# <span id="page-0-0"></span>Solutions of the Flavor Problem through Effective Theories

# Ioannis Plakias

Pôle Théorie IJCLab Université Paris-Saclay, CNRS December, 2022

Solutions of the Flavor Problem through Effective Theo[ri](#page-29-0)es 1  $/ 10$ 

The Standard Model (SM) is the most successful theory of particle physics. However, there are some questions that it cannot answer. In particular:

The Standard Model (SM) is the most successful theory of particle physics. However, there are some questions that it cannot answer. In particular:

#### Neutrino Masses and Oscillations

The Standard Model (SM) is the most successful theory of particle physics. However, there are some questions that it cannot answer. In particular:

- Neutrino Masses and Oscillations
- The Flavor Problem

In the SM, flavor is accommodated by the Yukawa interactions. However, there are many parameters in this sector (masses and quark mixing matrix) which need to be extracted from exp. data.



The solutions of these problems require physics beyond the SM. Therefore, it is fundamental to look for its effects in all possible ways:

- $\rightarrow$  Direct production of new particles.
- $\rightarrow$  Precision tests in low-energy observables allowed in the SM (e.g.  $b \rightarrow s\gamma$ ).
- $\rightarrow$  Search for processes which are forbidden in the SM (e.g. LFV).

The solutions of these problems require physics beyond the SM. Therefore, it is fundamental to look for its effects in all possible ways:

- $\rightarrow$  Direct production of new particles.
- $\rightarrow$  Precision tests in low-energy observables allowed in the SM (e.g.  $b \rightarrow s\gamma$ ).
- $\rightarrow$  Search for processes which are forbidden in the SM (e.g. LFV).

The main goal of my thesis is to exploit low-energy experimental data in order to search for New Physics (NP).

The solutions of these problems require physics beyond the SM. Therefore, it is fundamental to look for its effects in all possible ways:

- $\rightarrow$  Direct production of new particles.
- $\rightarrow$  Precision tests in low-energy observables allowed in the SM (e.g.  $b \rightarrow s\gamma$ ).
- $\rightarrow$  Search for processes which are forbidden in the SM (e.g. LFV).

The main goal of my thesis is to exploit low-energy experimental data in order to search for New Physics (NP). Two of the projects I have worked on so far are:

- 1) Impact of meson mixing on  $B_s \to \phi ee$  angular observables.[arXiv: 2210.11995]
- 2) LFV in semileptonic observables at 1-loop.[In preparation]

# **Motivation**

Further motivation to look for NP in the flavor sector arises from hints of Lepton Flavor Universality Violation (LFUV) in  $b \to s \ell \ell$  and  $b \to c \ell \nu$  observables:

$$
\left| R_{K^{(*)}} = \frac{\mathcal{B}(B \to K^{(*)} \mu \mu)}{\mathcal{B}(B \to K^{(*)} e e)} \right|
$$

$$
\boxed{\mathbf{R}_{\mathrm{D}^{(*)}}=\frac{\mathcal{B}(B\to D^{(*)}\tau\bar{\nu})}{\mathcal{B}(B\to D^{(*)}\ell\bar{\nu})}}
$$

## **Motivation**

Further motivation to look for NP in the flavor sector arises from hints of Lepton Flavor Universality Violation (LFUV) in  $b \rightarrow s \ell \ell$  and  $b \rightarrow c \ell \nu$  observables:

$$
\mathcal{R}_{\mathcal{K}^{(*)}} = \frac{\mathcal{B}(B \to K^{(*)} \mu \mu)}{\mathcal{B}(B \to K^{(*)} e e)} \Bigg| \qquad \qquad \mathcal{R}_{\mathcal{D}^{(*)}} = \frac{\mathcal{B}(B \to D^{(*)} \tau \bar{\nu})}{\mathcal{B}(B \to D^{(*)} \ell \bar{\nu})}
$$

These anomalies can be explained by NP in semileptonic operators above the  $\Lambda \sim 1$  TeV scale. There are concrete UV scenarios that predict contributions to these operators (e.g. Leptoquarks (LQ)).

#### **Motivation**

Further motivation to look for NP in the flavor sector arises from hints of Lepton Flavor Universality Violation (LFUV) in  $b \rightarrow s \ell \ell$  and  $b \rightarrow c \ell \nu$  observables:

$$
\mathcal{R}_{\mathcal{K}^{(*)}} = \frac{\mathcal{B}(B \to K^{(*)} \mu \mu)}{\mathcal{B}(B \to K^{(*)} e e)} \Bigg| \qquad \qquad \mathcal{R}_{\mathcal{D}^{(*)}} = \frac{\mathcal{B}(B \to D^{(*)} \tau \bar{\nu})}{\mathcal{B}(B \to D^{(*)} \ell \bar{\nu})}
$$

These anomalies can be explained by NP in semileptonic operators above the  $\Lambda \sim 1$  TeV scale. There are concrete UV scenarios that predict contributions to these operators (e.g. Leptoquarks (LQ)).



# Effective Field Theories

The main tool in my research is EFTs. Heavy NP can be in general parameterized in terms of non-renormalizable operators (bottom-up approach):

$$
\mathcal{L}_{\text{EFT}} = \sum_{i} \sum_{d>4} \frac{C_i^{(d)} \mathcal{O}_i^{(d)}}{\Lambda^{d-4}} , \quad \Lambda = \mathcal{O}(M) .
$$

 $\bullet$  *M* : Masses of the heavy degrees of freedom (dof).  $\mathcal{O}^{(d)}_i$  $i_i^{(u)}$ : d-dimension operators made of the light dof.  $C_i^{(d)}$  $i_i^{(u)}$ : Wilson coefficients (dimensionless constants).

# Effective Field Theories

The main tool in my research is EFTs. Heavy NP can be in general parameterized in terms of non-renormalizable operators (bottom-up approach):

$$
\mathcal{L}_{\text{EFT}} = \sum_{i} \sum_{d>4} \frac{C_i^{(d)} \mathcal{O}_i^{(d)}}{\Lambda^{d-4}} , \quad \Lambda = \mathcal{O}(M) .
$$

- $\bullet$  *M* : Masses of the heavy degrees of freedom (dof).
- $\mathcal{O}^{(d)}_i$  $i_i^{(u)}$ : d-dimension operators made of the light dof.
- $C_i^{(d)}$  $i_i^{(u)}$ : Wilson coefficients (dimensionless constants).

At low-energies  $(E \ll \Lambda)$  the EFT can be truncated at a given order in  $1/\Lambda$ , depending on the accuracy at which we want to compute an observable.

The LQ models that explain  $R_{K(*)}$  and  $R_{D(*)}$  predict large LFV in  $\tau \to \mu$  transitions.

The LQ models that explain  $R_{K(*)}$  and  $R_{D(*)}$  predict large LFV in  $\tau \to \mu$  transitions. One such model is the  $S_1 + S_3$  LQ model, which can be described by the SMEFT.

The LQ models that explain  $R_{K(*)}$  and  $R_{D(*)}$  predict large LFV in  $\tau \to \mu$  transitions. One such model is the  $S_1 + S_3$  LQ model, which can be described by the SMEFT. For example for  $\tau \to \mu \phi$ :



which we computed at 1-loop for the first time.

The LQ models that explain  $R_{K(*)}$  and  $R_{D(*)}$  predict large LFV in  $\tau \to \mu$  transitions. One such model is the  $S_1 + S_3$  LQ model, which can be described by the SMEFT. For example for  $\tau \to \mu \phi$ :



which we computed at 1-loop for the first time. From the best-fit values of the LQ masses and couplings, we obtain:

$$
\mathcal{B}(\tau \to \mu \phi)|_{\text{tree}} \approx 2 \times 10^{-11} , \quad \mathcal{B}(\tau \to \mu \phi)|_{\text{loop}} \approx 4 \times 10^{-9} ,
$$

The LQ models that explain  $R_{K(*)}$  and  $R_{D(*)}$  predict large LFV in  $\tau \to \mu$  transitions. One such model is the  $S_1 + S_3$  LQ model, which can be described by the SMEFT. For example for  $\tau \to \mu \phi$ :



which we computed at 1-loop for the first time. From the best-fit values of the LQ masses and couplings, we obtain:

$$
\mathcal{B}(\tau \to \mu \phi)|_{\text{tree}} \approx 2 \times 10^{-11} , \quad \mathcal{B}(\tau \to \mu \phi)|_{\text{loop}} \approx 4 \times 10^{-9} ,
$$

while the current exp. limit from Belle is  $\mathcal{B}(\tau \to \mu \phi)_{\rm expt.} < 8.4 \times 10^{-8}$  at 90% C.L. and will be improved by Belle-II.

The LQ models that explain  $R_{K(*)}$  and  $R_{D(*)}$  predict large LFV in  $\tau \to \mu$  transitions. One such model is the  $S_1 + S_3$  LQ model, which can be described by the SMEFT. For example for  $\tau \to \mu \phi$ :



which we computed at 1-loop for the first time. From the best-fit values of the LQ masses and couplings, we obtain:

$$
\mathcal{B}(\tau \to \mu \phi)|_{\text{tree}} \approx 2 \times 10^{-11} , \quad \mathcal{B}(\tau \to \mu \phi)|_{\text{loop}} \approx 4 \times 10^{-9} ,
$$

while the current exp. limit from Belle is  $\mathcal{B}(\tau \to \mu \phi)_{\rm expt.} < 8.4 \times 10^{-8}$  at 90% C.L. and will be improved by Belle-II.

The same pattern is observed for similar processes (e.g.  $\tau \to \mu \eta$ ).

 $R_K$  hints to NP in  $b \to s \ell \ell$  transitions. We can also search for NP in processes related to the  $b \to s\gamma$  transition, which probes the operators  $\mathcal{O}_{7^{(\prime)}}$ :

$$
\mathcal{O}_7=[\bar{s}\sigma_{\mu\nu}P_Rb]F^{\mu\nu}\ ,\quad \mathcal{O}_{7'}=[\bar{s}\sigma_{\mu\nu}P_Lb]F^{\mu\nu}\ .
$$

 $R_K$  hints to NP in  $b \to s \ell \ell$  transitions. We can also search for NP in processes related to the  $b \to s\gamma$  transition, which probes the operators  $\mathcal{O}_{7^{(\prime)}}$ :

$$
\mathcal{O}_7=[\bar{s}\sigma_{\mu\nu}P_Rb]F^{\mu\nu}\;,\quad \mathcal{O}_{7'}=[\bar{s}\sigma_{\mu\nu}P_Lb]F^{\mu\nu}\;.
$$

One of the main observables to constrain  $C_{7^{(\prime)}}$  is  $\underline{B} \to X_s \gamma$ 

 $R_K$  hints to NP in  $b \to s \ell \ell$  transitions. We can also search for NP in processes related to the  $b \to s\gamma$  transition, which probes the operators  $\mathcal{O}_{7^{(\prime)}}$ :

$$
\mathcal{O}_7=[\bar{s}\sigma_{\mu\nu}P_Rb]F^{\mu\nu}\;,\quad \mathcal{O}_{7'}=[\bar{s}\sigma_{\mu\nu}P_Lb]F^{\mu\nu}\;.
$$

One of the main observables to constrain  $C_{7^{(t)}}$  is  $\underline{B} \to X_s \gamma \Rightarrow$ sensitive only to the moduli of  $C_{7^{(t)}},$  not probing CP violation.

 $R_K$  hints to NP in  $b \to s \ell \ell$  transitions. We can also search for NP in processes related to the  $b \to s\gamma$  transition, which probes the operators  $\mathcal{O}_{7^{(\prime)}}$ :

$$
\mathcal{O}_7=[\bar{s}\sigma_{\mu\nu}P_Rb]F^{\mu\nu}\;,\quad \mathcal{O}_{7'}=[\bar{s}\sigma_{\mu\nu}P_Lb]F^{\mu\nu}\;.
$$

One of the main observables to constrain  $C_{7^{(t)}}$  is  $\underline{B} \to X_s \gamma \Rightarrow$ sensitive only to the moduli of  $C_{7^{(t)}},$  not probing CP violation.

We can also constrain  $C_{7^{(\prime)}}$  from angular observables of  $B_s \to \phi ee$  and  $\underline{B \to K^*ee}$  at low  $q^2$  ( $e^+e^-$  invariant mass), to probe both CP conserving and CP violating NP.

 $R_K$  hints to NP in  $b \to s \ell \ell$  transitions. We can also search for NP in processes related to the  $b \to s\gamma$  transition, which probes the operators  $\mathcal{O}_{7^{(\prime)}}$ :

$$
\mathcal{O}_7=[\bar{s}\sigma_{\mu\nu}P_Rb]F^{\mu\nu}\;,\quad \mathcal{O}_{7'}=[\bar{s}\sigma_{\mu\nu}P_Lb]F^{\mu\nu}\;.
$$

One of the main observables to constrain  $C_{7^{(t)}}$  is  $\underline{B} \to X_s \gamma \Rightarrow$ sensitive only to the moduli of  $C_{7^{(t)}},$  not probing CP violation.

We can also constrain  $C_{7^{(\prime)}}$  from angular observables of  $B_s \to \phi ee$  and  $\underline{B \to K^*ee}$  at low  $q^2$  ( $e^+e^-$  invariant mass), to probe both CP conserving and CP violating NP.

There are already exp. data for  $B \to K^*ee$  from LHCb [arXiv: 2010.06011] and there will be similar searches also for  $B_s \to \phi ee$ in the future.

For  $B_s \to \phi ee$  at LHCb, the  $B_s - \bar{B}_s$  mixing effects have to be taken into account since we cannot perform flavor tagging.

For  $B_s \to \phi ee$  at LHCb, the  $B_s - \bar{B}_s$  mixing effects have to be taken into account since we cannot perform flavor tagging. In [arXiv: 2210.11995] we studied the impact of mixing into the angular observables.

For  $B_s \to \phi ee$  at LHCb, the  $B_s - \bar{B}_s$  mixing effects have to be taken into account since we cannot perform flavor tagging. In [arXiv: 2210.11995] we studied the impact of mixing into the angular observables.

For two of the angular observables  $(A_T^{(2)})$  $T^{(2)}$  and  $A_T^{(Im);CF}$  $T^{(1m),cT}$ , we find the projected constraints on the effective couplings  $C_{7^{(\prime)}}$ :



Figure 1: Constraints with mixing (left) and without mixing (right)

My thesis is devoted to searching for NP in flavor observables.

- My thesis is devoted to searching for NP in flavor observables.
- The main tool I use for my research is EFTs to analyze exp. data.
- My thesis is devoted to searching for NP in flavor observables.
- The main tool I use for my research is EFTs to analyze exp. data.
- The projects I have worked on so far are:
	- 1) Impact of meson mixing on  $B_s \to \phi ee$  angular observables.[arXiv: 2210.11995]
	- 2) LFV in semileptonic observables at 1-loop.[In preparation]
	- 3) Constraints on LFV observables from LHC data.[In preparation]

# <span id="page-29-0"></span>Thank you!