

IRN Neutrino - Orsay

Leptonic CPV phases: impact for LFV Higgs & Z-boson decays and CP asymmetries

based on 2207.10109, with A. Abada, J. Kriewald, S. Rosauero and A. M. Teixeira



Emanuelle Pinsard

LPC - Clermont



Flavour: Interactions between fermion families

In the Standard Model, Yukawas encode the **flavour** dynamics (masses, mixings and CP violation)

Gauge interactions are **flavour universal**

Lepton sector: Neutrinos are strictly **massless** in the SM

↪ **Conservation** of total **lepton number** and **lepton flavours**

↪ **Lepton Flavour Universality** (only broken by Yukawas)

↪ **No source of CPV** (only in the quark sector, not enough for BAU ...)

Numerous **tensions** between SM and observation: $(g - 2)_\ell$, B-meson “anomalies”, ...

And **observational caveats** of the SM:

dark matter, neutrino oscillations, baryon asymmetry of the Universe

Neutrino oscillations: 1st laboratory **evidence of NP**

↪ neutrinos are **massive** & **leptons mix** $\mathcal{U}_{\alpha i}^{\text{PMNS}}$

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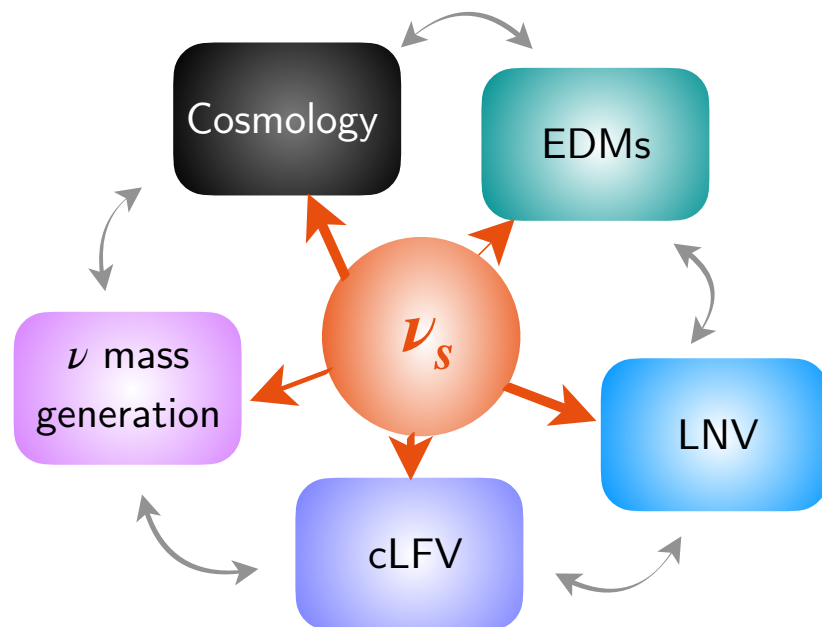
⇒ **Need New Physics**

↪ Need new fields: **Majorana**? **LNV**? New sources of **CPV**?

Which model? At which scale? ↪ **Searches for NP** in the **lepton sector**

Strong arguments in favour of **New Physics** involving (neutral) leptons!

Majorana sterile fermions are a **very appealing hypothesis**, motivated by extensive **theoretical** and **observational** arguments

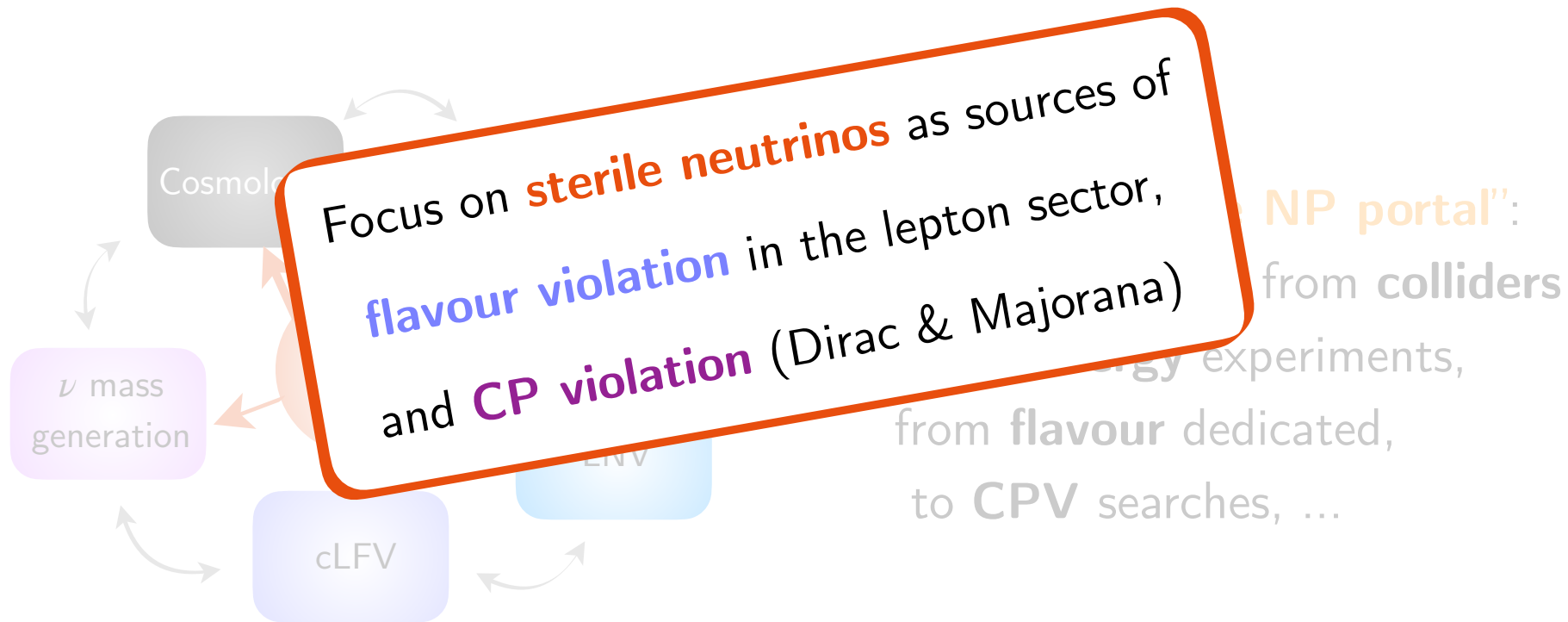


Potentially very **“visible NP portal”**:

↪ Extensive imprints from **colliders** to **low-energy** experiments, from **flavour** dedicated, to **CPV** searches, ...

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Sterile fermions ($\sim \nu_R$ RH neutrinos):

(minimal) SM extension to account for **neutrino masses & mixings**

- ↪ Interactions with SM fields only through **mixings** with active neutrinos
- ↪ No bound on the **number** and **mass scale** of the sterile states
- ↪ Common to numerous NP models, wide range of scales $\Lambda_{EW} \rightarrow \Lambda_{GUT}$

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High-scale type I seesaw

$$m_{\nu_s} = \mathcal{O}(10^{10-15} \text{ GeV})$$

“natural” $Y^\nu \sim 1$

Decoupled new states

“**Low-scale**” seesaw

$$m_{\nu_s} = \mathcal{O}(\text{MeV} - \text{TeV})$$

Finetune Y^ν

New states **within experimental reach!**

⇒ Phenomenological implications strongly depend on their **masses**

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Low-scale seesaw ↪ non-decoupled states, **modified lepton currents!**

Rich phenomenology at low-energies, high-intensity and colliders

EW precision tests, cLFV transitions and decays, $0\nu 2\beta$ decays, rare meson decays, ...

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Low-scale seesaw ↪ **enhanced lepton currents!**

Rich phenomenology: intensity and colliders

EW precision tests and decays, $0\nu 2\beta$ decays, rare meson decays, ...

Testability!

Constructing simplified models

Minimal “toy model” for **phenomenological** analyses: SM + ν_s

- ↪ **Ad-hoc** construction: extend SM with n_s **Majorana massive** states leading to **new mixings** and **CPV phases** (Dirac and Majorana)
 - ↪ No assumption on the **mass generation mechanism**
 - ↪ **Well-defined interactions** in physical basis
- ⇒ Explore the **low-energy phenomenology** common to complete models (type I seesaw, ISS, ...)

Focus on **sterile** fermions and **cLFV observables**

Minimal “toy model” for phenomenological analyses: SM + 1 ν_s

Active-sterile mixing $\mathcal{U}_{\alpha i}$

$$|n_L\rangle = \mathcal{U}_{4\times 4} |\nu_i\rangle$$

Left-handed lepton mixing $\tilde{\mathcal{U}}_{\text{PMNS}}$
3 × 3 sub-block, **non-unitary!**

⇒ **Modified** charged
& **neutral lepton currents!**

$$\mathcal{U}_{4\times 4} = \begin{pmatrix} \mathcal{U}_{e1} & \mathcal{U}_{e2} & \mathcal{U}_{e3} & \mathcal{U}_{e4} \\ \mathcal{U}_{\mu 1} & \mathcal{U}_{\mu 2} & \mathcal{U}_{\mu 3} & \mathcal{U}_{\mu 4} \\ \mathcal{U}_{\tau 1} & \mathcal{U}_{\tau 2} & \mathcal{U}_{\tau 3} & \mathcal{U}_{\tau 4} \\ \mathcal{U}_{s1} & \mathcal{U}_{s2} & \mathcal{U}_{s3} & \mathcal{U}_{s4} \end{pmatrix}$$

Physical parameters:

- ↪ 4 masses: 3 light mostly active & 1 heavy mostly sterile
- ↪ 6 mixing angles
- ↪ 6 CPV phases (3 Dirac δ_{ij} and 3 Majorana φ_i)

Minimal “toy model” for phenomenological analyses: **SM + 2** ν_s

Active-sterile mixing $\mathcal{U}_{\alpha i}$

Left-handed lepton mixing $\tilde{\mathcal{U}}_{\text{PMNS}}$

3 × 3 sub-block, **non-unitary!**

$$|n_L\rangle = \mathcal{U}_{5 \times 5} |\nu_i\rangle$$

$$\mathcal{U}_{5 \times 5} =$$

$$\begin{pmatrix} \mathcal{U}_{e1} & \mathcal{U}_{e2} & \mathcal{U}_{e3} & \mathcal{U}_{e4} & \mathcal{U}_{e5} \\ \mathcal{U}_{\mu1} & \mathcal{U}_{\mu2} & \mathcal{U}_{\mu3} & \mathcal{U}_{\mu4} & \mathcal{U}_{\mu5} \\ \mathcal{U}_{\tau1} & \mathcal{U}_{\tau2} & \mathcal{U}_{\tau3} & \mathcal{U}_{\tau4} & \mathcal{U}_{\tau5} \\ \mathcal{U}_{s1} & \mathcal{U}_{s2} & \mathcal{U}_{s3} & \mathcal{U}_{s4} & \mathcal{U}_{s5} \\ \mathcal{U}_{s'1} & \mathcal{U}_{s'2} & \mathcal{U}_{s'3} & \mathcal{U}_{s'4} & \mathcal{U}_{s'5} \end{pmatrix}$$

⇒ **Modified** charged
& **neutral lepton currents!**

⇒ **Sizeable contributions** to **cLFV observables** (already present in 3 + 1)

⇒ **Interference effects** between heavier states expected

Constructive & **destructive** interference effects

in cLFV leptonic and boson decays!

Minimal “toy model” for phenomenological analyses: SM + $2\nu_s$

Parametrise \mathcal{U} via a series of 10 (complex) **rotations** R_{ij} and a diagonal matrix including the **4 physical Majorana phases** φ_i

$$\mathcal{U}_{5 \times 5} = \begin{pmatrix} \mathcal{U}_{e1} & \mathcal{U}_{e2} & \mathcal{U}_{e3} & \mathcal{U}_{e4} & \mathcal{U}_{e5} \\ \mathcal{U}_{\mu1} & \mathcal{U}_{\mu2} & \mathcal{U}_{\mu3} & \mathcal{U}_{\mu4} & \mathcal{U}_{\mu5} \\ \mathcal{U}_{\tau1} & \mathcal{U}_{\tau2} & \mathcal{U}_{\tau3} & \mathcal{U}_{\tau4} & \mathcal{U}_{\tau5} \\ \mathcal{U}_{s1} & \mathcal{U}_{s2} & \mathcal{U}_{s3} & \mathcal{U}_{s4} & \mathcal{U}_{s5} \\ \mathcal{U}_{s'1} & \mathcal{U}_{s'2} & \mathcal{U}_{s'3} & \mathcal{U}_{s'4} & \mathcal{U}_{s'5} \end{pmatrix}$$

$$\mathcal{U} = R_{45} R_{35} R_{25} R_{15} R_{34} R_{24} R_{14} R_{23} R_{13} R_{12} \times \text{diag}(1, e^{i\varphi_2}, e^{i\varphi_3}, e^{i\varphi_4}, e^{i\varphi_5})$$

With, for illustration:

Mixing parameters θ_{ij}

$$R_{45} = \begin{pmatrix} \mathbb{1}_{3 \times 3} & \mathbb{0}_{2 \times 3} \\ \mathbb{0}_{3 \times 2} & \Theta_{45} \end{pmatrix} \quad \Theta_{45} = \begin{pmatrix} \cos \theta_{45} & \sin \theta_{45} e^{-i\delta_{45}} \\ -\sin \theta_{45} e^{i\delta_{45}} & \cos \theta_{45} \end{pmatrix}$$

Dirac phases δ_{ij}

Abada and Toma, [1511.03265]

Minimal “toy model” for **phenomenological** analyses: **SM** + **$2\nu_s$**

2 heavy sterile states with **masses m_4 and m_5** , **leptonic mixing $\mathcal{U}_{5\times 5}$**

CPV phases (**Dirac δ** and/or **Majorana φ**)

Simplified

Illustrative (*simplified*) approach

⇒ **No experimental constraint**

↪ Assume **degenerate masses**

$$m_4 = m_5$$

↪ Assume **degenerate mixing**

$$\text{angles } \theta_{\alpha 4} = \theta_{\alpha 5}$$

↪ Unconsidered **phases** set to 0

Illustrate the **impact** of **phases**

Full analysis

Take into account all available experimental constraints ⇒ **Full phenomenological study**

↪ Limits on **active-sterile mixings**

↪ Negative results of searches for **sterile** states

↪ **Electroweak precision** tests

↪ Bounds on searches for other **cLFV transitions**

⇒ **No assumptions** on **active-sterile mixings**
& all **CPV phases** randomly varied

Minimal “toy model” for phenomenological analyses: $\text{SM} + 2\nu_s$

2 heavy **sterile** states with masses m_4 and m_5 , leptonic mixing $\mathcal{U}_{5\times 5}$

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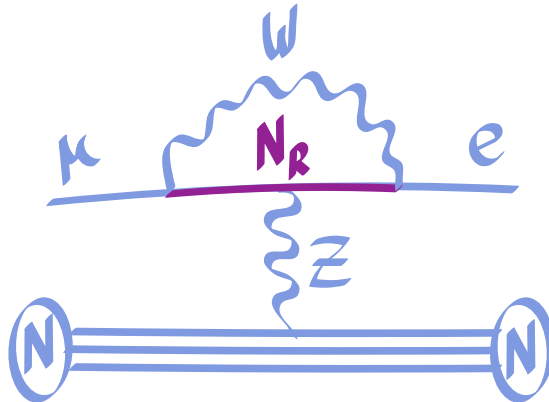
(Leptonic) cLFV
with CPV phases

cLFV: $\mu - e$ conversion in nuclei with **CPV Dirac and Majorana phases**

3 + 2 heavy ν_s : simplified approach $\sin \theta_{\alpha 4} = \sin \theta_{\alpha 5}$, $m_4 = m_5 = 1$ TeV

Simplified

Both **destructive** and **constructive** interference effects



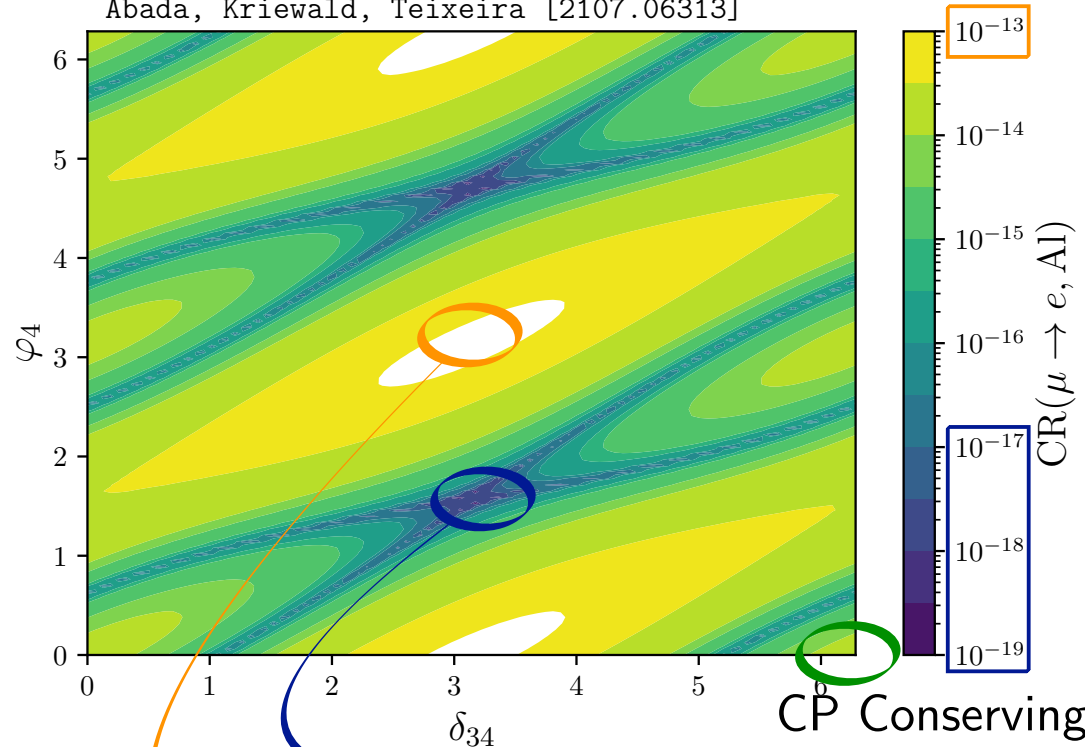
Joint effect of **Dirac** (δ_{34}) and **Majorana** (φ_4) **CPV phases**

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Abada, Kriewald, Teixeira [2107.06313]



Both **destructive** and **constructive** interference effects

Joint effect of **Dirac** (δ_{34}) and **Majorana** (φ_4) **CPV phases**

⇒ From beyond experimental sensitivity...
to within future reach...
and even **already excluded!**

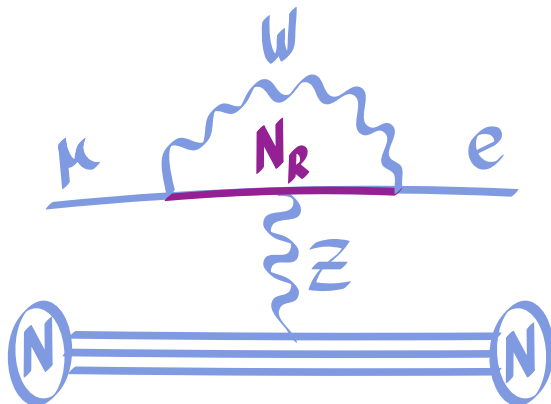
→ **CPV (max destructive interference)**

→ **CPV (max constructive interference)**

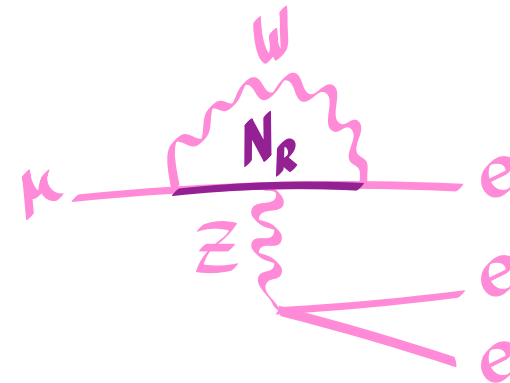
Ratios of cLFV observables to identify mediators & constrain their masses

An example: observables dominated by a **common topology** (Z-penguin)

$\mu - e$ conversion

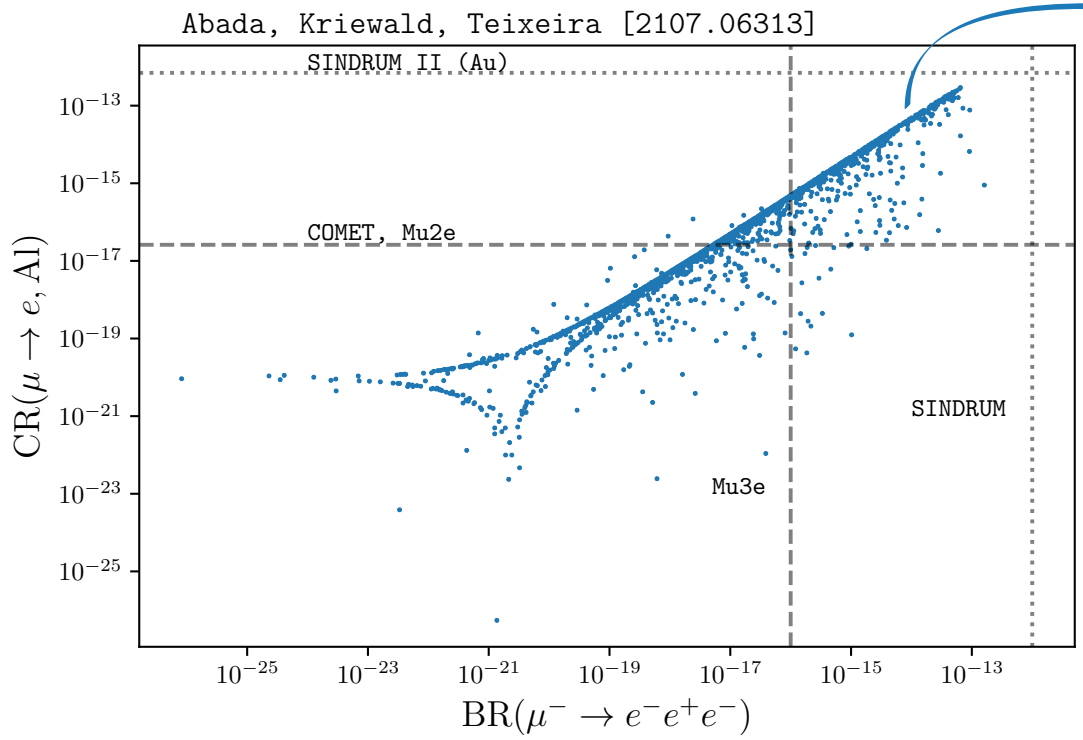


3-body muon decays ($\mu \rightarrow 3e$)



Ratios of cLFV observables to identify mediators & constrain their masses
Heavy sterile neutrino masses fixed to $m_4 = m_5 = 1 \text{ TeV}$

Full analysis



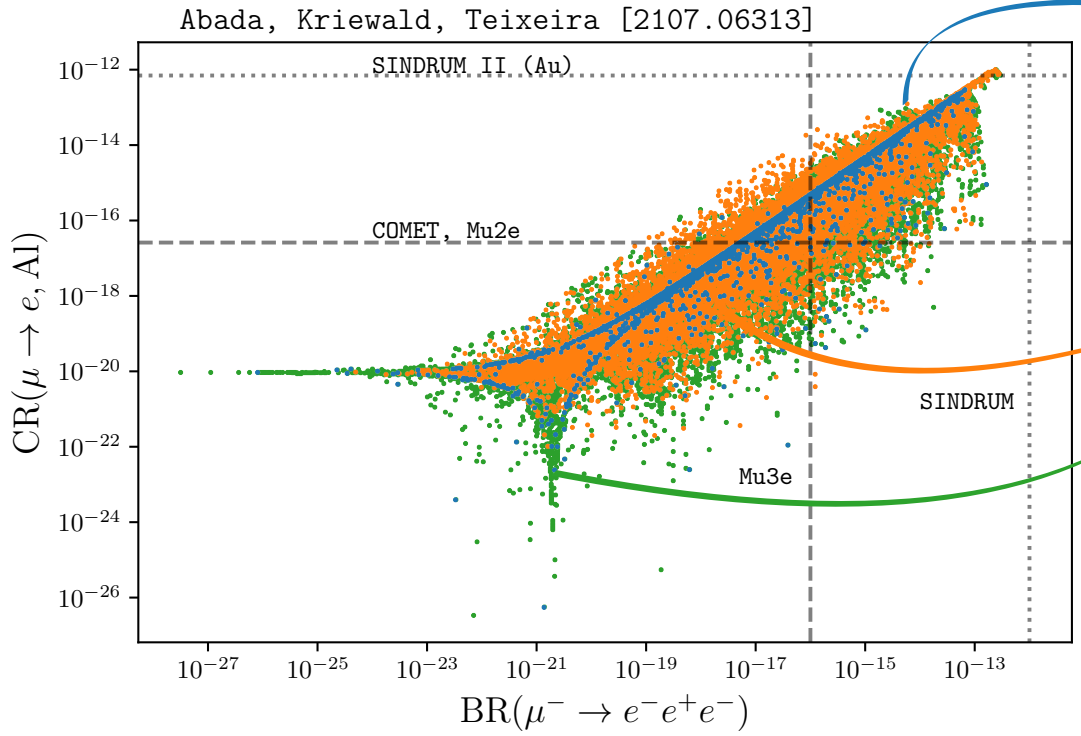
Strong correlation
(CP conserving)

Observation of $\mu \rightarrow 3e$
 \implies observation of
 $\mu - e$ conversion

Ratios of cLFV observables to identify mediators & constrain their masses

But **CP violating phases do matter!** And impact naïve expectations....

Full analysis



Strong correlation
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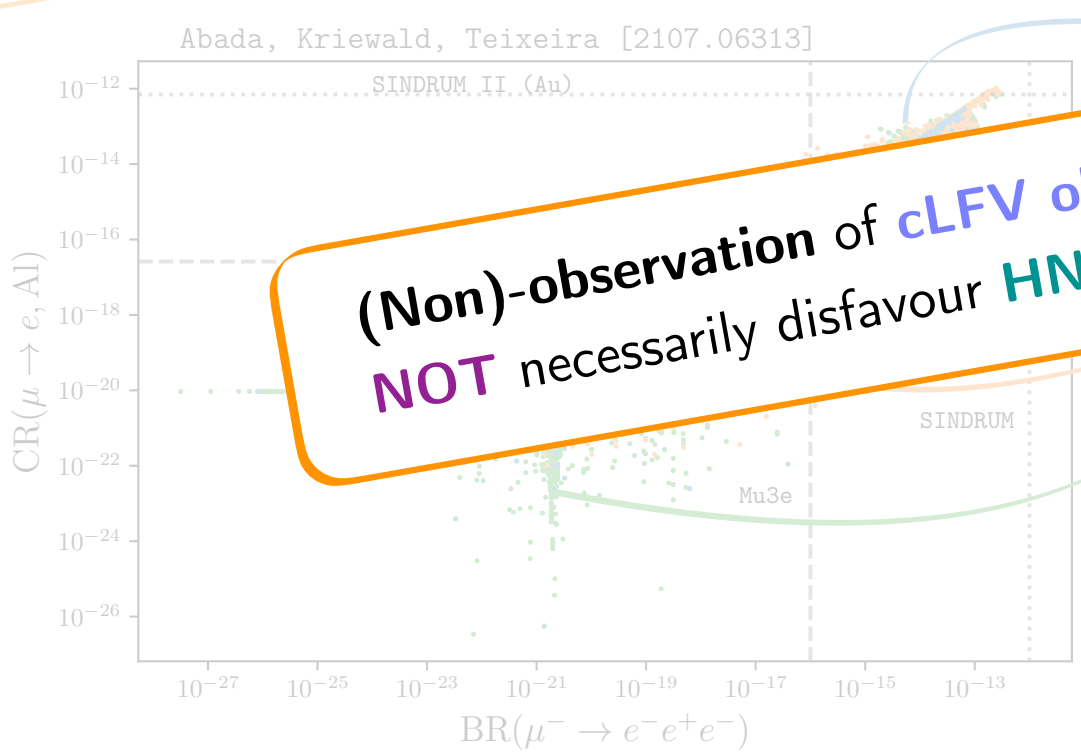
Loss of correlation!
(CP violating)

Observation of $\mu \rightarrow 3e$
 \Rightarrow observation of $\mu - e$ conversion

Ratios of cLFV observables to identify mediators & constrain their masses

But **CP violating phases do matter!** And impact naïve expectations....

Full analysis



(Non)-observation of cLFV observable(s)
NOT necessarily disfavour HNL extension!

Strong correlation

(conserving)

Loss of correlation!

(CP violating)

Observation of $\mu \rightarrow 3e$

\nRightarrow observation of

$\mu - e$ conversion

Impact of **CPV phases** regarding **experimental** prospects

Some *illustrative* benchmark points - **CP conserving** vs **CP violating**

Abada, Kriewald, Teixeira [2107.06313]

	BR($\mu \rightarrow e\gamma$)	BR($\mu \rightarrow 3e$)	CR($\mu - e, Al$)	BR($\tau \rightarrow 3\mu$)	BR($Z \rightarrow \mu\tau$)
P_1	3×10^{-16} ○	1×10^{-15} ✓	9×10^{-15} ✓	2×10^{-13} ○	3×10^{-12} ○
P'_1	1×10^{-13} ✓	2×10^{-14} ✓	1×10^{-16} ✓	1×10^{-10} ✓	2×10^{-9} ✓
P_2	2×10^{-23} ○	2×10^{-20} ○	2×10^{-19} ○	1×10^{-10} ✓	3×10^{-9} ✓
P'_2	6×10^{-14} ✓	4×10^{-14} ✓	9×10^{-14} ✓	8×10^{-11} ✓	1×10^{-9} ✓
P_3	2×10^{-11} ✗	3×10^{-10} ✗	3×10^{-9} ✗	2×10^{-8} ✓	8×10^{-7} ✓
P'_3	8×10^{-15} ○	1×10^{-14} ✓	6×10^{-14} ✓	2×10^{-9} ✓	1×10^{-8} ✓

○ beyond future reach
 ✓ within future sensitivity
 ✗ conflicts current bounds

P_2 : only $\mu - \tau$ cLFV within future reach, cLFV μ decays beyond sensitivity...

P'_2 : all considered **cLFV transitions within reach!**

Observation of cLFV observable(s)
NOT necessarily disfavour **HNL extension!**

Impact of **CPV phases** regarding **experimental** prospects

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○ beyond future reach
 ✓ within future sensitivity
 ✗ conflicts current bounds

P_3 : large active-sterile mixings, excluded due to bounds on cLFV μ decays

P'_3 : suppression of rates from **CPV phases: reconcile large mixings with observation!**

CPV phases matter and must be included!

cLFV boson decays and CPV

Gauge bosons (Z, W) and Higgs decays are sensitive to **New Physics** including heavy **sterile** states!

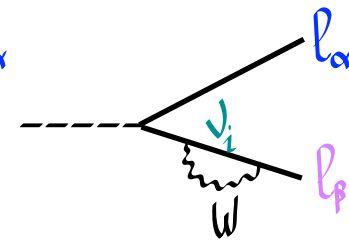
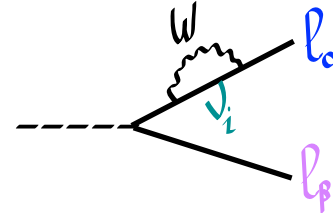
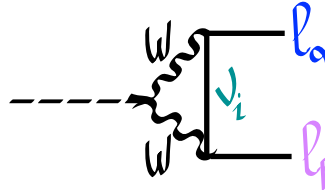
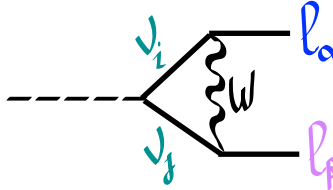
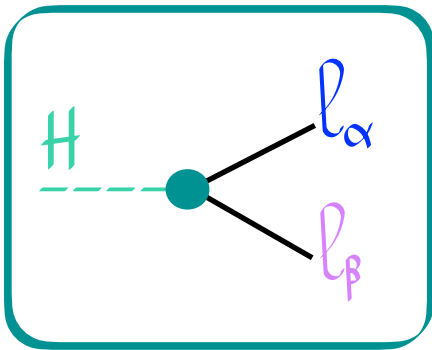
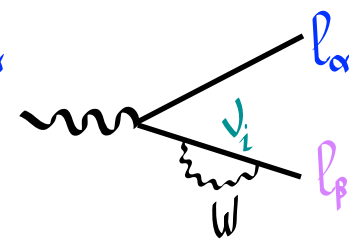
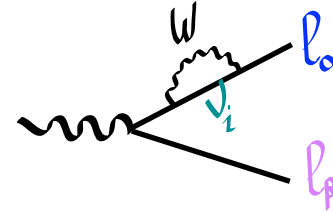
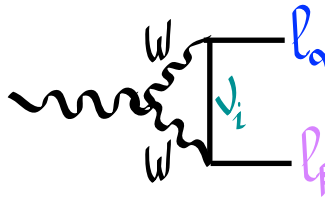
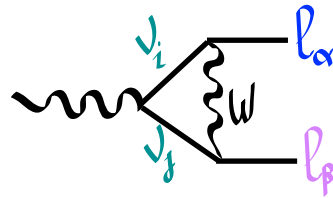
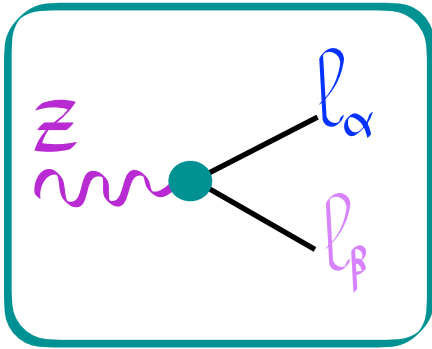
Significant contributions leading to **strong constraints**:

$$\Gamma_Z^{\text{inv}} \text{ and } Z \rightarrow \ell_\alpha \ell_\beta, H \rightarrow \ell_\alpha \ell_\beta$$

\implies What is the **impact** of **CPV Dirac & Majorana** phases on cLFV Z and Higgs decays?

Abada, Kriewald, EP, Rosauero, Teixeira [2207.10109]

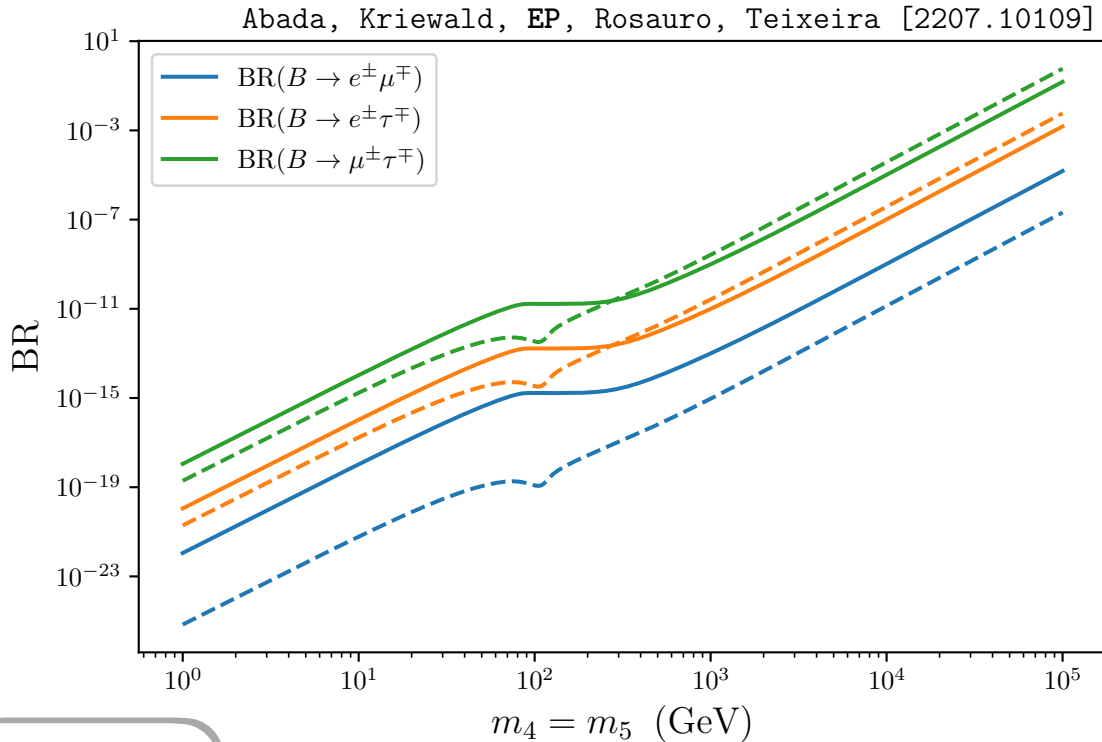
cLFV boson decays: $Z \rightarrow \ell_\alpha \ell_\beta$, $H \rightarrow \ell_\alpha \ell_\beta$
with **heavy sterile states** and **CPV phases**



Full computation of cLFV widths; both unitary & Feynman gauges for complete HNL models

$Z \rightarrow \ell_\alpha \ell_\beta$ and $H \rightarrow \ell_\alpha \ell_\beta$ with **heavy sterile states** (degenerate masses)

Simplified



No CPV phases included

— $Z \rightarrow \ell_\alpha \ell_\beta$
 - - - $H \rightarrow \ell_\alpha \ell_\beta$

Significant dependance of cLFV boson decays on **sterile masses**

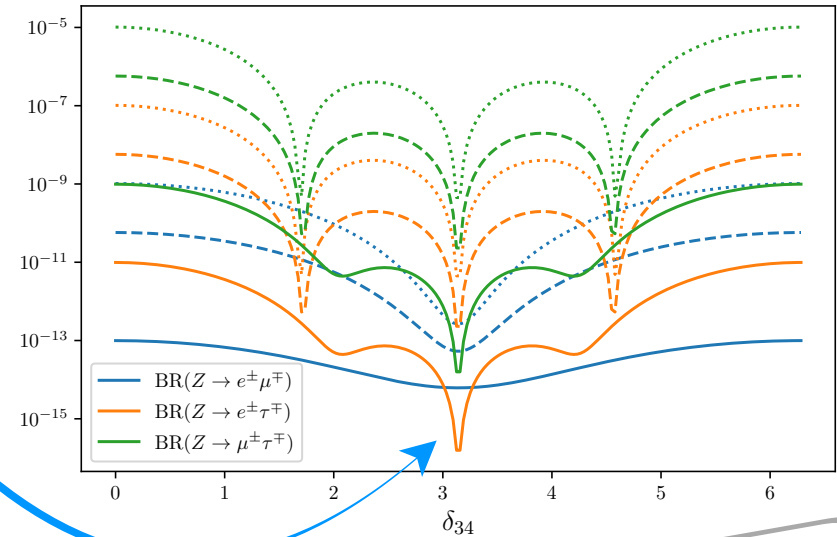
$\theta_{14} = \theta_{15} = 10^{-3}$
 $\theta_{24} = \theta_{25} = 10^{-2}$
 $\theta_{34} = \theta_{35} = 10^{-1}$

$$Z \rightarrow \ell_\alpha \ell_\beta$$

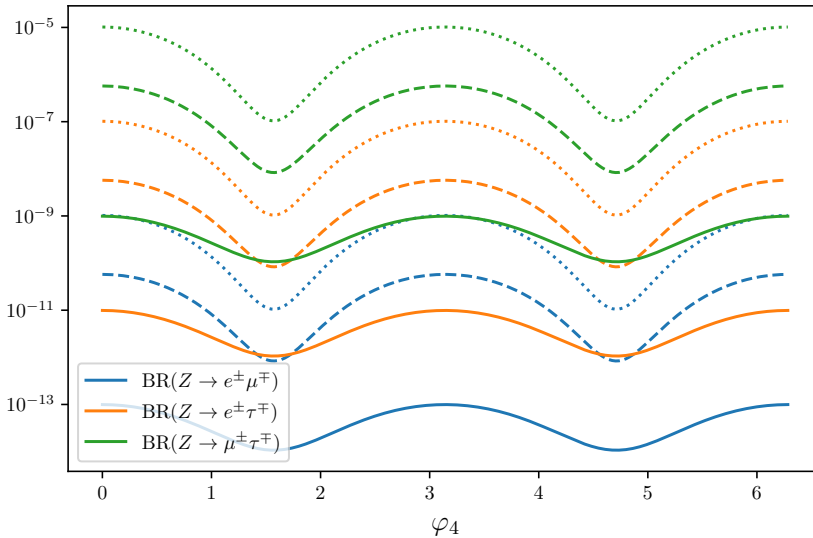
Strong dependence on **Dirac CPV phase** (δ_{34})

Possibility of decay rate **suppressions**

Abada, Kriewald, EP, Rosauero, Teixeira [2207.10109]



Abada, Kriewald, EP, Rosauero, Teixeira [2207.10109]



Simplified

- 1 TeV
- - 5 TeV
- 10 TeV

Sensitivity to **Majorana CPV phase** (φ_4)

These effects are amplified for larger HNL masses

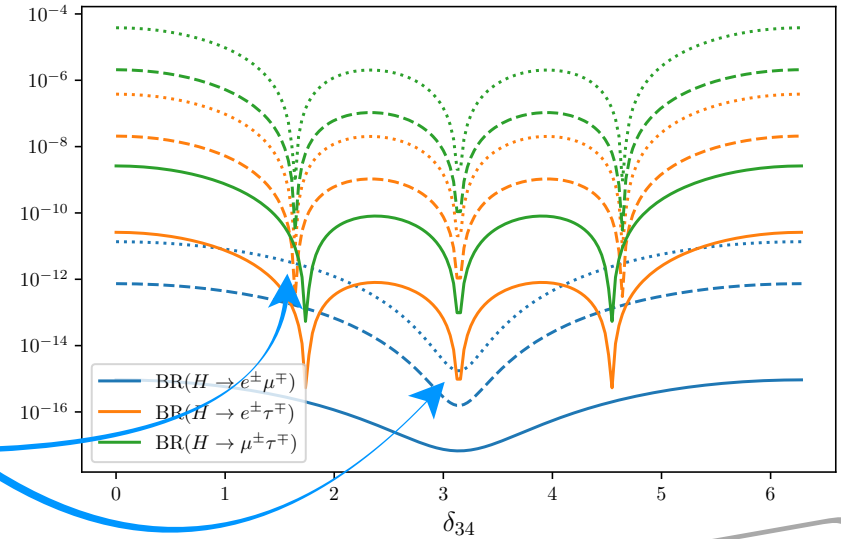
Including CPV phases

$$H \rightarrow \ell_\alpha \ell_\beta$$

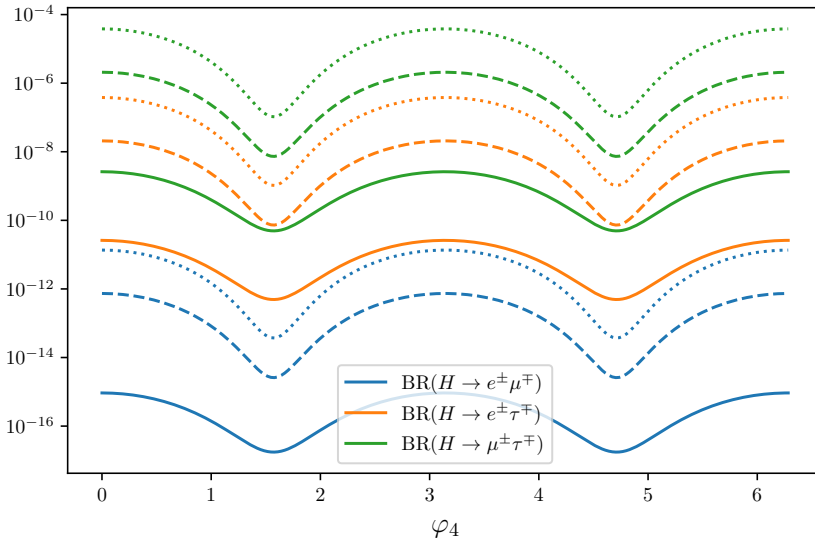
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Simplified

— 1 TeV
- - 5 TeV
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Varying all **CPV phases** associated with the **sterile states**

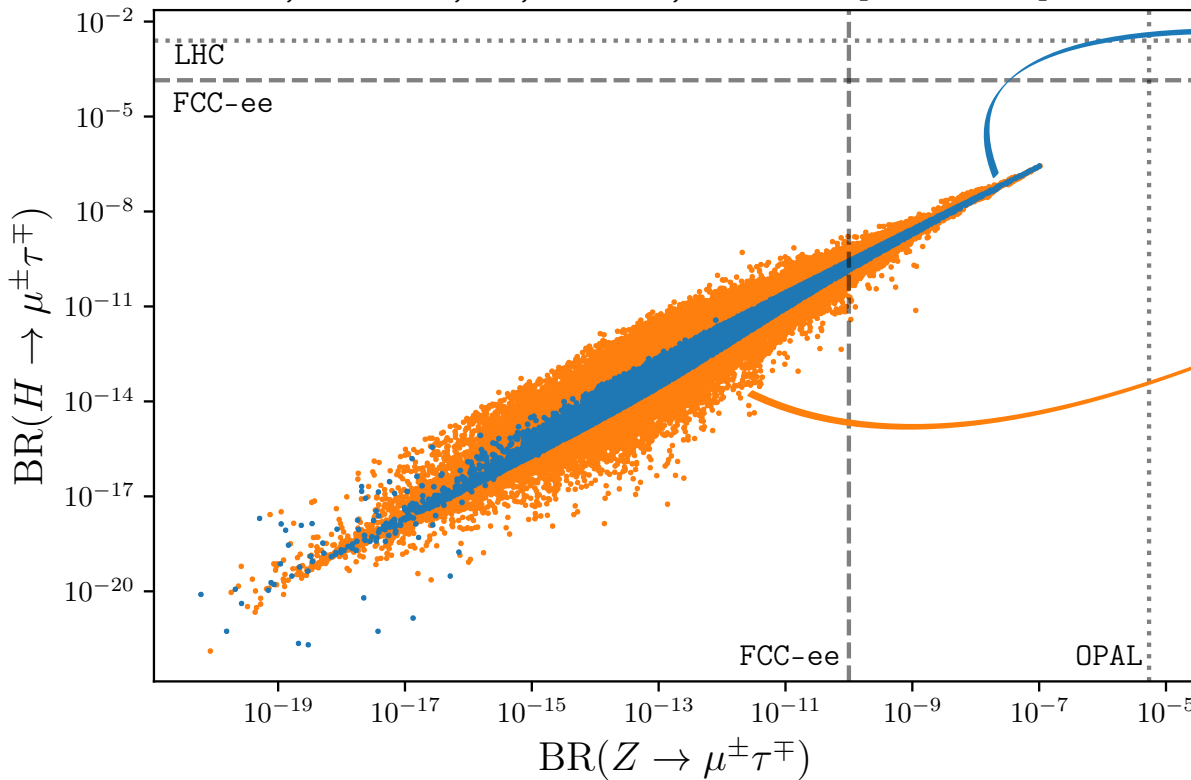
Full analysis

$$H \rightarrow \mu \tau \text{ vs. } Z \rightarrow \mu \tau$$

$$m_4 = 5 \text{ TeV}$$

$$m_5 - m_4 \in [10 \text{ MeV}, 1 \text{ TeV}]$$

Abada, Kriewald, EP, Rosauro, Teixeira [2207.10109]



Strong correlation
(CP conserving)

Loss of correlation
(CP violating)
Constructive & destructive

**$Z \rightarrow \mu \tau$ within
future sensitivity**

$H \rightarrow \mu \tau$ beyond future experimental sensitivity

**An interesting observable:
CP asymmetries in Z-decays**

$Z \rightarrow \mu \tau$ decays potentially observable **AND** impacted by **CPV** phases

\Rightarrow Consider **CP-asymmetries**

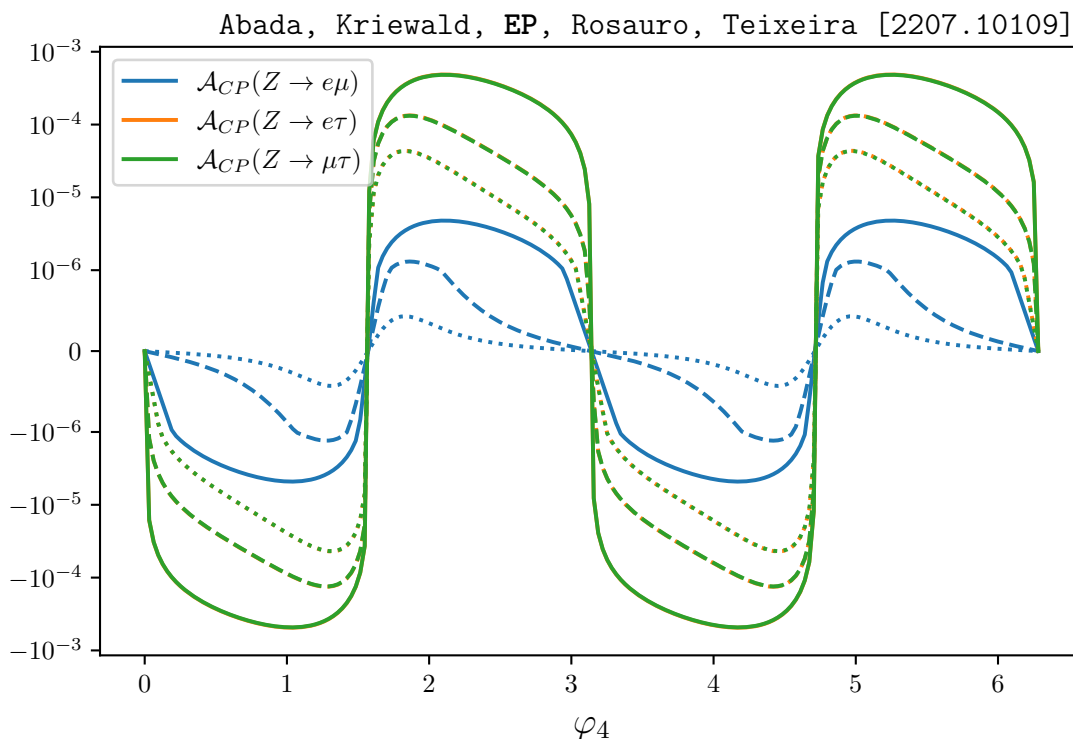
$$\mathcal{A}_{CP}(Z \rightarrow \ell_{\alpha} \ell_{\beta}) = \frac{\Gamma(Z \rightarrow \ell_{\alpha}^{-} \ell_{\beta}^{+}) - \Gamma(Z \rightarrow \ell_{\alpha}^{+} \ell_{\beta}^{-})}{\Gamma(Z \rightarrow \ell_{\alpha}^{-} \ell_{\beta}^{+}) + \Gamma(Z \rightarrow \ell_{\alpha}^{+} \ell_{\beta}^{-})}$$

Additional **observables** to ultimately **probe** the presence of **CPV**

- \rightsquigarrow Up to which extent can such a **minimal BSM** model be at the source of **non-vanishing contributions to CP-asymmetries?**
- \rightsquigarrow Contributions induced by *both* **Majorana** and **Dirac CPV** phases

Up to which extent can such a **minimal BSM** model be at the source of **non-vanishing contributions to CP-asymmetries?**

Simplified



Impact of **Majorana**
CPV phases
on **CP-asymmetry!**

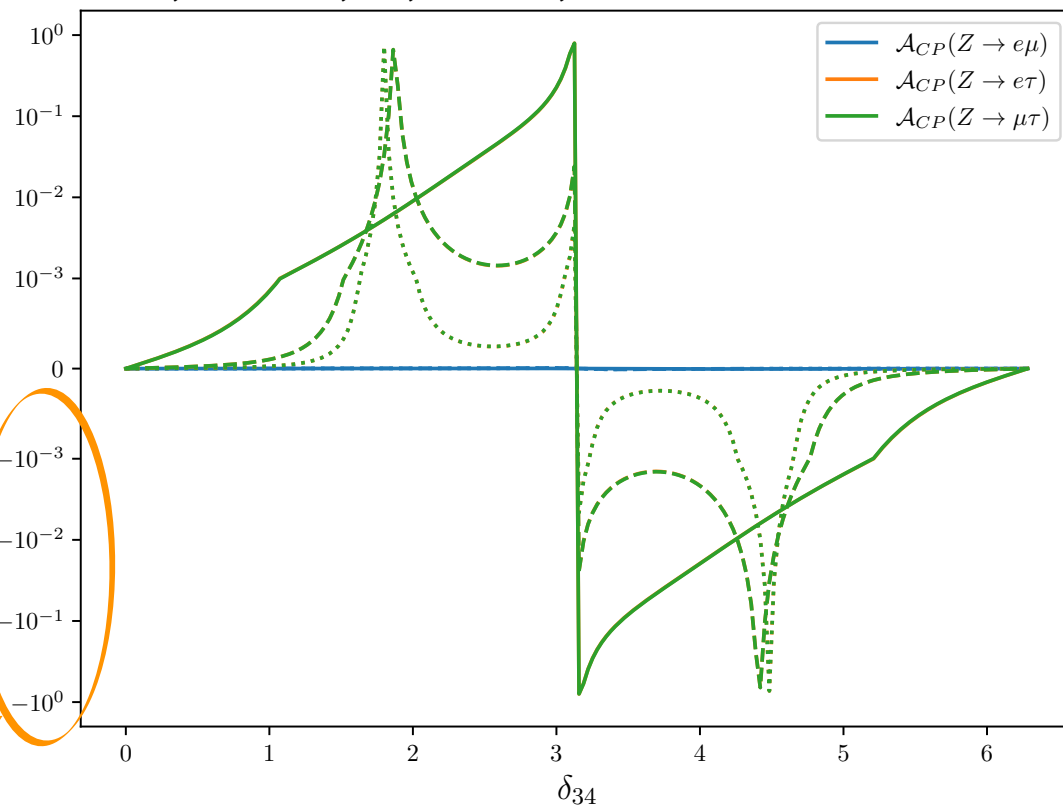
Up to which extent can such a **minimal BSM** model be at the source of **non-vanishing contributions to CP-asymmetries?**

Simplified

Impact of **Dirac CPV phases!**

Can lead to **very large CP-asymmetry!**

Abada, Kriewald, EP, Rosauero, Teixeira [2207.10109]

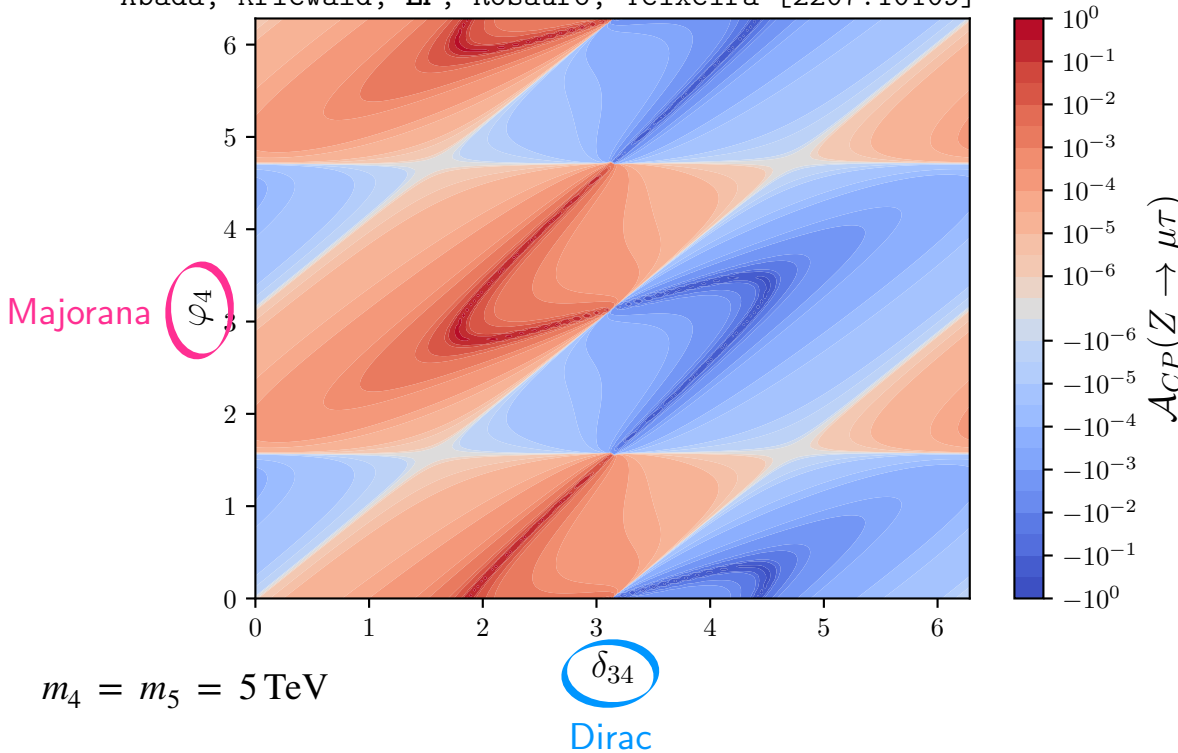


Up to which extent can such a **minimal BSM** model be at the source of **non-vanishing contributions to CP-asymmetries?**

Simplified

Contributions induced by *both* **Majorana** and **Dirac** CPV phases

Abada, Kriewald, EP, Rosauro, Teixeira [2207.10109]

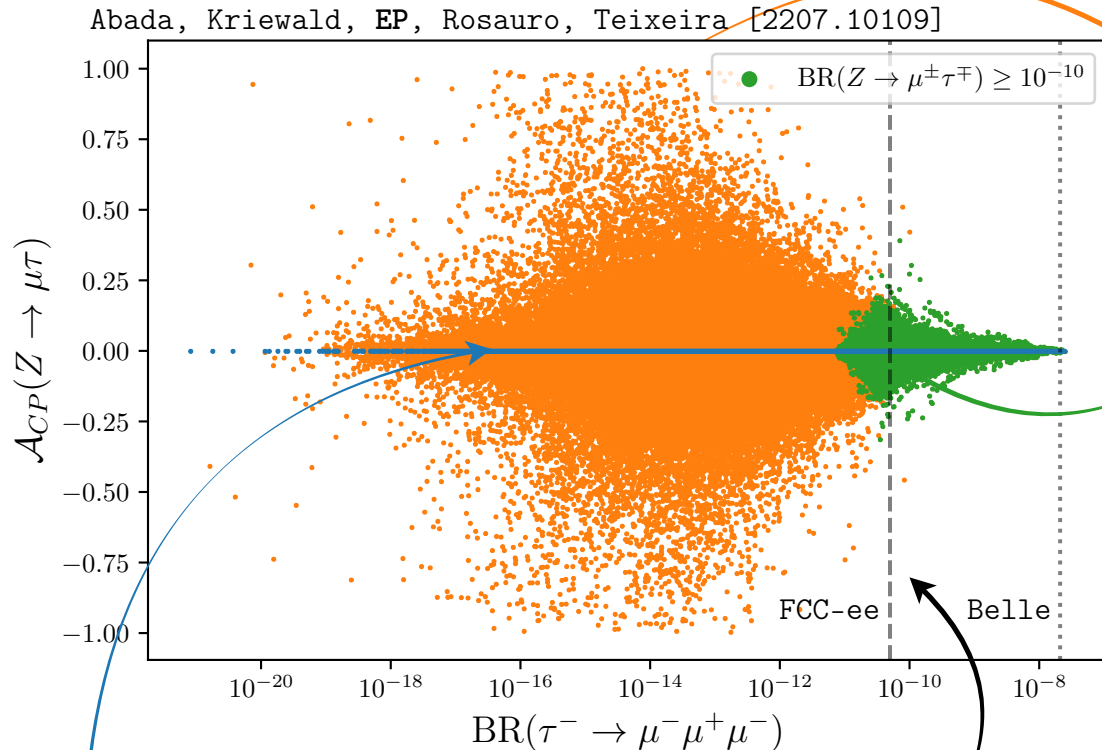


Provided that **Dirac** CPV phases are present
 \implies **Majorana** phases have a **significant impact on CP-asymmetries!**

CP-asymmetry vs. $\tau \rightarrow 3\mu$

Varying all **CPV phases** associated with the **sterile states** randomly

Full analysis



\mathcal{A}_{CP} can be as **large** as **100%**

$Z \rightarrow \mu\tau$

within **future sensitivity!**

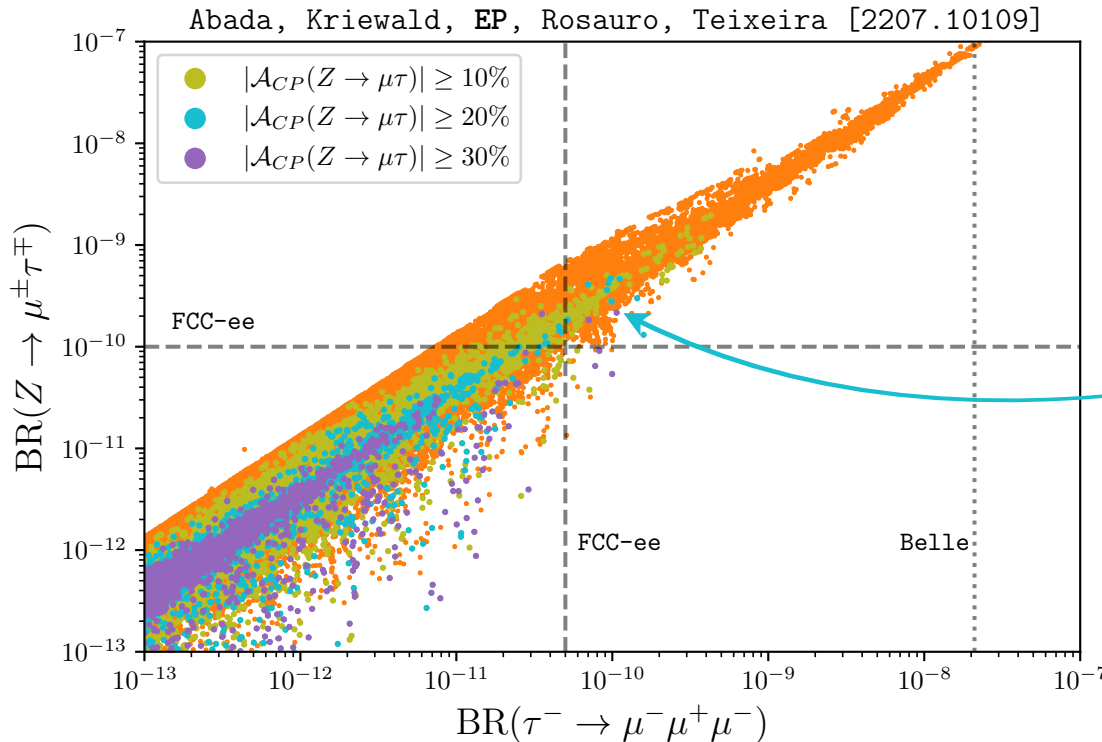
Associated \mathcal{A}_{CP}
up to **20%**

CP conserving

$\tau \rightarrow 3\mu$ also within **future reach**

Looking at the joint behaviour of $Z \rightarrow \mu\tau$, $\mathcal{A}_{CP}(Z \rightarrow \mu\tau)$ and $\tau \rightarrow 3\mu$

Full analysis



For $Z \rightarrow \mu\tau$ and $\tau \rightarrow 3\mu$
within future sensitivity

$|\mathcal{A}_{CP}(Z \rightarrow \mu\tau)|$
can reach $\geq 20\%$

If joint observation \implies highly suggestive of such an extension!

With at least 2 heavy Majorana fermions

Impact of (potential) measurement of the CP asymmetries

$$\begin{aligned} & m_4 = 5 \text{ TeV}, m_5 = 5.1 \text{ TeV}, \\ P_A \quad & s_{14} = -0.0028, s_{15} = 0.0045, s_{24} = -0.0052, s_{25} = -0.0037, s_{34} = -0.052, s_{35} = -0.028, \\ & \delta_{ij} = \varphi_i = 0, \end{aligned} \quad \text{CP Conserving}$$

$$\begin{aligned} & m_4 = 5 \text{ TeV}, m_5 = 5.1 \text{ TeV}, \\ P_B \quad & s_{14} = 0.00020, s_{15} = -7.1 \times 10^{-5}, s_{24} = -0.0024, s_{25} = 0.029, s_{34} = -0.073, s_{35} = -0.037, \\ & \delta_{14} = 0.71, \delta_{15} = 5.21, \delta_{24} = 2.06, \delta_{25} = 4.78, \delta_{34} = 3.80, \delta_{35} = 4.74, \varphi_4 = 1.77, \varphi_5 = 4.33. \end{aligned} \quad \text{CP Violating}$$

Both benchmark points P_A and P_B lead to **common cLFV predictions**:
with $\mu \rightarrow 3e$, $\mu - e$ conversion, $\tau \rightarrow 3\mu$ and $Z \rightarrow \mu\tau$ within future sensitivity

Indistinguishable if **cLFV** signals are observed

Impact of (potential) measurement of the CP asymmetries

$$m_4 = 5 \text{ TeV}, m_5 = 5.1 \text{ TeV},$$

$$P_A \quad s_{14} = -0.0028, s_{15} = 0.0045, s_{24} = 0.0028,$$

$$\delta_{ij} = \varphi_i = 0$$

BUT CP asymmetries in Z - boson decays offer a **clear distinction**:

$$P_B \text{ leads to } \mathcal{A}_{CP}(Z \rightarrow \mu\tau) = 30\%$$

\Rightarrow Can disentangle between **CP conserving** et **CPV** scenarios!

Both P_A and P_B lead to common **cLFV** predictions:
with $\mu \rightarrow 3e$, $\mu - e$ conversion, $\tau \rightarrow 3\mu$ and $Z \rightarrow \mu\tau$ within future sensitivity

Indistinguishable if **cLFV** signals are observed

In summary...

↪ **Minimal** and **simple** BSM construction:

SM + 2 **heavy Majorana sterile** states (no assumption on mass mechanism)

Low-energy phenomenology of complete (high-energy) models

↪ **Impact of the heavy steriles:**

Depends on **masses** & **mixings** with active states (**CPV**) \Rightarrow non unitary \tilde{U}_{PMNS}

↪ **cLFV: CPV phases** affect **correlations** & **interpretation** of exp data!

↪ **cLFV boson decays** sensitive to the presence of HNL: **CPV phases** have a **clear impact** on the decay rates (Dirac CPV striking reductions)

$Z \rightarrow \mu\tau$ within future sensitivity, large associated $\mathcal{A}_{CP} \Rightarrow$ Importance of taking **multiple observables** into account to **probe** CPV or CP conserving scenarios

CP asymmetry key to establish the presence of CP violation!

→ **Minimal** and **simple** BSM construction:

SM + 2 heavy **Majorana** sterile states (no assumption on mass mechanism)

Low-energy phenomenology of complete (high-energy) models

→ Impact of the **heavy steriles**:

Depends on **masses** & **mixings** with

→ **cLFV**: CP

→ **cLFV** bos

clear impact

$Z \rightarrow \mu\tau$ within future sensitivity, large associated $\mathcal{A}_{CP} \Rightarrow$ Importance of taking

multiple observables into account to **probe** CPV or CP conserving scenarios

CP violating phases do matter!
and should be taken into account

CP asymmetry key to establish the presence of **CP violation!**

Thank you for your attention

$$\mathcal{L}_{W^\pm} = -\frac{g_w}{\sqrt{2}} W_\mu^- \sum_{\alpha=1}^3 \sum_{j=1}^{3+n_S} \mathcal{U}_{\alpha j} \bar{l}_\alpha \gamma^\mu P_L \nu_j + \text{H.c.},$$

$$\mathcal{L}_{Z^0}^\nu = -\frac{g_w}{4 \cos \theta_w} Z_\mu \sum_{i,j=1}^{3+n_S} \bar{\nu}_i \gamma^\mu (P_L C_{ij} - P_R C_{ij}^*) \nu_j,$$

$$\mathcal{L}_{Z^0}^\ell = -\frac{g_w}{2 \cos \theta_w} Z_\mu \sum_{\alpha=1}^3 \bar{l}_\alpha \gamma^\mu (\mathbf{C}_V - \mathbf{C}_A \gamma_5) l_\alpha,$$

$$\mathcal{L}_{H^0} = -\frac{g_w}{4M_W} H \sum_{i \neq j=1}^{3+n_S} \bar{\nu}_i [C_{ij} (P_L m_i + P_R m_j) + C_{ij}^* (P_R m_i + P_L m_j)] \nu_j,$$

$$C_{ij} = \sum_{\rho=1}^3 u_{i\rho}^\dagger u_{\rho j}$$

Active mixings ($\theta_{\alpha\beta}$) and Dirac CPV δ_{13} : Central values of NuFIT 5.1 results

Active-sterile mixing angles $\theta_{\alpha 4,5}$ constrain from **low-** and **high-**energy observables:

(Semi-)leptonic τ decays
Light mesons leptonic decays } Construct ratios;
sensitivity to **modified** $W\ell\nu$ vertex

$$R_W^{\ell_1\ell_2} = \frac{\Gamma(W \rightarrow \ell_1\nu)}{\Gamma(W \rightarrow \ell_2\nu)} \quad \Gamma(Z \rightarrow \text{inv})$$

Upper bounds on the entries of η indirectly taking into account constrains from **modifications** of G_F , $\sin^2\theta_w$ and M_W

Bound on **HNL decay width** to comply with perturbative unitarity
 \implies bound on sterile **masses** and **couplings to active states**

$0\nu 2\beta$: upper limit on the **effective mass** m_{ee} from KamLAND-ZEN

For TeV-scale HNL, collider searches and cosmological bounds are not competitive

Heavier masses: assumed to be **sufficiently close** to allow for **interferences**

→ Fix m_4 and take random values of m_5 from half-normal distributions

(scale representative of the sterile states width)

Active-sterile mixing angles: **independently** varied & randomly varying signs

For $m_4 = 5 \text{ TeV}$, the range of parameters to be explored is:

$$\begin{aligned} m_5 - m_4 &\in [10 \text{ MeV}, 1 \text{ TeV}], \\ |\sin \theta_{14,15}| &\in [6.0 \times 10^{-5}, 6.0 \times 10^{-3}], \\ |\sin \theta_{24,25}| &\in [1.9 \times 10^{-4}, 0.036], \\ |\sin \theta_{34,35}| &\in [8.3 \times 10^{-4}, 0.13]. \end{aligned}$$

⇒ Correspond to regimes complying with **experimental data** for the **CP conserving case**

Analysis: Select randomly 10^4 points (consistent with experimental data), vary all **CPV phases** associated with sterile states $\delta_{\alpha 4,5}$, $\varphi_{4,5}$ for each tuple of mixing angles.

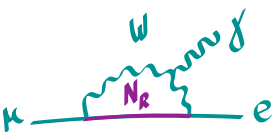
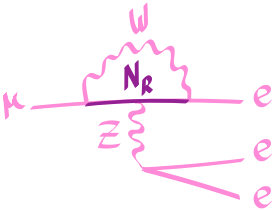
Consider only regimes that do not lead to **cLFV predictions** far away from the corresponding **future experimental sensitivity**

Synergy of cLFV observables very important: **Probe** different operators/topologies

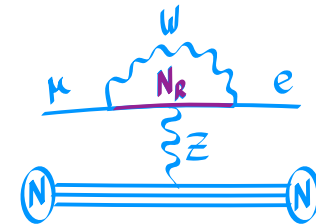
$BR(\mu \rightarrow e\gamma)$, $BR(\mu \rightarrow eee)$ and $CR(\mu - e, N)$ correlated by **common topologies**:

γ dipoles & anapoles, Z penguins, tree-level contributions, ... \Rightarrow 4-fermion operators

Model-dependent: certain topologies dominate, tree-level contributions might be present



Model	$\mu \rightarrow eee$	$\mu N \rightarrow eN$	$\frac{BR(\mu \rightarrow eee)}{BR(\mu \rightarrow e\gamma)}$	$\frac{CR(\mu N \rightarrow eN)}{BR(\mu \rightarrow e\gamma)}$
MSSM	Loop	Loop	$\approx 6 \times 10^{-3}$	$10^{-3} - 10^{-2}$
Type-I seesaw	Loop*	Loop*	$3 \times 10^{-3} - 0.3$	0.1-10
Type-II seesaw	Tree	Loop	$(0.1 - 3) \times 10^3$	$\mathcal{O}(10^{-2})$
Type-III seesaw	Tree	Tree	$\approx 10^3$	$\mathcal{O}(10^3)$
LFV Higgs	Loop [†]	Loop* [†]	$\approx 10^{-2}$	$\mathcal{O}(0.1)$
Composite Higgs	Loop*	Loop*	0.05 - 0.5	2 - 20

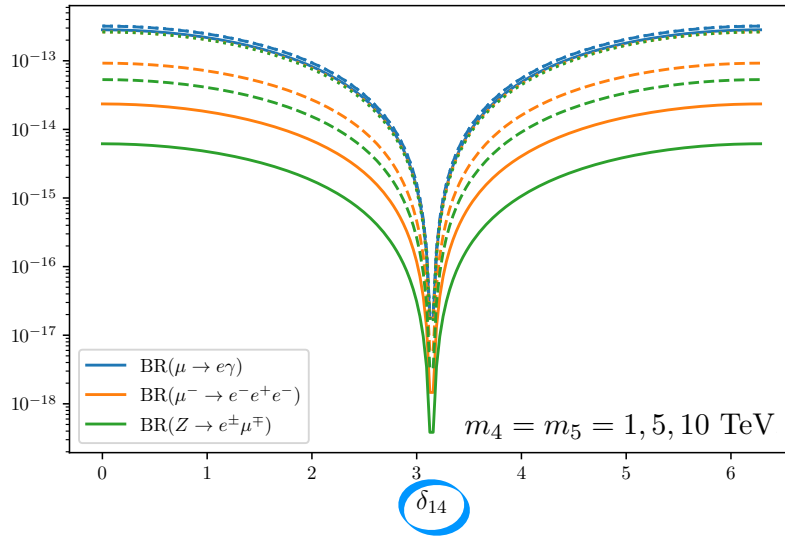


\Rightarrow study **correlations/ratios** of cLFV observables, might find **peculiar cLFV patterns**

\Rightarrow provide complementary information to direct searches

More on cLFV & CPV phases

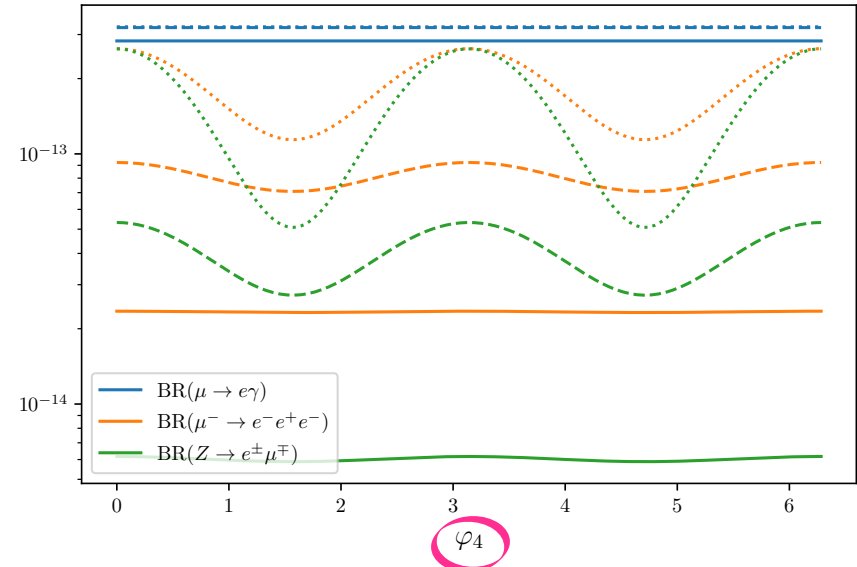
Abada, Kriewald, Teixeira [2107.06313]



Full cancellation of the rates for
 $\delta_{14} = \pi$
 Similar results for other **Dirac phases**

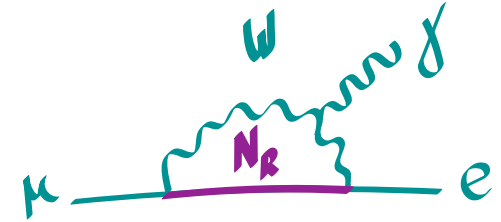
Milder dependence, γ -penguin
 independent of **Majorana phases**

Abada, Kriewald, Teixeira [2107.06313]



Radiative decays:

$$G_{\gamma}^{\mu e} = \sum_{i=4,5} \mathcal{U}_{ei} \mathcal{U}_{\mu i}^* G_{\gamma} \left(\frac{m_{N_i}^2}{m_W^2} \right)$$



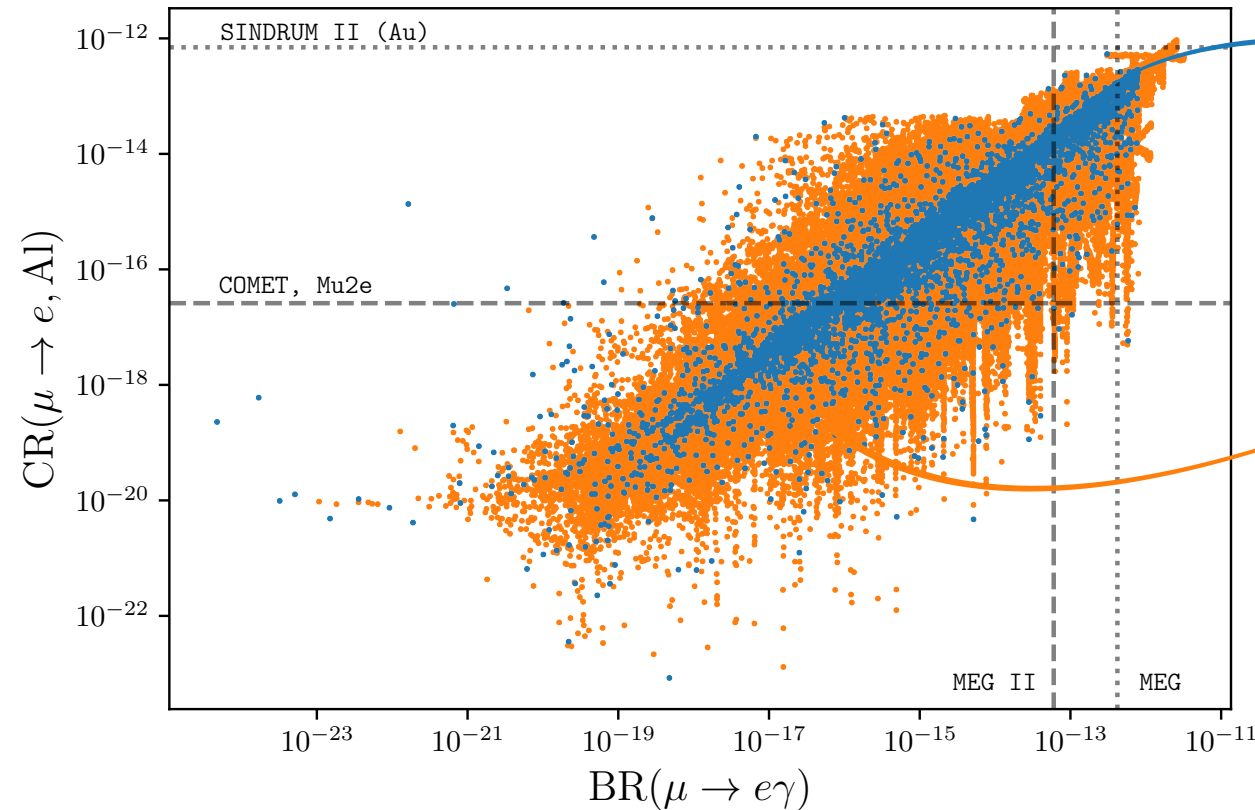
Assume (for *simplicity & illustrative purposes*): $m_4 \approx m_5$, $\sin \theta_{\alpha 4} \approx \sin \theta_{\alpha 5} \ll 1$

$$|G_{\gamma}^{\mu e}|^2 \approx 4 \sin^2 \theta_{e4} \sin^2 \theta_{\mu 4} \cos^2 \left(\frac{\delta_{14} + \delta_{25} - \delta_{15} - \delta_{24}}{2} \right) G_{\gamma} \left(\frac{m_{N_i}^2}{m_W^2} \right)$$

\Rightarrow **Radiative decays:** rate depends **only on Dirac phases**;
full cancellation for $\Sigma \delta = \pi$

More on cLFV & CPV phases

Abada, Kriewald, Teixeira [2107.06313]



Strong correlation
(CP conserving)

Loss of correlation!
(CP violating)