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1 Phenomenological stage

A substance presents itself, as it is, in a group of phenomena Ex : Resonance of hadrons

2 Substantialistic stage

Investigation of the structure of the substance, distinction from the phenomena

Ex : Discovery of quarks

3 Essentialistic stage

Dynamics is understood : interactions and laws of motion are clarified

Ex : Formulation of QCD

Quantum gravity needs to skip the first stage.

Space-time $(\mathcal{M}, g_{\mu\nu})$

Partially ordered set (\mathcal{M}, \prec)

$$
\forall x, y, z \in \mathcal{M}, \left\{ \begin{array}{l} x \prec x \text{ (Reflexivity)} \\ x \prec y \text{ and } y \prec x \implies x = y \text{ (Acyclicity)} \\ x \prec y \text{ and } y \prec z \implies x \prec z \text{ (Transitivity)} \end{array} \right.
$$

Causal set (c, \prec)

 (c, \prec) is a partially ordered set $\forall x, y \in c, |\{z \in c / x \prec z \prec y\}| < \infty$

HKMML theorem :

Let (M_1, g_1) and (M_2, g_2) be d-dimensional Lorentzian manifolds with $d > 2$ such that the chronological (or timelike) past and future of each point in space-time is unique.

If there exists a causal bijection between (M_1, g_1) and (M_2, g_2) .

(If $\exists f : (\mathcal{M}_1, \prec_1) \rightarrow (\mathcal{M}_2, \prec_2) \forall x, y \in \mathcal{M}_1, x \prec_1 y \Leftrightarrow f(x) \prec_2 f(y))$

Then (M_1, g_1) and (M_2, g_2) are conformally isometric.

This means that the causal structure determine not just one space-time, but the full conformal equivalence class of it.

If we know only (c_N, \prec) , with c_N a causal set of N elements, we cannot recover a volume of space-time out of it.

We need the information contained in ρ (or ε) such that :

$$
\frac{N}{\rho} = \int_{V} \sqrt{\det g} \, \mathrm{d}^d x.
$$

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Formalism

A causal set could be represented by a Hasse diagram.

- \bullet c \prec b \prec a
- \bullet {b, a} is a link or a 2-chain
- $\bullet \{c, b, a\}$ is a 3-chain
- \bullet {d, a} is a 2-antichain
- $\text{ Past}(a)$ is the set of all elements e such that $e \prec a$ (the green set)
- **o** This is a **Past-finite** causal set c : $\forall e \in c, |\text{Fast}(e)| < \infty.$

$(c, \prec) + \varepsilon \leftrightarrow (M, g)$

Faithful embedding map Φ : $c \rightarrow (\mathcal{M}, g)$

$$
\bullet \; x \prec_c y \Leftrightarrow \Phi(x) \prec_M \Phi(y)
$$

- Embedded points are distributed uniformly with unit density.
- The characteristic length over which the continuous geometry varies ≫ mean spacing between embedded points

Poissonian selection of random positions $(\mathcal{M}, g) \to c$

 \bullet The probability of finding *n* elements in a spacetime region of volume V is given by :

$$
P_V(n) = \frac{(\rho V)^n e^{-\rho V}}{n!}.
$$

 $\bullet \langle N \rangle = \rho V$ $\Delta N =$ √ $\overline{\rho}$ V If a causal set c could have arisen from a sprinkling process into (\mathcal{M}, g) "with relatively high probability" :

$$
(c, \prec) + \varepsilon \simeq (\mathcal{M}, g) \quad \left(\Leftrightarrow (c, \prec) + \varepsilon \quad \frac{\text{Spring}}{\text{Embedding}} \quad (\mathcal{M}, g) \right)
$$

The Hauptvermutung of CST :

$$
(c, \prec) + \varepsilon \simeq (\mathcal{M}, g) \& (c, \prec) + \varepsilon \simeq (\mathcal{M}', g')
$$

$$
\Rightarrow (\mathcal{M}, g) \simeq (\mathcal{M}', g')
$$

then (M, g) and (M', g') differ only at scale smaller than ρ .

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We interpret $Z=\int {\cal D}g e^{iS[g]}$ as $Z_{\Omega}\equiv \sum_{c\in \Omega} e^{i\frac{S(c)}{\hbar}}$ where Ω is the space of all past-finite causal set that we need to construct.

The standard action for causal set is the Benincasa-Dowker action defined as :

$$
S(c_n^i) = \frac{4}{\sqrt{6}} \left[n - N_0^{(n,i)} + 9N_1^{(n,i)} - 16N_2^{(n,i)} + 8N_3^{(n,i)} \right].
$$

Where $N_k^{(n,i)}$ $\kappa_k^{(n,r)}$ is the total number of *k*-element order interval in the causal set c_n^i . It gives the Einstein-Hilbert action in the continuum limit.

 $\begin{picture}(40,40) \put(0,0){\vector(0,1){30}} \put(15,0){\vector(0,1){30}} \put(15,0){\vector(0$

As $n \to \infty$, this growth process generates the sample space Ω of countable labelled past finite causal sets.

PHYSICAL REQUIREMENTS OF THE DYNAMICS :

- **•** Markov sum rule
- Internal temporality
- Discrete general covariance
- **•** Bell causality

Bell Causality

$$
\frac{\alpha(c_4 \rightarrow c_5^1)}{\alpha(c_4 \rightarrow c_5^2)} = \frac{\alpha(c_2 \rightarrow c_3^1)}{\alpha(c_2 \rightarrow c_3^2)}
$$

The continuum may be a mathematical construct which approximates an underlying physical discreteness.

$$
(\mathcal{M},g)\xrightarrow[\text{Sprinkling}]{\text{Embedding}}(c,\prec)+\varepsilon
$$

Causal set approach :

- ✓ cures divergences in QFT.
- \checkmark cures curvature singularity in GR.
- \checkmark cures infinite entanglement entropy of black holes.
- ✓ measures metric at sub-Planckian scale.
- \checkmark is compatible with Lorentz invariance.
- $\sqrt{\ }$ predicts the right magnitude of Λ .
- ? could give fruitful formulation of quantum fields dynamics.
- ? could solve the Hard Problem of Consciousness as a birth process happening in the brain.

Thank you for your attention

Additional slides

Let Ω be a non-empty set. $\mathfrak{A} \subset \mathcal{P}(\Omega)$ is an algebra if :

\n- (i)
$$
\Omega \in \mathfrak{A}
$$
\n- (ii) $A \in \mathfrak{A} \implies A^c \in \mathfrak{A}$
\n- (iii) $A_1, A_2, \ldots, A_n \in \mathfrak{A} \implies \bigcup_{k=1}^n A_k \in \mathfrak{A}$
\n- \mathfrak{A} is closed under **finite** unions
\n

 $\mathfrak A$ is a σ -algebra if is also closed under **countable** unions. $A \in \mathfrak{A}$ is a measurable set. (Ω, \mathfrak{A}) is a measurable space. A measure on a measurable space is a map satisfying :

(i) $\mu(\emptyset) = 0$ (ii) $A_1, ..., A_n \in \mathfrak{A}$ pairwise disjoint $\implies \mu(\bigcup_{k=1}^n A_k) = \sum_{k=1}^n \mu(A_k)$

 $(\Omega, \mathfrak{A}, \mu)$ is a *measure space* \rightarrow gives the CST dynamics.

CHK theorem

Caratheodary-Hahn-Kluvnek (CHK) theorem

If μ is bounded weakly countably additive vector measure over $\mathfrak A$: μ is strongly additive ⇔∃! countably additive extension of μ to $\mathfrak{S}_{\mathfrak{A}}$

• bounded

$$
\sup\nolimits_{x^*} {\left\{ \sup_{\pi} \sum_{\alpha_i \in \pi} ||x^*(\mu(\alpha_i))||; x^* \in \mathcal{H}^*, ||x^*|| \leq 1 \right\}} < \infty
$$

where the second supremum is taken over all partitions π of Ω

• weakly countably additive For every $x^* \in \mathcal{H}^*$, $x^*(\mu)$ countably additive.

$$
\Leftrightarrow x^* \left(\mu \left(\bigcup_i \alpha_i \right) \right) = \sum_i x^* (\mu(\alpha_i))
$$

for infinite sequence $\{\alpha_n\}$ of pairwise disjoint element of $\mathfrak A$

CHK theorem

Caratheodary-Hahn-Kluvnek (CHK) theorem

If μ is bounded weakly countably additive vector measure over $\mathfrak A$: μ is strongly additive ⇔∃! countably additive extension of μ to $\mathfrak{S}_{\mathfrak{A}}$

vector measure \bullet

> DEFINITION 1. A function F from a field F of subsets of a set Q to a Banach space X is called a *finitely additive vector measure*, or simply a *vector measure*, if whenever E_1 and E_2 are disjoint members of $\mathcal F$ then $F(E_1 \cup E_2) = F(E_1) +$ $F(E_2)$.

• strongly additive

$$
||\sum_{n=1}^{\infty}\mu(\alpha_n)||<\infty
$$

for every sequence $\{\alpha_n\}$ of pairwise disjoint element of $\mathfrak A$

CHK theorem simplified

Caratheodary-Hahn-Kluvnek (CHK) theorem

If μ is bounded weakly countably additive vector measure over \mathfrak{A} : μ is strongly additive ⇔∃! countably additive extension of μ to $\mathfrak{S}_{\mathfrak{A}}$

Variation of μ , $\forall \alpha \in \mathfrak{A}$, where π is a finite partition of α :

$$
|\mu|(\alpha) \equiv \sup_{\pi} \sum_{\alpha_i \in \pi} ||\mu(\alpha_i)||
$$

Measure μ is of bounded variation if $|\mu|(\Omega) < \infty$

Theorem : μ of bounded variation \implies μ strongly additive **Theorem** : μ strongly additive $\implies \mu$ bounded

Simplified CHK theorem

If μ is a weakly countably additive vector measure over \mathfrak{A} : µ of bounded variation \Rightarrow ∃! countably additive extension to $\mathfrak{S}_{\mathfrak{A}}$

Need for extension

$$
cylinder set: \operatorname{cyl}(c_n^i) \equiv \{c \in \Omega | c|_n = c_n^i\} = \bigcup_{j(i)} \operatorname{cyl}(c_{n+1}^{j(i)}) \subset \Omega
$$

Nesting property for $m > n$:

$$
\operatorname{cyl}(c_m^i) \cap \operatorname{cyl}(c_n^j) \neq 0 \implies \operatorname{cyl}(c_m^i) \subset \operatorname{cyl}(c_n^j)
$$

 $\text{cyl}(c_n^i) \subset \Omega$. $\mathfrak A$ is generated from the cylinder sets via finite unions, intersections and set differences.

$$
\mu(c_n^i) \equiv \mu(\text{cyl}(c_n^i))
$$

The event algebra $\mathfrak A$ does not suffice to be able to define covariant observables like the originary event : $\alpha_{\text{orig}} = \left(\bigcup_{n>1} \bigcup_{i \in \mathcal{I}_n} \text{cyl}(c_n^i)\right)^c$. \implies One needs to include countable set operations on \mathfrak{A} .

Covariant events $\in \mathfrak{S}_{\mathfrak{A}}/\sim$

 $c \sim c' \Leftrightarrow c, c'$ are *order-isomorphic* to each other

Transitive percolation

Generate a random causal set by the following algorithm :

- Start with n elements labeled $0, 1, 2, \dots, n 1$ $(n = \infty$ not excluded.)
- 2 With a fixed probability p, introduce a relation between every pair of points labeled *i* and *j*, where $i < j$.
- Form the transitive closure of these relations (e.g. if 2 \prec 5 and $5 \prec 8$ then enforce that $2 \prec 8$.)

The transition probability α_n from c_n^i to a specified child $c_{n+1}^{j(i)}$:

$$
\alpha_n^{(S)} = \rho^m (1-\rho)^{n-\varpi}
$$

 $m =$ number of maximal elements in the past S of the new element ϖ = size of the past S of the new element

PHYSICAL REQUIREMENTS OF THE DYNAMICS : \checkmark

Physical requirement for transitive percolation

✓ Internal temporality

Build into our definition of the growth process

✓ Discrete general covariance

Net probability of a given c_n^i in "manifestly covariant form" is $P(c_n^i) = Wp^Lq^{{n \choose 2}-R}$ where L is the number of links in c_n^i , R the number of relations, and W the number of (natural) labelings of c_n^i .

✓ Bell causality

Consider two different children, one with $(m, \varpi) = (m_1, \varpi_1)$ and the other with $(m, \varpi) = (m_2, \varpi_2)$

$$
\frac{\alpha_n^{(m_1,\varpi_1)}}{\alpha_n^{(m_2,\varpi_2)}}=\frac{\alpha_{n'}^{(m_1,\varpi_1)}}{\alpha_{n'}^{(m_2,\varpi_2)}}\Leftrightarrow\frac{p^{m_1}q^{n-\varpi_1}}{p^{m_2}q^{n-\varpi_2}}=\frac{p^{m_1}q^{n'-\varpi_1}}{p^{m_2}q^{n'-\varpi_2}}
$$

where $n' \leq n$ is the cardinality of the union of the precursor sets of the two transitions.

✓ Markov sum rule

Trivial in a well-defined probabilistic procedure.

Parameters of the growth $= q_n =$ probabilities to add a completely disconnected element at stage n.

$$
\alpha_n^{(S)} = \alpha_n^{(m,\varpi)} = \sum_{k=0}^m (-1)^k {m \choose k} \frac{q_n}{q_{\varpi-k}} = \frac{\sum_{l=m}^{\varpi} \left(\frac{\varpi - m}{\varpi - l} \right) t_l}{\sum_{j=0}^n {n \choose j} t_j}
$$

 $m =$ number of maximal elements in the past S of the new element ϖ = size of the past S of the new element

Alternative parameters : $t_n = \sum_{k=0}^n (-1)^{n-k} \binom{n}{k}$ $\binom{n}{k} \frac{1}{q}$ $\frac{1}{q_k}$.

PHYSICAL REQUIREMENTS OF THE DYNAMICS : \checkmark

Physical requirement for general transition probability

✓ Internal temporality

Build into our definition of the growth process

- ✓ Discrete general covariance Probability of a labeled causal set \tilde{c}_n^i : $P(\tilde{c}_n^i) = \prod_{i=0}^{N-1} \alpha(i, \varpi_i, m_i)$ This is a product over all elements $x \in c_n^i$ of poset invariant quantities that depends only on the structure of $past(x)$.
- ✓ Bell causality

$$
\frac{\alpha_n^{(m_1,\varpi_1)}}{\alpha_n^{(m_2,\varpi_2)}}=\frac{\alpha_{n'}^{(m_1,\varpi_1)}}{\alpha_{n'}^{(m_2,\varpi_2)}} \Leftrightarrow \frac{\sum_{k=0}^{m_1}(-1)^k\binom{m_1}{k}\frac{q_n}{q_{\varpi_1-k}}}{\sum_{k=0}^{m_2}(-1)^k\binom{m_2}{k}\frac{q_n}{q_{\varpi_2-k}}}=\frac{\sum_{k=0}^{m_1}(-1)^k\binom{m_1}{k}\frac{q_{n'}}{q_{\varpi_1-k}}}{\sum_{k=0}^{m_2}(-1)^k\binom{m_2}{k}\frac{q_{n'}}{q_{\varpi_2-k}}}
$$

The ratios depends only on precursor set structure.

✓ Markov sum rule

Impose a constraint :

$$
\sum_{i=0}^{N-1} \alpha(i, \varpi_i, m_i) = 1 \Leftrightarrow \sum_{S} \sum_{I} t_I \binom{|S| - m(S)}{1 - m(S)} = \sum_{j} t_j \binom{n}{j}
$$

$$
\Leftrightarrow \forall I, \sum_{S} \binom{|S| - m(S)}{1 - m(S)} = \binom{n}{I}
$$

Corresponding measure space

The dynamics is a specification of the measure over \mathfrak{A} .

$$
\mu(\mathrm{cyl}(c_n^i)) \equiv P(c_n^i) = \prod_{i=1}^n \alpha_i
$$

$$
\mu: {\mathfrak{A}} \rightarrow [0,1], \ \ \mu(\Omega)=\mu({\rm cyl}(c_1^1))=1
$$

 $\forall \alpha \in \mathfrak{A}$, there exists a smallest $n < \infty$ and a subset $S \subset \{1, 2, ..., |\Omega_n|\}$ such that $\alpha = \bigcup_{k \in S} \text{cyl}(c_n^k)$.

$$
\mu(\alpha) = \sum_{k \in S} P(c_n^k)
$$

 μ scalar real measure $\implies \mu$ extends to $\mathfrak{S}_{\mathfrak{A}}$.

Entropy catastrophe : KR posets in CSG

Lemma (Brightwell, Dowker, Garcia, Henson, Sorkin) : In the CSG dynamics with $t_k \neq 0$ for some $k > 1$, a causet containing an infinite level almost surely does not occur.

$$
\sum_{n=|S|+1}^{\infty} \alpha_n^{(S)} = \sum_{l=m}^{\infty} {\binom{\varpi-m}{\varpi-l}} t_l \sum_{n=|S|+1}^{\infty} \frac{1}{\sum_{j=0}^{n} {\binom{n}{j}} t_j} < \infty
$$

Straightforward generalisation of classical sequential growth, the transition amplitudes are :

$$
A_n^{(S)} = A_n^{(m,\varpi)} = \sum_{k=0}^m (-1)^k \binom{m}{k} \frac{q_n}{q_{\varpi-k}} = \frac{\sum_{l=m}^{\varpi} \binom{\varpi-m}{\varpi-l} t_l}{\sum_{j=0}^n \binom{n}{j} t_j}
$$

with now $q_n, t_n \in \mathbb{C}$.

The measure of $\text{cyl}(c_n^i) \in \mathfrak{A}$ yields :

$$
|c_n^i\rangle \equiv \mu\big(\text{cyl}(c_n^i)\big) \propto \prod_{m \text{ in branch}} A(c_m \to c_{m+1}) \in \mathbb{C}
$$

The product is over transition along the nodes from c_1^1 to c_n^i .

PHYSICAL REQUIREMENTS OF THE DYNAMICS : \checkmark

Physical requirement on **C**SG

✓ Internal temporality

Build into our definition of the growth process

✓ Discrete general covariance

$$
\implies |c_n^i\rangle = |c_n^j\rangle \text{ whenever } c_n^i \sim c_n^j
$$

✓ Bell causality :

$$
\frac{A_{n}^{(m_1,\varpi_1)}}{A_{n}^{(m_2,\varpi_2)}}=\frac{A_{n'}^{(m_1,\varpi_1)}}{A_{n'}^{(m_2,\varpi_2)}}\Leftrightarrow \frac{\sum_{k=0}^{m_1}(-1)^k\binom{m_1}{k}\frac{1}{q_{\varpi_1-k}}}{\sum_{k=0}^{m_2}(-1)^k\binom{m_2}{k}\frac{1}{q_{\varpi_2-k}}}=\frac{\sum_{k=0}^{m_1}(-1)^k\binom{m_1}{k}\frac{1}{q_{\varpi_1-k}}}{\sum_{k=0}^{m_2}(-1)^k\binom{m_2}{k}\frac{1}{q_{\varpi_2-k}}}
$$

✓ Markov sum rule

$$
|c_{n+1}^{(i)} \rangle = \hat{O}(c_n^j \to c_{n+1}^{(i)}) |c_n^j \rangle \& \text{cyl}(c_n^j) = \bigcup_{j(i)} \text{cyl}(c_{n+1}^{(i)})
$$

$$
\mu(\text{cyl}(c_n^i)) = \mu\Big(\bigcup_{j(i)} \text{cyl}(c_{n+1}^{j(i)})\Big) = \sum_{j(i)} \mu(\text{cyl}(c_{n+1}^{j(i)})) = \sum_{j(i)} \hat{O}(c_n^i \rightarrow c_{n+1}^{j(i)})|c_n^i >
$$

$$
\implies \sum_{j(i)} \hat{O}(c_n^j \rightarrow c_{n+1}^{j(i)}) = \mathbb{1}
$$

Simplified CHK theorem

If μ is a weakly countably additive vector measure over \mathfrak{A} : µ of bounded variation \Rightarrow ∃! countably additive extension to $\mathfrak{S}_{\mathfrak{N}}$

Define
$$
\zeta_n^i \ge 0
$$
 such as $\sum_{j(i)} |A(c_n^i \rightarrow c_{n+1}^{j(i)})| = 1 + \zeta_n^i \ge 1$ we have :

$$
\zeta_n^{\max} \equiv \max_{c_n^j \in \Omega_n} \zeta_n^i \stackrel{!}{=} \frac{\sum_{k=0}^n \binom{n}{k} |t_k|}{\left| \sum_{k=0}^n \binom{n}{k} t_k \right|} - 1 \stackrel{!}{=} \zeta_n^{\text{a}}
$$

$$
\zeta_n^{\min} \equiv \min_{c_n^j \in \Omega_n} \zeta_n^i \stackrel{!}{=} \sum_{\infty=1}^n \frac{\left| \sum_{k=1}^{\infty-1} \binom{\infty-1}{k-1} t_k \right|}{\left| \sum_{k=0}^n \binom{n}{k} t_k \right|} + \frac{|t_0|}{\left| \sum_{k=0}^n \binom{n}{k} t_k \right|} - 1 \stackrel{!}{=} \zeta_n^{\text{c}}
$$

Theorem (Surya, Zalel) : μ is of bounded variation if $\sum_{n=1}^{\infty} \zeta_{n}^{\text{max}}$ converges. μ is not of bounded variation if $\sum_{n=1}^{\infty} \zeta_{n}^{\min}$ diverges.

Entropy catastrophe : KR posets in **C**SG

In complex sequential growth :

$$
\sum_{n=1}^{\infty} \zeta_n^{\max} < \infty \implies \sum_{n=|S|+1}^{\infty} \alpha_n^{(S)} < \infty
$$

Personal work in progress.

Quantum sequential growth

Event = set of histories :
$$
E = \{\gamma_1, \gamma_2, ...\}
$$

$$
\mu(E) = \mu(\gamma_1) + \mu(\gamma_2) + \ldots + I(\gamma_1, \gamma_2) + I(\gamma_1, \gamma_3) + I(\gamma_2, \gamma_3) + \ldots = D(E, E)
$$

 $I(x, y) = D(x, y) + D(y, x)$ are interferences terms.

 $D: \mathfrak{A} \times \mathfrak{A} \to \mathbb{C}$ is the *decoherence functional* defined with :

- Hermiticity : $\forall \alpha, \beta \in \mathfrak{A}, D(\alpha, \beta) = D(\beta, \alpha)^*$
- **·** Linearity : $\forall \alpha, \beta, \delta \in \mathfrak{A}/\beta \cap \delta = \varnothing$, $D(\alpha, \beta \cup \delta) = D(\alpha, \beta) + D(\alpha, \delta)$
- Normalisation : $D(\Omega, \Omega) = 1$
- Strong positivity : for any $\{\alpha_i\}$ finite collection in \mathfrak{A} : $M_{ij}=D(\alpha_i,\alpha_j)$ has non-negative eigenvalues

$$
\exists \alpha, \beta \in \mathfrak{A}, \alpha \cap \beta = \varnothing / \mu(\alpha \cup \beta) \neq \mu(\alpha) + \mu(\beta)
$$

$$
\mu(\alpha \cup \beta \cup \delta) = \mu(\alpha \cup \beta) + \mu(\alpha \cup \delta) + \mu(\beta \cup \delta) - \mu(\alpha) - \mu(\beta) - \mu(\delta)
$$

GNS Construction \implies Quantum vector measure $\mu_v : \mathfrak{A} \to \mathcal{H}$.

$$
\forall \alpha, \beta \in \mathfrak{A}, \alpha \cap \beta = \varnothing / \mu(\alpha \cup \beta) = \mu(\alpha) + \mu(\beta)
$$

PHYSICAL REQUIREMENTS OF THE DYNAMICS : X Bell Causality ?

Physical requirement on QSG

- ✓ Internal temporality Build into our definition of the growth process
- ✓ Discrete general covariance

 $\implies |c_n^i>=|c_n^j>$ whenever $c_n^i\sim c_n^j$

- Bell causality : ???
- ✓ Markov sum rule

$$
|c_{n+1}^{(i)} \rangle = \hat{O}(c_n^j \to c_{n+1}^{(i)}) |c_n^j \rangle \& cyl(c_n^j) = \bigcup_{j(i)} cyl(c_{n+1}^{(i)})
$$
\n
$$
\mu(cyl(c_n^j)) = \mu\Big(\bigcup_{j(i)} cyl(c_{n+1}^{j(i)})\Big) = \sum_{j(i)} \mu(cyl(c_{n+1}^{j(i)}) = \sum_{j(i)} \hat{O}(c_n^j \to c_{n+1}^{j(i)}) |c_n^j \rangle
$$
\n
$$
\implies \sum_{j(i)} \hat{O}(c_n^j \to c_{n+1}^{j(i)}) = \mathbb{1}
$$

Complex percolation is a natural quantum generalisation of transitive percolation in which real probabilities are replaced by complex amplitudes.

$$
P(c_n^i) = Wp^Lq^{\binom{n}{2}-R} \rightarrow A(c_n^i)
$$

Lemma : The quantum vector measure of complex percolation is not of bounded variation when the parameter p is not real.

The construction of a Hilbert space from the event algebra $\mathfrak A$ and the decoherence functional D implies that the quantum measure is equivalent to a Hilbert space valued measure which is additive, unlike the quantum measure :

Vector pre-measure $\eta_v : \mathfrak{A} \to \mathcal{B}/\mathfrak{A}$

$$
\eta_{\mathsf{v}}\Big(\bigcup_{n=1}^N\alpha_n\Big)=\sum_{n=1}^N\eta_{\mathsf{v}}(\alpha_{\mathsf{v}})
$$

Vector measure $\eta_{\nu}: \mathfrak{S} \to \mathcal{B}/\mathfrak{S}$

$$
\bar{\eta}_v\Big(\bigcup_{n=1}^{\infty}\alpha_n\Big)=\sum_{n=1}^{\infty}\bar{\eta}_v(\alpha_v)
$$

Quantum Measure Theory is a formulation of quantum theory based on the path integral.

Systems described by a quantum measure space $(\Omega, \mathfrak{A}, \mu)$.

- Ω : sample space of histories γ or spacetime configurations.
- \mathfrak{A} : event algebra or set of proposition about the system.
- μ : $\boldsymbol{\mathsf{quantum}}$ pre-measure given by the path integral, $\mu: \mathfrak{A} \rightarrow \mathbb{R}^+.$

$$
\exists \alpha, \beta \in \mathfrak{A}, \alpha \cap \beta = \varnothing / \mu(\alpha \cup \beta) \neq \mu(\alpha) + \mu(\beta)
$$

 $\mu(\alpha \cup \beta \cup \delta) = \mu(\alpha \cup \beta) + \mu(\alpha \cup \delta) + \mu(\beta \cup \delta) - \mu(\alpha) - \mu(\beta) - \mu(\delta)$

To make predictions about infinite-time events \implies extension of the quantum pre-measure to a σ -algebra. Construction of the inner product vector space $(\mathcal{H}_1, +, \cdot, \langle \cdot, \cdot \rangle_1)$

 $\mathcal{H}_1 \equiv$ set of all complex-valued functions on $\mathfrak A$ which are non-zero only on a finite number of events.

$$
\forall u, v \in \mathcal{H}_1, \alpha \in \mathfrak{A}, \quad (u+v)(\alpha) \equiv u(\alpha) + v(\alpha)
$$

$$
\forall u \in \mathcal{H}_1, \lambda \in \mathbb{C}, \quad (\lambda \cdot u)(\alpha) \equiv \lambda u(\alpha)
$$

$$
\forall u, v \in \mathcal{H}_1, \quad \langle u, v \rangle_1 \equiv \sum_{\alpha \in \mathfrak{A}} \sum_{\beta \in \mathfrak{A}} u^*(\alpha) v(\beta) D(\alpha, \beta)
$$

Problem : The inner product is degenerate

Construction of the Hilbert space of histories $(\mathcal{H}_2, +, \cdot, \langle \cdot, \cdot \rangle_2)$

$$
\{u_n\} \sim \{v_n\} \Leftrightarrow \lim_{n \to +\infty} ||u_n - v_n||_1 = 0
$$

$$
\mathcal{H}_2\equiv\mathcal{H}_1/\sim
$$

 \sim equivalence class of a Cauchy sequence $\{u_n\}$ is denoted by $[u_n]$

$$
\forall [u_n], [v_n] \in \mathcal{H}_1, \alpha \in \mathfrak{A}, \quad [u_n] + [v_n] \equiv [u_n + v_n]
$$

$$
\forall [u_n] \in \mathcal{H}_1, \lambda \in \mathbb{C}, \quad \lambda \cdot [u_n] \equiv [\lambda u_n]
$$

$$
\forall [u_n], [v_n] \in \mathcal{H}_1, \quad \langle [u_n], [v_n] \rangle_2 \equiv \lim_{n \to +\infty} \langle u_n, v_n \rangle_1
$$

$$
\mu_{\mathsf{v}}(\alpha) \equiv [\chi_{\alpha}] \in \mathcal{H} \text{ with the indicator } \chi_{\alpha}(\beta) = \left\{ \begin{array}{ll} 1 & \text{if } \beta = \alpha, \\ 0 & \text{if } \beta \neq \alpha. \end{array} \right.
$$

If $\mathcal{H}=\mathbb{C}^n$, and $\mu_\mathsf{v}^{(i)}:\mathfrak{A}\to\mathbb{C},$ for $i=1,...,n$ are the components of μ_{ν} in an orthonormal basis :

 $\mu_{\mathbf v}$ is of bounded variation $\ \Leftrightarrow\ \mu_{\mathbf v}^{(i)}$ is of bounded variation

$$
\langle \mu_{\nu}(\alpha), \mu_{\nu}(\beta) \rangle = D(\alpha, \beta) \quad \langle \bar{\mu}_{\nu}(\alpha), \bar{\mu}_{\nu}(\beta) \rangle = \bar{D}(\alpha, \beta)
$$

$$
\bar{D} : \mathfrak{S}_{\mathfrak{A}} \times \mathfrak{S}_{\mathfrak{A}} \to \mathbb{C} \quad \bar{D}|_{\mathfrak{A}} = D
$$

Sorkin's slogan

 $ORDER + NUMBER = GEOMETRY$

CST dynamics is given by $(\Omega, \mathfrak{A}, \mu)$.

- Extension of measure for classical sequential growth.
- Condition on extension of measure for complex sequential growth.
- Definition of quantum vector measure.
- How to get Bell causality in quantum sequential growth?
- \Box What are conditions for extension in quantum sequential growth?