Dark matter indirect detection limits including complete annihilation patterns

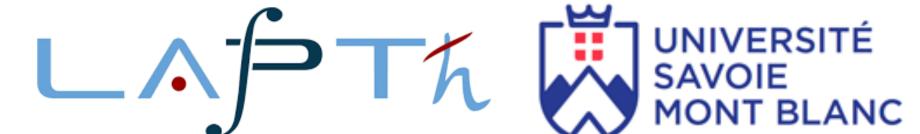
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Based on **2210.01220** [hep-ph]















DARK MATTER

85% of the total matter of our Universe

Relic density observed experimentally by Planck:

$$\Omega_{\chi} h^2 \simeq 0.1200 \pm 0.0012$$

Ref: Ade et al. 2016, Astrophys. 594, A13

Its identification would reveal new Physics

Proving its existence and nature would improve our understanding of the Universe

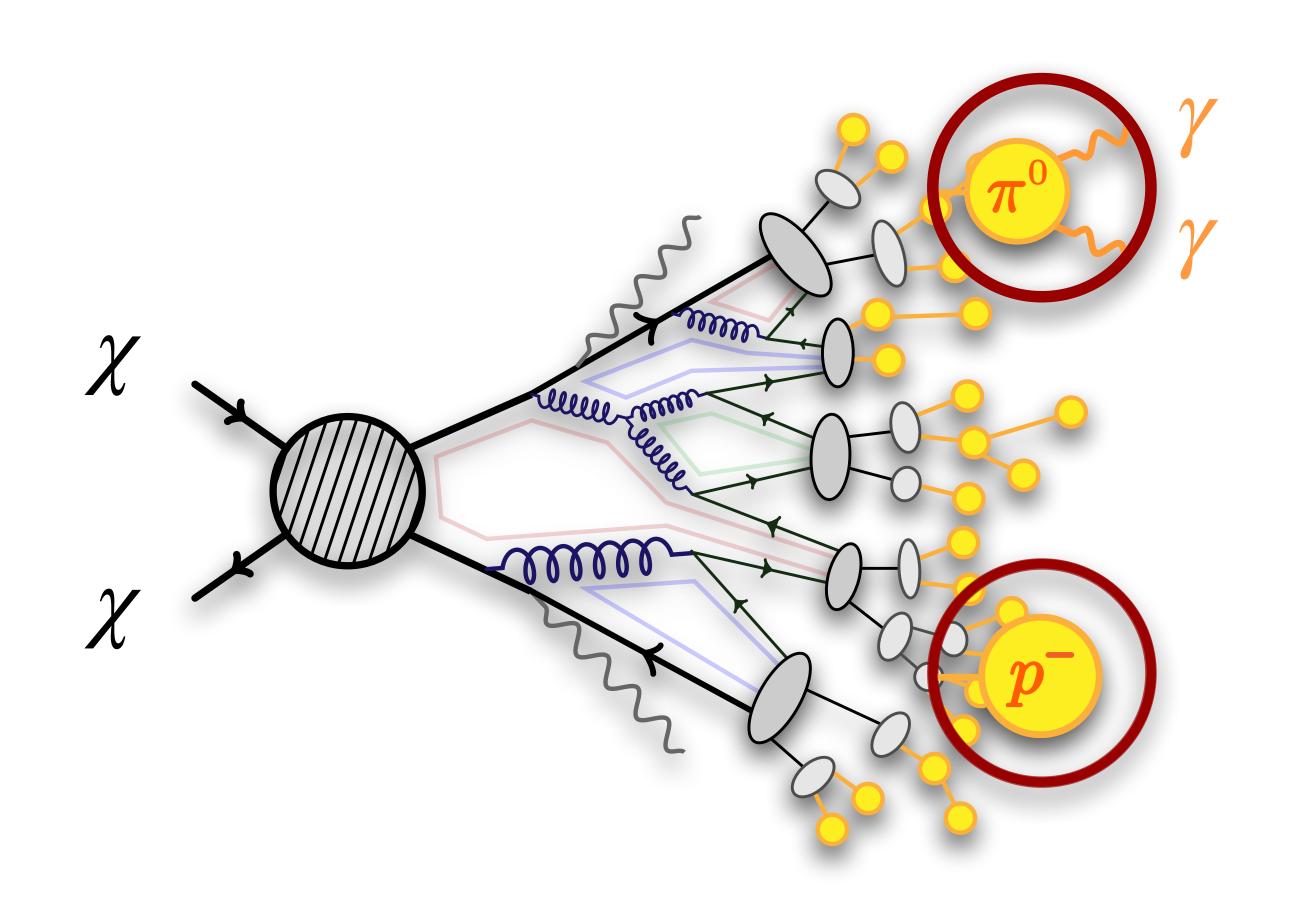
GOALS

Study of the impact of a more complete particle model

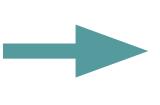
New prediction of DM upper limits with CTA mockdata of Sculptor

- Previously: use of individual annihilation channels
- This work: Collaboration with a theoretician to include a more complex and more complete model

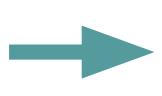
INDIRECTSEARCHES



Dark Matter (DM) annihilation



Standard Model particles (bosons, quarks, leptons)



Final state products such as γ rays

INDIRECTSEARCHES

Expected y-ray flux from DM annihilation

Astrophysical

J factor

$$\frac{d\Phi\left(\langle\sigma v\rangle,J\right)}{dE} = \frac{1}{4\pi} \frac{\langle\sigma v\rangle}{2m_{\chi}^2} \sum_{f} \mathsf{BR}_{f} \frac{\mathsf{d}N_{f}}{\mathsf{d}E} \times \int_{\Delta\Omega} \int_{\log} \rho_{\mathrm{DM}}^2 ds d\Omega$$

Particle Physics

factor

 $<\sigma v>$ = annihilation cross-section m_X = DM particle mass BR_f = branching ratio

 dN_f/dE = differential spectrum ρ_{DM} = DM density

where

STATISTICALANALYSIS

LOG-LIKELIHOOD RATIO TEST STATISTICS

Constrained

minimization

$$\Lambda = -2 \ln \frac{\mathcal{L}_{H_0}}{\mathcal{L}_{H_1}} = -2 \ln \frac{\mathcal{L}(\langle \sigma v \rangle_0 | \hat{N}_B, \hat{J})}{\mathcal{L}(\langle \hat{\sigma v} \rangle, \hat{N}_B, \hat{J})}$$

Global

minimization

Ref: Cowan et al, 2010 Eur.Phys.J.C71:1554,2011

<σv>

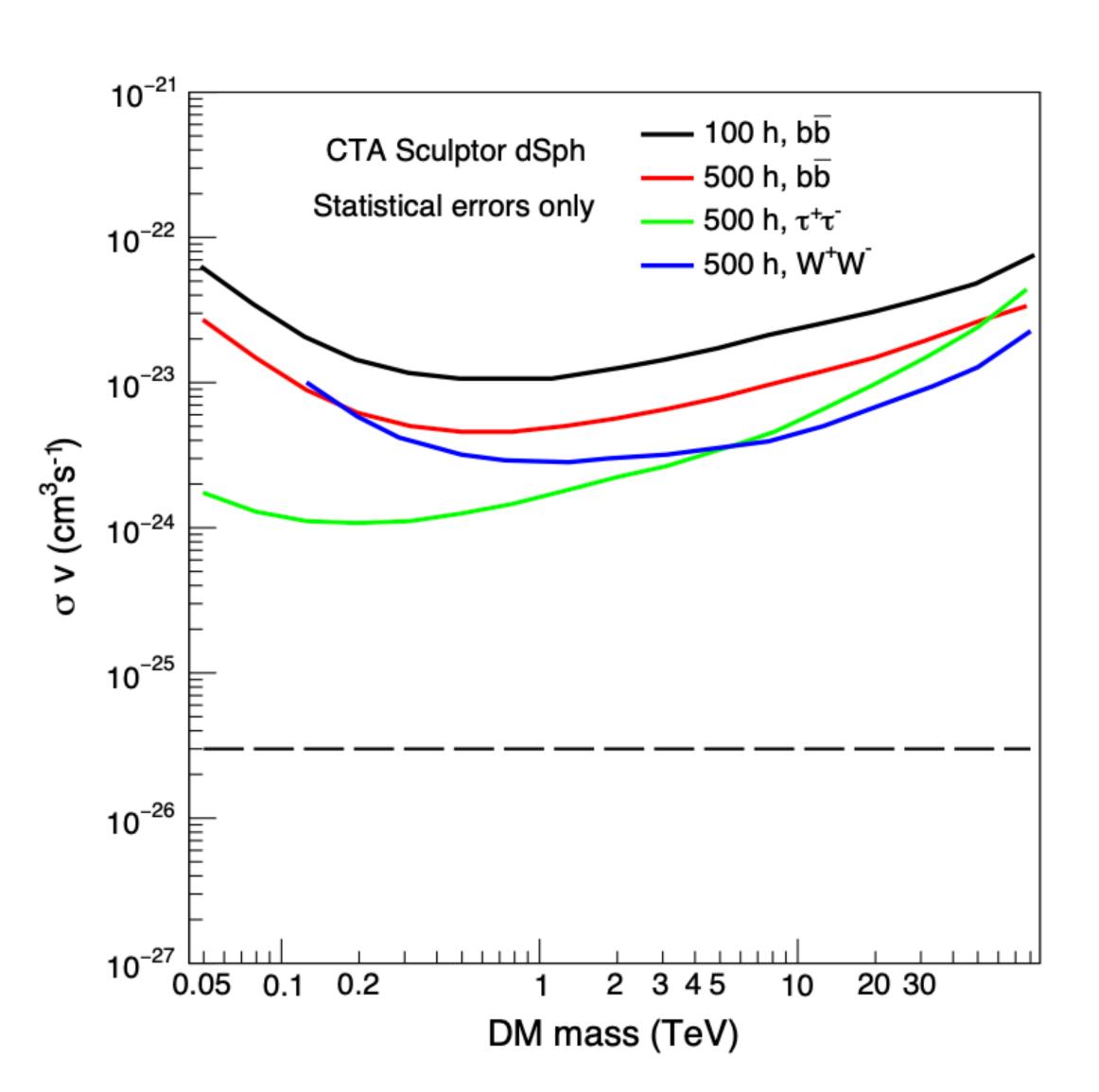
Parameter of interest

N_B, J

Nuisance parameters

2.71 at 95% Confidence Level

UPPER LIMITS



- Each annihilation channel treated independently
- Corresponding to a branching ratio of 100%
- Simplest model possible where all DM particles annihilate through the same channel

We change the particle physics model?

Standard model extended by an additional scalar field (DM)

$$V_{\text{scalar}} \supset 2\lambda_H v^2 h^2 + \frac{1}{2}\mu_S^2 S^2 + \frac{1}{4}\lambda_{SH} v^2 S^2 + \frac{1}{4}\lambda_{SH} v S^2 h + \lambda_{SH} S^2 h^2$$

DM mass

$$m_S^2 = \mu_S^2 + \frac{1}{2} \lambda_{SH} v^2$$

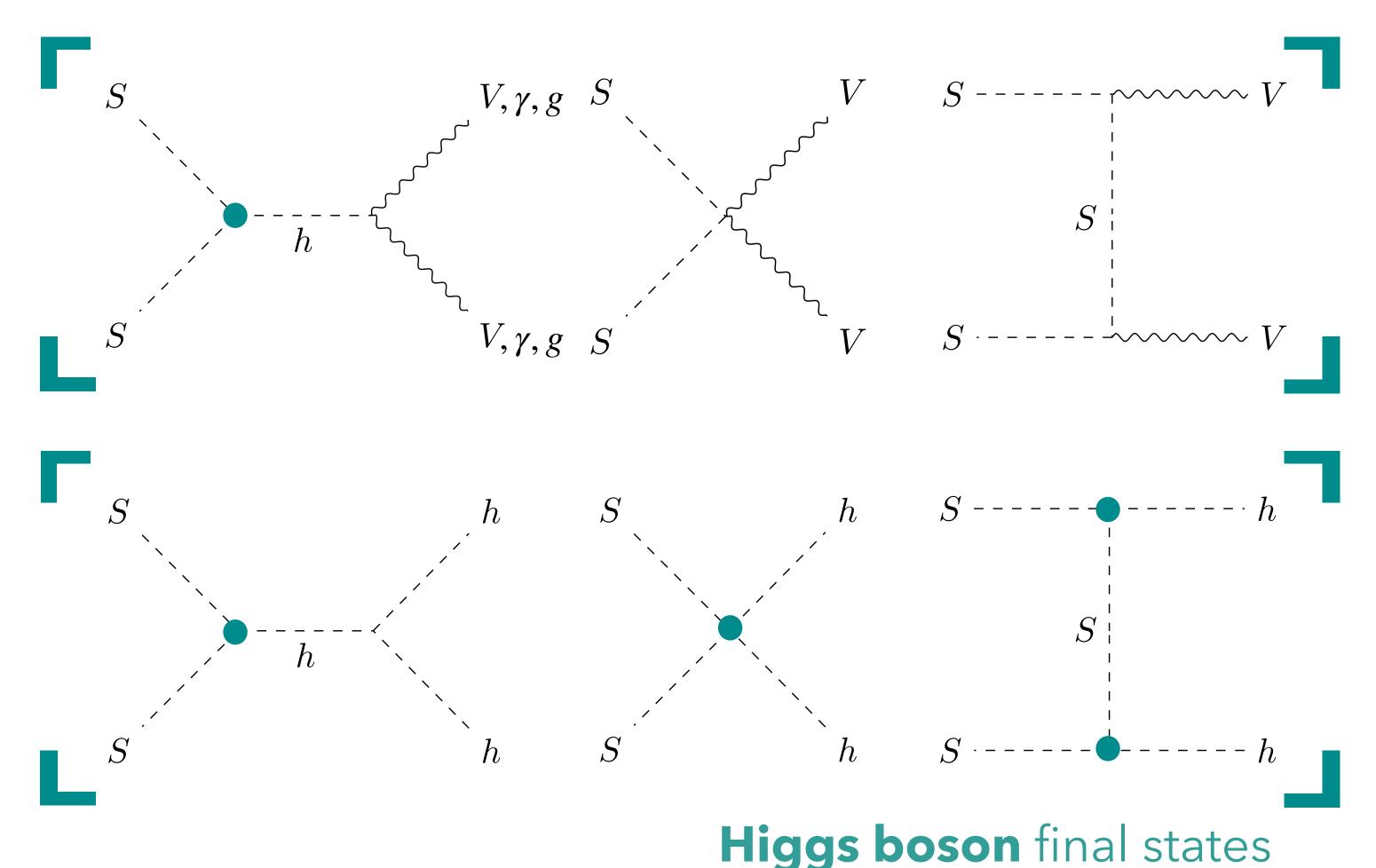
DM – Higgs interaction ("Higgs portal")

Phenomenology governed by

$$m_S$$
 (DM mass)

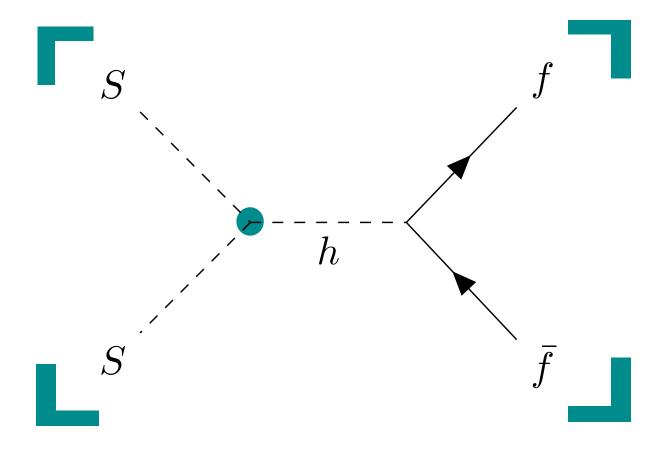
 λ_{SH} (DM coupling)

Possible dark matter annihilation channels (DM relic density + indirect detection)



Gauge boson final states

$$V = Z^0, W^{\pm}$$



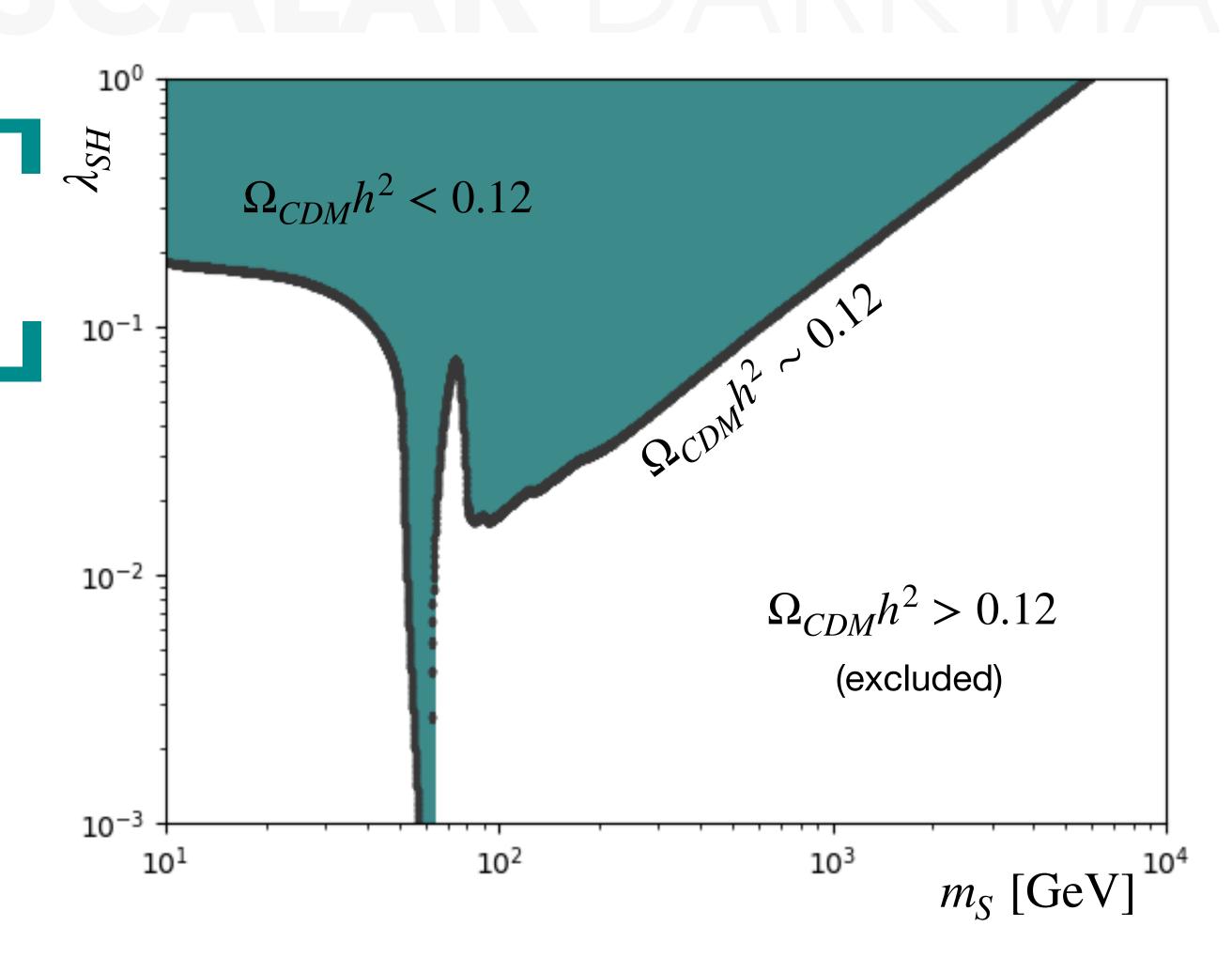
Quark or lepton final states

$$f = u, d, c, s, b, t, e, \mu, \tau$$

DM coupling vs DM mass

Relic density and branching ratio grid computed using micrOMEGAs

Ref: Bélanger, Pukhov et al. 2003 - 2022



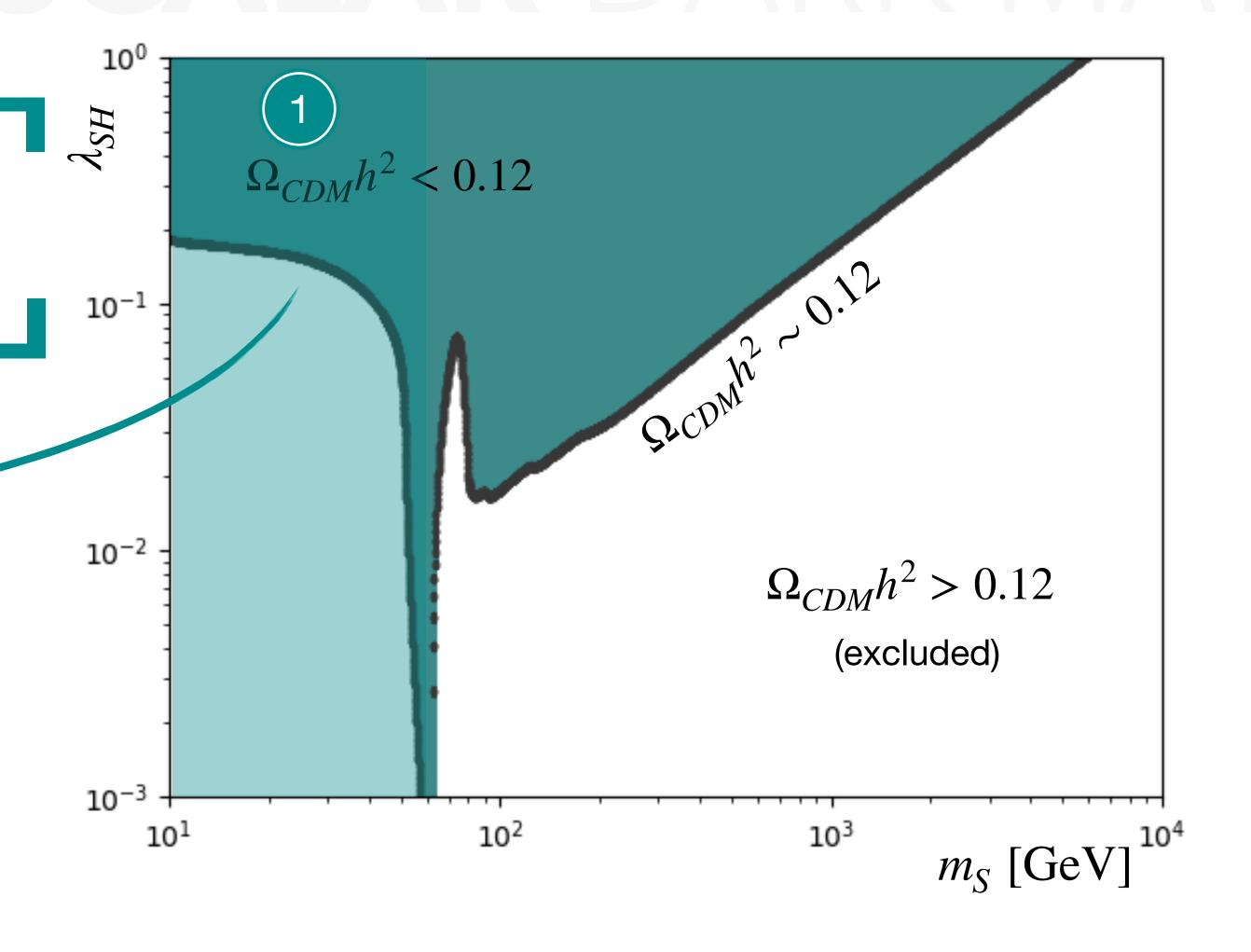
DM coupling vs **DM** mass

Relic density and branching ratio grid computed using micrOMEGAs

Ref: Bélanger, Pukhov et al. 2002 - 2022

$$SS \rightarrow b\bar{b} \sim 75 - 85\%$$

 $SS \rightarrow \tau^{+}\tau^{-} \sim 7 - 10\%$
 $SS \rightarrow c\bar{c} \sim 3 - 4\%$
 $SS \rightarrow gg \sim 3 - 15\%$



DM coupling vs **DM** mass

Relic density and branching ratio grid computed using micrOMEGAs

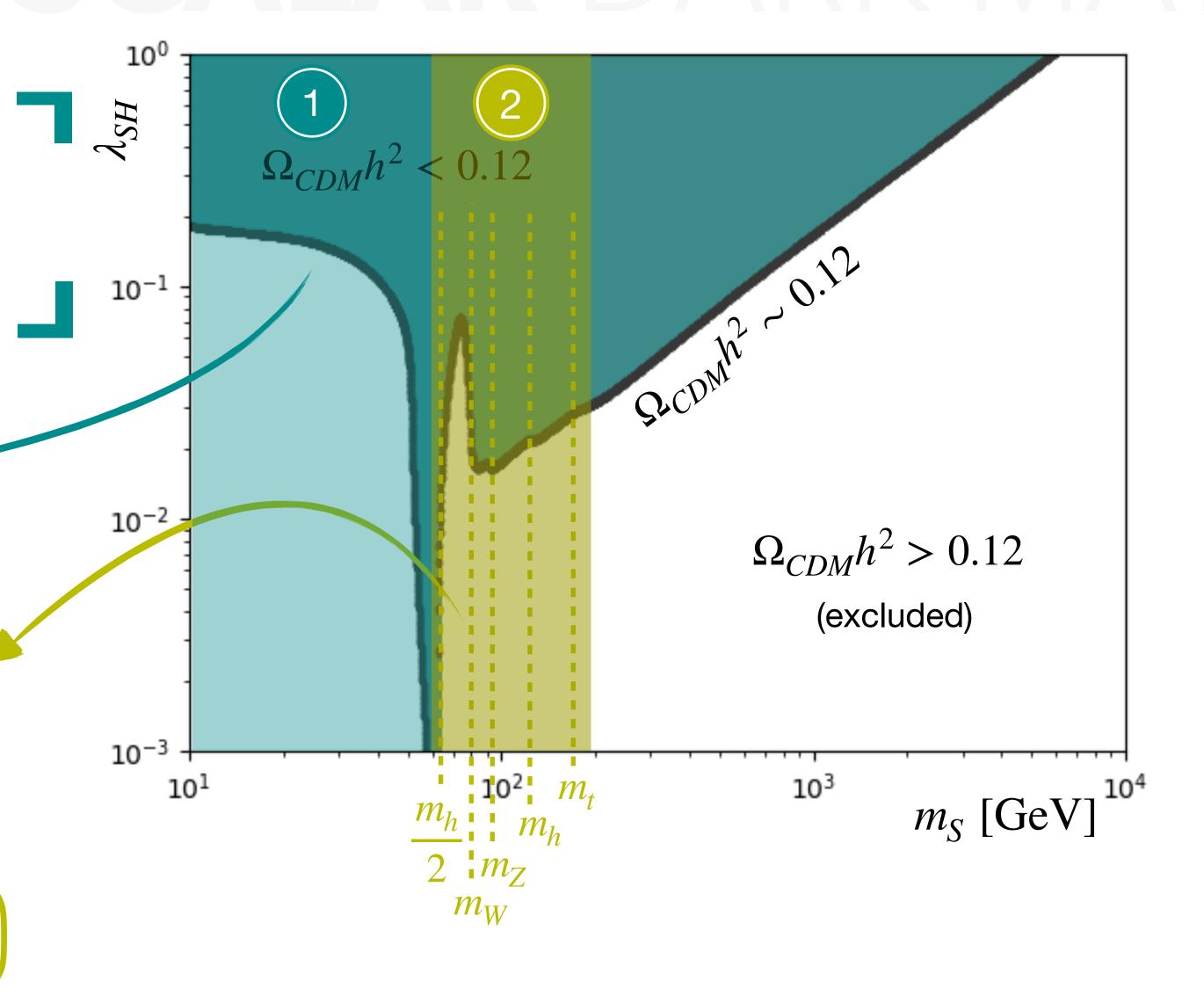
Ref: Bélanger, Pukhov et al. 2002 - 2022

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Region of frequent change of the dominant annihilation channel

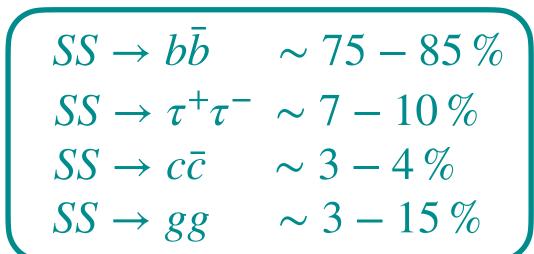
New annihilation channels open, Higgs resonance at $m_h/2$



DM coupling vs **DM** mass

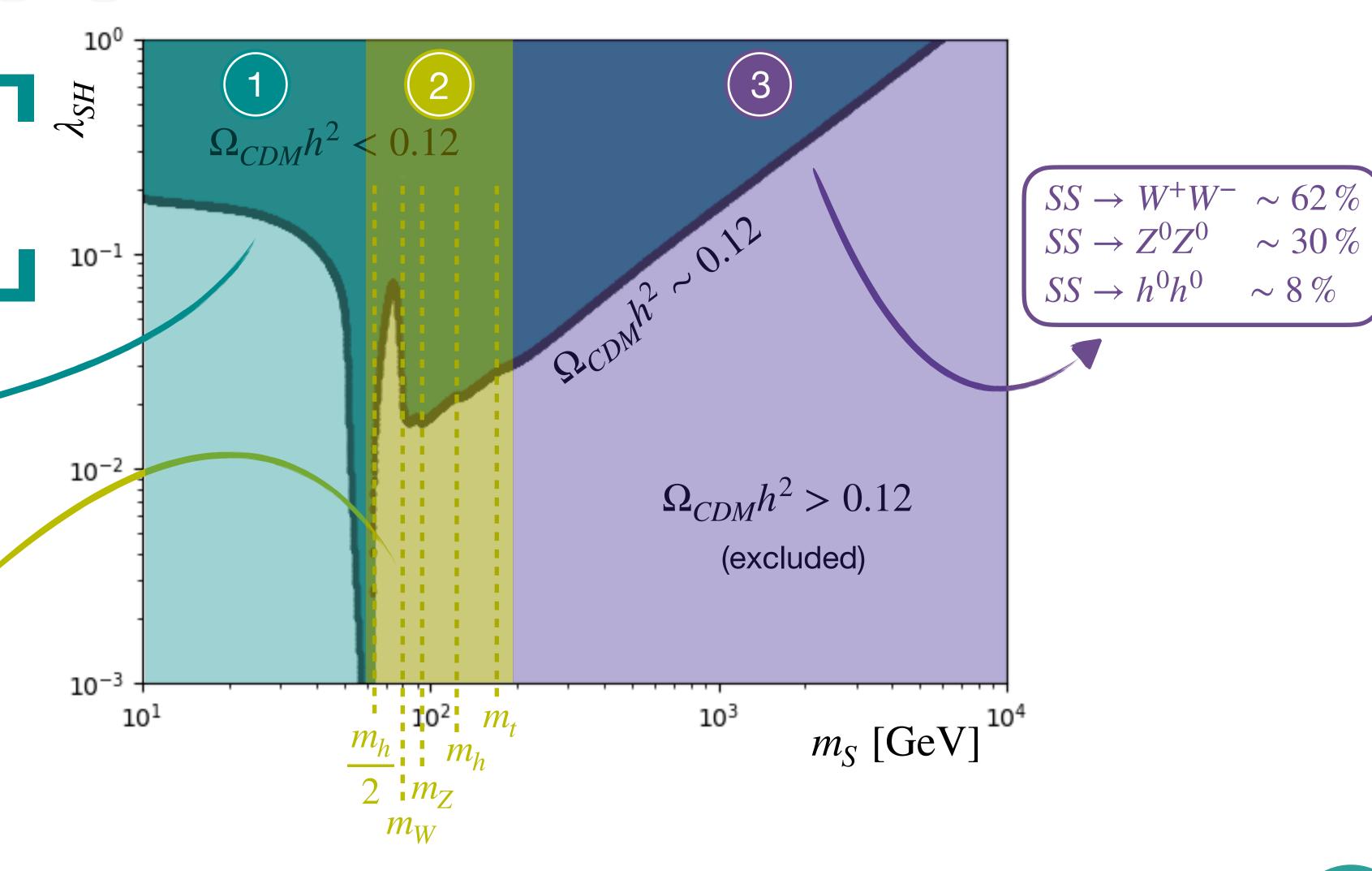
Relic density and branching ratio grid computed using micrOMEGAs

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Region of frequent change of the dominant annihilation channel

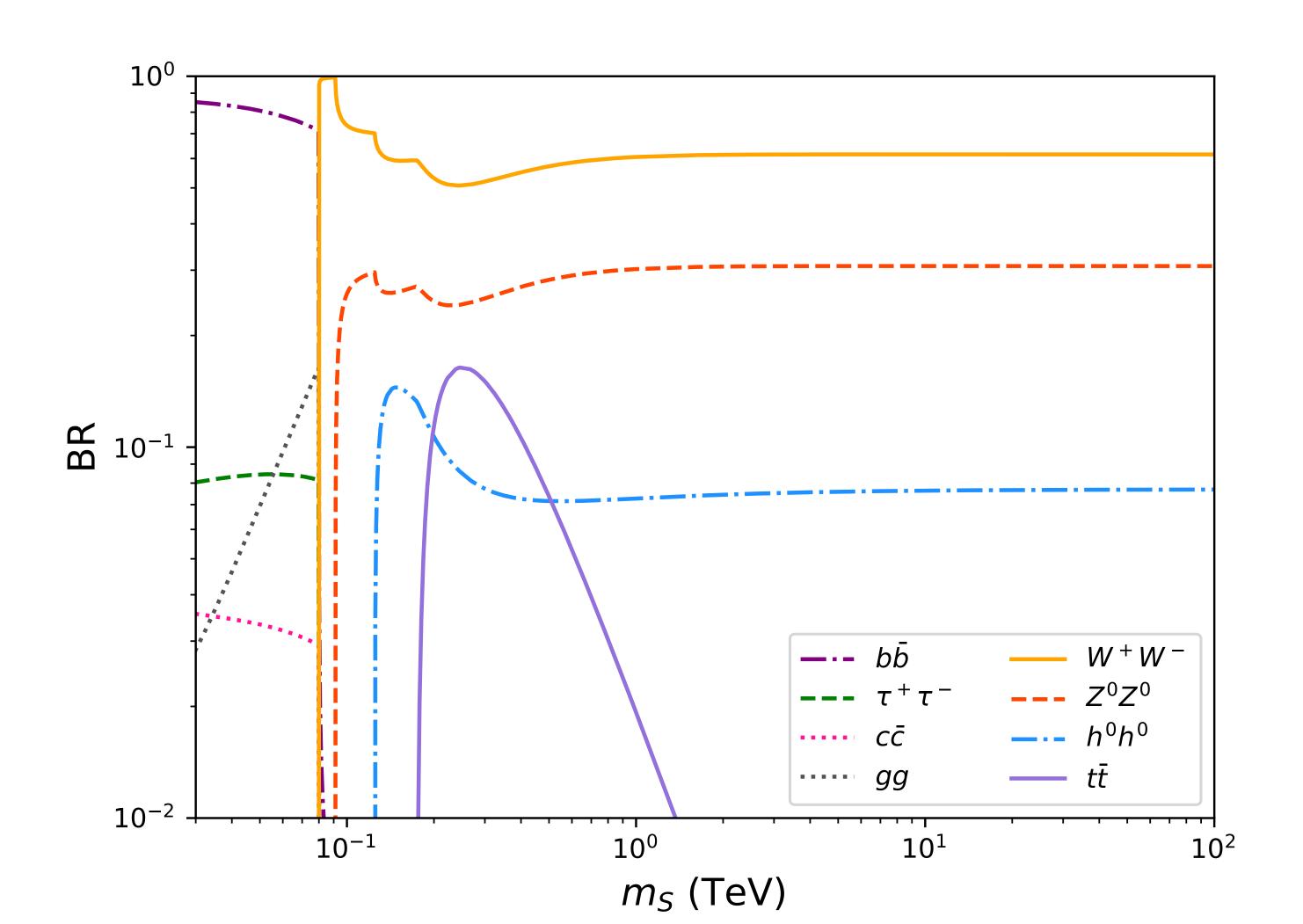
New annihilation channels open, Higgs resonance at $m_h/2$





All annihilation channels treated all together whose branching ratio varies with respect to the DM mass

Branching Ratio according to the relic density constraint



None of the annihilation channels are at 100% branching ratio over the full mass range

For the remaining part, we focus on the case where the relic density constraint is satisfied (black line in previous figure):

$$\Omega_{\chi} h^2 \simeq 0.1200 \pm 0.0012$$

Even in such a simple setup, the "100% hypothesis" is not justified...

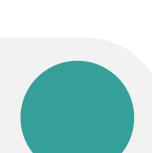
More complex models invoke an even richer phenomenology...

TARGET SOURGE

Dwarf galaxy selected for the CTA dark matter program



South Hemisphere $I = 287.62^{\circ}, b = -83.16^{\circ}$



Sculptor



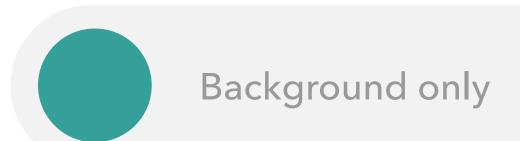
Mock data prepared with Gammapy 0.18.2

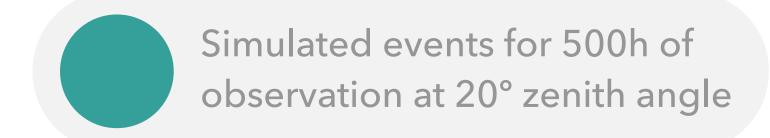
J factor $Log_{10} J_{0.1} = 18.3 \pm 0.3$



Ref: Bonnivard et al, 2015 ApJ 808 L3







NEW UPPER LIMITS

Computation of the predicted DM cross section VS DM particle mass

Expected limits - Sample of 500 Poisson realizations of the simulated background events



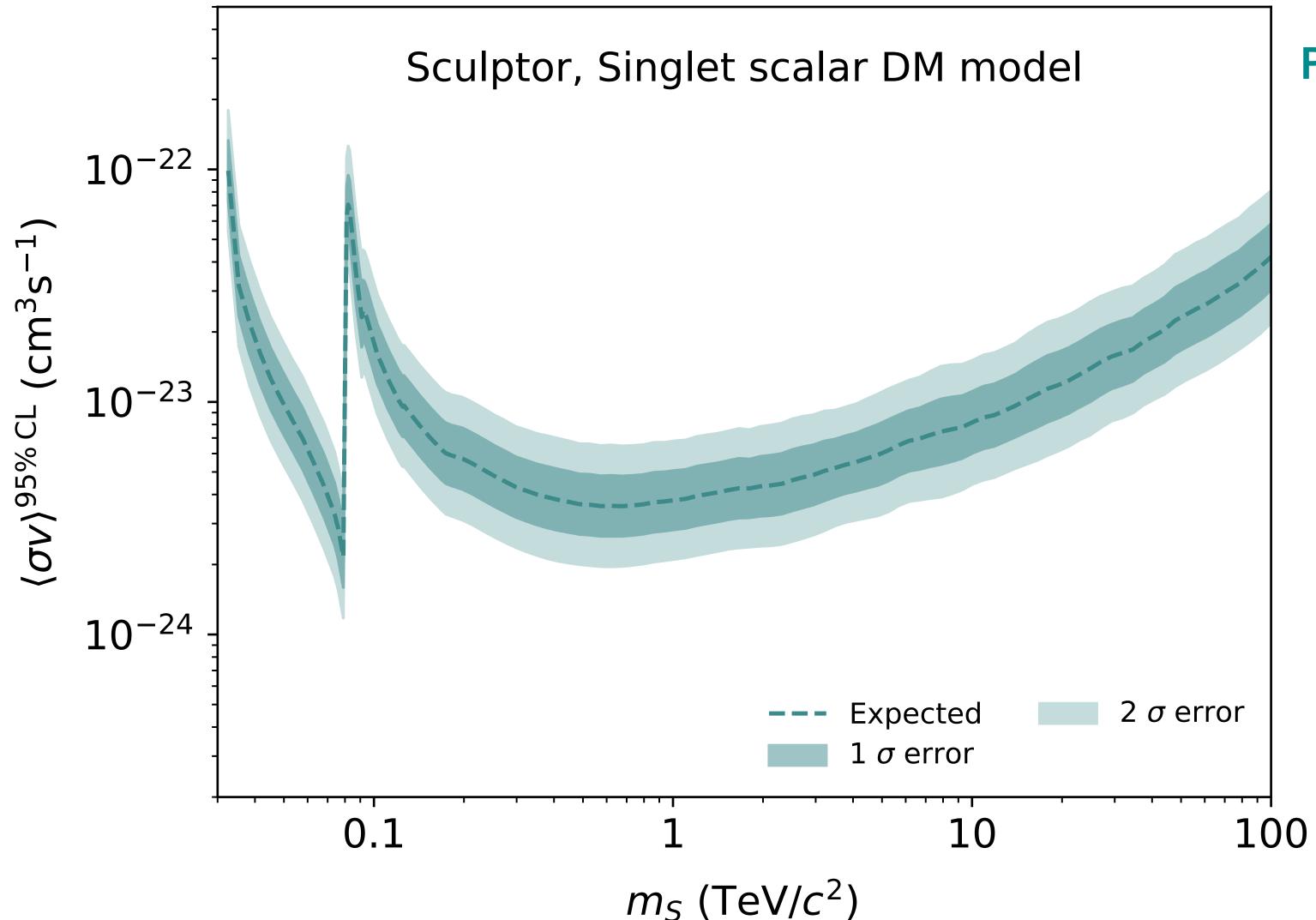


Mean expected limits

Mean of the derived <σν> distribution

Statistical uncertainty bands Standard deviation at 1 and 2σ

RESULTS



Predicted upper limit and uncertainties

Assuming a singlet scalar DM model γ-ray spectra taken from Cirelli et al. *JCAP* 03 (2011) 051

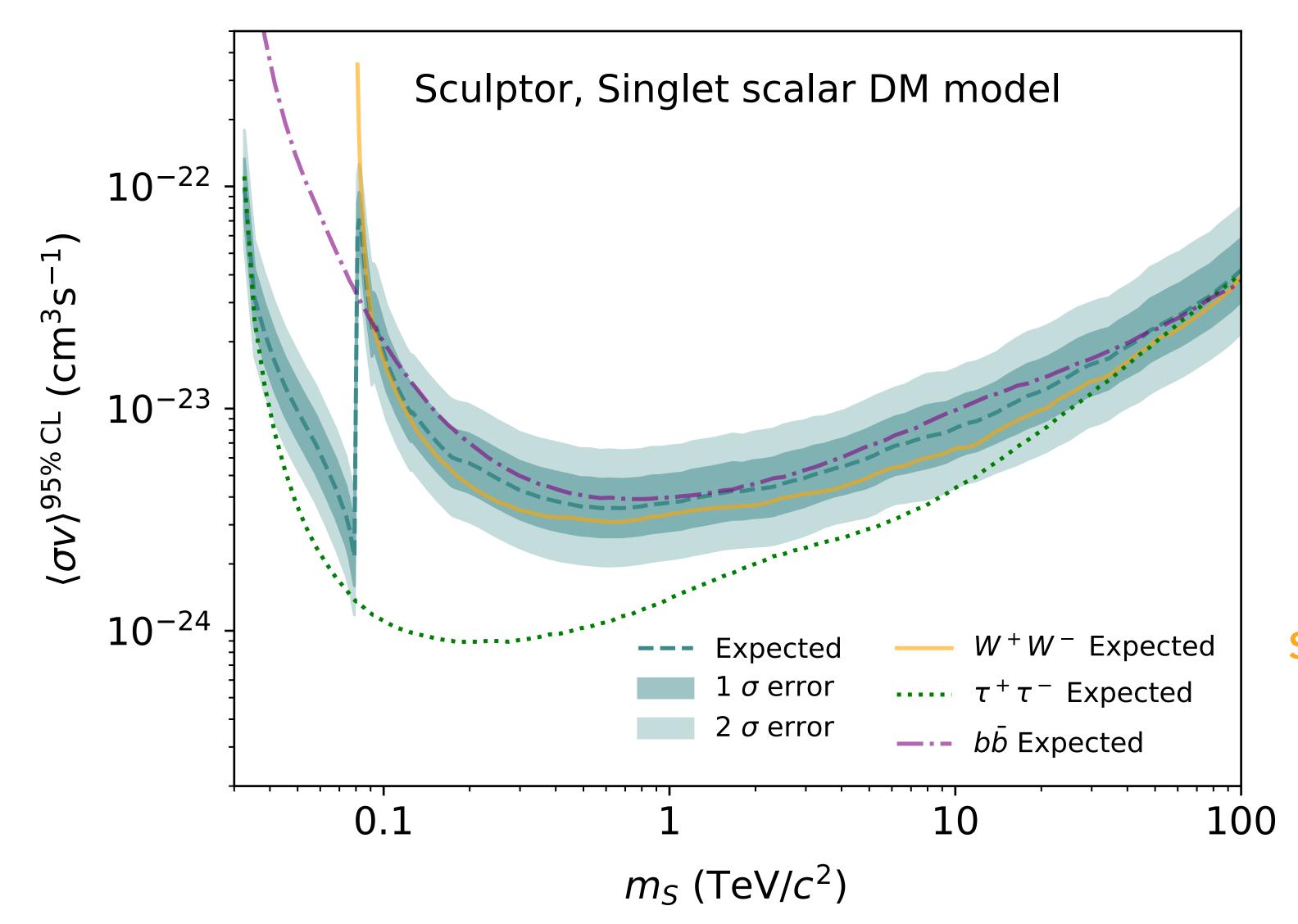
Inflection point

Due to the Higgs resonance

Sudden increase

Due to the opening of the WW channel

COMPARISONS



SINGLET SCALAR MODEL VS

100% W+W-

Below the W mass

No upper limit for 100% WW since the WW channel does not exist

Above the W mass

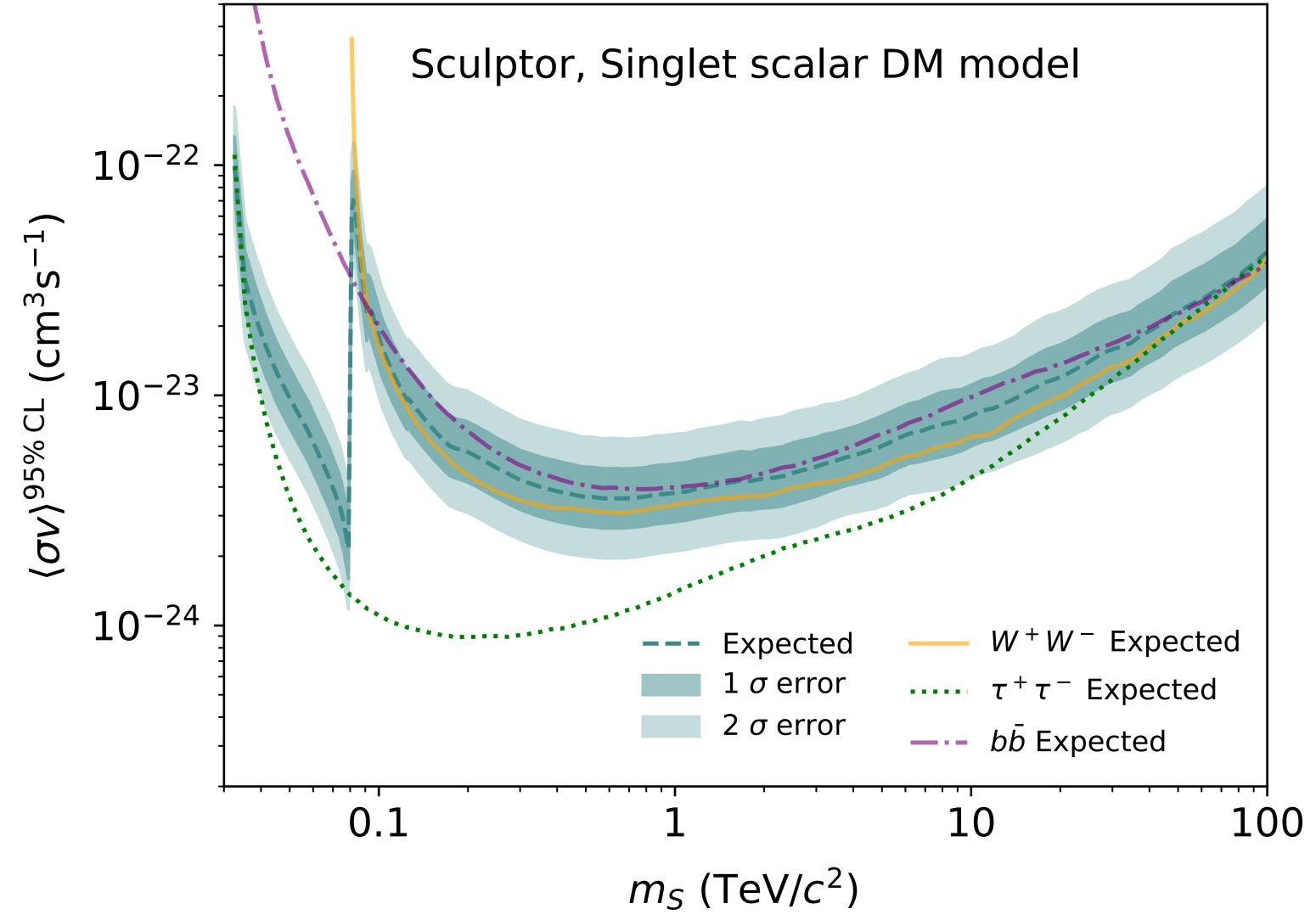
Additional contributions: ZZ, hh, tt which produce less y than the WW channel

Above 1 TeV

Stability with ~62% WW - 30% ZZ - 8% hh

More conservative limit with the singlet scalar DM model

COMPARISONS



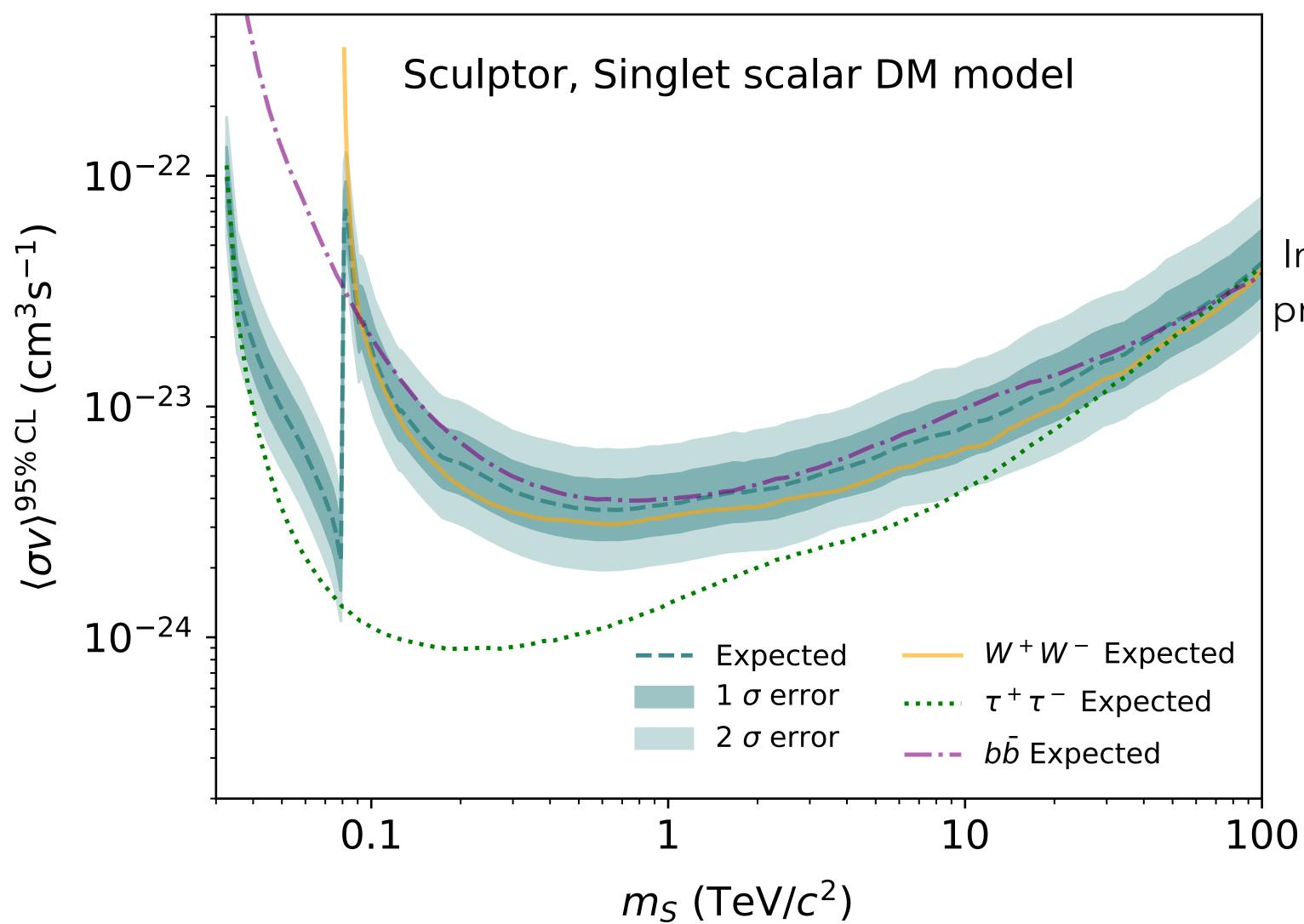
SINGLET SCALAR MODEL VS 100% T+T-

100% T+T- produces more γ rays
Leads to more constraining upper limits

However, in the singlet scalar model, this T+T-channel is never dominant

100% T+T = over estimation of the contribution

COMPARISONS



SINGLET SCALAR MODEL VS 100% bb

Below the W mass

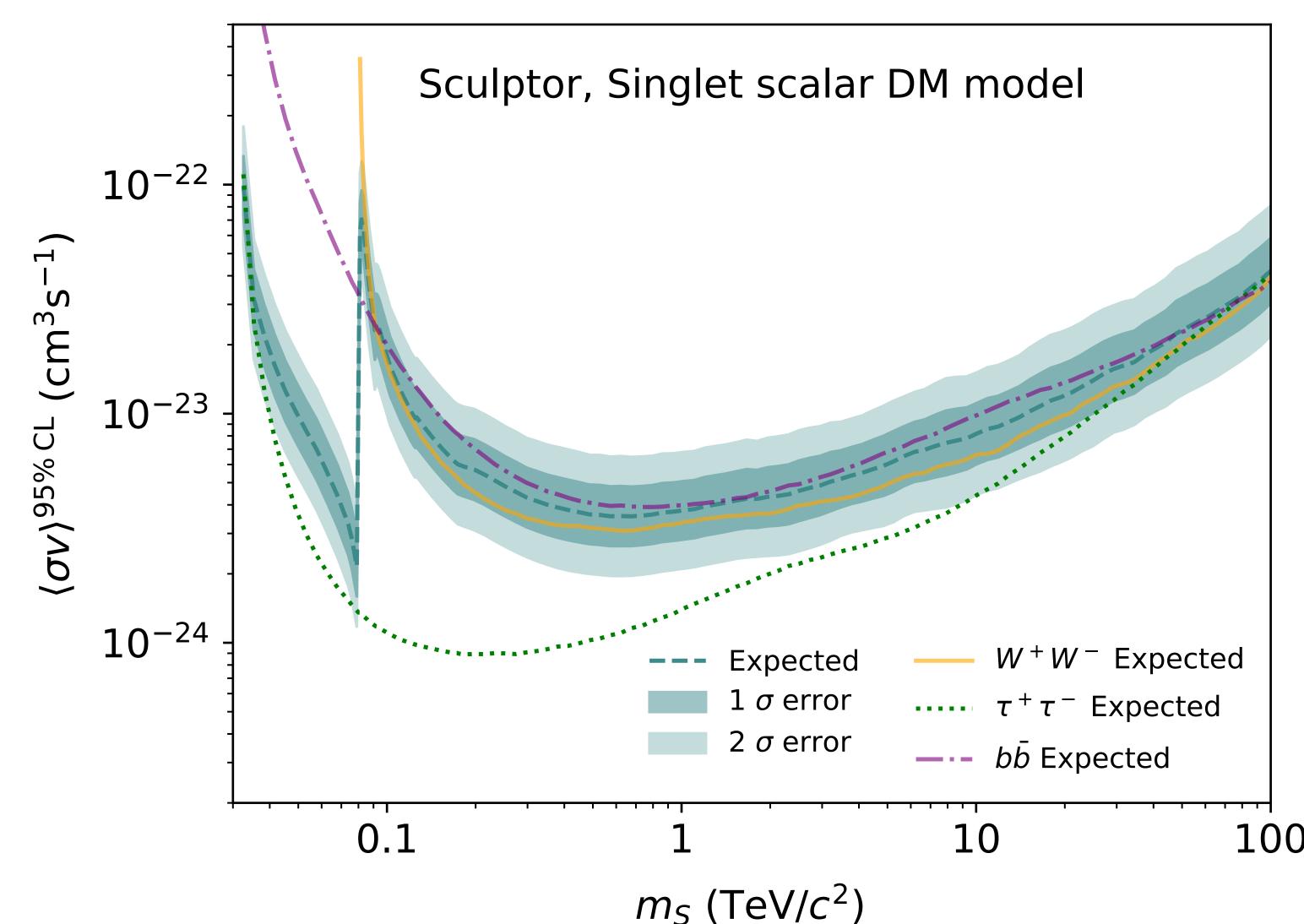
Singlet scalar much more constraining Important difference with 100% bb due to the presence of the subdominant channels $\tau^+\tau^-$, cc, and gg

Above the W mass

bb suppressed in the singlet scalar
bb produces less γ than the 4 contributing
channels WW, ZZ, hh, and tt (between the t
mass and 1 TeV)

More constraining limit with the singlet scalar DM model

RELATIVE ERRORS



WW

- Ranges between -22% and -6%
- Reaches values beyond 100% for masses at the W mass threshold

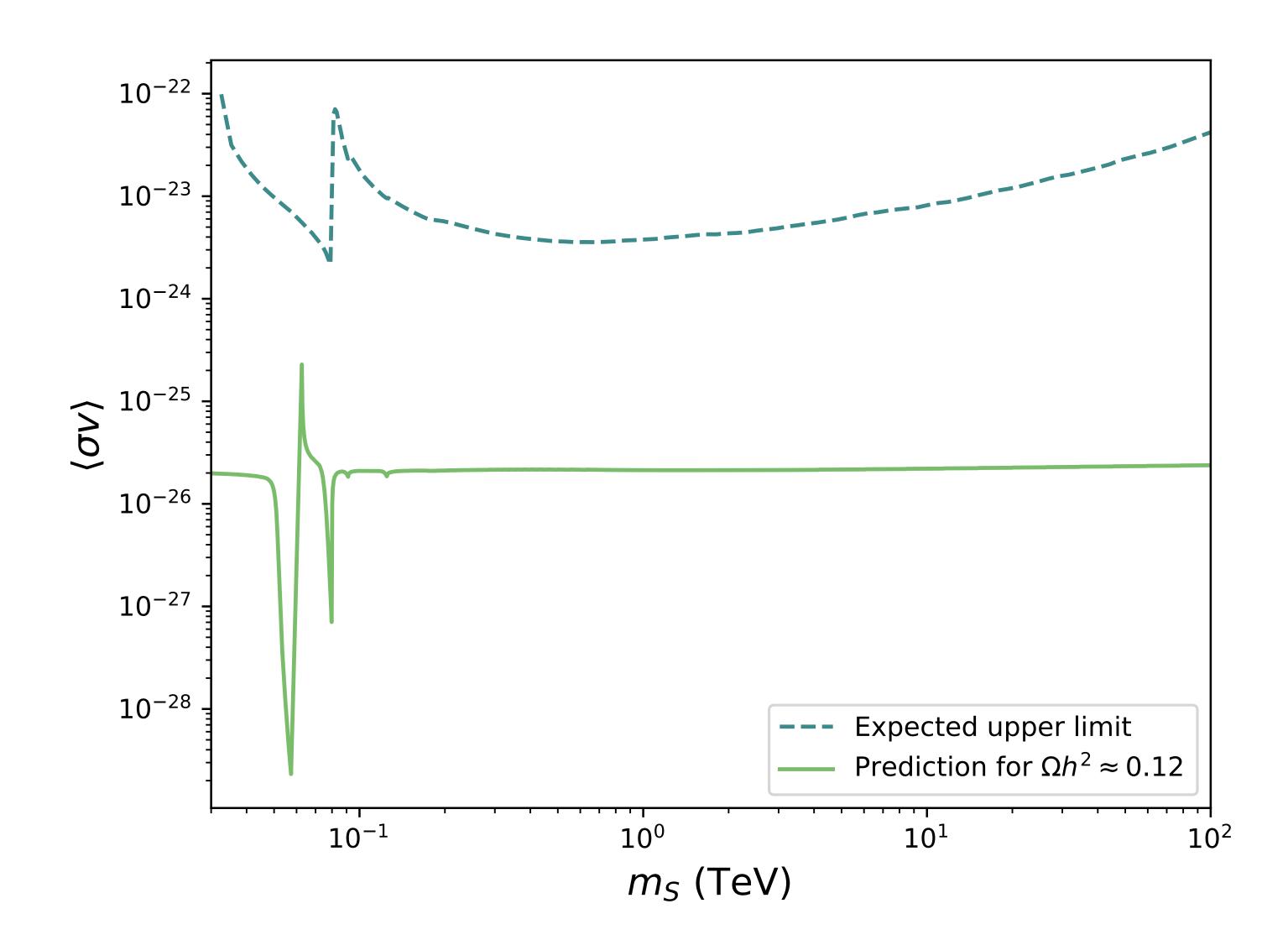
T+T-

- Ranges between -65% and +12% below the W mass threshold
- Around the W mass threshold reaches -98%
- Then decreases to -3% at 100 TeV

bb

- In the order of 1000% below the W mass threshold
- Then drops to -56% at the W mass threshold
- Then remains in the range of -10% to
 +38% after the threshold

EXCLUSION



Singlet scalar model **not excluded** by DM indirect detection

However, resonance and kinematical thresholds influence the exclusion curve together with the predicted annihilation cross-section

Generally, such fluctuations might lead to exclusion...

CONCLUSION & PERSPECTIVES

- Use of a more complex and more complete particle physics model
- Takes into account the full phenomenology with all annihilation channels at once
- Change of dominant annihilation channel(s) along with the DM particle mass
- Affects the predicted upper limits
- Feature can be expected in any particle physics model
- Derivation of a predicted upper limit and its 1σ and 2σ uncertainty bands
- Particle physics model could be used as well on the future data of CTA
- Paper submitted to JCAP & available on 2210.01220 [hep-ph]

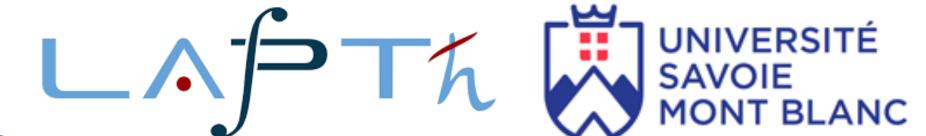
Thanks for your attention







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BACKUP



STATISTICALANALYSIS

Total likelihood

$$\mathcal{L}(\langle \sigma v \rangle, N_B, J) = \prod_{i=1}^{i} \mathcal{L}_{P_i}(\langle \sigma v \rangle, N_{B_i}, J | N_{\text{ON}_i}, N_{\text{OFF}_i}, \alpha) \mathcal{L}^J(J | \bar{J}, \sigma_J)$$

Poisson likelihood

Log-normal likelihood

Poisson likelihood for each energy bin

$$\mathcal{Z}_{i}^{P} = \frac{(N_{S_{i}} + N_{B_{i}})^{N_{\text{ON}_{i}}}}{N_{\text{ON}_{i}}!} e^{-(N_{S_{i}} + N_{B_{i}})} \cdot \frac{(\alpha N_{B_{i}})^{N_{\text{OFF}_{i}}}}{N_{\text{OFF}_{i}}!} e^{-\alpha N_{B_{i}}}$$

Log-normal likelihood to model the uncertainties of the J factor

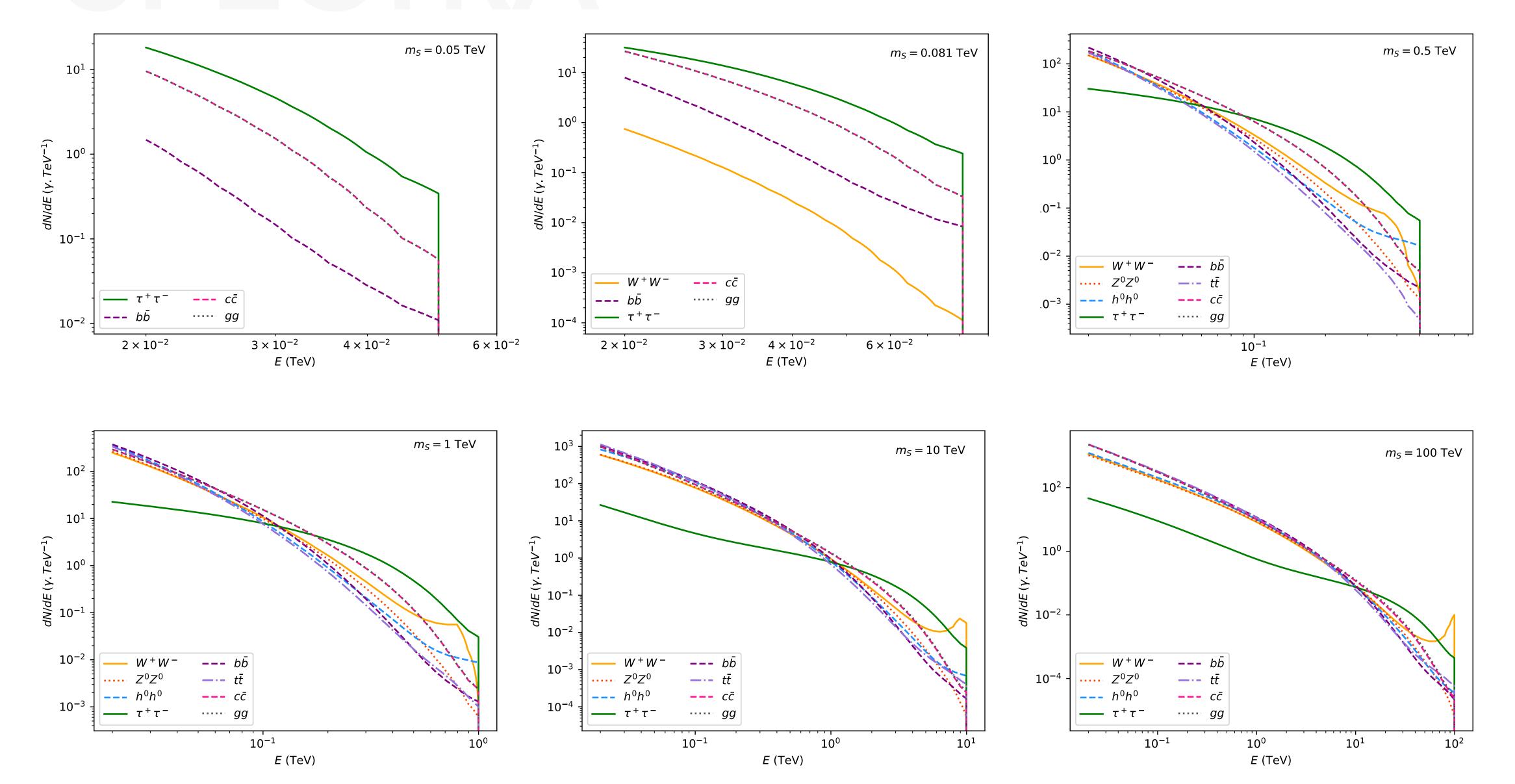
$$\mathcal{L}^{J} = \frac{1}{\ln(10)\sqrt{2\pi\sigma_{J}J}} \exp{-\frac{(\log_{10}J - \log_{10}\bar{J})^{2}}{2\sigma_{J}^{2}}}$$

ON REGION

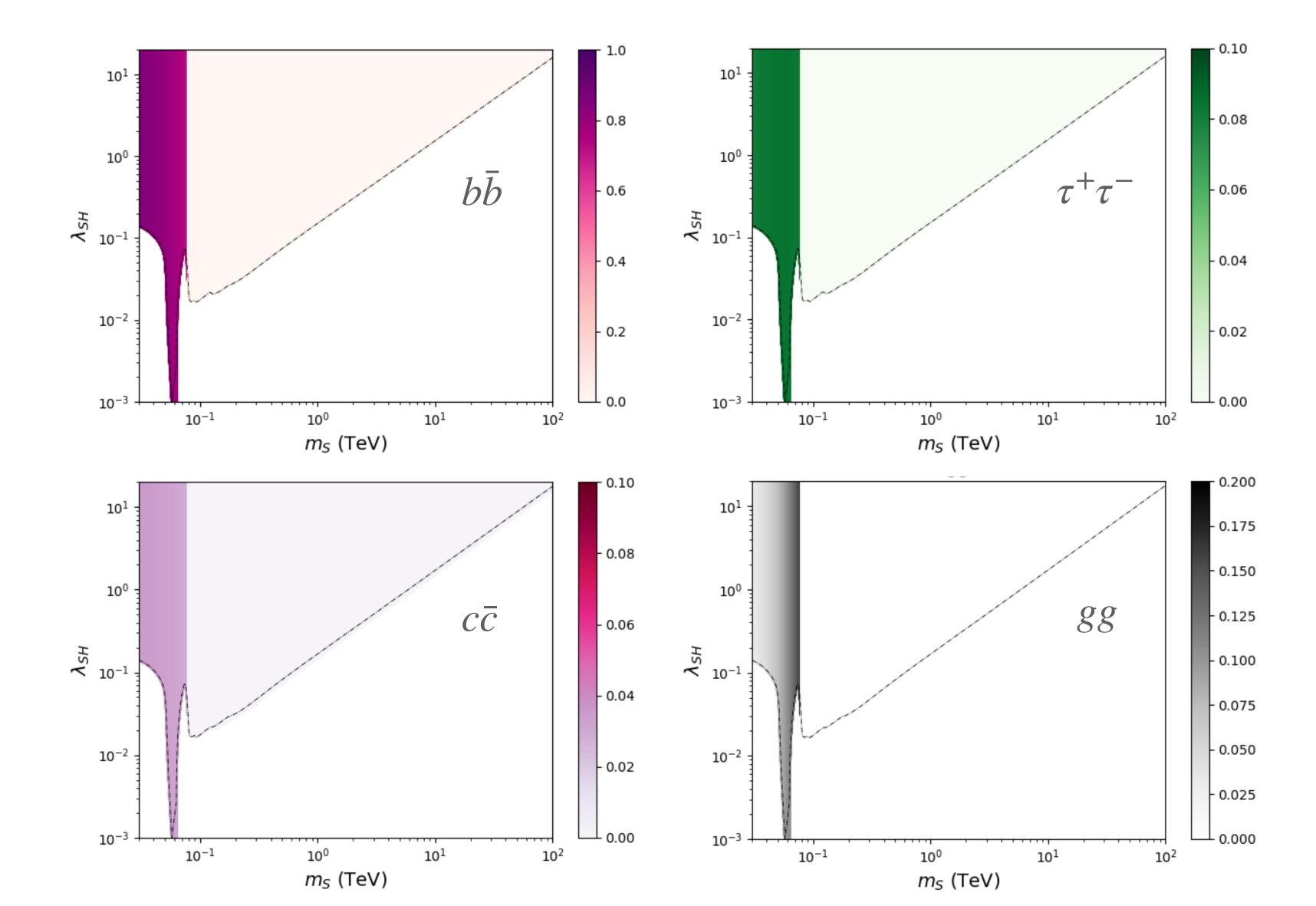
OFF REGION

SPECTRA

Ref: Cirelli et al. *JCAP* 03 (2011) 051

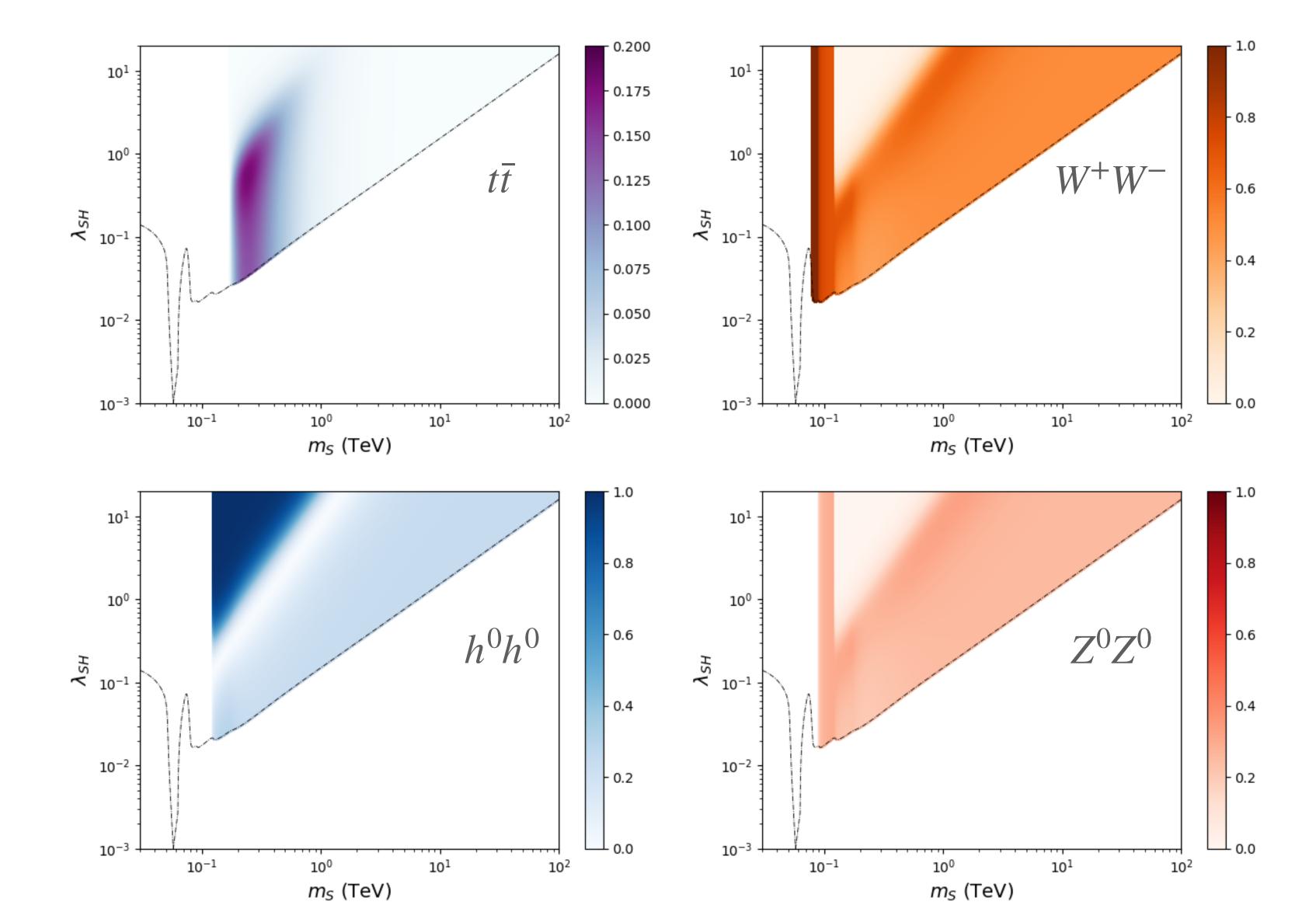


BRANCHING RATIOS



 $m_S \lesssim m_W$

BRANCHING RATIOS



$$m_S \gtrsim m_W$$