

Three-body contact of the resonant Fermi gas

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F. Werner and XL, arXiv:2211.09765
(submitted to Comptes rendus de l'Académie des sciences).

Outline

- 1 Introduction: three-body contact
- 2 Results for three-body contact in non-degenerate regime
- 3 Three-body contact and three-body losses
- 4 Conclusion-projects

1 Introduction: three-body contact

2 Results for three-body contact in non-degenerate regime

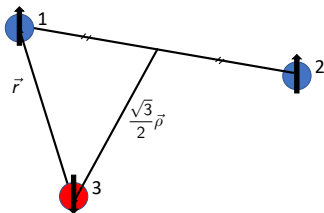
3 Three-body contact and three-body losses

4 Conclusion-projects

Introduction: three-body contact

System: Ultracold spin 1/2 interacting Fermi gases (BEC-BCS crossover)

- Jacobi coordinates:



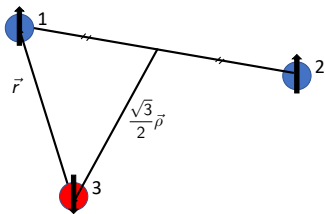
$\mathbf{R} = (\mathbf{r}, \rho)$ is hyperradius
(6D)

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Definition

$$N_3(\epsilon) \underset{\epsilon \rightarrow 0}{\sim} C_3 \epsilon^{2s+2}$$

C_3 is *three-body contact*.

$$s = 1.772724267\dots$$

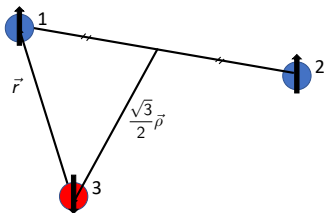
Efimov-like exponent for the 3-body problem for $l = 1$.

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Efimov-like exponent for the 3-body problem for $l = 1$.

Analogous to number of pairs $\uparrow\downarrow$ ($\bullet\bullet$)

$$\text{in } r < \epsilon : N_2(\epsilon) \underset{\epsilon \rightarrow 0}{\sim} C_2 \frac{\epsilon}{4\pi}$$

C_2 is Tan's (two-body) contact.

Introduction: three-body contact

Like Tan's contact C_2 , C_3 is due to *interactions*.

No interaction: $N_2^{(0)}(\epsilon) \propto \epsilon^3$, $N_3^{(0)}(\epsilon) \propto \epsilon^8$.

N_2 and N_3 are enhanced \Rightarrow Bunching effect.

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Zero range attractive interactions (scattering length a):
short range behavior of many-body wavefunction

2 or 3 particles close (r_{12} , $R \ll |a|$ or $(m|E|)^{-1/2}$):

V. Efimov 1970; D. Petrov, C. Salomon, G. Shlyapnikov 2004; S. Tan 2004.

$$\Psi(1, 2, \dots) \underset{r_{12} \rightarrow 0}{\simeq} \boxed{\frac{1}{r_{12}}} \times A(\dots) \Rightarrow N_2(\epsilon) \underset{\epsilon \rightarrow 0}{\sim} \epsilon$$

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$$\Psi(1, 2, 3 \dots) \underset{R \rightarrow 0}{\propto} \boxed{R^{s-2}} \times \left(\sum_{m=-1}^1 \Phi_m(\Omega) B_m(\dots) \right) \Rightarrow N_3(\epsilon) \underset{\epsilon \rightarrow 0}{\sim} \epsilon^{2s+2}$$

Introduction: three-body contact

Momentum distribution: $N_{\sigma}(\mathbf{k}) \underset{k \rightarrow +\infty}{\sim} \frac{C_2}{k^4}$

S. Tan 2005, 1D: M. Olchanyi, V. Dunjko 2003

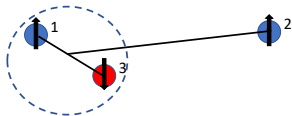
Number of pairs of opposite spin particles at a distance $< \epsilon$ of c.o.m. momentum \mathbf{K} :

Property

$$N_2(\epsilon, \mathbf{K}) \underset{\epsilon \rightarrow 0}{\sim} \frac{\epsilon}{4\pi} N_P(\mathbf{K})$$

$$N_P(\mathbf{K}) \underset{K \rightarrow +\infty}{\sim} \frac{C_3}{K^{2s+4}} \times (\text{constant}).$$

$$\Psi \sim \frac{1}{r} R^{s-1}$$



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Non degenerate regime

$n_\sigma \Lambda_T^3 \ll 1$, $\Lambda_T = \sqrt{\frac{2\pi\hbar^2}{mk_B T}}$ thermal wavelength: virial-like expansion

Virial expansion of three-body contact

$$\mathcal{C}_3 \underset{n_\sigma \Lambda_T^3 \rightarrow 0}{\sim} n^3 \left(\frac{\hbar^2}{mk_B T} \right)^{2-s} f\left(\frac{\lambda_T}{a}\right)$$

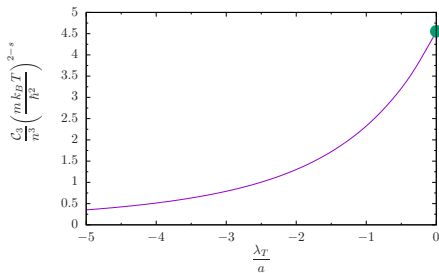
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Results :



Unitary limit:
exact result
(3-body problem)
 $f(0) = \frac{9\sqrt{3} \pi^3}{2^{2s+1} \Gamma(s+2)}$

Non degenerate regime, two approaches:

Wave-function:
(Unitary Limit)

UL : separability

$$\psi(\mathbf{R}) = \frac{F(R)}{R^2} \phi(\Omega)$$

exact results for F and ϕ .

$$F(R) \underset{R \rightarrow 0}{\sim} \text{const } R^s$$

Non degenerate regime, two approaches:

Green's functions:

3-body correlation function

Wave-function:
(Unitary Limit)

$$g_3(\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3) = \langle \psi_{\uparrow}^{\dagger}(\mathbf{r}_1) \psi_{\uparrow}^{\dagger}(\mathbf{r}_2) \psi_{\downarrow}^{\dagger}(\mathbf{r}_3) \psi_{\downarrow}(\mathbf{r}_3) \psi_{\uparrow}(\mathbf{r}_2) \psi_{\uparrow}(\mathbf{r}_1) \rangle$$

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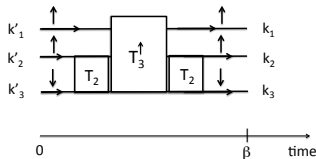
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(virial expansion with diagrams: XL 2011)

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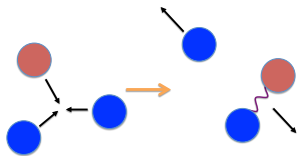
$$\Im(t_3^{\prime=1}(k, k')) \underset{k, k' \rightarrow +\infty}{\sim} \frac{B(E, a^{-1})}{(k k')^{s+1}}$$

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spin 1/2 (● ●) fermions
@ ultralow T (e. g. ${}^6\text{Li}$).

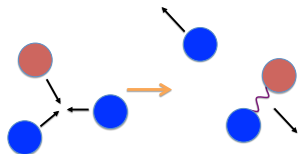


Three-body recombination \implies **losses**

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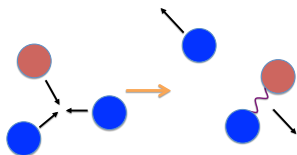
Three-body loss rate Γ_3 (s^{-1})

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Property

$$\Gamma_3 = -\frac{\hbar}{m} 4s(s+1) \text{Im}[a_3] C_3 \propto C_3$$

a_3 : "three-body parameter" for short-range physics

$$\Psi(\mathbf{R}, \dots) \underset{R \rightarrow 0}{\sim} \sum_{m=-1}^1 (R^s - a_3 R^{-s}) \frac{1}{R^2} \Phi_m(\boldsymbol{\Omega}) B(\dots)$$

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Conclusion-projects

- three-body contact: new observable, analogous to two-body Tan's contact, but for three-body correlations.
- analogous relations for short-range correlations $N_3(\epsilon)$, high-momentum tails of close pairs has $K^{-7.5454..}$ behavior.
- lowest order virial expansion of C_3 . This could be used has a calibration to measure C_3 , through three-body losses, at any T .
- j -particle contact ?

Derivative of the energy with respect to the three-body parameter a_3

Property

$$\left. \frac{\partial E}{\partial a_3} \right|_a = \frac{\hbar^2}{m} 2s(s+1) C_3.$$