

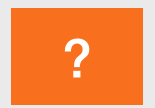
Physics at the High Energy Muon Collider

APR. 25 2023

ROBERTO FRANCESCHINI (ROMA 3 UNIVERSITY)



Open Questions on the “big picture” on fundamental physics circa 2020



• what is the dark matter in the Universe?



• why QCD does not violate CP?



• how have baryons originated in the early Universe?



• what originates flavor mixing and fermions masses?



• what gives mass to neutrinos?

EFT



• why gravity and weak interactions are so different?

EFT



• what fixes the cosmological constant?










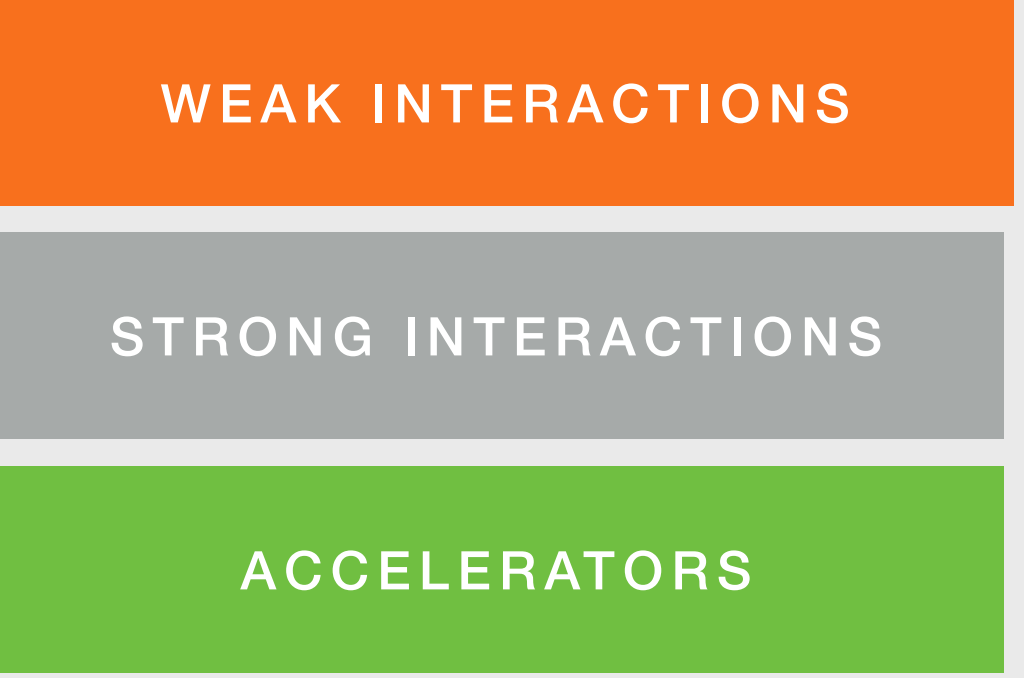
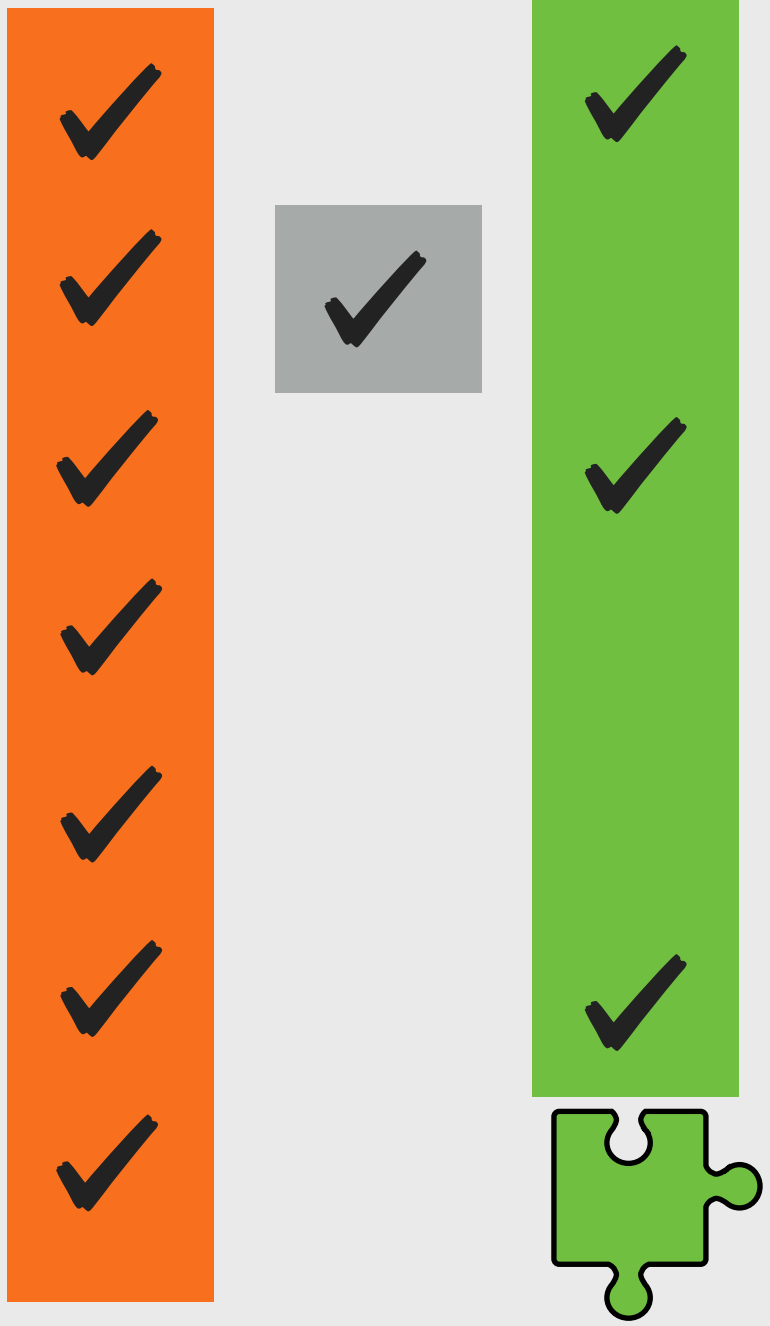
WEAK INTERACTIONS

STRONG INTERACTIONS

Accelerators are excellent probes

Open Questions on the “big picture” on fundamental physics circa 2020

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-  • what originates flavor mixing and fermions masses?
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- EFT*  • why gravity and weak interactions are so different?
- EFT*  • what fixes the cosmological constant?



Accelerators are excellent probes

Open Questions on the “big picture” on fundamental physics circa 2020

$\mu^+\mu^-$ sensitivity to weak interactions



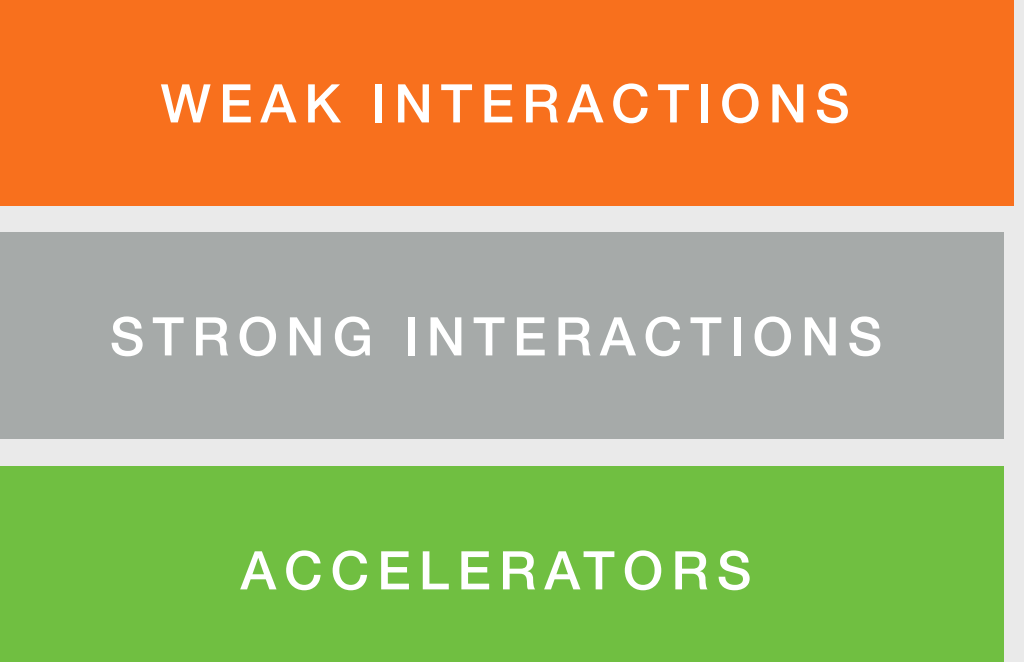
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EFT



EFT



Accelerators are excellent probes

Muon colliders

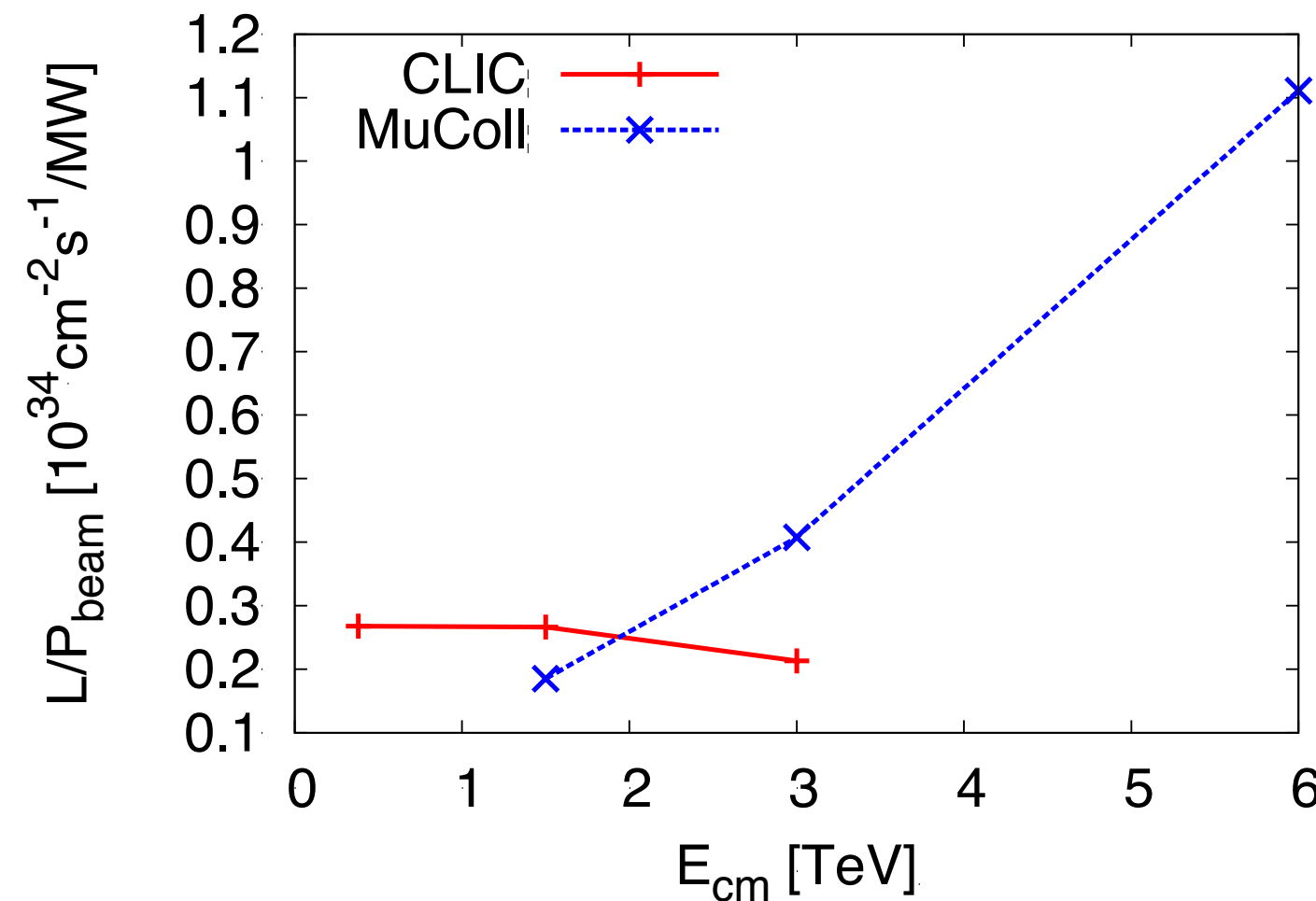
HIGHLY EFFICIENT

HIGH ENERGY COLLIDER

Luminosity Comparison

The luminosity per beam power is about constant in linear colliders

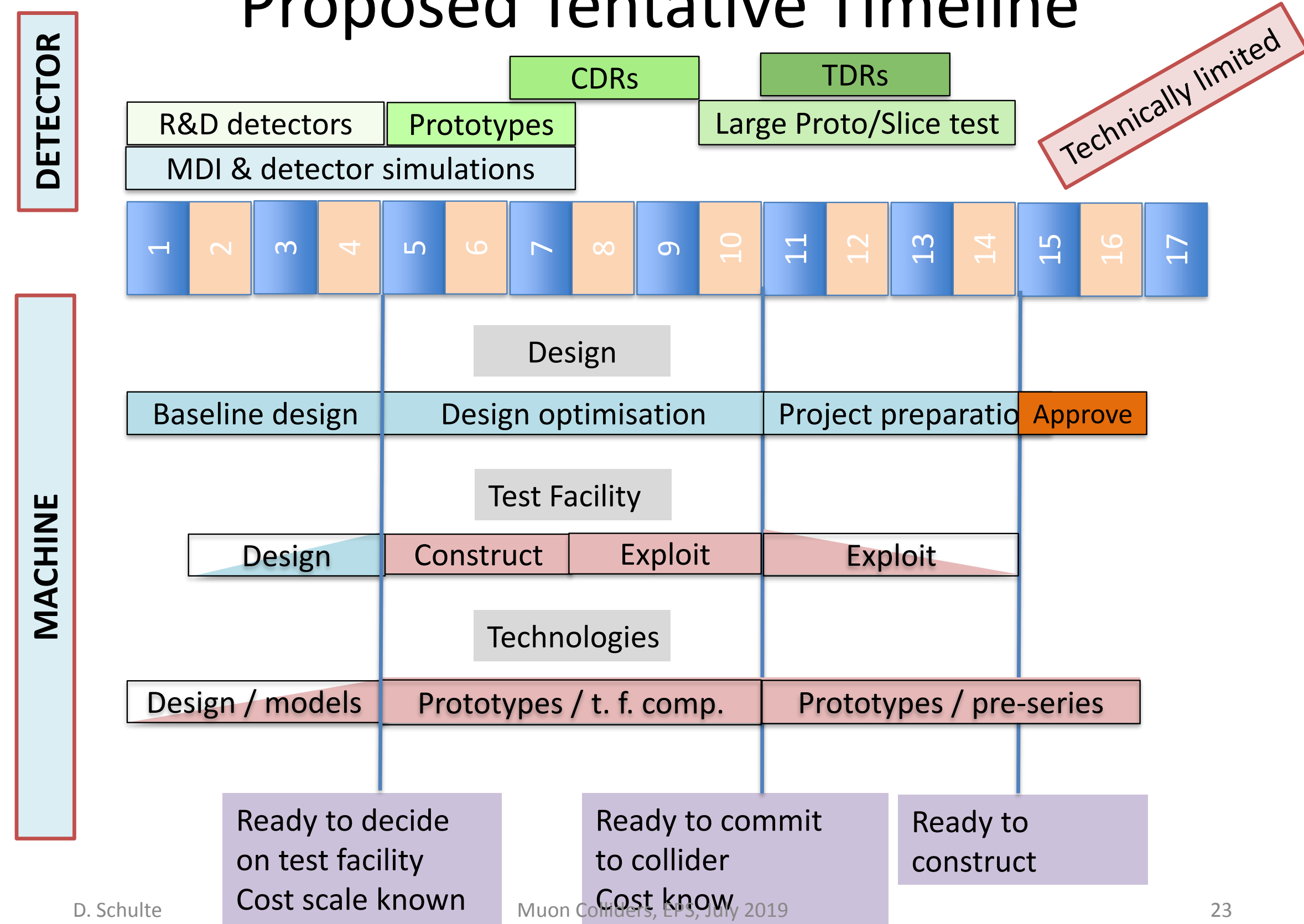
It can increase in proton-based muon colliders



Strategy CLIC:
 Keep all parameters at IP constant
 (charge, norm. emittances, betafunctions, bunch length)
 ⇒ Linear increase of luminosity with energy (beam size reduction)

Strategy muon collider:
 Keep all parameters at IP constant
 With exception of bunch length and betafunction
 ⇒ Quadratic increase of luminosity with energy (beam size reduction)

Proposed Tentative Timeline



Muon colliders

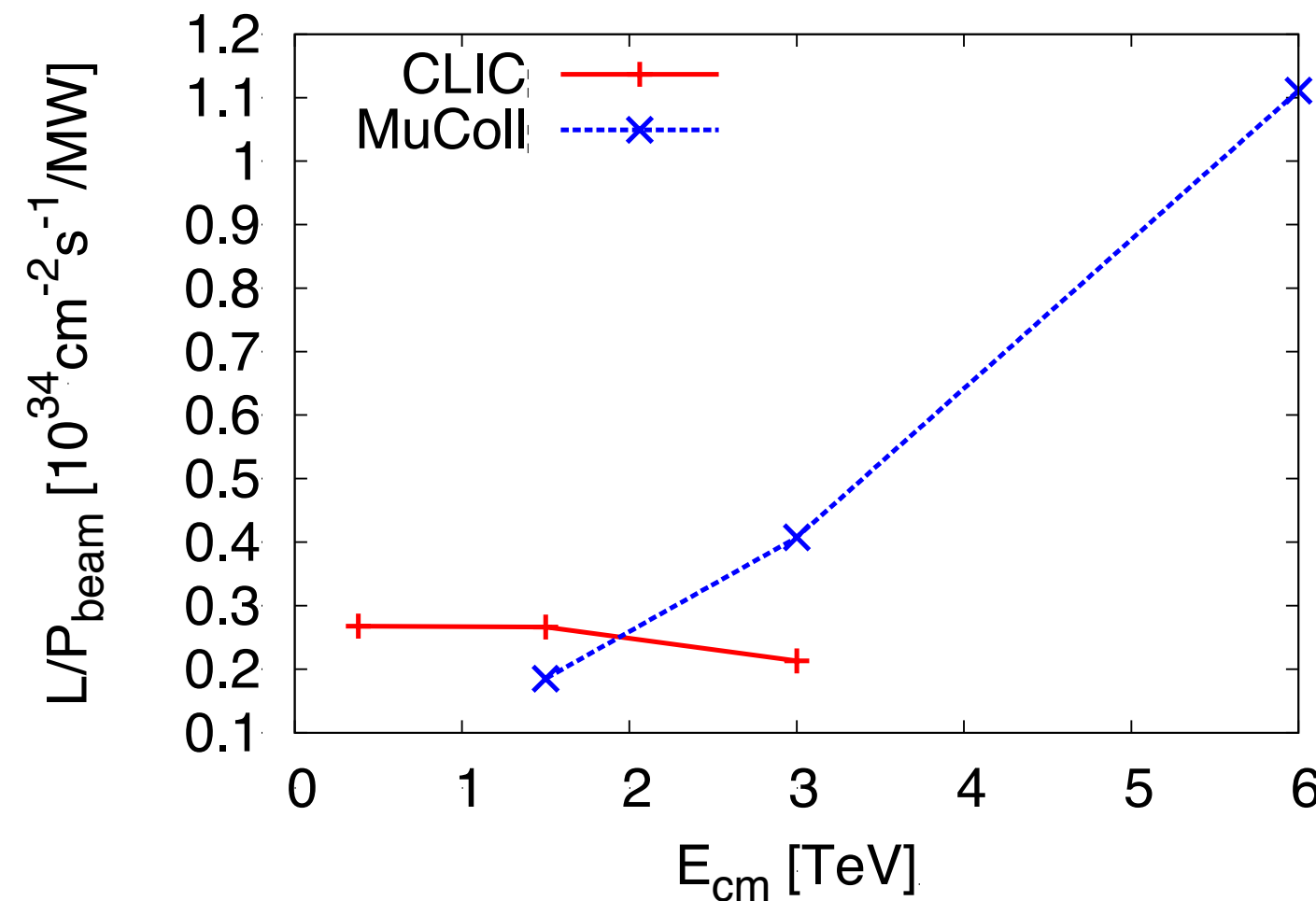
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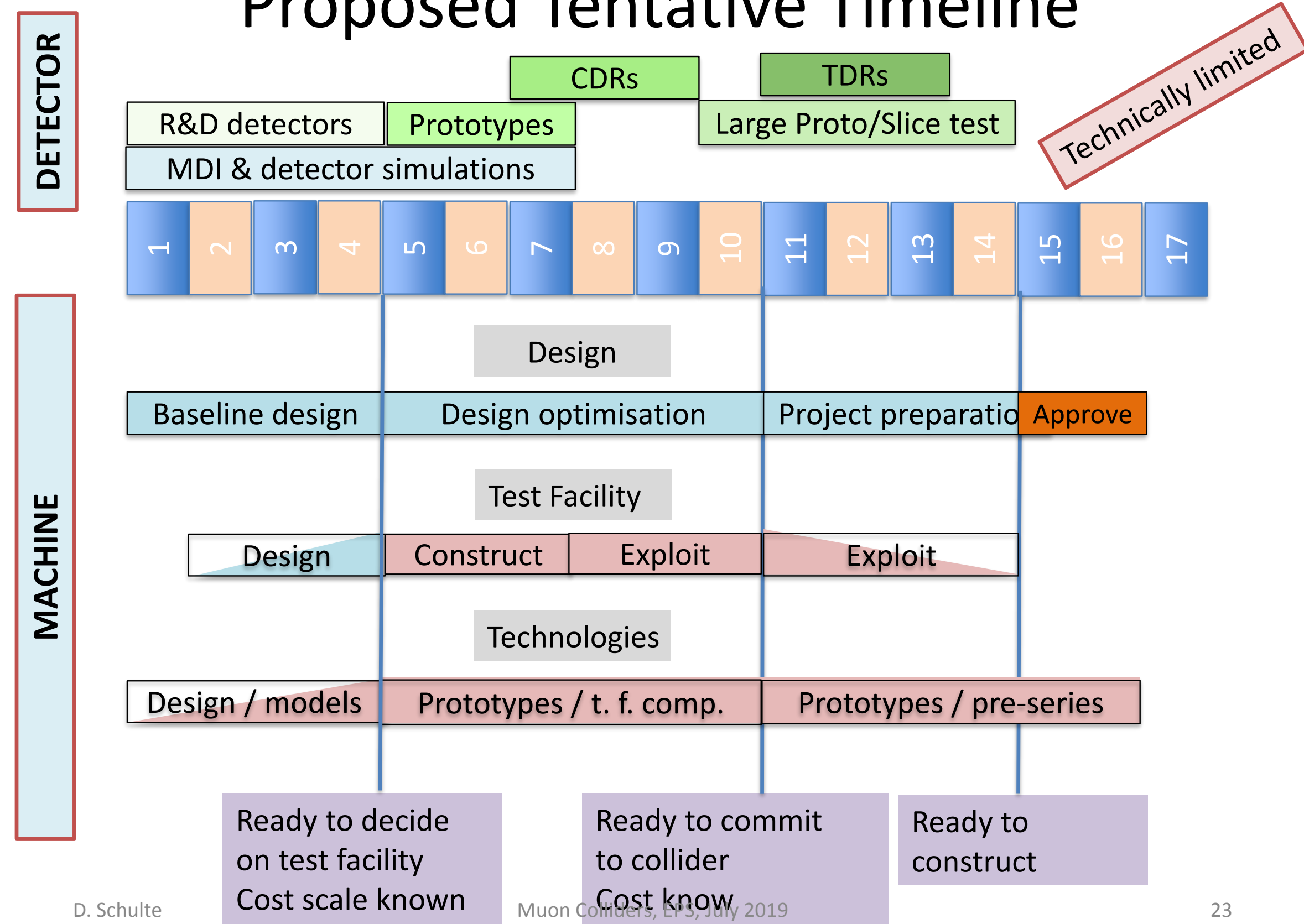
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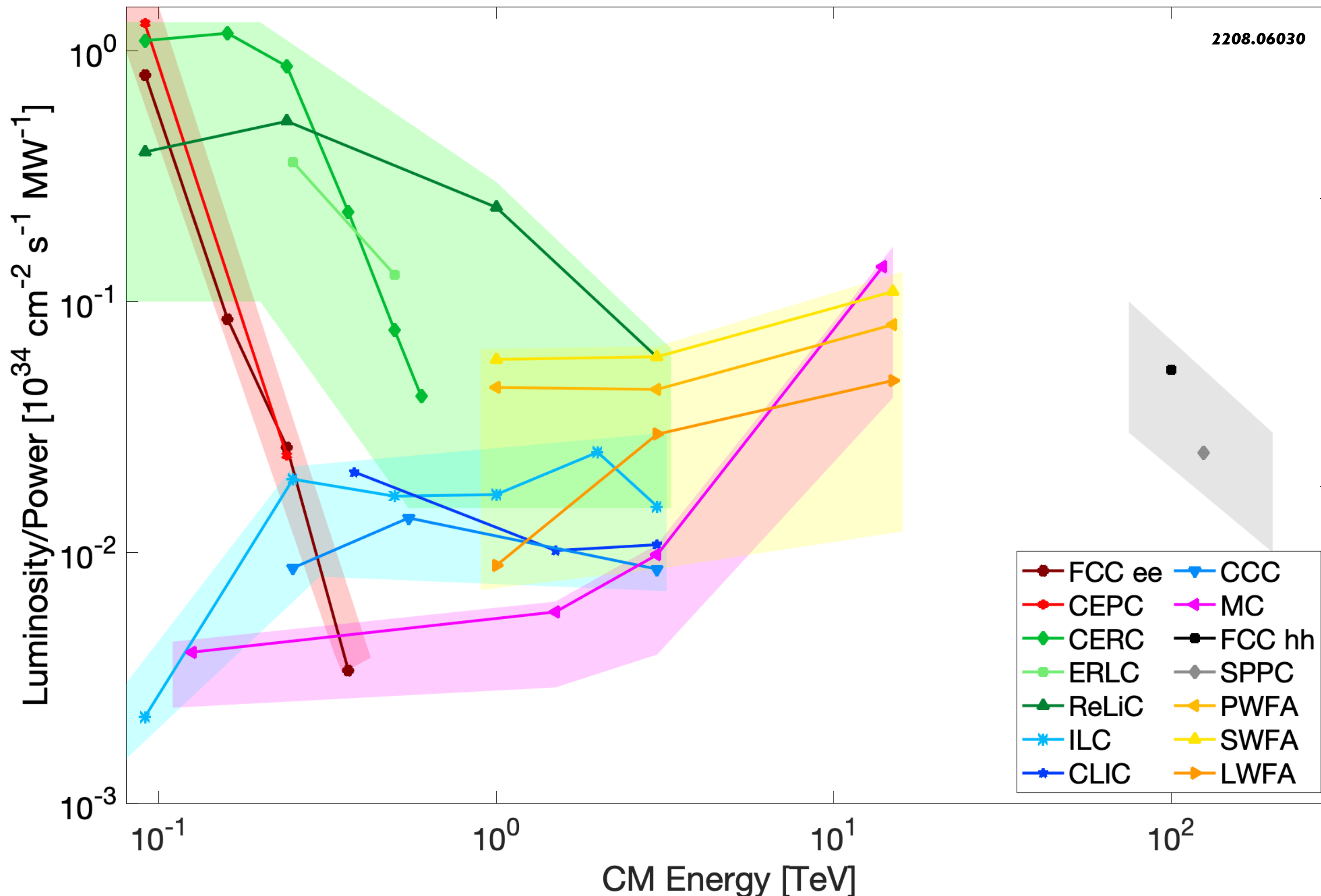
The luminosity per power is about comparable to linear colliders

It can increase in power based muon colliders

Strategy CLIC:
Keep all parameters constant (charge, norm. em) \Rightarrow Linear increase

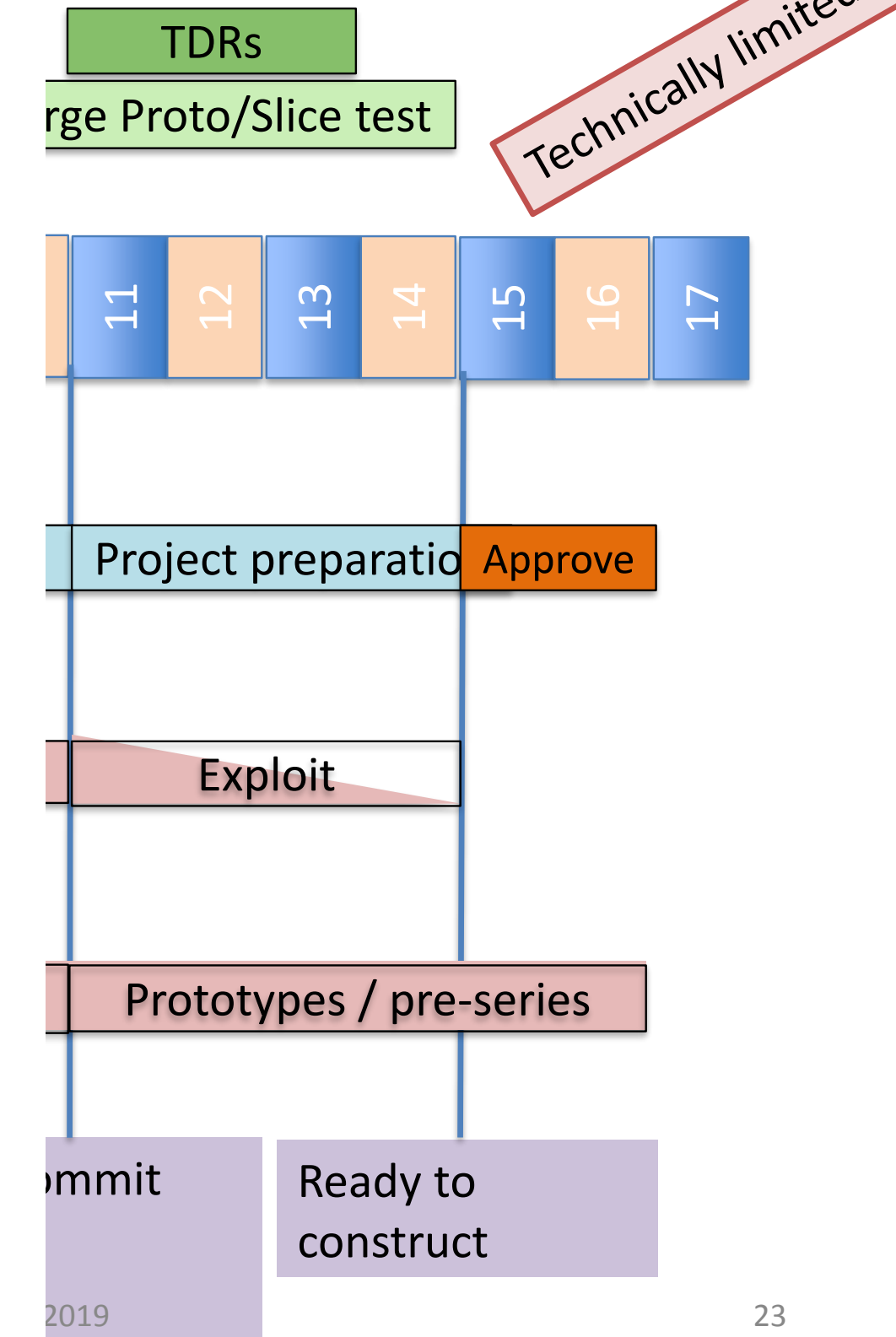
Strategy muon collider:
Keep all parameters constant With exception of \Rightarrow Quadratic increase

D. Schulte



Integrate Luminosity per Energy [$\text{ab}^{-1} \text{ TWh}^{-1}$]

Projective Timeline



Muon colliders

- International Muon Collider Collaboration formed to establish the physics case and the feasibility of a high energy muon collider



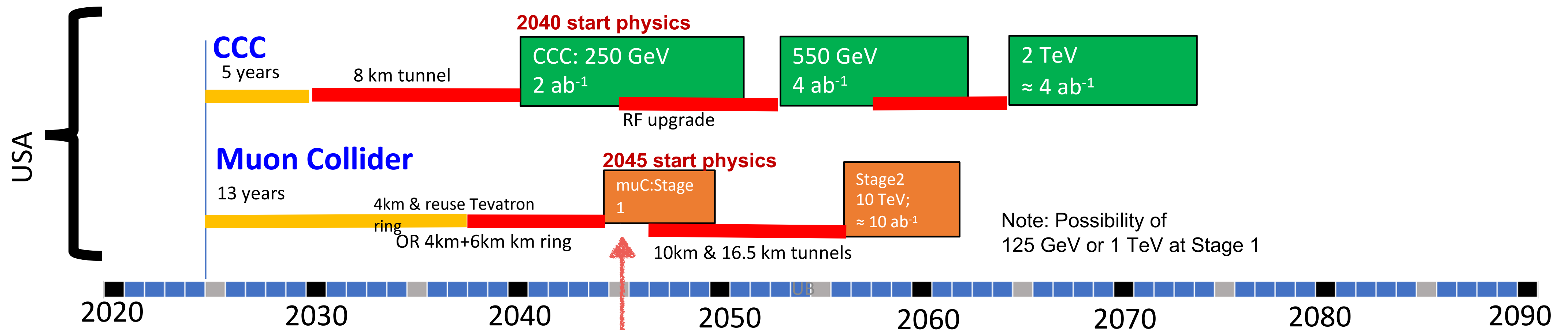
A Muon Collider in the US?

Possible scenarios of future colliders



Original from ESG by UB
Updated July 25, 2022 by MN

Proposals emerging from this Snowmass for a US based collider



- **Timelines technologically limited**
- Uncertainties to be sorted out
 - Find a contact lab(s)
 - Successful R&D and feasibility demonstration for CCC and Muon Collider
 - Evaluate CCC progress in the international context, and consider proposing an ILC/CCC [ie CCC used as an upgrade of ILC] or a CCC only option in the US.
 - International Cost Sharing

2045₀¹⁰ first $\mu^+\mu^-$ collisions pushing both the intensity and energy frontier

VERY SIGNIFICANT INTEREST FROM US COMMUNITY IN RECENT YEARS CULMINATED AT SNOWMASS 2022

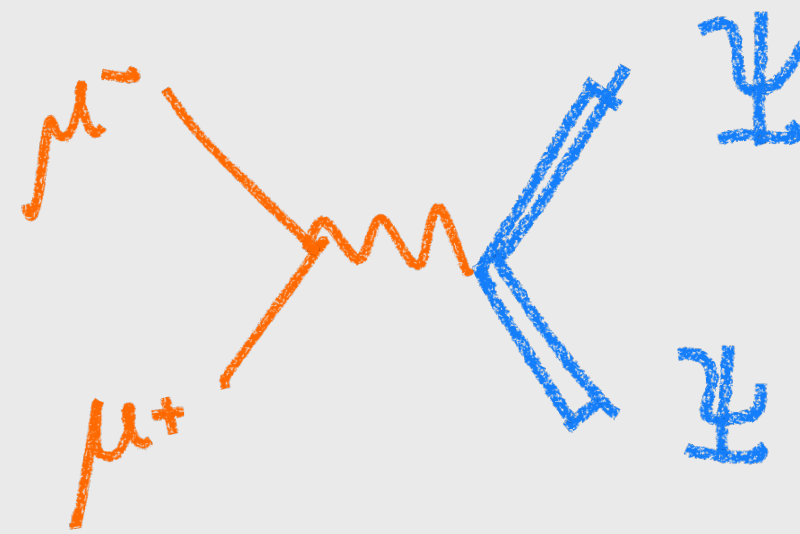
Muons collisions



$\mu^+ \mu^-$ collisions to probe fundamental physics

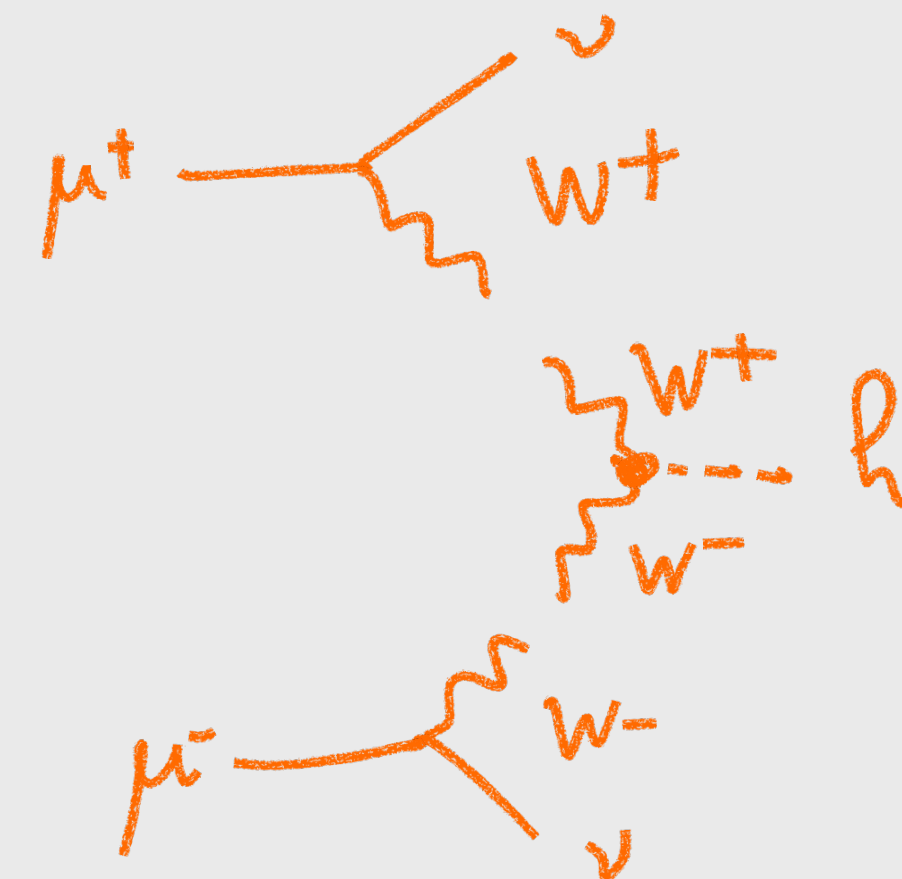
DIRECT SEARCHES

- production of SM and new physics in direct $\mu^+ \mu^-$ annihilation



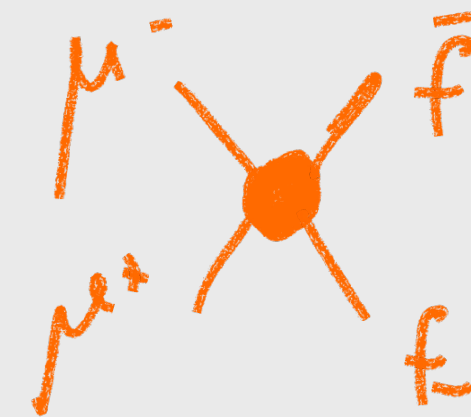
HIGH-INTENSITY PROBES

- production of SM and new physics using beam constituents (e.g. W bosons)



HIGH-ENERGY PROBES

- indirect probes of new physics in direct $\mu^+ \mu^-$ annihilation



$\sqrt{s} \gtrsim 3 \text{ TeV}$ center of mass brings significant extension compared to HL-LHC

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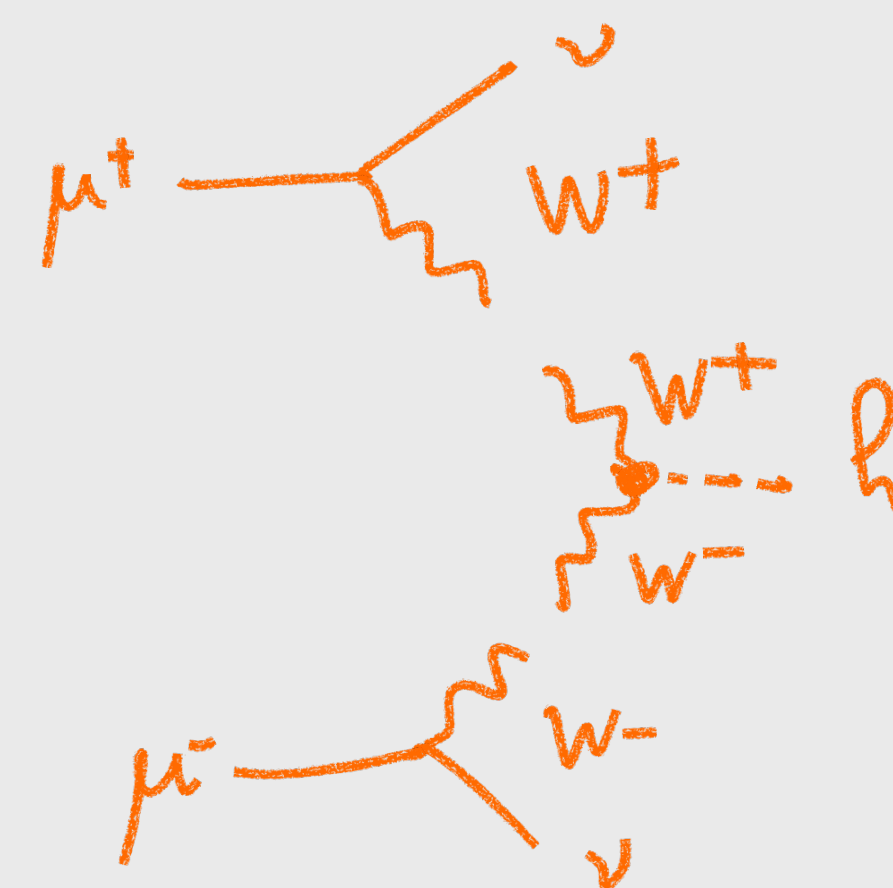
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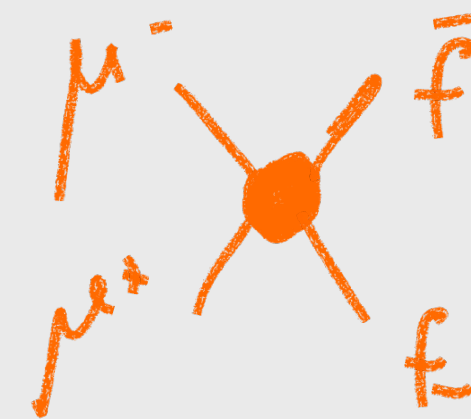
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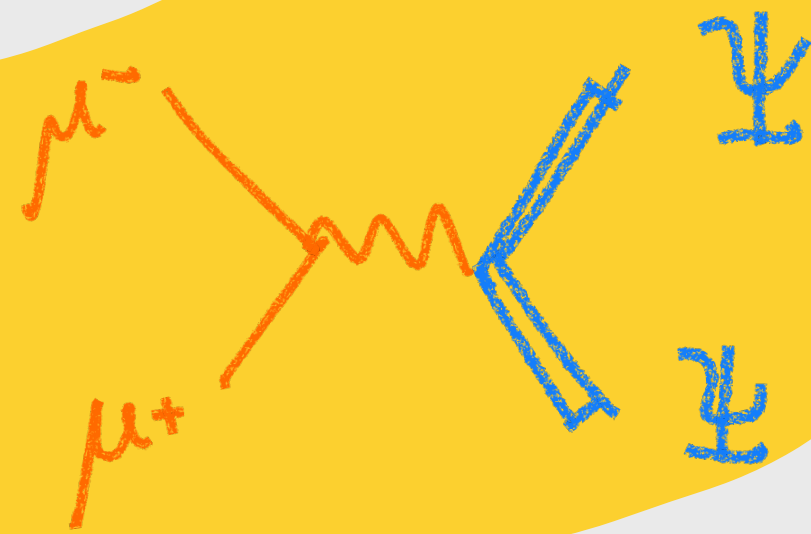


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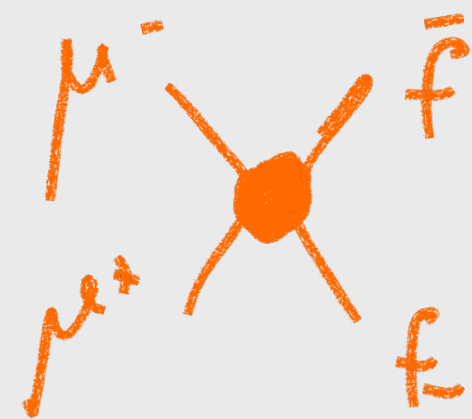
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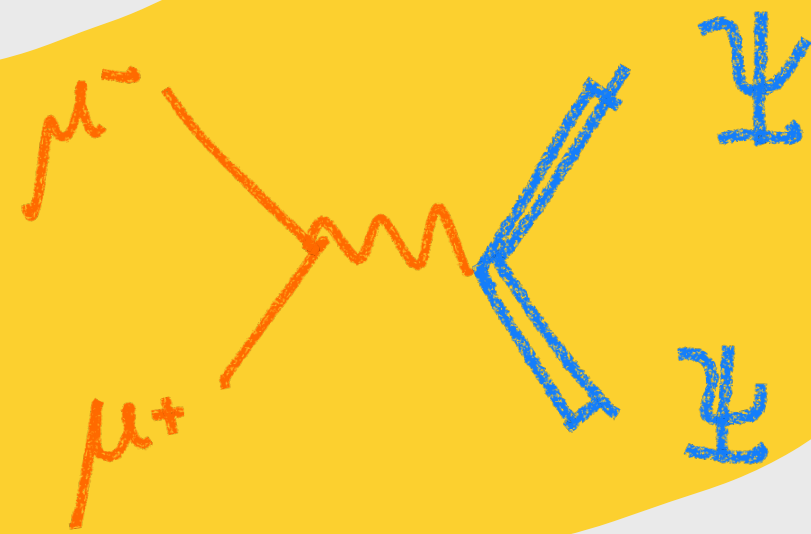


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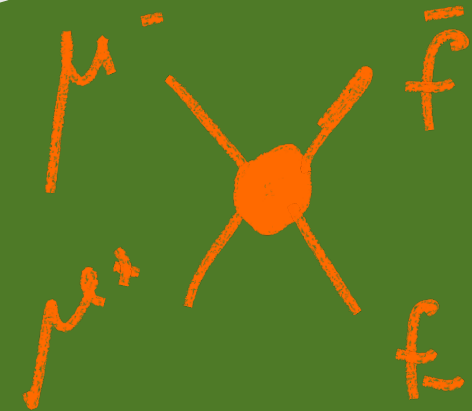
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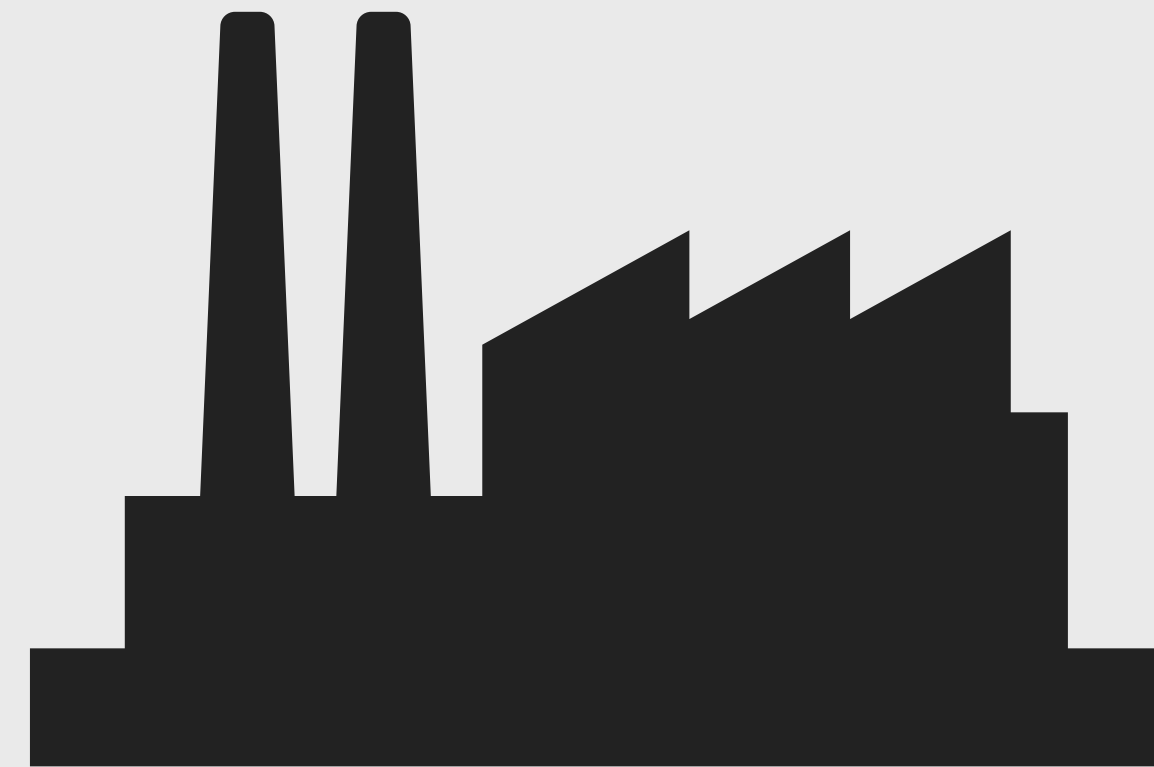


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Microscope & Factory



&



Microscope & Factory

Towards a Muon Collider

118 pages

e-Print: [2303.08533](https://arxiv.org/abs/2303.08533) [physics.acc-ph]

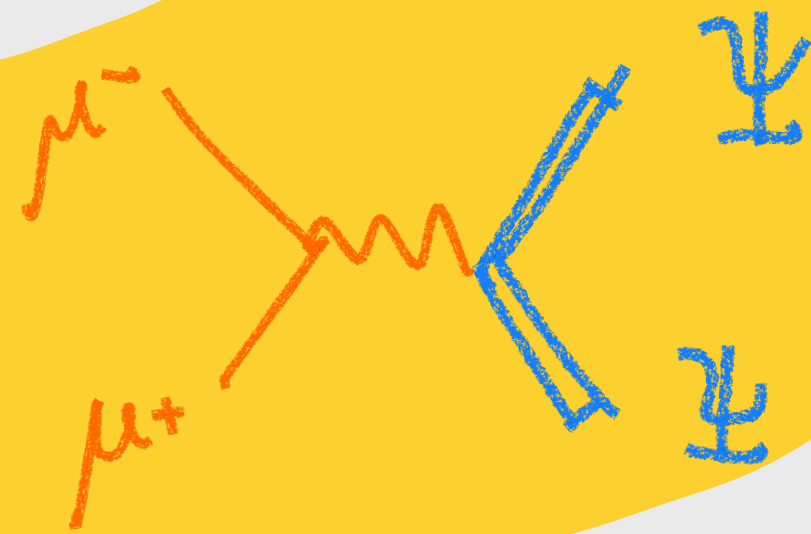
Report number: FERMILAB-PUB-23-123-AD-PPD-T

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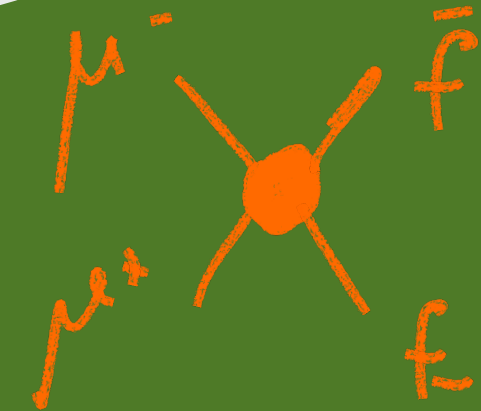
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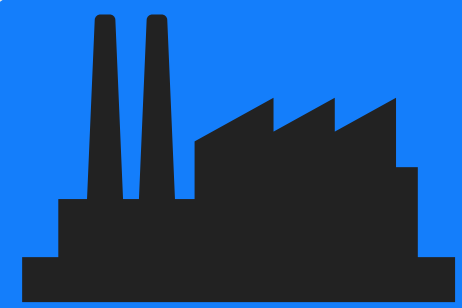
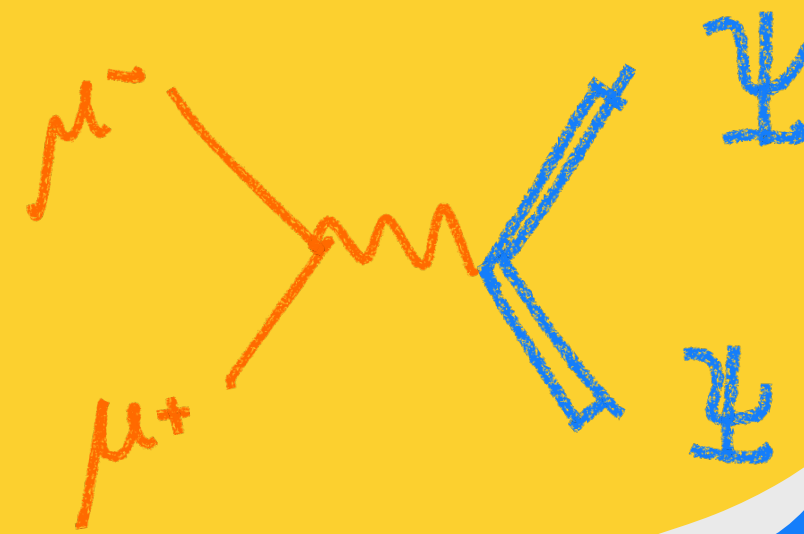
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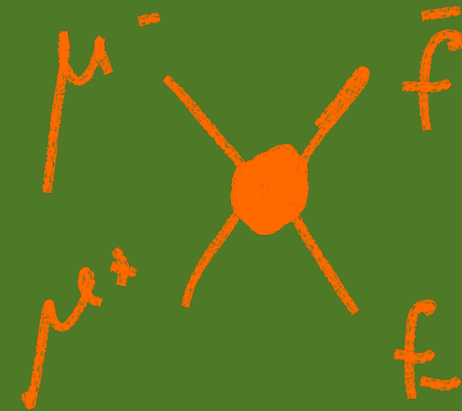
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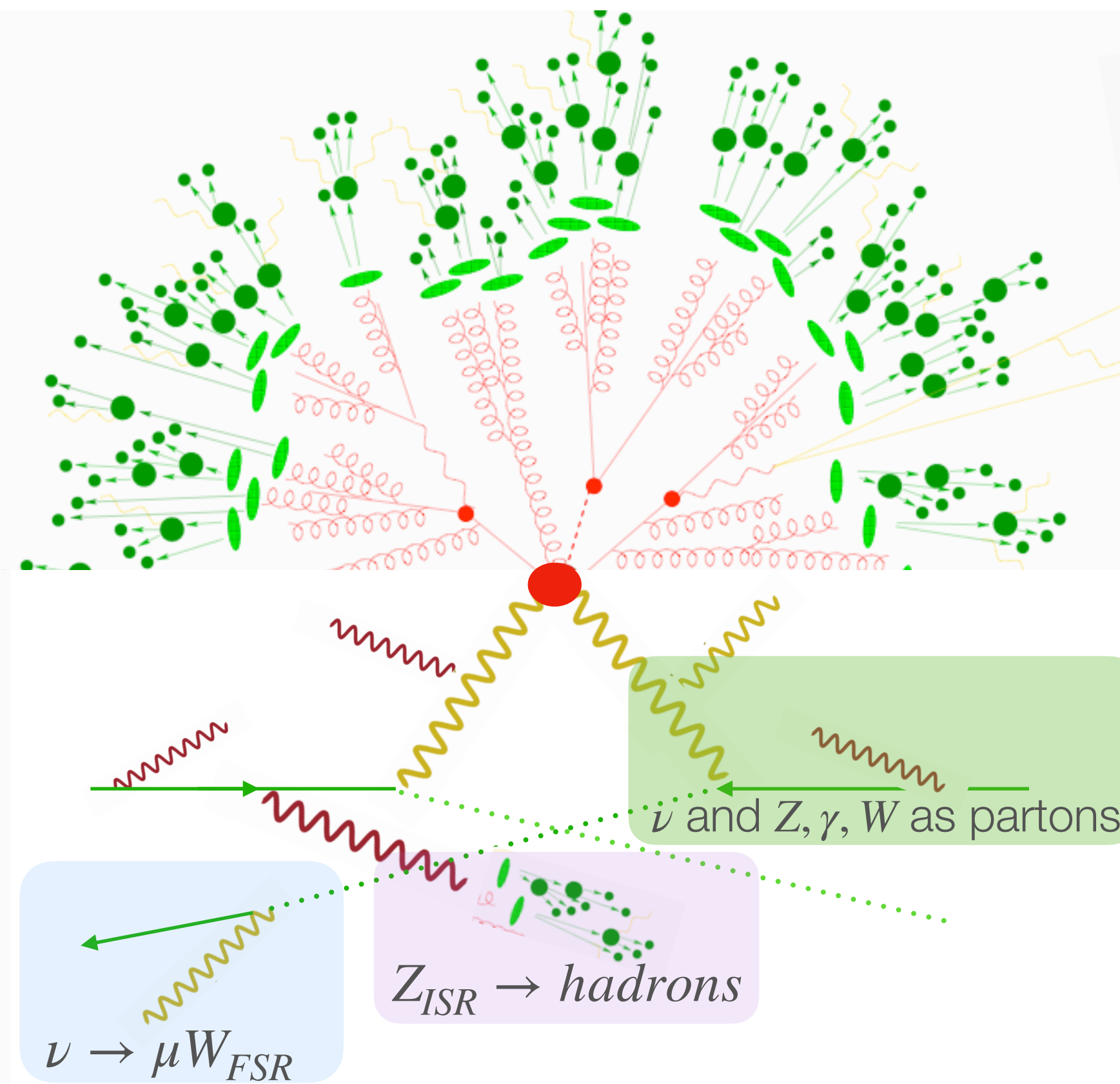
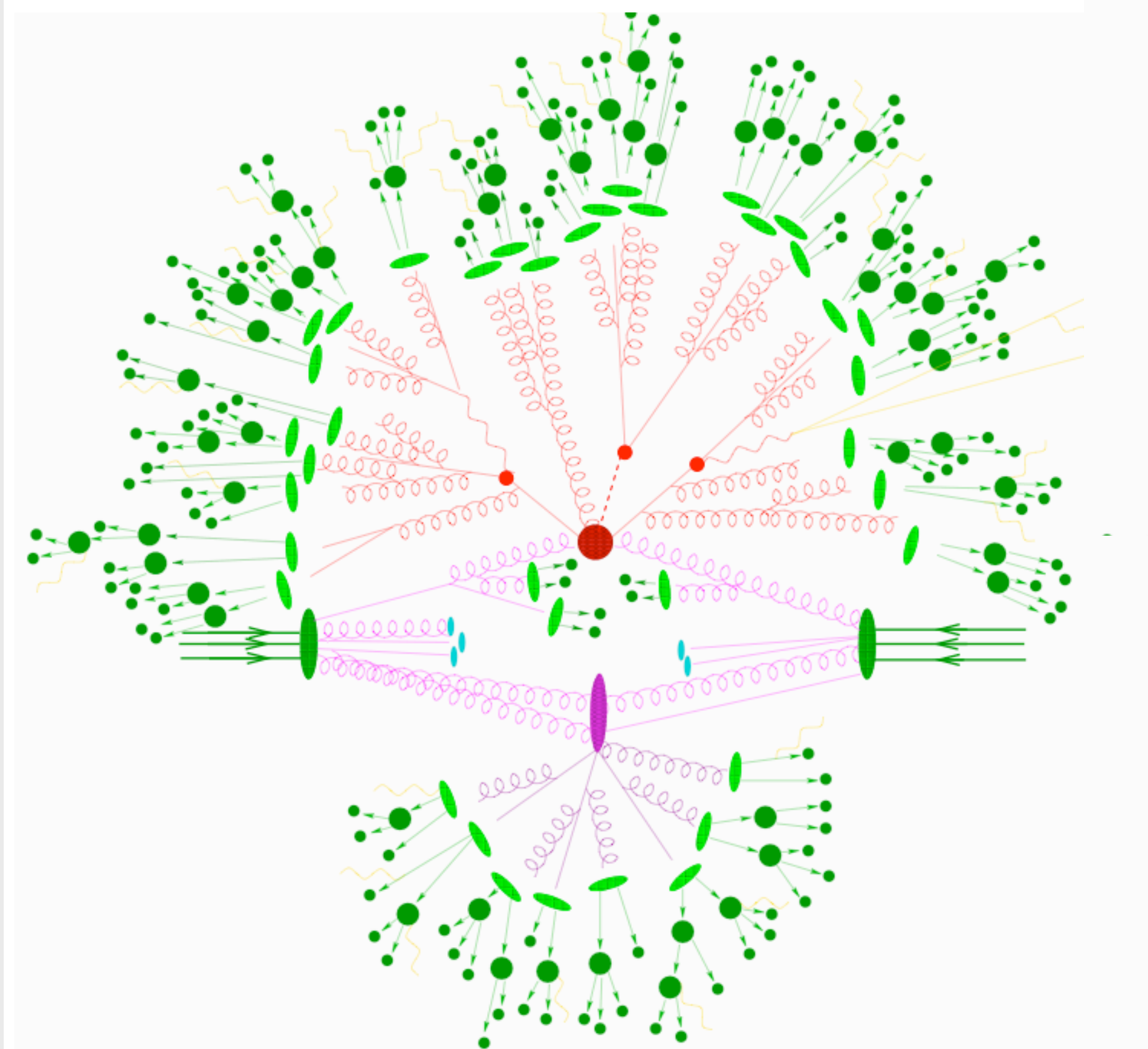
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STANDARD MODEL

“FACTORY”

tth production at the LHC (Fully hadronic)

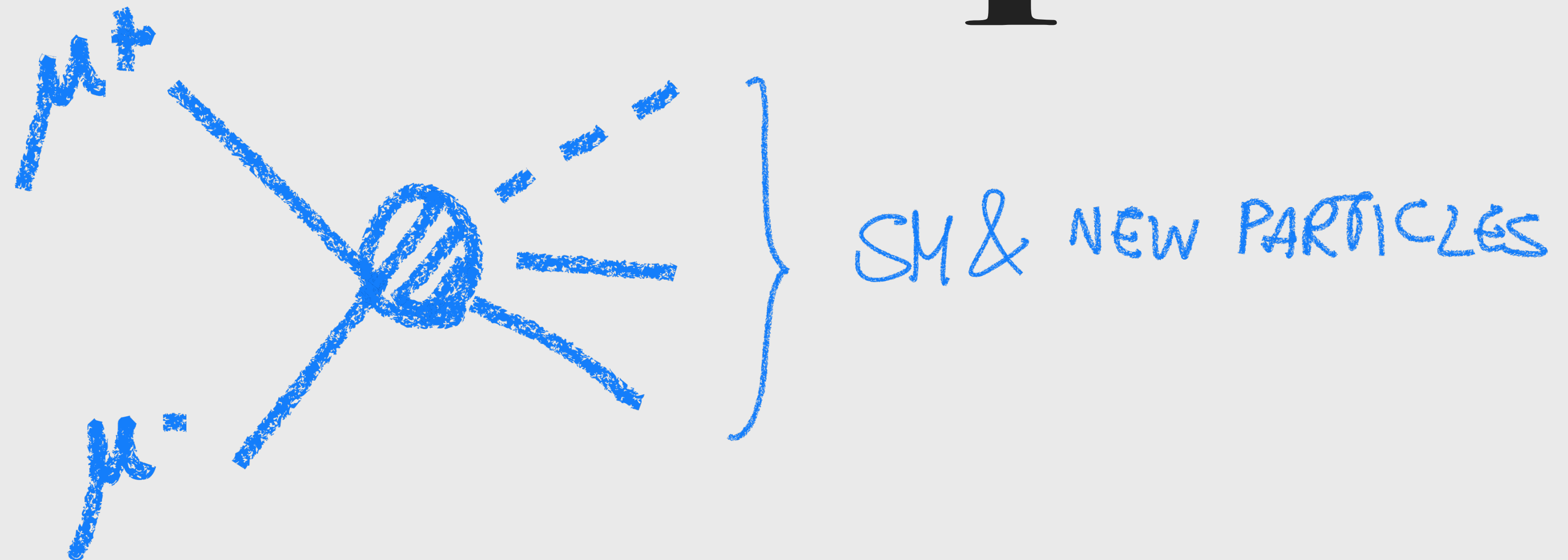
tth production at the muC 100 TeV (F. Maltoni)



NEW PHENOMENA AND
NEW REGIMES IN pQFT

- weak corrections become “ordinary”
- weak “partons”
- large EW logarithms

“Valence” Leptons

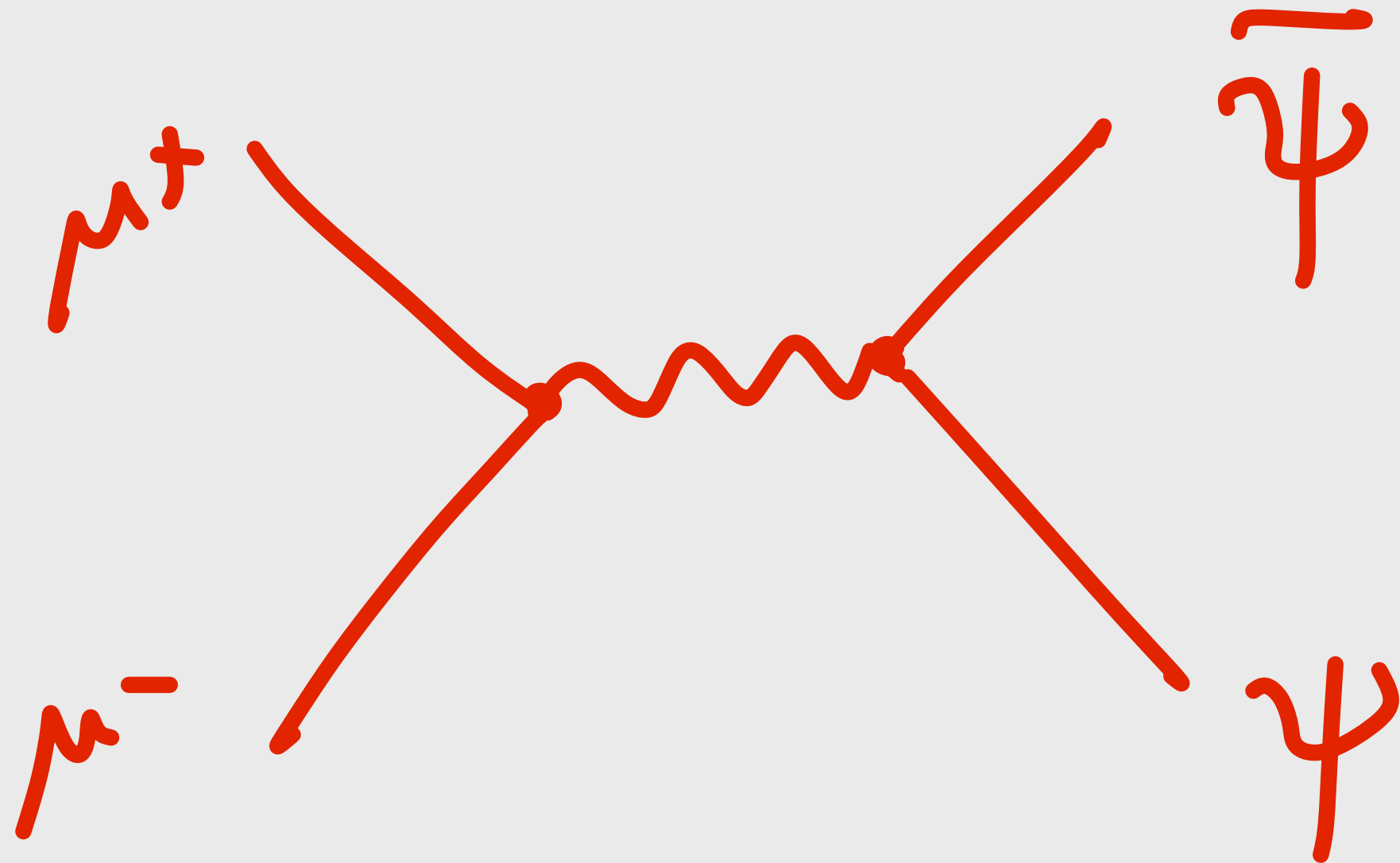


$\mu^+ \mu^- \rightarrow$ new physics

VALENCE

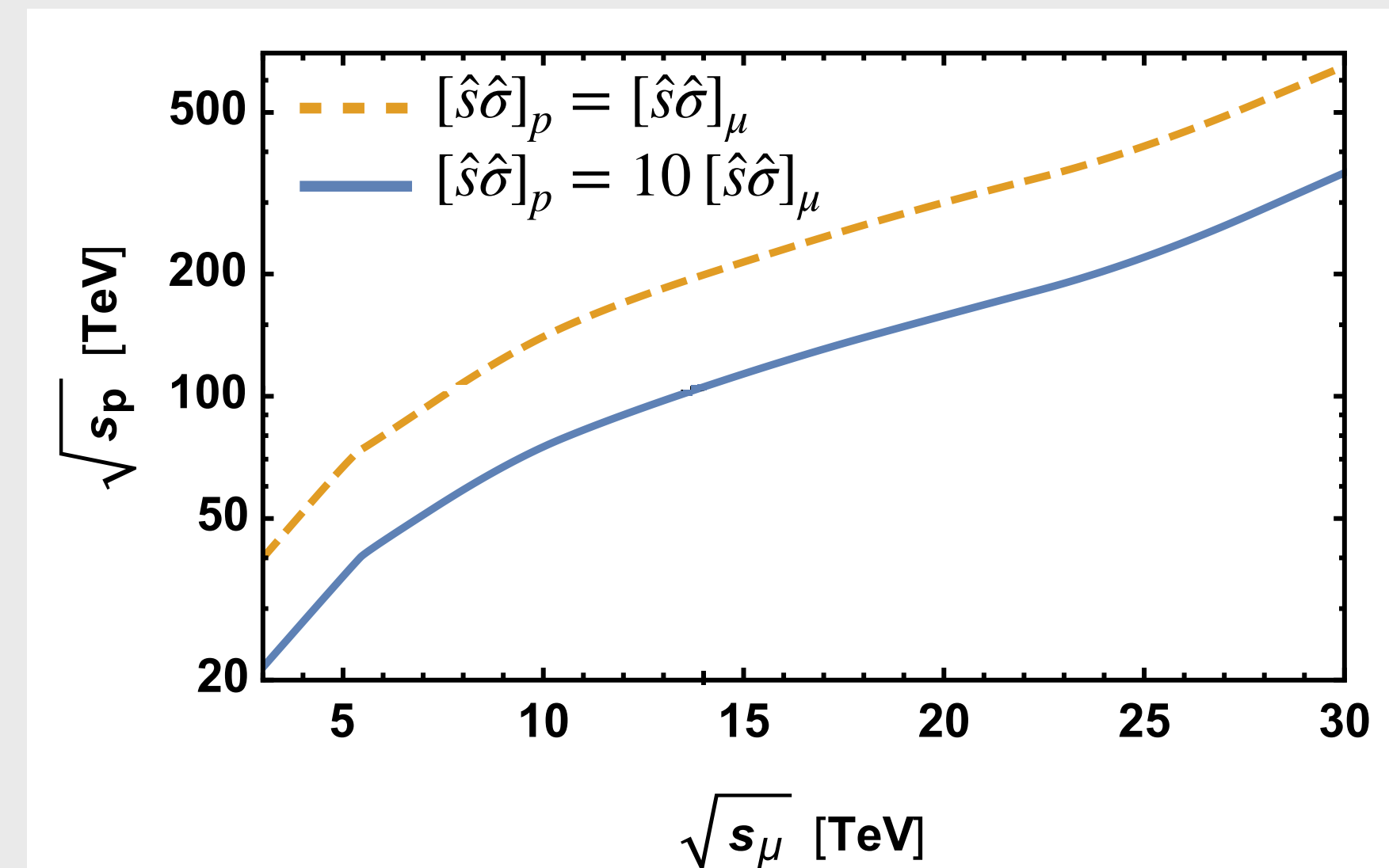
MUONS

Can produce heavy new physics (colored or not)



in principle can probe directly new states at $\frac{\sqrt{s}}{2}$ scale!

Compares pretty well with a pp collider 2005.10289
2209.01318

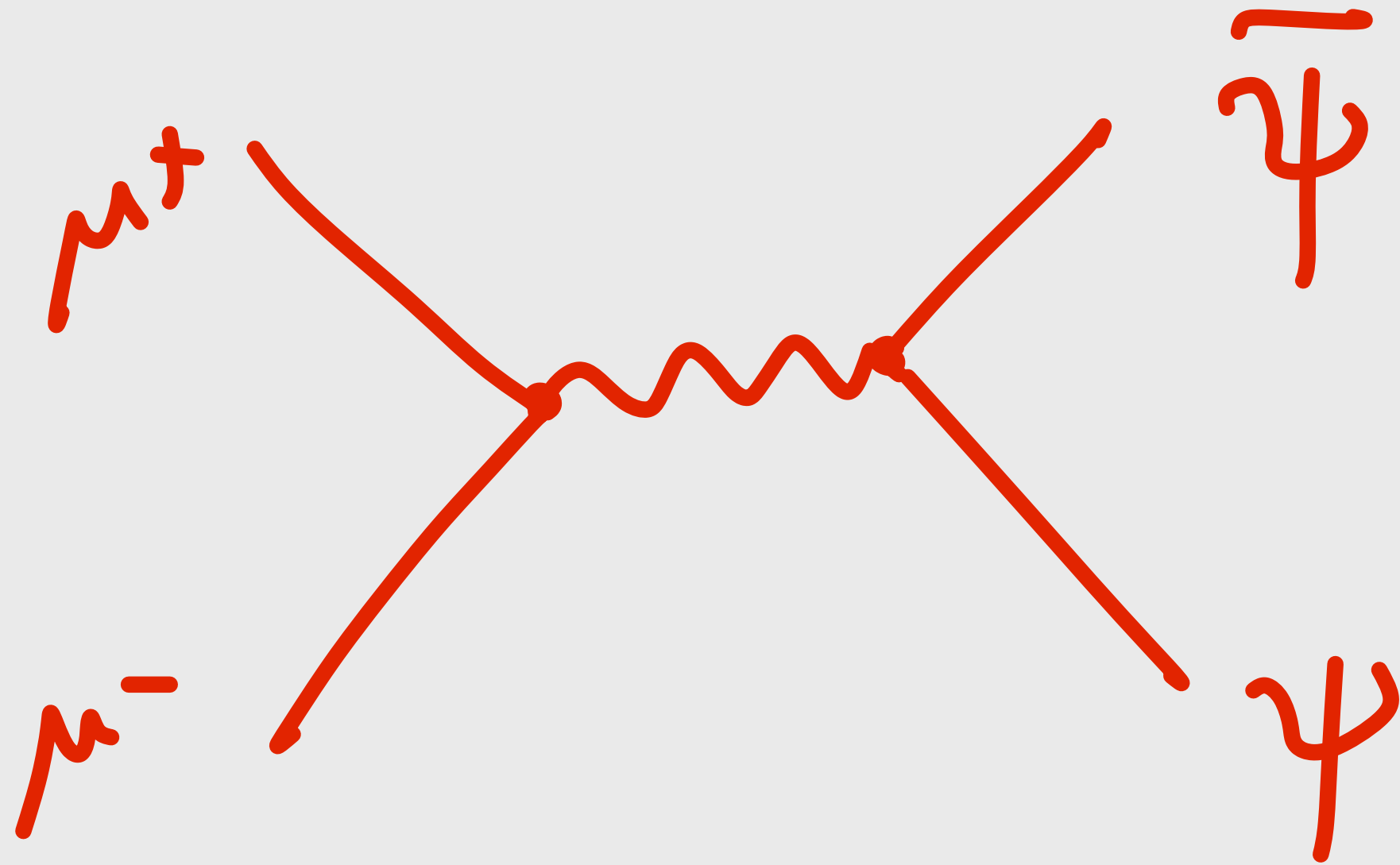


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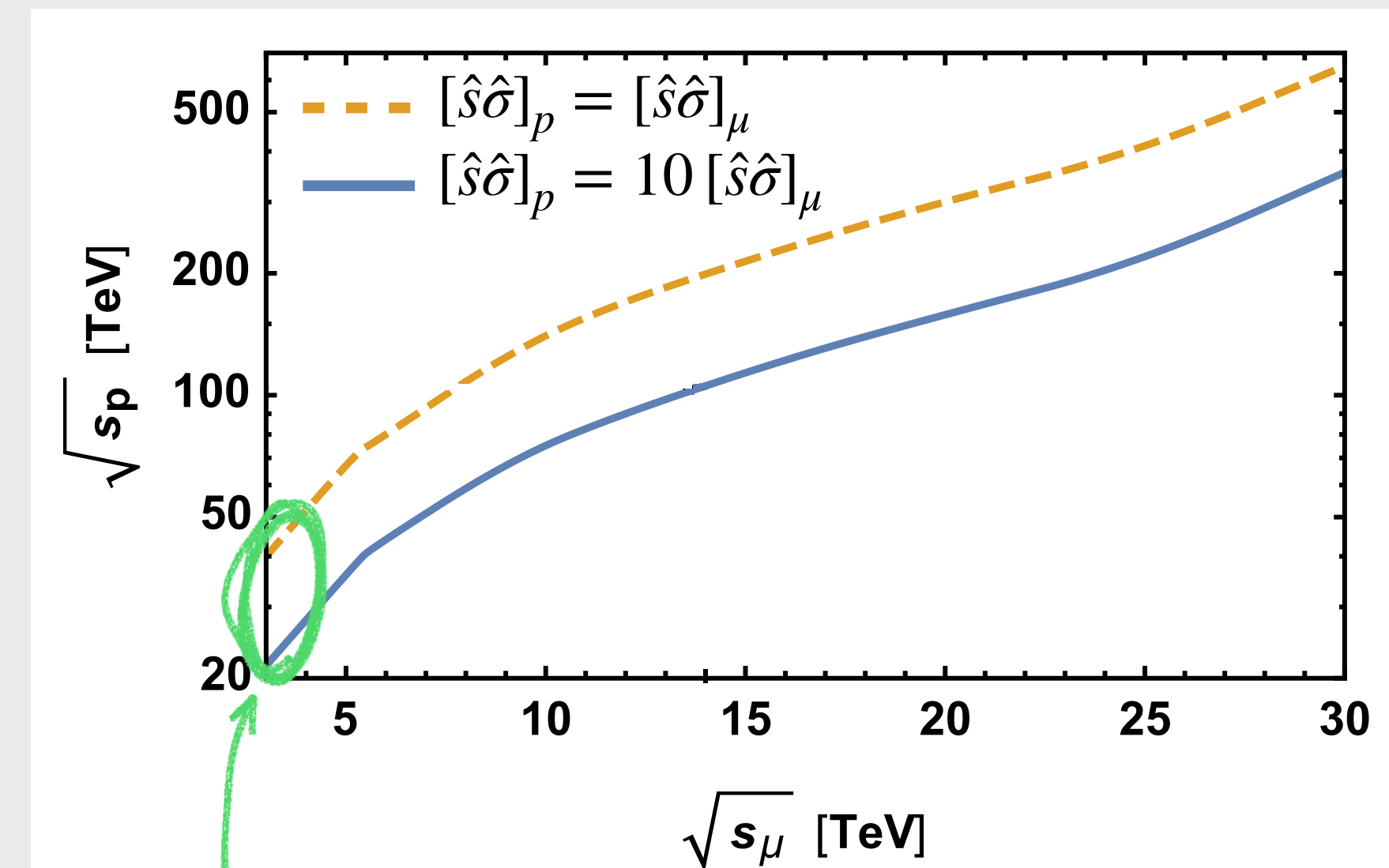
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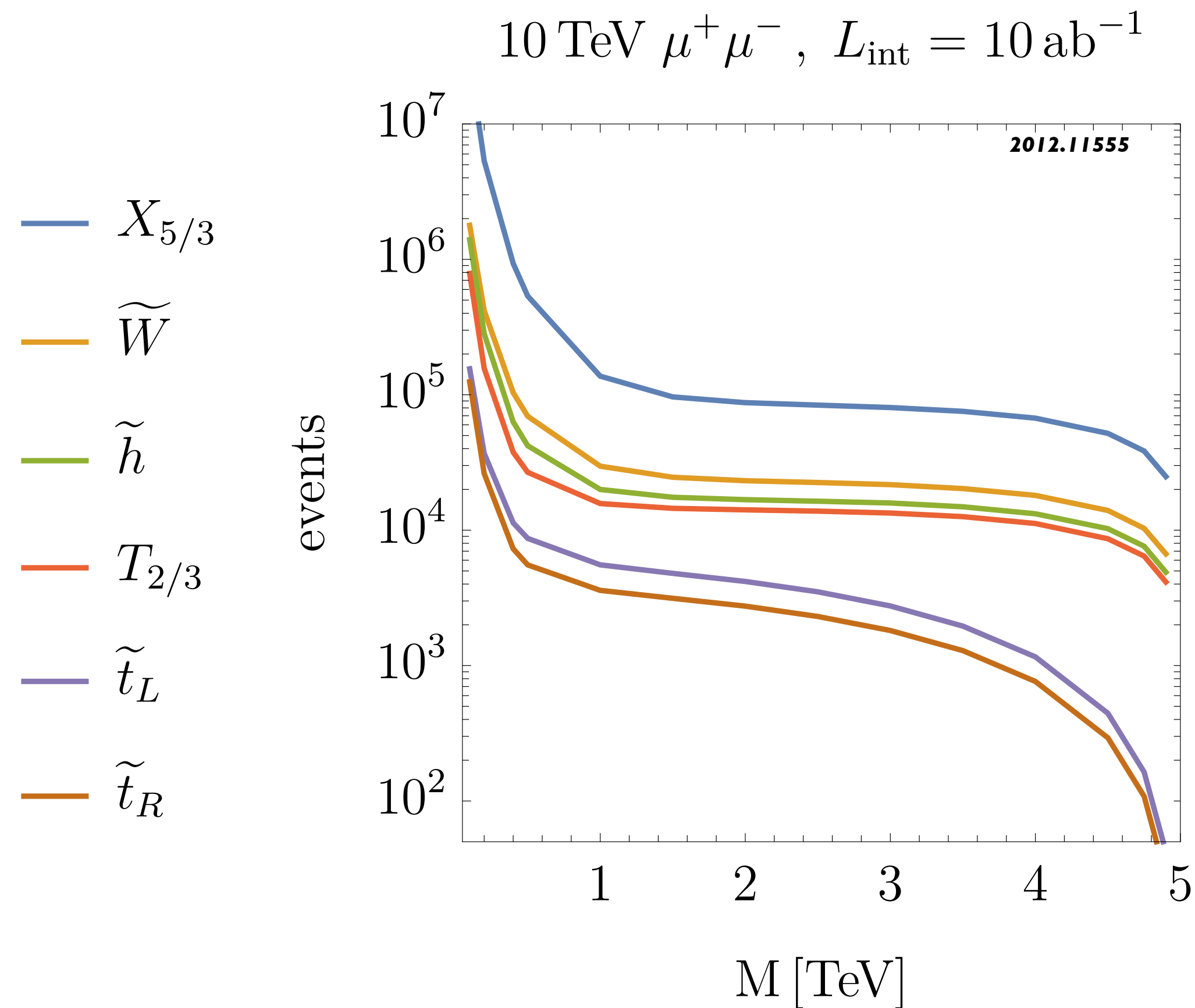


3 TeV $\mu^+ \mu^-$ roughly equivalent to 20+ TeV pp

$\mu^+ \mu^- \rightarrow$ new physics

VALENCE

MUONS



BEST POSITION TO OBSERVE ANY SIGN OF ELECTROWEAK NEW PHYSICS

(e.g. in the Higgs sector, or from new strong interactions at the TeV, fermions mass and mixing generation at the TeV)

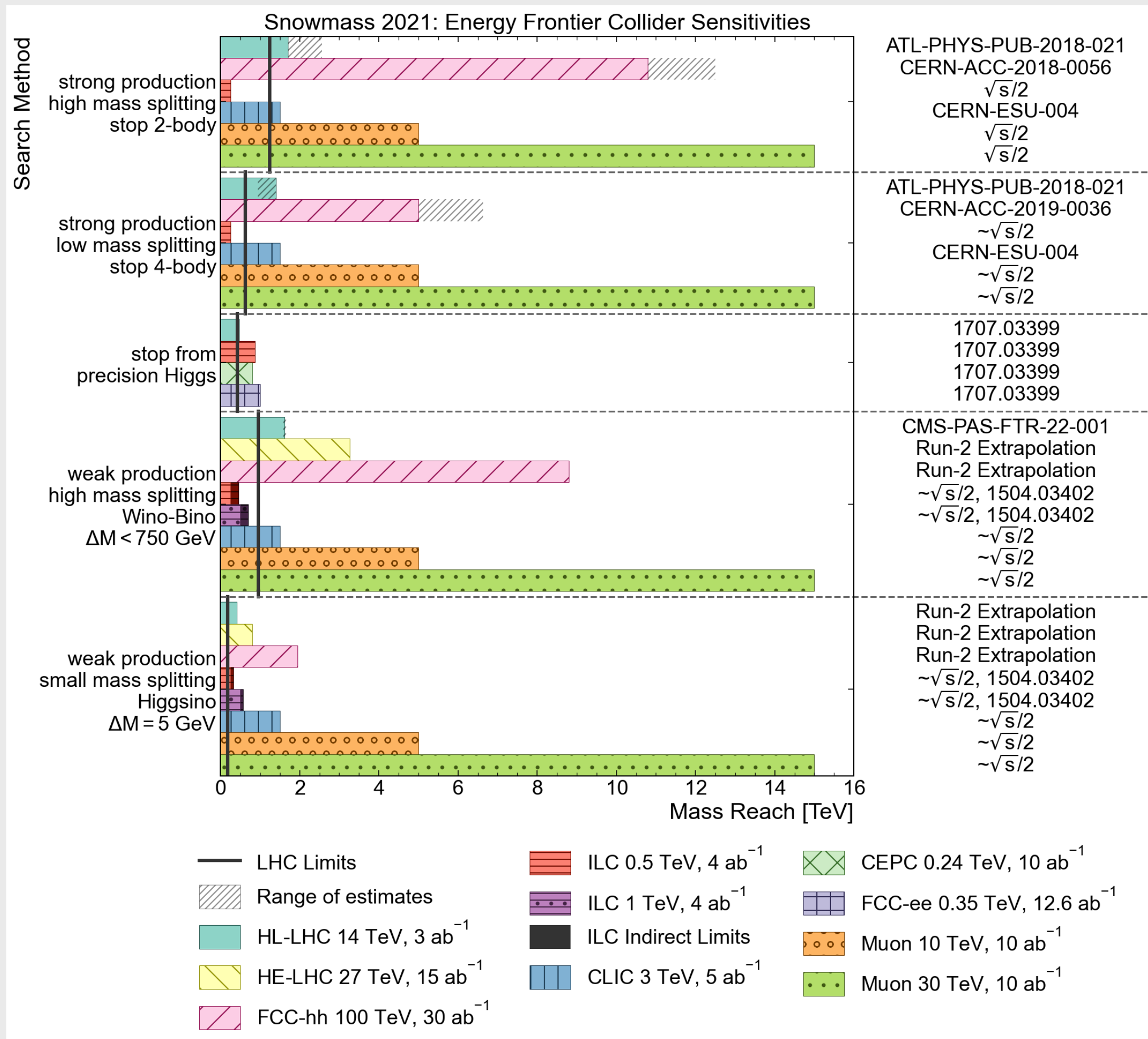
Any sign of SUSY below the TeV will be observable, no matter if the sparticles are colored or not.

$\mu^+ \mu^- \rightarrow$ new physics

VALENCE

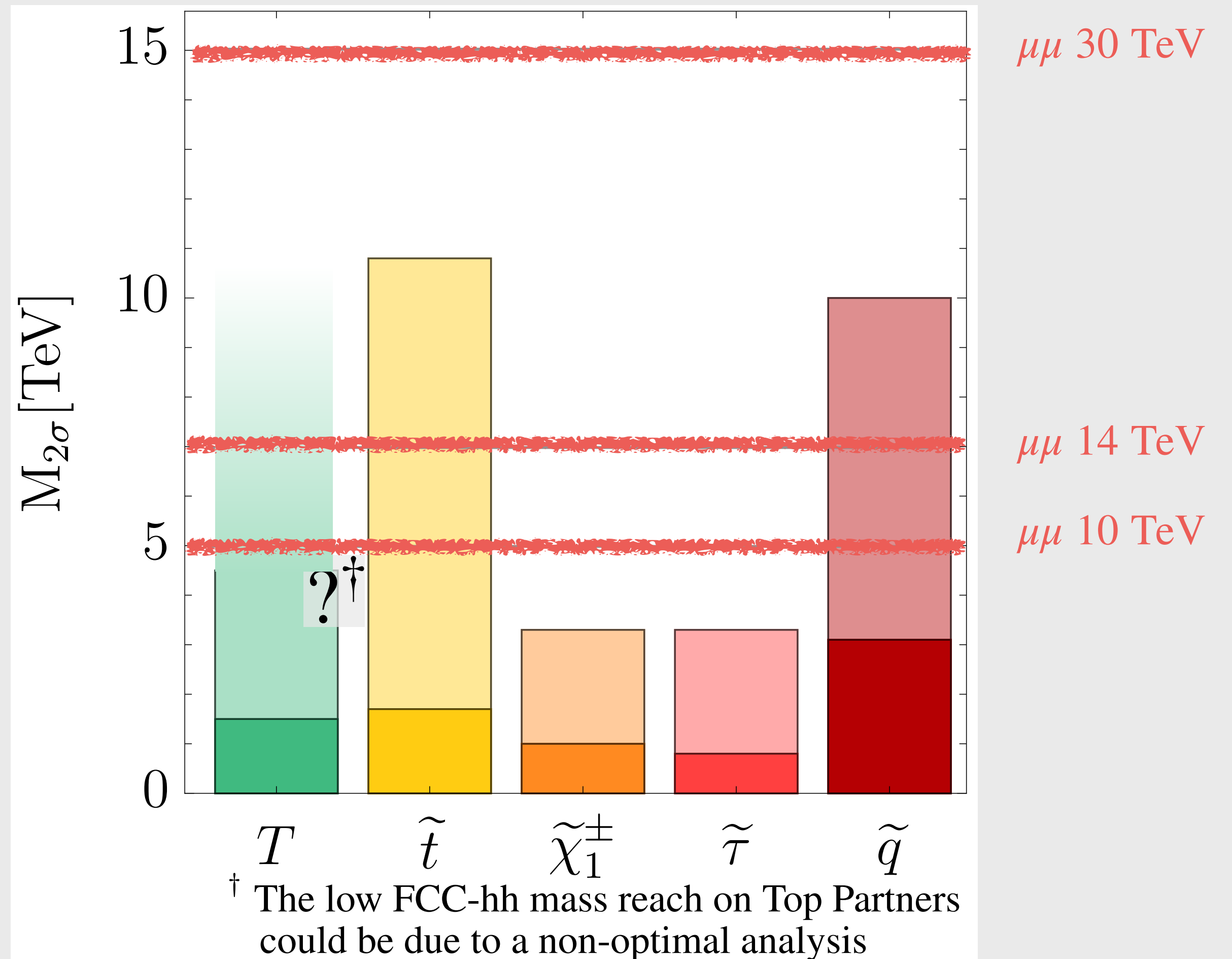
MUONS

2211.11084



FCChh
HL-LHC

2203.07256



2HDM

DOUBLETS

BIG JUMP AT FUTURE COLLIDERS

There is in general a weak sensitivity to new scalars, because of:

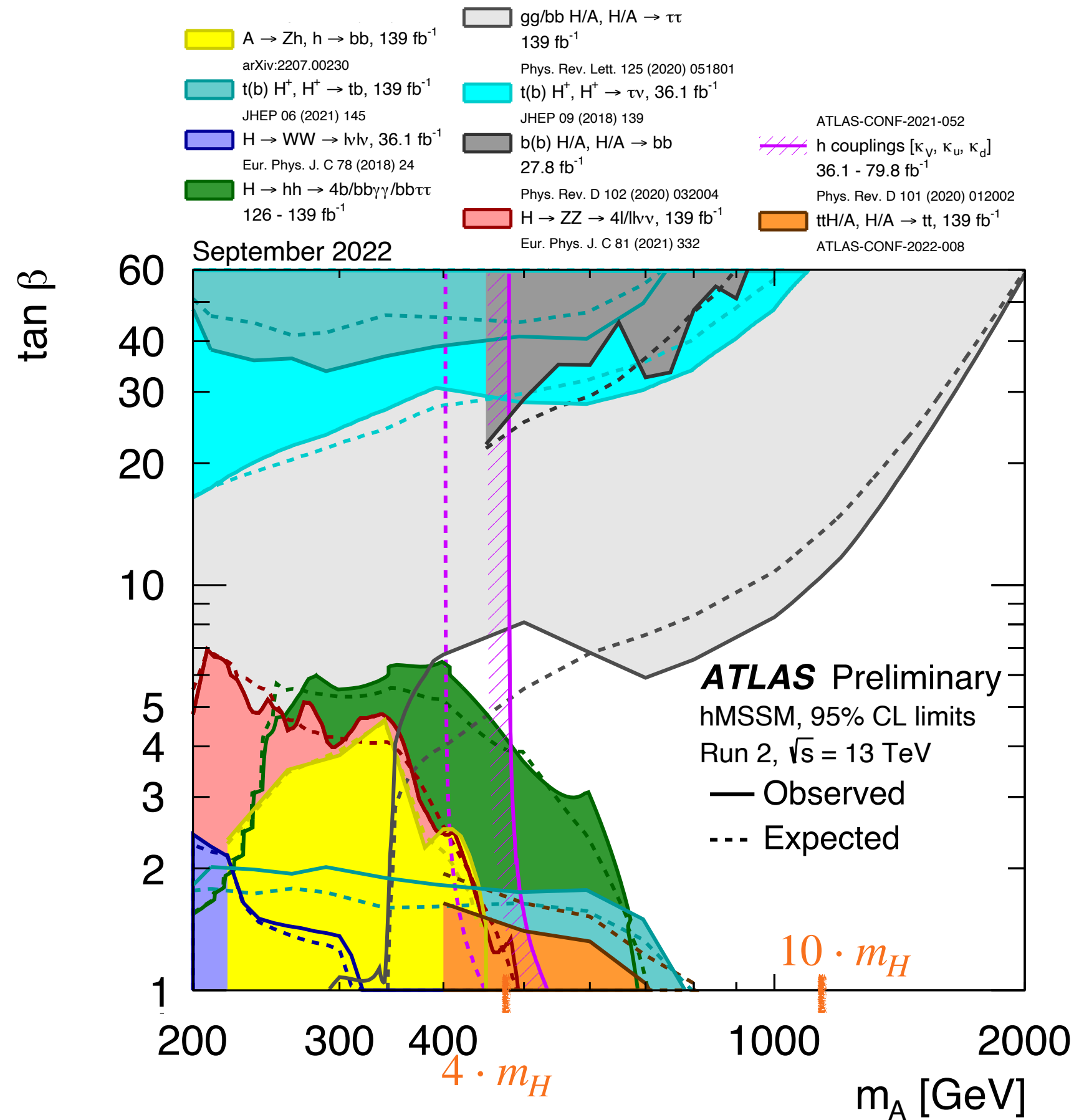
- **“small” cross-sections**
- **large backgrounds**

it is hard to explore the scalar sector and the only big discovery of the LHC may be left unmatched ... even if light scalars may exist.

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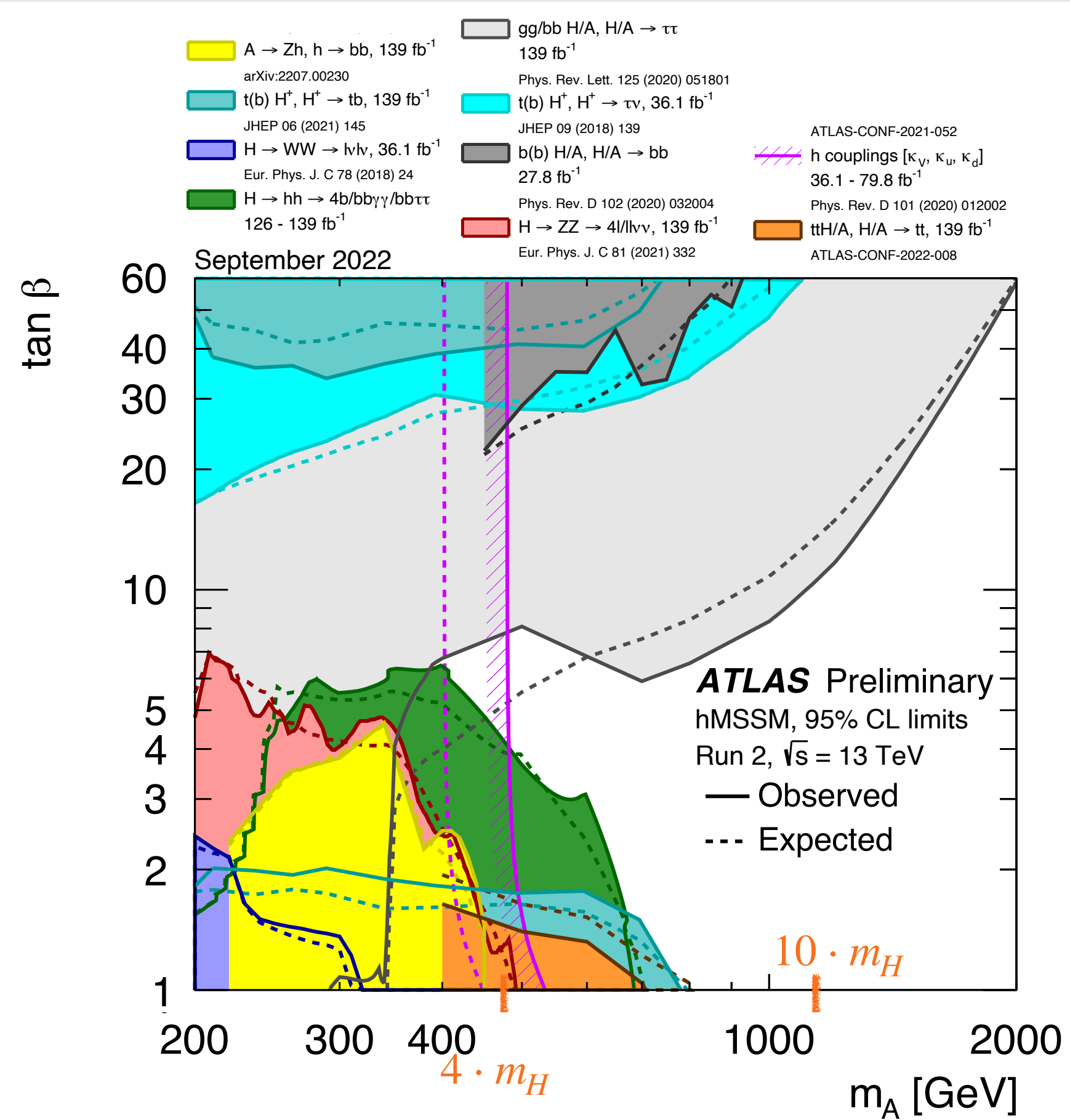
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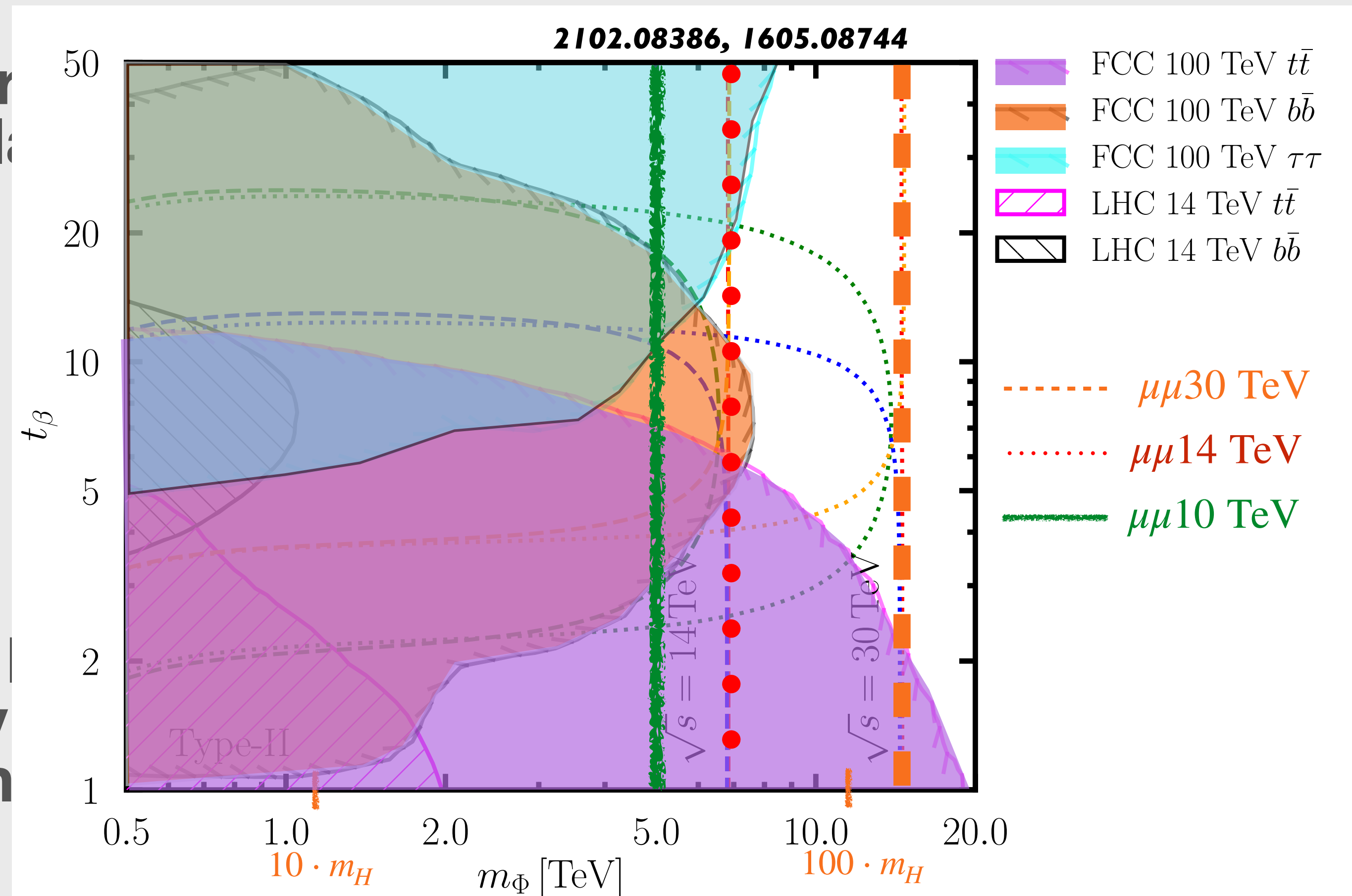
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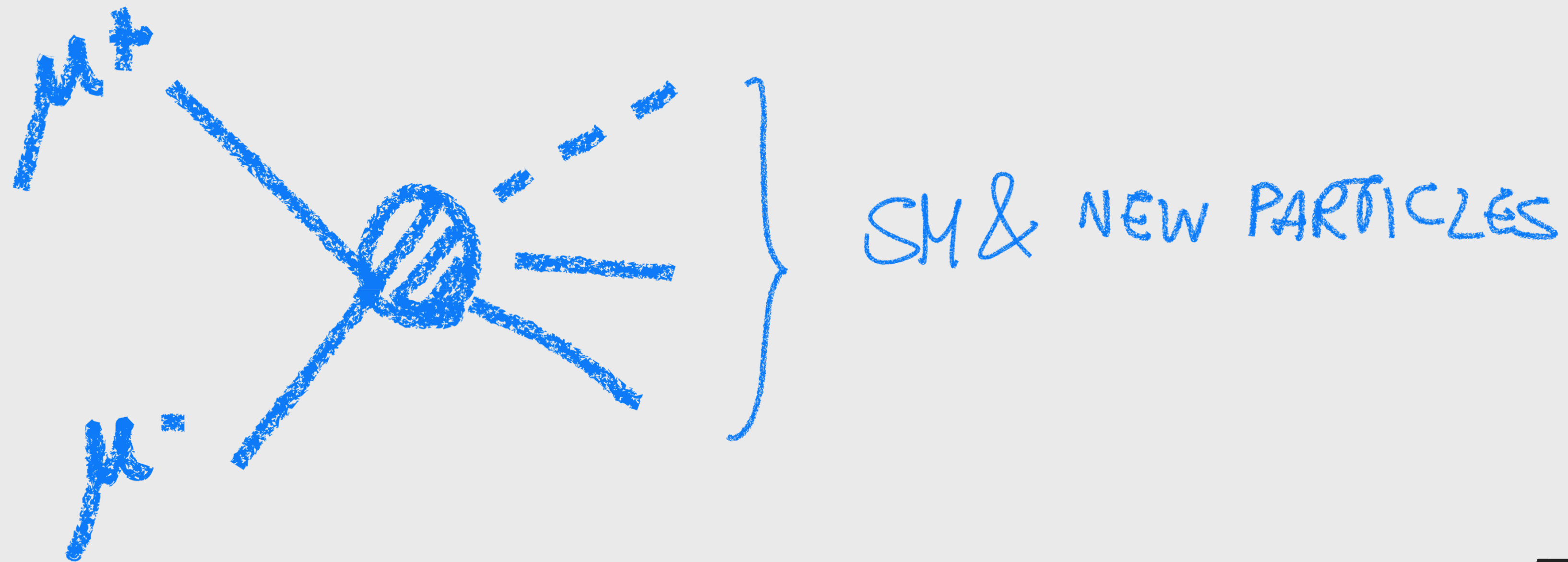


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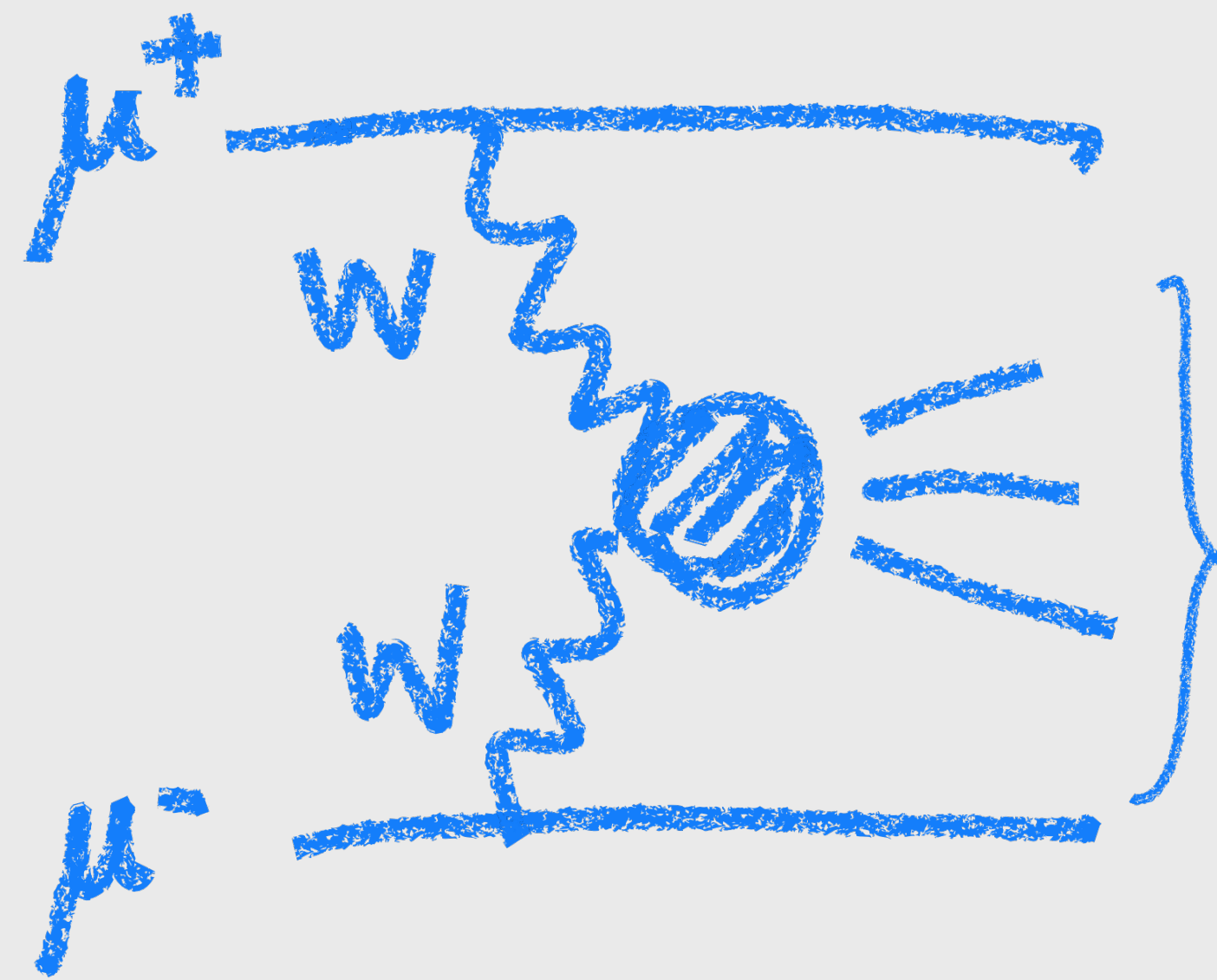
“Valence” Leptons



reach for new physics at $\sqrt{s}/2$

at $\sqrt{s} \gg 100 \text{ GeV}$

Weak Bosons collider



SM & NEW PARTICLES



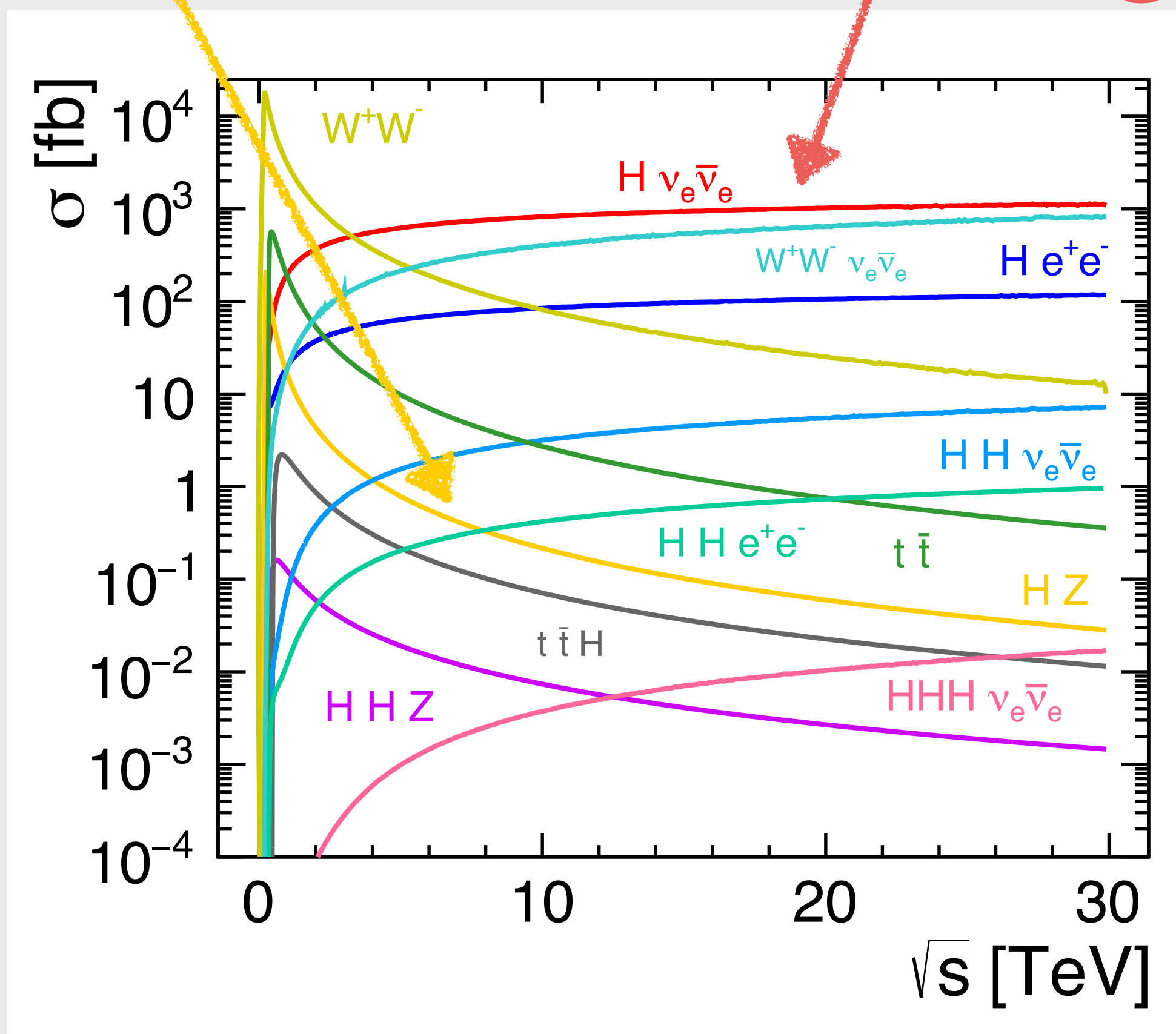
Higgs boson



10⁶ HIGGS BOSONS

MEGA-HIGGS FACTORY

$\sigma \sim 1/s$ $\sigma \sim \log(s)$



At 3 TeV the weak bosons are sufficiently light that can be radiated very efficiently

$$\sqrt{s} = 3 \text{ TeV}$$

$$\sigma \cdot \mathcal{L} \Rightarrow O(10^6) \text{ h}$$

- large number of Higgs bosons!

FURTHER OPPORTUNITIES

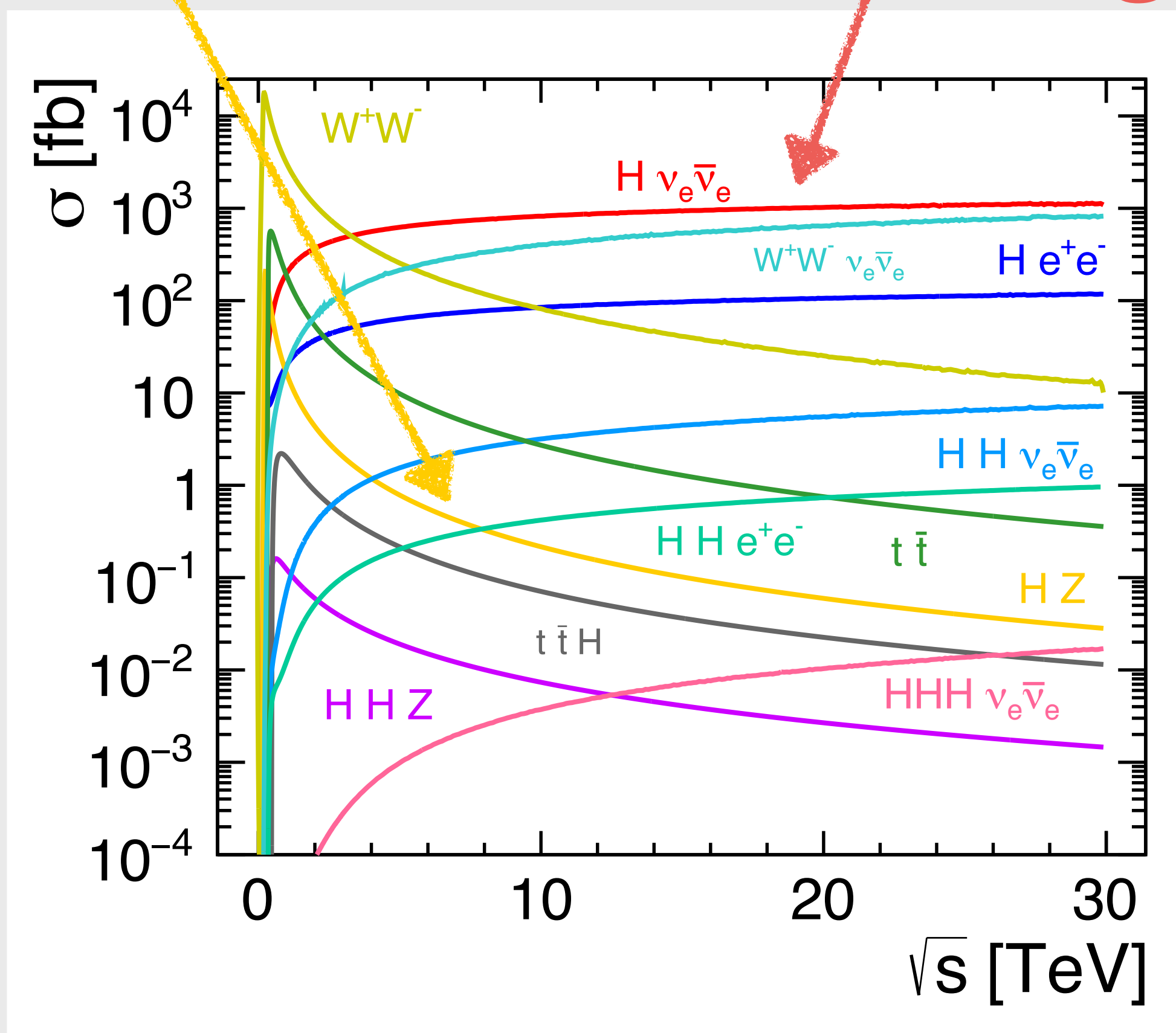
- ultra-rare Higgs decays
- differential distribution
- off-shell Higgs bosons
- rare production modes



10⁶ HIGGS BOSONS

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At 30 TeV the weak bosons are sufficiently light that can be radiated very efficiently

$$\sqrt{s} = 30 \text{ TeV}$$

$$\sigma \cdot \mathcal{L} \Rightarrow O(10^8) \text{ h}$$

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FURTHER OPPORTUNITIES

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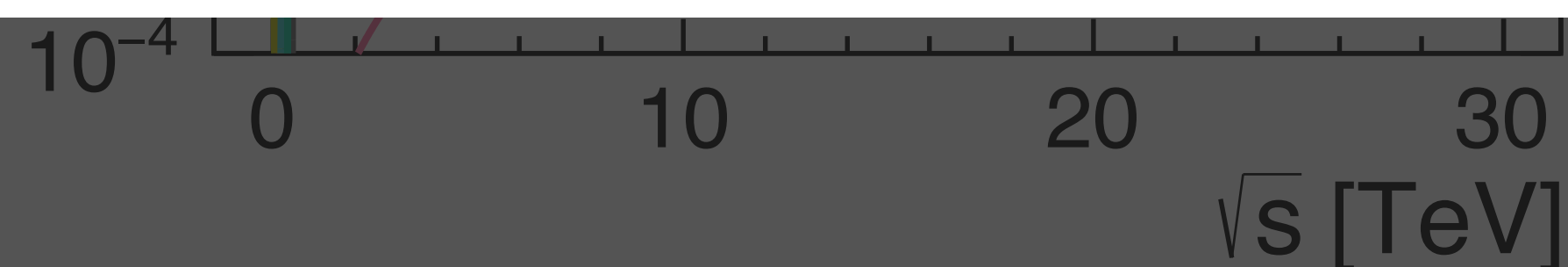
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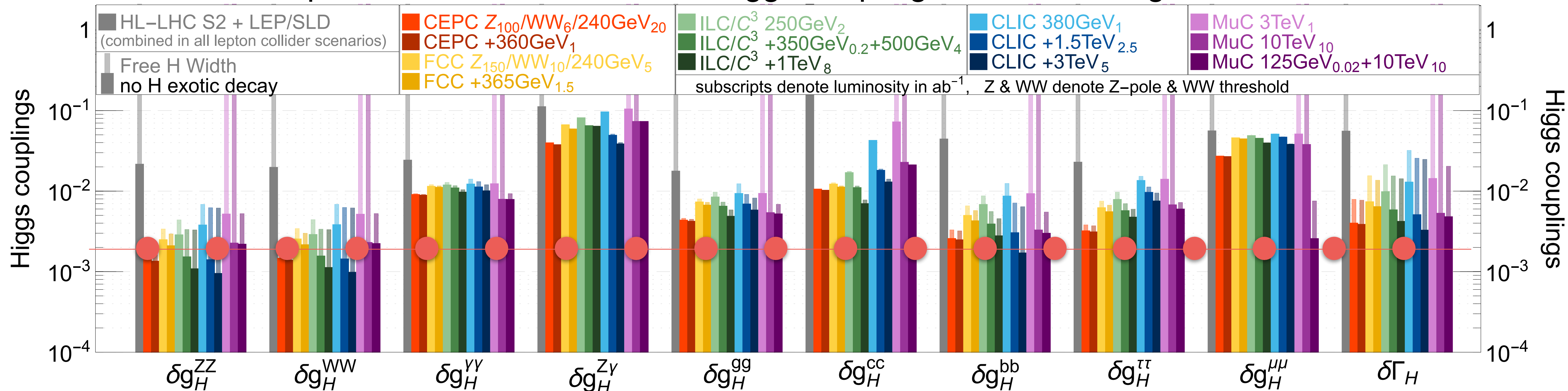
- Higgs factory at 3 TeV
- 10 × Higgs factory at 10 TeV
- 100 × Higgs factory at 30 TeV



- differential distribution
- off-shell Higgs bosons
- rare production modes

Summary: Higgs@FC (by couplings)

precision reach on effective Higgs couplings from SMEFT global fit



0.1% coupling precision, sensitivity to new physics at $10\text{ TeV} \simeq 100 \cdot m_h$

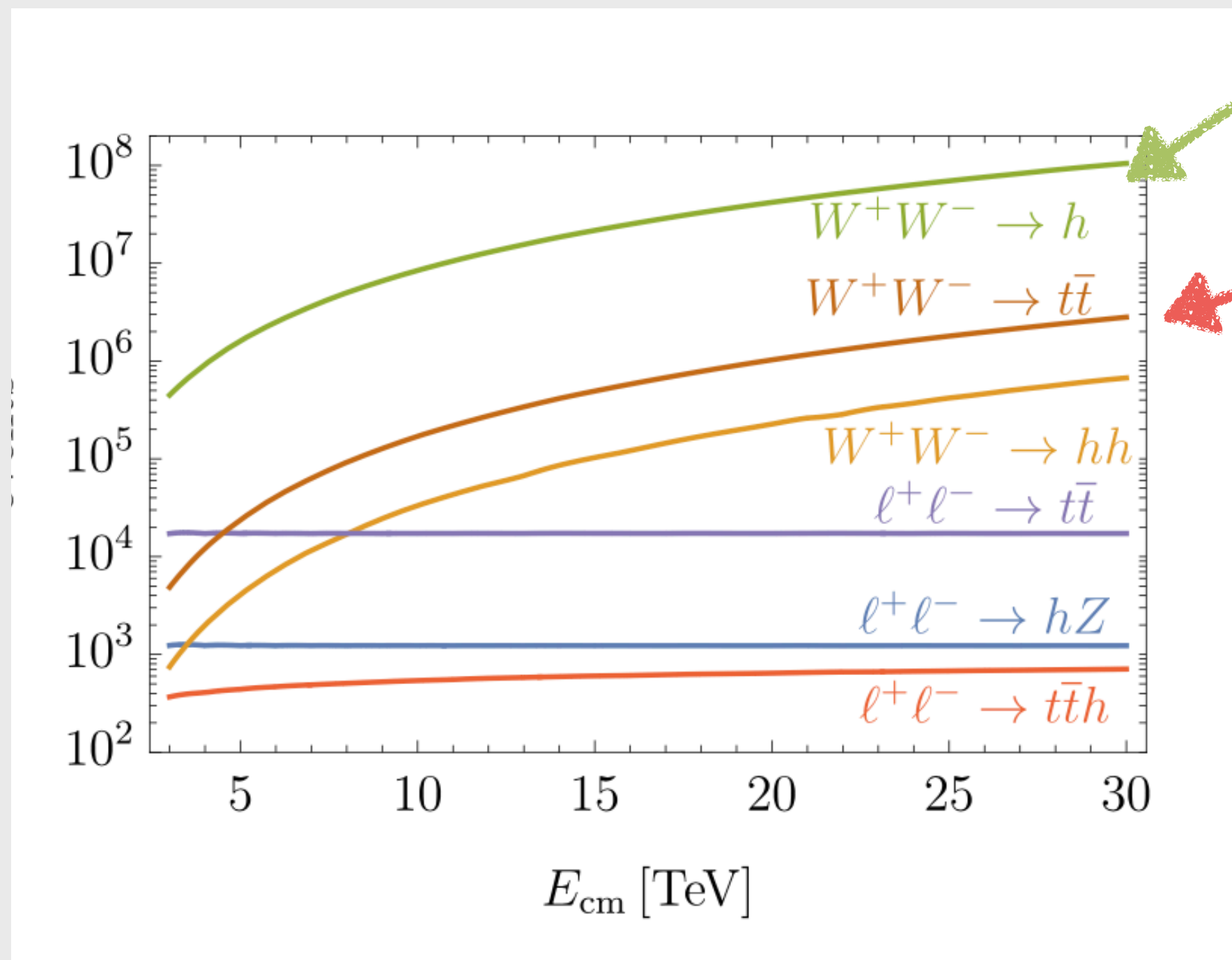
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STANDARD MODEL

“FACTORY”

$\sigma(h\nu\nu)$

$\sigma(tt\nu\nu)$



- large number of top quarks!
- rare top decays
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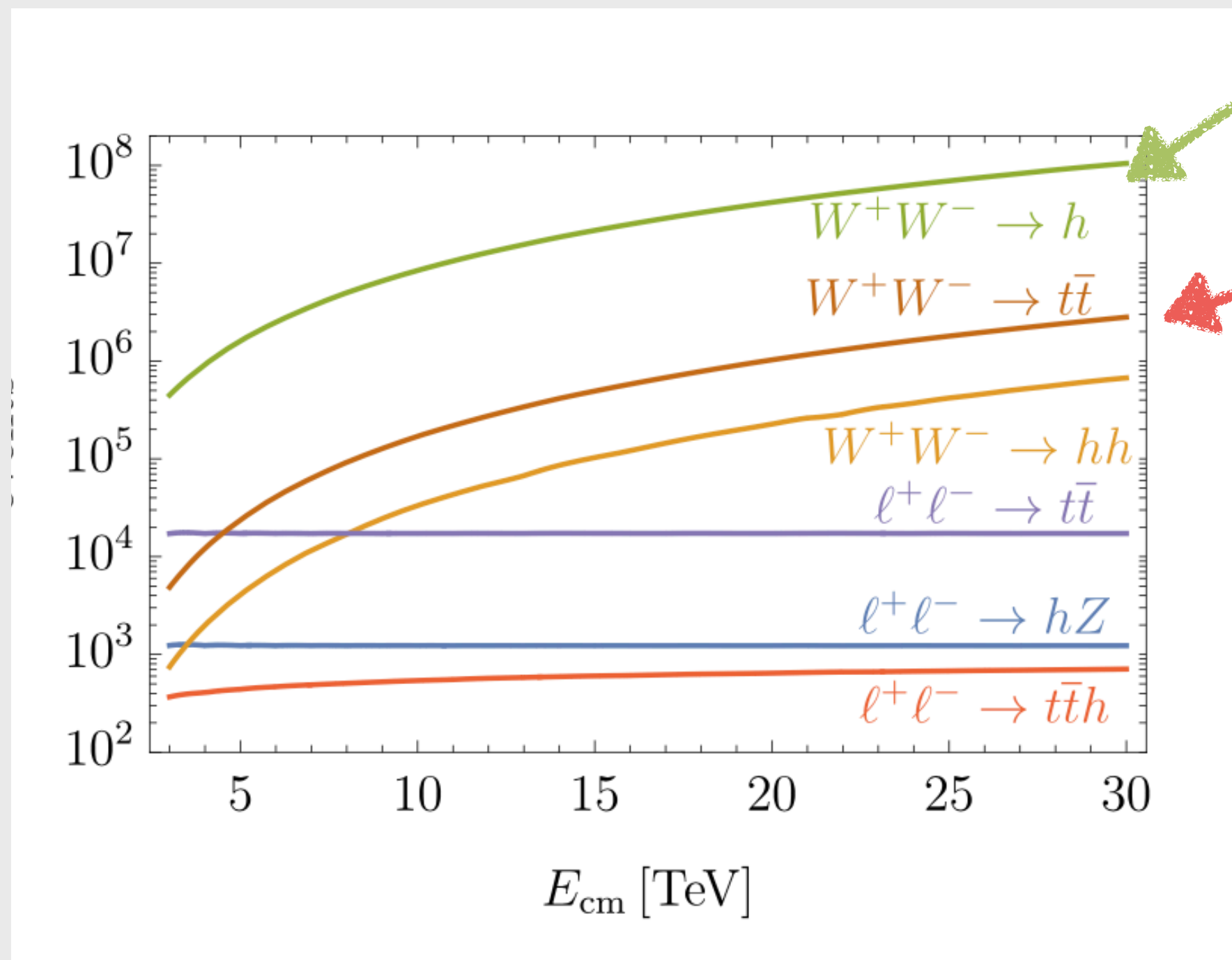
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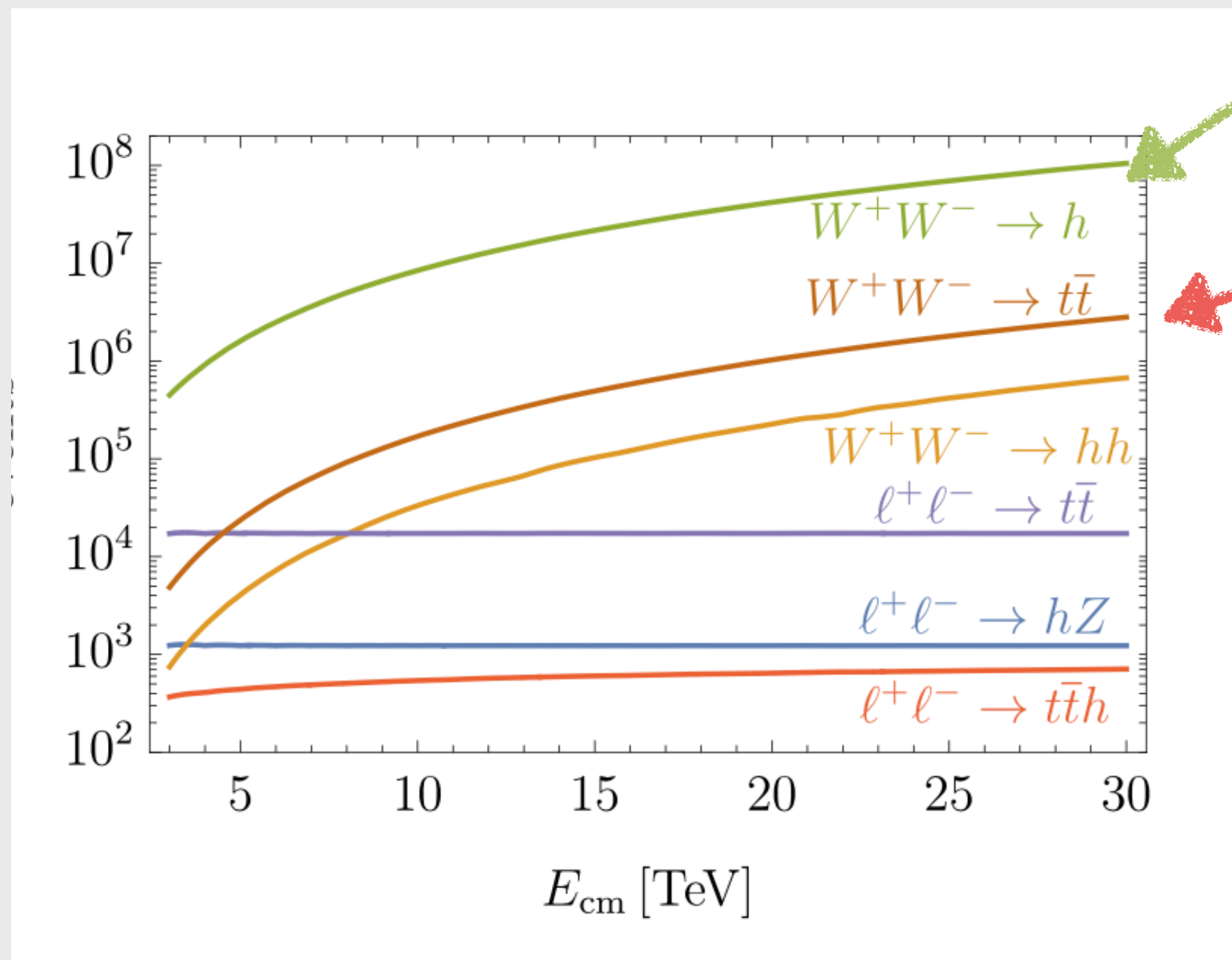
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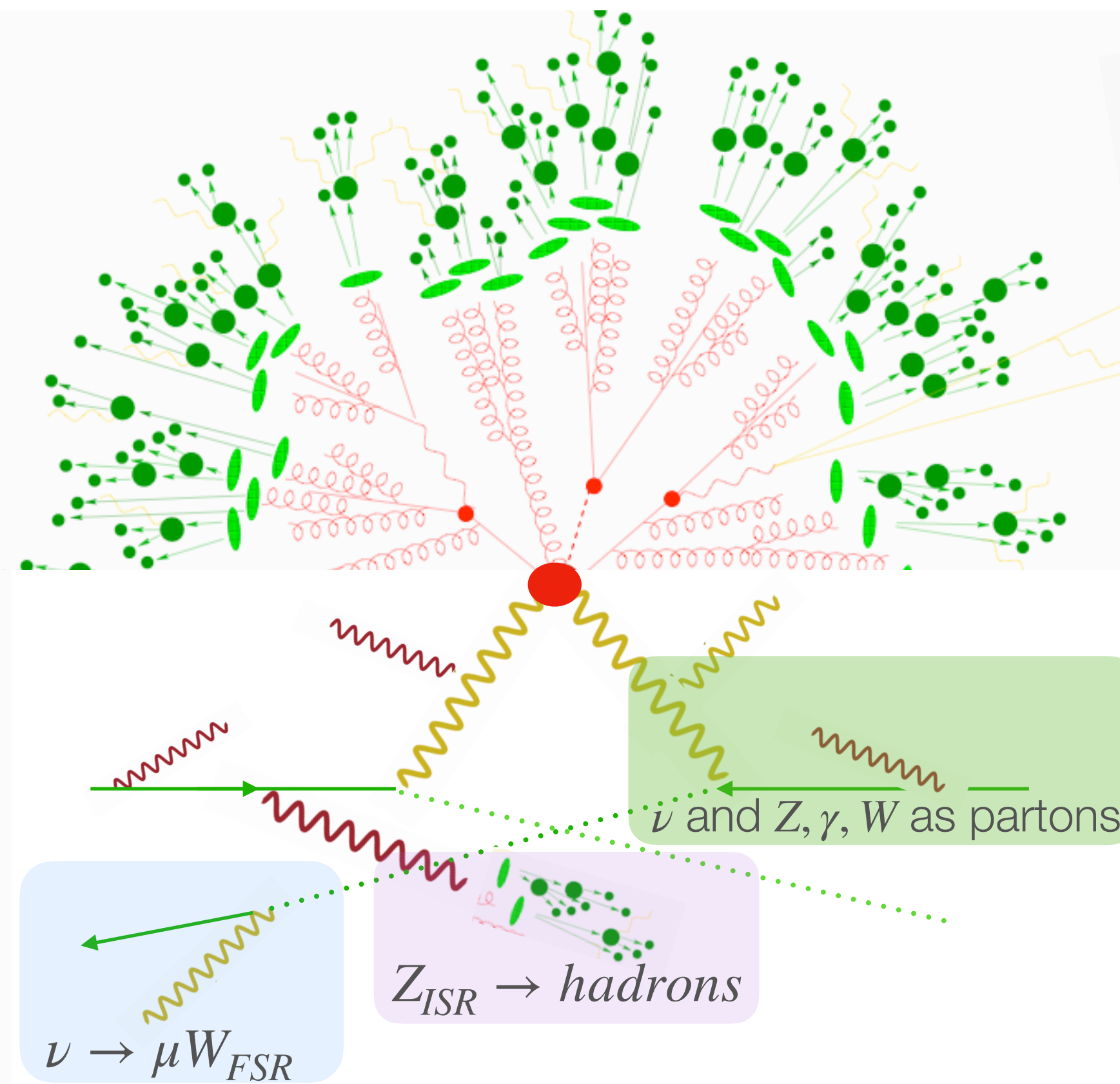
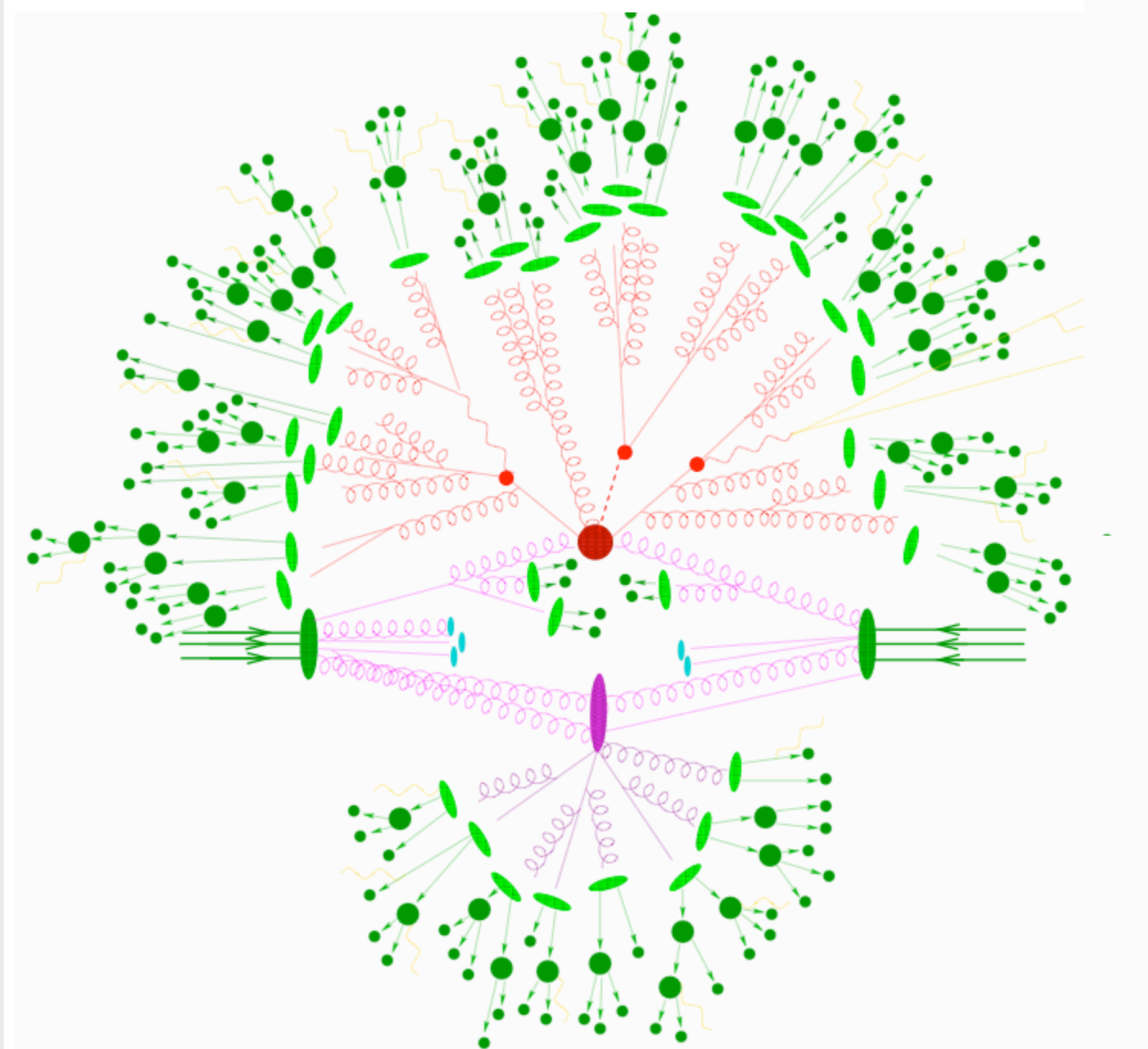
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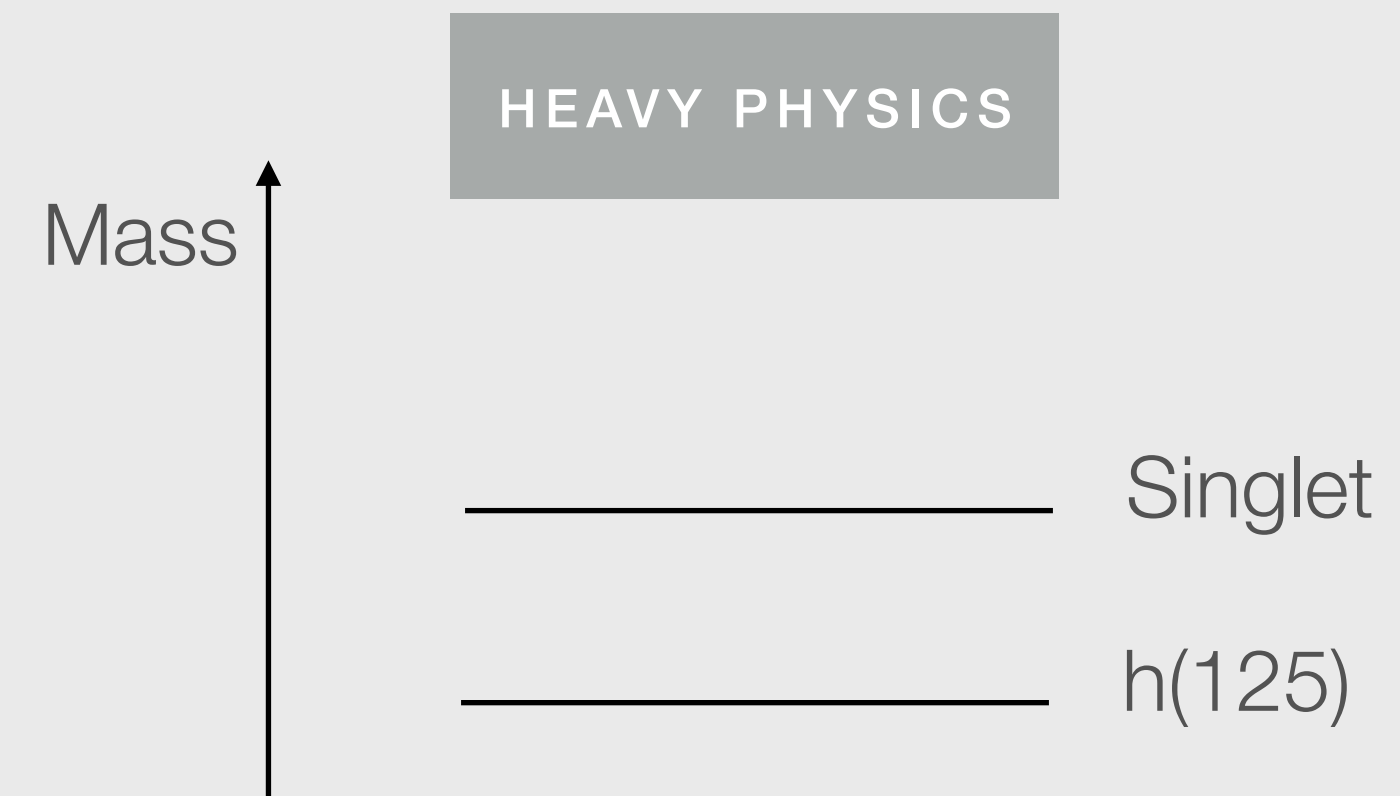
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Higgs + Singlet

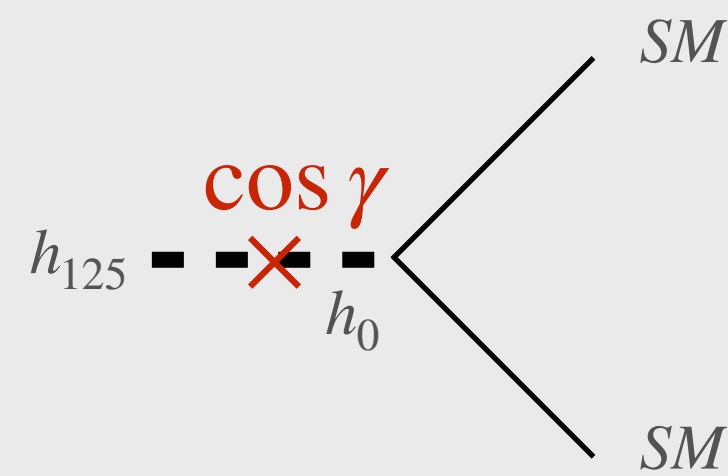


- Broad coverage of BSM scenarios: *(N)MSSM, Twin Higgs, Higgs portal, modified Higgs potential (Baryogenesis)*
- Phenomenology is also useful as “simplified model”

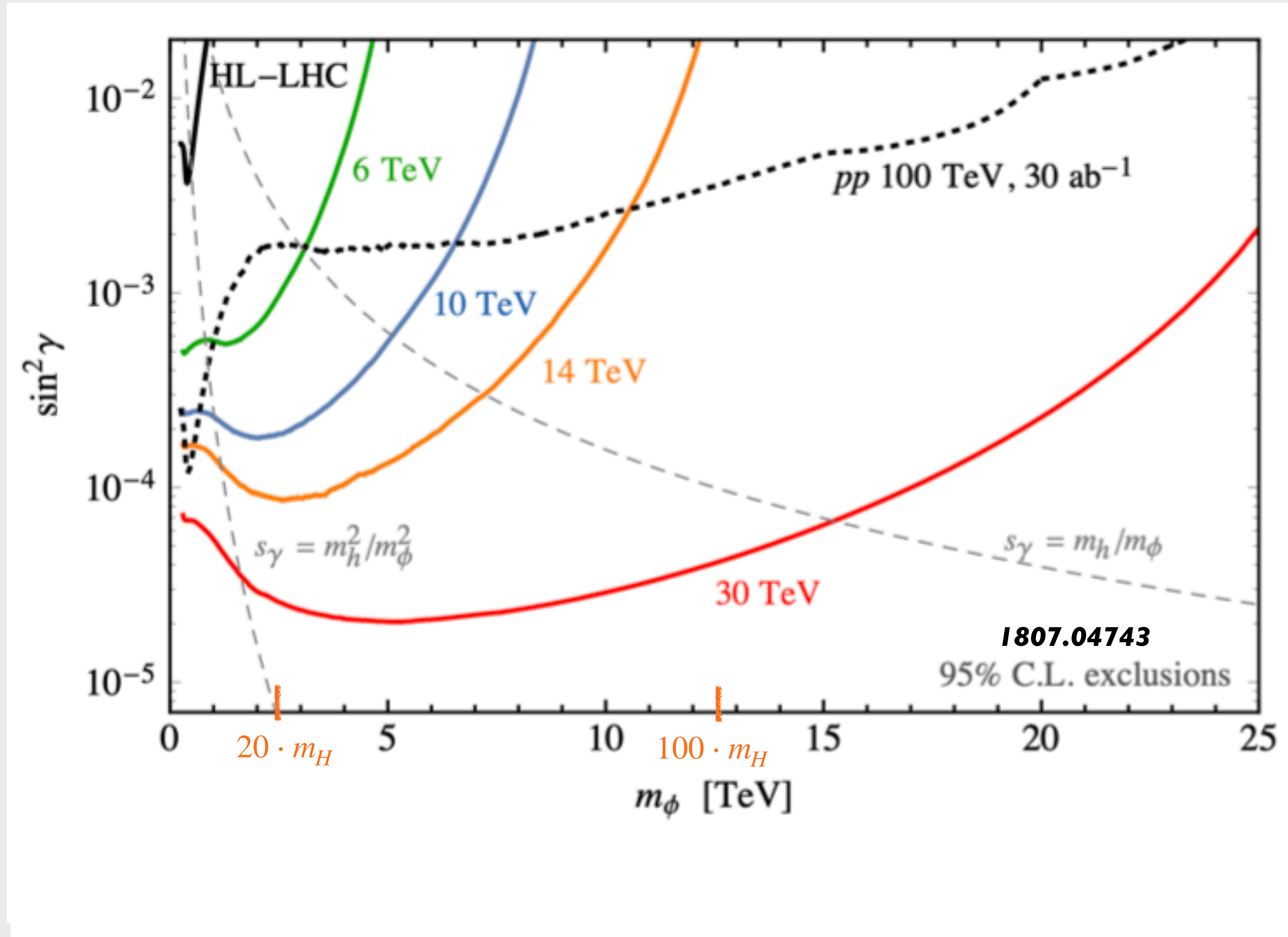
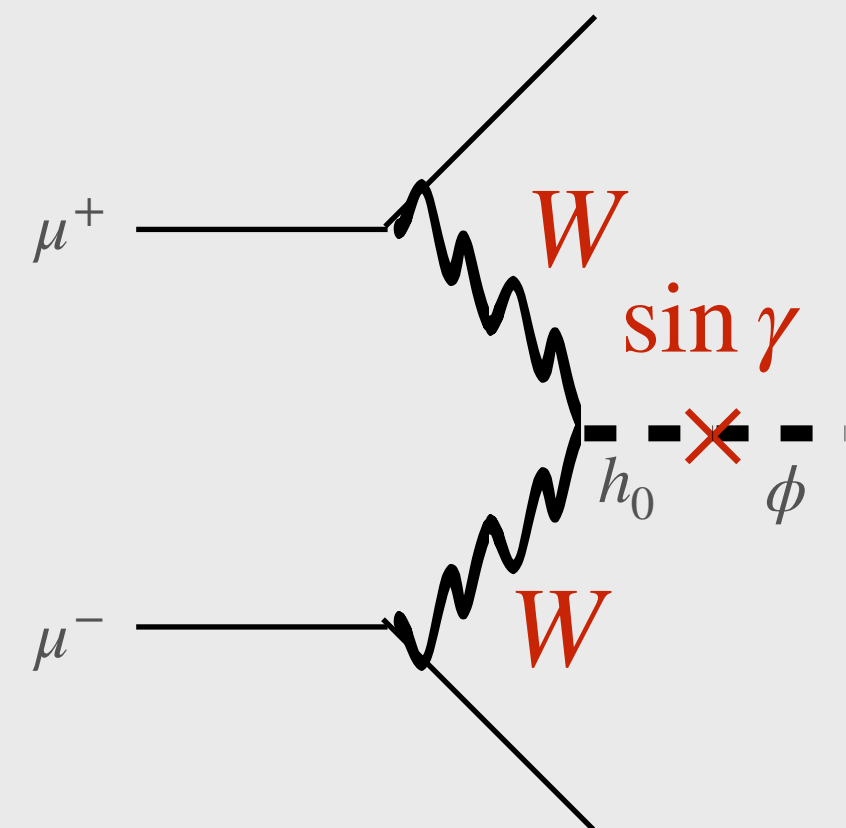
Higgs + Singlet

SINGELTS

BIG JUMP AT FUTURE COLLIDERS



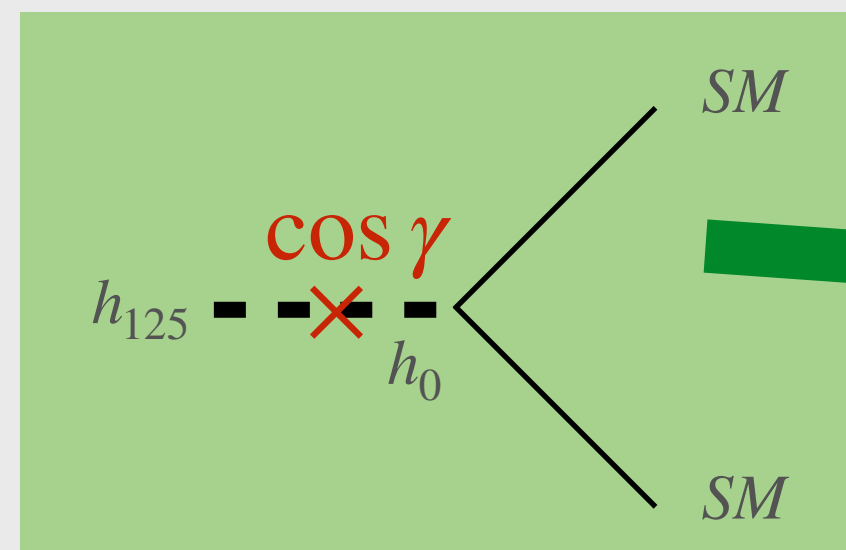
EXPLOIT ONCE MORE THE
W BOSON LUMINOSITY



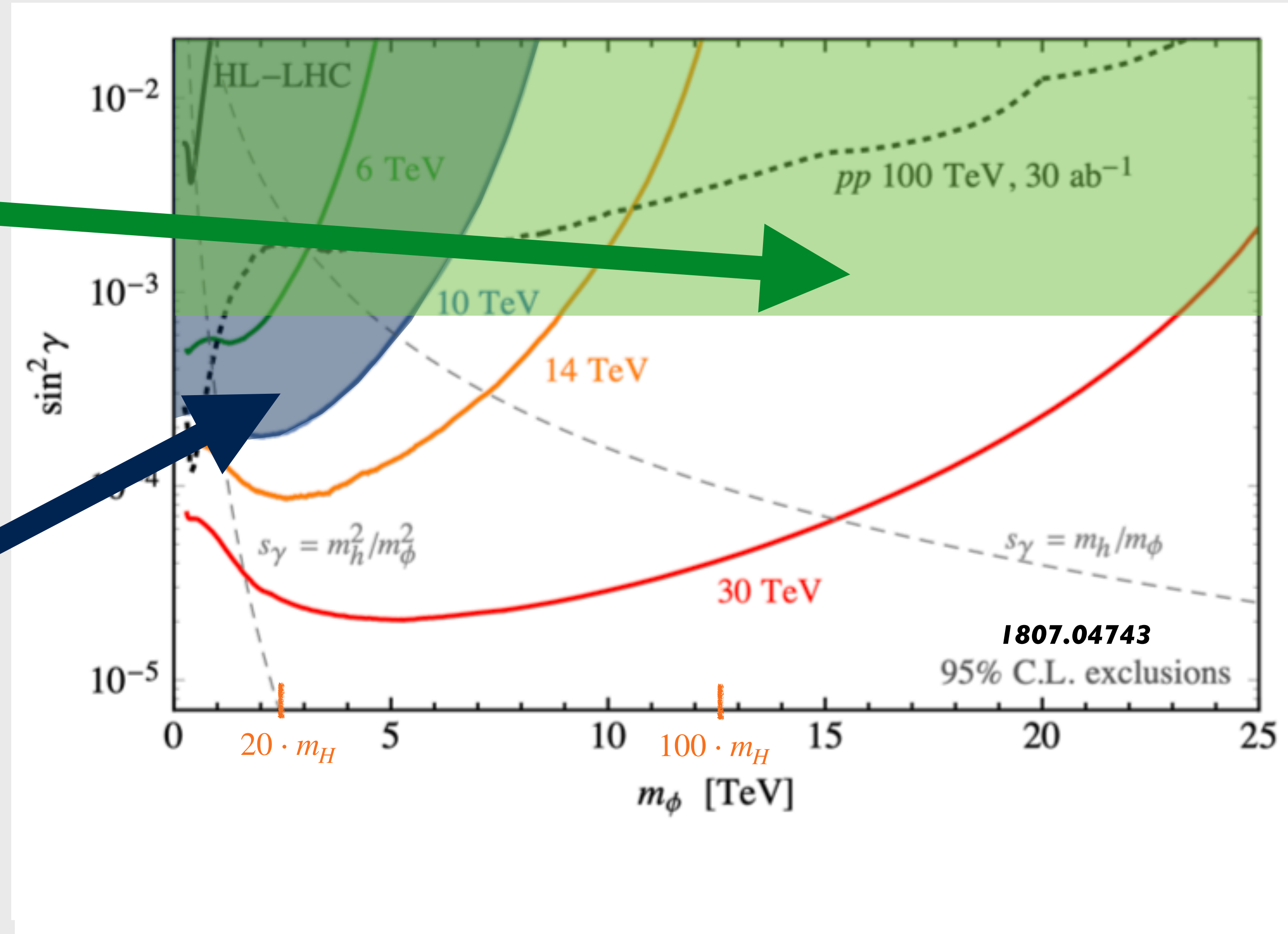
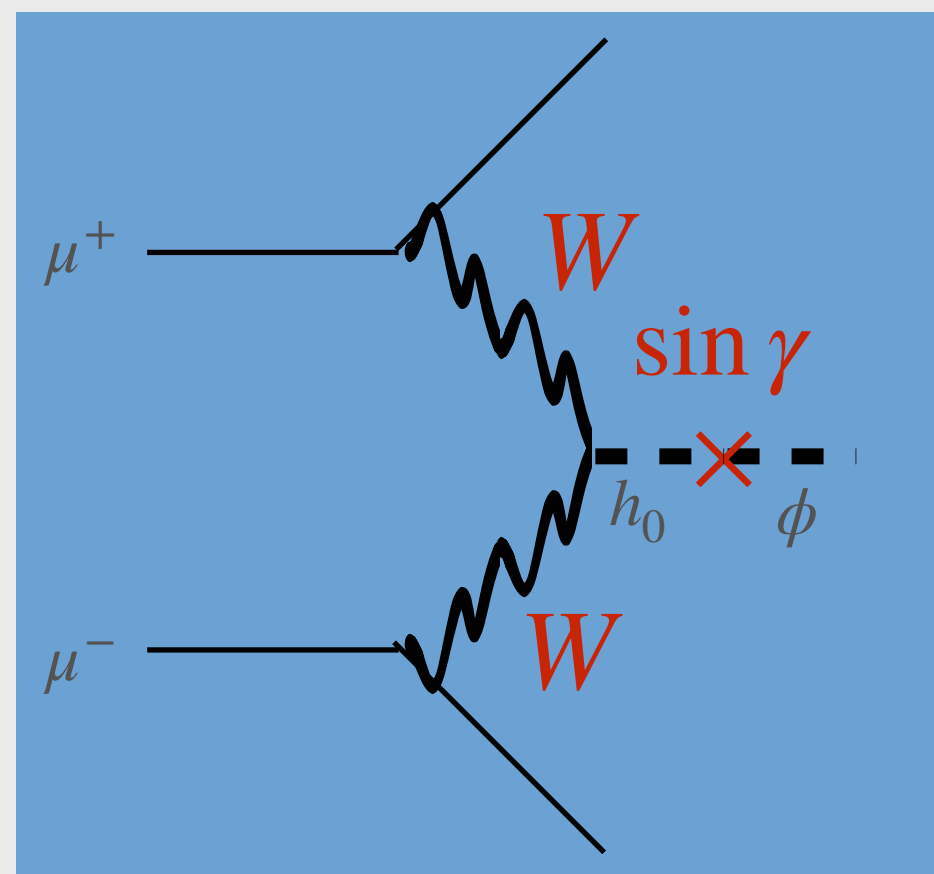
Higgs + Singlet

SINGELTS

BIG JUMP AT FUTURE COLLIDERS



EXPLOIT ONCE MORE THE
W BOSON LUMINOSITY



Higgs + Singlet

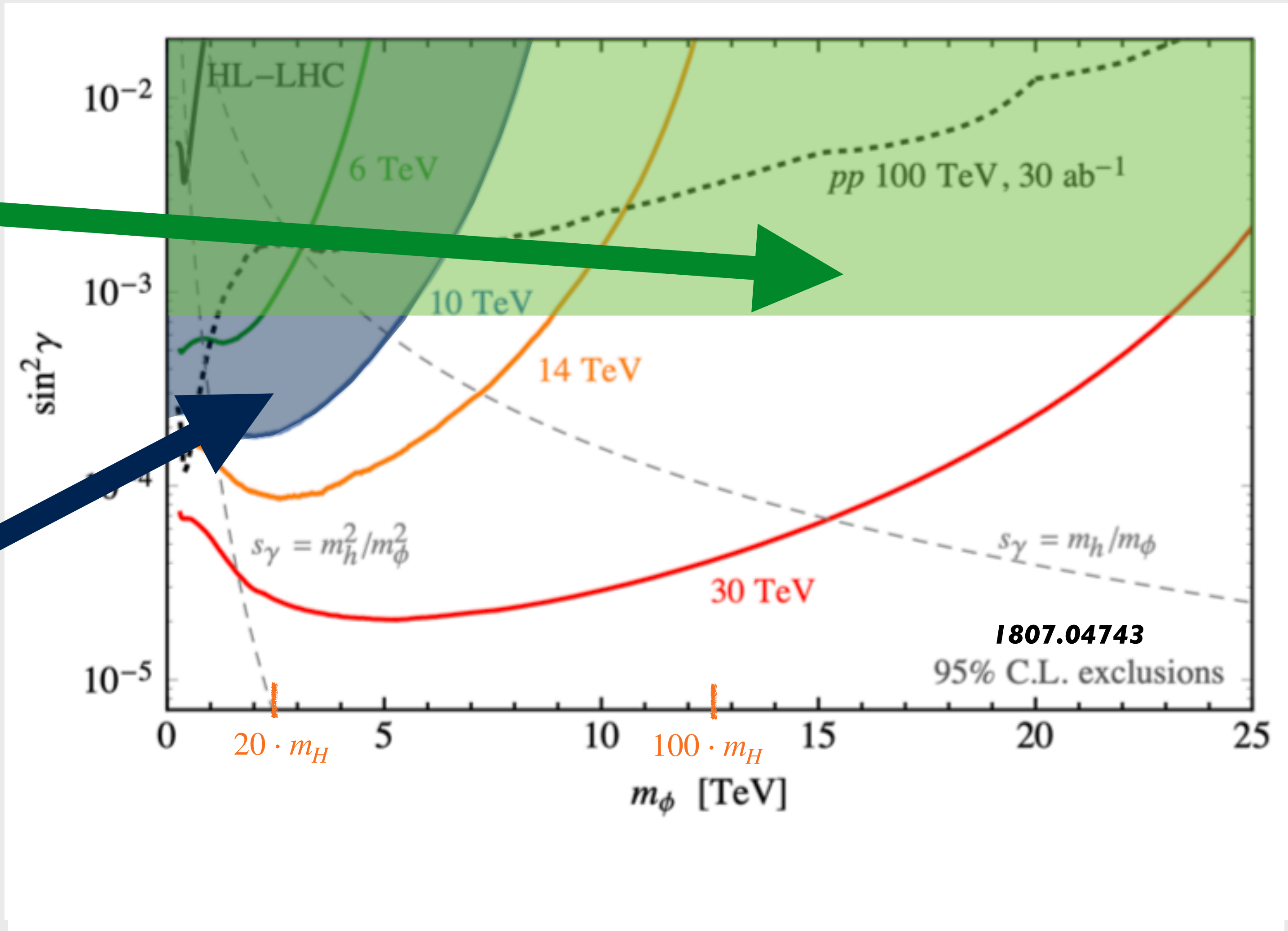
SINGELTS

BIG JUMP AT FUTURE COLLIDERS

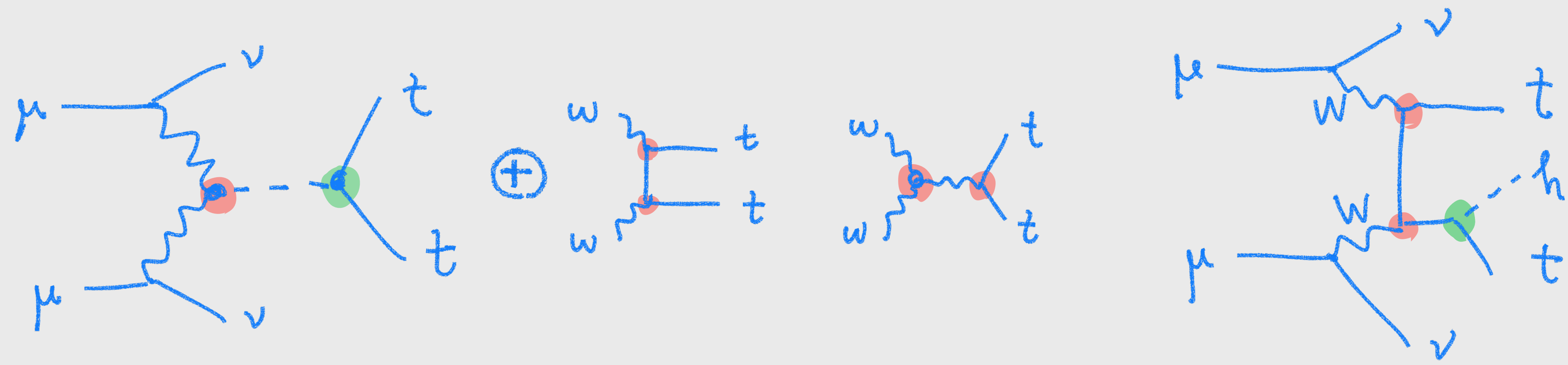
Probe new physics in two complementary ways at the same time



EX

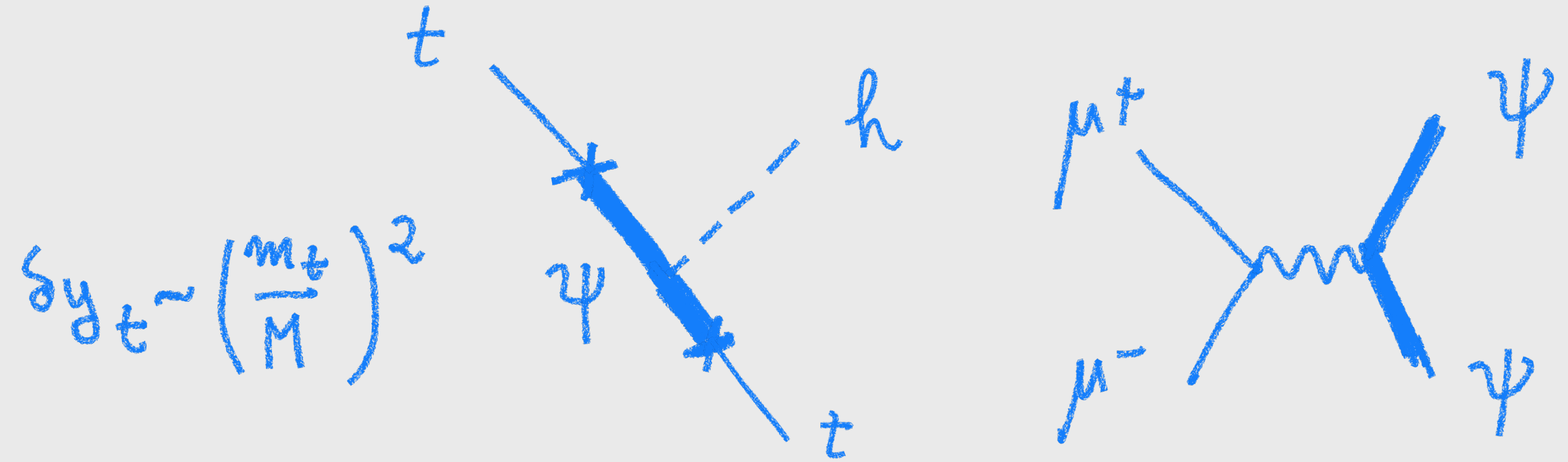


Joint-ventures on the Yukawa of the top quark



A global fit of “top \subset SM” to search for couplings deviations in $\mu\mu$ collider top factory”

The same muon collider that acts as “top factory” can also test directly the existence of new states responsible for the deviations in the couplings



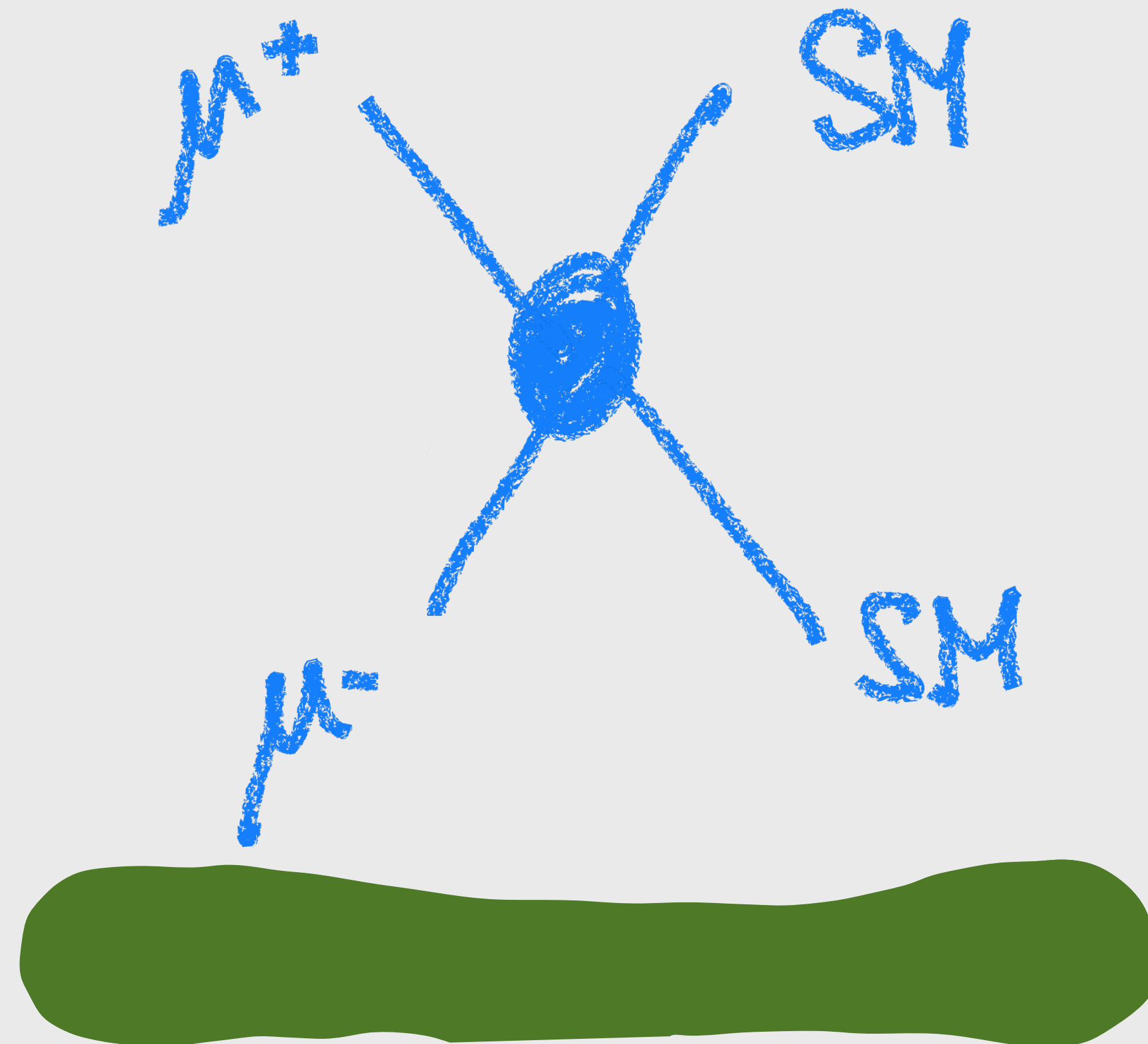
$$\delta y_t \sim \left(\frac{m_t}{M}\right)^2$$

TYPICALLY MUON COLLIDER EXCLUSION UP TO $S_{\mu\mu}/2$

same machine can perform complementary direct and indirect tests of BSM

at $\sqrt{s} \gg 100 \text{ GeV}$

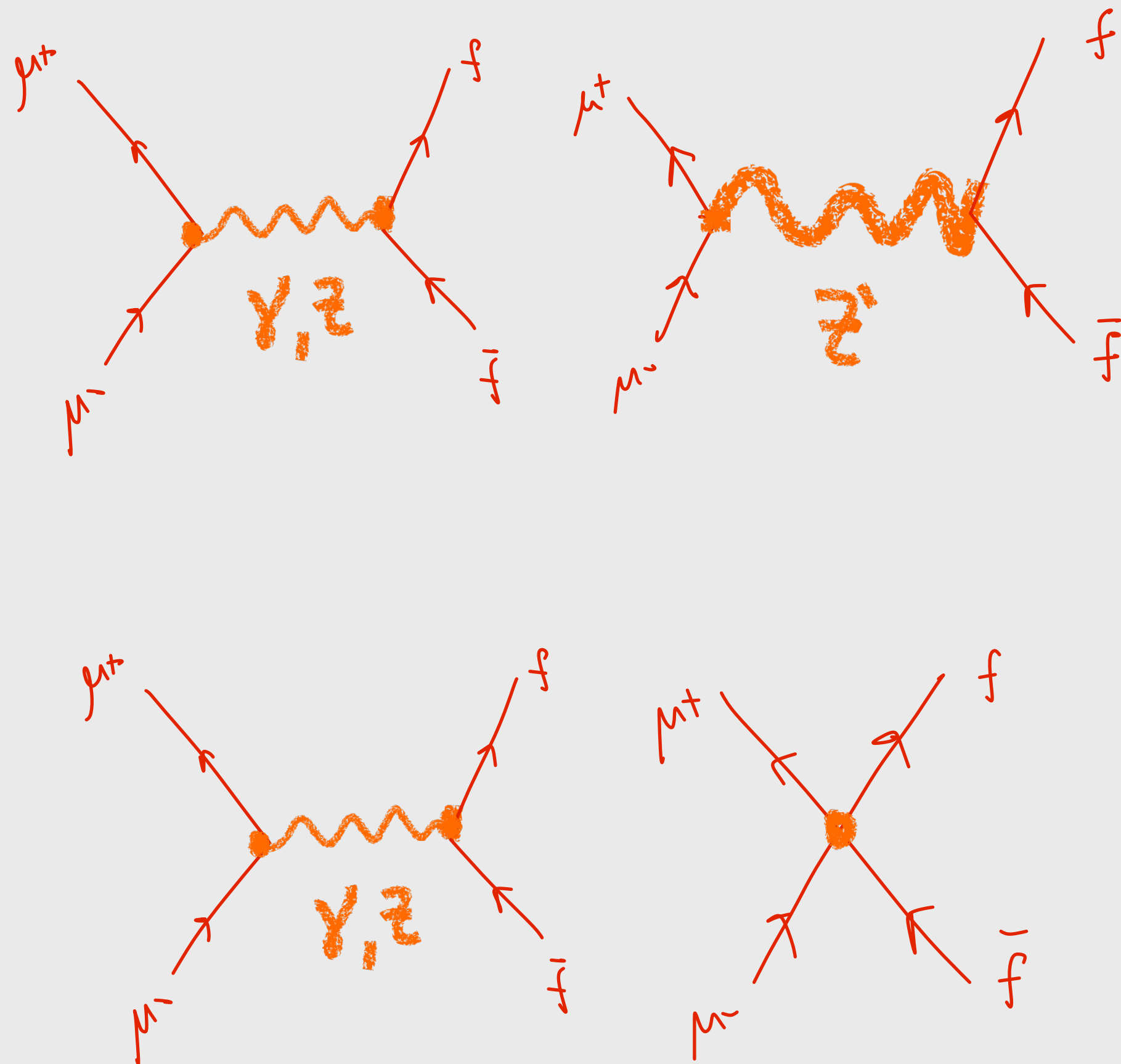
Indirect Effects



A heavy Z'

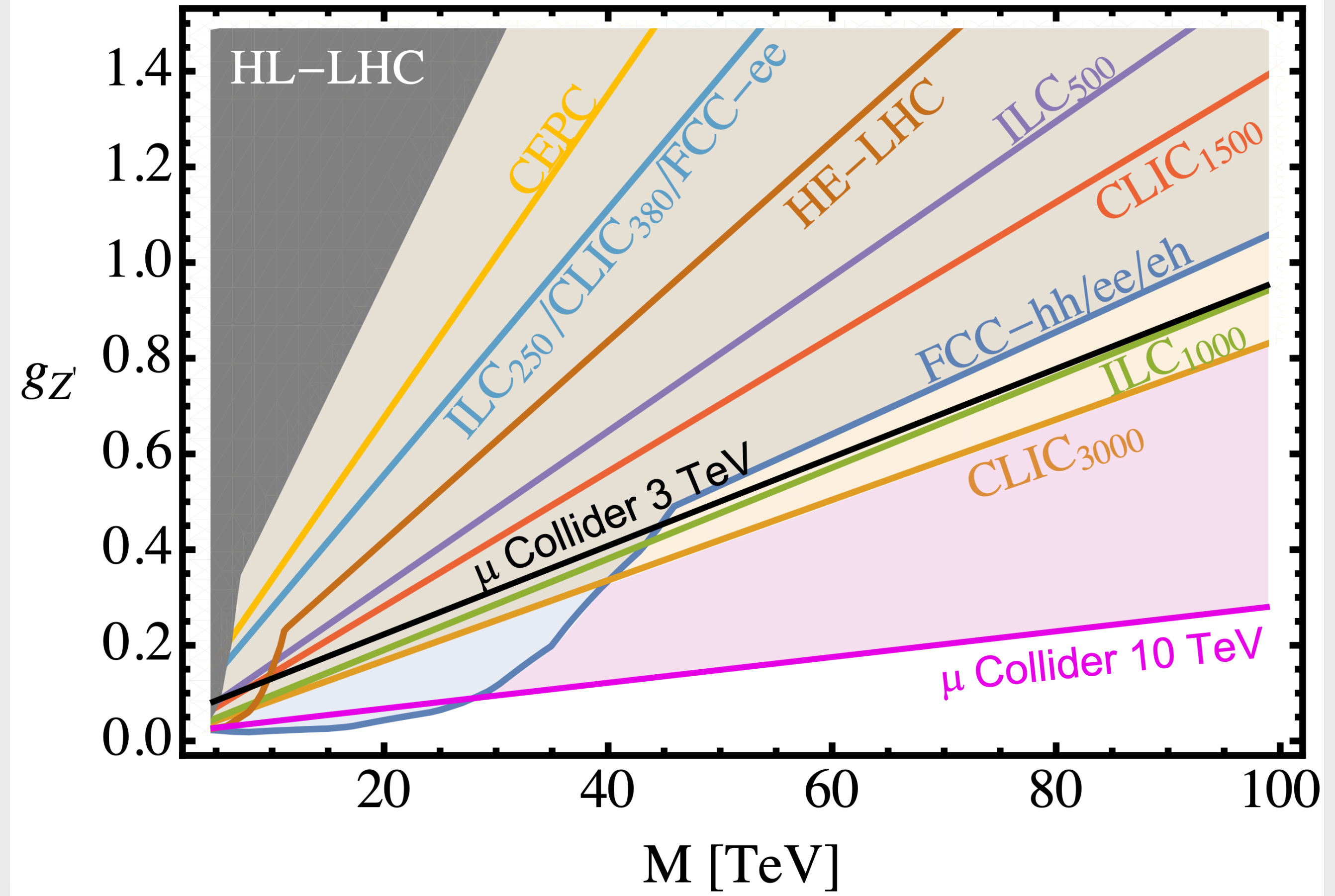
DRELL-YAN

RATES AND ANGULAR DISTRIBUTIONS



Y-Universal Z' , 2σ

2203.07256

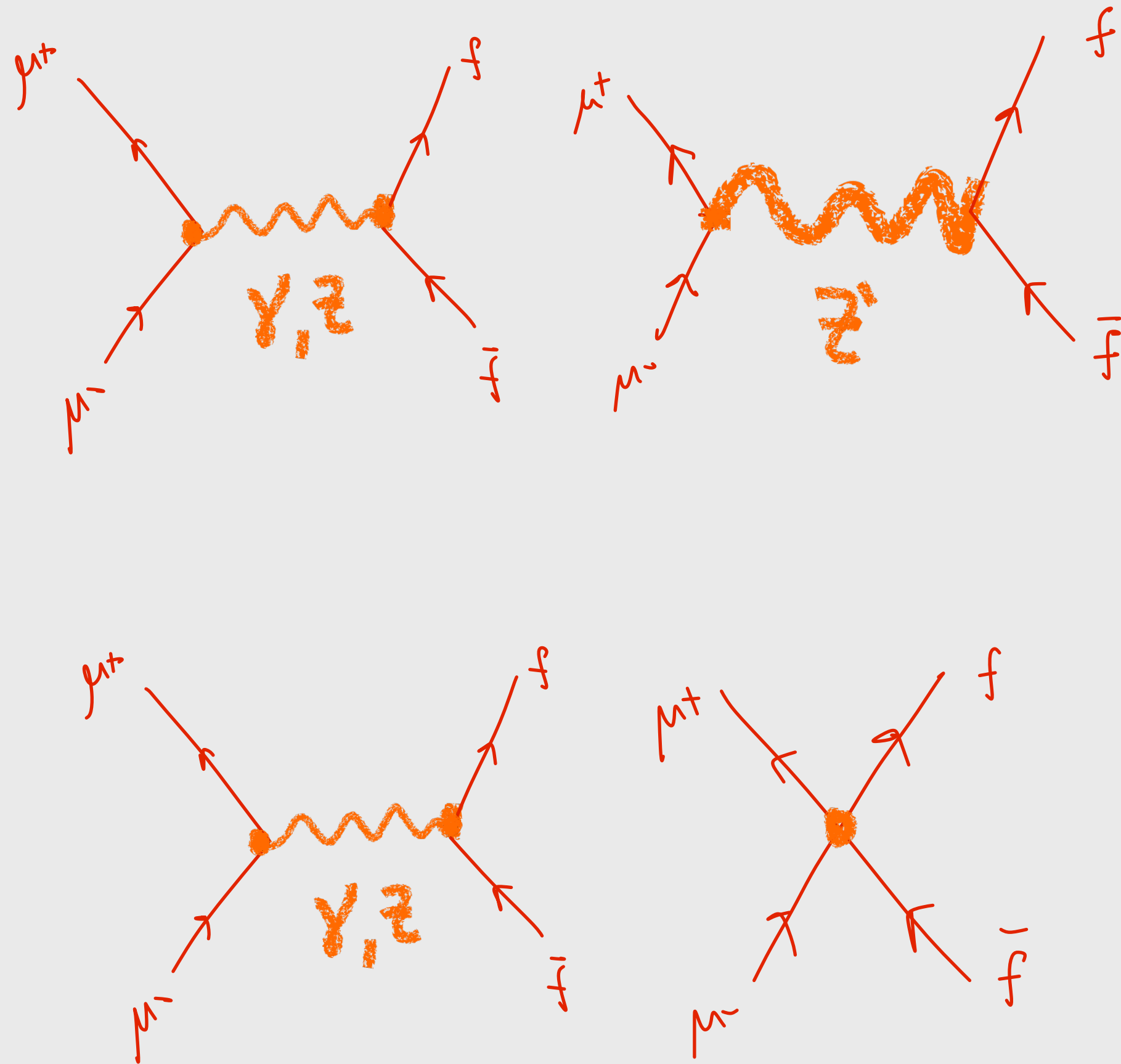


$\sqrt{s} \simeq 3 \text{ TeV}$ can probe 70+ TeV mass for $g_{Z'} \simeq g_{SM} \simeq 0.67$

A heavy Z'

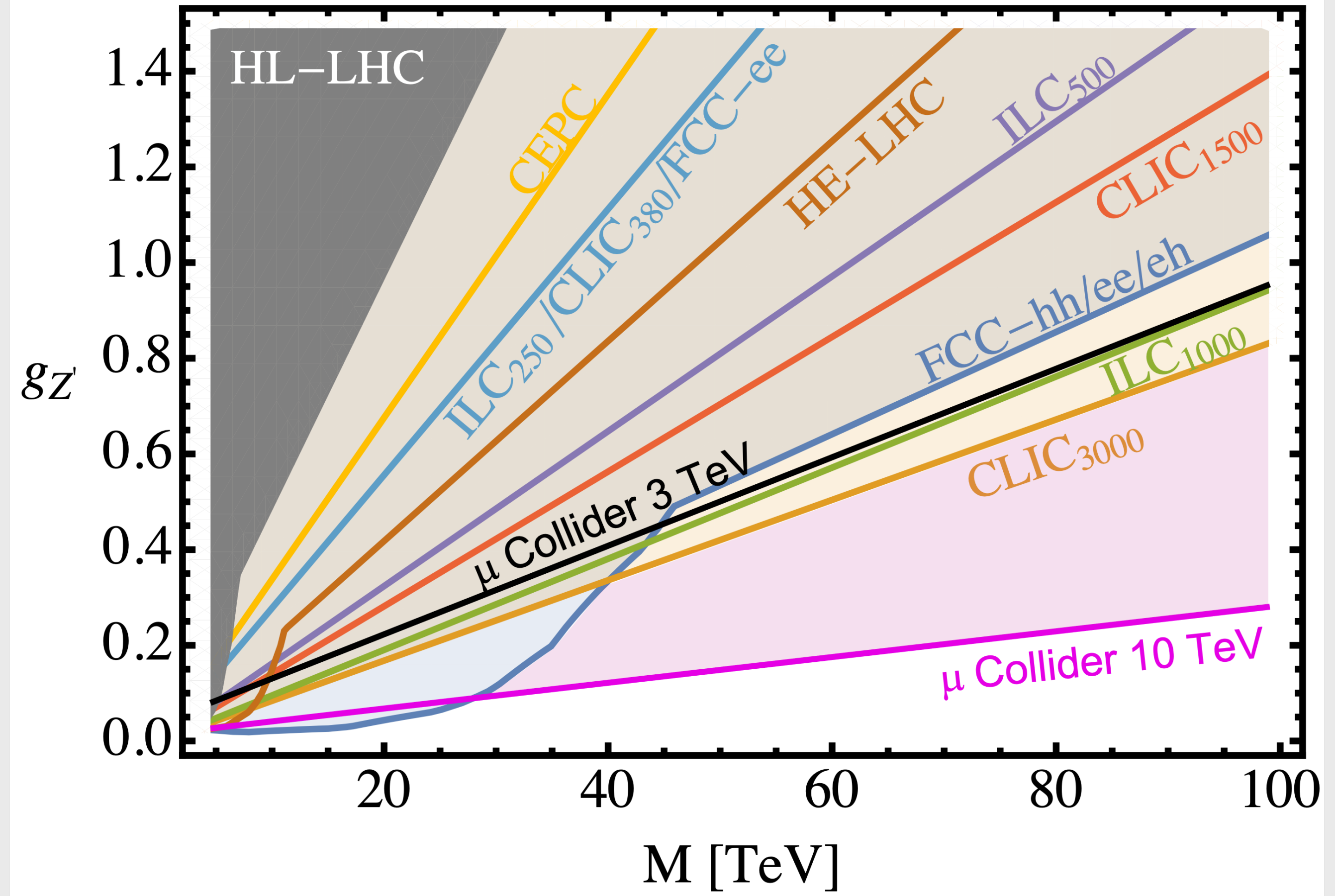
DRELL-YAN

RATES AND ANGULAR DISTRIBUTIONS



Y-Universal Z' , 2σ

2203.07256

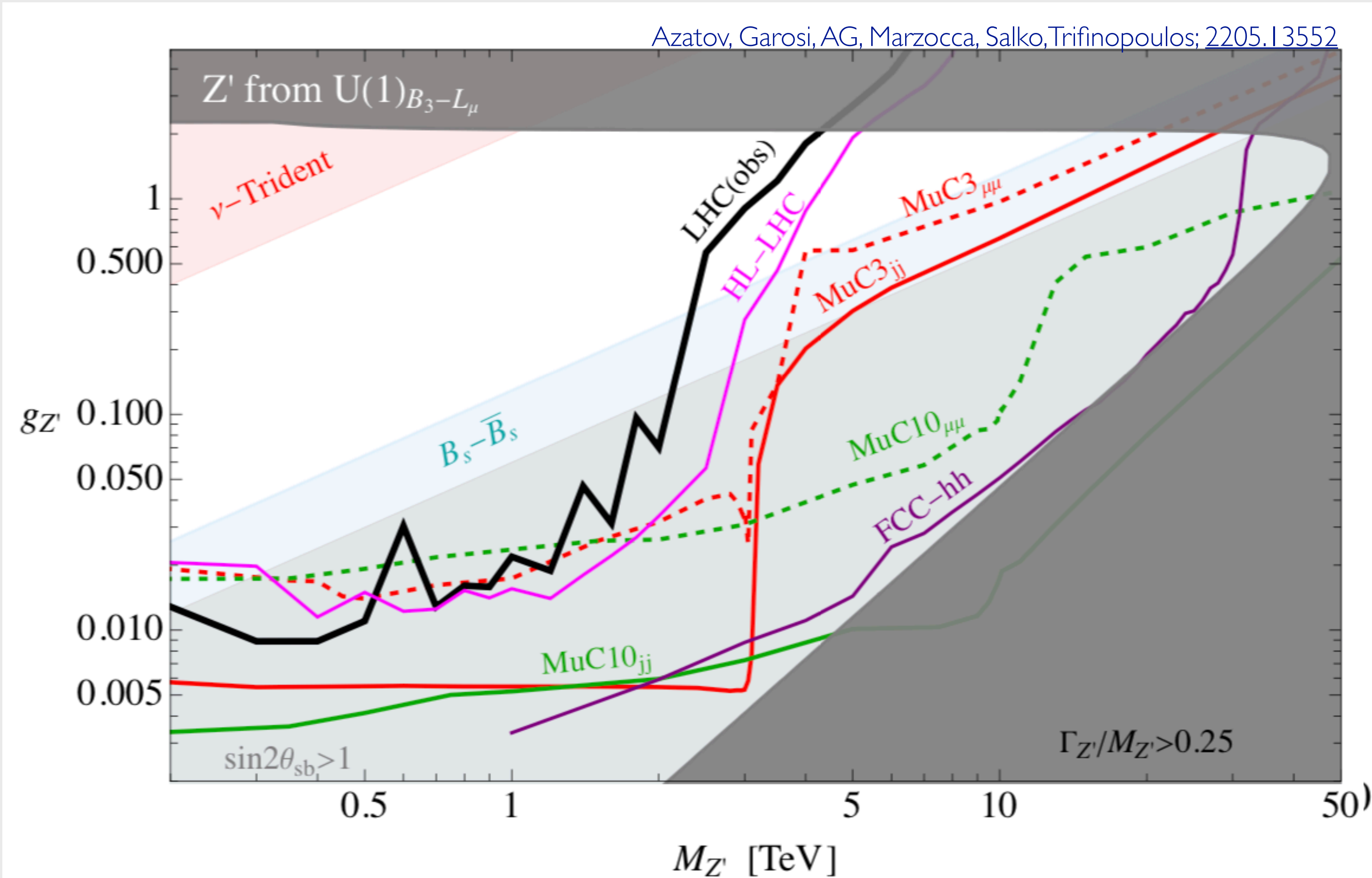
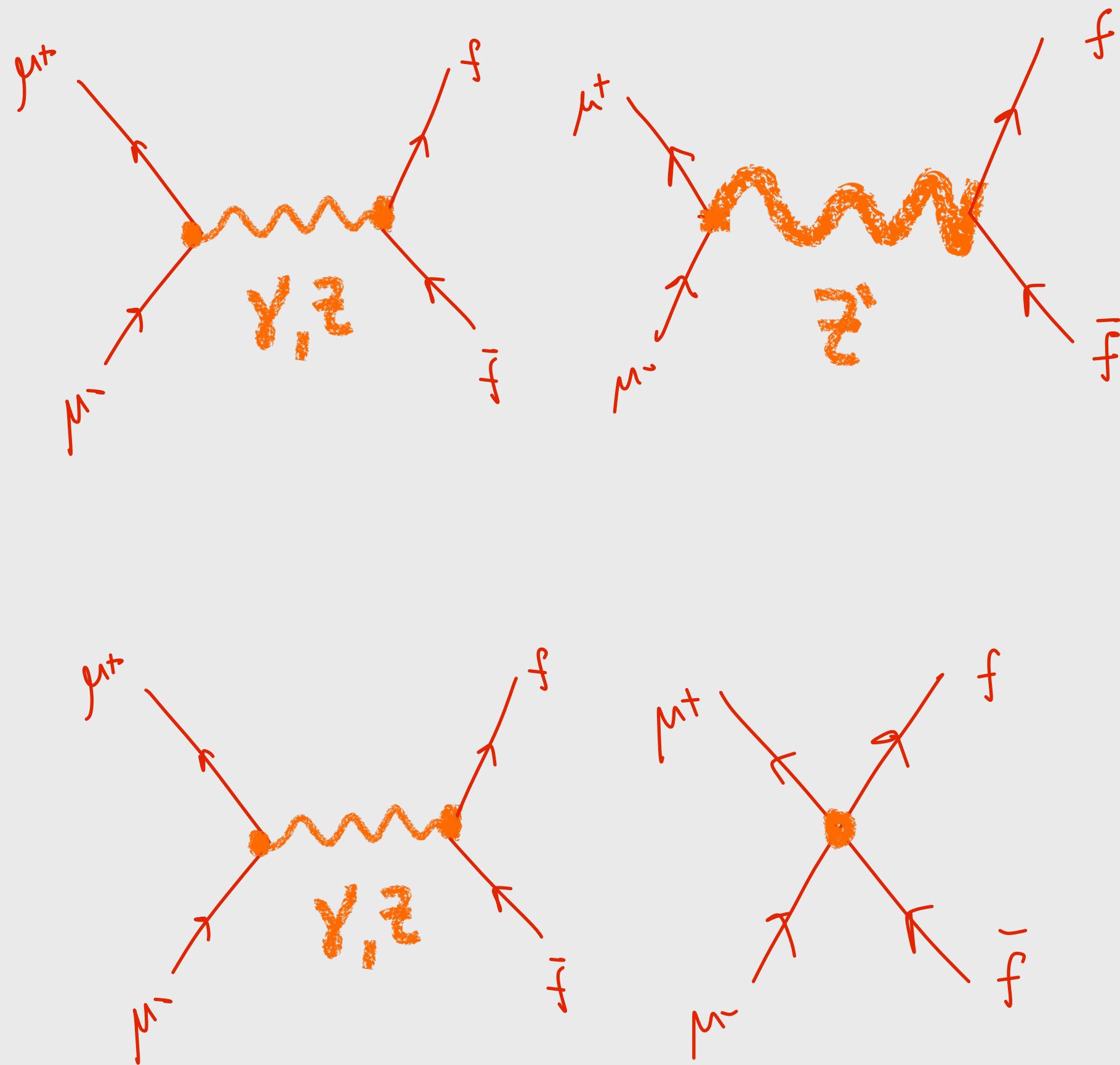


Another heavy Z'

DRELL-YAN

RATES AND ANGULAR DISTRIBUTIONS

$$bs\mu\mu : \epsilon_{sb} = -1.7 \times 10^{-3} \left(\frac{M_{Z'}}{g_{Z'} \text{TeV}} \right)^2$$



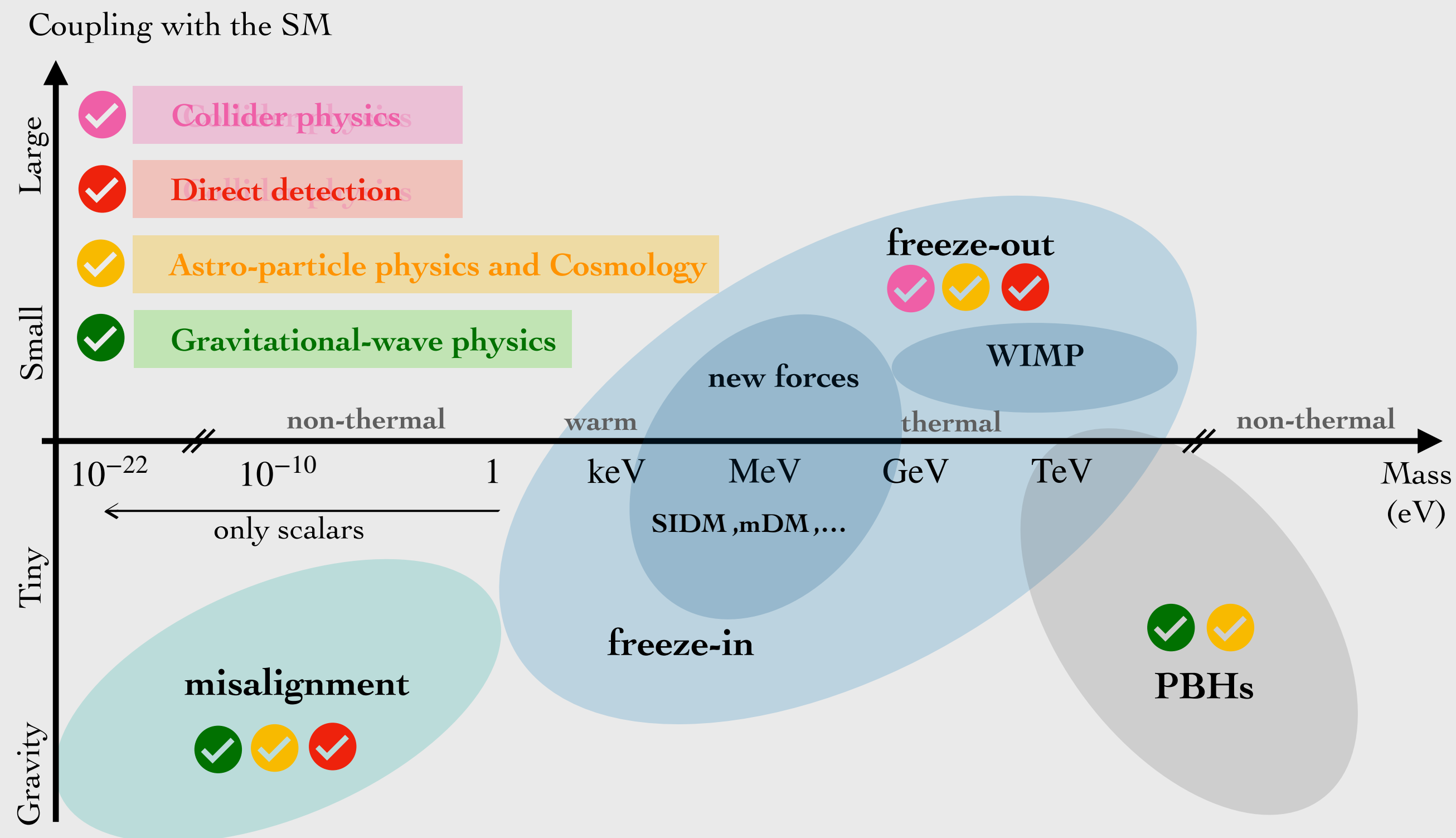
flashing concrete results for

Dark Matter at the weak scale



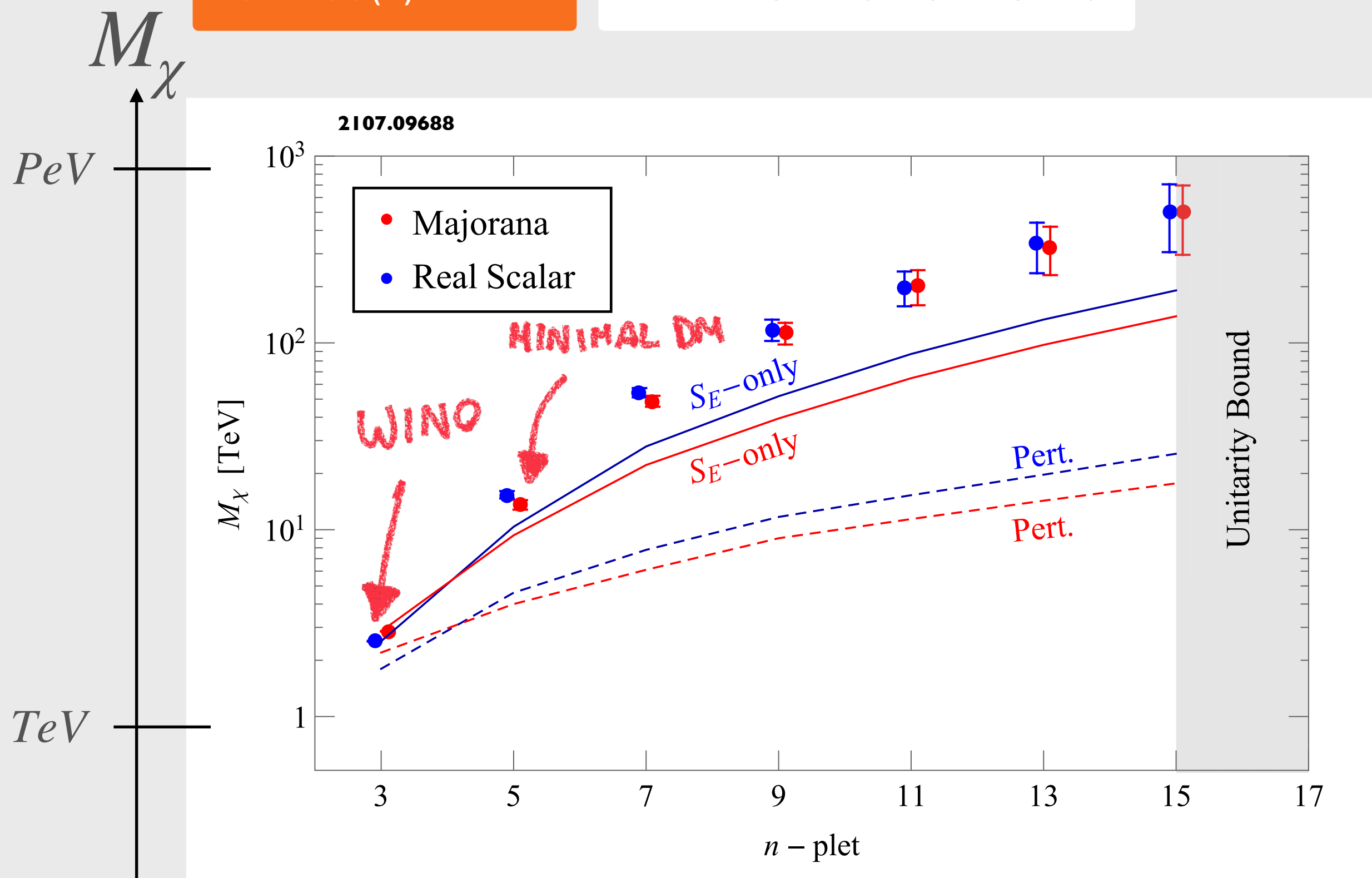
Electroweak Dark Matter: LSP (+NLSP)

- The chessboard of DM is very large!



- High energy colliders are excellent and very robust probes of WIMPs!

WIMP Dark Matter as $SU(2)$ n – plet

PURE $SU(2)$ N-PLETINTERPOLATOR UP TO PeV 

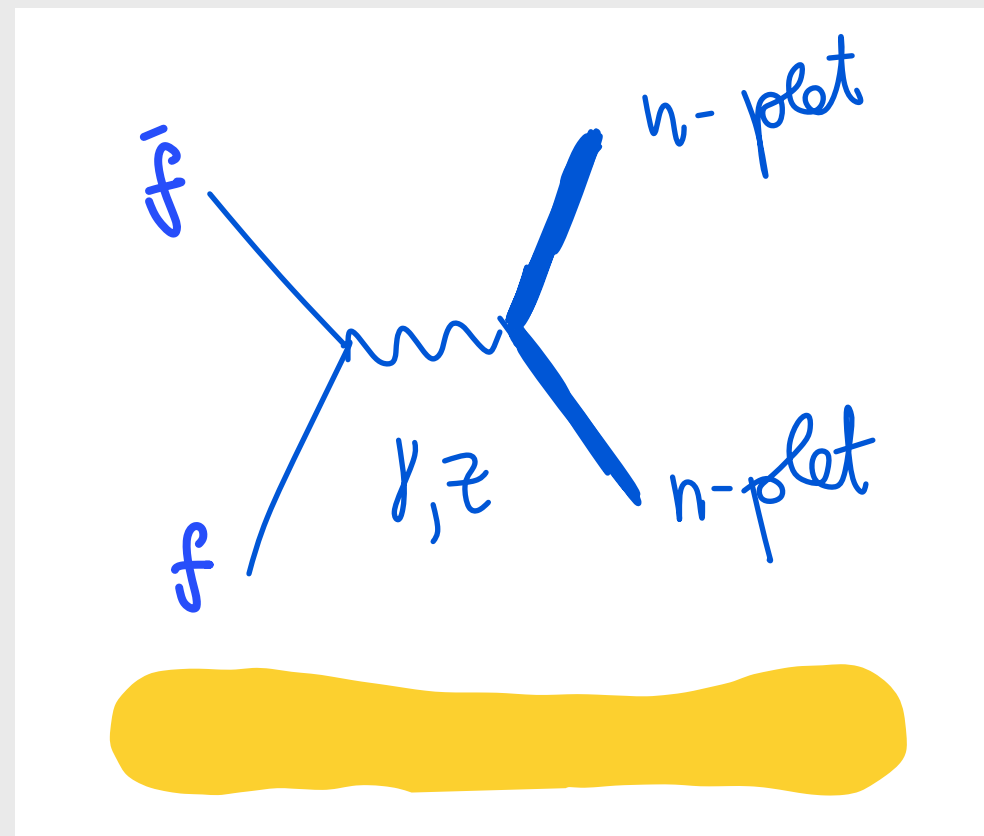
DM spin	EW n-plet	M_χ (TeV)	$(\sigma v)_{\text{tot}}^{J=0} / (\sigma v)_{\text{max}}^{J=0}$	$\Lambda_{\text{Landau}} / M_{\text{DM}}$	$\Lambda_{\text{UV}} / M_{\text{DM}}$
Real scalar	3	2.53 ± 0.01	–	3×10^{37}	$4 \times 10^{24*}$
	5	15.4 ± 0.7	0.002	5×10^{36}	2×10^{24}
	7	54.2 ± 3.1	0.022	2×10^{19}	2×10^{24}
	9	117.8 ± 15.4	0.088	3×10^3	2×10^{24}
	11	199 ± 42	0.25	20	3×10^{24}
	13	338 ± 102	0.6	3.5	3×10^{24}
Majorana fermion	3	2.86 ± 0.01	–	3×10^{37}	$8 \times 10^{12*}$
	5	13.6 ± 0.8	0.003	3×10^{17}	5×10^{12}
	7	48.8 ± 3.3	0.019	1×10^4	4×10^7
	9	113 ± 15	0.07	30	3×10^7
	11	202 ± 43	0.2	6	3×10^7
	13	324.6 ± 94	0.5	2.6	3×10^7

a “collection” of Dark Matter candidates with thermal mass from 1 TeV to fraction of PeV

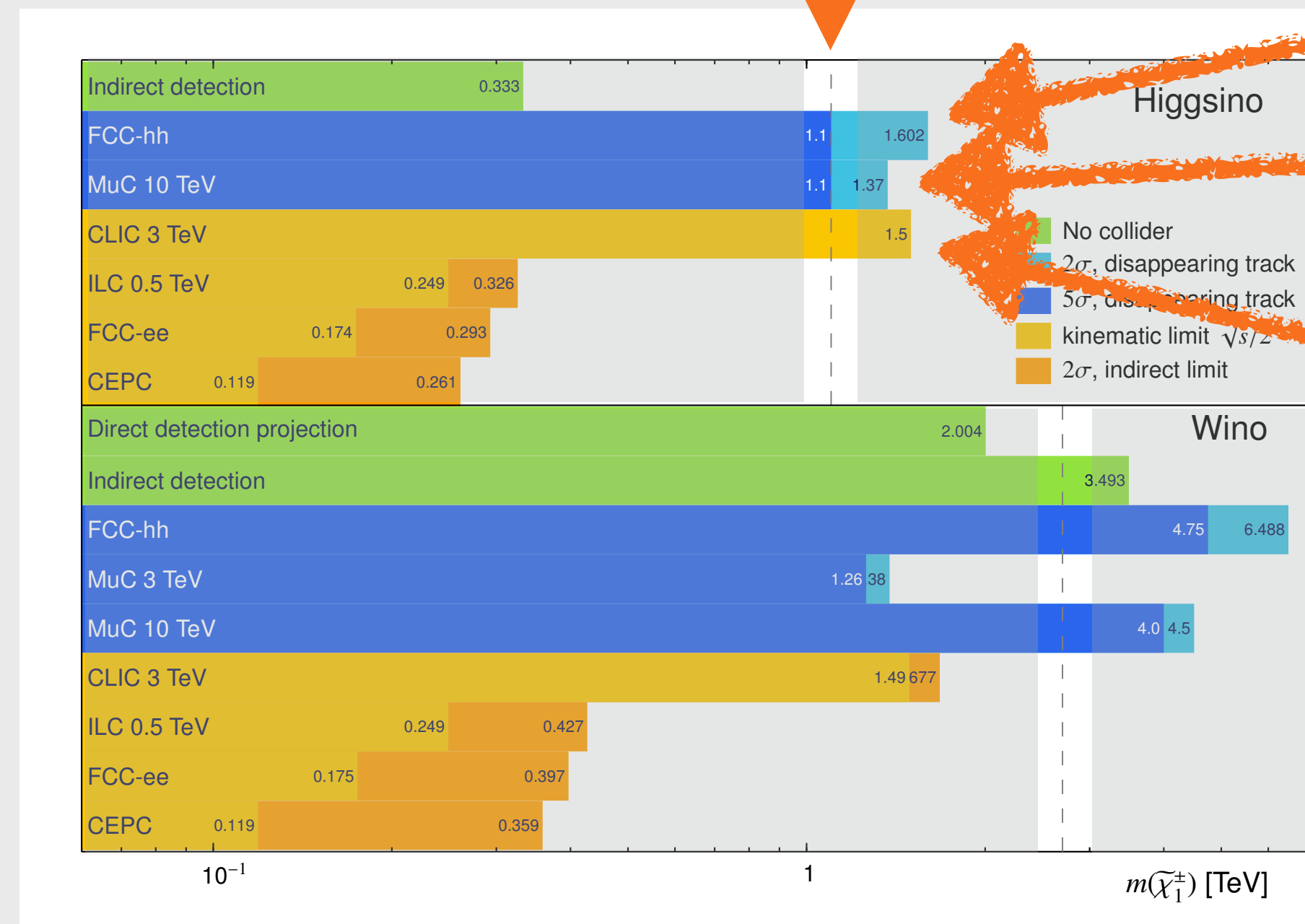
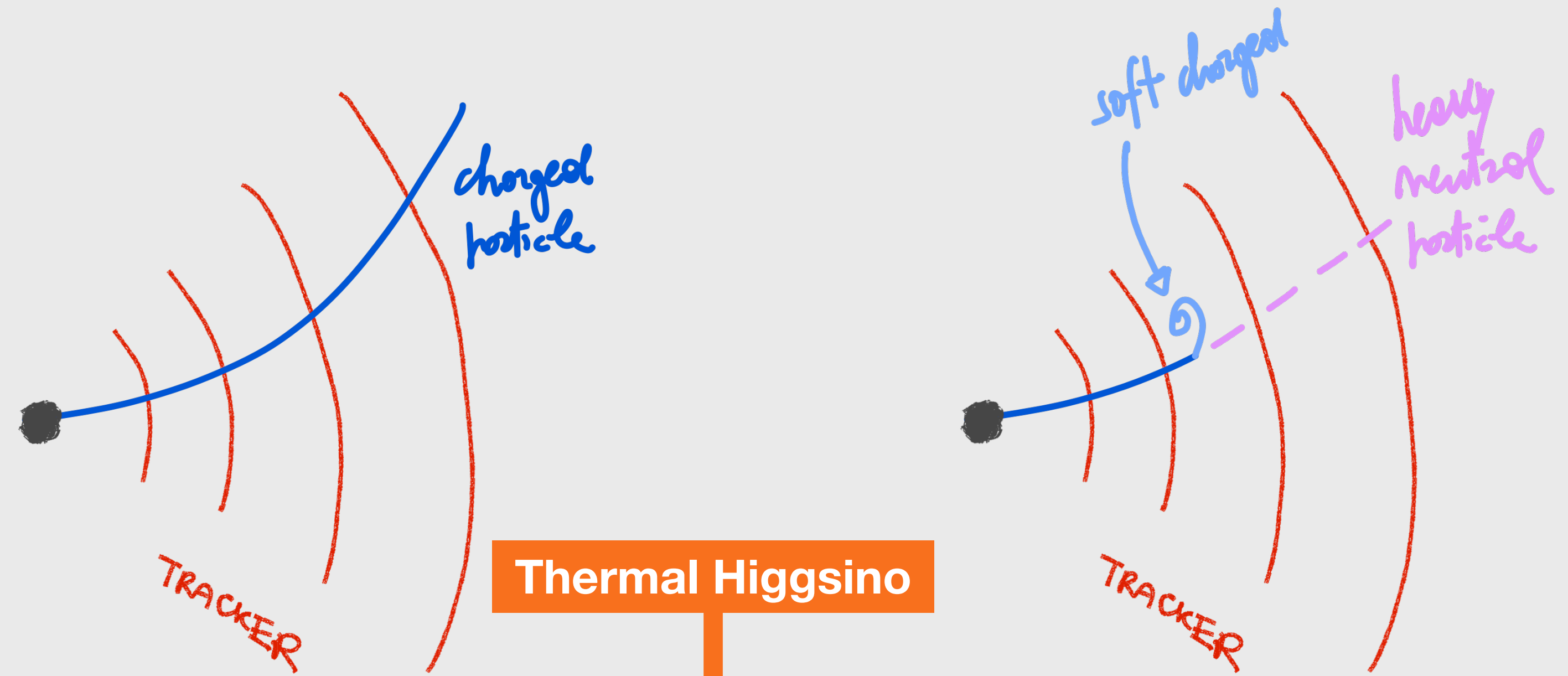
Higgsino DM

STUB-TRACKS EXOTIC SIGNAL

- Heavy n-plet of SU(2)
- Mass splitting $\sim \alpha_w m_W \sim 0.1 \text{ GeV} - \text{GeV}$



LARGE RATES, BUT NEEDS TO LIGHT UP THE DETECTOR IN A DISCERNIBLE WAY

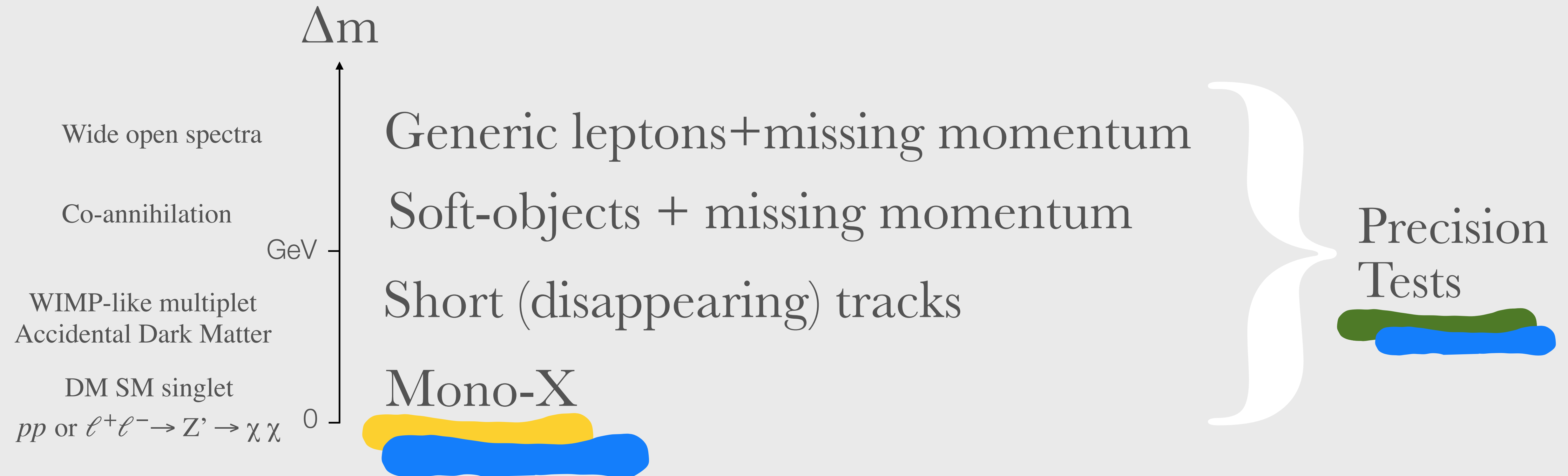


pp 100 TeV 30 ab^{-1}

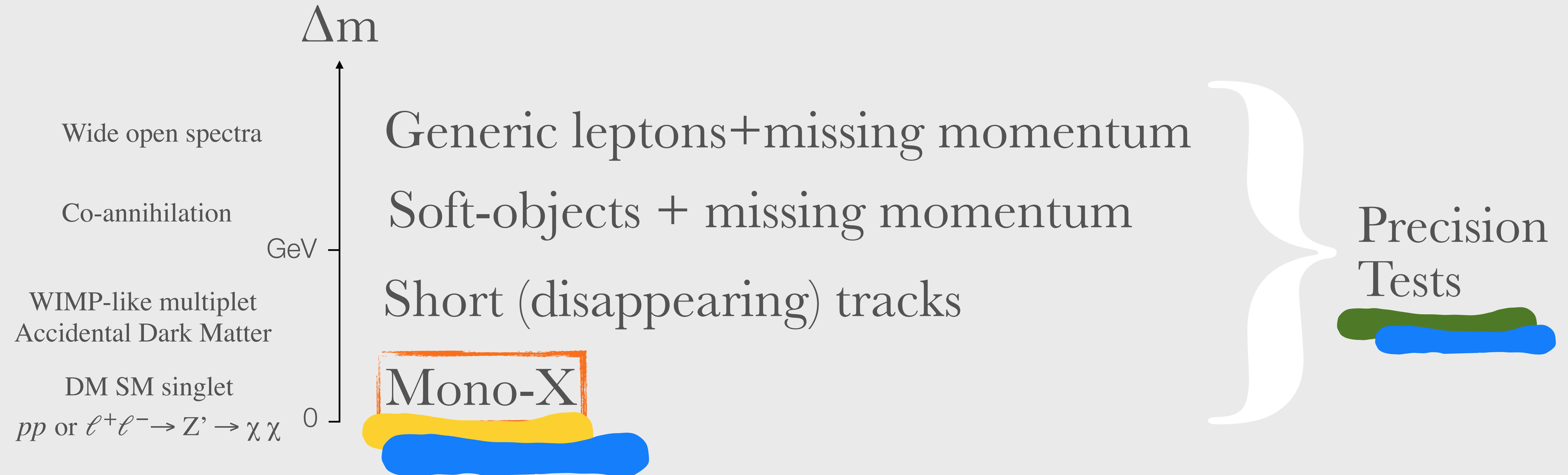
$\mu^+\mu^-$ 10 TeV 10 ab^{-1}

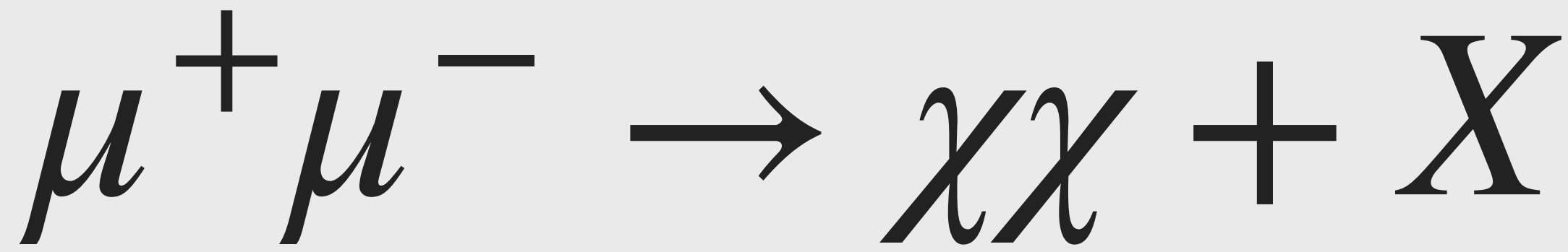
e^+e^- 3 TeV 5 ab^{-1}

Electroweak Dark Matter: LSP (+NLSP)



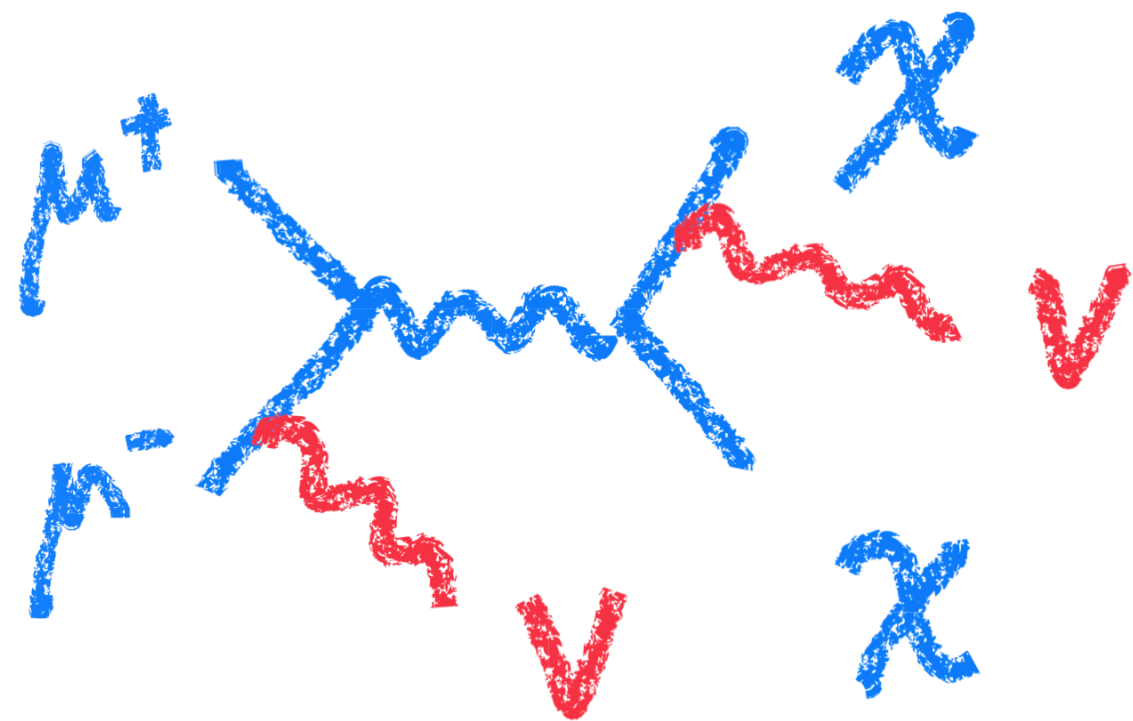
Electroweak Dark Matter: LSP (+NLSP)





MONO-X

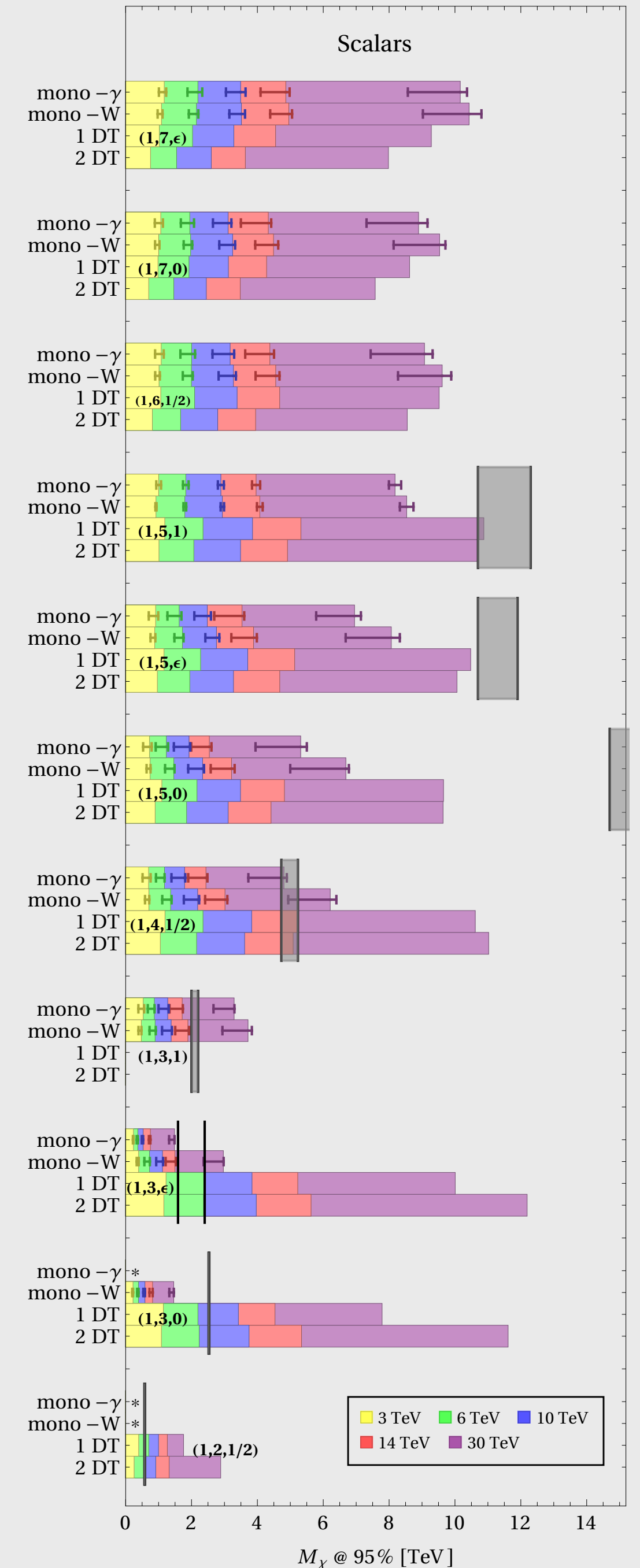
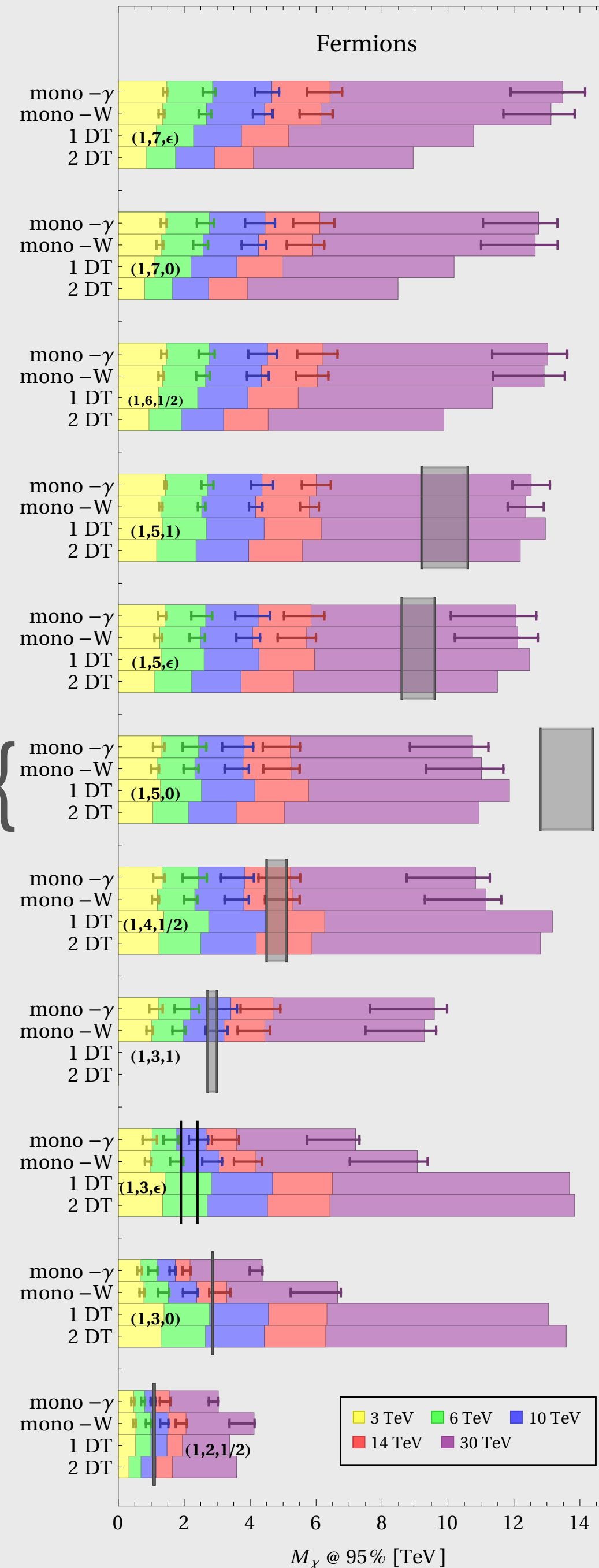
X = γ, W, \dots

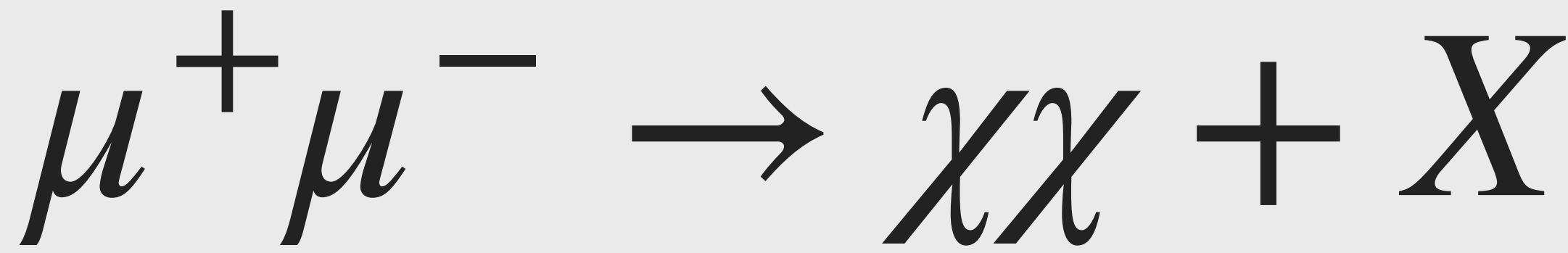


Large χ mass requires CoM energy!

Weak radiation yield the most constraining channel "mono-W"

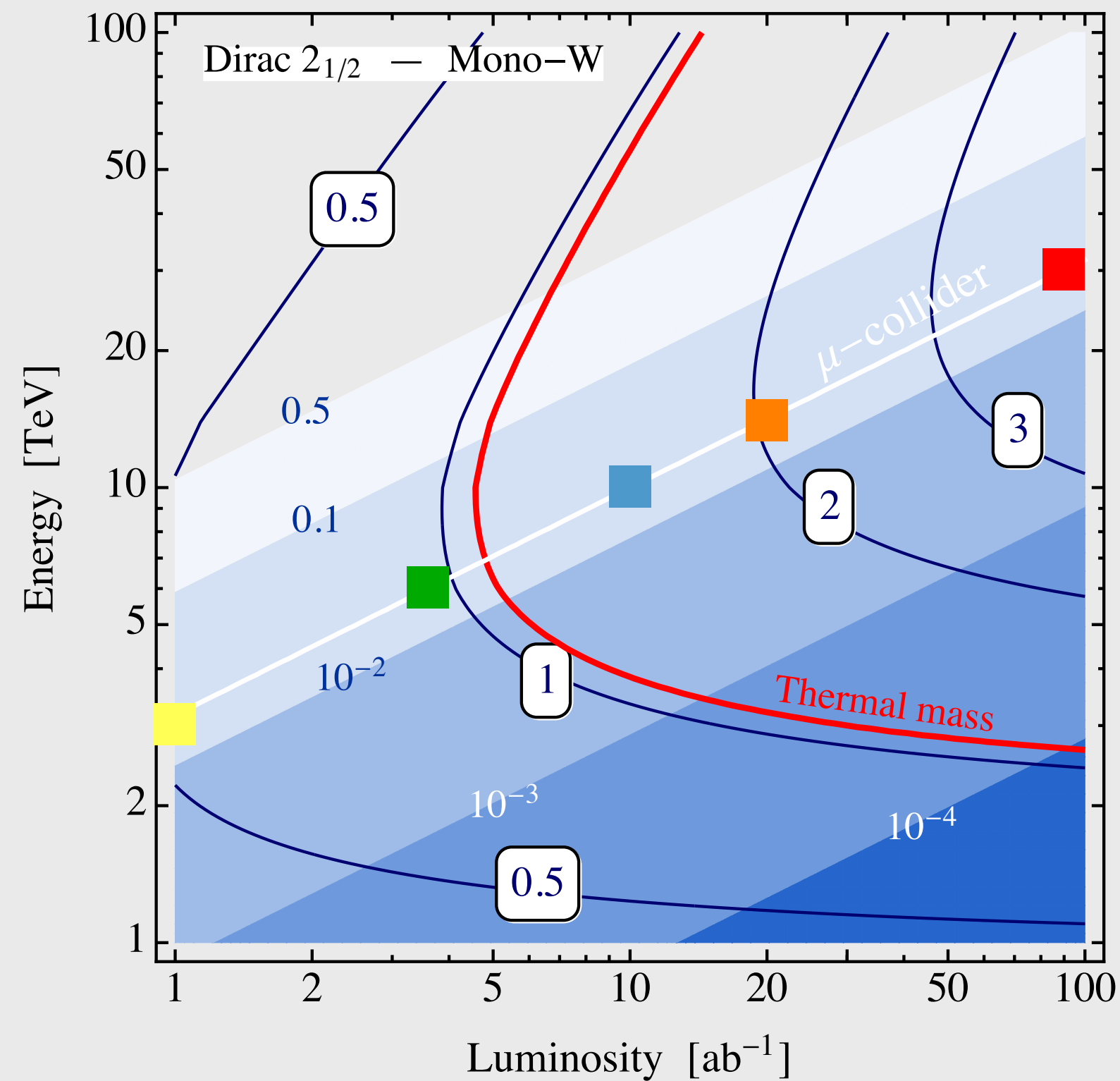
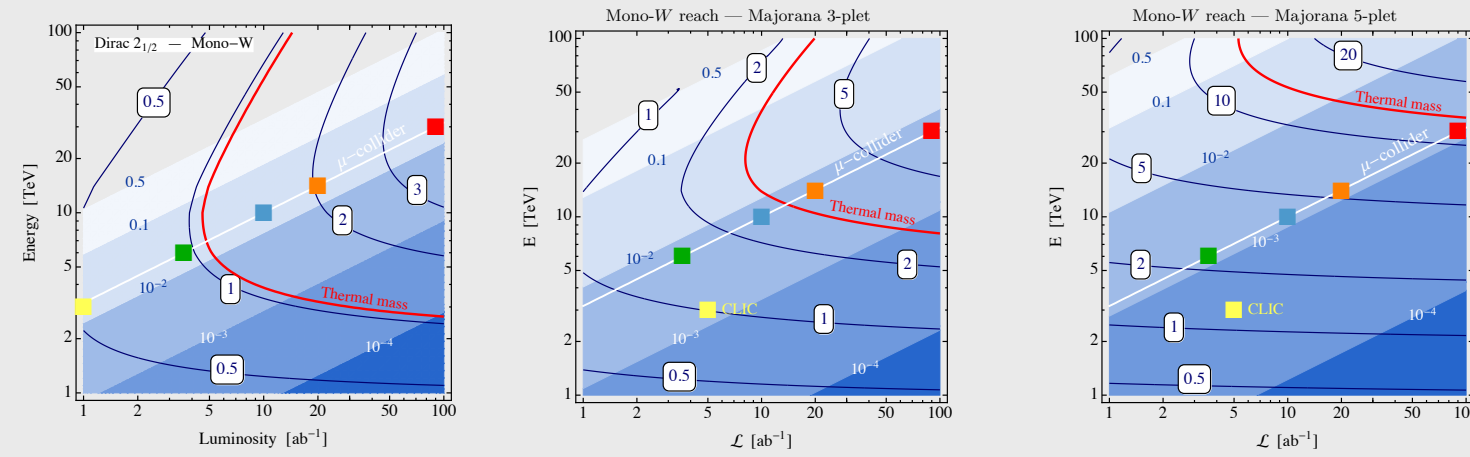
mono- γ
mono-W
tracklets





MONO-X

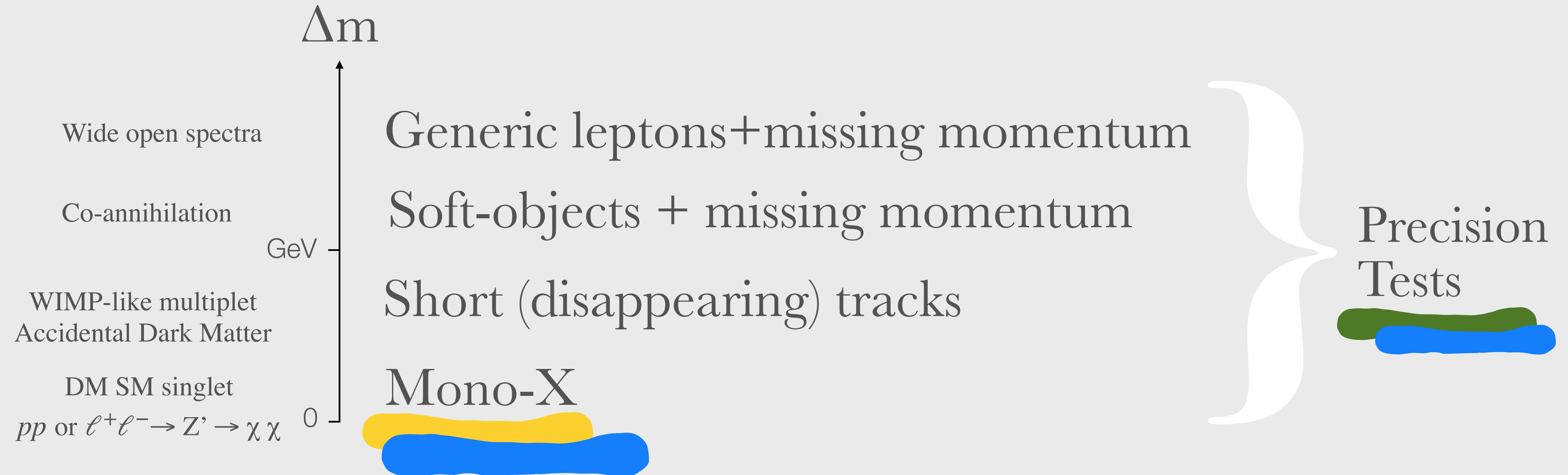
$X = \gamma, W, \dots$



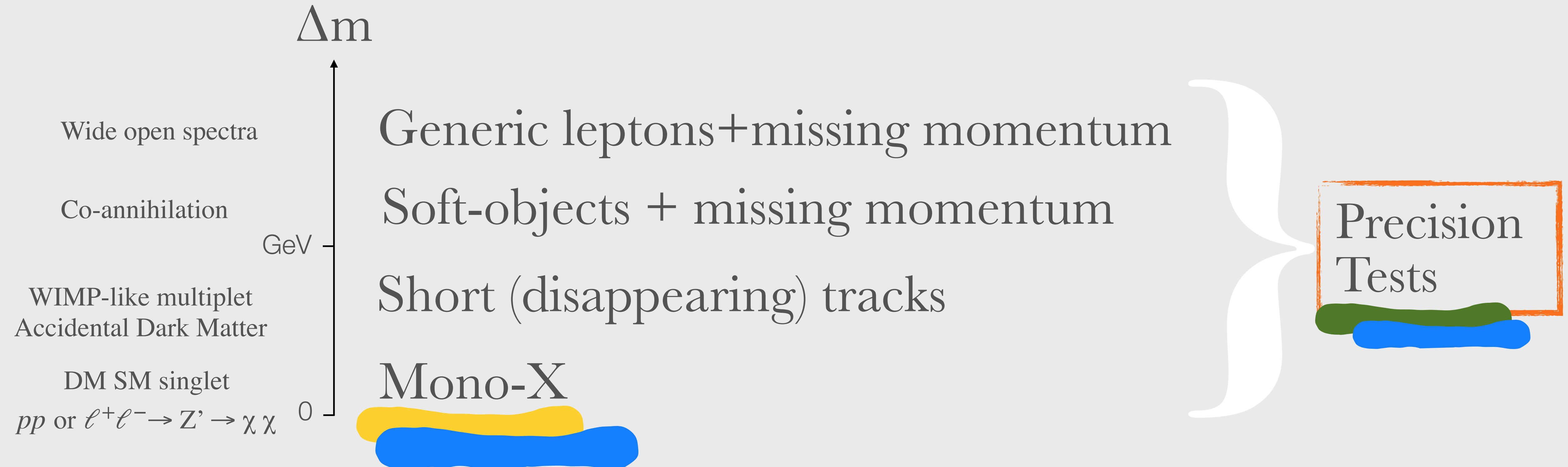
mono- γ
mono-W
tracklets



Electroweak Dark Matter: LSP (+NLSP)



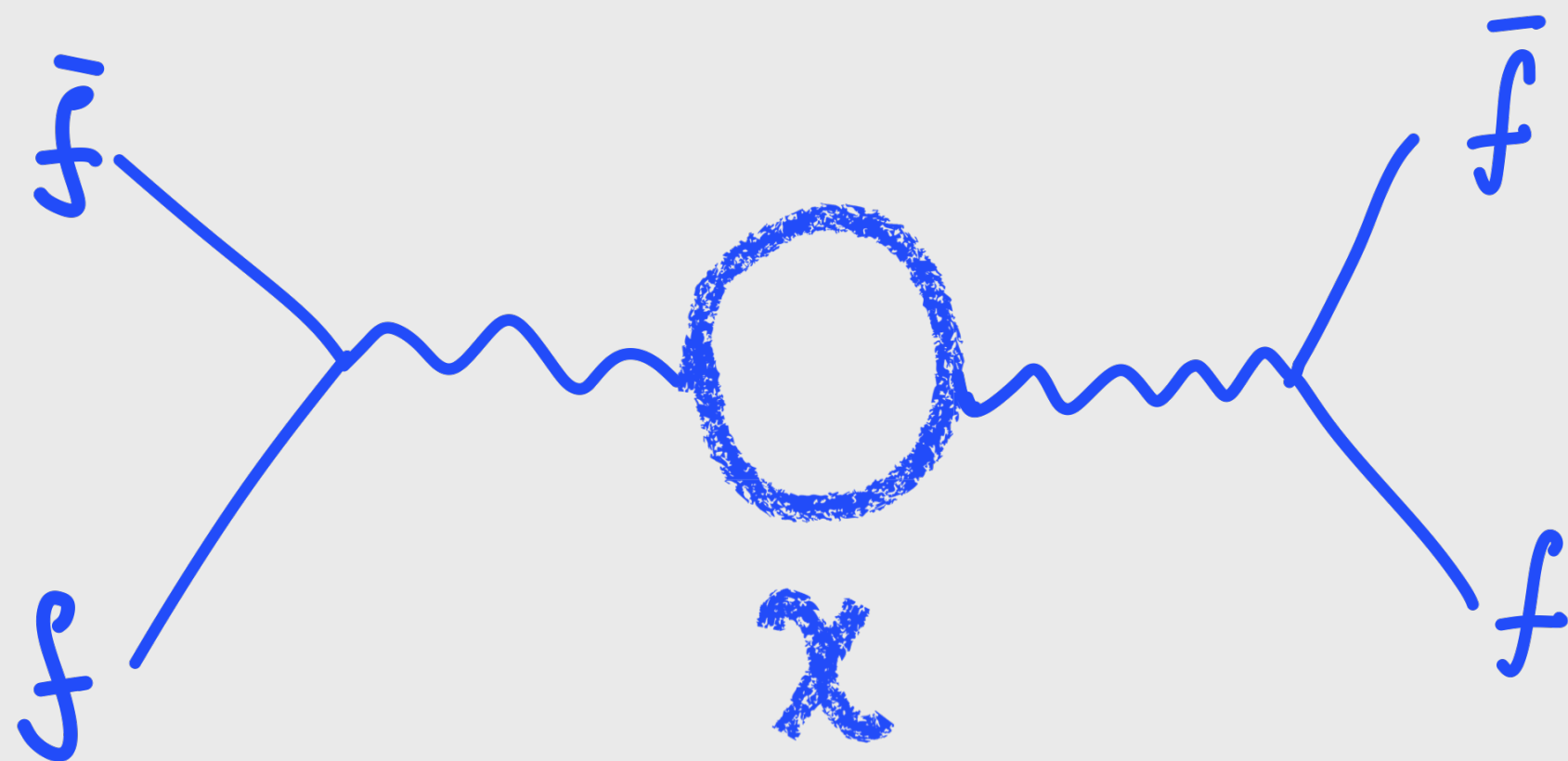
Electroweak Dark Matter: LSP (+NLSP)



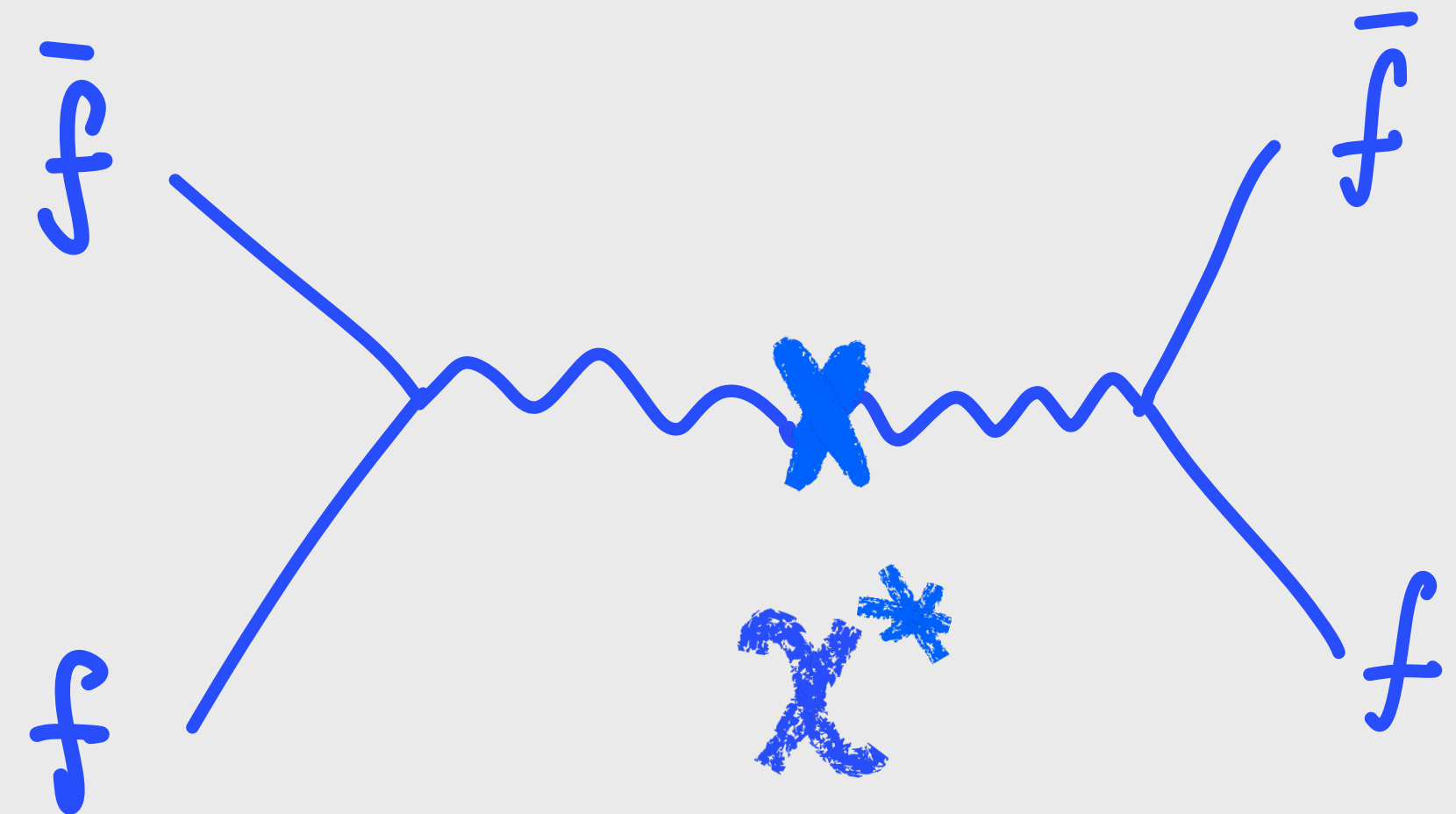
$$\mu^+ \mu^- \rightarrow f\bar{f}, W^+W^-, Zh, f\bar{f}'$$

PRECISION

TOTAL CROSS-SECTION

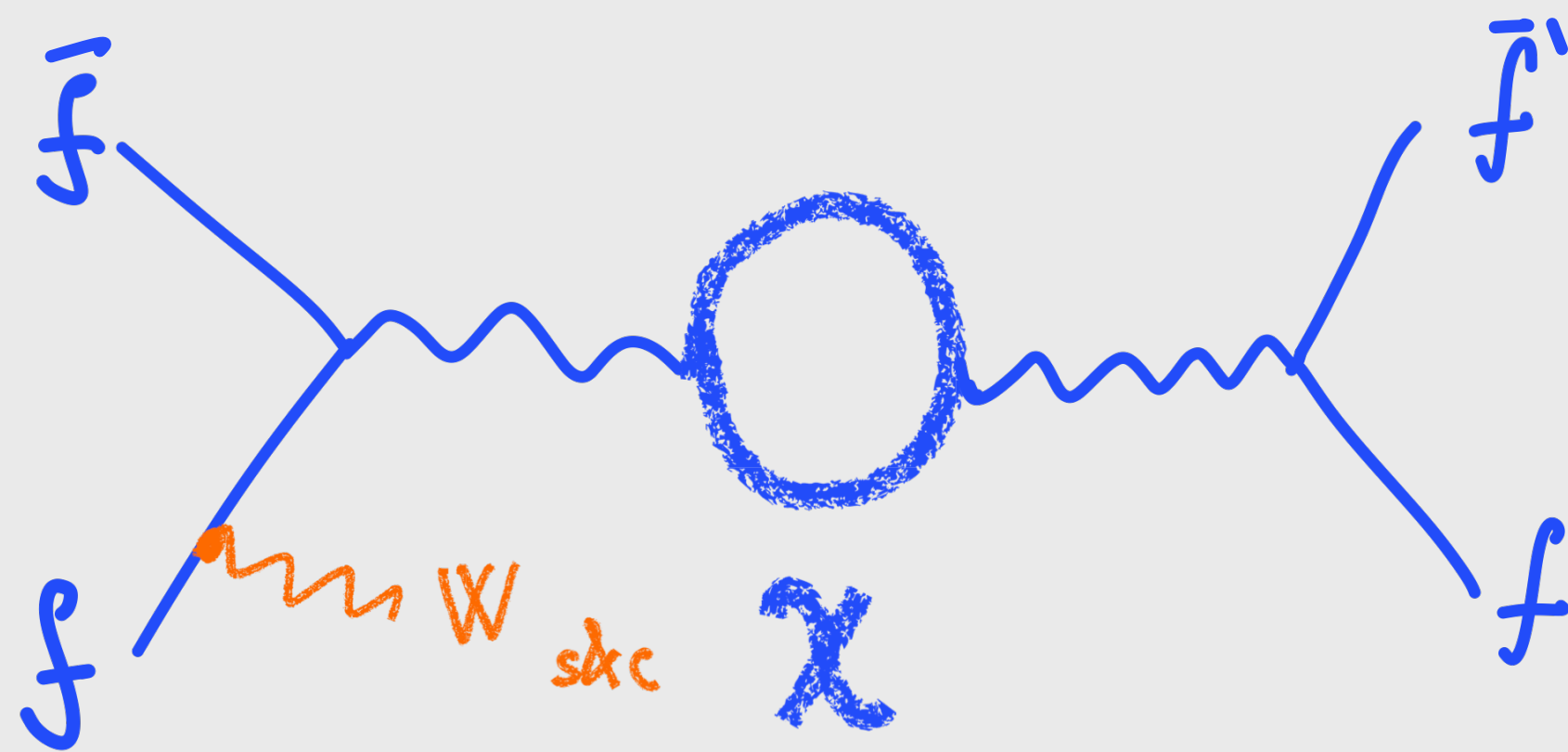


$m_\chi \gg E_{cm}$

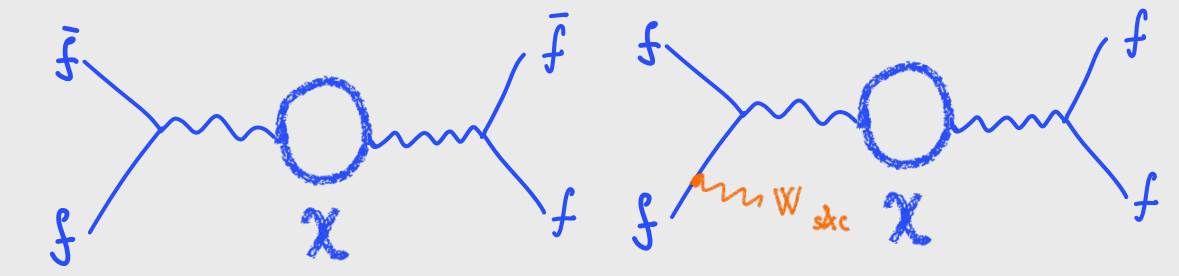


$$O_{2W} \propto (D_\mu W^{\mu\nu})^2$$

$E_{cm} \gg m_W$

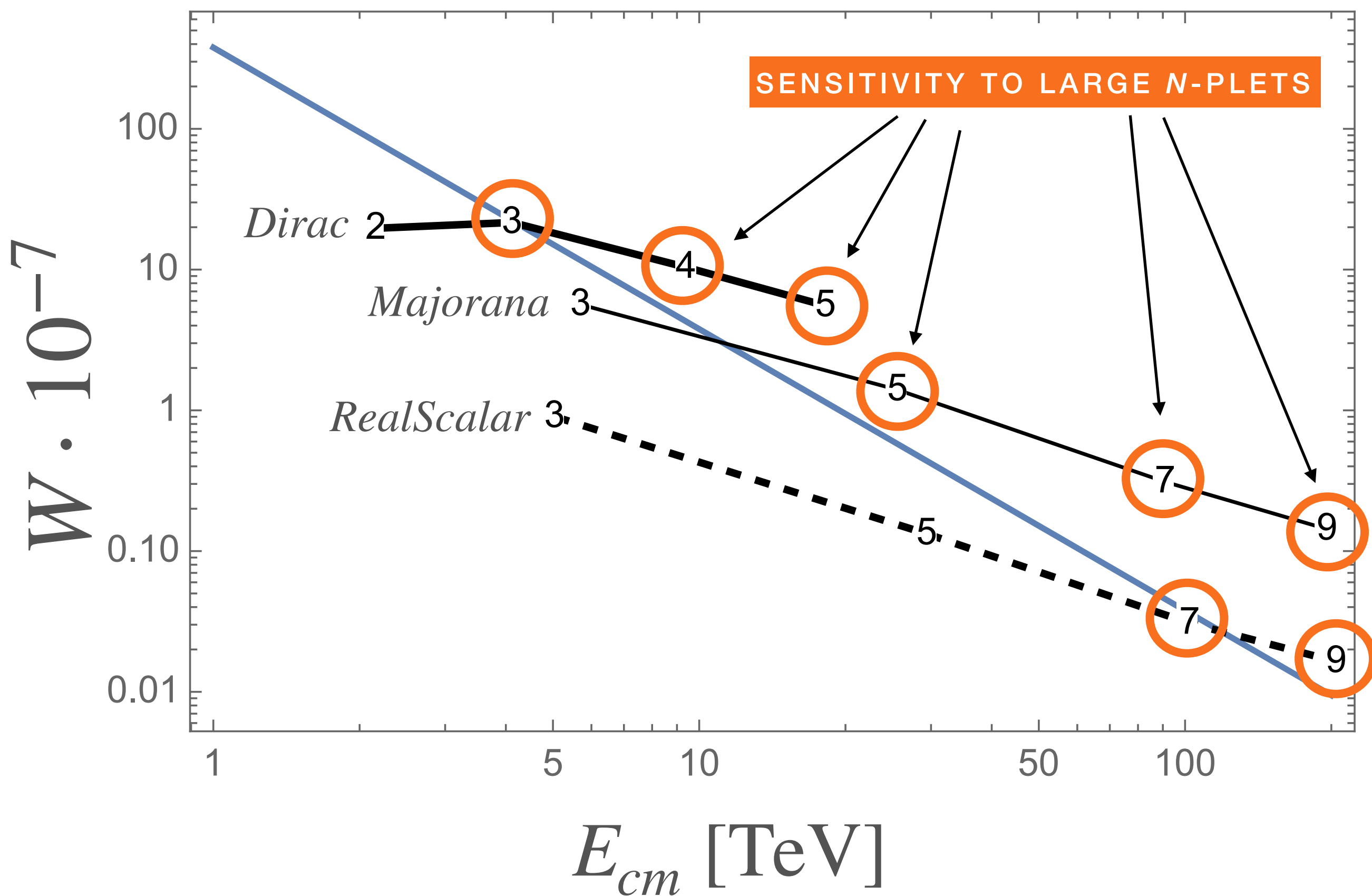


$$\ell^+ \ell^- \rightarrow f\bar{f}, Zh, W^+W^-, Wff'$$



PRECISION

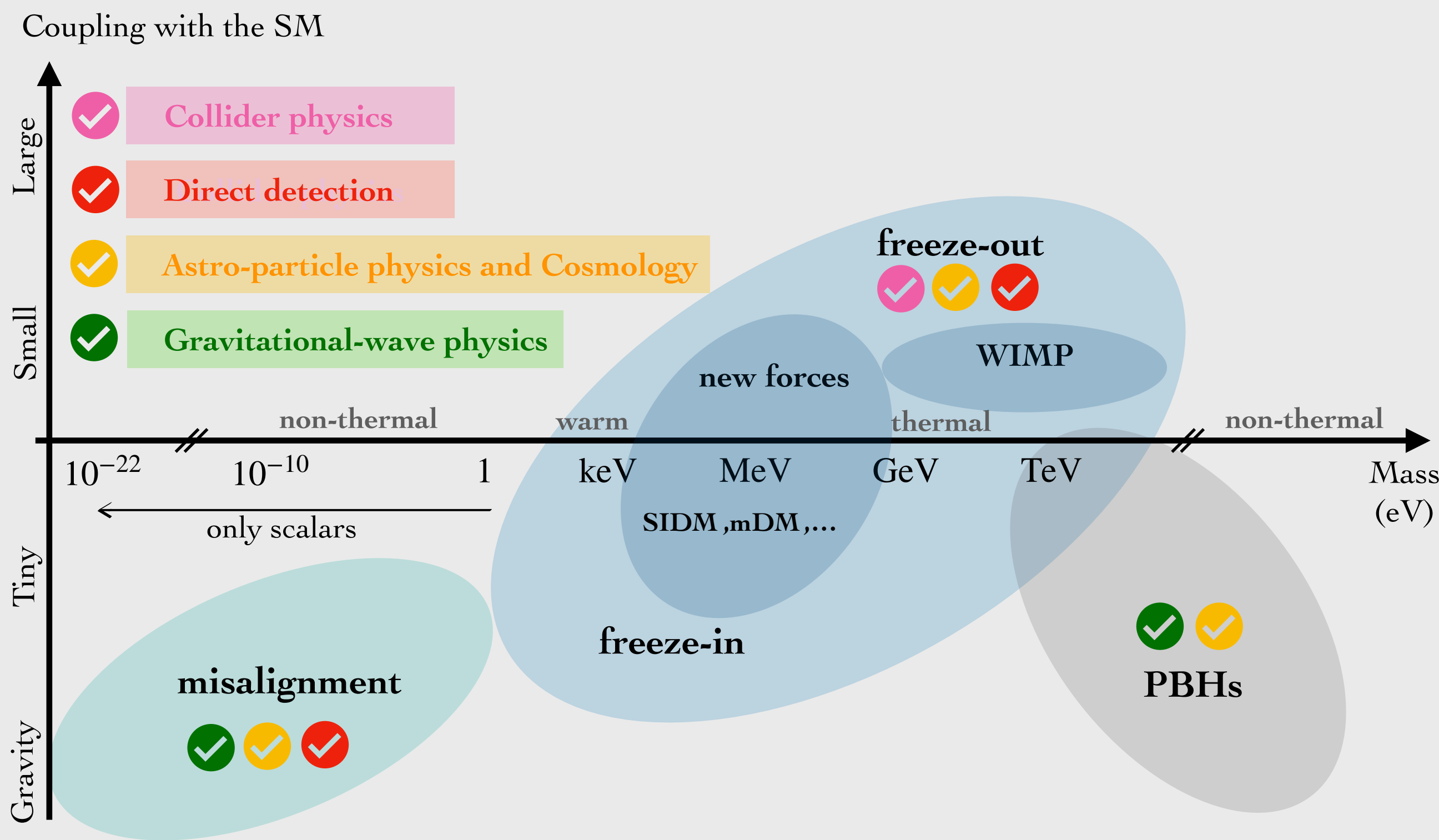
TOTAL CROSS-SECTION



2202.10509	exclusive	inclusive
	$W \times 10^7$	$W \times 10^7$
3 TeV	$[-53, 53]$	$[-41, 41]$
10 TeV	$[-5.71, 5.71]$	$[-3.71, 3.71]$
14 TeV	$[-3.11, 3.11]$	$[-1.90, 1.90]$
30 TeV	$[-0.80, 0.80]$	$[-0.42, 0.42]$

10 TeV	DL	e^{DL-1}	$SL(\frac{\pi}{2})$
$\ell_L \rightarrow \ell'_L$	-0.82	-0.56	0.33
$\ell_L \rightarrow q_L$	-0.78	-0.54	0.34
$\ell_L \rightarrow e_R$	-0.56	-0.43	0.17
$\ell_L \rightarrow u_R$	-0.48	-0.38	0.15
$\ell_L \rightarrow d_R$	-0.43	-0.35	0.13
$\ell_R \rightarrow \ell'_L$	-0.56	-0.43	0.17
$\ell_R \rightarrow q_L$	-0.53	-0.41	0.16
$\ell_R \rightarrow \ell'_R$	-0.30	-0.26	0.09
$\ell_R \rightarrow u_R$	-0.22	-0.20	0.07
$\ell_R \rightarrow d_R$	-0.17	-0.16	0.05

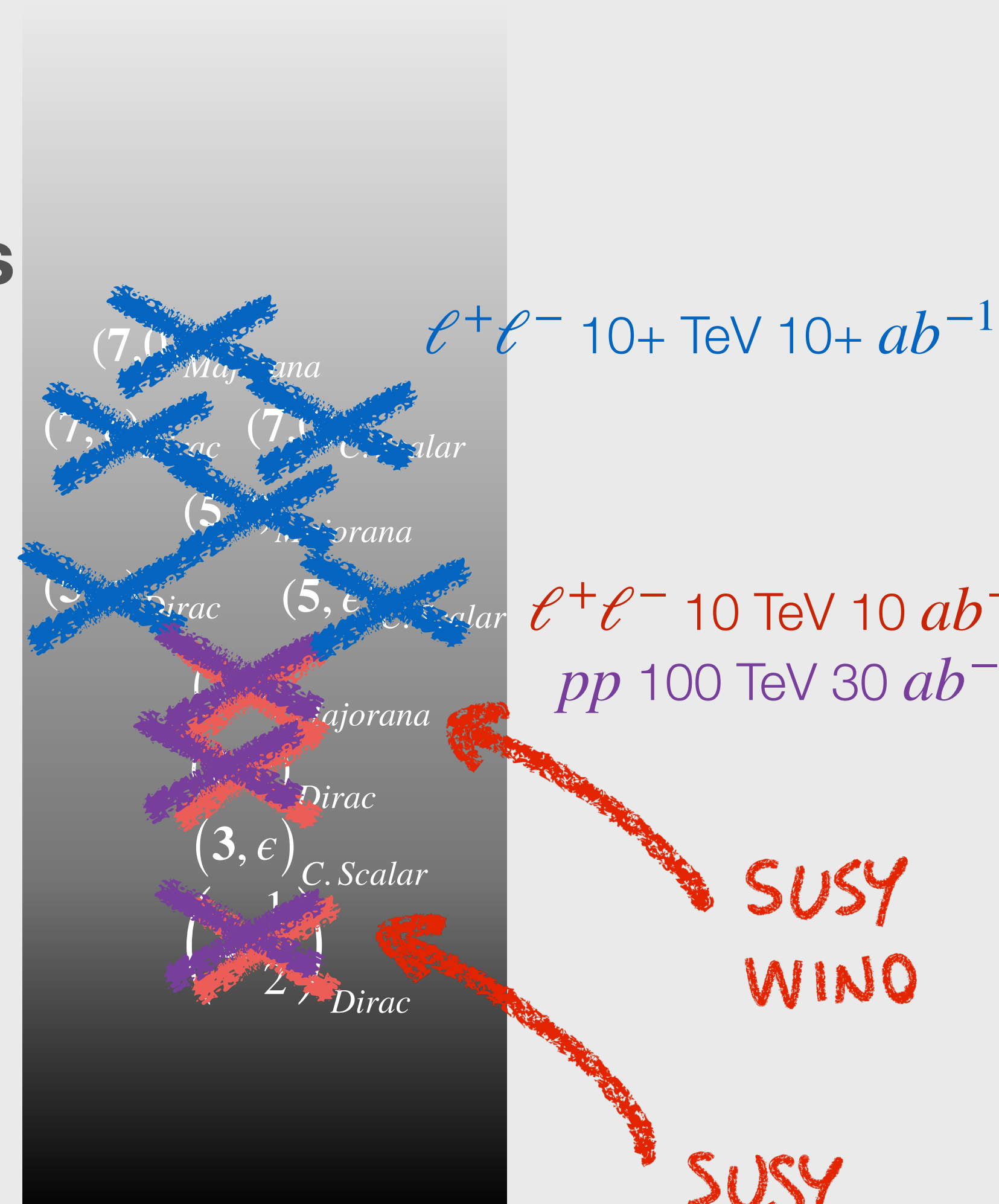
Electroweak Dark Matter: LSP (+NLSP)



Mass

50 TeV

1 TeV



Electroweak symmetry breaking

Big picture questions:

- Extended Higgs Sector
back to “valence” muon collisions
and direct production of new physics
- Higgs compositeness

Electroweak symmetry breaking

Big picture questions:

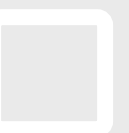
- Extended Higgs Sector
back to “valence” muon collisions
and direct production of new physics

- Higgs compositeness

“The size of the Higgs boson”

it matters because being “point-like” is the source of all the theoretical questions on the Higgs boson and weak scale

... and if it is not ... well, that is physics beyond the Standard Model!



Effects of the size of the Higgs boson

$h \sim \pi$

STRONGLY INTERACTING LIGHT HIGGS

$$\begin{aligned}
 \mathcal{L}_{universal}^{d=6} = & c_H \frac{g_*^2}{m_*^2} \mathcal{O}_H + c_T \frac{N_c \epsilon_q^4 g_*^4}{(4\pi)^2 m_*^2} \mathcal{O}_T + c_6 \lambda \frac{g_*^2}{m_*^2} \mathcal{O}_6 + \frac{1}{m_*^2} [c_W \mathcal{O}_W + c_B \mathcal{O}_B] \\
 & + \frac{g_*^2}{(4\pi)^2 m_*^2} [c_{HW} \mathcal{O}_{HW} + c_{HB} \mathcal{O}_{HB}] + \frac{y_t^2}{(4\pi)^2 m_*^2} [c_{BB} \mathcal{O}_{BB} + c_{GG} \mathcal{O}_{GG}] \\
 & + \frac{1}{g_*^2 m_*^2} [c_{2W} g^2 \mathcal{O}_{2W} + c_{2B} g'^2 \mathcal{O}_{2B}] + c_{3W} \frac{3! g^2}{(4\pi)^2 m_*^2} \mathcal{O}_{3W} \\
 & + c_{y_t} \frac{g_*^2}{m_*^2} \mathcal{O}_{y_t} + c_{y_b} \frac{g_*^2}{m_*^2} \mathcal{O}_{y_b}
 \end{aligned}$$

$$1/f \sim g_*/m_*$$

$$1/(g_* f) \sim 1/m_*$$

$$g_{SM}/(g_* f) \sim g_{SM}/m_*$$

$$\ell_{Higgs} \sim 1/m_*$$



$$\ell^+ \ell^- \rightarrow \mathbf{V}\mathbf{V} + \mathbf{X}$$

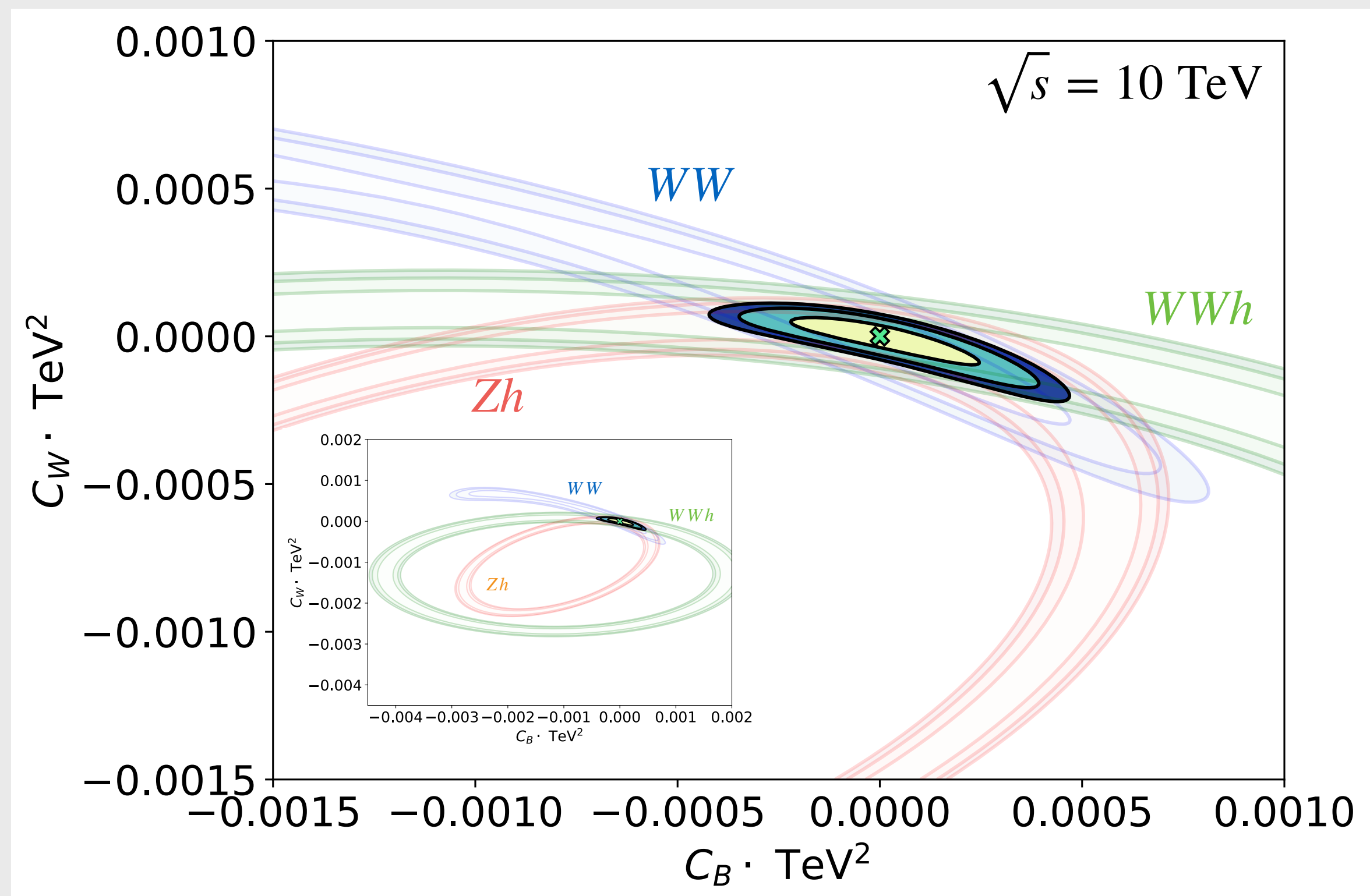
DI-BOSON

⊕ MULTI-BOSON

ZH: BSM and SM amplitudes have the same angular dependences, so the most powerful analysis is a simple cut-and-count.

WW: BSM and SM amplitudes **do not** have the same angular dependences, so the most powerful analysis is differential!

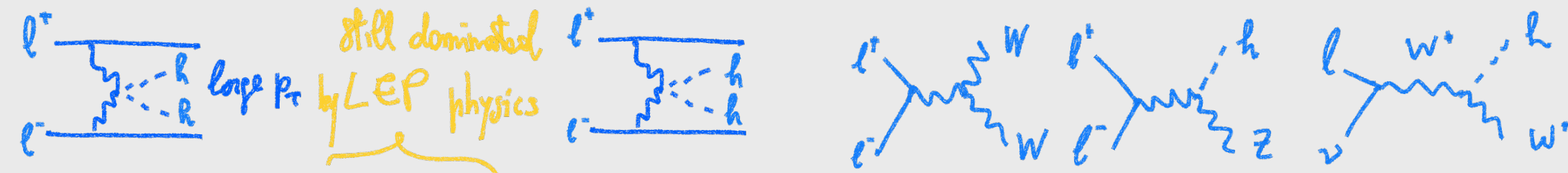
multi-body can contain hard sub-scattering with net electric charge, e.g. $e\nu \rightarrow Wh, WZ$ with new BSM couplings dependence



- ZH: elliptical belt in 2D BSM coupling space
- WW: basin in 2D BSM coupling space
- WWh: elliptical belt in 2D BSM coupling space

Effects of the size of the Higgs boson

STRONGLY INTERACTING TOP AND HIGGS

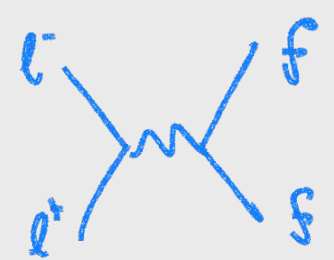


$$\mathcal{L}_{universal}^{d=6} = c_H \frac{g_*^2}{m_*^2} \mathcal{O}_H + c_T \frac{N_c \epsilon_q^4 g_*^4}{(4\pi)^2 m_*^2} \mathcal{O}_T + c_6 \lambda \frac{g_*^2}{m_*^2} \mathcal{O}_6 + \frac{1}{m_*^2} [c_W \mathcal{O}_W + c_B \mathcal{O}_B]$$

$$1/f \sim g_*/m_*$$

$$+ \frac{g_*^2}{(4\pi)^2 m_*^2} [c_{HW} \mathcal{O}_{HW} + c_{HB} \mathcal{O}_{HB}] + \frac{y_t^2}{(4\pi)^2 m_*^2} [c_{BB} \mathcal{O}_{BB} + c_{GG} \mathcal{O}_{GG}]$$

$$1/(g_* f) \sim 1/m_*$$

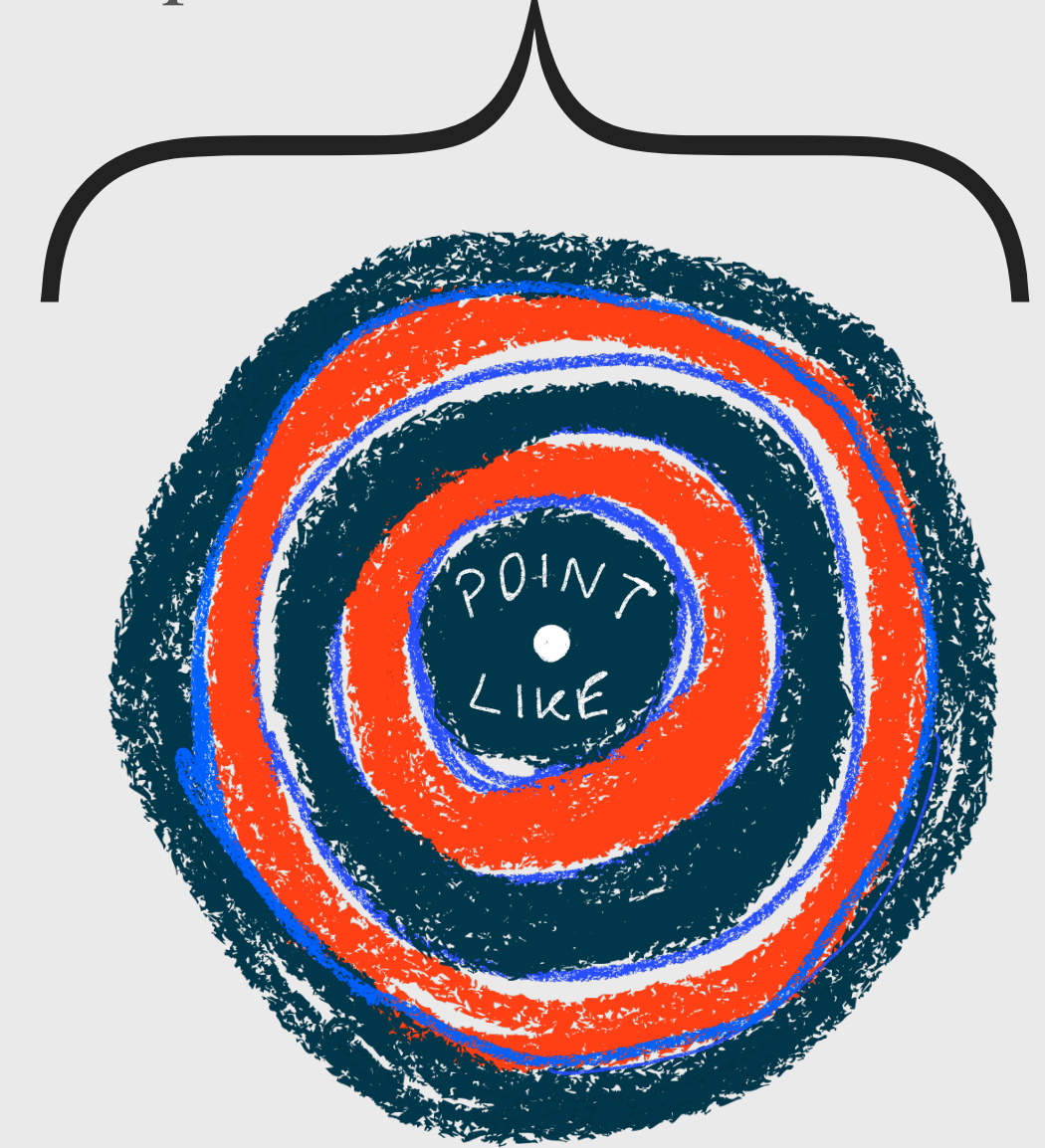


$$+ \frac{1}{g_*^2 m_*^2} [c_{2W} g^2 \mathcal{O}_{2W} + c_{2B} g'^2 \mathcal{O}_{2B}] + c_{3W} \frac{3! g^2}{(4\pi)^2 m_*^2} \mathcal{O}_{3W}$$

$$g_{SM}/(g_* f) \sim g_{SM}/m_*$$

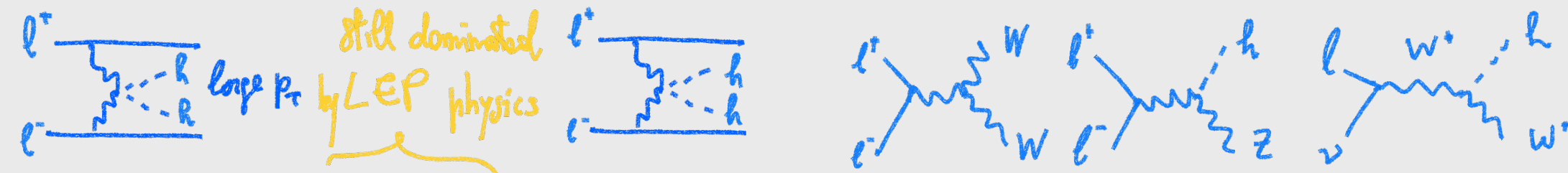
$$+ c_{y_t} \frac{g_*^2}{m_*^2} \mathcal{O}_{y_t} + c_{y_b} \frac{g_*^2}{m_*^2} \mathcal{O}_{y_b}$$

$$\ell_{top} \sim 1/m_* \sim \ell_{Higgs}$$



Effects of the size of the Higgs boson

STRONGLY INTERACTING TOP AND HIGGS

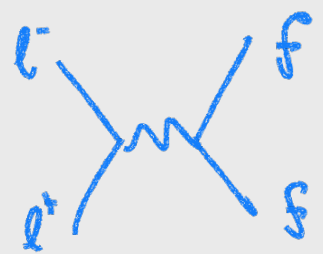


$$\mathcal{L}_{universal}^{d=6} = c_H \frac{g_*^2}{m_*^2} \mathcal{O}_H + c_T \frac{N_c \epsilon_q^4 g_*^4}{(4\pi)^2 m_*^2} \mathcal{O}_T + c_6 \lambda \frac{g_*^2}{m_*^2} \mathcal{O}_6 + \frac{1}{m_*^2} [c_W \mathcal{O}_W + c_B \mathcal{O}_B]$$

$$1/f \sim g_*/m_*$$

$$+ \frac{g_*^2}{(4\pi)^2 m_*^2} [c_{HW} \mathcal{O}_{HW} + c_{HB} \mathcal{O}_{HB}] + \frac{y_t^2}{(4\pi)^2 m_*^2} [c_{BB} \mathcal{O}_{BB} + c_{GG} \mathcal{O}_{GG}]$$

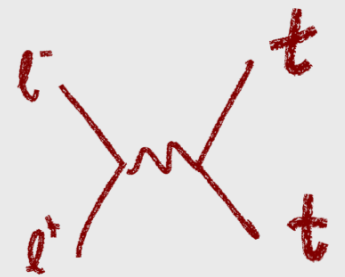
$$1/(g_* f) \sim 1/m_*$$



$$+ \frac{1}{g_*^2 m_*^2} [c_{2W} g^2 \mathcal{O}_{2W} + c_{2B} g'^2 \mathcal{O}_{2B}] + c_{3W} \frac{3! g^2}{(4\pi)^2 m_*^2} \mathcal{O}_{3W}$$

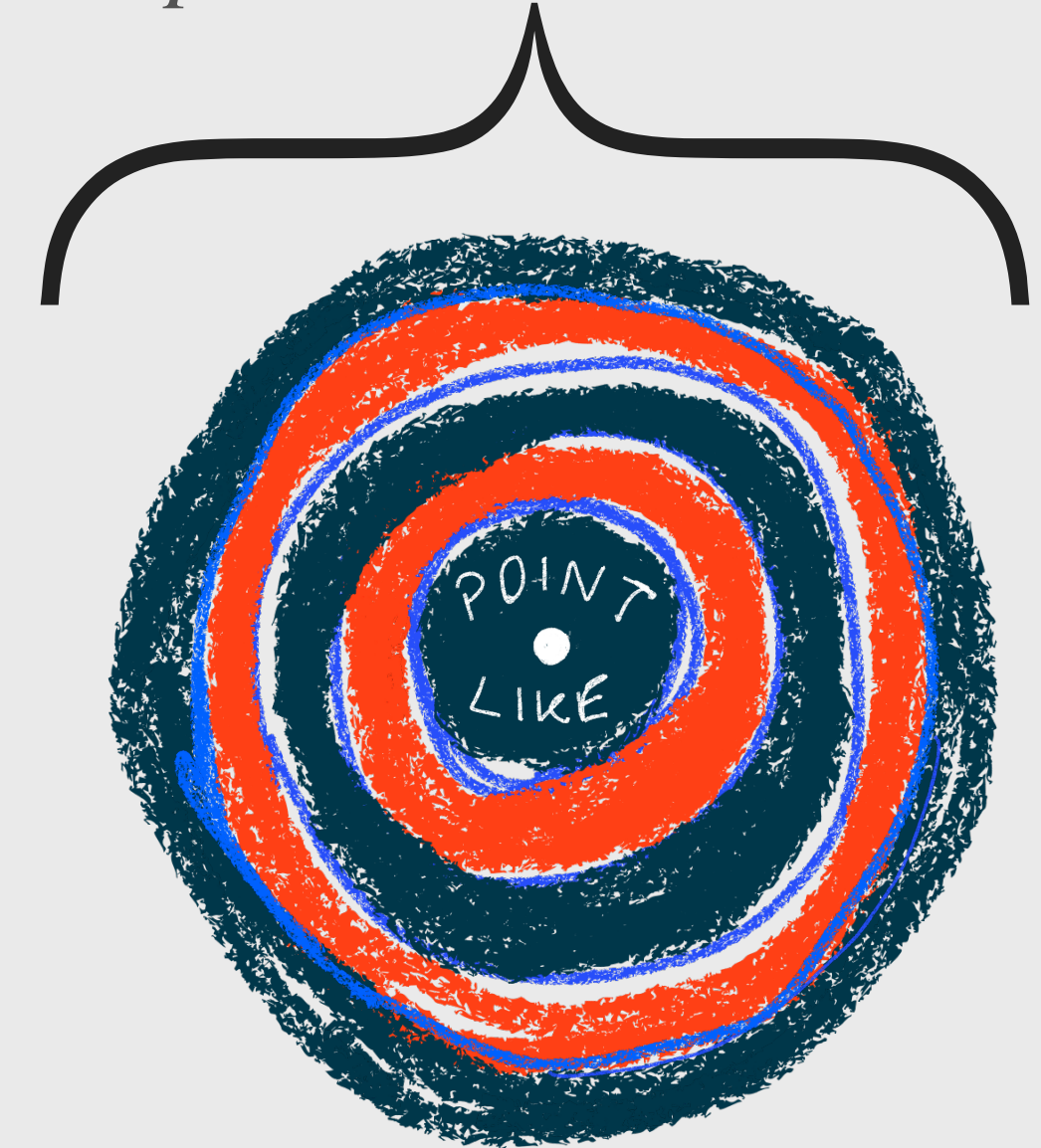
$$g_{SM}/(g_* f) \sim g_{SM}/m_*$$

$$+ c_{y_t} \frac{g_*^2}{m_*^2} \mathcal{O}_{y_t} + c_{y_b} \frac{g_*^2}{m_*^2} \mathcal{O}_{y_b}$$



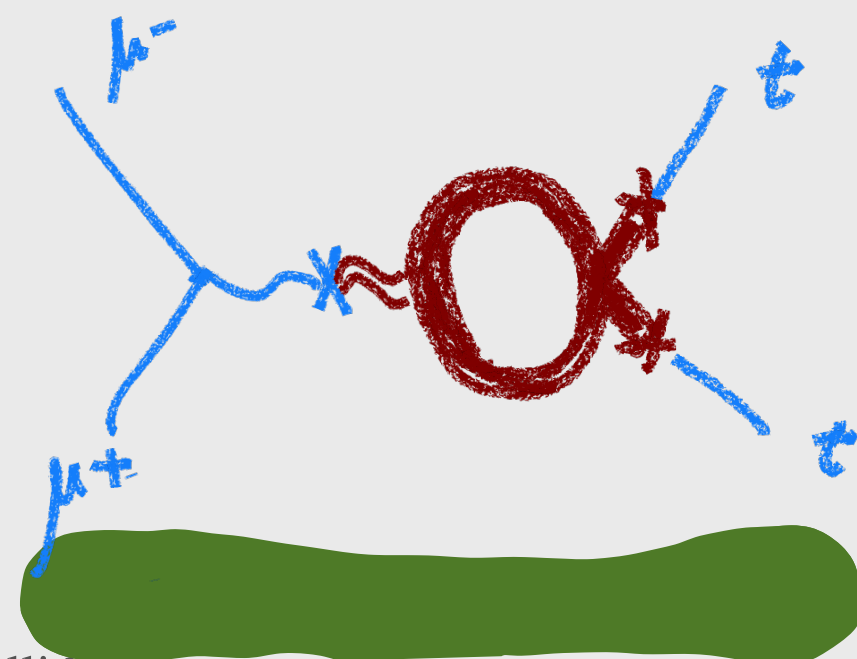
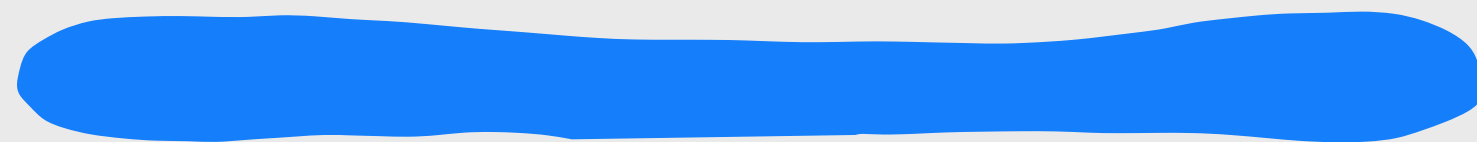
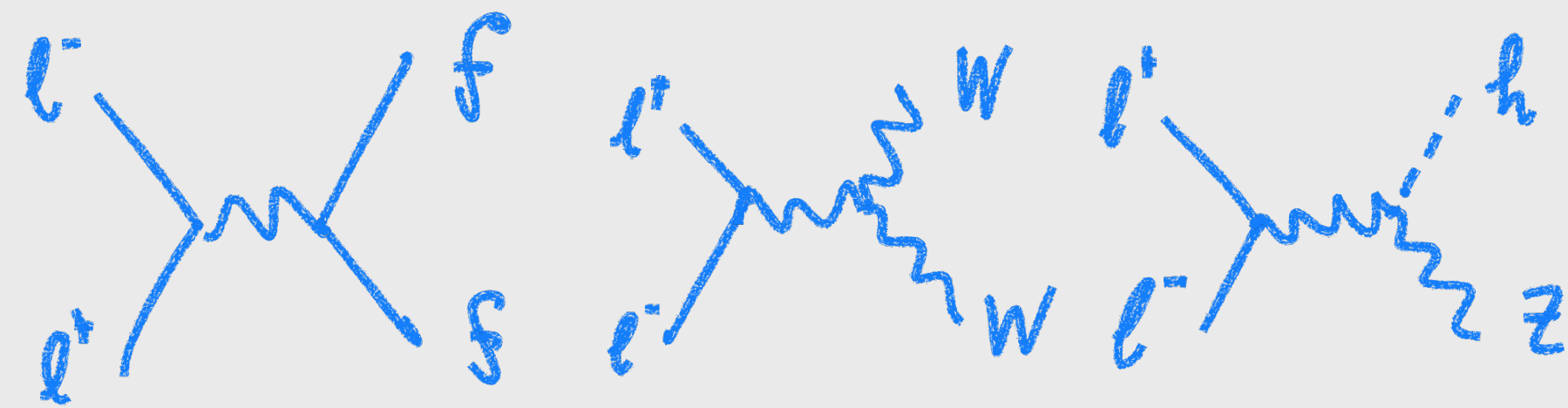
$$+ c_{tD} \frac{g_*^2}{m_*^2} \mathcal{O}_{tD}$$

$$\ell_{top} \sim 1/m_* \sim \ell_{Higgs}$$

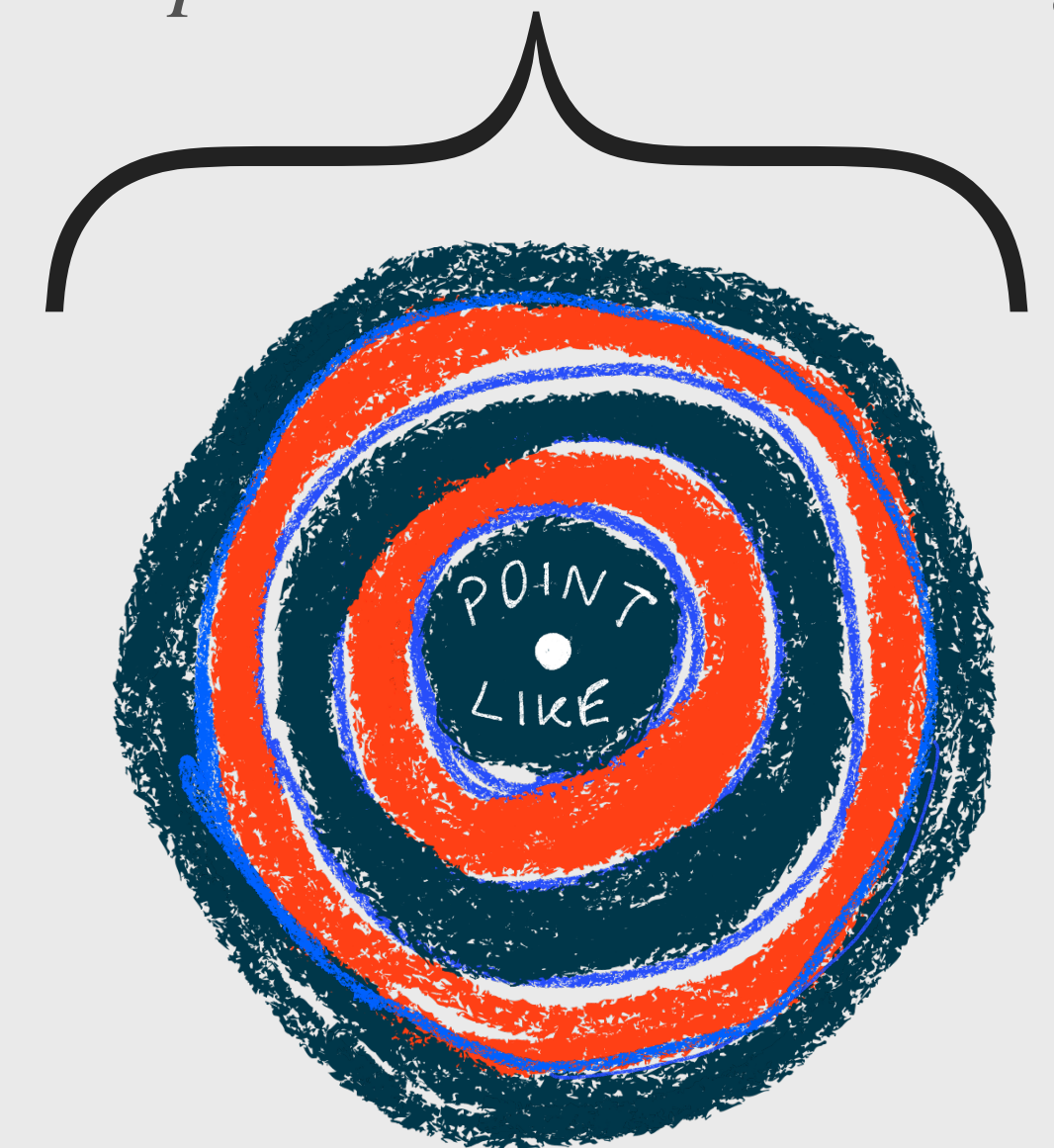


Effects of the size of the Higgs boson

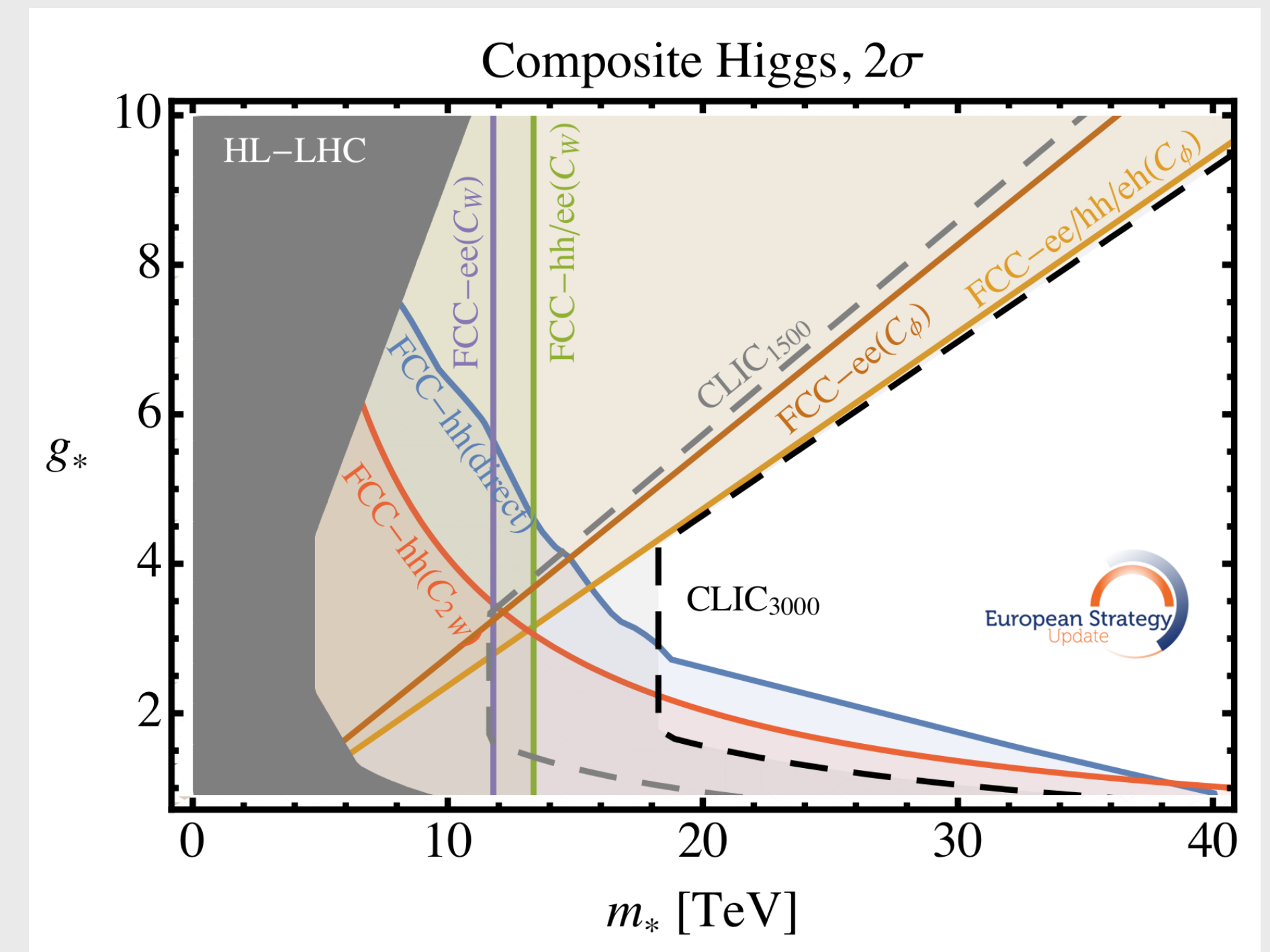
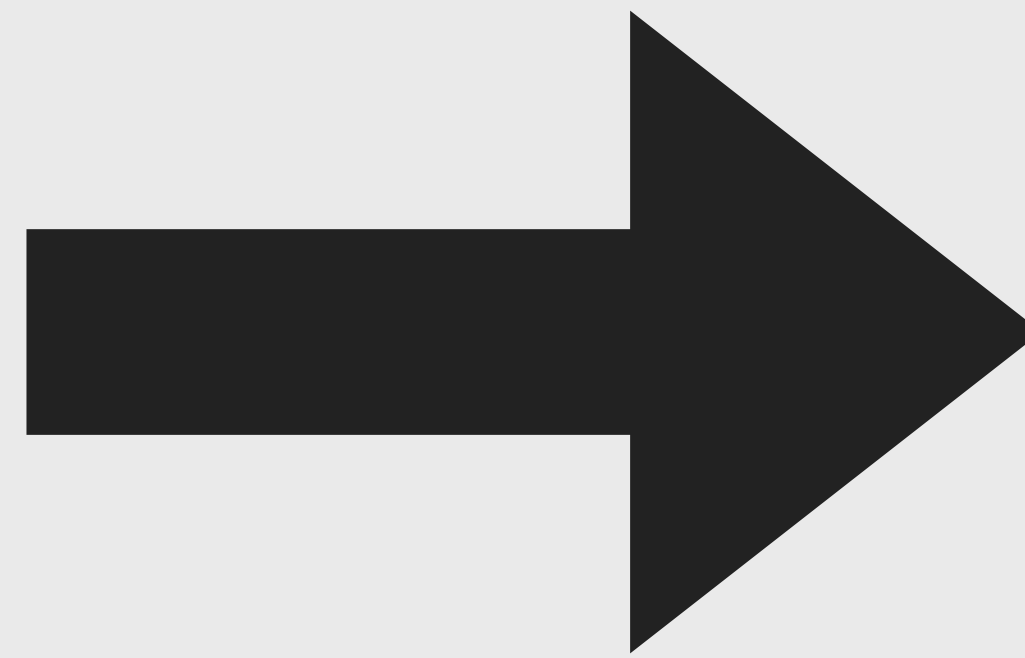
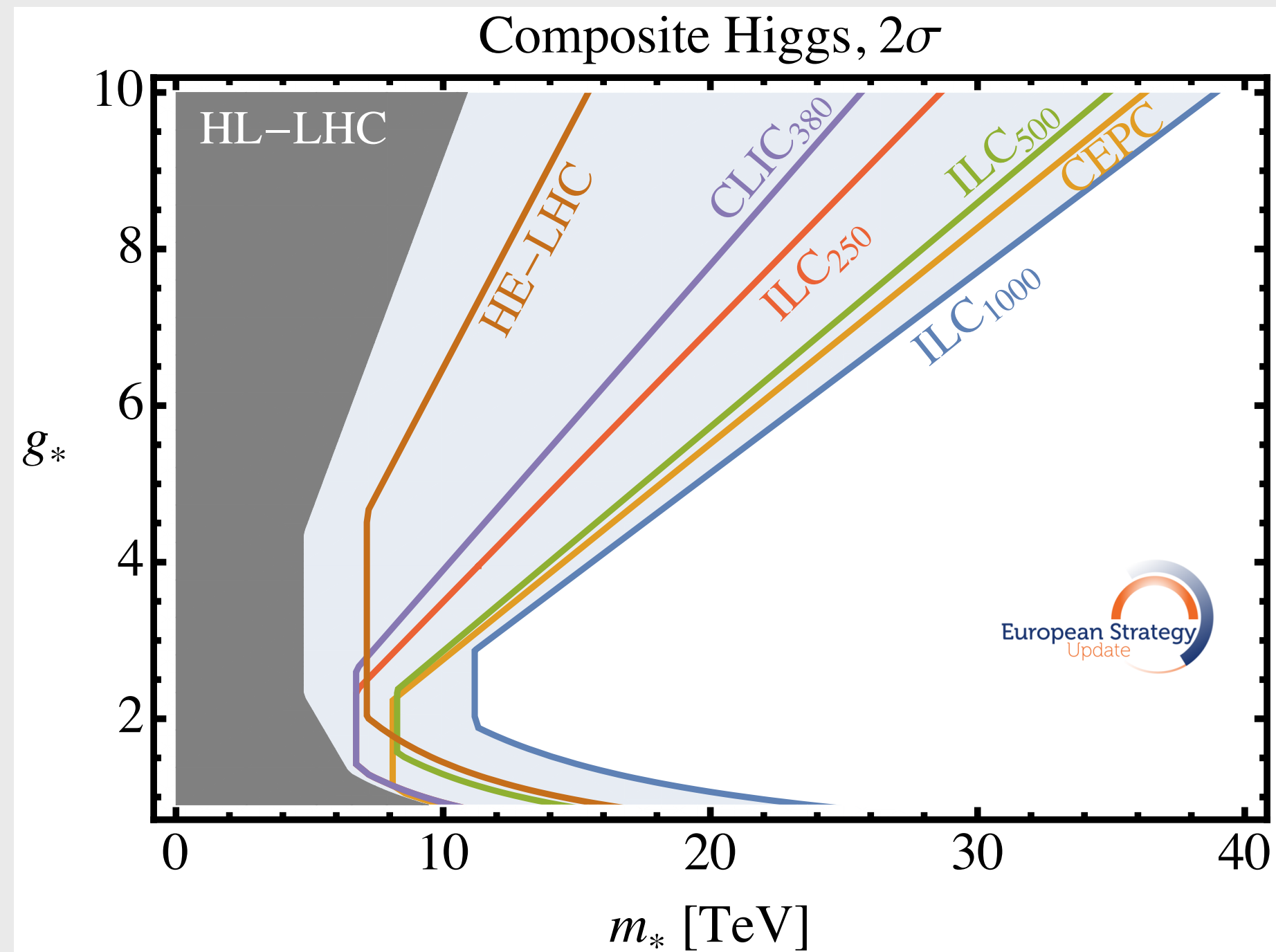
STRONGLY INTERACTING TOP AND HIGGS



$$\ell_{top} \sim 1/m_{\star} \sim \ell_{Higgs}$$

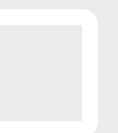


Looking ahead



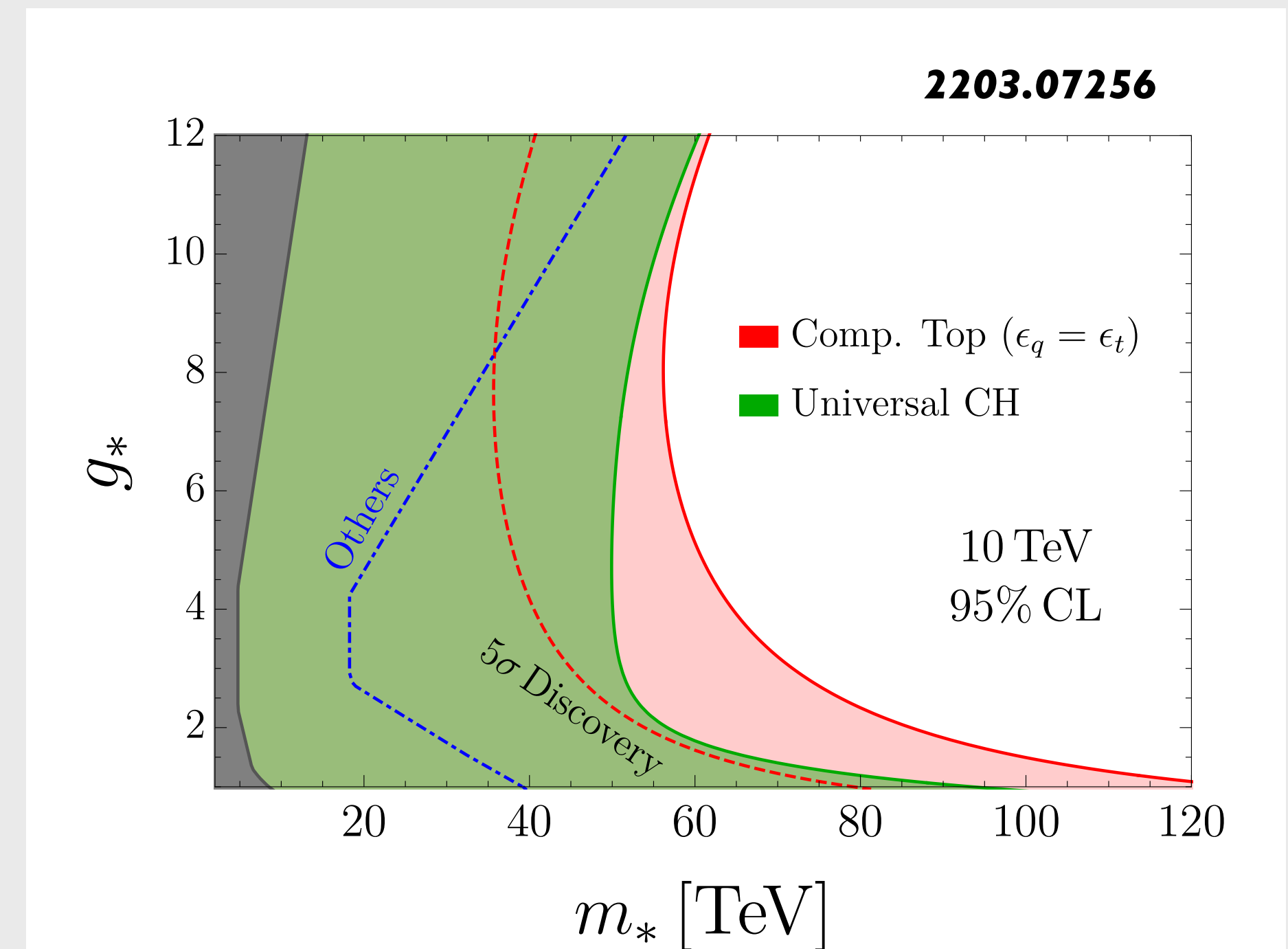
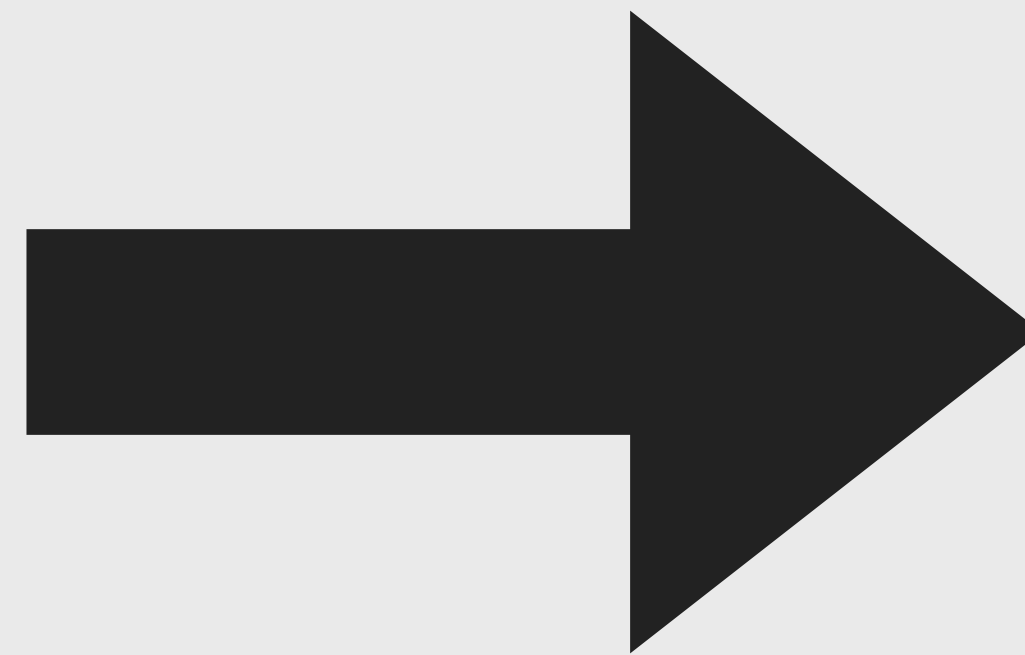
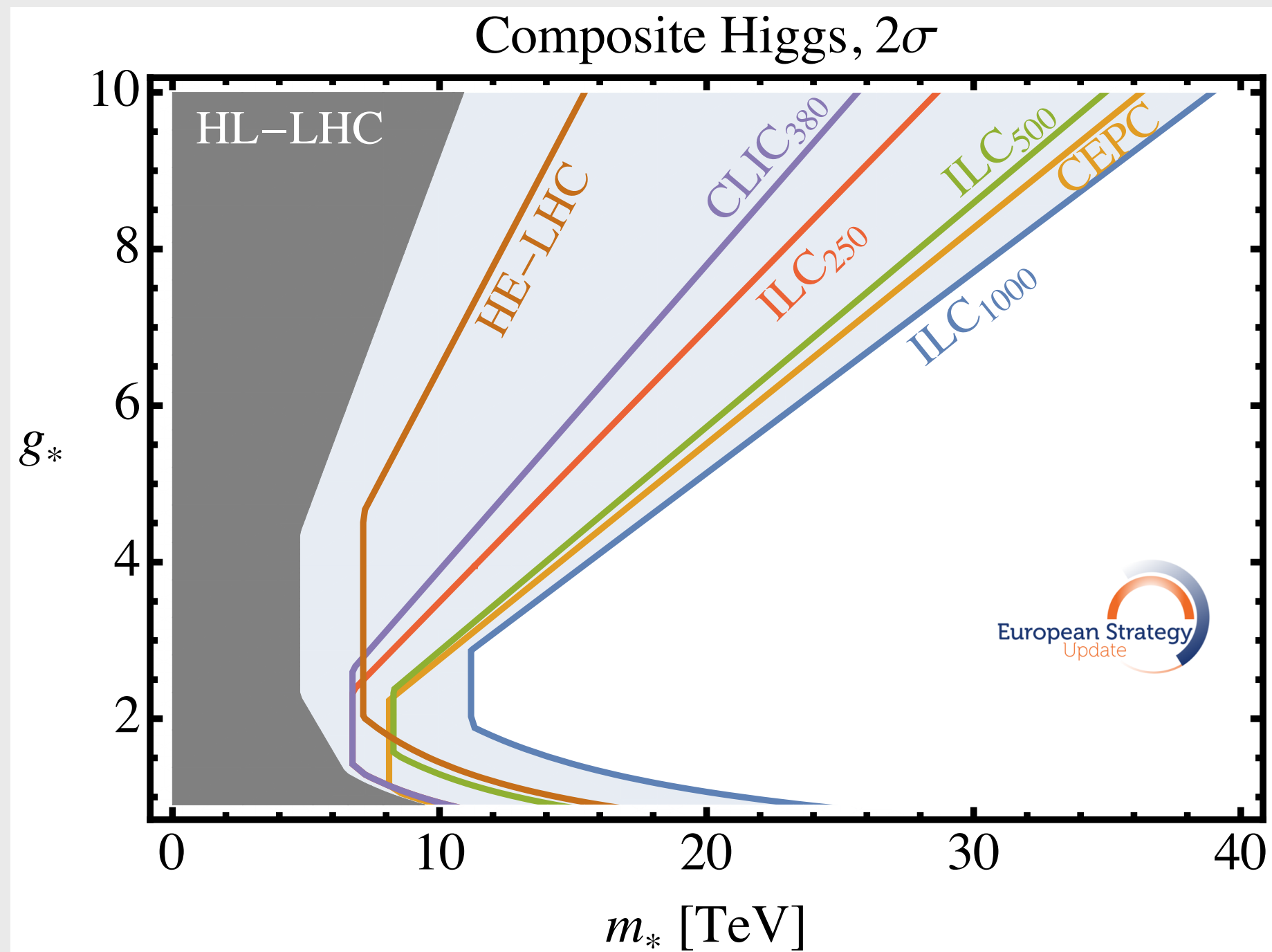
**compositeness at
few TeV @ HL-LHC**

**compositeness at
few 10 TeV**



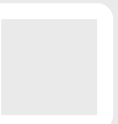
Looking ahead

UNIQUE AVENUE TO EXPLORE WEAK INTERACTIONS
FAR OFFSHORE FROM THE WEAK SCALE



compositeness at
few TeV @ HL-LHC

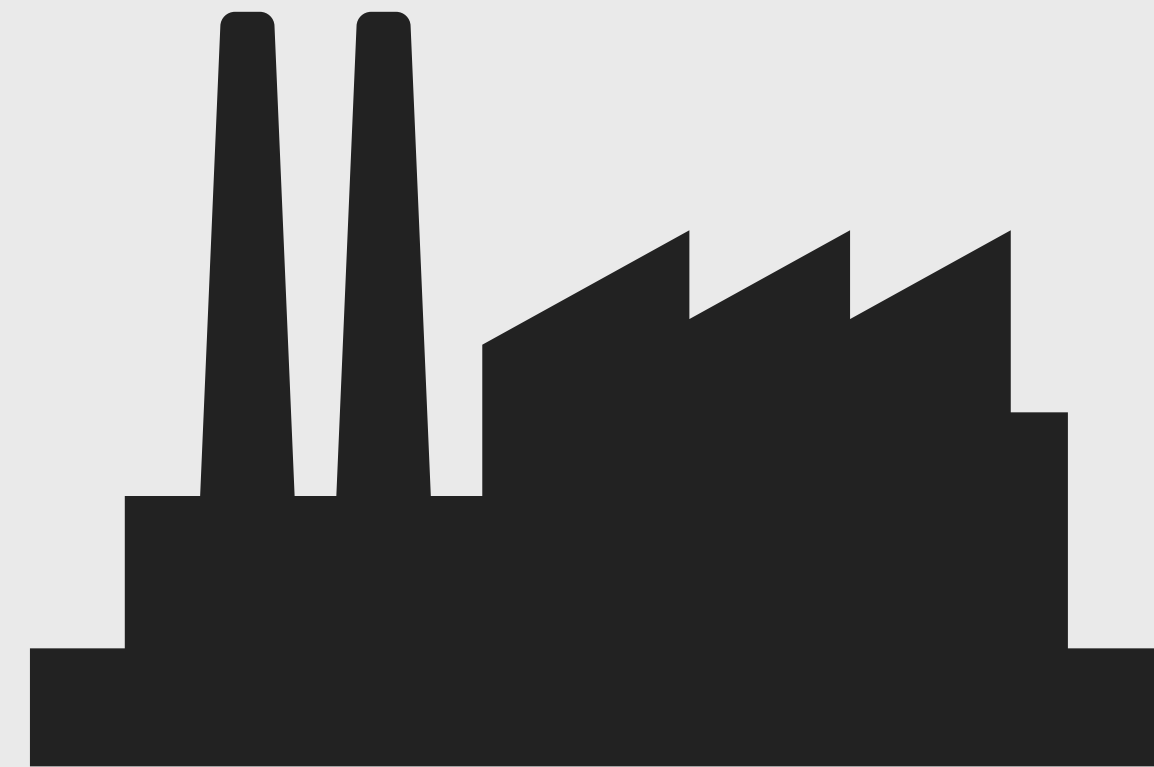
compositeness at
few 100 TeV



Microscope & Factory



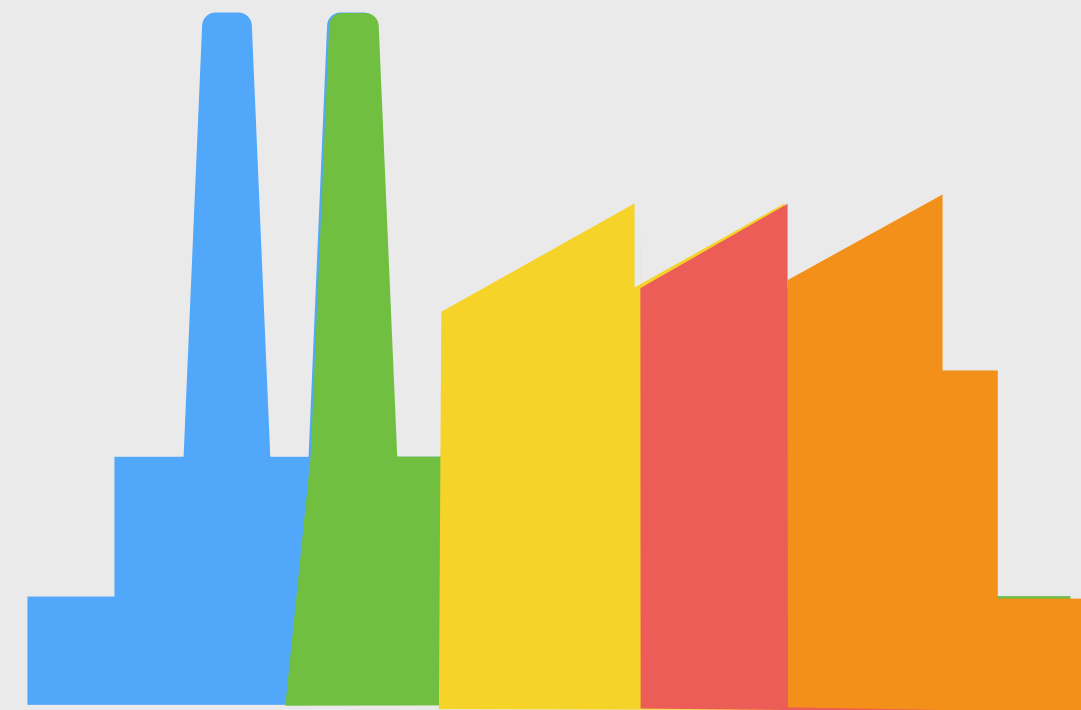
&



Microscope & Factory

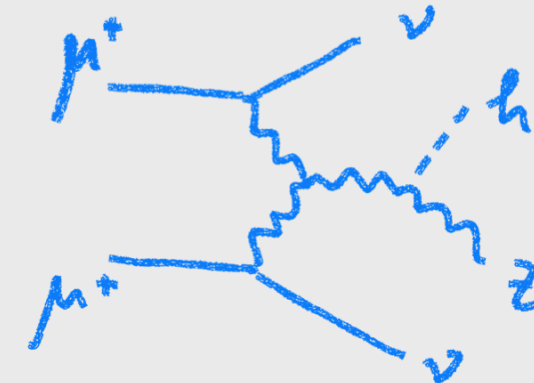
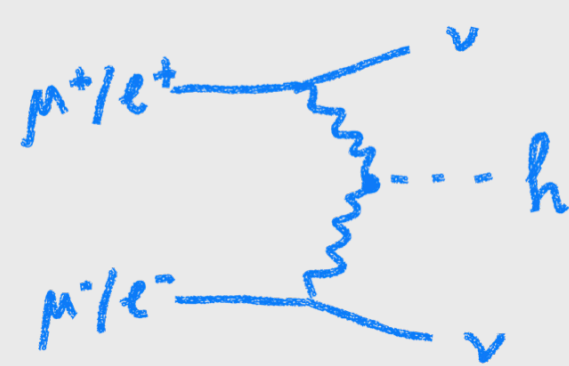
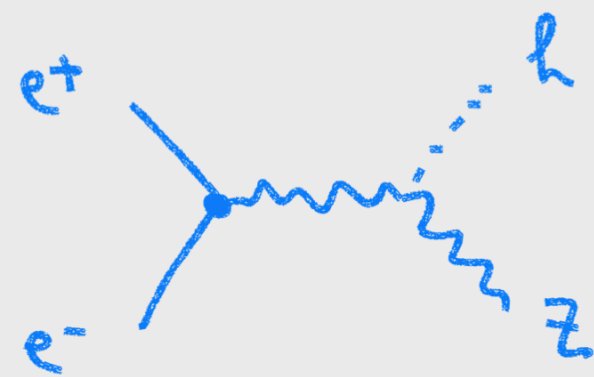


&



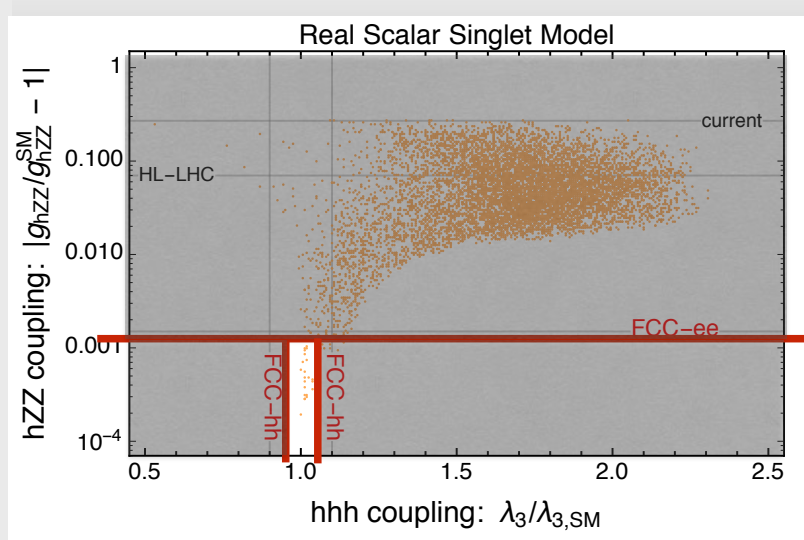
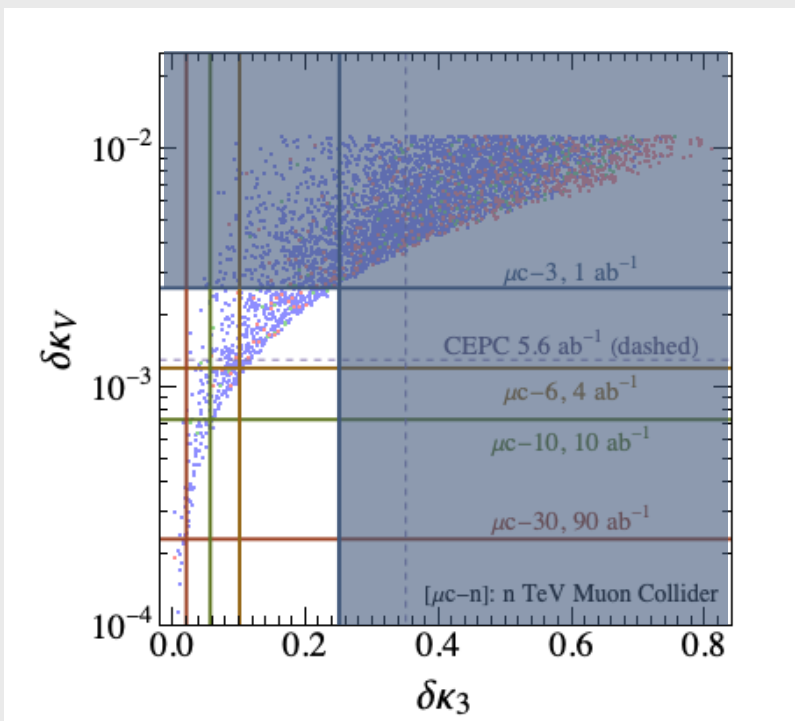
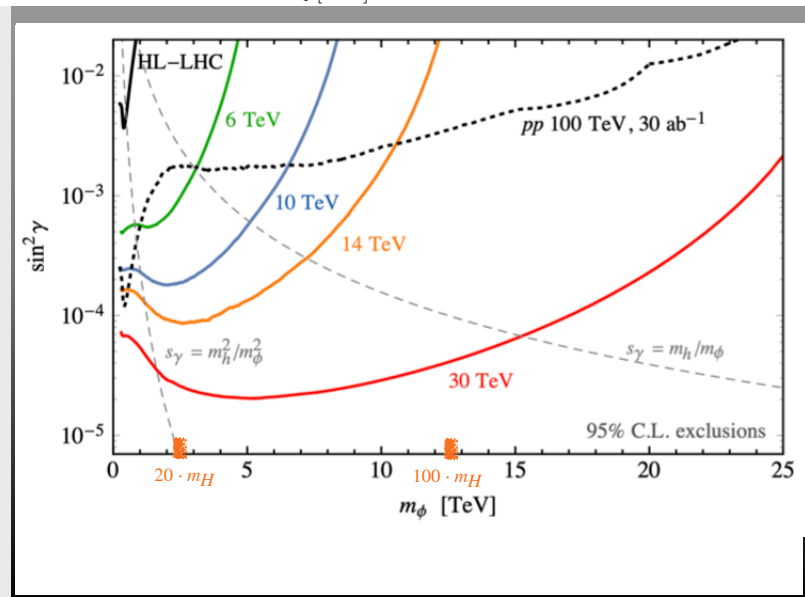
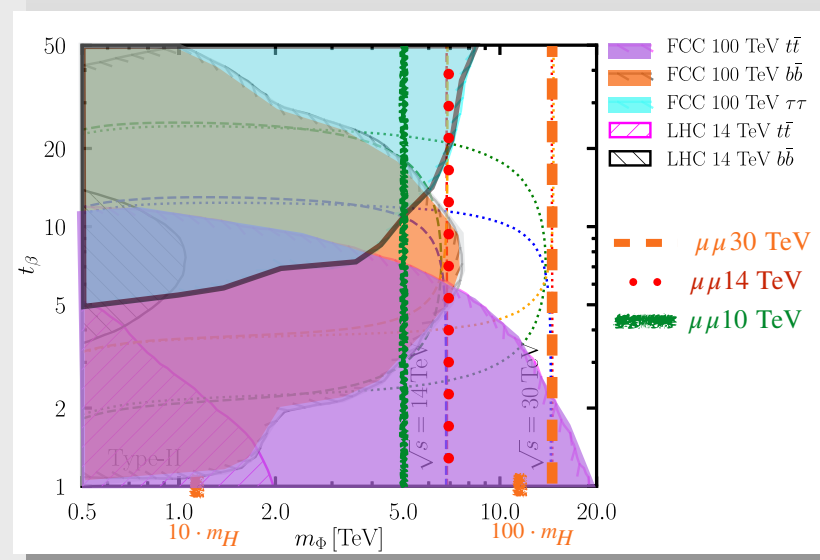
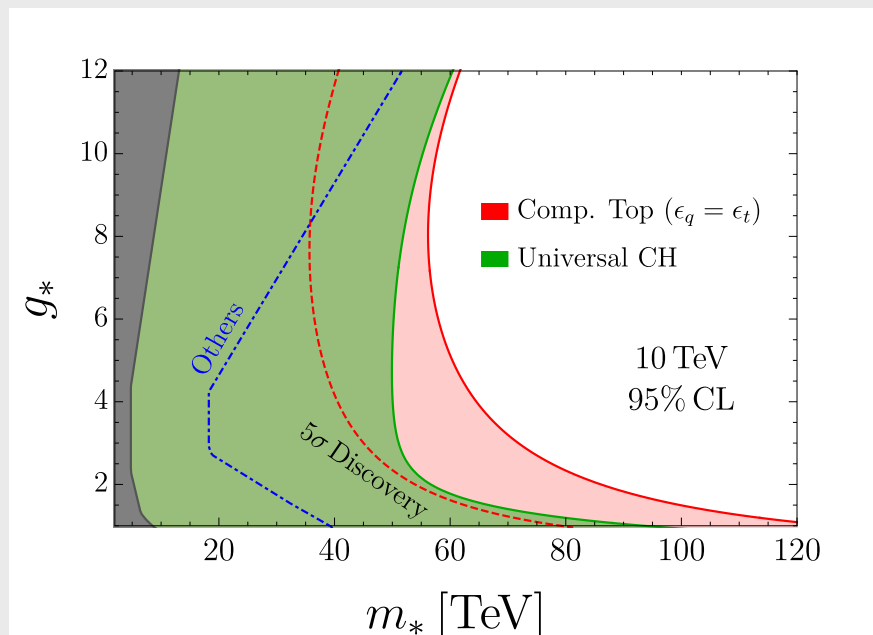
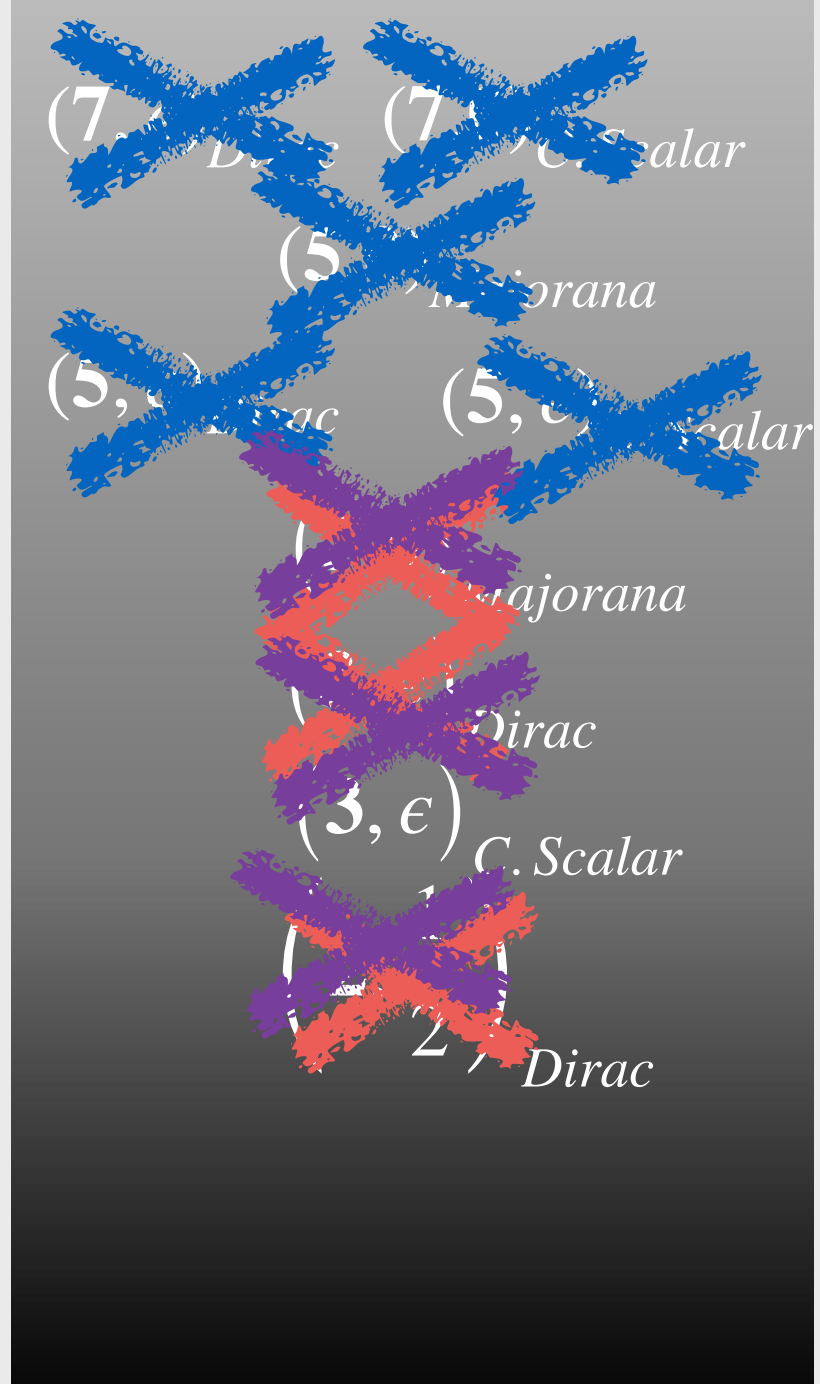
Summary and outlook: Higgs boson

- We have ambitious plans to thoroughly probe the Higgs boson



$$\frac{\delta g_h}{g_h} \simeq 10^{-3} \quad \rightarrow \quad \Lambda_{NP} \simeq 100 \cdot m_h$$

Summary and outlook: BSM



effective c.o.m. energy

100 TeV

10+ TEV

Extraordinary probes of Higgs boson, electroweak new physics and Dark Matter

3 TEV

Several important milestones: **full exploration of TeV EW states, EW phase transition, TeV Dark Matter**

1 TeV

200-300 GeV

Precision study of the Higgs boson



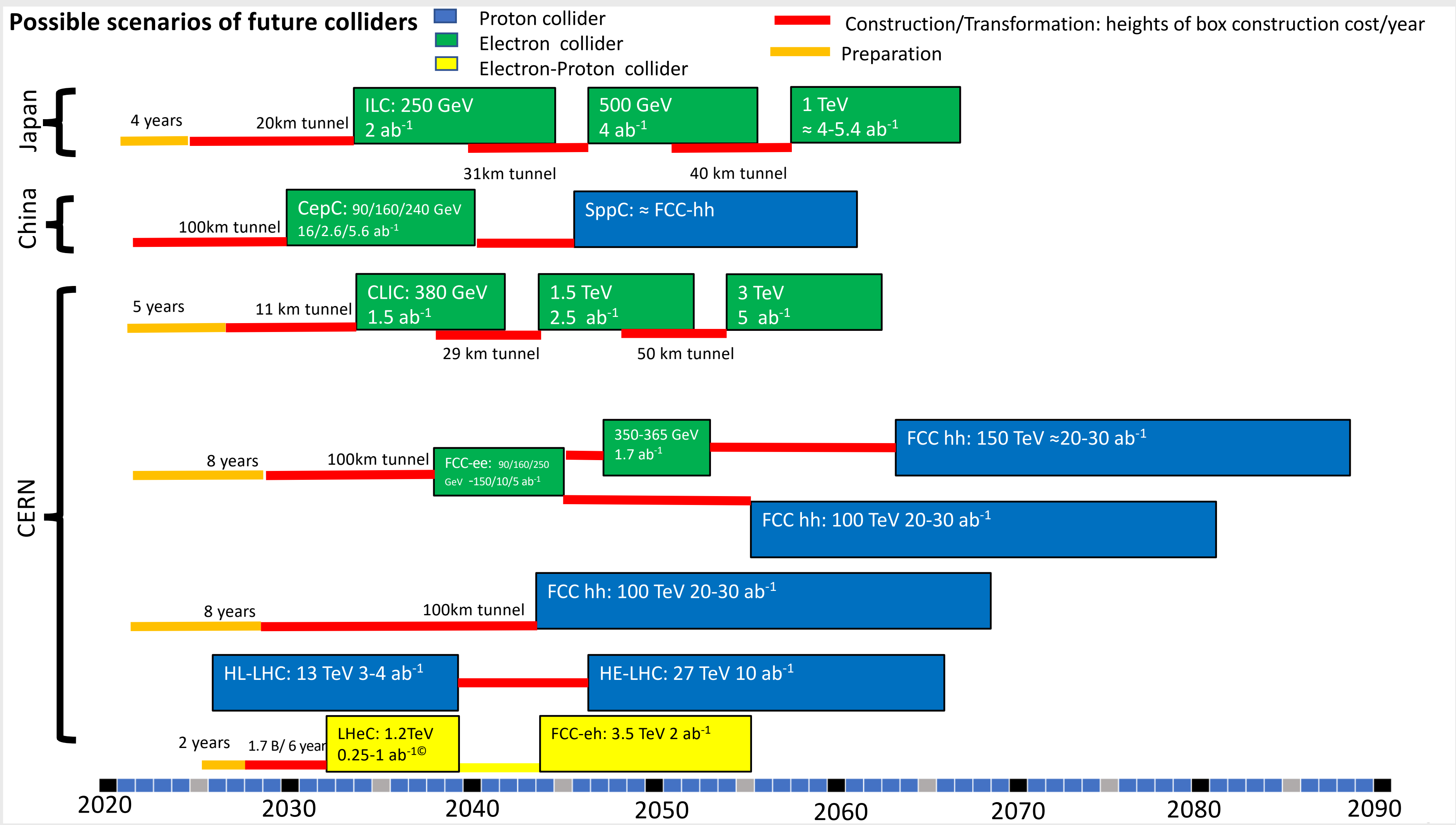
“WIMP” Dark Matter

EW symmetry breaking

EW phase transition
see backup



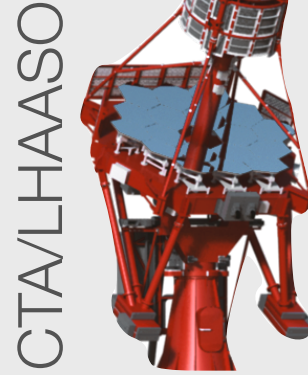
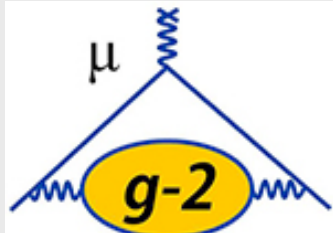
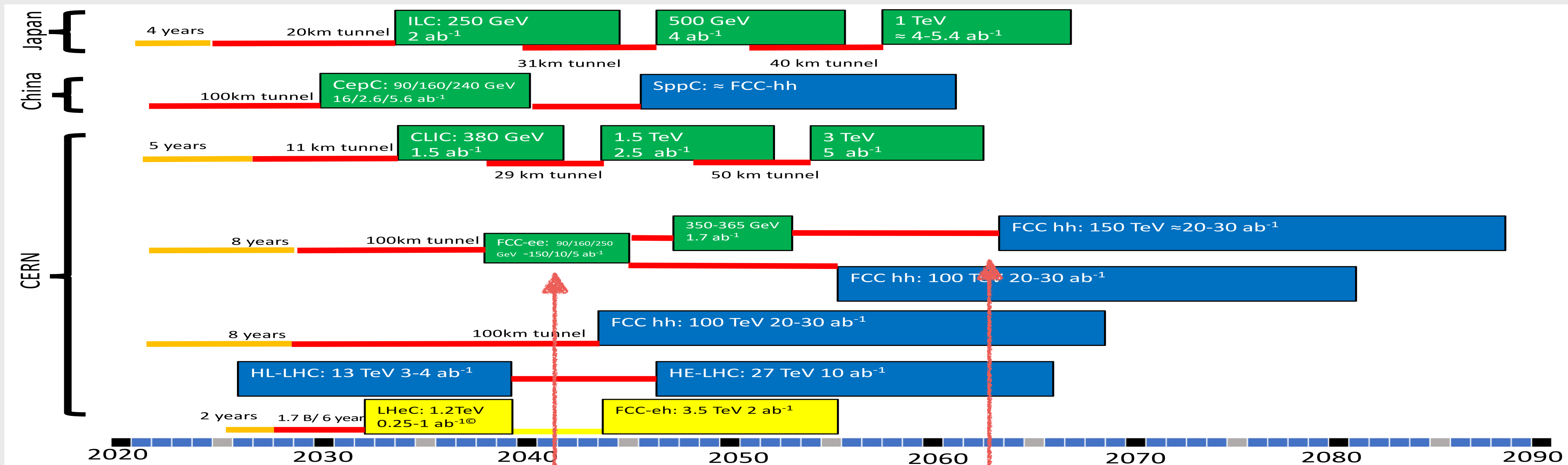
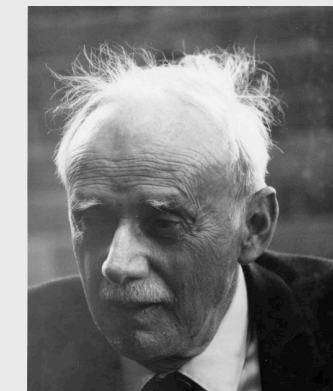
Are we ready?



Are we ready?



$$n_i = \frac{1}{e^{(\epsilon_i - \mu)/k_B T} + 1}$$



2060¹⁰ first *pp* collisions pushing the energy frontier

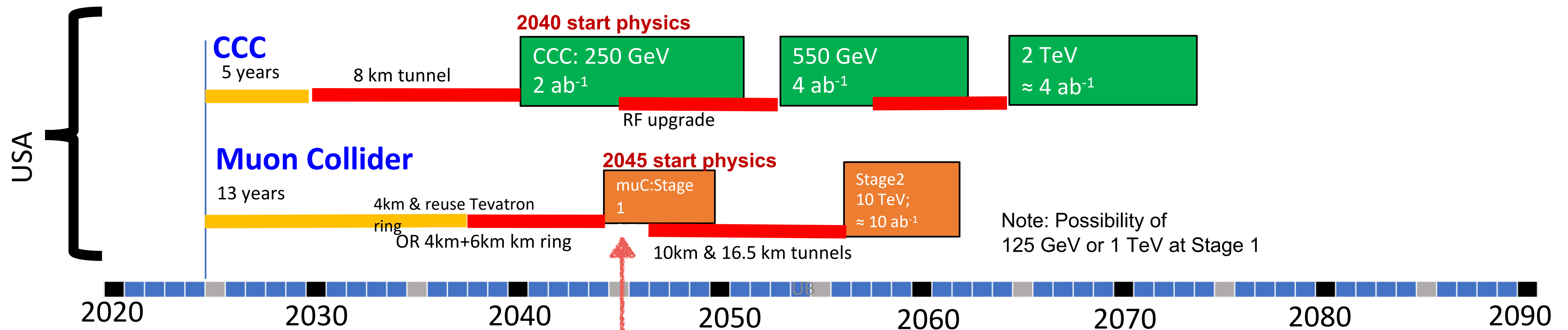
2045 ± 5 first *e⁺e⁻* collisions pushing the intensity frontier

Possible scenarios of future colliders



Original from ESG by UB
Updated July 25, 2022 by MN

Proposals emerging from this Snowmass for a US based collider



- **Timelines technologically limited**
- Uncertainties to be sorted out
 - Find a contact lab(s)
 - Successful R&D and feasibility demonstration for CCC and Muon Collider
 - Evaluate CCC progress in the international context, and consider proposing an ILC/CCC [ie CCC used as an upgrade of ILC] or a CCC only option in the US.
 - International Cost Sharing

2045₀¹⁰ first $\mu^+\mu^-$ collisions pushing both the intensity and energy frontier

VERY SIGNIFICANT INTEREST FROM US COMMUNITY IN RECENT YEARS CULMINATED AT SNOWMASS 2022

Are we ready?

TIME TO WORK ON MUON COLLIDER PHYSICS IS NOW!

Thank you!

TIME TO WORK ON MUON COLLIDER PHYSICS IS NOW!

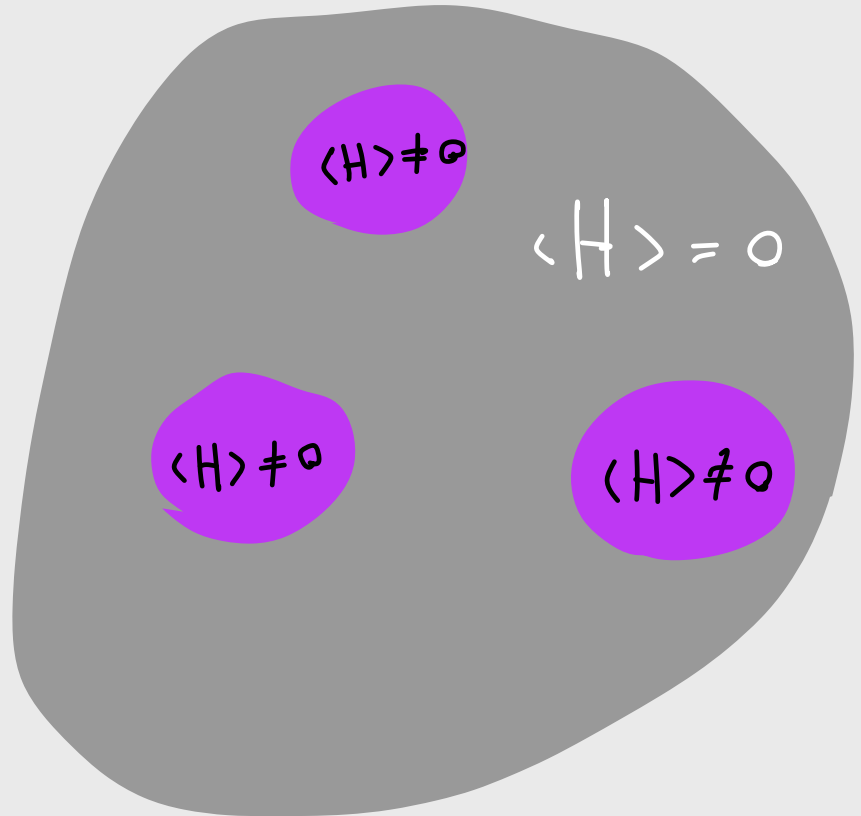
TALK TO YOU AT DINNER!

Thank you!

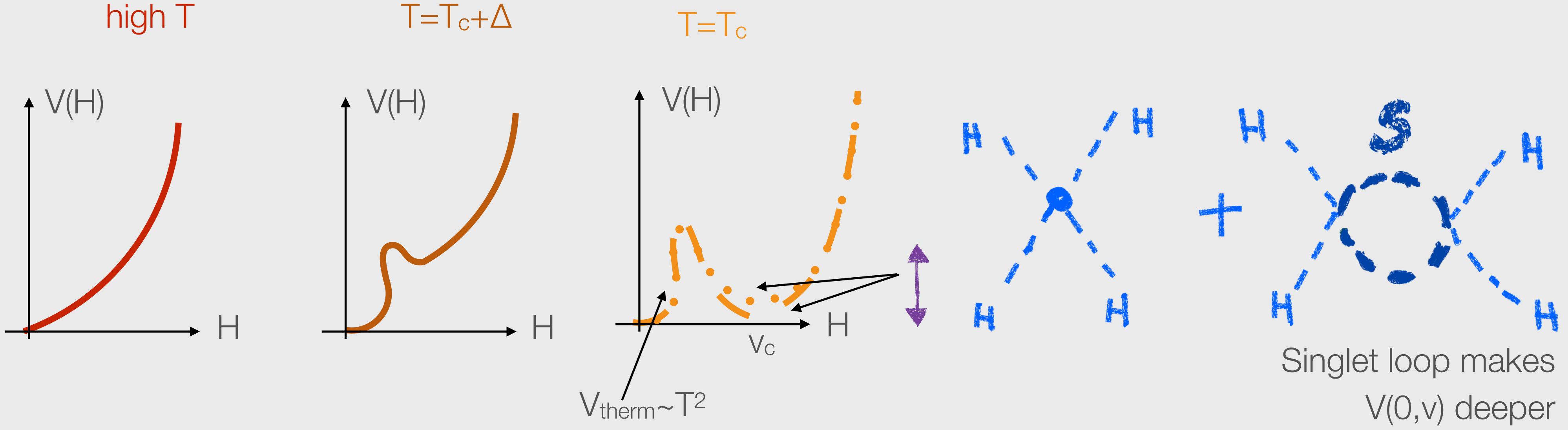
flashing concrete results for

Electroweak phase transition

Electroweak phase transition



- Modifications of the Higgs potential \Rightarrow Out of Equilibrium transition from one vacuum to a new energetically favorable one



Electroweak phase transition

- We need to study all possible new states that induce a change in the Higgs boson potential.
- For these new state to have sizable effects in the early Universe they must be light, around 1 TeV at most.
- All searches for new Higgs bosons (or general electroweak particles) probe such fundamental issue of the origin of matter in the early Universe!

$$V_{\text{therm}} \sim T^2$$

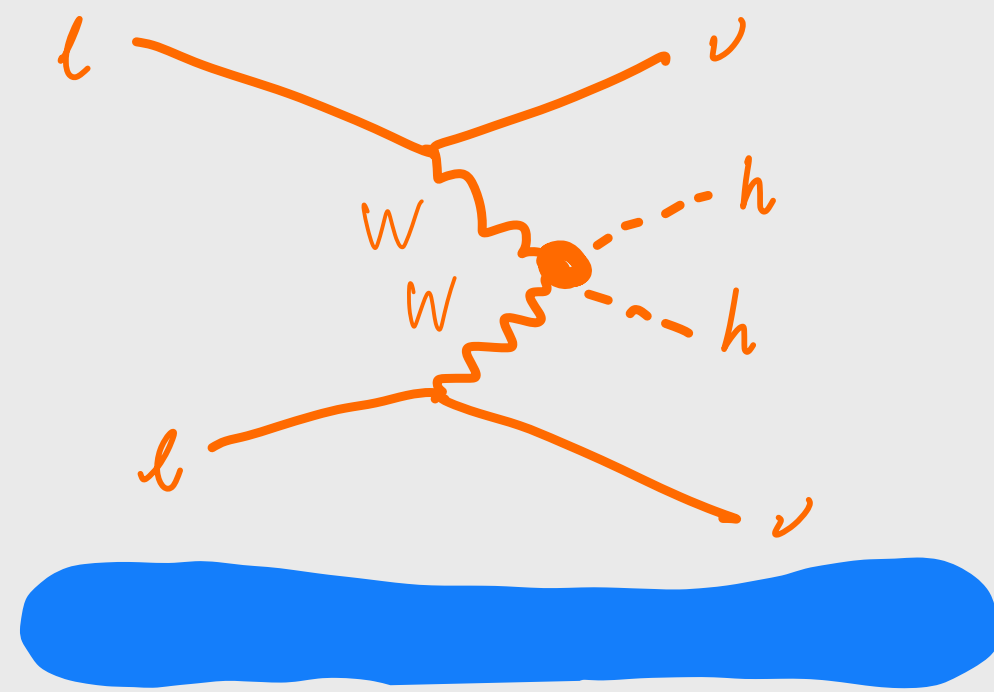
$V(0,v)$ deeper

pp or $\ell^+ \ell^- \rightarrow hh$

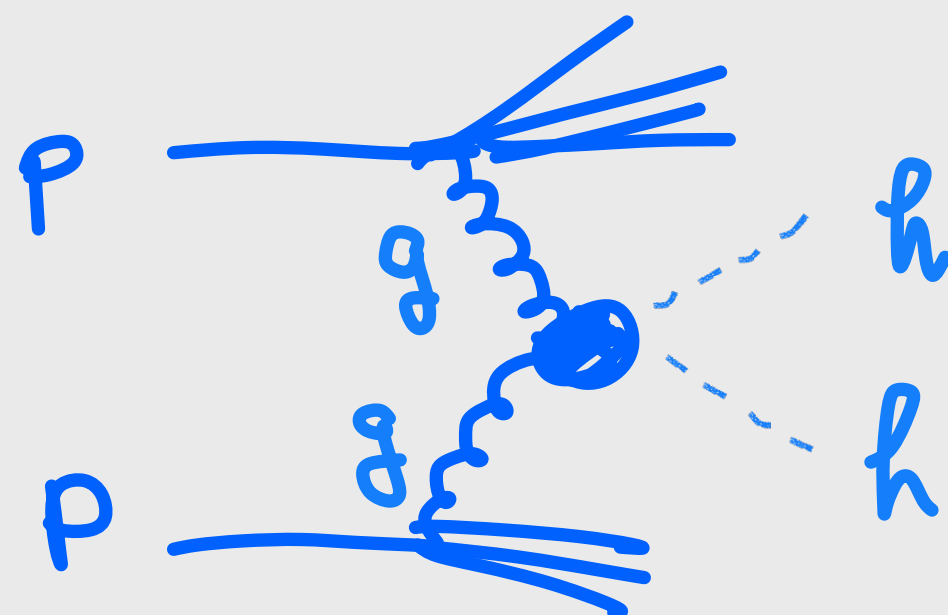
W BOSON

COLLIDER

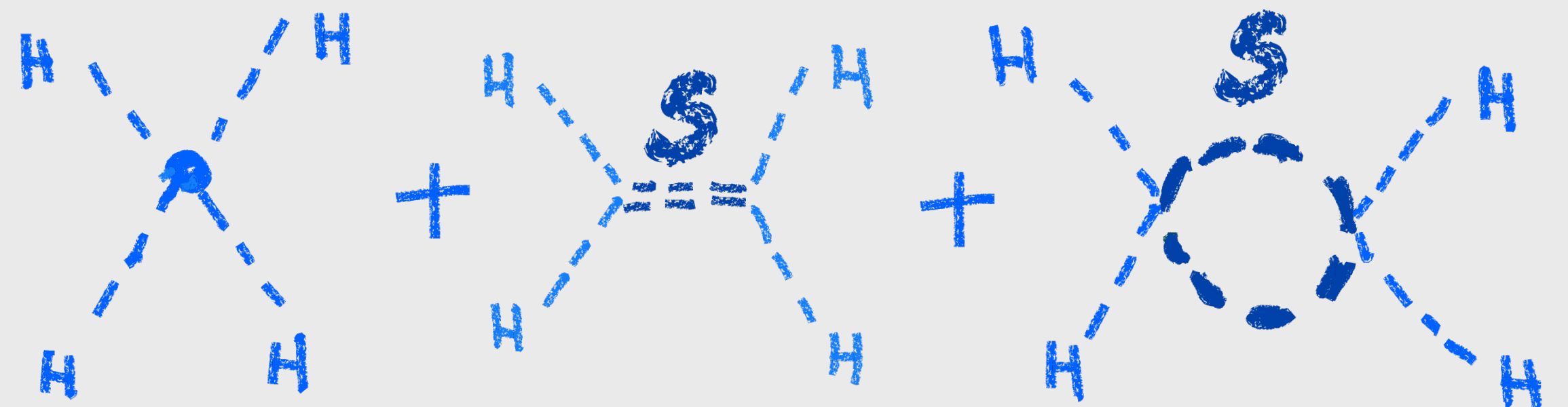
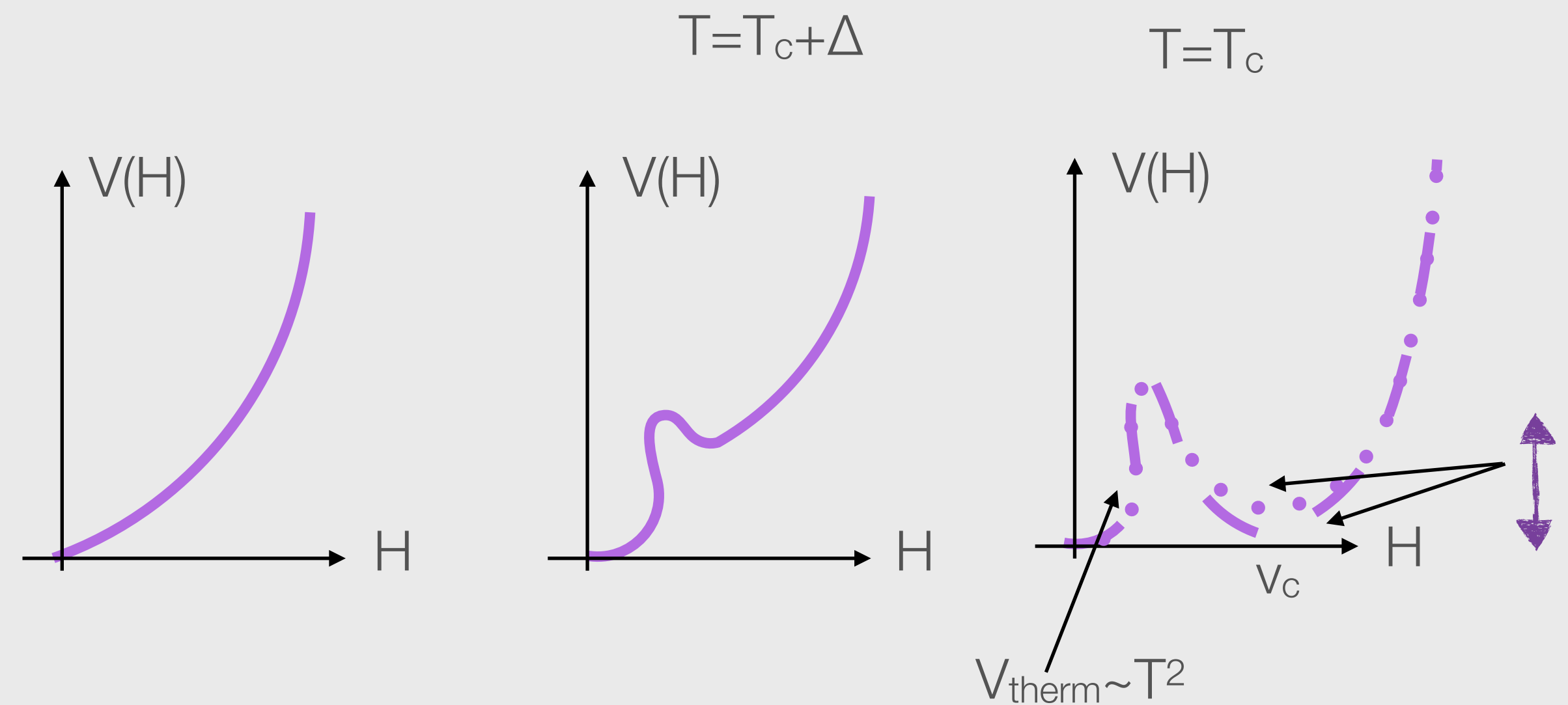
- High-Energy lepton collider has large flux of “partonic” W bosons



- gg collisions as usual



Electroweak phase transition



Singlet tree and loop makes $V(0,v)$ deeper

EW phase transition

DIRECT & INDIRECT

INTERPLAY

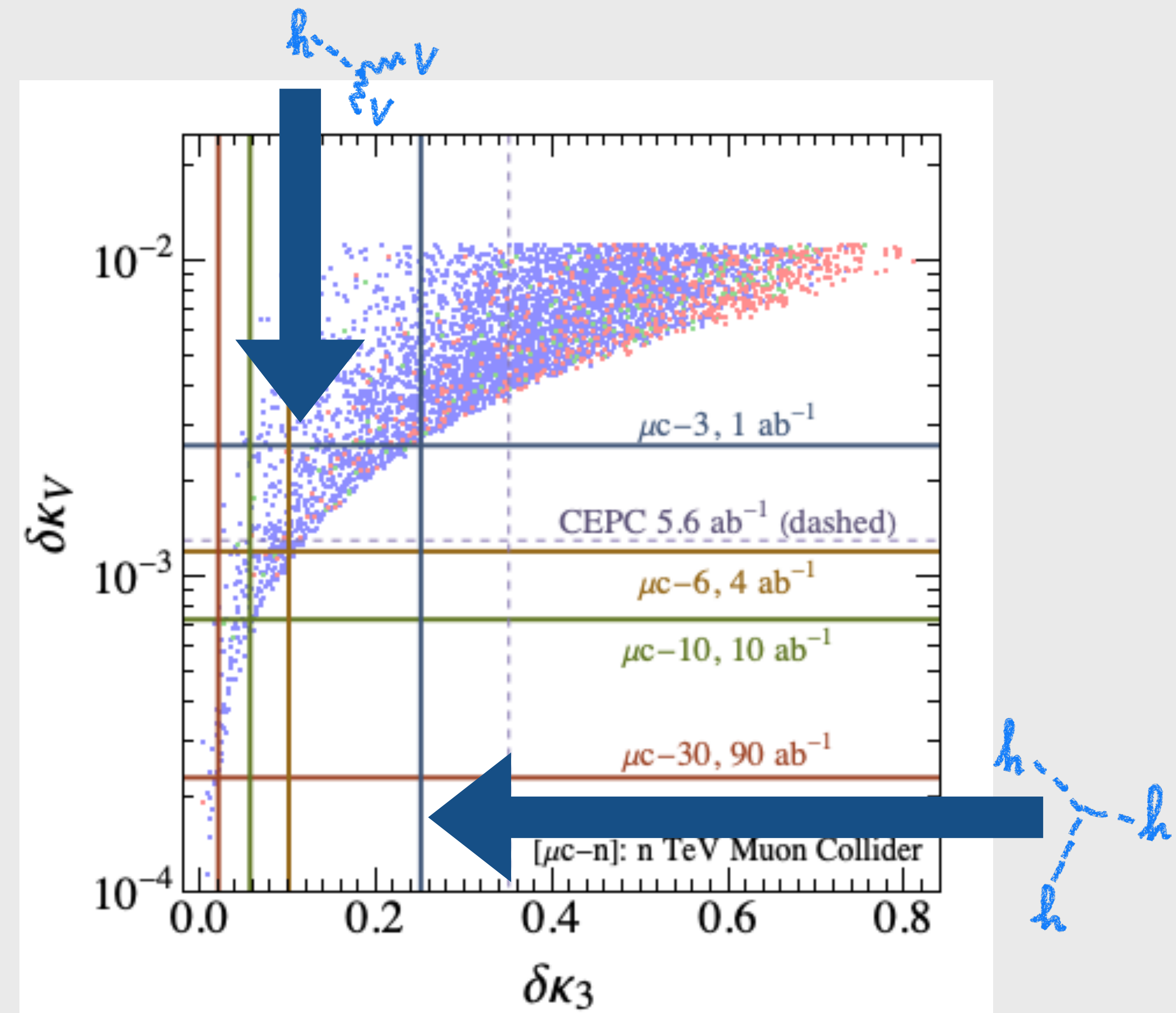
$$V(\Phi, S) = -\mu^2 (\Phi^\dagger \Phi) + \lambda (\Phi^\dagger \Phi)^2 + \frac{a_1}{2} (\Phi^\dagger \Phi) S + \frac{a_2}{2} (\Phi^\dagger \Phi) S^2 + b_1 S + \frac{b_2}{2} S^2 + \frac{b_3}{3} S^3 + \frac{b_4}{4} S^4.$$

independent parameters

$$\{M_{h_2}, \theta, v_s, b_3, b_4\}$$

strong First Order EW phase transition on all points

× ● ● → Gravity Wave SNR



EW phase transition

DIRECT & INDIRECT

INTERPLAY

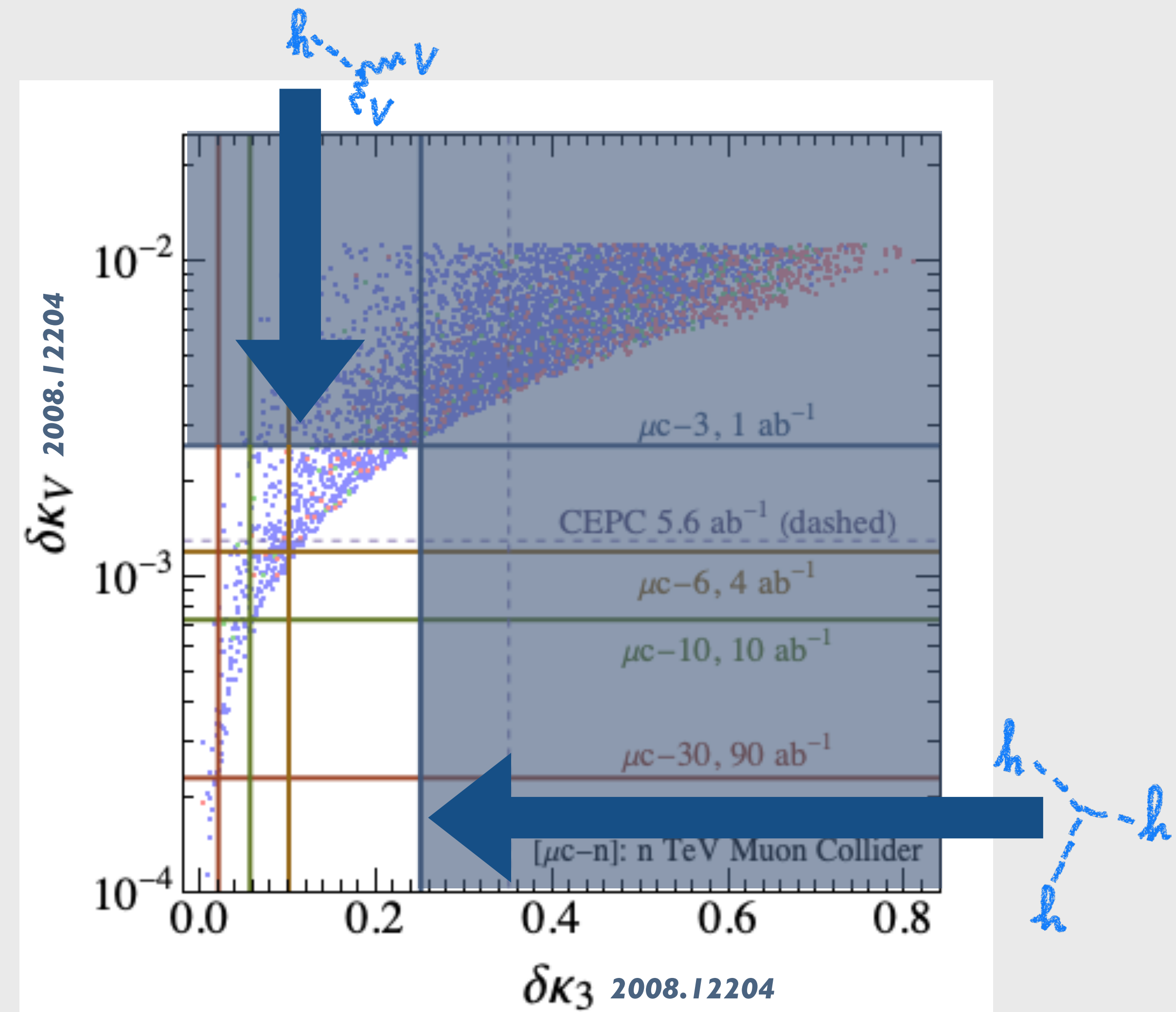
$$V(\Phi, S) = -\mu^2 (\Phi^\dagger \Phi) + \lambda (\Phi^\dagger \Phi)^2 + \frac{a_1}{2} (\Phi^\dagger \Phi) S + \frac{a_2}{2} (\Phi^\dagger \Phi) S^2 + b_1 S + \frac{b_2}{2} S^2 + \frac{b_3}{3} S^3 + \frac{b_4}{4} S^4.$$

independent parameters

$$\{M_{h_2}, \theta, v_s, b_3, b_4\}$$

strong First Order EW phase transition on all points

× ● ● → Gravity Wave SNR



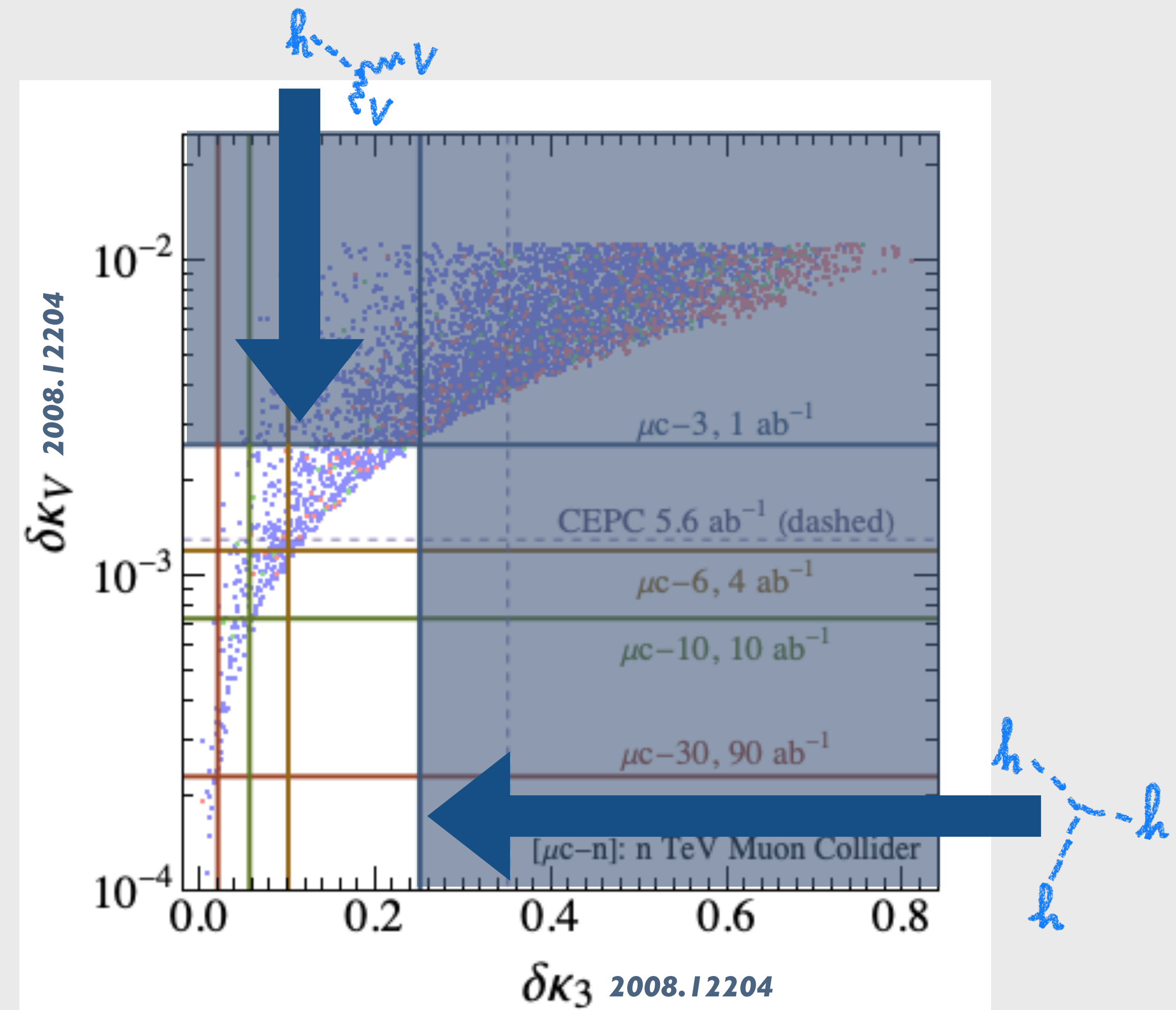
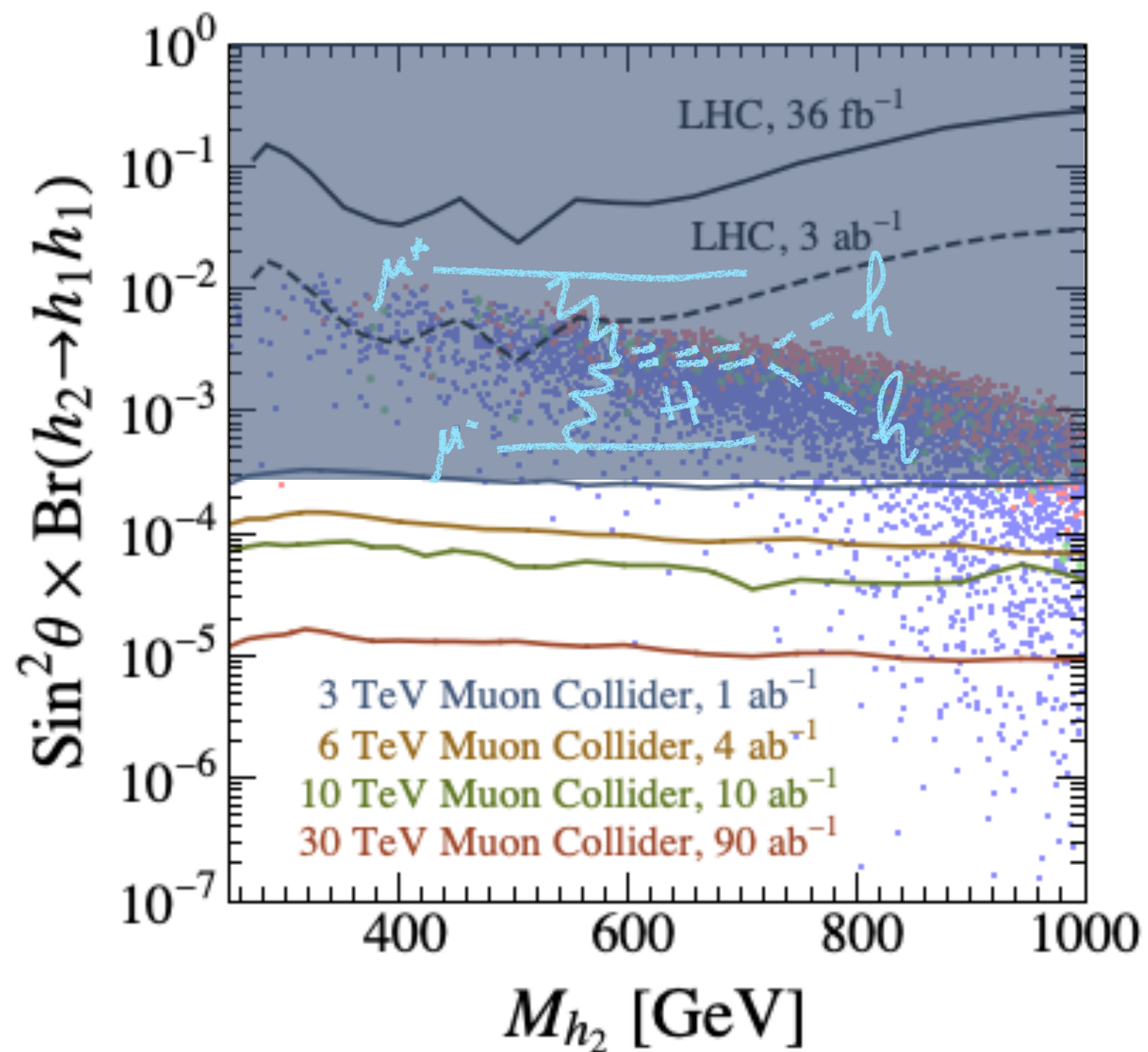
EW phase transition

strong First Order EW phase transition on all points

×
●
●
 → Gravity Wave SNR

DIRECT & INDIRECT

INTERPLAY



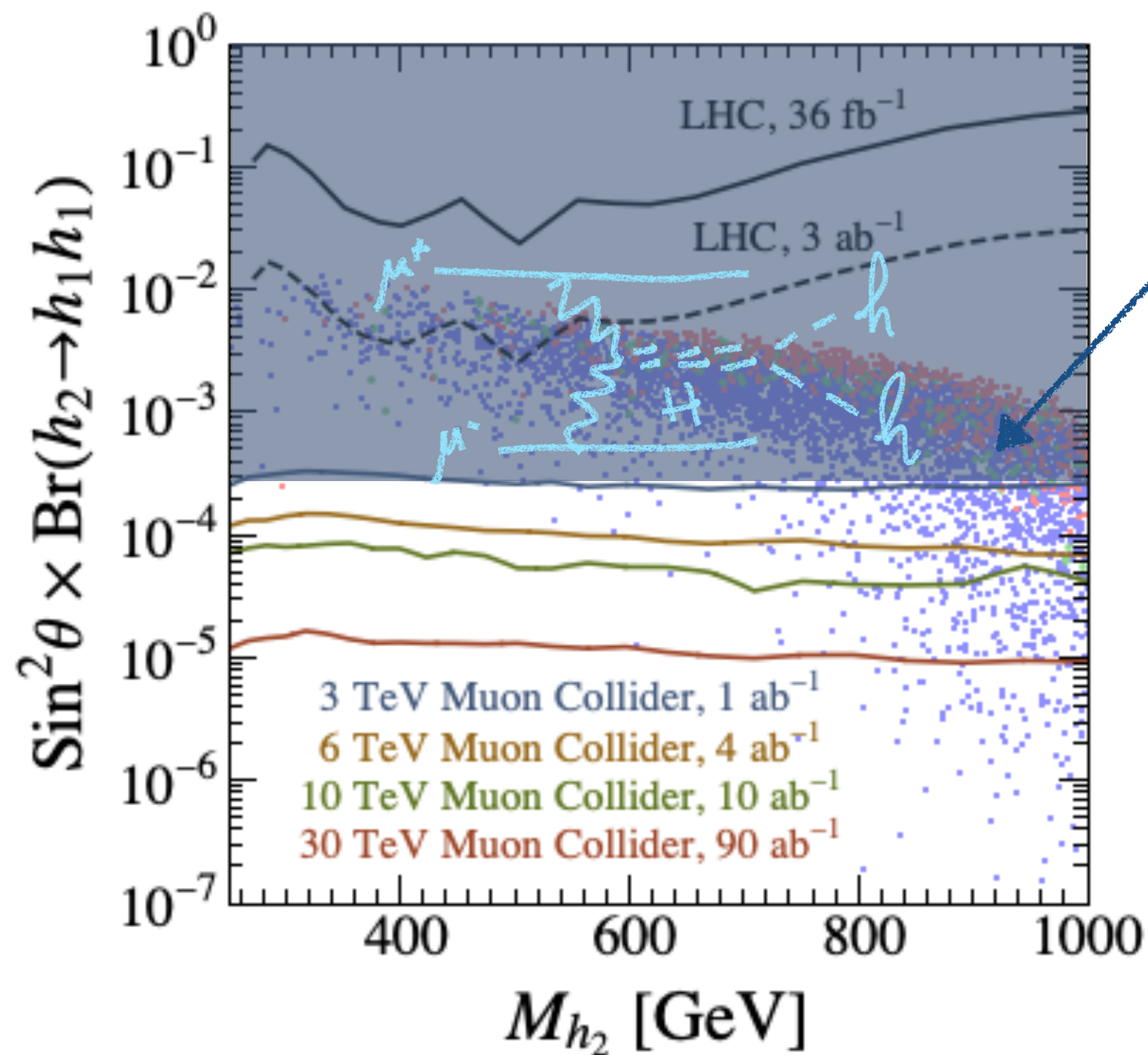
EW phase transition

strong First Order EW phase transition on all points

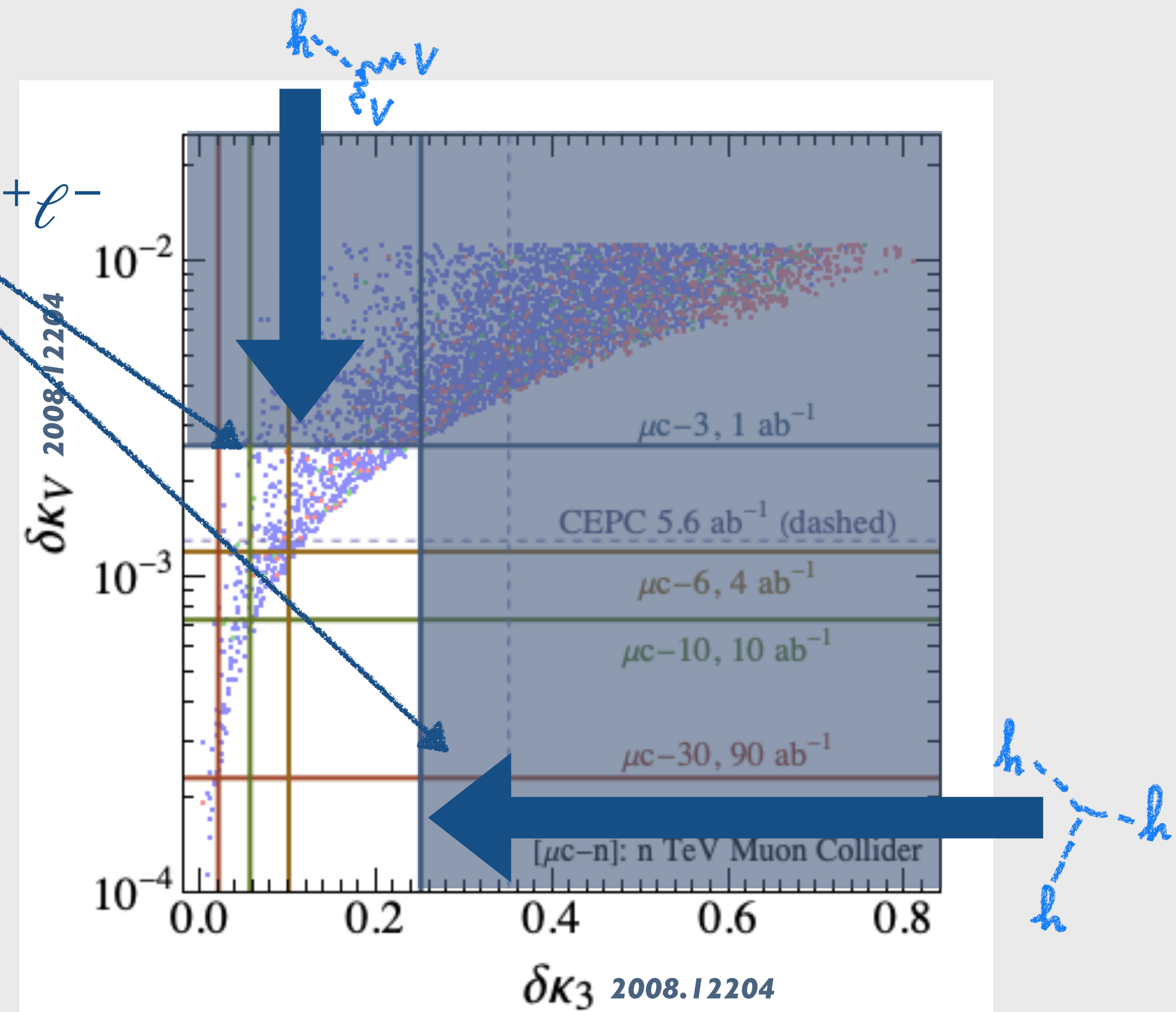
⊗ ⊙ ⊛ → Gravity Wave SNR

DIRECT & INDIRECT

INTERPLAY



3 TeV $\ell^+\ell^-$



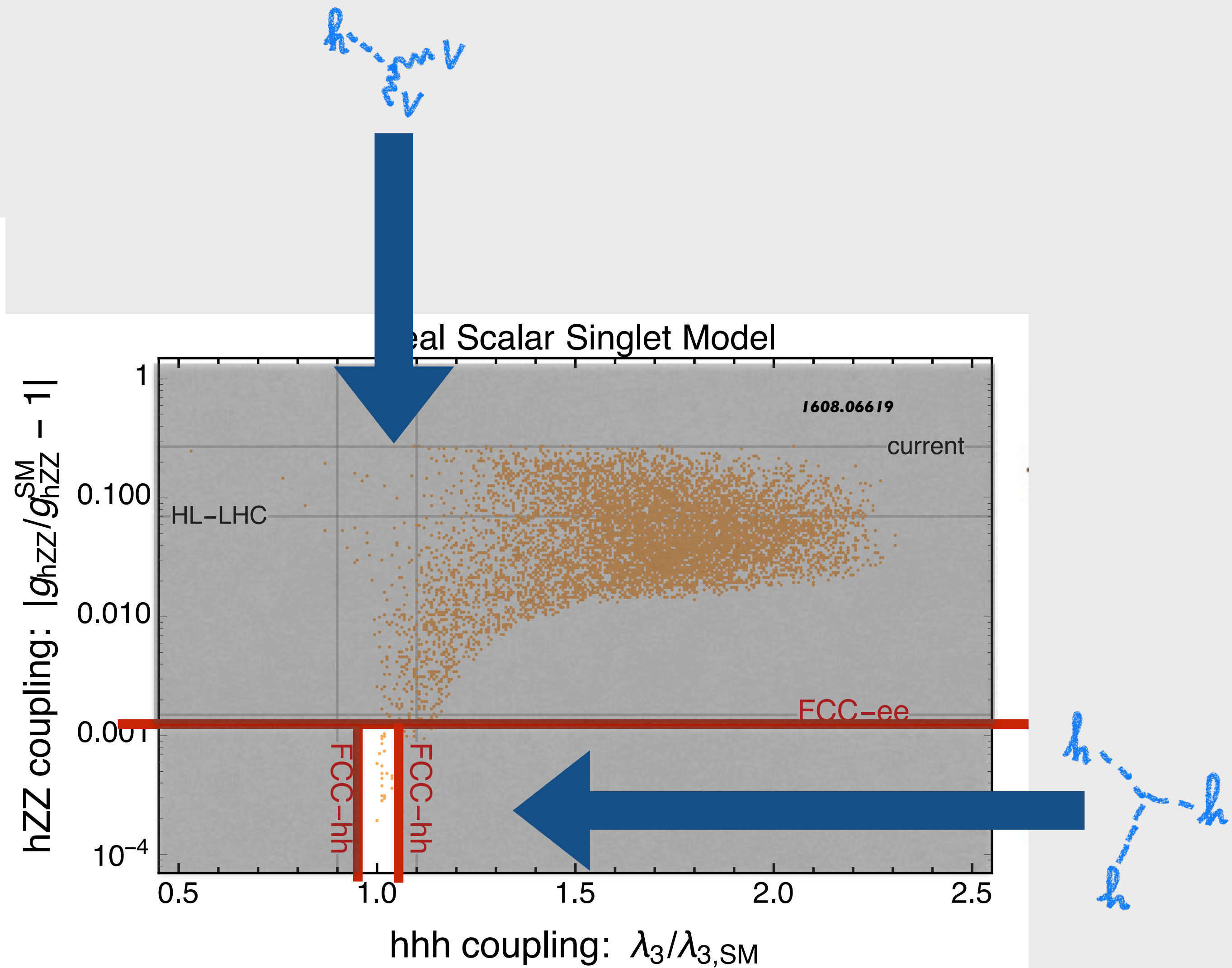
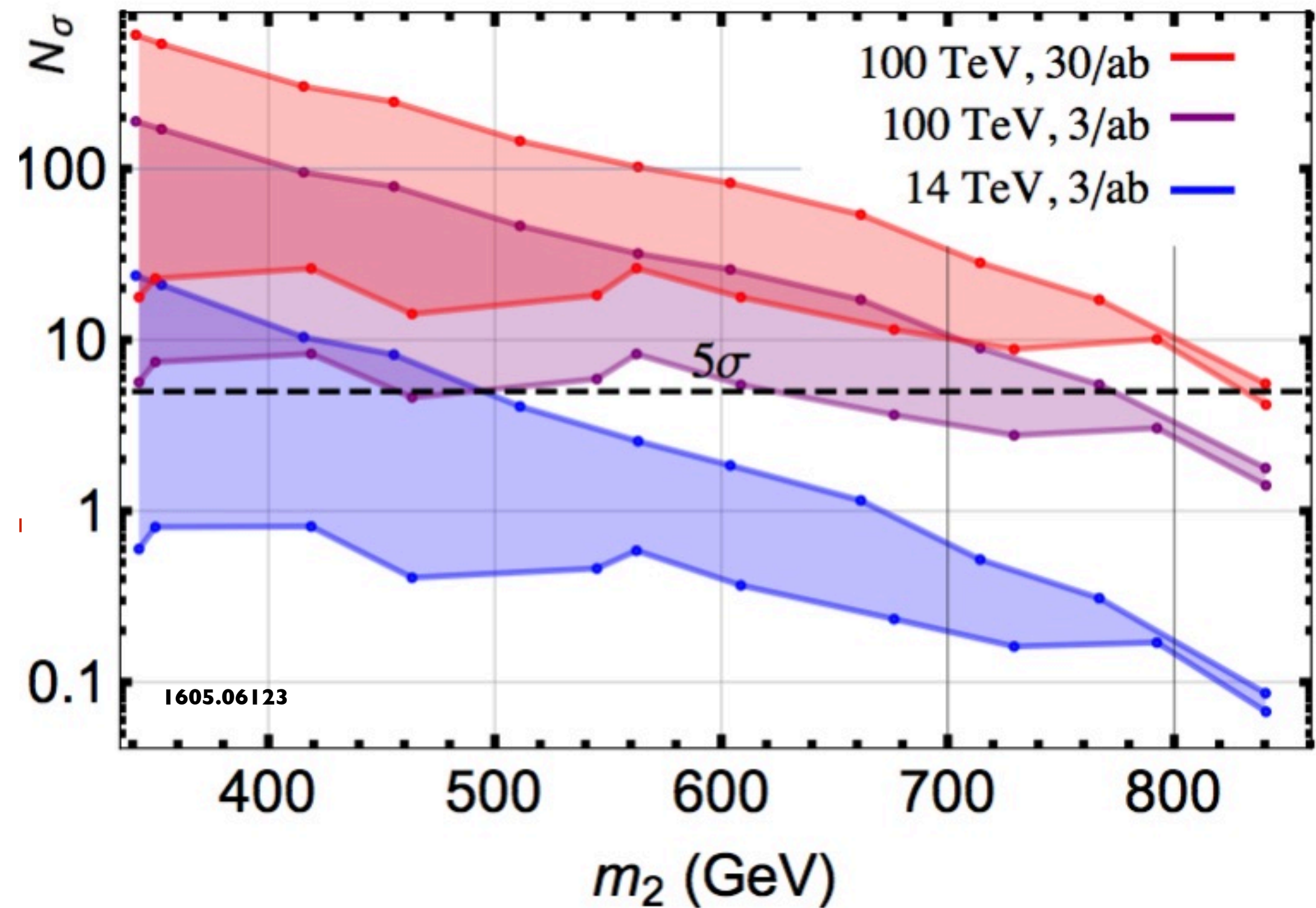
parameters space of 1st order phase transition accessible by **several measurements available at the 3 TeV $\ell^+\ell^-$ collider**

EW phase transition

DIRECT & INDIRECT

INTERPLAY

$$pp \rightarrow h_2 \rightarrow h^{(125)} h^{(125)}$$

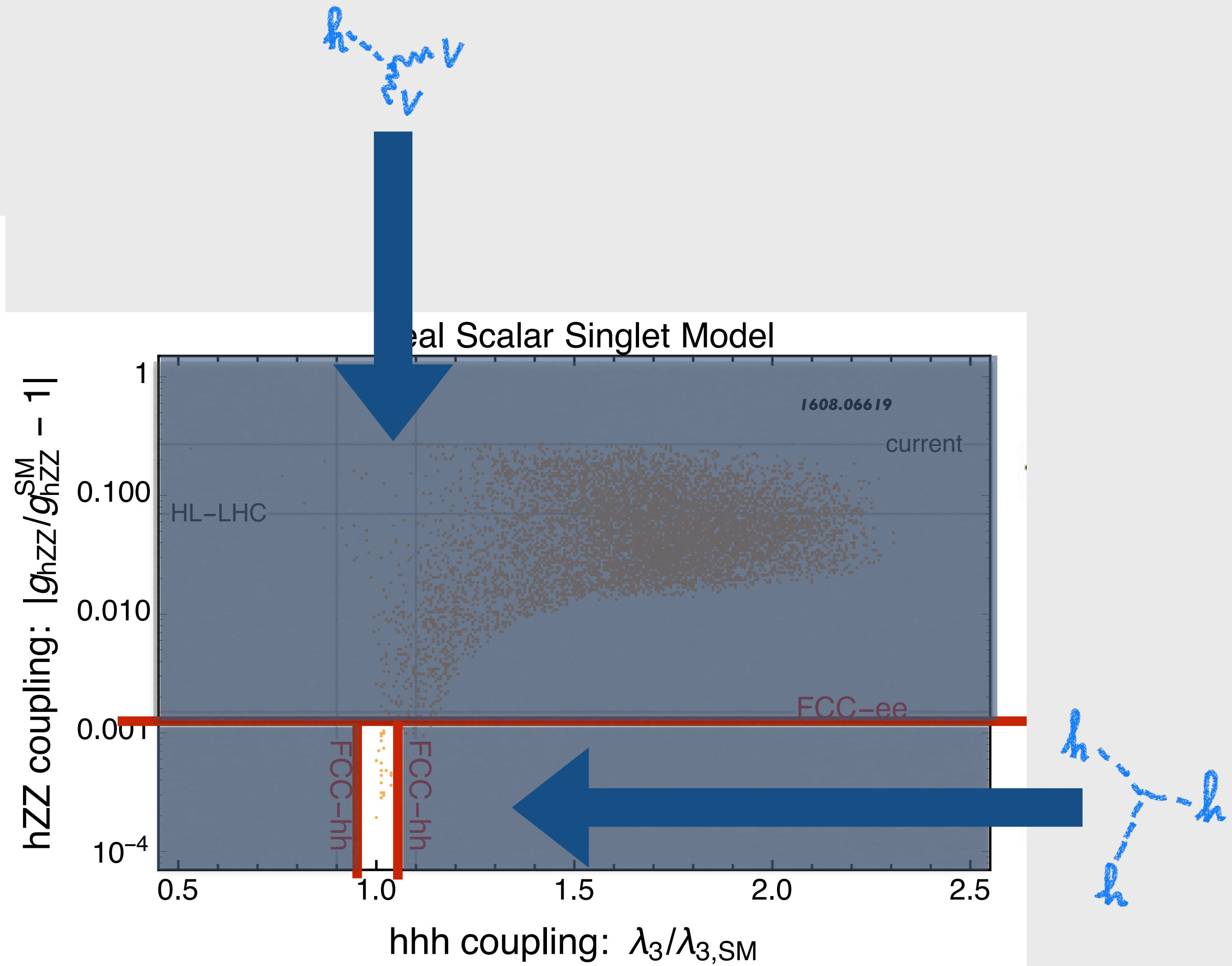
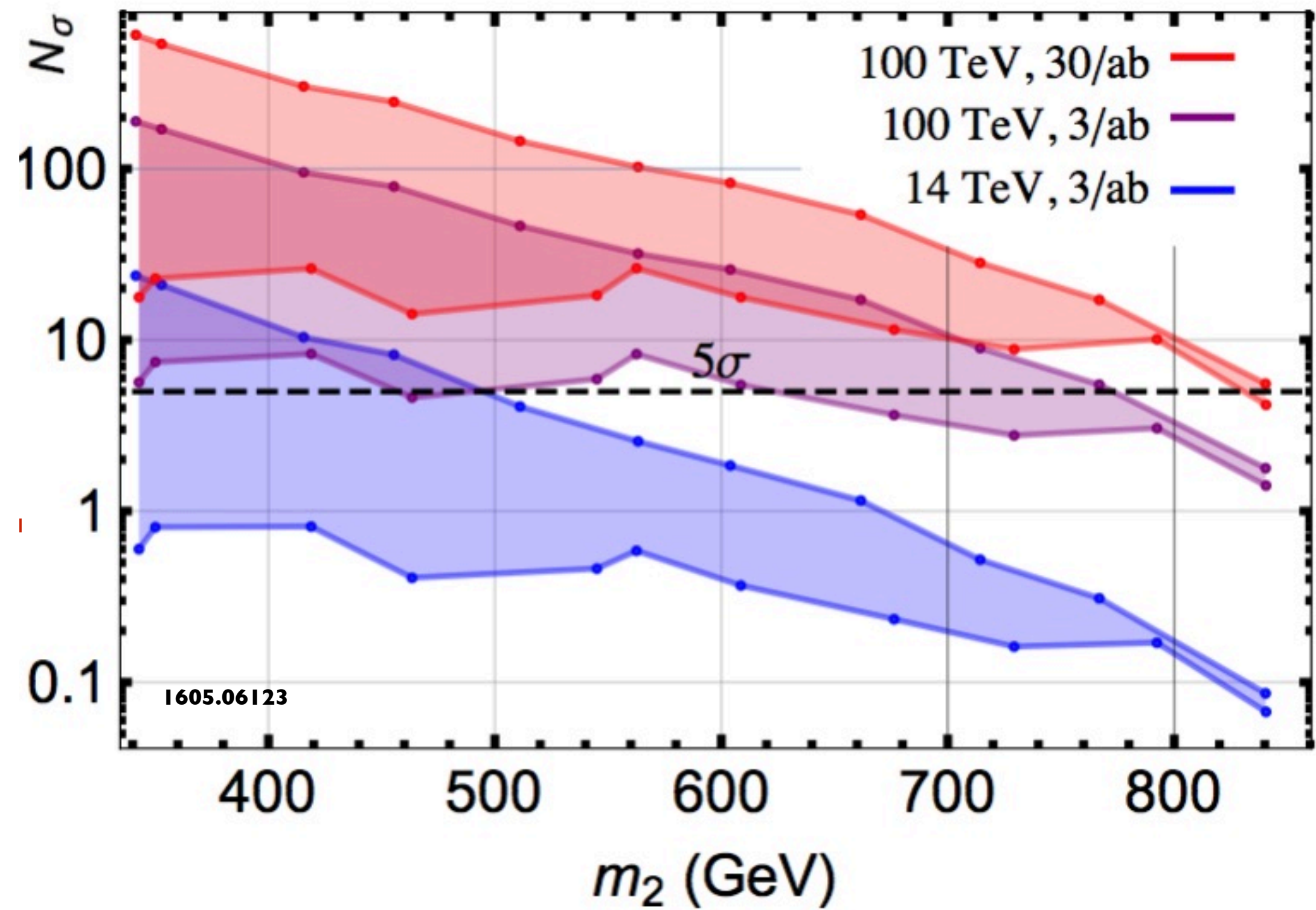


EW phase transition

DIRECT & INDIRECT

INTERPLAY

$$pp \rightarrow h_2 \rightarrow h^{(125)} h^{(125)}$$

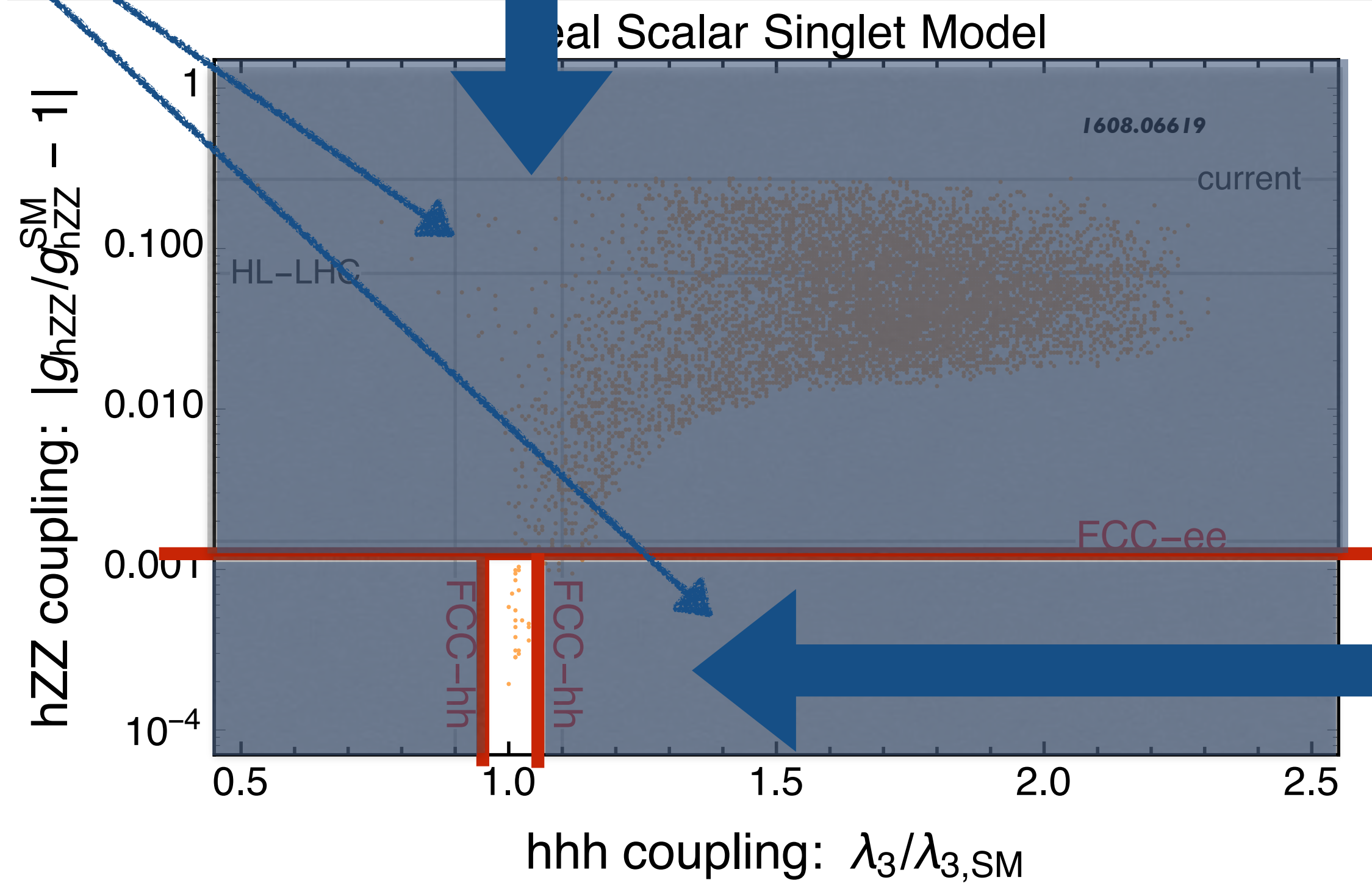
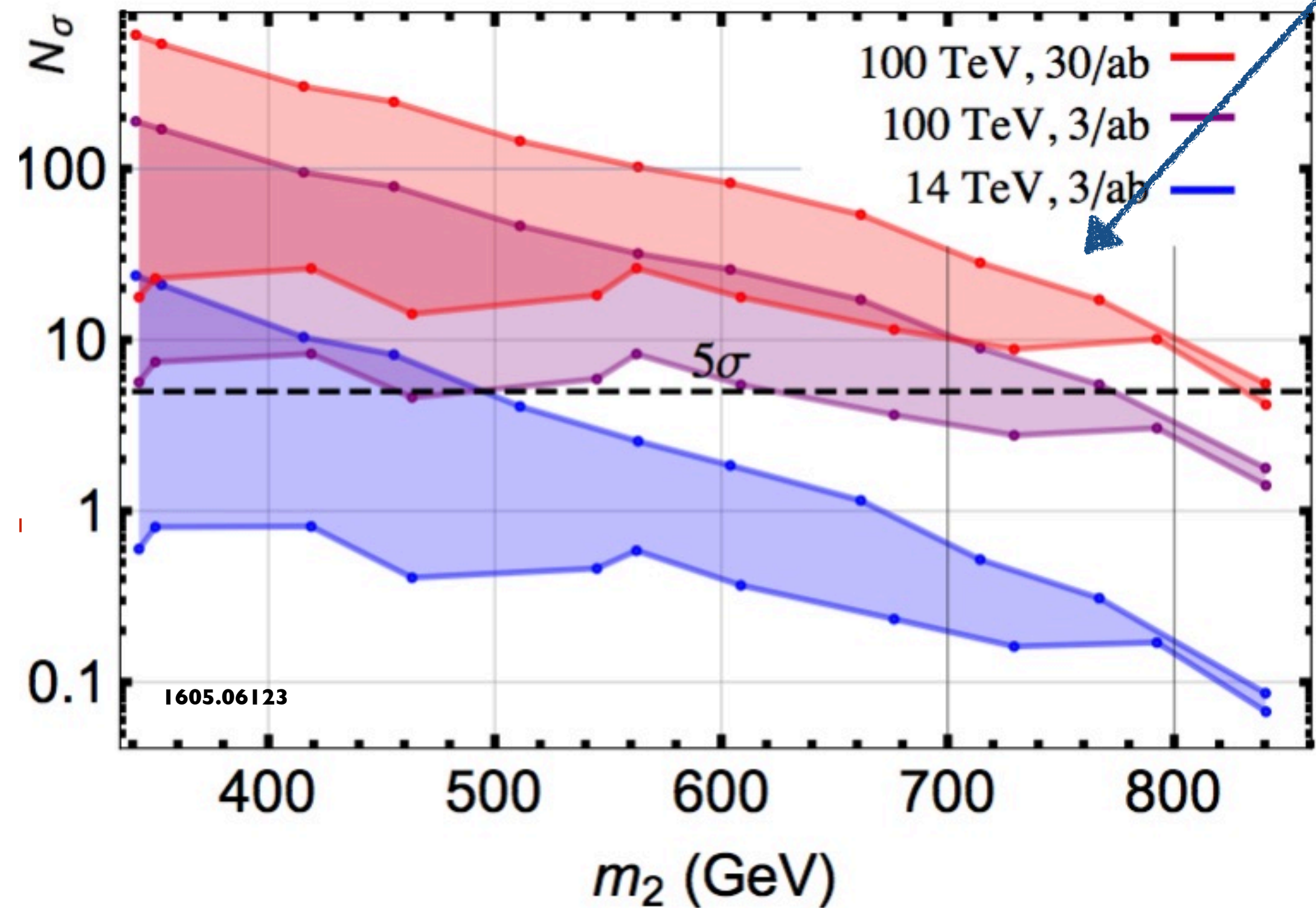


EW phase transition

DIRECT & INDIRECT

INTERPLAY

$$pp \rightarrow h_2 \rightarrow h^{(125)} h^{(125)} \quad 100 \text{ TeV } pp$$



parameters space of 1st order phase transition accessible by **several measurements available at the 100 TeV pp collider**