

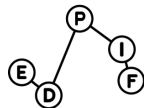
PERTURBATIVE QUANTUM FIELD THEORY TECHNIQUES APPLIED TO FLUCTUATING ELASTIC MEMBRANES.

RENCONTRES DES JEUNES PHYSICIEN(NE)S 2022 - COLLÈGE DE FRANCE

Simon Metayer

PhD advisor: Sofian Teber

November 2022



An overview on perturbative field theory

Field theory – What is it?

Field theory

Anything that can be described (directly or roughly) with waves moving in a field!

Electromagnetic waves, heat waves, mechanical waves, high energy particles ...

An outstanding example; Quantum Electrodynamics (QED)

Photons are obviously waves, but electrons are too! Let's make them interact:

$$\mathcal{L}_{\text{QED}} = \underbrace{\bar{\psi}(i\not{D} - m_e)\psi}_{\text{electrons+interactions}} - \frac{1}{4} \underbrace{F^{\mu\nu} F_{\mu\nu}}_{\text{Maxwell}}$$

QED is a QFT, which is classical field theory + special relativity + quantum mechanics.

FIELD THEORIES ARE INFINITELY COMPLICATED PROBLEMS IF WAVES CAN INTERACT WITH THEMSELVES AND/OR OTHER WAVES!

Field theory – What is it?

Field theory

Anything that can be described (directly or roughly) with waves moving in a field!

Electromagnetic waves, heat waves, mechanical waves, high energy particles ...

An outstanding example; Quantum Electrodynamics (QED)

Photons are obviously waves, but electrons are too! Let's make them interact:

$$\mathcal{L}_{\text{QED}} = \underbrace{\bar{\psi}(i\not{D} - m_e)\psi}_{\text{electrons+interactions}} - \frac{1}{4} \underbrace{F^{\mu\nu} F_{\mu\nu}}_{\text{Maxwell}}$$

QED is a QFT, which is classical field theory + special relativity + quantum mechanics.

FIELD THEORIES ARE INFINITELY COMPLICATED PROBLEMS IF WAVES CAN INTERACT WITH THEMSELVES AND/OR OTHER WAVES!

Field theory – What is it?

Field theory

Anything that can be described (directly or roughly) with waves moving in a field!

Electromagnetic waves, heat waves, mechanical waves, high energy particles ...

An outstanding example; Quantum Electrodynamics (QED)

Photons are obviously waves, but electrons are too! Let's make them interact:

$$\mathcal{L}_{\text{QED}} = \underbrace{\bar{\psi}(i\not{D} - m_e)\psi}_{\text{electrons+interactions}} - \frac{1}{4} \underbrace{F^{\mu\nu} F_{\mu\nu}}_{\text{Maxwell}}$$

QED is a QFT, which is classical field theory + special relativity + quantum mechanics.

FIELD THEORIES ARE INFINITELY COMPLICATED PROBLEMS IF WAVES CAN INTERACT WITH THEMSELVES AND/OR OTHER WAVES!

Field theory – What is it?

Field theory

Anything that can be described (directly or roughly) with waves moving in a field!

Electromagnetic waves, heat waves, mechanical waves, high energy particles ...

An outstanding example; Quantum Electrodynamics (QED)

Photons are obviously waves, but electrons are too! Let's make them interact:

$$\mathcal{L}_{\text{QED}} = \underbrace{\bar{\psi}(i\not{D} - m_e)\psi}_{\text{electrons+interactions}} - \frac{1}{4} \underbrace{F^{\mu\nu} F_{\mu\nu}}_{\text{Maxwell}}$$

QED is a QFT, which is classical field theory + special relativity + quantum mechanics.

FIELD THEORIES ARE INFINITELY COMPLICATED PROBLEMS IF WAVES CAN INTERACT WITH THEMSELVES AND/OR OTHER WAVES!

Field theory – Quantum perturbative approach

INFINITELY COMPLICATED PROBLEM \rightarrow INFINITE SUM OF “SIMPLE” PROBLEMS.

The famous Feynman diagrams

A visual representation of the computations ... An example of QED diagram:



Perturbative approach states that physical quantities (*probability amplitude*), e.g. how an electron truly propagates, can be computed precisely like:

$$X = \text{—} + \text{—} \text{---} \text{—} + \text{—} \text{---} \text{---} \text{—} + \dots$$

The equation shows a series of diagrams representing the perturbative expansion of a quantity X. The first term is a single solid line. The second term is a solid line with a self-energy loop (a wavy line connecting two vertices on the same line). The third term is a solid line with a more complex self-energy loop involving multiple vertices and wavy lines. The series continues with an ellipsis.

Summing all infinitely possible *self-interaction* loops gives the true X value ... Hopefully, the more complicated a diagram is, the smaller it contributes!

This is how we access very high precision prediction on quantities like:

“What is the probability of colliding particles A and B and get as an output particles X and Y?” or “How much, very precisely, is the magnetic moment of an electron?”.

Field theory – Quantum perturbative approach

INFINITELY COMPLICATED PROBLEM \rightarrow INFINITE SUM OF “SIMPLE” PROBLEMS.

The famous Feynman diagrams

A visual representation of the computations ... An example of QED diagram:



Perturbative approach states that physical quantities (*probability amplitude*), e.g. how an electron truly propagates, can be computed precisely like:

$$X = \text{—} + \text{—} \text{ with a self-energy loop} + \text{—} \text{ with a more complex self-energy loop} + \dots$$

The equation shows a series of terms representing the perturbative expansion of a quantity X. The first term is a simple horizontal line representing a free electron. The second term is a horizontal line with a self-energy loop (a cloud-like shape) attached to it. The third term is a horizontal line with a more complex self-energy loop (a cloud-like shape with two vertices) attached to it. The series continues with an ellipsis.

Summing all infinitely possible *self-interaction* loops gives the true X value ... Hopefully, the more complicated a diagram is, the smaller it contributes!

This is how we access very high precision prediction on quantities like:

“What is the probability of colliding particles A and B and get as an output particles X and Y?” or “How much, very precisely, is the magnetic moment of an electron?”.

Field theory – Quantum perturbative approach

INFINITELY COMPLICATED PROBLEM \rightarrow INFINITE SUM OF “SIMPLE” PROBLEMS.

The famous Feynman diagrams

A visual representation of the computations ... An example of QED diagram:



Perturbative approach states that physical quantities (*probability amplitude*), e.g. how an electron truly propagates, can be computed precisely like:

$$X = \text{—} + \text{—} \text{ with a self-energy loop} + \text{—} \text{ with a more complex self-energy loop} + \dots$$

The equation shows the true propagator X as a sum of terms. The first term is a simple straight line representing a free electron. The second term is a straight line with a self-energy loop (a loop of a photon and an electron) attached to it. The third term is a straight line with a more complex self-energy loop (a loop of a photon and an electron with an internal electron line). The sum continues with more terms indicated by an ellipsis.

Summing all infinitely possible *self-interaction* loops gives the true X value ... Hopefully, the more complicated a diagram is, the smaller it contributes!

This is how we access very high precision prediction on quantities like:

“What is the probability of colliding particles A and B and get as an output particles X and Y?” or “How much, very precisely, is the magnetic moment of an electron?”.

Field theory – Quantum perturbative approach

INFINITELY COMPLICATED PROBLEM \rightarrow INFINITE SUM OF “SIMPLE” PROBLEMS.

The famous Feynman diagrams

A visual representation of the computations ... An example of QED diagram:



Perturbative approach states that physical quantities (*probability amplitude*), e.g. how an electron truly propagates, can be computed precisely like:

$$X = \text{—} + \text{—} \text{ with a self-energy loop} + \text{—} \text{ with a more complex self-energy loop} + \dots$$

The equation shows the expansion of a quantity X. It starts with a single solid line representing a free electron. This is followed by a plus sign and a diagram of a self-energy loop (a wavy line forming a loop on a solid line). This is followed by another plus sign and a diagram of a more complex self-energy loop (a wavy line forming a loop with two vertices on a solid line). The sequence ends with a plus sign and an ellipsis, indicating an infinite series of terms.

Summing all infinitely possible *self-interaction* loops gives the true X value ... Hopefully, the more complicated a diagram is, the smaller it contributes!

This is how we access very high precision prediction on quantities like:

“What is the probability of colliding particles A and B and get as an output particles X and Y?” or “How much, very precisely, is the magnetic moment of an electron?”.

Field theory – Statistical perturbative approach

EXACTLY THE SAME FORMALISM IF FLUCTUATIONS ARE THERMAL OR STATISTICAL!

An example – From Ising model to ϕ^4 theory

Ising model is the simplest way to describe a magnet, i.e., ± 1 valued spins on a lattice. It is exactly solvable in 1&2 dimensions, but what to do for a physical 3D magnet?

Seen from far enough, Ising model may look like a field theory:

$$\underbrace{\mathcal{H}(S) = -J \sum_{\langle i,j \rangle} S_i S_j}_{\text{Ising discrete model}} \xrightarrow{\text{smoothing}} \underbrace{\mathcal{L}(\phi) = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{m^2}{2} \phi^2 - \frac{\lambda}{4!} \phi^4}_{\phi^4 \text{ field theory}}$$

Then (IR) quantities can be computed perturbatively via diagrammatic expansions too! :).

In fact, ϕ^4 model is an effective way to describe Ising model near its *critical point* ...

Field theory – Statistical perturbative approach

EXACTLY THE SAME FORMALISM IF FLUCTUATIONS ARE THERMAL OR STATISTICAL!

An example – From Ising model to ϕ^4 theory

Ising model is the simplest way to describe a magnet, i.e., ± 1 valued spins on a lattice. It is exactly solvable in 1&2 dimensions, but what to do for a physical 3D magnet?

Seen from far enough, Ising model may look like a field theory:

$$\underbrace{\mathcal{H}(S) = -J \sum_{\langle i,j \rangle} S_i S_j}_{\text{Ising discrete model}} \xrightarrow{\text{smoothing}} \underbrace{\mathcal{L}(\phi) = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{m^2}{2} \phi^2 - \frac{\lambda}{4!} \phi^4}_{\phi^4 \text{ field theory}}$$

Then (IR) quantities can be computed perturbatively via diagrammatic expansions too! :).

In fact, ϕ^4 model is an effective way to describe Ising model near its *critical point* ...

Field theory – Statistical perturbative approach

EXACTLY THE SAME FORMALISM IF FLUCTUATIONS ARE THERMAL OR STATISTICAL!

An example – From Ising model to ϕ^4 theory

Ising model is the simplest way to describe a magnet, i.e., ± 1 valued spins on a lattice. It is exactly solvable in 1&2 dimensions, but what to do for a physical 3D magnet?

Seen from far enough, Ising model may look like a field theory:

$$\underbrace{\mathcal{H}(S) = -J \sum_{\langle i,j \rangle} S_i S_j}_{\text{Ising discreet model}} \xrightarrow{\text{smoothing}} \underbrace{\mathcal{L}(\phi) = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{m^2}{2} \phi^2 - \frac{\lambda}{4!} \phi^4}_{\phi^4 \text{ field theory}}$$

Then (IR) quantities can be computed perturbatively via diagrammatic expansions too! :).

In fact, ϕ^4 model is an effective way to describe Ising model near its *critical point* ...

Field theory – Statistical perturbative approach

EXACTLY THE SAME FORMALISM IF FLUCTUATIONS ARE THERMAL OR STATISTICAL!

An example – From Ising model to ϕ^4 theory

Ising model is the simplest way to describe a magnet, i.e., ± 1 valued spins on a lattice. It is exactly solvable in 1&2 dimensions, but what to do for a physical 3D magnet?

Seen from far enough, Ising model may look like a field theory:

$$\underbrace{\mathcal{H}(S) = -J \sum_{\langle i,j \rangle} S_i S_j}_{\text{Ising discreet model}} \xrightarrow{\text{smoothing}} \underbrace{\mathcal{L}(\phi) = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{m^2}{2} \phi^2 - \frac{\lambda}{4!} \phi^4}_{\phi^4 \text{ field theory}}$$

Then (IR) quantities can be computed perturbatively via diagrammatic expansions too! :).

In fact, ϕ^4 model is an effective way to describe Ising model near its *critical point* ...

Field theory – Divergences ...

Sorry, I lied ... All graphs described above with at least one closed loop are $= \infty$...

Regularization & renormalization – taming infinities

Split divergent and convergent information, e.g.:


$$= \lim_{\varepsilon \rightarrow 0} \left(\frac{A}{\varepsilon} + B \right)$$

(some crazy looking divergent diagram)

- Convergent part (B) is the *renormalized probability amplitudes* info!
- Divergent part (A) is the *critical properties* info!

Critical properties

These are violent features of a model like **phases transitions!**

Described by *fixed points* and quantified by *critical exponents/anomalous dimensions*.

Allow to answer questions like

“What happen to the magnetization of a magnet near its melting point?” or
“How does specific heat of a liquid diverges near one of its phase transition”,
as well as describing superconductivity, superfluidity ...

Field theory – Divergences ...

Sorry, I lied ... All graphs described above with at least one closed loop are $= \infty$...

Regularization & renormalization – taming infinities

Split divergent and convergent information, e.g.:


$$\text{Diagram} = \lim_{\epsilon \rightarrow 0} \left(\frac{A}{\epsilon} + B \right)$$

(some crazy looking divergent diagram)

- Convergent part (B) is the *renormalized probability amplitudes* info!
- Divergent part (A) is the *critical properties* info!

Critical properties

These are violent features of a model like **phases transitions!**

Described by *fixed points* and quantified by *critical exponents/anomalous dimensions*.

Allow to answer questions like


“What happen to the magnetization of a magnet near its melting point?” or
“How does specific heat of a liquid diverges near one of its phase transition”,
as well as describing superconductivity, superfluidity ...

Field theory – Divergences ...

Sorry, I lied ... All graphs described above with at least one closed loop are $= \infty$...

Regularization & renormalization – taming infinities

Split divergent and convergent information, e.g.:


$$= \lim_{\varepsilon \rightarrow 0} \left(\frac{A}{\varepsilon} + B \right)$$

(some crazy looking divergent diagram)

- Convergent part (B) is the *renormalized probability amplitudes* info!
- Divergent part (A) is the *critical properties* info!

Critical properties

These are violent features of a model like **phases transitions!**

Described by *fixed points* and quantified by *critical exponents/anomalous dimensions*.

Allow to answer questions like


“What happen to the magnetization of a magnet near its melting point?” or
“How does specific heat of a liquid diverges near one of its phase transition”,
as well as describing superconductivity, superfluidity ...

Field theory – Divergences ...

Sorry, I lied ... All graphs described above with at least one closed loop are $= \infty$...

Regularization & renormalization – taming infinities

Split divergent and convergent information, e.g.:


$$= \lim_{\varepsilon \rightarrow 0} \left(\frac{A}{\varepsilon} + B \right)$$

(some crazy looking divergent diagram)

- Convergent part (B) is the *renormalized probability amplitudes* info!
- Divergent part (A) is the *critical properties* info!

Critical properties

These are violent features of a model like **phases transitions!**

Described by *fixed points* and quantified by *critical exponents/anomalous dimensions*.

Allow to answer questions like

“What happen to the magnetization of a magnet near its melting point?” or
“How does specific heat of a liquid diverges near one of its phase transition”,
as well as describing superconductivity, superfluidity ...

An intuitive field theory – Fluctuating flat membranes

Fluctuating flat membranes – Physical motivation

d -dim extended objects embedded in a larger D -dim space subject to small quantum and/or thermal fluctuations.

Applications:

- cond-mat: graphene, silicene, phosphorene ...
- bio: living cells surfaces (phospholipid bilayers)
- hep: worldsheet, branes ...

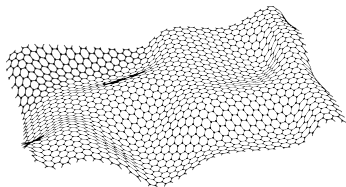


Figure: Fluctuating graphene

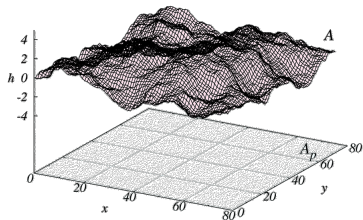


Figure: Generic fluctuating membrane

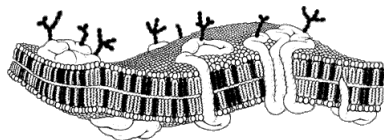


Figure: Cell bi-layered membrane

Fluctuating flat membranes – Physical motivation

d -dim extended objects embedded in a larger D -dim space subject to small quantum and/or thermal fluctuations.

Applications:

- cond-mat: graphene, silicene, phosphorene ...
- bio: living cells surfaces (phospholipid bilayers)
- hep: worldsheet, branes ...

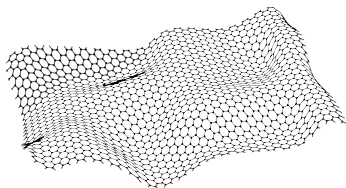


Figure: Fluctuating graphene

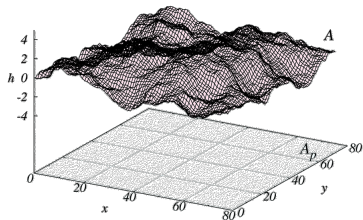


Figure: Generic fluctuating membrane

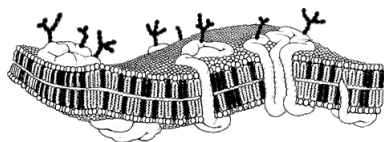
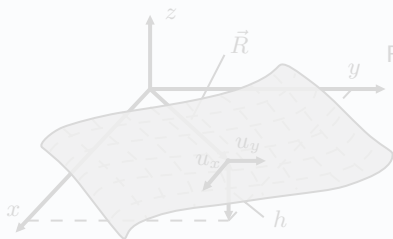


Figure: Cell bi-layered membrane

Fluctuating flat membranes – An intuitive field theory

Model parametrization



Fields parametrization:

- $\vec{u}(\vec{x}) \equiv$ longitudinal displ. (P-wave) (**phonon**)
- $h(\vec{x}) \equiv$ height displ. (S-wave) (**flexuron**)
- $\vec{R}(\vec{x}) = (\vec{x} + \vec{u}(\vec{x}), h(\vec{x})) \equiv$ coordinates with $\vec{x} = (x, y)$ and $\vec{u}(\vec{x}) = (u_x(\vec{x}), u_y(\vec{x}))$

Action

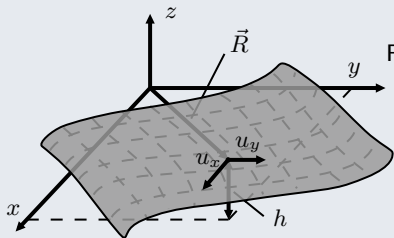
$$S[\vec{u}, h] = \int d^2x \left[\frac{1}{2} (\Delta h)^2 + \frac{\mu}{(4\pi)^2} (u_{ij})^2 + \frac{\lambda}{2(4\pi)^2} (u_{ii})^2 \right]$$

$$\text{with } u_{ij} \approx \frac{1}{2} [\partial_i u_j + \partial_j u_i + \partial_i h \partial_j h]$$

$u_{ij} \equiv$ stress tensor ; fluctuations with respect to the flat configuration $\vec{R}_0(\vec{x}) = (\vec{x}, 0)$
 $\lambda, \mu \equiv$ coupling constants \equiv Lamé coefficients

Fluctuating flat membranes – An intuitive field theory

Model parametrization



Fields parametrization:

- $\vec{u}(\vec{x}) \equiv$ longitudinal displ. (P-wave) (**phonon**)
- $h(\vec{x}) \equiv$ height displ. (S-wave) (**flexuron**)
- $\vec{R}(\vec{x}) = (\vec{x} + \vec{u}(\vec{x}), h(\vec{x})) \equiv$ coordinates with $\vec{x} = (x, y)$ and $\vec{u}(\vec{x}) = (u_x(\vec{x}), u_y(\vec{x}))$

Action

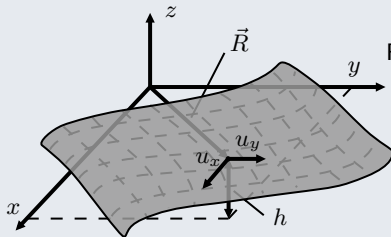
$$S[\vec{u}, h] = \int d^2x \left[\frac{1}{2} (\Delta h)^2 + \frac{\mu}{(4\pi)^2} (u_{ij})^2 + \frac{\lambda}{2(4\pi)^2} (u_{ii})^2 \right]$$

$$\text{with } u_{ij} \approx \frac{1}{2} [\partial_i u_j + \partial_j u_i + \partial_i h \partial_j h]$$

$u_{ij} \equiv$ stress tensor ; fluctuations with respect to the flat configuration $\vec{R}_0(\vec{x}) = (\vec{x}, 0)$
 $\lambda, \mu \equiv$ coupling constants \equiv Lamé coefficients

Fluctuating flat membranes – An intuitive field theory

Model parametrization



Fields parametrization:

- $\vec{u}(\vec{x}) \equiv$ longitudinal displ. (P-wave) (**phonon**)
- $h(\vec{x}) \equiv$ height displ. (S-wave) (**flexuron**)
- $\vec{R}(\vec{x}) = (\vec{x} + \vec{u}(\vec{x}), h(\vec{x})) \equiv$ coordinates with $\vec{x} = (x, y)$ and $\vec{u}(\vec{x}) = (u_x(\vec{x}), u_y(\vec{x}))$

Action

$$S[\vec{u}, h] = \int d^2x \left[\frac{1}{2} (\Delta h)^2 + \frac{\mu}{(4\pi)^2} (u_{ij})^2 + \frac{\lambda}{2(4\pi)^2} (u_{ii})^2 \right]$$

$$\text{with } u_{ij} \approx \frac{1}{2} [\partial_i u_j + \partial_j u_i + \partial_i h \partial_j h]$$

$u_{ij} \equiv$ stress tensor ; fluctuations with respect to the flat configuration $\vec{R}_0(\vec{x}) = (\vec{x}, 0)$
 $\lambda, \mu \equiv$ coupling constants \equiv Lamé coefficients

Fluctuating flat membranes – What to compute?

What to compute?

- $(\lambda, \mu) \equiv$ Lamé coefficients at a stable (scale invariant) fixed point.
- $\eta \equiv$ elastic critical exponent \equiv field anomalous dimension.

A papersheet subject to small deformations acquire anomalous rigidity and elasticity!

All other mechanical quantities are accessible with λ , μ and η only:

Quantities derived from λ and μ :

- Young modulus
- Poisson ratio (negative!)
- bulk modulus
- s-wave sound velocity ...

All crit. exponents depend only on η :

- bending/rigidity modulus
- Young modulus
- roughness exponent
- lower-crit dim ...

Fluctuating flat membranes – What to compute?

What to compute?

- $(\lambda, \mu) \equiv$ Lamé coefficients at a stable (scale invariant) fixed point.
- $\eta \equiv$ elastic critical exponent \equiv field anomalous dimension.

A papersheet subject to small deformations acquire anomalous rigidity and elasticity!

All other mechanical quantities are accessible with λ , μ and η only:

Quantities derived from λ and μ :

- Young modulus
- Poisson ratio (negative!)
- bulk modulus
- s-wave sound velocity ...

All crit. exponents depend only on η :

- bending/rigidity modulus
- Young modulus
- roughness exponent
- lower-crit dim ...

Fluctuating flat membranes – What to compute?

What to compute?

- $(\lambda, \mu) \equiv$ Lamé coefficients at a stable (scale invariant) fixed point.
- $\eta \equiv$ elastic critical exponent \equiv field anomalous dimension.

A papersheet subject to small deformations acquire anomalous rigidity and elasticity!

All other mechanical quantities are accessible with λ , μ and η only:

Quantities derived from λ and μ :

- Young modulus
- Poisson ratio (negative!)
- bulk modulus
- s-wave sound velocity ...

All crit. exponents depend only on η :

- bending/rigidity modulus
- Young modulus
- roughness exponent
- lower-crit dim ...

Fluctuating flat membranes – What to compute?

What to compute?

- $(\lambda, \mu) \equiv$ Lamé coefficients at a stable (scale invariant) fixed point.
- $\eta \equiv$ elastic critical exponent \equiv field anomalous dimension.

A papersheet subject to small deformations acquire anomalous rigidity and elasticity!

All other mechanical quantities are accessible with λ , μ and η only:

Quantities derived from λ and μ :

- Young modulus
- Poisson ratio (negative!)
- bulk modulus
- s-wave sound velocity ...

All crit. exponents depend only on η :

- bending/rigidity modulus
- Young modulus
- roughness exponent
- lower-crit dim ...

Fluctuating flat membranes – Computations

We computed the Lamé coefficients (μ, λ) and the critical exponent η **analytically** with high precision:


Flexuron propagator \equiv  **Phonon** propagator \equiv 

- **First order** (hand computations) [Aronovitz & Lubensky, '88] (4 integrals)



- **Second order** (partially automated) [Coquand, Mouhanna, Teber, '20] (318 integrals)

Fluctuating flat membranes – Computations




We computed the Lamé coefficients (μ, λ) and the critical exponent η **analytically** with high precision:

Flexuron propagator \equiv  **Phonon** propagator \equiv 

- **First order** (hand computations) [Aronovitz & Lubensky, '88] (4 integrals)

$$\Sigma_1 = \text{---} \bullet \text{---} \text{---} \bullet \text{---} \text{---} \quad \text{and} \quad \Pi_1 = \text{---} \bullet \text{---} \bullet \text{---} \text{---}$$


- **Second order** (partially automated) [Coquand, Mouhanna, Teber, '20] (318 integrals)

$$\Sigma_2 = \text{---} \bullet \text{---} \bullet \text{---} \text{---} + \text{---} \bullet \text{---} \text{---} \bullet \text{---} \text{---} + \text{---} \bullet \text{---} \text{---} \bullet \text{---} \text{---} + \text{---} \bullet \text{---} \text{---} \bullet \text{---} \text{---} + \text{---} \bullet \text{---} \text{---} \bullet \text{---} \text{---}$$


Fluctuating flat membranes – Computations

- **Third order** (highly automated) [Metayer, Mouhanna, Teber, '21]

Fluctuating flat membranes – Computations

- **Third order** (highly automated) [Metayer, Mouhanna, Teber, '21]

$$\Sigma_3 =$$
$$\Pi_3 =$$

(231939 integrals...)

Fluctuating flat membranes – Results

Analytical results exact order by order in perturbation theory:

$$\lambda = - \underbrace{\frac{4}{25}}_{1\text{st order}} + \underbrace{\frac{56}{3125}}_{2\text{nd order}} - \underbrace{\frac{2(1703808\zeta_3 - 3032351)}{48828125}}_{3\text{rd order}} + \dots$$
$$= -0.160 + 0.018 + 0.040 + \dots$$

$$\mu = + \underbrace{\frac{12}{25}}_{1\text{st order}} - \underbrace{\frac{88}{3125}}_{2\text{nd order}} + \underbrace{\frac{6(1847808\zeta_3 - 2076601)}{48828125}}_{3\text{rd order}} + \dots$$
$$= 0.480 - 0.028 + 0.018 + \dots$$

$$\eta = + \underbrace{\frac{24}{25}}_{1\text{st order}} - \underbrace{\frac{144}{3125}}_{2\text{nd order}} - \underbrace{\frac{4(1286928\zeta_3 - 568241)}{146484375}}_{3\text{rd order}} + \dots$$
$$= 0.960 - 0.0461 - 0.027 + \dots$$

Fluctuating flat membranes – Results

Analytical results exact order by order in perturbation theory:

$$\lambda = - \underbrace{\frac{4}{25}}_{1\text{st order}} + \underbrace{\frac{56}{3125}}_{2\text{nd order}} - \underbrace{\frac{2(1703808\zeta_3 - 3032351)}{48828125}}_{3\text{rd order}} + \dots$$
$$= -0.160 + 0.018 + 0.040 + \dots$$

$$\mu = + \underbrace{\frac{12}{25}}_{1\text{st order}} - \underbrace{\frac{88}{3125}}_{2\text{nd order}} + \underbrace{\frac{6(1847808\zeta_3 - 2076601)}{48828125}}_{3\text{rd order}} + \dots$$
$$= 0.480 - 0.028 + 0.018 + \dots$$

$$\eta = + \underbrace{\frac{24}{25}}_{1\text{st order}} - \underbrace{\frac{144}{3125}}_{2\text{nd order}} - \underbrace{\frac{4(1286928\zeta_3 - 568241)}{146484375}}_{3\text{rd order}} + \dots$$
$$= 0.960 - 0.0461 - 0.027 + \dots$$

Fluctuating flat membranes – Results

This leads the fundamental numbers:

$$\lambda = -0.102 \quad \mu = 0.470 \quad \eta = 0.888$$

From which we compute all desired property of the membrane, e.g.:

$$\text{Bulk modulus} \equiv K = \lambda + \mu = 0.368 \quad \text{Young modulus} \equiv E = \frac{4\mu(\lambda + \mu)}{\lambda + 2\mu} = 0.825$$

$$\text{Poisson ratio} \equiv \nu = \frac{\lambda}{\lambda + 2\mu} = -0.121 \quad \text{P-wave modulus} \equiv M = \lambda + 2\mu = 0.837$$

$$\text{Anomalous roughness} \equiv \zeta = \frac{2 - \eta}{2} = 0.556 \quad \text{Anomalous rigidity} \equiv \eta = 0.888 \quad \text{etc}$$

These are the experimental results that would be obtained by averaging the mechanical properties of a membrane over a large number of (fractal) configurations.

Fluctuating flat membranes – Results

This leads the fundamental numbers:

$$\lambda = -0.102 \quad \mu = 0.470 \quad \eta = 0.888$$

From which we compute all desired property of the membrane, e.g.:

$$\text{Bulk modulus} \equiv K = \lambda + \mu = 0.368 \quad \text{Young modulus} \equiv E = \frac{4\mu(\lambda + \mu)}{\lambda + 2\mu} = 0.825$$

$$\text{Poisson ratio} \equiv \nu = \frac{\lambda}{\lambda + 2\mu} = -0.121 \quad \text{P-wave modulus} \equiv M = \lambda + 2\mu = 0.837$$

$$\text{Anomalous roughness} \equiv \zeta = \frac{2 - \eta}{2} = 0.556 \quad \text{Anomalous rigidity} \equiv \eta = 0.888 \quad \text{etc}$$

These are the experimental results that would be obtained by averaging the mechanical properties of a membrane over a large number of (fractal) configurations.

Fluctuating flat membranes – Results

This leads the fundamental numbers:

$$\lambda = -0.102 \quad \mu = 0.470 \quad \eta = 0.888$$

From which we compute all desired property of the membrane, e.g.:

$$\begin{aligned} \text{Bulk modulus} \equiv K &= \lambda + \mu = 0.368 & \text{Young modulus} \equiv E &= \frac{4\mu(\lambda + \mu)}{\lambda + 2\mu} = 0.825 \\ \text{Poisson ratio} \equiv \nu &= \frac{\lambda}{\lambda + 2\mu} = -0.121 & \text{P-wave modulus} \equiv M &= \lambda + 2\mu = 0.837 \\ \text{Anomalous roughness} \equiv \zeta &= \frac{2 - \eta}{2} = 0.556 & \text{Anomalous rigidity} \equiv \eta &= 0.888 \quad \text{etc} \end{aligned}$$

These are the experimental results that would be obtained by averaging the mechanical properties of a membrane over a large number of (fractal) configurations.

Fluctuating flat membranes – Results

This leads the fundamental numbers:

$$\lambda = -0.102 \quad \mu = 0.470 \quad \eta = 0.888$$

From which we compute all desired property of the membrane, e.g.:

$$\begin{aligned} \text{Bulk modulus} \equiv K &= \lambda + \mu = 0.368 & \text{Young modulus} \equiv E &= \frac{4\mu(\lambda + \mu)}{\lambda + 2\mu} = 0.825 \\ \text{Poisson ratio} \equiv \nu &= \frac{\lambda}{\lambda + 2\mu} = -0.121 & \text{P-wave modulus} \equiv M &= \lambda + 2\mu = 0.837 \\ \text{Anomalous roughness} \equiv \zeta &= \frac{2 - \eta}{2} = 0.556 & \text{Anomalous rigidity} \equiv \eta &= 0.888 \quad \text{etc} \end{aligned}$$

These are the experimental results that would be obtained by averaging the mechanical properties of a membrane over a large number of (fractal) configurations.

Thank you for your attention :).