

# Recent results on neutrino oscillations from NOvA

Pierre Lasorak  
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- Neutrinos change flavour as they travel
  - Nobel Prize 2015 for Kajita-san (SK) and Art McDonald (SNO)
- Oscillations depend on
  - Travel distance and energy of neutrinos
  - 3 angles of “rotation”:  $\theta_{12}$ ,  $\theta_{13}$  and  $\theta_{23}$
  - 2 differences of mass squared:  $\Delta m_{32}^2$  and  $\Delta m_{12}^2$
  - 1 phase:  $\delta_{CP}$
  - Matter density (through the MSW effect)

- 3 neutrinos leading order disappearance probability:

$$P(\nu_\mu \rightarrow \nu_\mu) = 1 - \sin^2 2\theta_{23} \sin^2 \frac{1.27 \Delta m_{32}^2 L}{E}$$

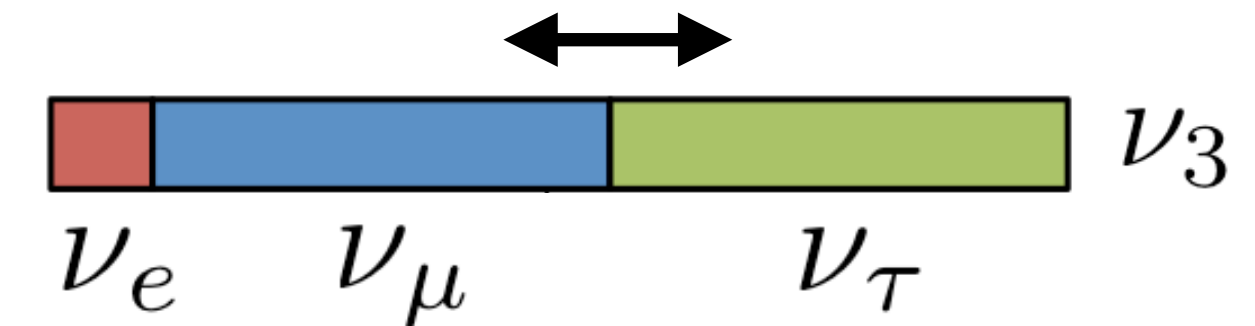
$$\Delta m_{32}^2 = m_2^2 - m_3^2$$

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \underbrace{\begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix}}_{\text{atmospheric, beam}} \underbrace{\begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta_{CP}} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta_{CP}} & 0 & c_{13} \end{pmatrix}}_{\text{reactor, beam}} \underbrace{\begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}}_{\text{solar, reactor}} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

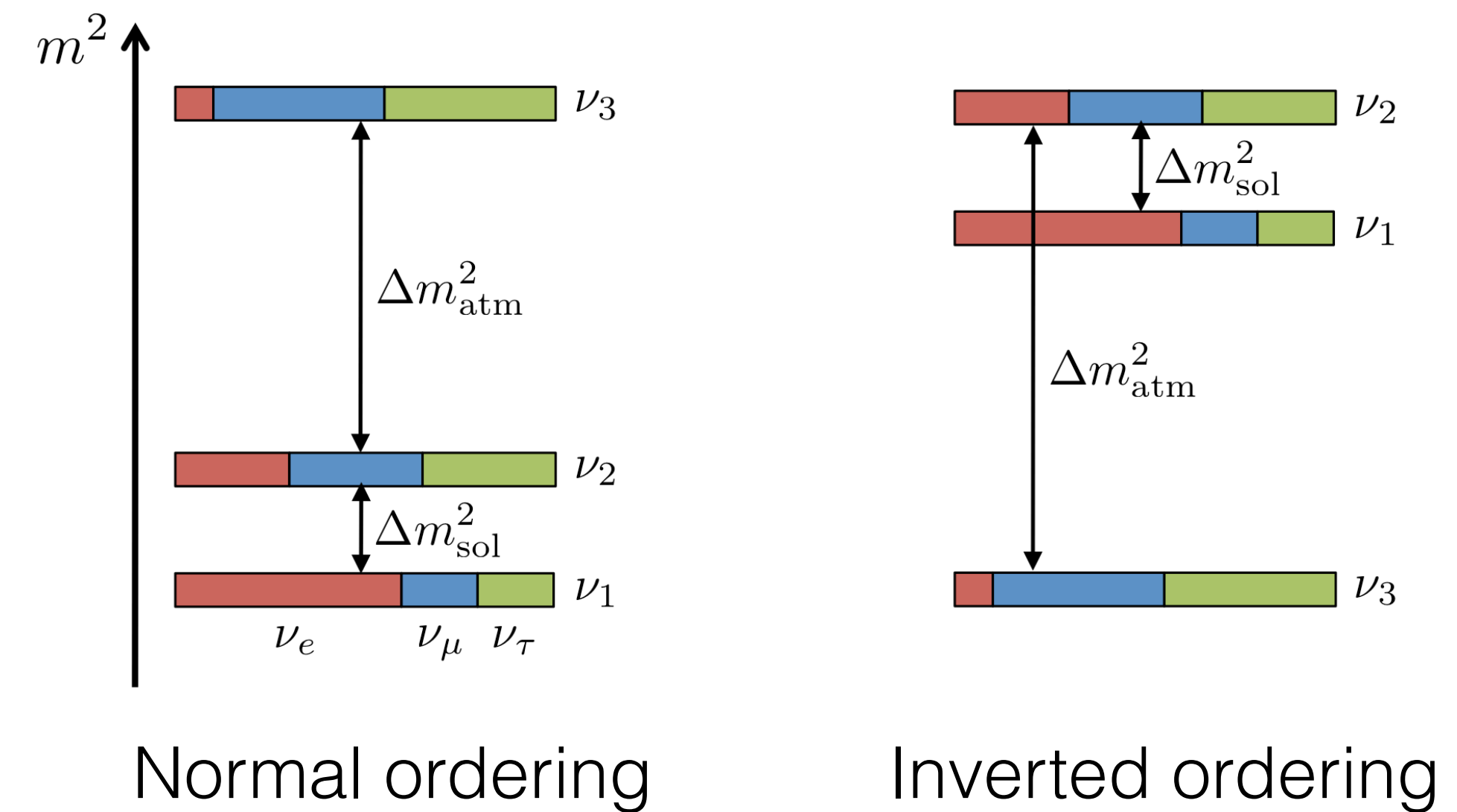
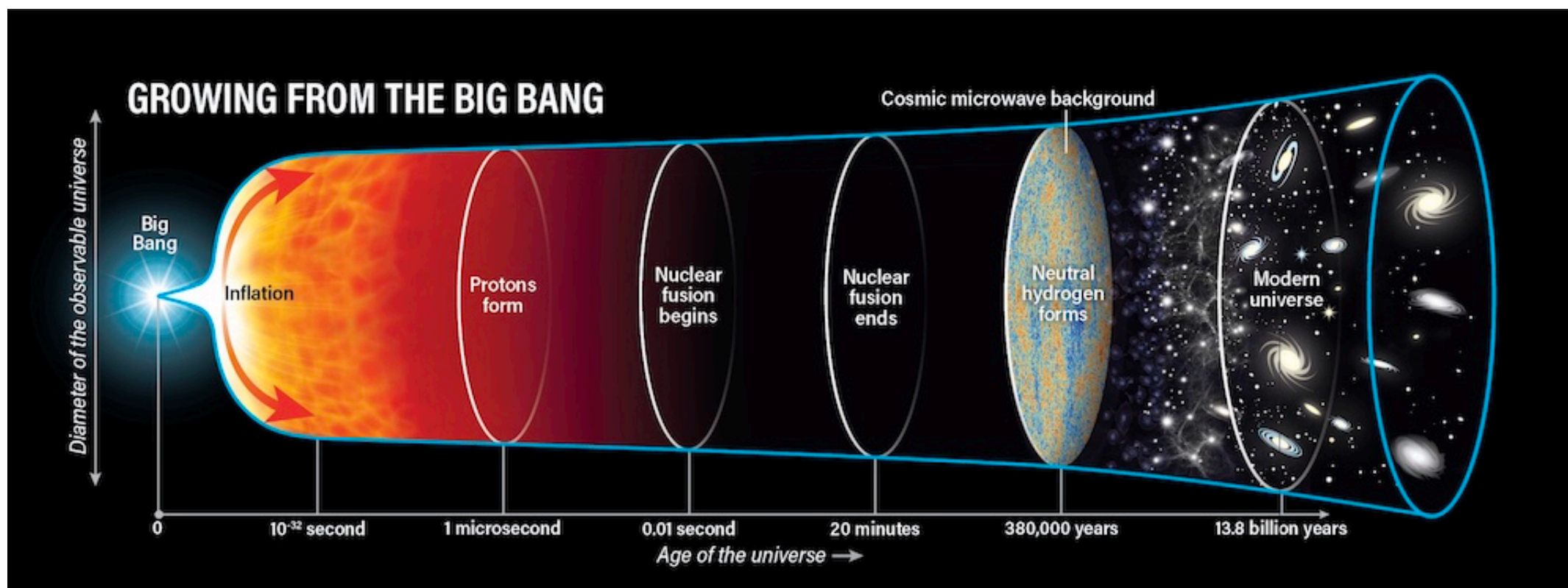
$$s_{ij} = \sin\theta_{ij}$$

$$c_{ij} = \cos\theta_{ij}$$

- Unknowns left to discover/measure
  - $\theta_{23}$  octant / maximal mixing: is  $\sin^2\theta_{23} = 0.5$ ?
  - Mass ordering: sign of  $\Delta m^2_{32}$ ?
  - Charge parity violation: is  $\delta_{CP} \neq 0$ ?



$$P(\nu_\mu \rightarrow \nu_\mu) \simeq 1 - \sin^2 2\theta_{23} \sin^2 \frac{1.27 \Delta m^2_{32} L}{E}$$



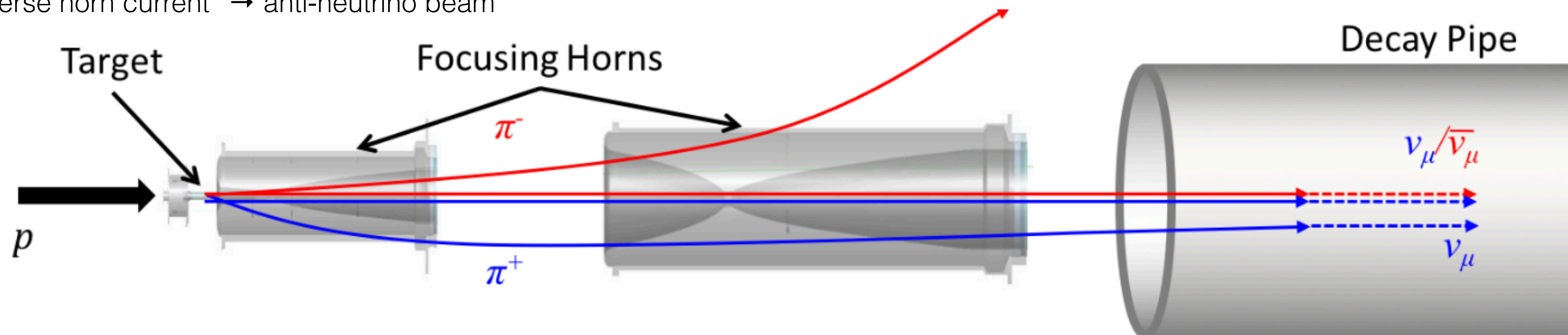
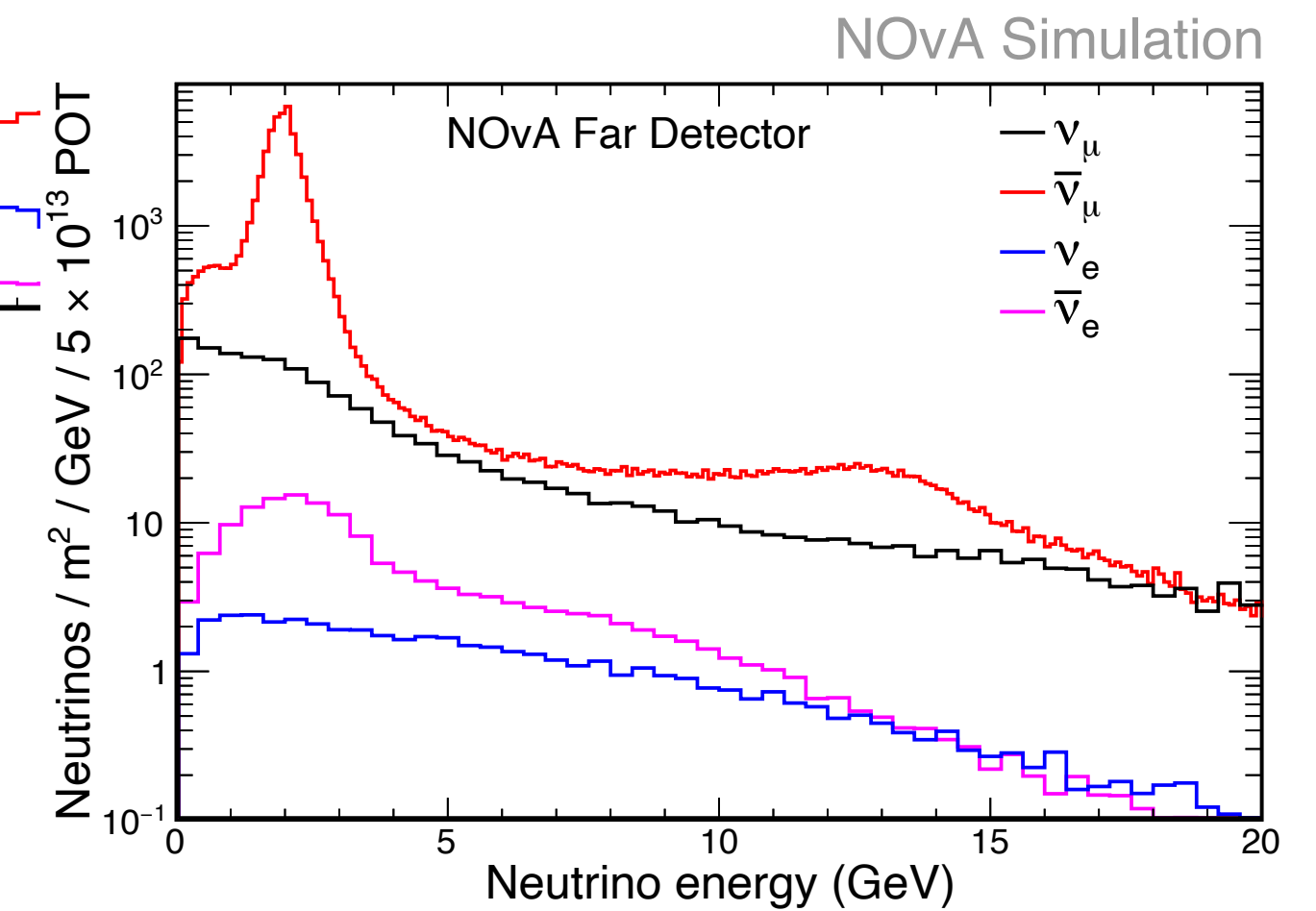
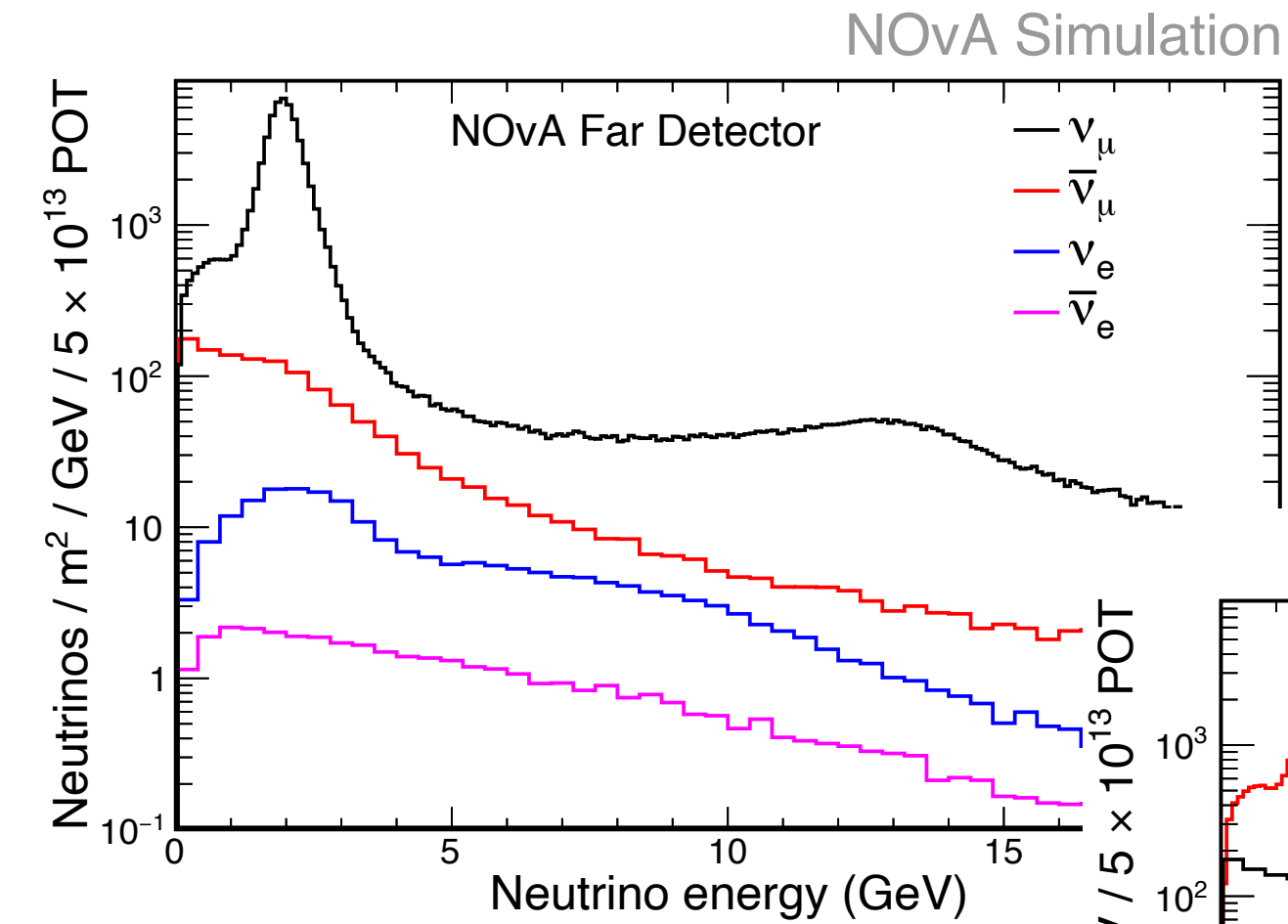


- NOvA (NuMI Off-axis electron- $\nu$  Appearance experiment)
- Long baseline neutrino experiment
- Taking data since 2014, ~250 physicists
- Aims to measure  $\nu_\mu$  disappearance and  $\nu_e$  appearance in a  $\nu_\mu$  beam
- 3 parts
  - NuMI proton accelerator, target and neutrino beam line producing a 2 GeV  $\nu_\mu$  beam (km 0) @ Fermilab
  - Near detector (300 tonnes of liquid plastic scintillator, 1 km away) @ Fermilab
  - Far detector (14 ktonnes, 810 km away) @ Ash River

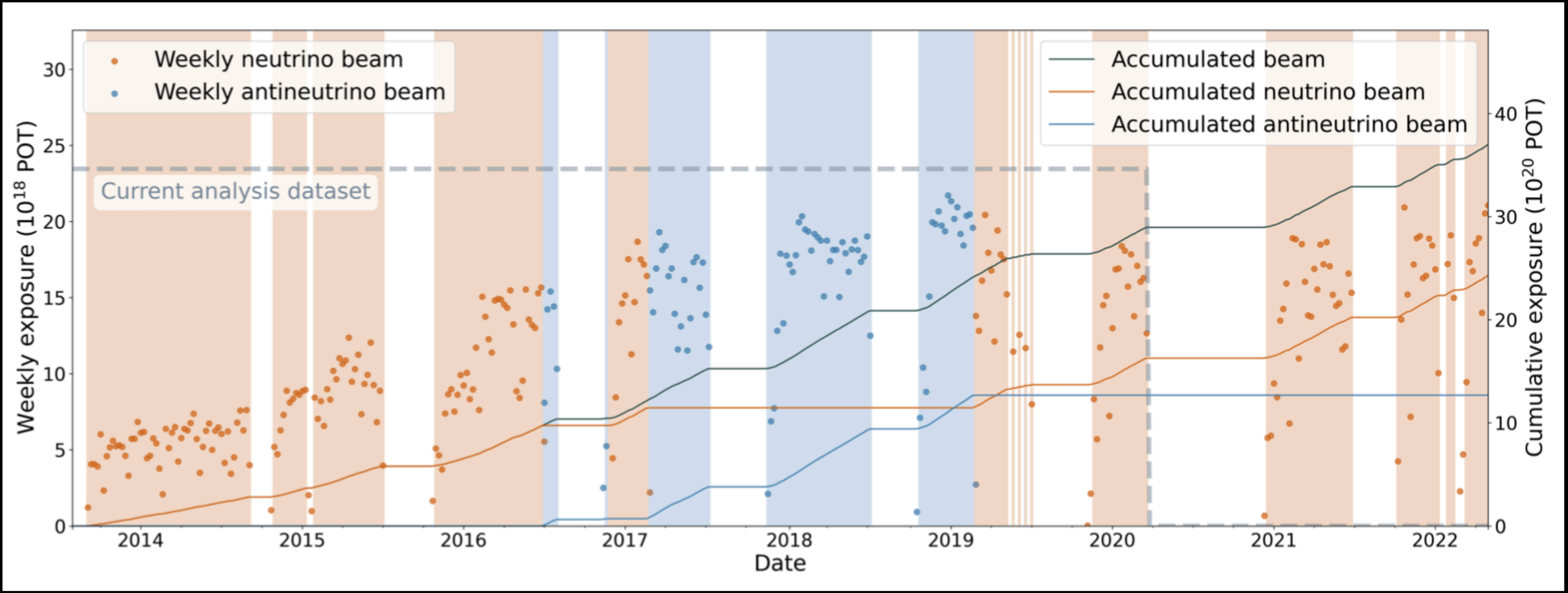




- NuMI proton accelerator
  - Capable of delivering ~850 kW (current record at 893 kW) of 120 GeV protons
  - Collimated to a 1m graphite target
- Produces a lot of  $\pi$ , which decay into  $\nu_\mu$ 
  - $\pi$  propagate in an intense magnetic field produced by the horns and decay in the 675 m decay pipe
- Detectors sit 14 mrad from the beam axis and see  $\nu_\mu$  of energy 2 GeV (off-axis effect)
- Horn current can be inverted:
  - FHC “forward horn current”  $\rightarrow$  neutrino beam
  - RHC “reverse horn current”  $\rightarrow$  anti-neutrino beam



- NuMI proton accelerator target and neutrino beam line



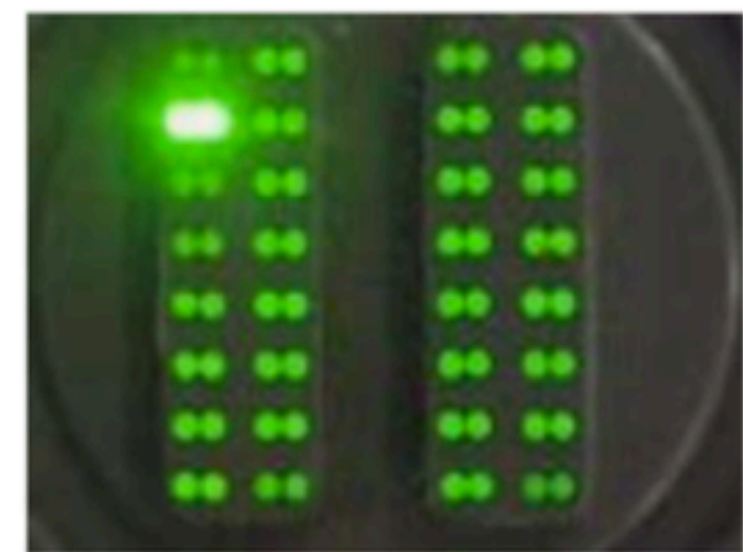
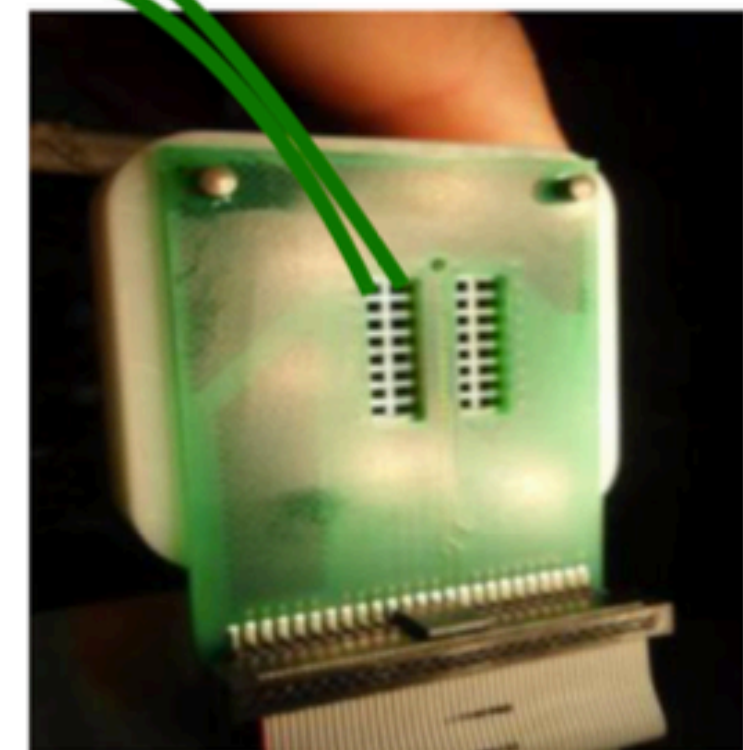
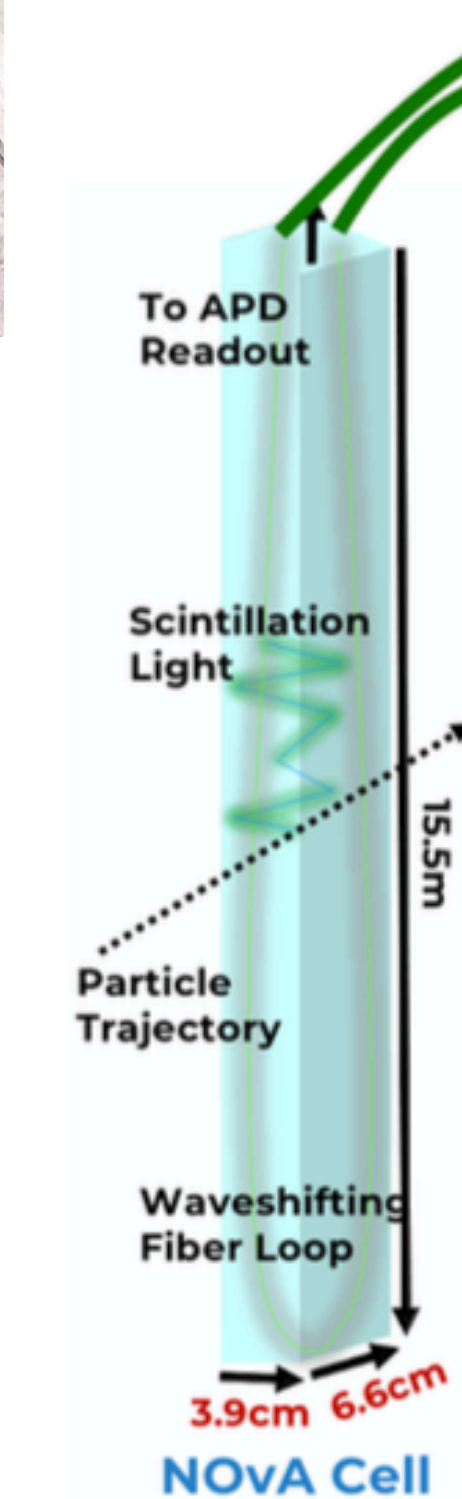
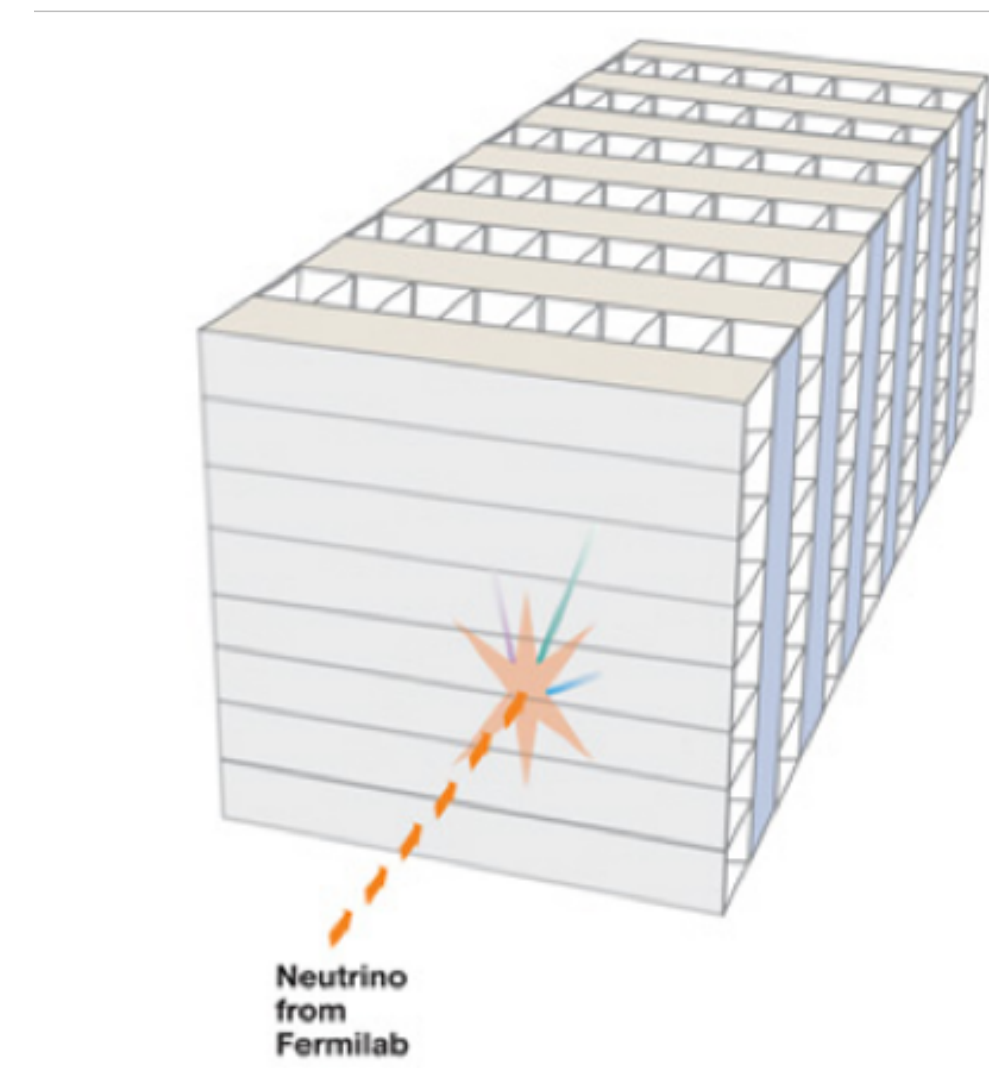
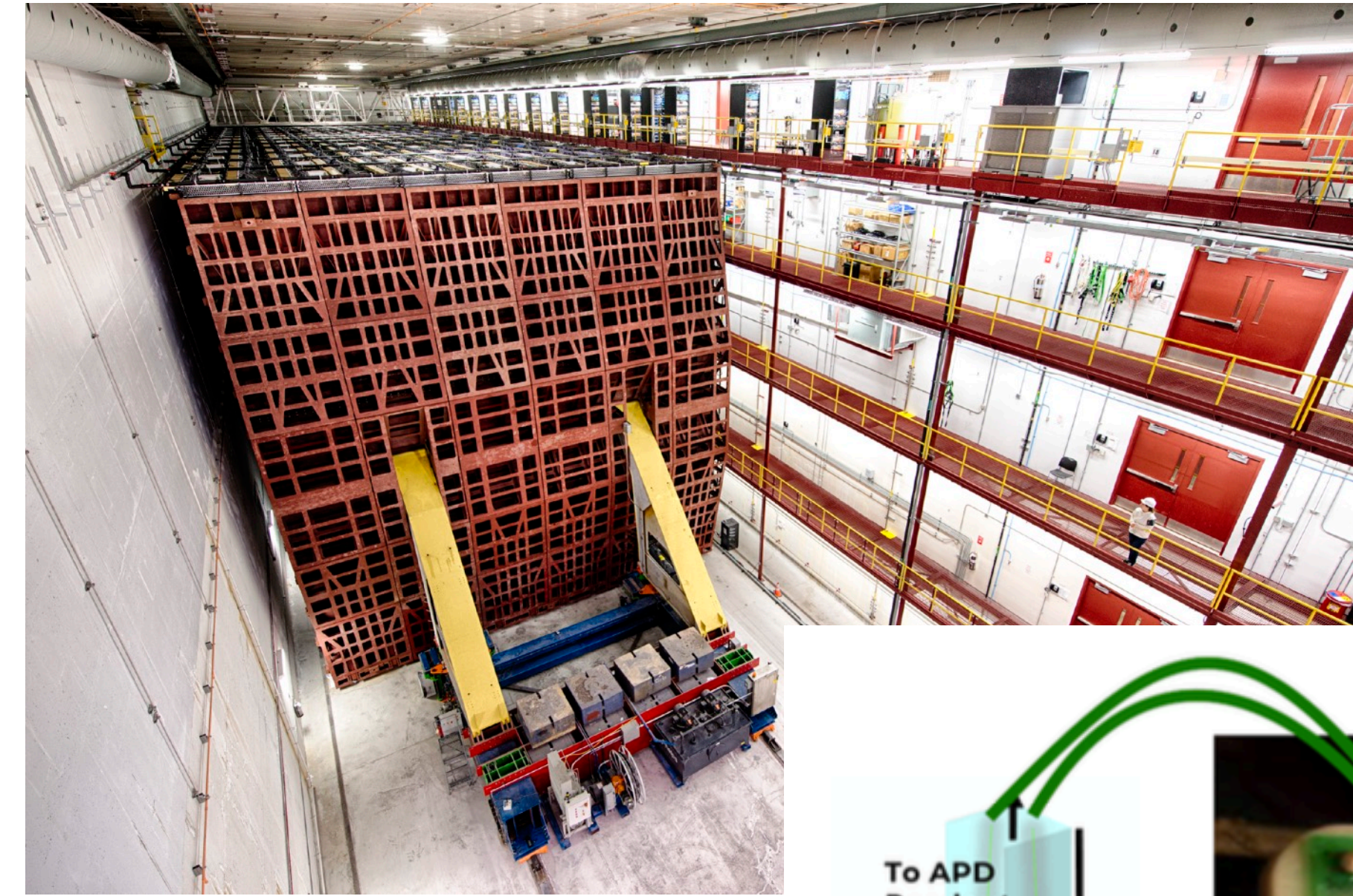
$p$

$\pi^+$

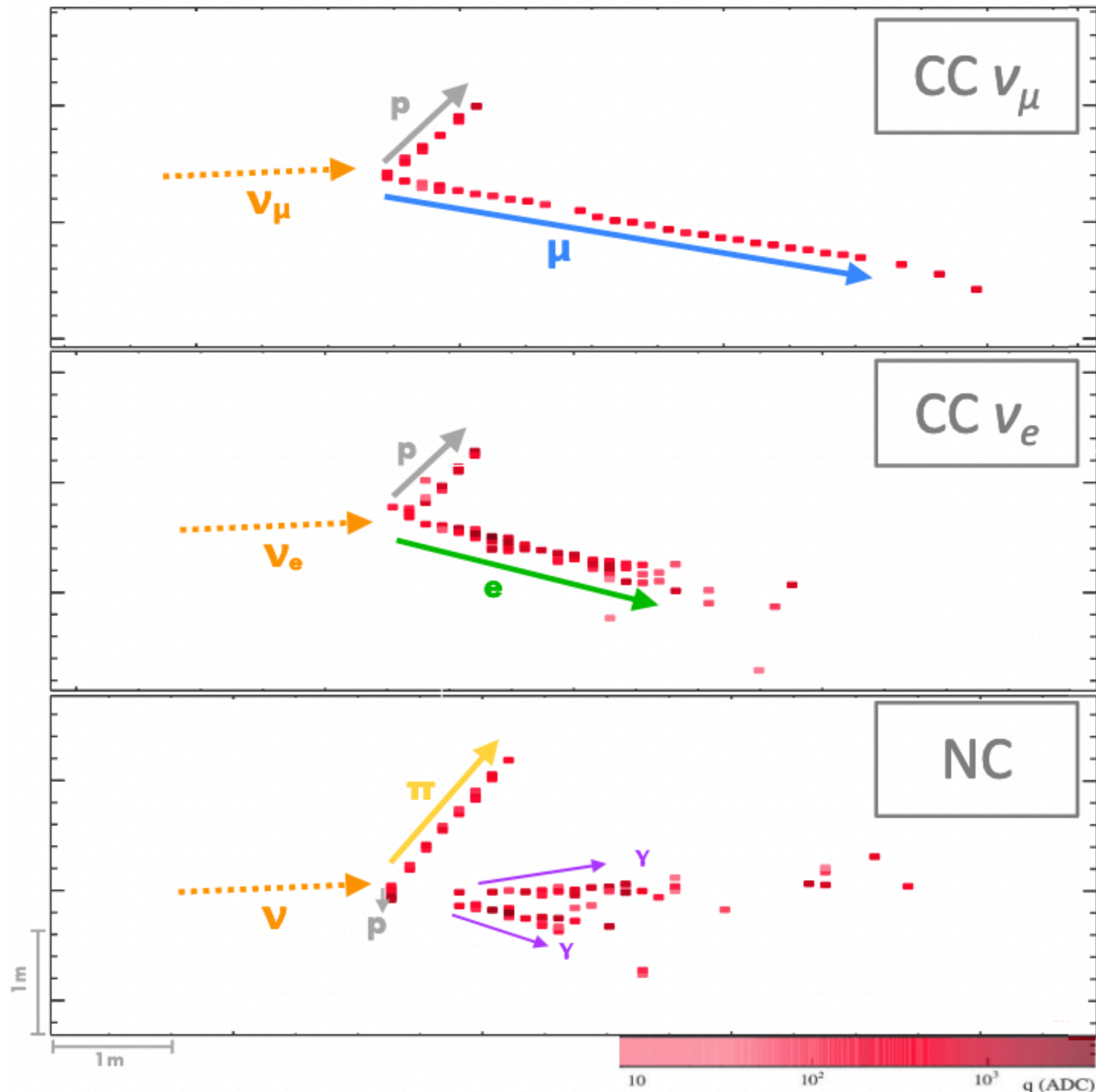
$\nu_\mu$



- Near and far detectors are functionally identical (km 1 and 810)
- Tracking calorimeters
  - Extruded plastic cells filled with liquid scintillator
  - Alternate in vertical and horizontal position
- Wavelength shifting fibre running through each cell
- Avalanche Photodiode readout





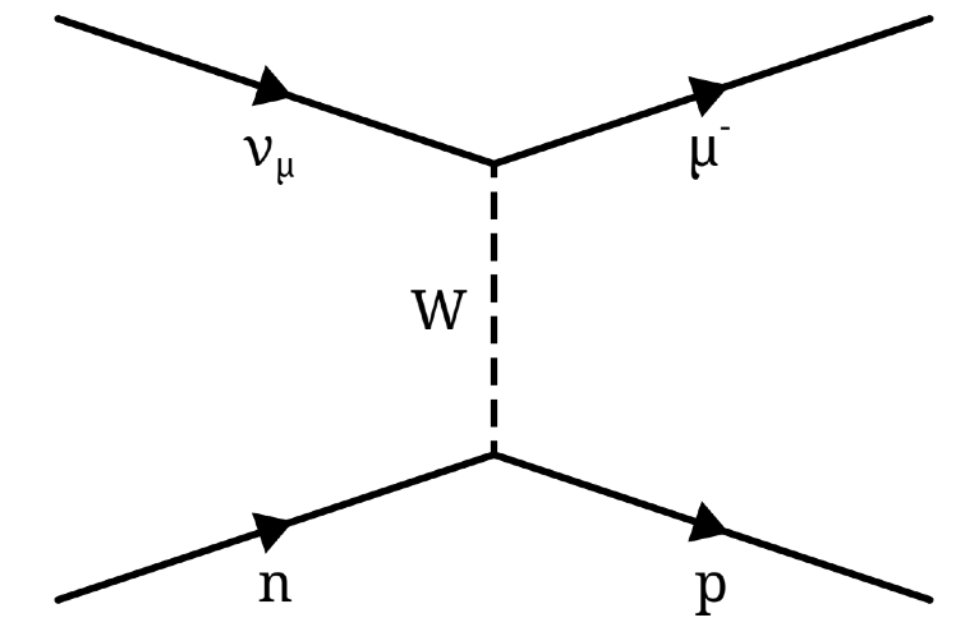


- Neutrino oscillation formula needs

- $\nu$  flavour

- L

- $E_\nu$



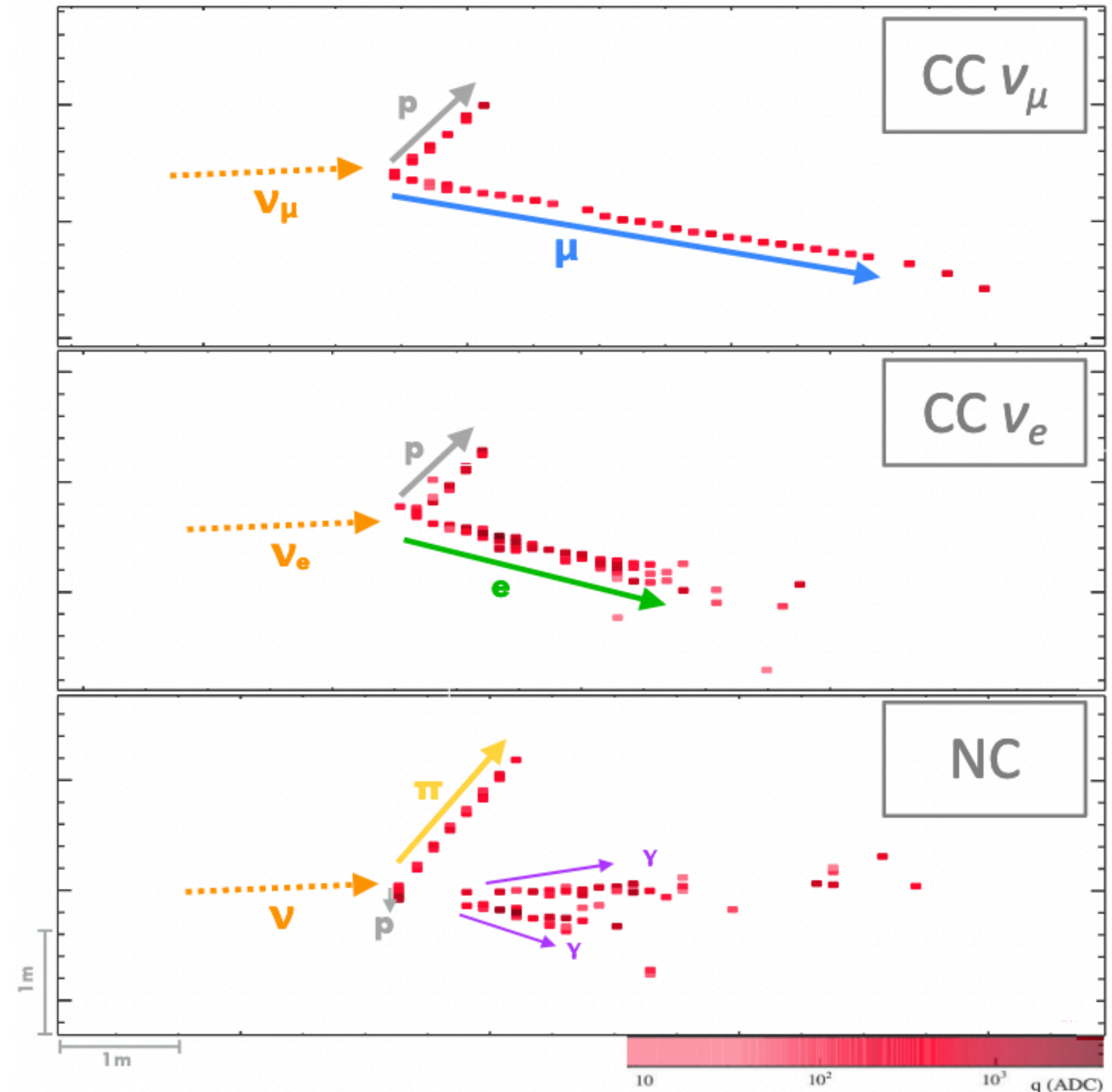
- Detailed images of the neutrino Charged Current (CC) events @ NOvA allows:

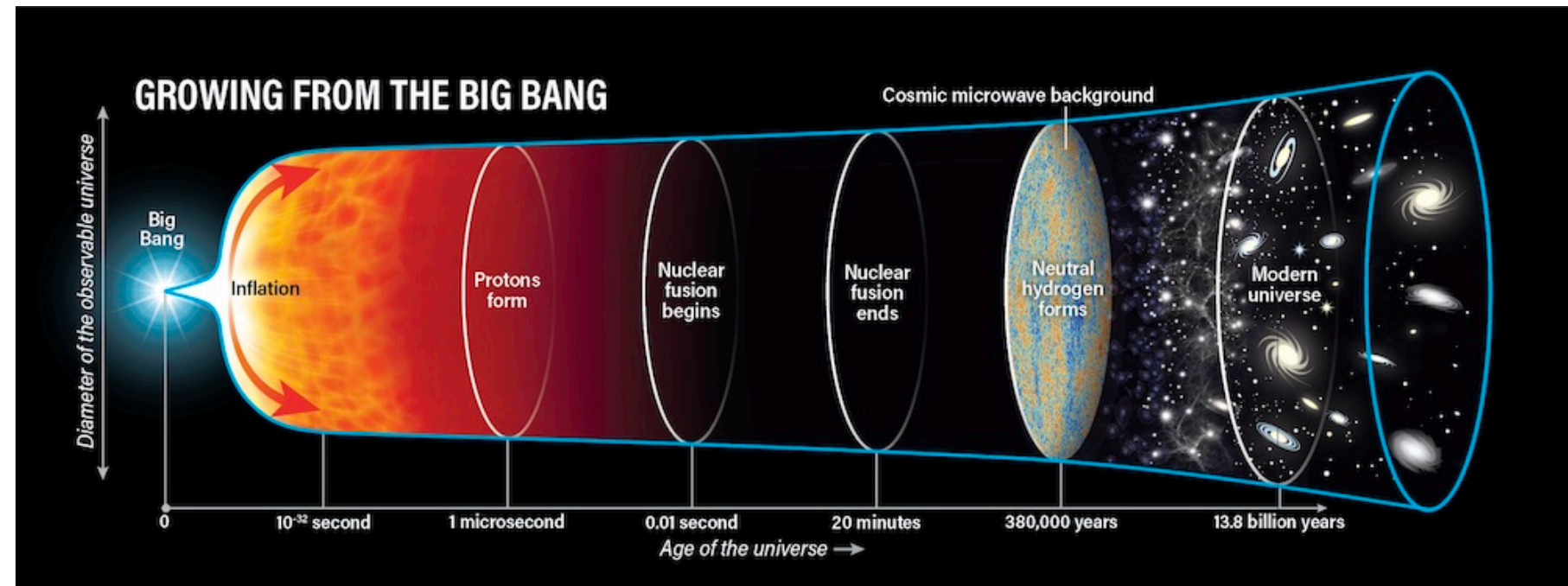
- Event classification  $\rightarrow$  Convolutional Neural Network

- Energy reconstruction



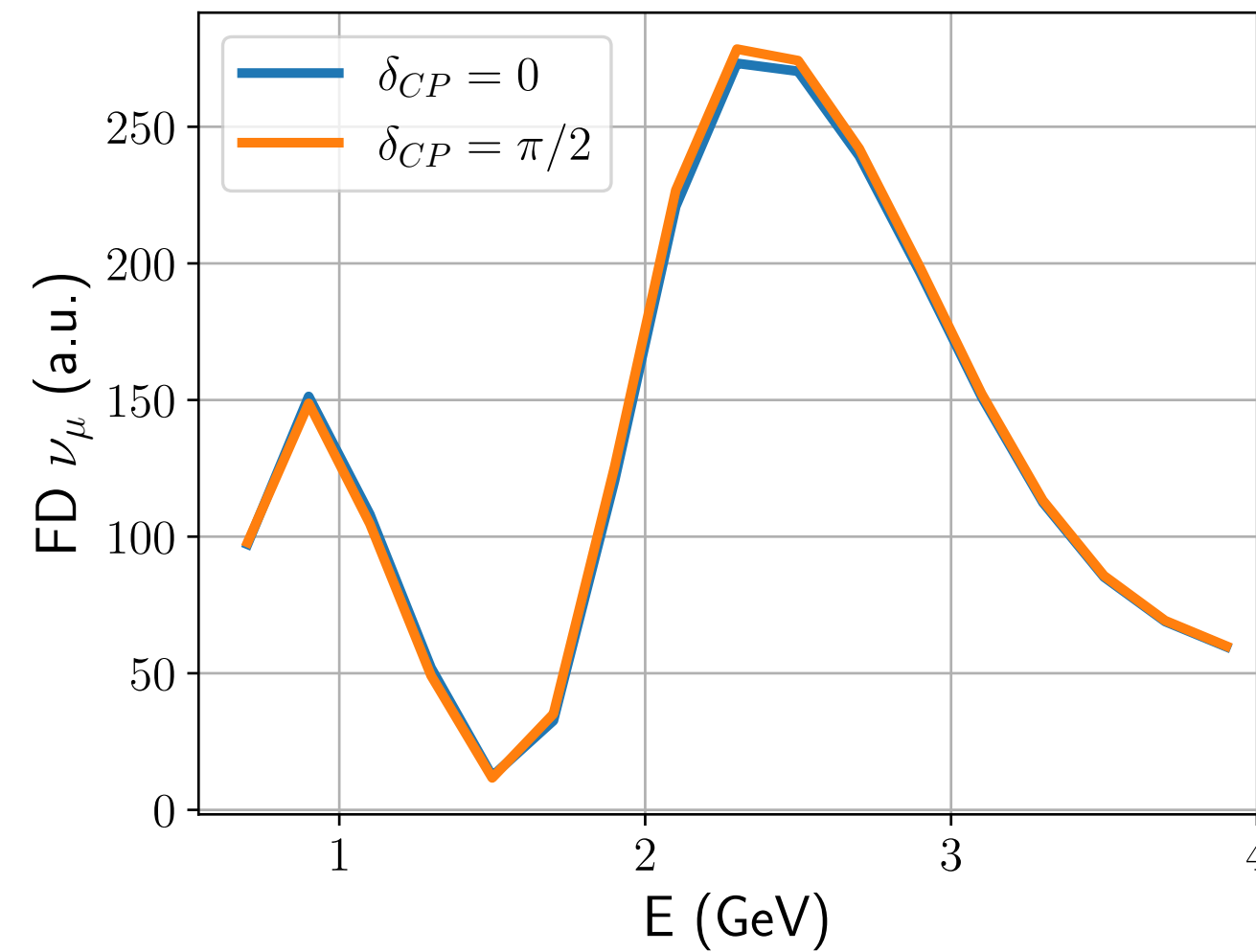
- Muon neutrinos CC:  $E_{\nu_\mu} = E_{Had} + E_\mu$   
(calorimetry and range) energy resolution: 9%
- Electron neutrinos CC:  $E_{\nu_e} = f(E_{Had}, E_{EM})$  (both with calorimetric estimation) average energy resolution: 11%
- NC interactions: Calorimetry (14~17% resolution)





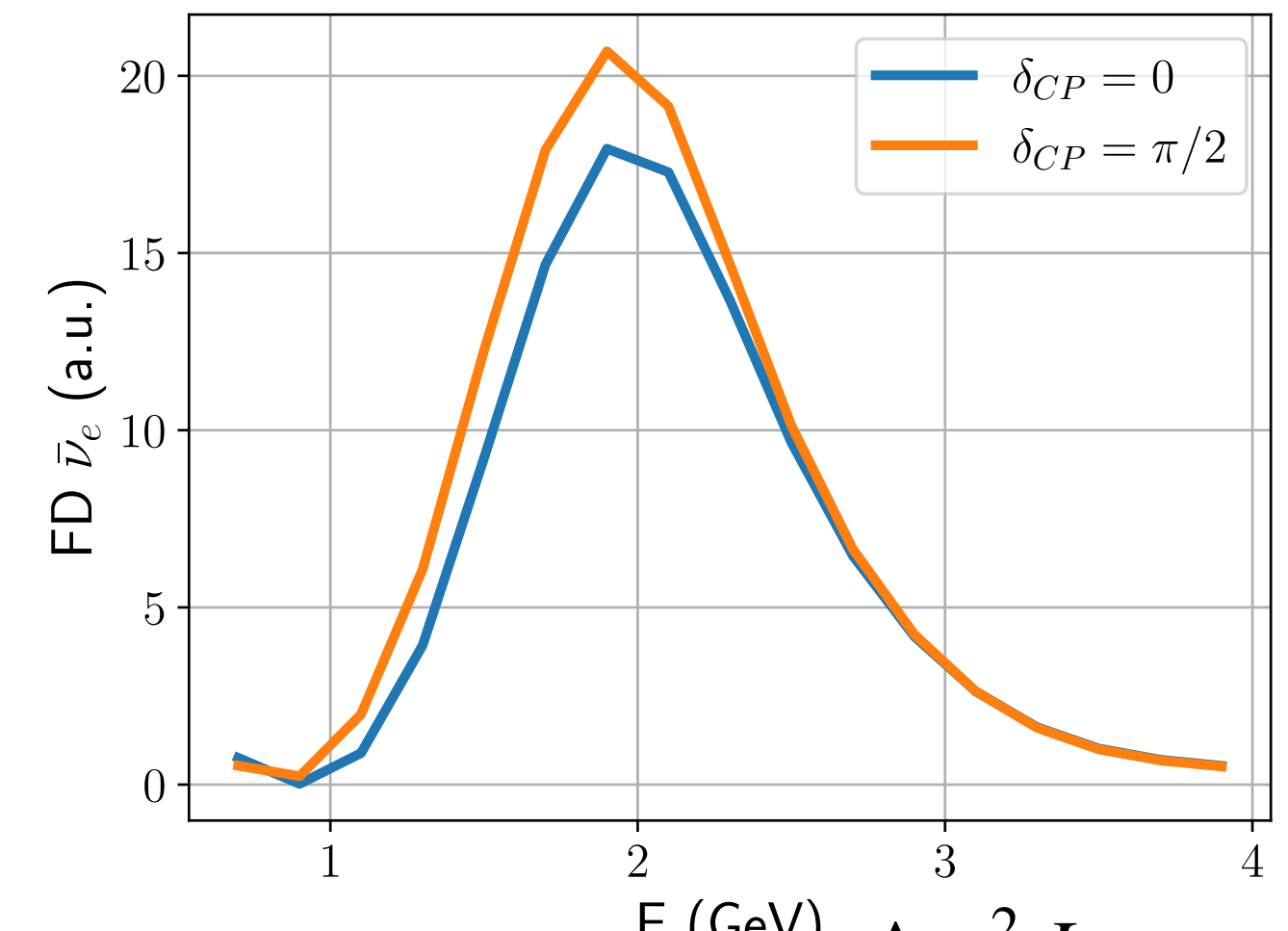
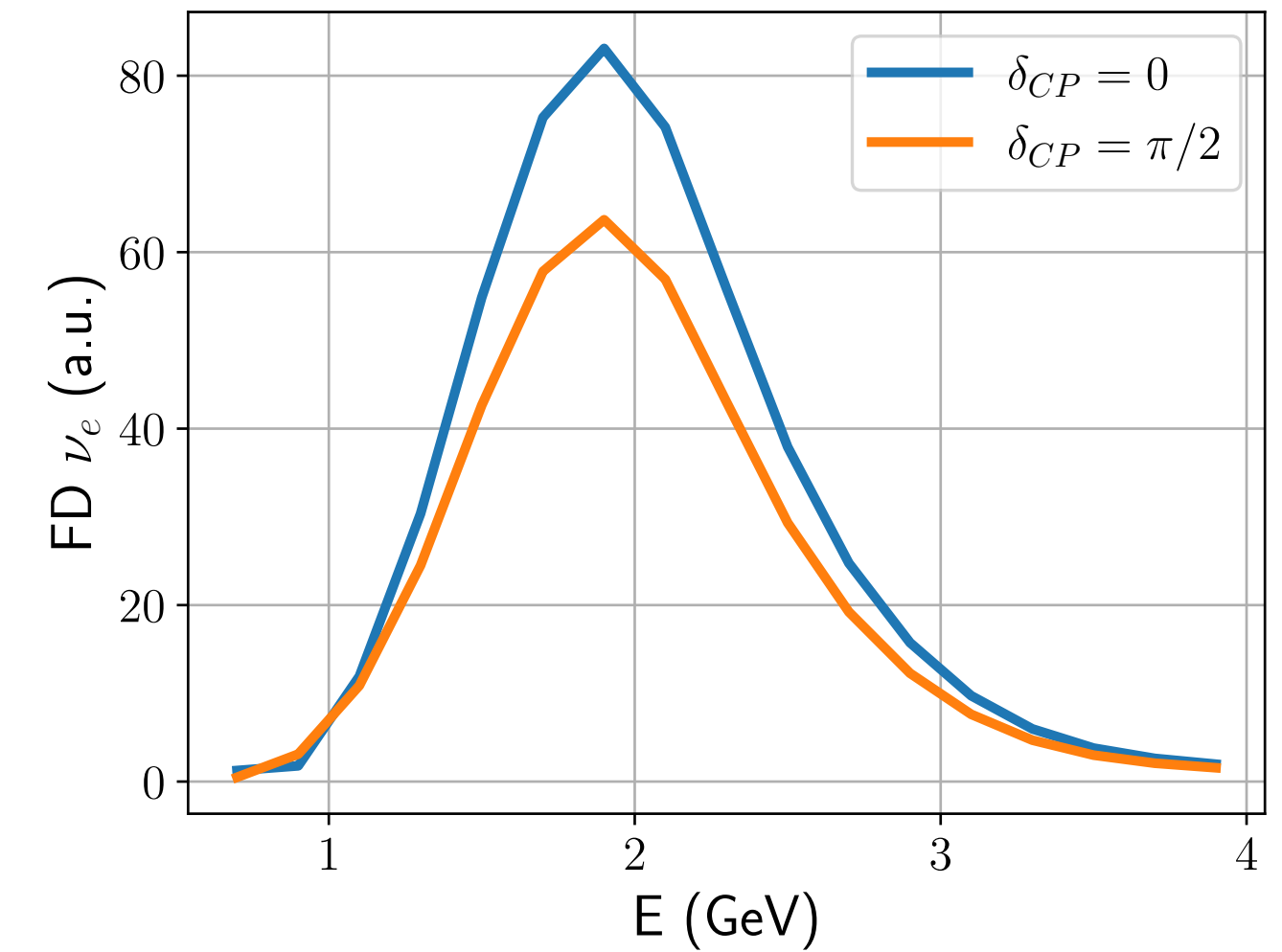
- Effect of  $\delta_{CP}$
- No effect on the  $\nu_\mu$  spectra
- Tiny effect due to matter effect
- $\nu_\mu \rightarrow \nu_e$  oscillations are reduced compared to the anti-neutrino equivalent at maximal CP violation

$$\Phi_{\nu_\mu} \times \sigma_{\nu_\mu} \times P(\nu_\mu \rightarrow \nu_\mu)$$



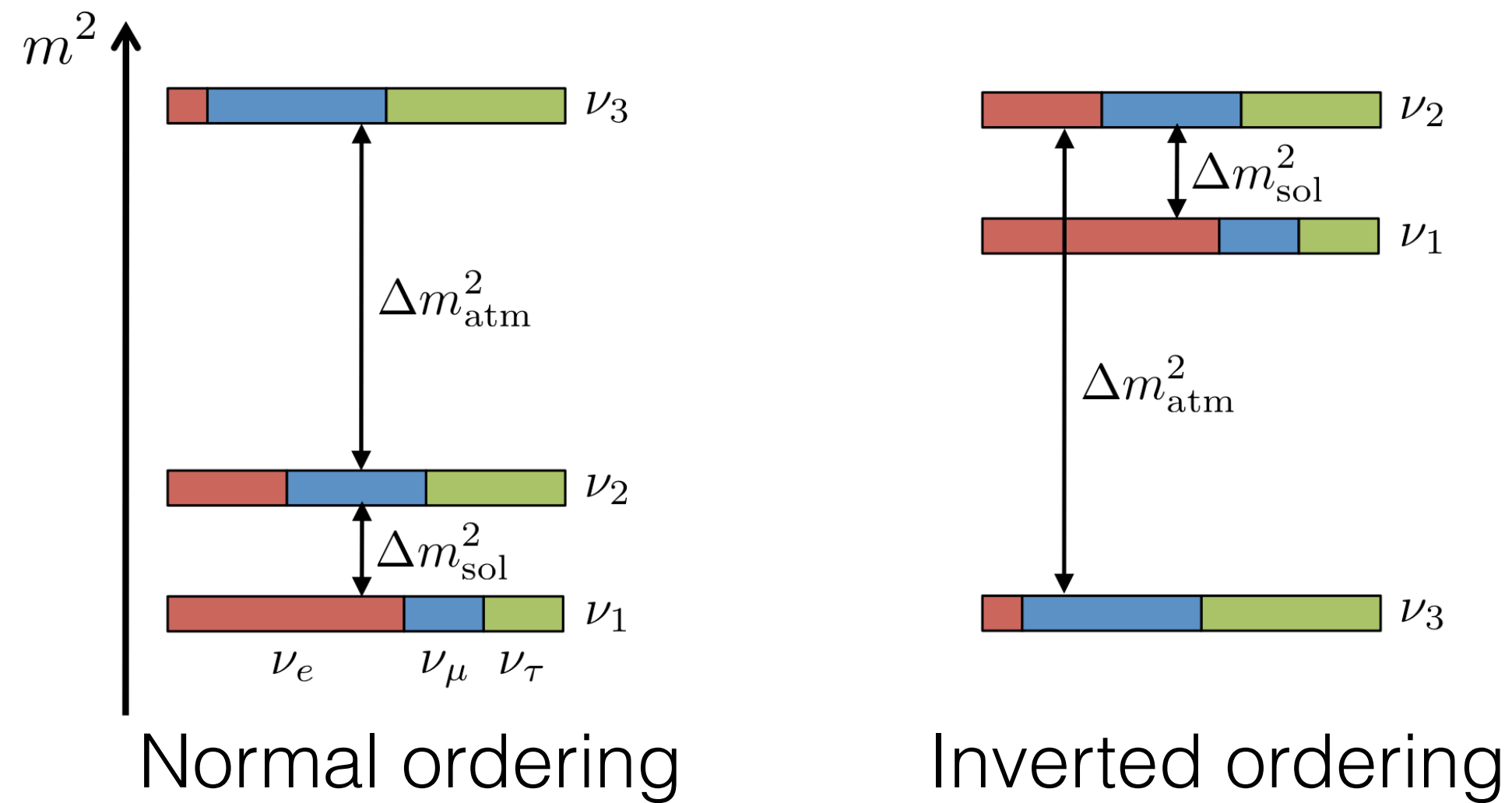
$$P(\nu_\mu \rightarrow \nu_\mu) \simeq 1 - \sin^2 2\theta_{23} \sin^2 \frac{1.27 \Delta m_{32}^2 L}{E}$$

$$\Phi_{\nu_\mu} \times \sigma_{\nu_e} \times P(\nu_\mu \rightarrow \nu_e)$$



$$P(\nu_\mu \rightarrow \nu_e) \simeq \sin^2 \theta_{23} \sin^2 2\theta_{13} \sin^2 \frac{\Delta m_{32}^2 L}{4E} + f(\delta_{CP})$$



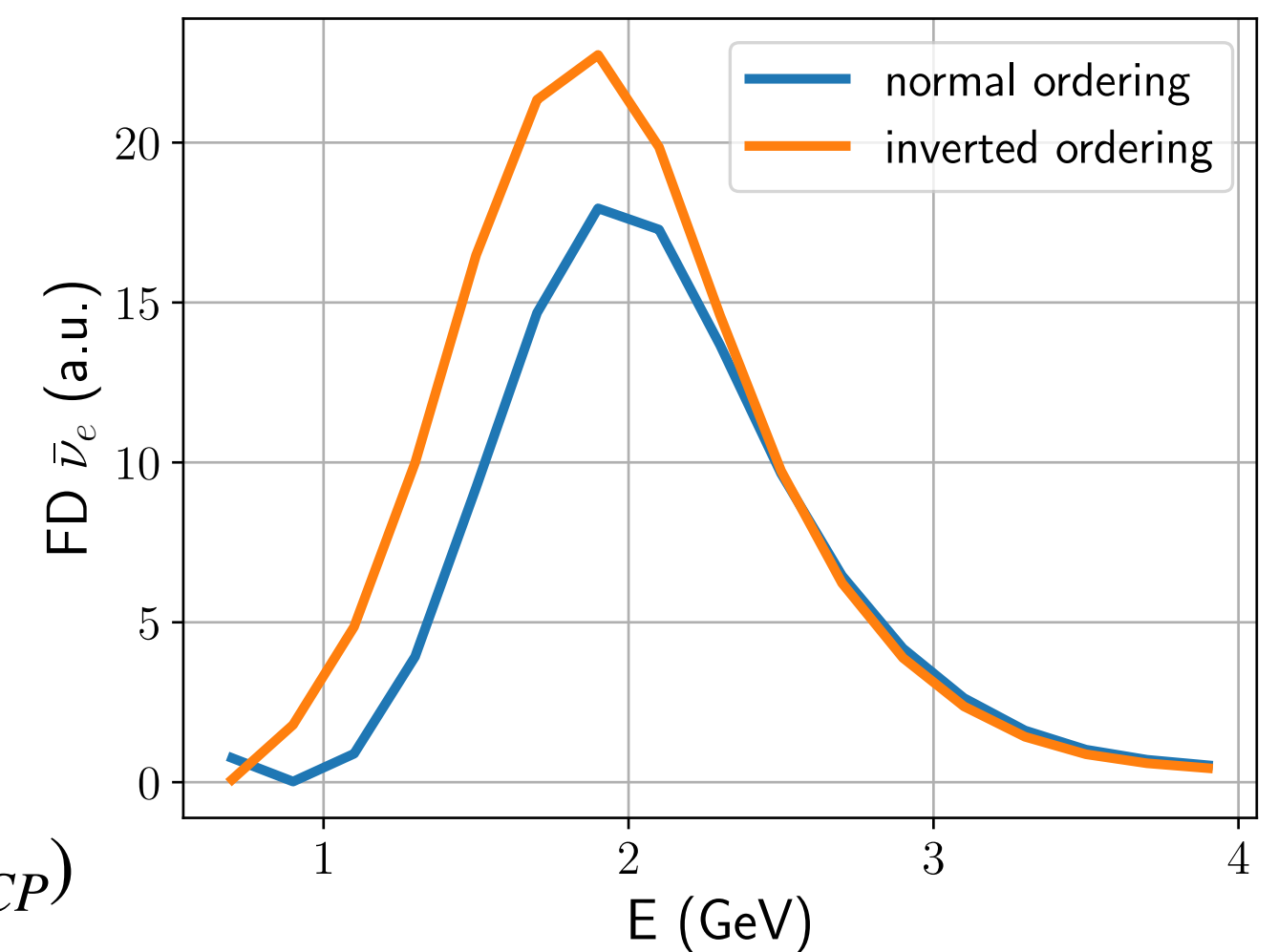
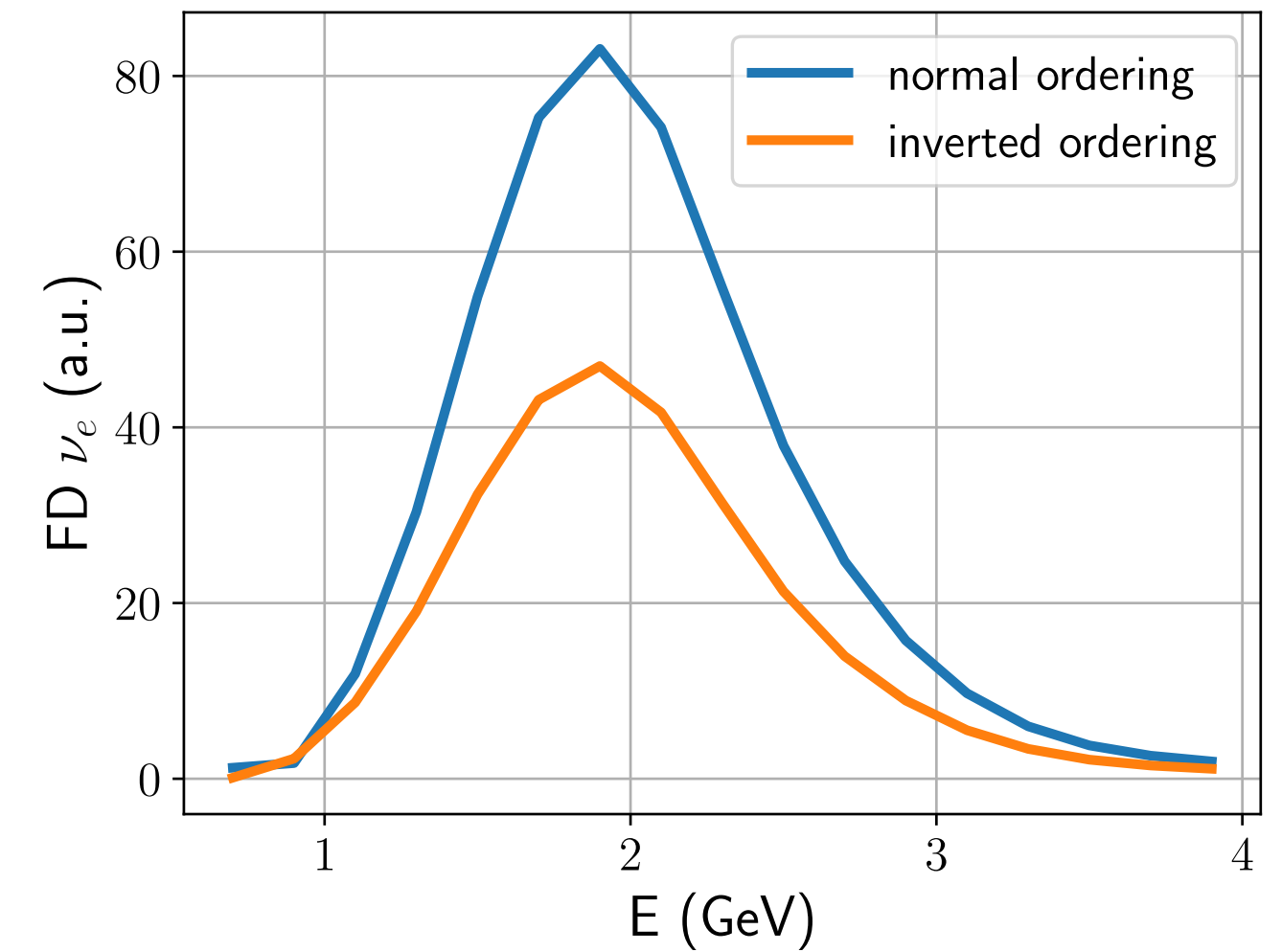
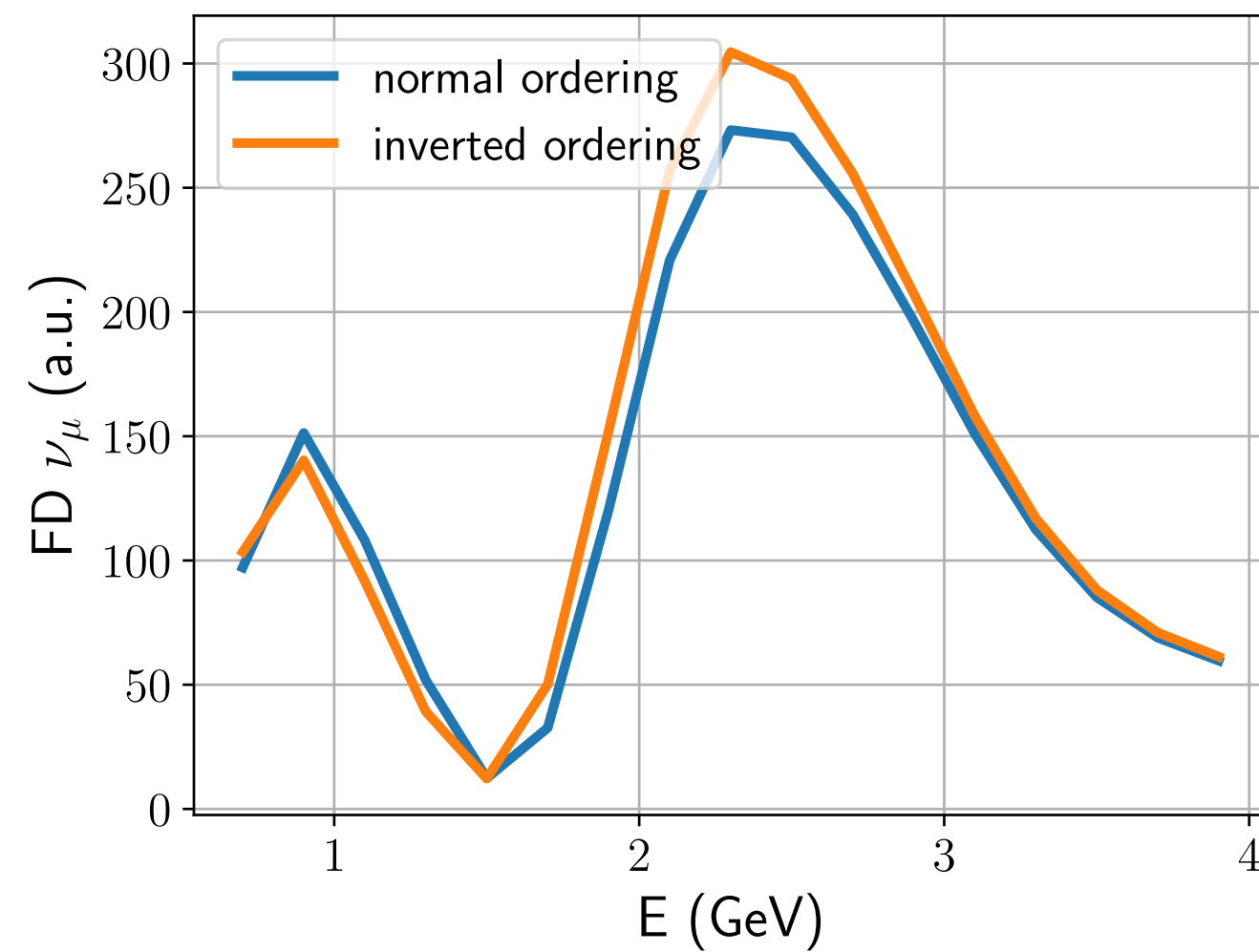


- Effect of the mass ordering

- $\nu_\mu$  channel (FHC/RHC): moves the oscillation dip a bit
- Disappearance sensitive to  $|\Delta m^2_{32}|$

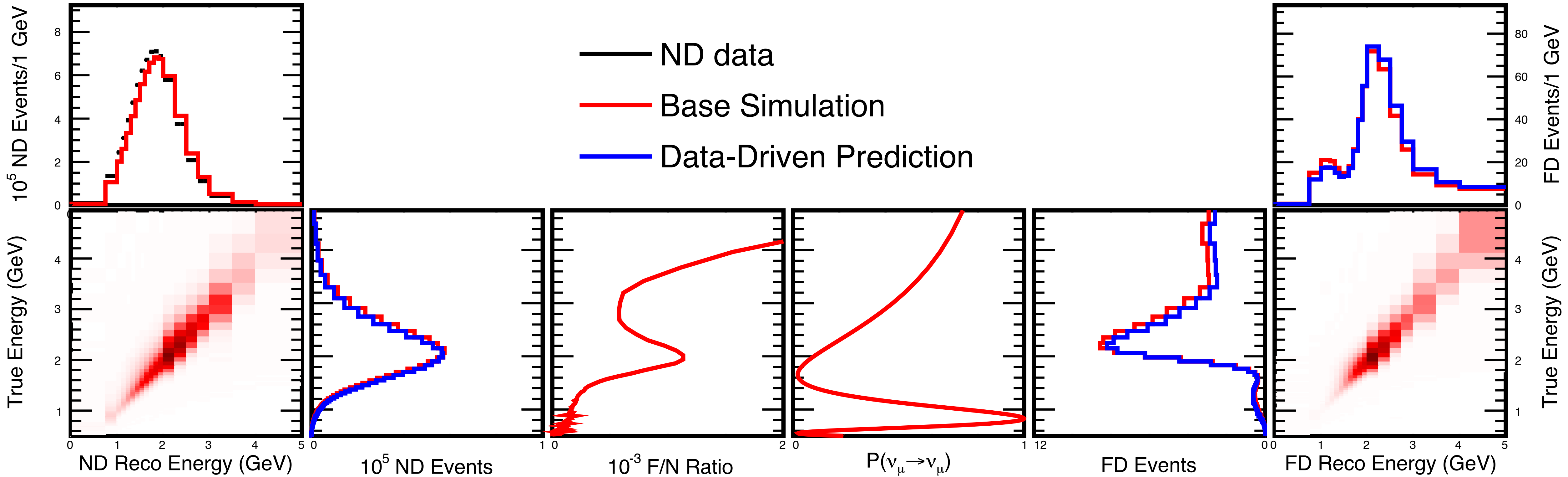
- $\nu_e$  channel: changes the normalisation in opposite way for FHC/RHC

- Appearance sensitive to  $\Delta m^2_{32}$  via matter effect  $P(\nu_\mu \rightarrow \nu_e) \simeq \sin^2 \theta_{23} \sin^2 2\theta_{13} \sin \frac{\Delta m^2_{32} L}{4E} + f(\delta_{CP})$



$$P(\nu_\mu \rightarrow \nu_\mu) \simeq 1 - \sin^2 2\theta_{23} \sin^2 \frac{1.27 \Delta m^2_{32} L}{E}$$

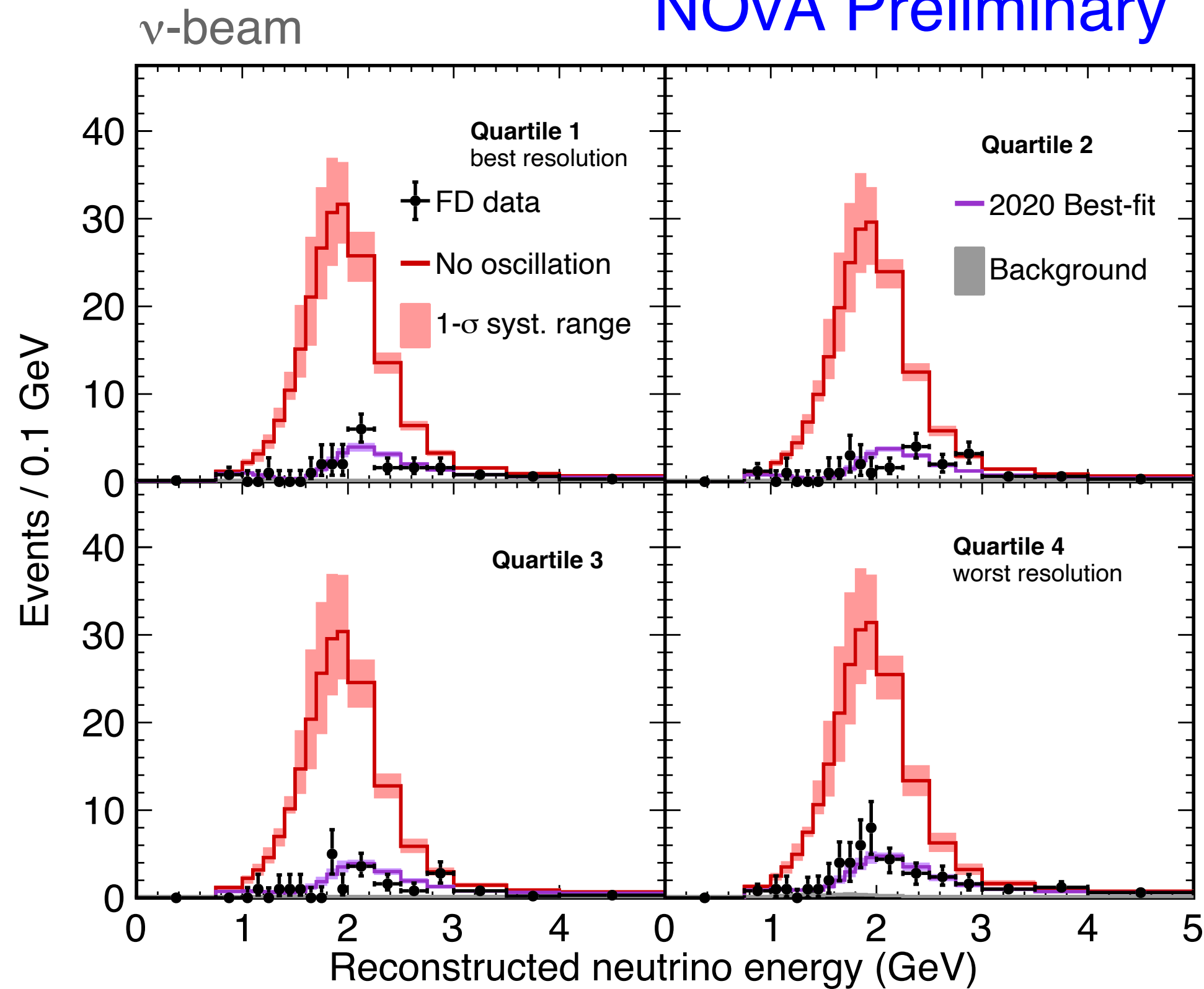
$$P(\nu_\mu \rightarrow \nu_e) \simeq \sin^2 \theta_{23} \sin^2 2\theta_{13} \sin \frac{\Delta m^2_{32} L}{4E} + f(\delta_{CP})$$



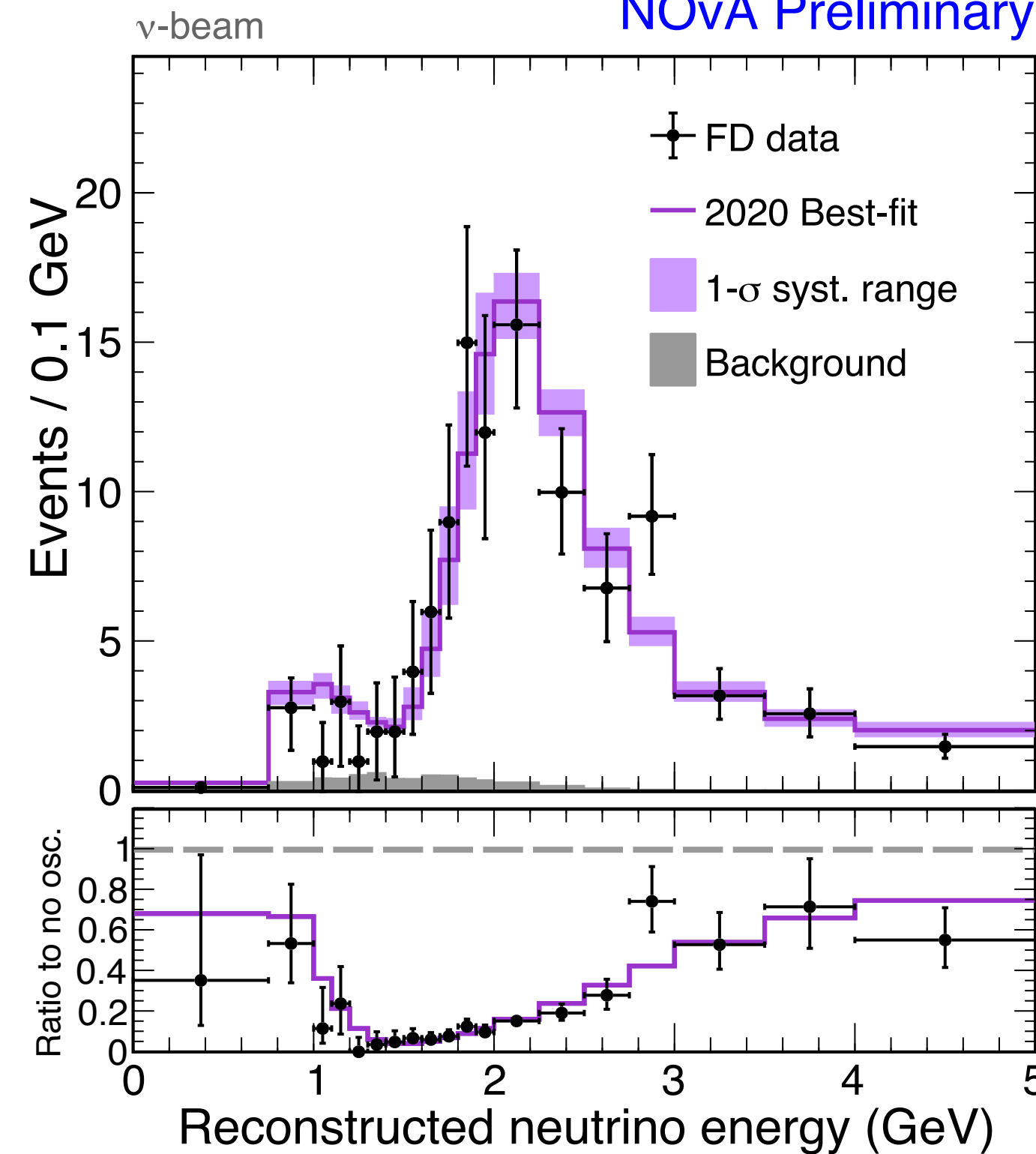
- Makes extensive use of the functionally identical detectors
  - Basic idea is to unfold the ND spectra to true energy, add oscillation, corrections (beam divergences etc.), and fold it again to the FD
- Able to drastically reduce systematic uncertainties



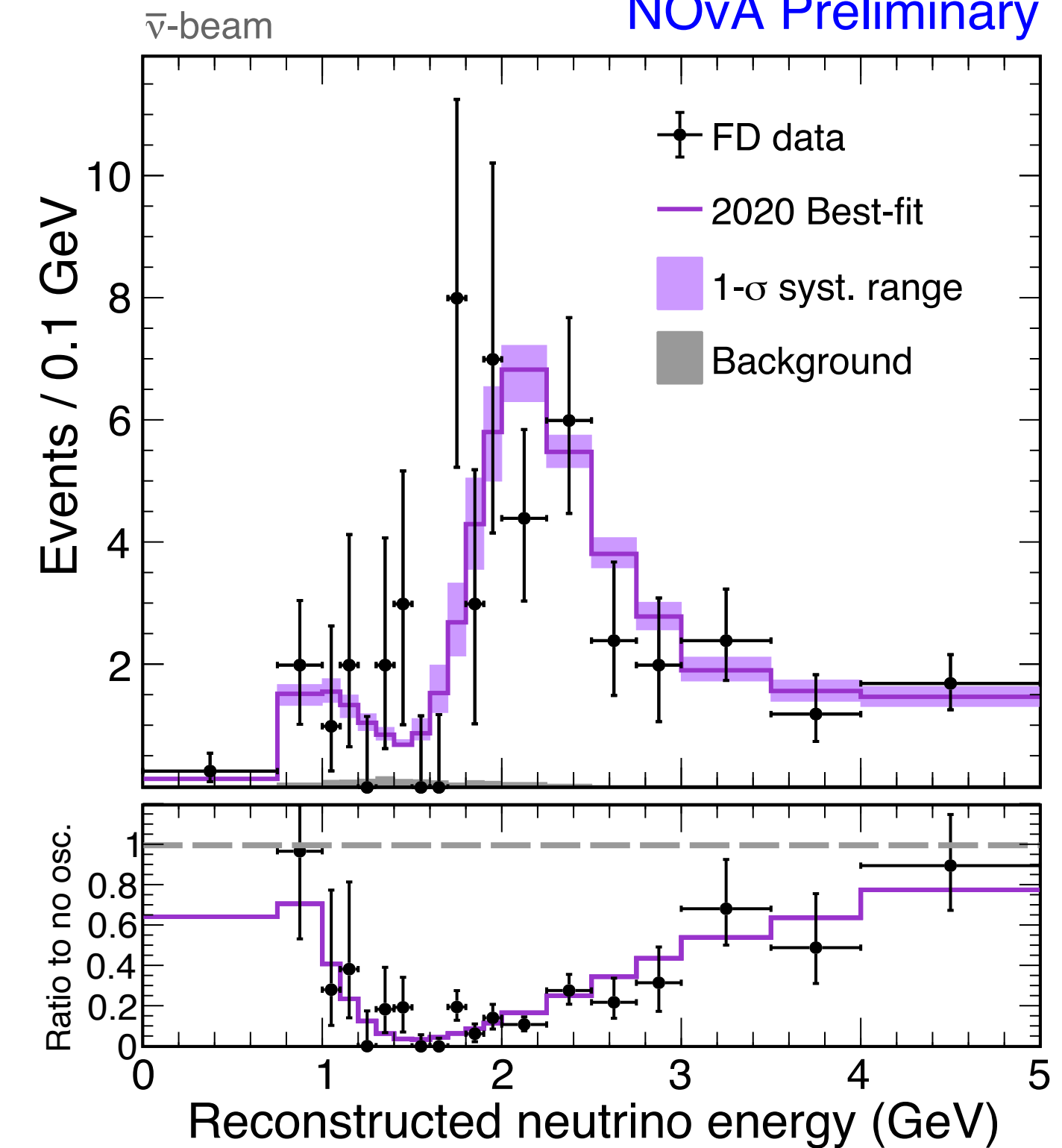
## NOvA Preliminary



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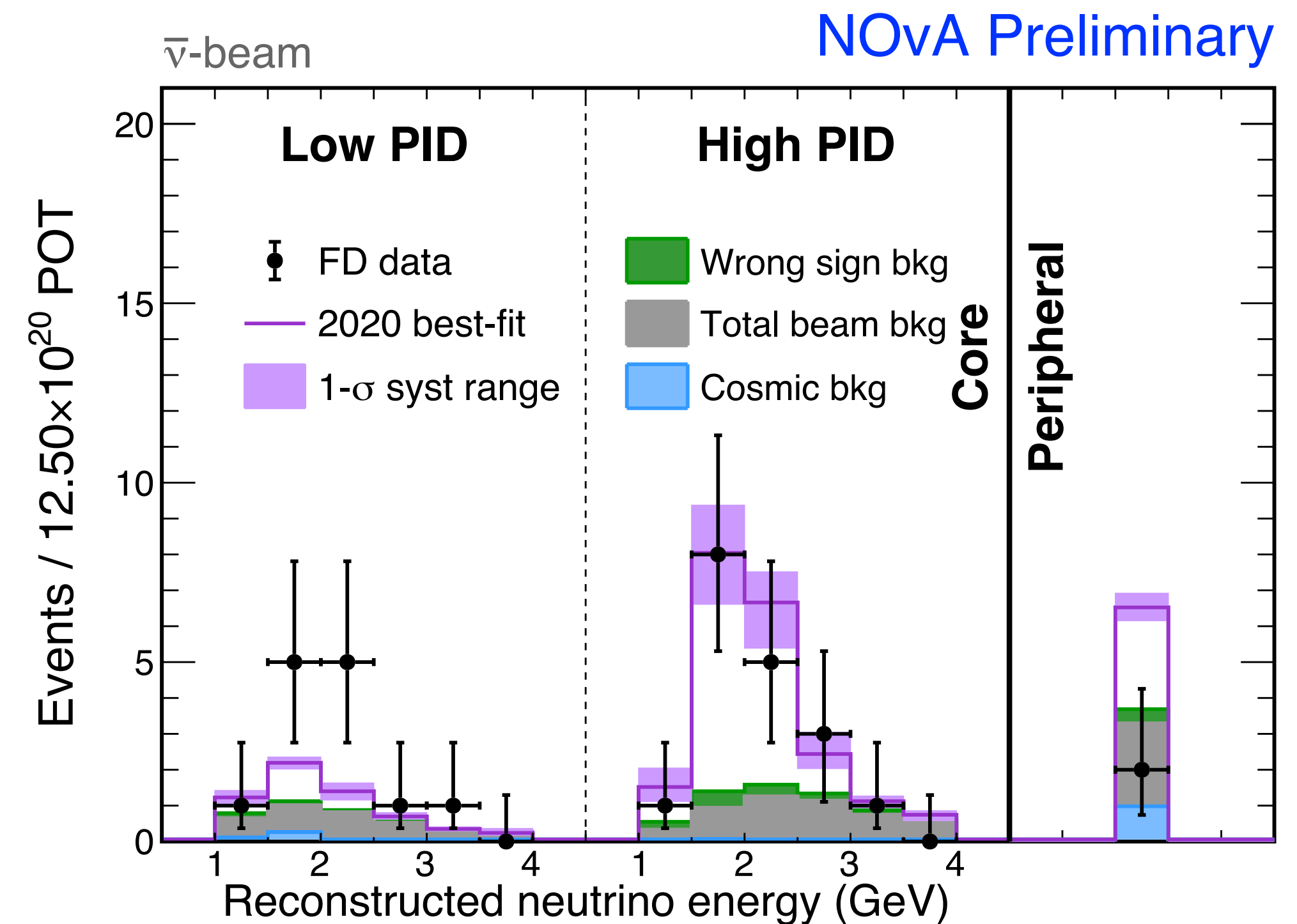
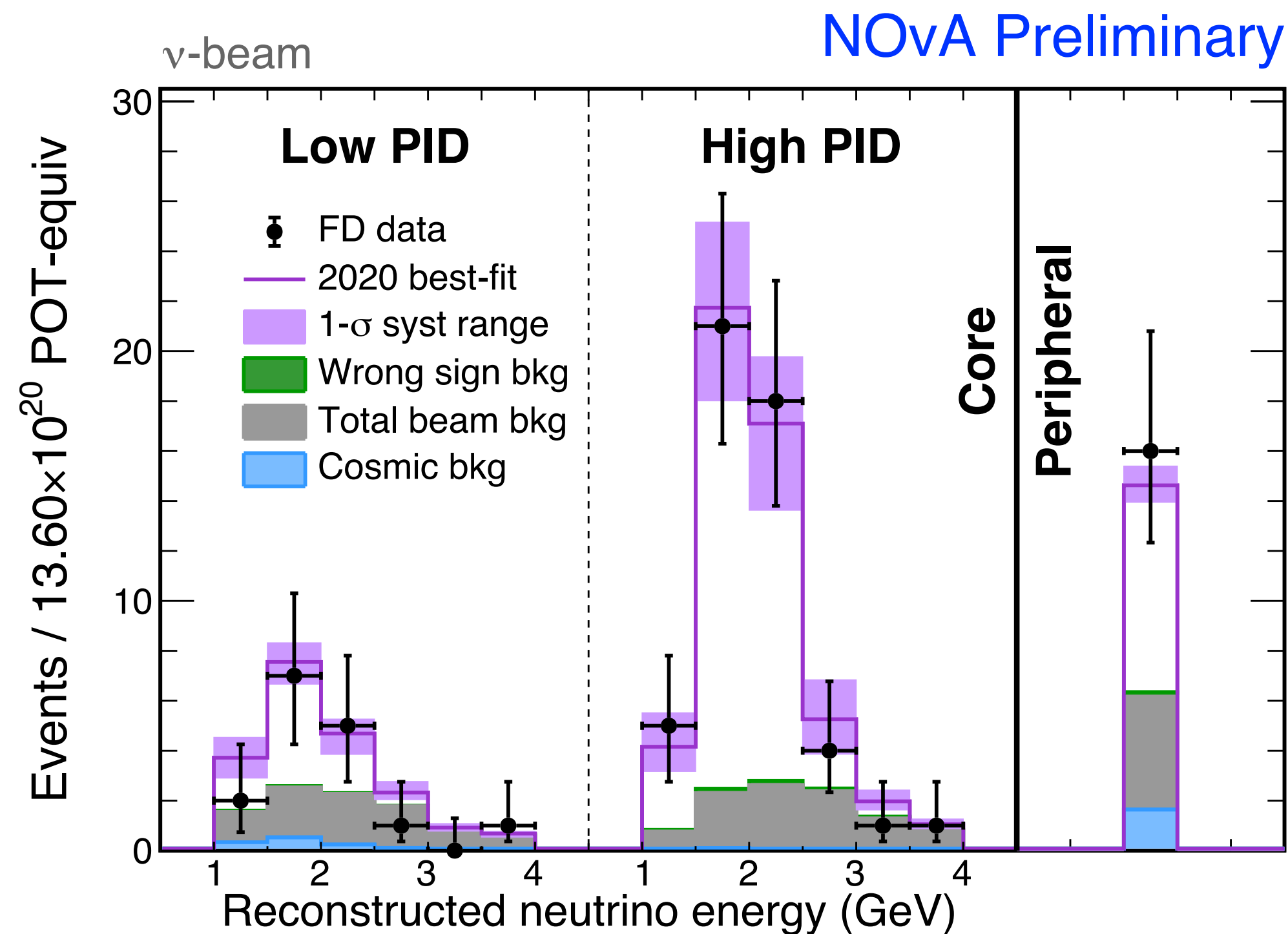


- Fitted:  $E_{\text{Had}}$  fraction quantile
- Muon  $|\text{pT}|$  bins are extrapolated separately

- Typical dip for the disappearance of (anti-) $\nu_\mu$ 
  - FHC observed: 211  $\nu_\mu$  (BG: 8.2)
  - RHC observed: 105  $\nu_\mu$  (BG: 2.1)

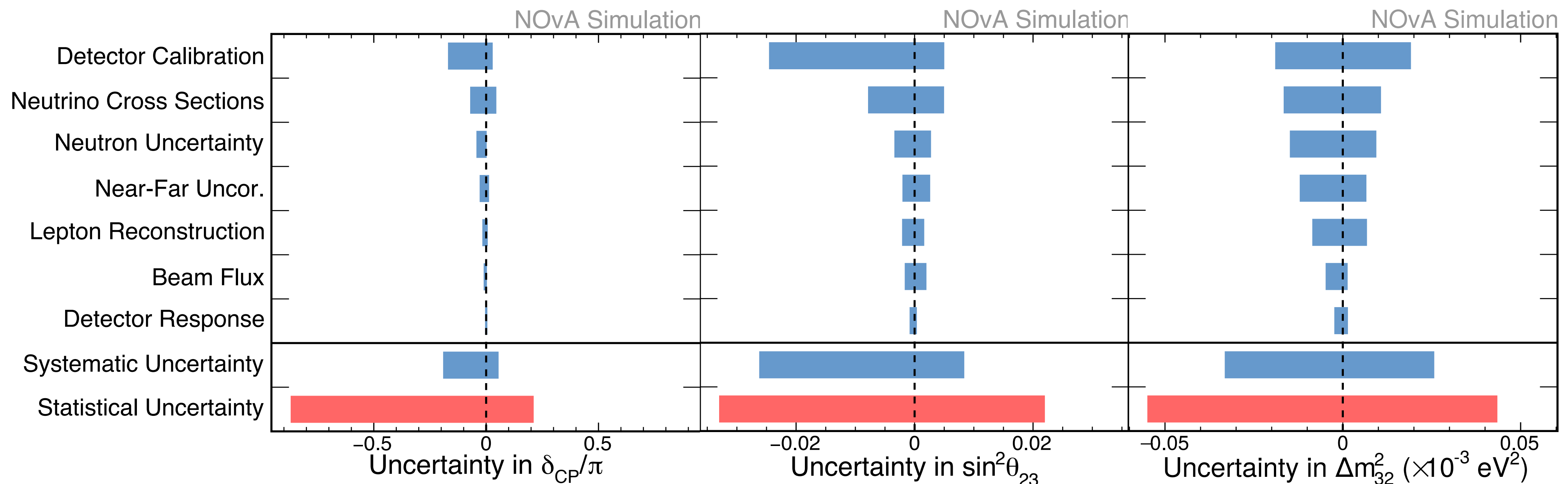


- Split by PID score from a convolutional neural network
- FHC observed: 82 (BG: 26.8)
- RHC observed: 33 (BG: 14)

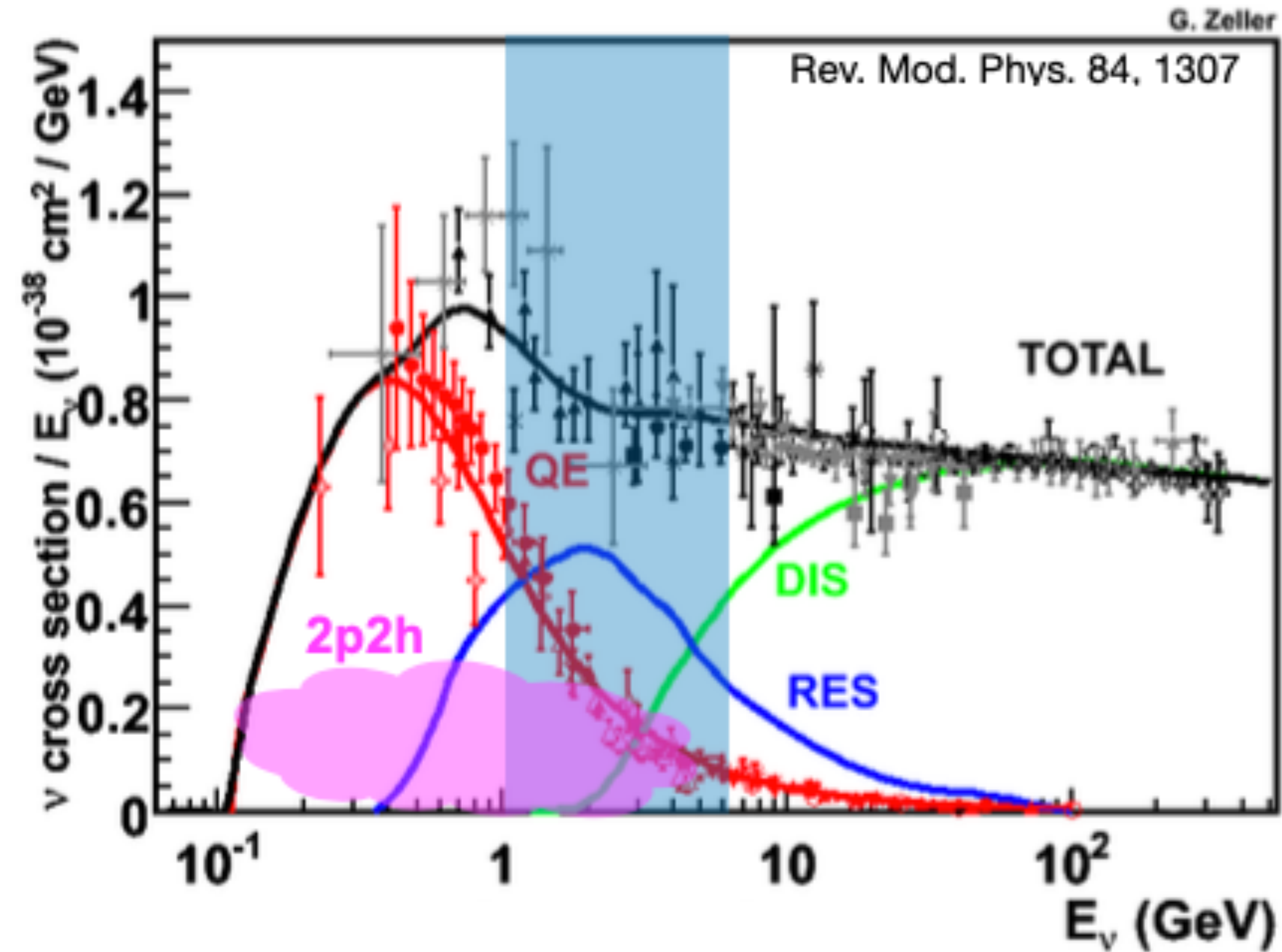
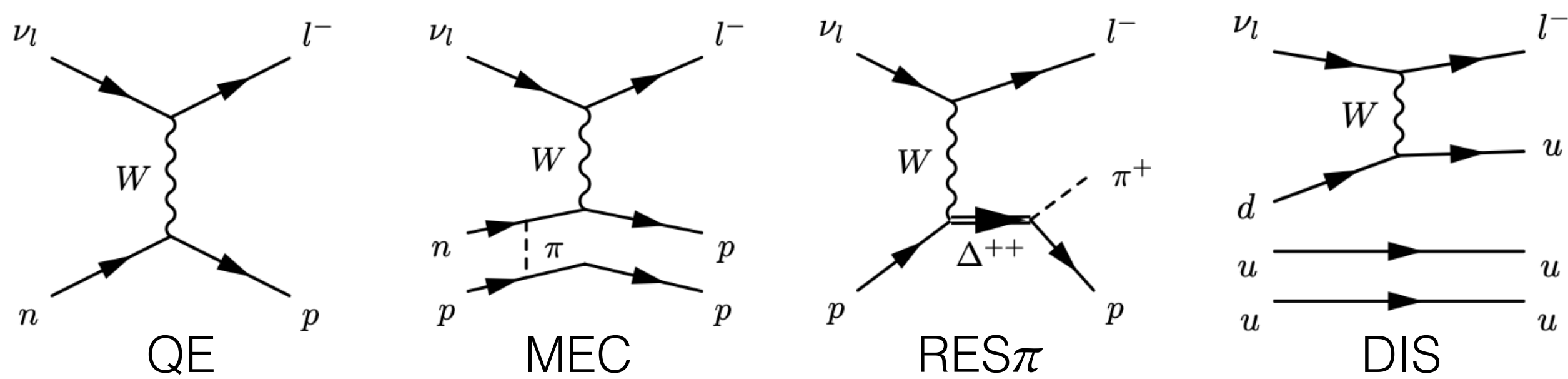
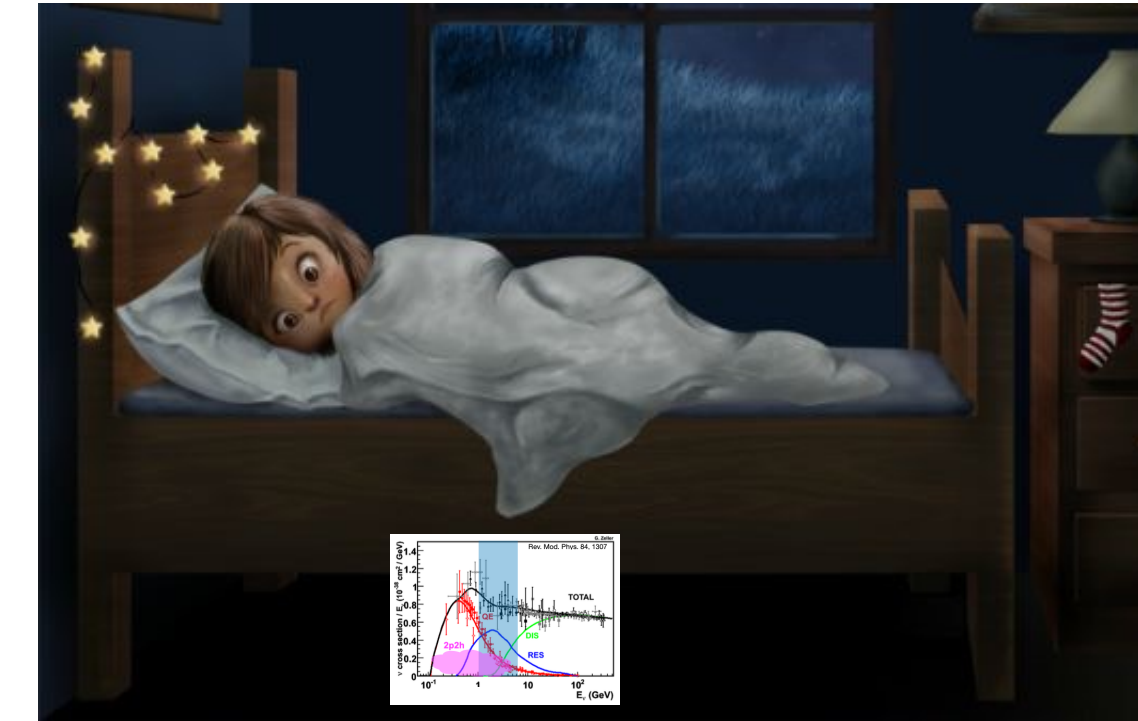




- $\nu_e$  samples are statistically limited  $\rightarrow$  effect of systematics on  $\delta_{CP}$  and  $\theta_{13}$  is negligible
- Main effects for  $\nu_\mu$  samples
- Calibration and neutrino cross section



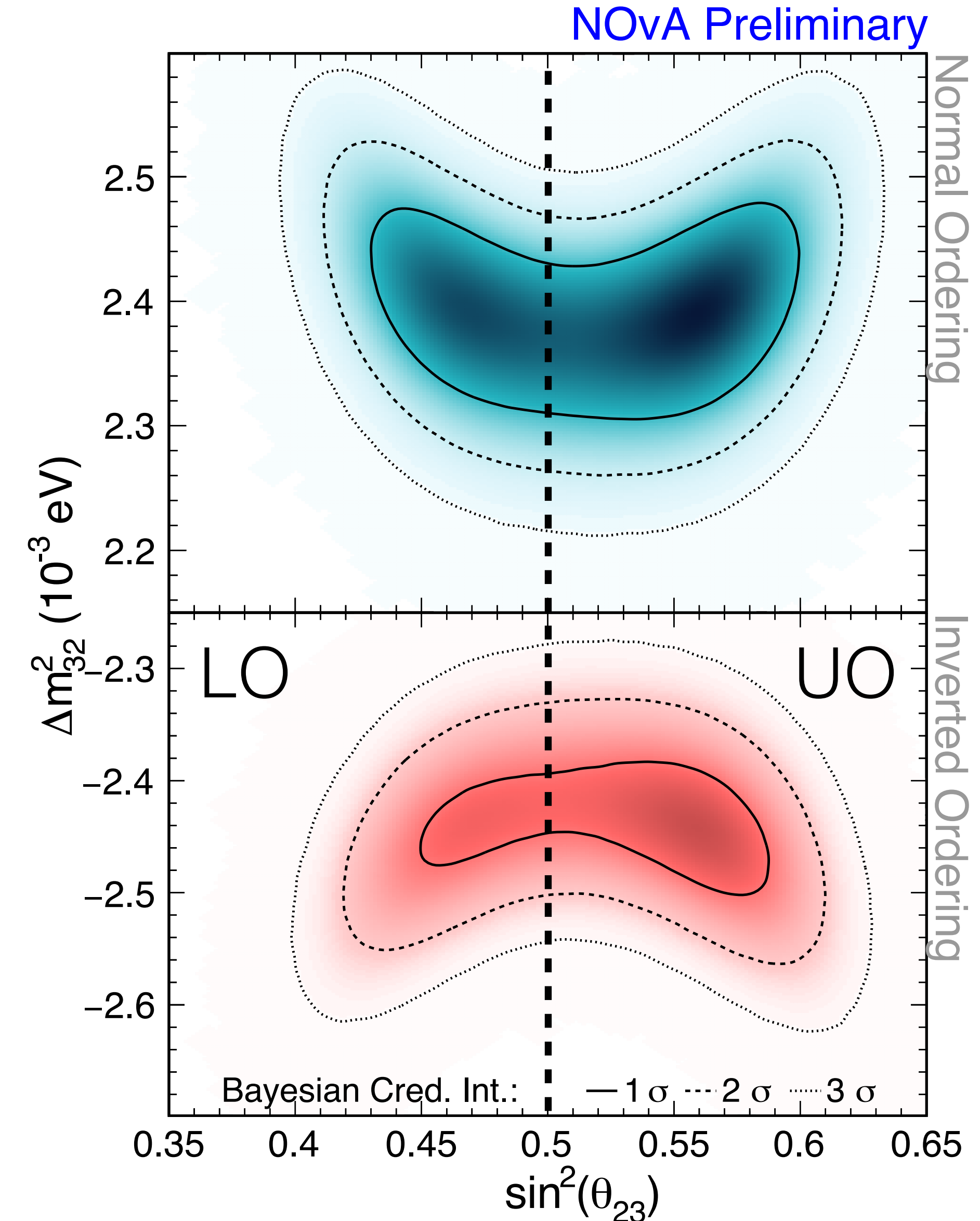
- (Should be) every LBL neutrino physicist's nightmare
  - A lot of unanswered questions that have a high impact on long baseline neutrino experiments
- NOvA uses GENIE v3.0.6 as neutrino nucleus interaction generator
  - Tuned by the GENIE collaboration and us (so-called N1810j\_0211a - "NINJA" tune)
    - Z-expansion Quasi-Elastic Form Factor
  - Extra systematics
    - Low  $Q^2$  suppression
    - Multi nucleon knockout (2p2h/MEC)
    - Final State Interactions (FSI)



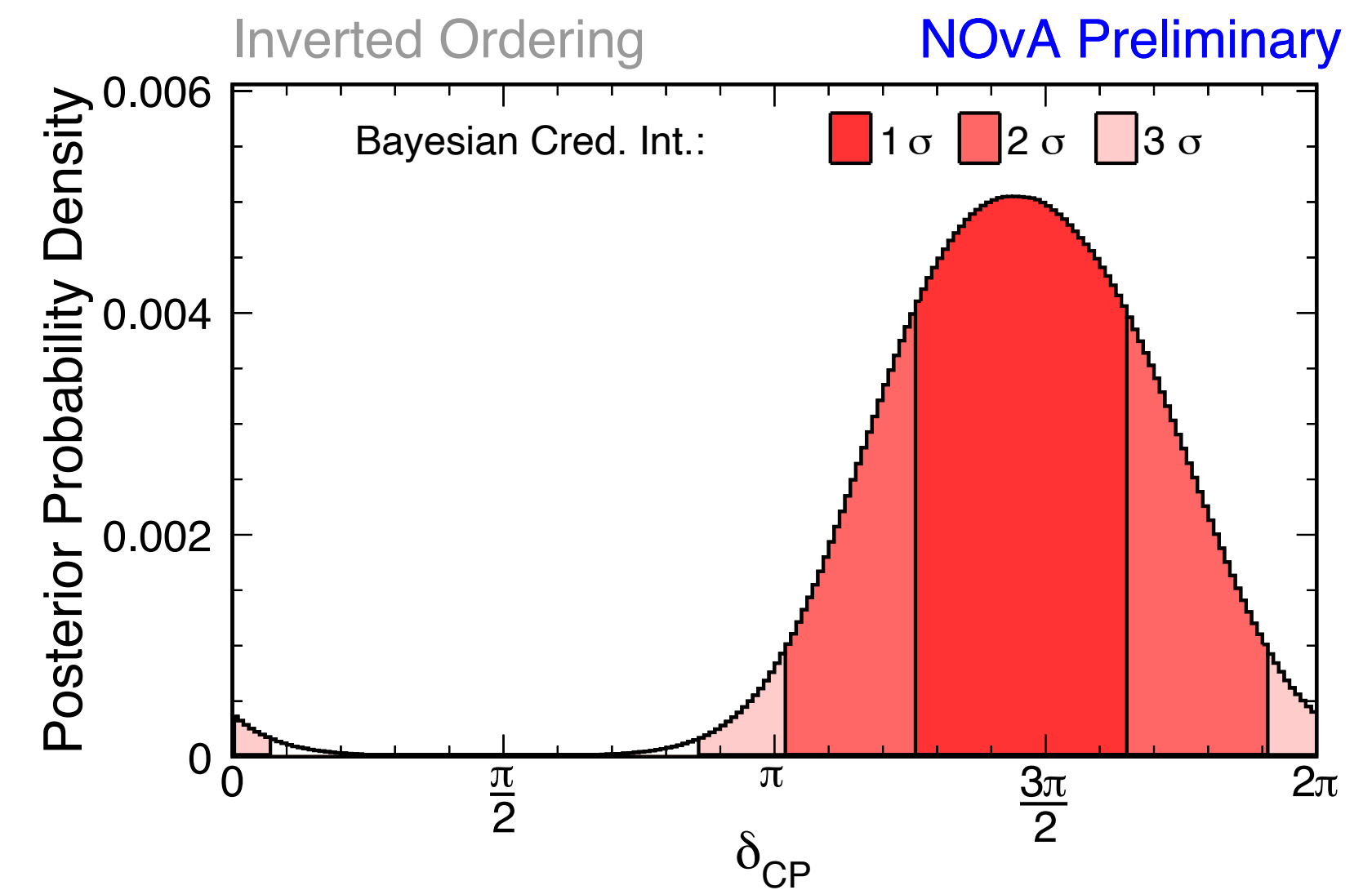
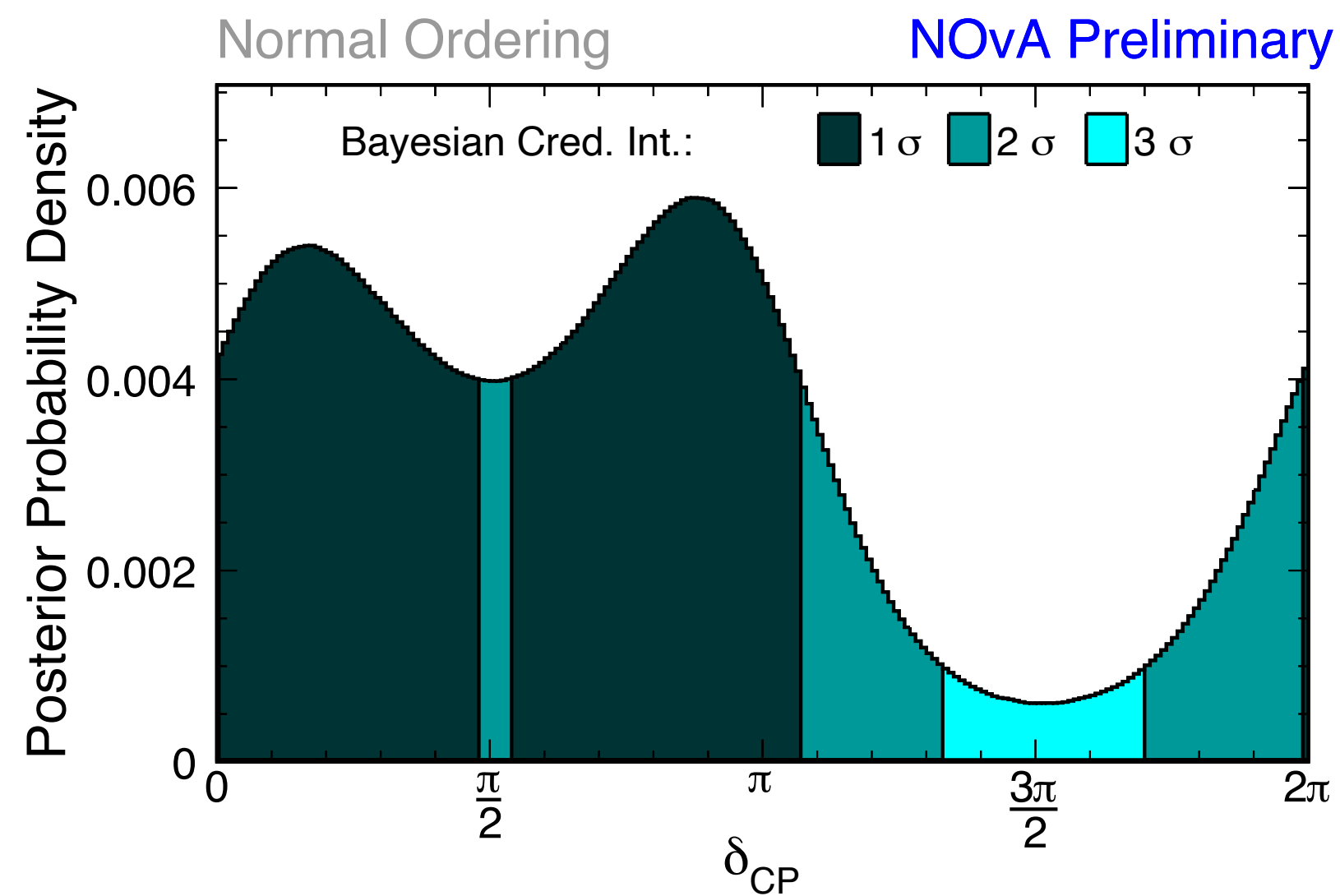


- “Disappearance parameters” =  $\theta_{23}$  and  $\Delta m^2_{32}$  (i.e. typical atmospheric parameters first seen by SK)
- Weak preference in upper octant normal mass ordering
- Bayes factor (ratio of steps):
  - $BF_{\frac{NO}{IO}} = 2.1$
  - $BF_{\frac{UO}{LO}} = 1.7$

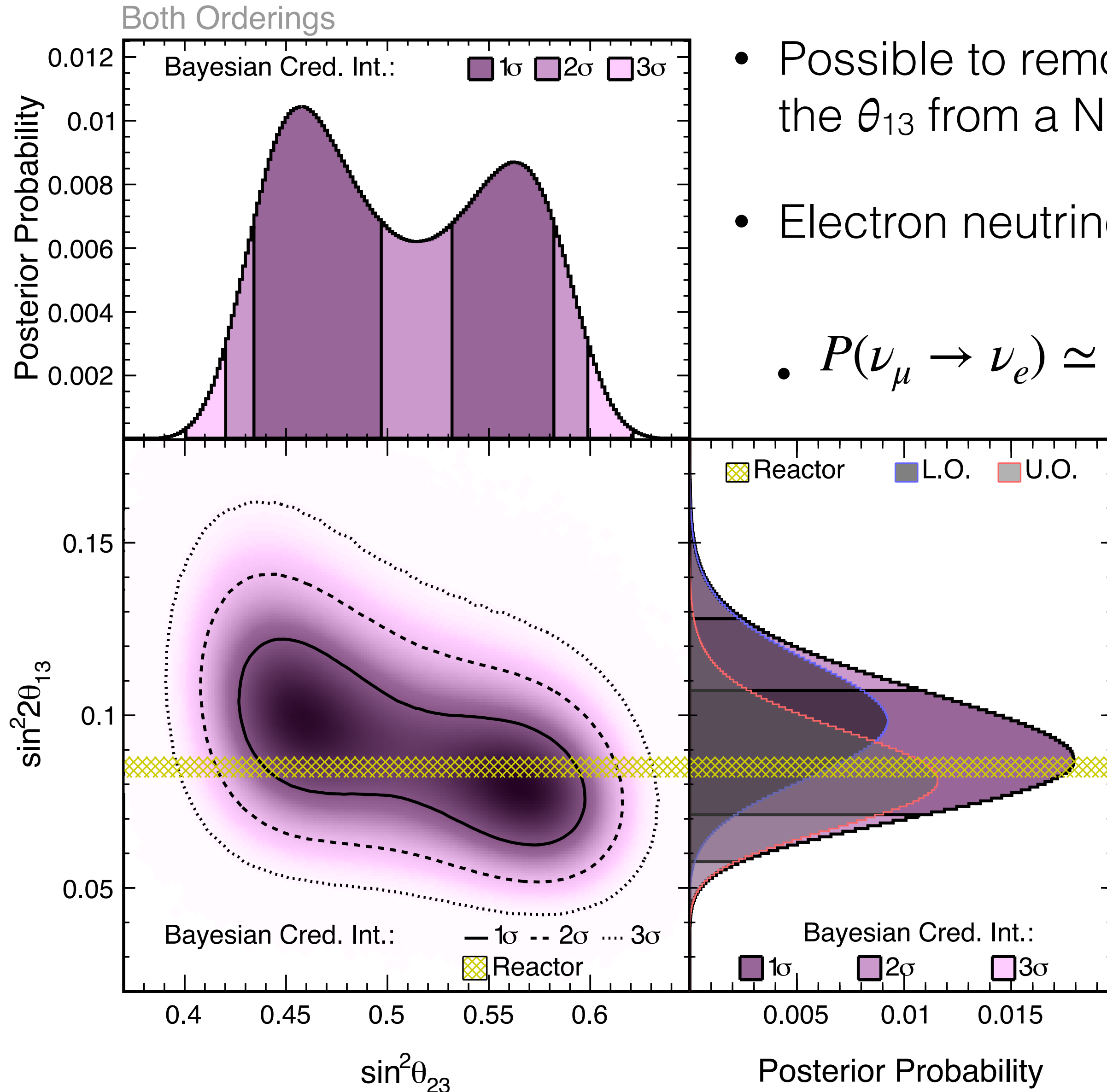
		Mass ordering		Total
		Normal	Inverted	
Octant	Upper	41.7%	20.9%	62.6%
	Lower	25.8%	11.5%	37.4%
Total		67.5%	32.5%	100%



- “Appearance parameters” =  $\theta_{13}$  and  $\delta_{CP}$ , although they can’t be observed without disappearance parameters...
- Strong correlation between mass hierarchy and  $\delta_{CP}$ 
  - Normal ordering:  $\delta_{CP} = 1.5 \pi$  outside the  $2 \sigma$  credible intervals (CI)
  - Inverted ordering:  $\delta_{CP} = 0.5 \pi$  outside the  $3 \sigma$  CI





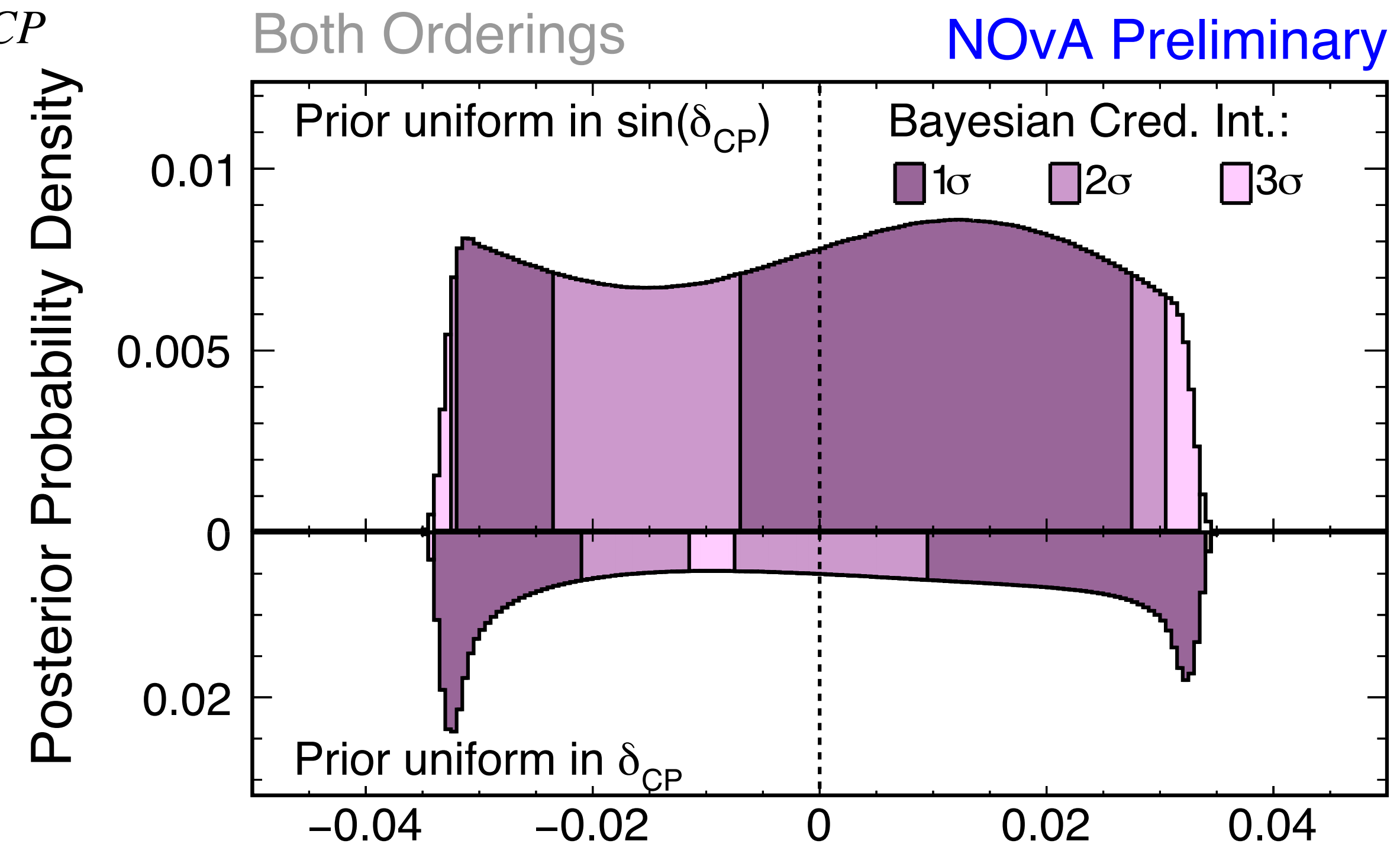


- Possible to remove the reactor constraint and check the posterior of the  $\theta_{13}$  from a NOvA only fit
- Electron neutrino appearance probability is:

$$P(\nu_\mu \rightarrow \nu_e) \simeq \sin^2 \theta_{23} \sin^2 2\theta_{13} \sin \frac{\Delta m_{32}^2 L}{4E} + f(\delta_{CP})$$

- So  $\theta_{23}$  (from  $\nu_\mu$  survival probability) and  $\theta_{13}$  (from  $\nu_e$  appearance probability) should be anti-correlated
- $\nu_e$  disappearance (reactor neutrino) and  $\nu_e$  appearance data (NOvA) have high degree of consistency

- Jarlskog invariant is a measure of CP violation, independent of the PMNS parameterisation
- $J = \cos \theta_{12} \cos^2 \theta_{13} \cos \theta_{23} \sin \theta_{12} \sin \theta_{13} \sin \theta_{23} \sin \delta_{CP}$
- Basic idea:
  - $J = 0 \rightarrow$  CP is conserved
  - $J \neq 0 \rightarrow$  CP is violated
- NOvA has weak sensitivity to  $J$ , so the choice of prior matters
  - Flat prior in  $\sin \delta_{CP}$  (related to what we see)
  - Flat prior in  $\delta_{CP} \rightarrow J \neq 0$  more preferred
- $\text{BF}_{\frac{J \neq 0}{J = 0}} = 1.5$  (from the Savage-Dickey method)





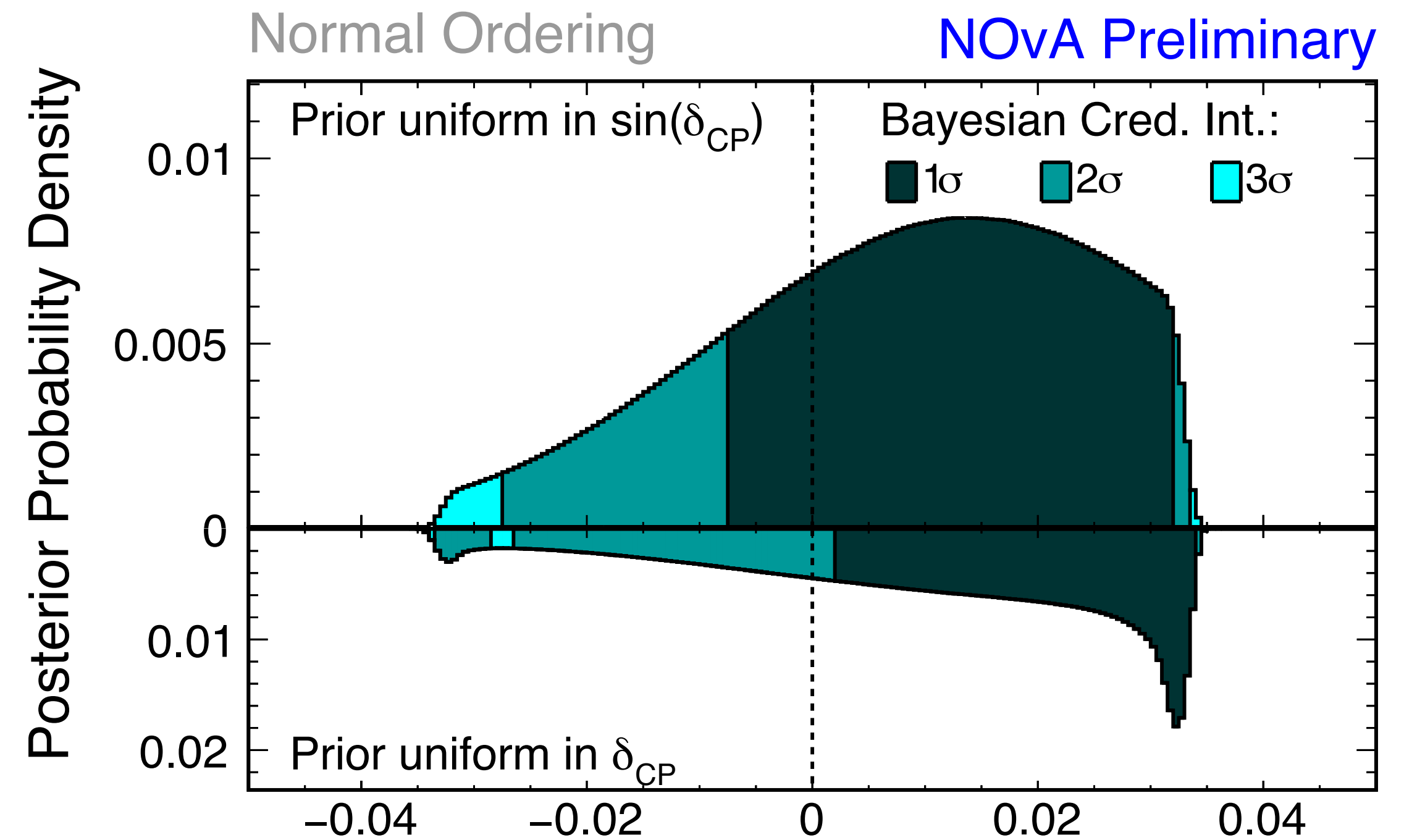
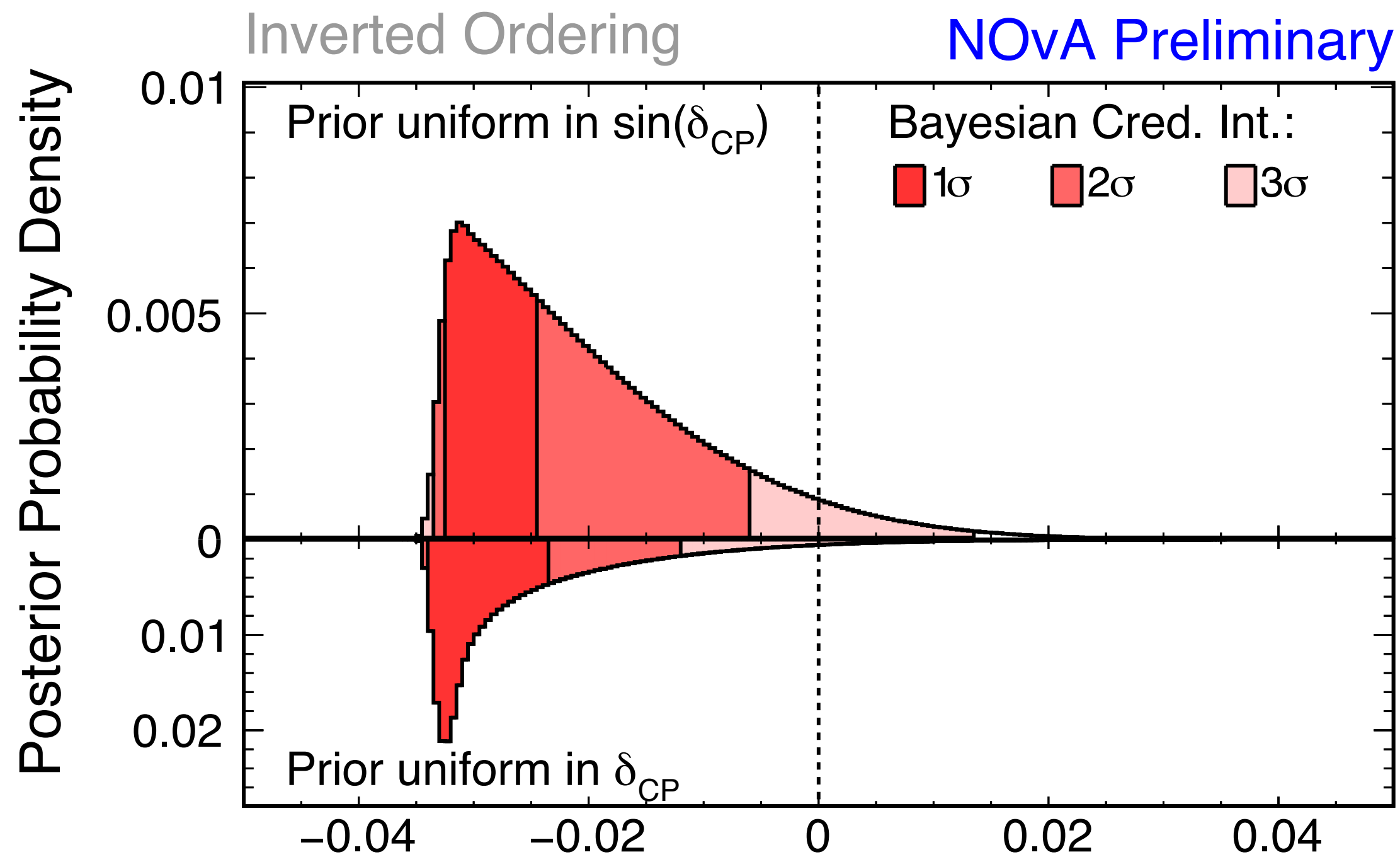
- Analysis of the NOvA data shows
  - No deviation from PMNS 3 neutrino flavour model
  - Marginal preference for the upper octant of  $\theta_{23}$  and normal neutrino mass ordering
  - No hints of CP violation in the neutrino sector
- Future of NOvA
  - More POT: we analysed  $26 \times 10^{20}$  POT, collected  $12 \times 10^{20}$  since 2020 dataset, and should get around  $60 \times 10^{20}$  by the end of NOvA
  - Test beam
  - Joint analysis with T2K

*Thanks for your attention!*

# Backup

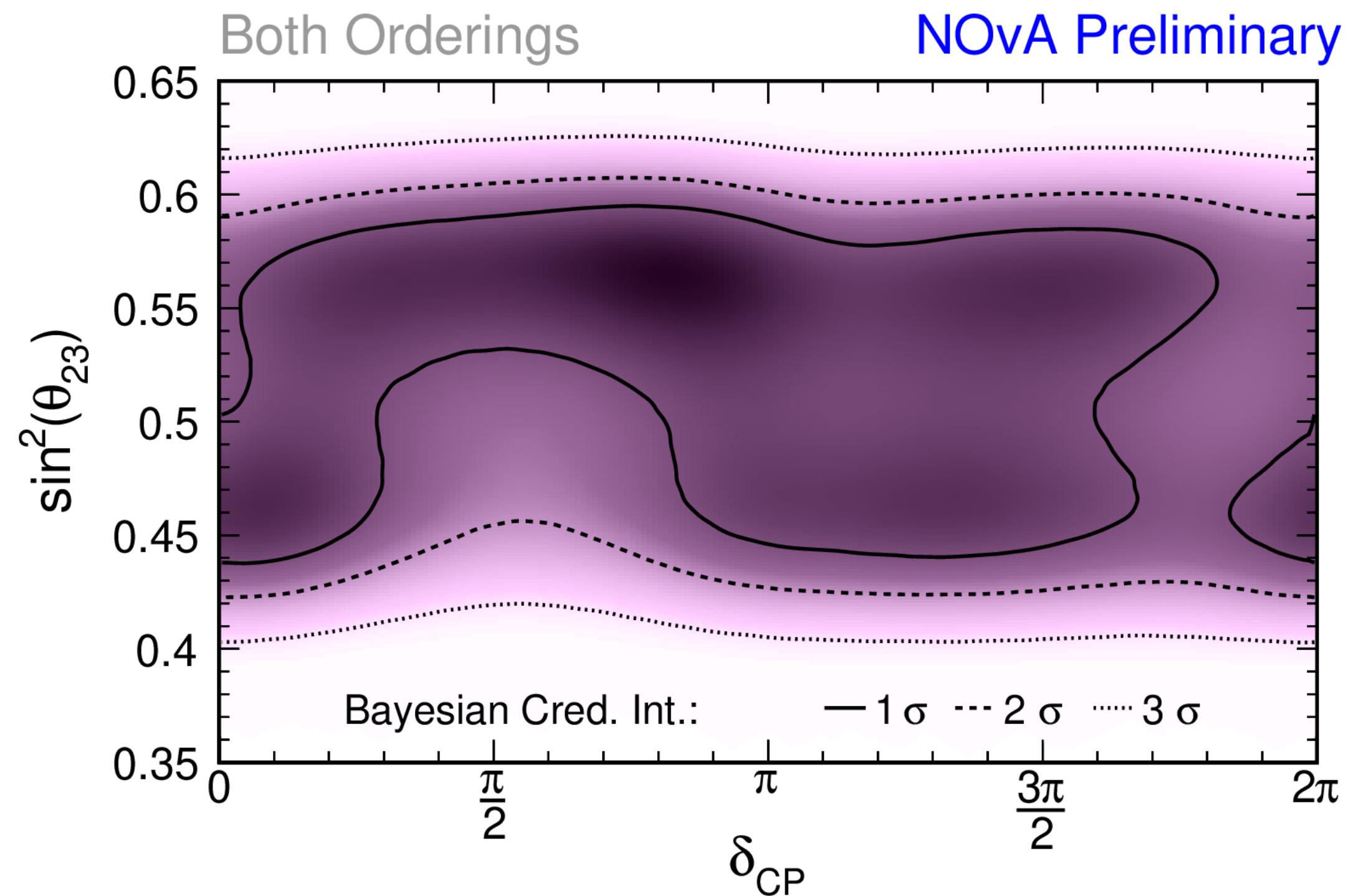


- Sign of  $J$  is governed by the sign of  $\sin \delta_{CP}$
- In inverted ordering:  $2\sigma$  exclusion of  $J = 0$
- In normal ordering no exclusion of  $J = 0$

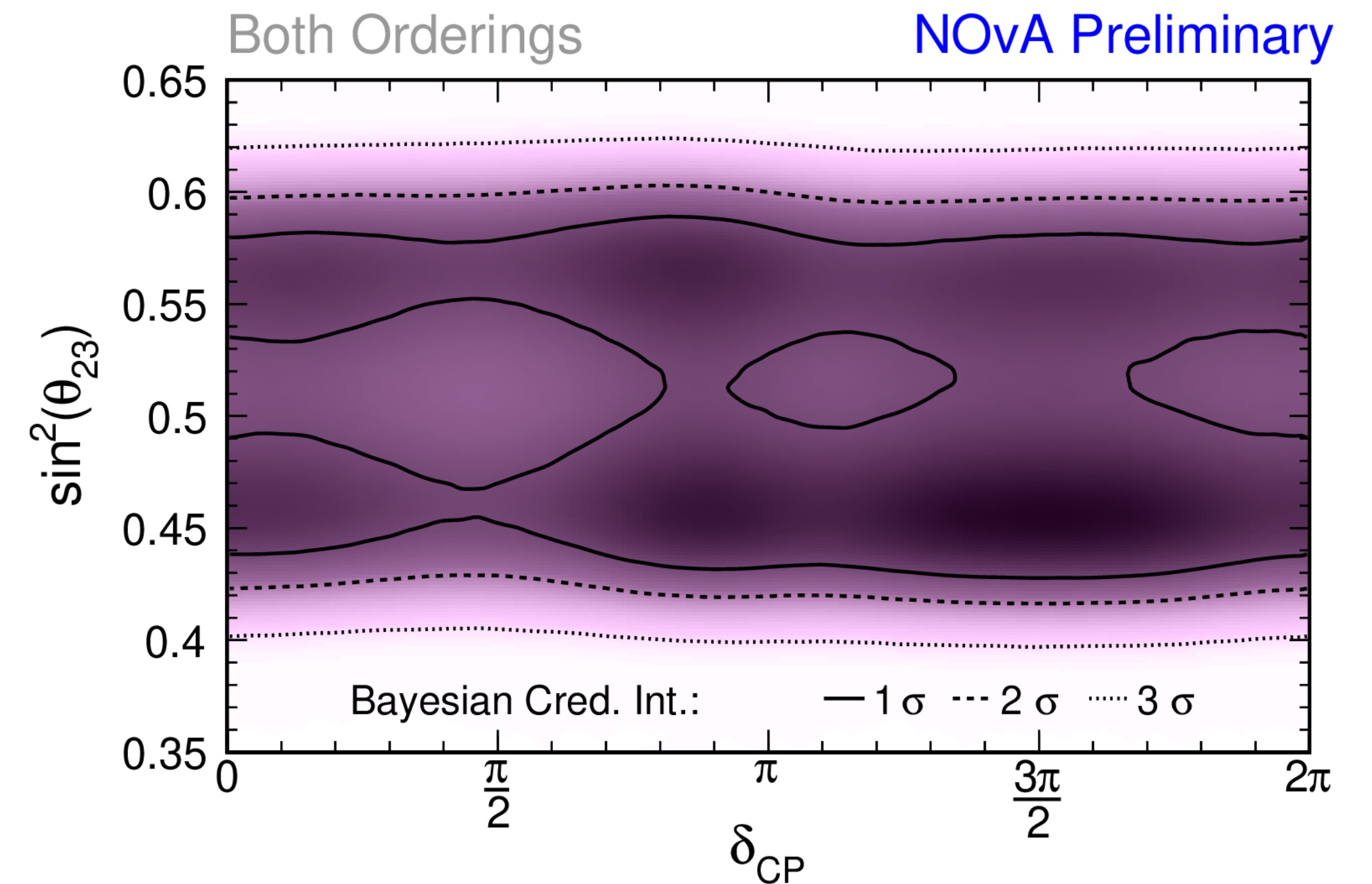


- Typically long baseline neutrino experiment use the reactor neutrino experiment's measurement on  $\theta_{13}$  to enhance their sensitivity to  $\delta_{CP}$

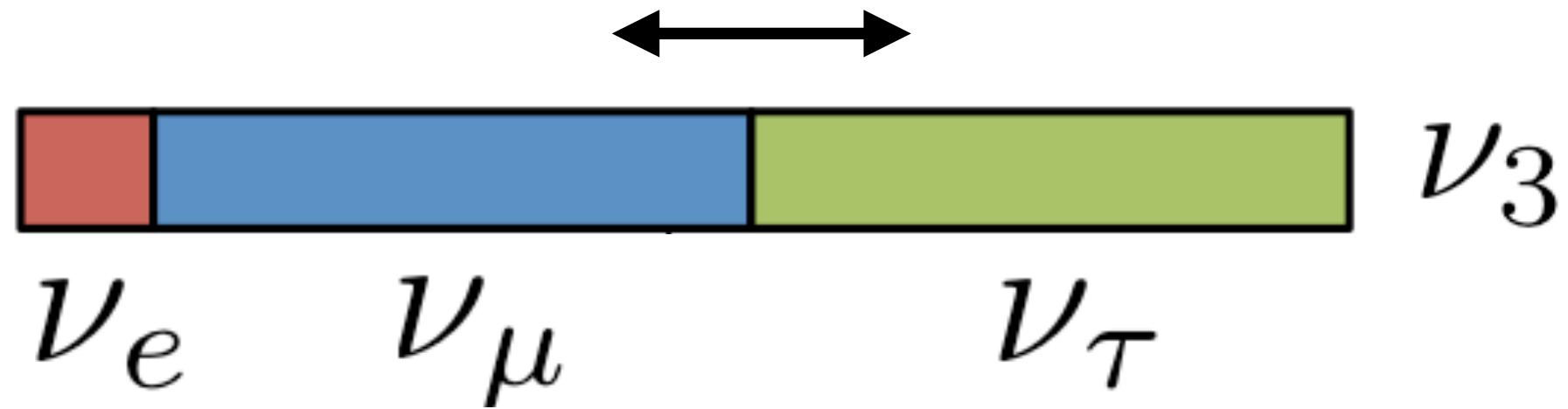
With reactor constraint



Without reactor constraint

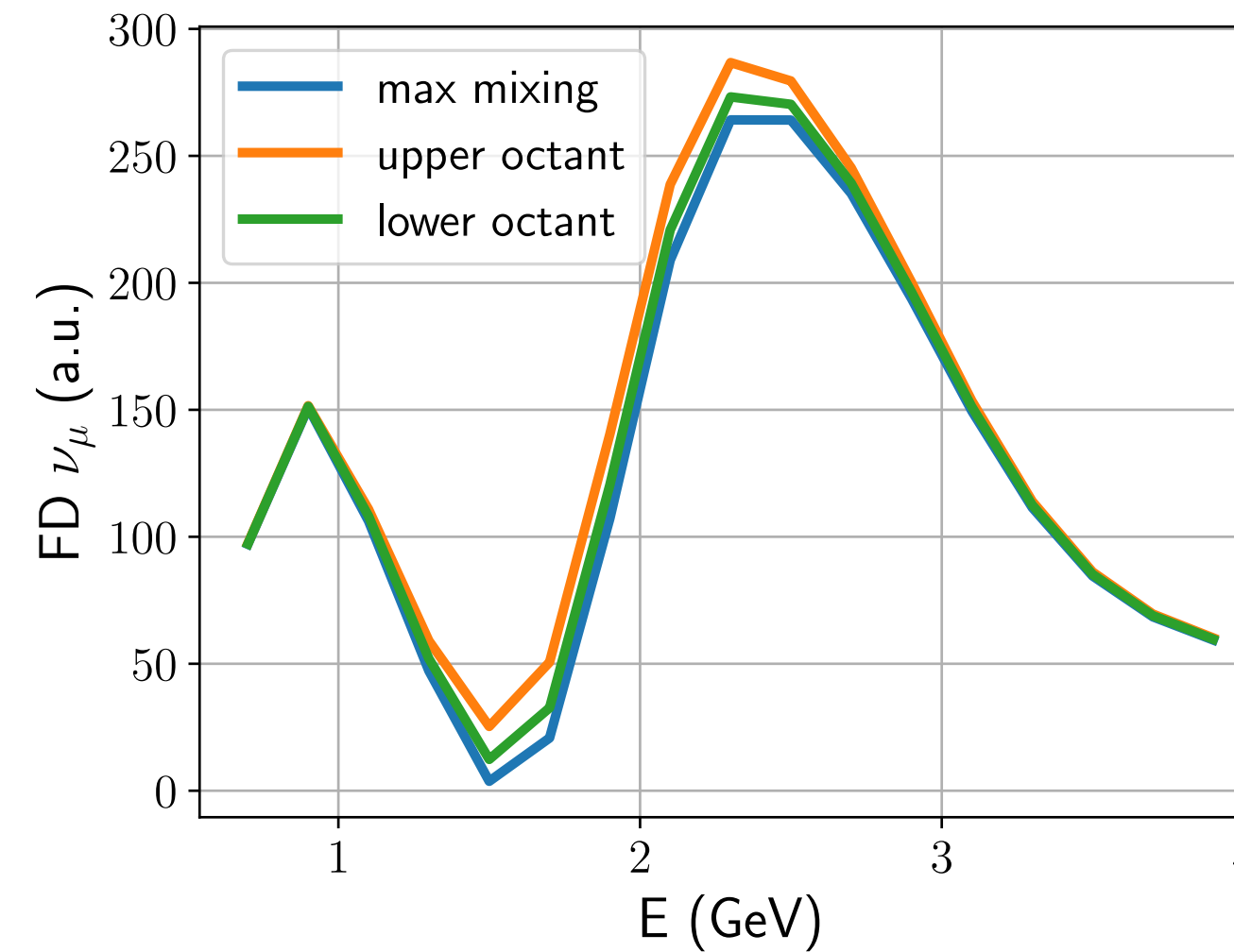






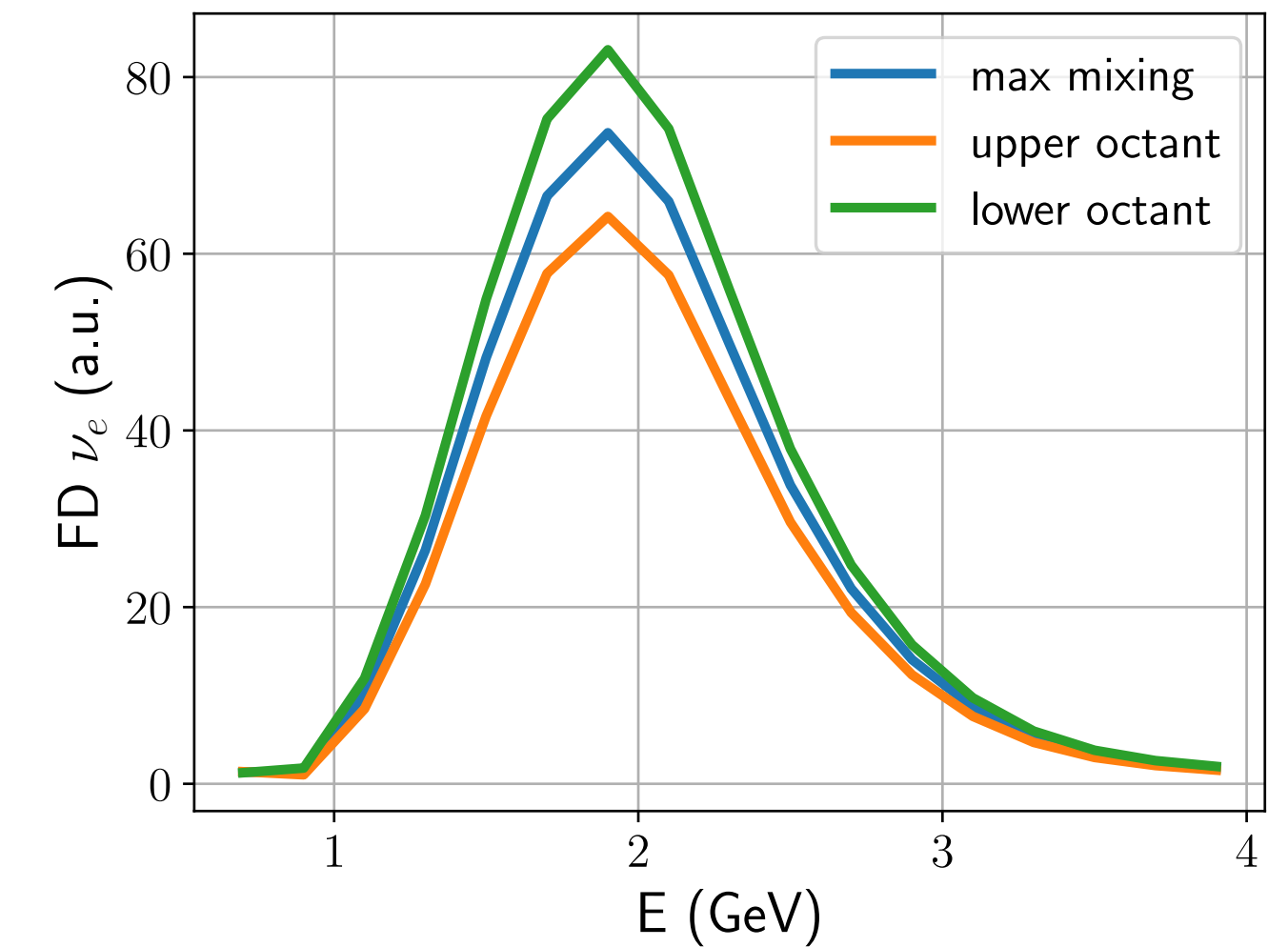
True variables

- Effect of the octant ( $\theta_{23}$  bigger or smaller than  $45^\circ$ )
  - Same effects in FHC/RHC
  - $\nu_\mu$  channel:
    - Maximal disappearance at  $\theta_{23} = 45^\circ$  any movement away from that  $\rightarrow \nu_\mu$  disappear less
      - Disappearance  $\propto \sin^2 2\theta_{23}$
  - $\nu_e$  channel:
    - Sensitive to octant
      - Appearance  $\propto \sin^2 \theta_{23}$



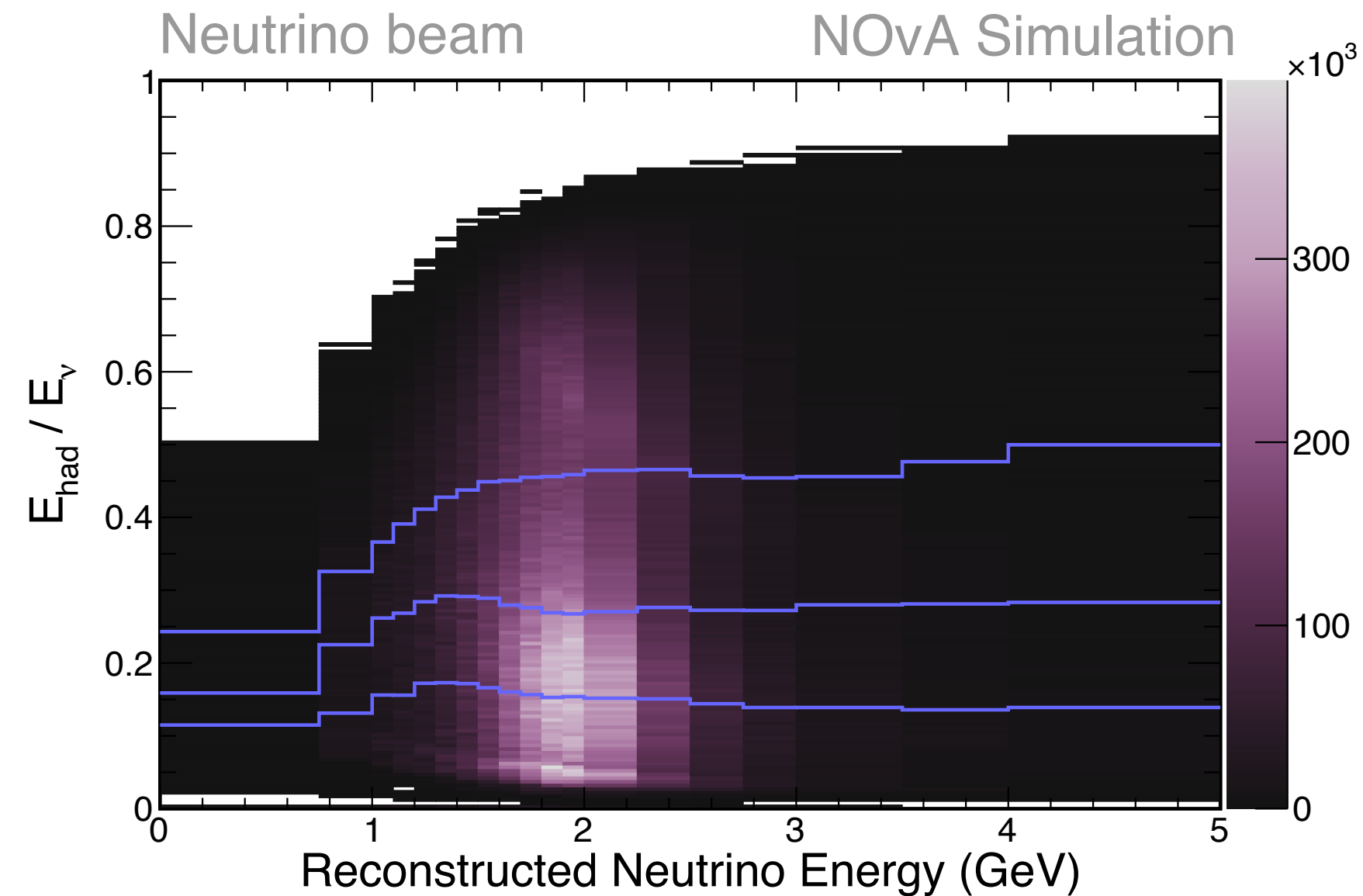
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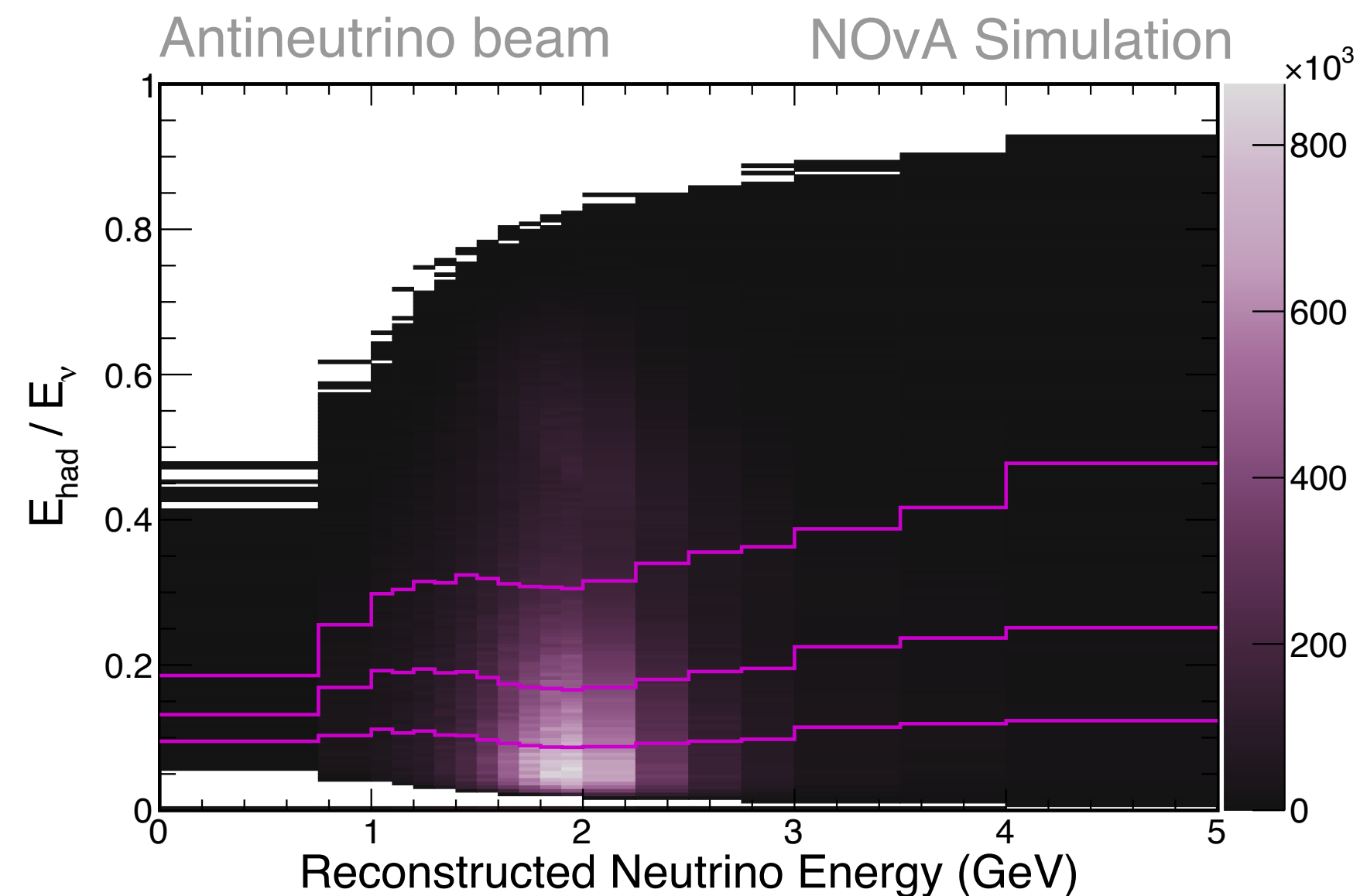


$$P(\nu_\mu \rightarrow \nu_e) \times \sigma_{\nu_e} \times \Phi_{\nu_\mu}$$

$$P(\nu_\mu \rightarrow \nu_e) \simeq \sin^2 \theta_{23} \sin^2 2\theta_{13} \sin \frac{\Delta m_{32}^2 L}{4E} + f(\delta_{CP})$$

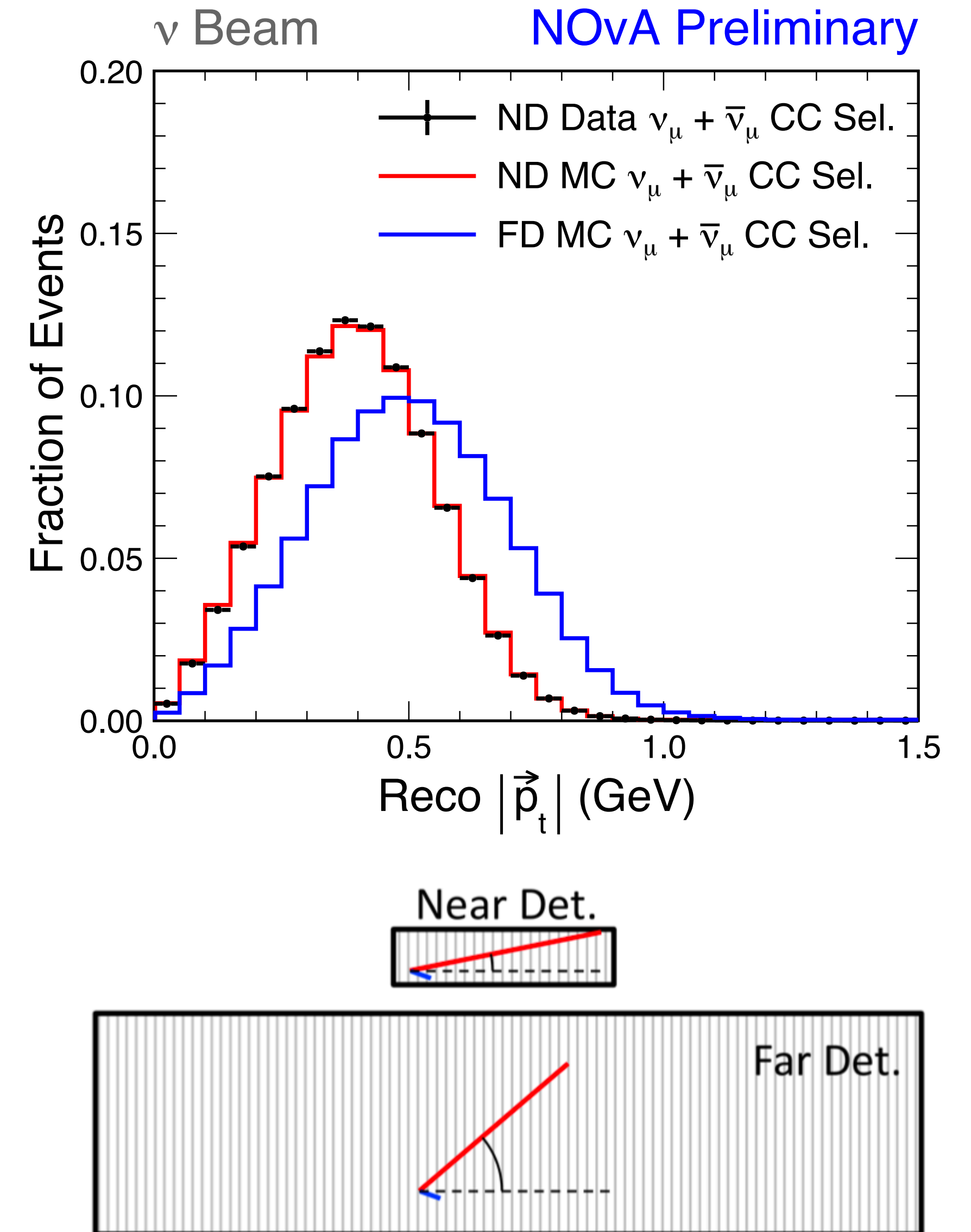
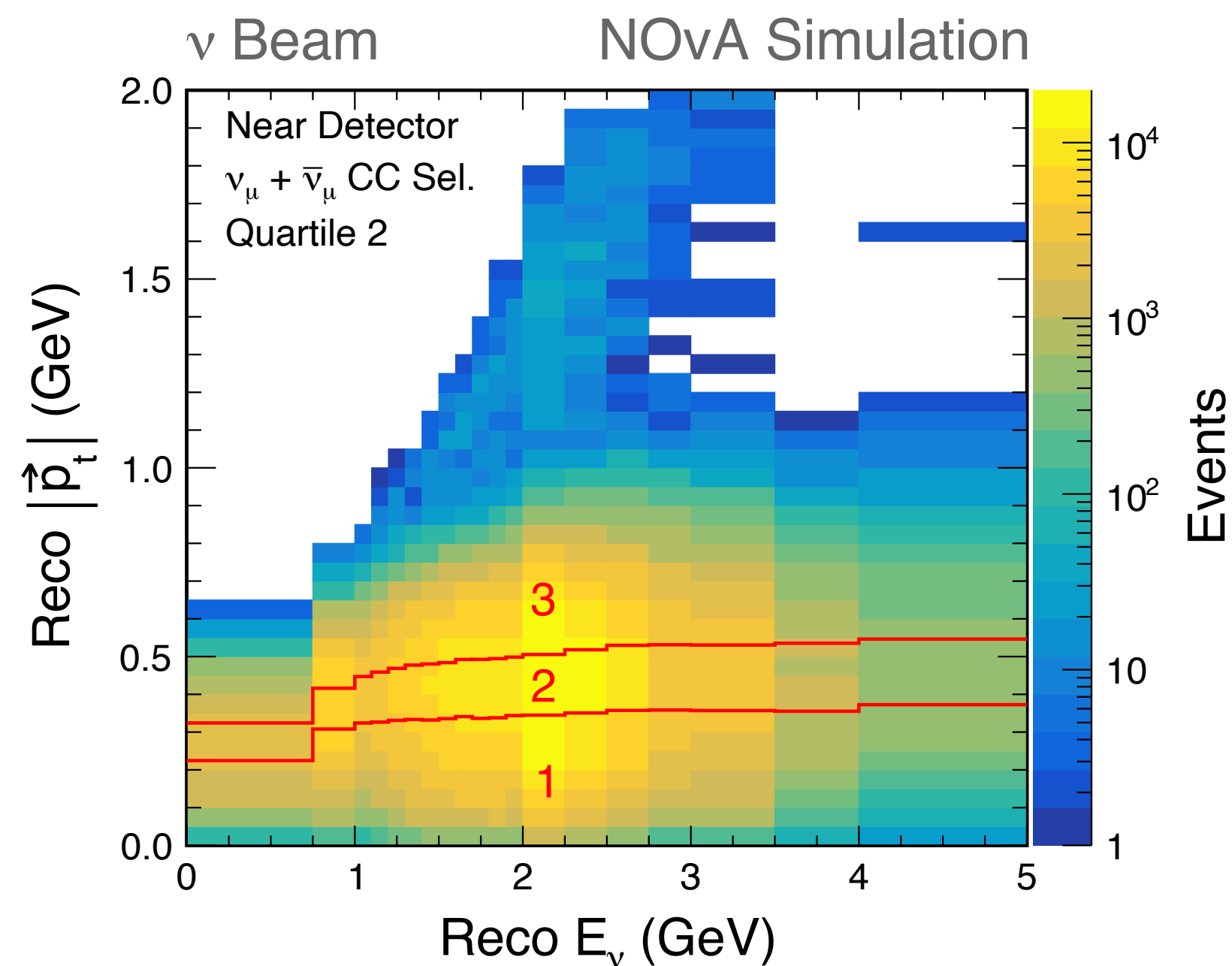


- Energy reconstruction depends on the hadronic energy fraction
- Split the hadronic energy fraction in quartiles
- Extrapolate ND  $\rightarrow$  FD each of them independently

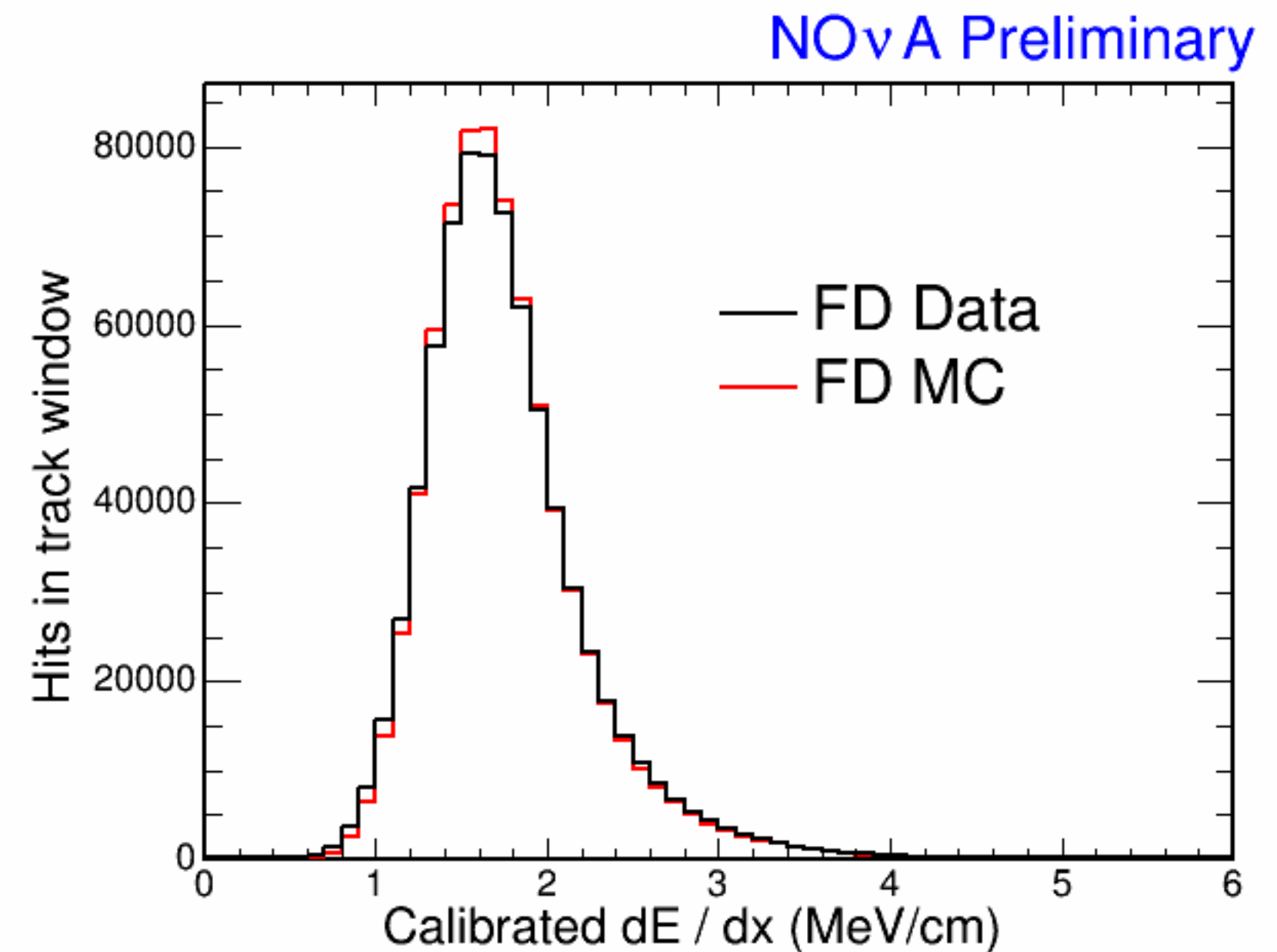
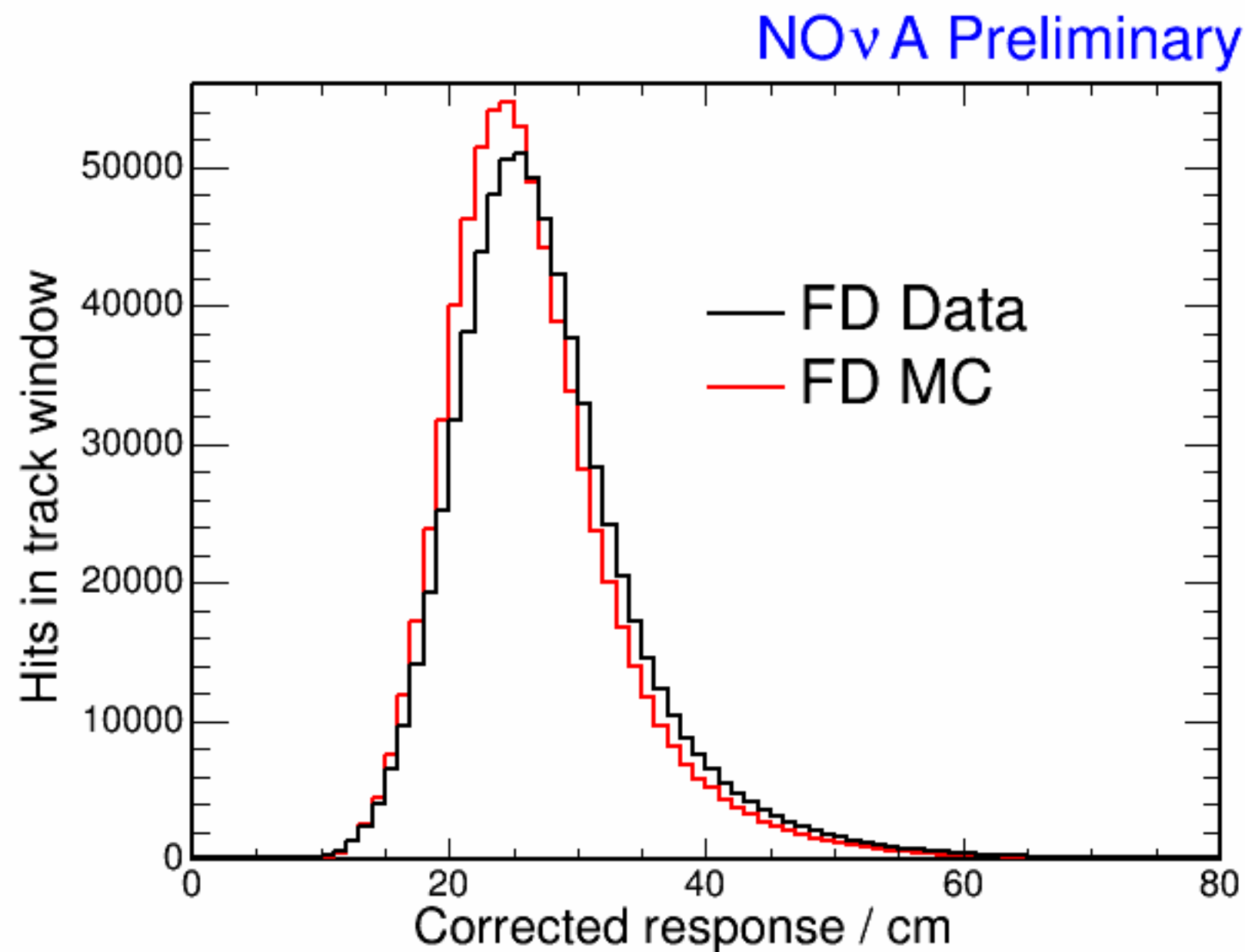




- ND and FD have different sizes
- Transverse muon momentum is different
- Separate into 3  $p_T$  bins and extrapolate each of them independently

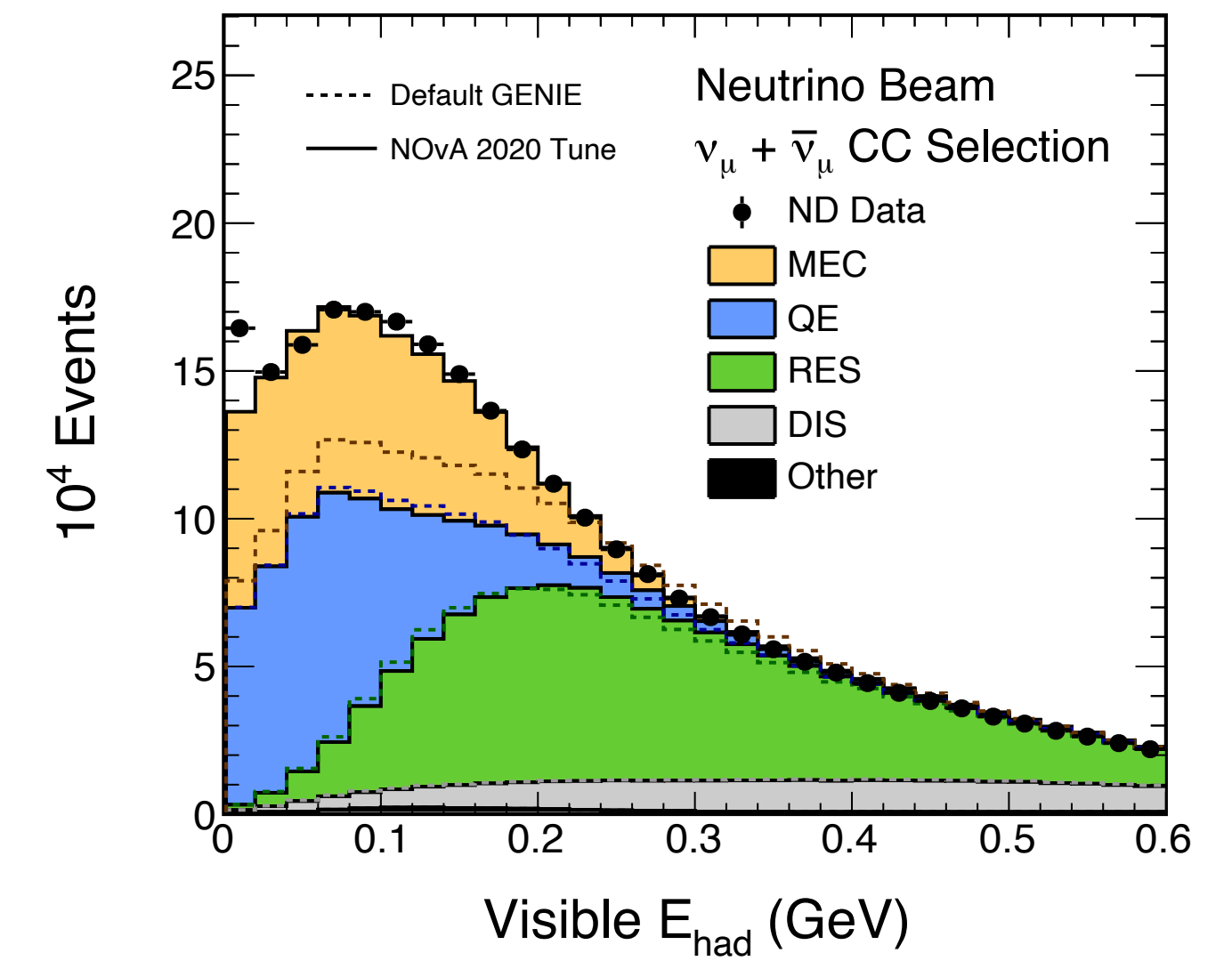
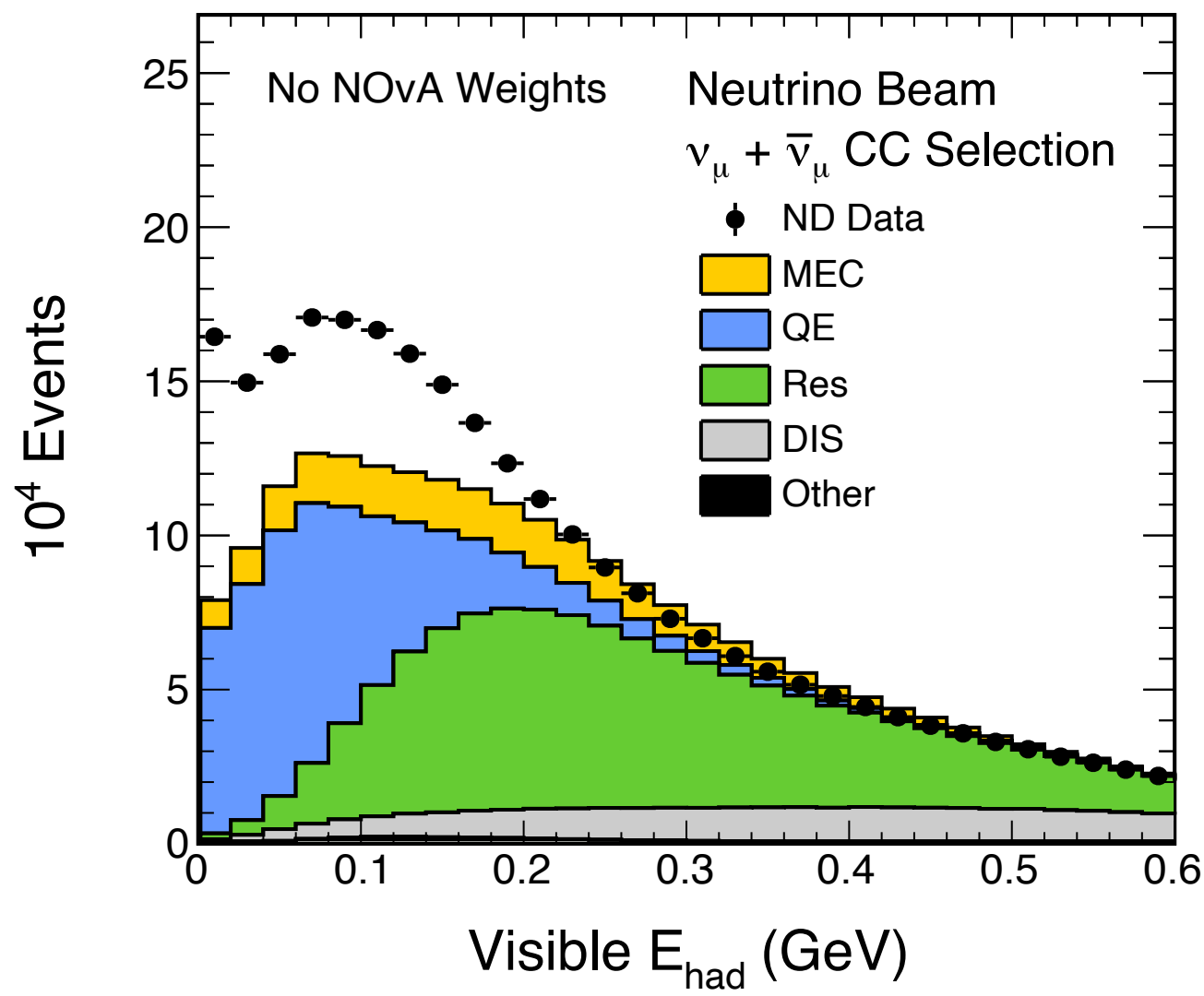
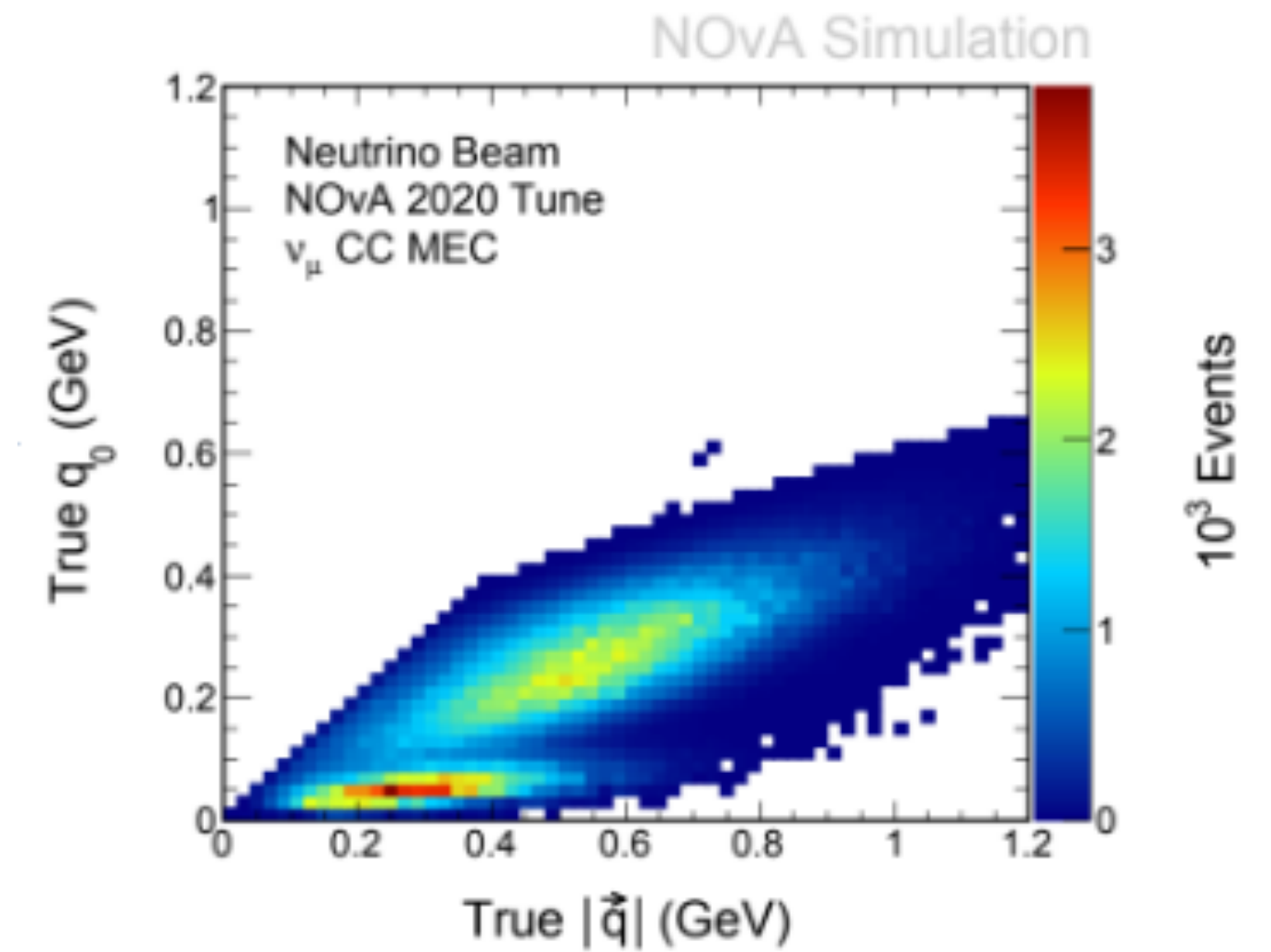
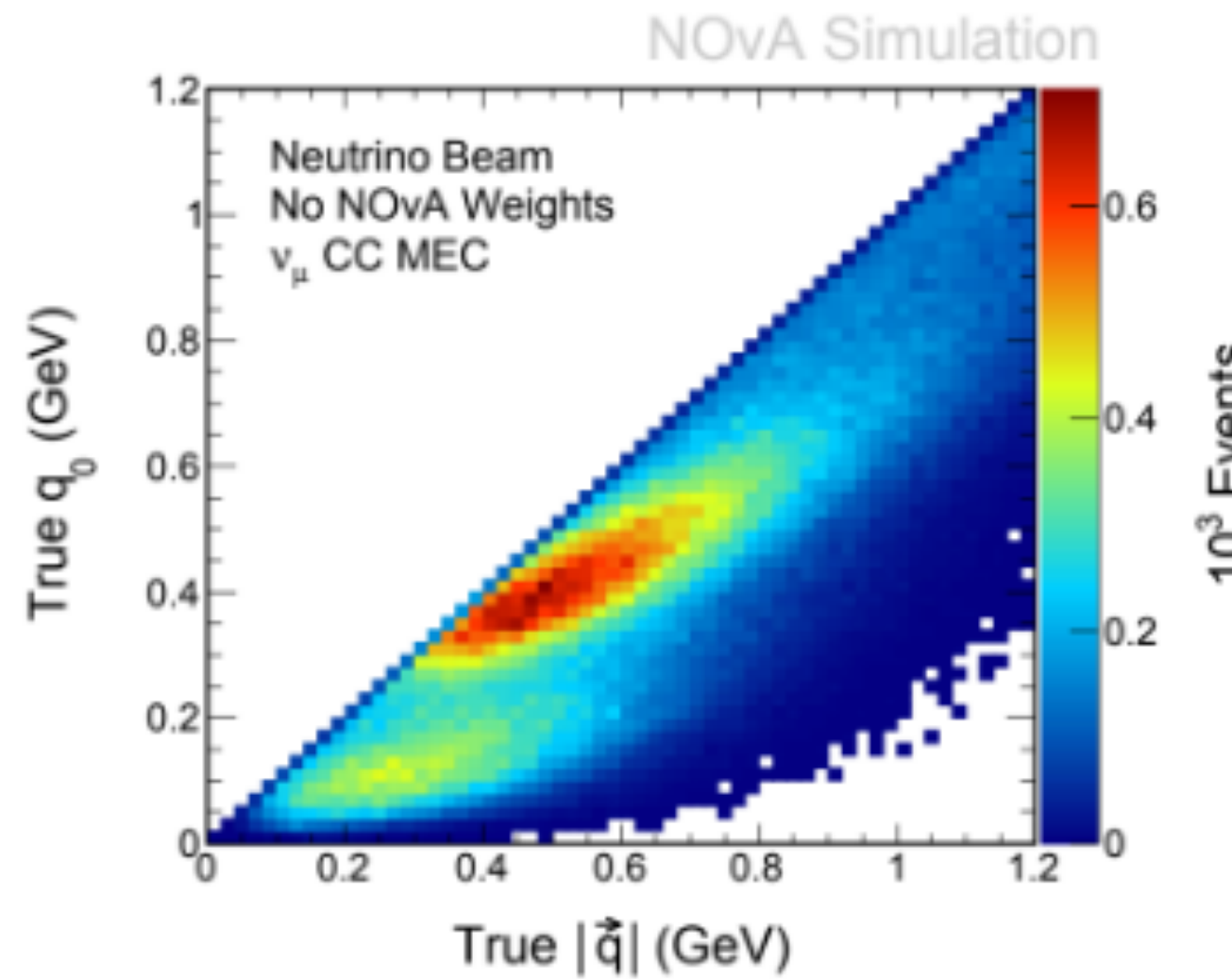
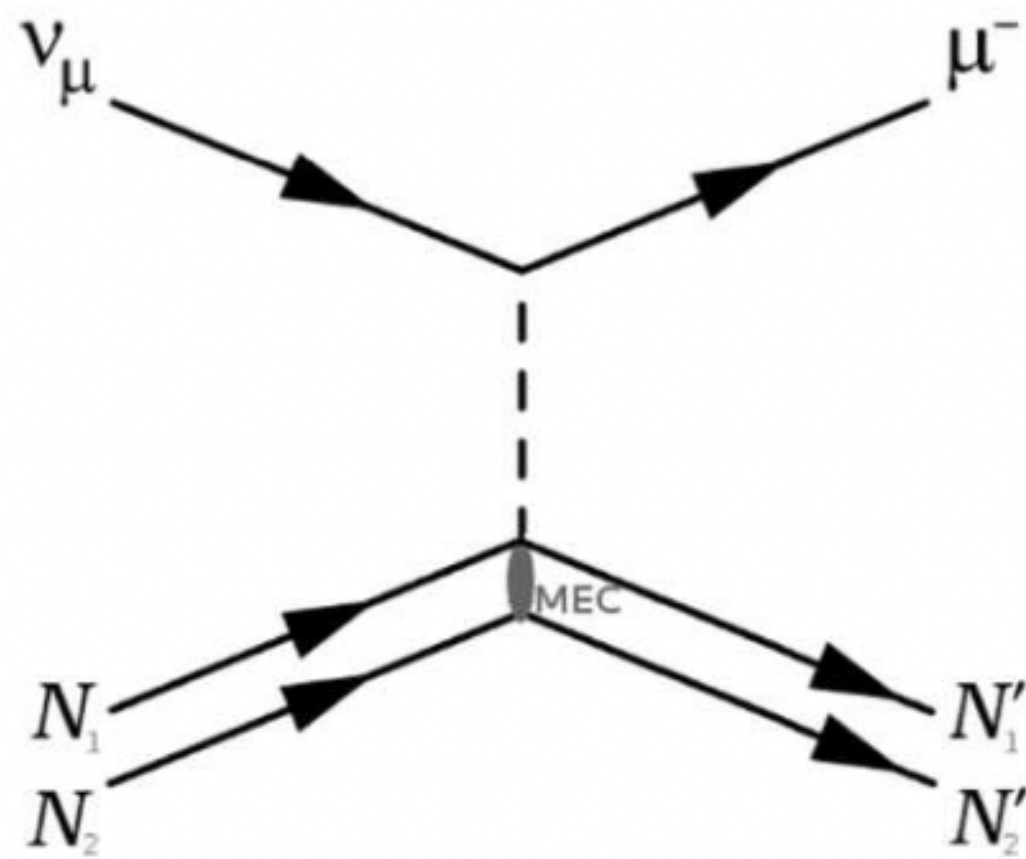


- Done with cosmic muons
  - Select stopping muons
  - Measure distance from the end of the track
  - Compare dEdx for MC and data and adjust
- Note these are plots from 2017! Got better since
- Limitation
  - Hadronic energy deposition is quite different from MIP!
- Hoping that the test beam can help us





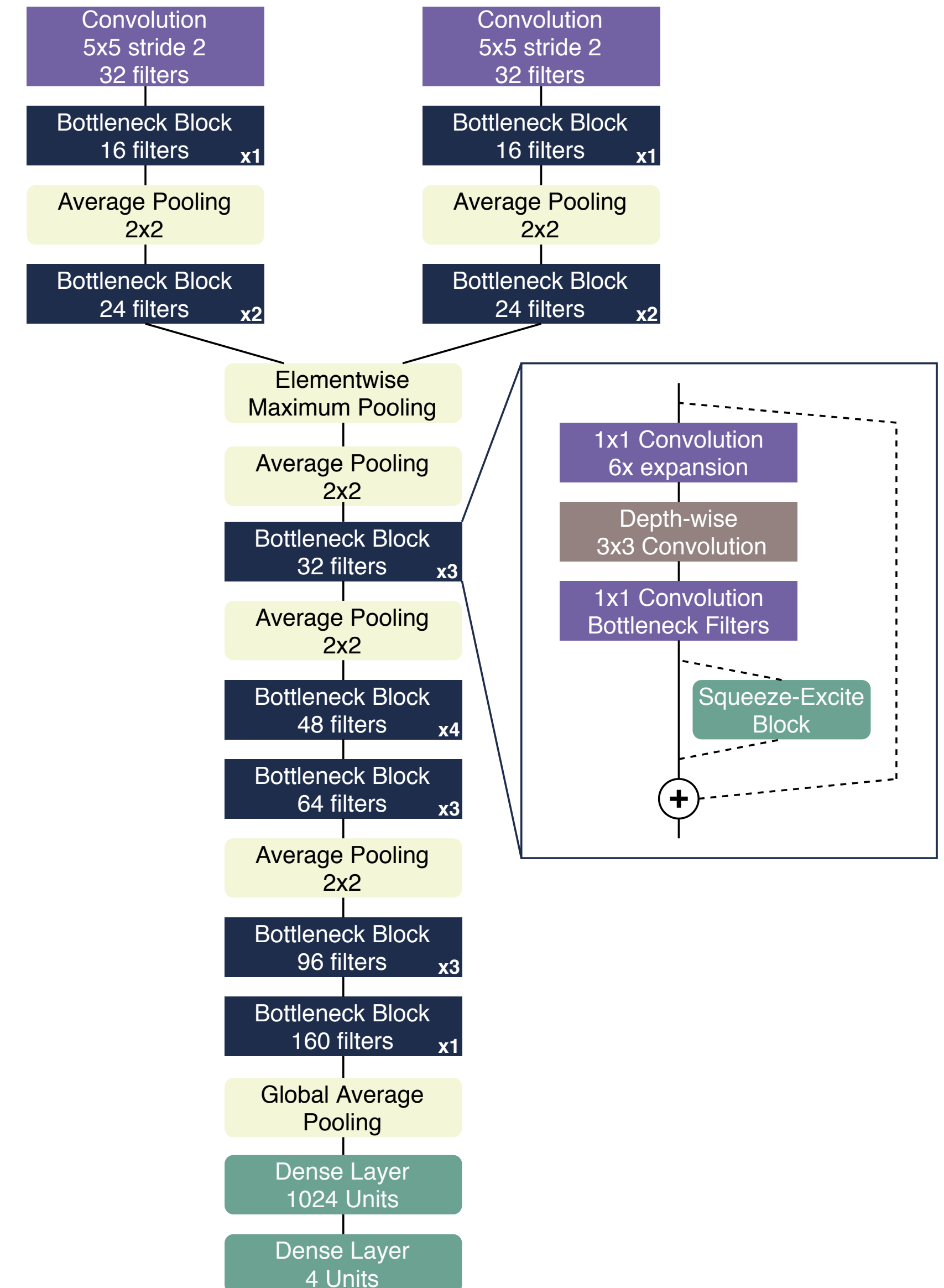
- $\sigma_{\text{MEC}}$  needs a huge increase to agree with ND data
- Not theory motivation
- “Good” answer requires a lot of work from both phenomenology and experiment data



GENIE

NOvA's GENIE

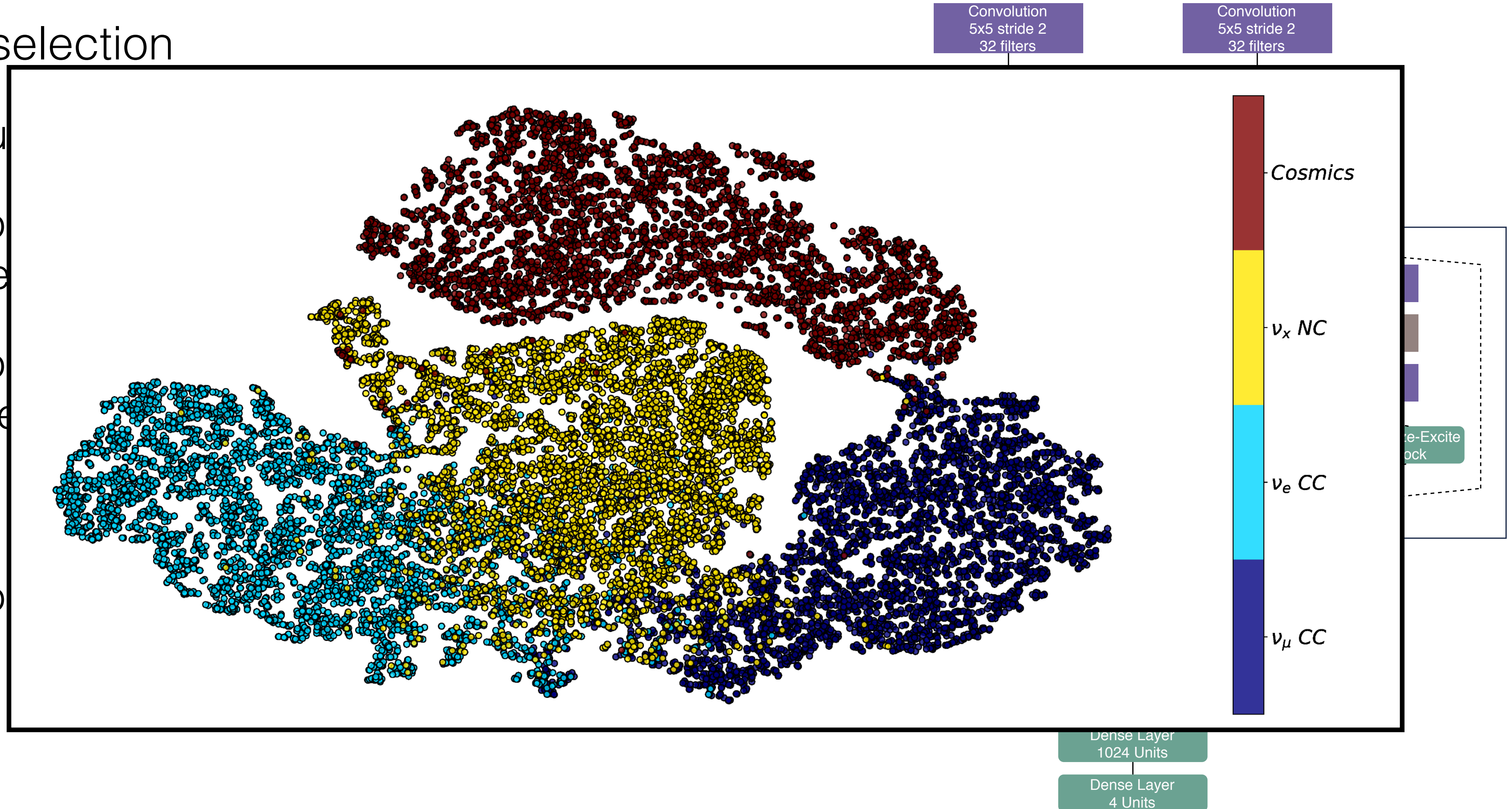
- Preselection
  - Quality cut (detector was running normally)
  - Contained cut (nothing seems to be exiting the detector)
  - Cosmic rejection Boosted Decision Tree (the event doesn't look like a cosmic ray)
- PID with Convolutional Neural Network (CNN)
  - Topology-based algorithm
  - High purity samples





- Preselection

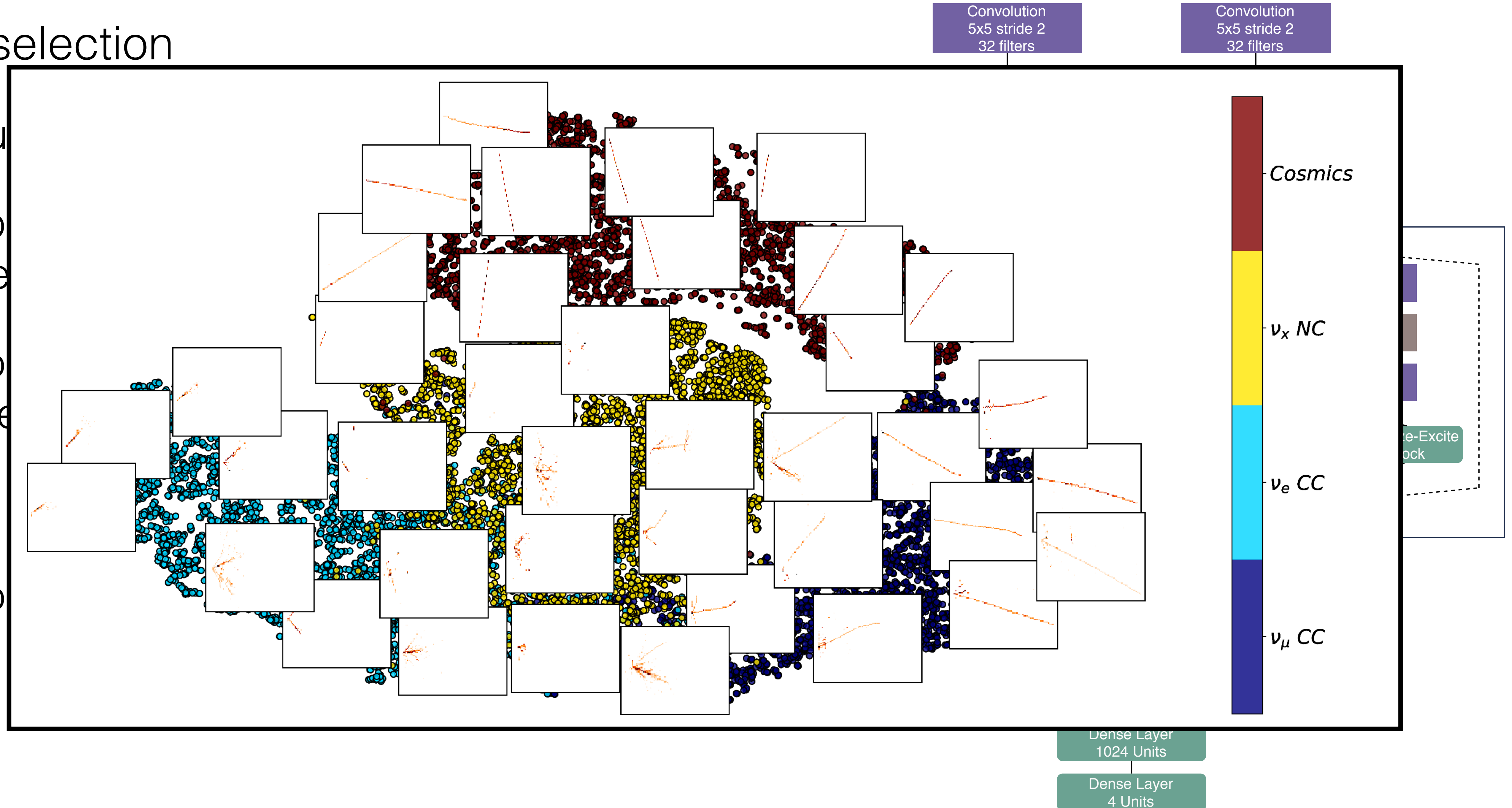
- Qu
- Co  
the
- Co  
eve
- PID
- Co





- Preselection

- Qu
- Co  
the
- Co  
eve
- PID
- Co



# CVN confusion

