

# Looking Beyond the Standard Model with neutrons and nuclei

GDR Intensity Frontier  
2-4 Nov. 2022

G. Pignol  
M. Versteegen

# Why use n and nuclei ?

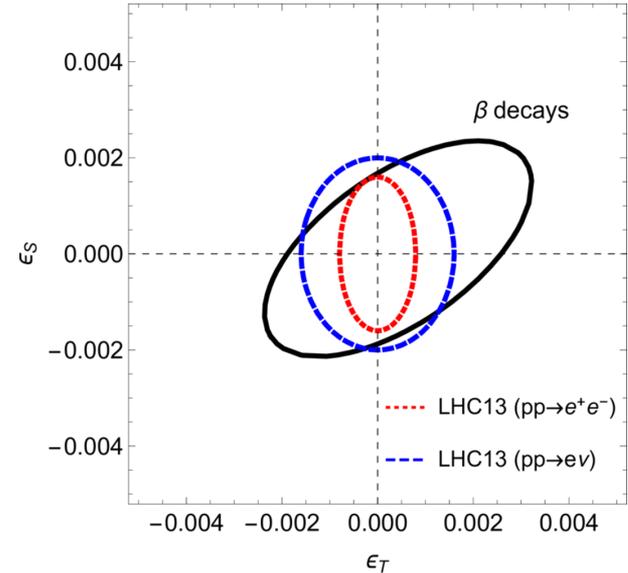
- **Beta decay**  
n and nuclei extensively used to establish the properties of the weak interaction in the framework of the SM
- **Effective Field Theory**  
Model independent approach : no assumption on NP origin

Wilson coefficients :

$$\epsilon_i \propto \left( \frac{m_W}{\Lambda} \right)^2 \sim 10^{-3}$$

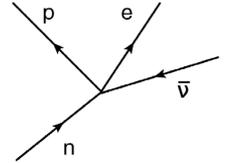
→ TeV NP scale

- ⇒ **Beta decay brings independent and competitive constraints to HEP in the weak sector when going to 0.1% level**



*M. González-Alonso, O. Naviliat-Cuncic, N. Severijns Prog. Part. Nucl. Phys. (2019)  
A. Falkowski, M. González-Alonso, O. Naviliat-Cuncic JHEP04 (2021)*

# Beta decay



- n beta decay Lagrangian

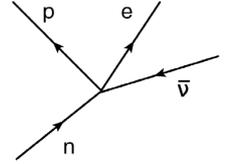
$$\begin{aligned} -\mathcal{L}_{LY} = & C_V \left( \bar{p} \gamma^\mu n + \frac{C_A}{C_V} \bar{p} \gamma^\mu \gamma_5 n \right) \times \bar{e} \gamma_\mu (1 - \gamma_5) \nu_e \\ & + C_S \bar{p} n \times \bar{e} (1 - \gamma_5) \nu_e + \frac{1}{2} C_T \bar{p} \sigma^{\mu\nu} n \times \bar{e} \sigma_{\mu\nu} (1 - \gamma_5) \nu_e + hc \\ & + \textit{right-handed neutrinos} \end{aligned}$$

SM “V-A” structure

Exotic currents : S and T  
P omitted

T.Lee, C-N Yang Phys. Rev. 104 (1956)  
M. González-Alonso, Colloque GANIL (2019)

# Beta decay



- n beta decay Lagrangian

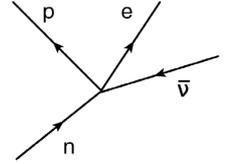
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 & \text{---} + \textit{right-handed neutrinos} \text{---} \quad \text{SM} \Rightarrow C_i = C_i'
 \end{aligned}$$

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# Beta decay



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 \end{aligned}$$

~~+ right-handed neutrinos~~ SM  $\Rightarrow C_i = C_i'$

**EFT**

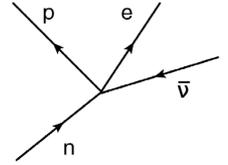
$$\begin{aligned}
 \bar{C}_V + \bar{C}'_V &= 2g_V (1 + \epsilon_L + \epsilon_R) \\
 \bar{C}_A + \bar{C}'_A &= -2g_A (1 + \epsilon_L - \epsilon_R) \\
 \bar{C}_S + \bar{C}'_S &= 2g_S \epsilon_S \\
 \bar{C}_P + \bar{C}'_P &= 2g_P \epsilon_P \\
 \bar{C}_T + \bar{C}'_T &= 8g_T \epsilon_T
 \end{aligned}$$

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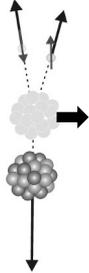
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SM “V-A” structure

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- **Decay rate distribution for polarized nuclei**



$$\frac{dW(\mathbf{J})}{dE_e d\Omega_e d\Omega_\nu} = dW_0 \times \xi \left\{ 1 + a \frac{\mathbf{p}_e \cdot \mathbf{p}_\nu}{E_e E_\nu} + b \frac{m_e}{E_e} + \frac{\langle \mathbf{J} \rangle}{J} \cdot \left( A \frac{\mathbf{p}_e}{E_e} + B \frac{\mathbf{p}_\nu}{E_\nu} + D \frac{\mathbf{p}_e \times \mathbf{p}_\nu}{E_e E_\nu} \right) \right\}$$

$a$   $\beta$ - $\nu$  correlation coefficient  
CP conserving  
Access to  $C_S$  and  $C_T$  quadratically

$b$  Fierz interference term  
CP conserving  
Access to  $C_S$  and  $C_T$  linearly

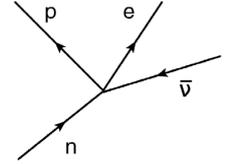
$D$  « D » coefficient  
CP violating  
Access to  $C_A, C_A', C_V, C_V'$  linearly

T.Lee, C-N Yang Phys. Rev. 104 (1956)

M. González-Alonso, Colloque GANIL (2019)

J.D Jackson, S.B Treiman, H.W Wylid Nuclear Phys 4 (1957)

# Beta decay



## ▪ n beta decay Lagrangian

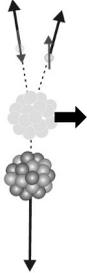
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## ▪ Decay rate distribution for polarized nuclei



$$\frac{dW(\mathbf{J})}{dE_e d\Omega_e d\Omega_\nu} = dW_0 \times \xi \left\{ 1 + a \frac{\mathbf{p}_e \cdot \mathbf{p}_\nu}{E_e E_\nu} + b \frac{m_e}{E_e} + \frac{\langle \mathbf{J} \rangle}{J} \cdot \left( A \frac{\mathbf{p}_e}{E_e} + B \frac{\mathbf{p}_\nu}{E_\nu} + D \frac{\mathbf{p}_e \times \mathbf{p}_\nu}{E_e E_\nu} \right) \right\}$$

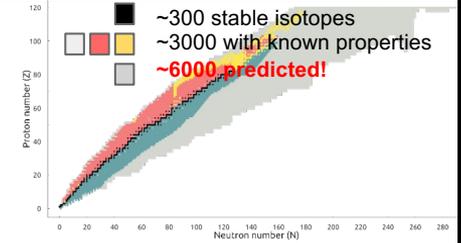
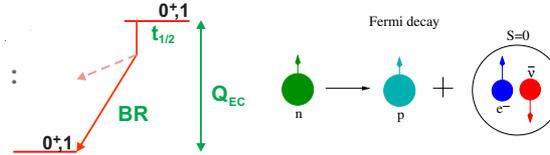
- Ft values :  $V_{ud}$ ,  $b$
- Beta spectrum shape :  $b$
- Correlation measurements :  $a$ ,  $b$ ,  $D$

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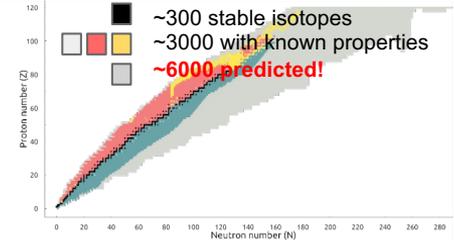
# Ft values : total decay rates

- Unitarity test of the CKM matrix 1<sup>st</sup> row :  $|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 1 + \text{NP}$

$0^+ \rightarrow 0^+$  superallowed Fermi transition :



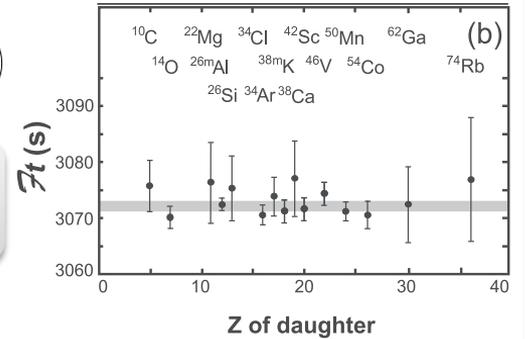
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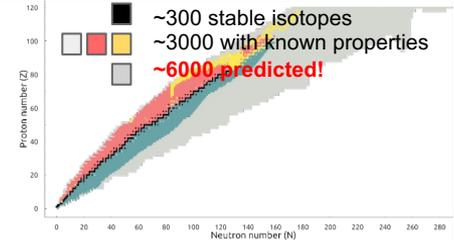


<p><b>Statistical rate function</b> <math>f \propto \int dW_0</math></p> <p>Trap : <math>Q_{EC}</math> <math>\Delta m \sim 10^{-8}</math></p>	<p><b>Partial half-life</b> <math>t = \frac{t_{1/2}}{BR}(1 + P_{EC})</math></p> <p>Beta counting and Ge with calibrated <math>\varepsilon</math> : <math>t_{1/2}</math> and BR <math>\Delta \varepsilon \sim 0.2\%</math></p>	<p><b>Corrections:</b> Radiative &lt; 1% Structure &lt; 1%</p> <p>Theoretical Calculations uncertainties &lt; 0.1% (except <math>^{62}\text{Ga}</math> &amp; <math>^{74}\text{Rb}</math>)</p>	<p><b><math>\mathcal{F}t = \frac{K}{2G_F^2 V_{ud}^2 (1 + \Delta V_R)}</math></b></p> <p>15 transitions with uncertainties &lt; 0.3%</p>
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222 individual measurements from 23 decays :  $|V_{ud}| = 0.97373 \pm 0.00031$

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$$\text{Statistical rate function } f \propto \int dW_0 \times \text{Partial half-life } t = \frac{t_{1/2}}{BR} (1 + P_{EC}) \times \text{Corrections: Radiative } < 1\% \text{ Structure } < 1\% = \mathcal{F}t = \frac{K}{2G_F^2 V_{ud}^2 (1 + \Delta V_R)}$$

Trap :  $Q_{EC}$   
 $\Delta m \sim 10^{-8}$

Beta counting and Ge with calibrated  $\varepsilon$  :  
 $t_{1/2}$  and BR  
 $\Delta \varepsilon \sim 0.2\%$

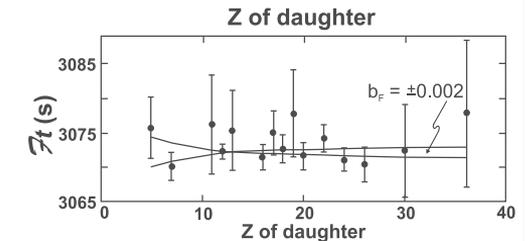
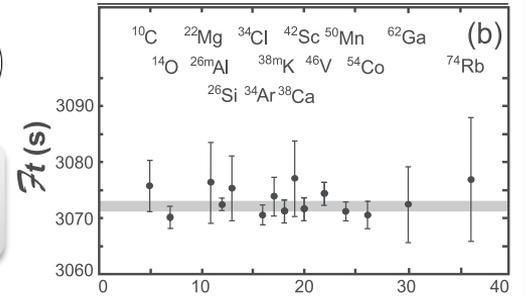
Theoretical Calculations  
Radiative < 1%  
Structure < 1%  
uncertainties < 0.1%  
(except  $^{62}\text{Ga}$  &  $^{74}\text{Rb}$ )

15 transitions with uncertainties < 0.3%

222 individual measurements from 23 decays :  $|V_{ud}| = 0.97373 \pm 0.00031$

- Sensitivity to exotic scalar currents :  $b = \pm 0.002$

If  $b \neq 0$  then  $f$  is affected :  $ft' = ft \times \frac{1}{1 + b < \frac{m_e}{E_e} >}$



JC. Hardy & IS. Towner Phys.Rev.C 102 (2020)

# Beta spectrum shape

- **Beta energy spectrum for non polarized nuclei :**

$$\frac{dW(\mathbf{J})}{dE_e d\Omega_e d\Omega_\nu} = dW_0 \times \xi \left\{ 1 + a \frac{\mathbf{p}_e \cdot \mathbf{p}_\nu}{E_e E_\nu} + b \frac{m_e}{E_e} + \frac{\langle \mathbf{J} \rangle}{J} \cdot \left( A \frac{\mathbf{p}_e}{E_e} + B \frac{\mathbf{p}_\nu}{E_\nu} + D \frac{\mathbf{p}_e \times \mathbf{p}_\nu}{E_e E_\nu} \right) \right\}$$

↪

$$dW = dW_0 \times \xi \left( 1 + b \frac{m}{E_e} \right)$$

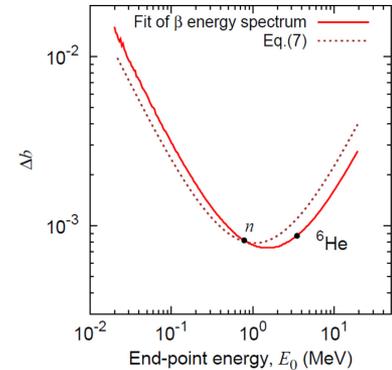
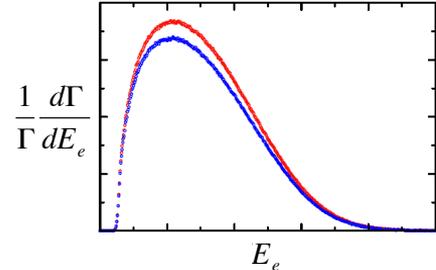
Highest sensitivity candidates due to kinematics : endpoint energy 1-4 MeV  
All theoretical corrections under control at 0.1% level

- **Experimental Challenges**

- Electron backscattering
- energy loss in source
- Detector dead layer

- **Set-ups : 4 $\pi$ , MWDC, CRES...**

Ongoing programs with  ${}^6\text{He}$ @LPC,  ${}^{114}\text{In}$ @Leuven,  ${}^{20}\text{F}$ ...



M. González-Alonso, O. Naviliat-Cuncic *Phys. Rev. C* 94 (2016)

L. Hayen et al, *Rev. Mod. Phys.* 90 (2018)

# Correlation measurements : WISARD

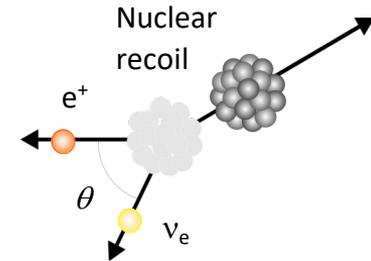


- Decay rate for non polarized nuclei

$$\frac{dW(\mathbf{J})}{dE_e d\Omega_e d\Omega_\nu} = dW_0 \times \xi \left\{ 1 + a \frac{\mathbf{p}_e \cdot \mathbf{p}_\nu}{E_e E_\nu} + b \frac{m_e}{E_e} + \frac{\langle \mathbf{J} \rangle \cdot \left( A \frac{\mathbf{p}_e}{E_e} + B \frac{\mathbf{p}_\nu}{E_\nu} + D \frac{\mathbf{p}_e \times \mathbf{p}_\nu}{E_e E_\nu} \right)}{J} \right\}$$



$$dW = dW_0 \times \xi \left( 1 + a \frac{\mathbf{p}_e \cdot \mathbf{p}_\nu}{E_e E_\nu} + b \frac{m_e}{E_e} \right)$$



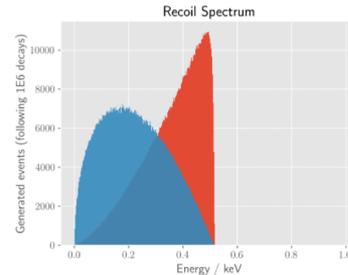
- Correlation measurement = recoil measurement

$a > 0$  :  $\theta = 0^\circ$  favored and large recoil

$a < 0$  :  $\theta = 180^\circ$  favored and small recoil

access to :  $\tilde{a} \sim \frac{a}{1 + b \frac{m_e}{E_e}}$

Beta nuclear recoil < 1 keV



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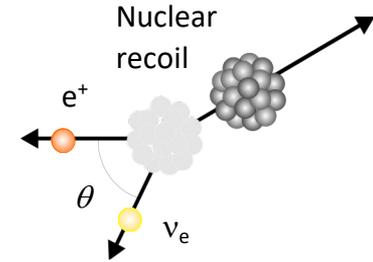


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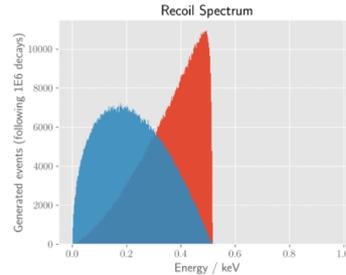


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Beta nuclear recoil < 1 keV



**Pure Fermi transition  $\Delta J=0 S=0$**

$$a_F \cong 1 - \frac{|C_S|^2 + |C'_S|^2}{|C_V|^2} = 1$$

$$b_F \approx \pm \text{Re} \left( \frac{C_S + C'_S}{C_V} \right) = 0$$

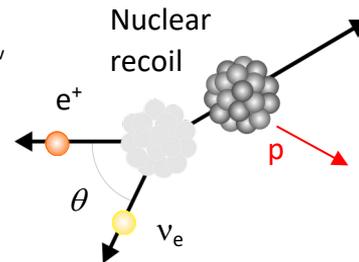
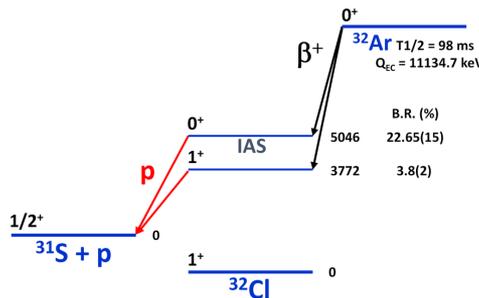
⇒ Best measurement  $a_F$  at 0.48%

# Correlation measurements : WISArD



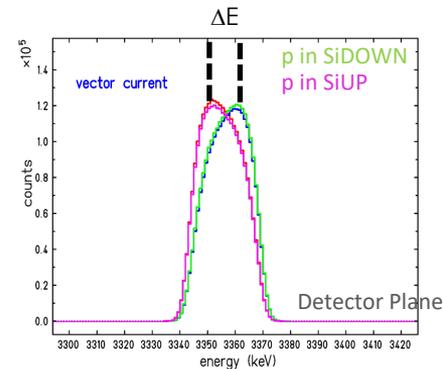
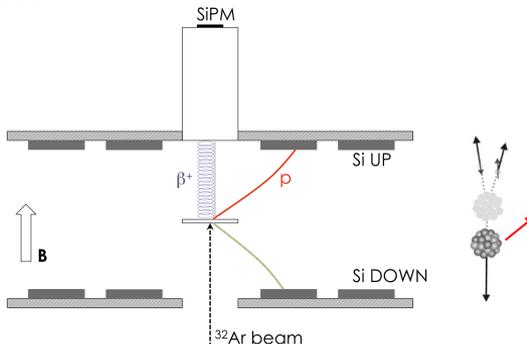
## β-delayed p emission in <sup>32</sup>Ar

- Fermi  $0^+ \rightarrow 0^+$  transition from GS to IAS
- Recoil energy  $\sim 100$ s eV
- Beta delayed p emission  $\sim 3$  MeV
- IAS :  $\Gamma \sim 20$  eV  $\Leftrightarrow T_{1/2} \sim 10^{-17}$  s
- ⇒ p emission in flight from the recoil

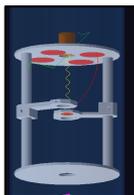
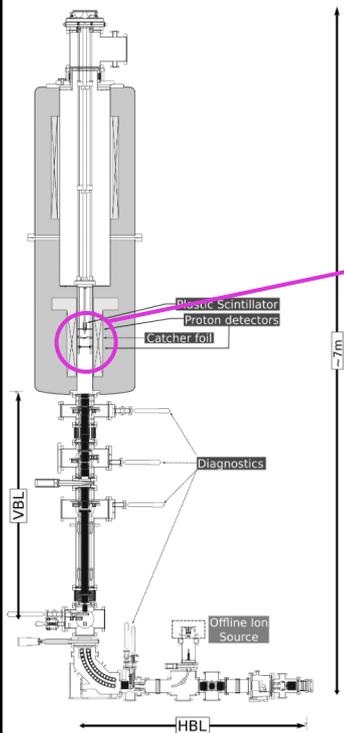


## β-p coincidence measurement

- Beta detector with detection threshold below 10 keV
- Strong magnetic field
- 2 symmetrical p detectors with resolution  $< 15$  keV and high solid angle



# WISArD at ISOLDE

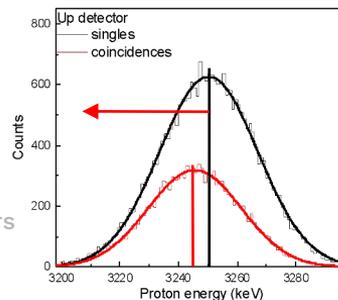


$\beta$  detector: plastic scintillator and SiPM

6  $\mu\text{m}$  mylar Catcher

2 x 4 proton detectors  
(300  $\mu\text{m}$  Si Detectors)

in B = 4 T  
(WITCH magnet) + FASTER DAQ



## Proof-of-principle (2018)

- Readily available  $\beta$  and p detectors
- $\sim 1700$  pps of  $^{32}\text{Ar}$  instead of 3000 nominal
- $\sim 35\text{h}$  of beamtime

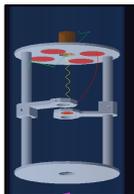
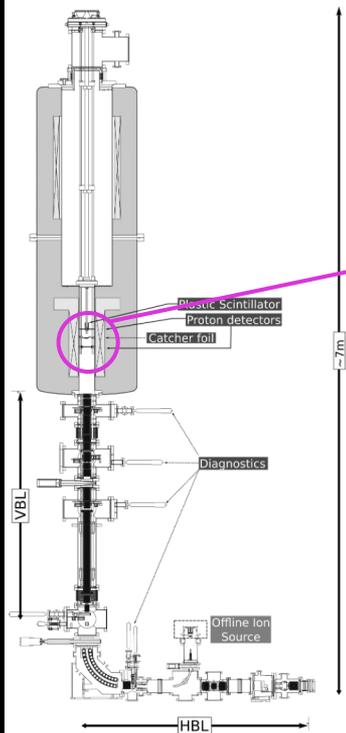
$$\Delta E_F = 4.49(3) \text{ keV}$$

$$\tilde{\alpha}_F = 1.007(32)_{\text{stat}}(25)_{\text{syst}}$$

◊ 3rd best result



# WISArD at ISOLDE



$\beta$  detector: plastic scintillator and SiPM

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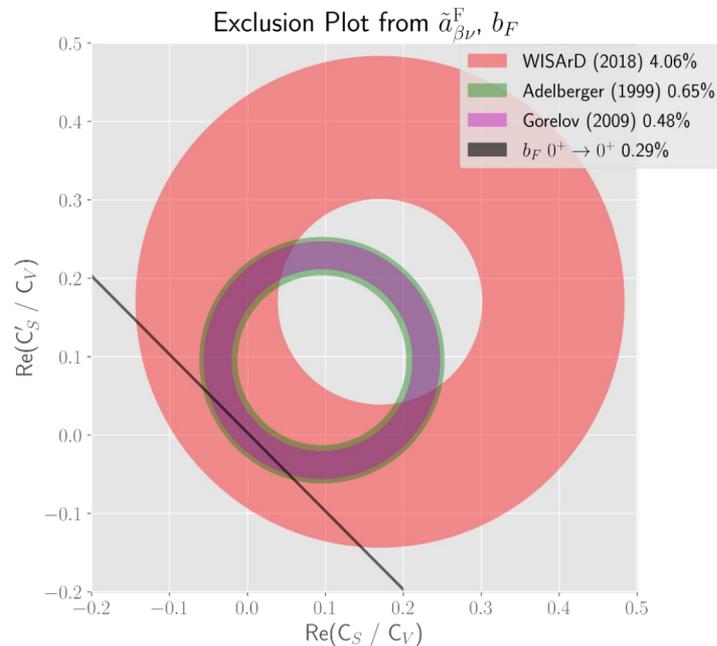
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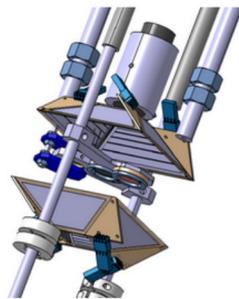
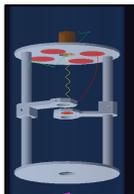
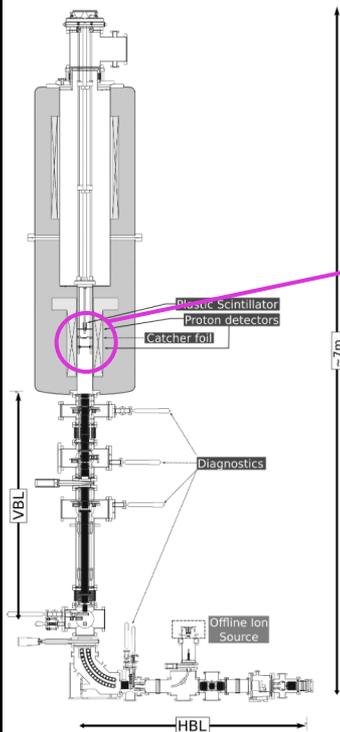
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↔ 3rd best result



Exclusion plot from D. Atanasov  
V. Araujo-Escalona et al. Phys. Rev. C 101 (2020)

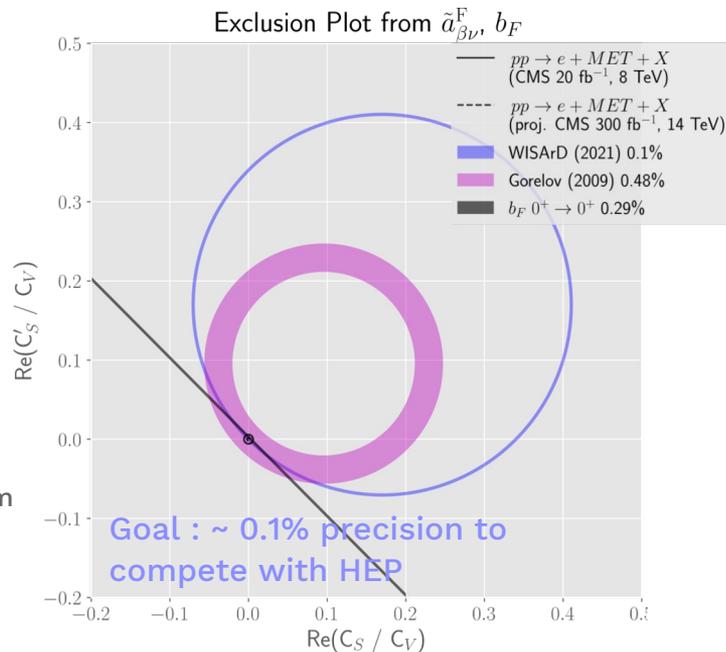
# WISArD at ISOLDE



## 2022 Upgrade

- **Ion beam transport**  
98% transmission (SIMION)
- **p detectors**  
40% solid angle + 10 keV resolution + 100 nm dead layer
- **Beta detector**  
Lower detection threshold + Validation of backscattering (GEANT4)

Next data taking : **spring 2023**



Exclusion plot from D. Atanasov

# Correlation measurements : many projects

$$\frac{dW(\mathbf{J})}{dE_e d\Omega_e d\Omega_\nu} = dW_0 \times \xi \left\{ 1 + a \frac{\mathbf{p}_e \cdot \mathbf{p}_\nu}{E_e E_\nu} + b \frac{m_e}{E_e} + \frac{\langle \mathbf{J} \rangle}{J} \cdot \left( A \frac{\mathbf{p}_e}{E_e} + B \frac{\mathbf{p}_\nu}{E_\nu} + D \frac{\mathbf{p}_e \times \mathbf{p}_\nu}{E_e E_\nu} \right) \right\}$$

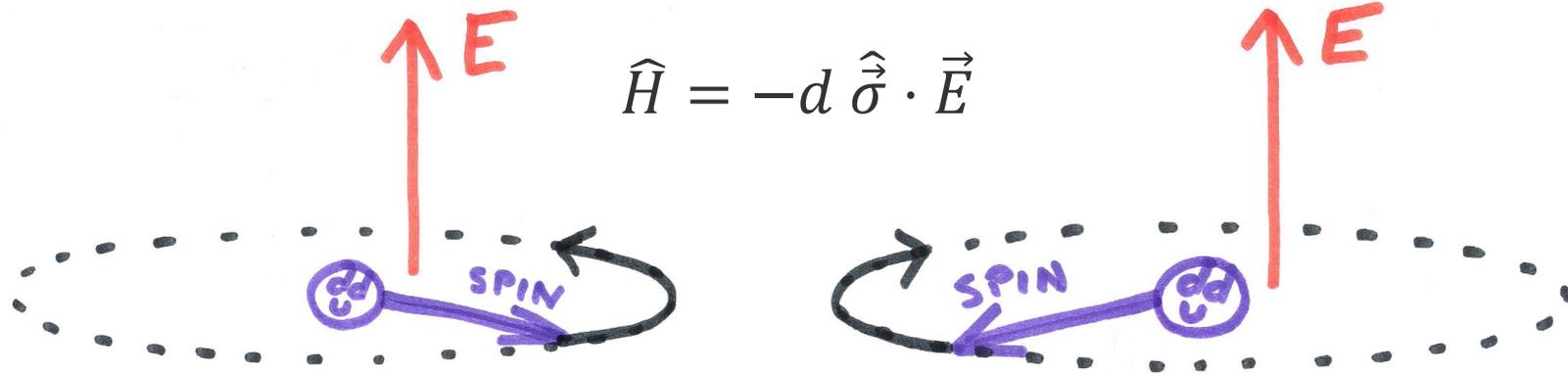
<sup>6</sup>He @ LPC (Paul trap)  
<sup>8</sup>Li @ ANL (Paul trap)  
<sup>6</sup>He @ ANL (MOT)  
<sup>32</sup>Ar @ Texas A&M (Penning)  
<sup>38m</sup>K @ TRIUMF (MOT)  
 n @ aSPECT  
 ...

<sup>114</sup>In @ ISOLDE  
<sup>6</sup>He @ LPC (bSTILED)  
<sup>6</sup>He @ NSCL  
 ...



See talk by N. Goyal

# EDMs: coupling between spin and E-field



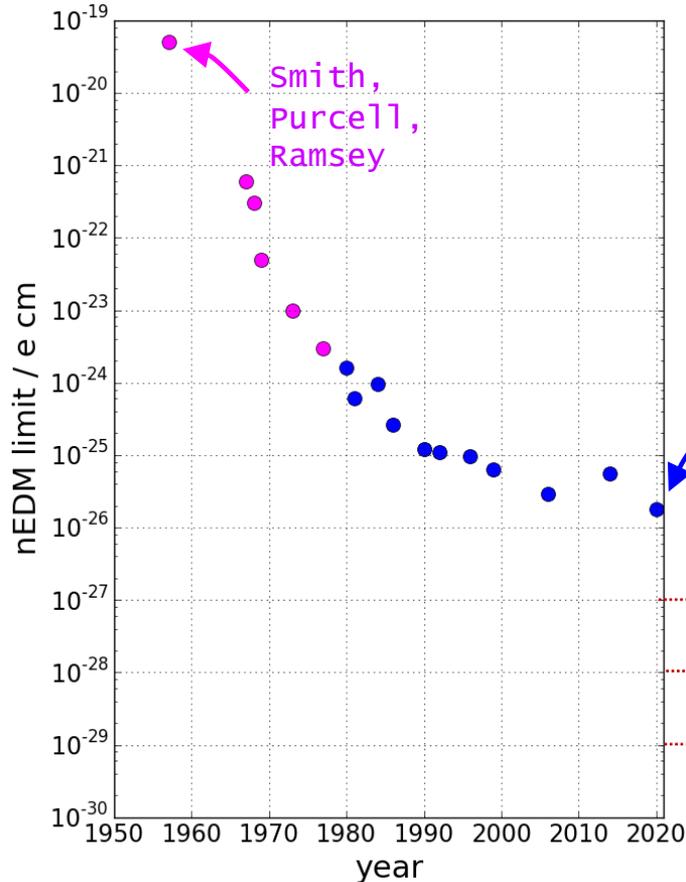
>> PLAY >>

<< REWIND <<

If  $d \neq 0$  the process and its time reversed version are different.

violation of T  $\xrightarrow{\text{CPT}}$  violation of CP

# The neutron EDM is still zero



Best limit (nEDM@PSI)

Abel et al, PRL 124, 081803 (2020)

$$|d_n| < 1.8 \times 10^{-26} \text{ e cm}$$

Design sensitivity range of 4 experiments

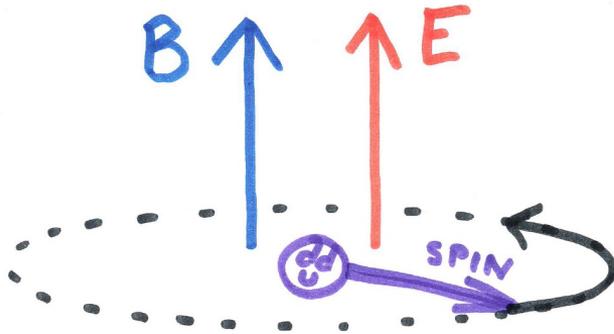
**n2EDM** @PSI, panEDM @ILL, LANL EDM, tucan @TRIUMF  
under construction now

Design sensitivity EDM@SNS, starting 2028

Possible reach with present neutron sources

CKM background uncertain, possibly  $10^{-31}$  e cm

# Basics of nEDM measurement



Larmor frequency  
 $f = 30 \text{ Hz @ } B = 1 \mu\text{T}$

$$2\pi f = \frac{2\mu}{\hbar} B \pm \frac{2d}{\hbar} |E|$$

If  $d = 10^{-26} e \text{ cm}$  and  $E = 11 \text{ kV/cm}$   
one full turn in a time

$$\frac{\pi\hbar}{dE} = 200 \text{ days}$$

To detect such a minuscule coupling

- Long interaction time
- High intensity/statistics
- Control the magnetic field

- Long interaction time
- High intensity/statistics
- Control the magnetic field

## Build colossal magnetic shields

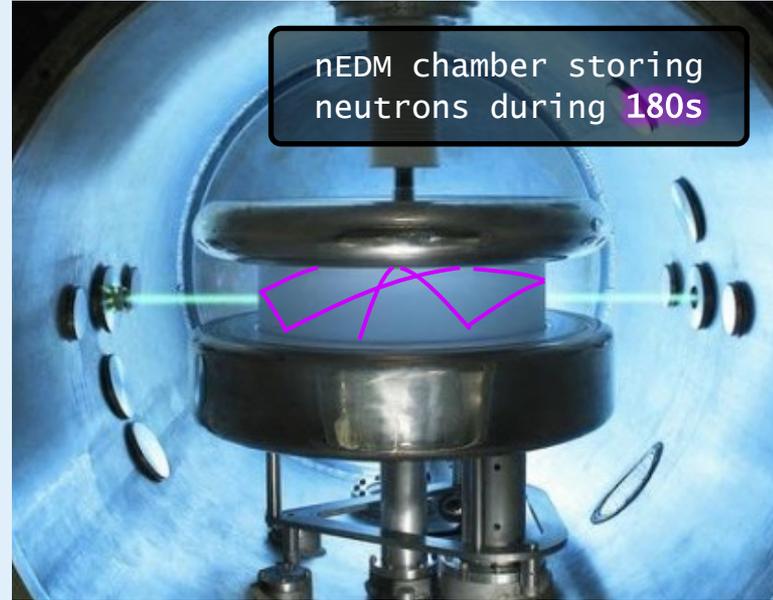


6 layers of mu-metal  
Shielding factor 100,000

+ Use quantum magnetometry

## Use Ultracold neutrons

Neutrons with velocity  $<5\text{m/s}$  ( $E < 200\text{neV}$ ) can undergo total reflection and be stored in material "bottles"



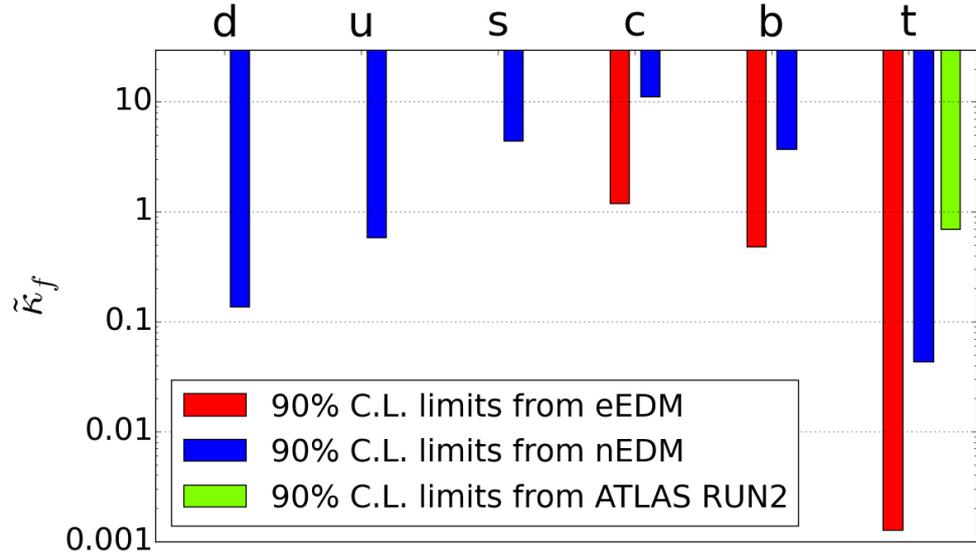
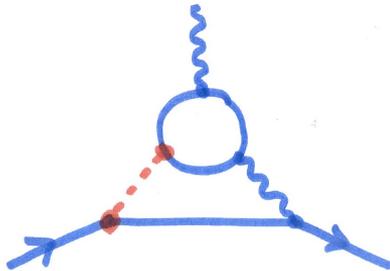
# EDMs beyond the SM: modified Higgs couplings

Modified Higgs-fermion Yukawa coupling

$$\mathcal{L} = -\frac{y_f}{\sqrt{2}} \left( \kappa_f \bar{f} f h + i \tilde{\kappa}_f \bar{f} \gamma_5 f h \right)$$

CP

Generates EDM at 2 loops  
Barr, Zee, PRL 65 (1990)



Brod, Haich, Zupan, 1310.1385  
 Brod, Stamou, 1810.12303  
 Brod, Skodras, 1811.05480  
 ATLAS, PRL 125, 061802 (2020)

# Concluding view: hunt for forbidden couplings...

Beyond SM physics =  
EFT with fundamental degrees of freedom:  
quarks & leptons, gauge, Higgs

EFT with  
nucleons,  
leptons and  
photons



observables

Lee-Yang operators:  
semileptonic, isospin changing

$$\begin{aligned}
 -\mathcal{L}_{LY} = & C_V \left( \bar{p} \gamma^\mu n + \frac{C_A}{C_V} \bar{p} \gamma^\mu \gamma_5 n \right) \times \bar{e} \gamma_\mu (1 - \gamma_5) \nu_e \\
 & + C_S \bar{p} n \times \bar{e} (1 - \gamma_5) \nu_e \\
 & + \frac{1}{2} C_T \bar{p} \sigma^{\mu\nu} n \times \bar{e} \sigma_{\mu\nu} (1 - \gamma_5) \nu_e + hc
 \end{aligned}$$

Nuclear Beta decay alphabet:  
a, b, A, B, D

$$\frac{dW(\mathbf{J})}{dE_e d\Omega_e d\Omega_\nu}$$

Isospin-diagonal, CPV operators

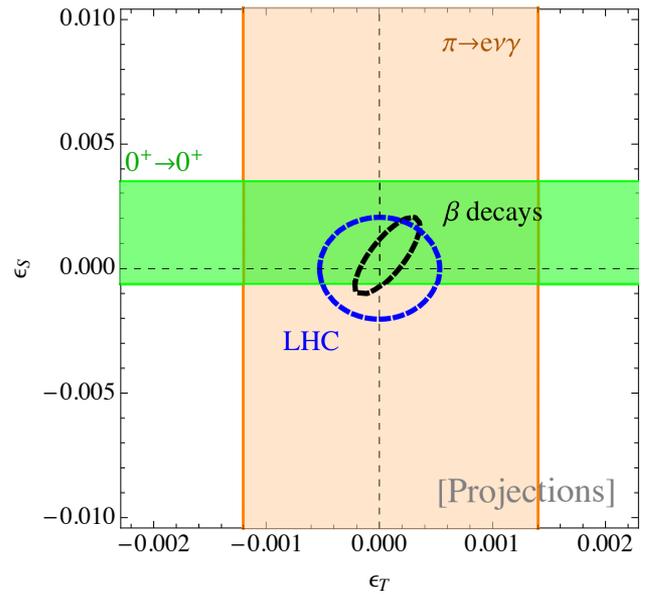
$$\begin{aligned}
 -\mathcal{L}_{EDM} \\
 = & \frac{1}{2} d_n \bar{n} \sigma_{\mu\nu} i \gamma_5 n F^{\mu\nu} \\
 & + \frac{G_F}{\sqrt{2}} C_S^0 \bar{n} n \bar{e} i \gamma_5 e + \dots
 \end{aligned}$$

EDMs of nucleons and atoms

$$\hat{H} = -d \hat{\vec{\sigma}} \cdot \vec{E}$$

**Thank you for your attention**

**Back up**

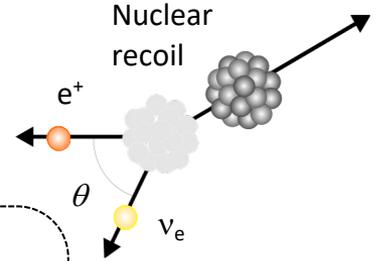


# Correlation measurements : WISARD



- Decay rate for non polarized nuclei :

$$dW = dW_0 \left( 1 + a \frac{\mathbf{p}_e \cdot \mathbf{p}_\nu}{E_e E_\nu} + b \frac{m_e}{E_e} \right)$$

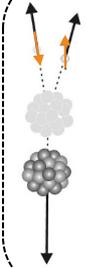


**Pure F transition  $\Delta J=0$  S=0**

$$a_F \cong 1 - \frac{|C_S|^2 + |C'_S|^2}{|C_V|^2} = 1$$

$$b_F \approx \pm \text{Re} \left( \frac{C_S + C'_S}{C_V} \right) = 0$$

Best measurement  $a_F$  at 0.45%



**Pure GT transition  $\Delta J=0$  or 1 S=1**

$$a_{GT} \cong -\frac{1}{3} \left( 1 - \frac{|C_T|^2 + |C'_T|^2}{|C_A|^2} \right) = -1/3$$

$$b_{GT} \approx \pm \text{Re} \left( \frac{C_T + C'_T}{C_A} \right) = 0$$

Best measurement  $a_{GT}$  at  $< \sim 1\%$

