Quark and gluon helicity evolution at small-x: Revised and updated

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Overview

Part 1: Proton spin puzzle

- 19 May 1988
- Theoretical prediction in the 70's
- The missing spin of the proton?

Part 2: Quark flavor-singlet helicity distribution

- Generalities
- Initial framework
- Let's do it again, new dipoles

Part 1: Proton spin puzzle

Proton spin puzzle / crisis.

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The spin asymmetry in deep inelastic scattering of longitudinally polarised muons by longitudinally polarised protons has been measured over a large x range $(0.01 < x < 0.7)$. The spin-dependent structure function $g_1(x)$ for the proton has been determined and its integral over x found to be $0.114 \pm 0.012 \pm 0.026$, in disagreement with the Ellis–Jaffe sum rule. Assuming the validity of the Bjorken sum rule, this result implies a significant negative value for the integral of g_1 for the neutron. These values for the integrals of g_1 lead to the conclusion that the total quark spin constitutes a rather small fraction of the spin of the nucleon.

Reminder

$$g_1^\gamma = \frac{1}{2} \sum_q e_q^2 \; (\Delta q + \Delta \bar{q}) \,, \qquad \Delta q = q^\uparrow - q^\downarrow \text{ w.r.t. the proton spin} \tag{1}$$

and they observed for the proton

$$\int_{0.01}^{0.7} g_1(x) \; \mathrm{d}x = 0.114 \pm 0.012(\text{stat.}) \pm 0.026(\text{syst.}) \tag{2}$$

Remarks

- In blue: finite range of integration. "... *the small x region is expected to make a large contribution to the integrals.*"
- In red: Ellis-Jaffe sum rule. $\rightarrow$ Theoretical understanding of the 70's.

$\rightarrow$ How do we understand this value?

Theoretical prediction in the 70's

$\rightarrow$ How do we understand this value? $0.114 \pm 0.012(\text{stat.}) \pm 0.026(\text{syst.})$

Ellis-Jaffe sum rule, assumptions

- Sea $q\bar{q}$: $\lambda^+(x) \simeq \lambda^-(x) \simeq \bar{\lambda}^+(x) \simeq \bar{\lambda}^-(x)$
- Ansatz $\Delta s \sim 0$ (no intrinsic strangeness)
- Belief that valence quarks carry the proton spin.

we obtain[11]

$$\int_0^1 d\xi\, g_1^{ep}(\xi) = \frac{g_A}{12}(1.78), \tag{6}$$

$$\int_0^1 d\xi\, g_1^{en}(\xi) = \frac{g_A}{12}(-0.22), \tag{7}$$

where $g_A = 1.248 \pm .010$.

Ellis-Jaffe sum rule prediction (70's): $0.185 \pm 0.0015 \longrightarrow$ Not compatible with 0.114

Where is the missing spin ?

Old fundamental problem ($\sim 30y$) $\rightarrow$ looking at small number adding up to $1/2$.

- There are progresses $\rightarrow$ Still don't understand the spin of the proton in term of QCD dof.

The missing spin of the proton?

A more recent picture of the proton spin.

Spin sum rule (Jaffe Manohar decomposition [Nucl. Phys. B337, 509 (1990)])

$$\underbrace{\frac{1}{2}}_{\text{Proton spin}} = \underbrace{\frac{1}{2}\Sigma_q}_{\text{Quark spin}} + \underbrace{\Sigma_g}_{\text{Gluon spin}} + \overbrace{L_q}^{\text{Quark OAM}} + \overbrace{L_g}^{\text{Gluon OAM}}$$

Possibilities:

- Gluon spin
- Quark/Gluon angular orbital momentum

Problems?

- Gauge-dependence and non-unicity of the decomposition.
  (e.g. Ji-decomposition has no identification of gluon spin/OAM)

Large and low x region. Experiments only access a finite range of x...

$$\Sigma_q = \int_0^1 \mathrm{d}x \left(q^{\uparrow}(x) - q^{\downarrow}(x)\right) \tag{3}$$

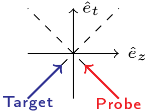
Possibilities

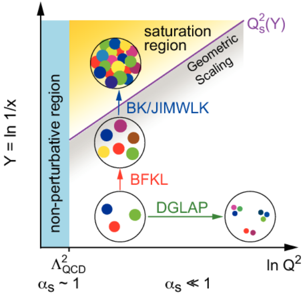
- Large-x?
- Small-x?

Part 2: Quark flavor-singlet helicity distribution

Comments on the framework (1/3)

Using lc coordinates $a^{\pm} = \frac{1}{\sqrt{2}}(a^0 \pm a^3)$
Frame choice: $\rightarrow$ Probe minus-moving, target plus-moving.





- Aim: Contribution to the spin using small-x asymptotic.
  $\longrightarrow$ Evolution in rapidity.
- Approach: Take a TMD,
  $\longrightarrow$ Simplify / Evolve / Solve.
- Equations in the spirit of BK-evolution.
  Initiated by [Kovchegov, Pitonyak, and Sievert].

Rmk: $\exists$ other frameworks for g_1 at small-x, such as Bartel, Ermolaev, and Ryskin [BER] - (1996).

Def: Wilson Lines for any irreducible representation (irrep) are

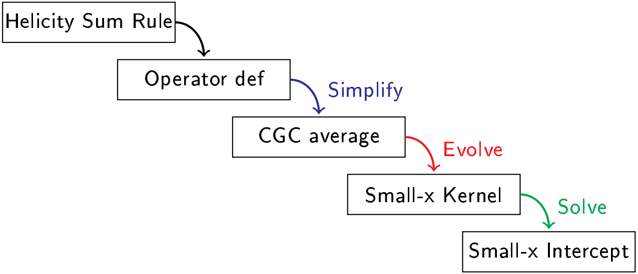
$$W^{(R)}_{\underline{x}}[b^-, a^-] \equiv \mathrm{P}\ \exp\left\{ ig \int_{a^-}^{b^-} dx^-\ t^a_R\ A^+ a(x^+ = 0, x^-, \underline{x}) \right\}. \qquad (4)$$

$\Rightarrow$ Depends only on the background field A^+ (Lorentz Gauge).

Notation: we use V for fundamental WL, and U for adjoint WL.

Yuri's (and *al.*) approach "Simplify, Evolve, and Solve"

[Y.V Kovchegov, D. Pitonyak, and M. D. Sievert 2016 2017] [Y.V Kovchegov, and M. D. Sievert 2018]



Mixed space:

- $z \equiv k^+/P^+$
- $\underline{x} = \mathbf{x} = \vec{x}_\perp$

Fourier phase

- $e^{-i\underline{k}\cdot(\underline{\zeta}-\underline{y})}$

Dipole model

- Splitting factor $\underline{x}/\underline{x}^2$

Helicity distributions (flavor-singlet)

$$g_{1L}^S(x, k_T^2) = \frac{8N_c}{(2\pi)^6} \int d^2\underline{\zeta}\, d^2\underline{w}\, d^2\underline{y}\, e^{-i\underline{k}\cdot(\underline{\zeta}-\underline{y})} \int_{\Lambda^2/s}^{1} \frac{dz}{z} \frac{\underline{\zeta}-\underline{w}}{|\underline{\zeta}-\underline{w}|^2} \cdot \frac{\underline{y}-\underline{w}}{|\underline{y}-\underline{w}|^2} G_{\underline{w},\underline{\zeta}}(zs) \tag{5}$$

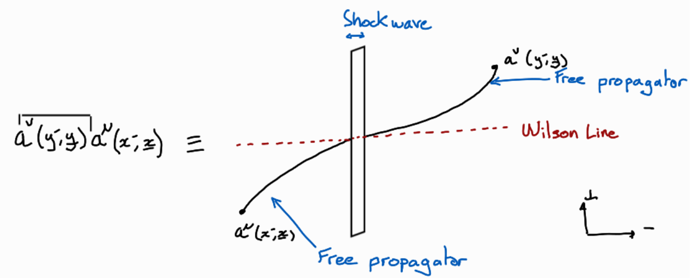
where

$$G_{\underline{w},\underline{\zeta}}(zs) = \frac{k_1^- p^+}{N_c} \mathrm{Re} \left\langle \mathrm{T\,tr}\left[V_{\underline{\zeta}}\, V_{\underline{w}}^{pol\,\dagger}\right] + \mathrm{T\,tr}\left[V_{\underline{w}}^{pol}\, V_{\underline{\zeta}}^{\dagger}\right]\right\rangle \tag{6}$$

Think of it as a regular dipole amplitude (for the moment)

Recipe:

- Split the background field A^μ into a new background A^μ and a quantum field a^μ (according to their longitudinal momentum fraction)
- Integrate out quantum fields a^μ.
- Require the propagator in the new background. Use shockwave approximation.
- Pull out the corresponding kernel for one step of evolution.



Remarks

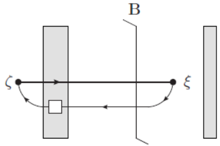
- In our case, we go beyond eikonal approximation since helicity-dependence is a genuine subeikonal effect.
- Introduce Wilson line and polarized Wilson lines, up to subeikonal level.

Situation prior to this contribution (1/3)

Consider the quark helicity TMD [Kovchegov et al. 2018]

$$g_{1L}^q(x,k_T^2) = \frac{1}{(2\pi)^3}\frac{1}{2}\sum_{S_L} S_L \int \mathrm{d}^2\underline{r}dr^- \; e^{ik\cdot r}\langle p,S_L|\bar{\psi}(0)U[0,r]\frac{\gamma^+\gamma^5}{2}\psi(r)|p,S_L\rangle. \qquad (7)$$

- Gauge-link $U[0,r]$ is process dependent, SIDIS $\rightarrow$ forward staple.
- Simplify at small-x, remaining diagram is B.



After some algebra...

$$g_{1L}^q(x,k_T^2) = -\frac{2p^+}{(2\pi)^3}\int \mathrm{d}^2\zeta \mathrm{d}^2 w \frac{\mathrm{d}^2 k_1 dk_1^-}{(2\pi)^3} e^{i(\underline{k}_1+\underline{k})\cdot(\underline{w}-\underline{\zeta})}\theta(k_1^-)\sum_{\sigma_1,\sigma_2}$$
$$\times\, \bar{v}_{\sigma_2}(k_2)\tfrac{1}{2}\gamma^+\gamma^5 v_{\sigma_1}(k_1) 2\sqrt{k_1^- k_2^-} \times \left\langle \mathrm{T} V_{\underline{\zeta}}^{ij}\left(\bar{v}_{\sigma_1}(k_1)\hat{V}_{\underline{w}}^{\dagger ji} v_{\sigma_2}(k_2)\right)\right\rangle$$
$$\times\, \frac{1}{[2k_1^- xP^+ + \underline{k}_1 - i\epsilon k_1^-][2k_1^- xP^+ + \underline{k}^2 + i\epsilon k_1^-]}\Bigg|_{k_2^-=k_1^-,\underline{k}_2=-\underline{k}} + c.c. \qquad (8)$$

Situation prior to this contribution (2/3)

The previous green operator reads

$$\left(\bar{v}_\sigma(p)\hat{V}_{\underline{x}}^\dagger v_{\sigma'}(p')\right) = 2\sqrt{p^- p'^-}\delta_{\sigma\sigma'}\left(V_{\underline{x}}^\dagger - \sigma V_{\underline{x}}^{pol\dagger} + \cdots\right). \tag{9}$$

Recall the flavor-singlet contribution simplified at small-x gives

$$g_{1L}^S(x, k_T^2) = \frac{8N_c i}{(2\pi)^5}\int \mathrm{d}^2\zeta \mathrm{d}^2\underline{w}\; e^{-i\underline{k}\cdot(\underline{\zeta}-\underline{w})}\int_{\Lambda^2/s}^{1}\frac{dz}{z}\frac{\underline{\zeta}-\underline{w}}{(\underline{\zeta}-\underline{w})^2)}\cdot\frac{\underline{k}}{\underline{k}^2}G_{\underline{w},\underline{y}}(zs). \tag{10}$$

The dipole operator $G_{\underline{w},\underline{y}}(zs)$ is

$$G_{\underline{w},\underline{y}}(zs) = \frac{k_1^- p^+}{N_c}\mathrm{Re}\left\langle \mathrm{T}\,\mathrm{tr}\left[V_{\underline{x}}V_{\underline{w}}^{pol\dagger}\right] + \mathrm{T}\,\mathrm{tr}\left[V_{\underline{w}}^{pol}V_{\underline{x}}^\dagger\right]\right\rangle, \tag{11}$$

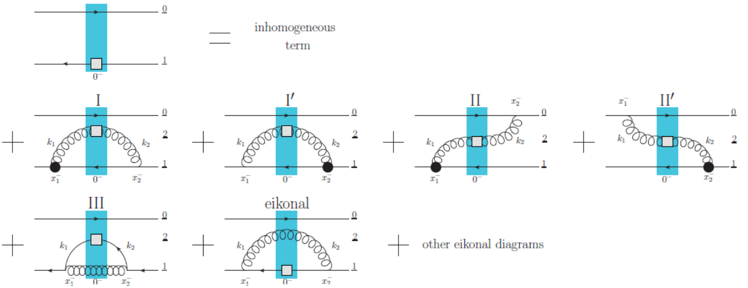
where the polarized Wilson line reads

$$\begin{aligned} V_{\underline{x}}^{pol} &= ig\frac{p^+}{s}\int dx^- V_{\underline{x}}[\infty, x^-]F^{12}V_{\underline{x}}[x^-, -\infty] \qquad (12) \\ &- g^2\frac{p^+}{s}\int dx_1^-\int_{x_1^-} dx_2^- V_{\underline{x}}[\infty, x_2^-]t^b\psi_\beta(x_2^-, \underline{x}U_{\underline{x}}^{ba}[x_2^-, x_1^-]\left[\tfrac{1}{2}\gamma^+\gamma^5\right]_{\alpha\beta}\bar{\psi}_\alpha(x_1^-, \underline{x})t^a V_{\underline{x}}[x_1^-, -\infty]. \end{aligned}$$

Remarks

- "Dressed dipoles" involve polarized WL. Obtained as sub-eikonal corrections to the scattering of a quark on a target.
- Corrections are proportional to $\sigma\delta_{\sigma\sigma'}$ in helicity basis (Brodsky-Lepage spinors in the minus direction).

Evolution (DLA, Involves the same WL at different coordinates $\longrightarrow \sigma\delta_{\sigma\sigma'}$)



Solve

Intercept in the pure glue case is $\Delta\Sigma \sim \Delta G \sim g_1 \sim (1/x)^{\alpha_h^q}$ with $\alpha_h \sim 2.31\sqrt{\frac{\alpha_s N_c}{2\pi}}$.

✗ Disagreement with BER pure glue intercept $\alpha_h \sim 3.66\sqrt{\frac{\alpha_s N_c}{2\pi}}$

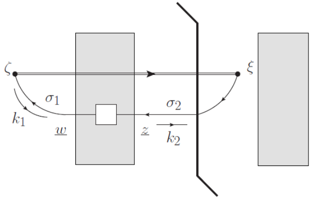
Quark flavor-singlet helicity TMD - New dipole (1/2)

Let us start again from the quark helicity TMD: [2004.11898]

$$g_{1L}^{q}(x,k_T^2) = -\frac{2p^+}{(2\pi)^3}\int \mathrm{d}^2\zeta \mathrm{d}^2 w \mathrm{d}^2 z \frac{\mathrm{d}^2 k dk^-}{(2\pi)^3} e^{i\underline{k}_1\cdot(\underline{w}-\underline{\zeta})+i\underline{k}\cdot(\underline{z}-\underline{\zeta})}\theta(k_1^-)\sum_{\sigma_1,\sigma_2}$$
$$\times \bar{v}_{\sigma_2}(k_2)\tfrac{1}{2}\gamma^+\gamma^5 v_{\sigma_1}(k_1) 2\sqrt{k_1^- k_2^-} \times \left\langle \mathrm{Ttr}\left[V_{\underline{\zeta}} V^{\dagger}_{\underline{z},\underline{w};\sigma_2,\sigma_1}\right]\right\rangle$$
$$\times \left. \frac{1}{[2k_1^- xP^+ + \underline{k}_1 - i\epsilon k_1^-][2k_1^- xP^+ + \underline{k}^2 + i\epsilon k_1^-]}\right|_{k_2^-=k_1^-,\underline{k}_2=-\underline{k}} + c.c. \qquad (13)$$

Remarks

- $V_{\underline{z},\underline{w};\sigma',\sigma}$ is the quark S-matrix for a quark-target scattering in helicity-basis.
- Allows for non locality before and after the shock wave.



Wilson lines and eikonal expansion

At sub-eikonal order:

$$
\begin{aligned}
V_{\underline{x},\underline{y};\sigma',\sigma} &= V_{\underline{x}}\,\delta^2(\underline{x}-\underline{y})\,\delta_{\sigma,\sigma'} \\
&+ \frac{i\,P^+}{s}\int_{-\infty}^{\infty} dz^-\, d^2z\; V_{\underline{x}}[\infty,z^-]\,\delta^2(\underline{x}-\underline{z})\left[-\delta_{\sigma,\sigma'}\overleftarrow{D}^i\,D^i + g\,\sigma\,\delta_{\sigma,\sigma'}\,F^{12}\right](z^-,\underline{z})\,V_{\underline{y}}[z^-,-\infty]\,\delta^2(\underline{y}-\underline{z}) \\
&- \frac{g^2 P^+}{2\,s}\delta^2(\underline{x}-\underline{y})\int_{-\infty}^{\infty} dz_1^- \int_{z_1^-}^{\infty} dz_2^-\, V_{\underline{x}}[\infty,z_2^-]\,t^b\,\psi_\beta(z_2^-,\underline{x})\,U_{\underline{x}}^{ba}[z_2^-,z_1^-]\left[\delta_{\sigma,\sigma'}\,\gamma^+ - \sigma\,\delta_{\sigma,\sigma'}\,\gamma^+\gamma^5\right]_{\alpha\beta} \\
&\qquad \times \bar{\psi}_\alpha(z_1^-,\underline{x})\,t^a\,V_{\underline{x}}[z_1^-,-\infty],
\end{aligned}
\tag{14}
$$

Remarks

- Blue $\longrightarrow$ Already used in previous V^{pol}. Label *of the first kind*; notation $V^{pol\,[1]}$. Proportional to $\sigma\delta_{\sigma\sigma'}$.
- Red $\longrightarrow$ "NEW" (in our framework). Label *of the second kind*; notation $V^{pol\,[2]}$. Proportional to $\delta_{\sigma\sigma'}$.

Picture?

For the quark S-matrix at sub eikonal order, see also:

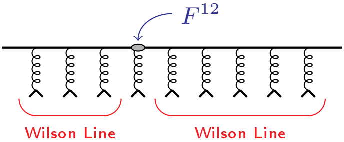
- Balitsky and Tarasov, e.g. [1505.02151]
- Chirilli, e.g. [1807.11435]
- Altinoluk et al., e.g. [2012.03886]
- Kovchegov et al., e.g. [1808.09010] [2108.03667]

Polarized WL,

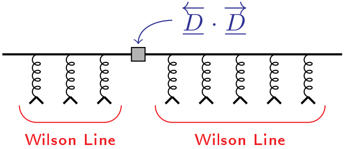
$$V^{\text{pol}[1]}_{\underline{x}} = \underbrace{V^{\text{G}[1]}_{\underline{x}} + V^{\text{q}[1]}_{\underline{x}}}_{\sigma\,\delta_{\sigma\sigma'}}, \quad V^{\text{pol}[2]}_{\underline{x},\underline{y}} = \underbrace{V^{\text{G}[2]}_{\underline{x},\underline{y}} + V^{\text{q}[2]}_{\underline{x}}\,\delta^2(\underline{x}-\underline{y})}_{\delta_{\sigma\sigma'}}.$$

can be represented as

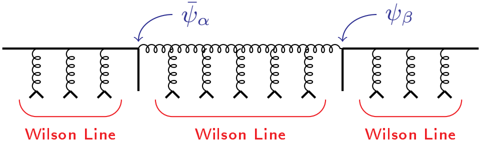
$V^{G[1]}_{\underline{x}} =$



$V^{G[2]}_{\underline{x}} =$



$V^{q[1/2]}_{\underline{x}} =$



Contraction with $(\gamma^+\gamma^5)_{\alpha\beta} \times \sigma\delta_{\sigma\sigma'}$ or $\gamma^+_{\alpha\beta} \times \delta_{\sigma\sigma'}$

Quark flavor-singlet helicity TMD - New dipole (2/2)

Simplified at small-x, the quark flavor-singlet helicity TMD reads

$$g_{1L}^S(x,k_T^2) = \frac{8\,N_c\,N_f}{(2\pi)^5} \int\limits_{\Lambda^2/s}^{1} \frac{dz}{z} \int d^2x_{10}\, e^{i\underline{k}\cdot\underline{x}_{10}} \left[i\,\frac{\underline{x}_{10}}{x_{10}^2} \cdot \frac{\underline{k}}{\underline{k}^2} \,\left[Q(x_{10}^2,zs) + G_2(x_{10}^2,zs)\right] - \frac{(\underline{k}\times\underline{x}_{10})^2}{\underline{k}^2\,x_{10}^2}\, G_2(x_{10}^2,zs) \right], \qquad (15)$$

The new dipole G_2 is defined with

$$G_{10}^j(zs) \equiv \frac{1}{2N_c} \left\langle\!\!\left\langle \mathrm{tr}\left[V_{\underline{\zeta}}^\dagger\, V_{\underline{\xi}}^{j\;\mathrm{G}[2]} + \left(V_{\underline{\xi}}^{j\;\mathrm{G}[2]}\right)^\dagger V_{\underline{\zeta}} \right] \right\rangle\!\!\right\rangle \qquad (16)$$

$$\int d^2\left(\frac{x_1+x_0}{2}\right)\, G_{10}^i(zs) = (x_{10})_\perp^i\, G_1(x_{10}^2,zs) + \epsilon^{ij}\,(x_{10})_\perp^j\, G_2(x_{10}^2,zs). \qquad (17)$$

Remarks

- Dependence on previously used dipole $Q(x_{10}^2,zs)$.
- The previously missing dependence is proportional to $G_2(x_{10}^2,zs)$.
- New contribution depends on the sub-eikonal operator $\overleftrightarrow{\underline{D}}$, related to the Jaffe-Manohar polarized gluon distribution.

Evolution, revised and updated

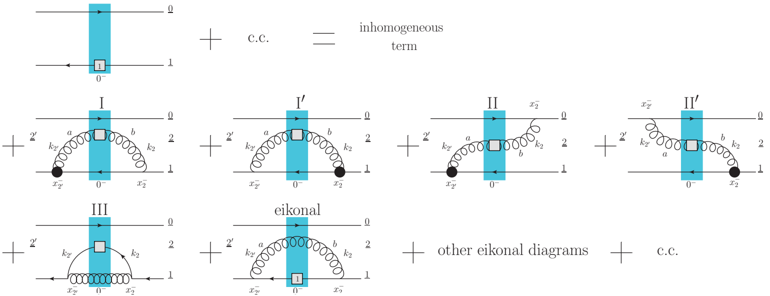
One step of evolution reads the formal form

$$\hat{\mathcal{O}}_i = \hat{\mathcal{O}}_i^{(0)} + \sum_j \mathcal{K}_{ij} \otimes \hat{\mathcal{O}}_j \tag{18}$$

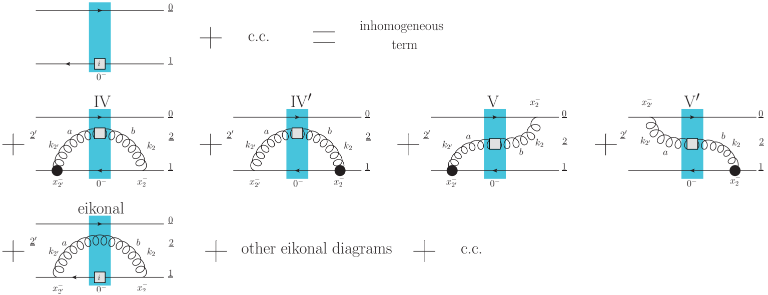
- Mixing to operators involving Wilson lines of first and/or second kind.
- Kernel involves transverse and longitudinal logarithmic integrals. The evolution is DLA, as opposed to the unpolarized one being SLA.
- Lifetime ordering is explicit $\theta(z\underline{x}_{10}^2 - z'\underline{x}_{21}^2)$.
- Similar to the Balitsky hierarchy, equation are not closed.
- Can be closed in the 't Hooft large N_c-limit or Veneziano large $N_c\&N_f$-limit.

Results

- In the pure glue sector, the intercept becomes $\alpha_h \sim 3.66\sqrt{\frac{\alpha_s N_c}{2\pi}}$. In complete agreement with BER result.
- Iterating this kernel, one recover the small-x spin-dependent DGLAP kernel.



$$
\begin{aligned}
&\frac{1}{2N_c}\boxed{\left\langle\!\!\left\langle \mathrm{tr}\left[V_{\underline{0}}\, V_{\underline{1}}^{\mathrm{pol}[1]\,\dagger}\right] + \mathrm{c.c.}\right\rangle\!\!\right\rangle(zs)} = \frac{1}{2N_c}\left\langle\!\!\left\langle \mathrm{tr}\left[V_{\underline{0}}\, V_{\underline{1}}^{\mathrm{pol}[1]\,\dagger}\right] + \mathrm{c.c.}\right\rangle\!\!\right\rangle_0(zs) \\
&+\frac{\alpha_s N_c}{2\pi^2}\int\limits_{\Lambda^2/s}^{z}\frac{dz'}{z'}\int d^2x_2\left\{\left[\frac{1}{x_{21}^2}-\frac{\underline{x}_{21}}{x_{21}^2}\cdot\frac{\underline{x}_{20}}{x_{20}^2}\right]\frac{1}{N_c^2}\boxed{\left\langle\!\!\left\langle \mathrm{tr}\left[t^b V_{\underline{0}} t^a V_{\underline{1}}^{\dagger}\right]\left(U_{\underline{2}}^{\mathrm{pol}[1]}\right)^{ba} + \mathrm{c.c.}\right\rangle\!\!\right\rangle(z's)}\right. \\
&\left.+\left[2\frac{\epsilon^{ij}\,x_{21}^j}{x_{21}^4}-\frac{\epsilon^{ij}\,(x_{20}^j+x_{21}^j)}{x_{20}^2\,x_{21}^2}-\frac{2\,\underline{x}_{20}\times\underline{x}_{21}}{x_{20}^2\,x_{21}^2}\left(\frac{x_{21}^i}{x_{21}^2}-\frac{x_{20}^i}{x_{20}^2}\right)\right]\frac{1}{N_c^2}\boxed{\left\langle\!\!\left\langle \mathrm{tr}\left[t^b V_{\underline{0}} t^a V_{\underline{1}}^{\dagger}\right]\left(U_{\underline{2}}^{i\,\mathrm{G}[2]}\right)^{ba} + \mathrm{c.c.}\right\rangle\!\!\right\rangle(z's)}\right\} \\
&+\frac{\alpha_s N_c}{4\pi^2}\int\limits_{\Lambda^2/s}^{z}\frac{dz'}{z'}\int\frac{d^2x_2}{x_{21}^2}\left\{\frac{1}{N_c^2}\boxed{\left\langle\!\!\left\langle \mathrm{tr}\left[t^b\, V_{\underline{0}}\, t^a\, V_{\underline{2}}^{\mathrm{pol}[1]\,\dagger}\right]\,U_{\underline{1}}^{ba}\right\rangle\!\!\right\rangle(z's)}+2\,\frac{\epsilon^{ij}\,x_{21}^j}{x_{21}^2}\,\frac{1}{N_c^2}\boxed{\left\langle\!\!\left\langle \mathrm{tr}\left[t^b\, V_{\underline{0}}\, t^a\, V_{\underline{2}}^{i\,\mathrm{G}[2]\,\dagger}\right]\,U_{\underline{1}}^{ba}\right\rangle\!\!\right\rangle(z's)}+\mathrm{c.c.}\right\} \\
&+\frac{\alpha_s N_c}{2\pi^2}\int\limits_{\Lambda^2/s}^{z}\frac{dz'}{z'}\int d^2x_2\,\frac{x_{10}^2}{x_{21}^2\,x_{20}^2}\left\{\frac{1}{N_c^2}\boxed{\left\langle\!\!\left\langle \mathrm{tr}\left[t^b\, V_{\underline{0}}\, t^a\, V_{\underline{1}}^{\mathrm{pol}[1]\,\dagger}\right]\,U_{\underline{2}}^{ba}\right\rangle\!\!\right\rangle(z's)}-\frac{C_F}{N_c^2}\boxed{\left\langle\!\!\left\langle \mathrm{tr}\left[V_{\underline{0}}\, V_{\underline{1}}^{\mathrm{pol}[1]\,\dagger}\right]\right\rangle\!\!\right\rangle(z's)}+\mathrm{c.c.}\right\}.
\end{aligned}
\tag{95}
$$



$$
\begin{aligned}
&\frac{1}{2N_c}\boxed{\left\langle\!\!\left\langle \mathrm{tr}\left[V_{\underline{0}}\, V_{\underline{1}}^{i\,\mathrm{G}[2]\,\dagger}\right] + \mathrm{c.c.}\right\rangle\!\!\right\rangle(zs)} = \frac{1}{2N_c}\left\langle\!\!\left\langle \mathrm{tr}\left[V_{\underline{0}}\, V_{\underline{1}}^{i\,\mathrm{G}[2]\,\dagger}\right] + \mathrm{c.c.}\right\rangle\!\!\right\rangle_0(zs) \qquad (106)\\
&+\frac{\alpha_s N_c}{4\pi^2}\int\limits_{\frac{\Lambda^2}{s}}^{z}\frac{dz'}{z'}\int d^2x_2\Bigg\{\left[\frac{\epsilon^{ij}x_{21}^j}{x_{21}^2}-\frac{\epsilon^{ij}x_{20}^j}{x_{20}^2}+2x_{21}^i\,\frac{\underline{x}_{21}\times\underline{x}_{20}}{x_{21}^2\,x_{20}^2}\right]\frac{1}{N_c^2}\boxed{\left\langle\!\!\left\langle \mathrm{tr}\left[t^b V_{\underline{0}} t^a V_{\underline{1}}^{\dagger}\right]\left(U_{\underline{2}}^{\mathrm{pol}[1]}\right)^{ba} + \mathrm{c.c.}\right\rangle\!\!\right\rangle(z's)}\\
&+\left[\delta^{ij}\left(\frac{3}{x_{21}^2}-2\frac{\underline{x}_{20}\cdot\underline{x}_{21}}{x_{20}^2\,x_{21}^2}-\frac{1}{x_{20}^2}\right)-2\frac{x_{21}^i\,x_{20}^j}{x_{21}^2\,x_{20}^2}\left(2\frac{\underline{x}_{20}\cdot\underline{x}_{21}}{x_{20}^2}+1\right)+2\frac{x_{21}^i\,x_{21}^j}{x_{21}^2\,x_{20}^2}\left(2\frac{\underline{x}_{20}\cdot\underline{x}_{21}}{x_{21}^2}+1\right)+2\frac{x_{20}^i\,x_{20}^j}{x_{20}^4}-2\frac{x_{21}^i\,x_{21}^j}{x_{21}^4}\right]\\
&\times\frac{1}{N_c^2}\boxed{\left\langle\!\!\left\langle \mathrm{tr}\left[t^b V_{\underline{0}} t^a V_{\underline{1}}^{\dagger}\right]\left(U_{\underline{2}}^{j\,\mathrm{G}[2]}\right)^{ba} + \mathrm{c.c.}\right\rangle\!\!\right\rangle(z's)}\Bigg\}\\
&+\frac{\alpha_s N_c}{2\pi^2}\int\limits_{\frac{\Lambda^2}{s}}^{z}\frac{dz'}{z'}\int d^2x_2\,\frac{x_{10}^2}{x_{21}^2\,x_{20}^2}\left\{\frac{1}{N_c^2}\boxed{\left\langle\!\!\left\langle \mathrm{tr}\left[t^b\, V_{\underline{0}}\, t^a\, V_{\underline{1}}^{i\,\mathrm{G}[2]\,\dagger}\right]\left(U_{\underline{2}}\right)^{ba}\right\rangle\!\!\right\rangle(z's)}-\frac{C_F}{N_c^2}\boxed{\left\langle\!\!\left\langle \mathrm{tr}\left[V_{\underline{0}}\, V_{\underline{1}}^{i\,\mathrm{G}[2]\,\dagger}\right]\right\rangle\!\!\right\rangle(z's)}+\mathrm{c.c.}\right\}.
\end{aligned}
$$

A very last slide

A quick conclusion

- Small-x evolution equations for helicity distributions at DLA.
- Involve G_2 operator.
- Numerical agreement with the intercept found by BER.

Some Prospects

- Solving those equation in the Veneziano limit (oscillations?).
- Going beyond the DLA limit. Resumming IR-log, and thus interfacing with full spin-dependent DGLAP.
- Fixing helicity-JIMWLK.
- Phenomenology using the JAM framework.

Extra

- Gluon helicity and Lipatov vertex
- Large N_c limit
- TMD's
- g_1
- Recovering small-x Pol DGALP

Gluon helicity

From the Jaffe-Manohar (JM) gluon helicity PDF

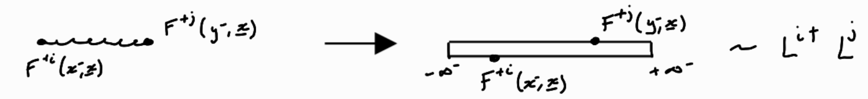
$$\Delta G(x,Q^2) = \int^{Q^2} d^2k\, g_{1L}^{G\,dip}(x,k_T^2) = \frac{-2i}{x\,P^+}\frac{1}{4\pi}\frac{1}{2}\sum_{S_L} S_L \int_{-\infty}^{\infty} d\xi^-\, e^{ixP^+\xi^-}$$
$$\times \langle P,S_L|\,\epsilon^{ij}\, F^{a+i}(0^+,0^-,\underline{0})\, U_{\underline{0}}^{ab}[0,\xi^-]\, F^{b+j}(0^+,\xi^-,\underline{0})\,|P,S_L\rangle\,, \tag{19}$$

Identify after some algebra the dipole gluon helicity TMD

$$g_{1L}^{G\,dip}(x,k_T^2) = \frac{-2i}{x\,P^+\,V^-}\frac{1}{(2\pi)^3}\frac{1}{2}\sum_{S_L} S_L\ \langle P,S_L|\,\epsilon^{ij}\,\mathrm{tr}\Big[L^{i\,\dagger}(x,\underline{k})\,L^j(x,\underline{k})\Big]\,|P,S_L\rangle \tag{20}$$

where we define the Lipatov vertex:

$$L^j(x,\underline{k}) \equiv \int_{-\infty}^{\infty} d\xi^-\, d^2\xi\, e^{ixP^+\xi^- - i\underline{k}\cdot\underline{\xi}}\ V_{\underline{\xi}}[\infty,\xi^-]\,(\partial^j A^+ + ixP^+A^j)\,V_{\underline{\xi}}[\xi^-,-\infty] \tag{21}$$



Gluon helicity

Expanding the Lipatov vertex in eikonality (i.e. Bjorken x)

$$L^j(x,\underline{k}) = \int_{-\infty}^{\infty} d\xi^- \, d^2\xi \, e^{-i\underline{k}\cdot\underline{\xi}} \, V_{\underline{\xi}}[\infty,\xi^-] \left[\partial^j A^+ + ixP^+ \left(\xi^- \, \partial^j A^+ + A^j\right) + \mathcal{O}(x^2)\right] V_{\underline{\xi}}[\xi^-,-\infty], \tag{22}$$

which we can write

$$L^j(x,\underline{k}) = -\frac{k^j}{g} \int d^2\xi \, e^{-i\underline{k}\cdot\underline{\xi}} \, V_{\underline{\xi}} - \frac{xP^+}{2g} \int d^2\xi \, e^{-i\underline{k}\cdot\underline{\xi}} \int_{-\infty}^{\infty} dz^- \, V_{\underline{\xi}}[\infty,z^-] \left[D^j - \overleftarrow{D}^j\right] V_{\underline{\xi}}[z^-,-\infty] \tag{23}$$

Performing the helicity dependent "CGC average"

$$g_{1L}^{G\,dip}(x,k_T^2) = \frac{-4i}{g^2\,(2\pi)^3} \, \epsilon^{ij} \, k^i \int d^2\zeta \, d^2\xi \, e^{-i\underline{k}\cdot(\underline{\xi}-\underline{\zeta})} \underbrace{\left\langle\!\!\left\langle \mathrm{tr}\left[V_{\underline{\zeta}}^\dagger \, V_{\underline{\xi}}^{j\,\mathrm{G}[2]} + \left(V_{\underline{\xi}}^{j\,\mathrm{G}[2]}\right)^\dagger V_{\underline{\zeta}}\right]\right\rangle\!\!\right\rangle}_{=2N_c G^j_{\underline{\xi},\underline{\zeta}}(zs)}, \tag{24}$$

with a polarized Wilson line of the second kind (more on this soon)

$$V_{\underline{z}}^{i\,\mathrm{G}[2]} \equiv \frac{P^+}{2s} \int_{-\infty}^{\infty} dz^- \, V_{\underline{z}}[\infty,z^-] \left[D^i(z^-,\underline{z}) - \overleftarrow{D}^i(z^-,\underline{z})\right] V_{\underline{z}}[z^-,-\infty]. \tag{25}$$

$\Longrightarrow$ We call $G^j_{\underline{\xi},\underline{\zeta}}(zs)$ a Polarized dipole amplitude of the second kind.

Solving, 't Hooft limit - Color

In the large N_c-limit (drop quarks t-channel exchanges)

$$U_{\underline{x}}^{\mathrm{pol}[1]} \to U_{\underline{x}}^{\mathrm{G}[1]} \tag{26}$$

Replace adjoint WL using:

$$(U_{\underline{x}})^{ba} = 2\,\mathrm{tr}[t^b V_{\underline{x}} t^a V_{\underline{x}}^{\dagger}]. \tag{27}$$

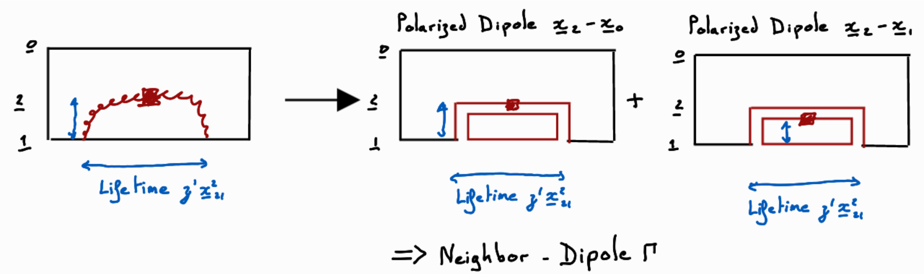
and

$$\left(U_{\underline{x}}^{\mathrm{G}[1]}\right)^{ba} = 2 \times \left\{2\,\mathrm{tr}\left[t^b\, V_{\underline{x}}\, t^a\, V_{\underline{x}}^{\mathrm{G}[1]\,\dagger}\right] + 2\,\mathrm{tr}\left[t^b\, V_{\underline{x}}^{\mathrm{G}[1]}\, t^a\, V_{\underline{x}}^{\dagger}\right]\right\}. \tag{28}$$

$$\left(U_{\underline{x}}^{i\,\mathrm{G}[2]}\right)^{ba} = 2\,\mathrm{tr}\left[t^b\, V_{\underline{x}}\, t^a\, V_{\underline{x}}^{i\,\mathrm{G}[2]\,\dagger}\right] + 2\,\mathrm{tr}\left[t^b\, V_{\underline{x}}^{i\,\mathrm{G}[2]}\, t^a\, V_{\underline{x}}^{\dagger}\right] \tag{29}$$

Notice the factor 2 in the former. A gluon has twice the spin of a quark.

Solving, 't Hooft limit - Lifetime



After Fiertzing arround, introduce neighbor dipole amplitude Γ to enforce lifetime ordering at each step of the evolution.

Solving, 't Hooft limit - Equation and intercept

$$G(x_{10}^2, zs) = G^{(0)}(x_{10}^2, zs) + \frac{\alpha_s N_c}{2\pi} \int\limits_{\frac{1}{s x_{10}^2}}^{z} \frac{dz'}{z'} \int\limits_{\frac{1}{z's}}^{x_{10}^2} \frac{dx_{21}^2}{x_{21}^2} \Big[\Gamma(x_{10}^2, x_{21}^2, z's) + 3\, G(x_{21}^2, z's) + 2\, G_2(x_{21}^2, z's) + 2\, \Gamma_2(x_{10}^2, x_{21}^2, z's) \Big], \tag{30a}$$

$$\Gamma(x_{10}^2, x_{21}^2, z's) = G^{(0)}(x_{10}^2, z's) + \frac{\alpha_s N_c}{2\pi} \int\limits_{\frac{1}{s x_{10}^2}}^{z'} \frac{dz''}{z''} \int\limits_{\frac{1}{z''s}}^{\min\left[x_{10}^2, x_{21}^2 \frac{z'}{z''}\right]} \frac{dx_{32}^2}{x_{32}^2} \Big[\Gamma(x_{10}^2, x_{32}^2, z''s) + 3\, G(x_{32}^2, z''s) + 2\, G_2(x_{32}^2, z''s) + 2\, \Gamma_2(x_{10}^2, x_{32}^2, z''s) \Big], \tag{30b}$$

$$G_2(x_{10}^2, zs) = G_2^{(0)}(x_{10}^2, zs) + \frac{\alpha_s N_c}{\pi} \int\limits_{\frac{\Lambda^2}{s}}^{z} \frac{dz'}{z'} \int\limits_{\max\left[x_{10}^2, \frac{1}{z's}\right]}^{\min\left[\frac{z}{z'} x_{10}^2, \frac{1}{\Lambda^2}\right]} \frac{dx_{21}^2}{x_{21}^2} \Big[G(x_{21}^2, z's) + 2\, G_2(x_{21}^2, z's) \Big], \tag{30c}$$

$$\Gamma_2(x_{10}^2, x_{21}^2, z's) = G_2^{(0)}(x_{10}^2, z's) + \frac{\alpha_s N_c}{\pi} \int\limits_{\frac{\Lambda^2}{s}}^{z' \frac{x_{21}^2}{x_{10}^2}} \frac{dz''}{z''} \int\limits_{\max\left[x_{10}^2, \frac{1}{z''s}\right]}^{\min\left[\frac{z'}{z''} x_{21}^2, \frac{1}{\Lambda^2}\right]} \frac{dx_{32}^2}{x_{32}^2} \Big[G(x_{32}^2, z''s) + 2\, G_2(x_{32}^2, z''s) \Big]. \tag{30d}$$

Numerical solution for the intercept:

$$\Delta\Sigma(x, Q^2) \sim \Delta G(x, Q^2) \sim g_1(x, Q^2) \sim (1/x)^{3.66 \sqrt{\frac{\alpha_s N_c}{2\pi}}}. \tag{31}$$

		Quark polarization		
		Unpolarized (U)	Longitudinally Polarized (L)	Transversely Polarized (T)
Nucleon Polarization	U	$f_1 =$	*	$h_1^{\perp} = \quad - \quad$
	L	*	$g_1 = \quad - \quad$	$h_{1L}^{\perp} = \quad - \quad$
	T	$f_{1T}^{\perp} = \quad - \quad$	$g_{1T} = \quad - \quad$	$h_1 = \quad - \quad$ $h_{1T}^{\perp} = \quad - \quad$

From "QCD2019 Workshop Summary"

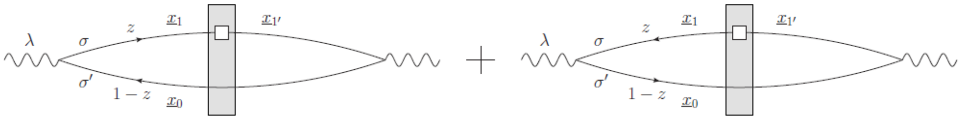
Getting g_1 - short recap

From the antisym hadronic tensor (e.g. [PDG] [Lampe and Reya 2000])

$$W^{[\mu\nu]} \sim i\epsilon_{\mu\nu\rho\sigma} \frac{q^\rho}{M_p P \cdot q} \left[S^\sigma g_1(x, Q^2) + \left(S^\sigma - \frac{Q \cdot q}{P \cdot q} P^\sigma \right) g_2(x, Q^2) \right] \tag{32}$$

DIS pol Scattering cross section is

$$\sigma^{\gamma^* p}(\lambda, \Sigma) = -\frac{8\pi^2 \alpha_{EM} x}{Q^2} \lambda\Sigma \left[g_1(x, Q^2) - \frac{4x^2 M_p^2}{Q^2} g_2(x, Q^2) \right] \tag{33}$$



One finally obtain (take the DLA limit)

$$g_1(x, Q^2) = -\sum_f \frac{Z_f^2}{2} \frac{N_c}{2\pi^3} \int_{\Lambda^2/s}^{1} \frac{dz}{z} \int_{1/zs}^{min\{1/zQ^2, 1/\Lambda^2\}} \frac{dx_{10}^2}{x_{10}^2} \left[Q(x_{10}^2, zs) + 2G_2(x_{10}^2, zs) \right] \tag{34}$$

Recovering small-x pol DGLAP

Pol DGLAP splitting function at small-x is

$$\Delta P_{gg}(z) \to \frac{\alpha_s}{2\pi} 4N_c + \left(\frac{\alpha_s}{2\pi}\right)^2 4N_c^2 \ln^2 z + \left(\frac{\alpha_s}{2\pi}\right)^3 \frac{7}{3} N_c^3 \ln^4 z \tag{35}$$

From the large N_C equations, start evolution with

$$G^{(0)}(x_1 0^2, zs) = 0, \qquad G_2^{(0)}(x_1 0^2, zs) = 1 \tag{36}$$

iterate three times, one finds

$$\Delta G^{(3)}(x, Q^2) = \left(\frac{\alpha_s}{\pi}\right)^3 \left[\frac{7}{120} \ln^5\left(\frac{1}{x}\right) \ln\left(\frac{Q^2}{\Lambda^2}\right) + \frac{1}{6} \ln^4\left(\frac{1}{x}\right) \ln^2\left(\frac{Q^2}{\Lambda^2}\right) + \frac{2}{9} \ln^3\left(\frac{1}{x}\right) \ln^3\left(\frac{Q^2}{\Lambda^2}\right)\right]$$

where

$$1/x_{10}^2 \to Q^2, \qquad zsz_{10}^2 \to 1/x \tag{37}$$