



Dark Matter: status and prospects at FCC

Dario Buttazzo



Istituto Nazionale di Fisica Nucleare

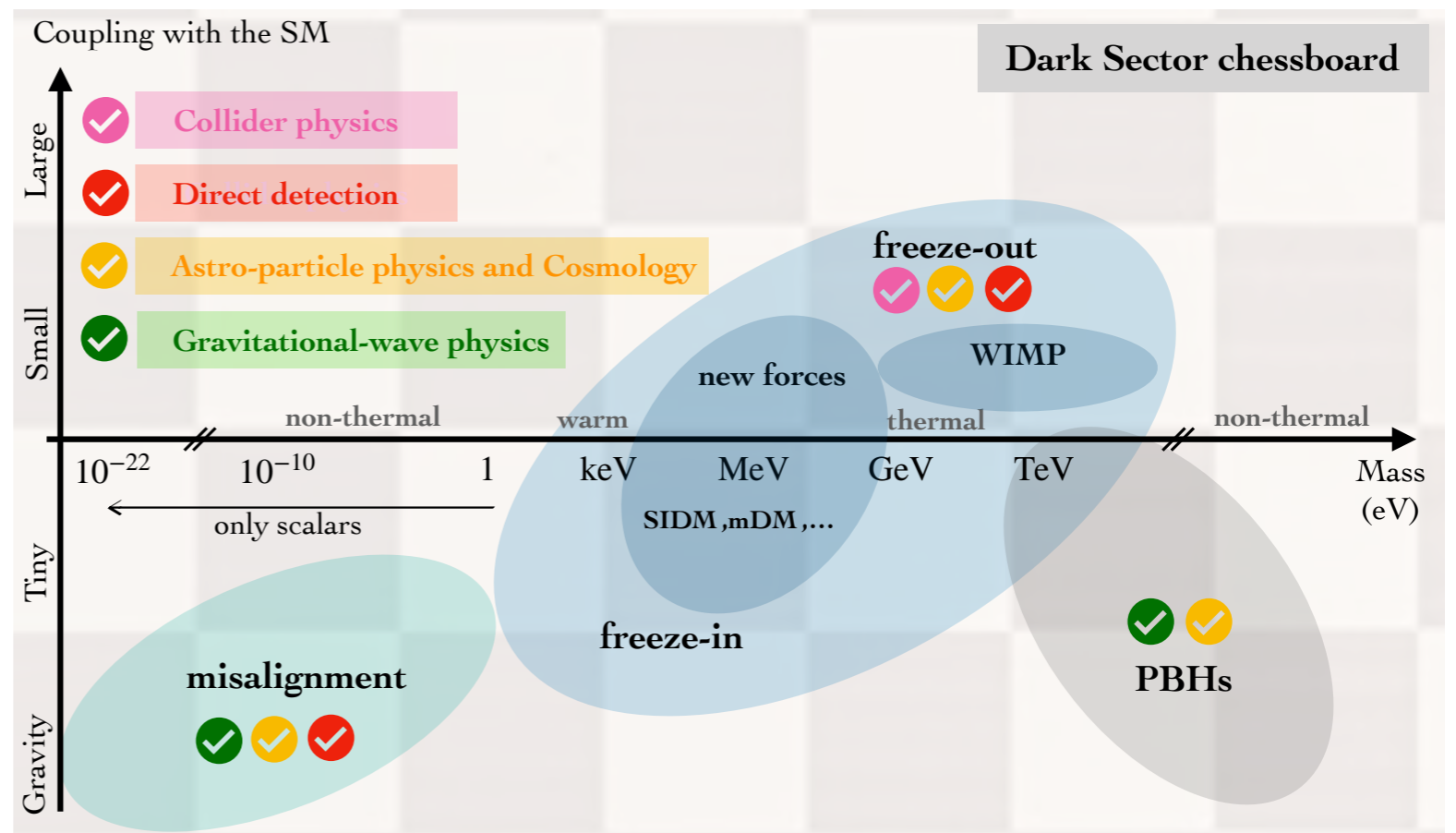
Sezione di Pisa

Motivation

- ◆ Dark Matter *exists*: the only direct evidence of physics beyond the Standard Model so far!
- ◆ One of the fundamental problems that the next large collider should be able to address

- ◆ Its fundamental nature is completely unknown: can span tens of orders of magnitude in mass and coupling

- ◆ The only measured quantity is the cosmological abundance $\Omega_{\text{DM}} \sim 0.26$



Motivation

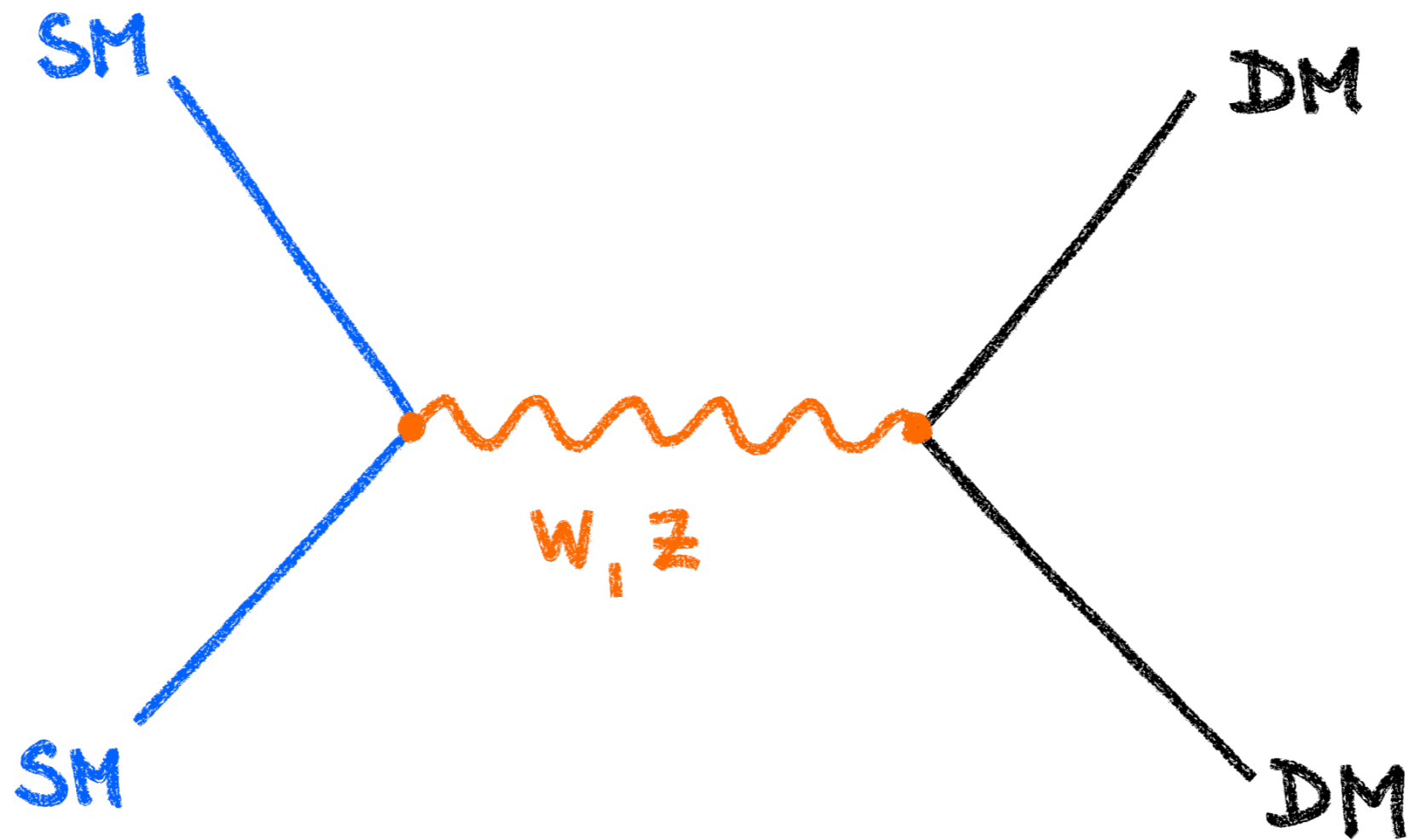
LOOKING UNDER THE LAMPPOST

LOST YOUR KEYS?

YEAH,
I LOST THEM OVER
THERE BUT THE
LIGHT'S BETTER HERE

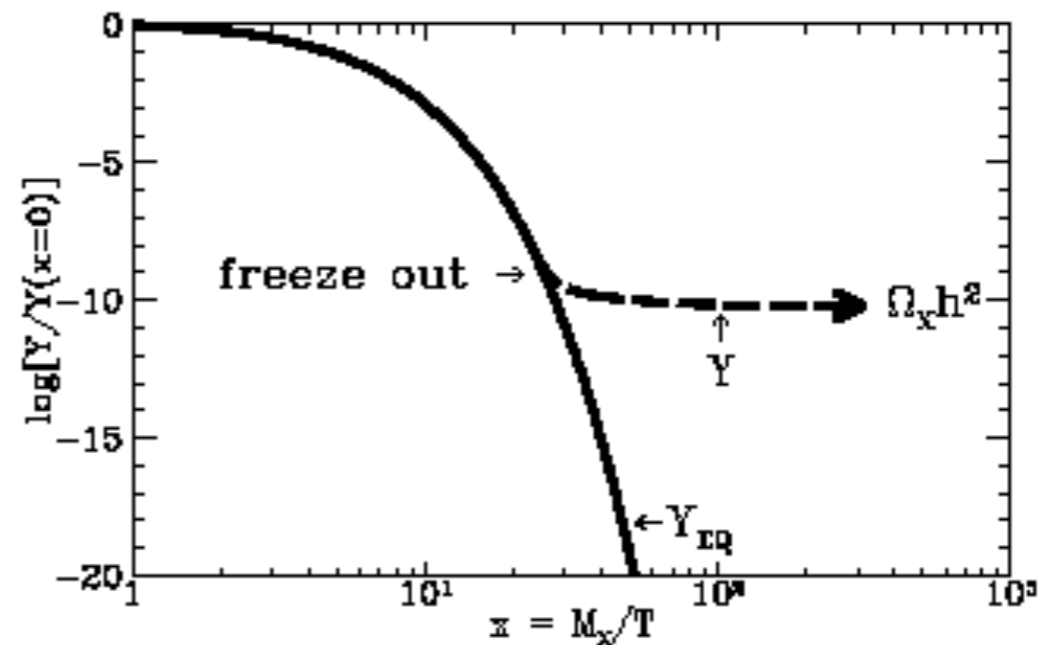
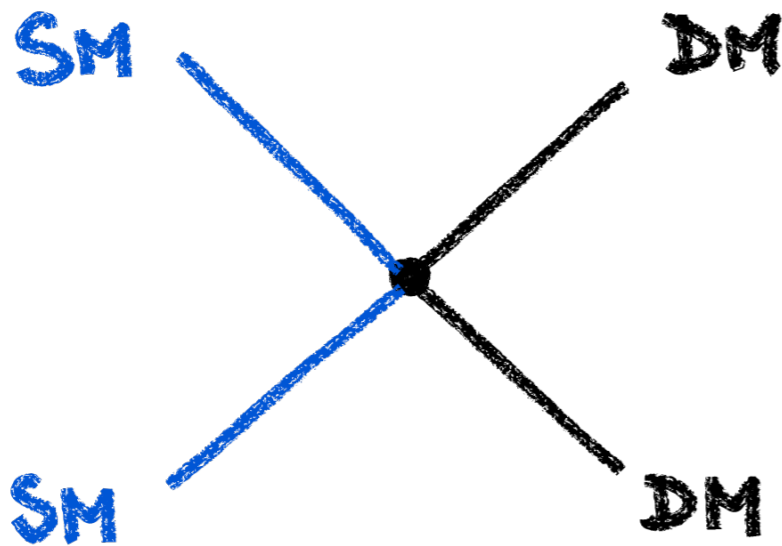


WIMP



The case for WIMPs

- ♦ **WIMP miracle:** offers at the same time a simple explanation for the observed Dark Matter abundance ($\Omega_{\text{DM}} \sim 0.26$) and a connection to naturalness of electroweak scale.
- ♦ Production in early Universe: thermal freeze-out of $2 \rightarrow 2$ scatterings



- ♦ For each value of the DM-SM coupling g_* DM mass is predicted.

☞ $g_* \sim g_{\text{EW}} \Rightarrow M_{\text{DM}} \sim \text{TeV}$

Ideal target for colliders!

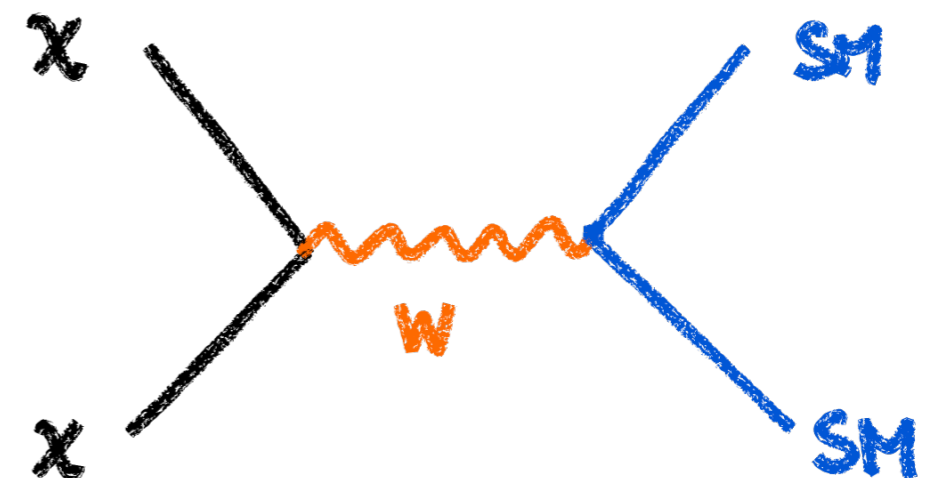
Which WIMP?

- ◆ Consider generic EW multiplet: interacts w/ SM through W, Z

DM is the neutral component $\chi_n = (\dots, \chi^-, \chi^0, \chi^+, \dots)$

“Minimal Dark Matter”: Cirelli, Fornengo, Strumia 2005

- ◆ DM needs to be stable
- ◆ Strong bounds from Direct Detection: No Z coupling @ tree-level
 - ▶ Real multiplet: $Y = 0, n$ odd
 - ▶ Complex multiplet: $Y \neq 0,$
(mass splittings from higher-dimensional operators needed)
- ◆ Single parameter sets the DM abundance:
mass M_{DM}



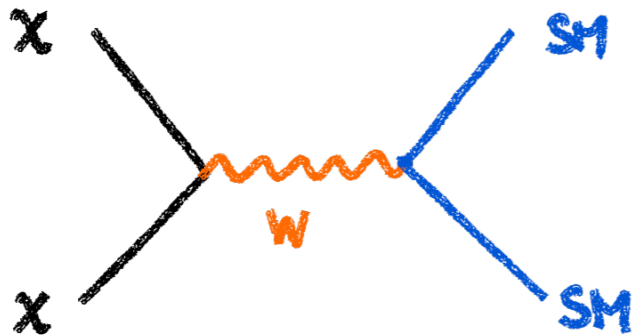
Which WIMP?

- ◆ Consider generic EW multiplet: interacts w/ SM through W, Z

$$\frac{dY}{dx} \propto \langle \sigma v \rangle (Y^2 - Y_{\text{eq}}^2)$$

which cross-section?

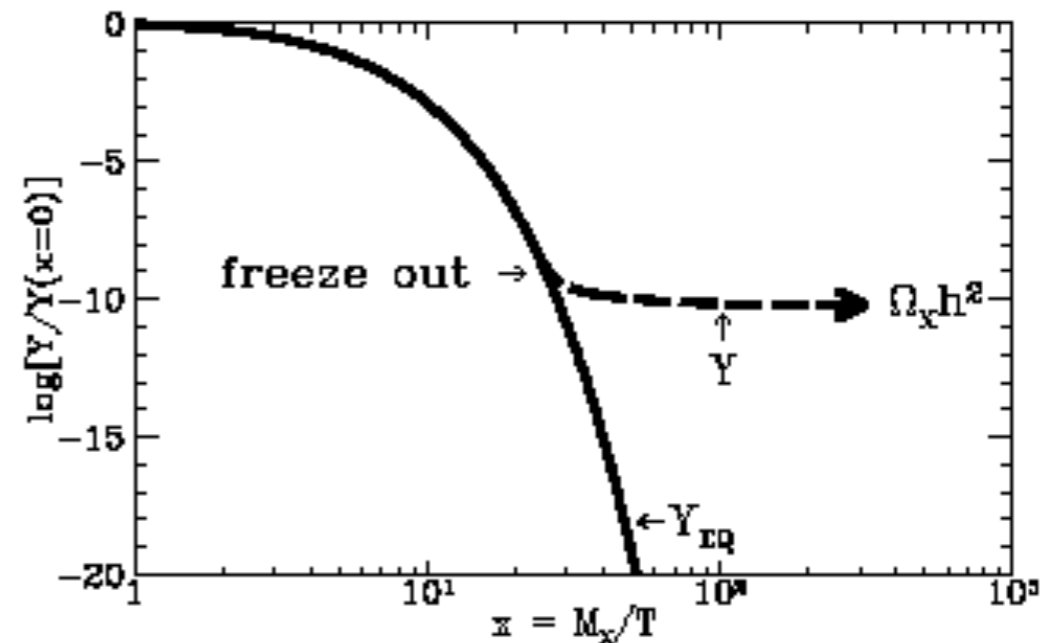
- ◆ Tree-level EW cross-section...



$$\langle \sigma v \rangle_0 = \frac{\pi \alpha_2^2 (2n^4 + 17n^2 - 19)}{16g_\chi M_\chi^2}$$

... is inaccurate!

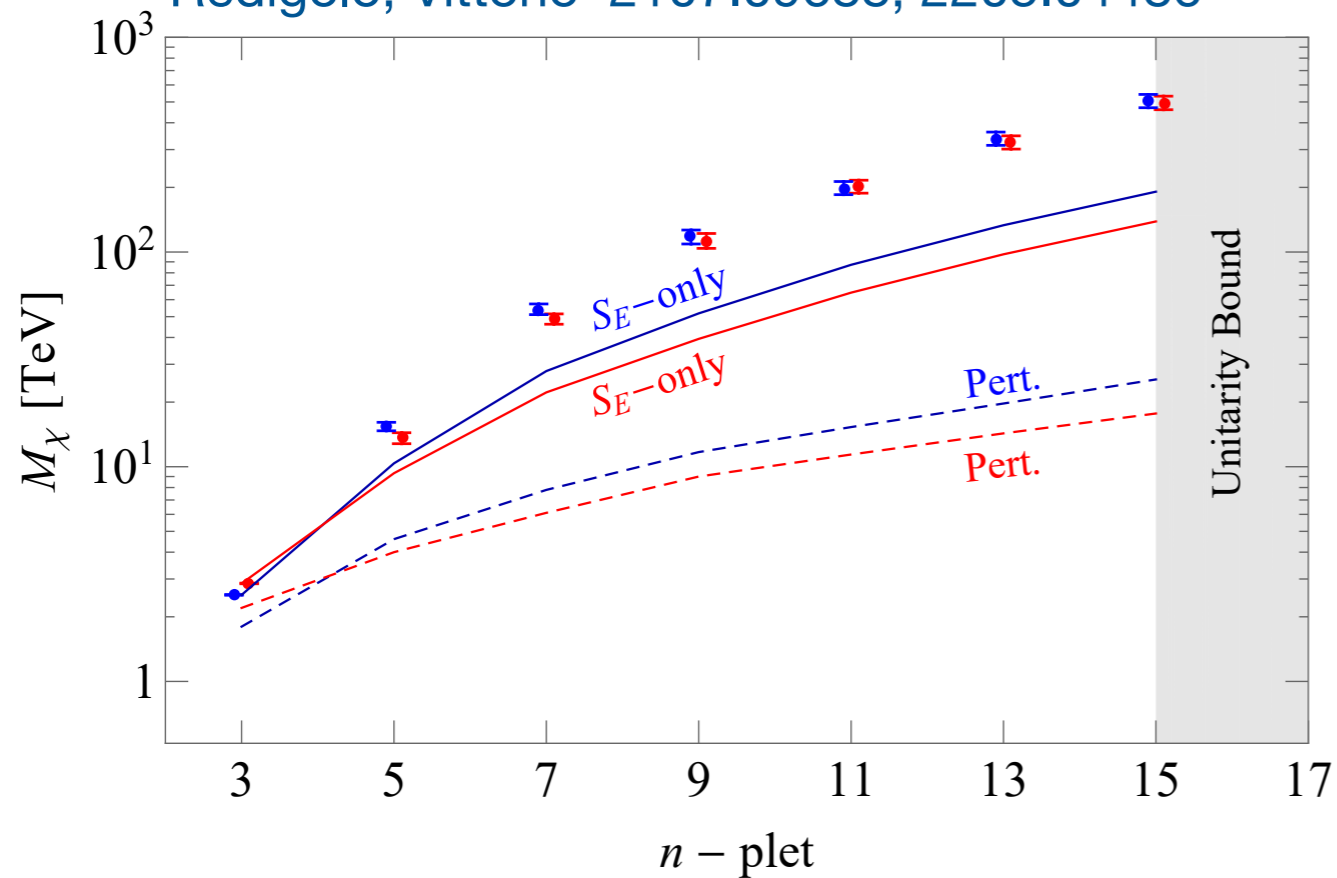
- ▶ Sommerfeld enhancement
- ▶ Bound states formation



Large non-perturbative,
non-relativistic effects

Thermal freeze-out masses

Bottaro, DB, Costa, Franceschini, Panci,
Redigolo, Vittorio 2107.09688, 2205.04486

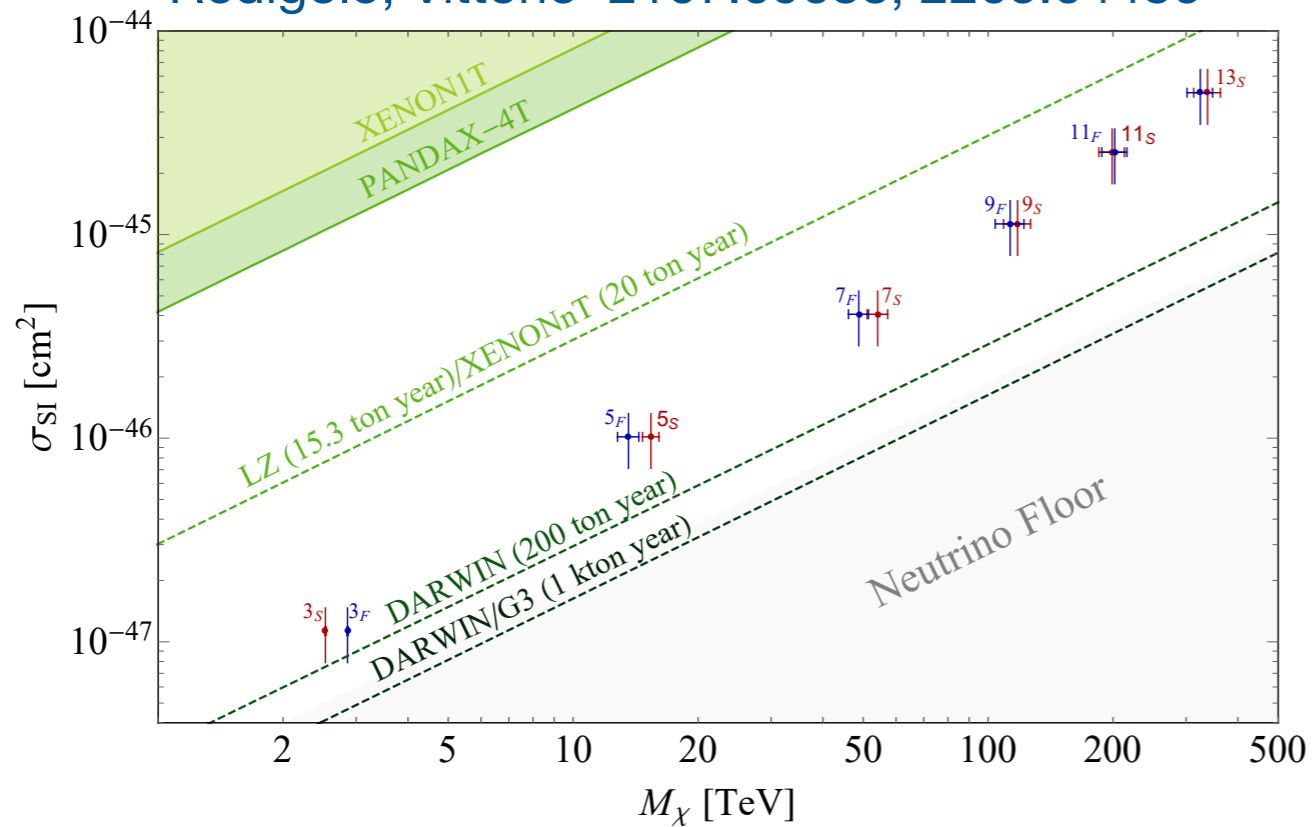


(and similar for scalars)

| | EW n -plet | Mass [TeV] |
|------------------|--------------|------------|
| Majorana fermion | 3_0 | 2.86 |
| | 5_0 | 13.6 |
| | 7_0 | 48.8 |
| | 9_0 | 113 |
| | 11_0 | 202 |
| | 13_0 | 324.6 |
| Dirac fermion | $2_{1/2}$ | 1.08 |
| | 3_1 | 2.85 |
| | $4_{1/2}$ | 4.8 |
| | 5_1 | 9.9 |
| | $6_{1/2}$ | 31.8 |
| | $8_{1/2}$ | 82 |
| | $10_{1/2}$ | 158 |
| | $12_{1/2}$ | 253 |

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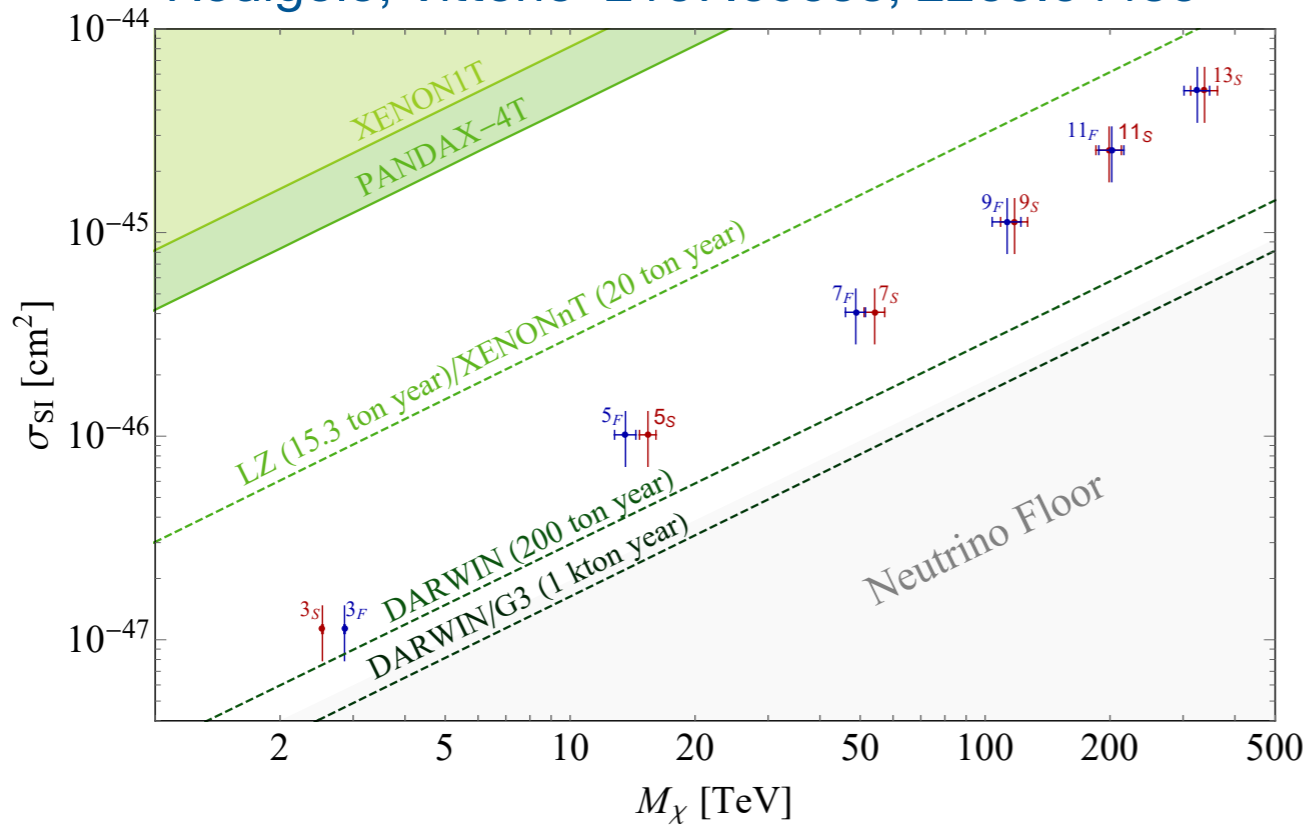


Direct Detection challenging...

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Thermal freeze-out masses

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Direct Detection challenging...

Potentially in reach of
a high-energy collider!

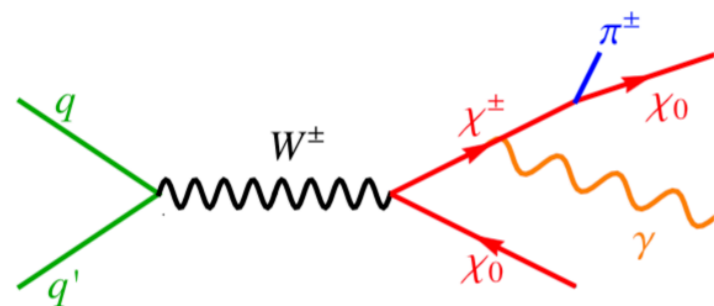
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Electroweak triplet: missing energy searches

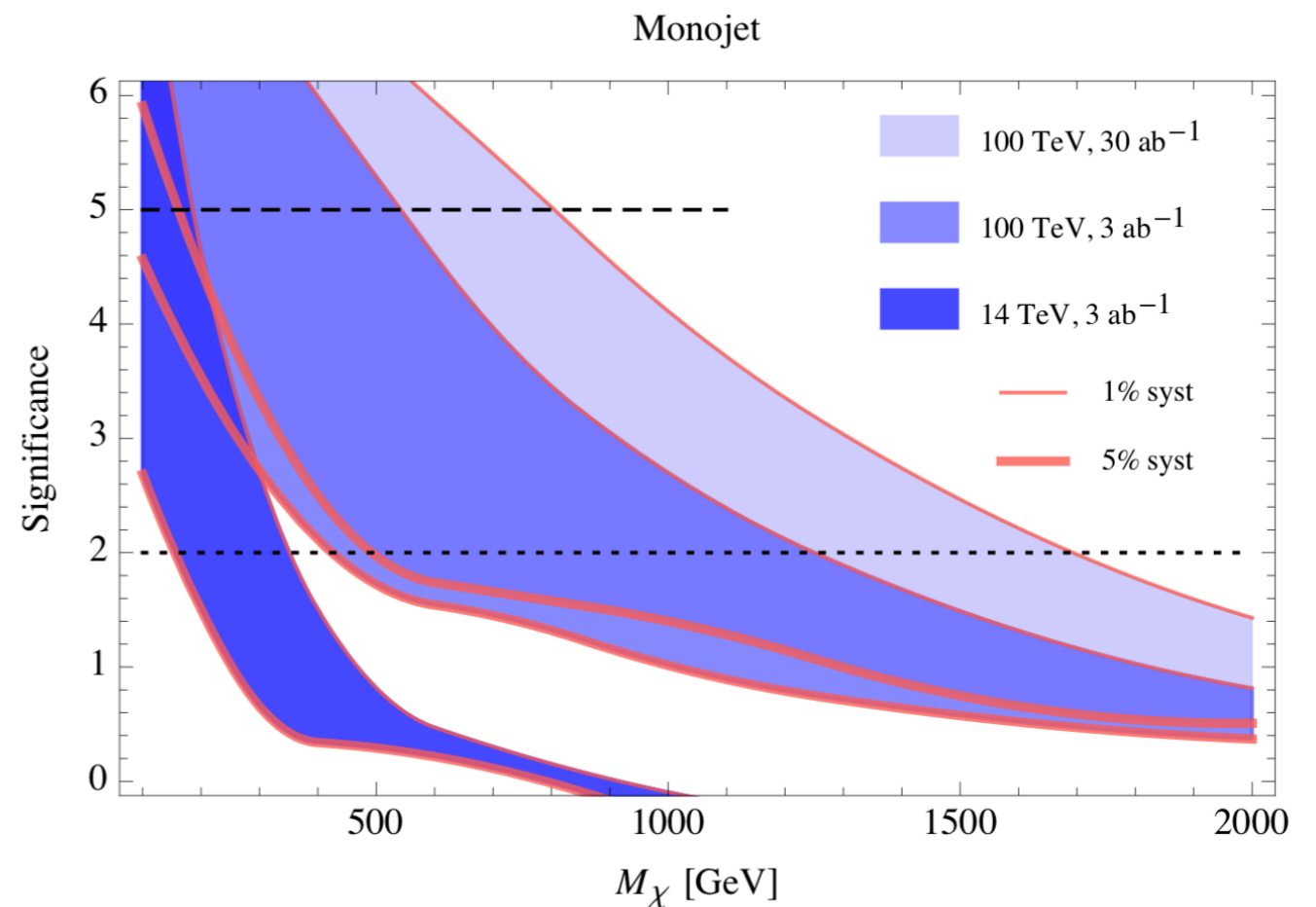
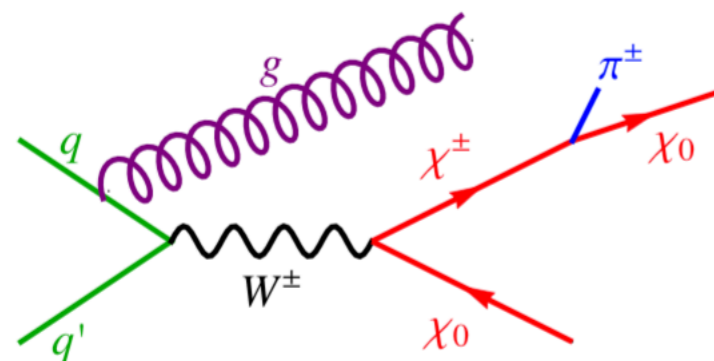
$$\mathcal{L} = \frac{1}{2} \bar{\chi} \left(iD_{\mu} \gamma^{\mu} - M_{\chi} \right) \chi, \quad \chi = (\chi^{-}, \chi^0, \chi^{+}) \sim \mathbf{3}$$

- ◆ $2 \rightarrow 2$ production of invisible χ pair + event tag

- ▶ Mono-photon:



- ▶ Mono-jet



Cirelli, Sala, Taoso 1407.7058

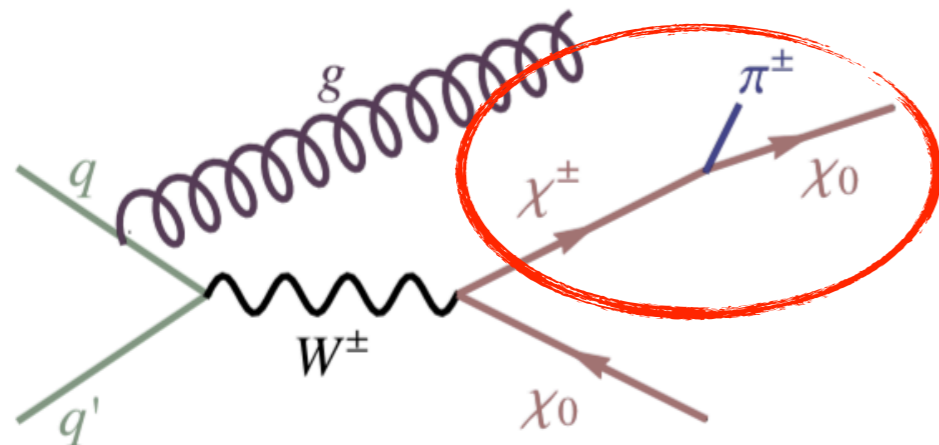
- ◆ Main background from $Z \rightarrow$ neutrinos
- ◆ Requires good control of systematics

Thermal freeze-out
mass: 2.86 TeV ✗

Electroweak triplet: disappearing tracks

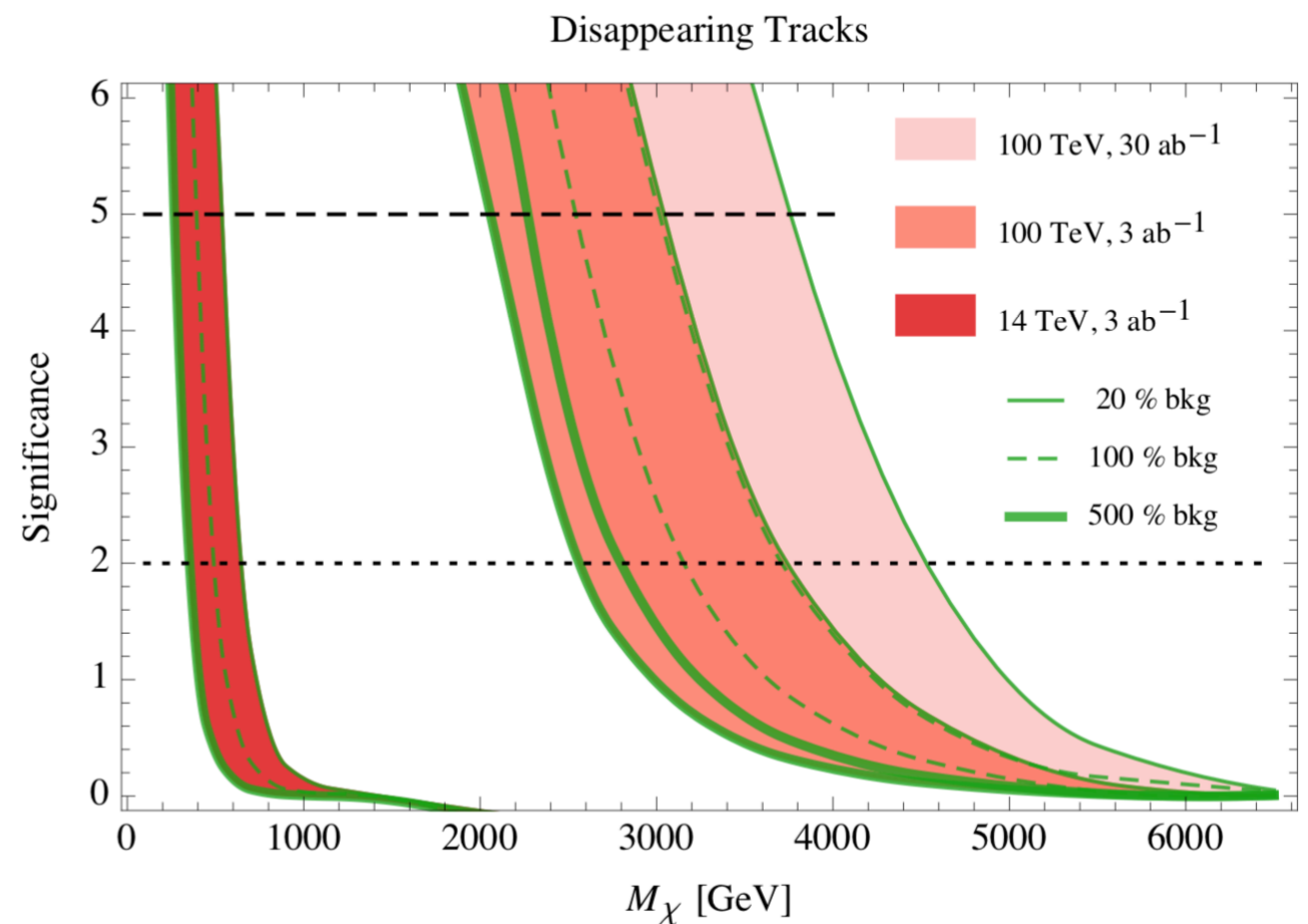
$$\mathcal{L} = \frac{1}{2} \bar{\chi} \left(iD_{\mu} \gamma^{\mu} - M_{\chi} \right) \chi, \quad \chi = (\chi^{-}, \chi^0, \chi^{+}) \sim \mathbf{3}$$

- ◆ 2 → 2 production of invisible χ pair + event tag
- ◆ Can look for the tracks of charged χ^{\pm} decaying to χ^0 DM:



- ▶ Real WIMPS: lifetime is fixed by EW interactions

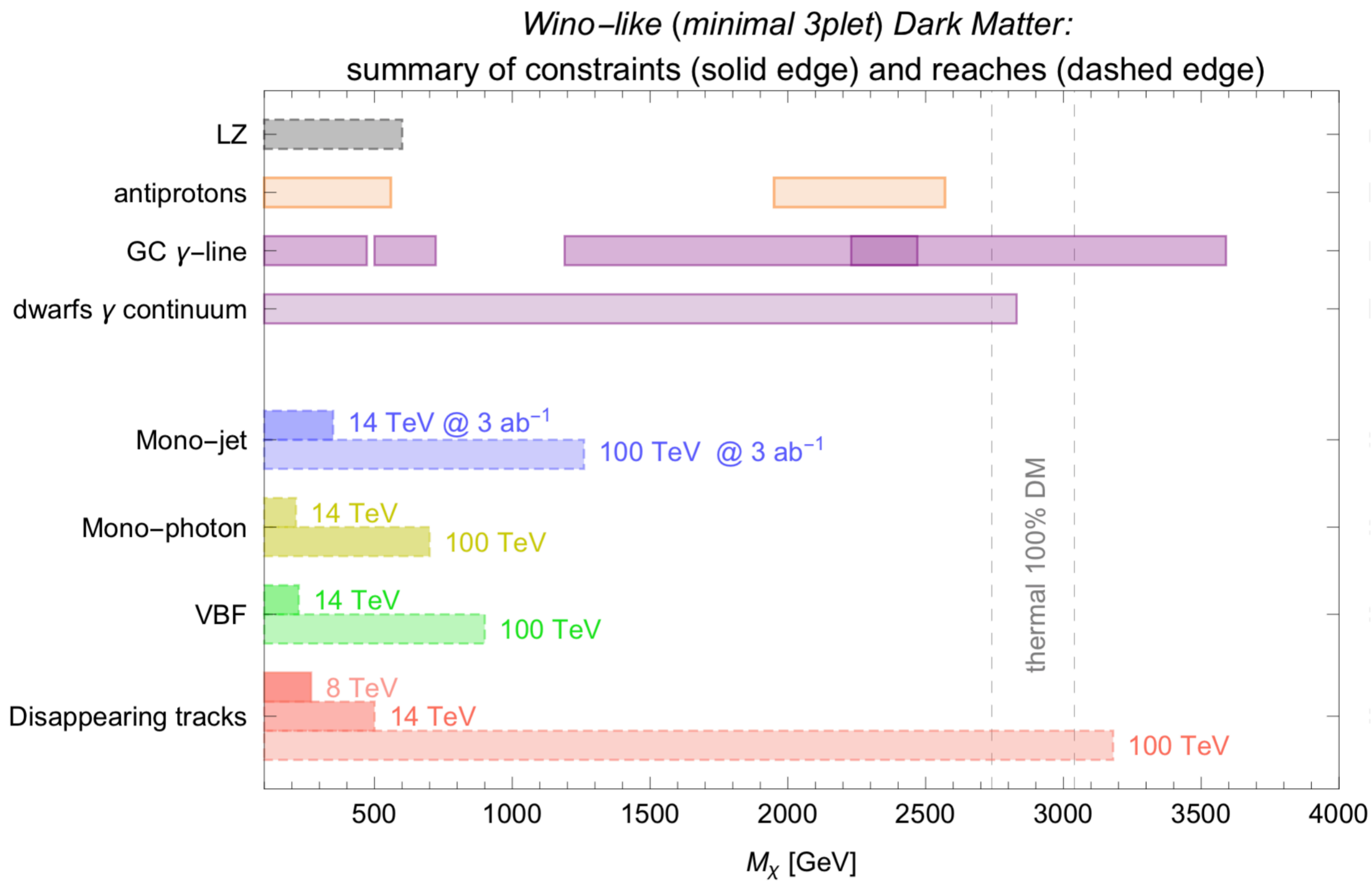
$$c\tau \approx 50 \text{ cm}/(n^2 - 1)$$



Thermal freeze-out mass: 2.86 TeV ✓

Cirelli, Sala, Taoso 1407.7058

Electroweak triplet @ FCC-hh

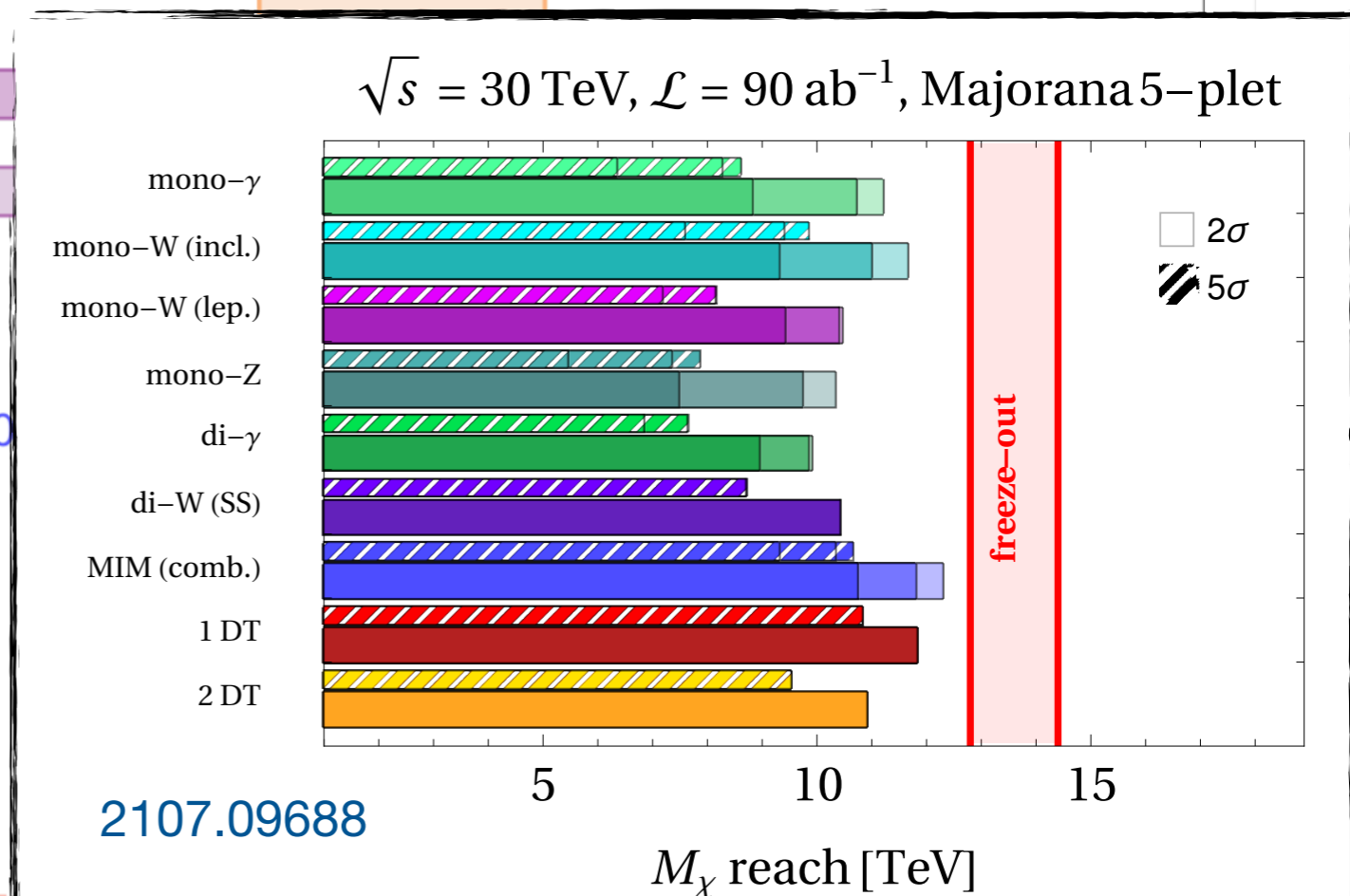
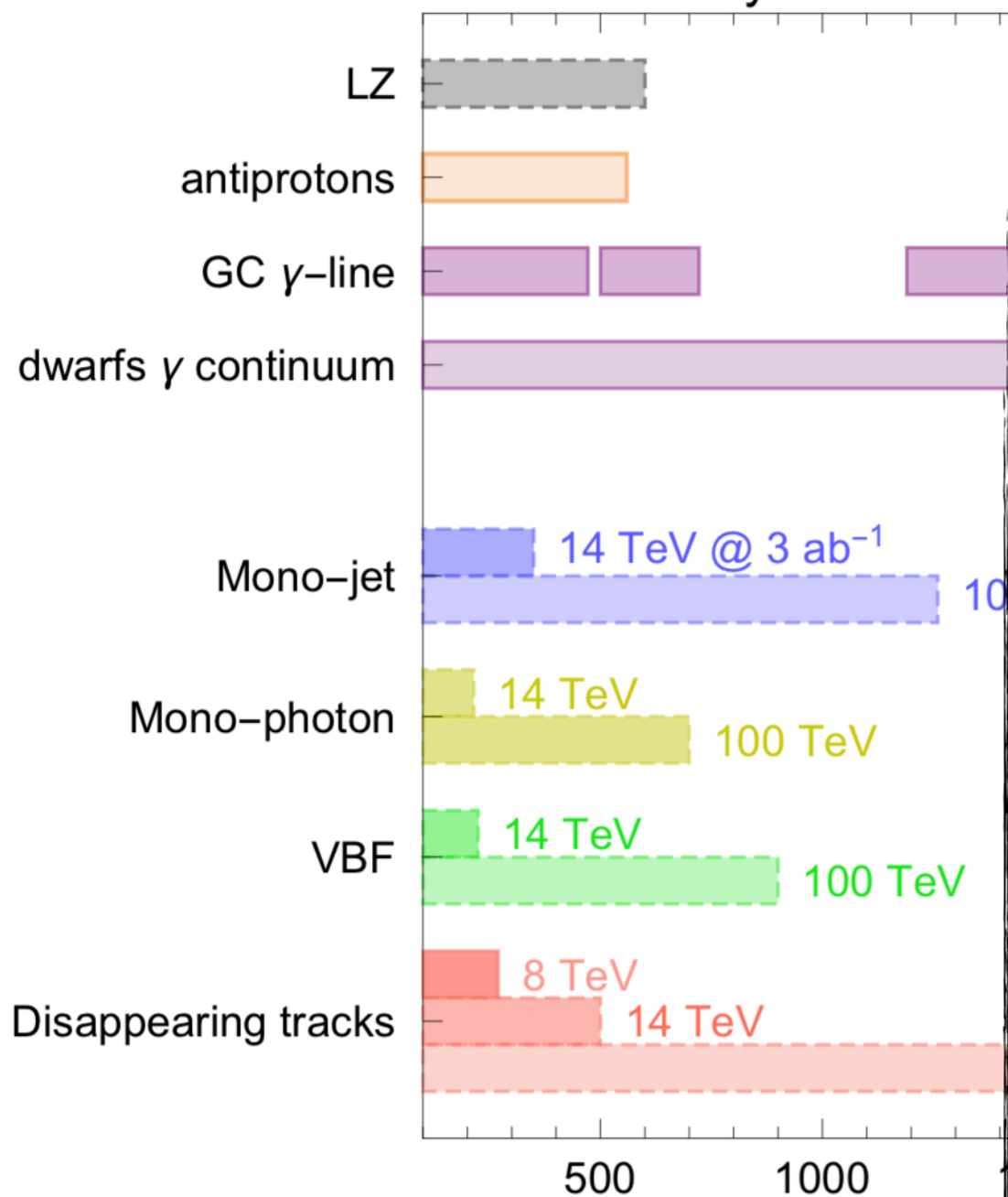


Cirelli, Sala, Taoso 1407.7058

FCC Physics Opportunities

Electroweak triplet @ FCC-hh

Wino-like (minimal 3plet) Dark Matter:
summary of constraints (solid edge) and reaches (dashed edge)



- ◆ Many other channels possible (EW radiation important at high energy)
- ◆ Can be combined to improve limits and reduce impact of systematics

Other EW multiplets

- ◆ Doublet (Higgsino) and triplet (Wino) fully probed at FCC-hh

- ▶ Results strongly depend on the detector layout!

Terashi, Sawada, Saito, Asai (2018)

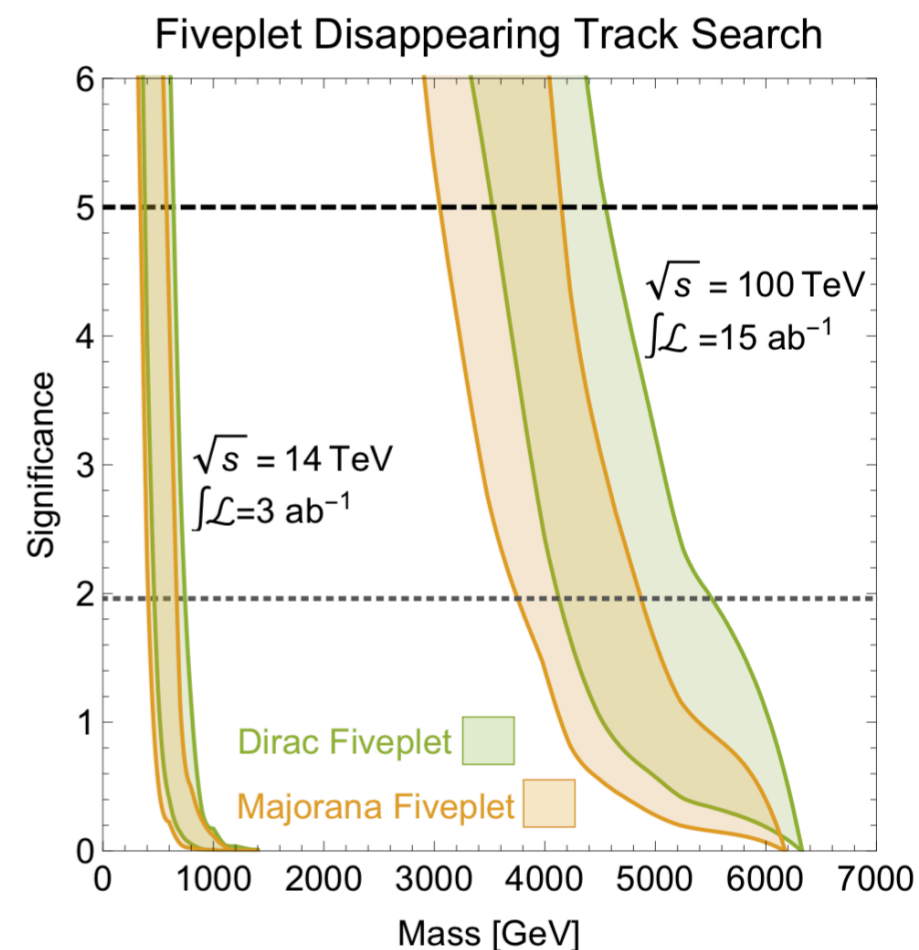
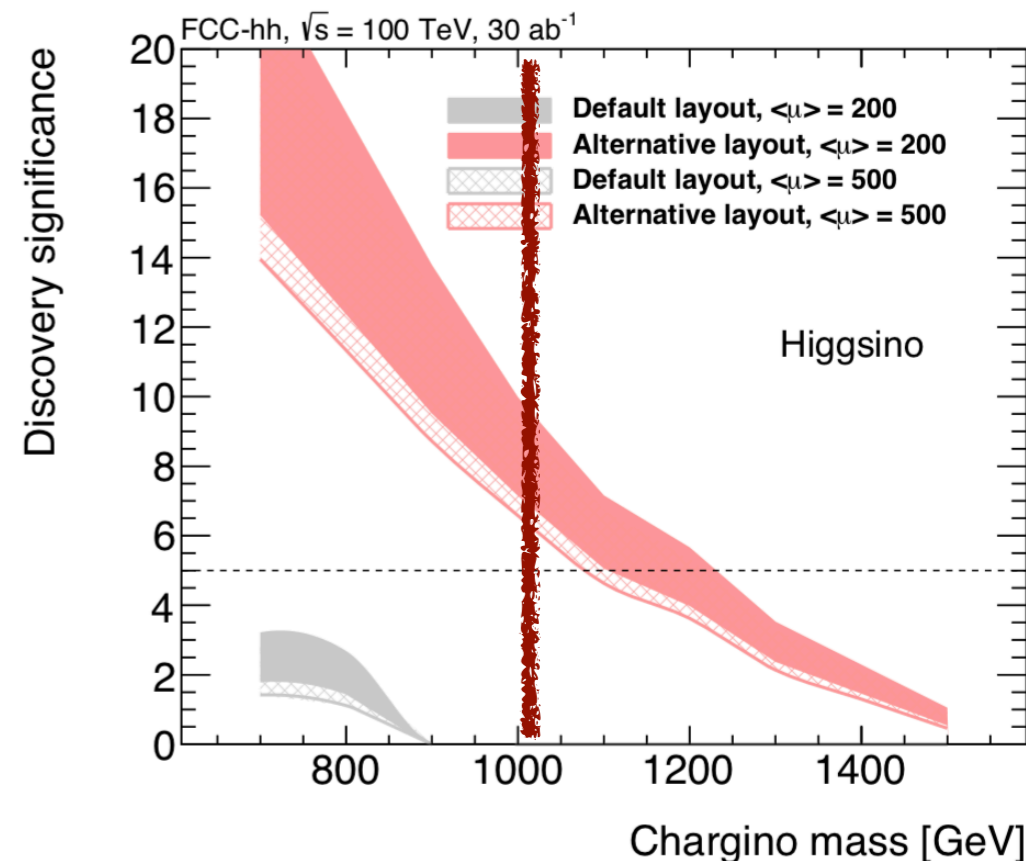
- ◆ Larger EW multiplets:

😊 Larger cross-section due to big EW charges

😞 Much larger thermal masses, $M_{DM} \sim n^{5/2}$

😞 Shorter lifetime of charged components, disappearing tracks less effective

$n \geq 5$ not testable at FCC



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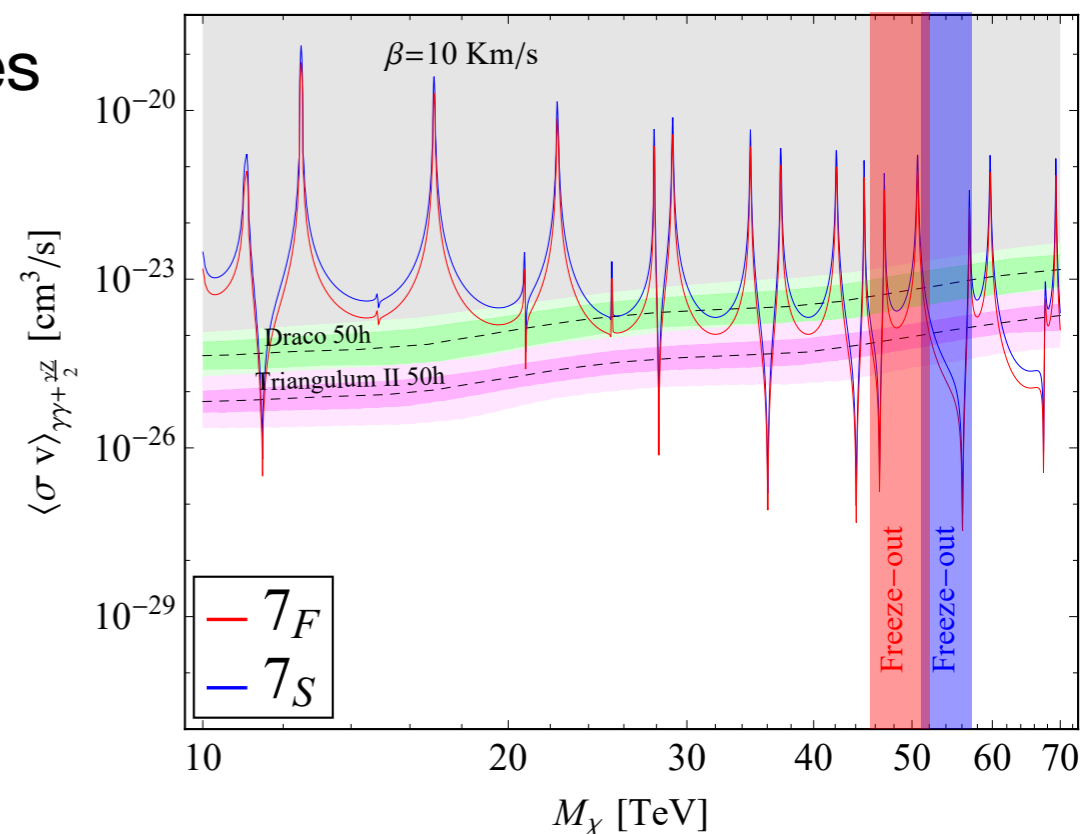
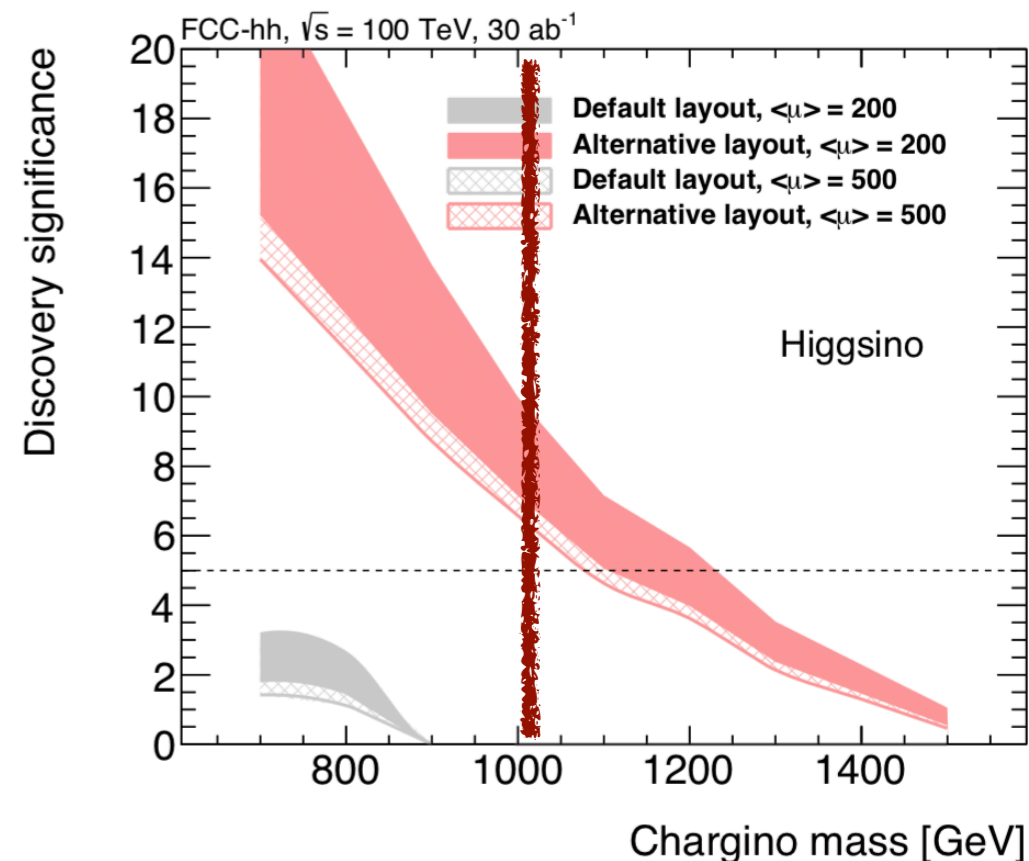
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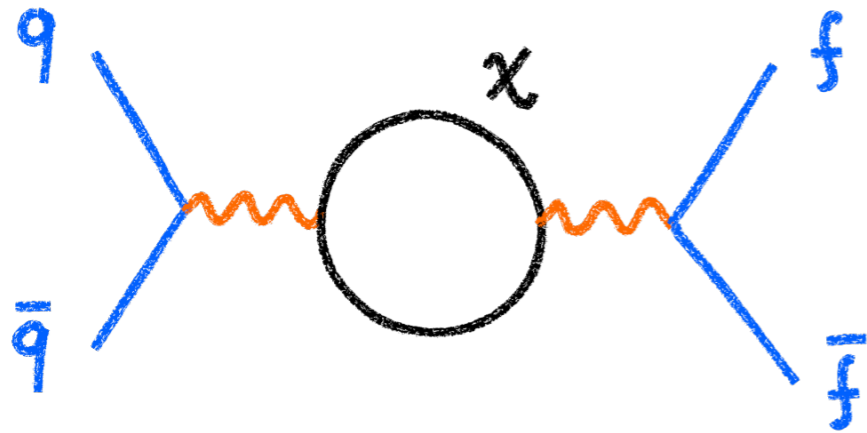
$n \geq 5$ not testable at FCC

Signals in Indirect Detection!



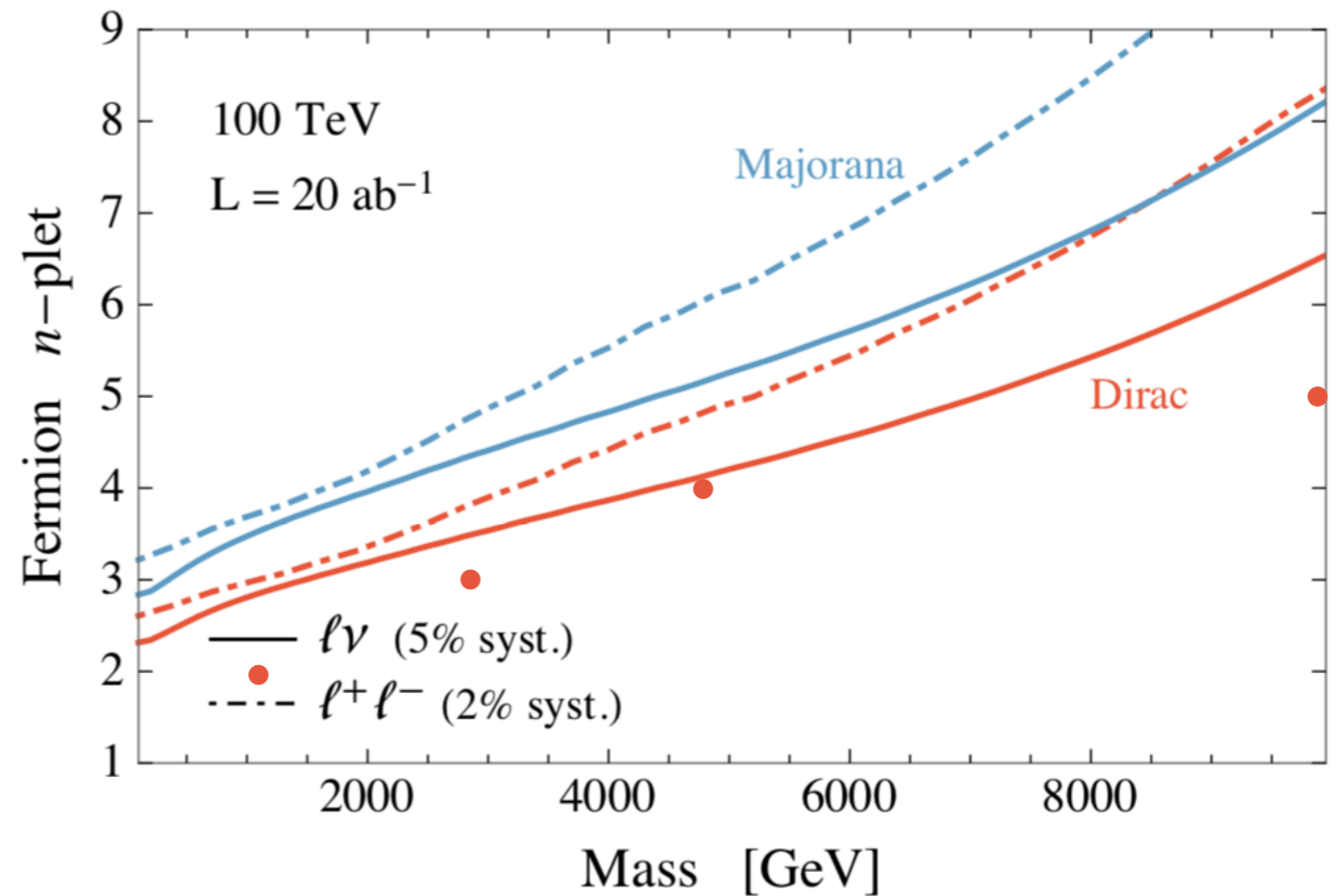
Indirect effects at colliders

- ♦ All EW multiplets give universal contributions to high-energy $2 \rightarrow 2$ fermion scattering (W, Y parameters)



These effects grow with energy, can be tested at FCC-hh

Di Luzio, Gröber, Panico 1810.10993



$$\hat{W} \approx 10^{-7} \times \left(\frac{1 \text{ TeV}}{M_{\text{DM}}} \right)^2 n^3 \propto 1/n^2$$

$$\hat{Y} \approx 10^{-7} \times \left(\frac{1 \text{ TeV}}{M_{\text{DM}}} \right)^2 Y^2 n \propto 1/n^4$$

Indirect effects at colliders

- ◆ Complex multiplets need mass splittings from higher dim. operators

- ▶ To make neutral component stable: $(\bar{\chi} T^a \chi) (H^\dagger \sigma^a H)$

- ▶ To suppress Z-induced scattering (DD): $(\bar{\chi} (T^a)^{2Y} \chi^c) (H^{\dagger c} \sigma^a H)^{2Y}$

- ◆ Contribution to S, T parameters

$$\hat{S} \approx 10^{-5} \times \left(\frac{1 \text{ TeV}}{M_{\text{DM}}} \right) \left(\frac{\delta M}{10 \text{ GeV}} \right) n^3$$

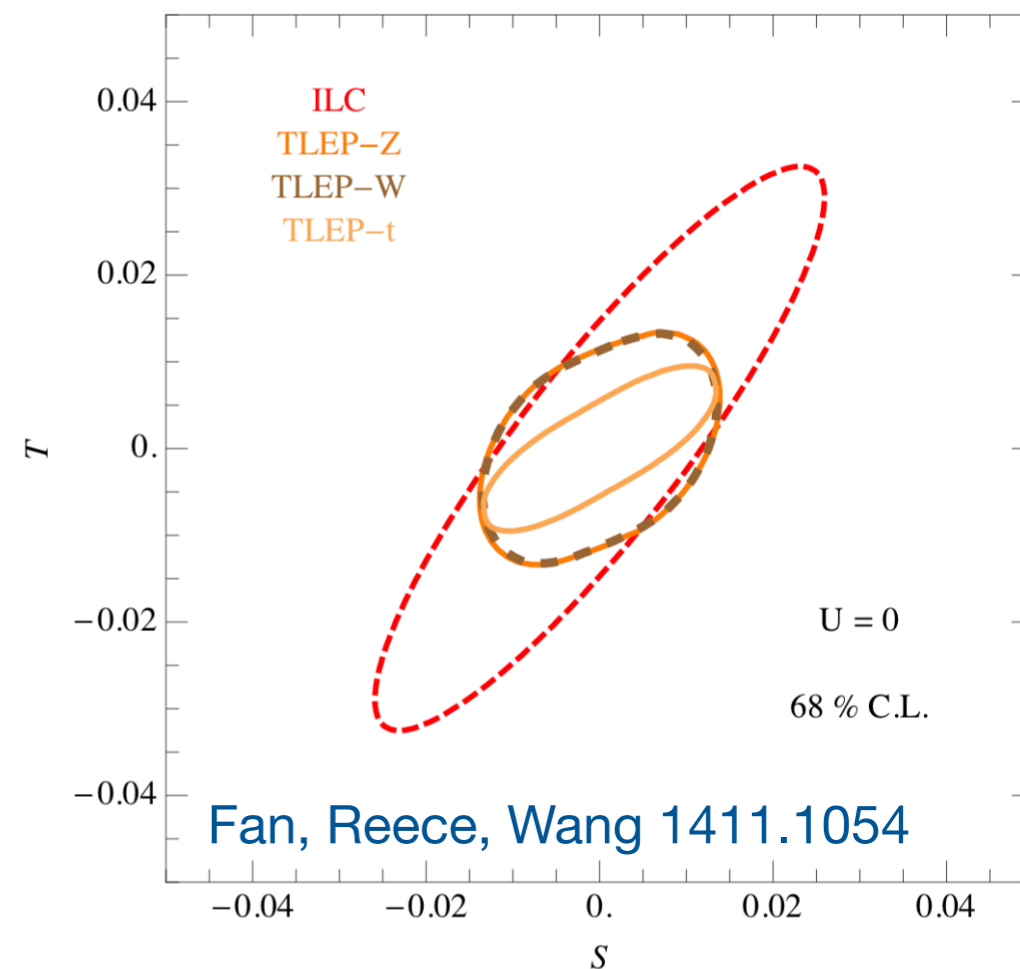
$$\hat{T} \approx 10^{-5} \times \left(\frac{\delta M}{10 \text{ GeV}} \right)^2 n^3$$

Bottaro, DB, Costa, Franceschini,
Panci, Redigolo, Vittorio 2205.04486

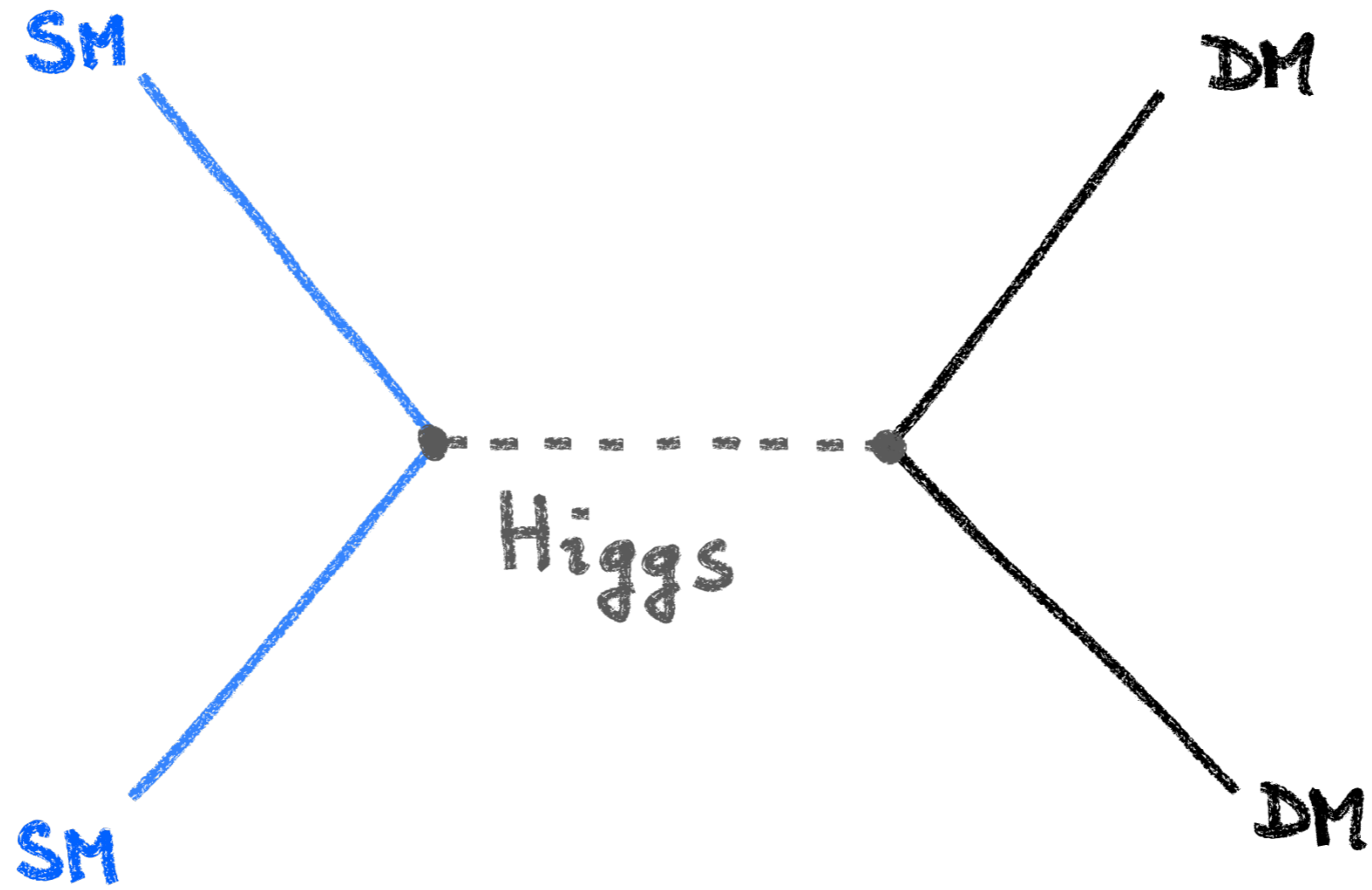
Di Luzio, Gröber, Kamenik, Nardecchia 1505.00359

- ◆ FCC-ee sensitivity: $\hat{S}, \hat{T} \lesssim \text{few} \times 10^{-5}$

could probe larger EW multiplets



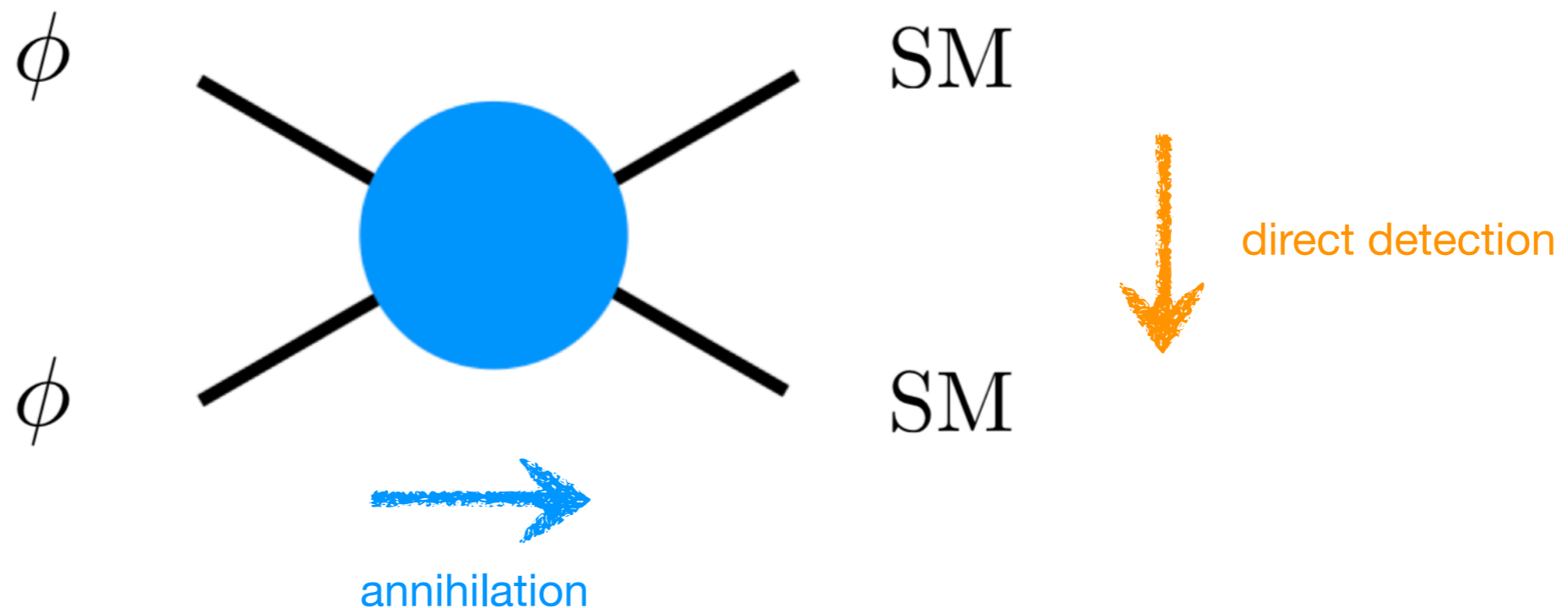
Higgs portal



Higgs portal

- ◆ Singlet scalar coupled through Higgs portal:
one of the simplest thermal DM scenarios beyond EW interactions

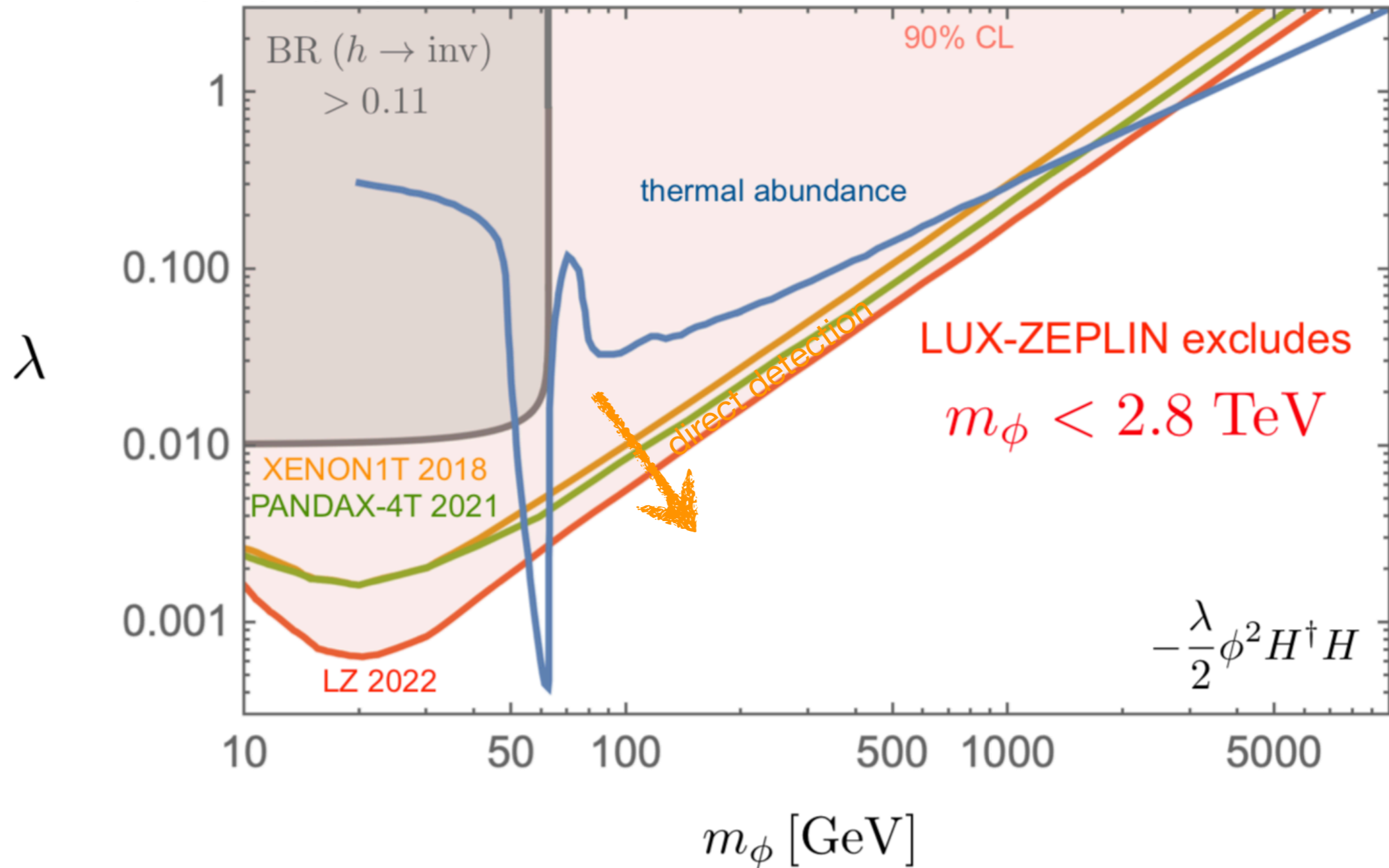
$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{1}{2} M_\phi^2 \phi^2 - \frac{\lambda}{2} \phi^2 H^\dagger H$$



Freeze-out: $\Omega_\phi h^2 \propto \frac{1}{\langle \sigma v \rangle} \approx 0.1 \left(\frac{0.1}{\lambda} \right)^2 \left(\frac{M_\phi}{100 \text{ GeV}} \right)^2$

Higgs portal

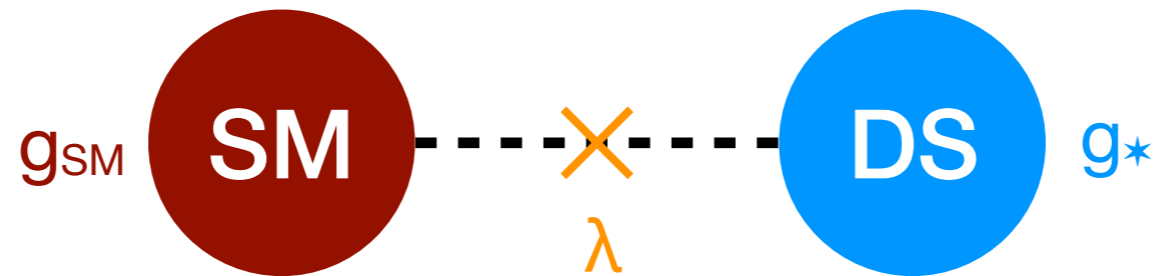
- ◆ Masses below few TeV strongly constrained by Direct Detection



Good example of the “WIMP vs direct detection” tension

Higgs portal

- ◆ Higgs portal to a *Dark Sector*: decouple freeze-out and direct detection



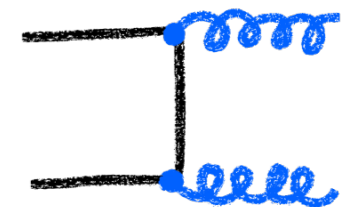
- ▶ freeze-out is controlled by a coupling g_* in the dark sector
- ▶ portal coupling $\lambda |H|^2 |\phi|^2$ can be made small

☞ Unitarity: if $g \sim 4\pi$, $m_{\text{DM}} \sim 100 \text{ TeV}$

Examples

- ◆ Dark Matter + dark gauge field

$$\mathcal{L} = -\frac{1}{4} \mathcal{F}_{\mu\nu}^a \mathcal{F}_{\mu\nu}^a + |D_\mu \phi|^2 - m_\phi^2 |\phi|^2 - \lambda \phi^2 (H^\dagger H)$$



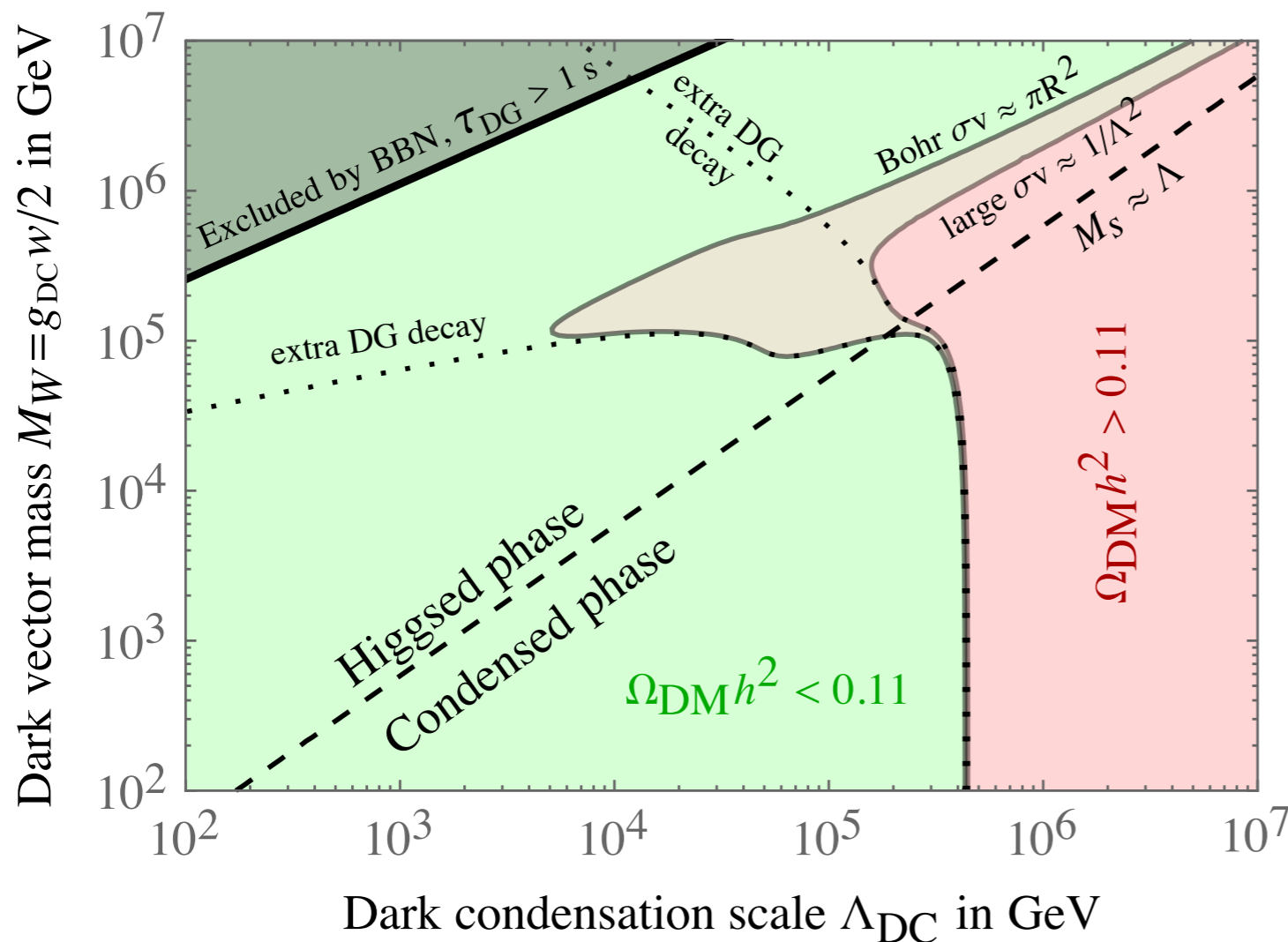
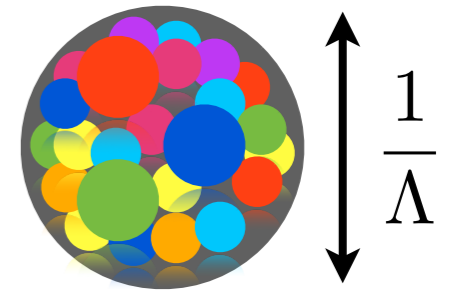
- ◆ Majorana fermion + singlet scalar

$$\mathcal{L} = \frac{1}{2} (\partial_\mu \phi)^2 - \frac{m_\phi^2}{2} \phi^2 + \bar{\chi} (i\partial - m_\chi) \chi - g_\chi \phi \chi \chi - \lambda \phi^2 (H^\dagger H)$$



Composite Dark Matter

- ◆ If dark group confines, Dark Matter is bound state of strong “dark force” at $\Lambda \sim \text{few} \times (1 - 100 \text{ TeV})$
- ◆ Stable thanks to accidental “dark baryon number”



constituents
can be fermions

Strassler, Zurek 2006
Antipitin, Redi, Strumia, Vigiani 2015
... many more...

or scalars

Hambye 2008
DB, Di Luzio, Landini,
Strumia, Teresi 2019

✗ Thermal mass $\sim \text{few} \times 100 \text{ TeV}$: not directly testable at colliders

Scalar singlets

- Can test the portal interaction with the singlet sector

$$\mathcal{L}_\phi = -\frac{1}{2}m_\phi^2\phi^2 - V(\phi) - \mu\phi H^\dagger H - \frac{\lambda}{2}\phi^2 H^\dagger H$$

induces mixing between
Higgs and scalar singlet

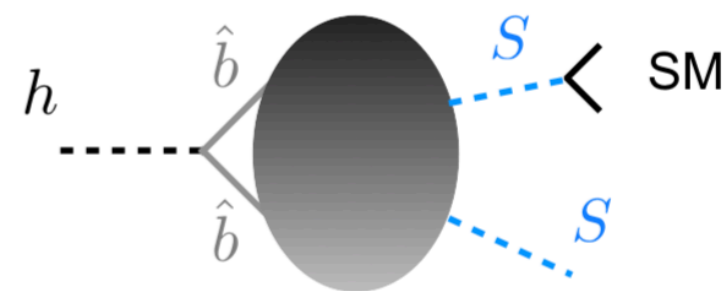
($\phi \rightarrow -\phi$ symmetry is not exact)

controls BR($h \rightarrow \phi\phi$),
or BR($\phi \rightarrow hh$)

- ϕ can be e.g. Goldstone, or glueball of dark gauge interactions...

- ϕ can mediate DM freeze-out if $M_{\text{DM}} > M_\phi$

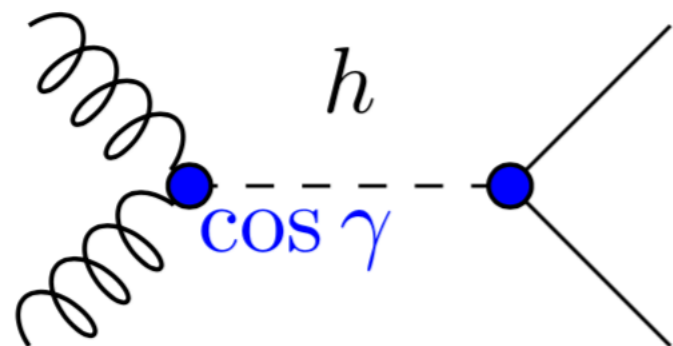
$$\mathcal{L}_{\text{DM}} = y_\chi \phi \chi \chi$$



Scalar singlets @ FCC

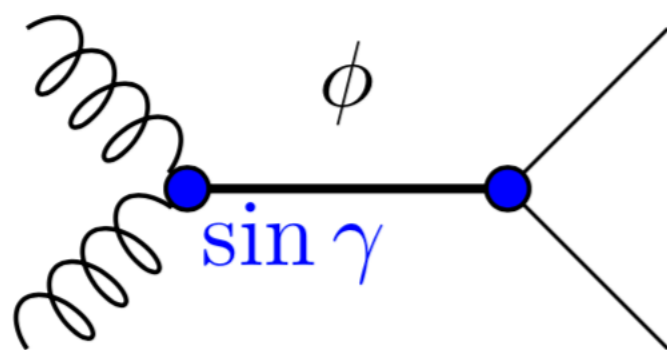
- Two complementary ways to look for the singlet at colliders:

Higgs signal strengths



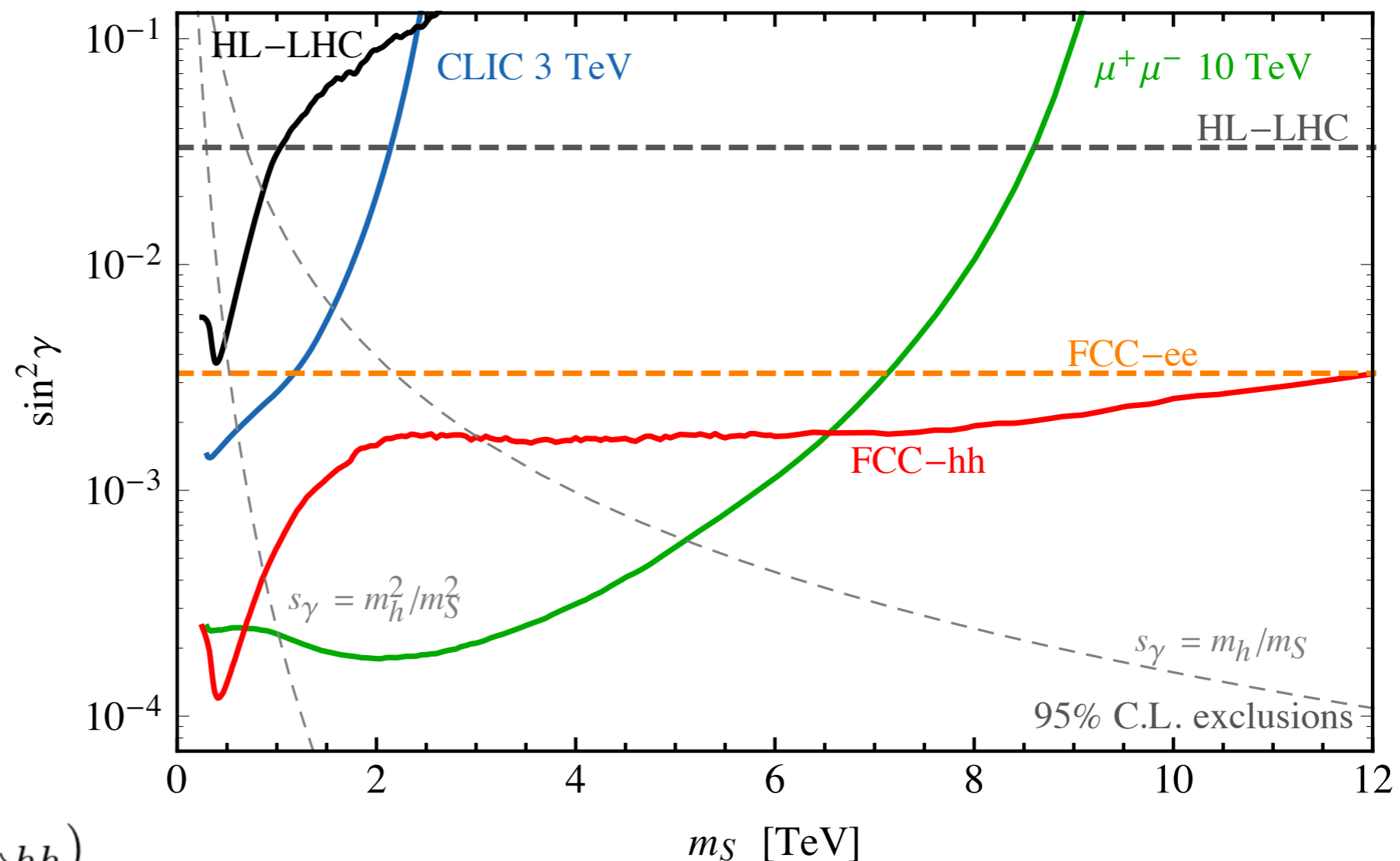
$$\mu_h = c_\gamma^2 \mu_{\text{SM}}$$

Direct searches



$$\mu_\phi = s_\gamma^2 \mu_{\text{SM}} \times (1 - \text{BR}_{\phi \rightarrow hh})$$

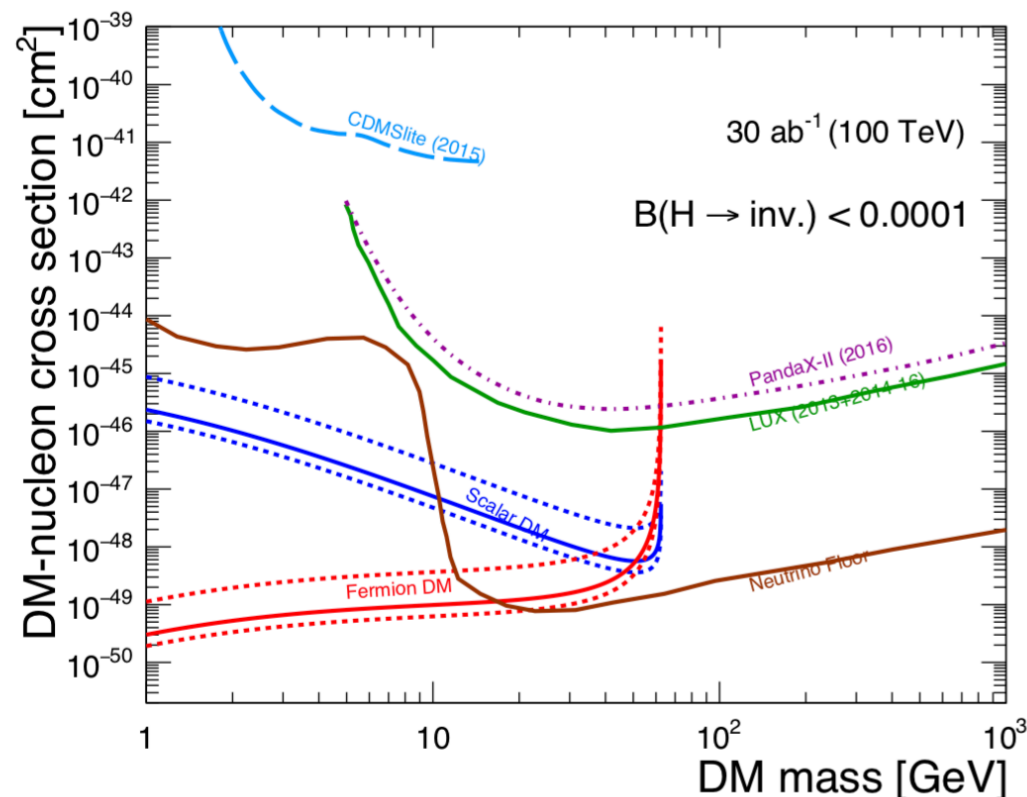
DB, Sala, Tesi 1505.05488
DB, Redigolo, Sala, Tesi 1807.04743



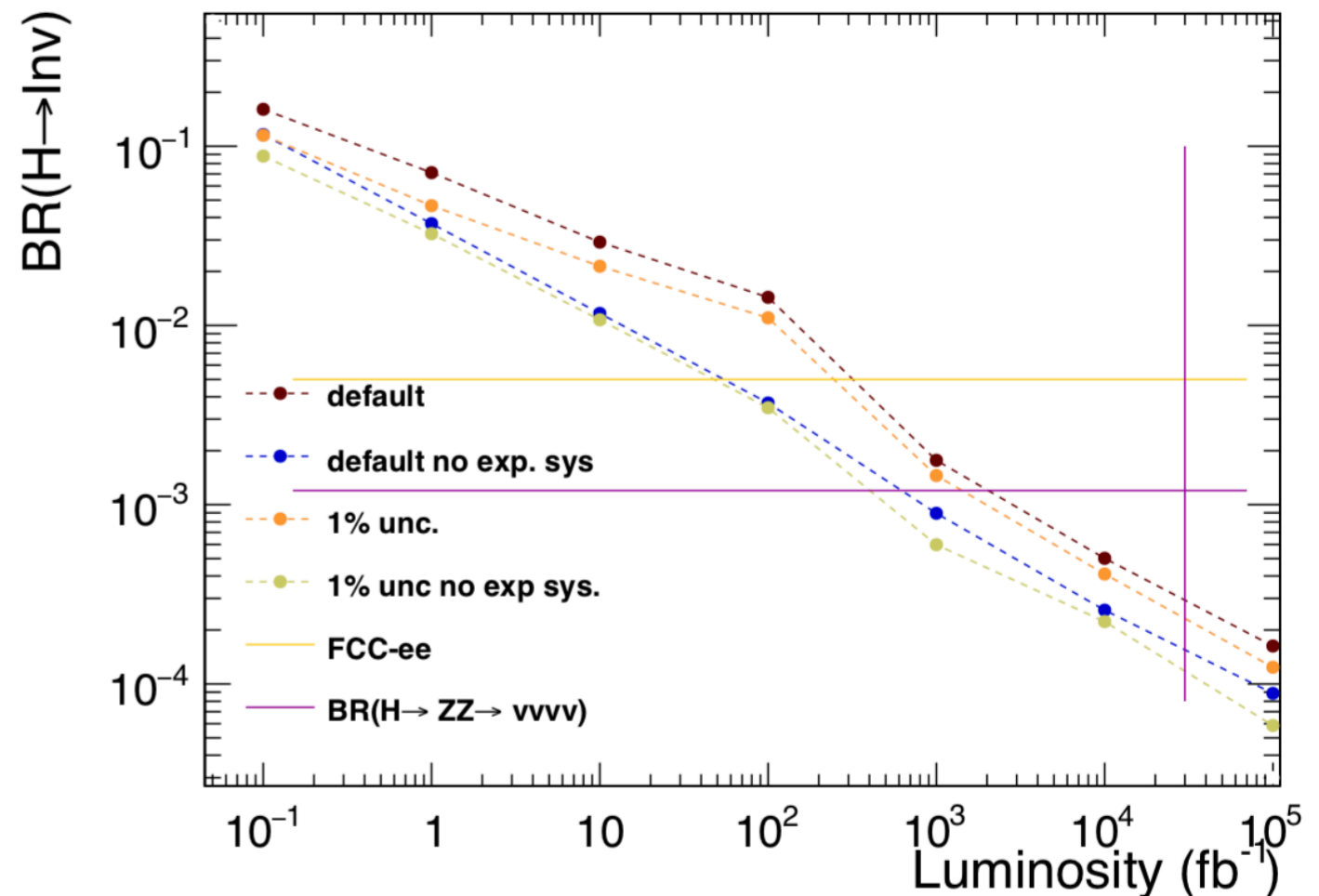
Light singlets: invisible Higgs

- ◆ In the zero-mixing limit the singlet is invisible

- ▶ $BR(h \rightarrow inv) < \text{few} \times 10^{-3}$
(missing mass @ FCC-ee)
- ▶ $BR(h \rightarrow inv) < \text{few} \times 10^{-4}$
from Higgs p_T distribution
in VBF or ttH @ FCC-hh



CERN-ACC-2018-0045



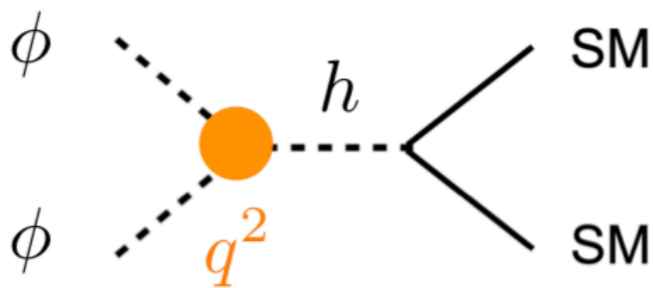
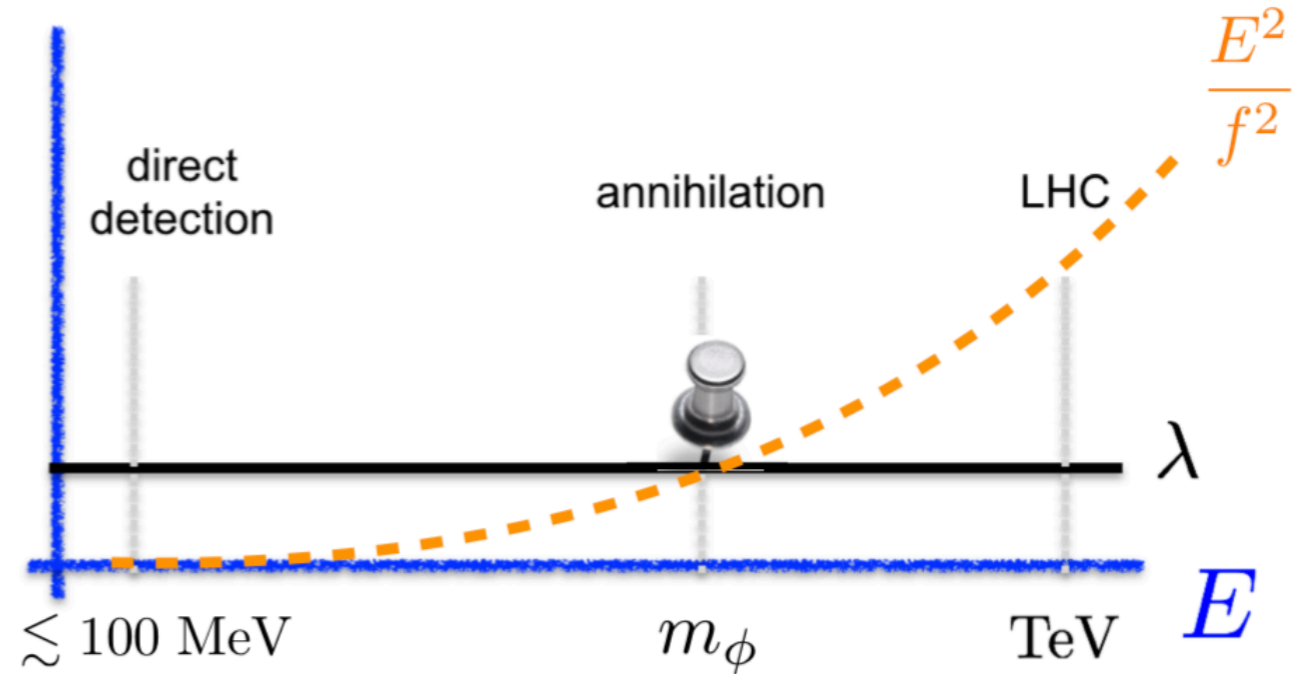
- ◆ No Higgs-singlet mixing: singlet is stable, can be identified with DM
stronger than Direct Detection prospects for low mass DM

Derivative Higgs portal

A viable alternative:

$$\mathcal{L} \supset \frac{1}{2f^2} \partial_\mu(\phi^2) \partial^\mu(H^\dagger H)$$

$$g_{\text{DM-SM}}^2(E)$$



annihilation $q^2 \sim m_{\text{DM}}^2$



scattering
 $q^2 \sim m_{\text{DM}}^2 v_{\text{DM}}^2$

$$v_{\text{DM}} \sim 10^{-3}$$

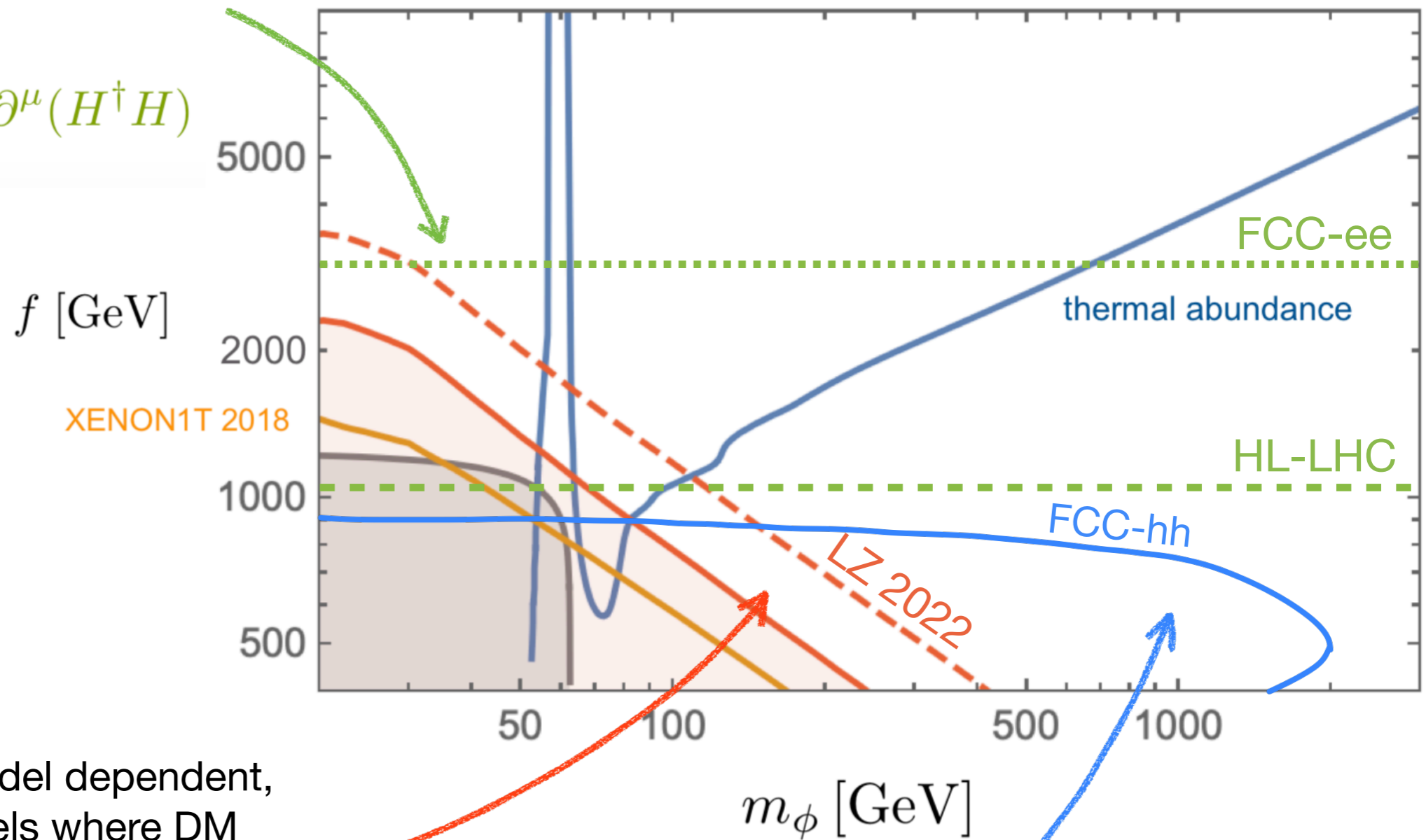
For the derivative Higgs portal,
 direct detection is automatically suppressed

Derivative Higgs portal

Higgs couplings: model dependent, e.g. composite models where DM and Higgs are pNGB

$$\mathcal{L} \supset \frac{1}{2f^2} \partial_\mu (H^\dagger H) \partial^\mu (H^\dagger H)$$

Ruhdorfer, Salvioni, Weiler 1910.04170

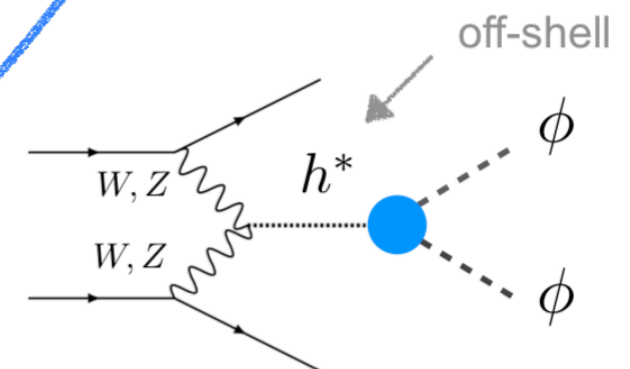


Direct Detection: model dependent, e.g. composite models where DM mass arises from down quark loops

$$\mathcal{L} \supset y_d \bar{q}_L H d_R \frac{\phi^2}{2f^2}$$

Vector boson fusion

$$\mathcal{L} \supset \frac{1}{2f^2} \partial_\mu (\phi^2) \partial^\mu (H^\dagger H)$$

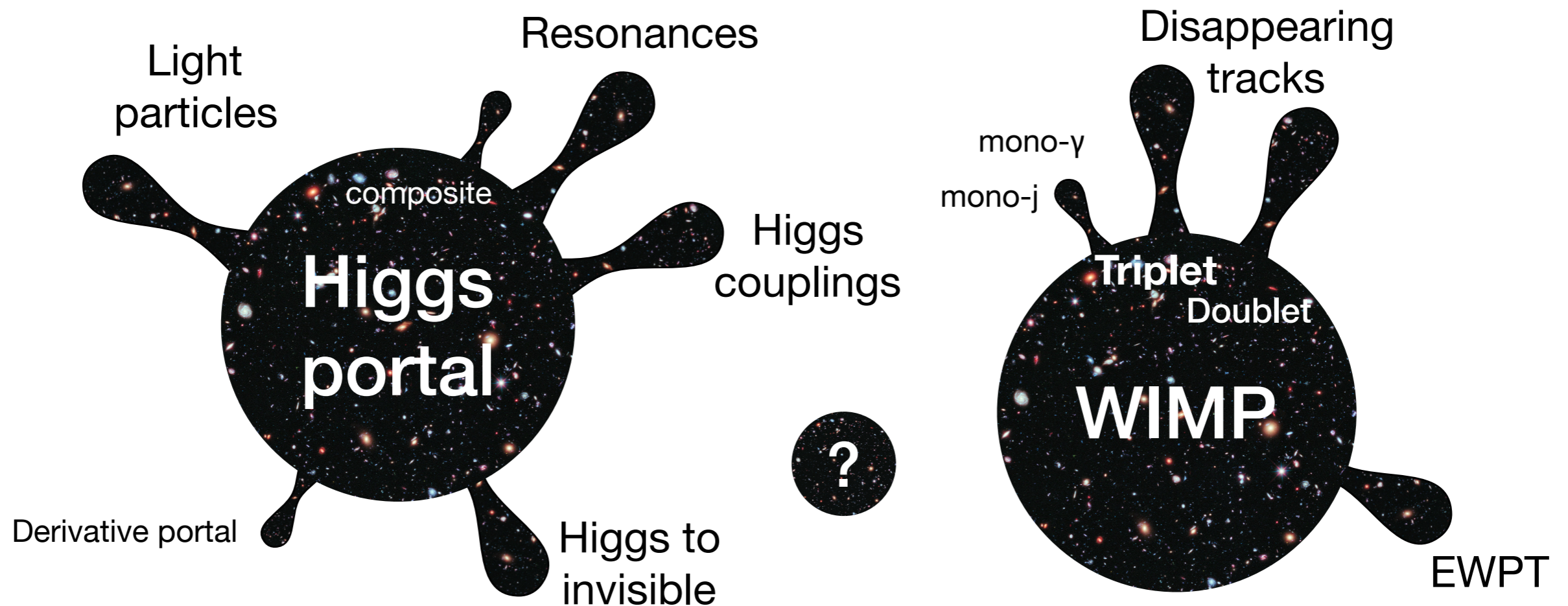


Summary

Searches for motivated DM candidates are difficult (EW states, singlets...)

Many viable models... Not everything can be probed at FCC.

High energy + precision: excellent prospects in many interesting cases.



Complementary information will come from ID and DD in the coming years