

# Theoretical Overview of Double Beta Decay

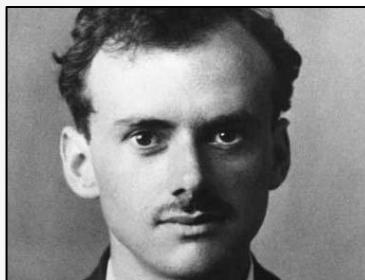
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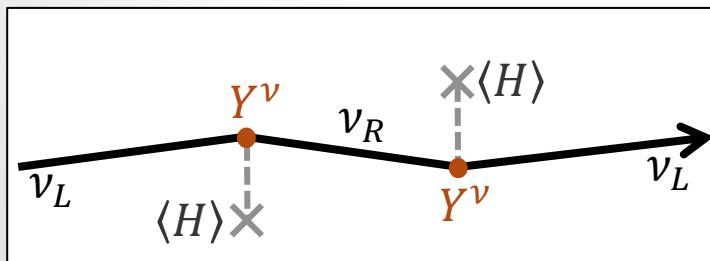
University College London

# Dirac versus Majorana

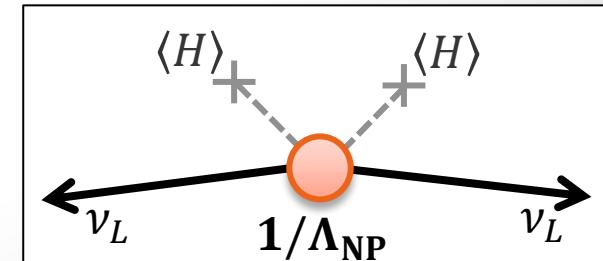
- ▶ Origin of neutrino masses beyond the Standard Model
- ▶ Two possibilities to define neutrino mass



Dirac mass analogous to other fermions but with  $m_\nu / \Lambda_{EW} \approx 10^{-12}$  couplings to Higgs



Majorana mass, using only a left-handed neutrino  
 → Lepton Number Violation



# Dirac versus Majorana

- ▶ Origin of neutrino masses beyond the Standard Model
- ▶ Crucial role of total lepton number  $L$  symmetry
  - Arises accidentally as global  $U(1)_L$  in SM from particle content and gauge symmetry
  - $L$  broken non-perturbatively but  $B - L$  conserved
  - Global symmetries expected to be broken gravitational effects?

$$m_\nu \approx \frac{v^2}{M_{\text{Planck}}} \approx 10^{-5} \text{ eV}$$

- Too small to explain oscillations but too large as subdominant splitting
- Connection to matter–antimatter asymmetry
  - Leptogenesis through heavy neutrinos in Type-I Seesaw

# Beta Decays and Neutrinos

## ► Single beta decay

$$(A, Z) \rightarrow (A, Z + 1) + e^- + \bar{\nu}_e$$

- Kinematic neutrino mass measurement

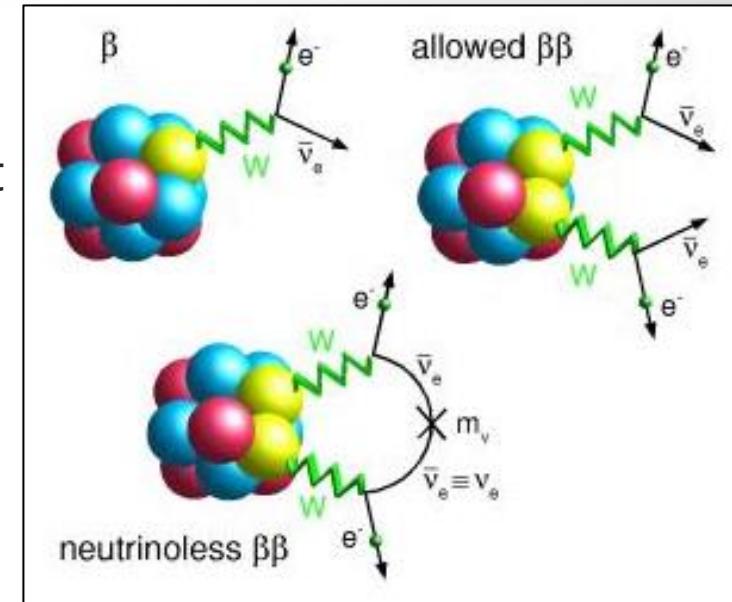
## ► Allowed double beta ( $2\nu\beta\beta$ ) decay

$$(A, Z) \rightarrow (A, Z + 2) + 2e^- + 2\bar{\nu}_e$$

## ► Neutrinoless double beta ( $0\nu\beta\beta$ ) decay

$$(A, Z) \rightarrow (A, Z + 2) + 2e^-$$

- Violation of lepton number
- Mediated by Majorana neutrinos
- Alternatives:
  - $0\nu\beta^+\beta^+$ :  $(A, Z) \rightarrow (A, Z - 2) + 2e^+$
  - $0\nu\beta^+EC$ :  $(A, Z) + e^- \rightarrow (A, Z - 2) + e^+$
  - $0\nuECEC$ :  $(A, Z) + 2e^- \rightarrow (A, Z - 2)$

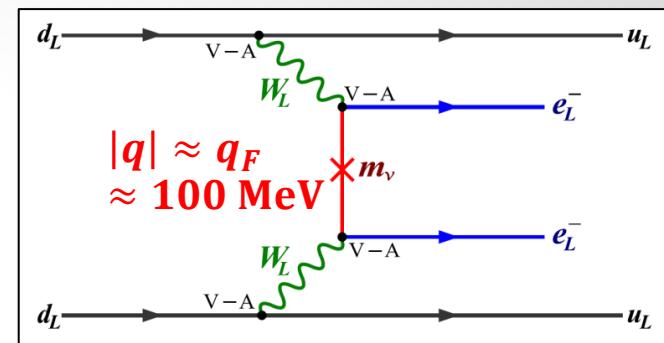


# Neutrinoless Double $\beta$ Decay

## ► Half-life

$$T_{1/2}^{-1} = |\mathbf{m}_{\beta\beta}|^2 \mathbf{G}^{0\nu} |\mathbf{M}^{0\nu}|^2$$

## ► Particle Physics



$$\mathcal{A}_{\mu\nu}^{lep} = \frac{1}{4} \sum_{i=1}^3 U_{ei}^2 \gamma_\mu (1 + \gamma_5) \frac{\cancel{q} + m_{\nu_i}}{q^2 - m_{\nu_i}^2} \gamma_\nu (1 - \gamma_5) \approx \frac{\gamma_\mu (1 + \gamma_5) \gamma_\nu}{4q^2} \sum_{i=1}^3 U_{ei}^2 m_{\nu_i} \rightarrow \mathbf{m}_{\beta\beta}$$

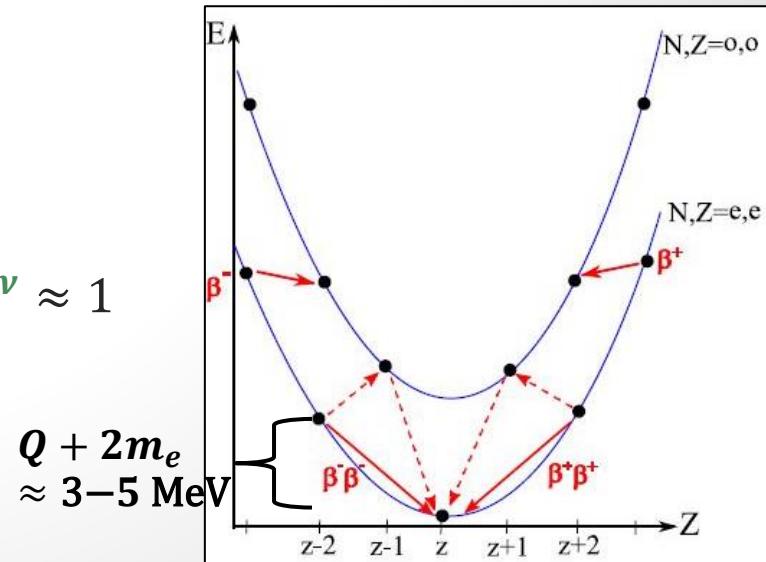
## ► Atomic Physics

- Leptonic phase space  $\mathbf{G}^{0\nu} \propto Q^5$

## ► Nuclear Physics

- Nuclear transition matrix element  $\mathbf{M}^{0\nu} \approx 1$   
but large uncertainties, factor 2–3

$$\frac{10^{25} \text{ y}}{T_{1/2}} \approx \left( \frac{|m_{\beta\beta}|}{\text{eV}} \right)^2$$

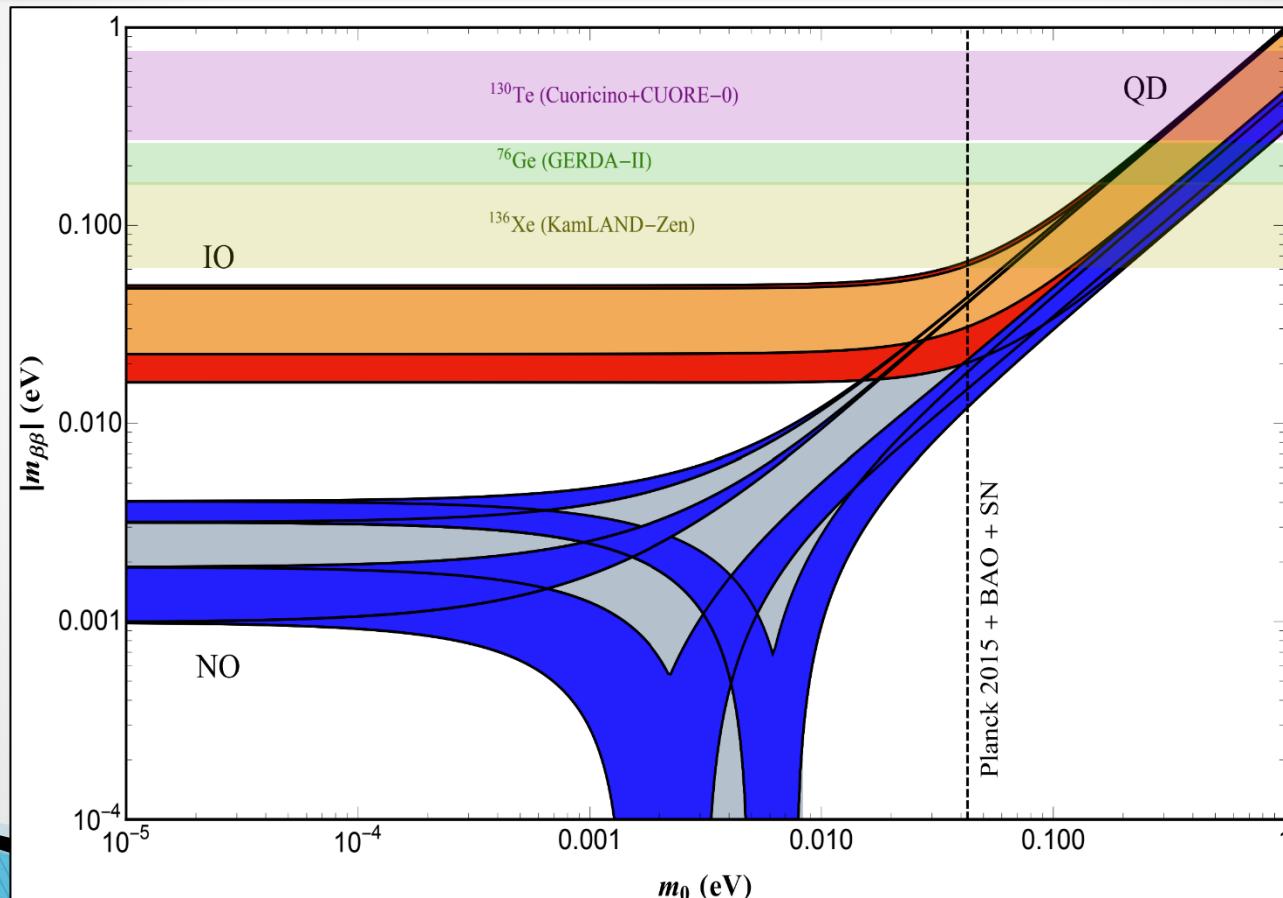


# Three Active Neutrinos

## ► Effective $0\nu\beta\beta$ Mass

degenerate &  $\theta_{13} \approx 0$

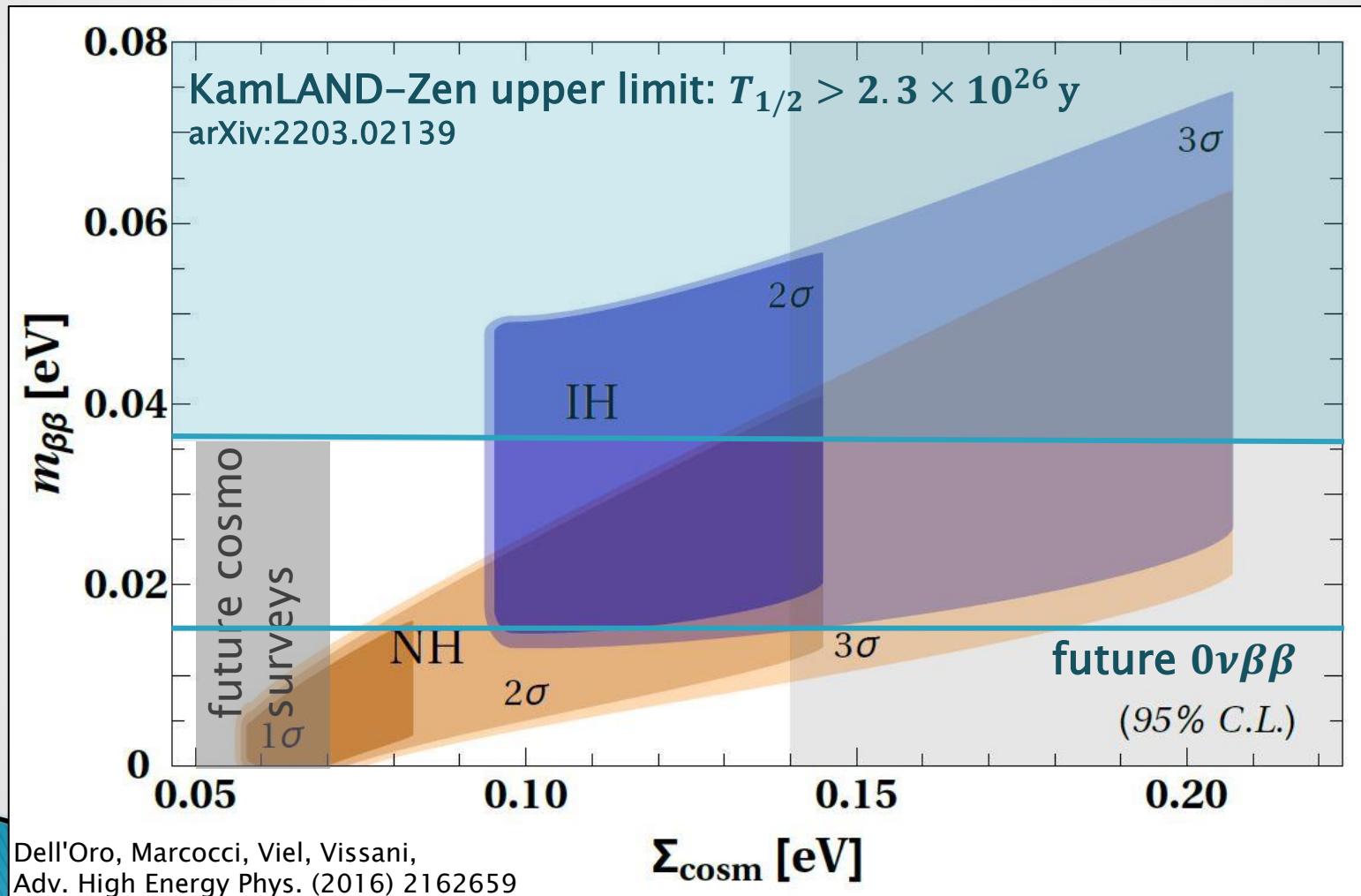
$$|m_{\beta\beta}| = |c_{12}^2 c_{13}^2 m_{\nu_1} + s_{12}^2 c_{13}^2 m_{\nu_2} e^{i\phi_{12}} + s_{13}^2 m_{\nu_3} e^{i\phi_{13}}| \approx m_\nu \sqrt{1 - \sin^2(2\theta_{12}) \sin^2(\phi_{12}/2)}$$



Hagedorn et al.,  
 Int.J.Mod.Phys.A 33 (2018) 05n06, 1842006

# Three Active Neutrinos

- ▶ Effective  $0\nu\beta\beta$  Mass



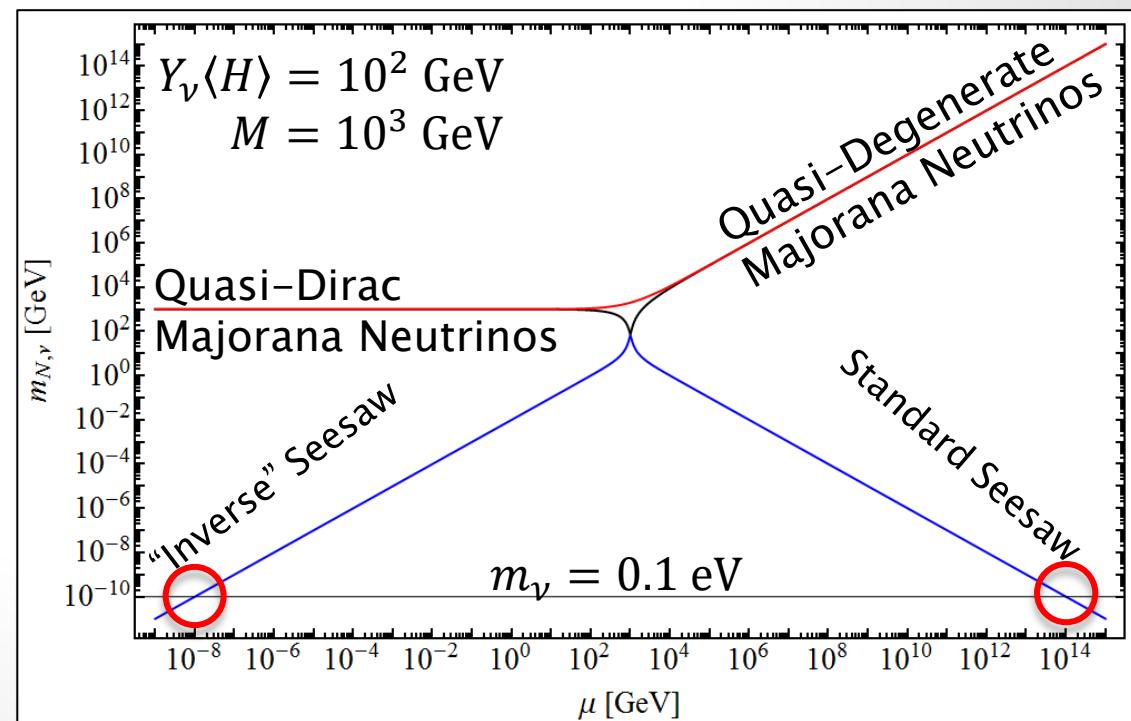
# Heavy Sterile Neutrinos

- ▶ Correct light neutrino masses for TeV scale heavy neutrinos
  - Seesaw Mechanism with TeV scale heavy neutrinos
    - Standard Seesaw with small Yukawa couplings
    - CLFV remains small
  - “Bent” Seesaw mechanisms
    - Decouple  $\Lambda_{\text{LNV}}$  from heavy neutrino mass
    - Example

$$\mathcal{M} = \begin{pmatrix} 0 & Y_\nu \langle H \rangle & 0 \\ Y_\nu \langle H \rangle & \mu & M \\ 0 & M & \mu \end{pmatrix}$$

- Potentially large CLFV
- In the limit  $\mu \rightarrow 0$ , no LNV but CLFV

$$V^{LR} \approx Y_\nu \approx 10^{-6} \sqrt{M_N/\text{TeV}}$$



# Sterile Neutrinos

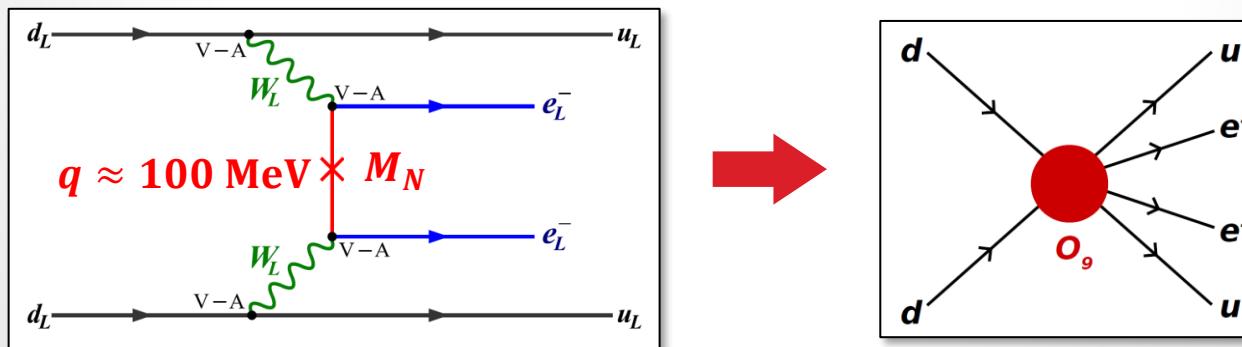
- ▶ Masses lighter than  $\approx 100$  MeV

$$|m_{\beta\beta}| = |c_{12}^2 c_{13}^2 m_{\nu_1} + s_{12}^2 c_{13}^2 m_{\nu_2} e^{i\phi_{12}} + s_{13}^2 m_{\nu_3} e^{i\phi_{13}} + s_{14}^2 m_{\nu_4} e^{i\phi_{14}} + \dots|$$

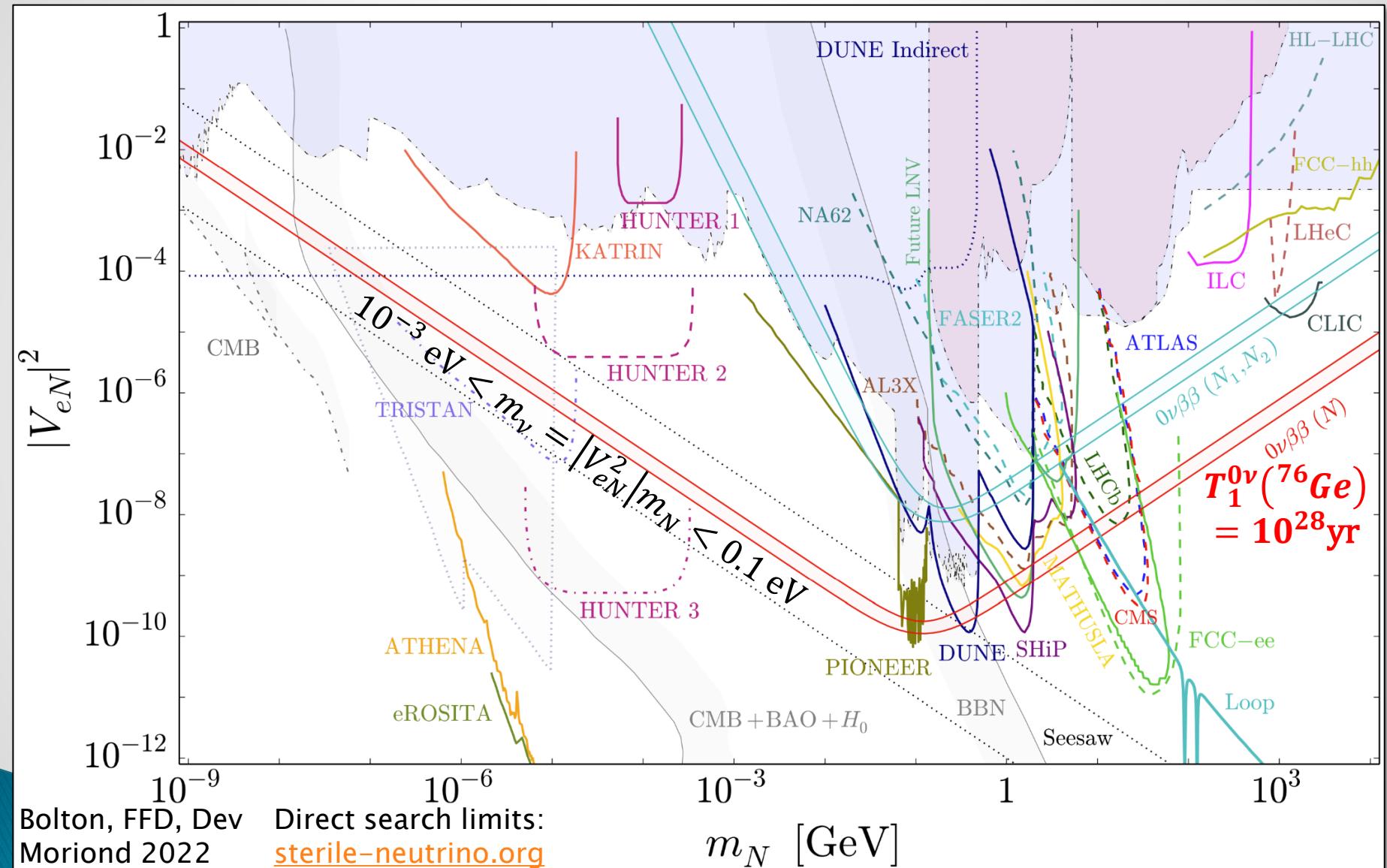
- ▶ Masses heavier than  $\approx 100$  MeV

$$\mathcal{A}_{\mu\nu}^{lep} = \frac{1}{4} \sum_{i=1}^3 V_{ei}^2 \gamma_\mu (1 + \gamma_5) \frac{\cancel{q} + M_{N_i}}{q^2 - M_{N_i}^2} \gamma_\nu (1 - \gamma_5) \approx \frac{-\gamma_\mu (1 + \gamma_5) \gamma_\nu}{4} \sum_{i=1}^3 \frac{V_{ei}^2}{M_{N_i}} \rightarrow \left\langle \frac{1}{M_N} \right\rangle_{\beta\beta}$$

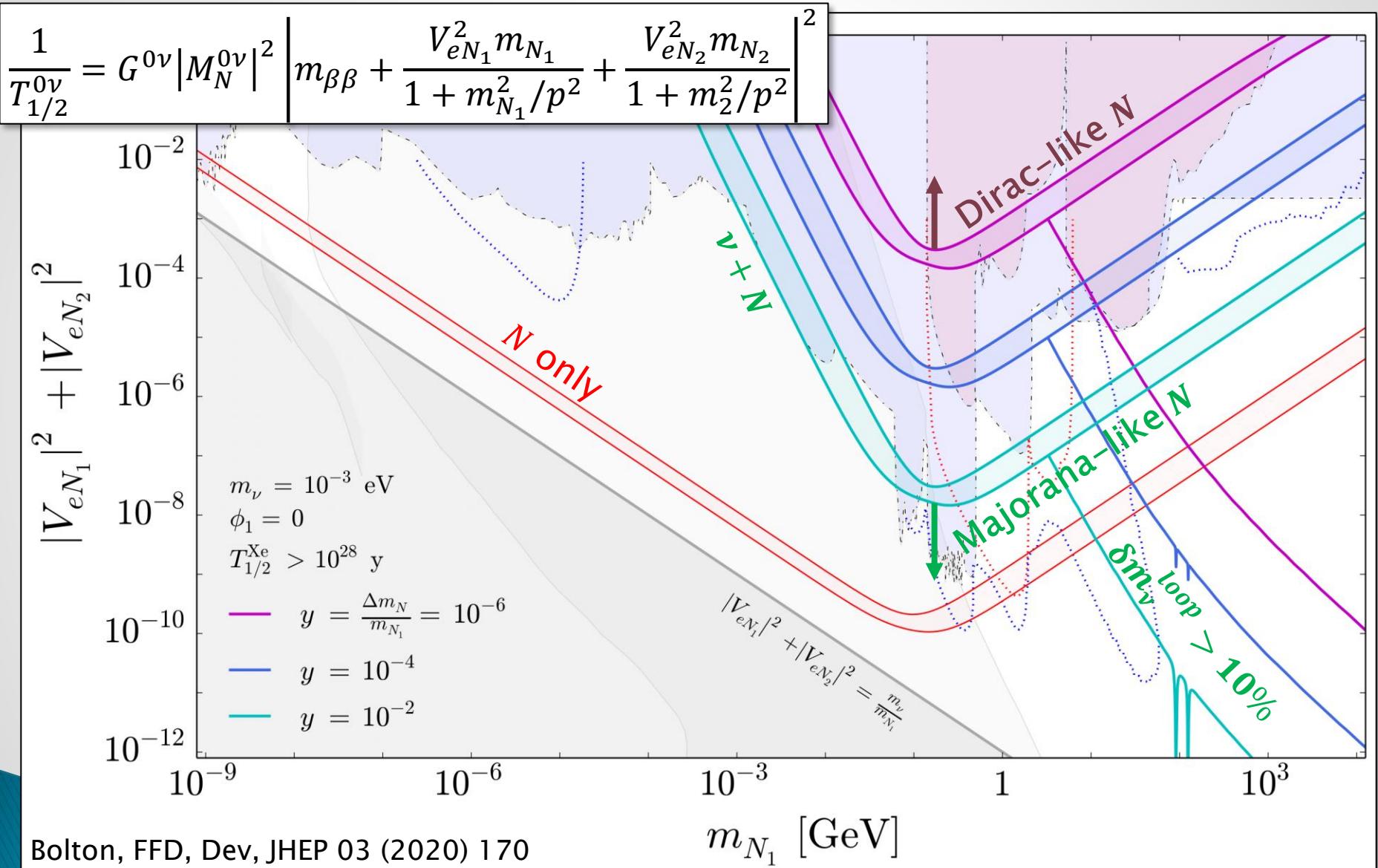
- ▶ Short-distance on nuclear scale



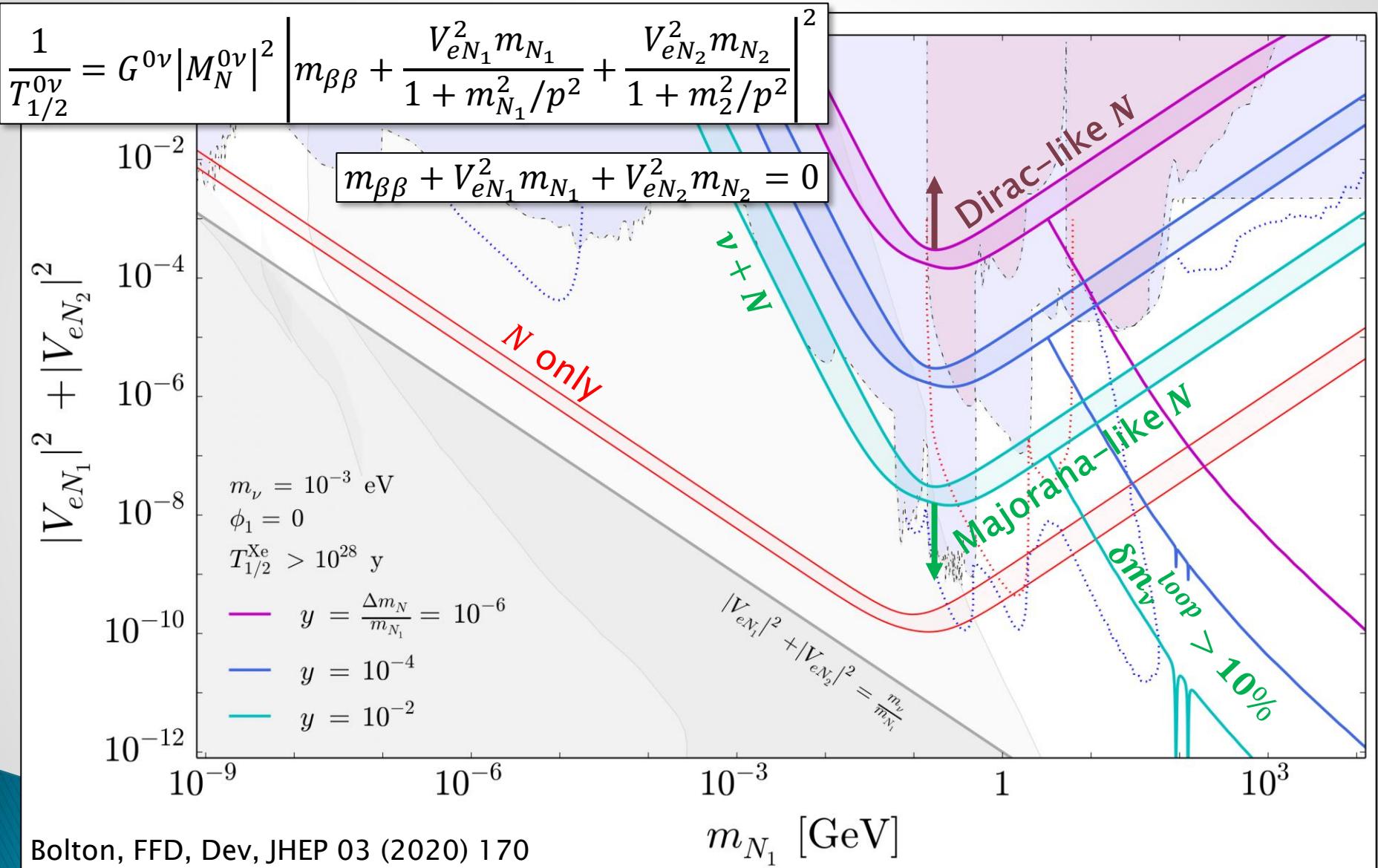
# HNL - Future Sensitivities



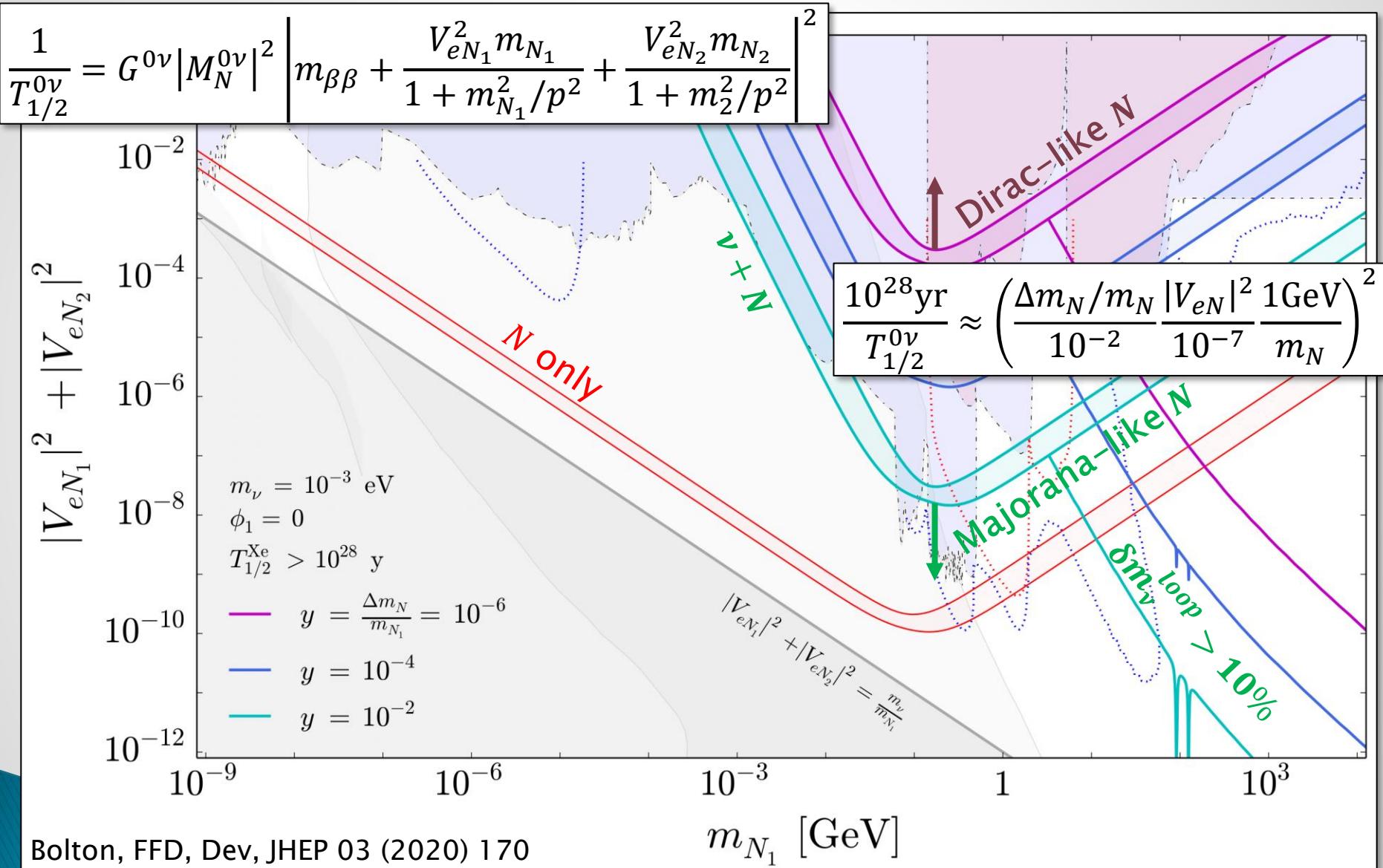
# HNL - Comparison with $0\nu\beta\beta$



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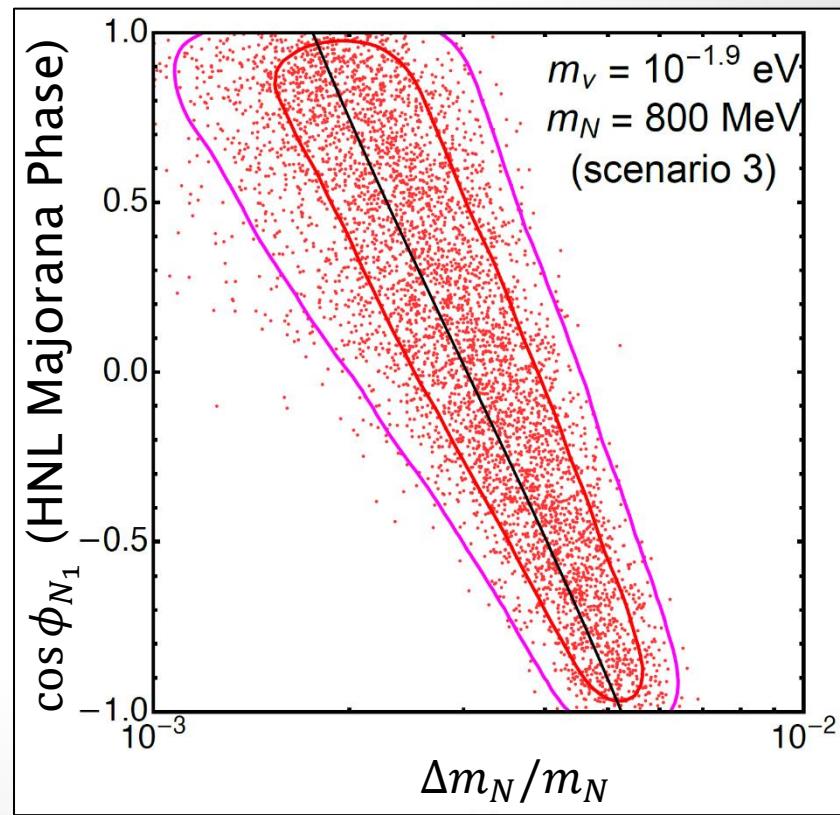
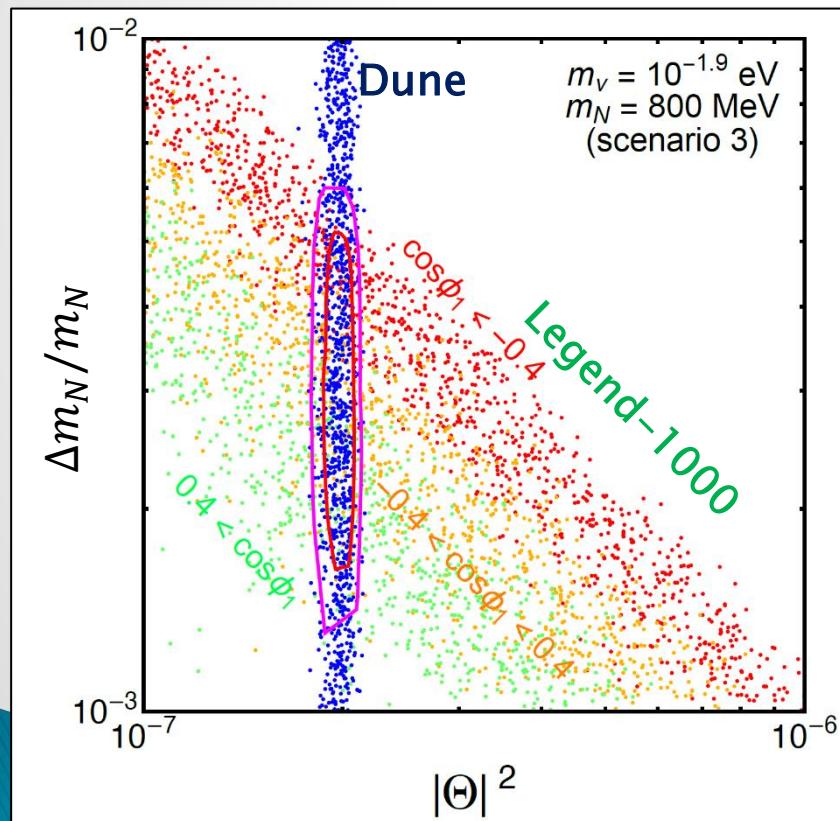
# HNL - Comparison with $0\nu\beta\beta$



# Complementarity

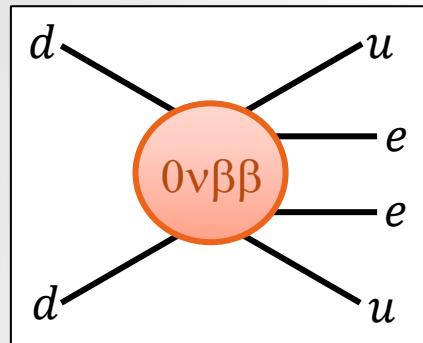
Patrick Bolton, FFD, Mudit Rai, Zhong Zhang, 2212.14690

- ▶ Between direct searches and  $0\nu\beta\beta$ 
  - Simulation of DUNE-like setup
  - Measurement of events at “DUNE” and  $0\nu\beta\beta$  decay at LEGEND-1000 near expected sensitivity

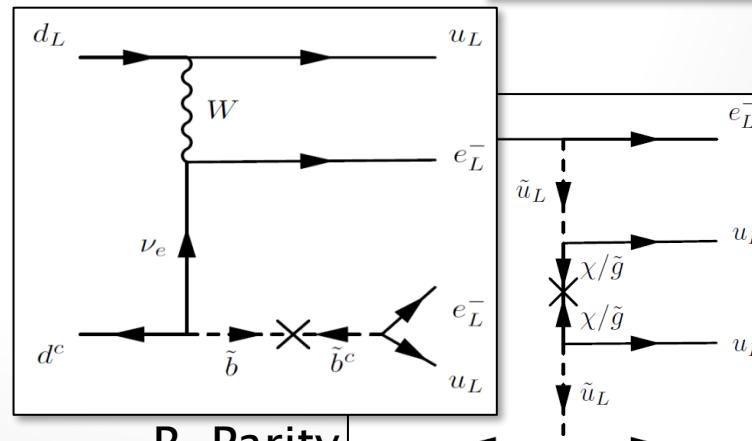
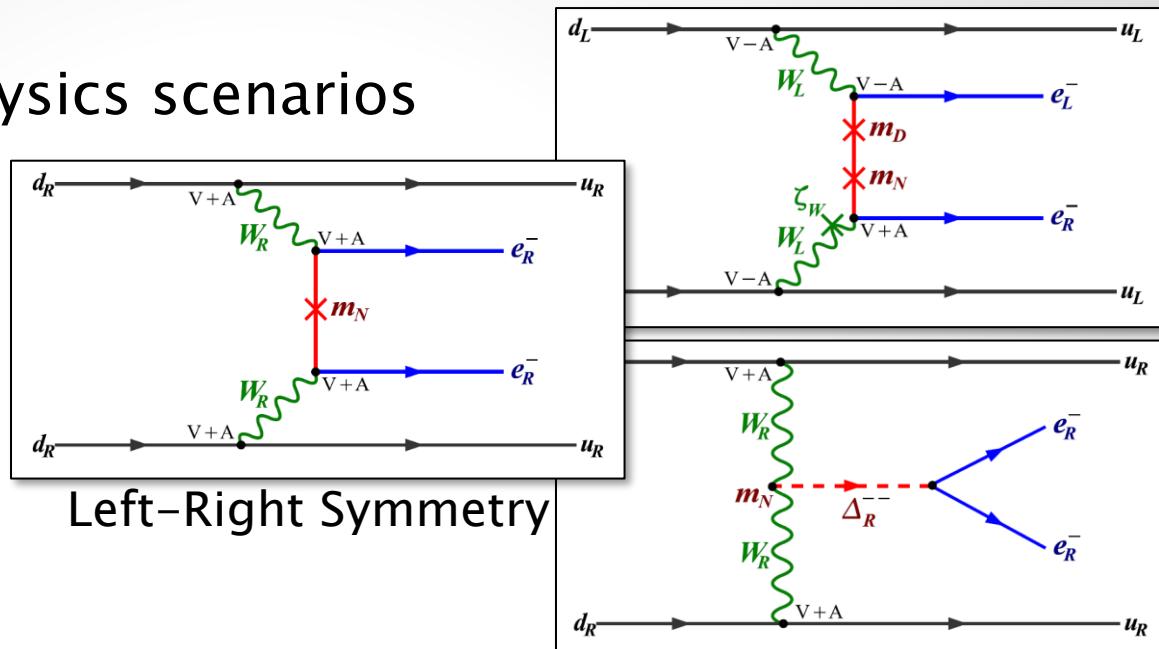


# New Physics and $0\nu\beta\beta$

## ► Plethora of New Physics scenarios



$$T_{1/2}^{-1} = \epsilon_{NP}^2 G_{NP}^{0\nu} |M_{NP}^{0\nu}|^2$$



R-Parity  
Violating SUSY

Extra Dimensions

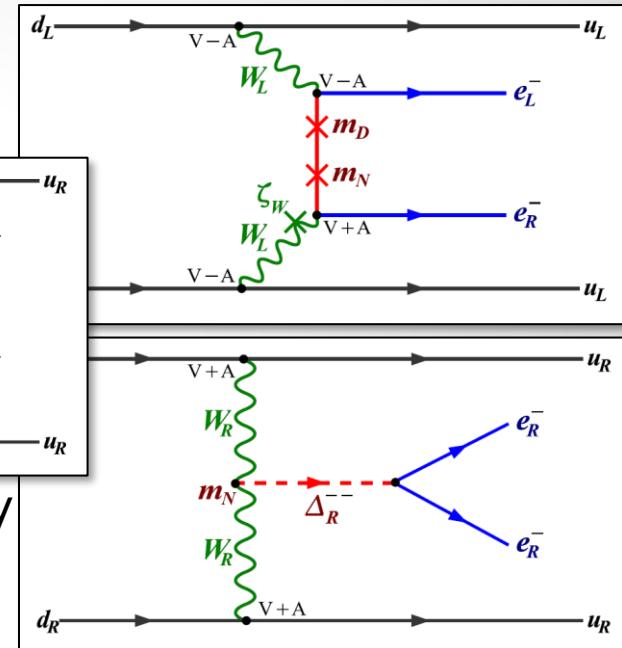
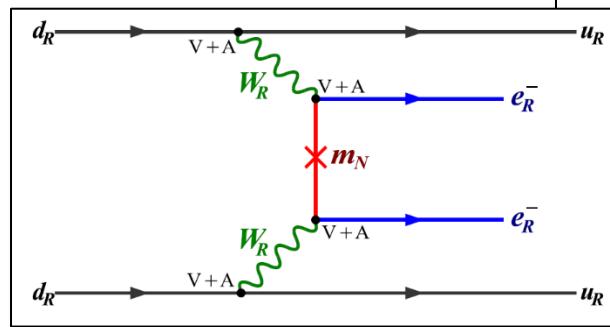
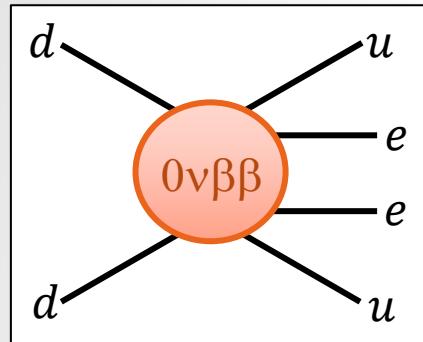
Majorons

Leptoquarks

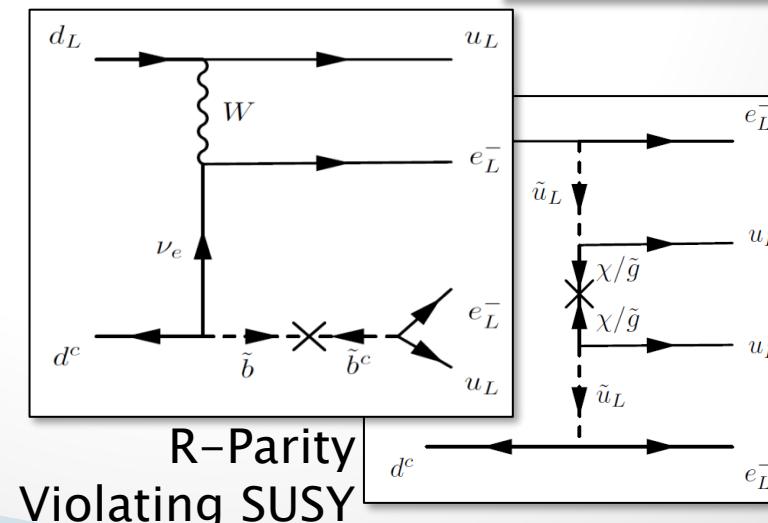
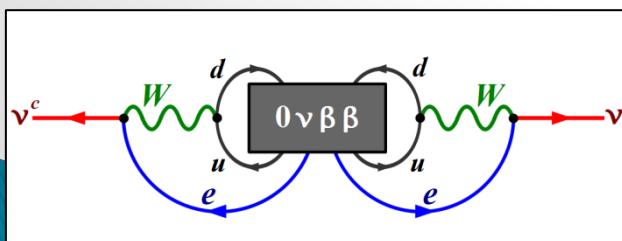
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# New Physics and $0\nu\beta\beta$

## ► Plethora of New Physics scenarios



► Neutrinos still Majorana  
 Schechter, Valle  
 Phys. Rev. D25 (1982) 2951



Extra Dimensions

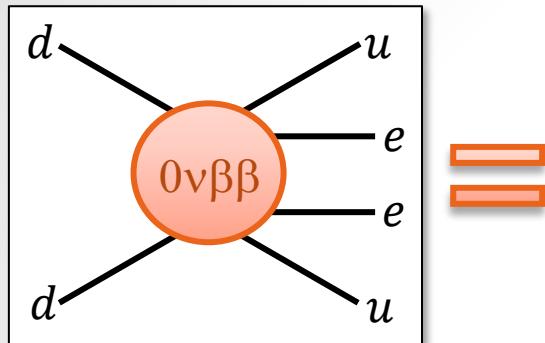
Majorons

Leptoquarks

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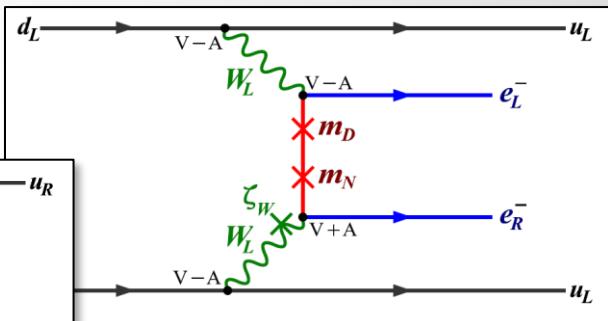
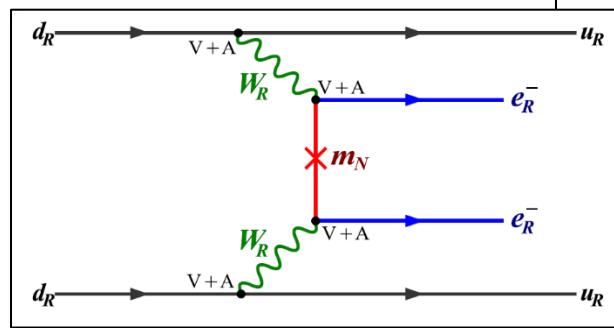
# New Physics and $0\nu\beta\beta$

## ► Examples in Left-Right Symmetry



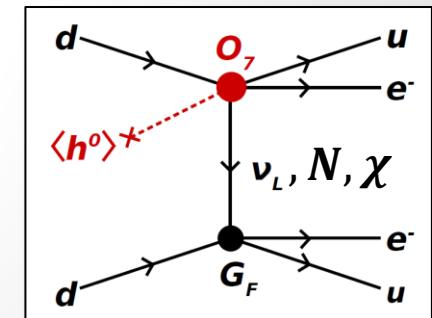
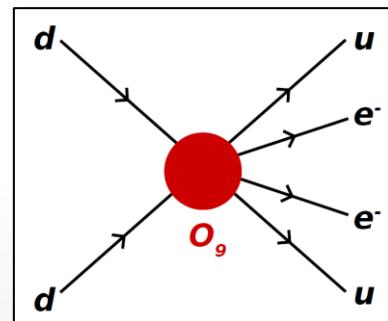
$$T_{1/2}^{-1} = \epsilon_{NP}^2 G_{NP}^{0\nu} |M_{NP}^{0\nu}|^2$$

►  $0\nu\beta\beta$  probes LNV at the TeV scale and above



$$\begin{aligned} \epsilon_{V-A}^{V+A} &= \sum_{i=1}^3 U_{ei} W_{ei} \tan \zeta_W \\ &\approx \frac{10^{-9}}{(\Lambda/10 \text{ TeV})^3} \end{aligned}$$

$$\begin{aligned} \epsilon_3^{RRZ} &= \sum_{i=1}^3 V_{ei}^2 \frac{m_p}{m_N} \frac{m_W^4}{m_{W_R}^4} \\ &\approx \frac{10^{-8}}{(\Lambda/1 \text{ TeV})^5} \end{aligned}$$

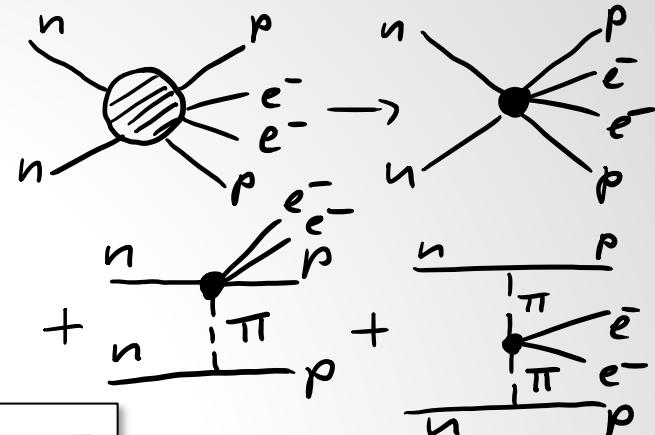
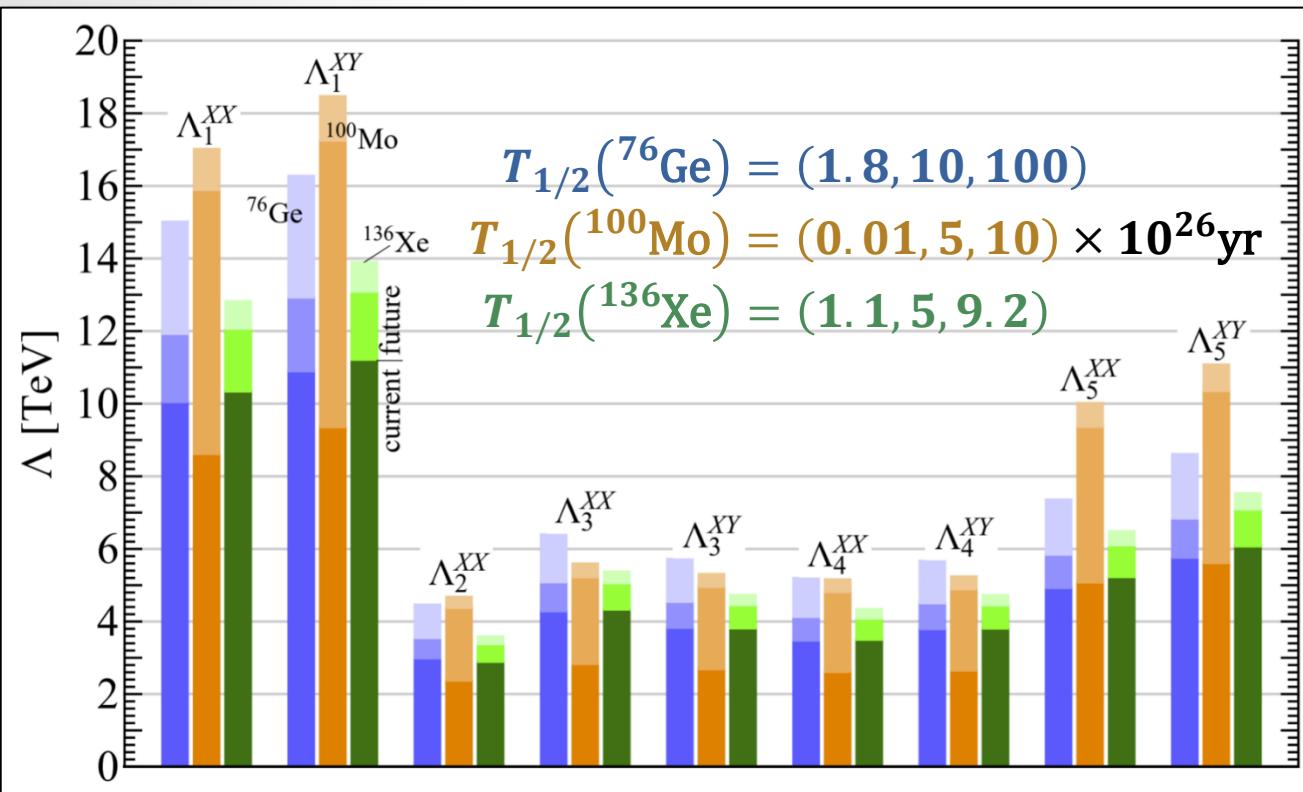


# Heavy New Physics

FFD, Graf, Iachello, Kotila, PRD 102 (2020)

- ▶ Limits on short-range operators

- NMEs from IBM-2 with  $g_A = 1.0$  and short-range correlations in Argonne parametrization



## Pion-mediated contributions

- ▶ R-parity violating SUSY (Faessler, Kovalenko, Simkovic, Schwieger, Phys.Rev.Lett. 78 (1997) 183)
- ▶ Chiral EFT with Pion operators from Lattice QCD (Cirigliano, Dekens, de Vries, Graesser, Mereghetti, JHEP 1812 (2018) 097)

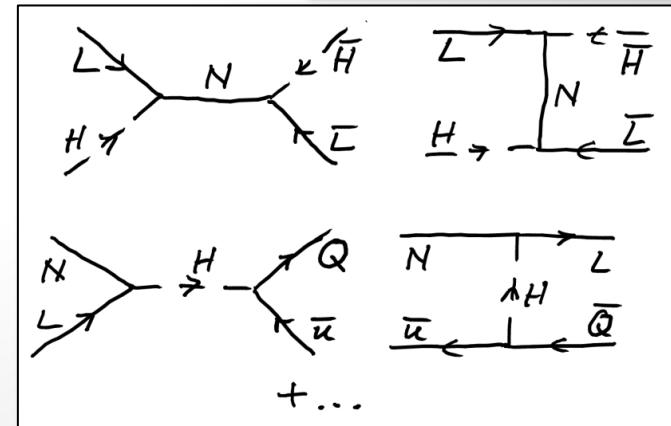
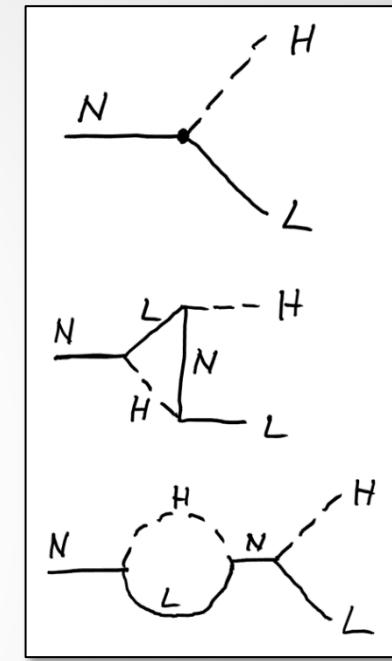
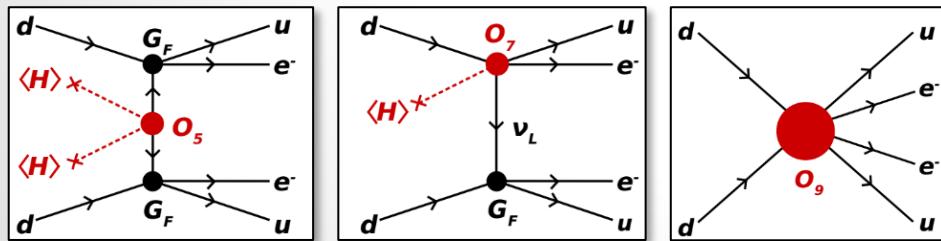
# Falsifying Baryogenesis

## ► Classic Example: High-Scale Leptogenesis

- Generation via heavy neutrino decays
- Competition with LNV washout processes
- Conversion to baryon asymmetry
  - EW sphaleron processes at  $T \approx 100$  GeV
  - Observed asymmetry

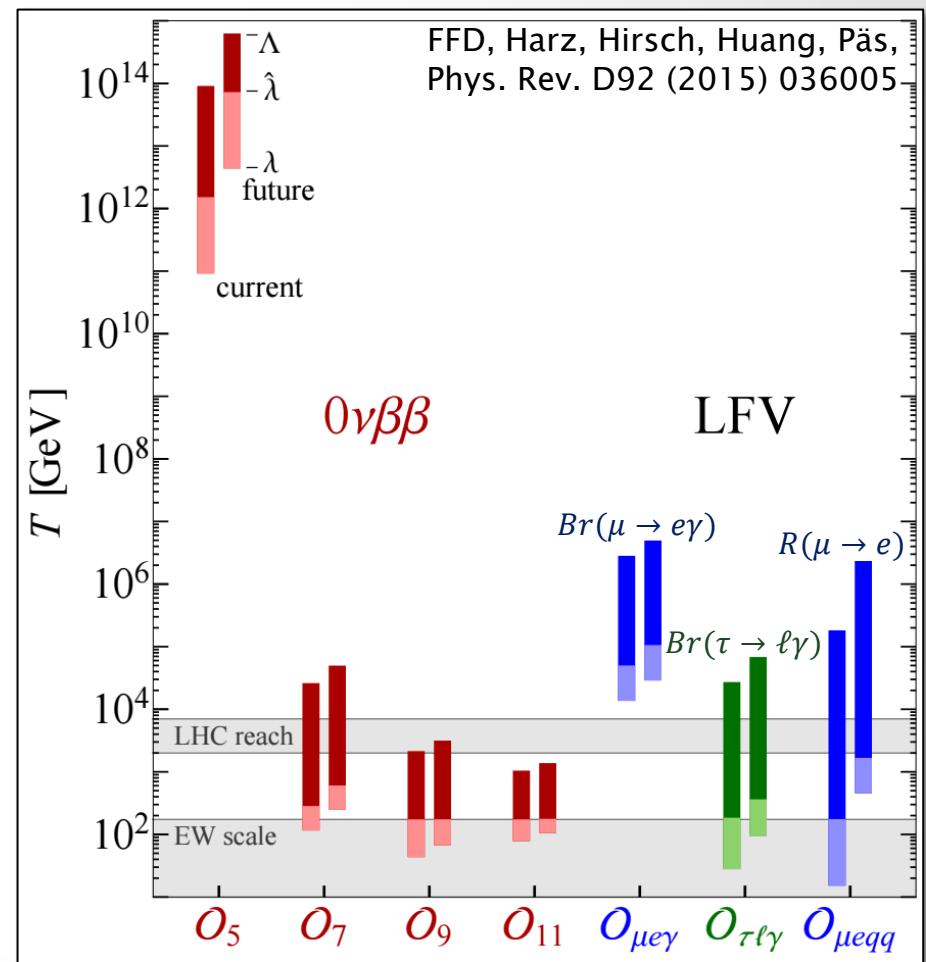
$$\eta_B \equiv \frac{n_B - n_{\bar{B}}}{n_\gamma} = (6.20 \pm 0.15) \times 10^{-10}$$

## ► What if we observe lepton number violating processes in $0\nu\beta\beta$ ?



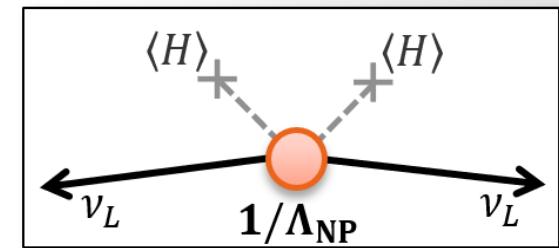
# Falsifying Baryogenesis

- ▶ Temperature ranges of strong equilibration
  - Assumes observation of corresponding process!
- ▶ Observation of LNV
  - gives information at what temperatures operators are in equilibrium
  - **can falsify high-scale baryogenesis scenarios**



# Conclusion

- ▶ **Neutrinos much lighter than other fermions**
  - Dirac or Majorana? Lepton Number Violation?
  - Determination of absolute mass scale
- ▶  **$0\nu\beta\beta$  crucial probe for BSM physics**
  - Universal probe of LNV physics
    - LNV physics near GUT scale
    - Direct sensitivity to LNV physics at scales  $m_N \approx 1 \text{ eV} - 100 \text{ TeV}$
    - Light exotic particles
  - Sensitive to CP properties of light and GeV-scale  $\nu$
- ▶  **$2\nu\beta\beta$  sensitive to New Physics**
  - Ongoing and future searches probe  $2\nu\beta\beta$  decay with high statistics
  - E.g., exotic (right-handed) currents,  $\nu$  self-interactions, sterile  $\nu$  mass endpoint



$$\frac{T_{1/2}^{0\nu\beta\beta}}{10^{28} \text{ y}} \approx \left( \frac{\Lambda_{\text{NP}}}{10^{15} \text{ GeV}} \right)^2$$