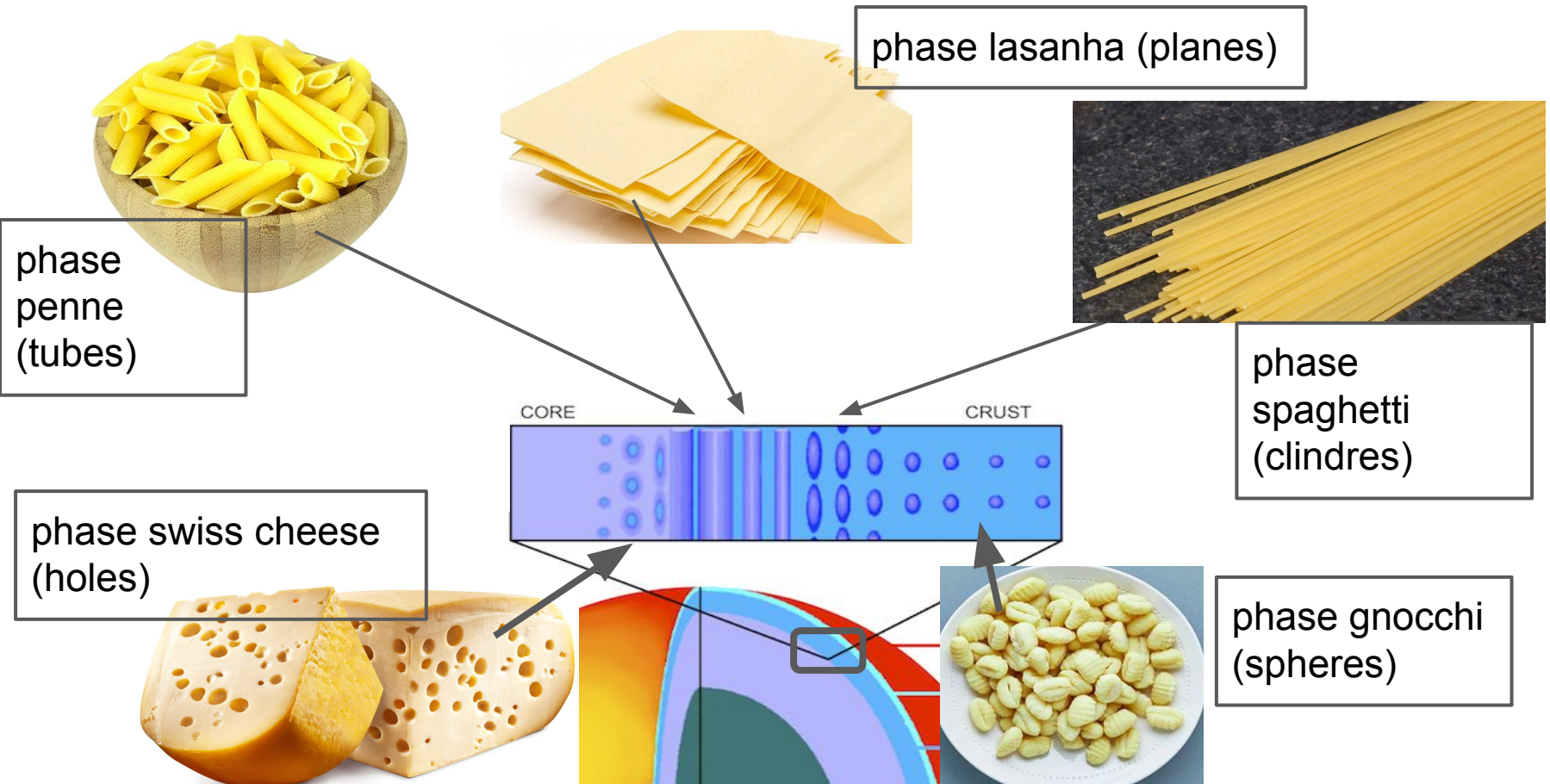
From BSkG series

- Improves the description of **masses** of nuclei in NS **outer-crust**.
- Nuclear **deformations** are important at the deep layers of the NS inner-crust, where **pasta phases** could be formed.

- 3D coordinate-space on Lagrangian mesh (high numerical accuracy);
- triaxial deformation;
- breaks time-reversal symmetry (time-odd terms);
- fission properties included in the fit.



Pasta phases : see talk of N. Shchepochin.



The BSkG3 functional

Connects the best features of Brussels functionals.

From BSk/HFB series

- Skyrme EDF + HFB method;
- fitted to essentially all experimental masses;
- constraints on infinite nuclear matter properties;
- extended Skyrme functional to obtain stiff neutron matter equation of state;
- realistic pairing with self-energy corrections.

From BSkG series

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The BSkG3 functional

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- Skyrme EDF + HFB method;
- fitted to essentially all experimental masses;
- constraints on infinite nuclear matter
- extended Skyrme functional to obtain stiff neutron matter equation of state;
- realistic pairing with self-energy corrections.

Not present in previous BSkG models.

From BSkG series

- 3D coordinate-space on Lagrangian mesh (high numerical accuracy);
- triaxial deformation;
- breaks time-reversal symmetry (time-odd terms);
- fission properties included in the fit.

The BSkG3 mass model

Reproduction of known nuclear masses.

Root-mean-square
(rms) deviation and
mean deviation
calculated from the
sum over: $O_{\text{exp}} - O_{\text{th}}$

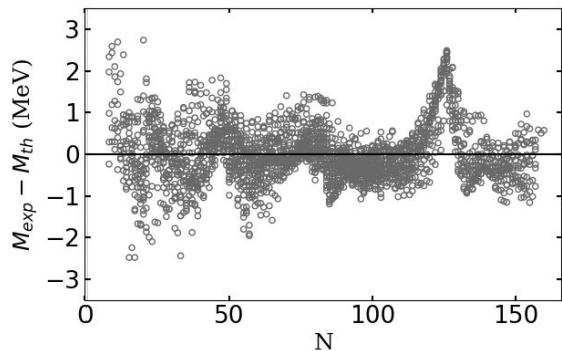
$$\sigma(M) = 0.681 \text{ MeV}$$

$$\varepsilon(M) = -0.041 \text{ MeV}$$

$$\sigma(R_C) = 0.0250 \text{ MeV}$$

$$\varepsilon(R_C) = -0.0045 \text{ MeV}$$

PRELIMINARY!



All 2457 nuclei
in the AME2020
nuclear chart.

The BSkG3 mass model

Reproduction of known nuclear masses.

Root-mean-square (rms) deviation and mean deviation calculated from the sum over: $O_{\text{exp}} - O_{\text{th}}$

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Comparison:

BSkG2

$$\sigma (M) = 0.678 \text{ MeV}$$

BSkG1

$$\sigma (M) = 0.741 \text{ MeV}$$

SLy4

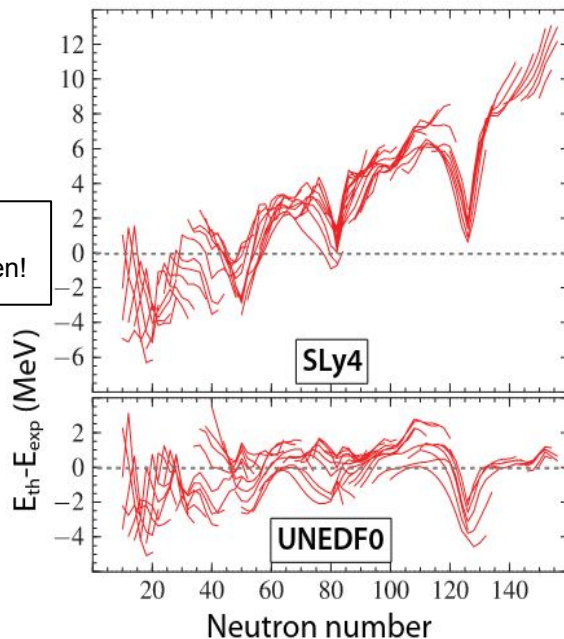
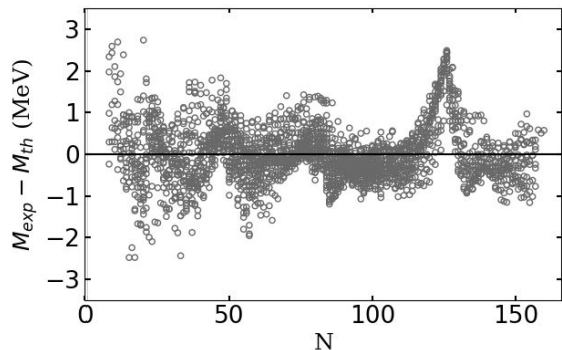
$$\sigma (M) = 4.80 \text{ MeV}$$

UNEDF0

$$\sigma (M) = 1.45 \text{ MeV}$$

Just 520 even-even!

All 2457 nuclei in the AME2020 nuclear chart.



The BSkG3 mass model

Reproduction of known nuclear masses.

Root-mean-square (rms) deviation and mean deviation calculated from the sum over: $O_{\text{exp}} - O_{\text{th}}$

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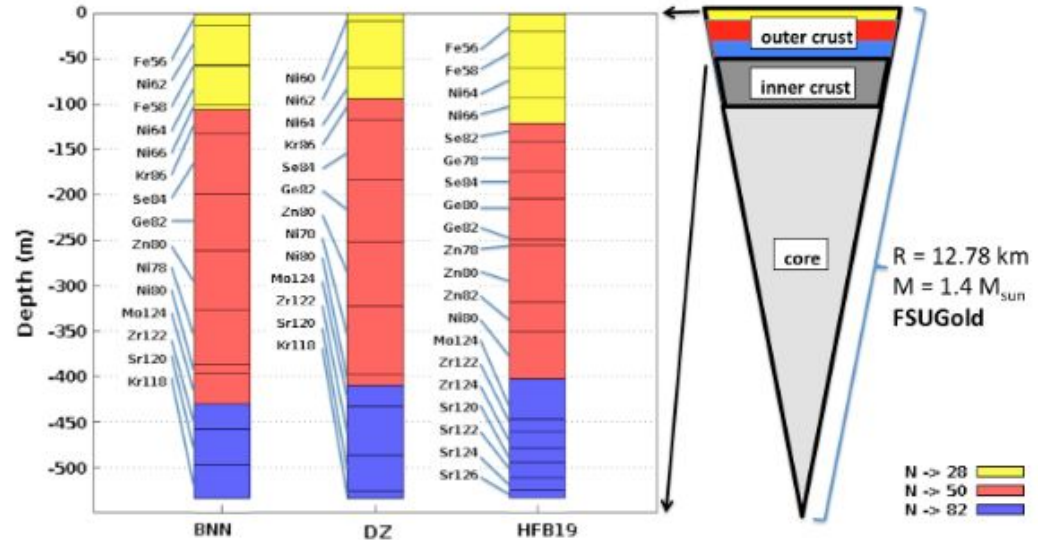
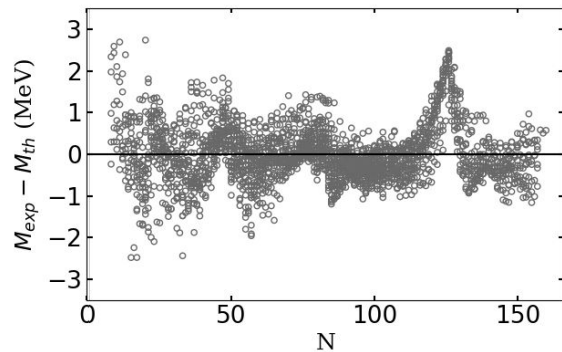
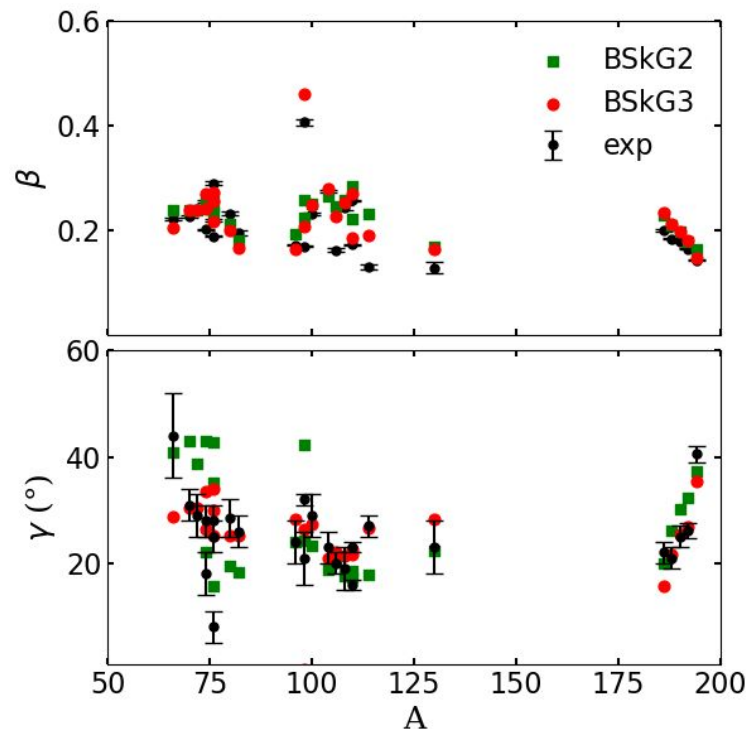
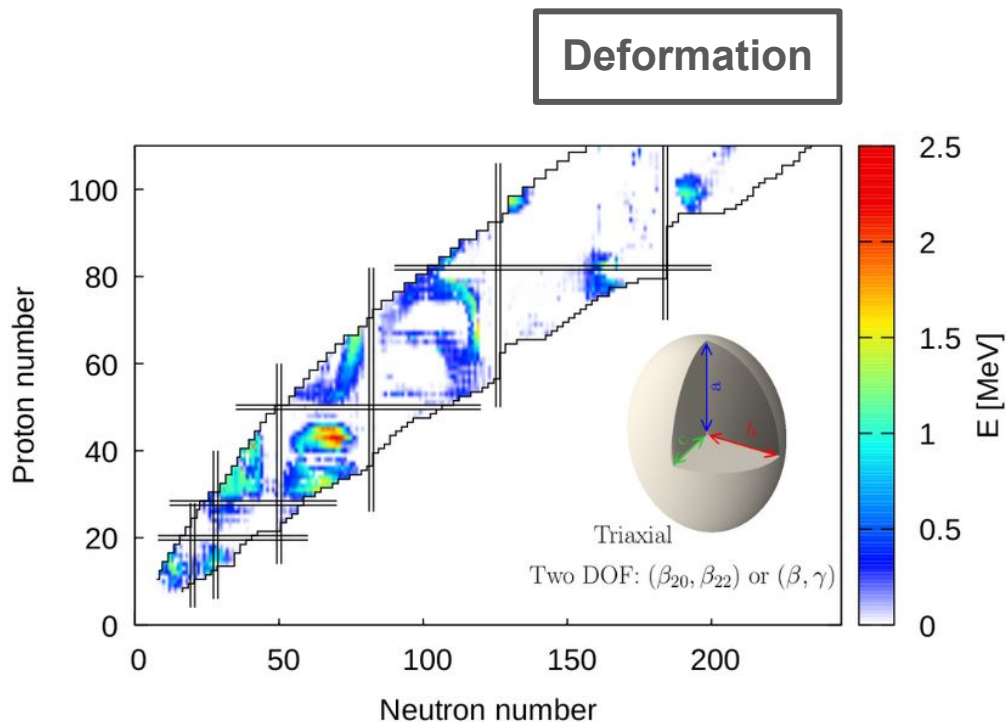


Fig. 4. Composition of a canonical $1.4 M_{\odot}$ neutron star with a 12.78 km radius as predicted by three mass models: “BNN-world”, DZ, and HFB19.

Figure from: J. Piekarewicz & R. Utama, Acta Phys. Pol. B 47, 659 (2016)

The BSkG3 mass model



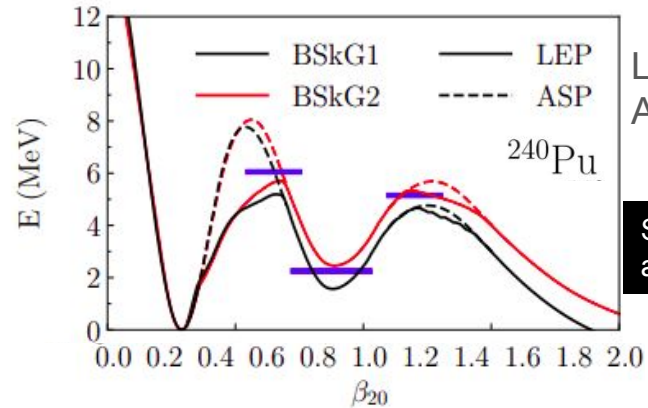
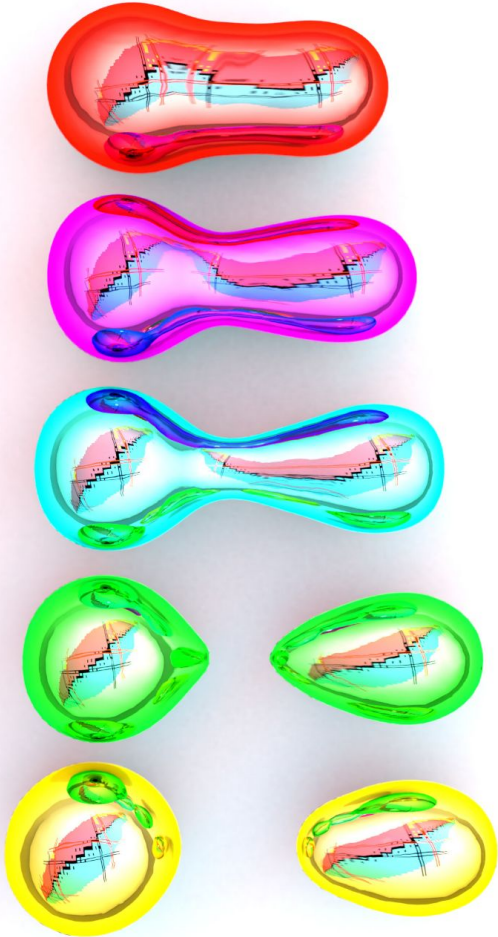
Binding energy difference between **unrestricted** calculations and calculations **restricted** to axial quadrupole deformation with BSkG1.

Quadrupole **deformation** (top) and **triaxiality** angle (bottom) for BSkG2 (green), BSkG3 (red) and experimental data (black).

The BSkG3 mass model

Fission barriers

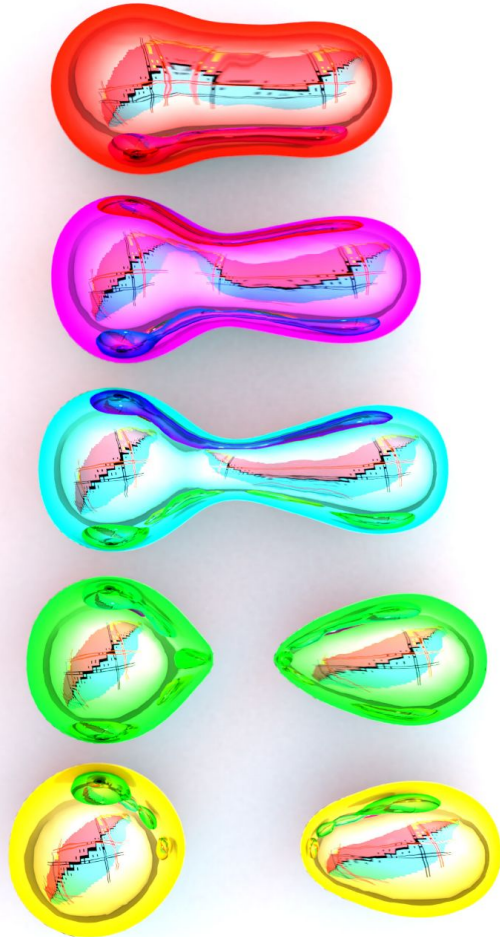
Primary (E_{\perp}), secondary (E_{\parallel}) **fission barrier** heights and fission isomer excitation energies (E_{iso}) of actinide nuclei.



LEP = Low Energy Path
ASP = Axial Symmetric Path

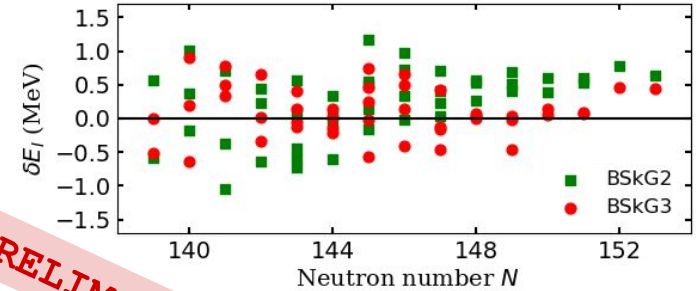
See W. Ryssens et al.
arXiv: 2302.03097

The BSkG3 mass model



Fission barriers

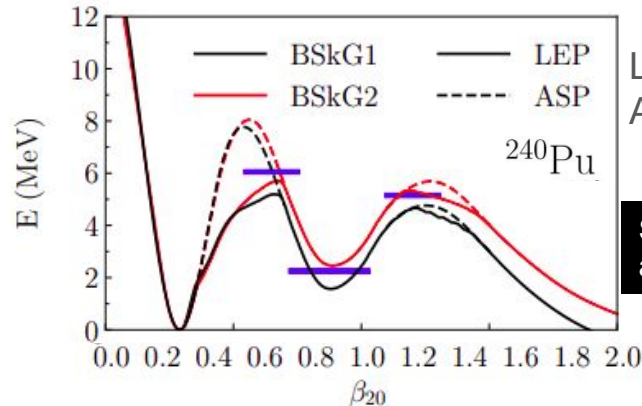
Primary (E_I), secondary (E_{II}) *fission barrier* heights and fission isomer excitation energies (E_{iso}) of actinide nuclei.



PRELIMINARY!

$$\sigma(E_I) = 0.48 \text{ MeV}$$

$$\varepsilon(E_I) = +0.10 \text{ MeV}$$

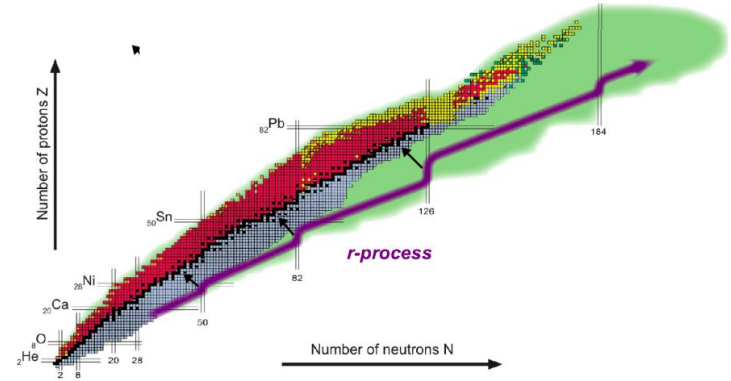
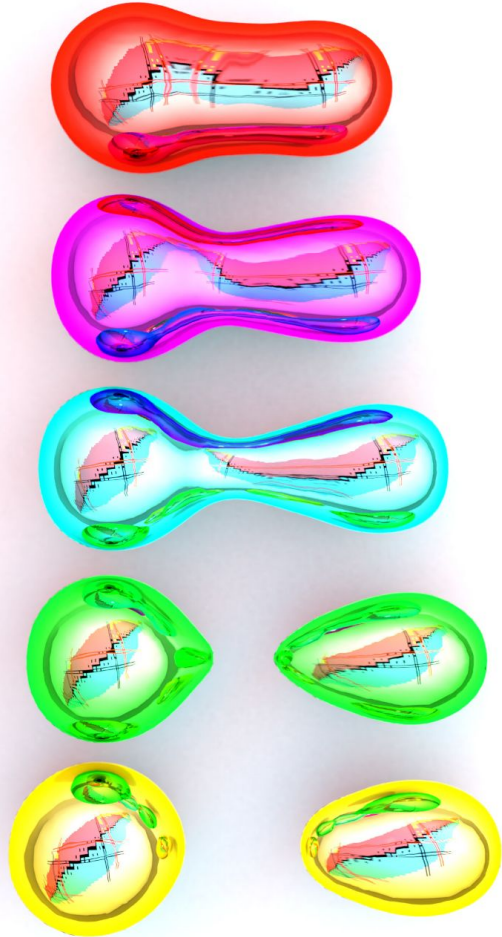


LEP = Low Energy Path
ASP = Axial Symmetric Path

See W. Ryssens et al.
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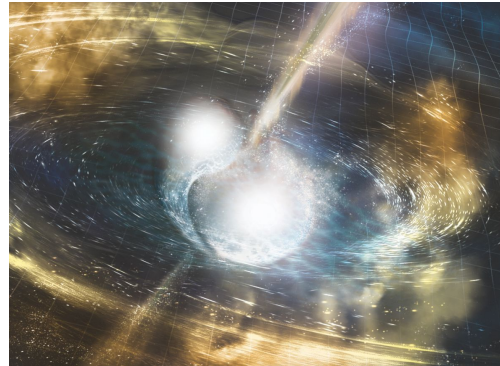
The BSkG3 mass model

Fission barriers



Fission properties impact several aspects of the r-process such as:

- "fission recycling";
- the **r-process abundances** in the $110 \leq A \leq 170$ region;
- the production of cosmic chronometers such as Th and U;
- the **heating rate of kilonovae**.

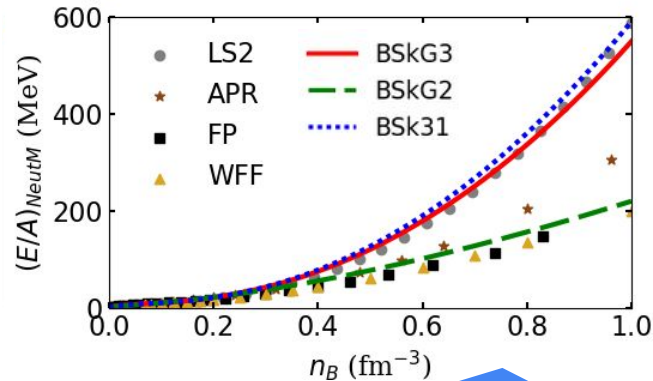


The BSkG3 interaction

Neutron matter (NeutM) energy

Constraint in *high density* neutron matter energy included in the fit protocol

Neutron matter energy



important for describing massive neutron star

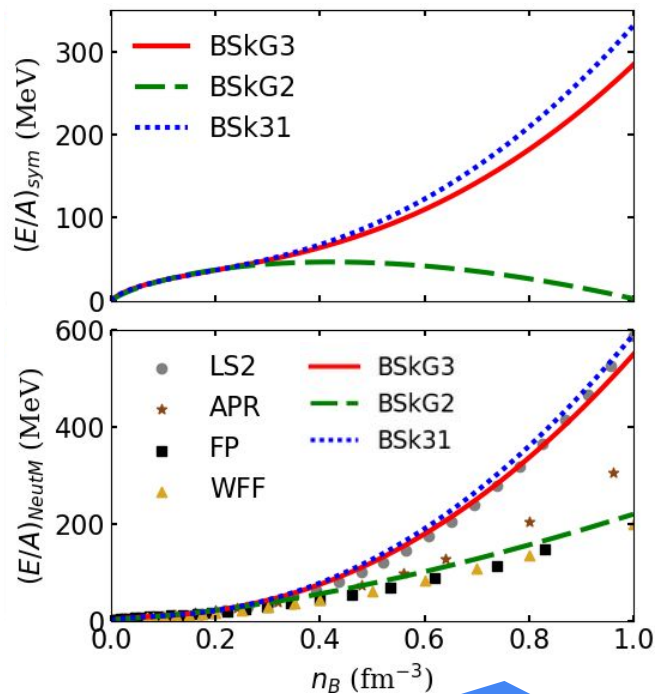
The BSkG3 interaction

Neutron matter (NeutM) & symmetry energy

Symmetry energy:

$$e_{\text{NeutM}} - e_{\text{SM}}$$

Neutron matter energy



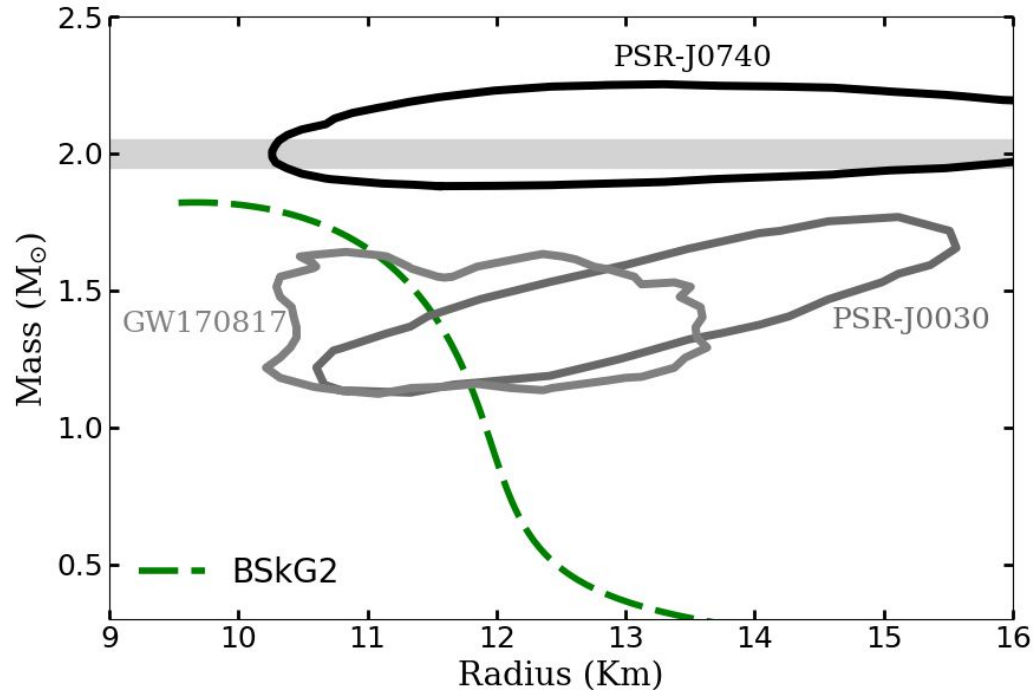
important for describing massive neutron star

The BSkG3 interaction

nuclear
EoS

TOV equations
(General Relativity)

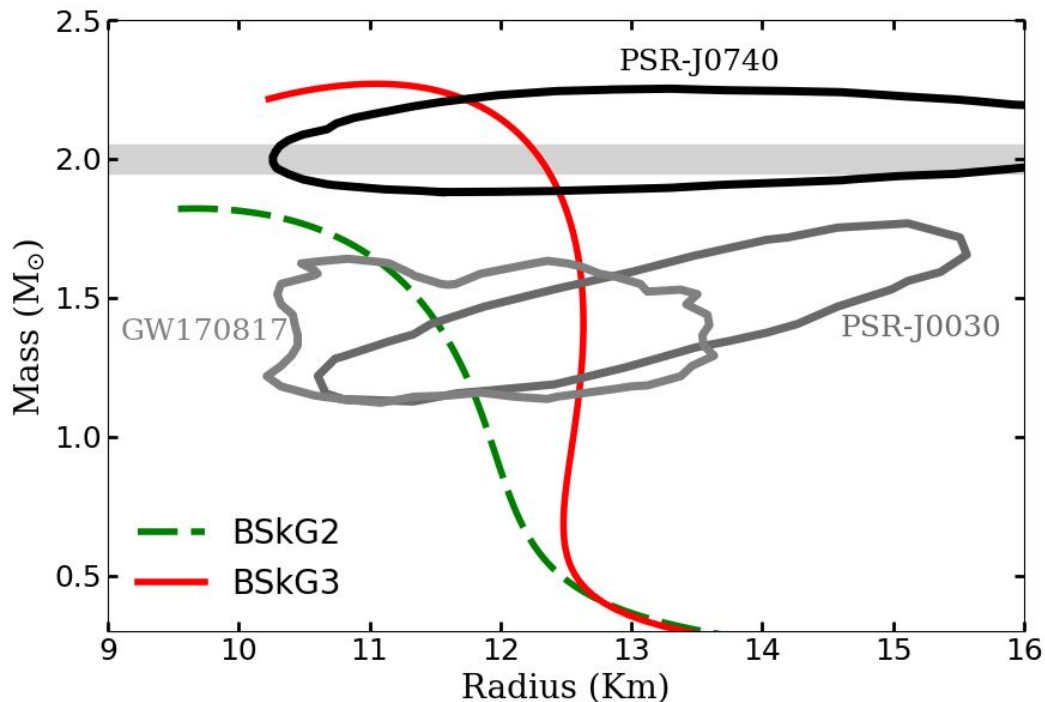
Neutron star properties



BSkG2
does not reproduce
massive pulsars

The BSkG3 interaction

Neutron star properties



$$R_{1.4} = 12.6 \text{ km}$$

$$M_{\text{max}} = 2.3 \text{ Msun}$$

$$R_{M_{\text{max}}} = 11.1 \text{ km}$$

$$M_{\text{dUrca}} = 1.5 \text{ Msun}$$

$$n_{\text{dUrca}} = 0.43 \text{ fm}^{-3}$$

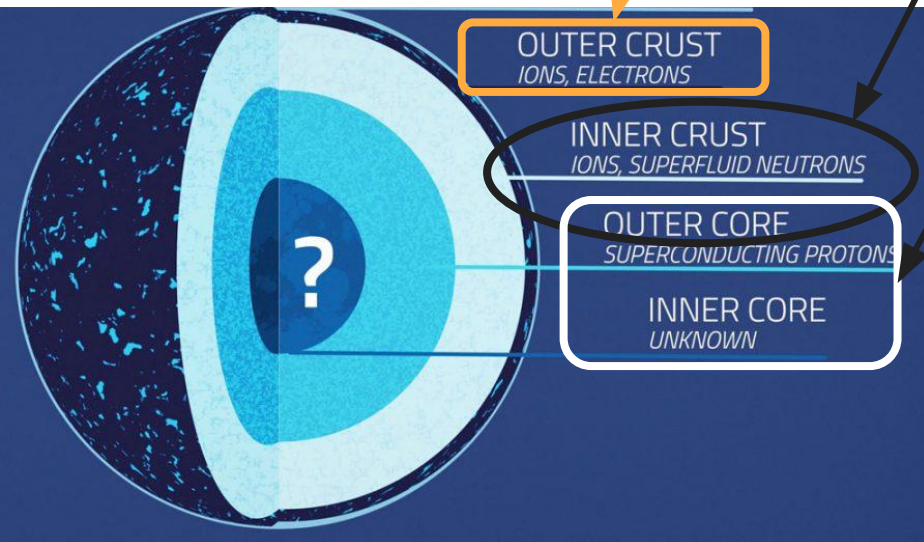
Summary

Root-mean-square on **experimental masses**

and fission barriers:

$$\sigma(M) < 800 \text{ KeV}$$

$$\sigma(E_f) < 500 \text{ KeV}$$



OUTER CRUST
IONS, ELECTRONS

INNER CRUST
IONS, SUPERFLUID NEUTRONS

OUTER CORE
SUPERCONDUCTING PROTONS

INNER CORE
UNKNOWN

- new **pairing** prescription with self-energy corrections (important for neutron and proton **superfluids** in NS **crust and core**);

- the **stiff neutron matter EoS** at **high densities**, which avoids the collapse of NS and allows the description of **heavy pulsars**;

The road ahead:

- systematic exploration of symmetry energy;
- move to finite temperature EoS;
- Fine-tune time-odd terms.

Acknowledgements

I thank:

- My collaborators: Wouter Ryssens, Guillaume Scamps, Stephane Goriely and Nicolas Chamel.
- The computational resources provided by the *Consortium des Équipements de Calcul Intensif* (CECI).
- The funding agencies FNRS and FWO.
- The EVEREST and MANASLU EOS projects.

• Thank you for the attention!



extra slides

The BSkG3 functional

Skyrme EDF

Standard
Skyrme used
in BSkG1 and
BSkG2

Energy

Coupling constants

Local densities ρ and kinetic density τ of a HFB state

$$E^{\text{Skyrme}} = C_1(\rho)\rho^2 + C_2\rho\tau + C_3\nabla\rho\nabla\rho + \dots$$

The BSkG3 functional

Skyrme EDF

Standard
Skyrme used
in BSkG1 and
BSkG2

Energy

Coupling constants

Local densities ρ and kinetic density τ of a HFB state

$$E^{\text{Skyrme}} = C_1(\rho)\rho^2 + C_2\rho\tau + C_3\nabla\rho\nabla\rho + \dots$$

Extended
Skyrme used
in BSkG3

$$E^{\text{BSkG3}} = C_1(\rho)\rho^2 + C_2(\rho)\rho\tau + C_3(\rho)\nabla\rho\nabla\rho + \dots$$

$$C_2(\rho) = C_2(t_1, x_1, t_2, x_2) + C_2(t_4, x_4)\rho^\beta + C_2(t_5, x_5)\rho^\gamma$$

where $t_1, x_1, \dots, t_5, x_5$ are
Skyrme parameters.

Ingredients of the mass model

We represent an atomic nucleus with N neutrons and Z protons with a many-body state of the Bogoliubov type, whose total energy E_{tot} we define as:

$$E_{\text{tot}} = E_{\text{HFB}} + E_{\text{corr}}. \quad (1)$$

We call E_{HFB} the mean-field energy and E_{corr} is a set of corrections that account (approximately) for correlations that cannot be captured by single mean-field reference state constructed from separate neutron and proton orbitals. More precisely, the mean-field energy consists of five parts:

$$E_{\text{HFB}} = E_{\text{kin}} + E_{\text{Sk}} + E_{\text{pair}} + E_{\text{Coul}} + E_{\text{cm}}^{(1)}, \quad (2)$$

which are, respectively, the kinetic, Skyrme, pairing and the Coulomb energy and the one-body part of the centre-of-mass correction [24]. The correction energy E_{corr} is written in terms of three parts:

$$E_{\text{corr}} = E_{\text{rot}} + E_{\text{cm}}^{(2)} + E_{\text{W}}, \quad (3)$$

$$E_{\text{Sk}} = \int d^3\mathbf{r} \sum_{t=0,1} [\mathcal{E}_{t,e}(\mathbf{r}) + \mathcal{E}_{t,o}(\mathbf{r})],$$

$$\begin{aligned} \mathcal{E}_{t,e}(\mathbf{r}) = & C_t^{\rho\rho}(\rho_0) \rho_t^2(\mathbf{r}) + C_t^{\rho\tau}(\rho_0) \rho_t(\mathbf{r}) \tau_t(\mathbf{r}) \\ & + C_t^{\rho\nabla J} \rho_t(\mathbf{r}) \nabla \cdot \mathbf{J}_t(\mathbf{r}) \\ & + C_t^{\rho\Delta\rho} \rho_t(\mathbf{r}) \Delta\rho_t(\mathbf{r}) \\ & + C_t^{\nabla\rho\nabla\rho}(\rho_0) \nabla\rho_t(\mathbf{r}) \cdot \nabla\rho_t(\mathbf{r}) \\ & + C_t^{\rho\nabla\rho\nabla\rho}(\rho_0) \rho_t(\mathbf{r}) \nabla\rho_0(\mathbf{r}) \cdot \nabla\rho_t(\mathbf{r}) \end{aligned} \quad (6)$$

$$\begin{aligned} \mathcal{E}_{t,o}(\mathbf{r}) = & C_t^{ss}(\rho_0) \mathbf{s}_t(\mathbf{r}) \cdot \mathbf{s}_t(\mathbf{r}) + C_t^{jj}(\rho_0) \mathbf{j}_t(\mathbf{r}) \cdot \mathbf{j}_t(\mathbf{r}) \\ & + C_t^{j\nabla s} \mathbf{j}_t(\mathbf{r}) \cdot \nabla \times \mathbf{s}_t(\mathbf{r}). \end{aligned} \quad (7)$$

$$\begin{aligned} C_t^{\rho\tau}(\rho_0) = & + \frac{1}{2} C_{0t}^+(t_1, x_1) + \frac{1}{2} C_{0t}^-(t_2, x_2) \\ & + \frac{1}{2} C_{0t}^+(t_4, x_4) \rho_0^\beta + \frac{1}{2} C_{0t}^-(t_5, x_5) \rho_0^\gamma, \end{aligned} \quad (8)$$

$$C_t^{\nabla\rho\nabla\rho}(\rho_0) = + \frac{3}{8} C_{0t}^+(t_4, x_4) \rho_0^\beta - \frac{1}{8} C_{0t}^-(t_5, x_5) \rho_0^\gamma, \quad (9)$$

$$C_t^{\rho\nabla\rho\nabla\rho}(\rho_0) = - \frac{1}{2} C_{0t}^+(t_4, x_4) \rho_0^{\beta-1}. \quad (10)$$

The BSkG3 interaction

Pairing interaction:

$$E_{\text{pair}} = \frac{1}{4} \sum_{q=p,n} \int d^3\mathbf{r} g_q(\rho_n, \rho_p) \tilde{\rho}_q^*(\mathbf{r}) \tilde{\rho}_q(\mathbf{r}),$$

Different from BSkG1 and BSkG2

$$g_q(\rho_n, \rho_p) = V_q(\rho_n, \rho_p) [1 + \kappa_q (\nabla \rho_0)^2]$$

The BSkG3 interaction

Pairing interaction:

$$E_{\text{pair}} = \frac{1}{4} \sum_{q=p,n} \int d^3\mathbf{r} g_q(\rho_n, \rho_p) \tilde{\rho}_q^*(\mathbf{r}) \tilde{\rho}_q(\mathbf{r}),$$

Designed to match the pairing gaps in INM deduced from *ab-initio* calculations.

L. G. Cao et al., PRC, 73, 014313, 2006.
Brueckner calculations with realistic two-body and three-body forces and account for self-energy corrections.

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$$g_q(\rho_n, \rho_p) = V_q(\rho_n, \rho_p) [1 + \kappa_q (\nabla \rho_0)^2]$$

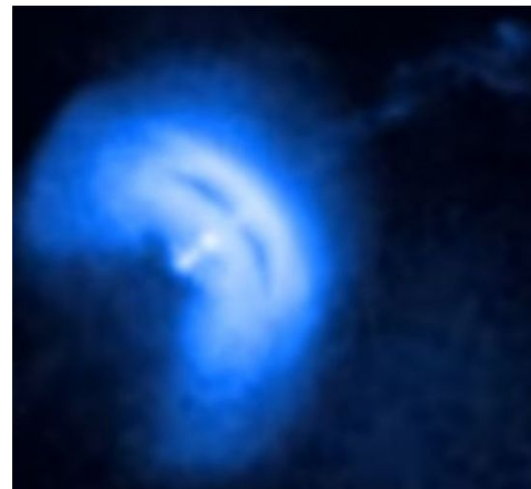
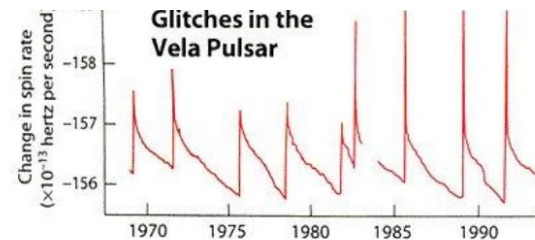
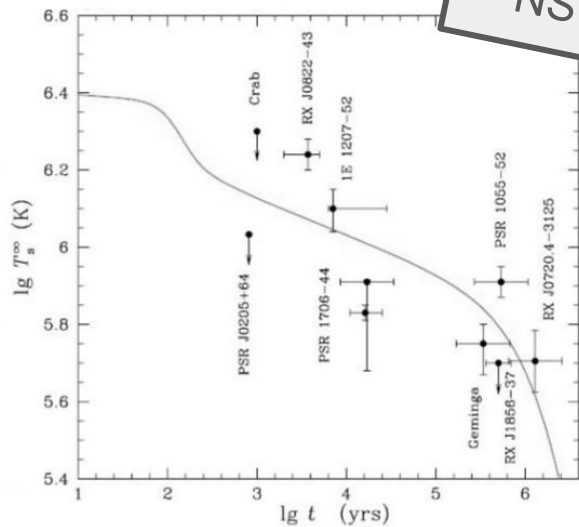
- Parameters $\kappa_{n/p}$ **adjusted** to reproduce **experimental** neutron and proton **pairing** gaps.
- $\kappa_{n/p}$ **vanish at INM** since it multiplies the density gradients.

The BSkG3 interaction

Pairing interaction:

Important to describe
superfluids in NSs.

NS cooling



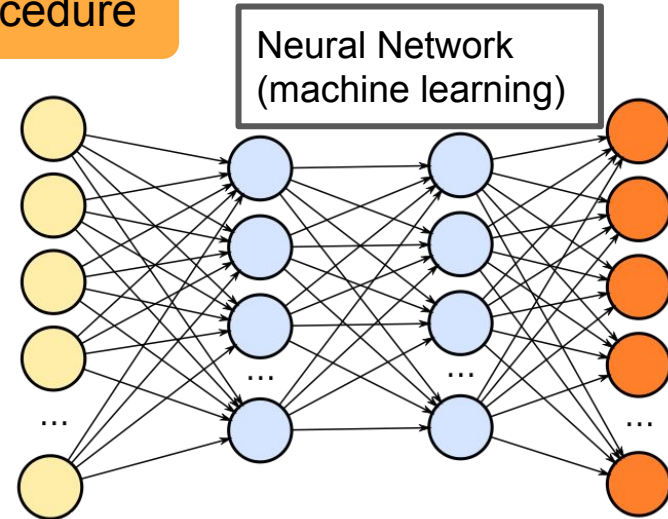
The BSkG3 functional

The fitting procedure

MOCCa

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- Fitted to experimental masses, charge radii, proton and neutron pairing gaps,
- constrain infinite nuclear matter properties.

Deformed nuclei are important to describe pasta configurations in NS

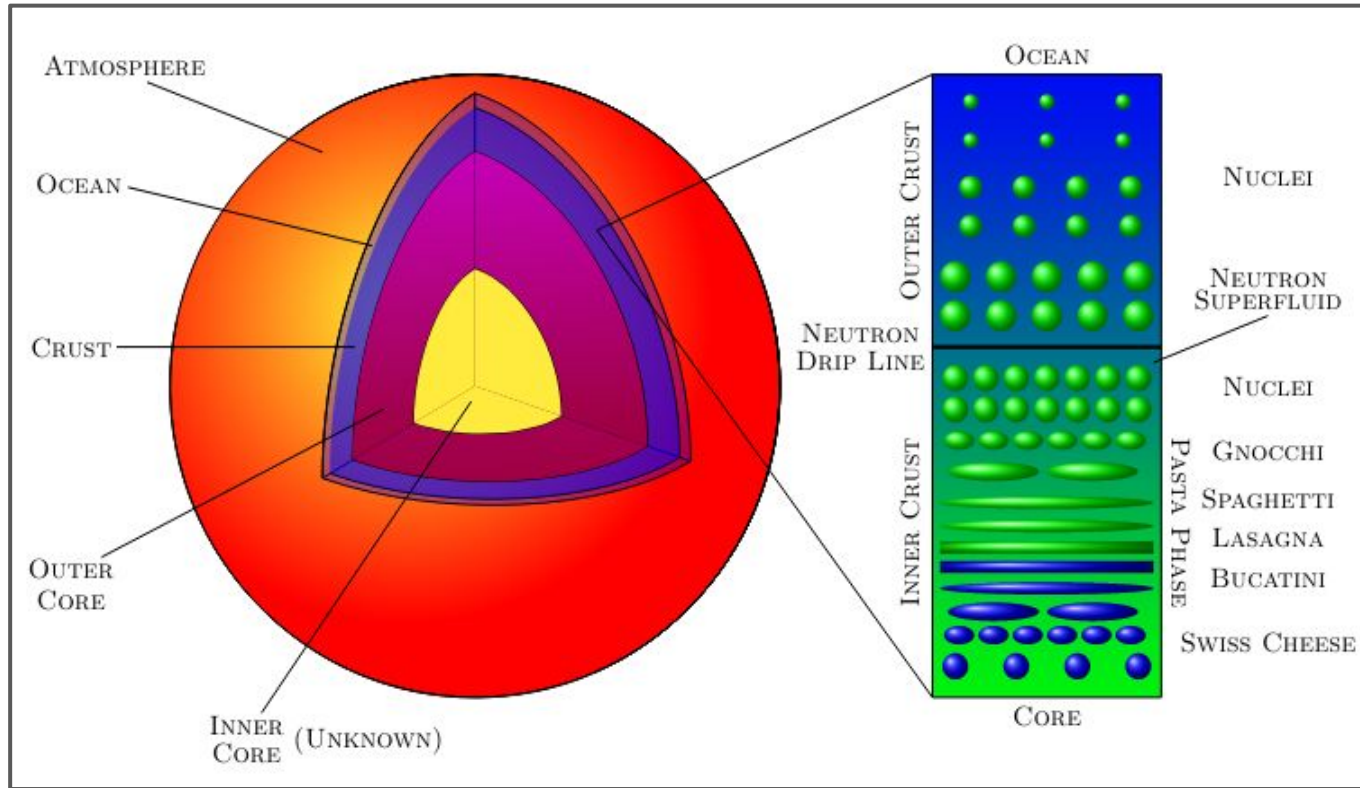
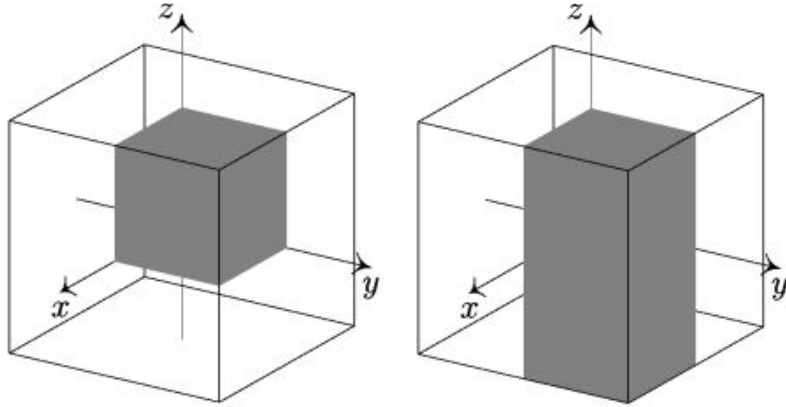


Figure from: Javier F. Acevedo et al JCAP 03, 038 (2020).

HFB solver

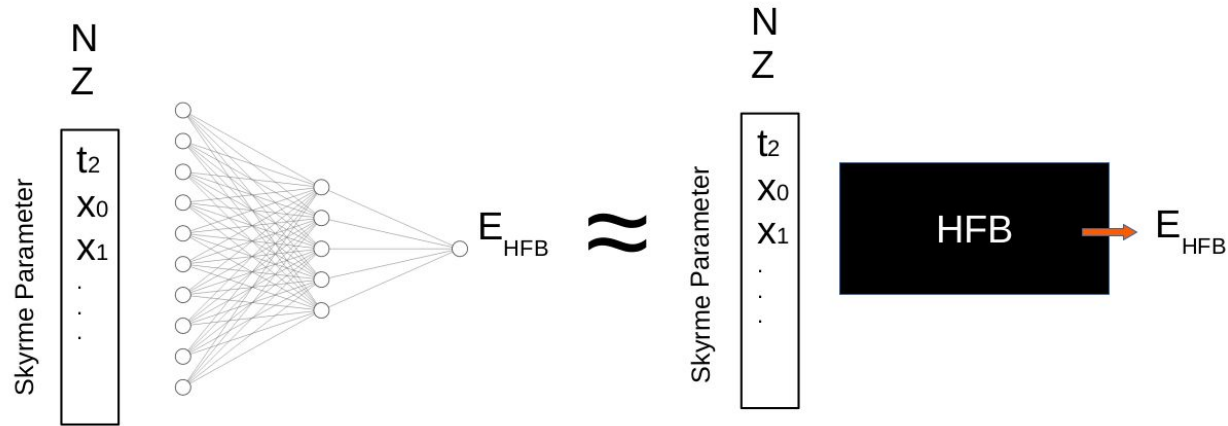
Modular Cranking Code = MOCCa

HFB solver at **3D coordinate-space** on **Lagrangian mesh** from W. Ryssens et al, W. R. PhD Thesis, ULB (2016).



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Machine learning method to fit



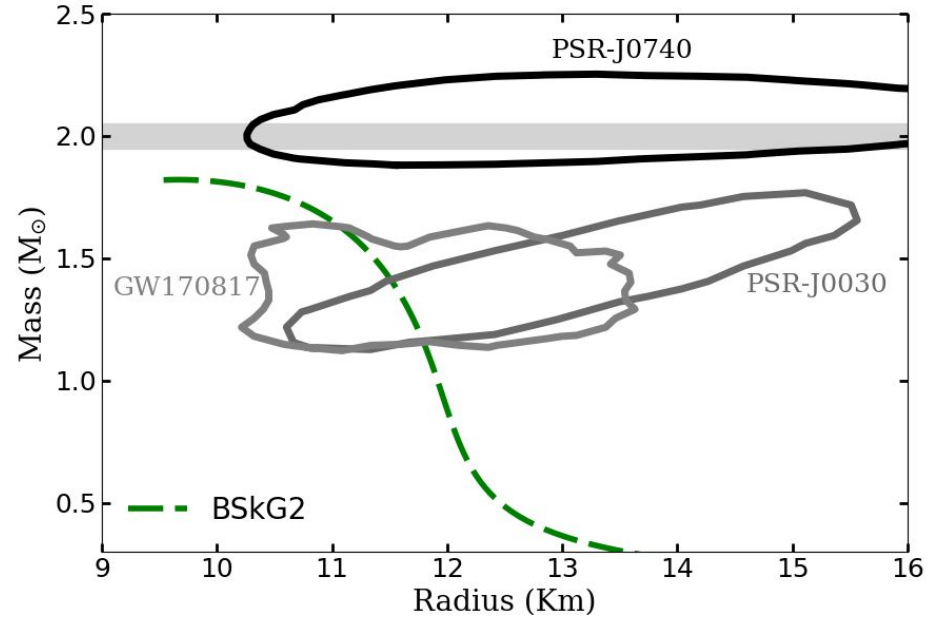
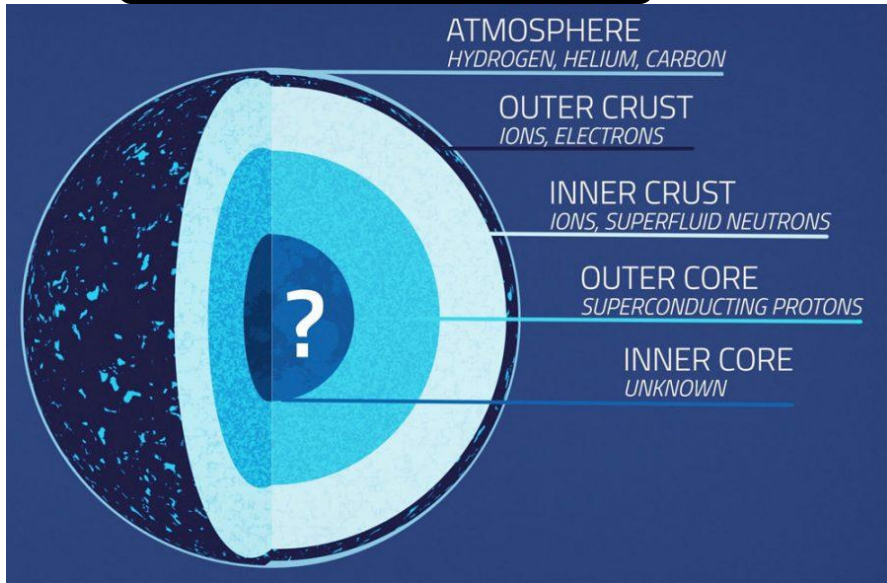
1. Create a **dataset of MOCCa results**.
2. Adjust a **machine learning system** to create an **emulator of the HFB solver**.
3. Adjust the mass model parameters using the emulator.

Machine learning prediction for one nucleus \sim a few seconds.
MOCCa predictions for one nucleus \sim 20 minutes.

4. **The mass table is computed with MOCCa.**

Motivation

Neutron star physics



Previous BSkG forces failed to reproduce massive neutrons stars.

The BSkG3 functional

The fitting procedure

1. Deduce INM equations, e.g., symmetry energy, energy per particle, compressibility, effective mass, etc.

$$\mathbf{J} = \mathbf{J}(t_2, \mathbf{x}_0, \mathbf{x}_1, \dots), \mathbf{K}_V = \mathbf{K}_V(t_2, \mathbf{x}_0, \mathbf{x}_1, \dots), m_{s/v}^* = m_{s/v}^*(t_2, \mathbf{x}_0, \mathbf{x}_1, \dots),$$

2. Impose constraints in the INM; obtain Skyrme parameters from the fit.

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2. Impose constraints in the INM; obtain Skyrme parameters from the fit.

1. Novelty at BSkG3:

impose constraint at **neutron matter energy (NeutM) at high density** (1 fm^{-3})

$$\mathbf{E}_{\text{NeutM}}(\text{at } 1 \text{ fm}^{-3}) = \mathbf{E}_{\text{NeutM}}(\mathbf{t}_4, \mathbf{t}_5, \dots, \mathbf{t}_2, \mathbf{x}_0, \mathbf{x}_1, \dots) \in [550, 600 \text{ MeV}].$$

2. Possible thanks to the **extended Skyrme** form which contains \mathbf{t}_4 and \mathbf{t}_5 terms (density dependent extension of \mathbf{t}_1 and \mathbf{t}_2).
3. Ensures a **stiff** neutron matter **equation of state** (essential for **neutron stars**).

The BSkG functionals

BSkG2

EPJA 58, 246 (2022)

Skyrme–Hartree–Fock–Bogoliubov mass models on a 3D mesh: II. Time-reversal symmetry breaking

Wouter Ryssens^{1,a}, Guillaume Scamps^{1,2}, Stephane Goriely¹, Michael Bender³

¹ Institut d'Astronomie et d'Astrophysique, Université Libre de Bruxelles, Campus de la Plaine CP 226, 1050 Brussels, Belgium

² Department of Physics, University of Washington, Seattle, WA 98195-1560, USA

³ Université de Lyon, Université Claude Bernard Lyon 1, CNRS / IN2P3, IP2I Lyon, UMR 5822, 69622 Villeurbanne, France

- **Allows for time-reversal symmetry breaking.**

Inclusion of '*time-odd*' terms in the Skyrme EDF.

No more Equal Filling Approximation (EFA).

- Incorporated information on the **fission properties** of twelve actinide nuclei in the fitting protocol.

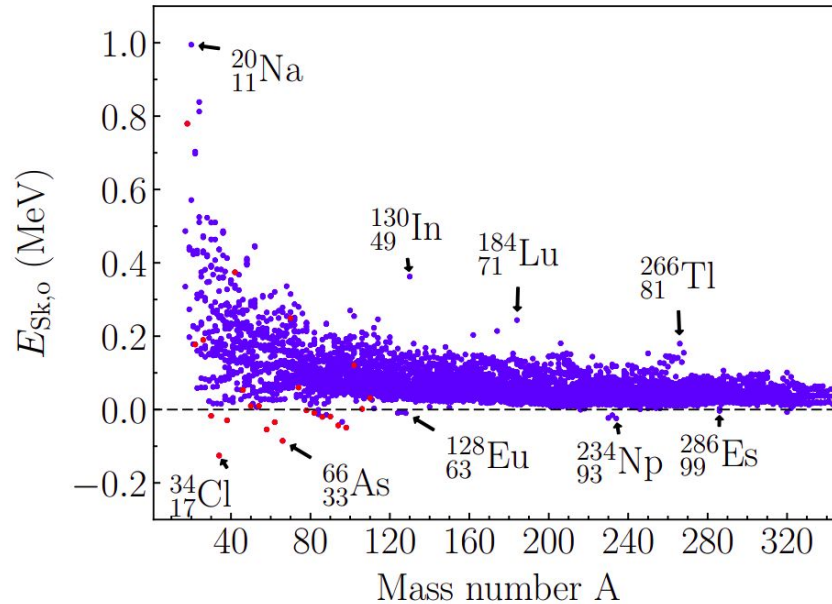
Access to more observables:

- magnetic moments;
- rotational bands.



The BSkG3 mass model

- **Allows for time-reversal symmetry breaking.**
Inclusion of '*time-odd*' terms in the Skyrme EDF - instead of Equal Filling Approximation (EFA) of previous model.
- Incorporated information on the **fission properties** of twelve actinide nuclei in the fitting protocol.



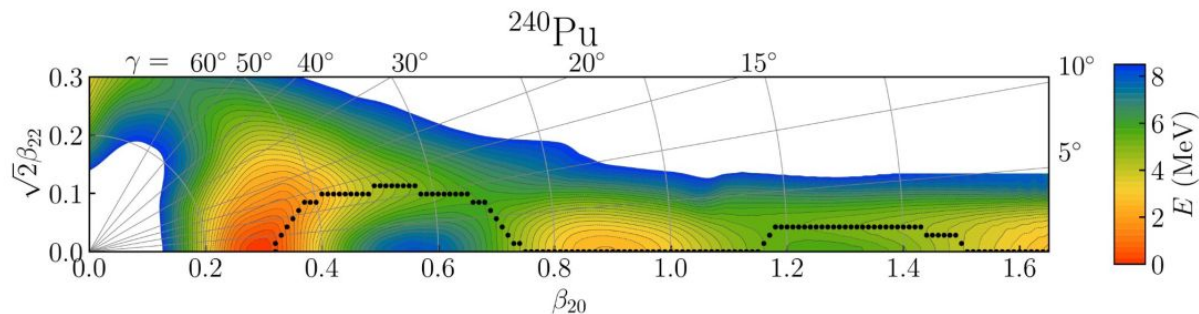
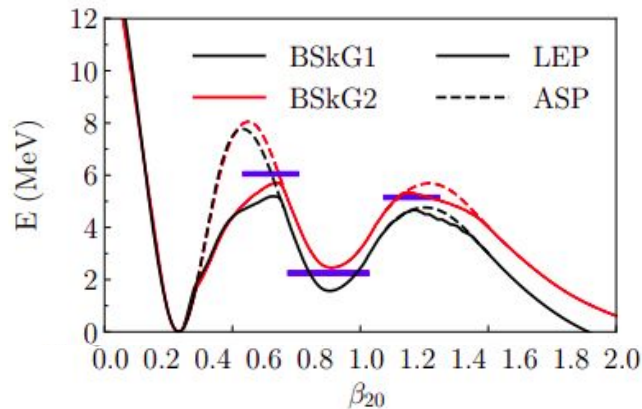
The BSkG functionals

BSkG2: fission properties

Skyrme-Hartree-Fock-Bogoliubov mass models on a 3D mesh. IIb. Fission properties of BSkG2.

arXiv: 2302.03097

Wouter Ryssens ^{a,1}, Guillaume Scamps ^{1,2}, Stephane Goriely ¹, Michael Bender ³



ASP = Axial Symmetric Path
LEP = Low Energy Path

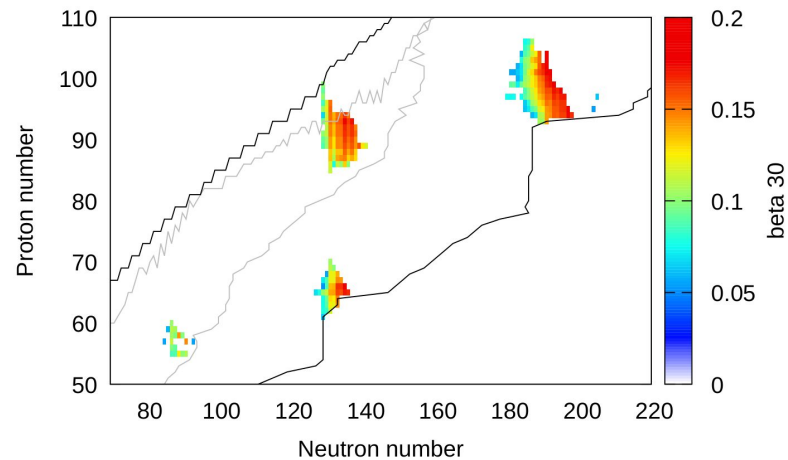
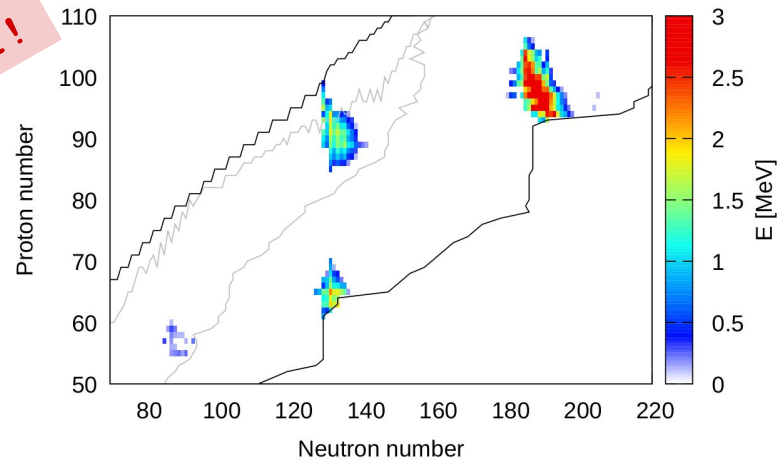
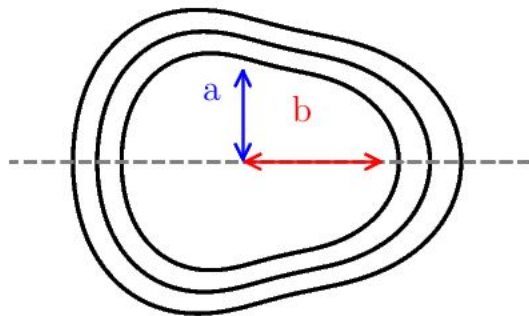


The BSkG3 mass model

Octupole shapes
in ground state

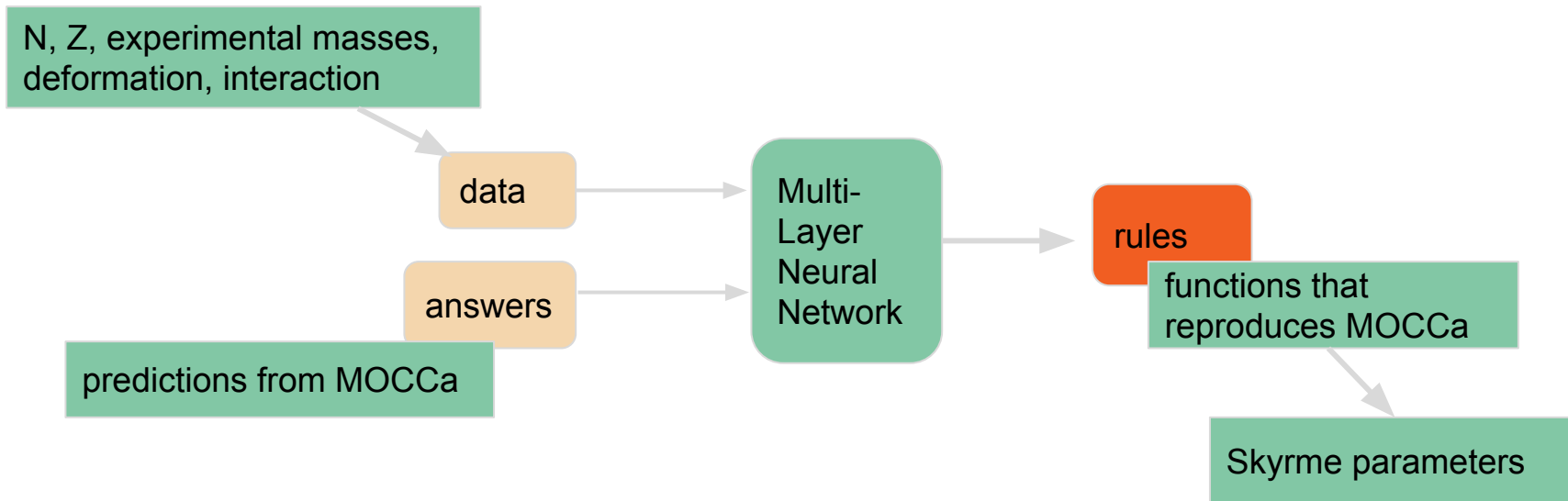
PRELIMINARY!

Reflection-asymmetric (RA)



workflow

The fitting procedure



How to describe an atomic nucleus?

macroscopic
description,
e.g., LDM: liquid
drop model

ab-initio: from the
bare nucleon-nucleon
interaction.

How to describe an atomic nucleus?

macroscopic description, e.g., LDM: liquid drop model.

too simple. (usually with no shell structure, neutron skin, pairing correlations...)

ab-initio: from the bare nucleon-nucleon interaction.

not (yet) feasible to describe thousands of nuclei along the nuclear chart.

A good compromise:

Energy density functional (EDF): an effective description of nuclei based on one-body densities.

Allows for predictions across the entire nuclear chart, firmly founded on a **microscopic description** of the nucleus.

The road to here

Brussels
mass models

AN ENERGY DENSITY NUCLEAR MASS FORMULA

(I). Self-consistent calculation for spherical nuclei

F. TONDEUR

Physique Nucléaire Théorique, Université Libre de Bruxelles, Campus de la Plaine, Cp 229, 1050 Bruxelles

Received 2 December 1977

HFB-1 mass model
BSk-1 nuclear force

A Hartree–Fock–Bogoliubov mass formula

M. Samyn^a, S. Goriely^{a,*}, P.-H. Heenen^b, J.M. Pearson^c, F. Tondeur^d

^a Institut d'Astronomie et d'Astrophysique, ULB-CP226, 1050 Brussels, Belgium

^b Service de Physique Nucléaire Théorique, ULB-CP229, 1050 Brussels, Belgium

^c Dépt. de Physique, Université de Montréal, Montréal, PQ, Canada H3C 3J7

^d Institut Supérieur Industriel de Bruxelles, 1000 Brussels, Belgium

Received 28 June 2001; revised 11 September 2001; accepted 25 September 2001

- EDF to calculate self-consistently the nuclei binding energies.
- Skyrme nucleon-nucleon interaction + Hartree-Fock many-body method.
- **Fitted to a large number of experimental masses.**

- **Pairing** channel first computed with BCS approximation (Tondeur 1977).
From HFB-1 models the **Bogoliubov** method is applied.
- **Constrained** on infinite **nuclear matter** properties (saturation density, symmetry energy, etc).
- Axial symmetry + harmonic oscillator basis.
- Equal Filling Approximation (EFA).

The road to here

Brussels
mass models

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Received 28 June 2001; revised 11 September 2001; accepted 25 September 2001

Skyrme-Hartree-Fock-Bogoliubov mass models on a 3D mesh: effect of triaxial shape

Guillaume Scamps^{1,a}, Stephane Goriely¹, Erik Olsen¹, Michael Bender^{2,b}, Wouter Ryssens^{1,3}

¹ Institut d'Astronomie et d'Astrophysique, Université Libre de Bruxelles, Campus de la Plaine CP 226, 1050 Brussels, Belgium

² Université de Lyon, Université Claude Bernard Lyon 1, CNRS, IP2I Lyon / IN2P3, UMR 5822, 69622 Villeurbanne, France

³ Center for Theoretical Physics, Sloane Physics Laboratory, Yale University, New Haven, CT 06520, USA

Received: 7 September 2021 / Accepted: 26 November 2021

- BSkG1: coordinate space representation on a **3D mesh**, allows for **triaxial** nuclei. *Machine learning tools as emulator* for the parameter adjustment.
- BSkG2: **break time-reversal symmetry**: includes time-odd terms on EDF (no more EFA) + information on **fission barriers** included on the fit procedure.

Coordinate space
representation
opens the door for
applications such
fission process
and **nuclear pasta**
on *neutron stars*
and *supernovae*.

The BSkG3 mass model

Skyrme-HFB
on the Grid

- **Nuclear interaction** described by an extended **Skyrme force**;
- **Many body problem** solved with **Hartree-Fock-Bogoliubov** approximation.
- **3D coordinate space** representation. **Triaxial** deformation.

Difference from previous BSkGs
to BSkG3:

The BSkG3 mass model

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- **Many body problem** solved with **Hartree-Fock-Bogoliubov** approximation.
- **3D coordinate space** representation. **Triaxial** deformation.

Difference from previous BSkGs
to BSkG3:

Improvement on *pairing* channel.
Functional fitted to ab-initio calculations of pairing
gaps with self-energy corrections.

Important for the neutron **superfluid** present on
neutron star inner-crust.

PHYSICAL REVIEW C **93**, 034337 (2016)

Further explorations of Skyrme-Hartree-Fock-Bogoliubov mass formulas. XVI. Inclusion of self-energy effects in pairing

S. Goriely,¹ N. Chamel,¹ and J. M. Pearson²

¹*Institut d'Astronomie et d'Astrophysique, CP-226, Université Libre de Bruxelles, 1050 Brussels, Belgium*

²*Département de Physique, Université de Montréal, Montréal (Québec), H3C 3J7 Canada*

(Received 20 January 2016; published 31 March 2016)

The BSkG3 mass model

Skyrme-HFB on the Grid

- Nuclear interaction described by an extended **Skyrme force**;
- Many body problem solved with **Hartree-Fock-Bogoliubov approximation**.
- **3D coordinate space** representation. **Triaxial** deformation.

PHYSICAL REVIEW C **80**, 065804 (2009)

Further explorations of Skyrme-Hartree-Fock-Bogoliubov mass formulas. XI **Stabilizing neutron stars** against a ferromagnetic collapse

N. Chamel,¹ S. Goriely,¹ and J. M. Pearson²

¹*Institut d'Astronomie et d'Astrophysique, CP226, Université Libre de Bruxelles, B-1050 Brussels, Belgium*

²*Département de Physique, Université de Montréal, Montréal, Québec H3C 3J7, Canada*

(Received 11 September 2009; published 17 December 2009)

We construct a new Hartree-Fock-Bogoliubov (HFB) mass model, labeled HFB-18, with a generalized Skyrme force. The additional terms that we have introduced into the force are density-dependent generalizations of the

PHYSICAL REVIEW C **82**, 035804 (2010)

Further explorations of Skyrme-Hartree-Fock-Bogoliubov mass formulas. XII. Stiffness and **stability of neutron-star matter**

S. Goriely,¹ N. Chamel,¹ and J. M. Pearson²

¹*Institut d'Astronomie et d'Astrophysique, CP-226, Université Libre de Bruxelles, B-1050 Brussels, Belgium*

²*Département de Physique, Université de Montréal, Montréal (Québec), H3C 3J7 Canada*

(Received 10 June 2010; published 24 September 2010)

We construct three new Hartree-Fock-Bogoliubov (HFB) mass models, labeled HFB-19, HFB-20, and HFB-21,

Difference from previous BSkGs
to BSkG3:

Two extra density depend
terms on EDF.

Already present on BSk-18 to
BSk-32, it helps to obtain a *stiff
neutron matter equation of
state*, and maintained a good
reproduction of experimental
masses.

The BSkG3 functional

Skyrme EDF

Standard Skyrme used in BSkG1 and BSkG2

Energy

Coupling constants

Local densities ρ and kinetic density τ of a HFB state

$$E^{Skyrme} = C_1(\rho)\rho^2 + C_2\rho\tau + C_3\nabla\rho\nabla\rho + \dots$$

Extended Skyrme used in BSkG3

$$E^{BSkG3} = C_1(\rho)\rho^2 + C_2(\rho)\rho\tau + C_3(\rho)\nabla\rho\nabla\rho + \dots$$

$$C_2(\rho) = C_2(t_1, x_1) + C_2^\beta(t_4, x_4)\rho^\beta$$

$$C_3(\rho) = C_3(t_2, x_2) + C_3^\gamma(t_5, x_5)\rho^\gamma$$